



Experiments with Rare Isotopes
- Recent (selected) Highlights -

R. Kanungo

Saint Mary's University & TRIUMF , Canada

Experiments with Rare Isotopes - Recent (selected) Highlights -

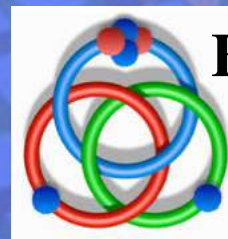
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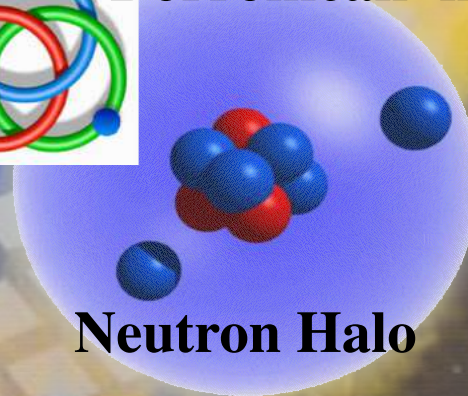
Proton Number



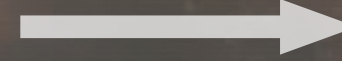
Stable Nucleus



Borromean nucleus



Neutron Halo



Neutron-rich matter

Rare Isotope
weak binding, $N/Z \gg 1$

Neutron Number



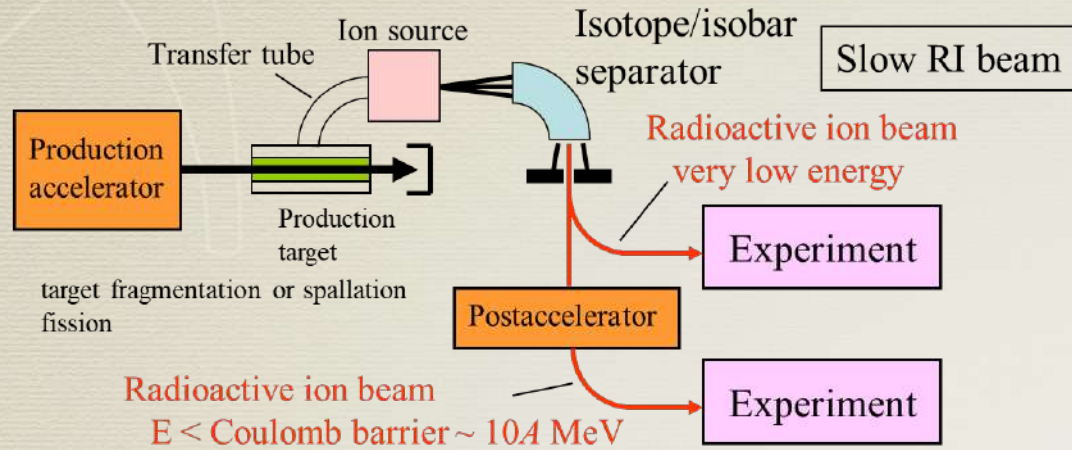
Why explore the rare isotopes ?

- **What new features emerge with neutron-proton asymmetry ?**
New structures - Halo, Skin
New Excitation modes
Change of shells
- **What is the nature of the nuclear force ?**
Tensor force
Three-body force
Pairing Interaction
- **How do rare isotopes shape our Universe ?**
Nucleosynthesis (*Talk : Hendrik Schatz*)
Structure information needed to constrain reaction rates
Equation of state of asymmetric nuclear matter
- **Rare isotopes test fundamental symmetries**

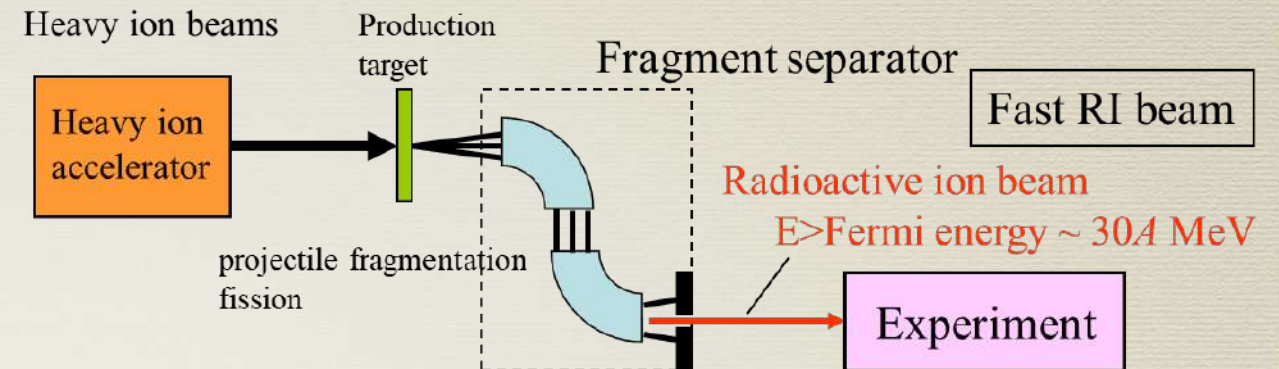
Rare Isotopes Facilities

Isotope Separator Online (ISOL)

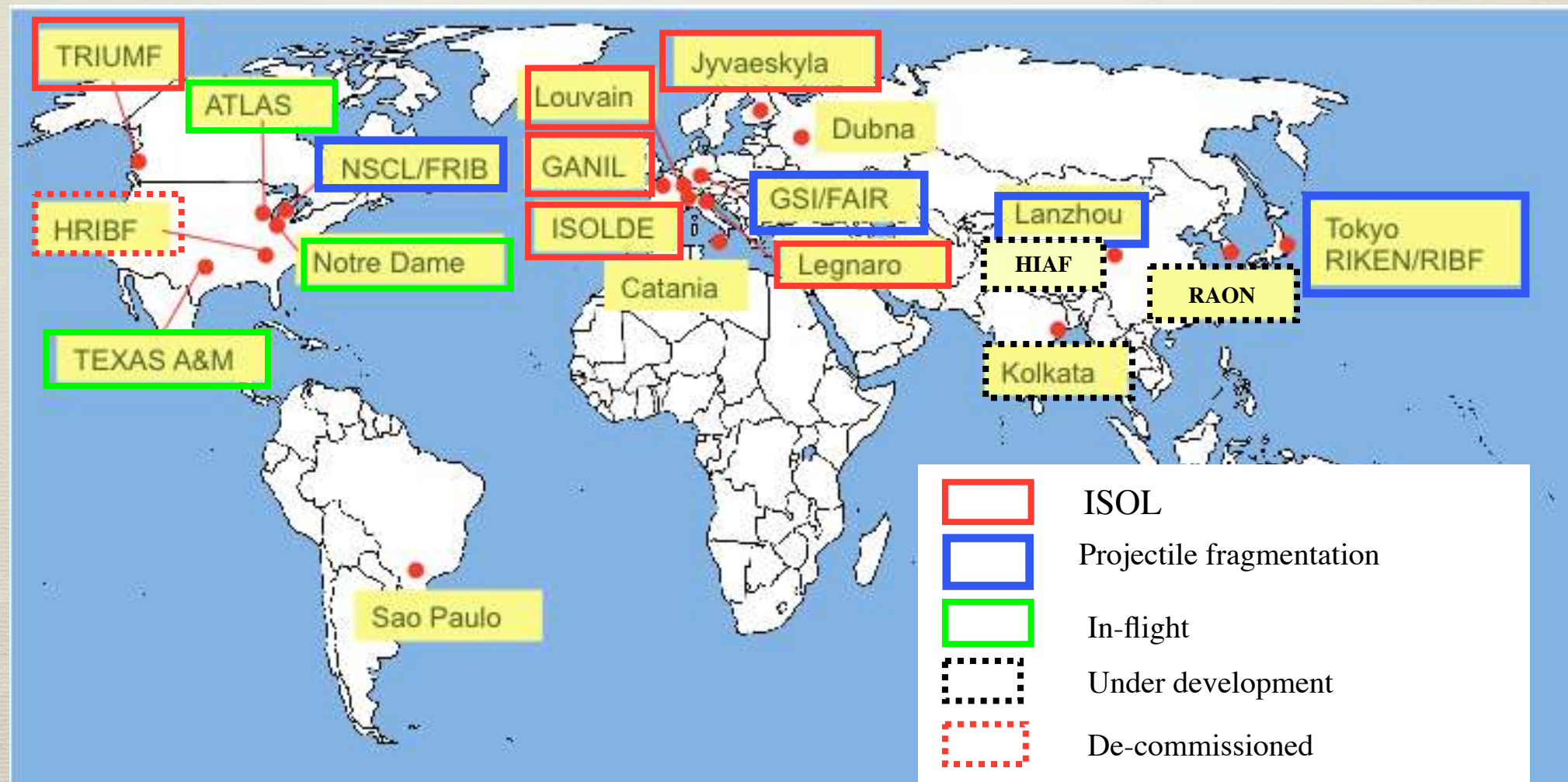
ISOL facility



In-flight - Projectile Fragmentation



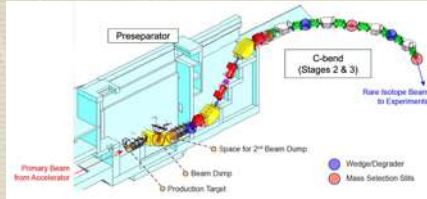
Courtesy : H. Sakurai



State-of-the-art instruments peek into rare isotopes

Fragment Separators

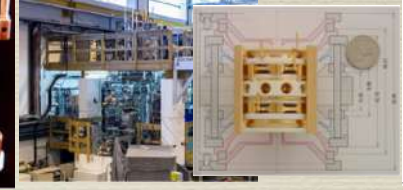
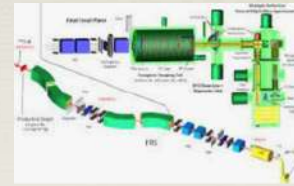
BigRIPS (RIKEN), FRS (GSI), ARIS+A1900 (FRIB)



Mass Measurements

Penning Traps: TITAN (TRIUMF), ISOLTRAP (CERN), CPT (ANL), LEBIT (FRIB), JYFLTRAP (Jyväskylä)

MR-TOF: RIKEN, GSI, TRIUMF



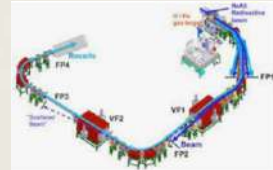
Storage Ring: ESR (GSI), Rare-RI Ring (RIKEN), HIRFL-CSR (Lanzhou)



High Rigidity Spectrometers & Mass Separators

High Rigidity Spectrometers: SAMURAI (RIKEN), GLAD-R³B (GSI/FAIR), S800 (FRIB)

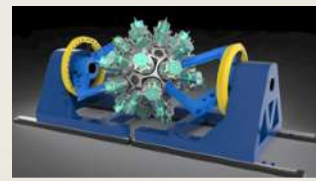
Mass Separators/Spectrometers: DRAGON, EMMA (TRIUMF), SECAR (FRIB), SHARAQ (RIKEN)



γ - spectroscopy

High Resolution: GRETINA/GRETA, CLARION (USA), AGATA (Europe), GRIFFIN, TIGRESS (Canada)

High Efficiency: DALI2 (RIKEN)

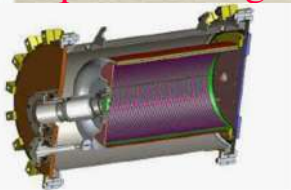


High Luminosity Targets

Active Targets: MAYA, ACTAR (GANIL), AT-TPC (FRIB), TEXAT (TAMU), SpecMAT (Leuven), MAIKo (Kyoto/RCNP)

Solid H₂ Target: IRIS (TRIUMF)

Liquid H₂ Target: MINOS (RIKEN)



High resolution charged particle spectroscopy

HELIOS (ANL), Isolde Solenoidal Spectrometer (CERN), SOLARIS (FRIB)



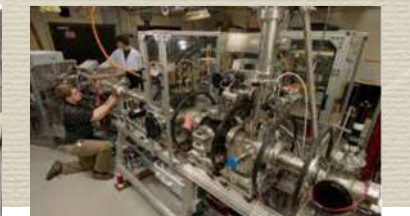
Neutron detectors

MoNA (FRIB), NeuLAND (GSI), NEBULA (RIKEN), BELEN (Europe), VANDLE (FRIB), DESCANT (TRIUMF), TexNEUT (TAMU)

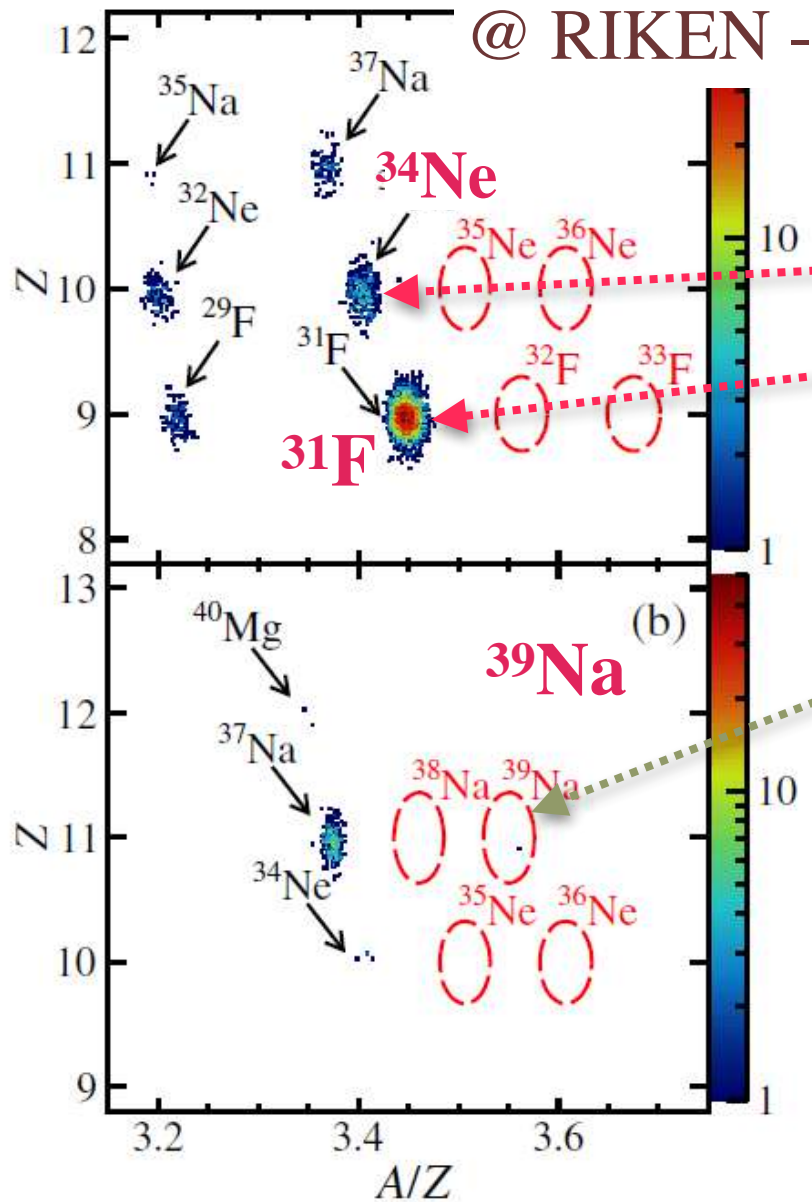


Laser Spectroscopy

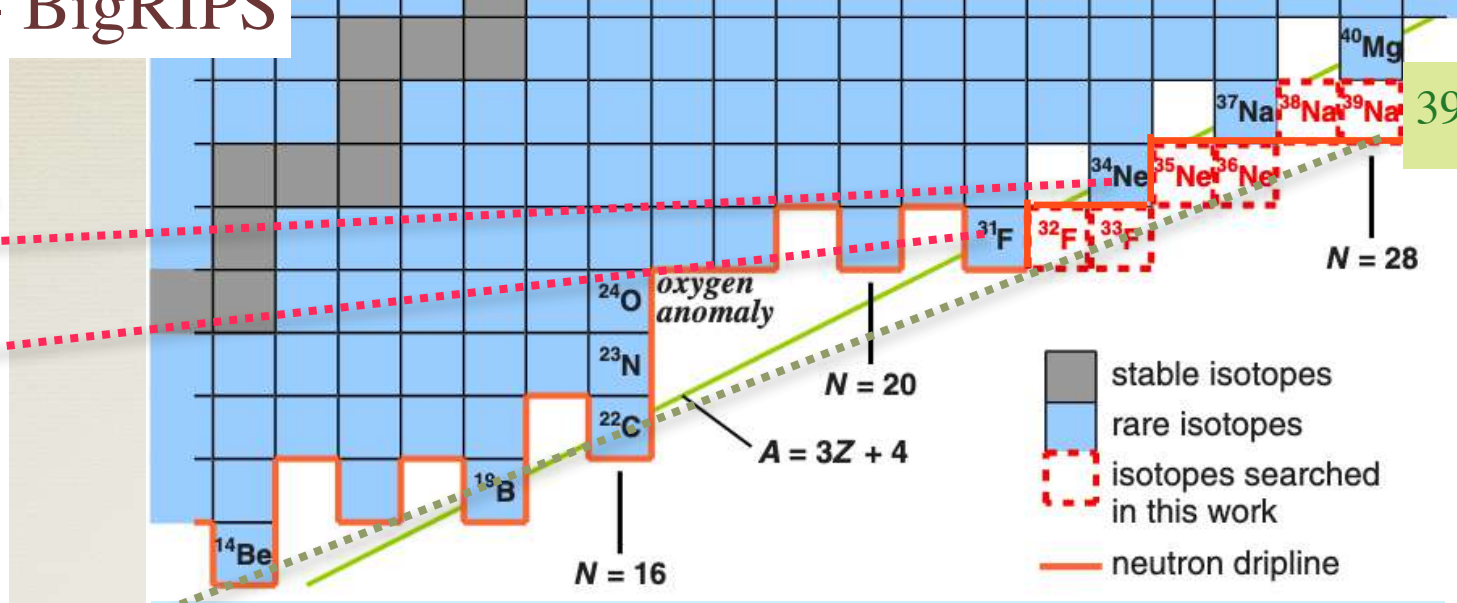
COLLAPS, CRIS (CERN), BECOLA (FRIB)



Hunt for the nuclear landscape boundary



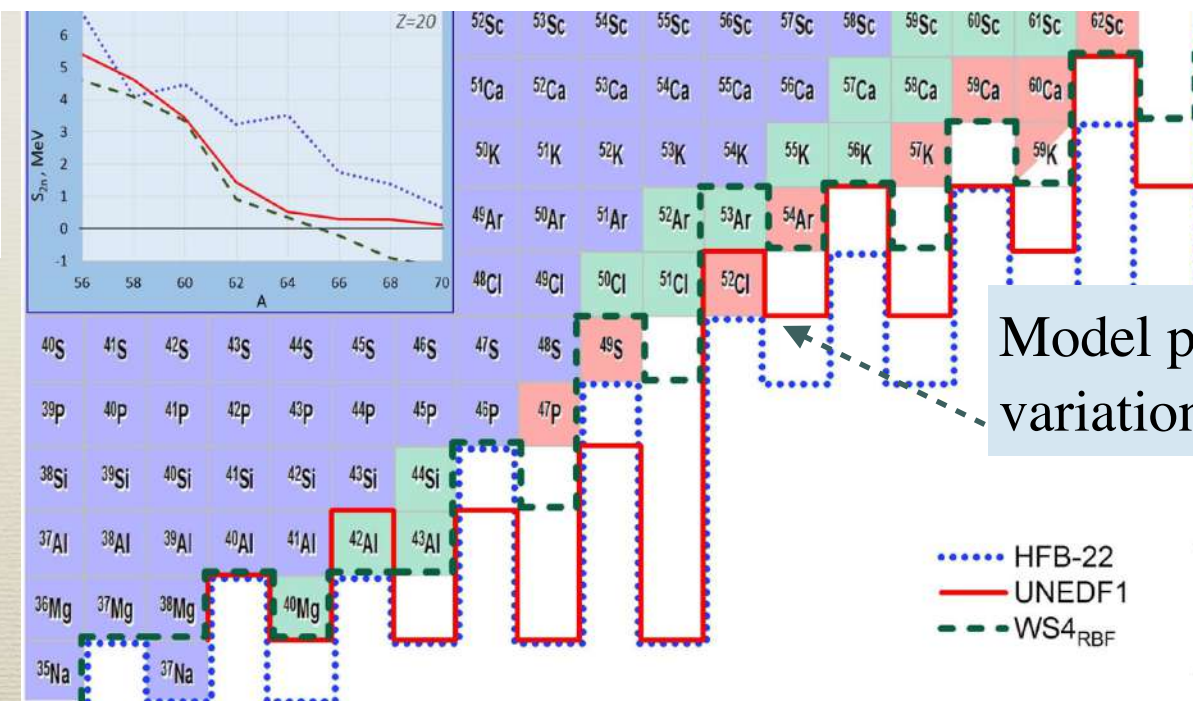
D. Ahn *et al.*, Phys. Rev. Lett. 123 (2019) 212501



^{39}Na is bound !

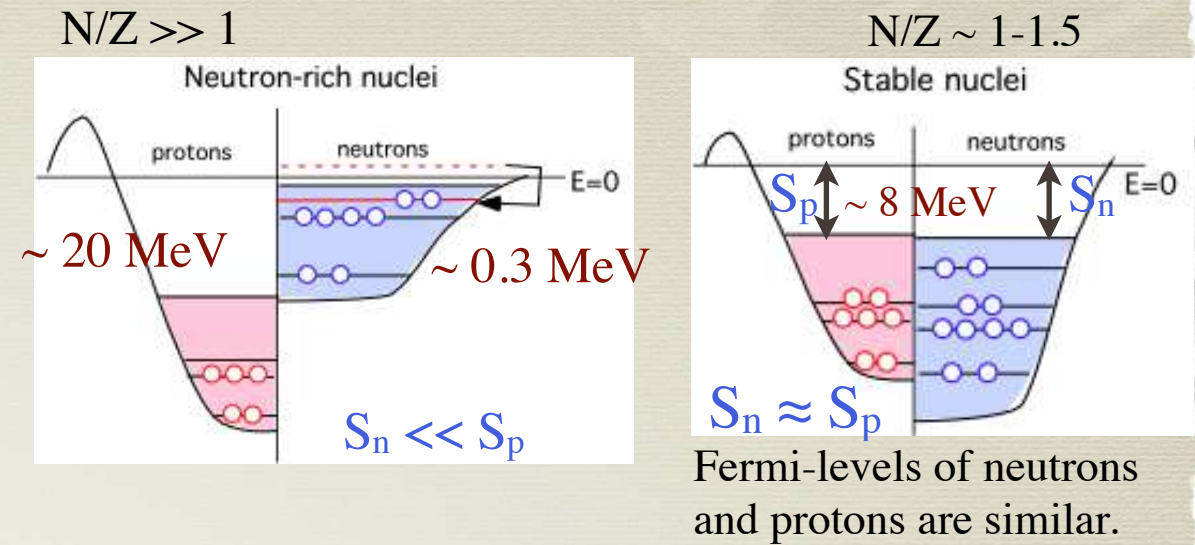
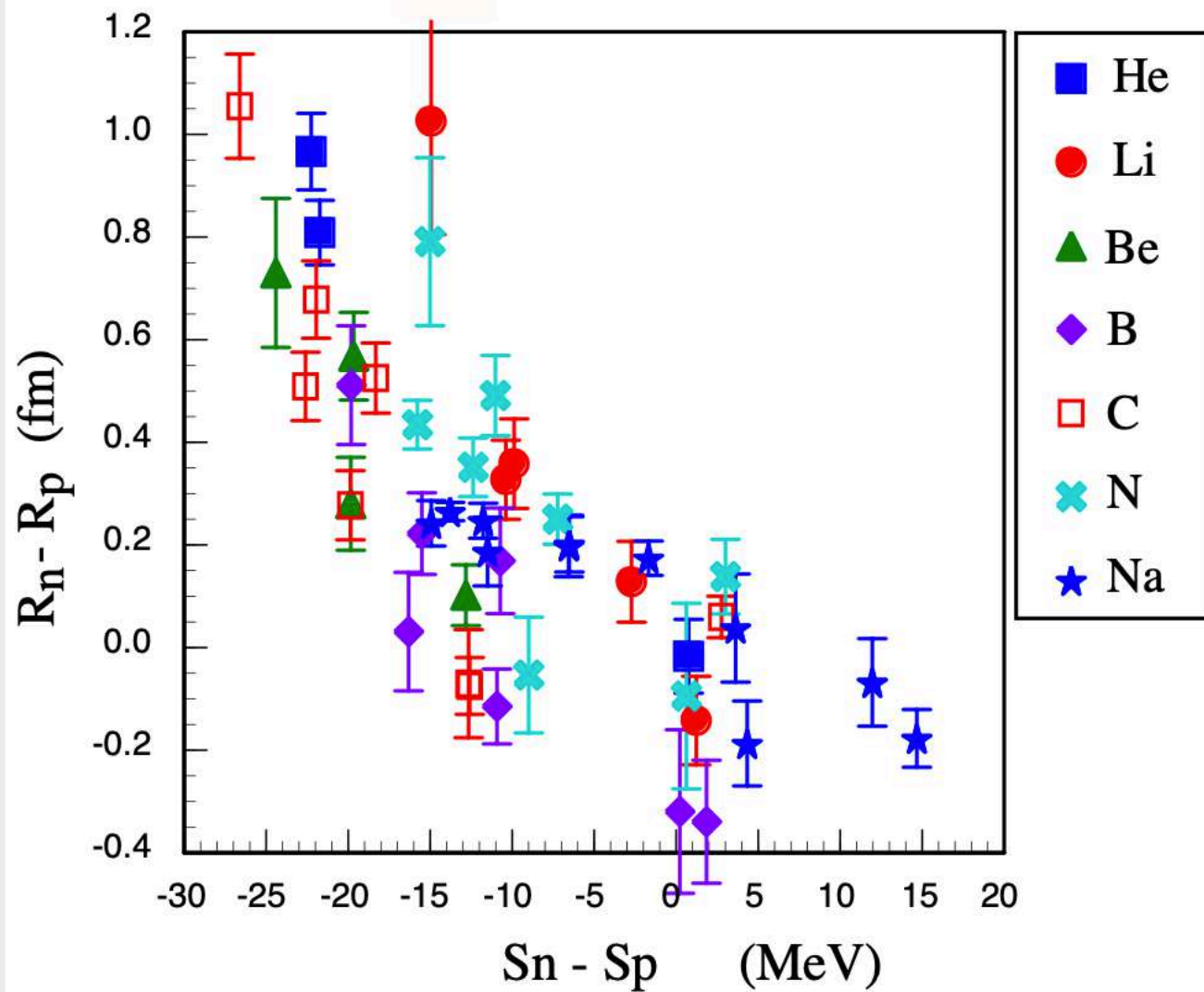
Neutron drip-line established for F and Ne
 Drip line @ ^{31}F signals it is deformed

O. B. Tarasov *et al.*, Phys. Rev. Lett. 121 (2018) 022501



Model predictions have a wide variation in predicting the drip-line

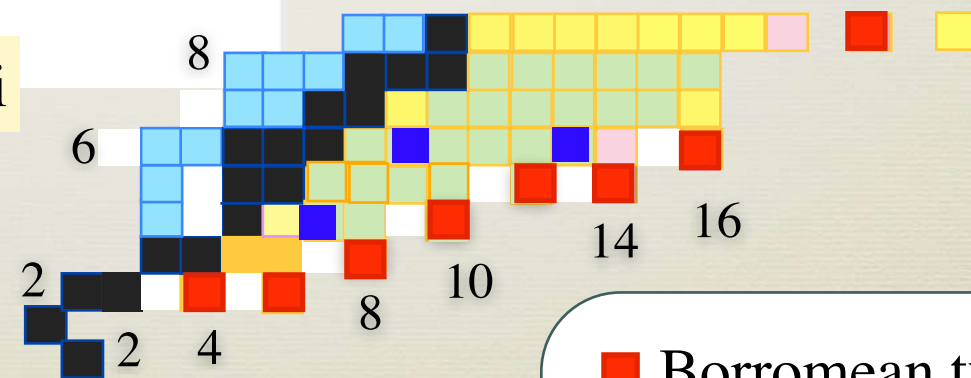
Neutron balloon in neutron-rich nuclei



❖ Weak binding & $n-p$ Fermi level difference

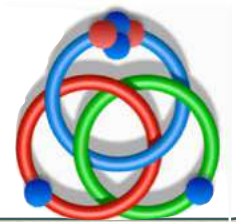
Neutron-rich nuclei

Stable nuclei



■ Borromean two-neutron halos
Core+n+n

■ One-neutron halos

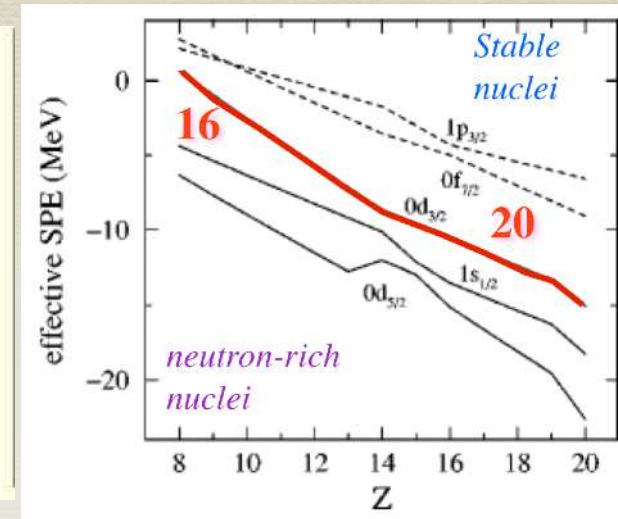
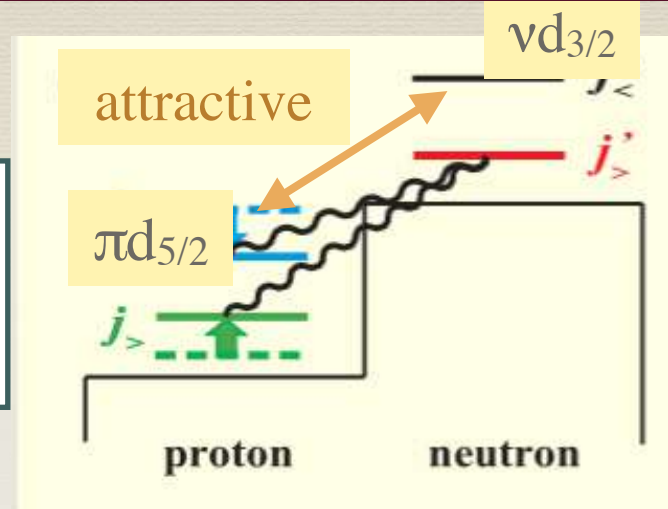


Exotic phenomena & nuclear force

❖ Tensor force

$p(j_< = l-1/2) - n(j_> = l+1/2)$ attractive
 $p(j_> = l+1/2) - n(j_> = l+1/2)$ repulsive

T. Otsuka *et al.*, Phys. Rev. Lett. 2005

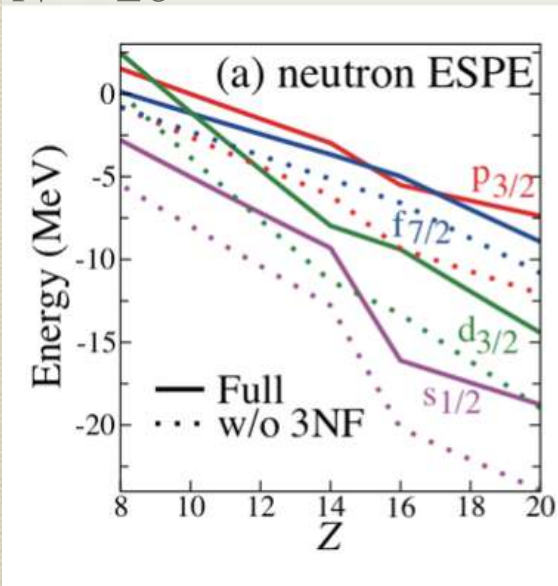


❖ Three-nucleon force

$N = 20$

N. Tsunoda *et al.*,
 Phys. Rev. C (R) 2017

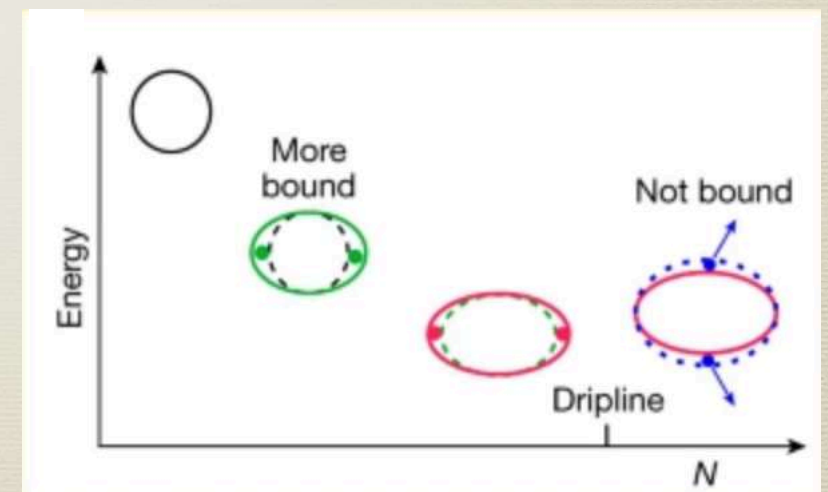
Neutron drip-line in O isotopes



❖ Deformation

N. Tsunoda, T. Otsuka, K. Takayanagi *et al.*, Nature, 2020

F to Mg : Strongly correlated valence neutrons \rightarrow ellipsoidal shape saturation marks the drip-line



❖ Pairing Interaction

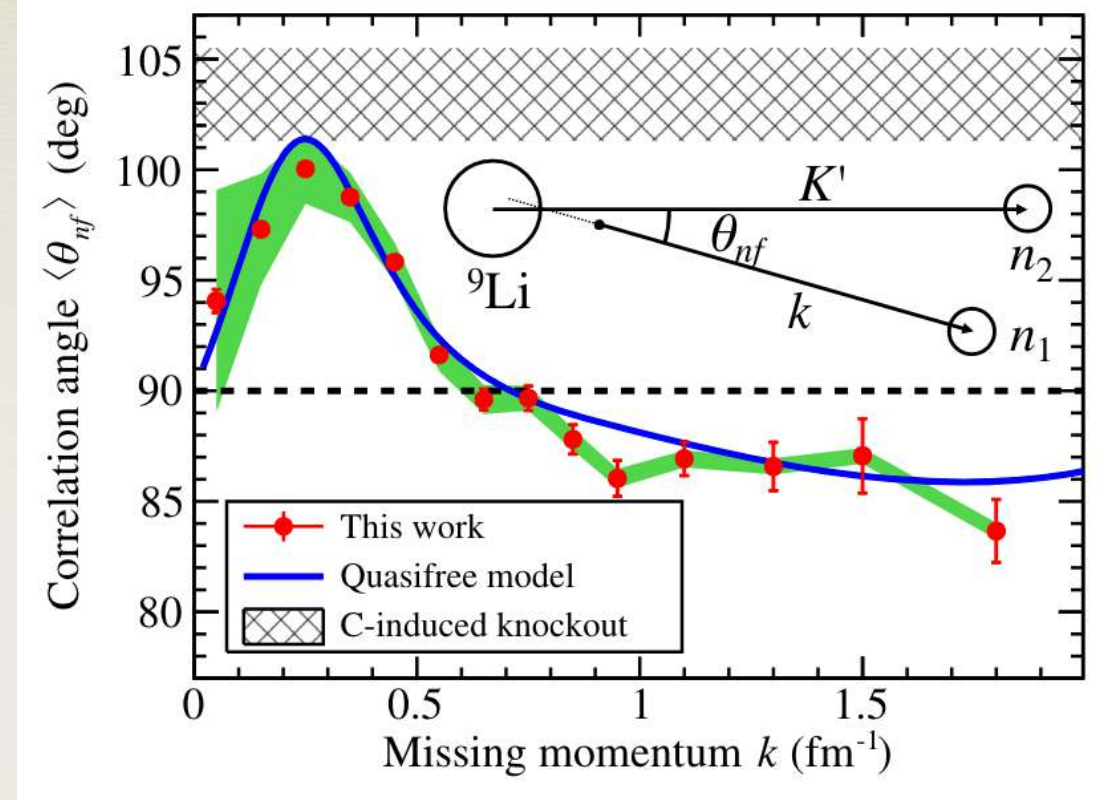
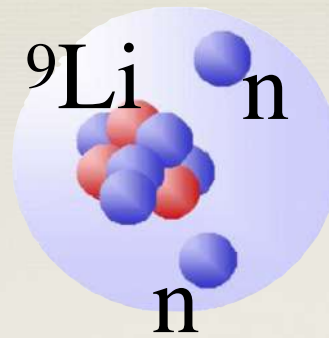
Neutron Halo features & Disappearance of known nuclear shell gaps

^{11}Li : Halo n-n correlation

@ RIKEN-SAMURAI $E/A = 246 \text{ MeV}$

$^{11}\text{Li}(p, pn)^{10}\text{Li}$

Y. Kubota, A. Corsi, G. Authelet *et al.*, Phys. Rev. Lett. 125 (2020) 252501



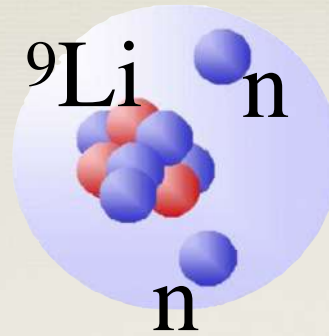
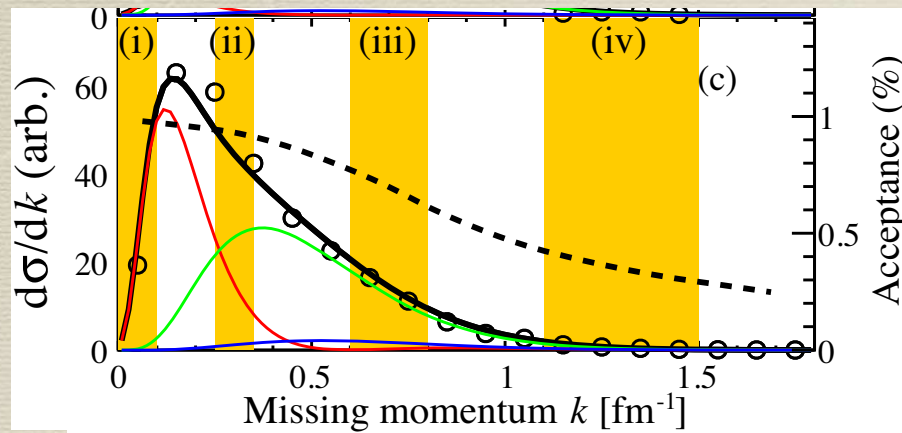
Di-neutron localized at the surface
i.e. Core - nn distance of $\sim 3.6 \text{ fm}$

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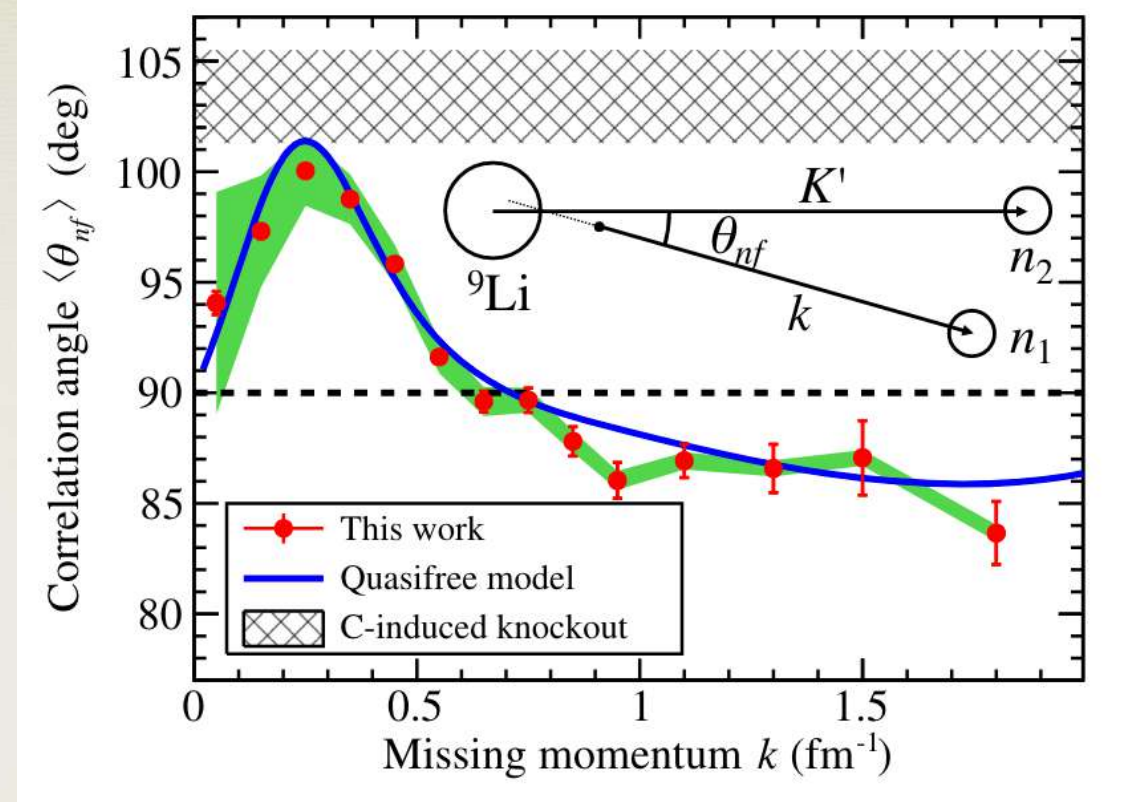
$2s_{1/2} : 35(4)\%$

