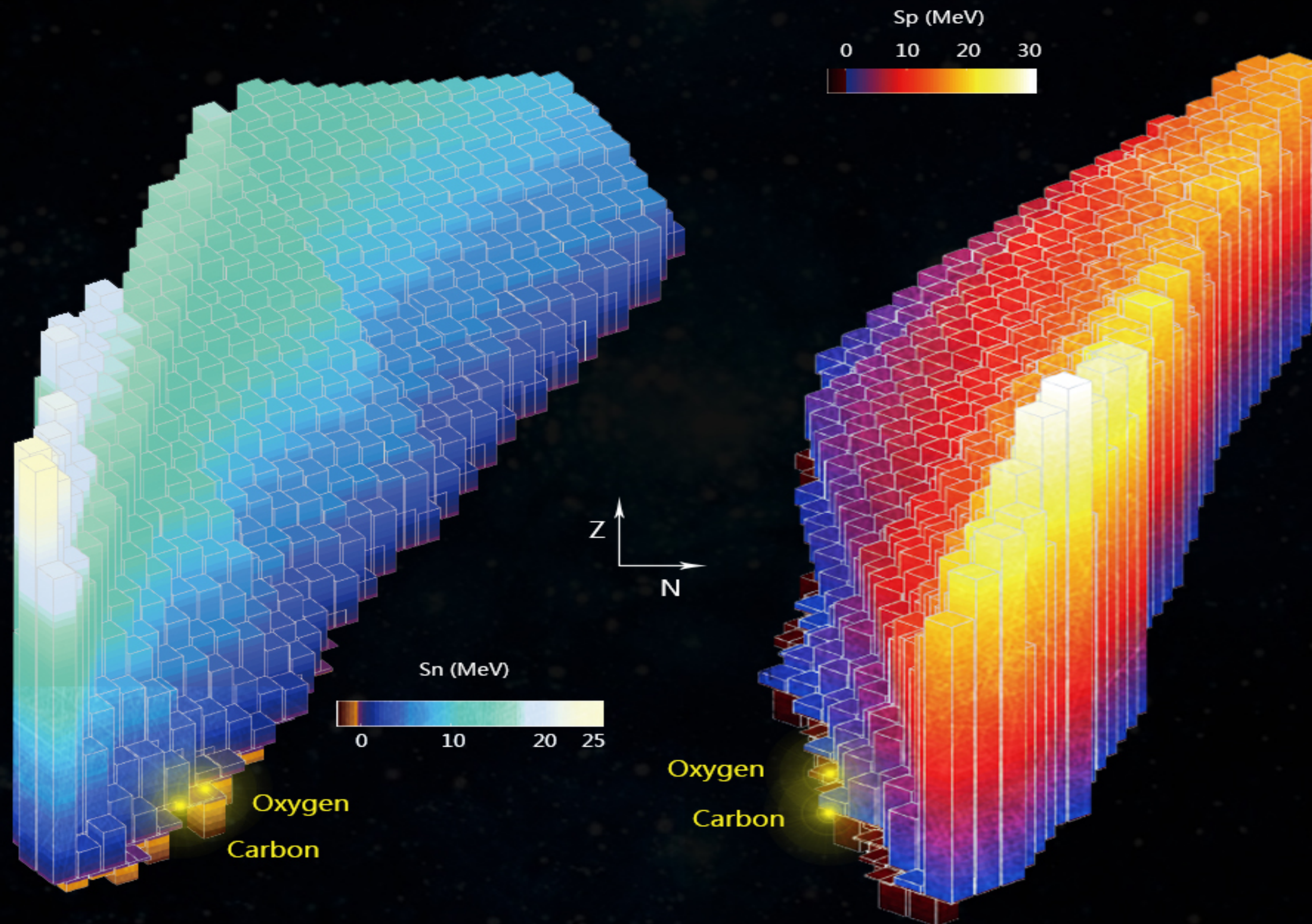


# In-medium similarity renormalization group with resonance and continuum



Bai-Shan Hu (胡柏山)  
Aug 9, 2022  
@ TRIUMF

# Outline

**I. What is IMSRG?**

**II. Why continuum?**

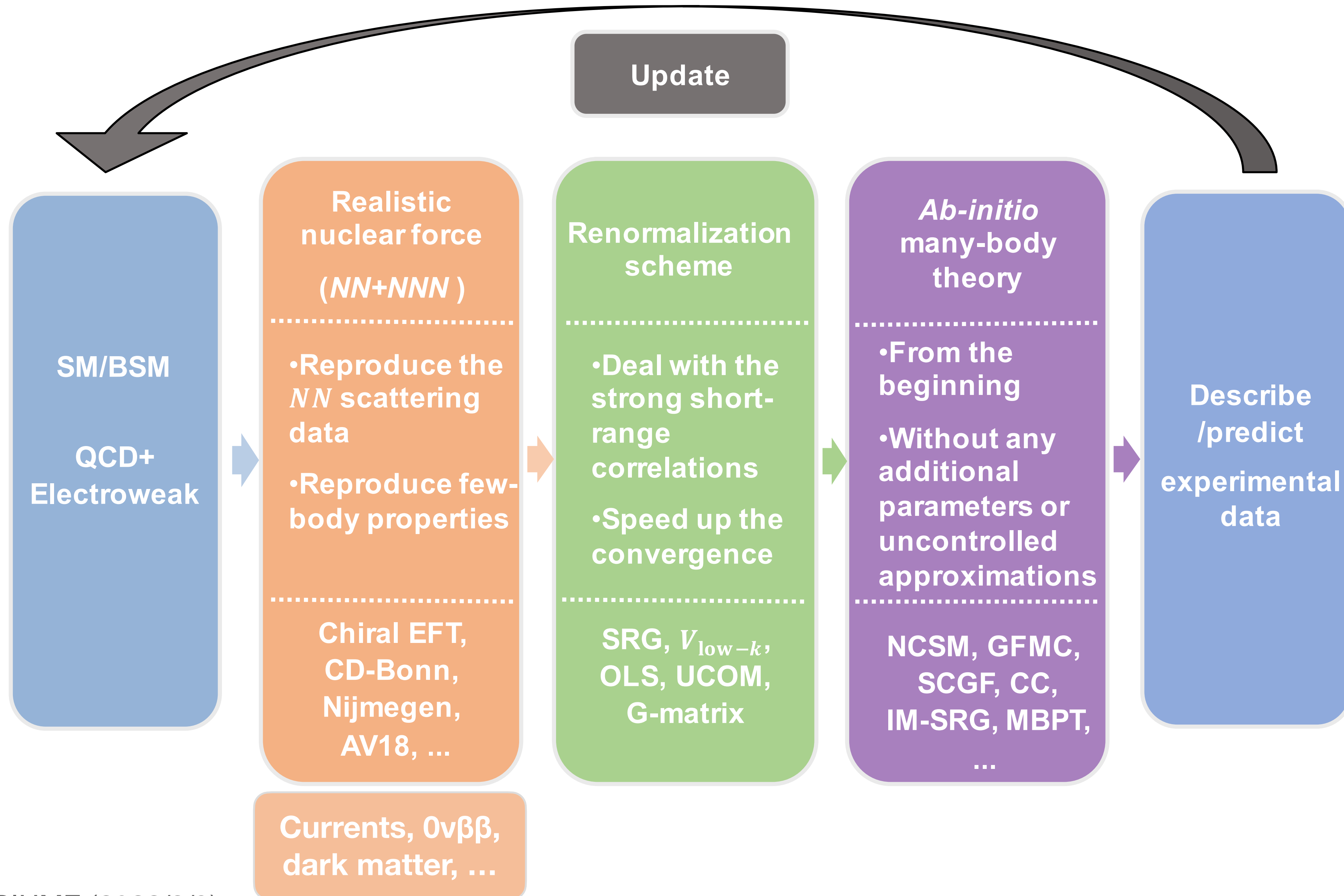
**III. How to include continuum into IMSRG**

 **resonance**

 **halo**

**IV. Summary**

# Workflow of *ab-initio* nuclear calculation



# Similarity Renormalization Group

drive the Hamiltonian towards a band- or block-diagonal form via continuous unitary transformation

$$U(s) \cdot U^\dagger(s) = U(s) \cdot U^{-1}(s) = 1 \quad \frac{dH(s)}{ds} = [\eta(s), H(s)] \quad \eta(s) = \frac{dU(s)}{ds} U^\dagger(s) = -\eta^\dagger(s)$$
$$H(s) = U(s)H(0)U^\dagger(s) \quad O(s) = U(s)OU^\dagger(s)$$

## Developed by Glazek and Wilson and by Wegner around in 1993

S. D. Glazek and K. G. Wilson, Phys. Rev. D 48, 5863 (1993)

F. Wegner, Ann. Phys. (Leipzig) 3, 77 (1994)

## Soften nucleon-nucleon interaction from 2007

S.K. Bogner, R.J. Furnstahl, R.J. Perry, Phys. Rev. C 75, 061001 (2007)

H. Hergert and R. Roth, Phys. Rev. C 75, 051001(R) (2007)

## In-medium SRG as a nuclear many-body method

Tsukiyama, S. K. Bogner, and A. Schwenk, PRL106, 222502 (2011); PRC85, 061304(R) (2012)

H. Hergert, *et al.*, PRL110, 242501 (2013)

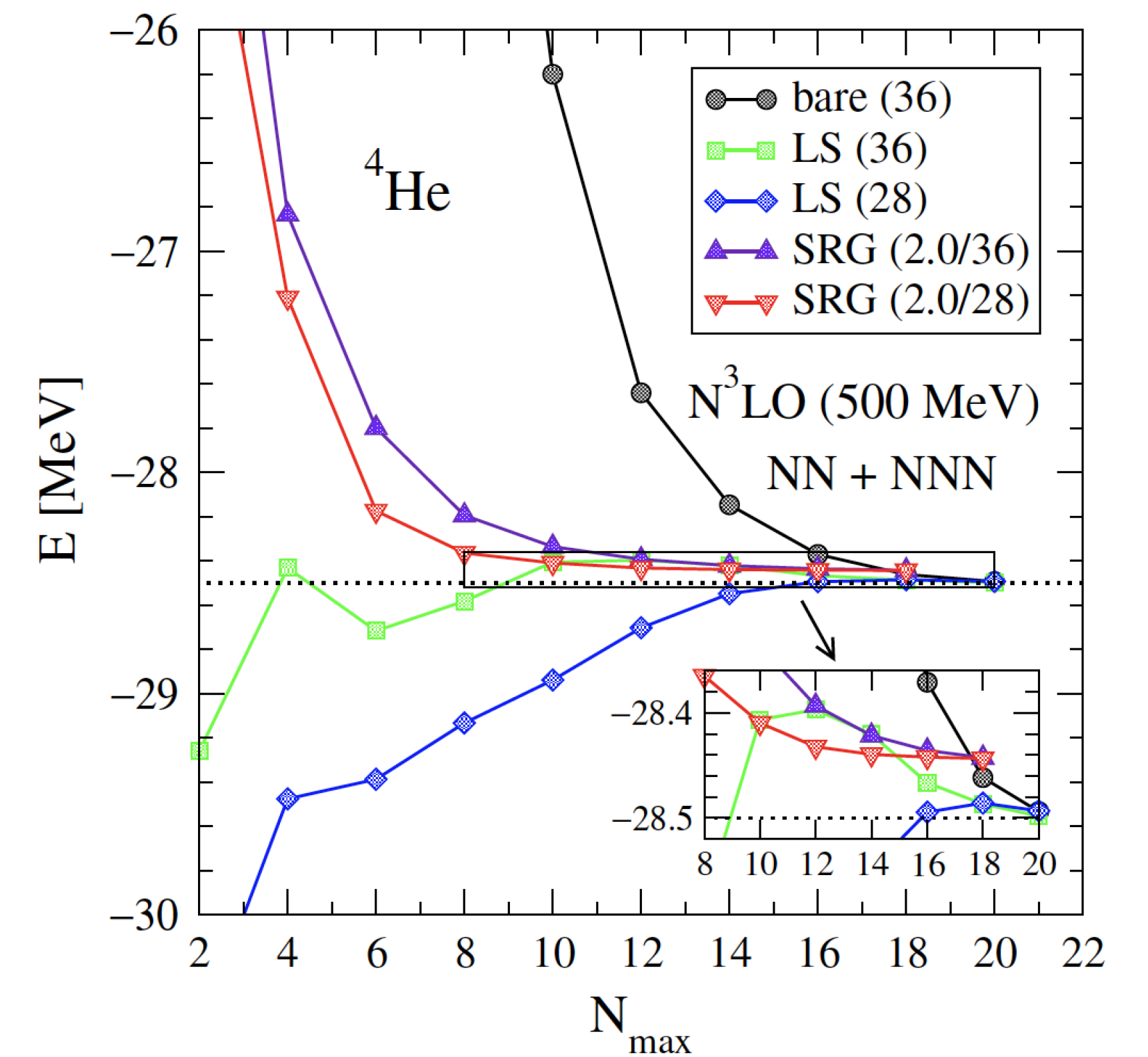
S.R. Stroberg, *et al.*, Phys. Rev. Lett. 118, 032502 (2017)

# SRG: soften nuclear interaction

$$H(s=0) = T_{\text{rel}} + V \quad \eta(s) = [T_{\text{rel}}, H(s)]$$

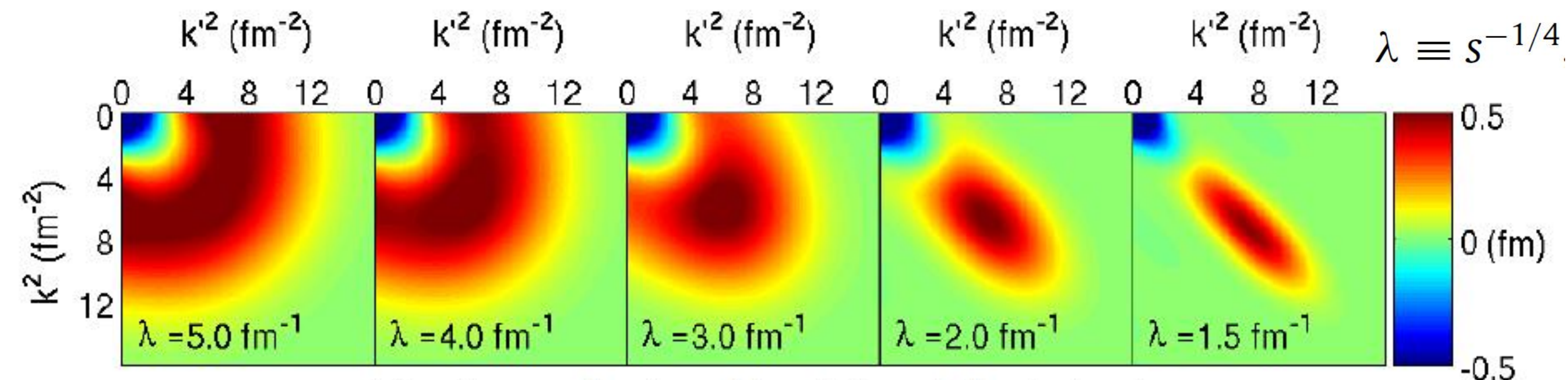
$$\frac{dH(s)}{ds} = [\eta(s), H(s)] \quad \frac{dH(s)}{ds} = [[T_{\text{rel}}, H(s)], H(s)]$$

$$\frac{dV_s(k, k')}{ds} = - (k^2 - k'^2)^2 V_s(k, k') + \frac{2}{\pi} \int_0^\infty q^2 dq (k^2 + k'^2 - 2q^2) V_s(k, q) V_s(q, k')$$



E. D. Jurgenson, P. Navrátil, and R. J. Furnstahl, PRL103, 082501 (2009)

<sup>1</sup>S<sub>0</sub> from N<sup>3</sup>LO (500 MeV) of Entem/Machleidt



<http://www.physics.ohio-state.edu/~ntg/srg/>

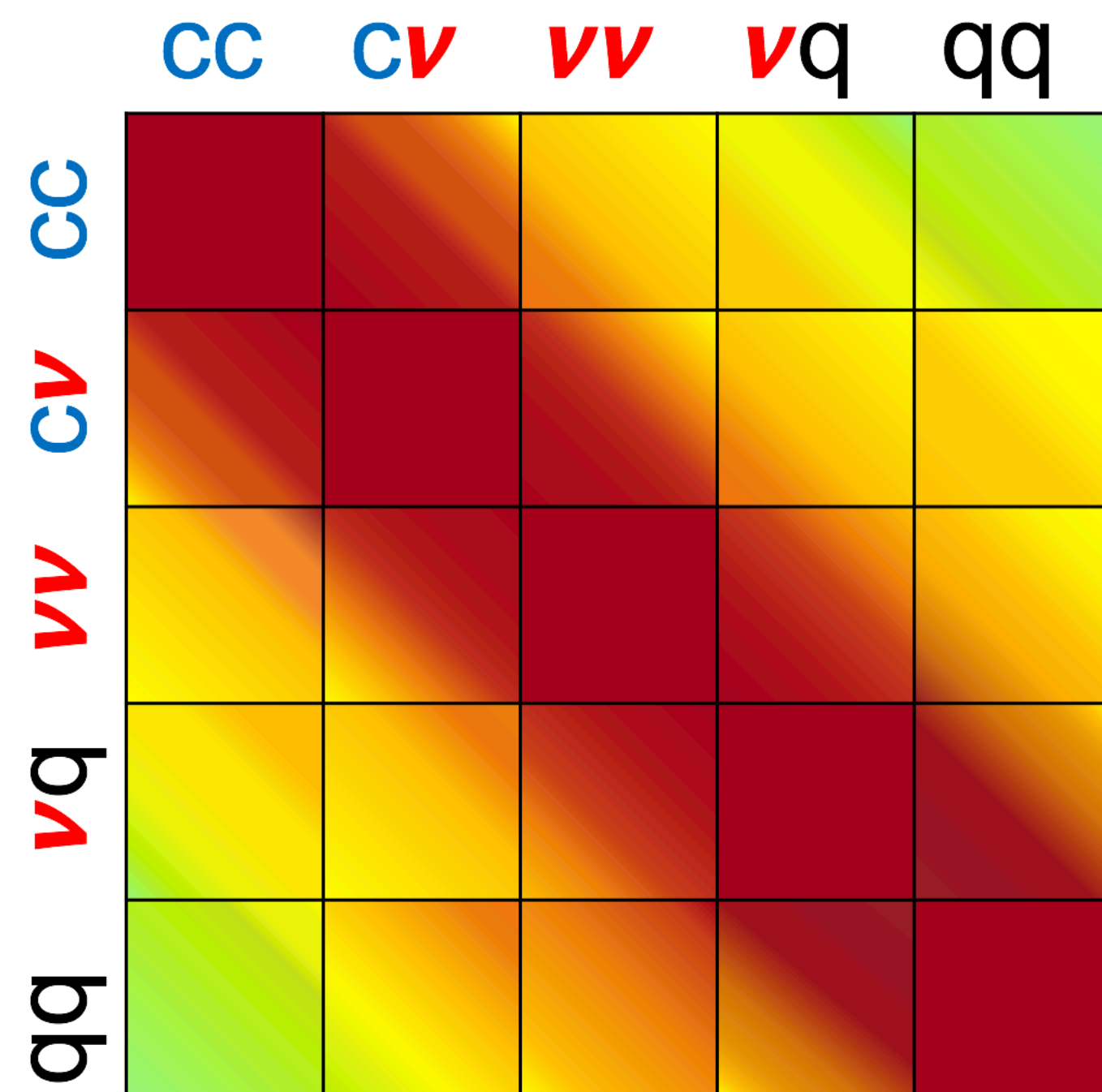
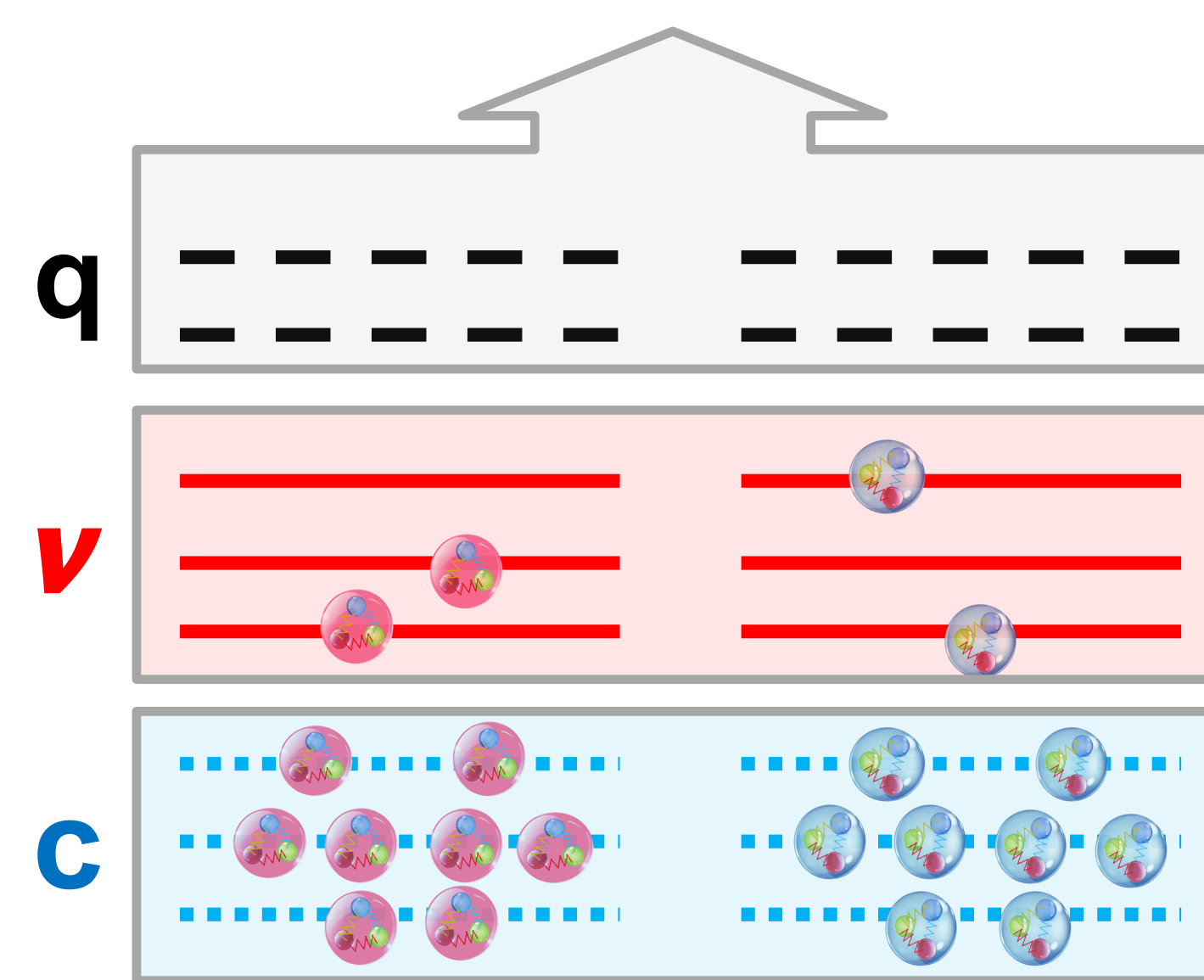
# In-Medium SRG (IM-SRG)

