

A new probe into three-nucleon-force effects on reaction observables

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Microscopic approach to many-body reactions

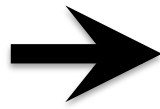
- ✓ Many-body nuclear direct reactions
- ✓ Microscopic reaction theory based on *NN effective interactions (g-matrix)*

founded on multiple scattering theory

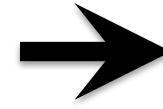
K. M. Watson, RMP30, 565 (1958).

M. Yahiro et al., PTP120, 767 (2008).

Realistic force
(Chiral EFT)



NN g-matrix
(Brueckner theory)



Many-body reactions
(folding model, ...)

N^3LO 2NF

N^2LO 3NF

E. Epelbaum et al., NPA747, 362 (2005); RMP81, 1773 (2009)

There is no ad hoc parameter!

Microscopic reaction theory is enable us to investigate exotic structures and (effective) interactions.

Elastic and inelastic scattering

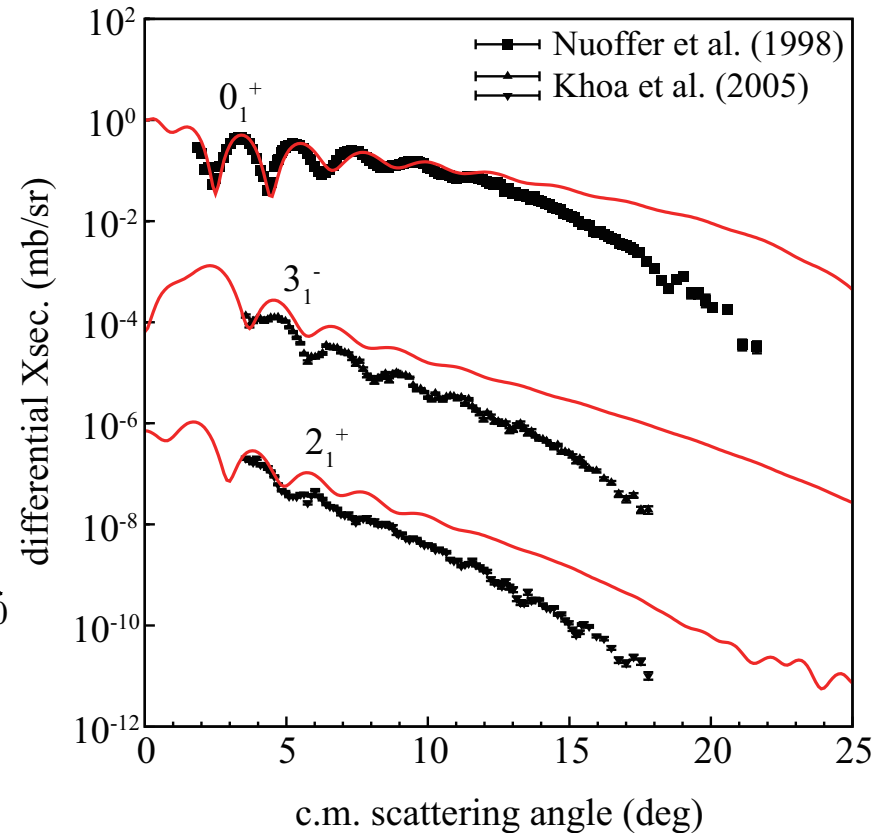
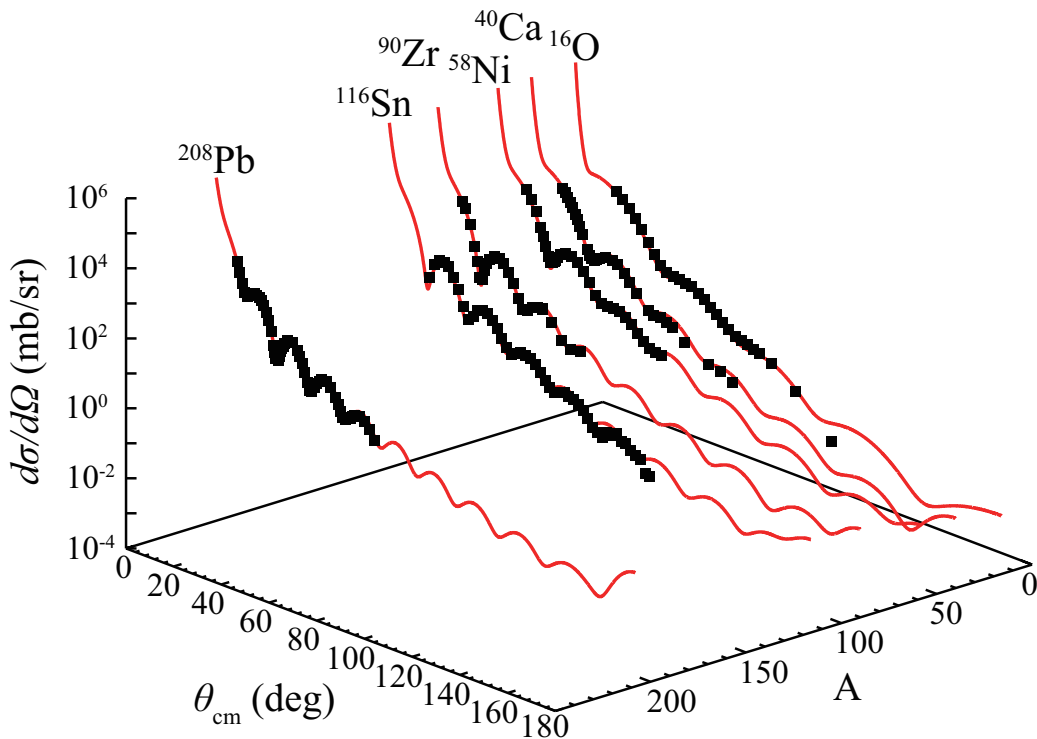
Folding (microscopic coupled-channels) calculations