$1p_{1/2} : 59(1)\%$

$1d_{5/2} : 6(4)\%$



$N = 8$ shell vanishes



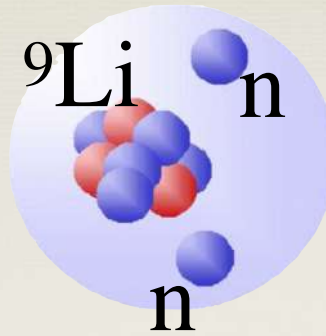
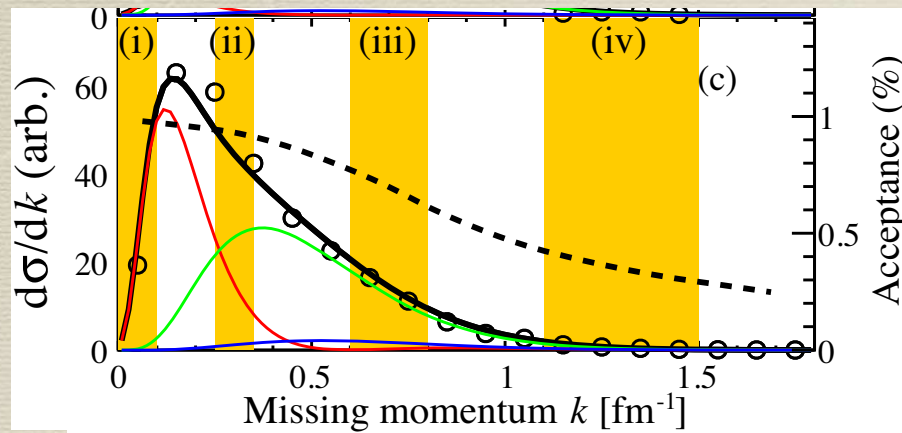
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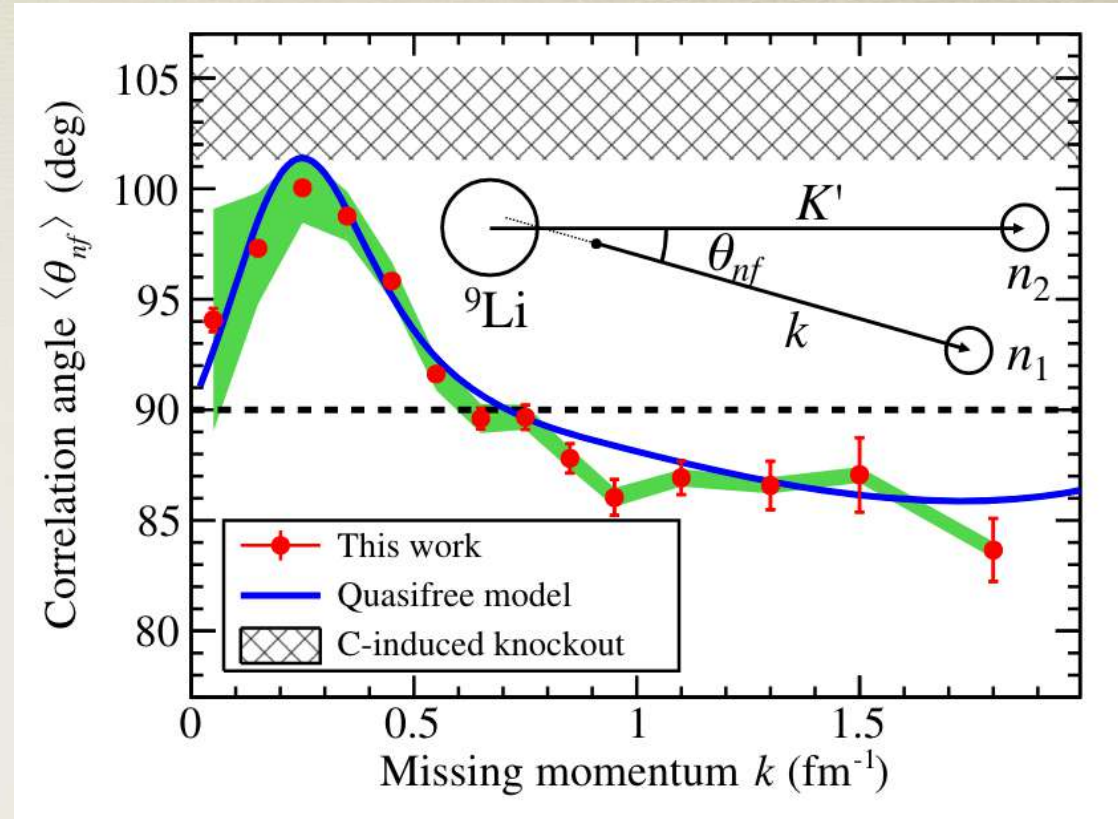
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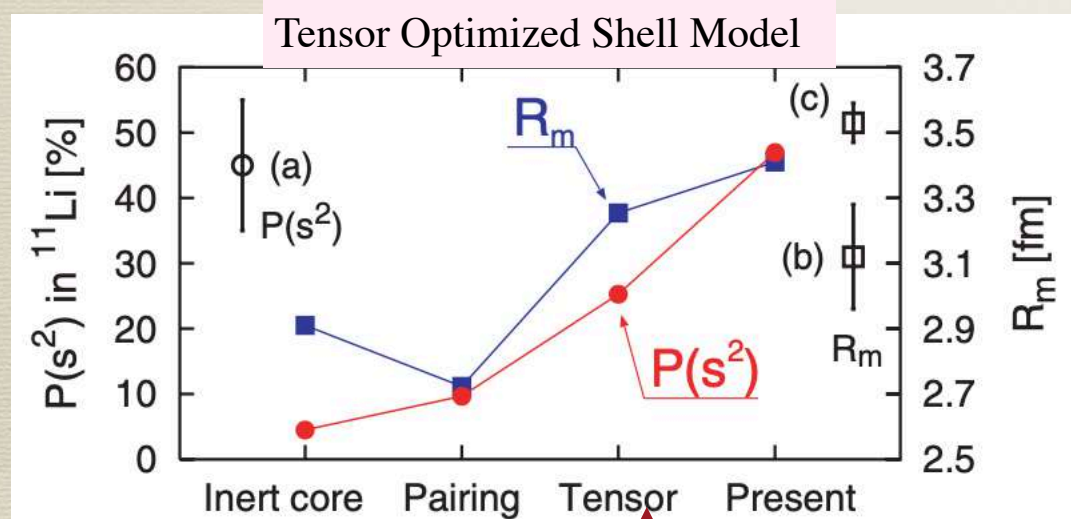
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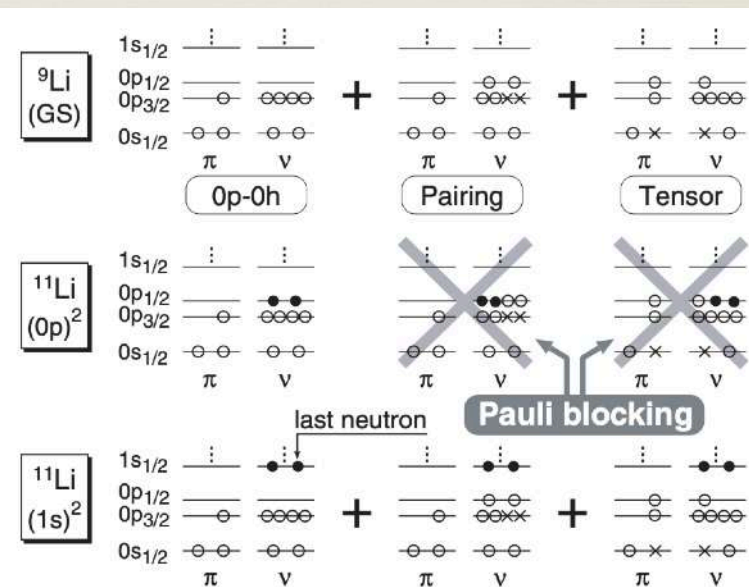


Di-neutron localized at the surface
i.e. Core - nn distance of $\sim 3.6 \text{ fm}$



T. Myo *et al.*, PRC (2007)

$2s_{1/2}$ lowering : predicted due to Tensor force - 2π exchange



^{29}F : Borromean Neutron halo

@ RIKEN-BigRIPS+ZDS

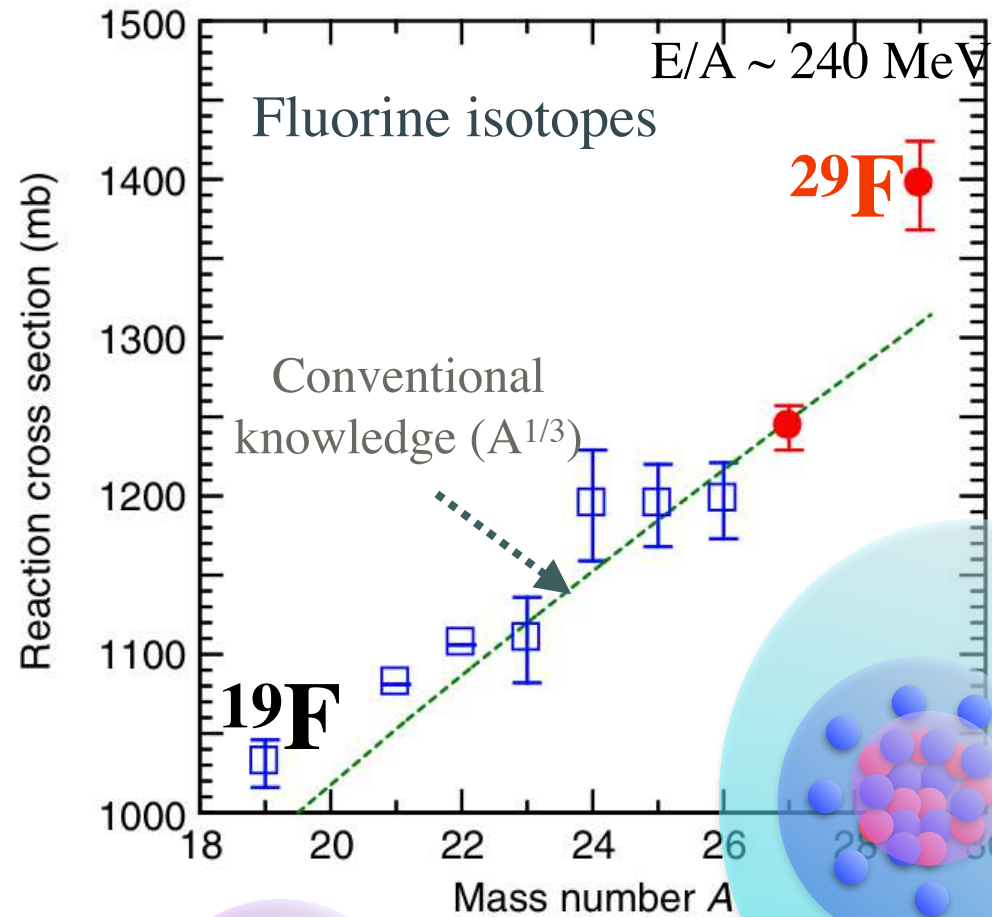
$$S_{2n} = 1.4(6) \text{ MeV}$$

	^{28}Al	^{30}Al	^{32}Al
12	^{28}Mg	^{30}Mg	^{32}Mg
10	^{26}Na	^{28}Na	^{30}Na
8	^{24}F	^{26}F	^{29}F
6	^{22}N		

Neutrons (N) →

18 20

^{22}C



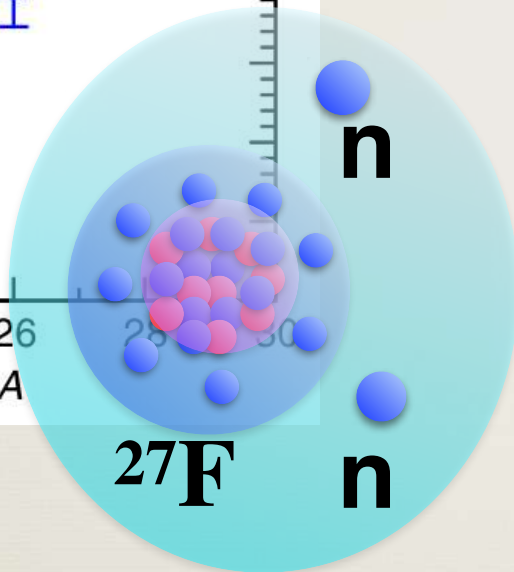
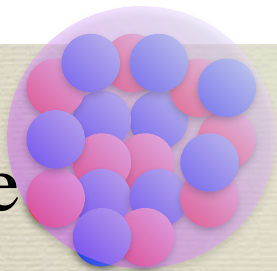
^{29}F Large increase in

$\sigma_R \rightarrow$ matter radius

$$\Delta R_m [^{29}\text{F} - ^{27}\text{F}] \sim 0.39(18) \text{ fm}$$

Two neutron halo

Stable



^{29}F

S. Bagchi et al., Phys. Rev. Lett. 124 (2020) 222504

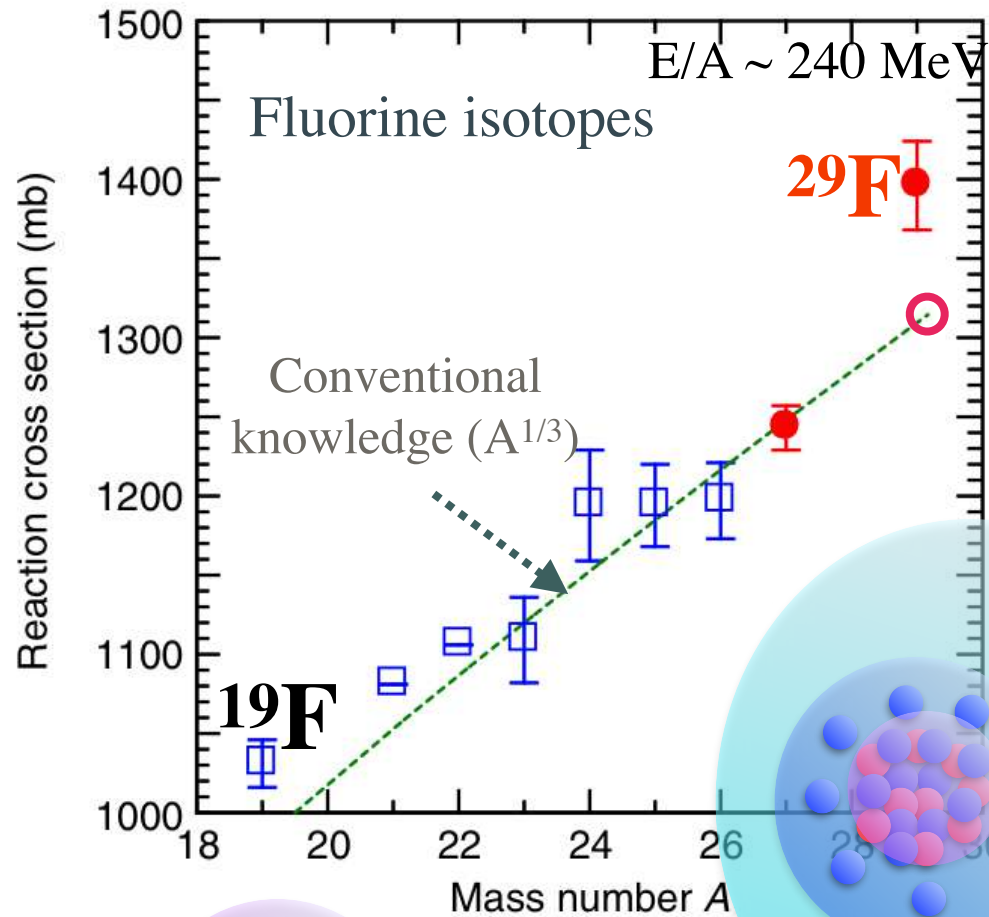
■ A. Honma et al., JPS Proc. 14 (2017) 021010



^{29}F : Borromean Neutron halo

@ RIKEN-BigRIPS+ZDS

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^{29}F Large increase in

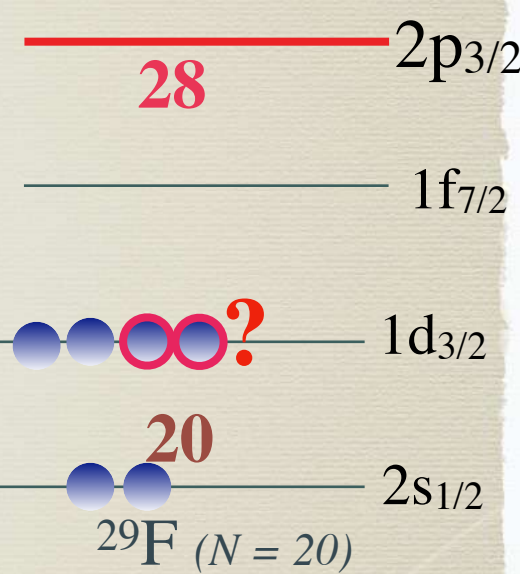
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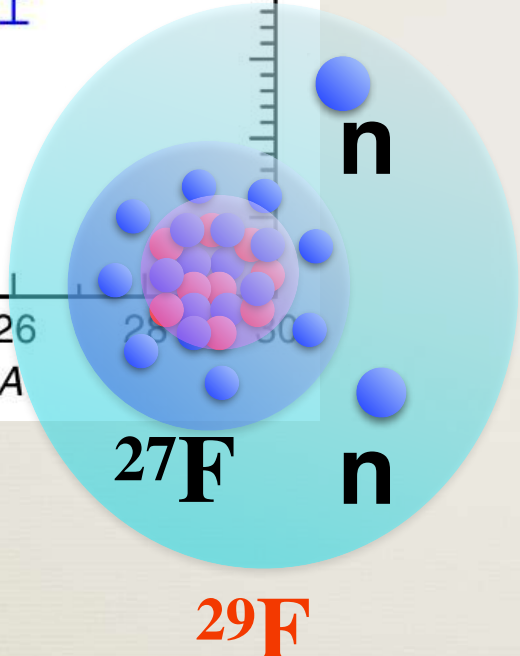
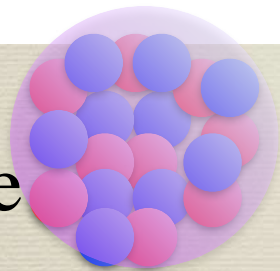
Two neutron halo

	^{28}Al	^{30}Al	^{32}Al	
12		^{28}Mg	^{30}Mg	^{32}Mg
	^{26}Na		^{28}Na	^{30}Na
10		^{26}Ne	^{28}Ne	^{30}Ne
	^{24}F		^{26}F	^{28}F
8			18	20
	^{22}N			
6				

Neutrons (N) \rightarrow



Stable



S. Bagchi et al., Phys. Rev. Lett. 124 (2020) 222504

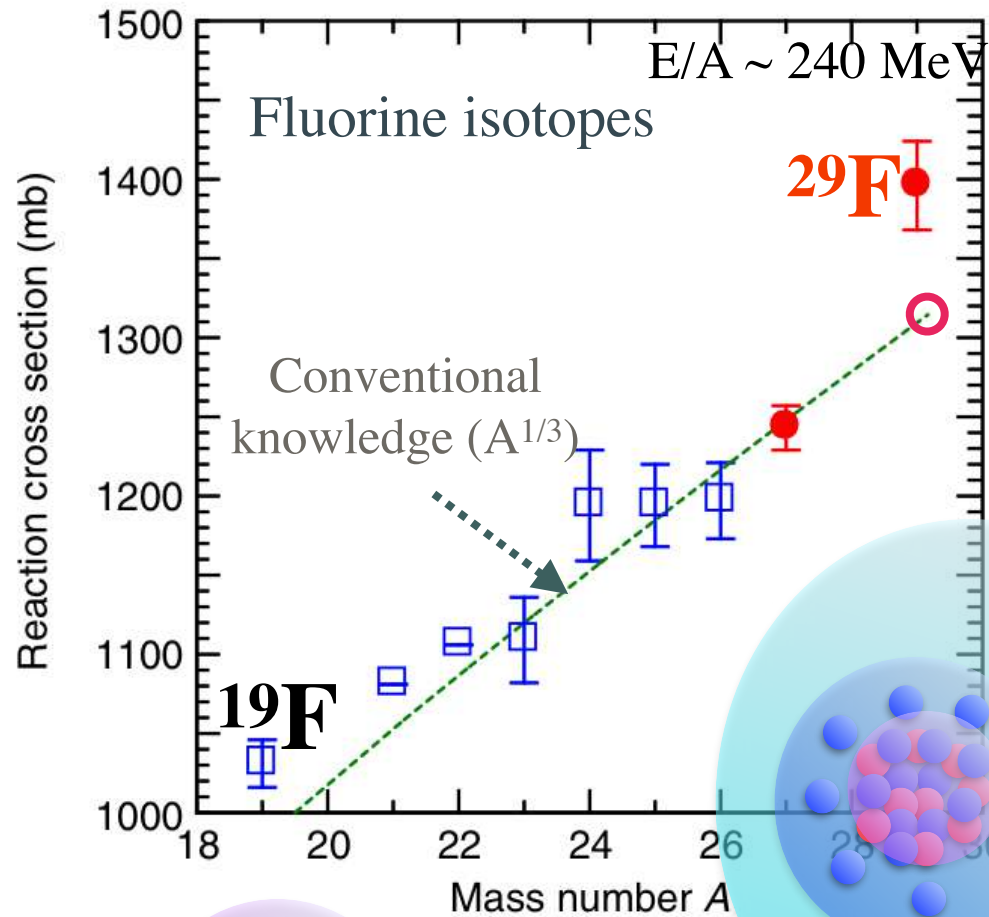
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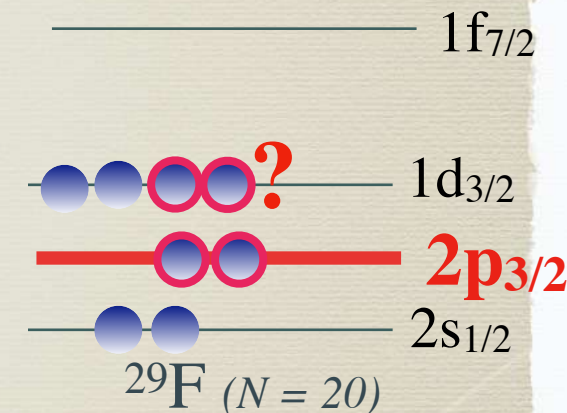
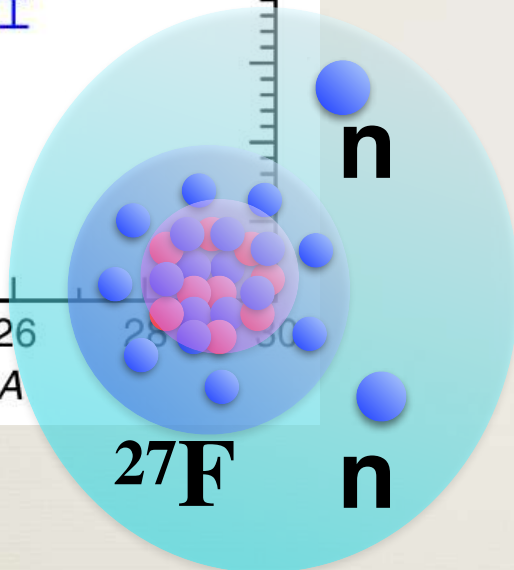
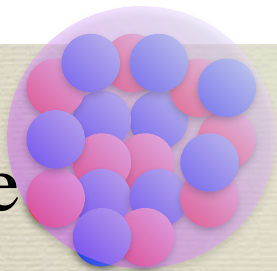
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Neutrons (N) \rightarrow

^{22}C

Stable



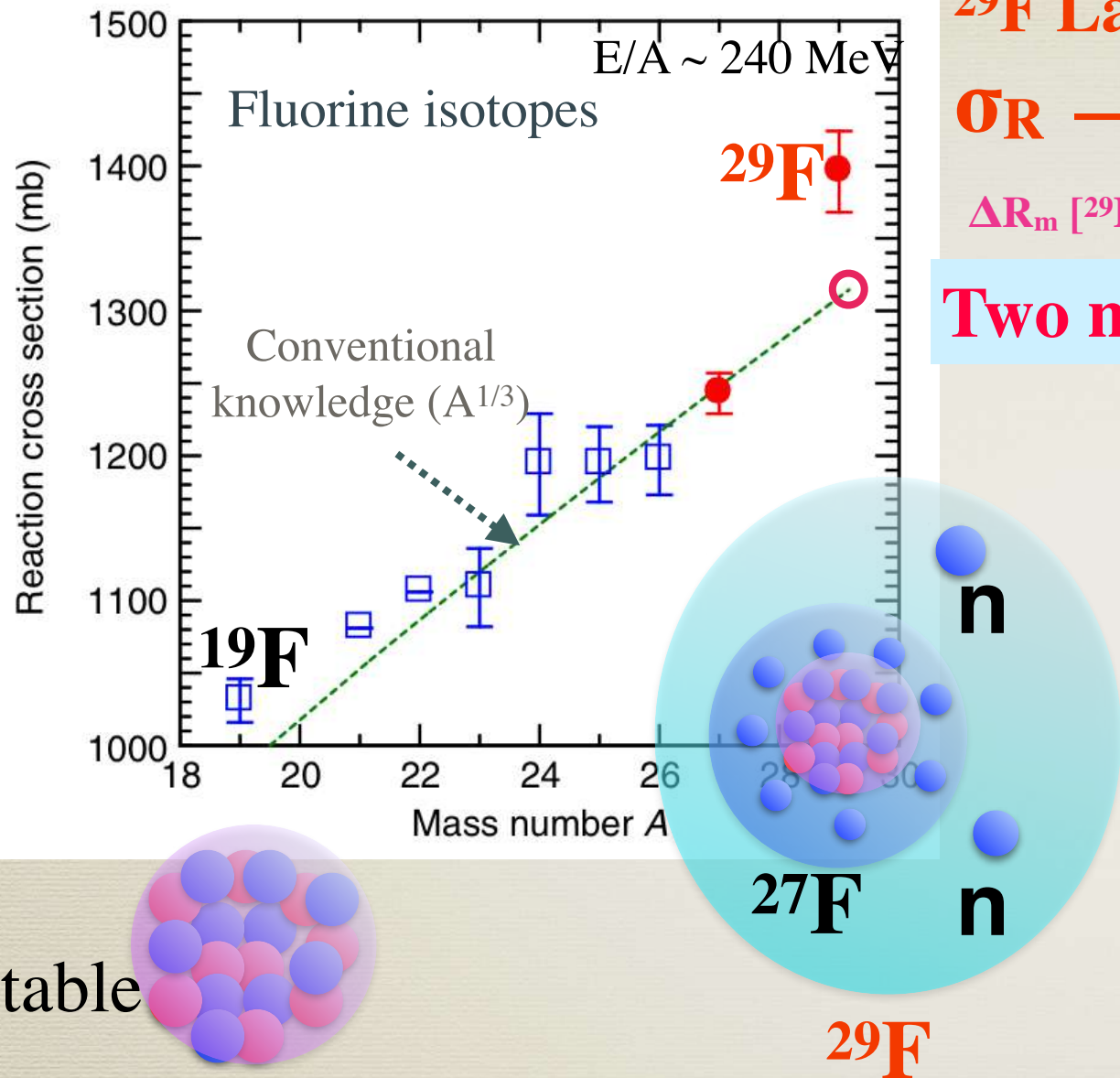
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^{29}F : Borromean Neutron halo

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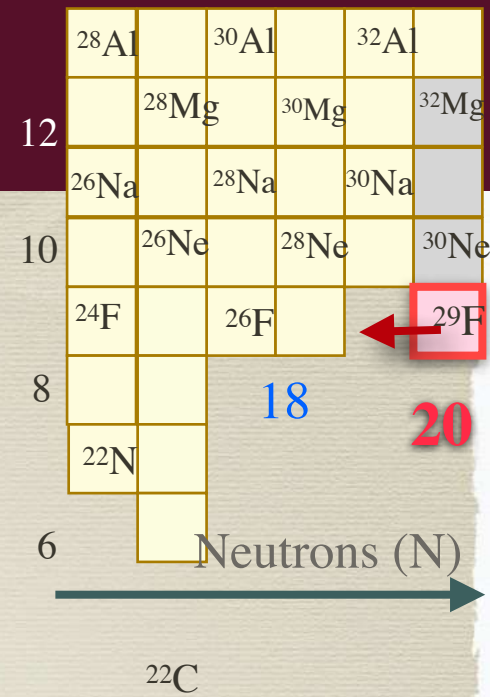


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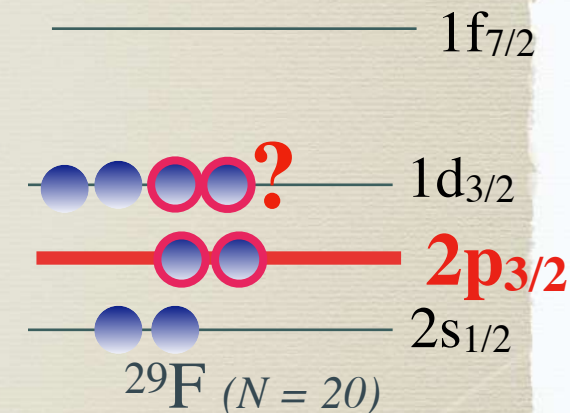
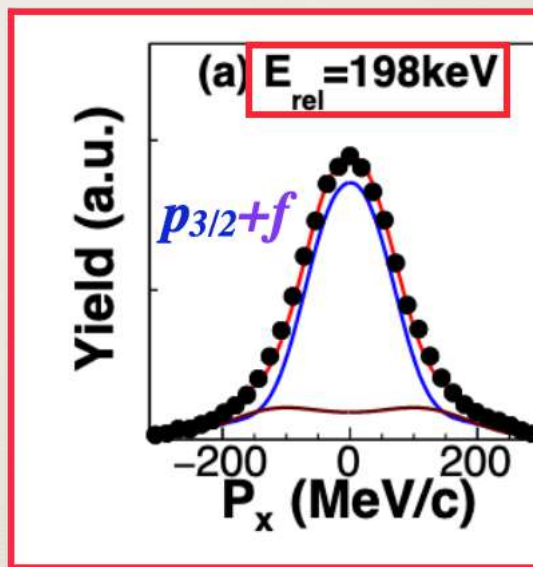
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Two neutron halo

$$S_{2n} = 1.4(6) \text{ MeV}$$



@ RIKEN-SAMURAI



Narrow momentum distribution \rightarrow $2p_{3/2}$ orbital

S. Bagchi et al., Phys. Rev. Lett. 124 (2020) 222504

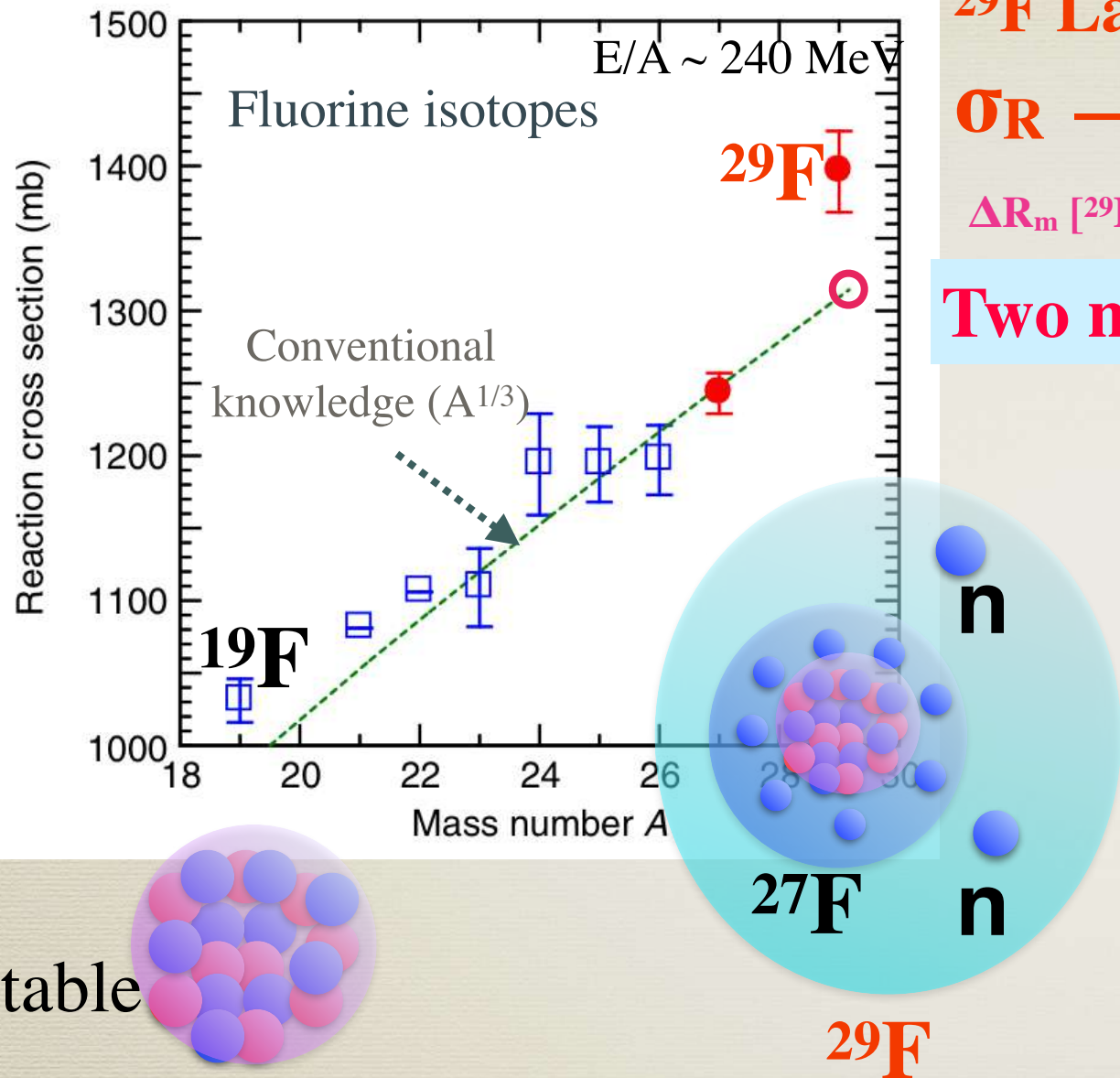
A. Honma et al., JPS Proc. 14 (2017) 021010

A. Revel et al., Phys. Rev. Lett. 124 (2020) 152502



^{29}F : Borromean Neutron halo

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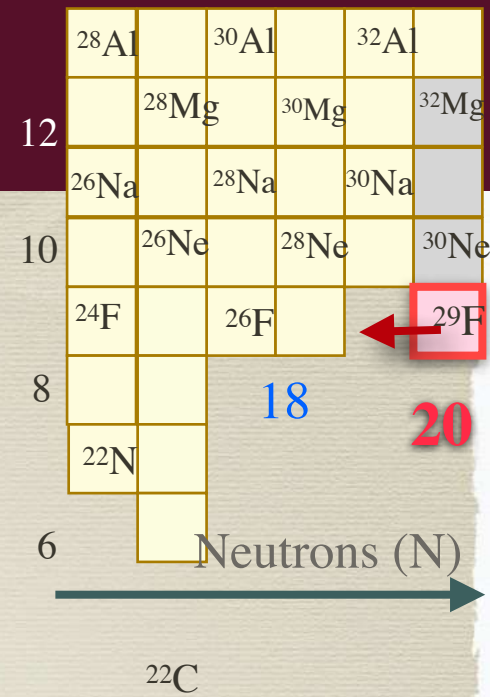


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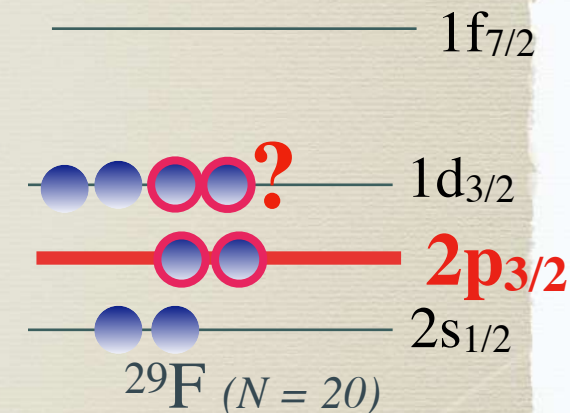
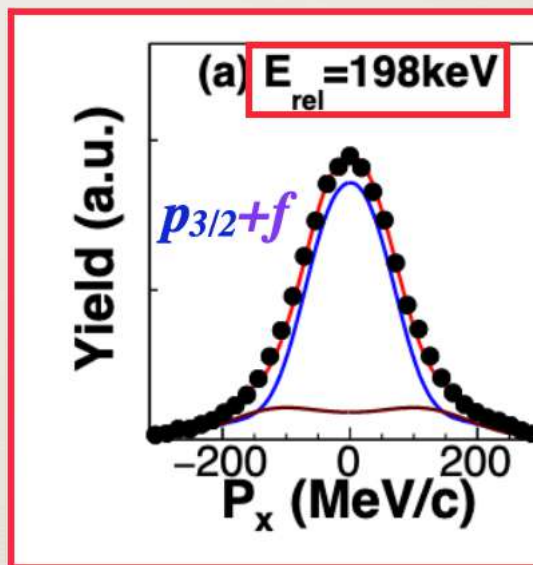
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S. Bagchi et al., Phys. Rev. Lett. 124 (2020) 222504

A. Honma et al., JPS Proc. 14 (2017) 021010

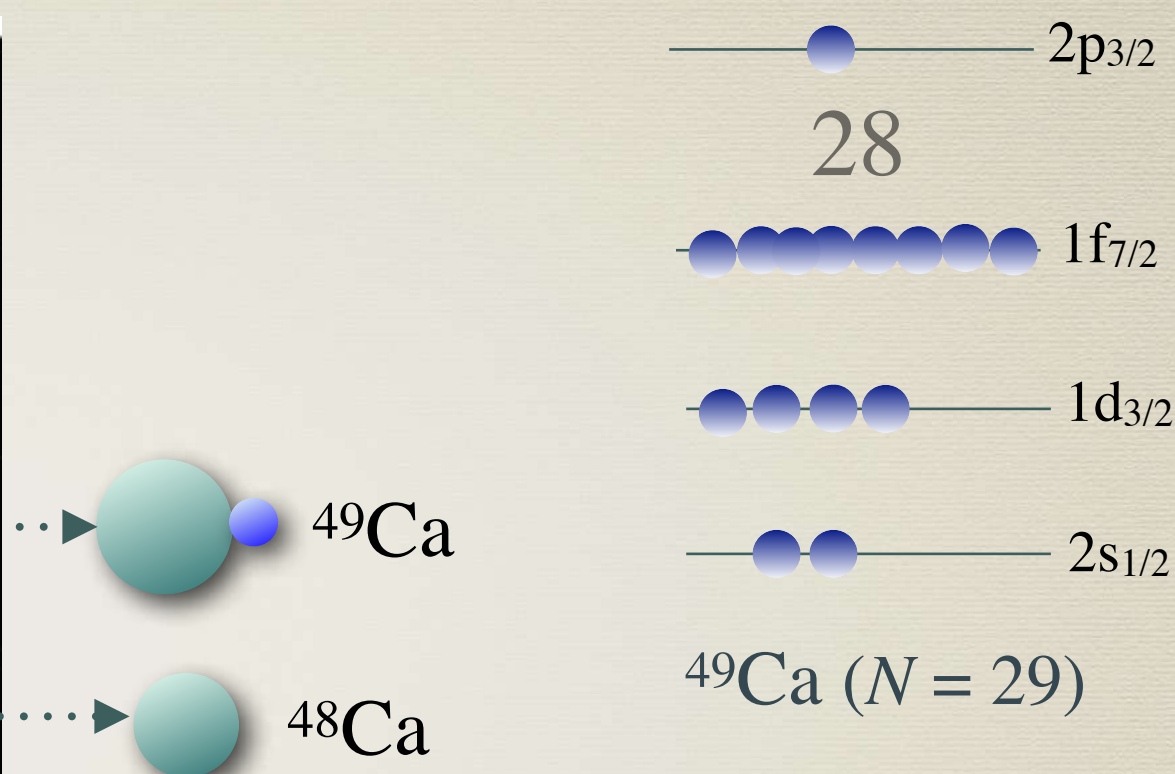
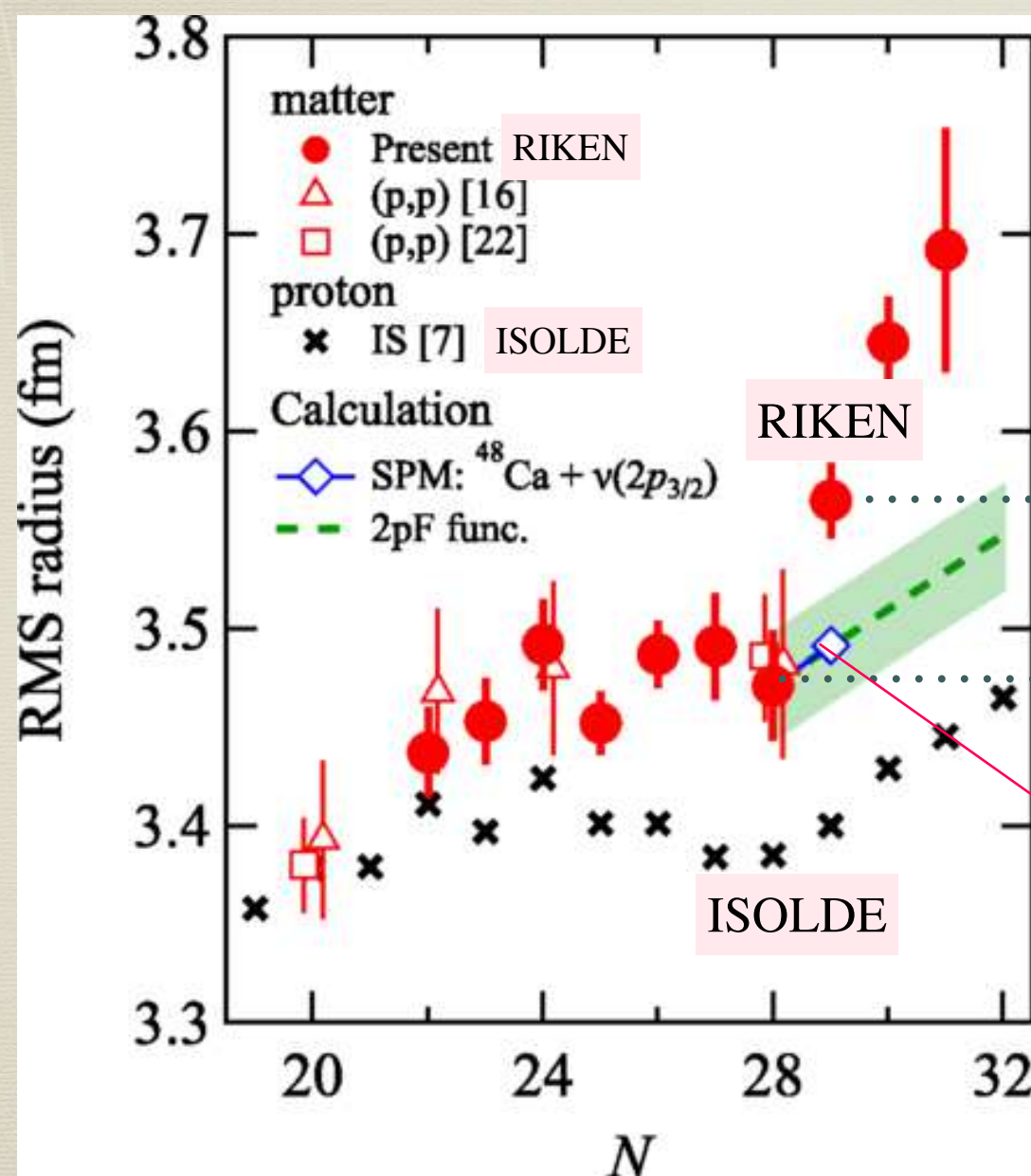
A. Revel et al., Phys. Rev. Lett. 124 (2020) 152502