$H$  is normal ordered with a finite-density reference state:

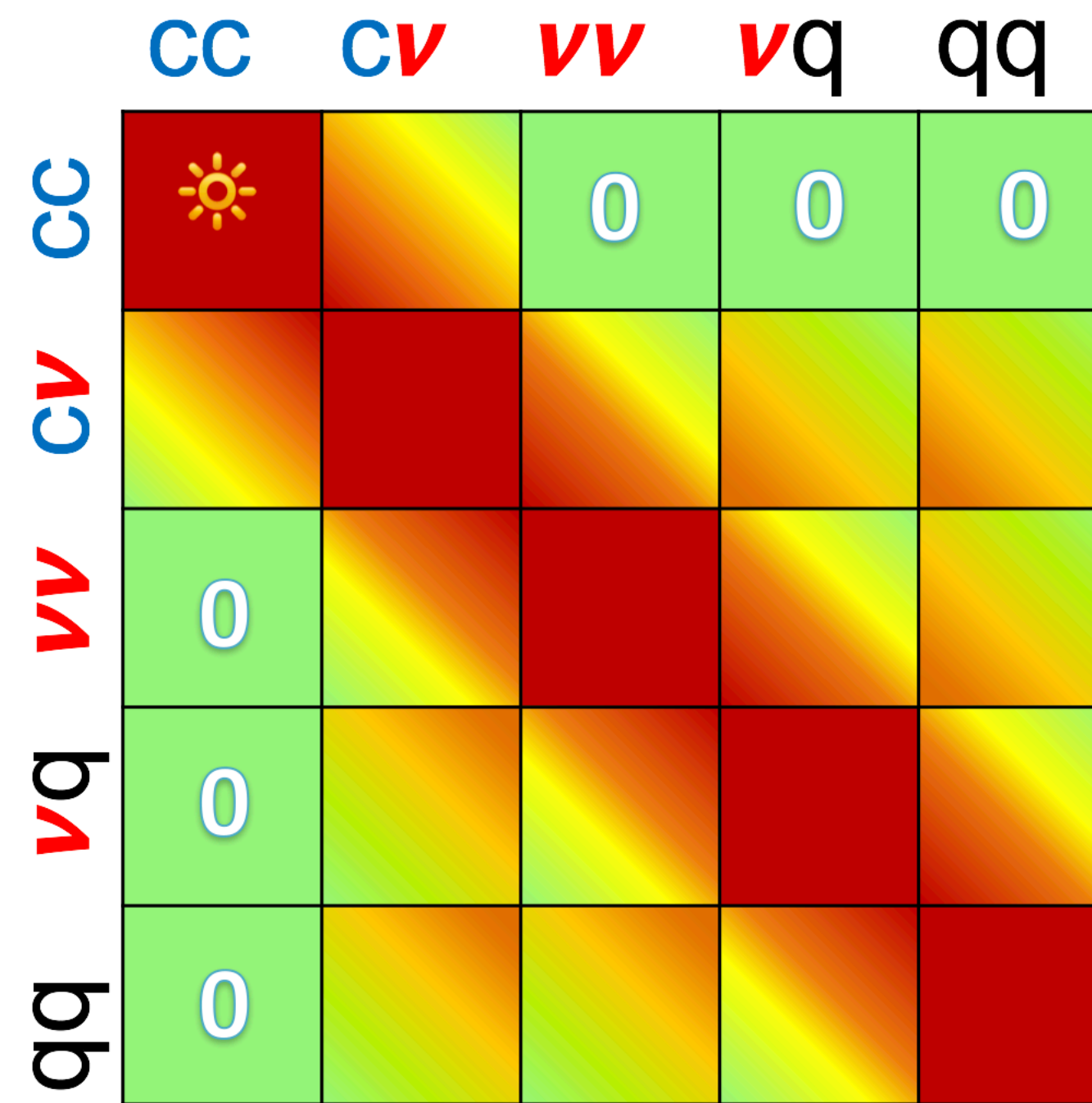
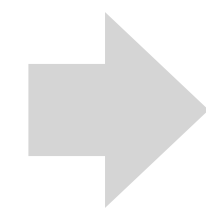
$$H = E_0 + \sum_{ij} f_{ij} \{a_i^\dagger a_j\} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}$$

Generator:

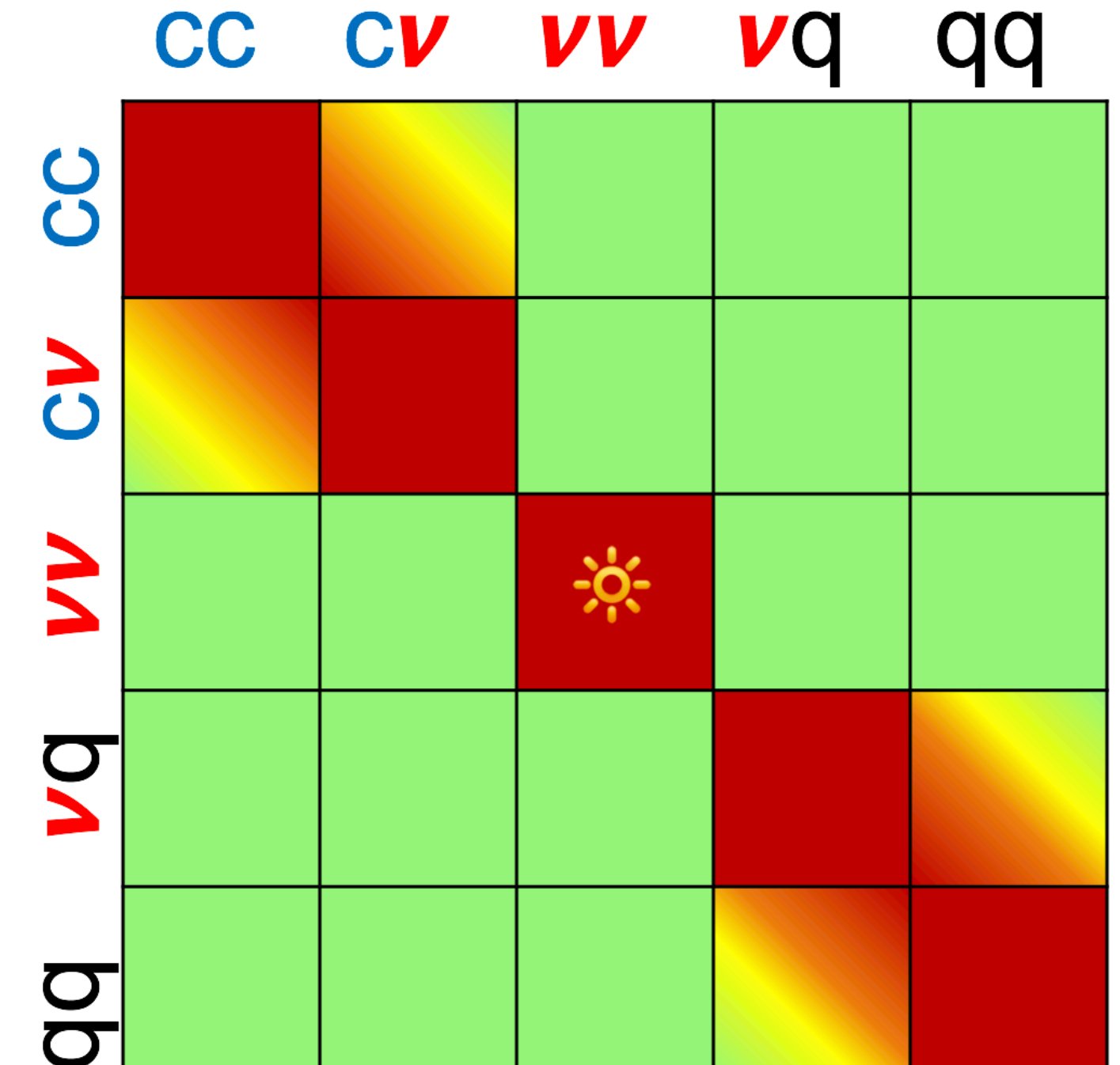
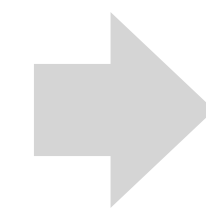
e.g., 
$$\eta(s) \equiv \sum_{ph} \frac{f_{ph}(s)}{\Delta_{ph}(s)} \{a_p^\dagger a_h\} + \sum_{pp'hh'} \frac{\Gamma_{pp'hh'}(s)}{\Delta_{pp'hh'}(s)} \{a_p^\dagger a_{p'}^\dagger a_h a_{h'}\} - H.c.$$



$\langle ij | H(s=0) | kl \rangle$

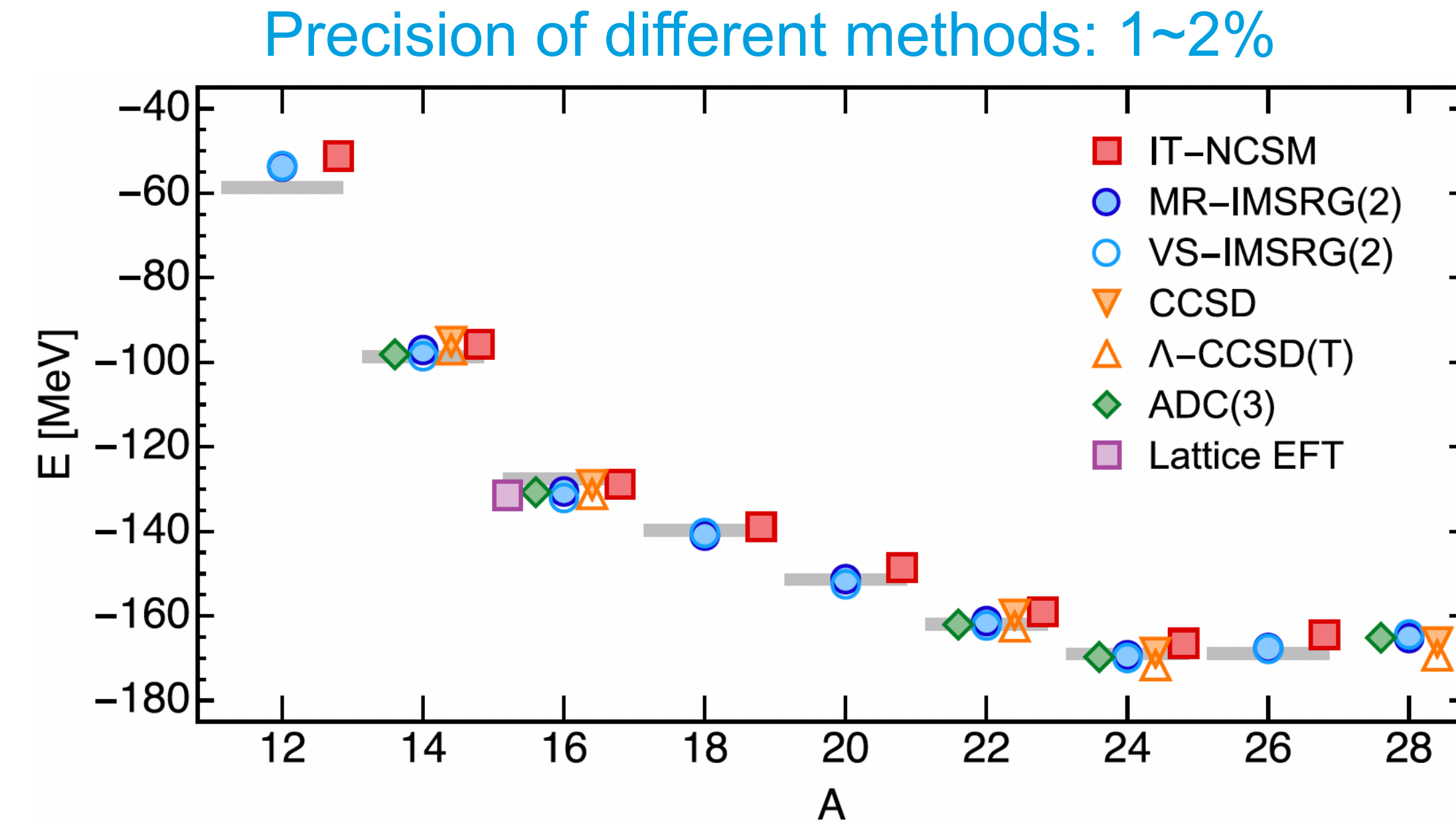
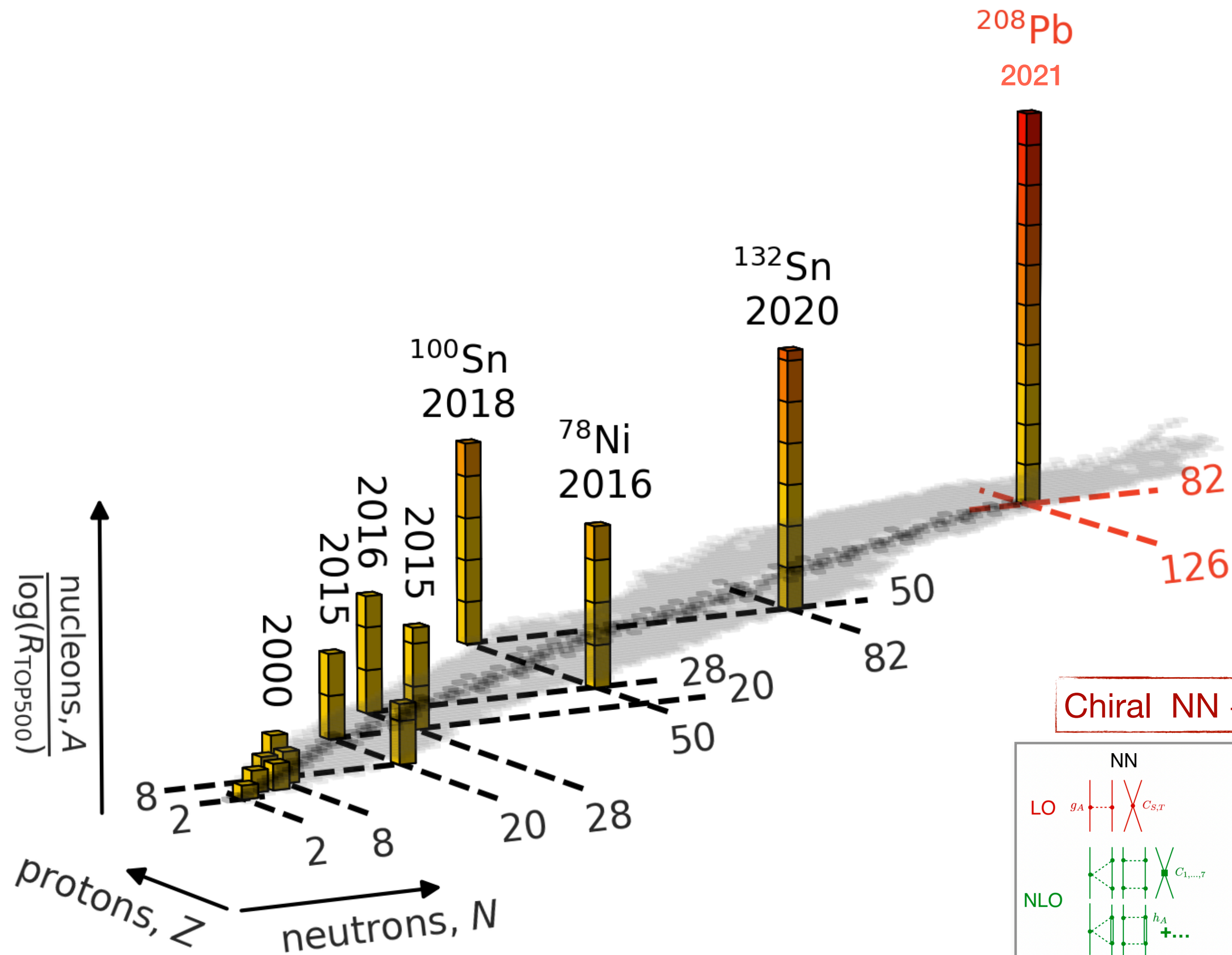


