



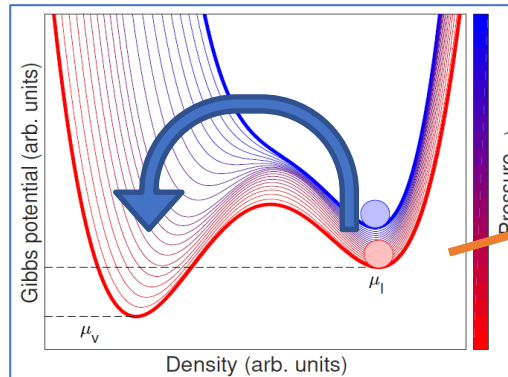
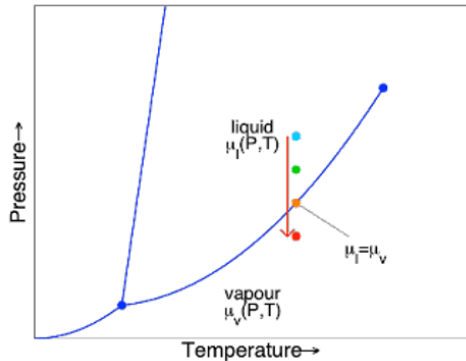
# Toward understanding the nuclear efficiency threshold of bubble chamber detectors

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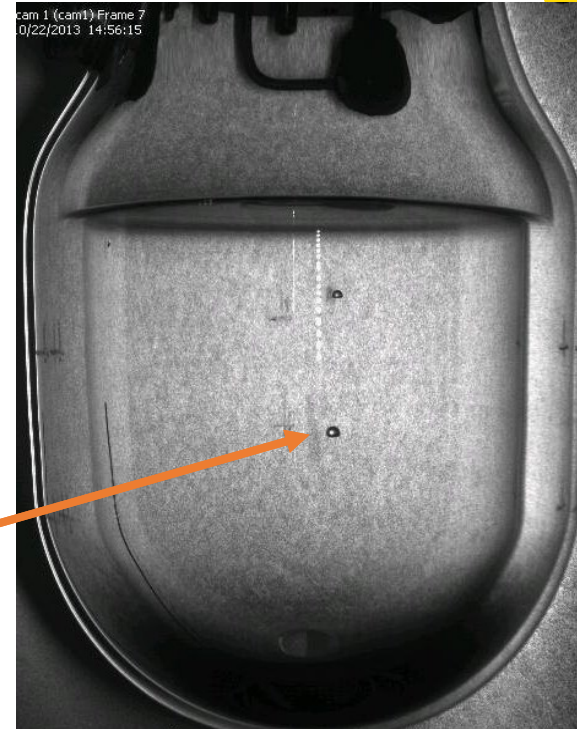


## How does it work?

- Particles interact with target nuclei causes nuclear recoils (NR).
- A small deposition of energy causes a phase change in superheated fluid.
- Nucleation can be observed optically, acoustically, or barometrically.



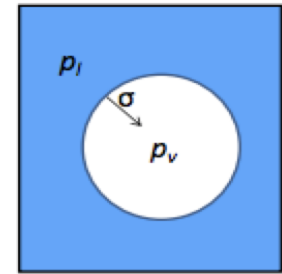
PICO 2L [1]



## Current theory: the heat spike model

The latent heat of evaporation.

The Seitz model [2]:  $E_{th} = \frac{4\pi}{3} R_c^3 \Delta p + \frac{4\pi}{3} R_c^3 \rho_v h_{lv} + \frac{4\pi}{3} R_c^3 \left( \sigma - T \frac{\partial \sigma}{\partial T} \right) + W_{irr}$



The work needed to form the bubble surface.

Combination of the free energy of surface and heat absorption from surrounding.

Irreversible work is small that we usually neglect.