$N = 20, 28$ shells vanish in a Borromean halo

Enlarged size of Ca isotopes beyond $N = 28$

@ RIKEN - BigRIPS

M. Tanaka, M. Takechi, A. Homma *et al.*, Phys. Rev. Lett. 124 (2020) 102501



R_m for ^{49}Ca is not explained by $R_m(^{48}\text{Ca}) + n(2p_{3/2})$

Dip in the proton distribution radius for ^{48}Ca shows $N = 28$ is closed shell

The ^{48}Ca core is enlarged in neutron-rich Ca isotopes : $p - n$ attractive force

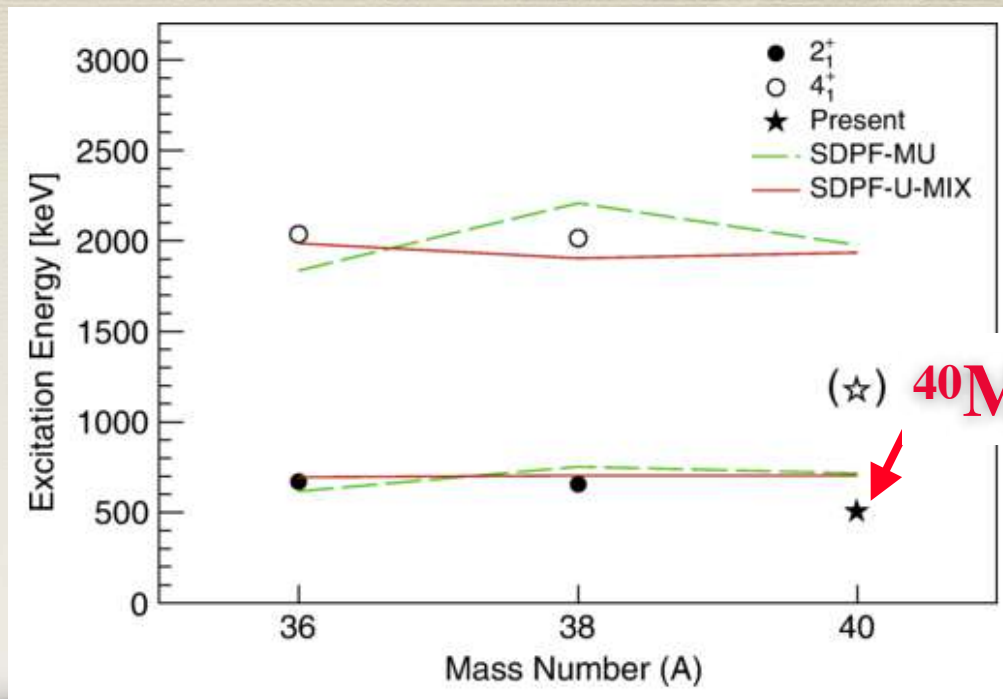
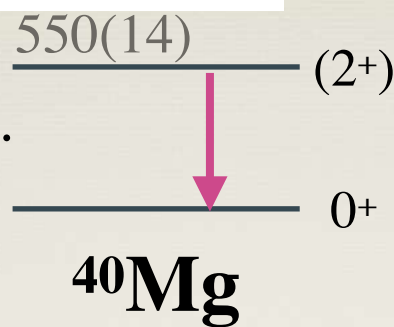
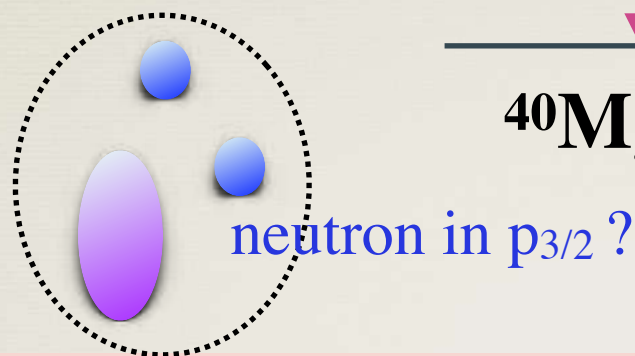
$N = 28$ shell closure breakdown

$^{41}\text{Al} \rightarrow ^{40}\text{Mg}^*$ @ RIKEN - BigRIPS & DALI2

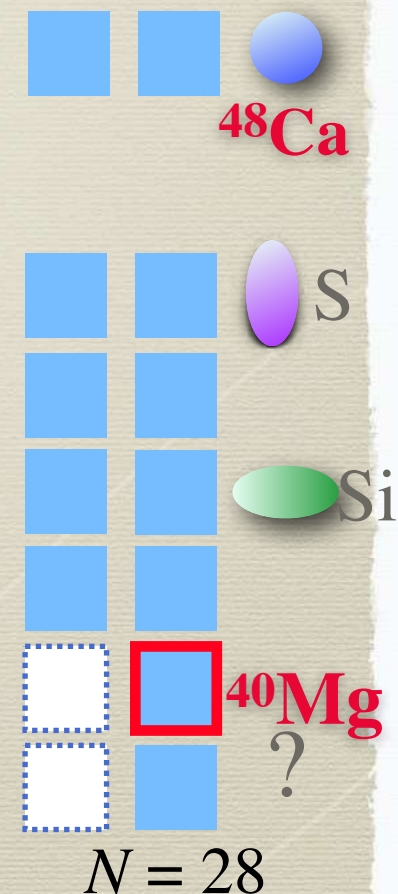
H. L. Crawford et al., Phys. Rev. Lett. 122 (2019) 052501

(2⁺) Excited state energy decreases in ^{40}Mg .

Not explained by shell model interactions



$N = 28$ shell gap disappears in ^{40}Mg .
Is ^{40}Mg a Borromean halo ?



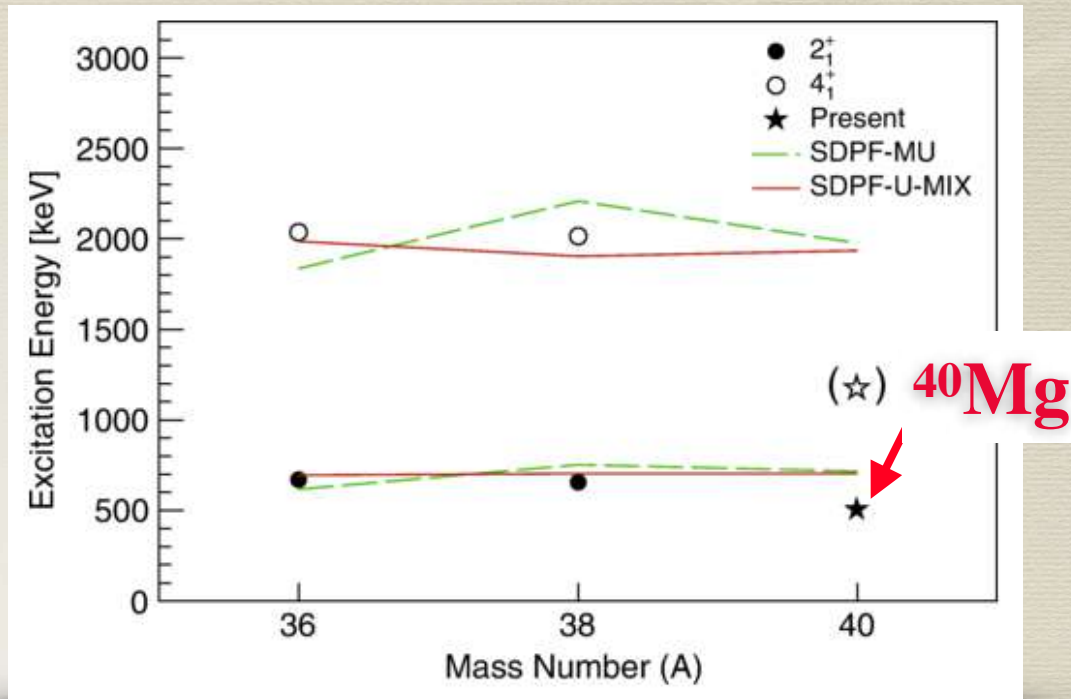
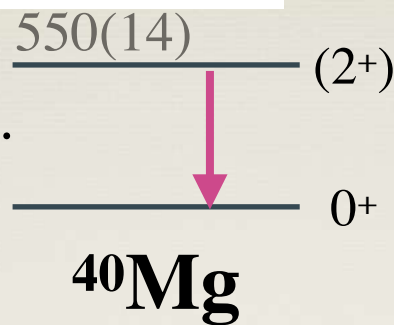
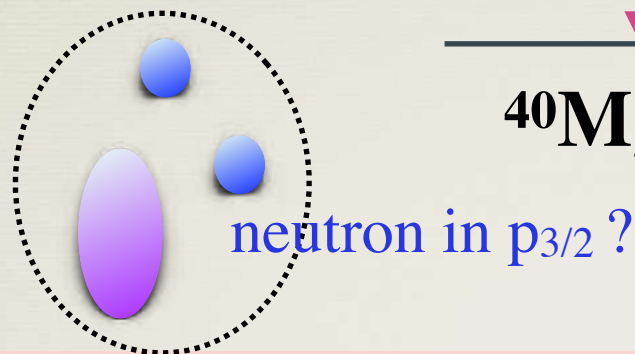
$N = 28$ shell closure breakdown

$^{41}\text{Al} \rightarrow ^{40}\text{Mg}^*$ @ RIKEN - BigRIPS & DALI2

H. L. Crawford et al., Phys. Rev. Lett. 122 (2019) 052501

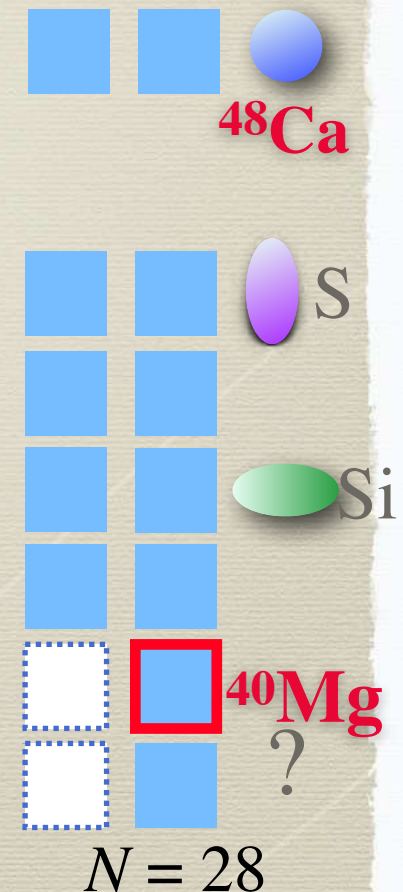
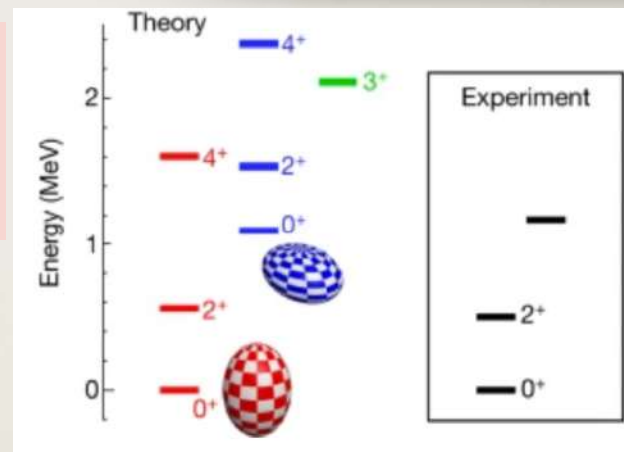
(2⁺) Excited state energy decreases in ^{40}Mg .

Not explained by shell model interactions



$N = 28$ shell gap disappears in ^{40}Mg .
Is ^{40}Mg a Borromean halo ?

N. Tsunoda, et al., Nature, 2020



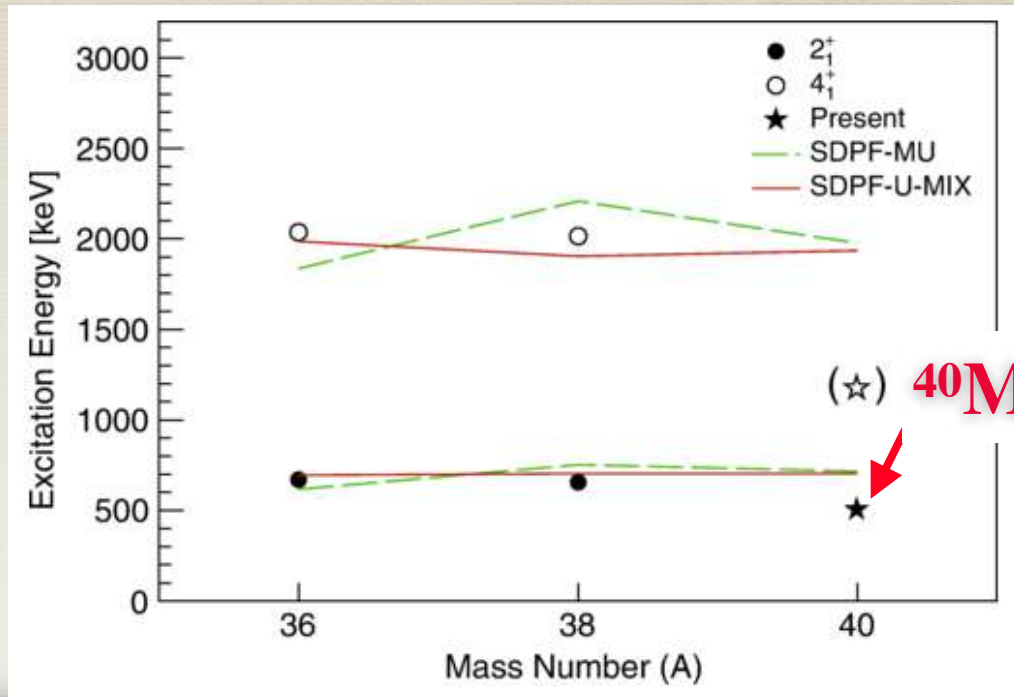
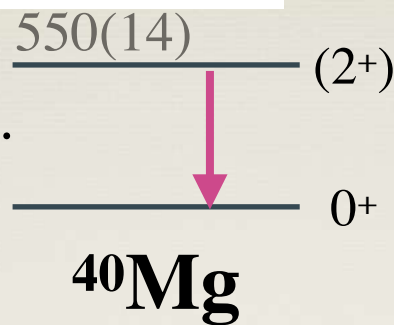
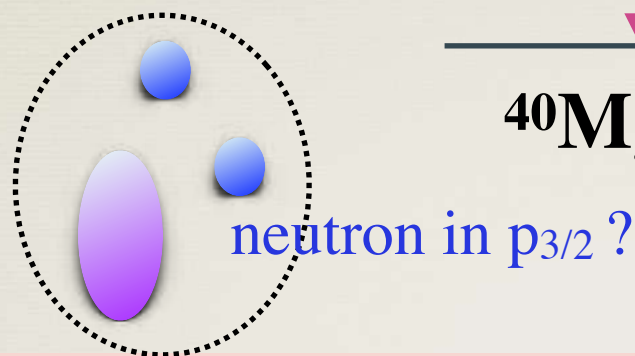
N = 28 shell closure breakdown

$^{41}\text{Al} \rightarrow ^{40}\text{Mg}^*$ @ RIKEN - BigRIPS & DALI2

H. L. Crawford et al., Phys. Rev. Lett. 122 (2019) 052501

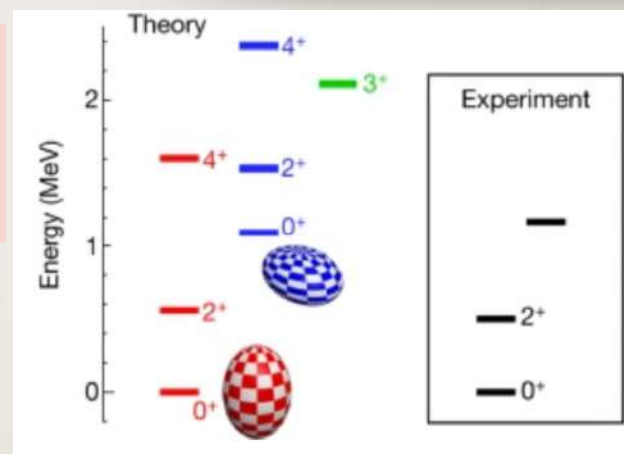
(2+) Excited state energy decreases in ^{40}Mg .

Not explained by shell model interactions



**N = 28 shell gap disappears in ^{40}Mg .
Is ^{40}Mg a Borromean halo ?**

N. Tsunoda, et al., Nature, 2020



^{43}S ($^{43}\text{S}^*m$) + ^{209}Bi (Coulex)

@ NSCL - CAESAR

B. Longfellow, D. Weisshaar, A. Gade et al., Phys. Rev. Lett. 125 (2020) 232501

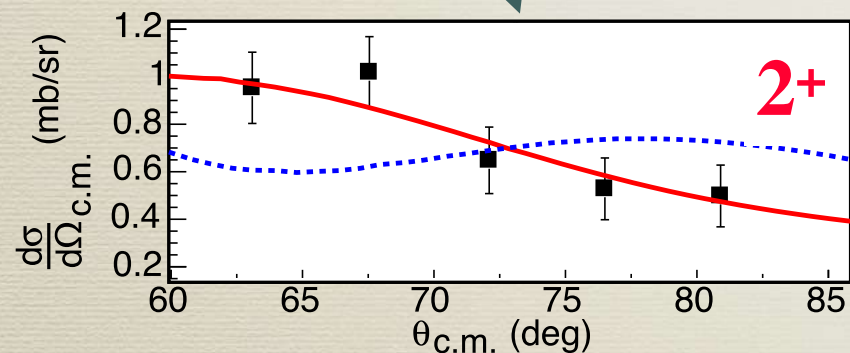
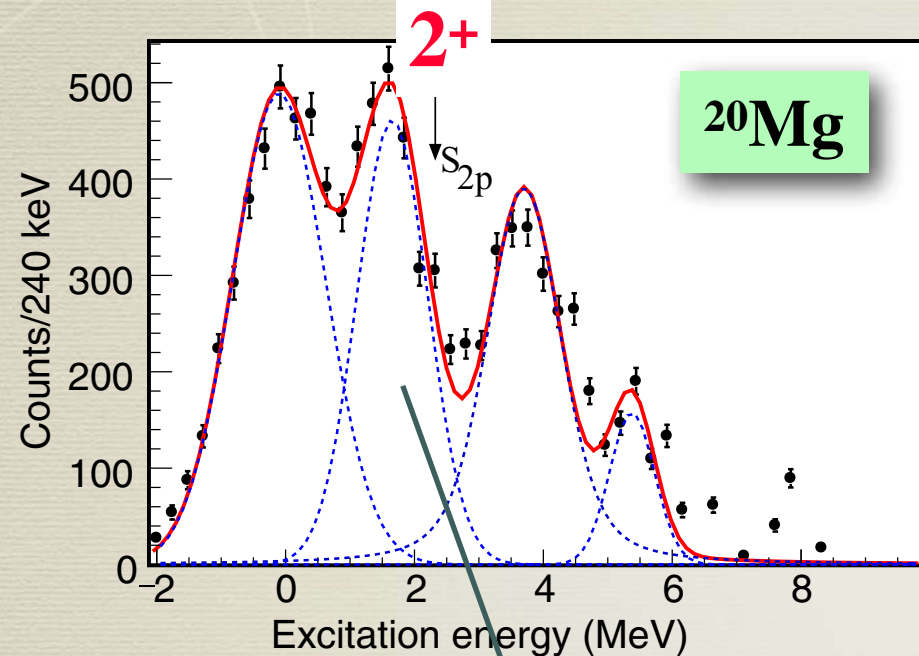
Shape co-existence in ^{43}S - degeneracy of $2p_{3/2}$ and $1f_{7/2}$ orbitals



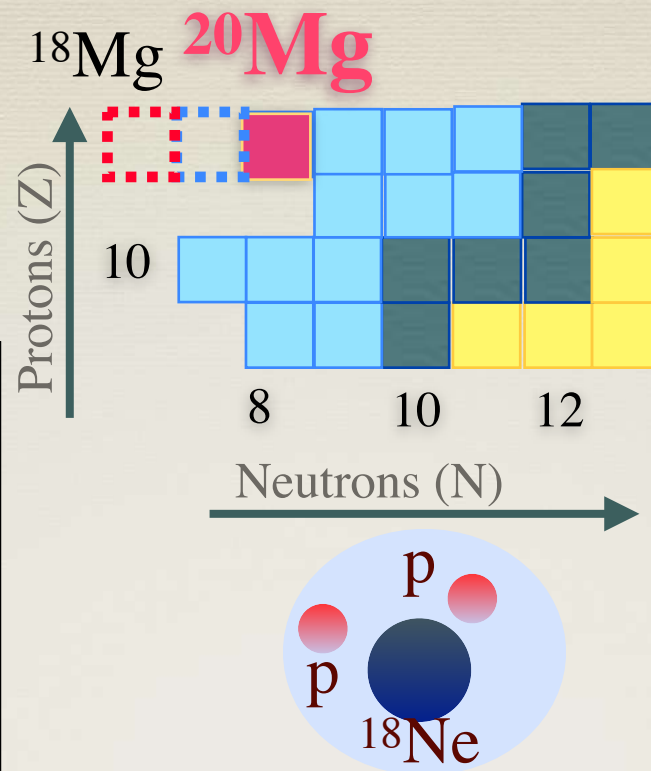
$N = 8$ vanishing @ Proton Drip Line ?

@ TRIUMF - IRIS

J.S. Randhawa, R.K., M. Holl et al.
Phys. Rev. C 99 (2019) 021301(R)

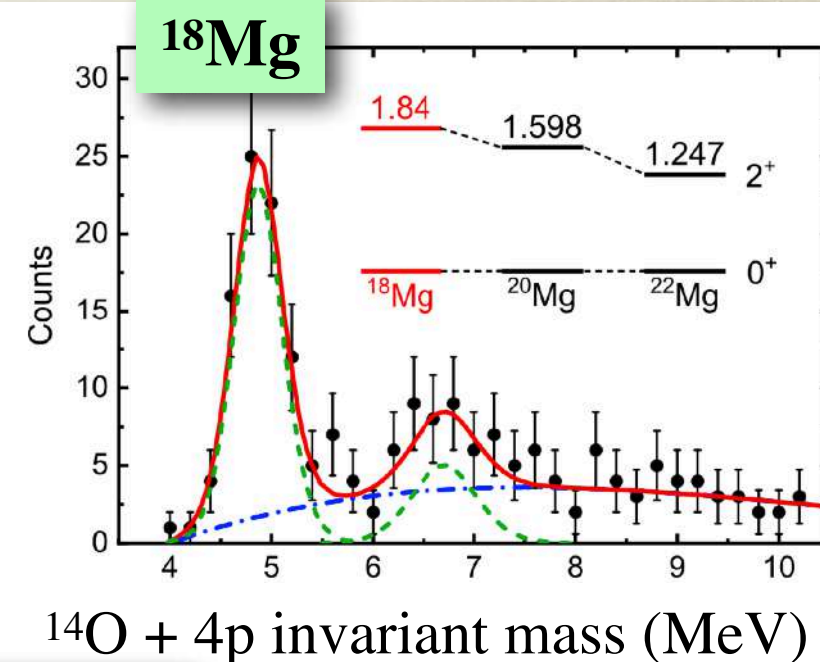


$N = 8$ Large quadrupole deformation $\beta = 0.44(4)$



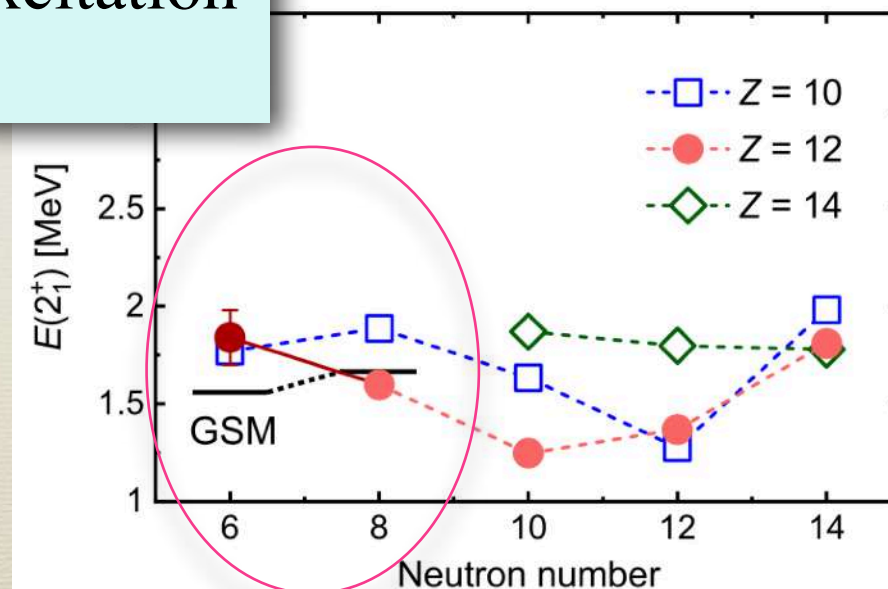
@ NSCL- HiRA

Y. Jin, C.Y. Niu, K.W. Brown *et al.*
Phys. Rev. Lett. 127 (2021) 262502



$N = 10$ Higher 2^+ excitation energy than $N = 8$

Challenge for theory



$N = 8$ shell gap quenching hinted at the proton drip-line from large deformation & lower energy of excited state

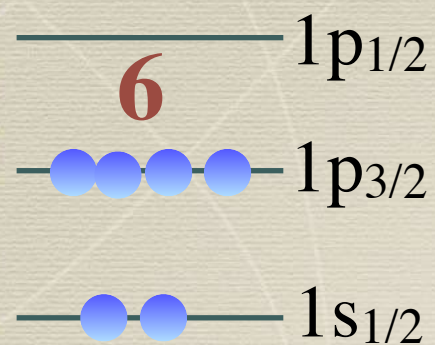
New shell gaps



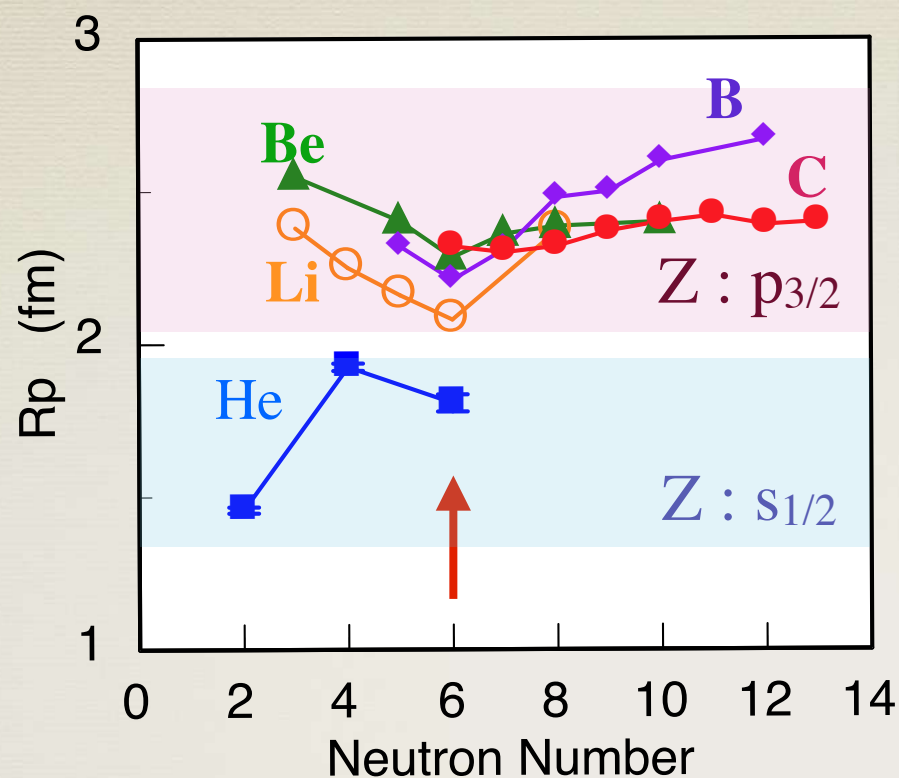
^8He : $N = 6$ new sub-shell

@ TRIUMF - IRIS

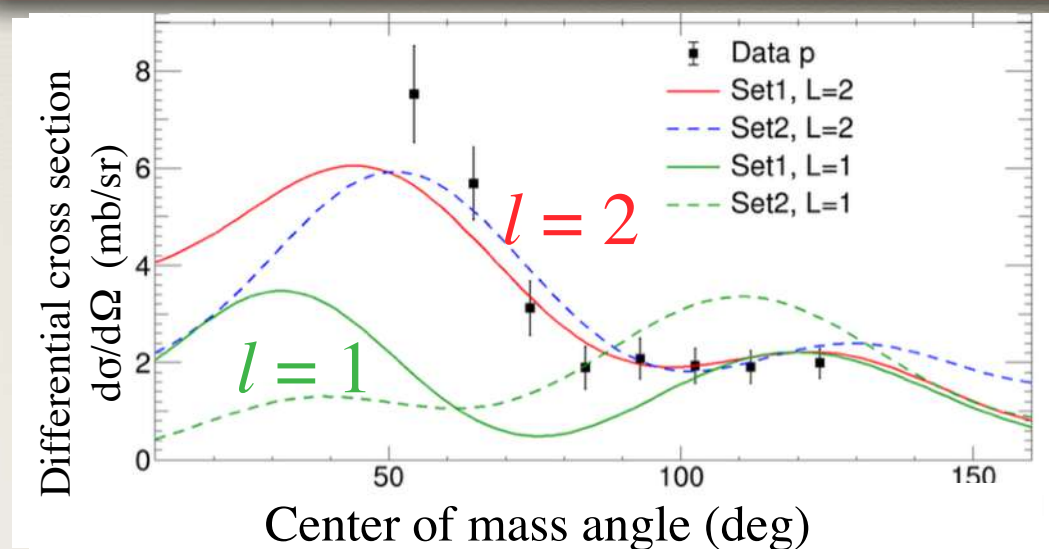
M. Holl, R.K., Z.H, Sun *et al.* Phys. Lett. B 822 (2021) 135748



Dip in R_p : Closed shell $N = 6$

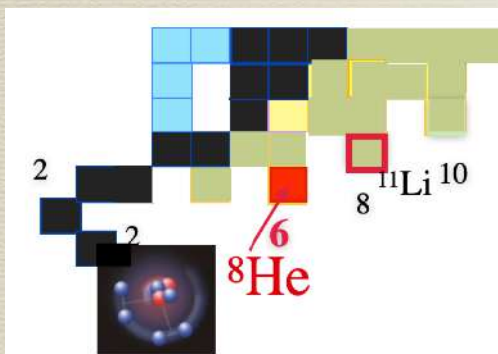


Doubly Magic - but deformed !



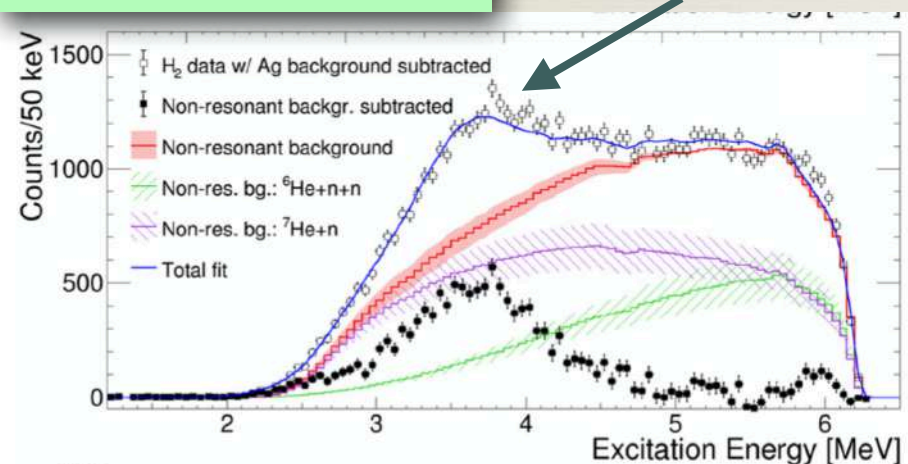
■ **Resonance (2^+) shows deformation**

$$\beta_2 = 0.40 \pm 0.03$$



High Excitation energy $\sim E_x \sim 3.5$ MeV : closed shell $N = 6$

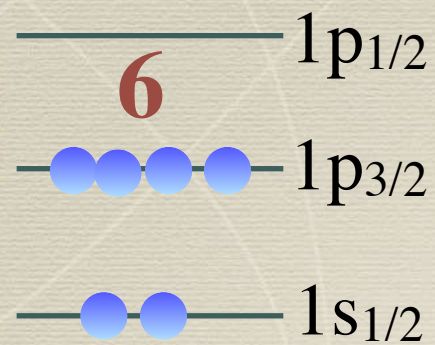
$^8\text{He}(p, p')^8\text{He}^*$



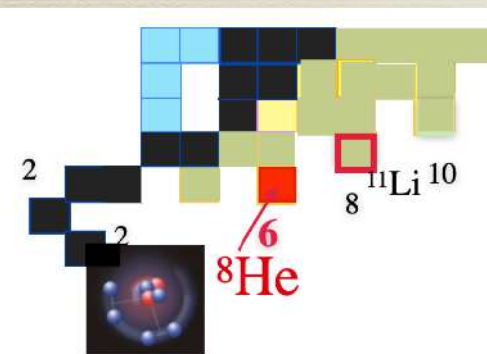
^8He : $N = 6$ new sub-shell

@ TRIUMF - IRIS

M. Holl, R.K., Z.H, Sun *et al.* Phys. Lett. B 822 (2021) 135748

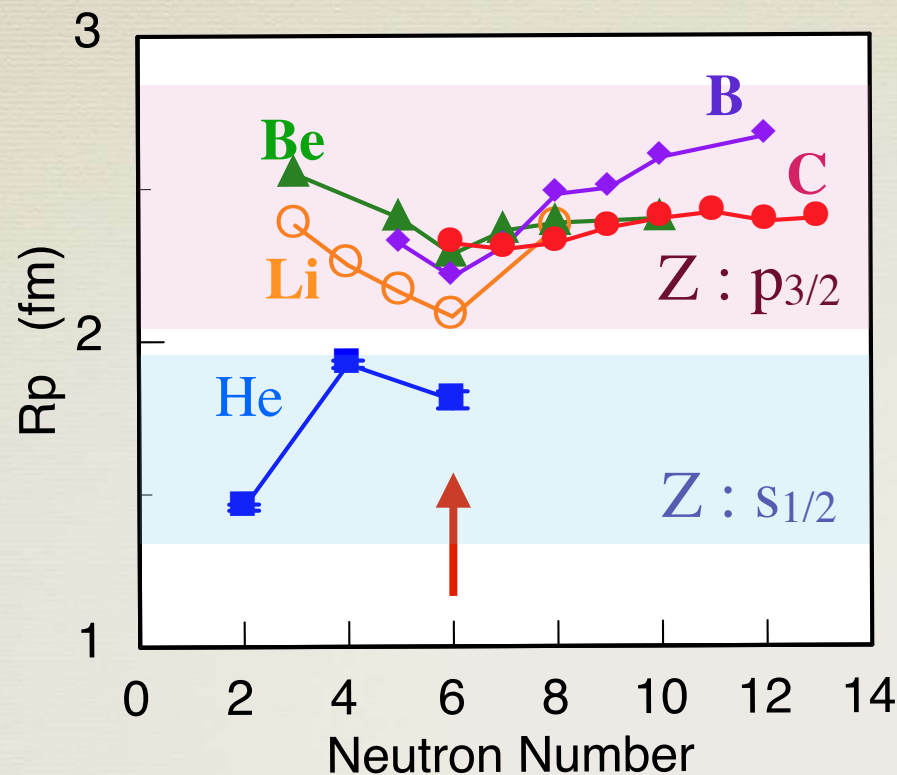
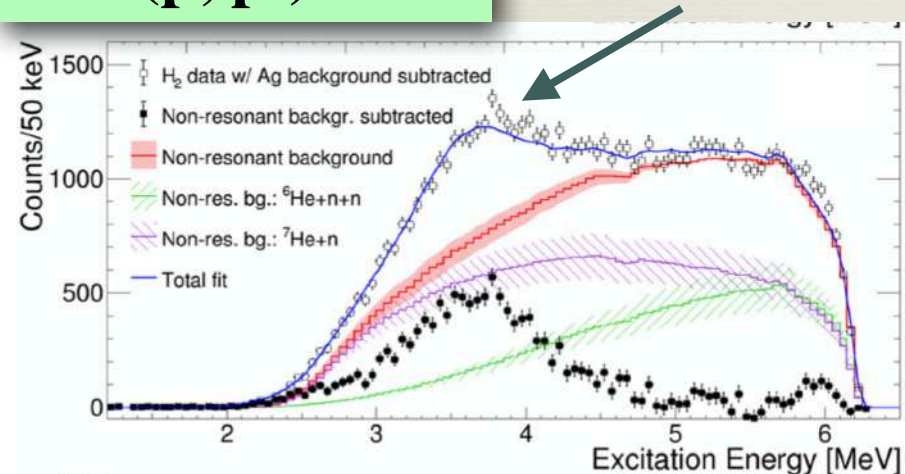


Dip in R_p : Closed shell $N = 6$

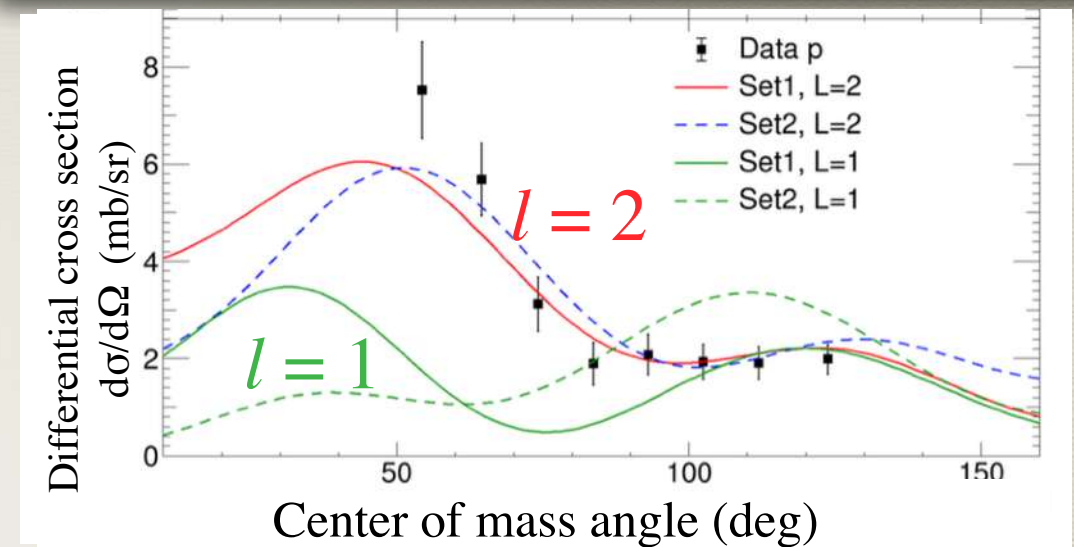


$^8\text{He}(p, p')^8\text{He}^*$

High Excitation energy $\sim E_x \sim 3.5$ MeV :
 closed shell $N = 6$



Doubly Magic - but deformed !