The nonlocality coming from knockon exchange process is localized.

KM et al., JPG37, 085011 (2010).

^{16}O - ^{16}O @70MeV/nucl.

p -scattering@65MeV



Note: “usual” folding potential does not work at **lower energies**.

cf.) Dispersive folding model J. Mueller et al., PRC83, 064605 (2011).

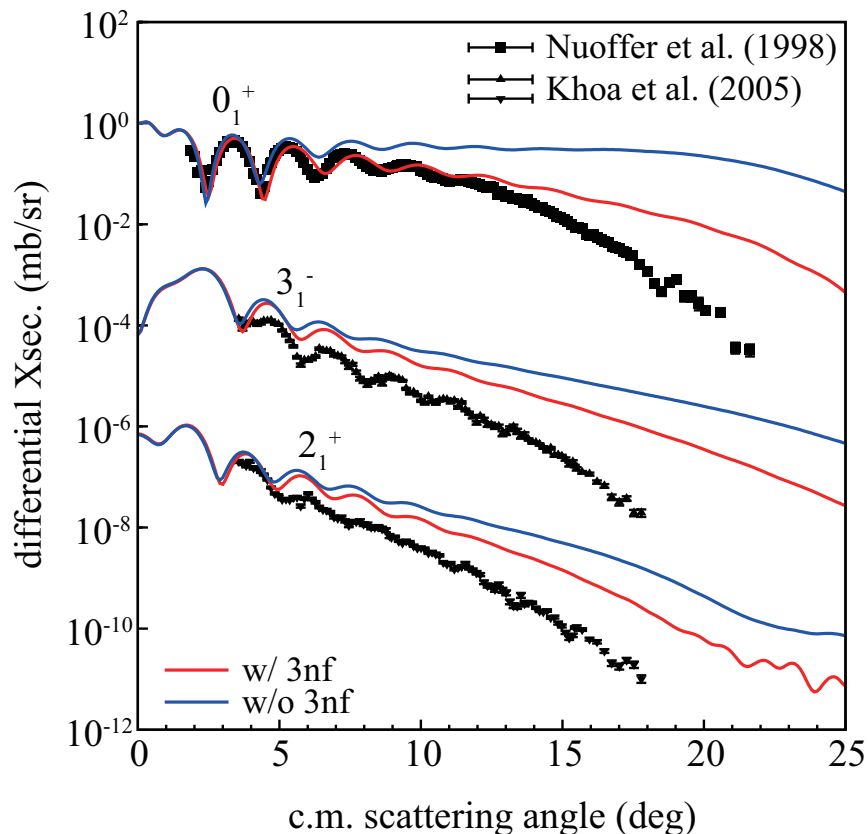
3NF effects on scattering observables

Nucleus-nucleus scattering

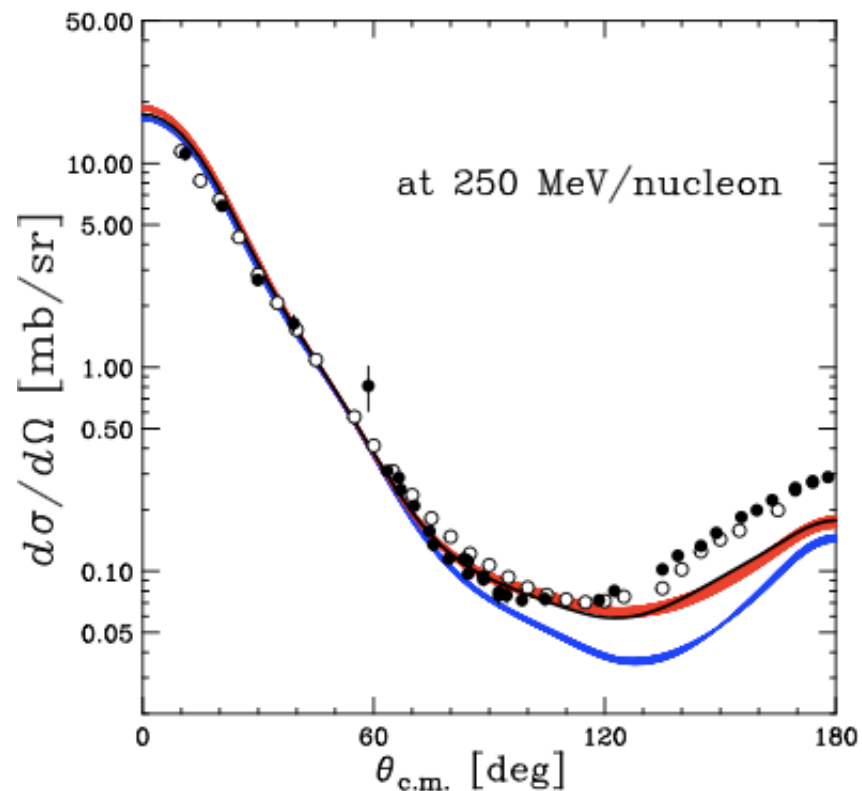
dp scattering

with the frozen density approx. ($\rho = \rho_P + \rho_T$)

KM et al., PRC93, 014607 (2016).



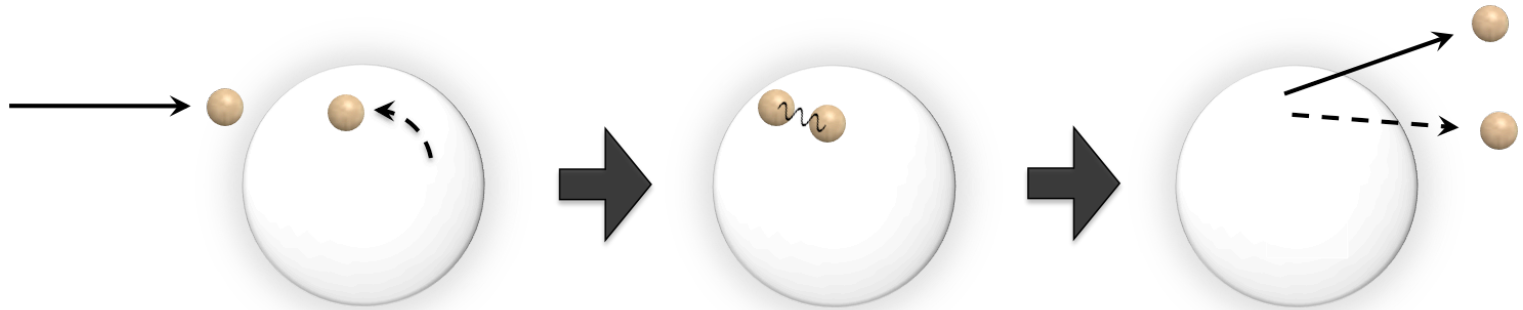
K. Sekiguchi et al., PRC89, 064007 (2014).



Can we investigate 3NF effects by using other reactions?

Knockout reactions as a probe into 3NF

Proton knockout reaction ($p,2p$)



$(p,2p)$ reaction occurs in nuclear interior so that 3NF effects are probed through the density dependence of g -matrix.

Motivations

- ✓ Examine the microscopic approach to knockout reactions
- ✓ Investigate the 3NF effects on many-body reactions

Microscopic DWIA

- ✓ Transition matrix element in the distorted wave Impulse Approx.

$$T = \langle \chi_{1,\mathbf{k}_1}^{(-)} \chi_{2,\mathbf{k}_2}^{(-)} | g(\kappa', \kappa, \theta; E, \rho) | \chi_{0,\mathbf{k}_0}^{(+)} \varphi_{nlj} \rangle$$

$g(\kappa', \kappa, \theta; E, \rho)$: **chiral g-matrix**

$\chi_{0,\mathbf{k}_0}^{(+)}$, $\chi_{1,\mathbf{k}_1}^{(-)}$, and $\chi_{2,\mathbf{k}_2}^{(-)}$: distorted waves

φ_{nlj} : single particle wave function

Microscopic DWIA

- ✓ Transition matrix element in the distorted wave Impulse Approx.

$$T = \langle \chi_{1,\mathbf{k}_1}^{(-)} \chi_{2,\mathbf{k}_2}^{(-)} | g(\kappa', \kappa, \theta; E, \rho) | \chi_{0,\mathbf{k}_0}^{(+)} \varphi_{nlj} \rangle$$