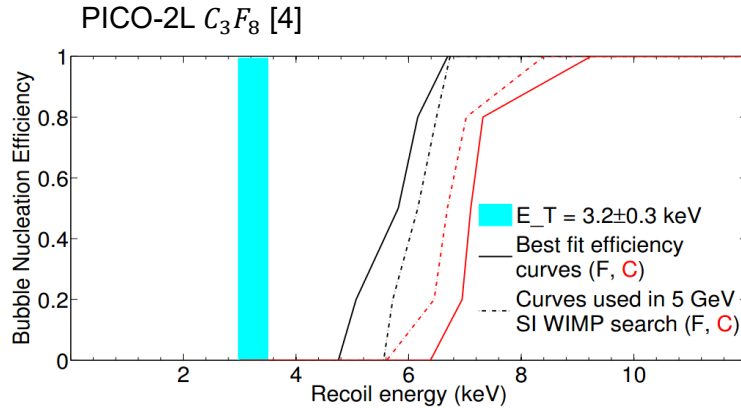
- $\Delta p$ : Pressure between bubble and liquid.
- $\rho_v$ : Density of the bubble.
- $\sigma$ : Surface tension.
- $\Delta H$ : Enthalpy change.
- $T$ : Temperature.
- Where  $R_c = \frac{2\sigma}{\Delta p}$ , is called critical radius.

### Condition for phase transition:

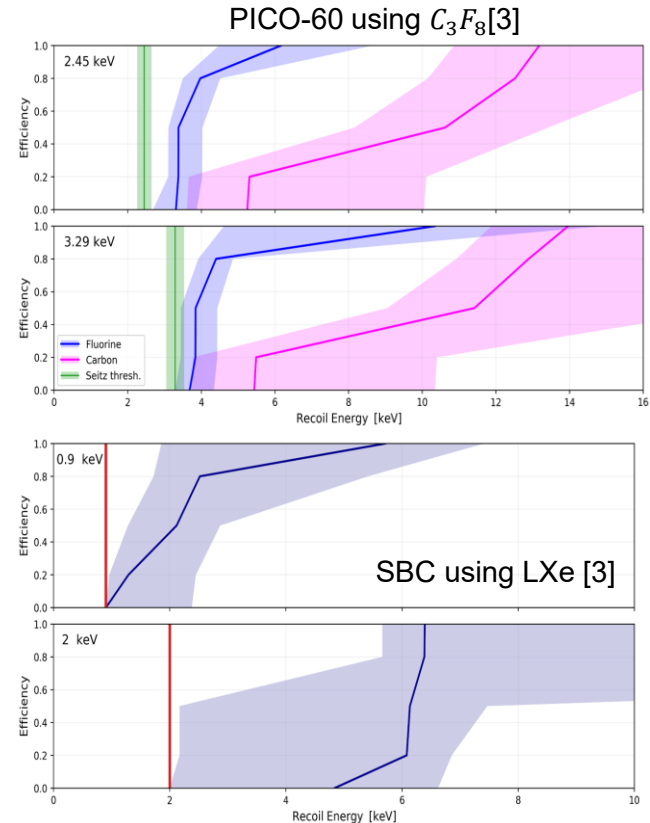
$$E_{dep} = \int_0^{l_c} \frac{dE}{dx} dx \geq E_{th}$$

- $l_c$ : The critical length or track length.
- Where  $l_c = bR_c$
- $b$  is a value obtained experimentally.

## Experimental discrepancy:



- It's well known that the threshold for nucleation deviates from the predicted Seitz threshold. (PICO-2L, PICO-60, SBC, etc.)
- Recent new fitting results [3]:  
→ See Daniel Durnford's talk



## Goals:

- Understand why the response is not a step function.
- Explore why the response is delayed and what factors cause the efficiency curve pattern.
- Establish a theory/model that one can extrapolate to other recoil threshold energies (in the range from 1 – 200 keV).

## Study methods:

- Molecular dynamic (MD) to study the bubble formation and growth in microscopic scale.
- SRIM Monte Carlo simulation to study the energy transfer (depth penetration, losses, etc.).

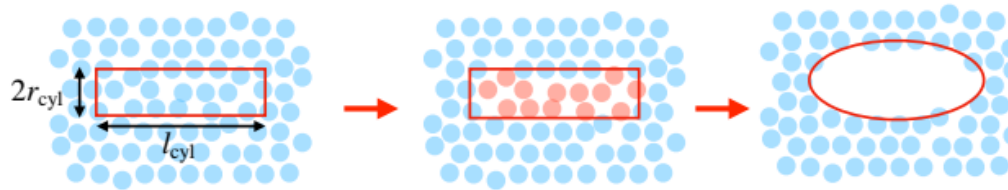
## Molecular dynamic simulation in LAMMPS [5]:

- In LAMMPS, the molecular or atoms are interacting with each other by the Lennard-Jones potential.

$$u_{LJ}(r) = 4\epsilon \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right]$$

- $\sigma$ : size of the atom/molecular.
- $\epsilon$ : minimum potential energy.
- $r$ : distance.