■ **Resonance (2^+) shows deformation**

$$\beta_2 = 0.40 \pm 0.03$$

No Core Shell Model with chiral force
 consistent with experimental observations

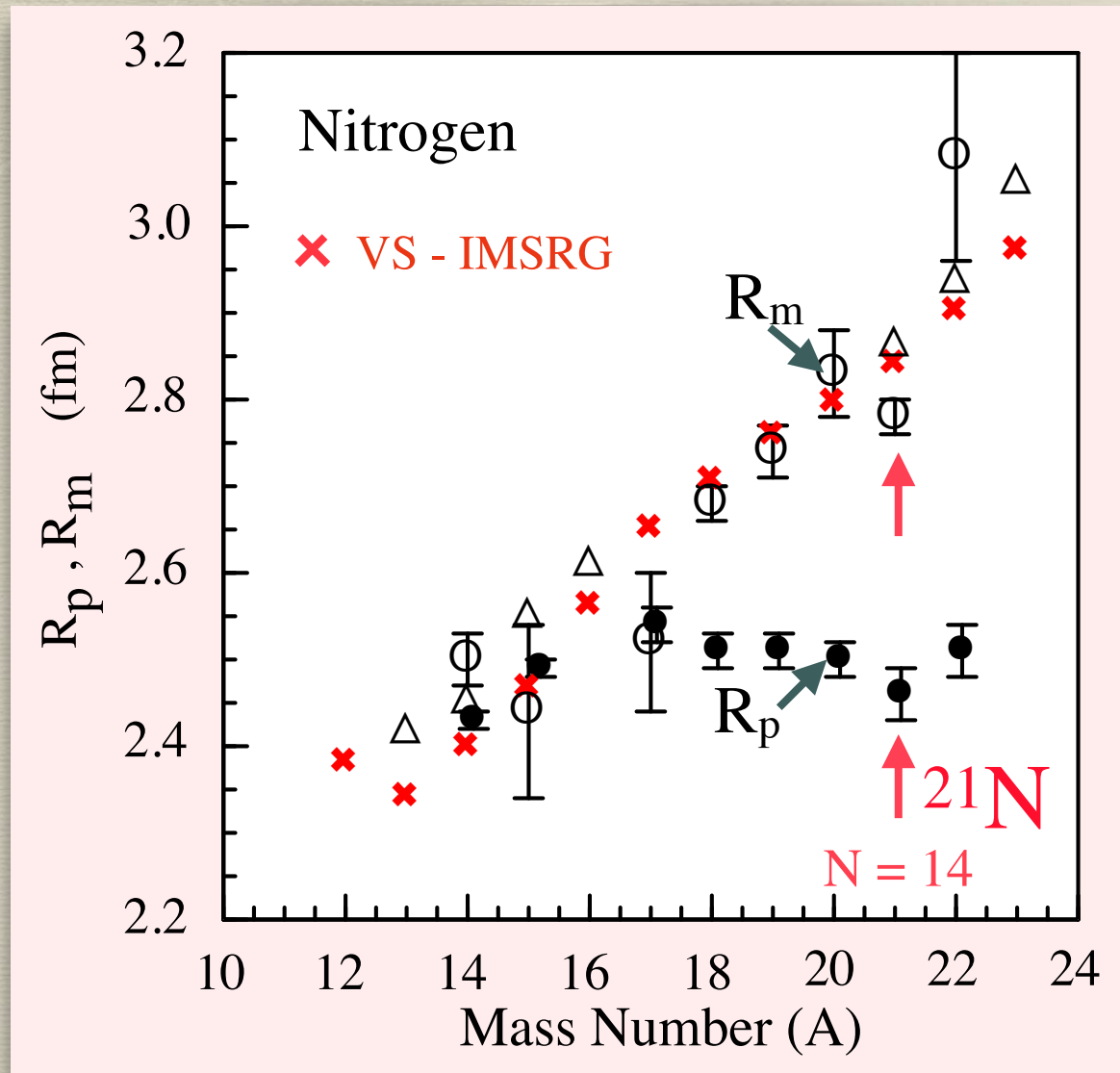
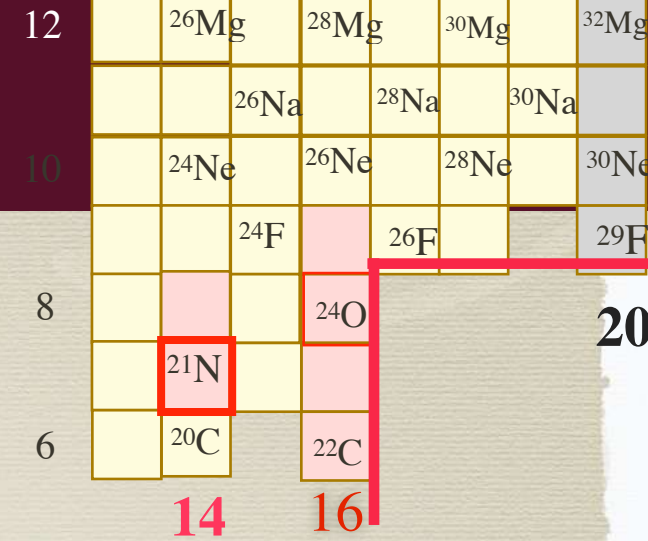
$$^8\text{He} (2^+) : \quad Q_n = 6.15 \text{ efm}^2, \quad Q_p = 0.60 \text{ efm}^2$$

^8He : spherical in protons and deformed in neutrons

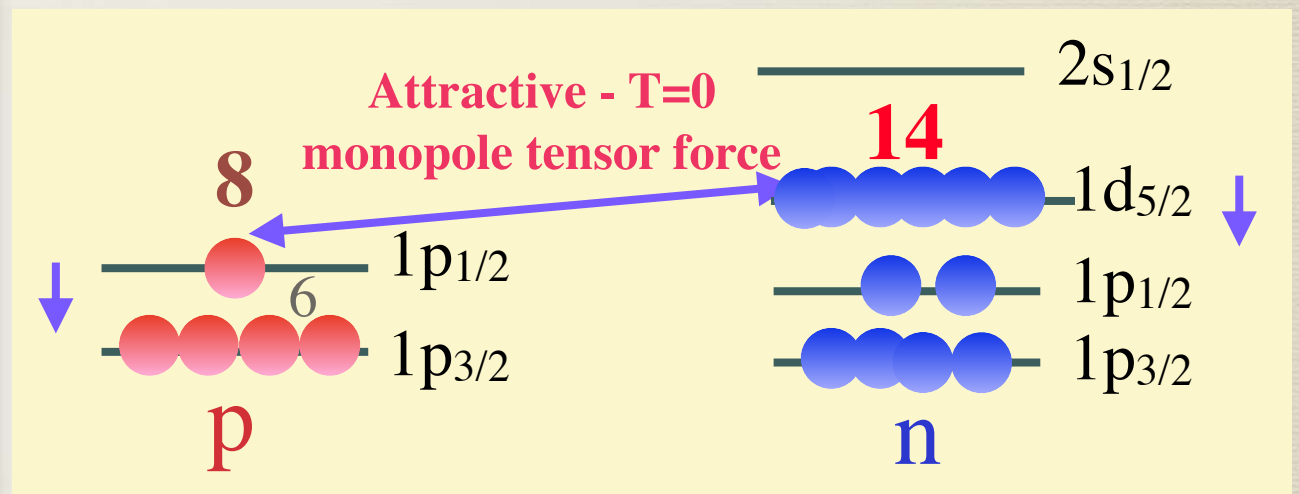
$N = 14$ new sub-shell

@ GSI - FRS

S. Bagchi, R.K., W. Horiuchi *et al.* Phys. Lett. B 790 (2019) 251



^{21}N : Dip in R_p and R_m directly shows effect of tensor force



Proton $p_{1/2}$ more strongly bound -> Smaller R_p . & $Z = 6$ gap reduced

Neutron $d_{5/2}$ more strongly bound -> Smaller R_m . & new gap @ $N = 14$

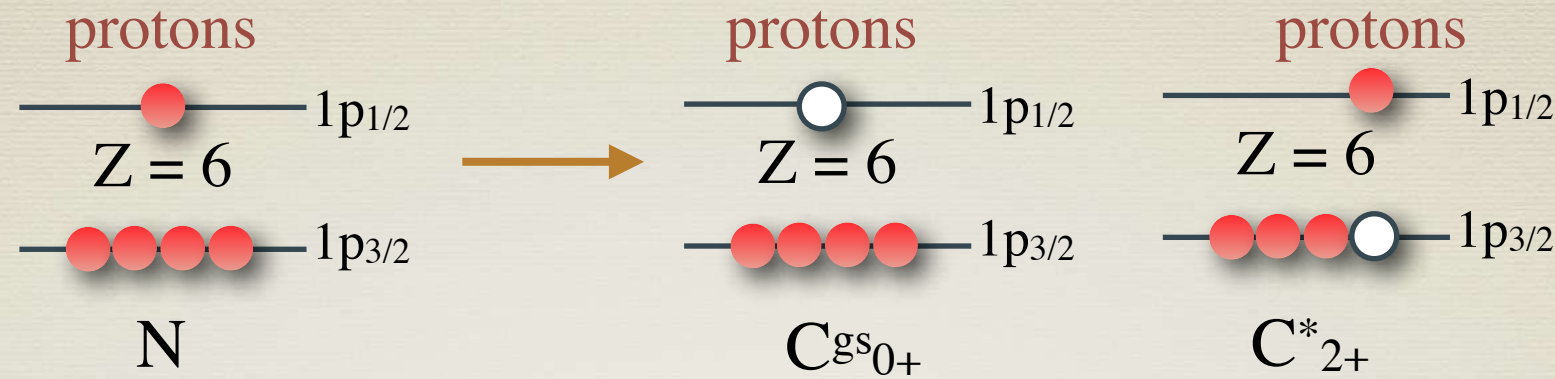
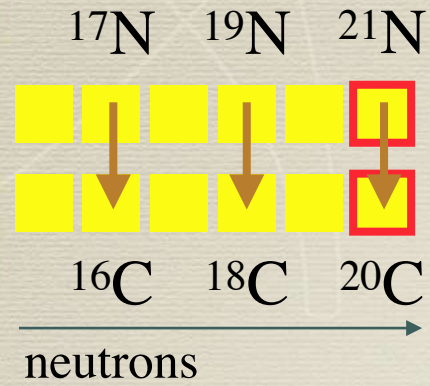
Dip in proton distribution radius shows new sub-shell gap @ $N = 14$

Quenching of proton sub-shell $Z = 6$

@ GSI - R³B/LAND

$A\text{N}(p, 2p)A-1\text{C}$

I. Syndicus, M. Petri, A.O. Macchiavelli *et al.* Phys. Lett. B 809 (2020) 135748

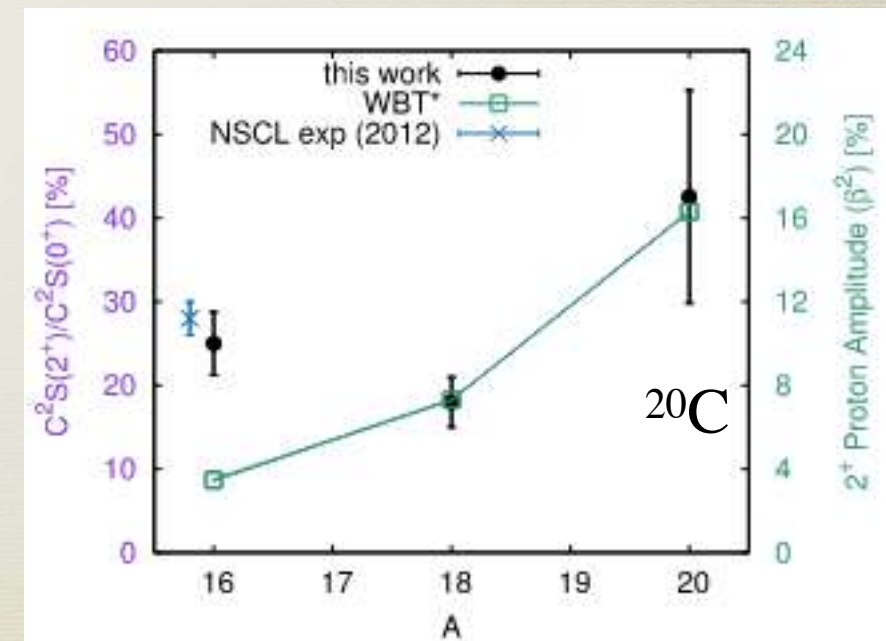
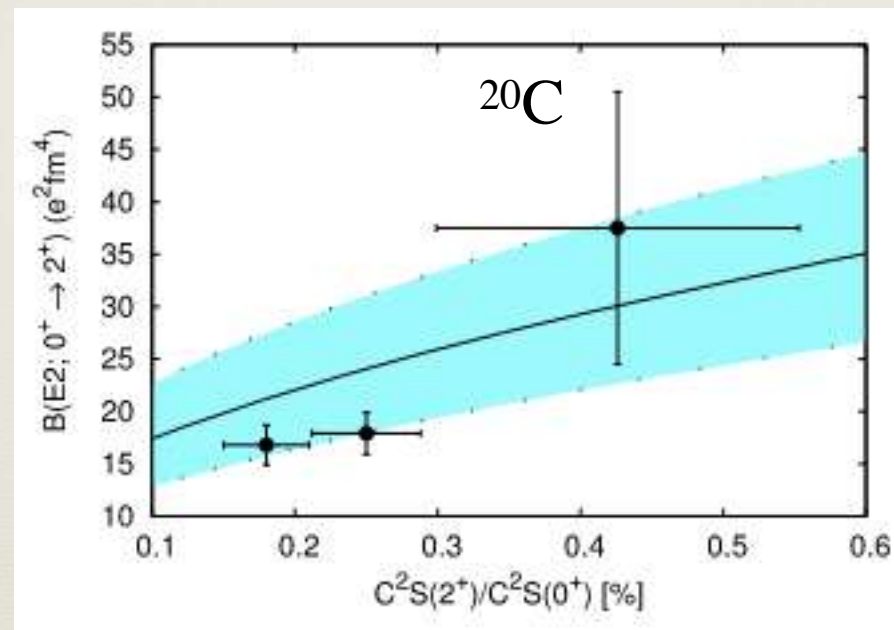


$$|2^+; A-1\text{C}\rangle \approx \alpha |v(sd)^n; J=2\rangle \otimes |\pi(p_{3/2})^4; J=0\rangle$$

proton excitation component: $2^+ + \beta |v(sd)^n; J=0\rangle \otimes |\pi(p_{3/2})^3(p_{1/2})^1; J=2\rangle$

Large quadrupole deformation \rightarrow [B(E2)] in ^{20}C

Deformation in ^{20}C correlated to proton excitation contribution



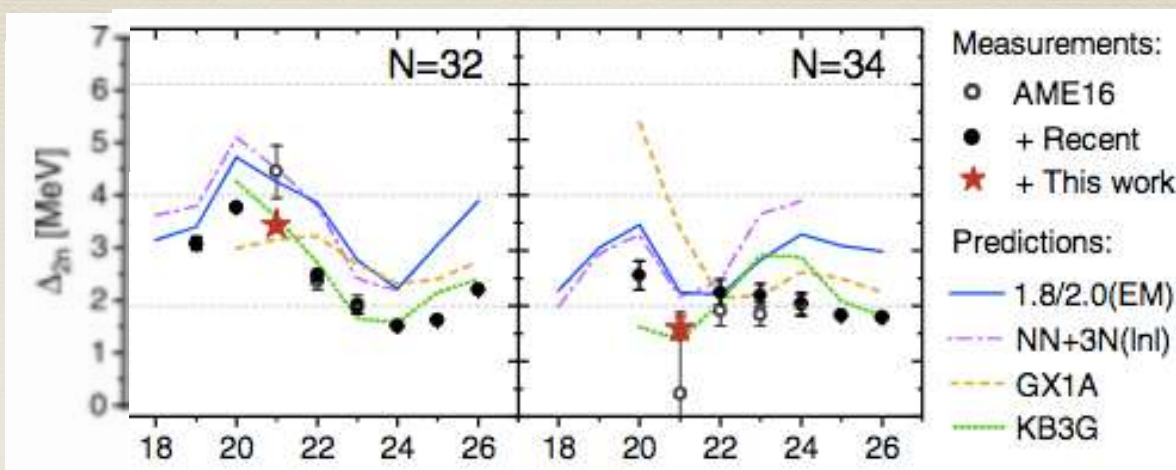
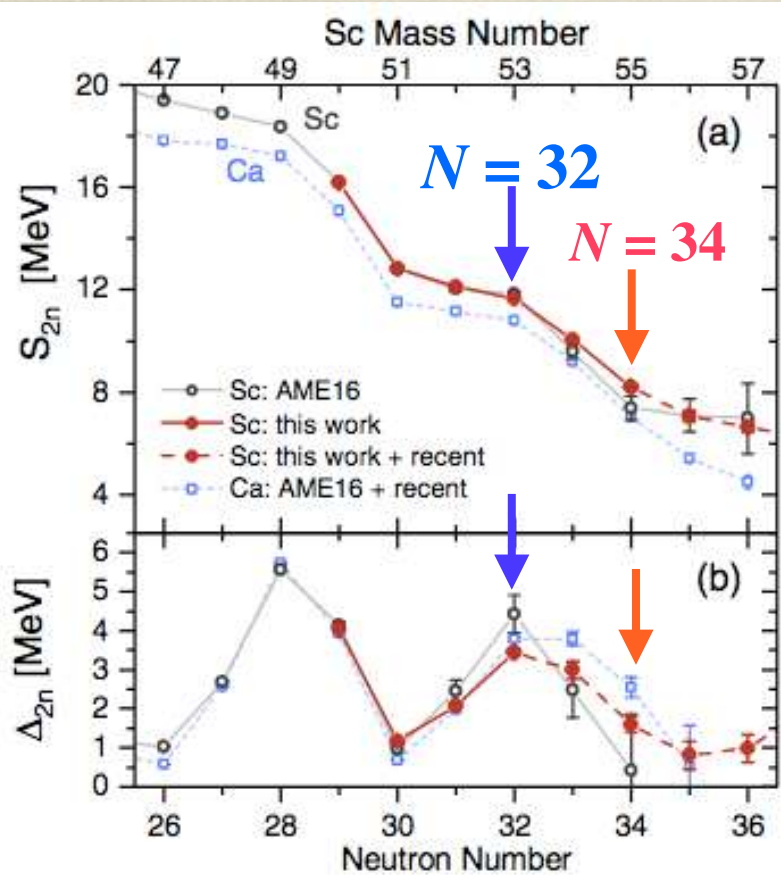
Hint of reduction in the $Z = 6$ sub-shell gap in neutron-rich nuclei

New shells $N = 32, 34$

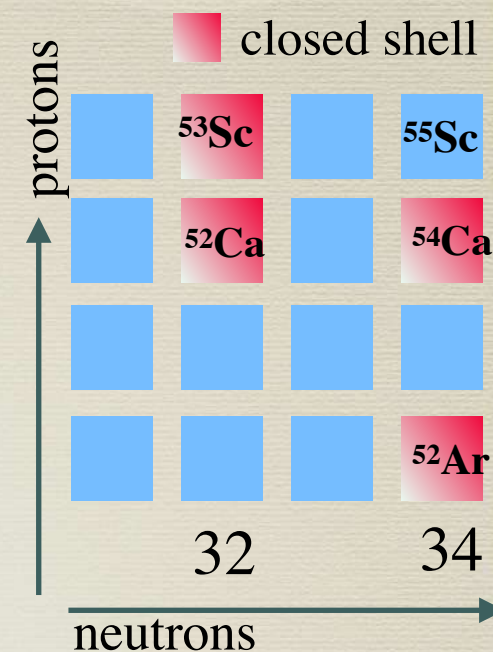
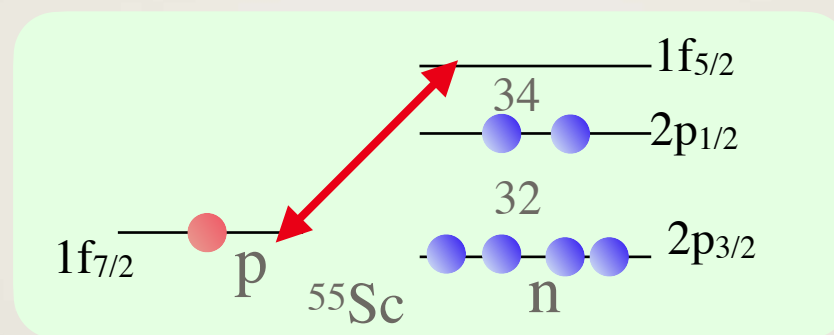
50-53Sc @ NSCL - LEBIT

54-55Sc @ TRIUMF - TITAN

E. Leistenschneider, E. Dunling, G. Bollen *et al.* Phys. Rev. Lett. 126 (2021) 042501



Theory predictions have large variation for $N = 34$

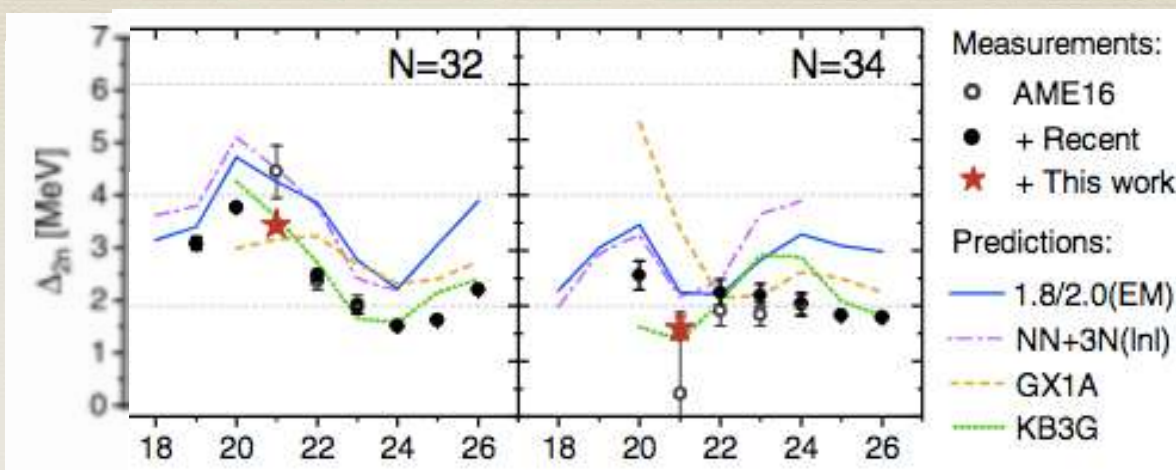
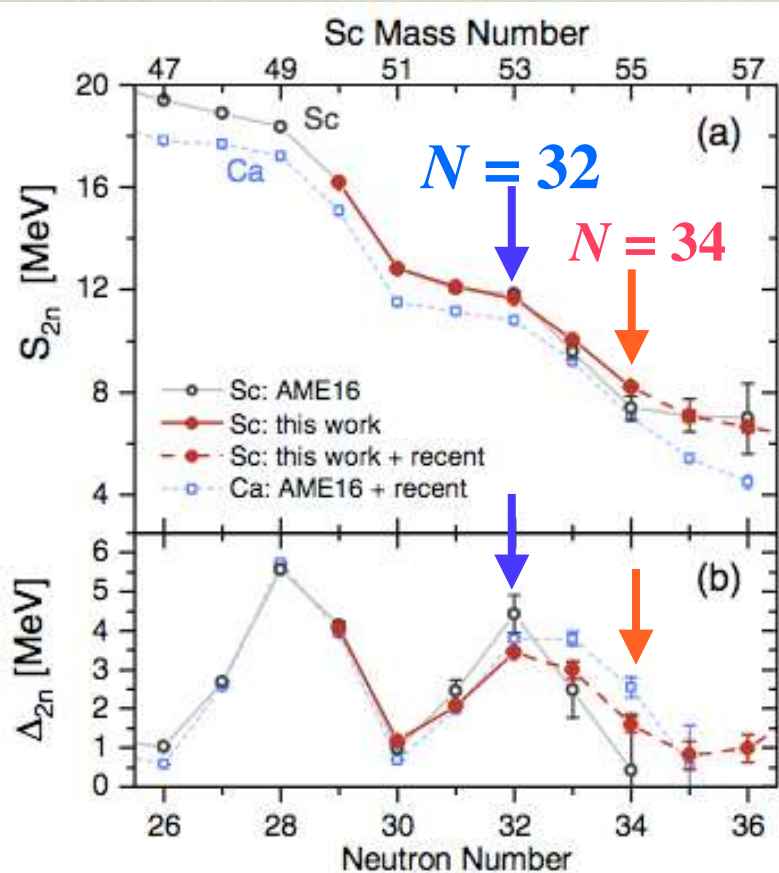


$N = 32$ shell gap seen in ^{53}Sc .
 No signature of $N = 34$ shell gap in ^{55}Sc

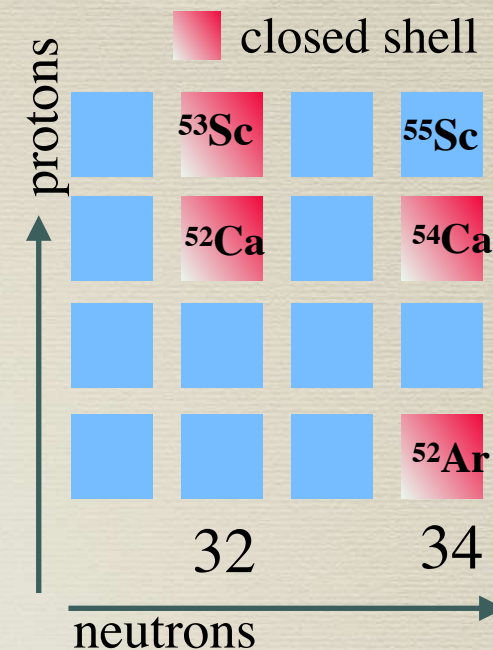
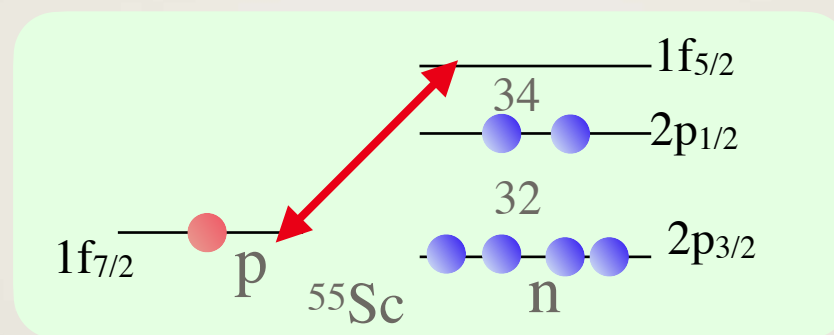
New shells $N = 32, 34$

$50-53\text{Sc}$ @ NSCL - LEBIT $54-55\text{Sc}$ @ TRIUMF - TITAN

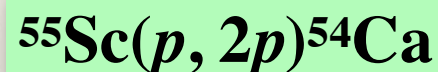
E. Leistenschneider, E. Dunling, G. Bollen *et al.* Phys. Rev. Lett. 126 (2021) 042501



Theory predictions have large variation for $N = 34$

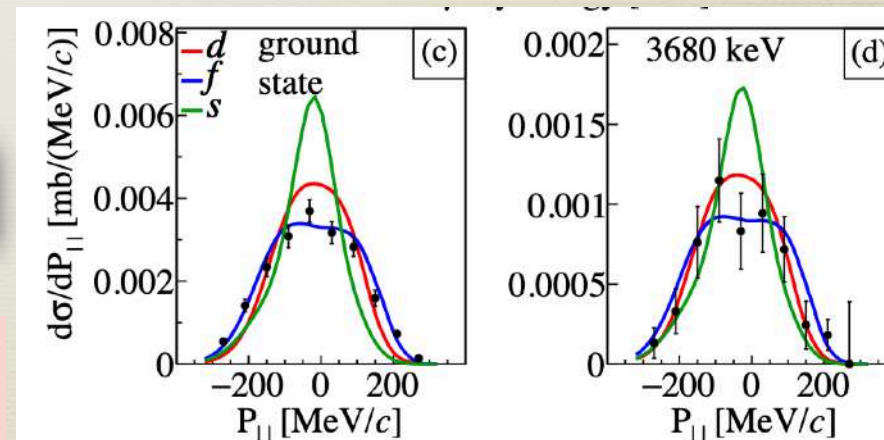


$N = 32$ shell gap seen in ^{53}Sc .
No signature of $N = 34$ shell gap in ^{55}Sc



$1f_{7/2}$ proton removal from ^{55}Sc dominantly populates $^{54}\text{Ca}_{\text{gs}}$. : Pairing force effect

^{54}Ca @ RIBF - BigRIPS



$$C^2S_{\text{g.s.}} = (\alpha_{0p0h}\beta_{0p0h} + \alpha_{2p2h}\beta_{2p2h})^2.$$

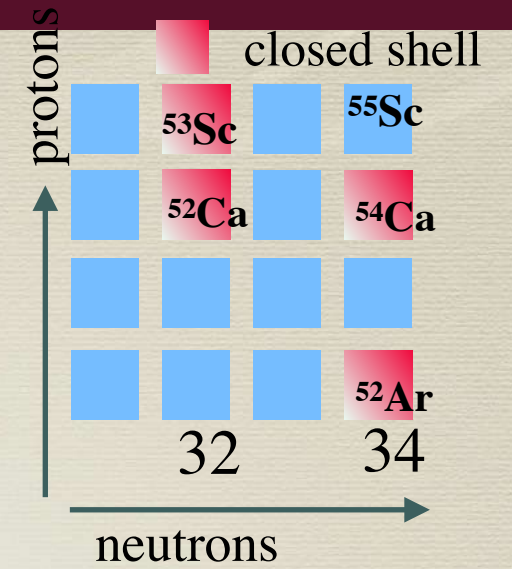
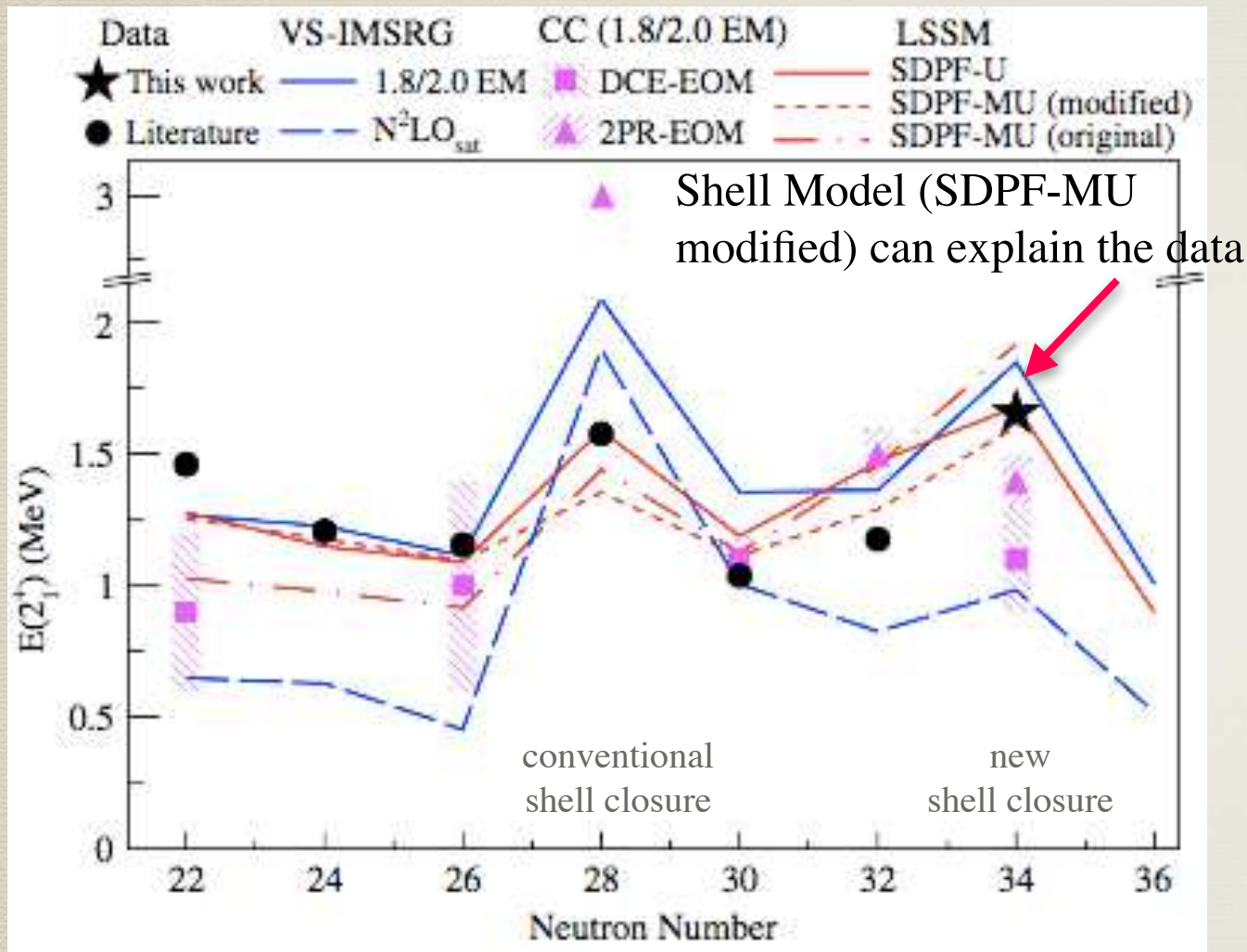
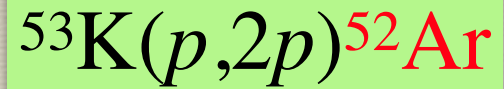
F. Browne, S. Chen, P. Doornenbal *et al.*, Phys. Rev. Lett. 126 (2021) 252501



New shells $N = 32, 34$ ^{52}Ar

@ RIBF - MINOS+SAMURAI

H. Liu, A. Obertelli, P. Doornenbal et al., Phys. Rev. Lett. 122 (2019) 072502.



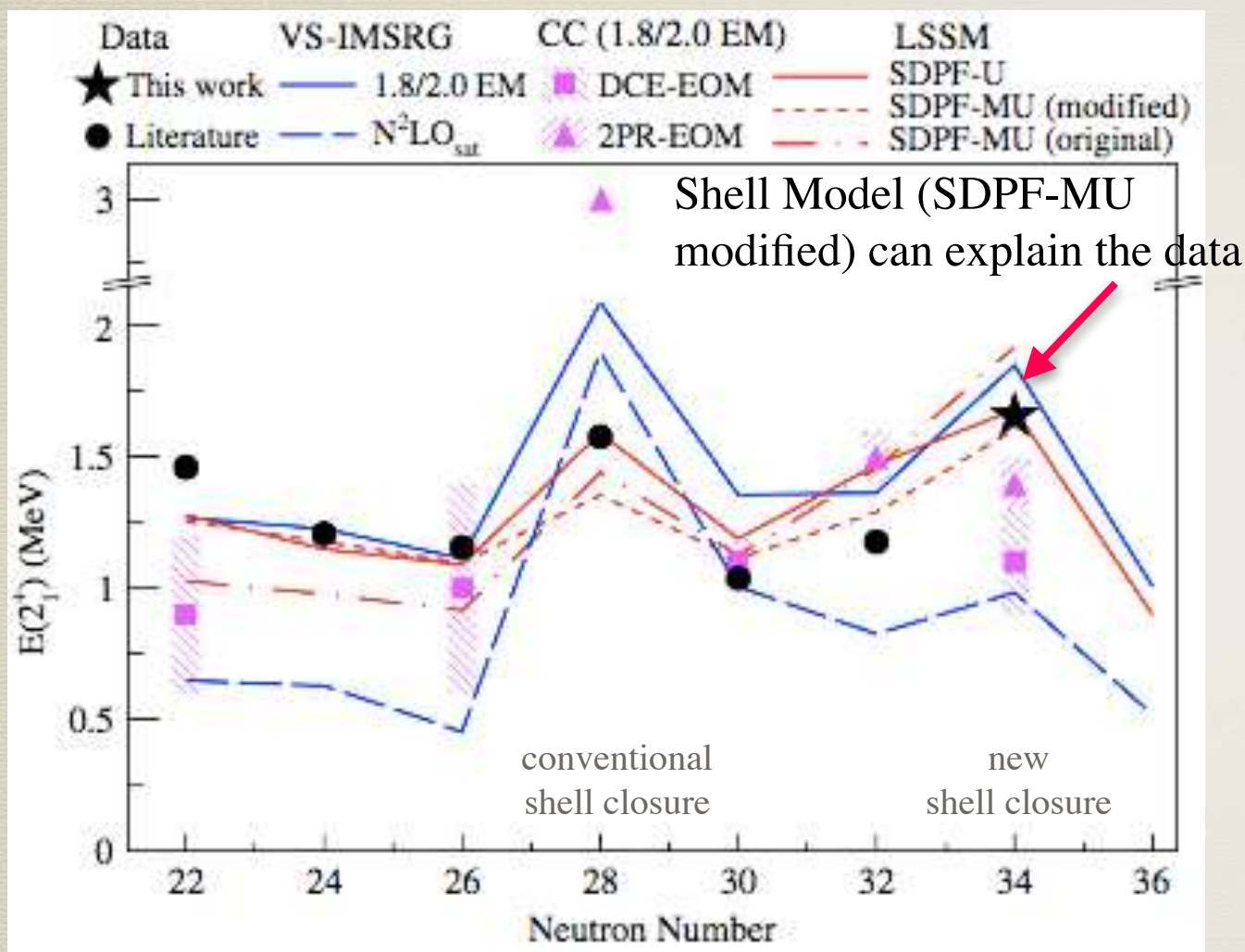
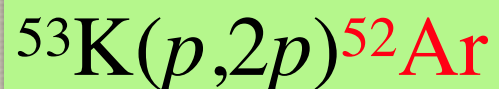
$N = 32$ does not show closed shell behaviour

$N = 34$ shell gap seen in ^{52}Ar from the high excitation energy

New shells $N = 32, 34$ ^{52}Ar

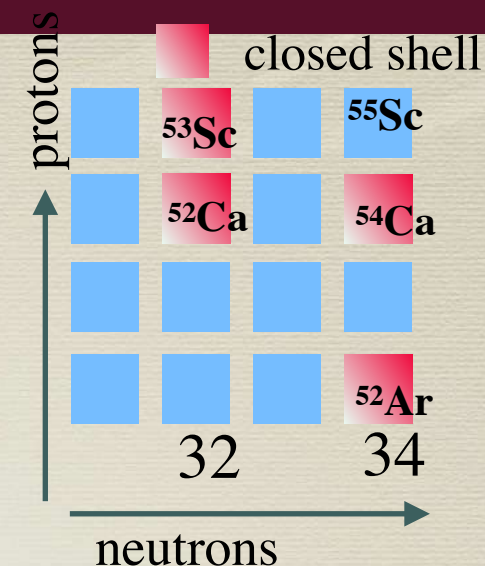
@ RIBF - MINOS+SAMURAI

H. Liu, A. Obertelli, P. Doornenbal et al., Phys. Rev. Lett. 122 (2019) 072502.



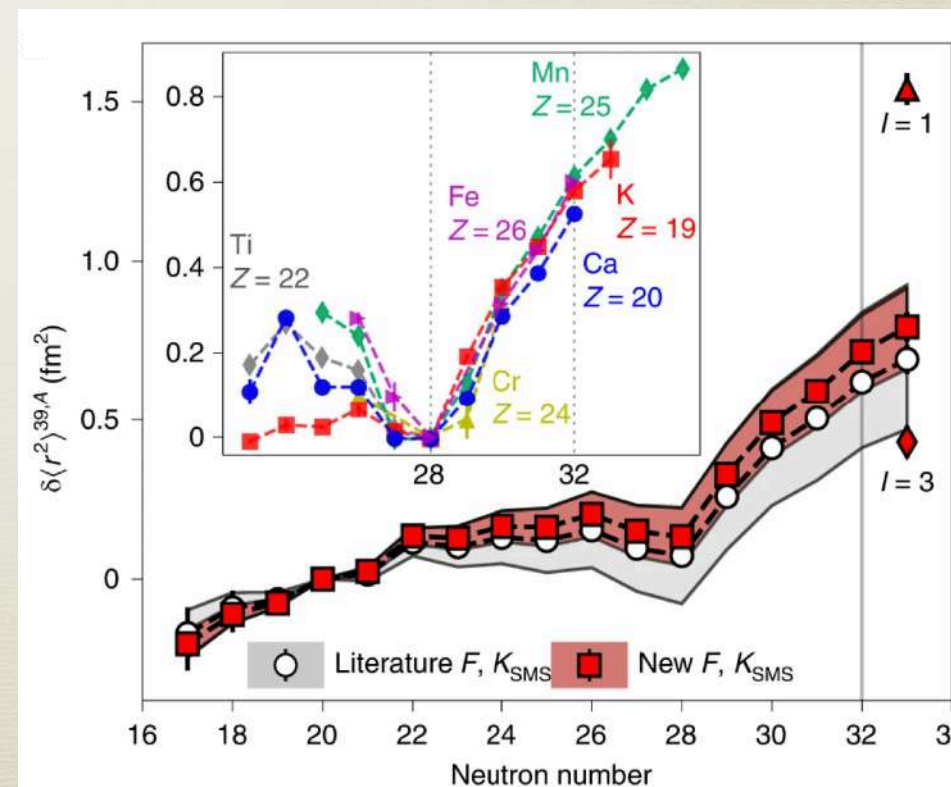
$N = 32$ does not show closed shell behaviour