$\langle ij | H(s) | kl \rangle$



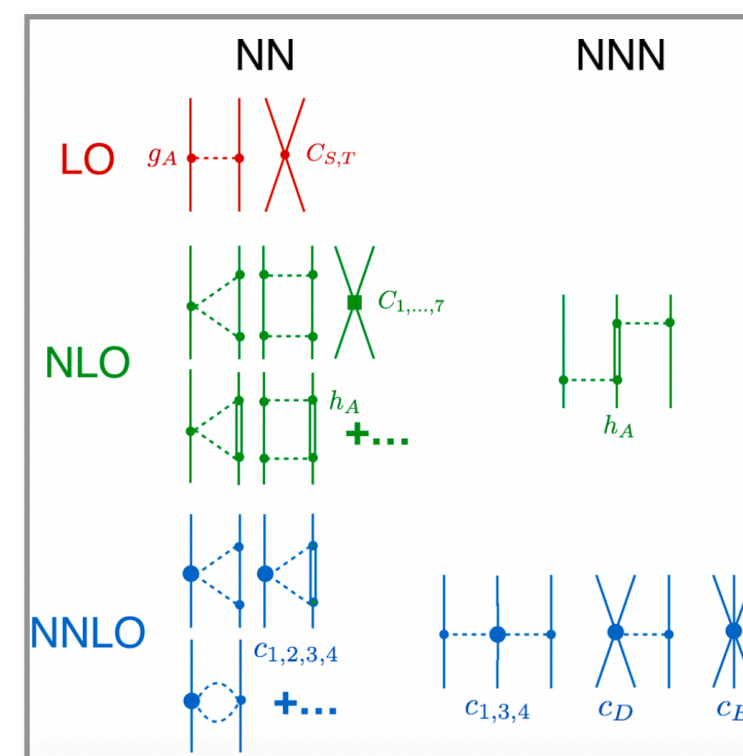
$\langle ij | H(s) | kl \rangle$

# Ab initio results for $^{208}\text{Pb}$ region



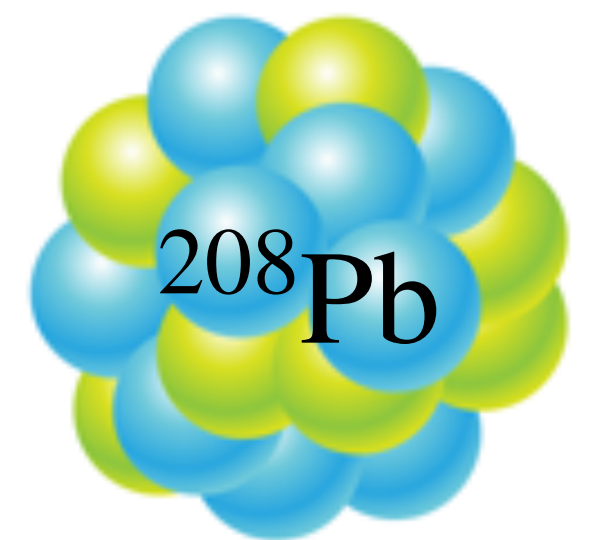
H. Hergert, Front. Phys. 8 (2020) 379

## Chiral NN + 3N

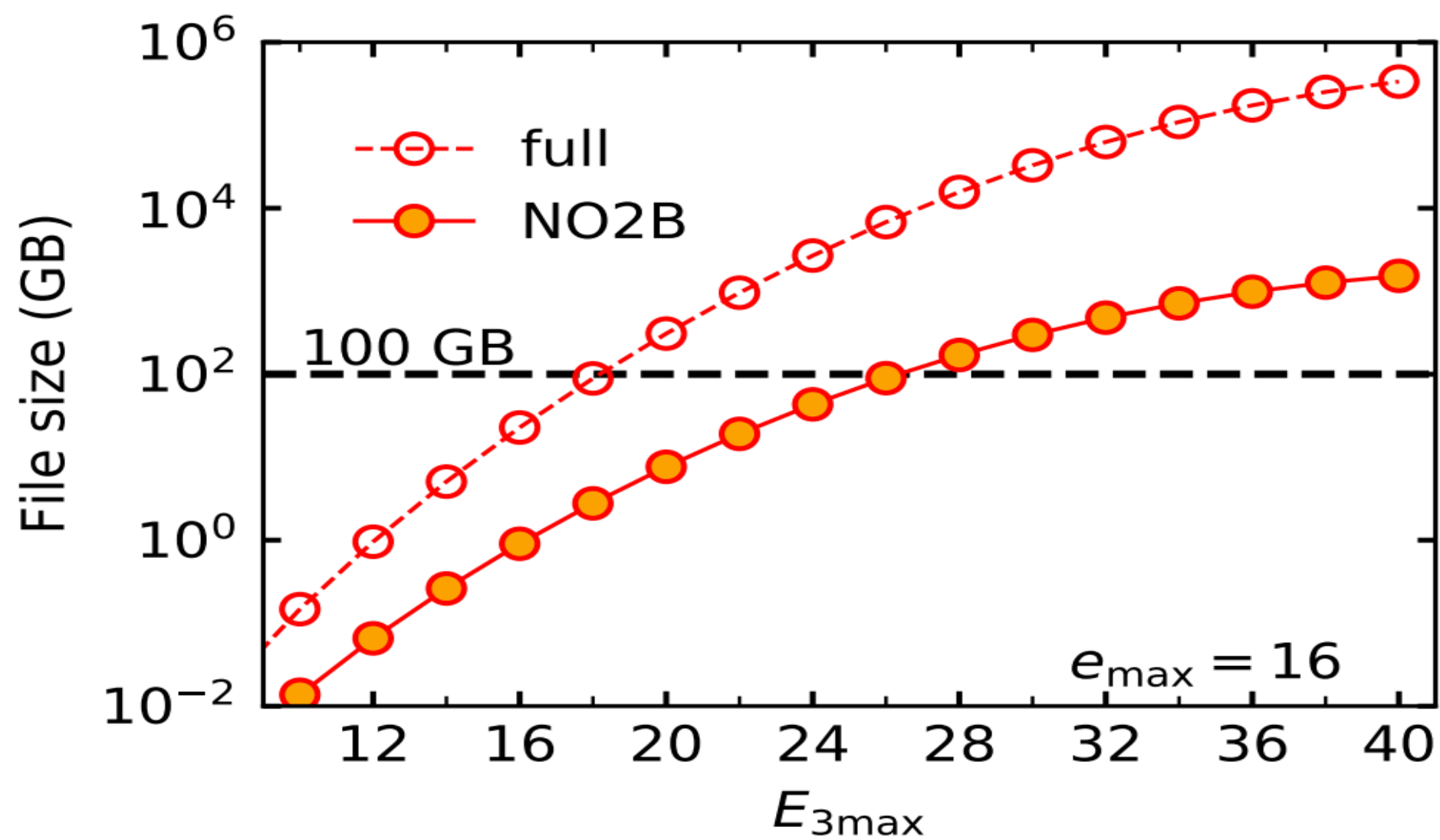
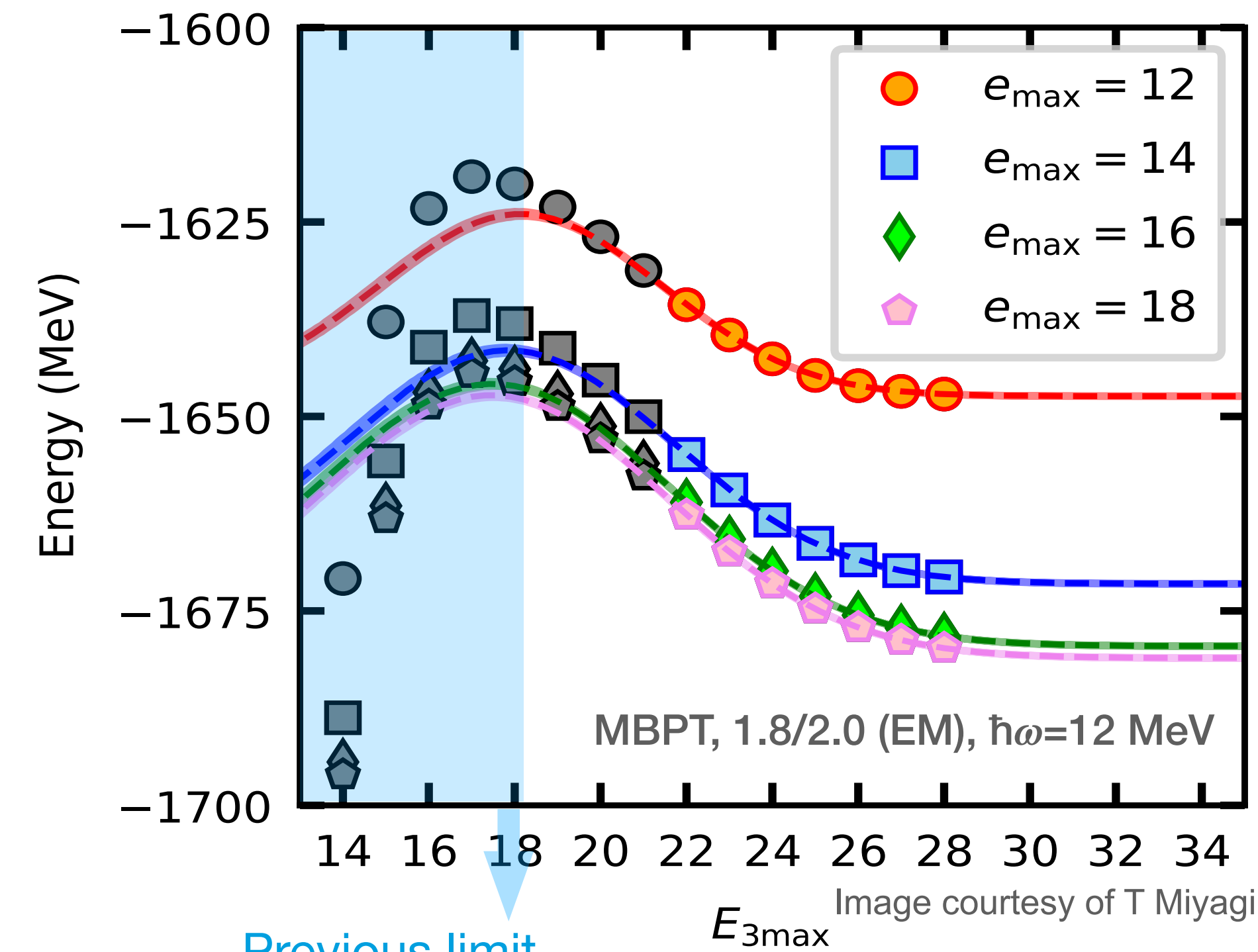


History matching  
 Emulator technology  
 Statistical tools

New 3N storage scheme  
 IMSRG, CC

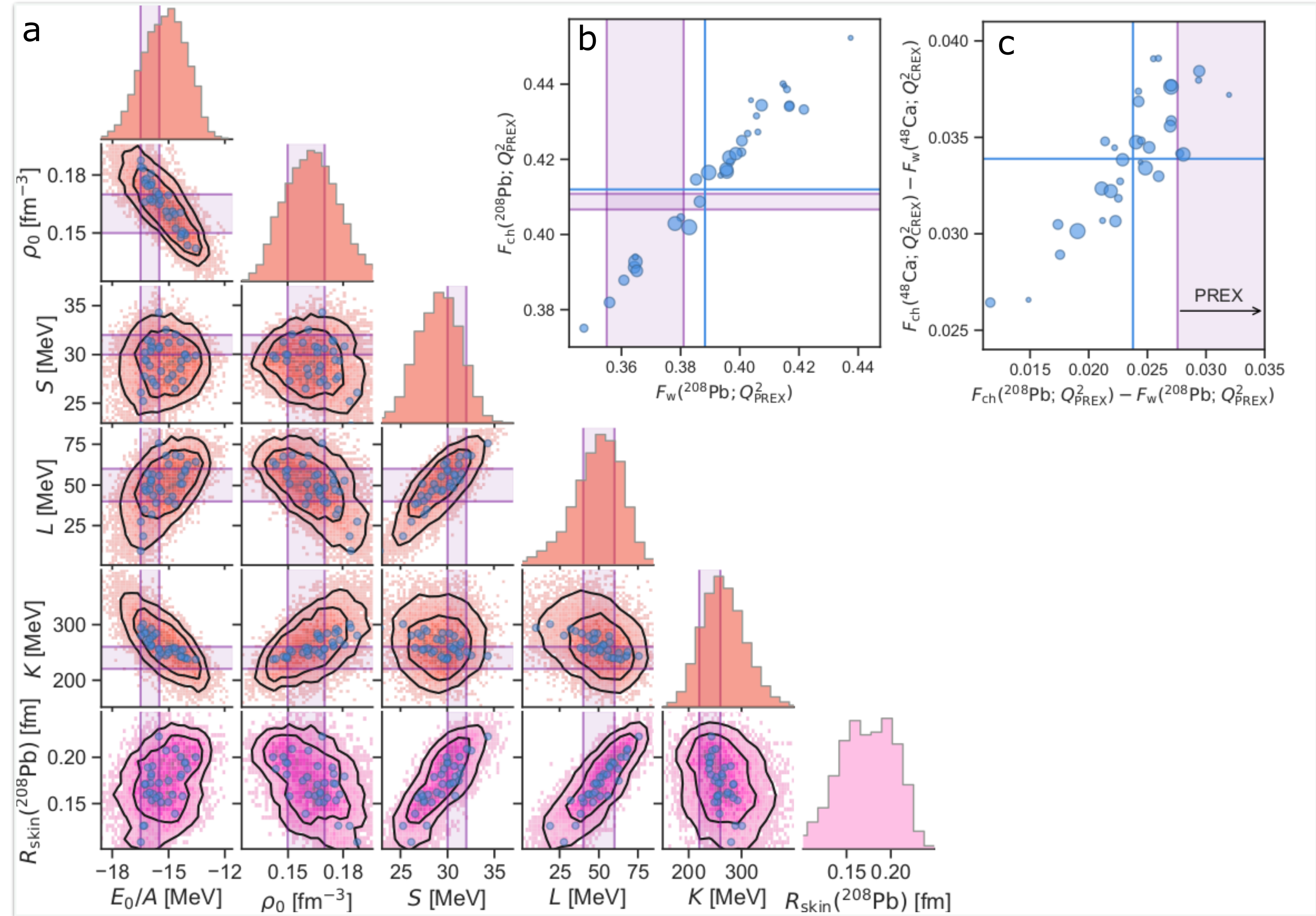


# Ab initio calculation of $^{208}\text{Pb}$



T Miyagi, *et al.*, PRC105 (2022) 014302

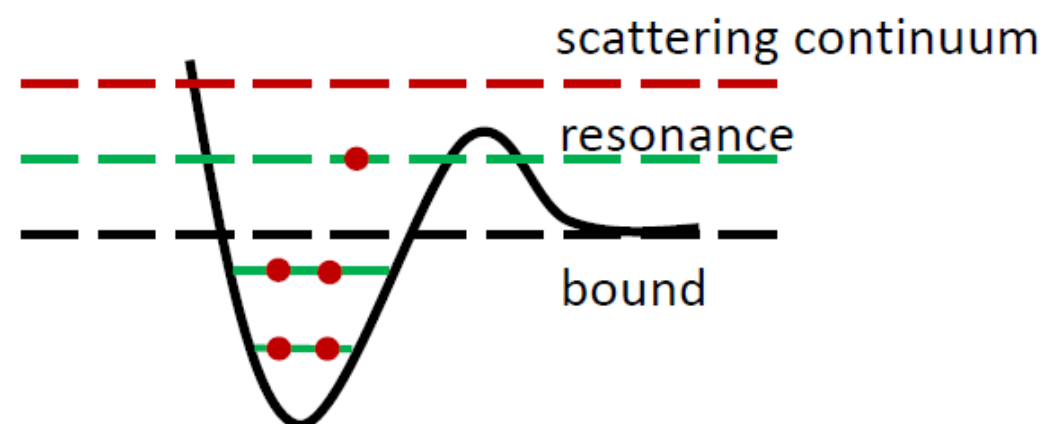
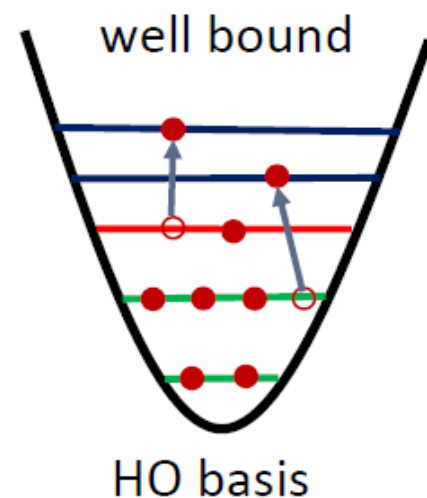
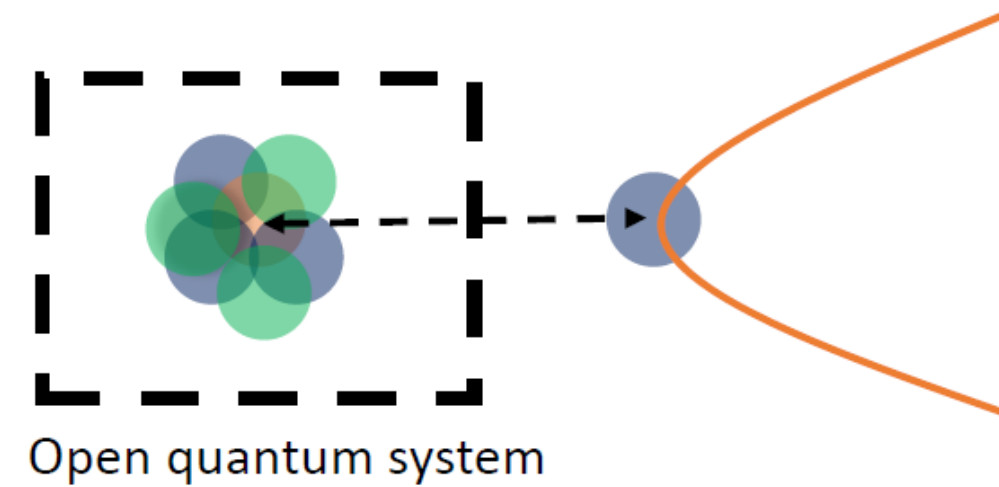
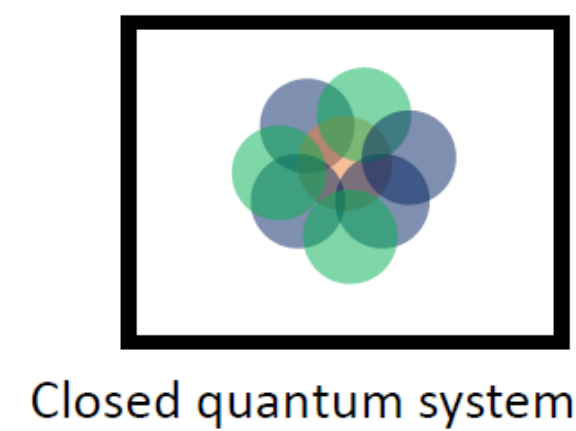
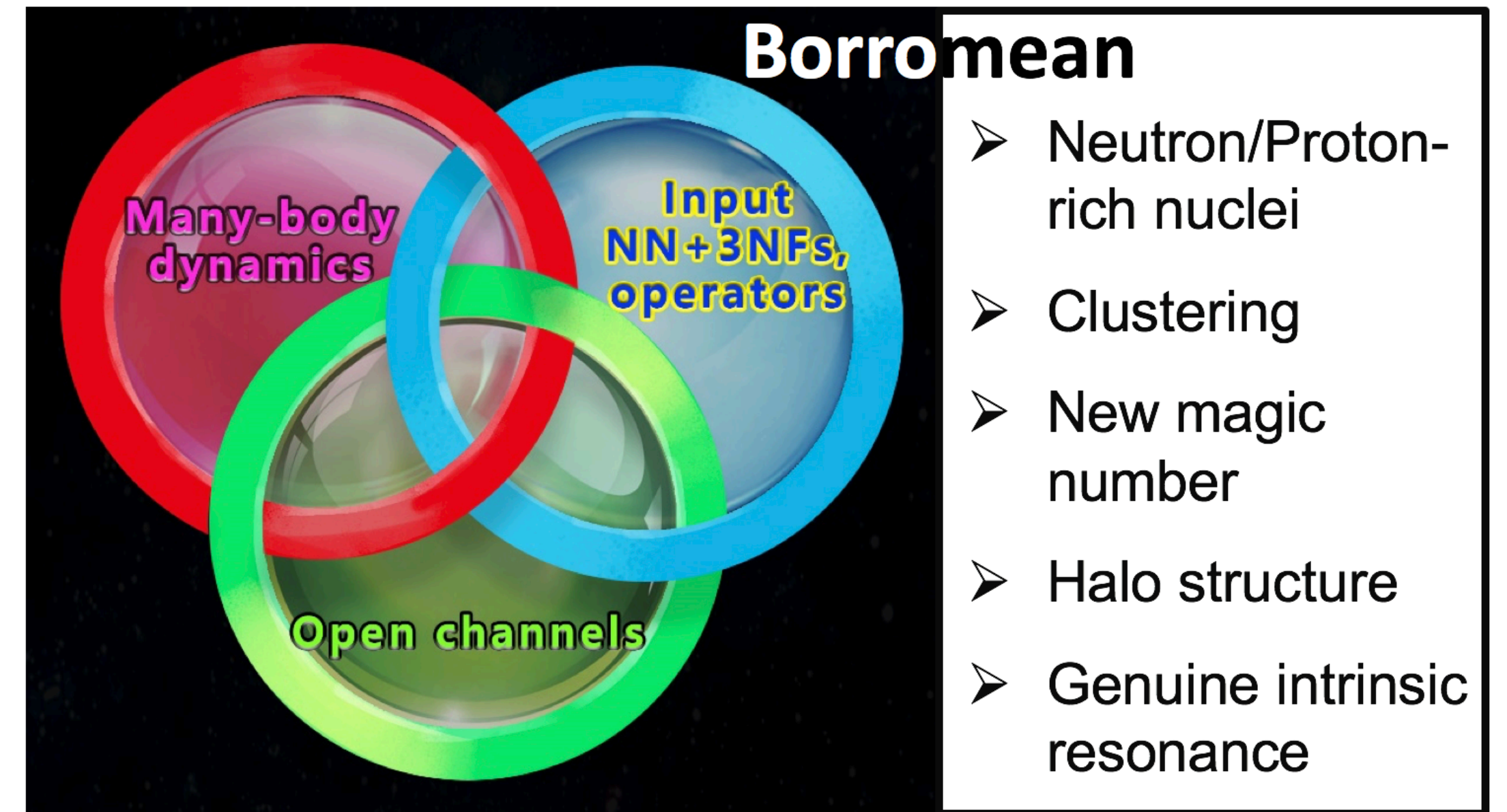
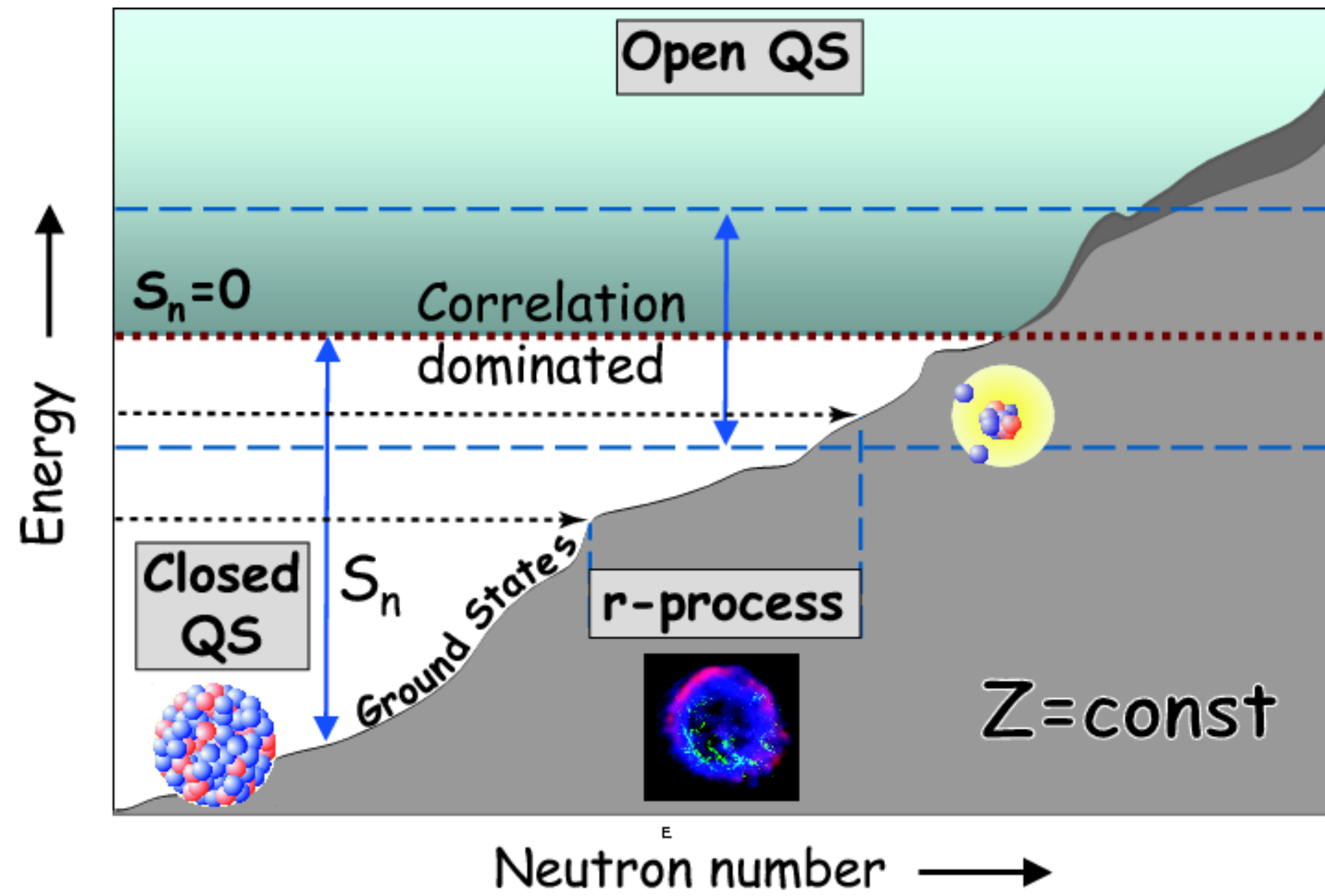
Baishan Hu - TRIUMF (2022/8/9)