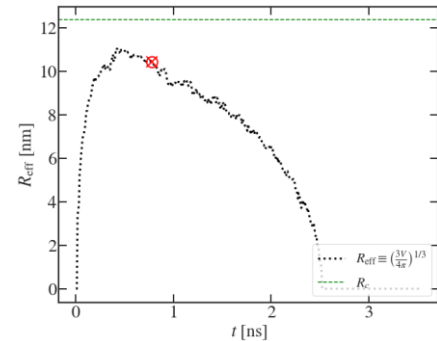
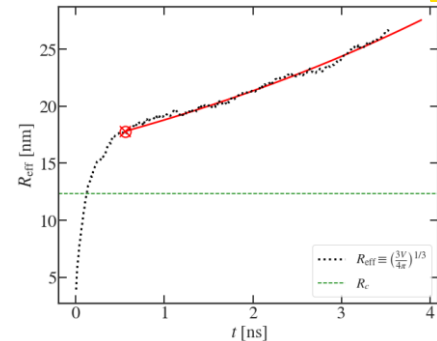
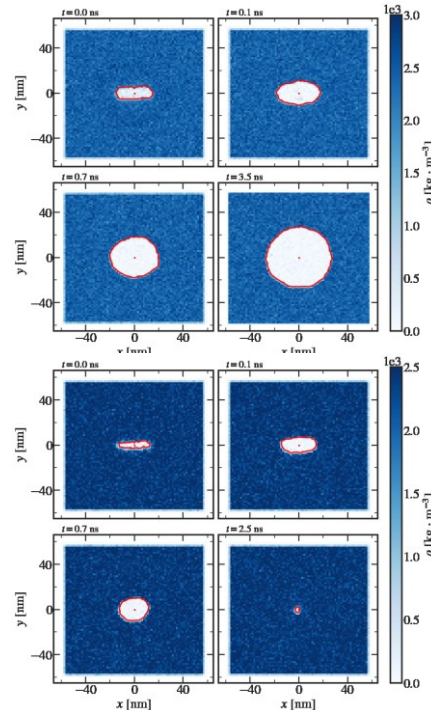
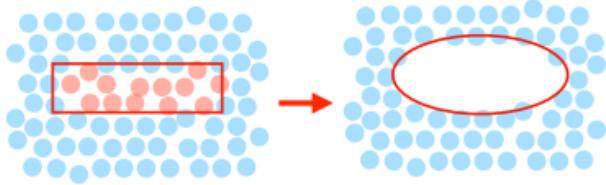
- An incoming particle deposits energy that is simulated by rescaling energy inside a cylinder.
- The resulting bubble will either grow or collapse depending on the Seitz energy threshold or energy density  $dE/dx$ .



$$l_{\text{cyl}} = l_c = bR_c$$

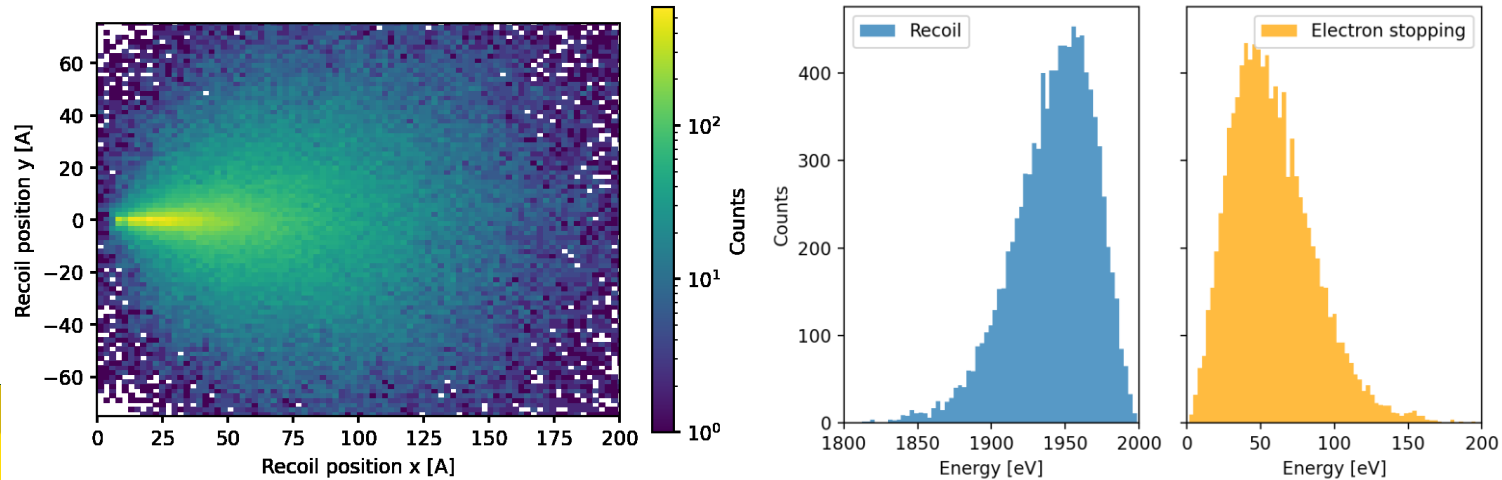
## Molecular dynamic simulation in LAMMPS [5]:

- $E_{th} = 0.9$  keV,  $T = -43$  °C,  
 $p = 25$  psia.
- 2 keV and 1 keV energy deposition in liquid xenon as an example
- The red fitting curve represents the bubble growth expansion described by Rayleigh-Plesset-Zwicky equation



## Monte Carlo simulation in SRIM [6]:

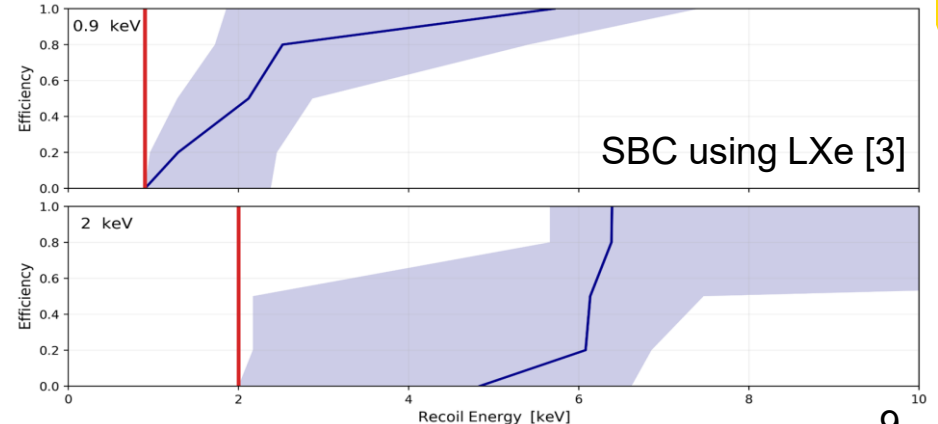
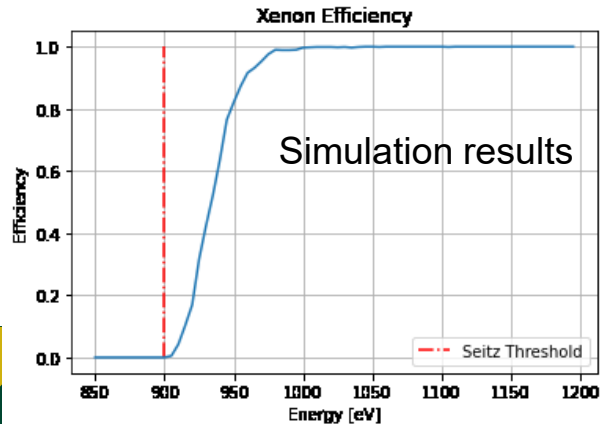
- Able to track how much energy of nuclear recoil is transferred into energy in heat spike.
- The graphs below show the 2 keV nuclear recoil events in liquid xenon.