$N = 34$ shell gap seen in ^{52}Ar from the high excitation energy



Charge Radius @ ISOLDE - CRIS

A. Koszorus, et al., Nature 17 (2021) 439



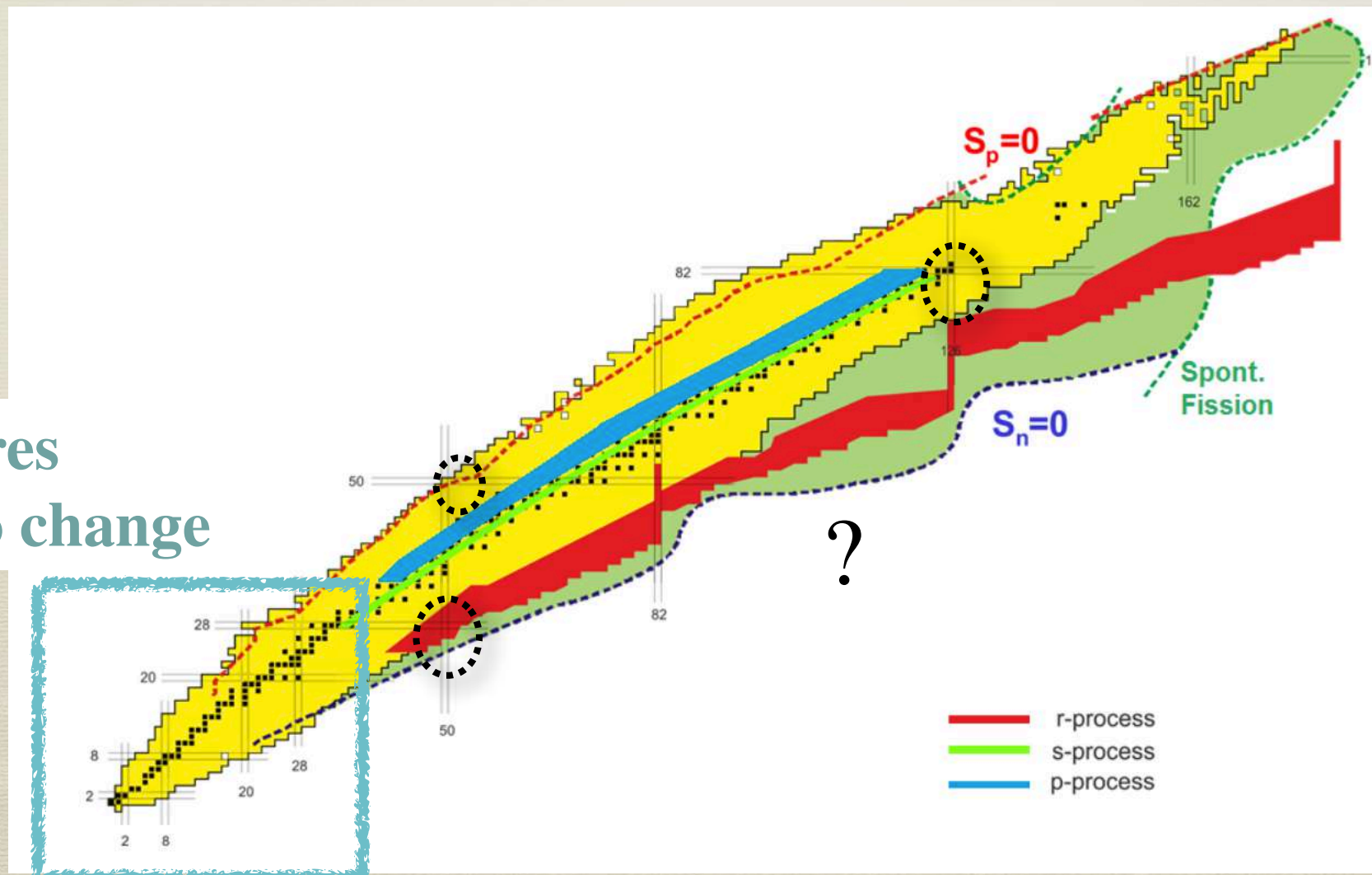
Charge radius does not show dip @ $N = 32$



Nuclear structure and shell evolution impacts heavy element synthesis

$N = 50$ & 126 conventional shells

Shell closures
observed to change



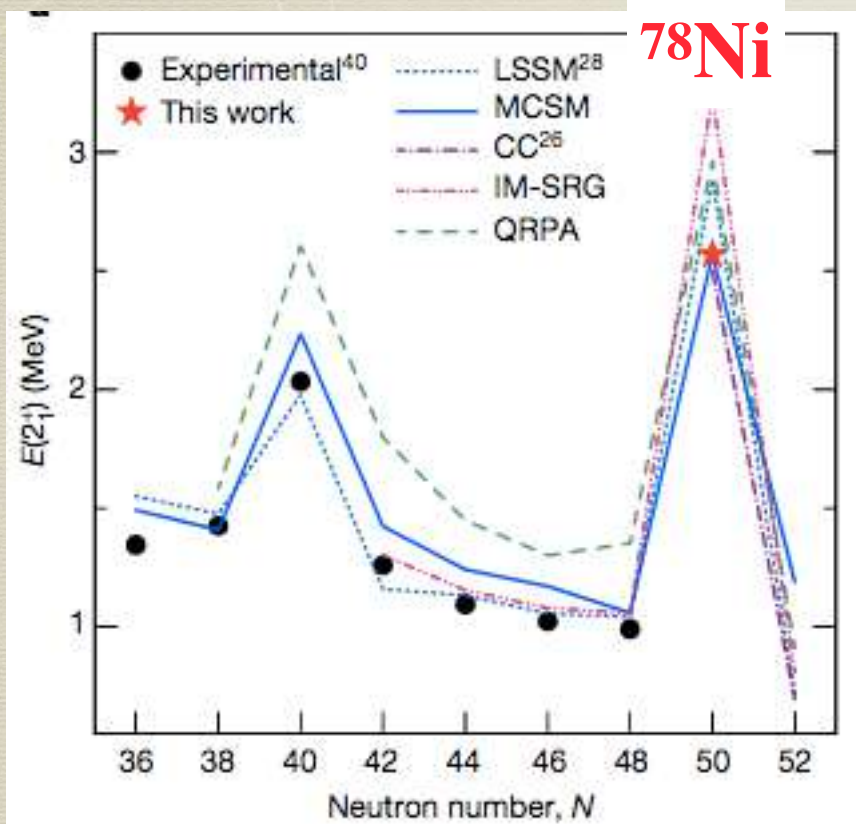
Doubly magic ^{78}Ni ($N = 50$)

@ RIKEN - BigRIPS

$^{80}\text{Zn}(p, 3p)^{78}\text{Ni}$

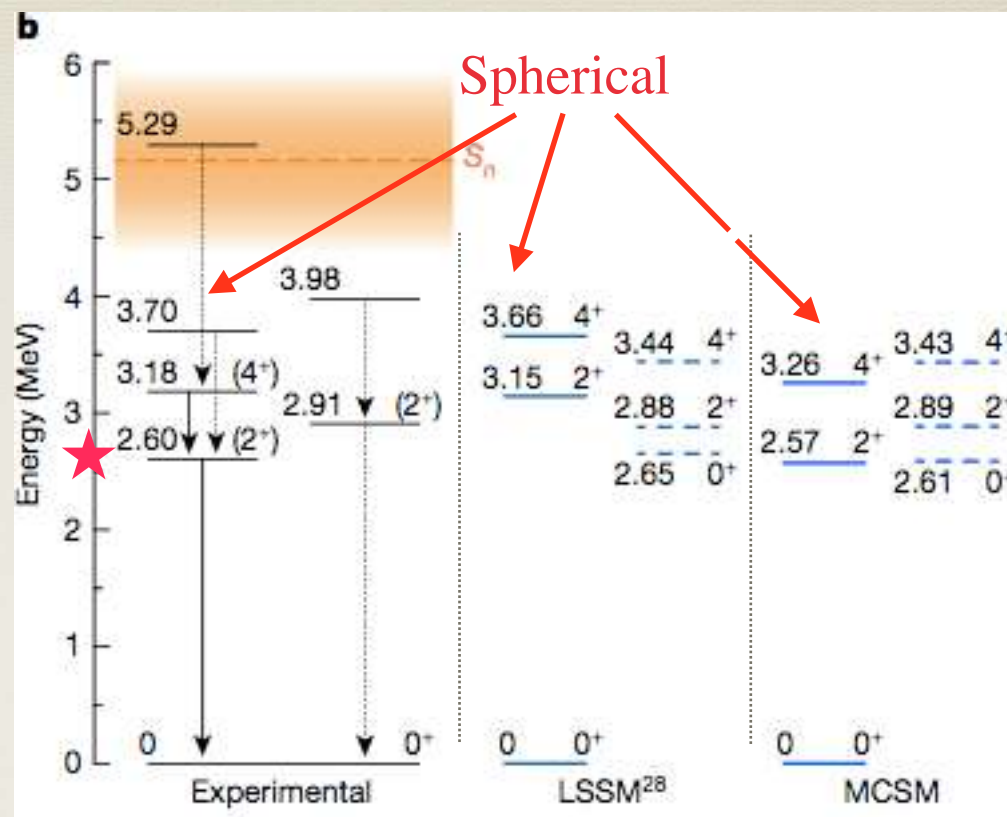
R. Taniuchi, C. Santamaria, P. Doornenbal *et al.* Nature 569 (2019) 53

$^{79}\text{Cu}(p, 2p)^{78}\text{Ni}$

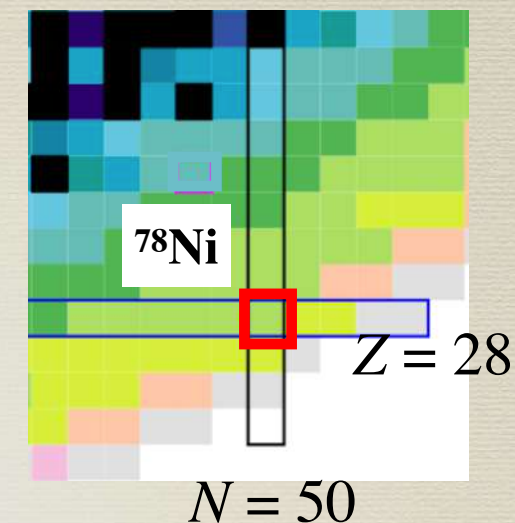


High excitation energy of first 2^+ state : $N = 50$ closed shell

Calculations with three nucleon force explain data



Second 2^+ state located close to first 2^+ state : Close lying spherical and deformed configurations



^{78}Ni is doubly magic - $N = 50$ shell closure persists

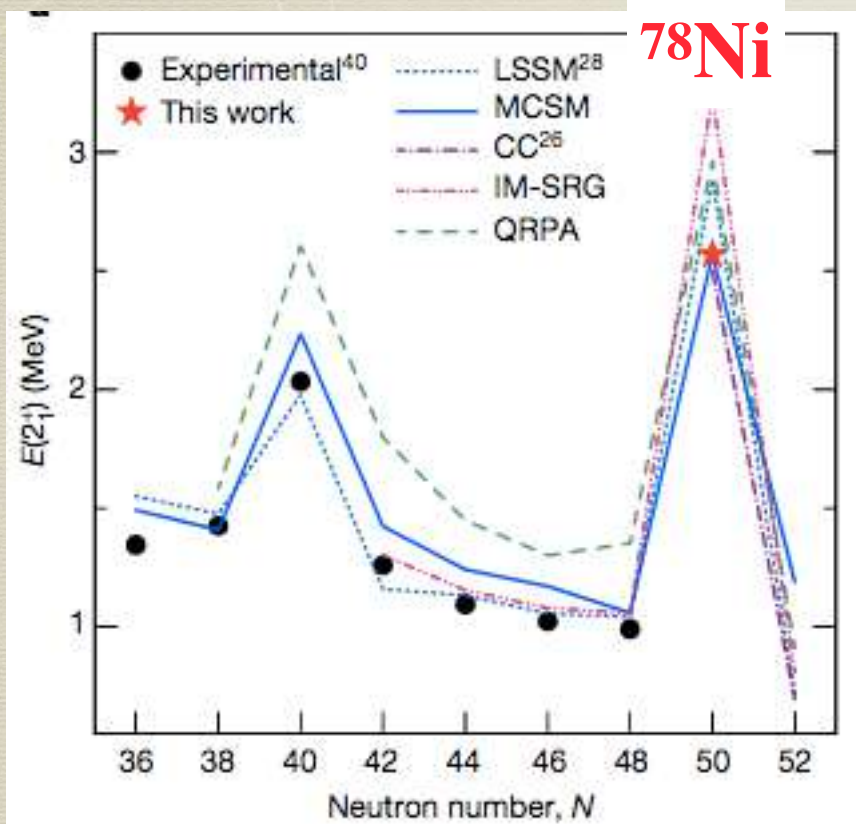
Doubly magic ^{78}Ni ($N = 50$)

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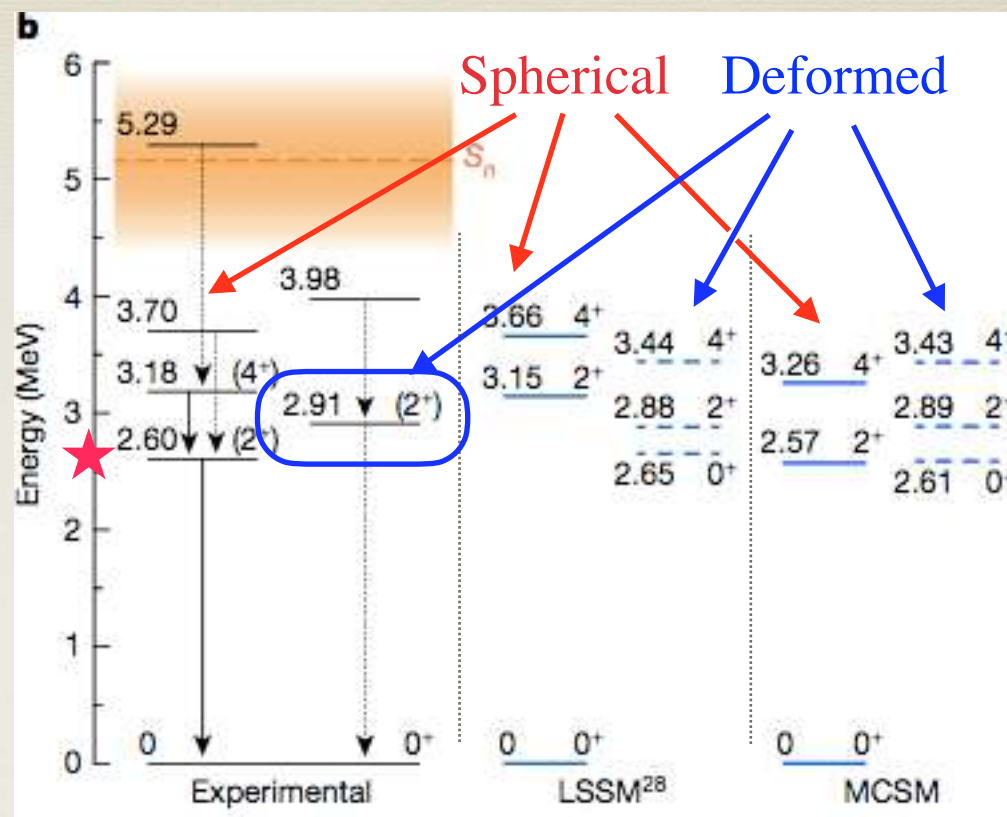
R. Taniuchi, C. Santamaria, P. Doornenbal *et al.* Nature 569 (2019) 53

$^{79}\text{Cu}(p, 2p)^{78}\text{Ni}$

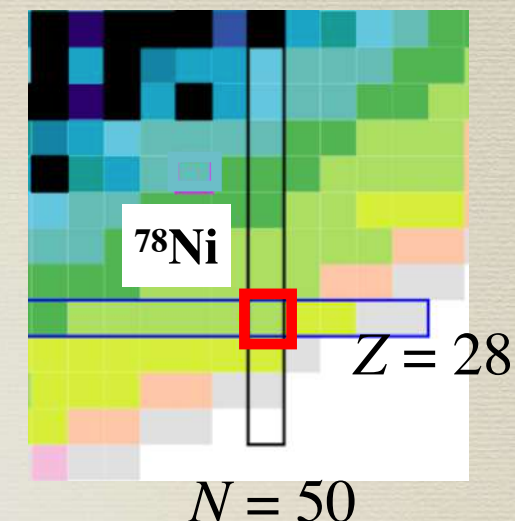


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$N = 50$ shell gap quenching beyond ^{78}Ni hinted - competing deformed & spherical shapes

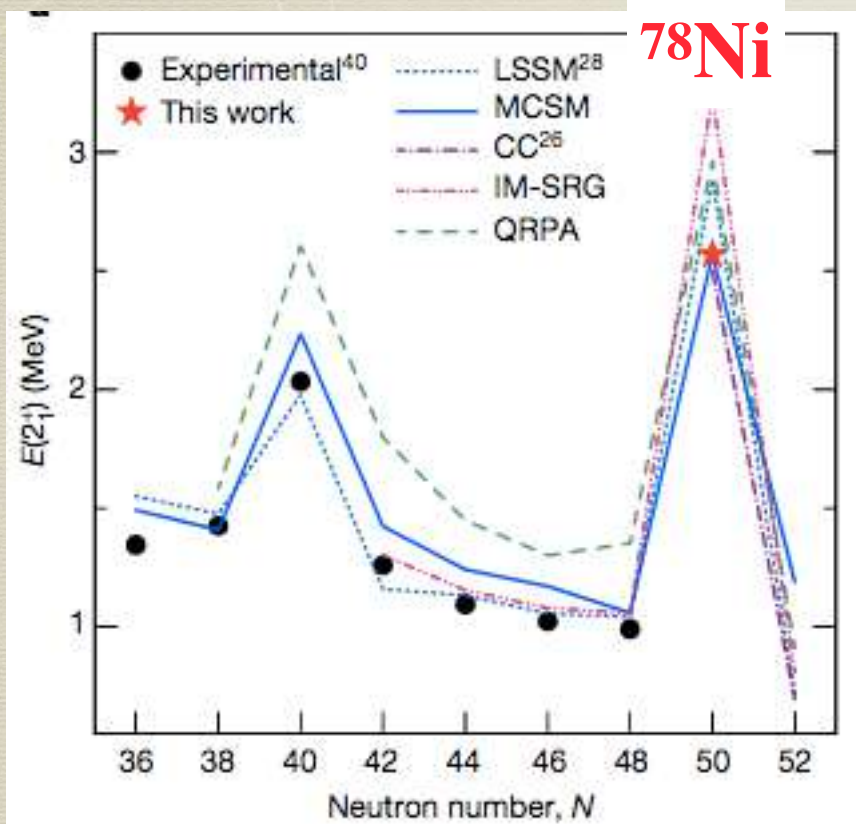
Doubly magic ^{78}Ni ($N = 50$)

@ RIKEN - BigRIPS

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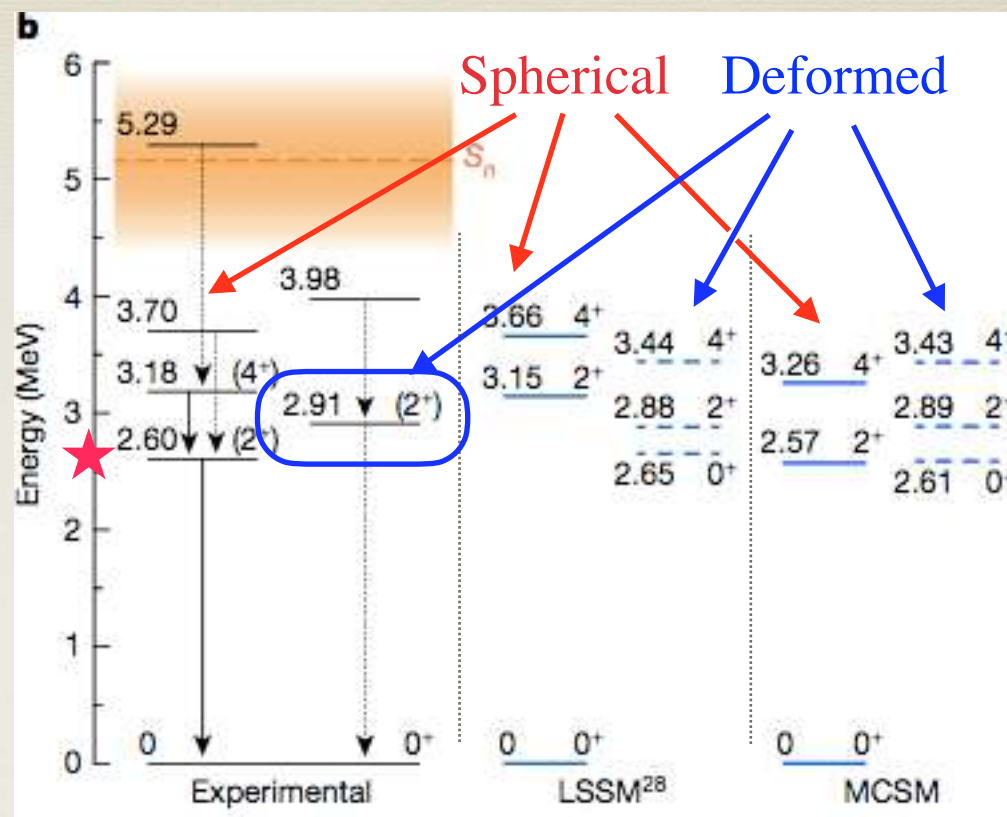
R. Taniuchi, C. Santamaria, P. Doornenbal *et al.* Nature 569 (2019) 53

$^{79}\text{Cu}(p, 2p)^{78}\text{Ni}$

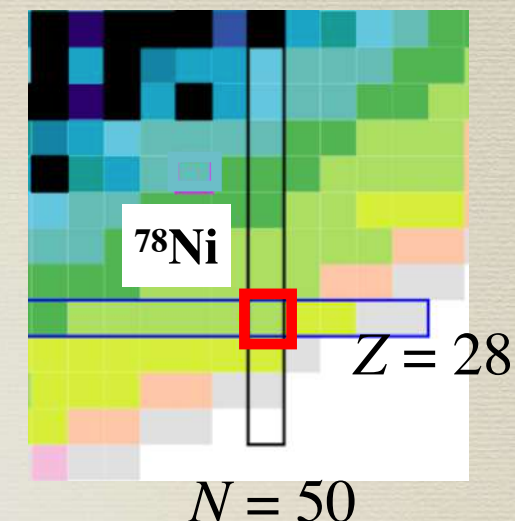


High excitation energy of first 2^+ state : $N = 50$ closed shell

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Second 2^+ state located close to first 2^+ state : Close lying spherical and deformed configurations



^{78}Ni is doubly magic - $N = 50$ shell closure persists

$N = 50$ shell gap quenching beyond ^{78}Ni hinted - competing deformed & spherical shapes

Future studies on ^{76}Fe and ^{74}Cr needed to search for erosion for $N = 50$ shell gap

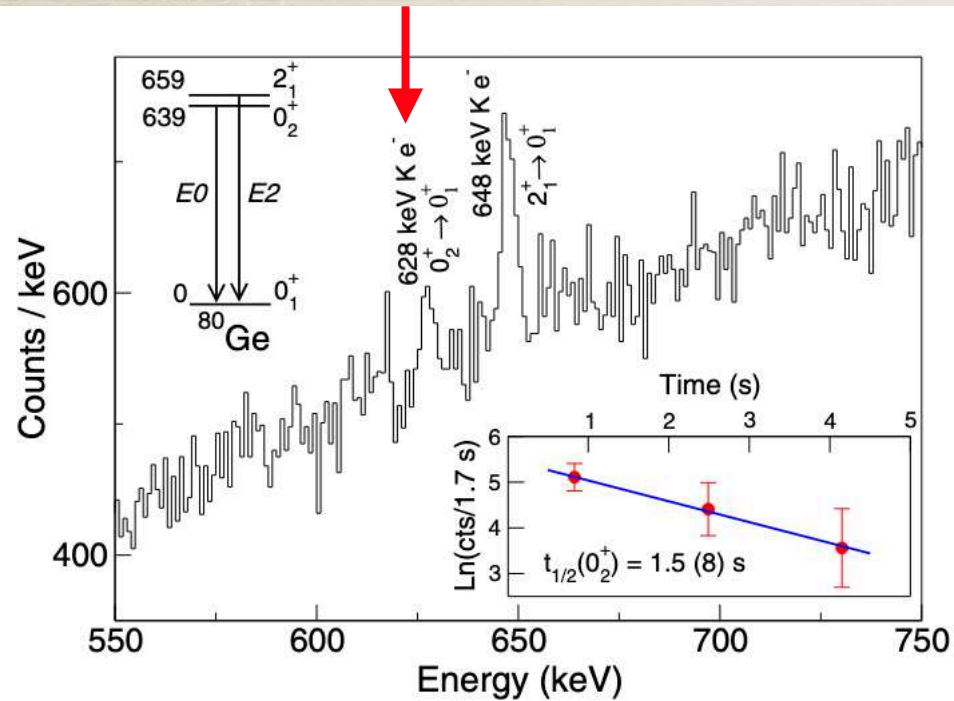


^{80}Ge ($N = 48$) : Shape coexistence or not ? controversy resolved

^{80}Ga β^- decay

@ ALTO

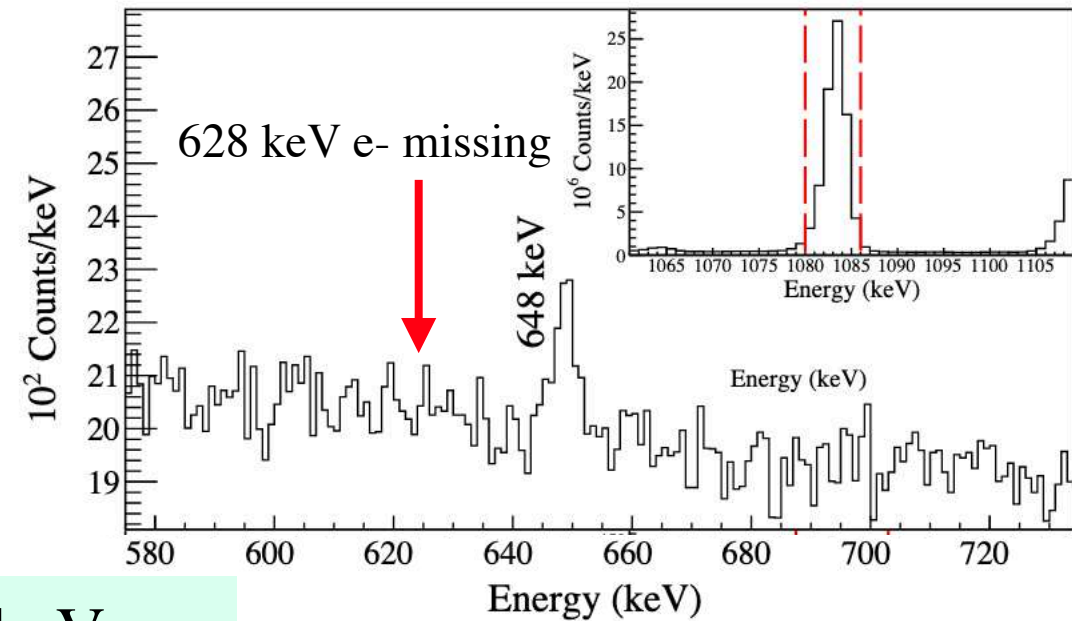
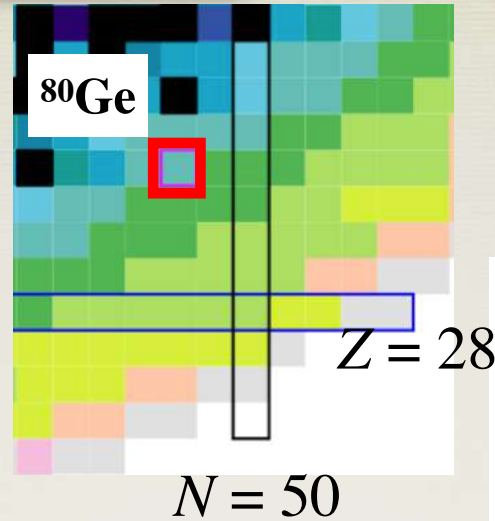
A. Gottardo et al., Phys. Rev. Lett. 116 (2016) 182501



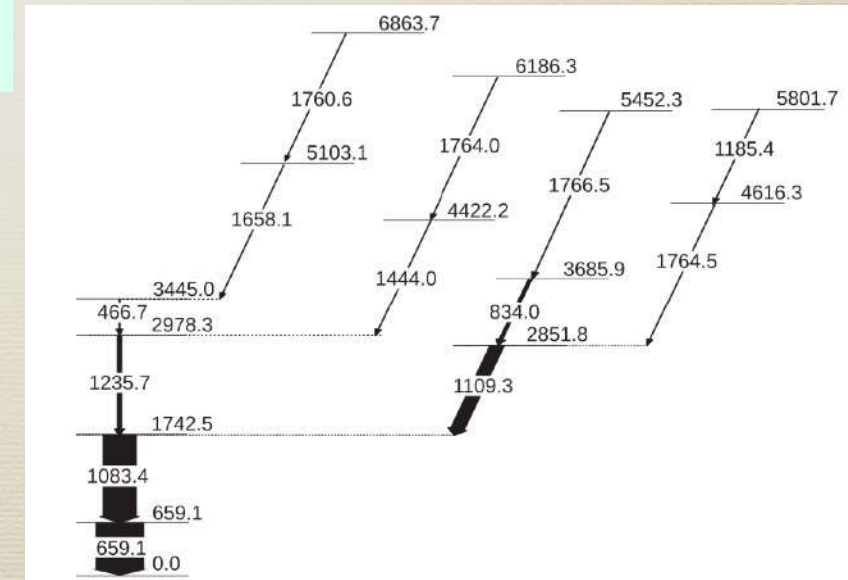
0_2^+ state @ 639 keV below 2^+ state reported as evidence of shape coexistence

@ TRIUMF - GRIFFIN

F. Garcia et al., Phys. Rev. Lett. 125 (2020) 172501



0_2^+ : 638 keV state not found !



No shape coexistence in ^{80}Ge

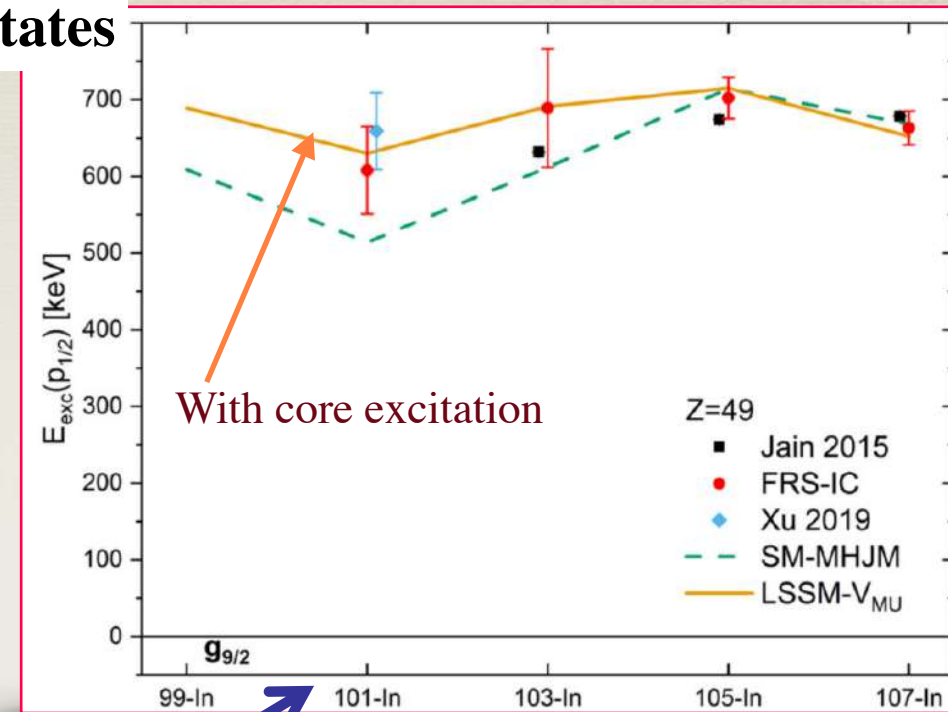
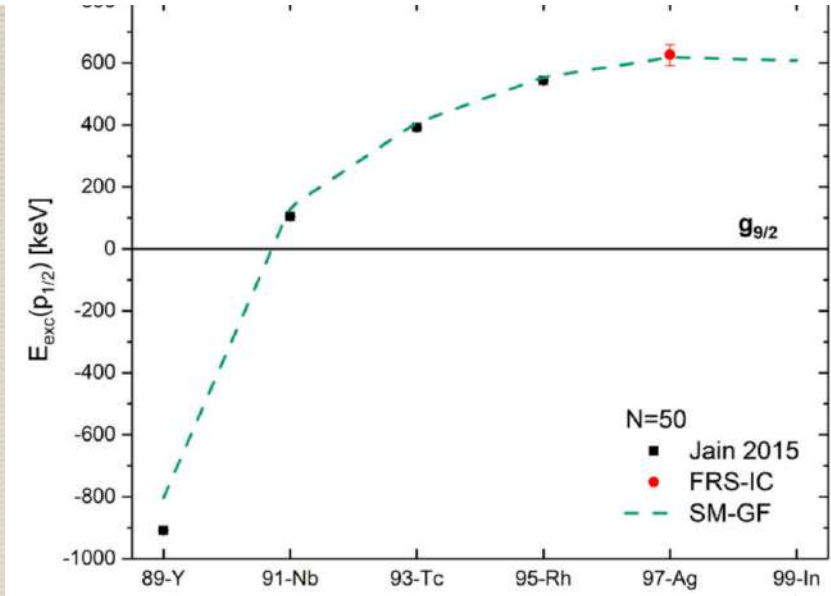
Isomer states found around ^{100}Sn

@ GSI - FRS Ion Catcher & MR-TOF

$N = Z = 50$ region

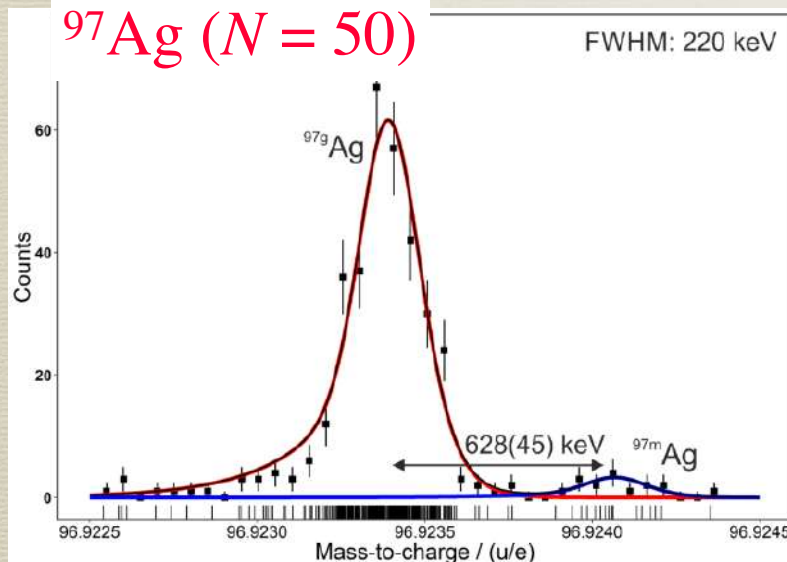
C. Homung *et al.* Phys. Lett. B 802 (2020) 135200

High precision mass measurements identify isomeric states

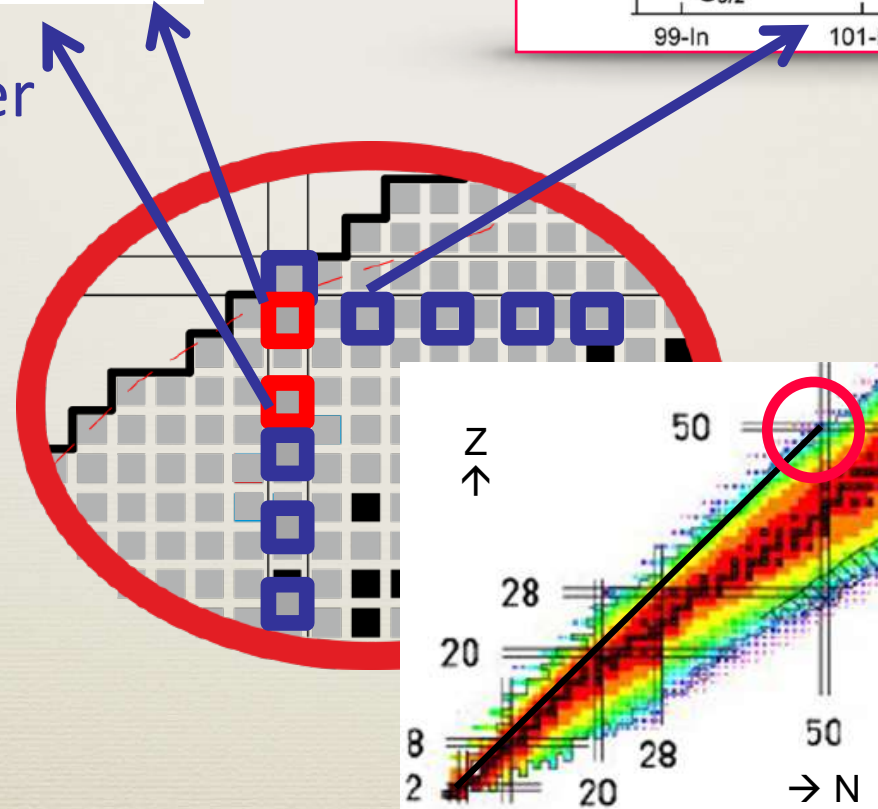


^{97}Ag , new (1/2-) isomer

^{97}Ag ($N = 50$)

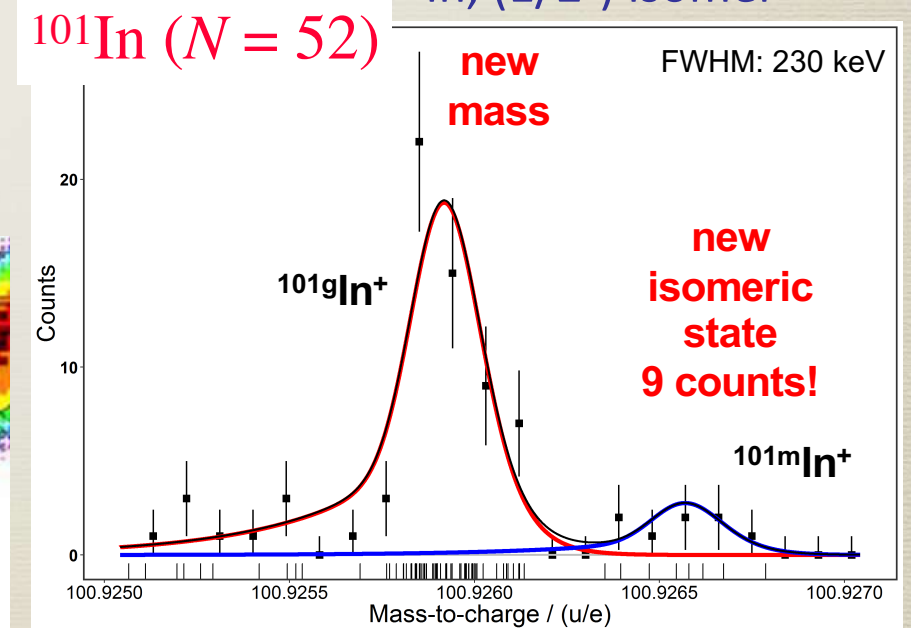


Courtesy : T. Dickel



^{101}In , (1/2-) isomer

^{101}In ($N = 52$)



Data shows the need for core excitation across $N = 50$: Large Scale Shell model calculations