BS Hu, WG Jiang, T Miyagi, ZH Sun, *et al.*, Nat. Phys. (2022), in press.  
 arXiv:2112.01125v1 (2021)



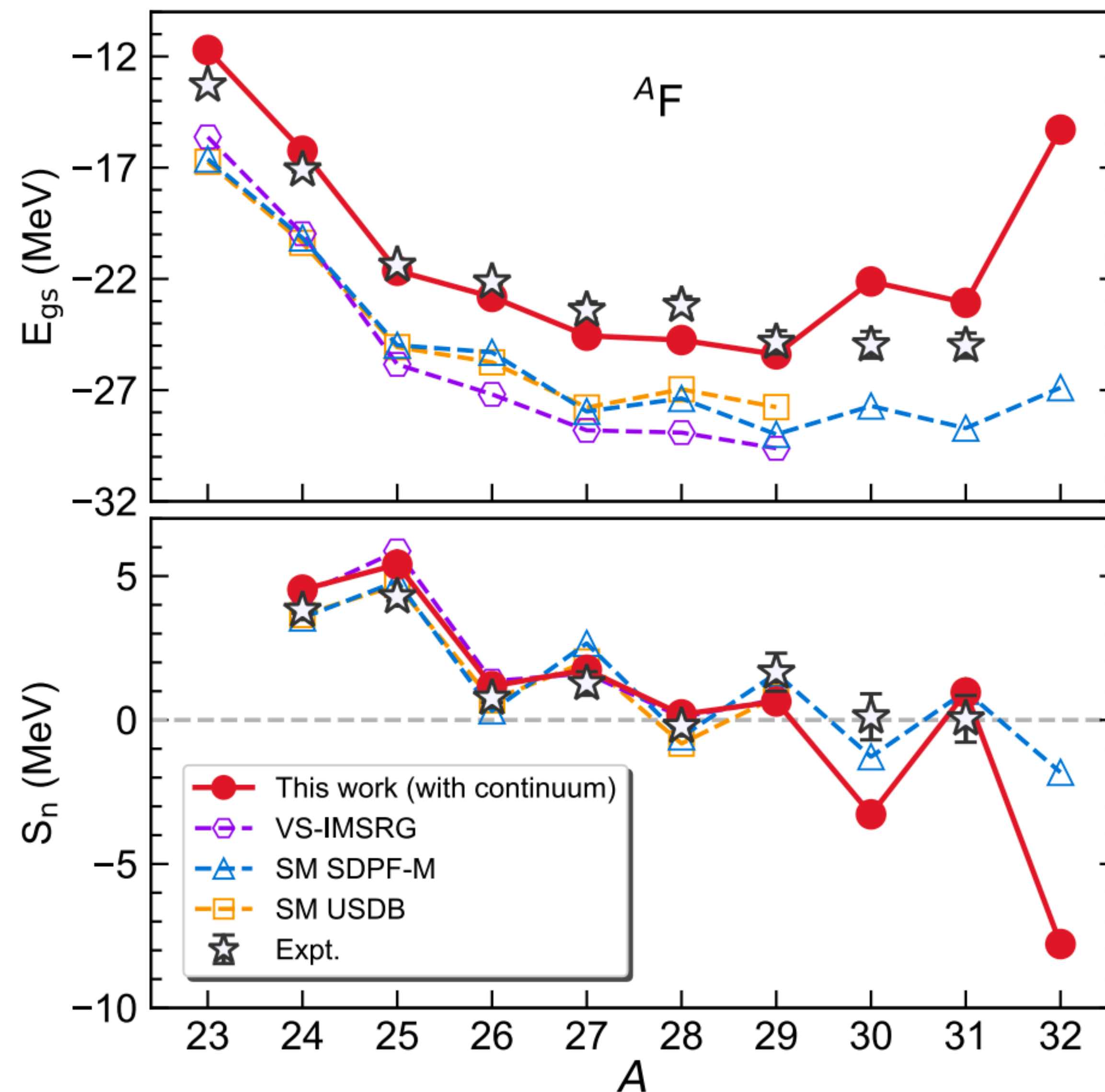
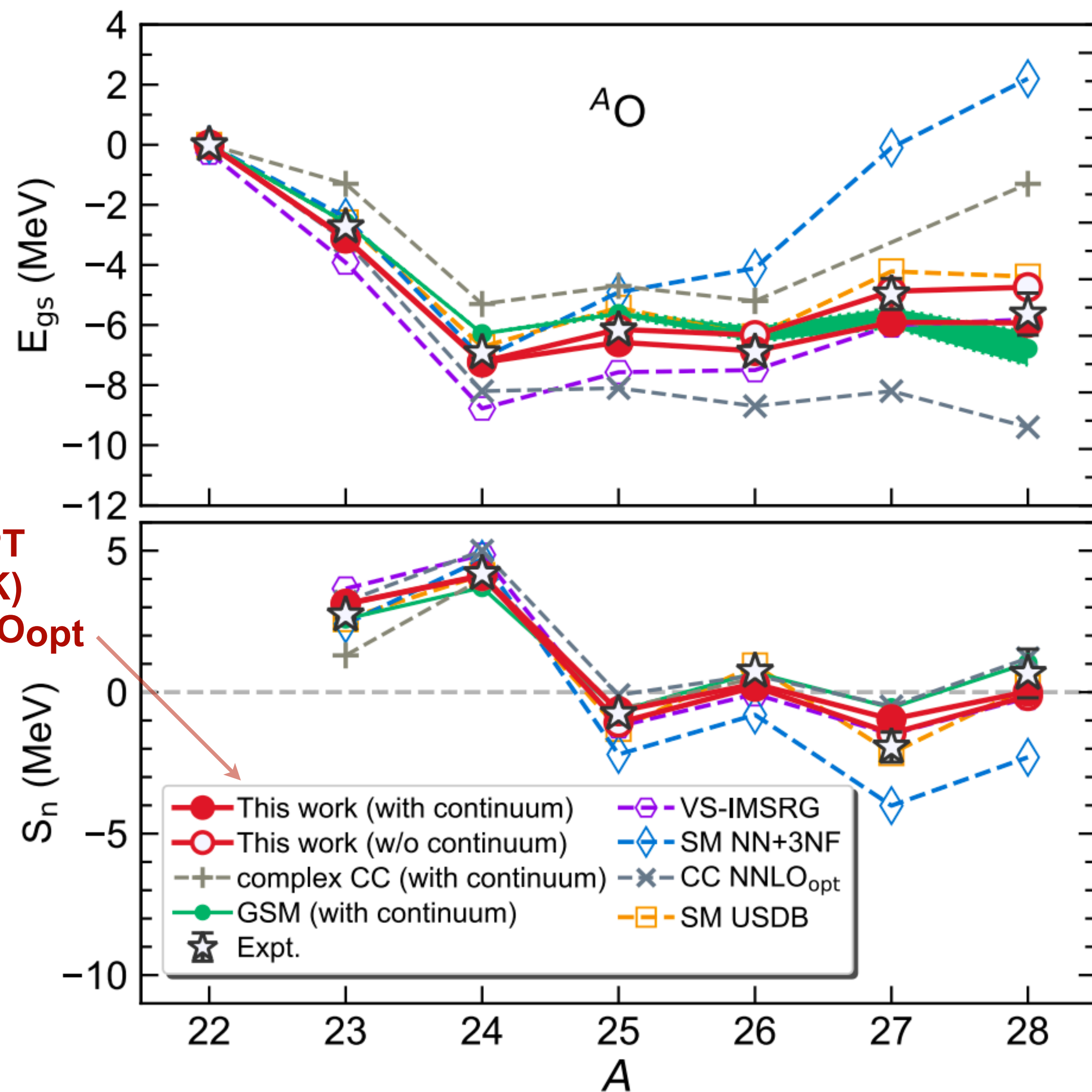
# Nucleus as an open quantum system



N. Michel *et al.*, JPG: NPP36 (2009) 013101

- **Bound, resonant and scattering states may be strongly coupled**
- **Need *ab initio* nuclear theory including the continuum**

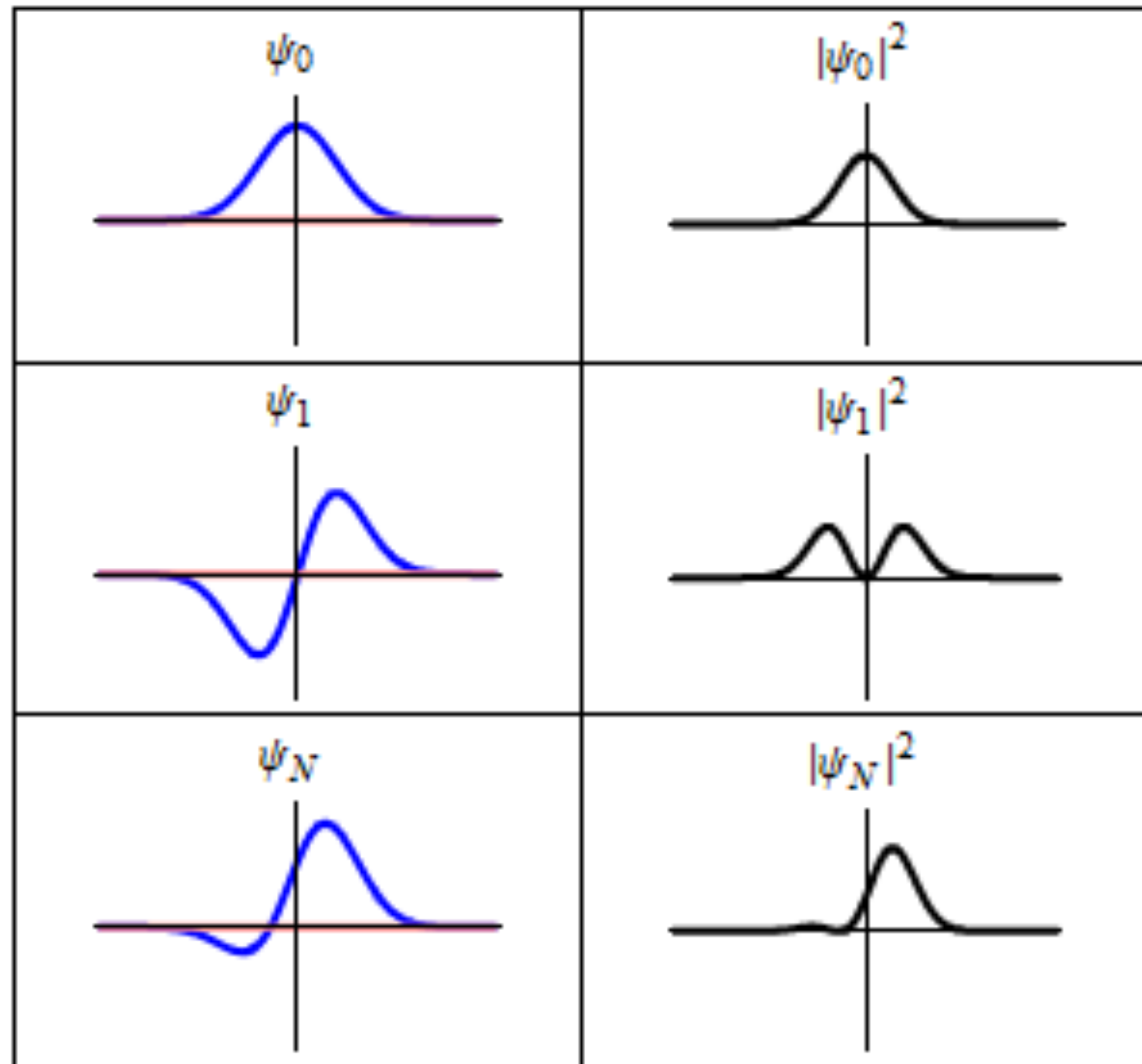
# Why consider the continuum effect?



BShu, Q. Wu, et al., PLB **802** (2020) 135206; arXiv:2001.02832

# Resonances

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = H |\psi\rangle$$



From Wikipedia

Time-independent:

$$H |\psi\rangle = E |\psi\rangle \quad \psi(t, \mathbf{r}) = \exp\left(-\frac{iE}{\hbar}t\right) \psi(0, \mathbf{r})$$

probability at  $\mathbf{r}$  is unchanged over time

Complex energy:  $E = E_0 - i\frac{\Gamma}{2}$

$$\begin{aligned} |\psi(t, \mathbf{r})|^2 &= \left| \exp\left(-\frac{iE_0}{\hbar}t\right) \exp\left(-\frac{\Gamma}{2\hbar}t\right) \psi(0, \mathbf{r}) \right|^2 \\ &= \exp\left(-\frac{\Gamma}{\hbar}t\right) |\psi(0, \mathbf{r})|^2 \end{aligned}$$

Describe a resonance decaying exponentially  
with half-life  $t_{1/2} = \hbar \ln 2 / \Gamma$

# Gamow-Berggren basis

- The wave function of a resonance with a peak at energy  $e_0$  and a width  $\gamma$

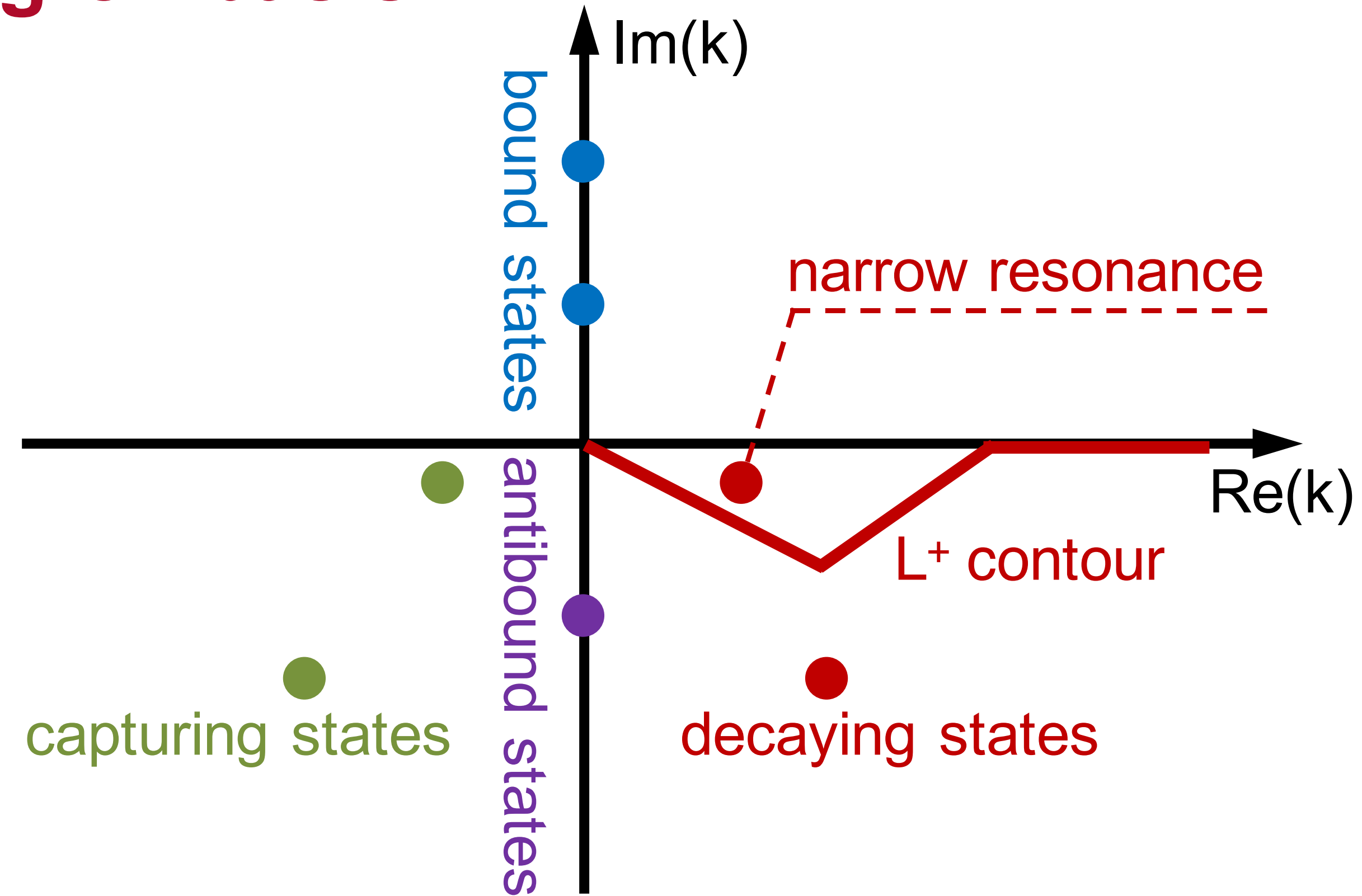
$$\Phi(e, \mathbf{r}) = \sqrt{\frac{\gamma/2}{\pi \left[ (e - e_0)^2 + (\gamma/2)^2 \right]}} \Psi(\mathbf{r})$$

- Through Fourier transformation, we obtain time evolution of the resonance

$$\Phi(t, \mathbf{r}) = \Psi(\mathbf{r}) e^{-i\tilde{e}t/\hbar} \quad \boxed{\tilde{e}_n = \frac{\hbar^2 k_n^2}{2m} = e_n - i\frac{\gamma_n}{2}} \quad t_{1/2} = \frac{\hbar \ln 2}{\gamma}$$

**Gamow state: complex energy**

G. Gamow, Z. Phys. 51 (1928) 204



**Berggren basis in complex- $k$  plane, describing bound, resonance and scattering on equal footing**

T. Berggren, Nucl. Phys. A109 (1968) 265

**Orthogonality and completeness:**

$$\delta(r - r') = \sum_n u_n(\tilde{e}_n, r) u_n(\tilde{e}_n, r') + \int_{L^+} d\tilde{e} u(\tilde{e}, r) u(\tilde{e}, r')$$

bound, resonance
scattering (discretized)

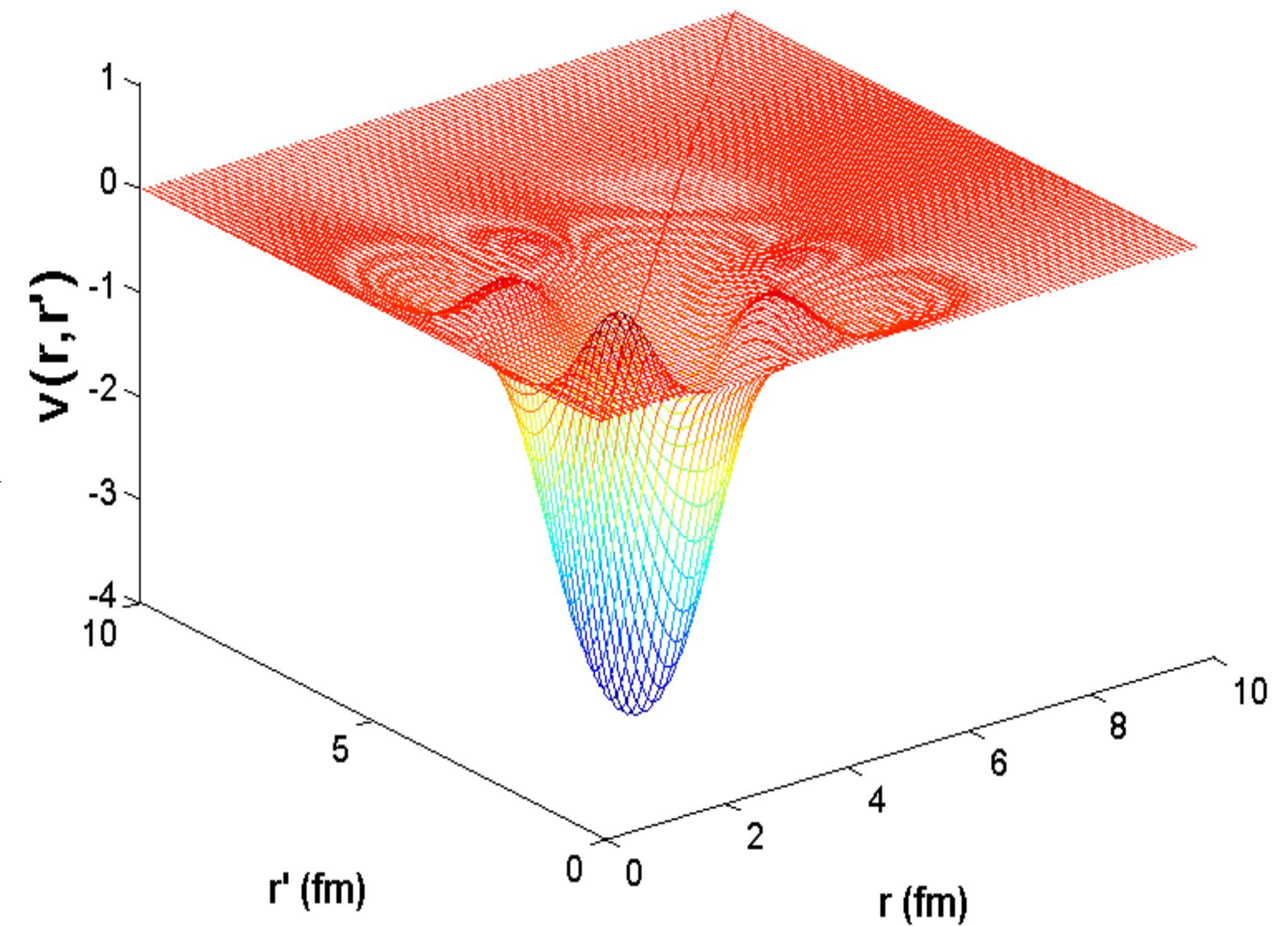
# Gamow Hartree-Fock

**Step 1:** Solve the HF equations in HO representation

$$H_{\text{int}} = \sum_{i=1}^A \left(1 - \frac{1}{A}\right) \frac{\vec{p}_i^2}{2m} + \sum_{i<j}^A \left( V_{NN,ij} - \frac{\vec{p}_i \cdot \vec{p}_j}{mA} \right) + \sum_{i<j<k}^A V_{NNN,ijk}$$