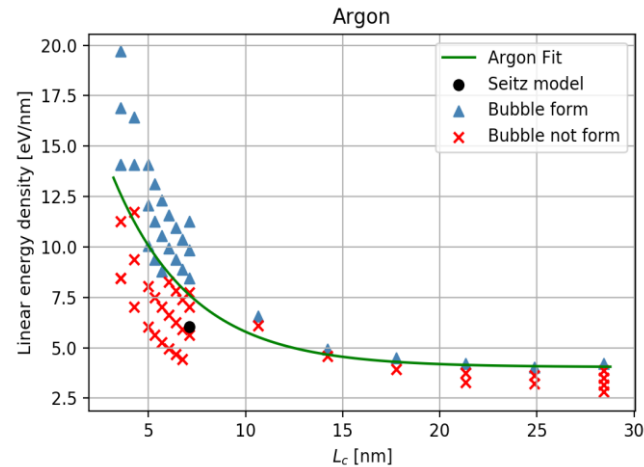
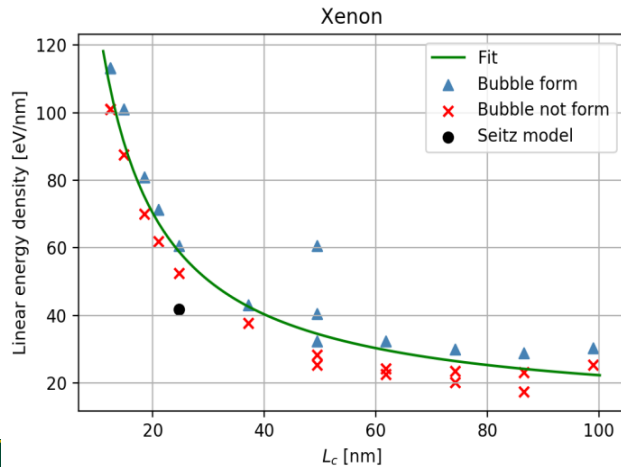
## Method 1

- We are using SRIM to study the bubble nucleation efficiency.
- Nuclear recoil energy contributes to the heat spike.
- Electron stopping does not contribute to heat spike.
- The assumption is that if the energy is larger than the Seitz threshold, the bubble can form.



## Bubble formation with respect to the track length in MD

- The fit function represents the separation between the bubble formed and not formed.
- The  $l_c$  is the length of energy deposition defined manually in MD simulation.

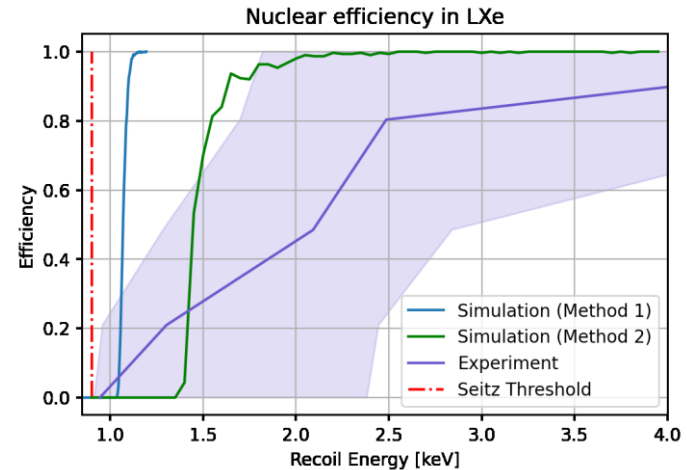
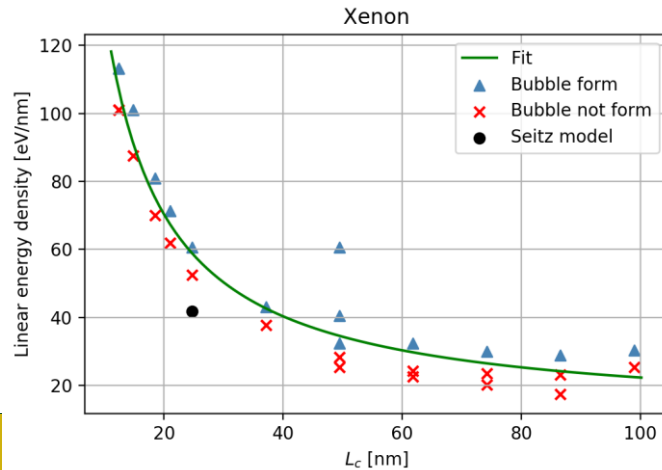


$$\text{Fitting function: } f(x) = \frac{a}{x^2} + \frac{b}{x} + c$$

## Method 2

- Assuming the linear energy density larger than the fitting curve the bubble will form.
- Results compared with nucleation efficiency for SBC using LXe [3].

→ See Daniel Durnford's slides.