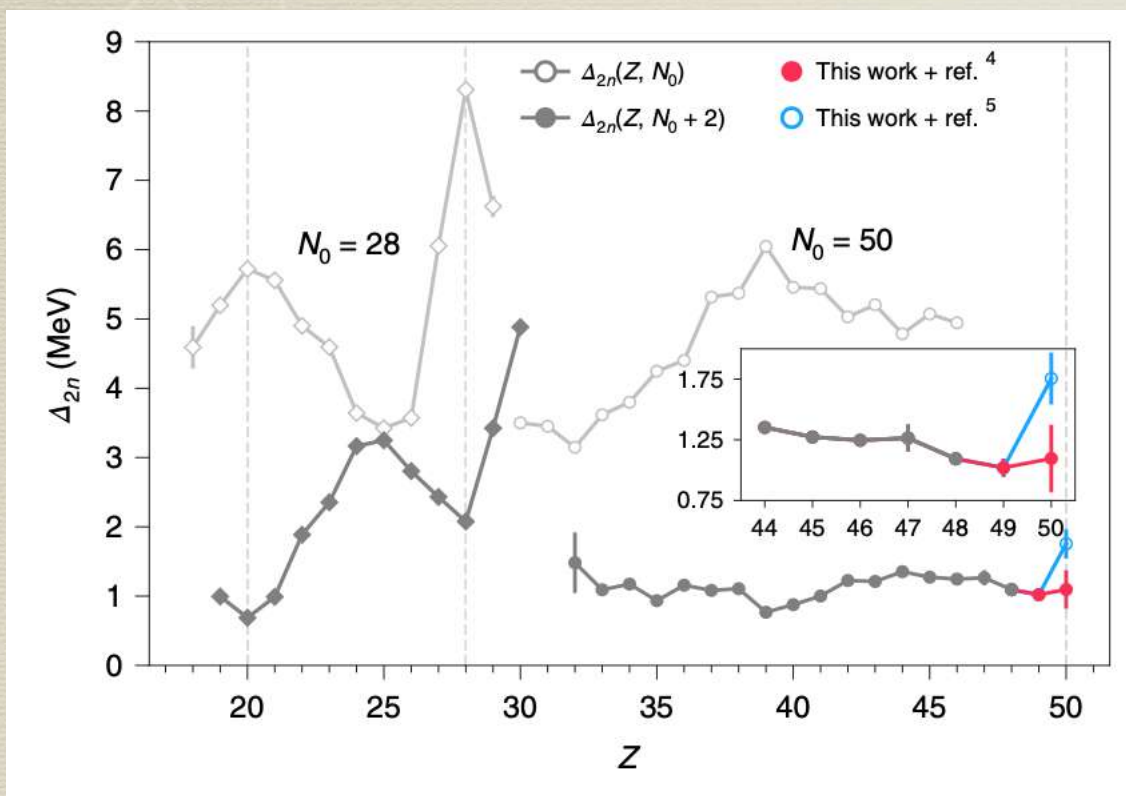


Mass $^{99-101}\text{In}$ towards ^{100}Sn

@ ISOLDE - ISOLTRAP & MR-TOF

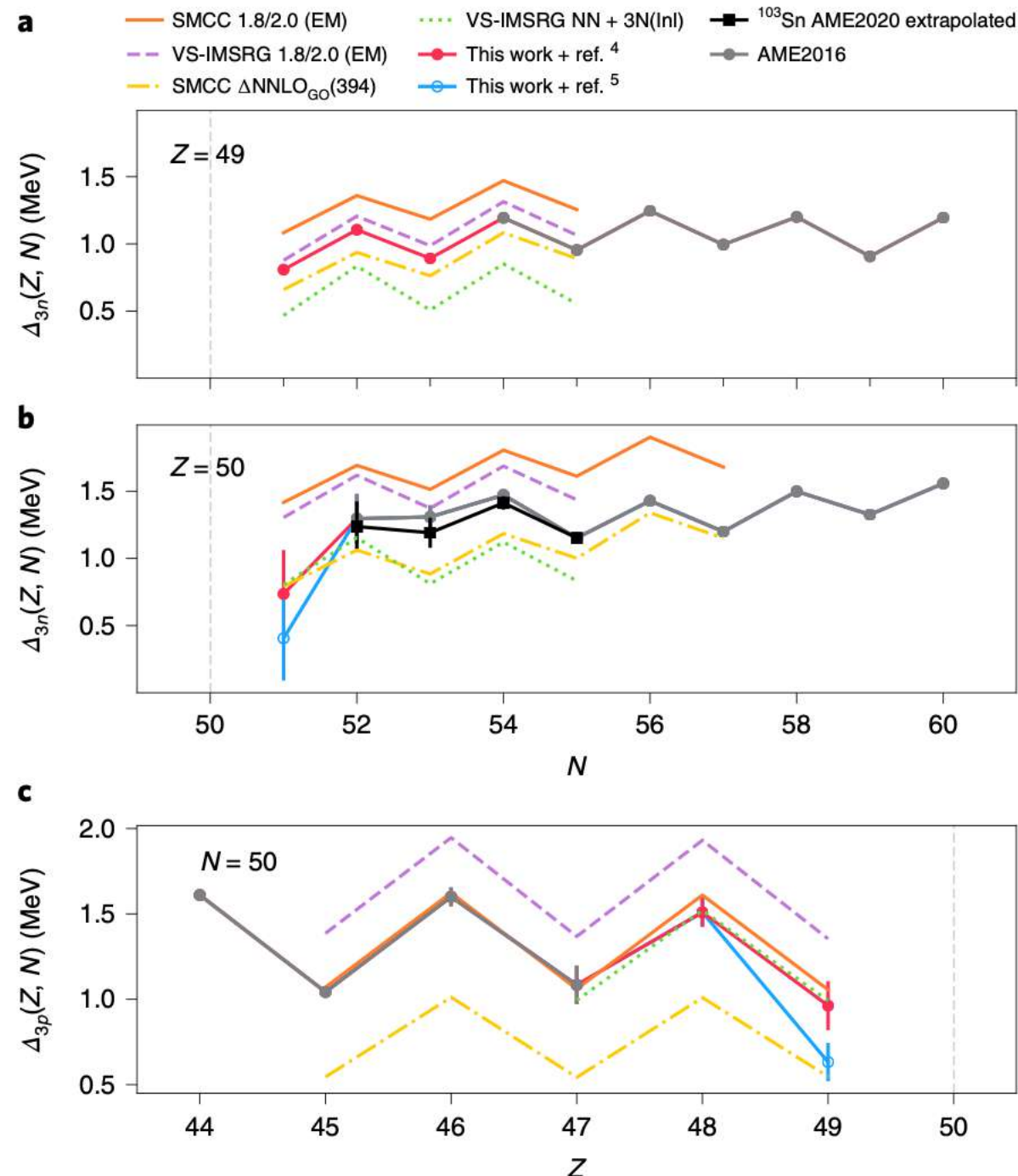
$N = Z = 50$ region

M. Mougeot et al., Nat. Phys. 17 (2021) 1099



$$\Delta_{2n}(Z, N_0) = M_E(Z, N_0 - 2) - 2M_E(Z, N_0) + M_E(Z, N_0 + 2)$$

High precision mass measurements show trend of shell closure towards $N = Z = 50$

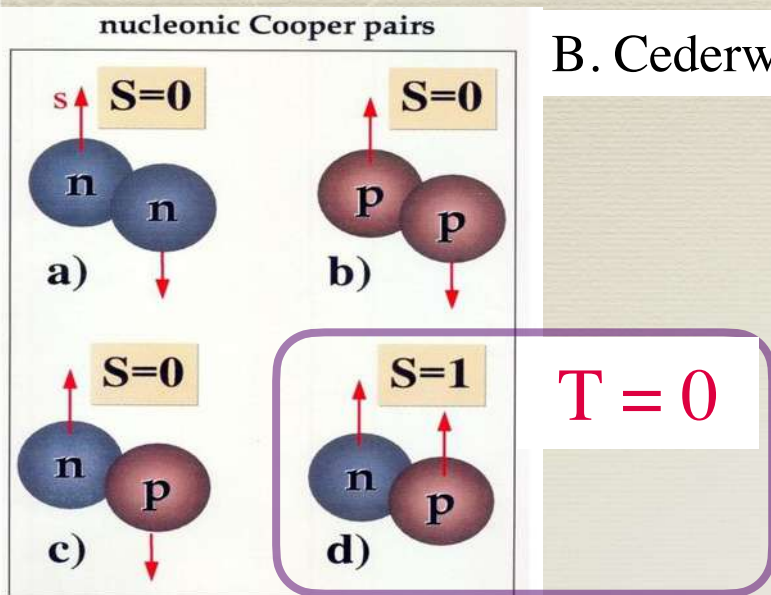


Ab initio predictions of odd-even staggering overall align well with data.

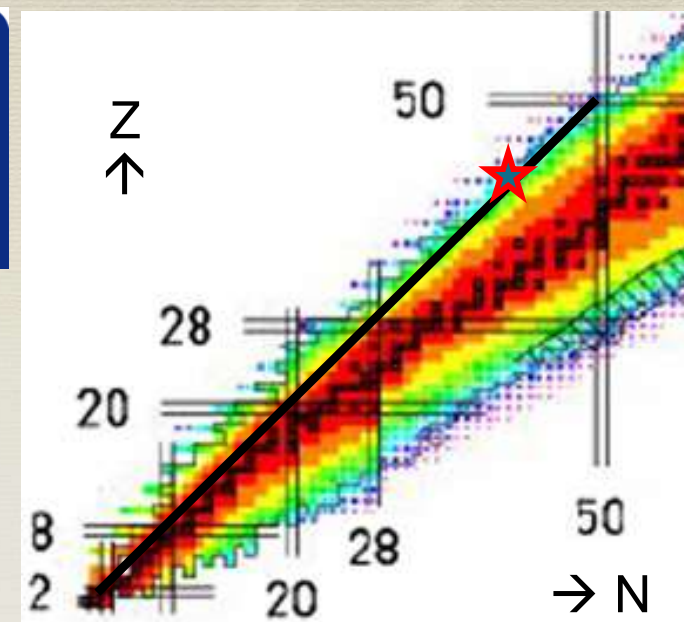
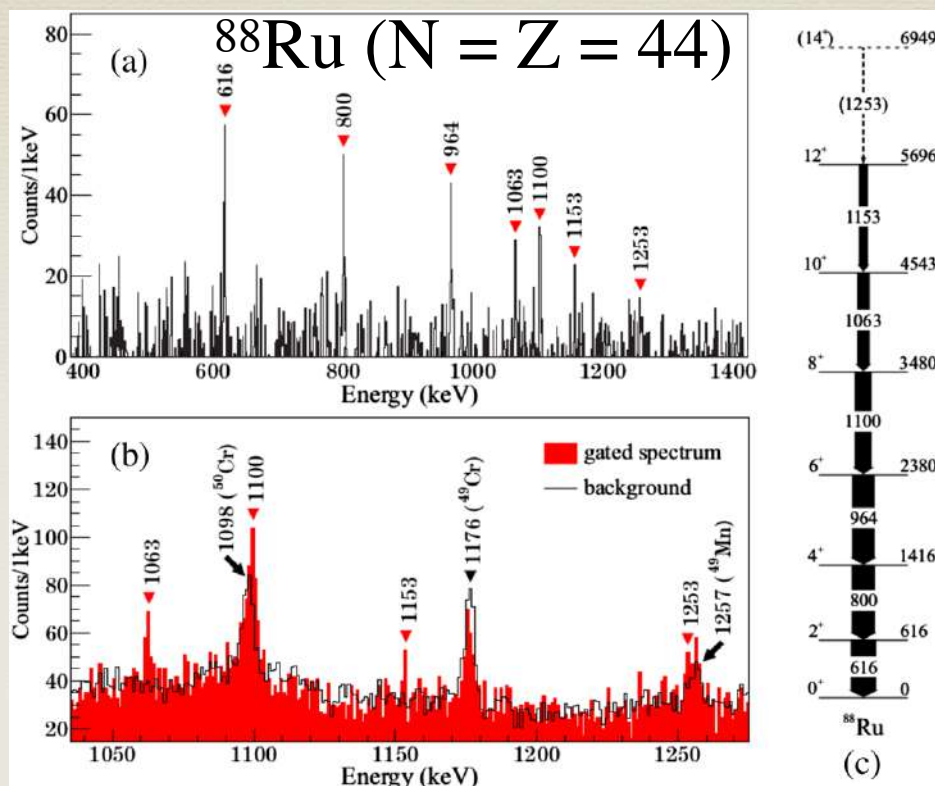


$N = Z$: Neutron-Proton (p - n) Isoscalar Pairing

@ GANIL - AGATA



B. Cederwall *et al.* Phys. Rev. Lett. 124 (2020) 062501

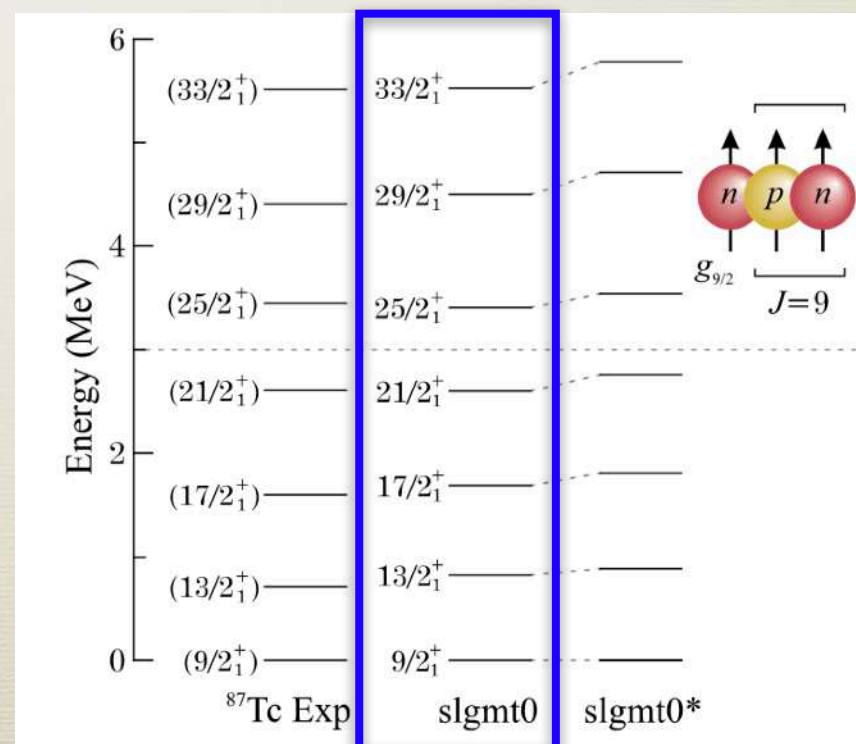


^{87}Tc ($N = 44, Z = 43$)

X. Liu *et al.* Phys. Rev. C Lett. 104 (2021) L021302

$^{56}\text{Fe}(^{36}\text{Fe}, 2n1p)^{87}\text{Tc}$

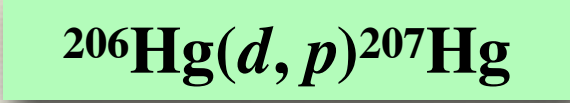
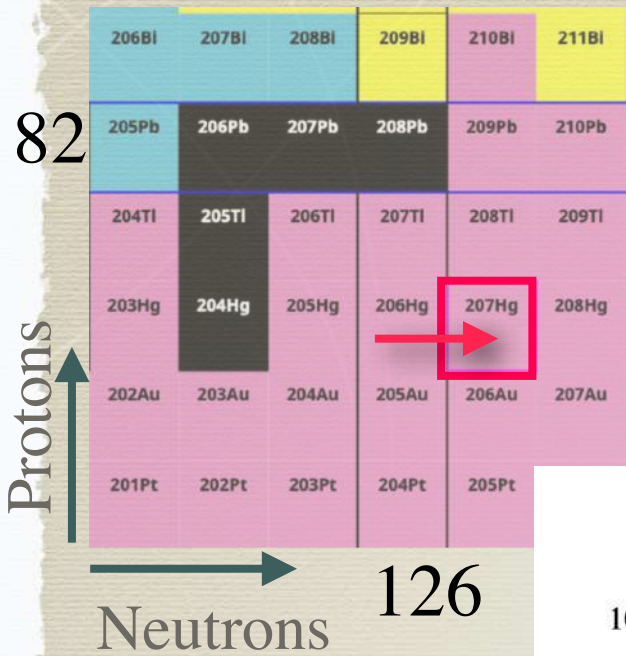
Observed rotation spectrum supports predictions with strong n - p pairing ($T=0$) correlation



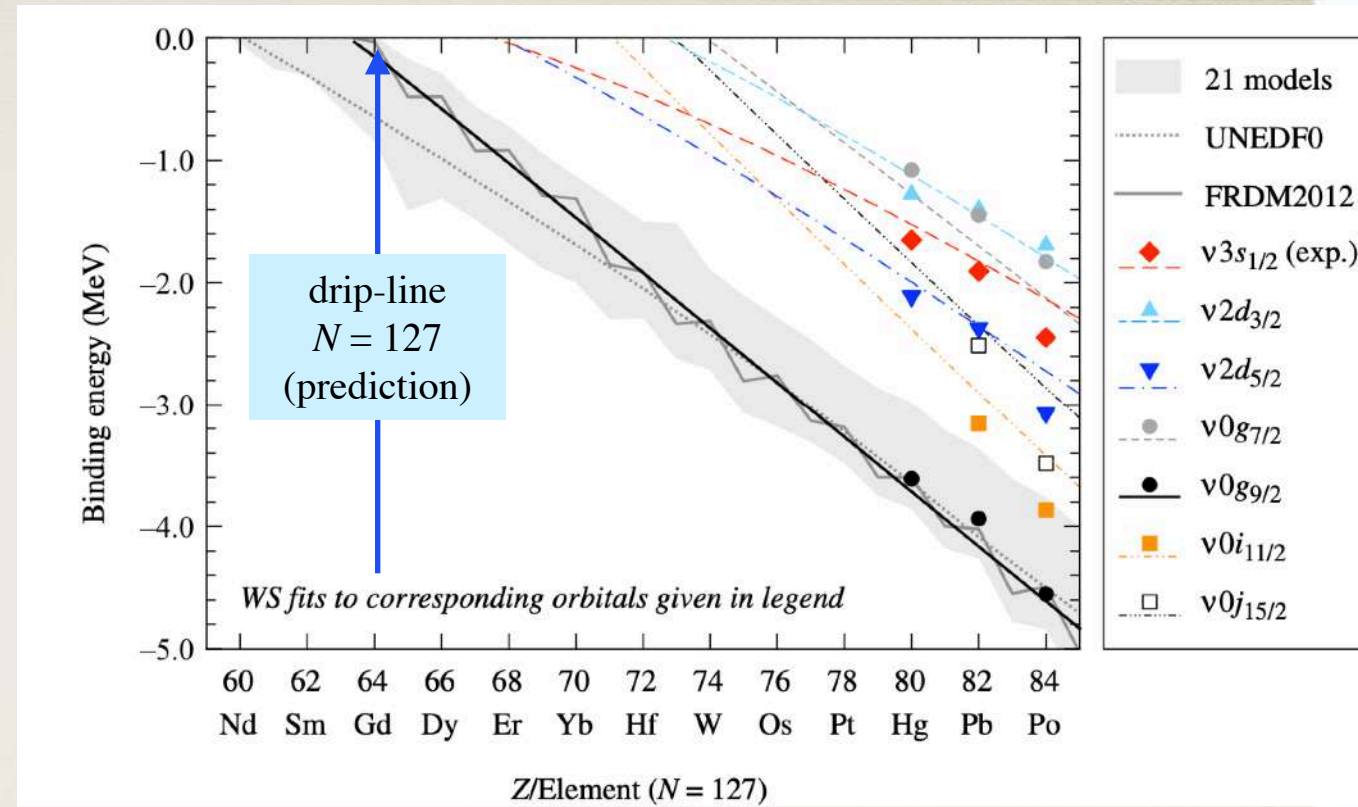
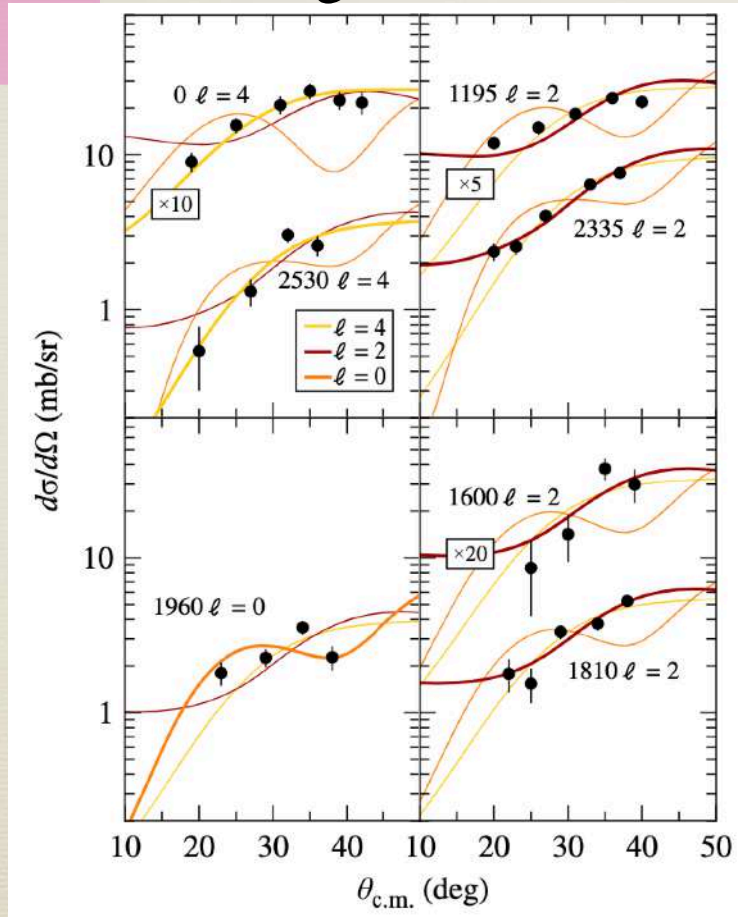
Shell structure beyond $N = 126$

@ ISOLDE - Solenoidal Spectrometer

T.L. Tang, B.P. Kay, C.R. Hofmann *et al.* Phys. Rev. Lett. 124 (2020) 062502



^{207}Hg excited states



^{207}Hg states show normal shell configuration. Binding energy aligns with predictions.

$Z = 64$ predicted to be the drip line of $N = 127$.

r-process neutron capture improbable with low angular momentum orbitals being unbound

Assumption : $N = 126$ shell closure persists. Future experiments will inform on this.



Unbound Nuclei - open quantum systems

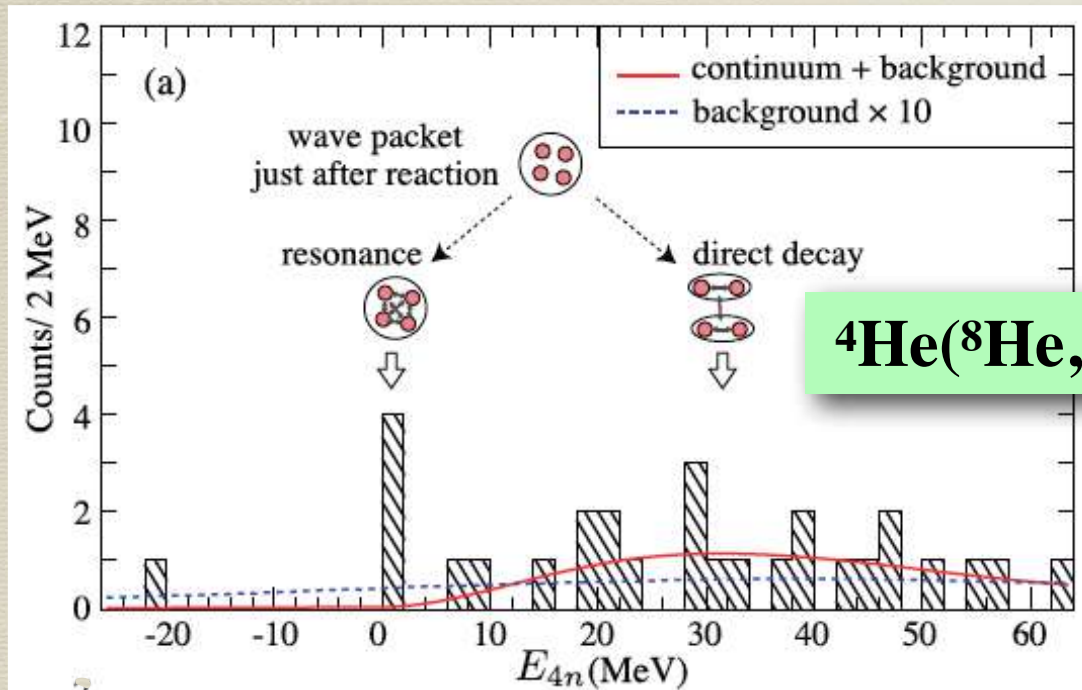
The hunt for tetraneutron

@ RIKEN - SHARAQ

K. Kisamori, S. Shimoura, H. Miya *et al.*, Phys. Rev. Lett. 116 (2016) 052501



The first result of 4n from SHARAQ



Resonance reported ~ 1 MeV

Theoretical Predictions

- Hiyama *et al.* No
(too strong 3N force is needed)
- Shirokov *et al.* Yes
(NSCM with JISP16 interaction)
- Gandolfi *et al.* Yes
(QMC with chiral interaction)

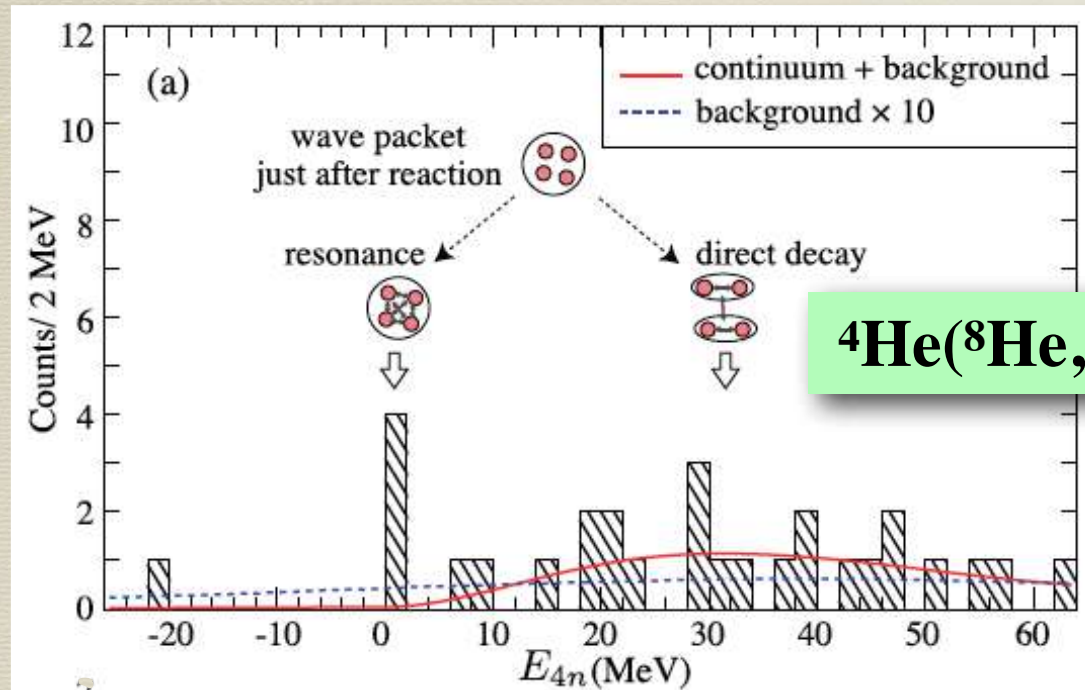
The hunt for tetra neutron

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The question is still open on the existence of narrow resonance for tetra neutron



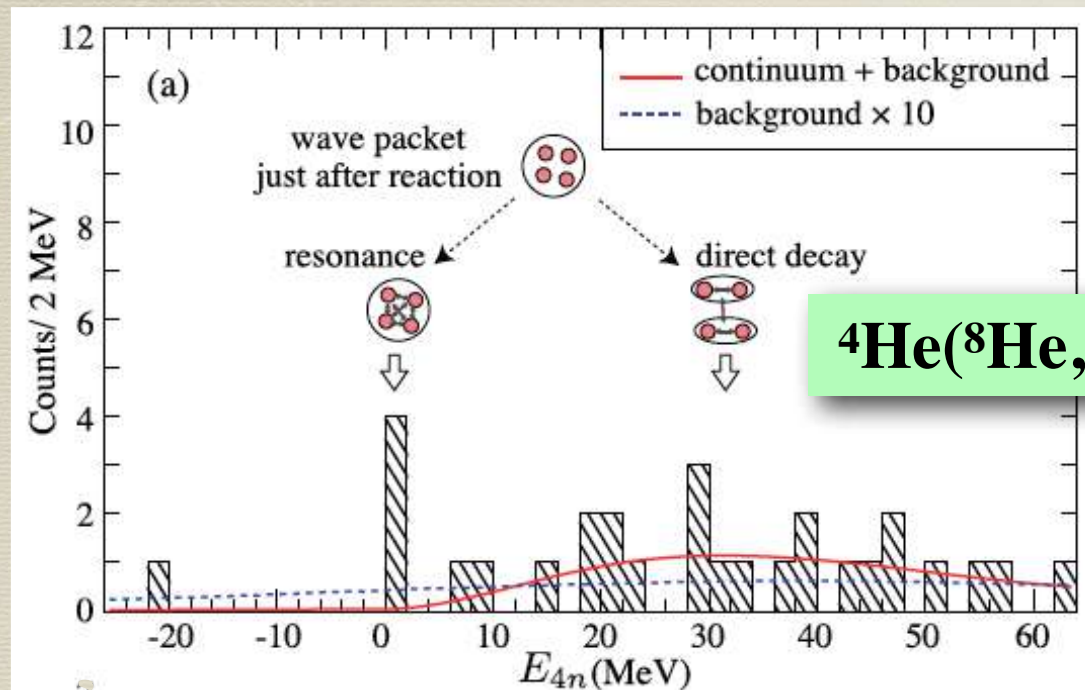
The hunt for tetra-neutron

@ RIKEN - SHARAQ

K. Kisamori, S. Shimoura, H. Miya *et al.*, Phys. Rev. Lett. 116 (2016) 052501



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New Experiments to study exotic system
@ RIKEN – RIBF

SHARAQ

- Shimoura *et al.*
Revisit $4\text{He}({}^8\text{He}, {}^8\text{Be})4n$
- Miki *et al.*
 ${}^3\text{H}({}^3\text{H}, {}^3\text{He})3n$

SAMURAI

- Rossi *et al.*
 ${}^8\text{He}(p, p\alpha)4n$
- Yang and Marques *et al.*
 ${}^8\text{He}(p, 2p){}^7\text{H} \rightarrow t + 4n$
- Beaumel *et al.*
 ${}^{14}\text{Be}(p, p\alpha)6n + \alpha$

Courtesy : H. Sakurai

The question is still open on the existence of narrow resonance for tetra-neutron

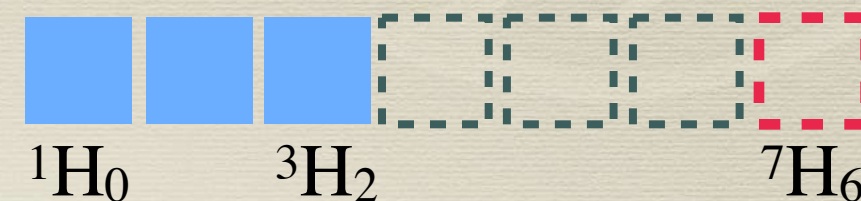


Superheavy hydrogen ${}^7\text{H}$

@ GANIL - MAYA

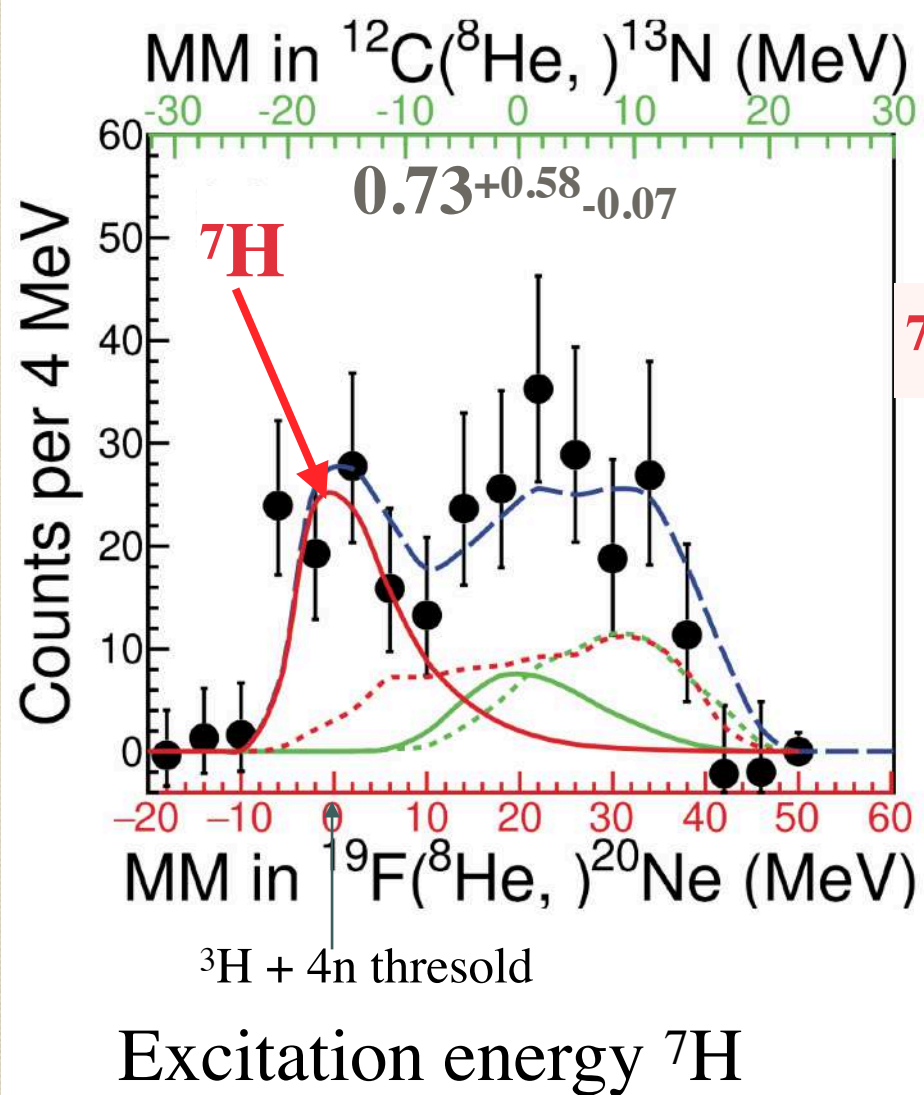
M. Cammano, T. Roger, T. Moro *et al.* Phys. Lett. B 829 (2022) 137067

Active Target gas He + CF₄



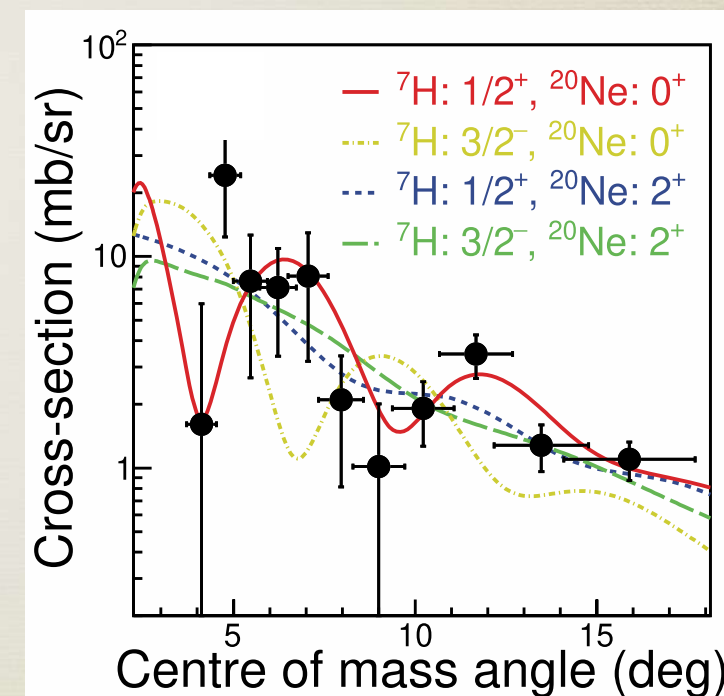
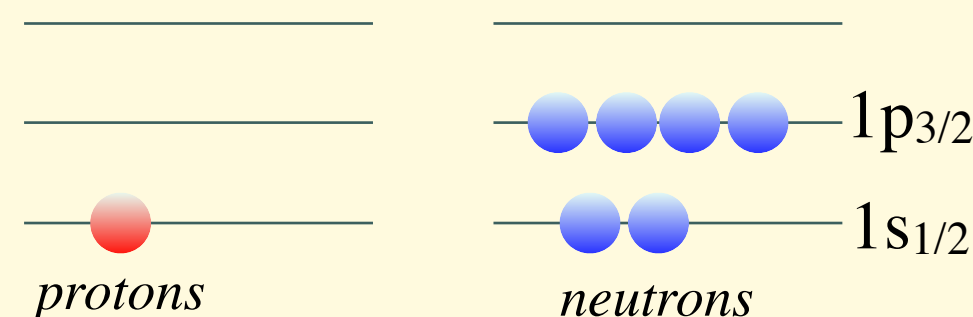
${}^{12}\text{C}({}^8\text{He}, {}^7\text{H}){}^{13}\text{N}$

${}^{19}\text{F}({}^8\text{He}, {}^7\text{H}){}^{20}\text{Ne}$



${}^7\text{H}$ Resonance energy

$0.73^{+0.58}_{-0.07}$ MeV



Resonance spin : $1/2^+$!

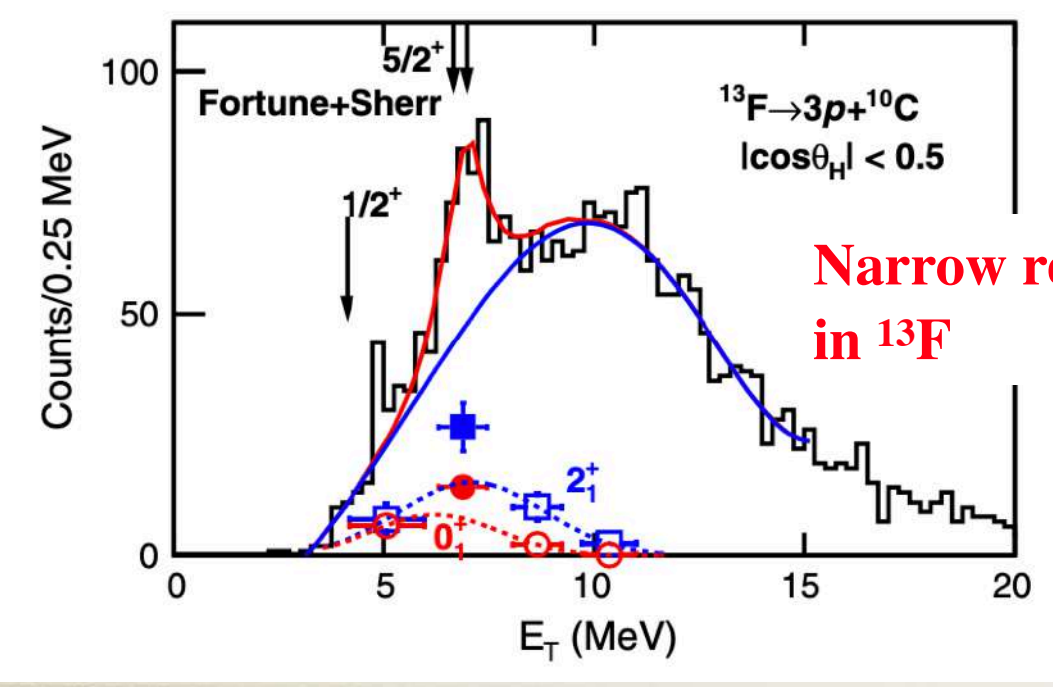
${}^7\text{H}$ resonance observed ~ 0.7 MeV. ${}^3\text{H} + 4n$ structure deduced.