**Step 2:** Extract the non-local HF potential  $v(r,r')$

$$h_{ij}^{\text{HF}} = \langle i | t | j \rangle + \langle i | v | j \rangle = \langle i | t | j \rangle + \sum_{k=1}^A \langle ik | V | jk \rangle$$



**Step 3:** Obtain the radial wave function  $u(r)/r$  in complex- $k$  plane

$$u''(r) = \left[ \frac{l(l+1)}{r^2} + v^{(\text{loc})}(r) - k^2 \right] u(r) + \int_0^{+\infty} v^{(\text{non-loc})}(r, r') u(r') dr'$$

$$\begin{cases} u(\tilde{e}, r) \sim C_0 r^{l+1} & r \rightarrow 0 \\ u(\tilde{e}, r) \sim C^+ H_{l\eta}^+(kr) + C^- H_{l\eta}^-(kr) & r \rightarrow +\infty \end{cases}$$

$$u_n(\tilde{e}_n, r) \sim O_l(k_n r) \sim e^{ik_n r}$$

**Resonance:  
outgoing**

$$\tilde{e}_n = \frac{\hbar^2 k_n^2}{2m} = e_n - i \frac{\gamma_n}{2}$$

G. Gamow, Z. Phys. 51 (1928) 204

**exterior complex scaling**

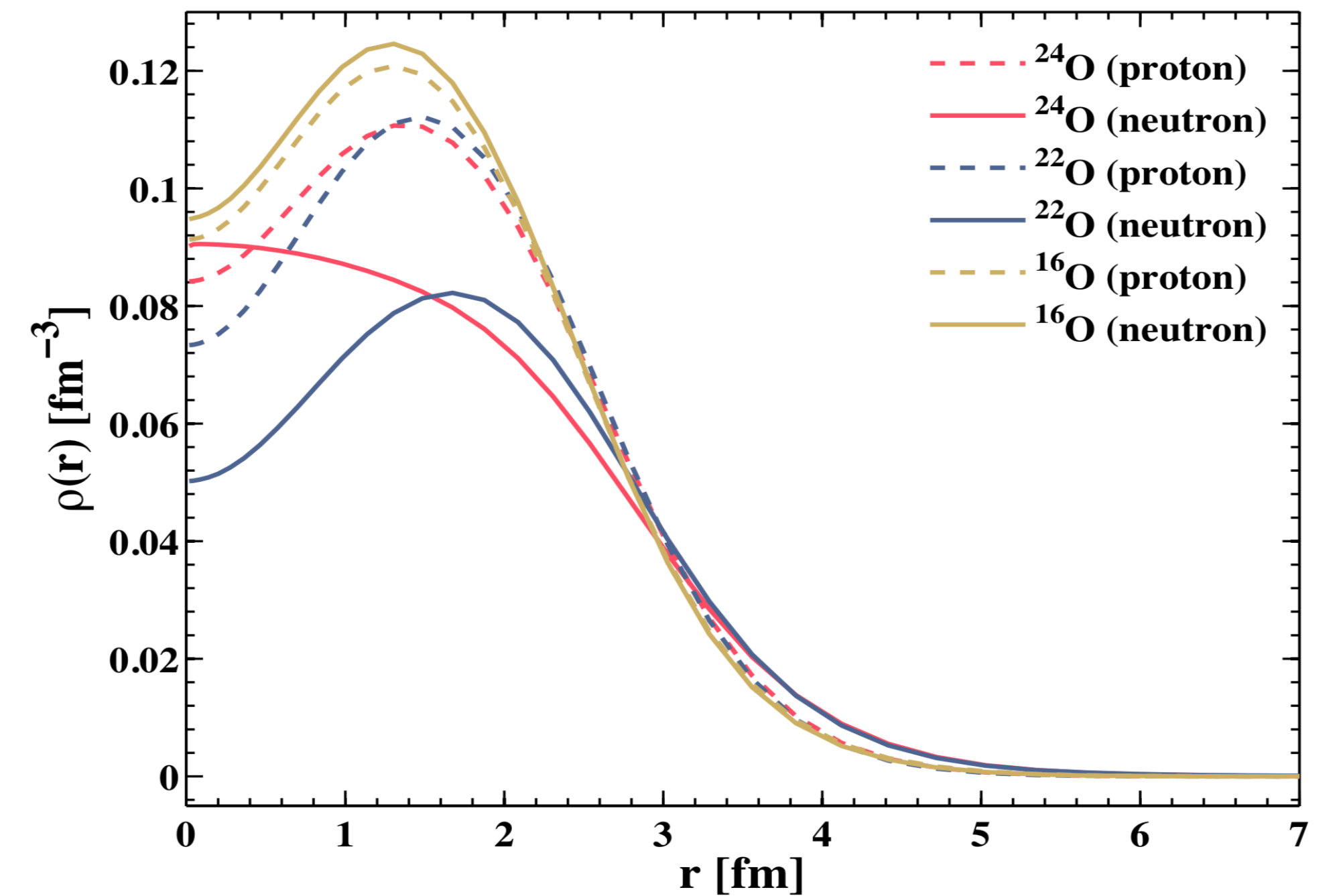
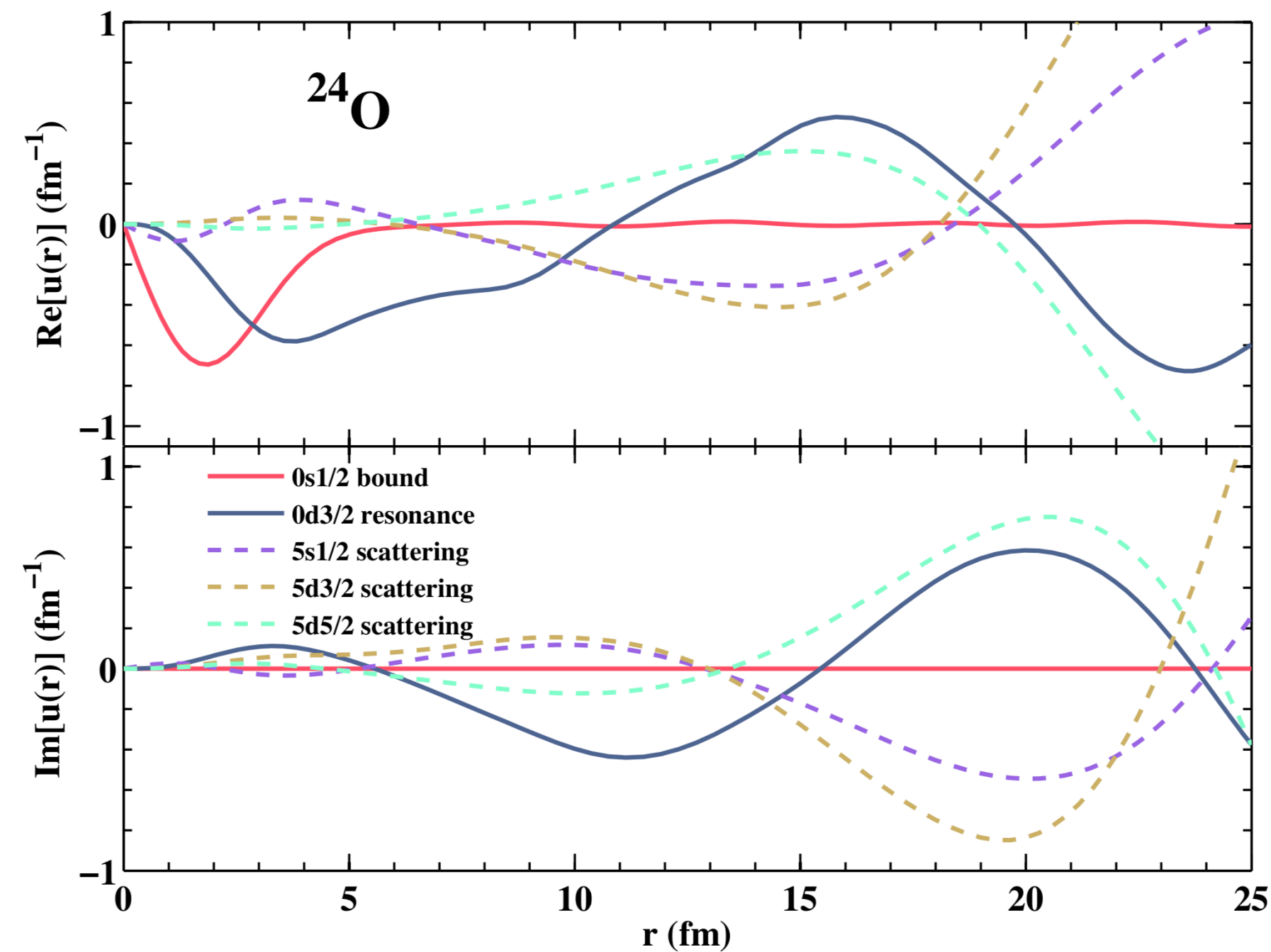
$$\int_0^{+\infty} u(\tilde{e}, r)^2 dr = \int_0^R u(\tilde{e}, r)^2 dr + (C^+)^2 \int_R^{+\infty} H_{l\eta}^+(kr)^2 dr$$

$$= \int_0^R u(\tilde{e}, r)^2 dr + (C^+)^2 \int_0^{+\infty} H_{l\eta}^+(kR + kxe^{i\theta})^2 e^{i\theta} dx$$

# Results of GHF

sp energies	$^{16}\text{O}$		$^{22}\text{O}$		$^{24}\text{O}$		$^{28}\text{O}$		MeV
	Re(E)	Im(E)	Re(E)	Im(E)	Re(E)	Im(E)	Re(E)	Im(E)	
$0s_{1/2}$	-48.858	0.000	-57.720	0.000	-59.313	0.000	-55.076	0.000	
$0p_{3/2}$	-22.735	0.000	-27.729	0.000	-28.132	0.000	-28.101	0.000	
$0p_{1/2}$	-13.863	0.000	-23.501	0.000	-22.669	0.000	-21.674	0.000	
$0d_{5/2}$	—	—	- 3.251	0.000	-3.993	0.000	-6.687	0.000	
$1s_{1/2}$	—	—	- 0.964	0.000	-2.374	0.000	-3.978	0.000	
$0d_{3/2}$	—	—	3.014	-0.626	<b>2.312</b>	<b>-0.368</b>	<b>1.088</b>	<b>-0.081</b>	

**sp resonance**







**NCSM**  
**CC**  
**GFMC**  
**IM-SRG**

**Included continuum coupling**



**NCSMC, SS-HORSE**  
**complex CC**  
**GFMC**  
**???**

# IM-SRG with resonance and continuum

-  **Include continuum degree of freedom into IM-SRG via Gamow-Berggren ensemble (Gamow Hartree-Fock)**
-  **Extend the continuous unitary transformation (Hermitian) to the orthogonal transformation (complex symmetric)**
-  **Describe bound, resonance and continuum states in a unified framework**
-  **Combine with resonating group method (RGM) for nuclear reactions**



# IMSRG framework

IMSRG Hermitian (HO/HF basis)		Gamow IMSRG Complex symmetric (Berggren basis)
$\langle a   H   b \rangle = \langle b   H   a \rangle^*$ $H(s) = U(s)H(0)U^\dagger(s)$ $U(s) \cdot U^\dagger(s) = U(s) \cdot U^{-1}(s) = 1$ $\eta(s) = \frac{dU(s)}{ds}U^\dagger(s) = -\eta^\dagger(s)$	$\frac{dH(s)}{ds} = [\eta(s), H(s)]$ $H(s) = U(s)HU^{-1}(s)$ $O(s) = U(s)OU^{-1}(s)$	$\langle \tilde{\alpha}   H   \beta \rangle = \langle \tilde{\beta}   H   \alpha \rangle^* = \langle \beta   H   \tilde{\alpha} \rangle$ $H(s) = U(s)H(0)U^T(s)$ $U(s) \cdot U^T(s) = U(s) \cdot U^{-1}(s) = 1$ $\eta(s) = \frac{dU(s)}{ds}U^T(s) = -\eta^T(s)$

## Magnus expansion:

T.D. Morris, N.M. Parzuchowski, and S.K. Bogner, PRC92 (2015) 034331

$$U(s) = e^{\Omega(s)}$$

$$\frac{d\Omega(s)}{ds} = \eta(s) + \frac{1}{2}[\Omega(s), \eta(s)] + \frac{1}{12}[\Omega(s), [\Omega(s), \eta(s)]] + \dots$$

$$H(s) = e^{\Omega(s)}He^{-\Omega(s)} = H + [\Omega(s), H] + \frac{1}{2}[\Omega(s), [\Omega(s), H]] + \dots$$

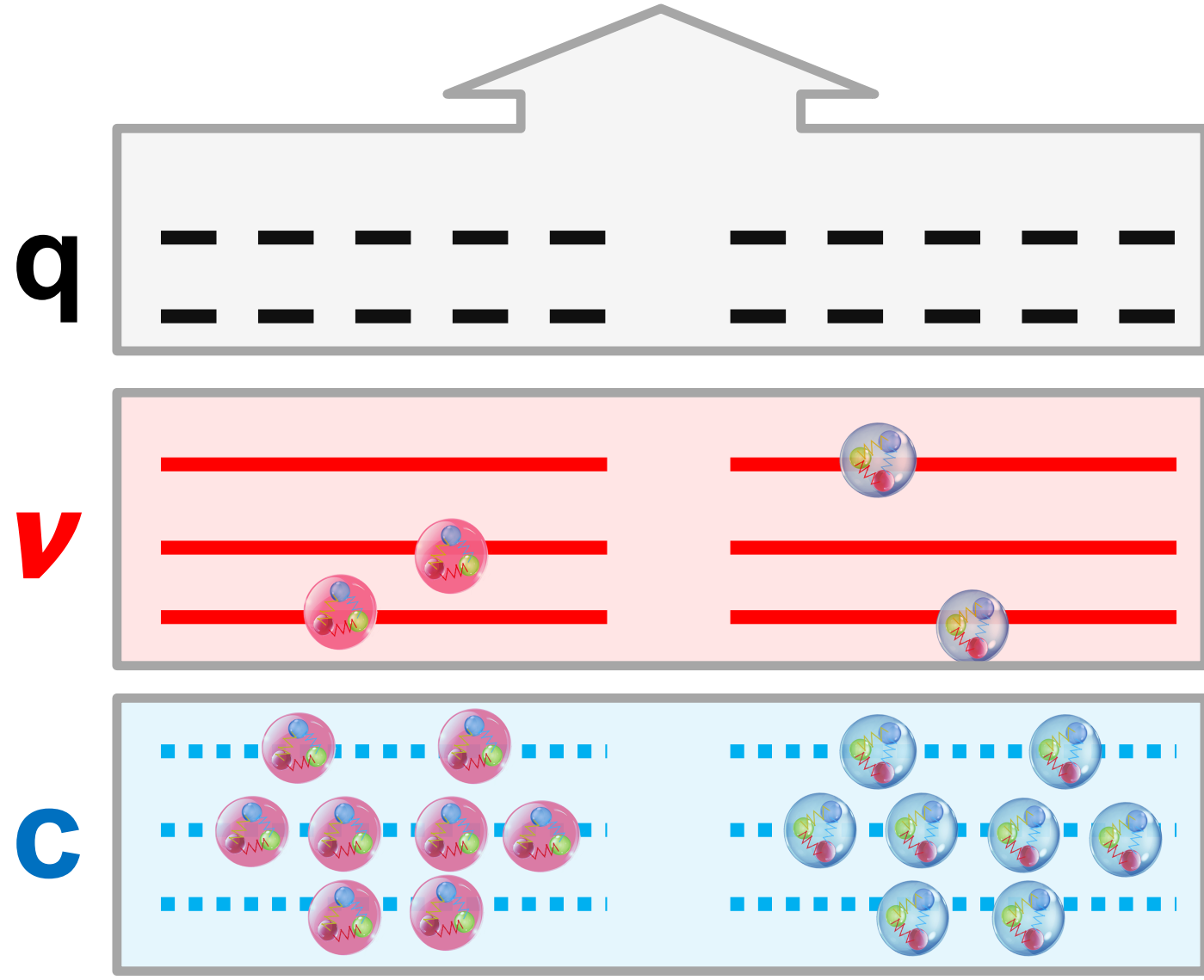
$$O(s) = e^{\Omega(s)}Oe^{-\Omega(s)} = O + [\Omega(s), O] + \frac{1}{2}[\Omega(s), [\Omega(s), O]] + \dots$$

## e.g., White generator $\eta$ :

$$\eta_{12} = \frac{f_{12}}{f_{11} - f_{22} + \Gamma_{1212}}$$

$$\eta_{1234} = \frac{\Gamma_{1234}}{f_{11} + f_{22} - f_{33} - f_{44} + A_{1234}}$$

$$A_{1234} = \Gamma_{1212} + \Gamma_{3434} - \Gamma_{1313} - \Gamma_{2424} - \Gamma_{1414} - \Gamma_{2323}$$



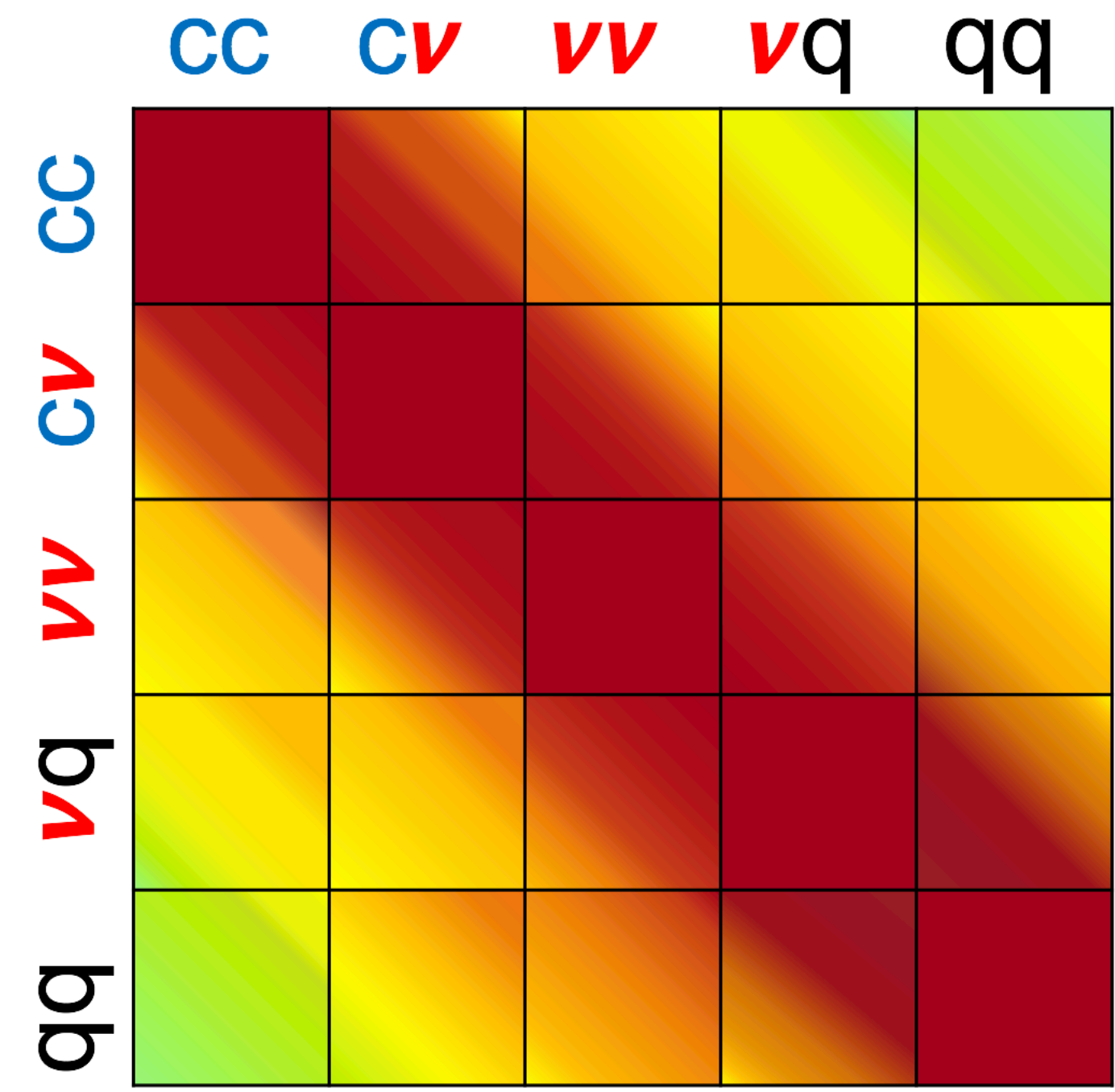
$$H = E_0 + \sum_{ij} f_{ij} \{ a_i^\dagger a_j \} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \{ a_i^\dagger a_j^\dagger a_l a_k \} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \{ a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l \}$$