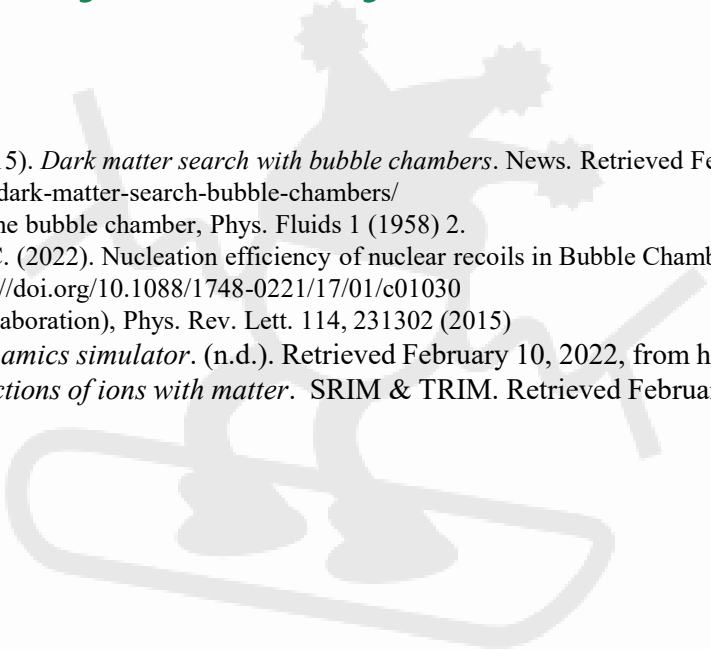
## Summary:

- We perform molecular dynamics and Monte Carlo simulation to study the nucleation efficiency discrepancy between the Seitz's model and experimental result.
- The preliminary are promising and seems to better reproduce the experimental data.

## Next step:

- Use other gas species (Ar,  $C_3F_8$ , etc.)
- Explore how to simulate multi-compounds atoms in molecular dynamics.
- Generalize the result to construct a real physics model to explain the discrepancy observed between the experimental results and the current Seitz model.

# Thank you for your attention!

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- [1] Amole, C. (2020, January 15). *Dark matter search with bubble chambers*. News. Retrieved February 10, 2022, from <https://news.fnal.gov/2016/07/dark-matter-search-bubble-chambers/>
- [2] F. Seitz, On the theory of the bubble chamber, *Phys. Fluids* 1 (1958) 2.
- [3] Durnford, D., & Piro, M.-C. (2022). Nucleation efficiency of nuclear recoils in Bubble Chambers. *Journal of Instrumentation*, 17(01). <https://doi.org/10.1088/1748-0221/17/01/c01030>
- [4] Amole, C et al. (PICO Collaboration), *Phys. Rev. Lett.* 114, 231302 (2015)
- [5] *LAMMPS molecular dynamics simulator*. (n.d.). Retrieved February 10, 2022, from <https://www.lammps.org/>
- [6] Ziegler, J. (n.d.). *Interactions of ions with matter*. SRIM & TRIM. Retrieved February 10, 2022, from <http://www.srim.org/>

Bubble expansion can be described by the equation ¶ :

$$\frac{dR}{dt} = - \left[ \frac{A^2 \sqrt{t - t_s}}{B} + \frac{2\nu_1}{R} \right] + \sqrt{A^2 - \frac{2\gamma}{\rho_l R} + \left( \frac{2\nu_1}{R} + \frac{A^2 \sqrt{t - t_s}}{B} \right)^2}$$

$A$  — speed of expansion in the linear growth phase [m/s]

$B$  — characteristic rate of expansion in the thermal growth phase [m/s<sup>2</sup>]

$\nu_1$  — Kinematic viscosity [m<sup>2</sup>/s]

$\gamma$  — Surface tension [N/m]

$\rho_l$  — Fluid density [kg/m<sup>3</sup>]

The initial condition is provided by the existing (MD) simulations.

¶ T. Kozynets, S. Fallows and C. Krauss, "Modeling emission of acoustic energy during bubble expansion in PICO bubble chambers", *Physical Review D*, vol. 100, no. 5, 2019. Available: [10.1103/physrevd.100.052001](https://doi.org/10.1103/physrevd.100.052001).

The linear energy density as function of track length is fit with:

$$f(x) = \frac{a}{x^2} + \frac{b}{x} + c$$

Where  $a = 16.88$  eV/nm,  $b = 79.41$  eV/nm,  $c = 13.40$  eV/nm.

$x$  is defined as  $l_c/R_c$ , where  $R_c = 12.36$  nm.