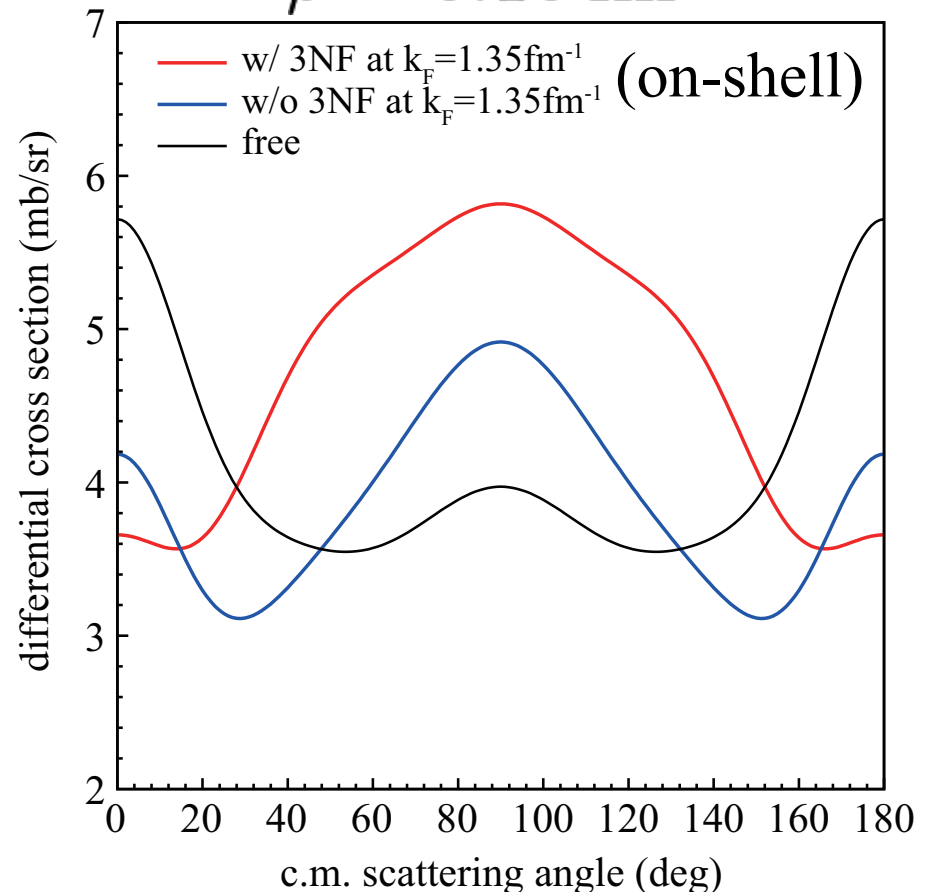
$g(\kappa', \kappa, \theta; E, \rho)$: **chiral g-matrix**

$$\rho = 0.16 \text{ fm}^{-3}$$

In medium pp scattering @ 200 MeV

$$\frac{d\sigma_{pp}}{d\Omega} \propto \left| g(\kappa', \kappa, \theta; E, \rho) \right|^2$$

3NF effects increase the pp cross section depending on the density.



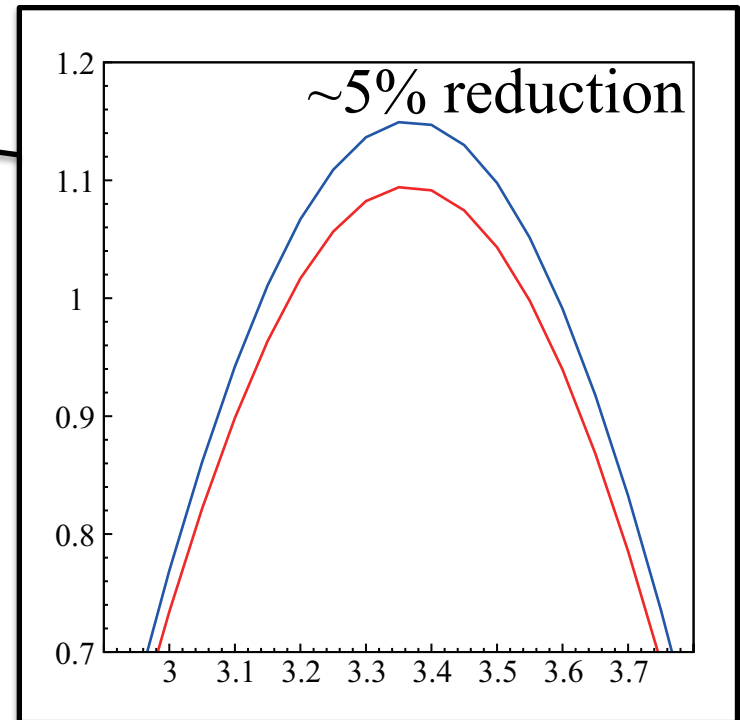
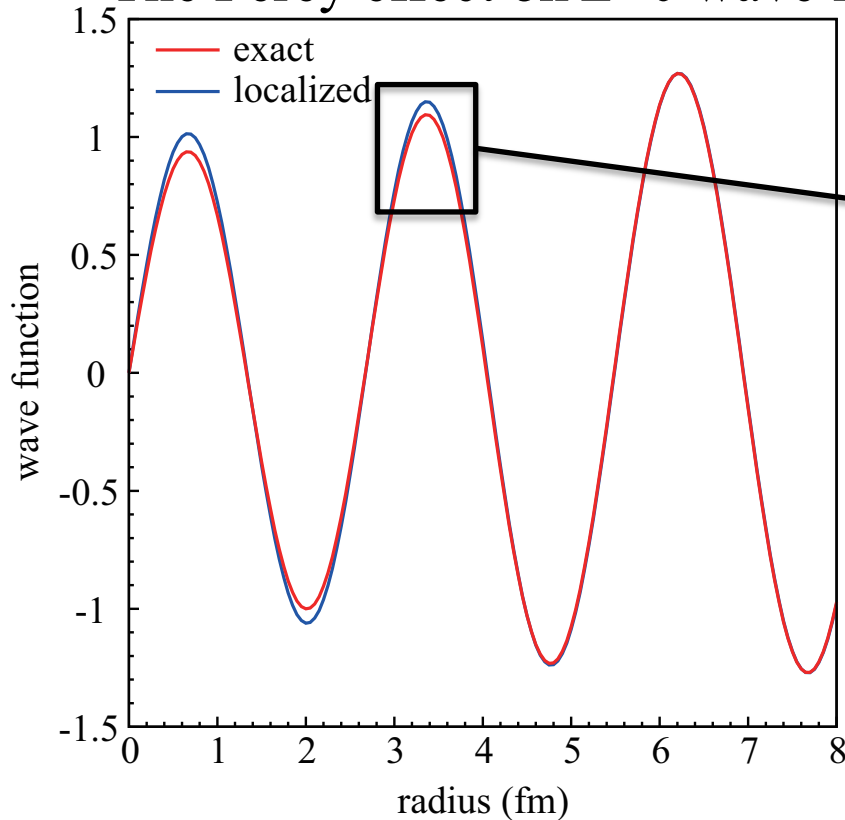
Microscopic DWIA