### EOM-IMSRG closed-shell nuclei

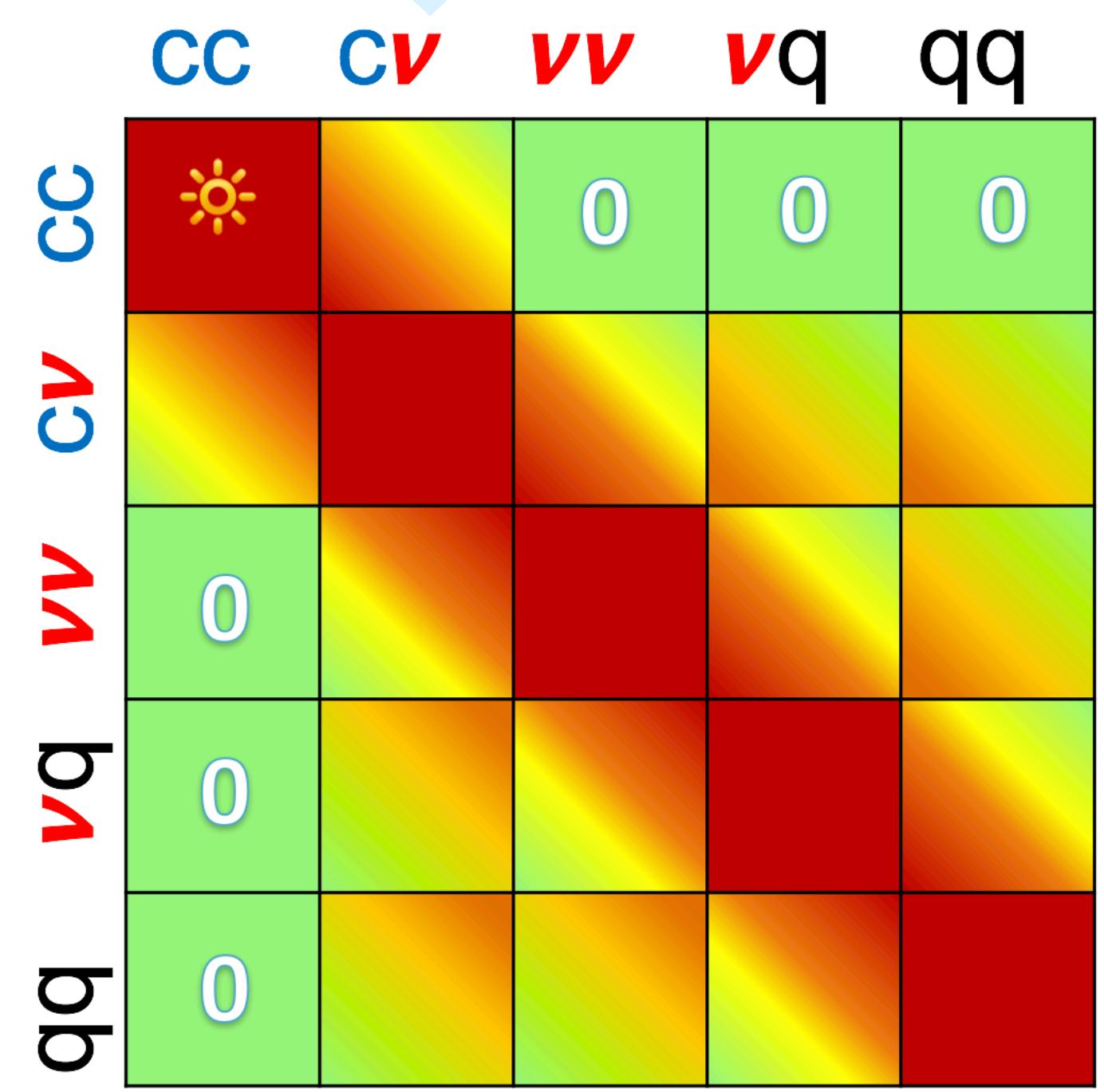
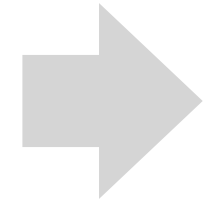
$$H^{\text{od}} \equiv \langle v | H | c \rangle + \langle q | H | c \rangle + \langle vv | H | cc \rangle + \langle qv | H | cc \rangle + \langle qq | H | cc \rangle$$

### VS-IMSRG shell-model effective interaction

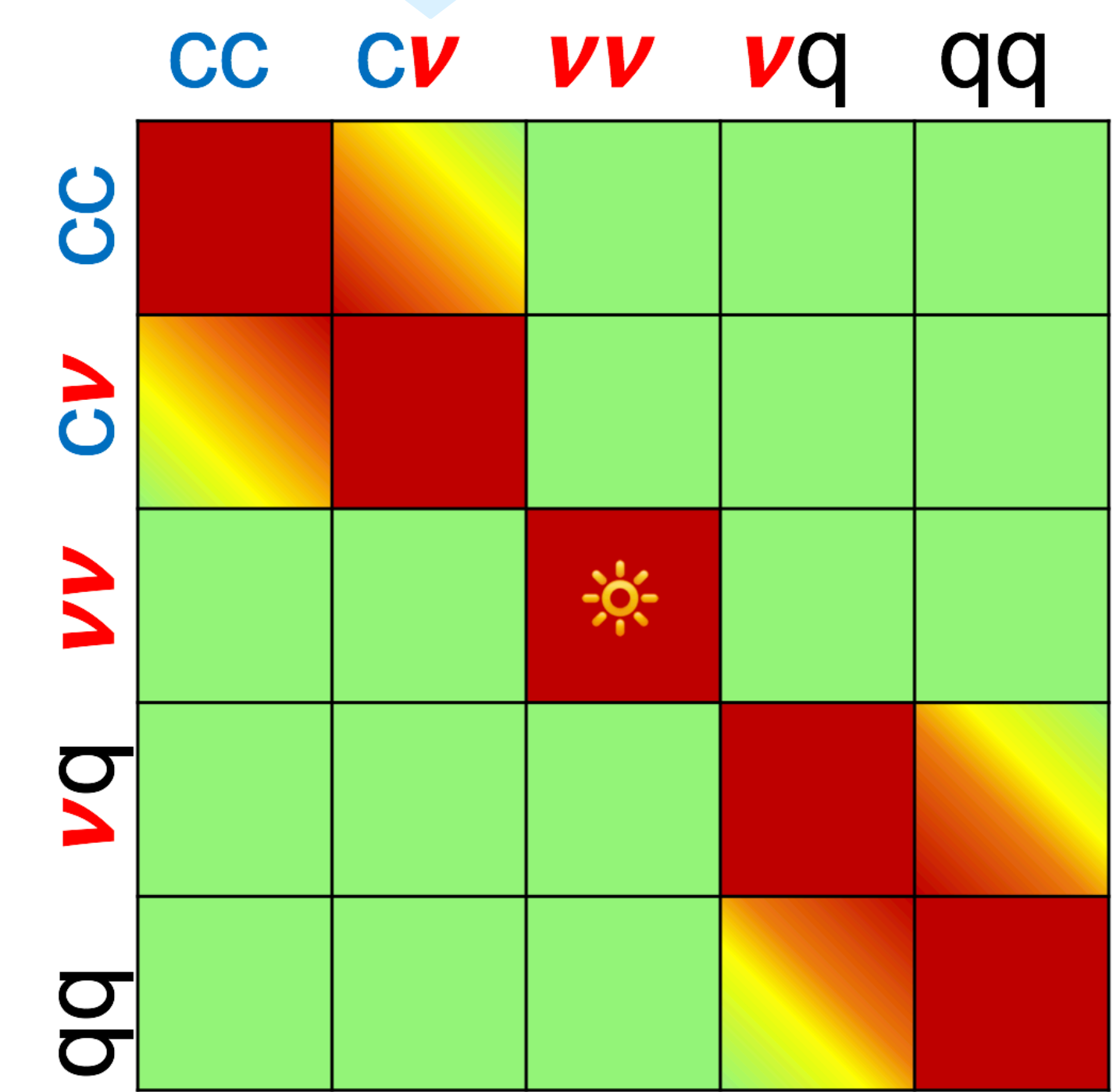
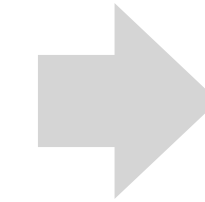
$$H^{\text{od}} \equiv \langle v | H | c \rangle + \langle q | H | c \rangle + \langle q | H | v \rangle + \langle vv | H | cc \rangle + \langle qv | H | cc \rangle + \langle qq | H | cc \rangle + \langle vv | H | vc \rangle + \langle qv | H | vc \rangle + \langle qq | H | vc \rangle + \langle qv | H | vv \rangle + \langle qq | H | vv \rangle$$



$\langle ij | H(s = 0) | kl \rangle$



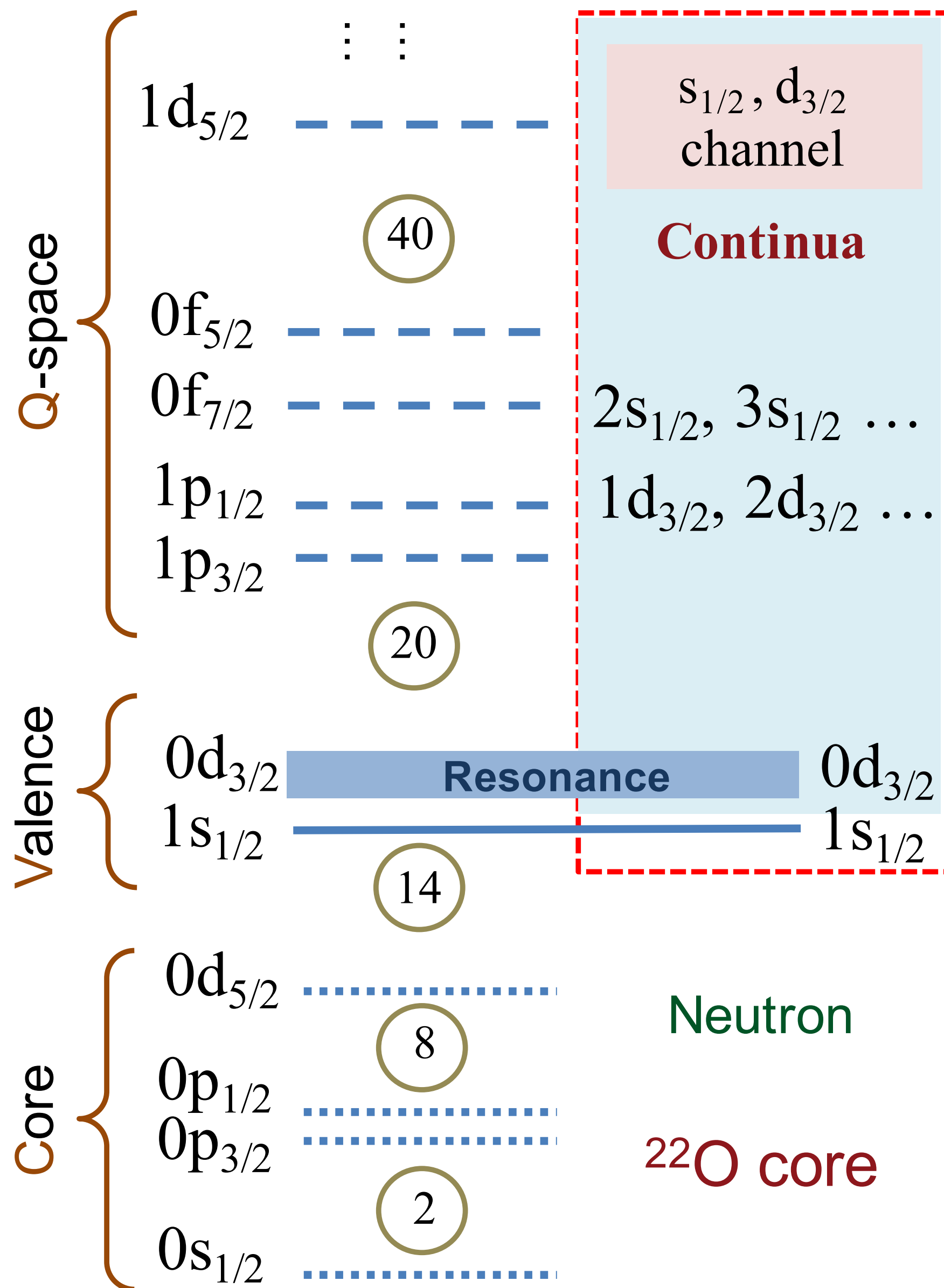
$\langle ij | H(s) | kl \rangle$



$\langle ij | H(s) | kl \rangle$

# Difficulties caused by continuum states

$$\langle vv | H^{od} | vq \rangle !!!$$



$$H(s) = e^{\eta(s)} H e^{-\eta(s)} = H + [\eta(s), H] + \frac{1}{2} [\eta(s), [\eta(s), H]] + \dots$$

$$\langle v | \eta^{\text{White}} | q \rangle = \frac{\langle v | H^{od} | q \rangle}{\langle v | H^d | v \rangle - \langle q | H^d | q \rangle}$$

**$\langle v | H^d | v \rangle$  close to  $\langle q | H^d | q \rangle$  : generates high-order terms in the transformation of H**

$$\theta = \frac{1}{2} \tan^{-1} [2a / (E_l - E_r)]. \quad (10)$$

magnitude. Note that the degenerate case  $E_l = E_r$  is nonsingular, generating an angle of  $\pm \pi/4$  (either angle can be chosen). Such a large transformation angle should be avoided if possible, however, since it generates high-order terms in the transformation of  $H$ .

S.R. White, J. Chem. Phys. 117 (2002) 7472

**IMSRG(3) or ...**

# Similar to many-multi-shells effective interaction

① 
$$\langle v | \eta^{\text{White}} | q \rangle = \frac{\langle v | H^{od} | q \rangle}{\langle v | H^d | v \rangle - \langle q | H^d | q \rangle}$$

$$\langle v | \eta^{\text{White}} | q \rangle = \frac{\langle v | H^{od} | q \rangle}{\langle v | H^d | v \rangle - \langle q | H^d | q \rangle + \Delta}$$

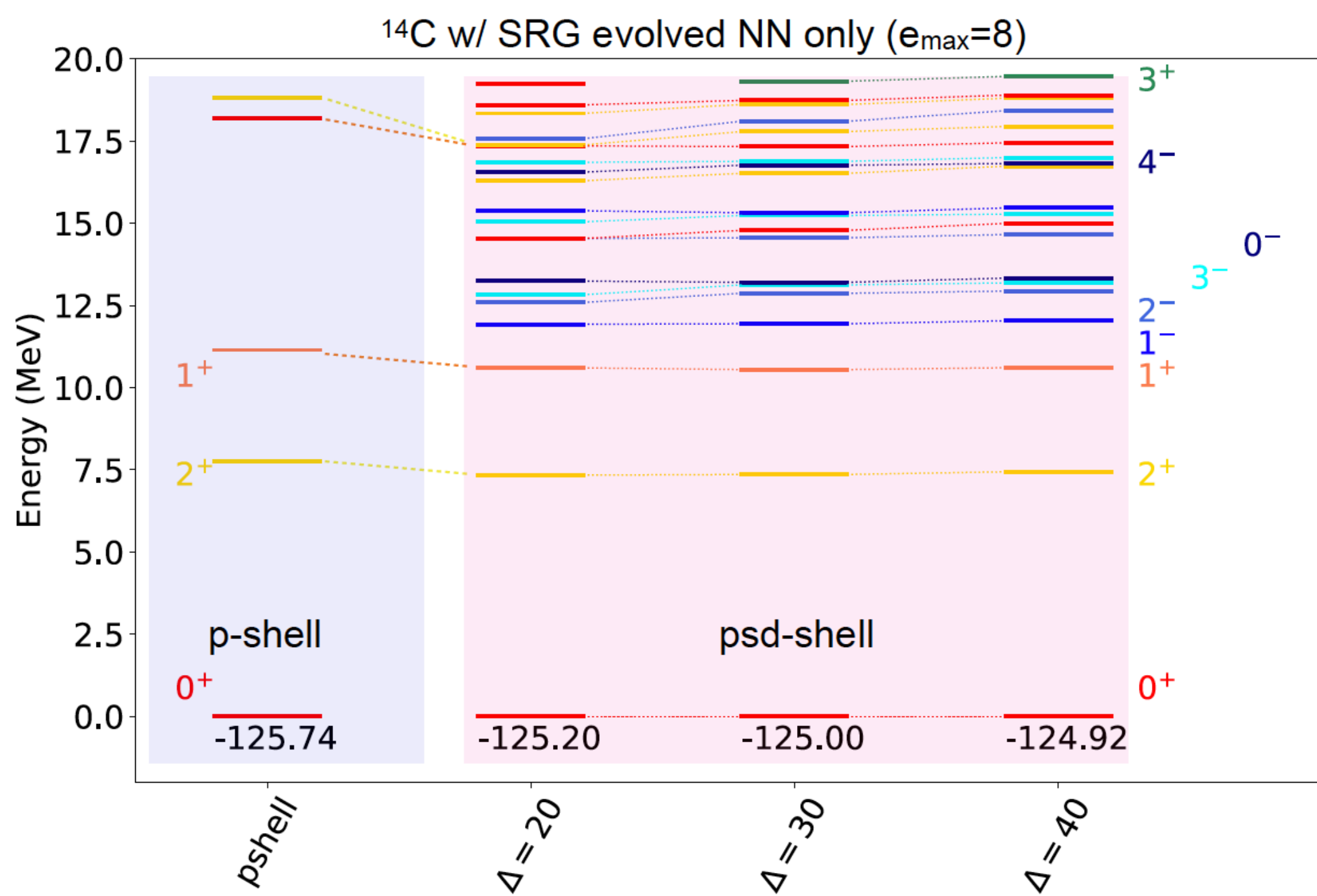
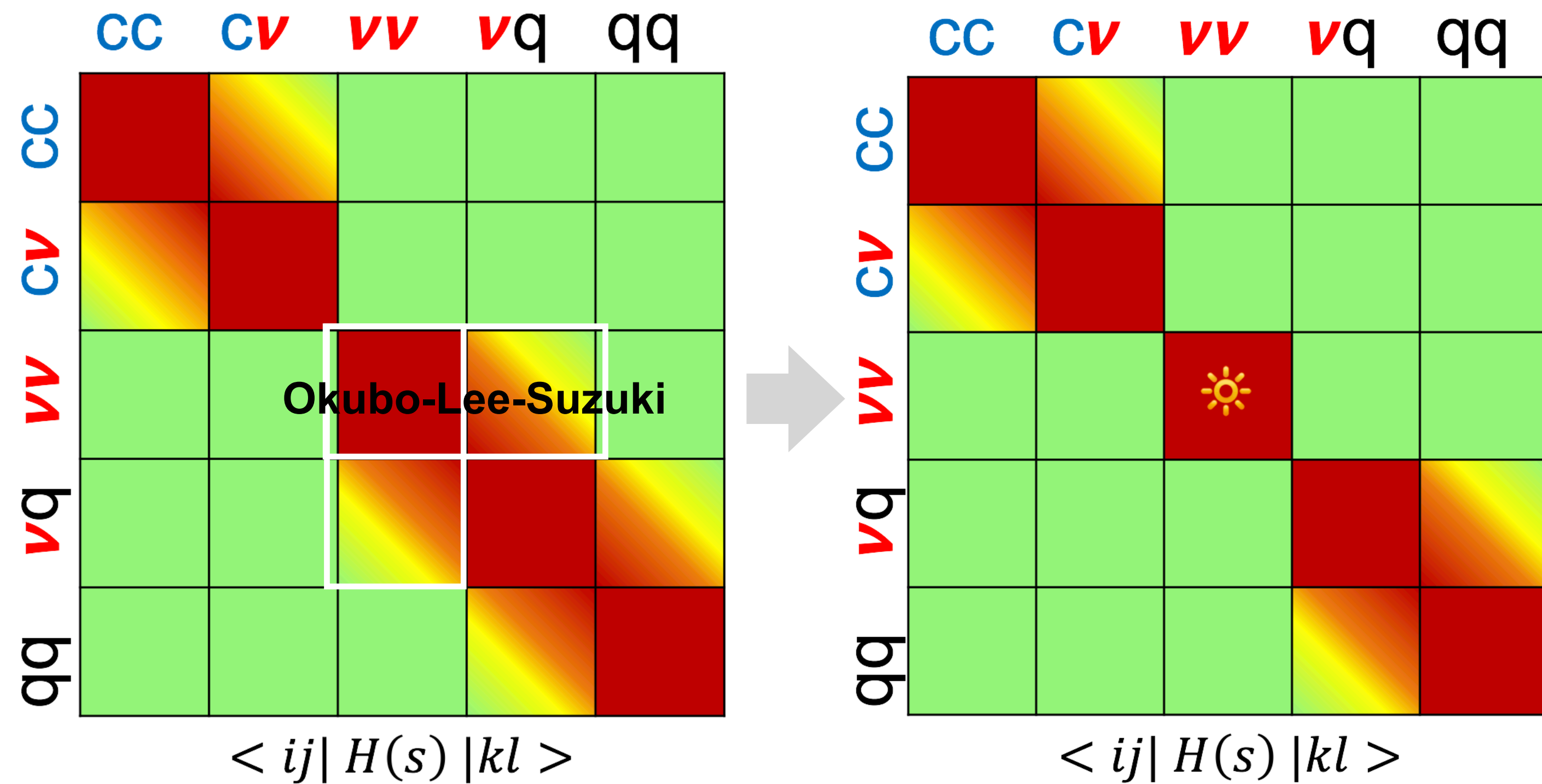