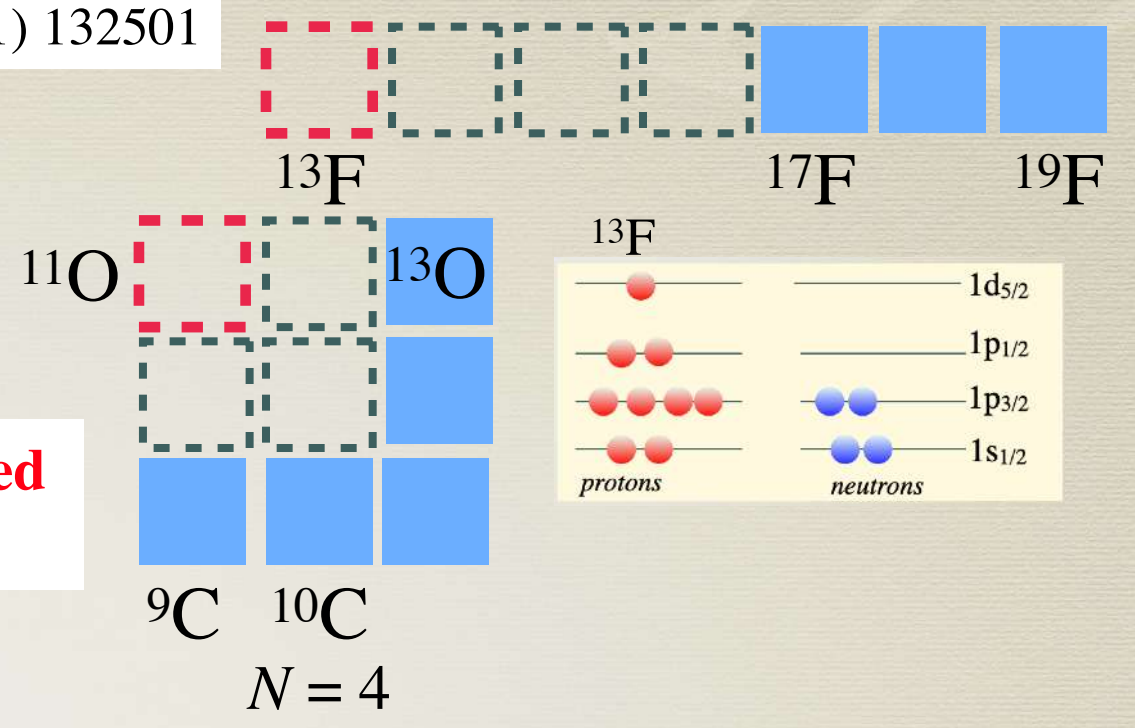
Beyond the proton drip-line ^{13}F , ^{11}O

@ NSCL - HIRA

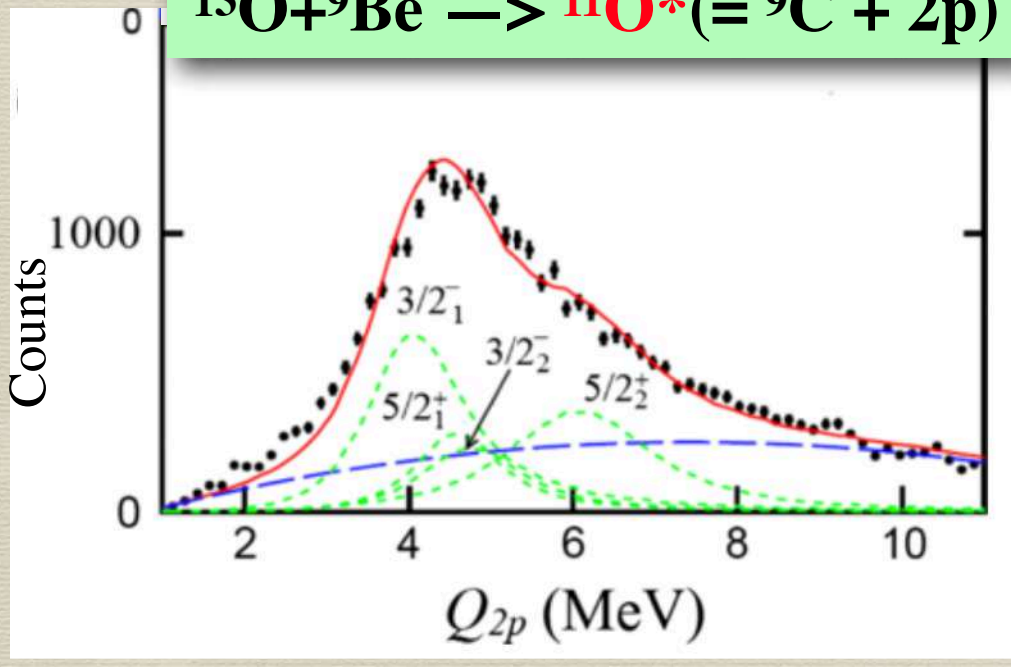
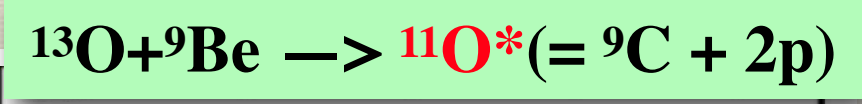
R.J. Charity, T.B. Webb, J.M. Elson *et al.* Phys. Rev. Lett. 126 (2021) 132501



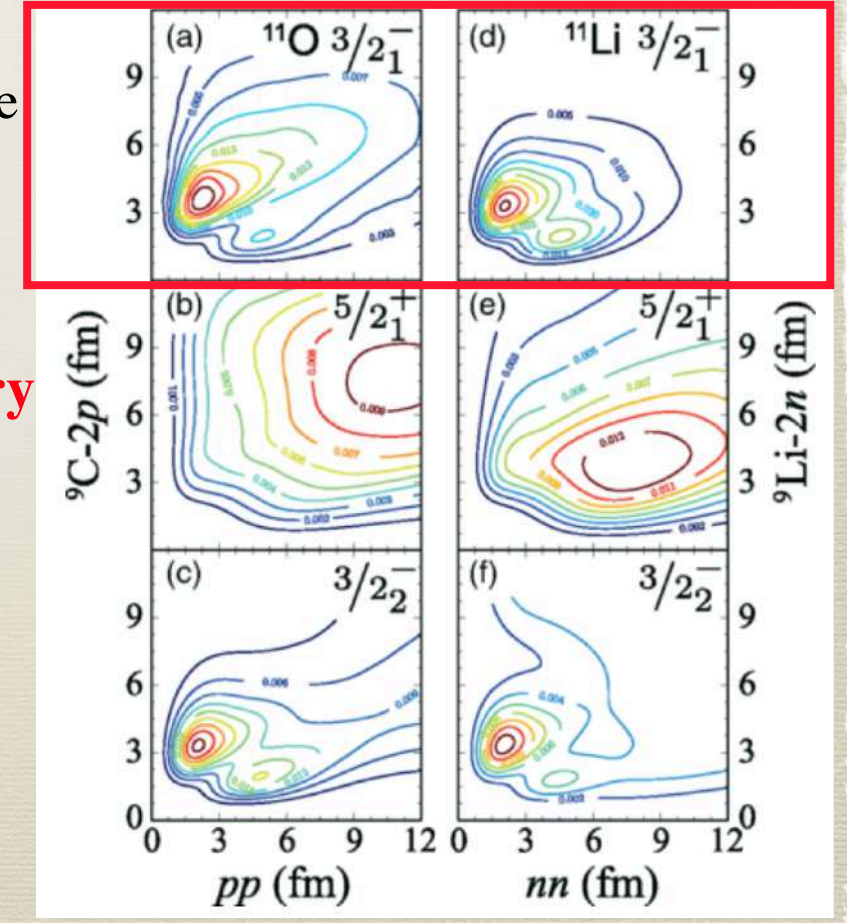
Narrow resonance sighted in ^{13}F



T. Webb *et al.* Phys. Rev. Lett. 122 (2019) 122501



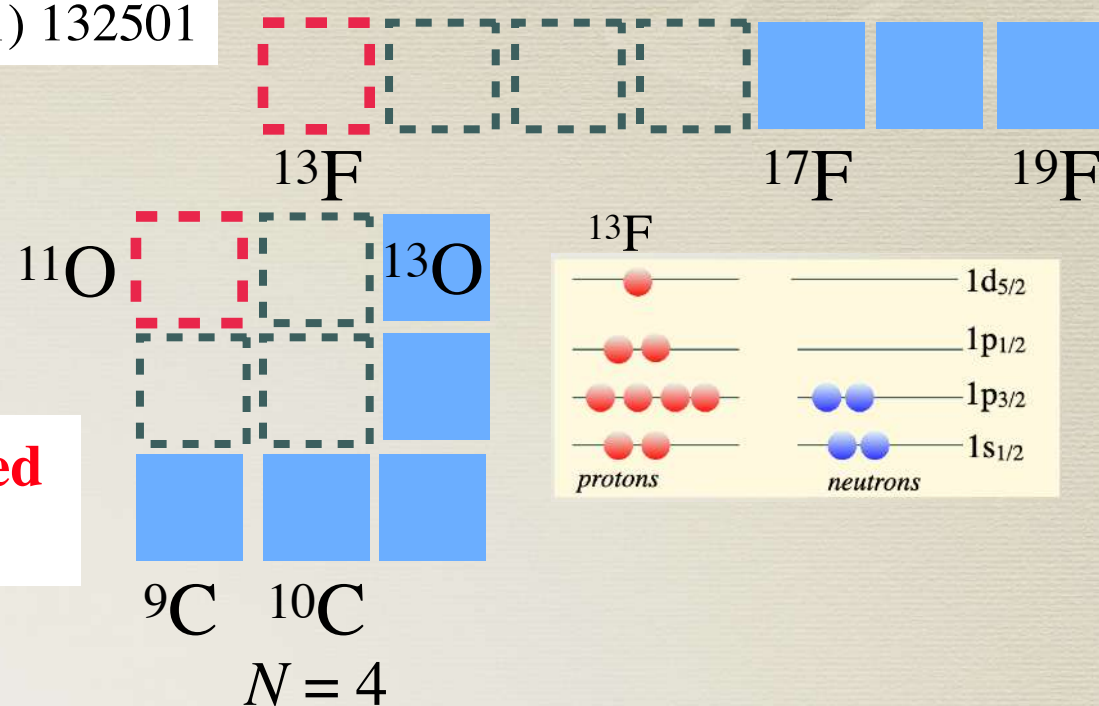
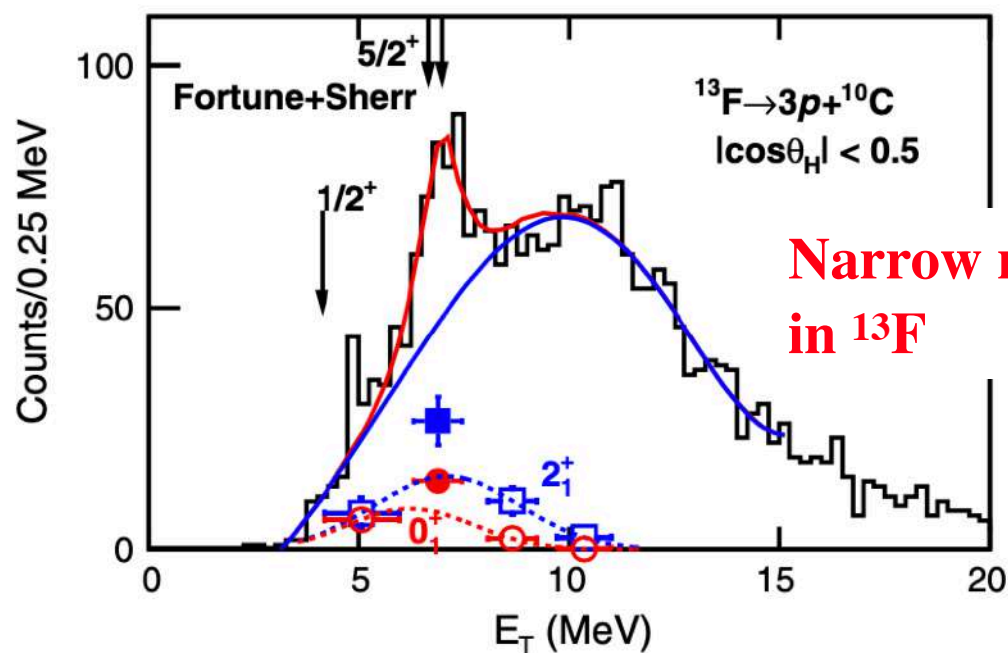
**^{11}O is mirror to ^{11}Li
Very small isospin asymmetry seen between ^{11}O and ^{11}Li**



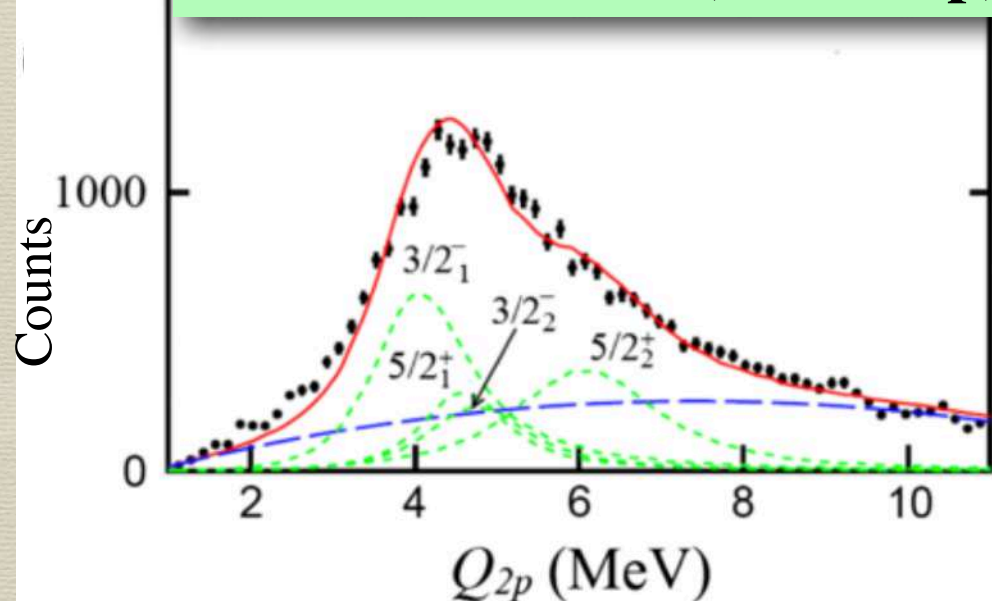
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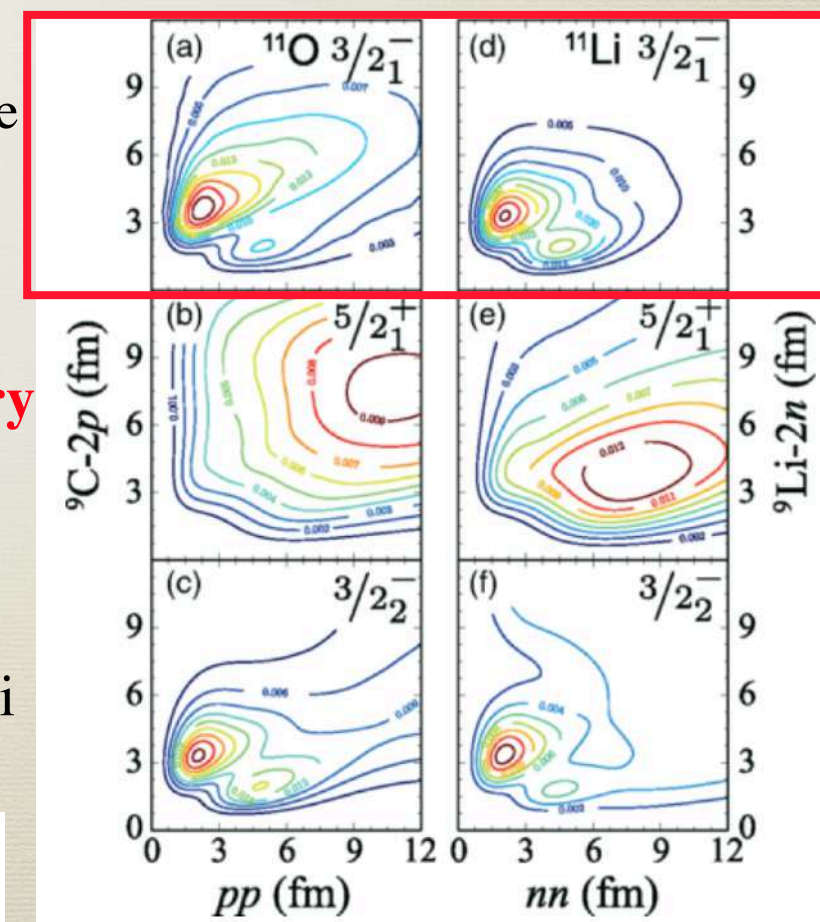
T. Webb *et al.* Phys. Rev. Lett. 122 (2019) 122501



^{11}O is mirror to ^{11}Li
Very small isospin asymmetry
seen between ^{11}O and ^{11}Li

Shape changes in mirror nuclei
 ^{70}Kr & ^{70}Se

K. Wimmer, W. Korten, P. Doornenbal
et al. Phys. Rev. Lett. 126 (2021)
072501

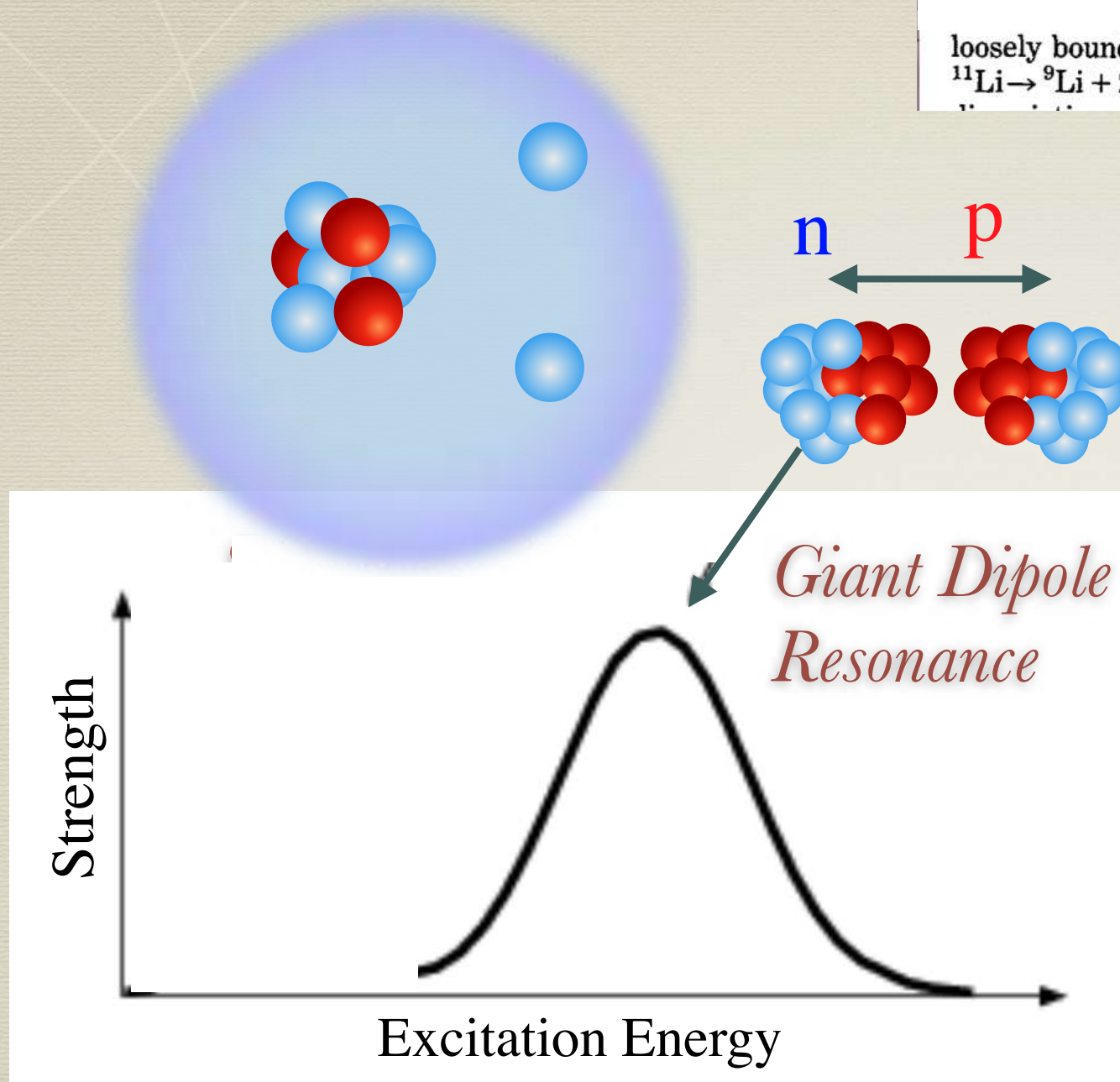


Effects of neutron skin

Equation of state of asymmetric nuclear matter

Neutron skin/halo oscillation

it is tempting to speculate that a loosely bound nucleus such as ^{11}Li will have a **soft electric-dipole mode** so that the reaction $^{11}\text{Li} \rightarrow ^9\text{Li} + 2n$ can be excited in Coulomb collisions at relatively low energy.



(Non-resonant) dipole (E1)
enhancement due halo density tail.

Soft dipole resonance
Halo oscillation

“Pygmy” dipole resonance
Neutron skin oscillation

(traditionally oscillation of neutrons outside $N = Z$ core)

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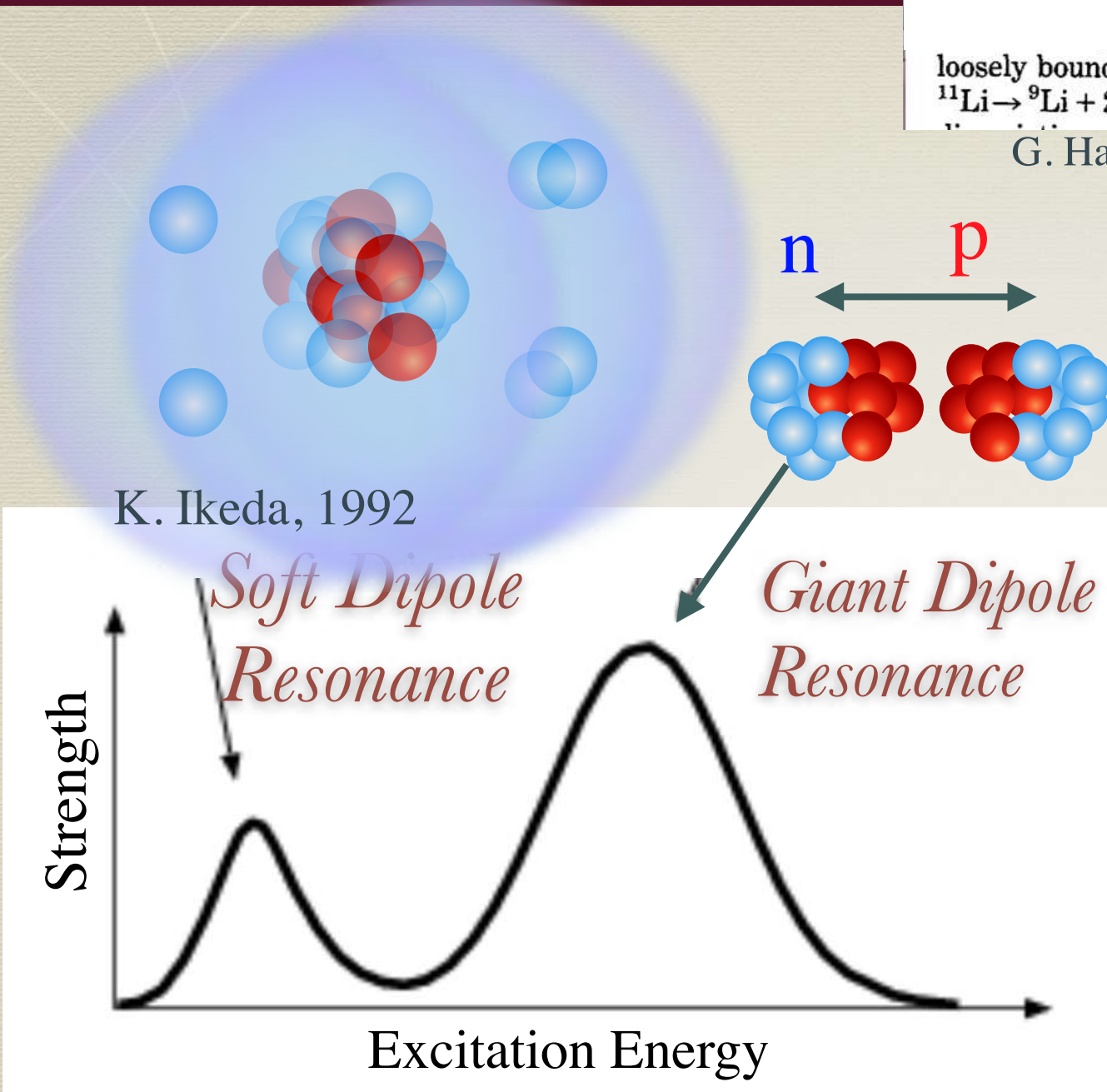
G. Hansen and B. Jonson, 1987

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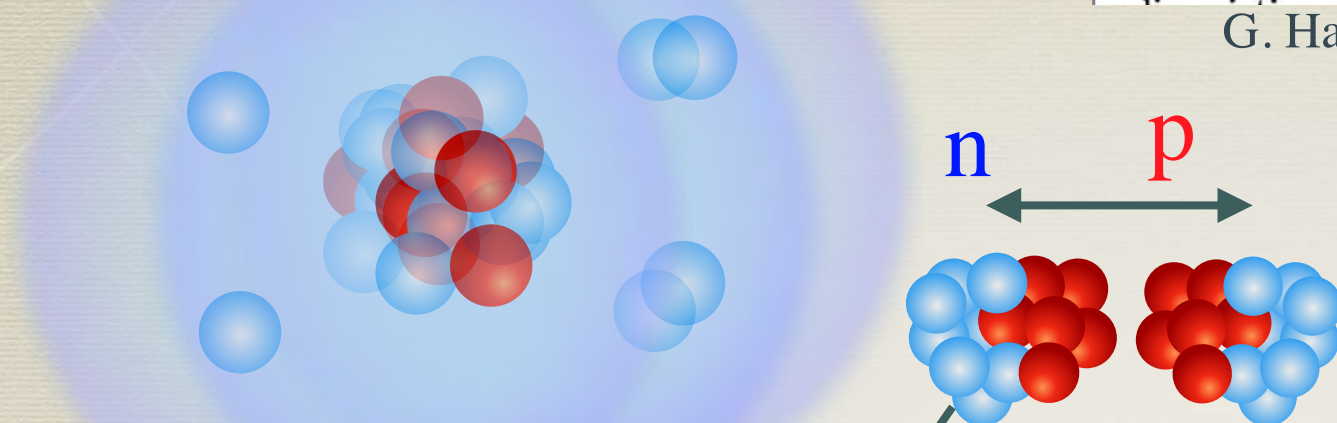
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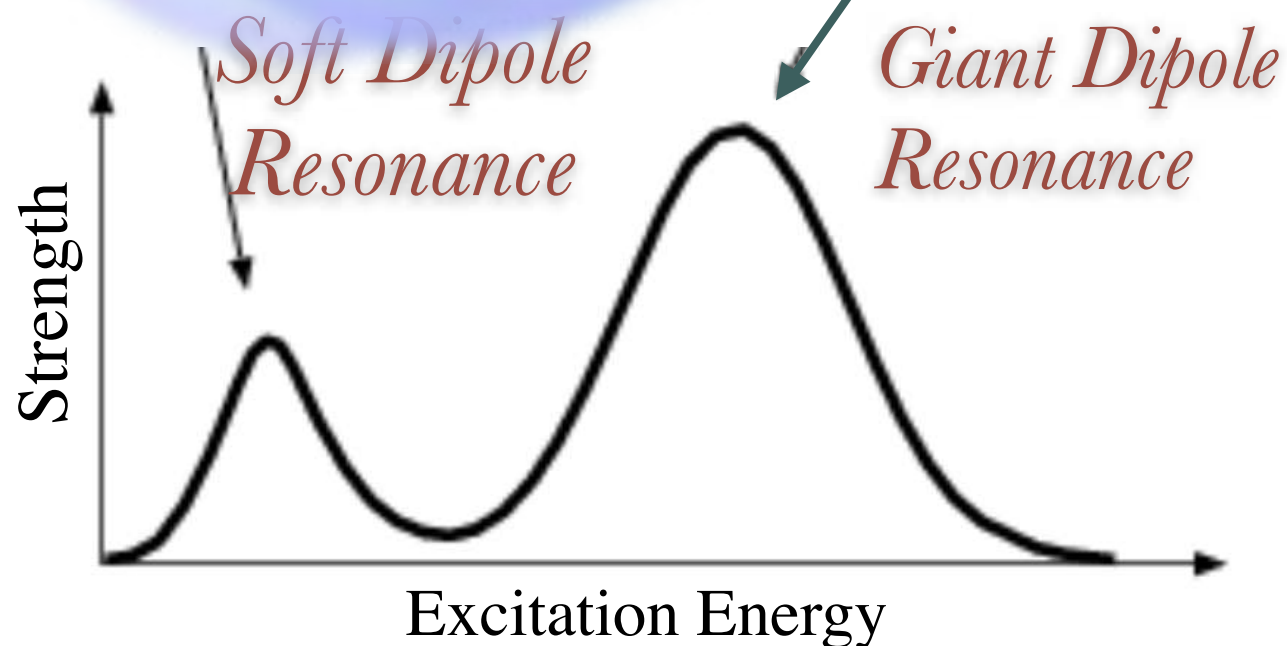
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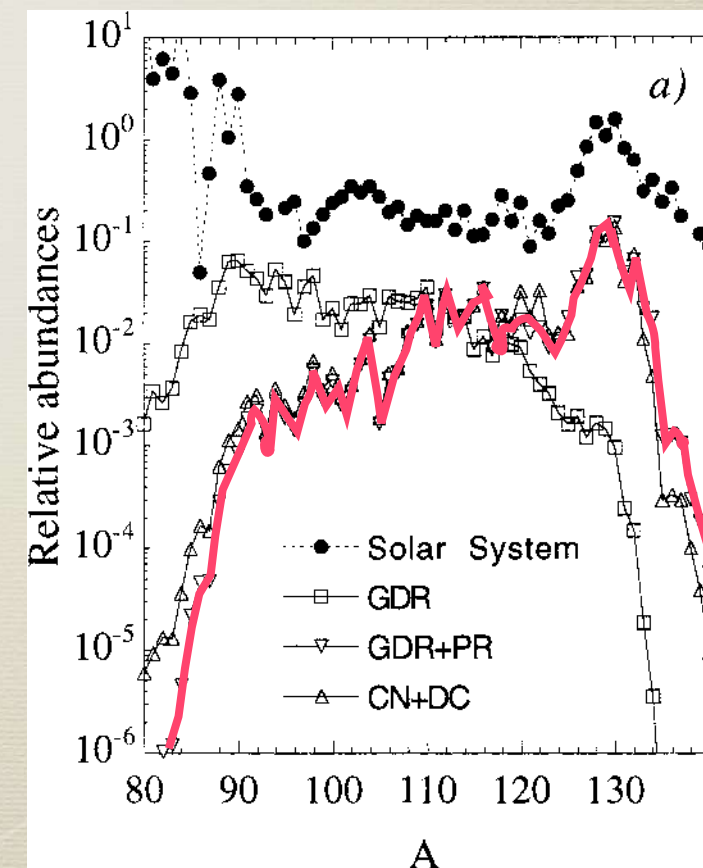


K. Ikeda, 1992



Low-energy dipole resonances ($\sim 1 - 2$ MeV above S_n) can enhance neutron capture rates in r -process

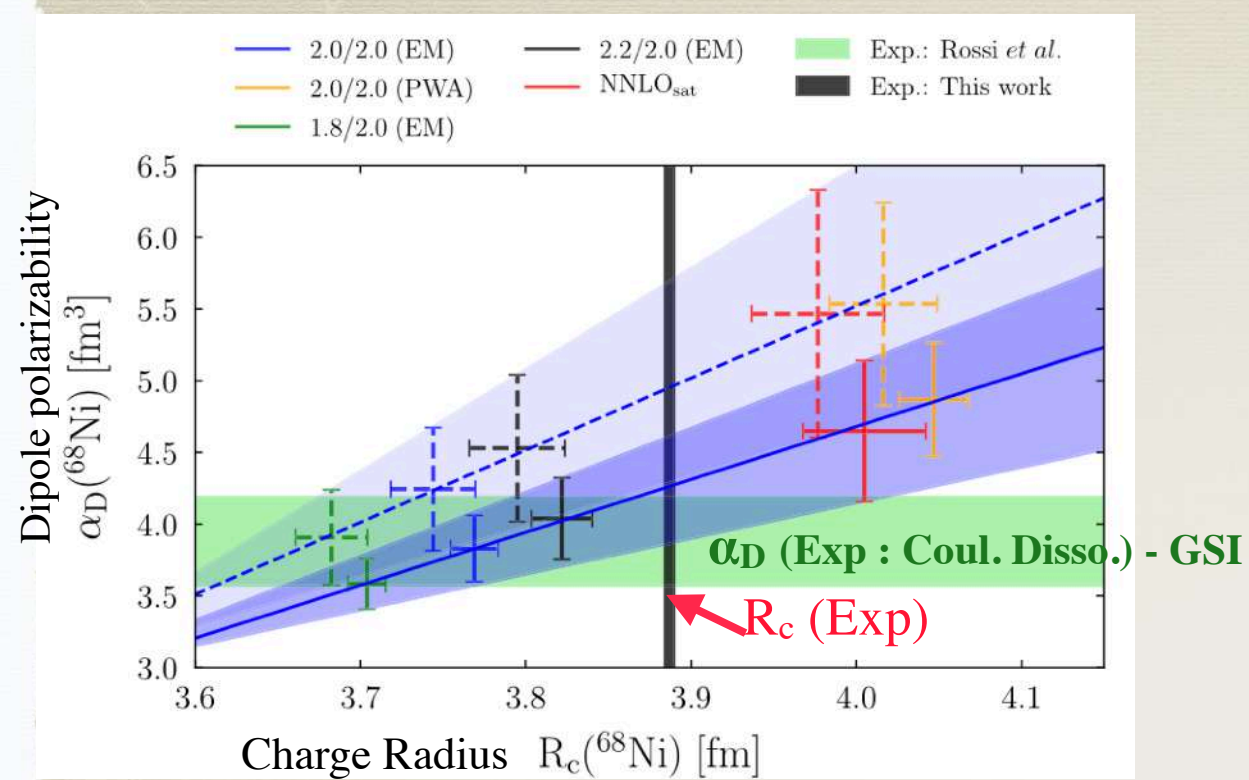
S. Goriely, '2001



Charge radius ^{68}Ni ($N = 40$)

@ ISOLDE - COLLAPS

S. Kaufmann, J. Simonis, S. Bacca *et al.* Phys. Rev. Lett. 124 (2020) 132502



Measured charge radius : challenge for chiral interactions

Ab initio calculations show a correlation between dipole polarizability and charge radius. 3p-3h correlation explains α_D & R_c .

Pygmy dipole ^{68}Ni

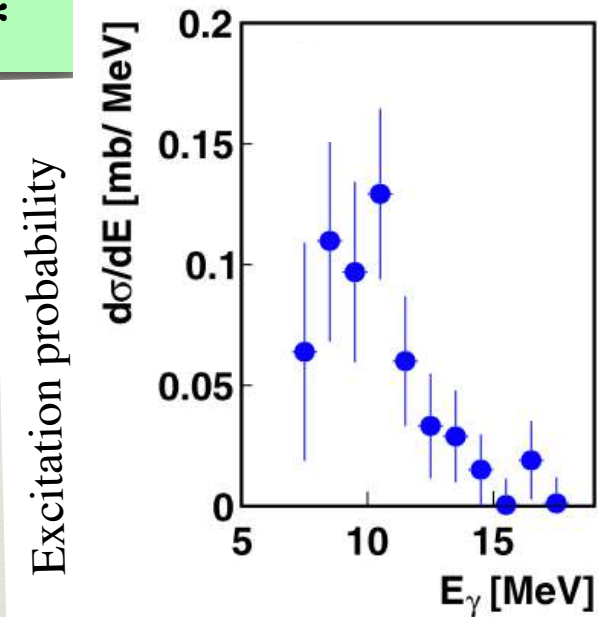
@ INFN - LNS

N.S. Martorana, G. Cardella, E.G. Lanza *et al.* Phys. Lett. B 782 (2018) 112



Isoscalar pygmy dipole ~ 10 MeV.

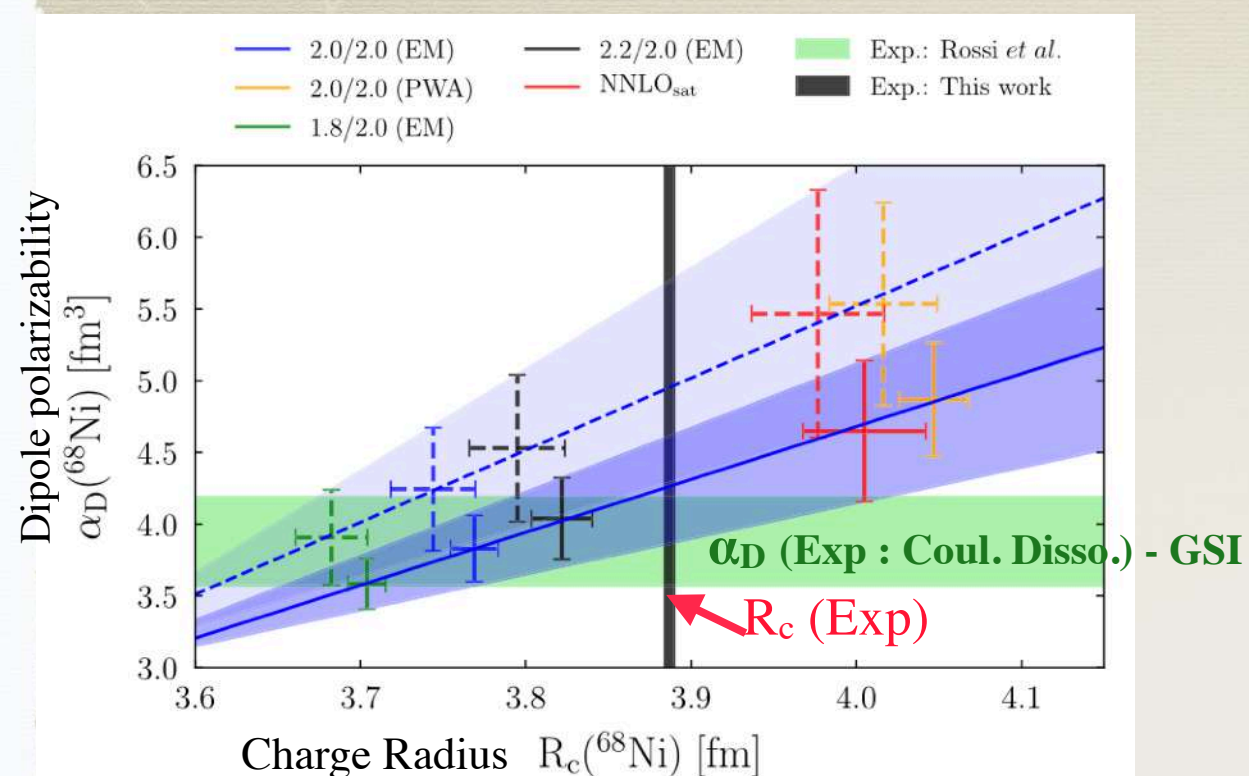
Consistent energy with Coulomb excitation @ GSI - LAND



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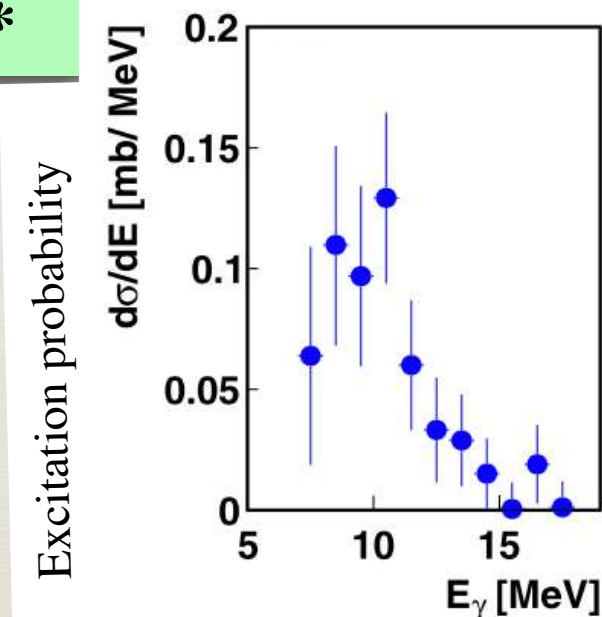
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Dipole Strength $^{62,64}\text{Fe}$ ($N = 36,38$)

@ GSI - AGATA

R. Avigo, O. Wieland, A. Bracco *et al.* Phys. Lett. B 811 (2020) 135951

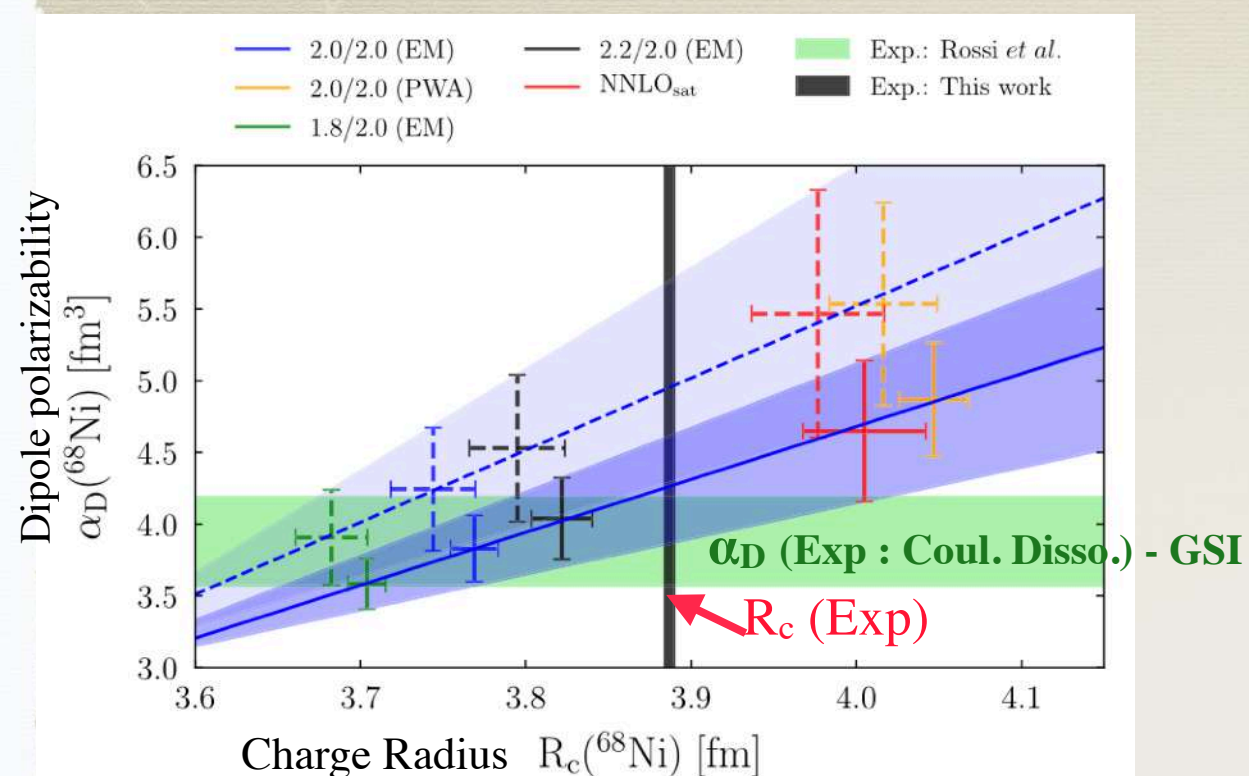
E1 strength at high E_γ increases with increasing neutron number and has complex 3p-3h structure