- ✓ Transition matrix element in the distorted wave Impulse Approx.

$$T = \langle \chi_{1,\mathbf{k}_1}^{(-)} \chi_{2,\mathbf{k}_2}^{(-)} | g(\kappa', \kappa, \theta; E, \rho) | \chi_{0,\mathbf{k}_0}^{(+)} \varphi_{nlj} \rangle$$

$\chi_{0,\mathbf{k}_0}^{(+)}$, $\chi_{1,\mathbf{k}_1}^{(-)}$, and $\chi_{2,\mathbf{k}_2}^{(-)}$: distorted waves calculated with
nonlocal microscopic optical potentials

The Perey effect on $L=0$ wave function for p - ^{40}Ca @100MeV



Microscopic DWIA

- ✓ Transition matrix element in the distorted wave Impulse Approx.

$$T = \langle \chi_{1,\mathbf{k}_1}^{(-)} \chi_{2,\mathbf{k}_2}^{(-)} | g(\kappa', \kappa, \theta; E, \rho) | \chi_{0,\mathbf{k}_0}^{(+)} \varphi_{nlj} \rangle$$

$g(\kappa', \kappa, \theta; E, \rho)$: **chiral g-matrix**

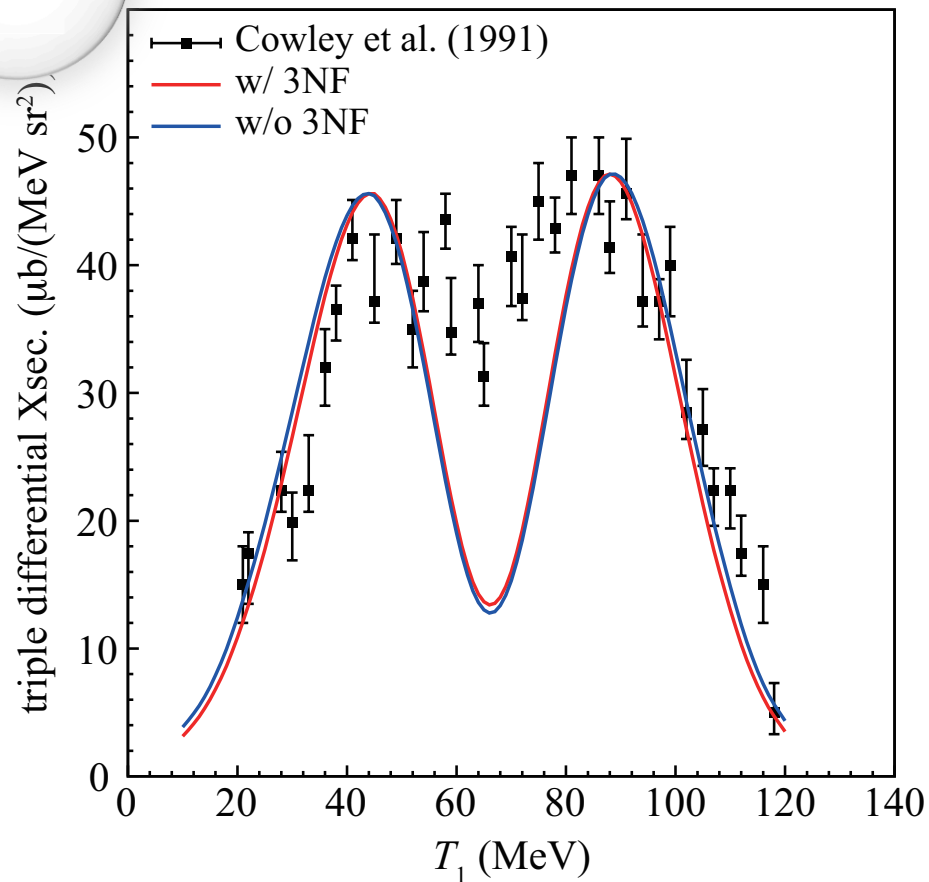
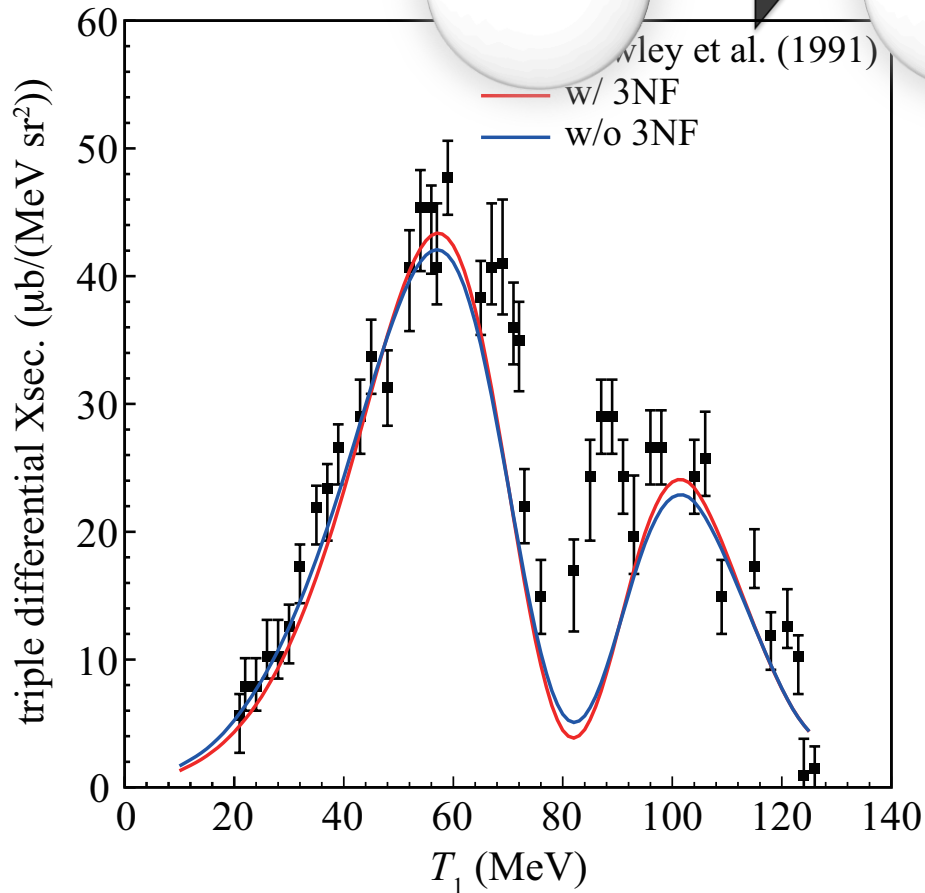
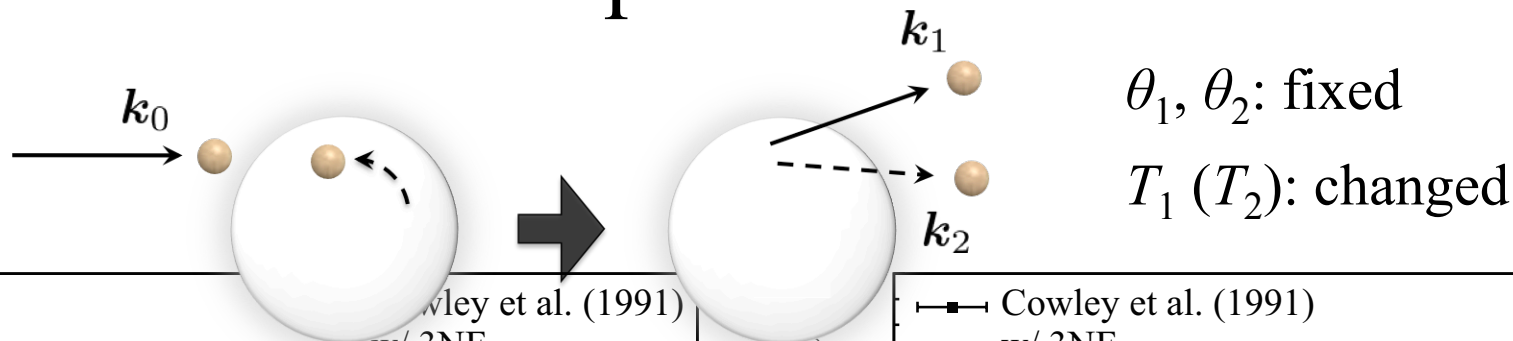
- ✓ medium effect
- ✓ off-shell properties

$\chi_{0,\mathbf{k}_0}^{(+)}$, $\chi_{1,\mathbf{k}_1}^{(-)}$, and $\chi_{2,\mathbf{k}_2}^{(-)}$: distorted waves calculated with **nonlocal microscopic optical potentials**

- ✓ Perey effect coming from knockon exchange process

φ_{nlj} : single particle wave function calculated by the Hartree-Fock method with the Gogny D1S force