Figure credit: T. Miyagi

② Use Okubo-Lee-Suzuki to decouple  $vv$  from  $vq$



$\omega = Q\omega P$

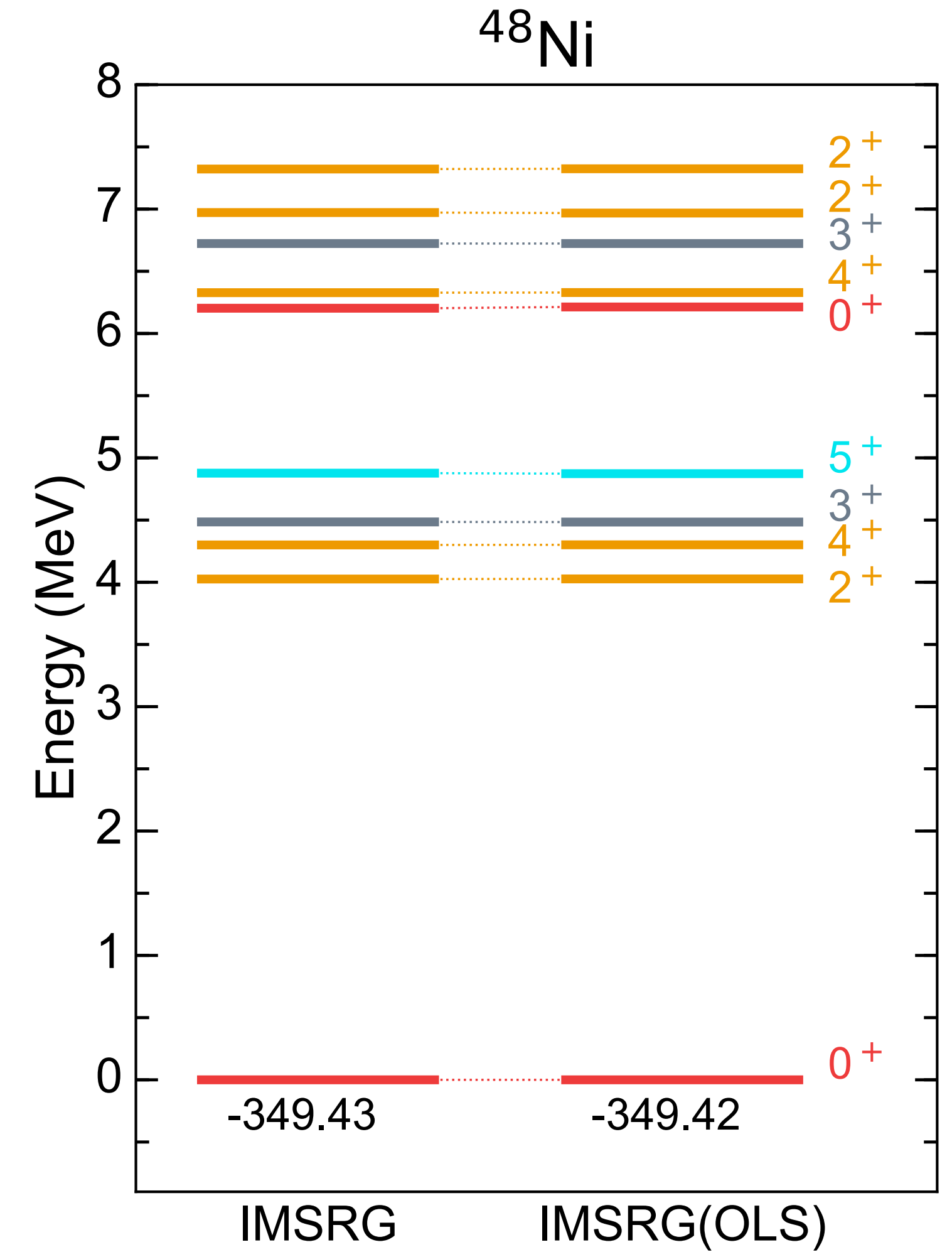
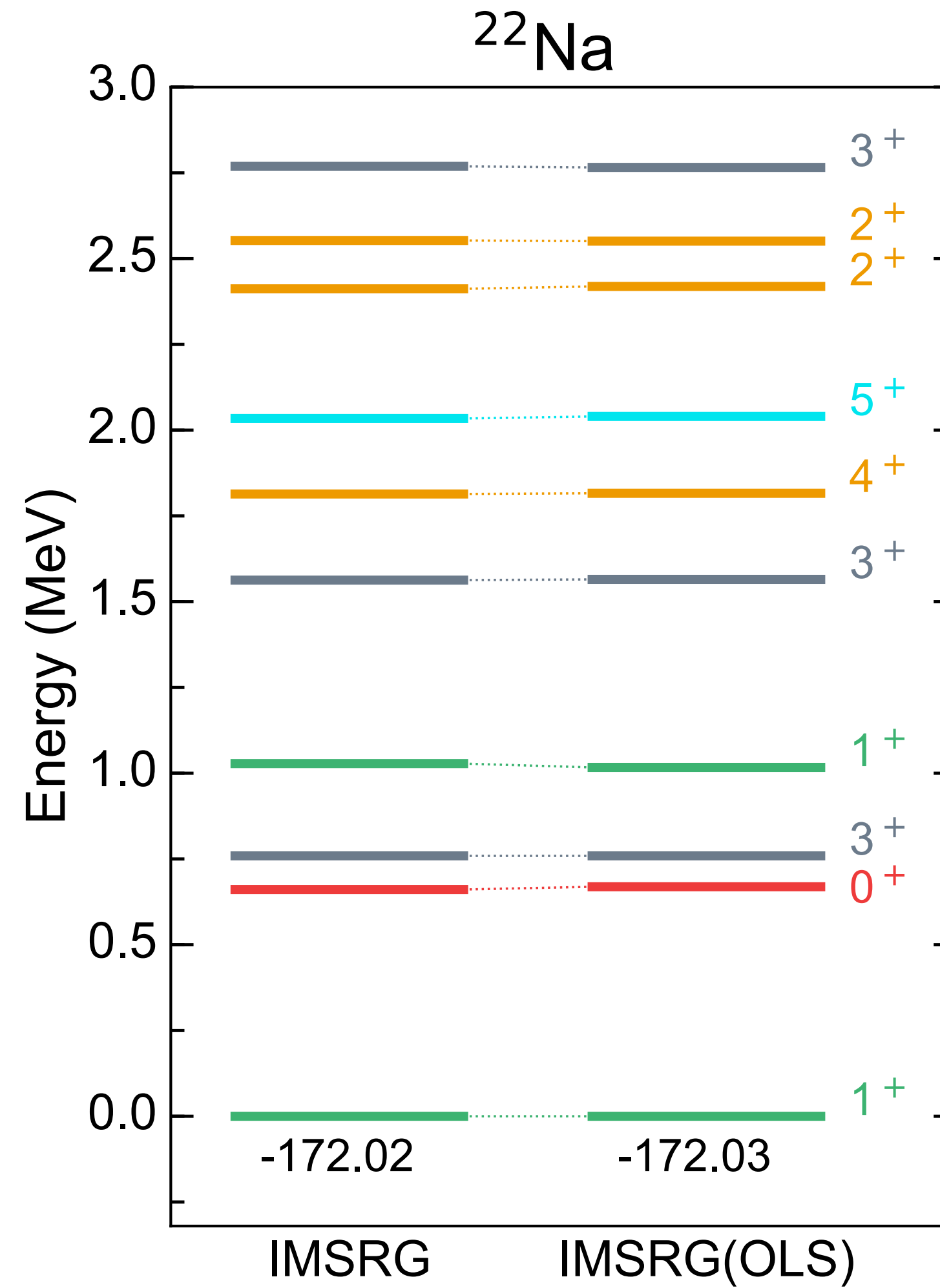
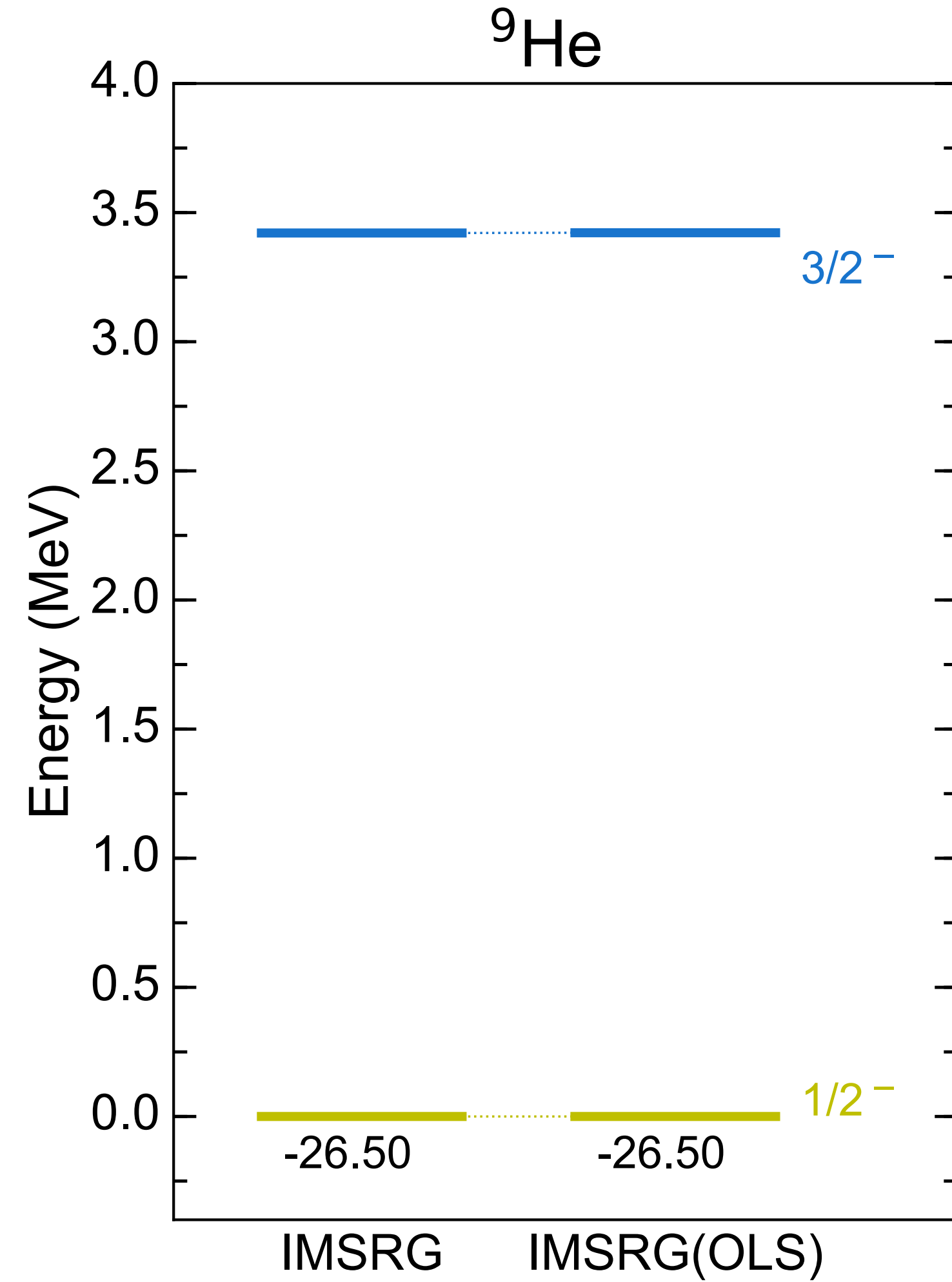
$$\langle \alpha_Q | \omega | \alpha_P \rangle = \sum_{k \in \kappa} \langle \alpha_Q | k \rangle \langle \tilde{k} | \alpha_P \rangle$$

$$\bar{H}_{\text{eff}} = [P(1 + \omega^T \omega)P]^{-1/2} (P + P\omega^T Q) H (P + Q\omega P) [P(1 + \omega^T \omega)P]^{-1/2}$$

# VS-IMSRG with OLS

EM1.8/2.0(NN + 3N),  $e_{\text{max}} = 12$ ,  $e_{3\text{max}} = 16$ ,  $\hbar\omega = 16$  MeV

Using ensemble normal ordering (ENO)  
Without continuum

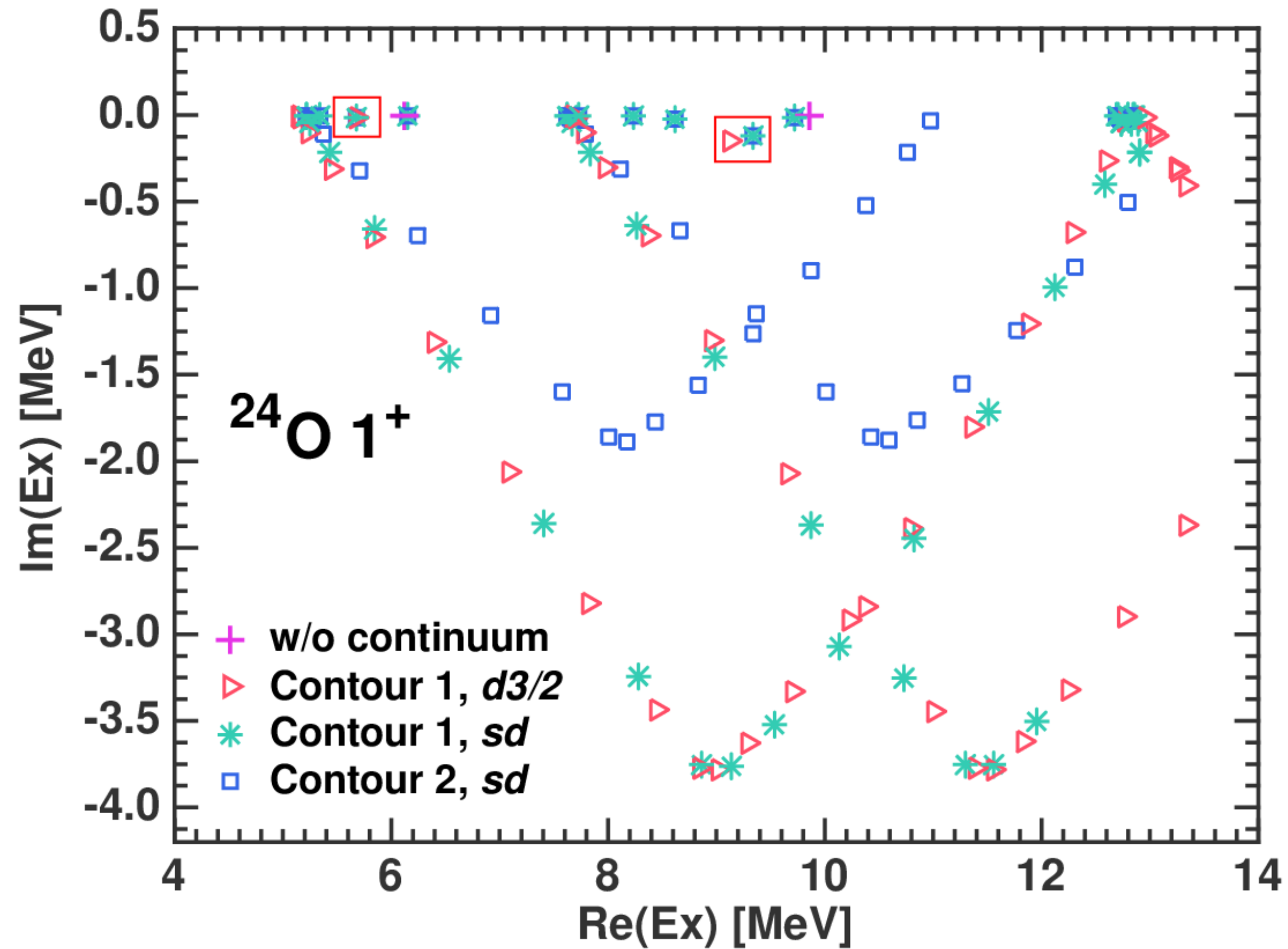


BS Hu, et al., In preparation (2022)

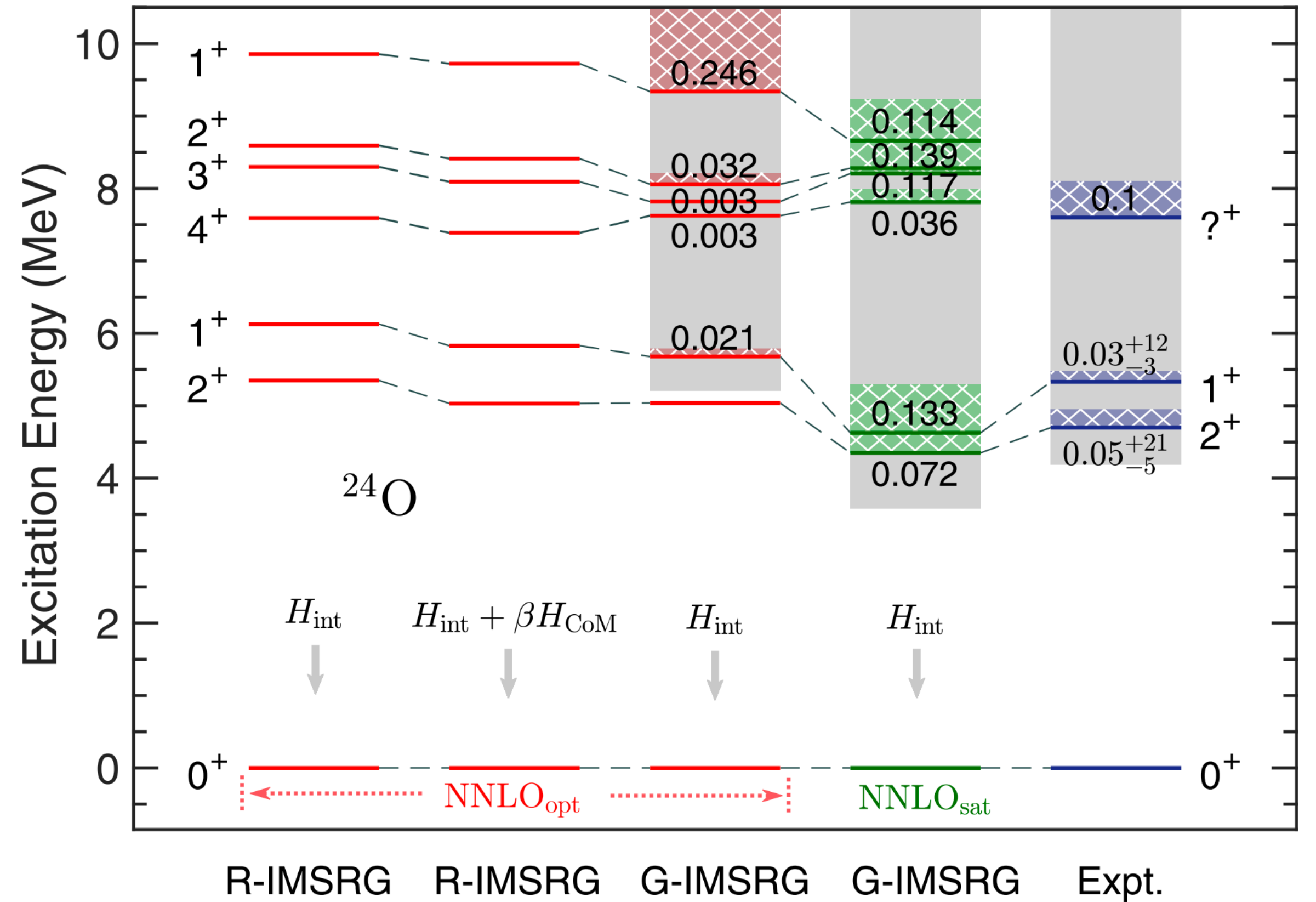
# Gamow EOM-IMSRG results

BShu, Q. Wu, et al., PRC **99** (2019) 061302(R); arXiv:1906.10539

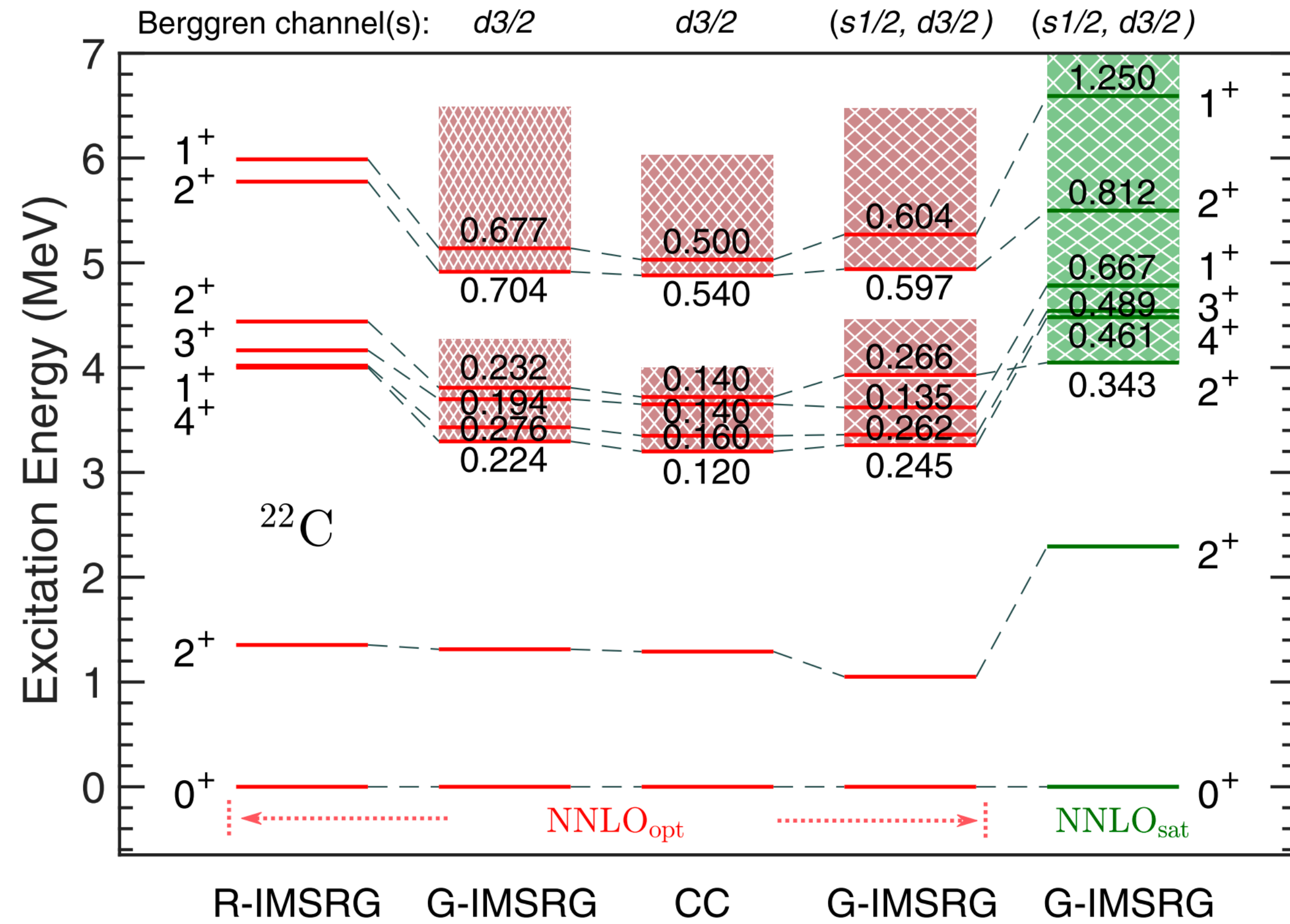
Convergence against different contour  
and discretization number