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@ ISOLDE - COLLAPS

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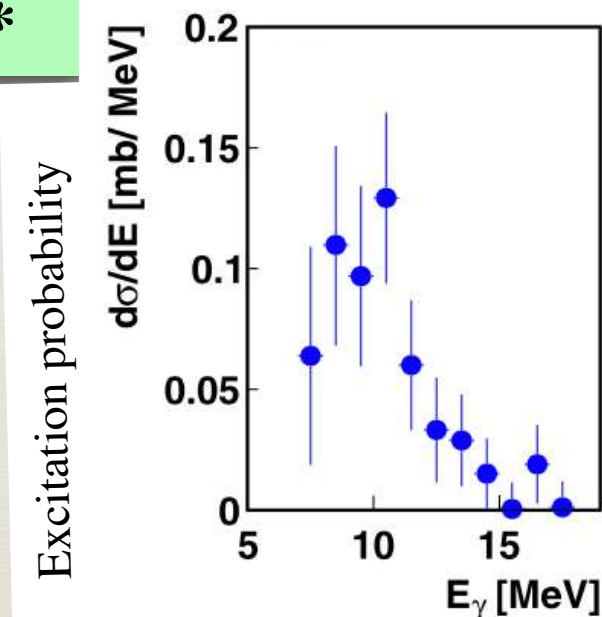
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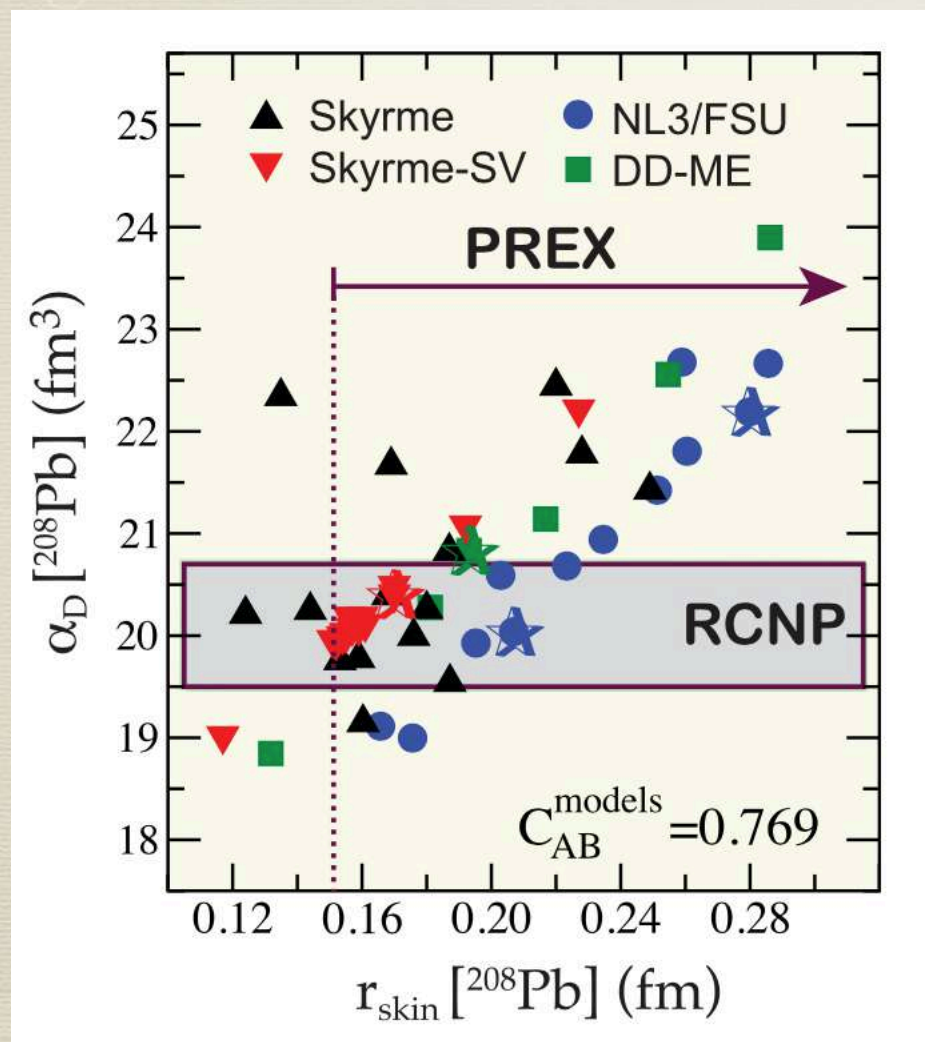
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Low-energy dipole resonances in neutron-rich heavy nuclei yet to be found

Dipole polarizability & neutron skin

Dipole polarizability (α_D)

$$\alpha_D = \frac{\hbar c}{2\pi^2} \int \frac{\sigma_{\text{abs}}}{E_x^2} dE_x = \frac{8\pi}{9} \int \frac{B(E1)}{E_x} dE_x$$



J. Piekarewicz (2012)

Neutron skin : bridge from neutron-rich nuclei to neutron stars

Equation of state of asymmetric nuclear matter

$$e(\rho, \delta) = e(\rho, 0) + c_{\text{sym}}(\rho)\delta^2 + \mathcal{O}(\delta^4) \quad \delta = \frac{\rho_n - \rho_p}{\rho}$$

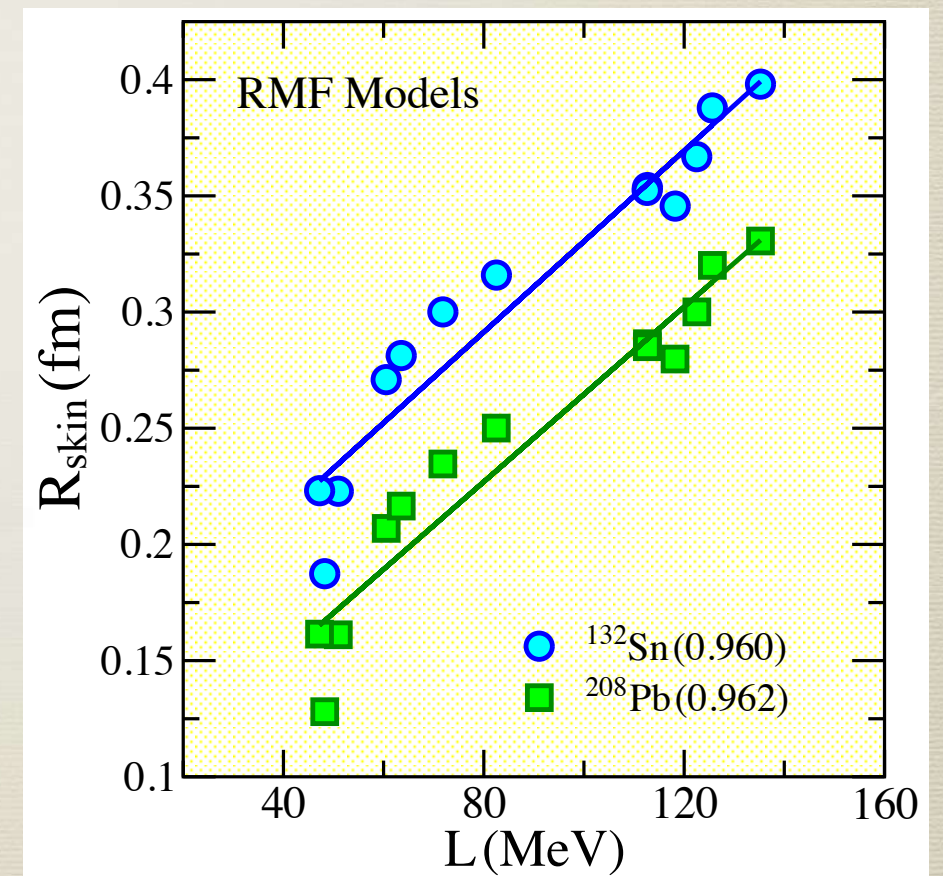
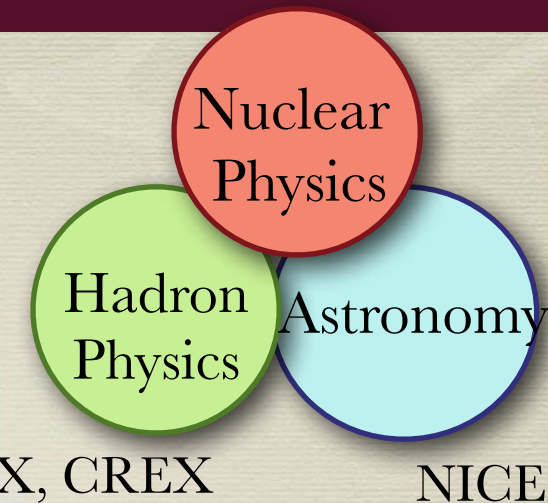
e = energy per particle

Symmetry Energy is poorly constrained

$$c_{\text{sym}}(\rho) = J - L\epsilon + \frac{1}{2}K_{\text{sym}}\epsilon^2 + \mathcal{O}(\epsilon^3) \quad \epsilon = (\rho_0 - \rho)/(3\rho_0)$$

$$L = 3\rho \partial c_{\text{sym}}(\rho) / \partial \rho |_{\rho_0}$$

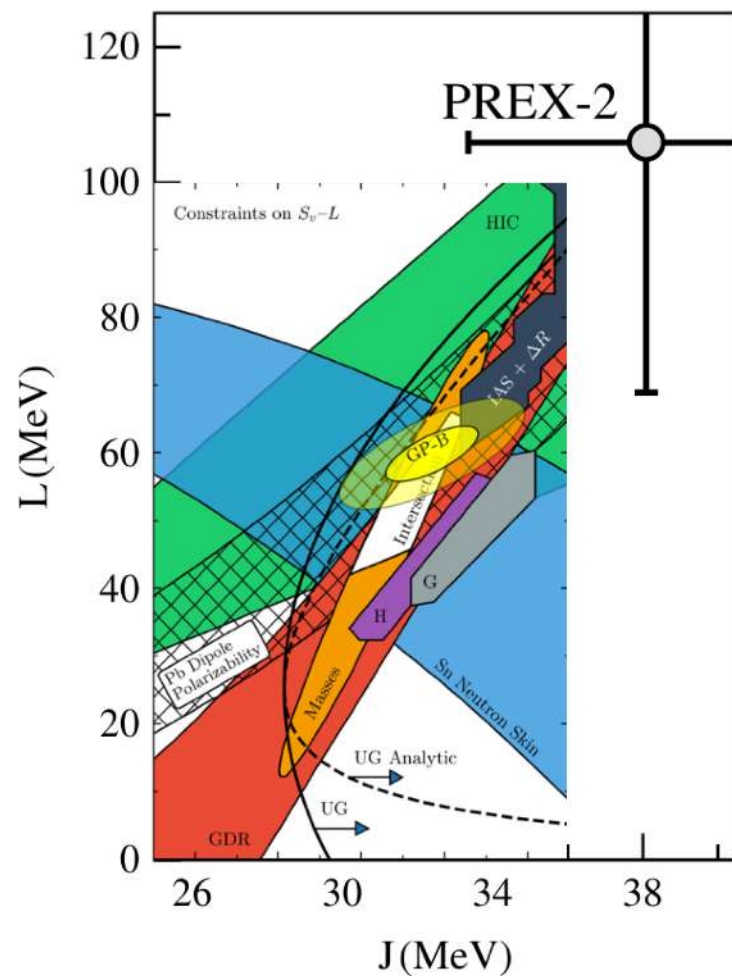
Neutron skin is strongly correlated with L



J. Piekarewicz

Neutron skin (PREX & CREX @ JLab) : Symmetry energy

Parity violating electron scattering



Reed et al., PRL (2021)

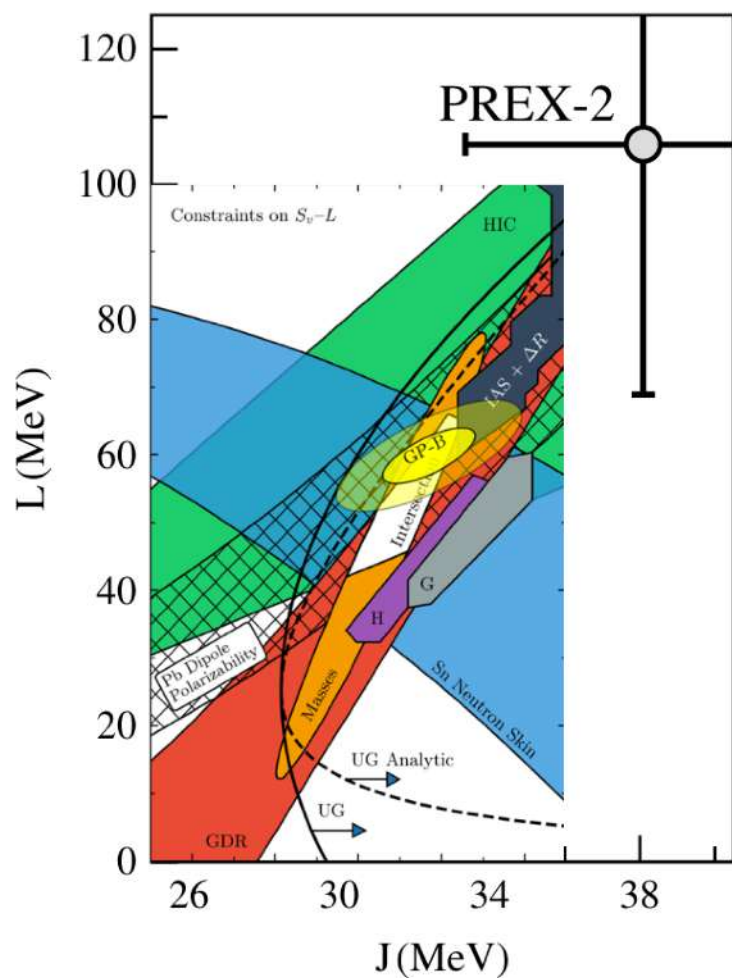
PREX value of neutron skin of ^{208}Pb is higher than other measurements

Higher value of L - Stiffer EOS

D. Adhikari *et al.*, Phys. Rev. Lett. 126 (2021) 172502

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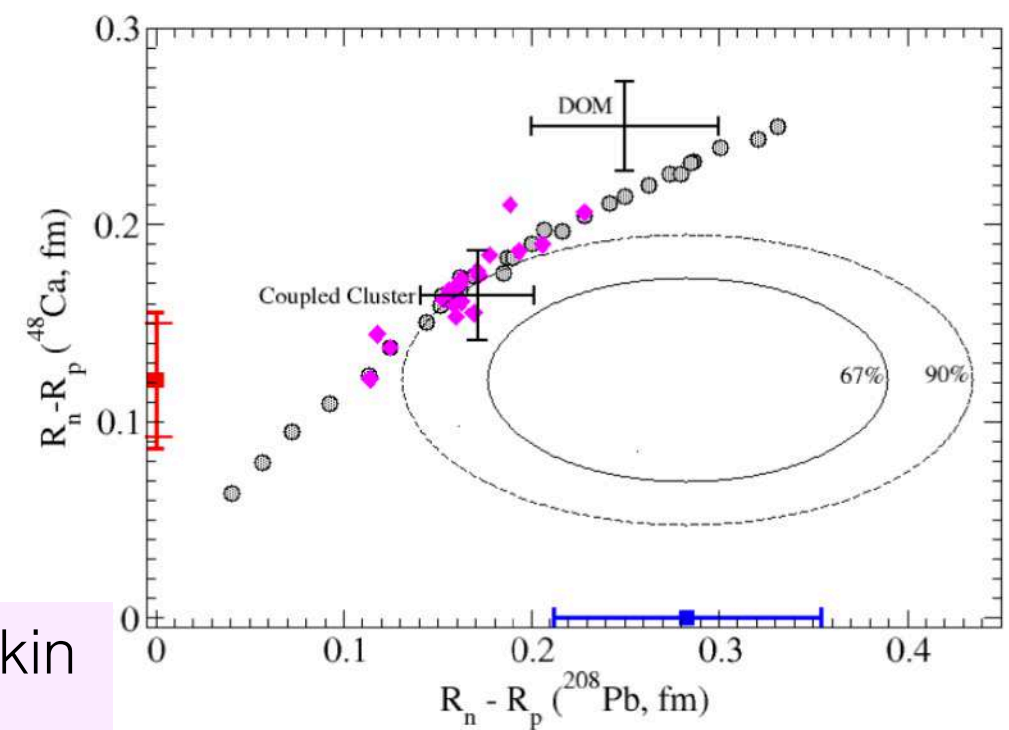


Reed et al., PRL (2021)

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D. Adhikari *et al.*, Phys. Rev. Lett. 126 (2021) 172502



D. Adhikari *et al.*, arXiv:2205.11593v1

CREX finds thin neutron skin

Lower value of L

^{48}Ca neutron skin

Dipole Polarizability

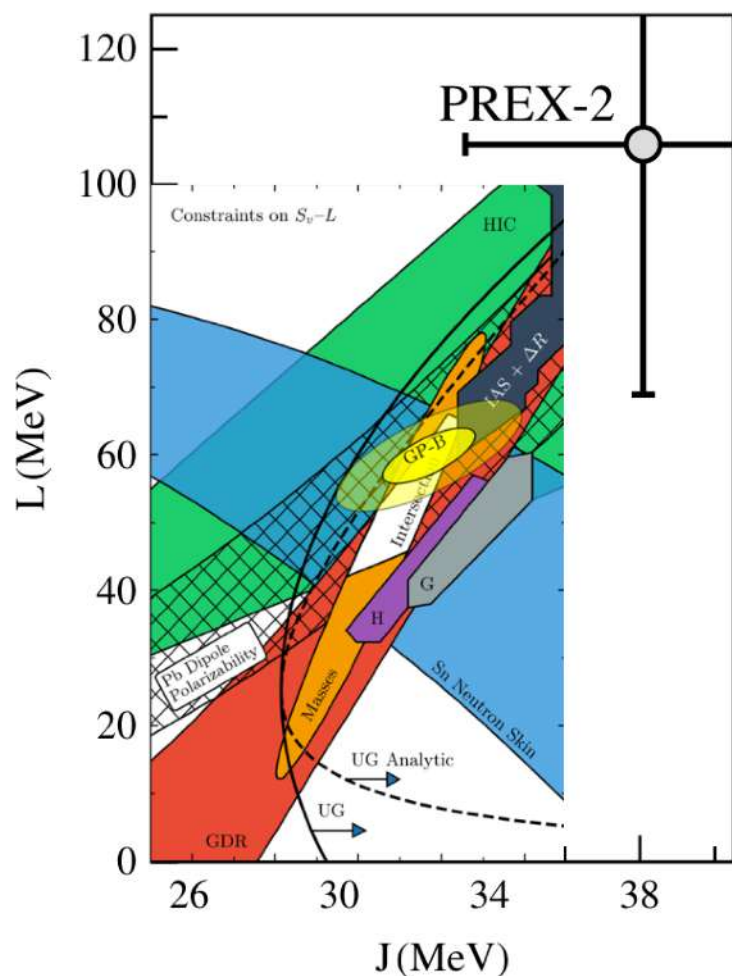
Reaction Cross Section

CREX

} Agree

Neutron skin (PREX & CREX @ JLab) : Symmetry energy

Parity violating electron scattering

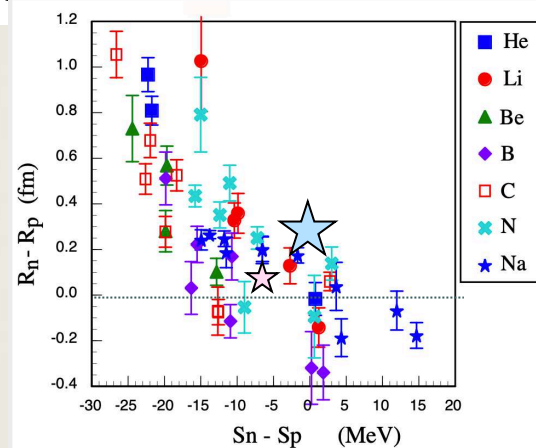


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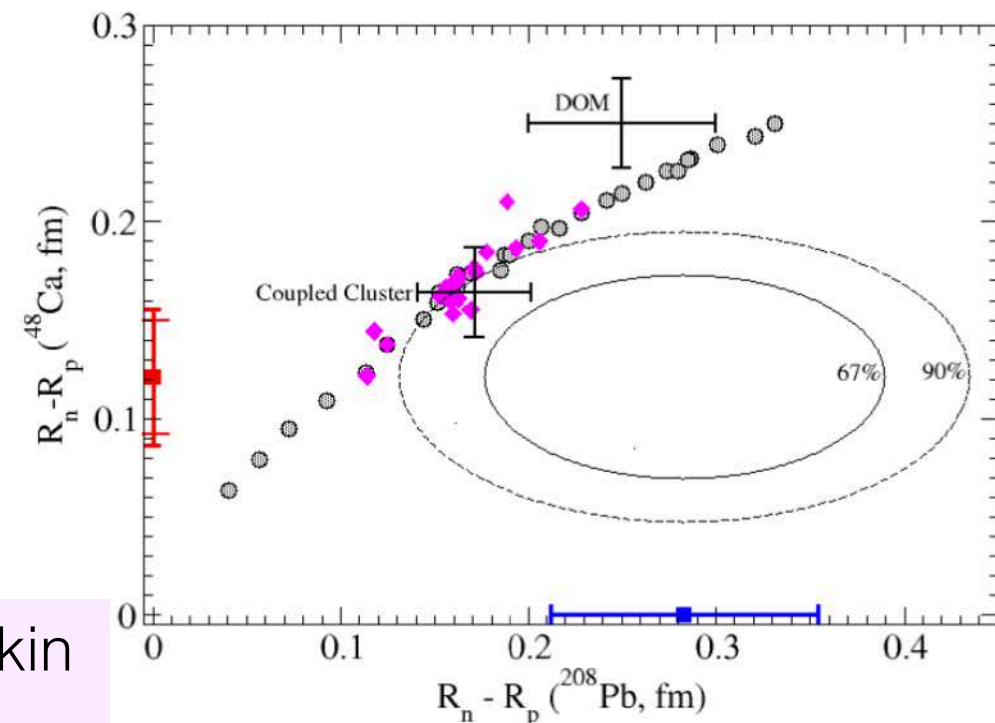
D. Adhikari *et al.*, Phys. Rev. Lett. 126 (2021) 172502



CREX finds thin neutron skin

Lower value of L

^{48}Ca neutron skin



D. Adhikari *et al.*, arXiv:2205.11593v1

Dipole Polarizability

Reaction Cross Section

CREX

} Agree

Rare isotopes with thicker skins will be more sensitive constraints on ' L '

Symmetry energy constraint at supra-saturation density

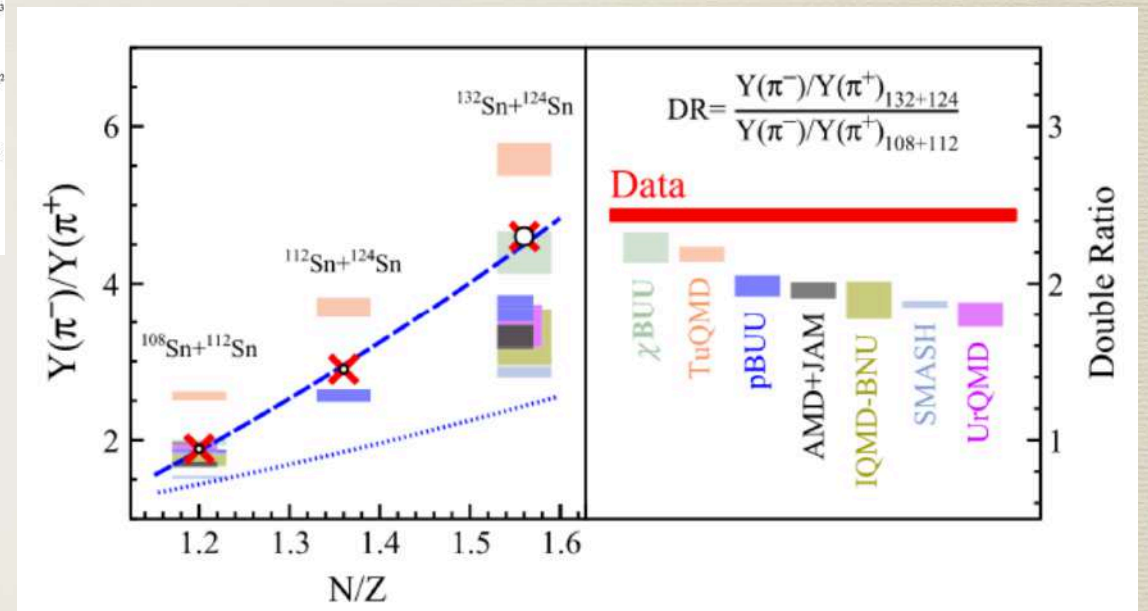
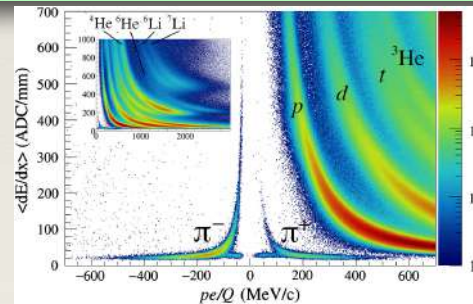
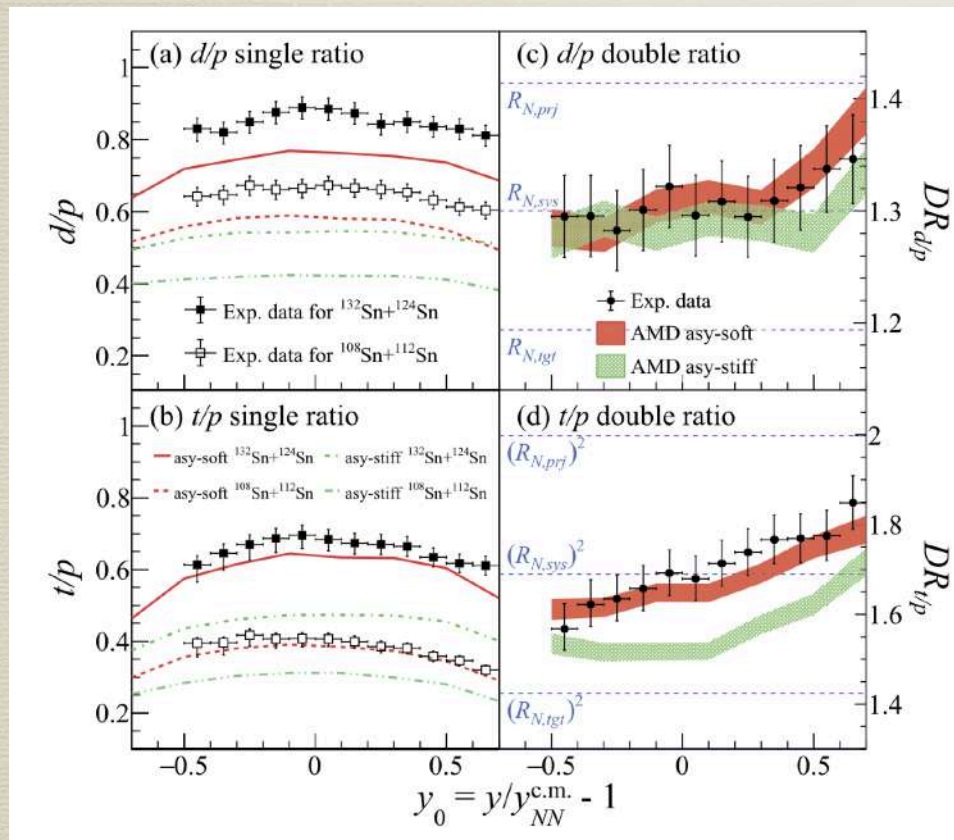
M. Kaneko et al., Phys. Lett. B. 822 (2021) 136681

G. Jhang et al., Phys. Lett. B. 813 (2021) 136016

$^{132}\text{Sn} + ^{124}\text{Sn}$ & $^{108}\text{Sn} + ^{112}\text{Sn}$

p, d, t

π^+, π^-



Small variation in predicted ratio for large (stiff) and small (soft) values of L .

$$DR_{d/p} = \frac{d/p(^{132}\text{Sn} + ^{124}\text{Sn})}{d/p(^{108}\text{Sn} + ^{112}\text{Sn})}, \quad DR_{t/p} = \frac{t/p(^{132}\text{Sn} + ^{124}\text{Sn})}{t/p(^{108}\text{Sn} + ^{112}\text{Sn})}$$

Measured Double Ratio t/p agrees with AMD predictions of soft EOS ($L \sim 46$ MeV)

The differences of transport models make it difficult to place a constraint on L .

Summary & Outlook

Rare Isotopes are enabling to unveil the unknown fundamentals of visible matter in the Universe

- ❖ Exotic forms of nuclei - unique quantum systems emerge far from the valley of stability
 - Do nuclear halos occur in heavy nuclei?
 - What new features of nucleon-nucleon pairing correlation emerge in neutron-rich nuclei?
 - What new phenomena surface with nuclear halo & skin?
 - Electron - RI Scattering
- ❖ Nuclear shells are mutating
 - How do nuclear change in heavy nuclei?
 - What is their influence on heavy element synthesis?
 - What features of the nuclear force drive the shell evolution?
- ❖ Neutron rich nuclei bring laboratory access to study behaviour of neutron-rich matter (EOS)
 - Connecting neutron skin driven effects to constrain EOS of asymmetric nuclear matter
- ❖ Constraining the nuclear force - defining from first principles : Excited States and Radii
 - Dynamical reaction probes
- ❖ Search for new physics - fundamental symmetries in nature (not covered)
 - Radioactive Molecules as probes of symmetry violation
 - Electric Dipole Moment measurements (Ra, Rn)
 - Beta - neutrino correlation
 - Unitarity of the CKM matrix element
- ❖ Rare strange matter : RI Hypernuclei (not covered)



Summary & Outlook

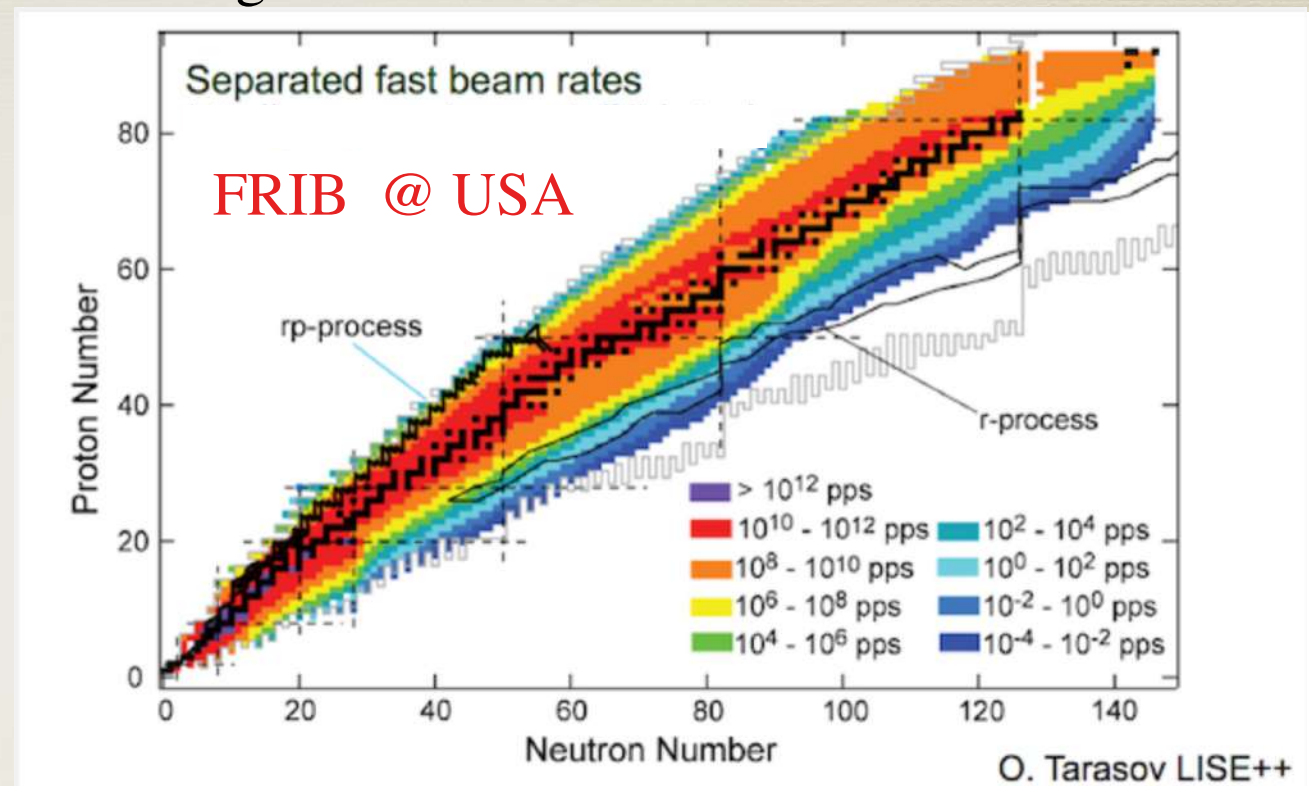
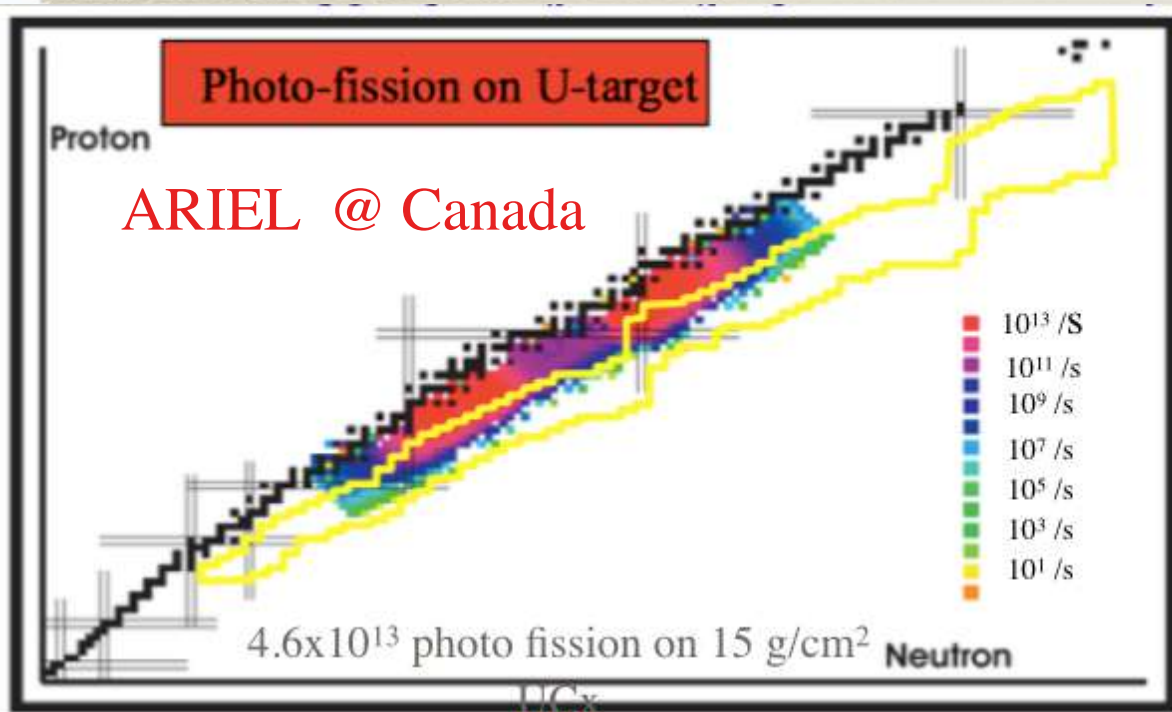
An era of new discoveries awaits in the horizon

New generation facilities bring access to rare isotopes in colliding stars

ISOL

North America

In-flight



- Masses & Half-lives
- Decay spectroscopy
- Charged particle spectroscopy - shell structure
- Safe Coulex - shell structure
- Direct capture

- New isotope search
- Masses, Half-lives, Radii, Decay & In-beam γ spectroscopy
- Knockout, Coulomb Disso. - shell structure
- Transfer reactions - shell structure



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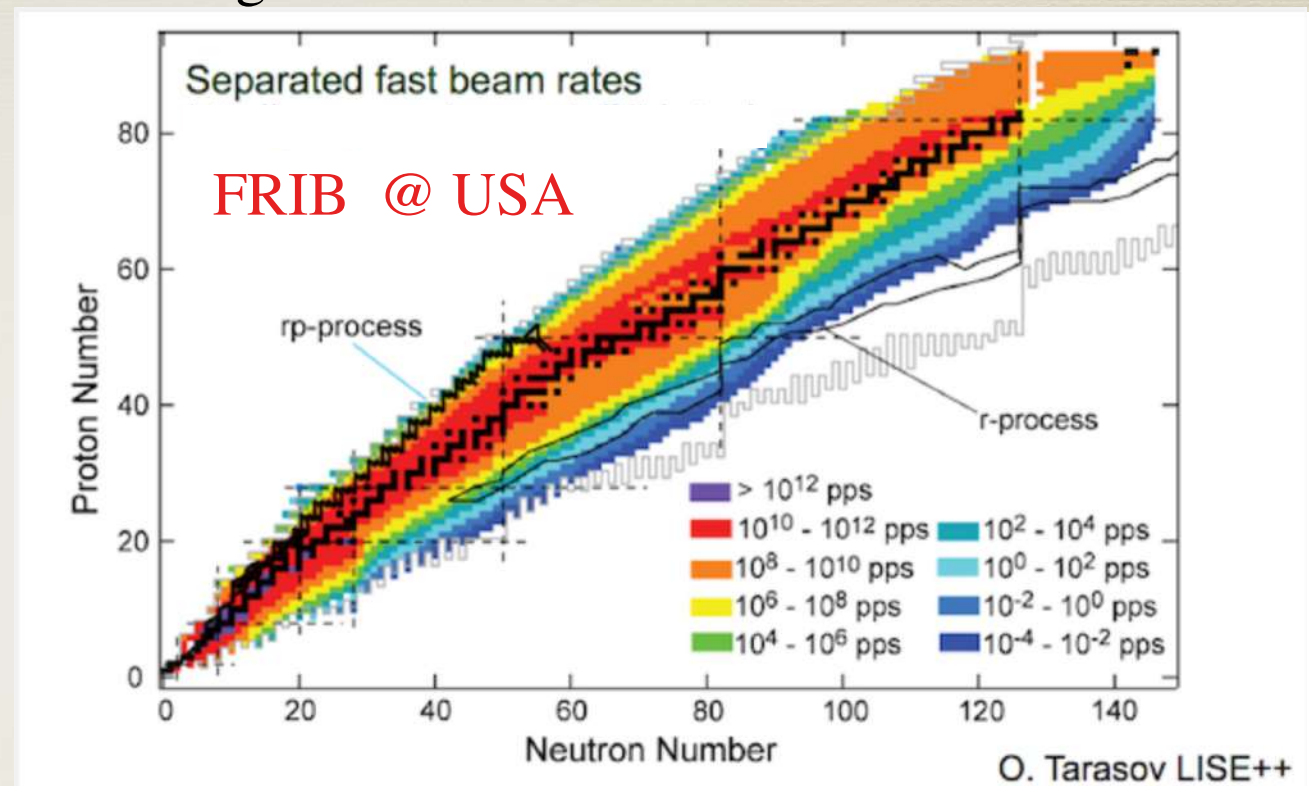
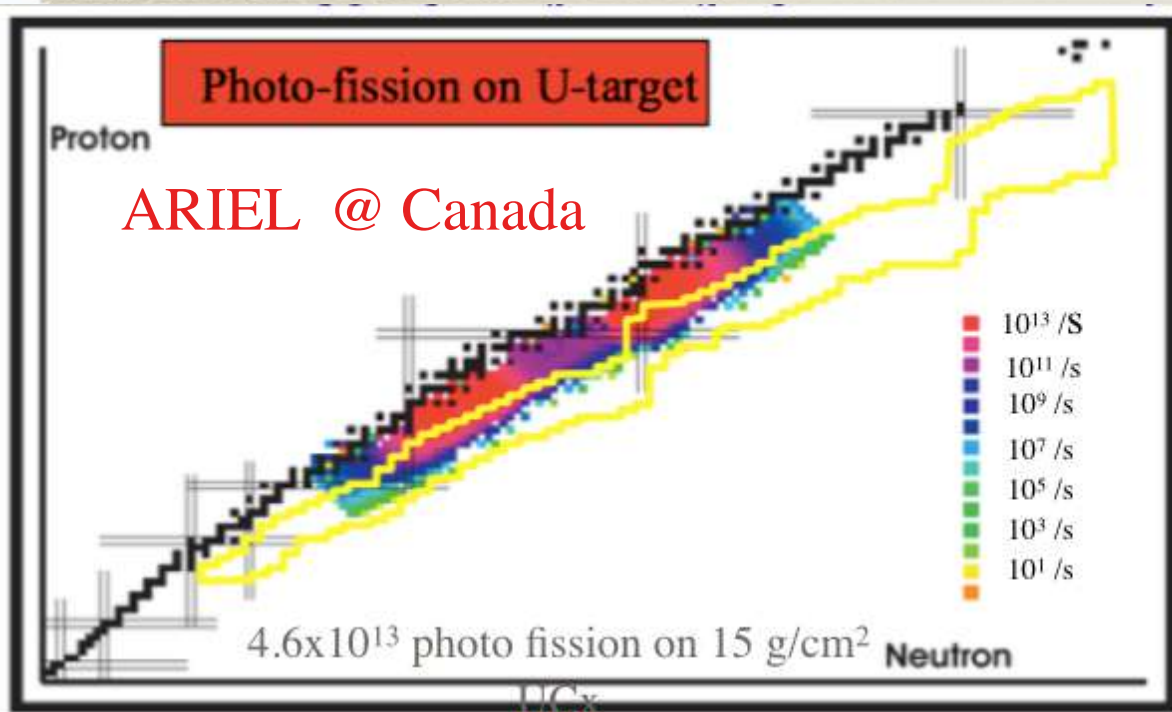
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Thanks : RI Beam Facilities and the funding agencies for enabling the discoveries.

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Thank you for your attention