Test of the present framework

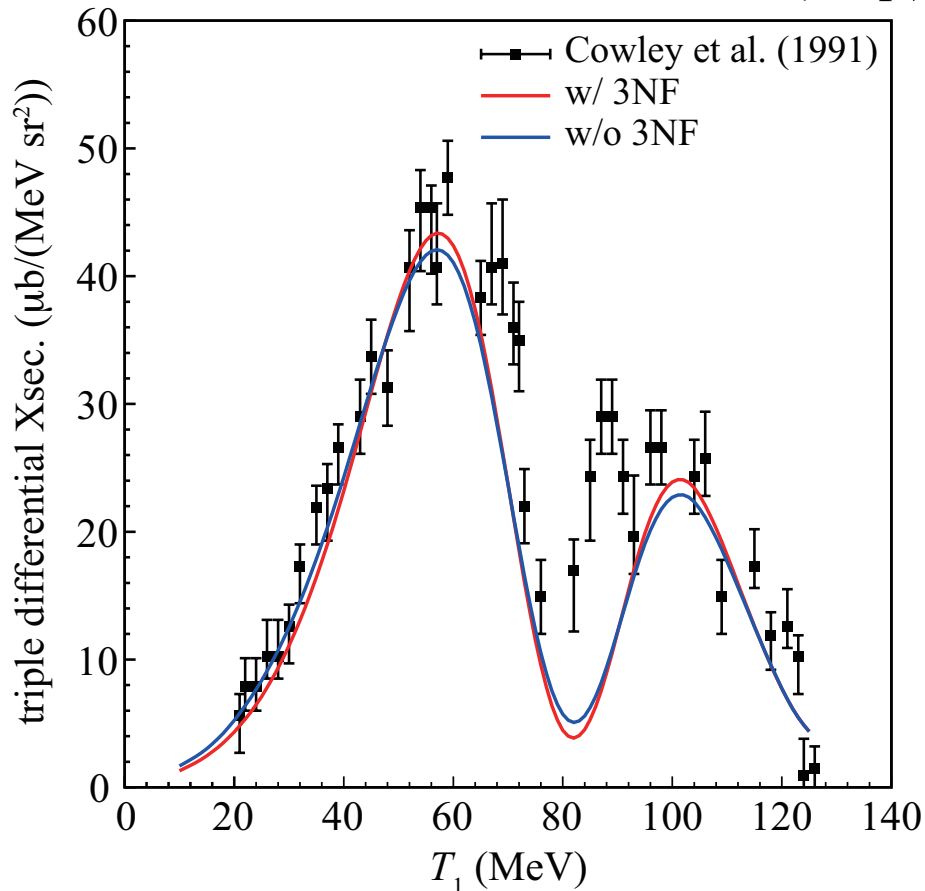


- ✓ The microscopic framework well reproduces the data.
- ✓ The 3NF effects are negligibly small.

Test of the present framework

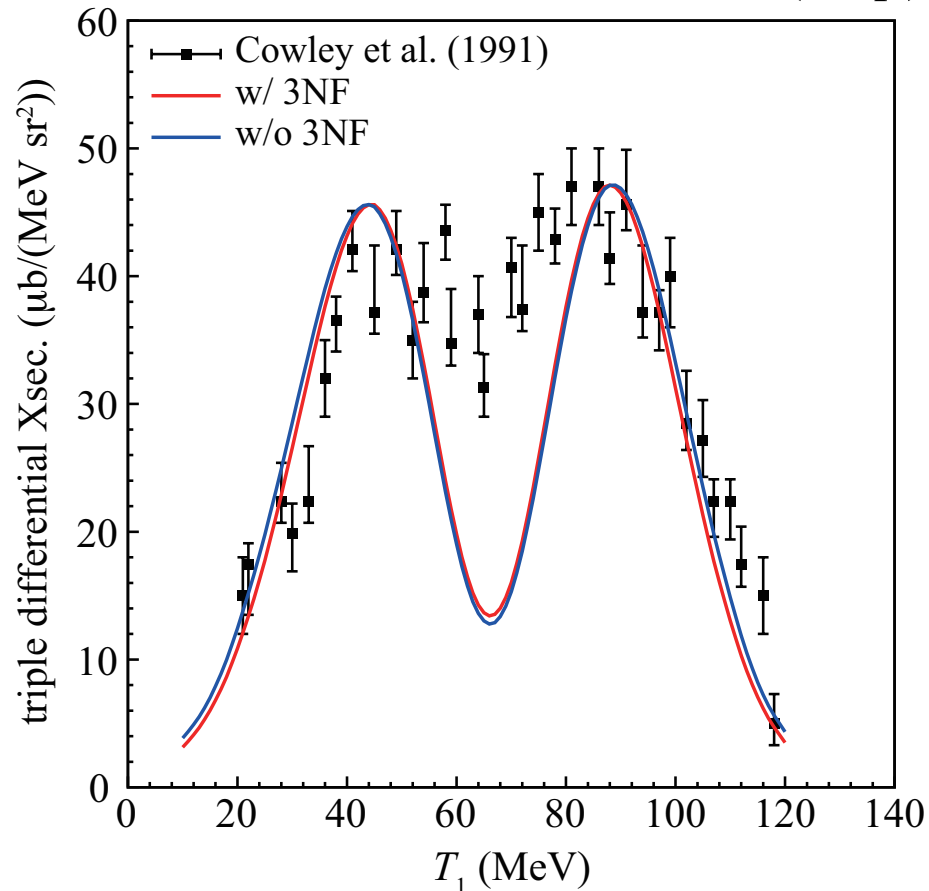
$^{16}\text{O}(p,2p)^{15}\text{N}_{\text{g.s.}}@151\text{MeV}$
(0p1/2 orbit)

$C^2S = 1.27 \pm 0.13$ extracted from $(e,e'p)$



$^{16}\text{O}(p,2p)^{15}\text{N}^*_{6.32\text{MeV}}@151\text{MeV}$
(0p3/2 orbit)

$C^2S = 2.25 \pm 0.22$ extracted from $(e,e'p)$

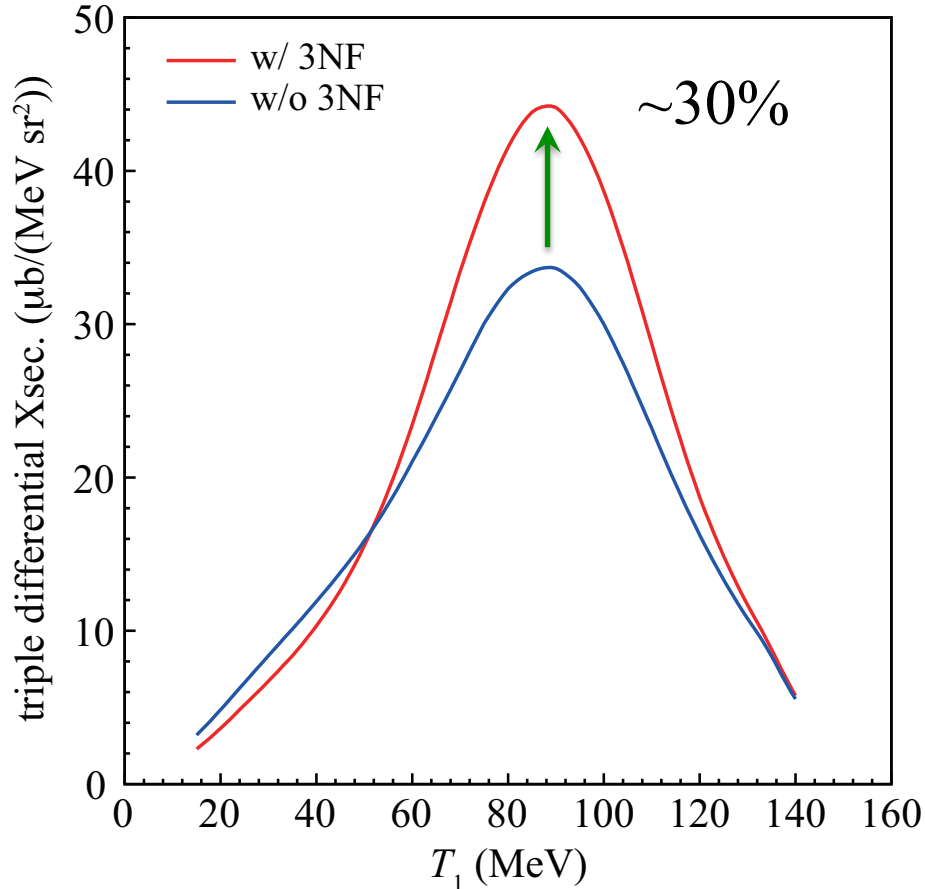


- ✓ The microscopic framework well reproduces the data.
- ✓ The 3NF effects are negligibly small.

Probing the 3NF effects

$^{40}\text{Ca}(p,2p)^{39}\text{K}^*$ @200MeV
(0s1/2 orbit)

$C^2S = 1$



✓ The 3NF effects enhance the cross sections near the peak.

✓ FWHM

w/ 3NF: 57 MeV

w/o 3NF: 67 MeV

cf.)

$C^2S = 1.56 \pm 0.28$ extracted from $(p,2p)$
Y. Yasuda et al., PRC81, 044315 (2010).

$C^2S = 1.50$ extracted from $(e,e'p)$
J. Mougey et al., NPA262, 461 (1976).

$C^2S = 1.10$

A. Fabrocini et al., PRC63, 044319 (2001).

$C^2S = 1.56$

C. Bisconti et al., PRC75, 054302 (2007).

$(p,2p)$ reactions can be used to probe 3NF effects!

Summary and perspective

Summary

- ✓ Microscopic approach to many-body direct reactions based on chiral interactions
- ✓ Microscopic DWIA framework
- ✓ A possibility of probing the 3NF effects by using $(p,2p)$ reaction

Perspective

- ✓ Spin observables $\frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$
- ✓ 3NF effects for $T=3/2$ state

Thank you very much for your attention!