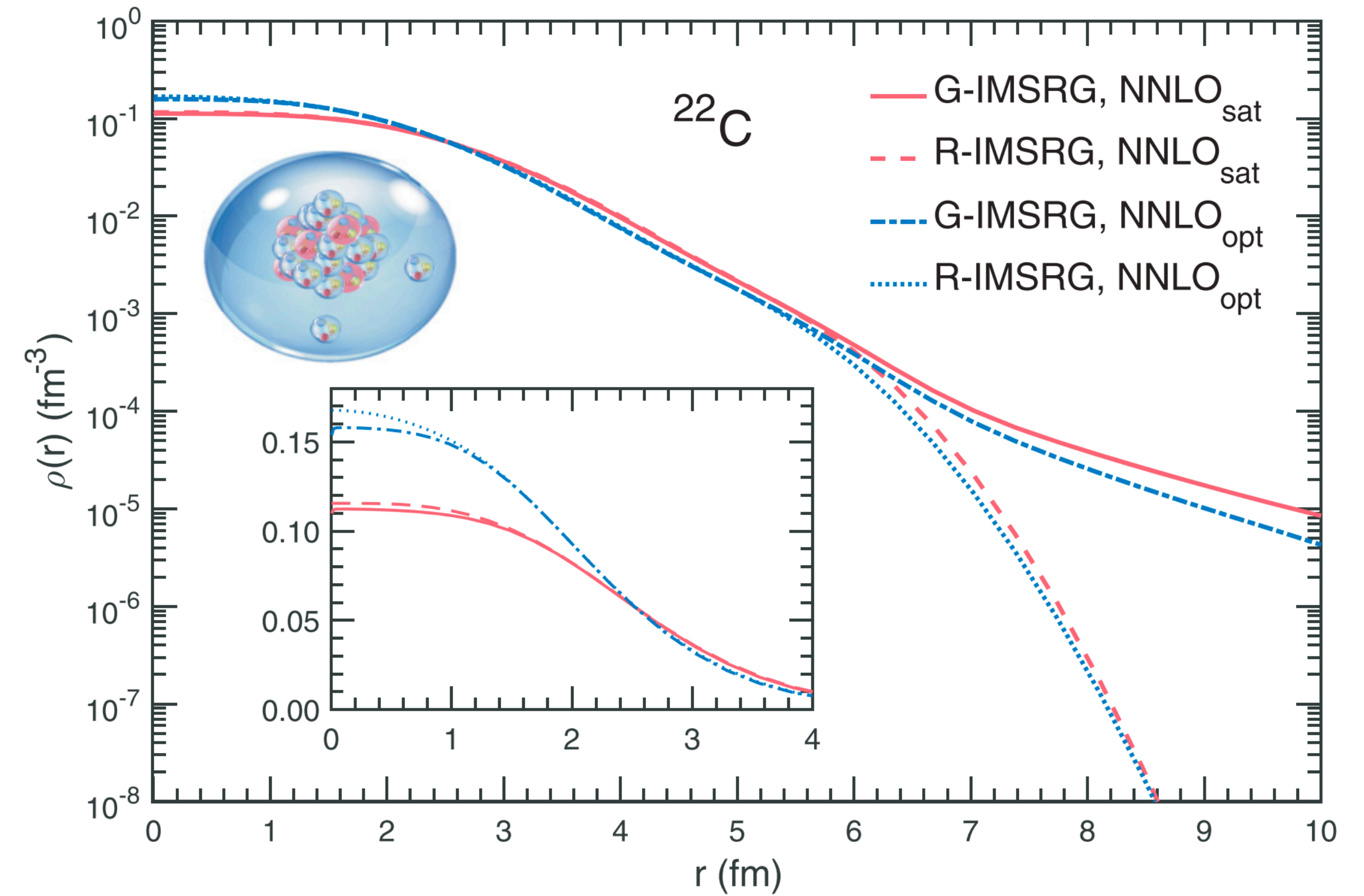
Center-of-mass spurious excitation



# EOM-IMSRG for Borromean $^{22}\text{C}$



BShu, Q. Wu, et al., PRC **99** (2019) 061302(R); arXiv:1906.10539



NNLO <sub>sat</sub>	R-IMSRG	G-IMSRG cont. $s, d$ waves	Expt. Estimated
matter radius (fm)	2.98	3.14	$3.44 \pm 0.08$ <sup>①</sup> $3.38 \pm 0.10$ <sup>②</sup>

①: Y. Togano *et al.*, PLB**761** (2016) 412

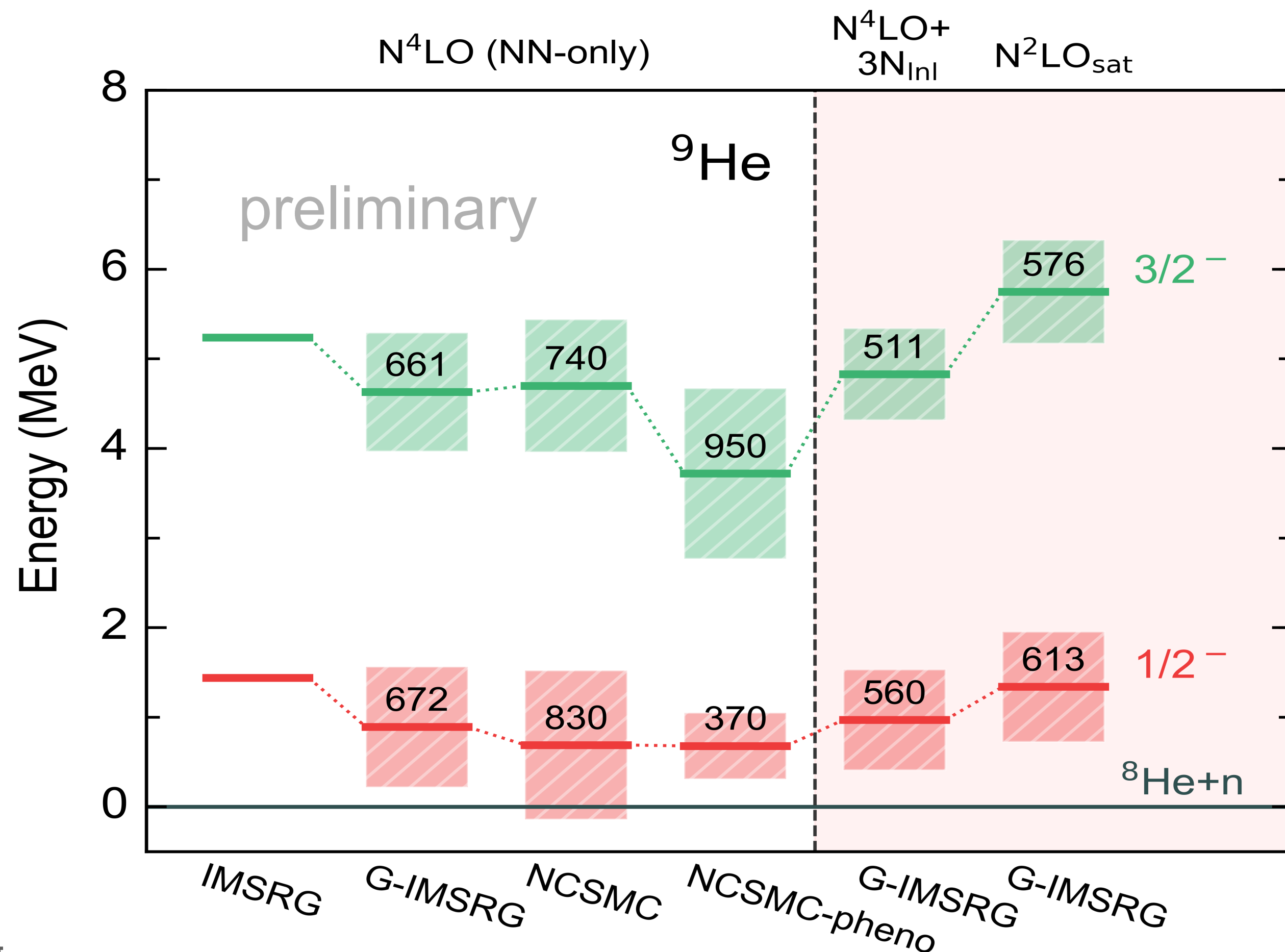
②: T. Nagahisa and W. Horiuchi,  
PRC**97** (2018) 054614

# Benchmark with NCSMC

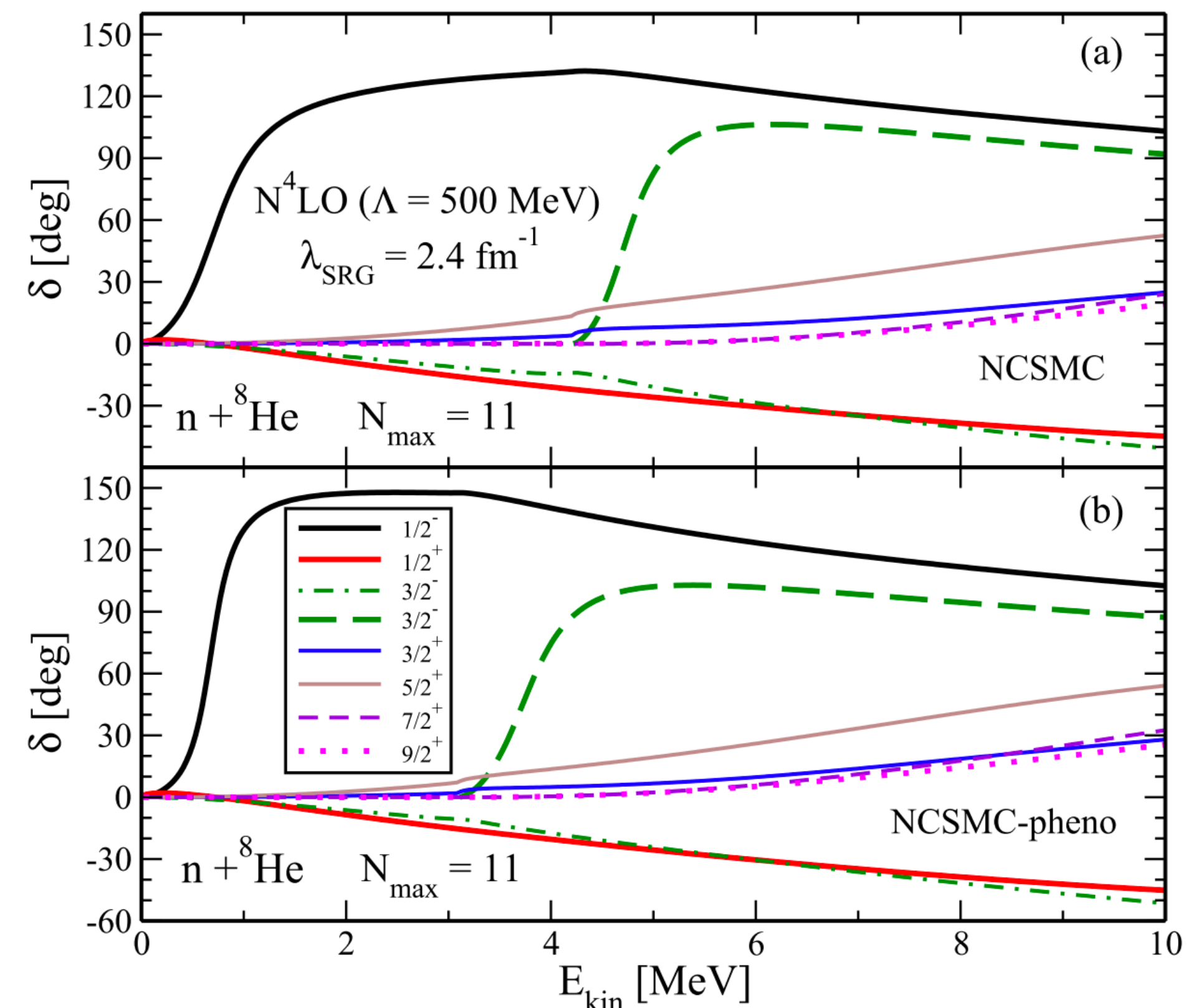
Core:  $^4\text{He}$

Valence space: neutron  $p_{1/2}$ ,  $p_{3/2}$  resonances and  $s_{1/2}$ ,  $p_{1/2}$ ,  $p_{3/2}$ ,  $d_{5/2}$  continua

$N^4\text{LO}(500)$ ,  $\lambda_{\text{SRG}} = 2.4 \text{ fm}^{-1}$ ,  $\hbar\Omega = 20 \text{ MeV}$

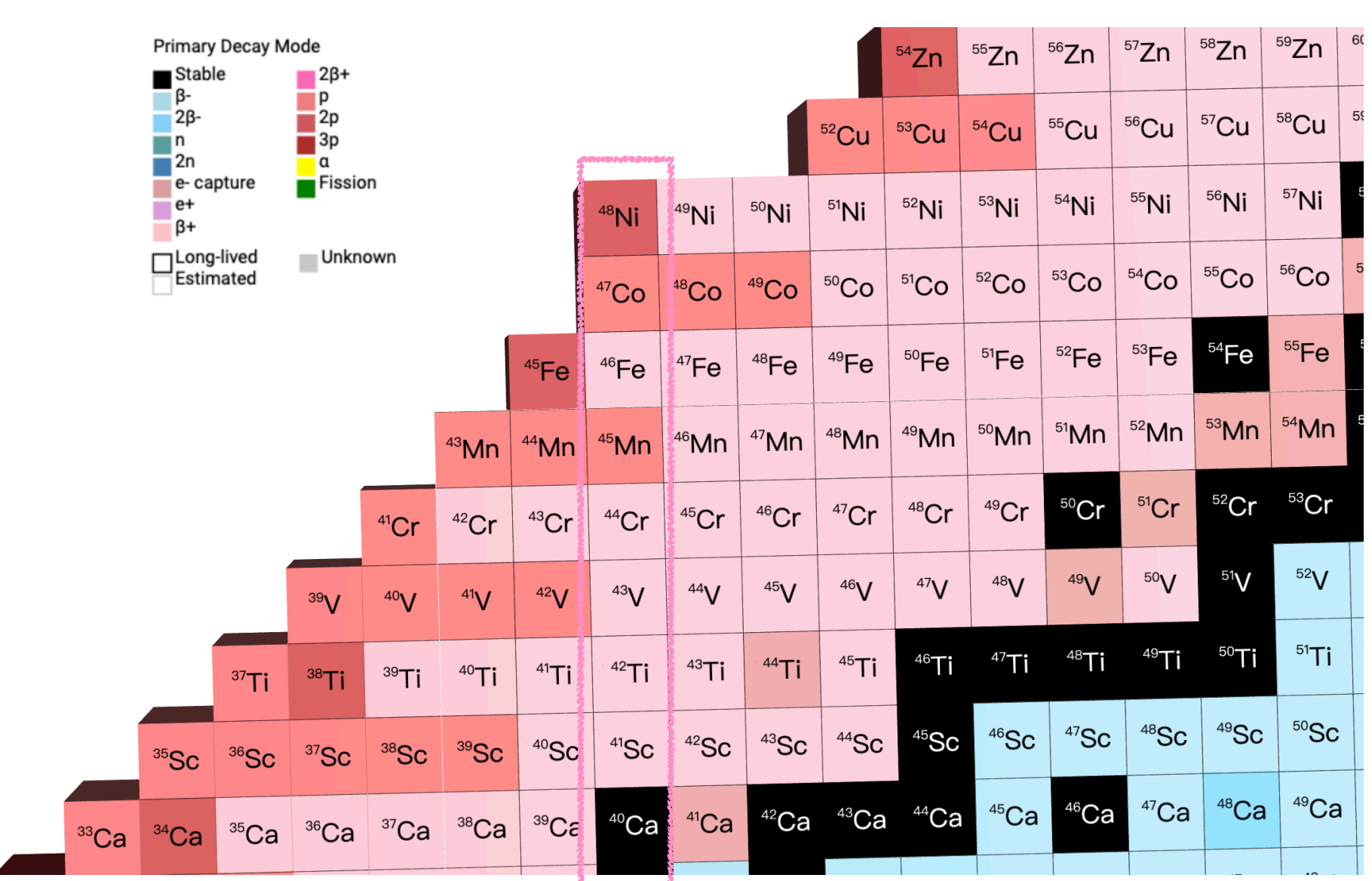
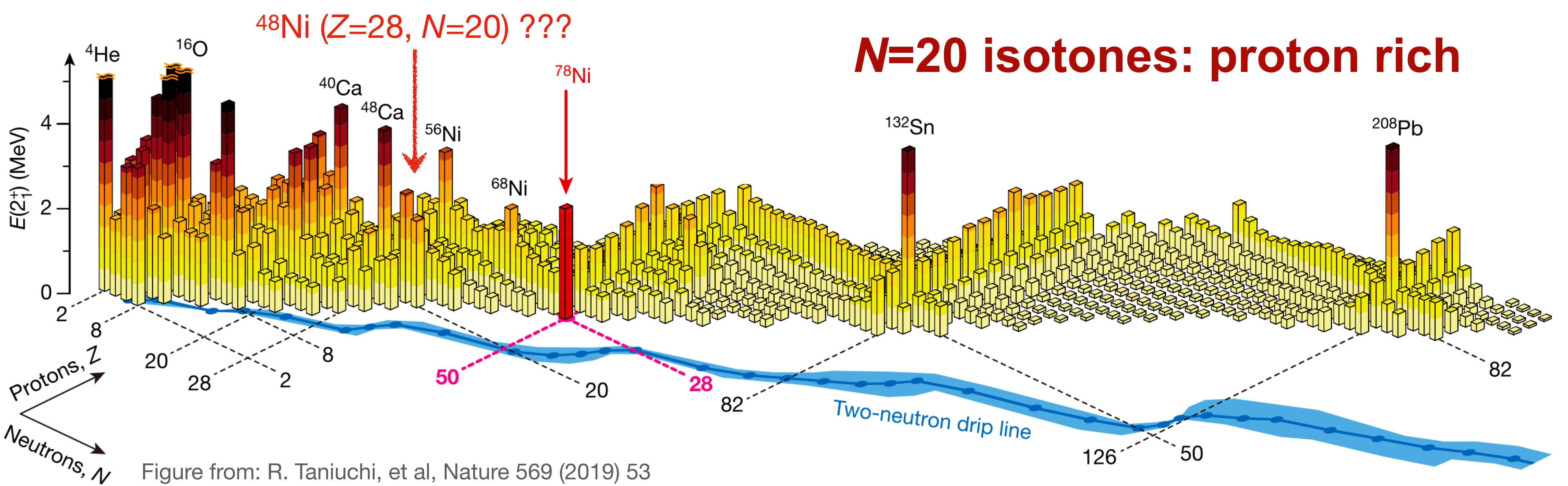


$E_{\text{g.s.}}$ (MeV)	$^4\text{He}$	$^6\text{He}$	$^8\text{He}$
IMSRG(ENO)	-28.50	-28.25	-29.47
NCSM	-28.36	-28.94(20)	-30.23(30)
Expt.	-28.30	-29.27	-31.41

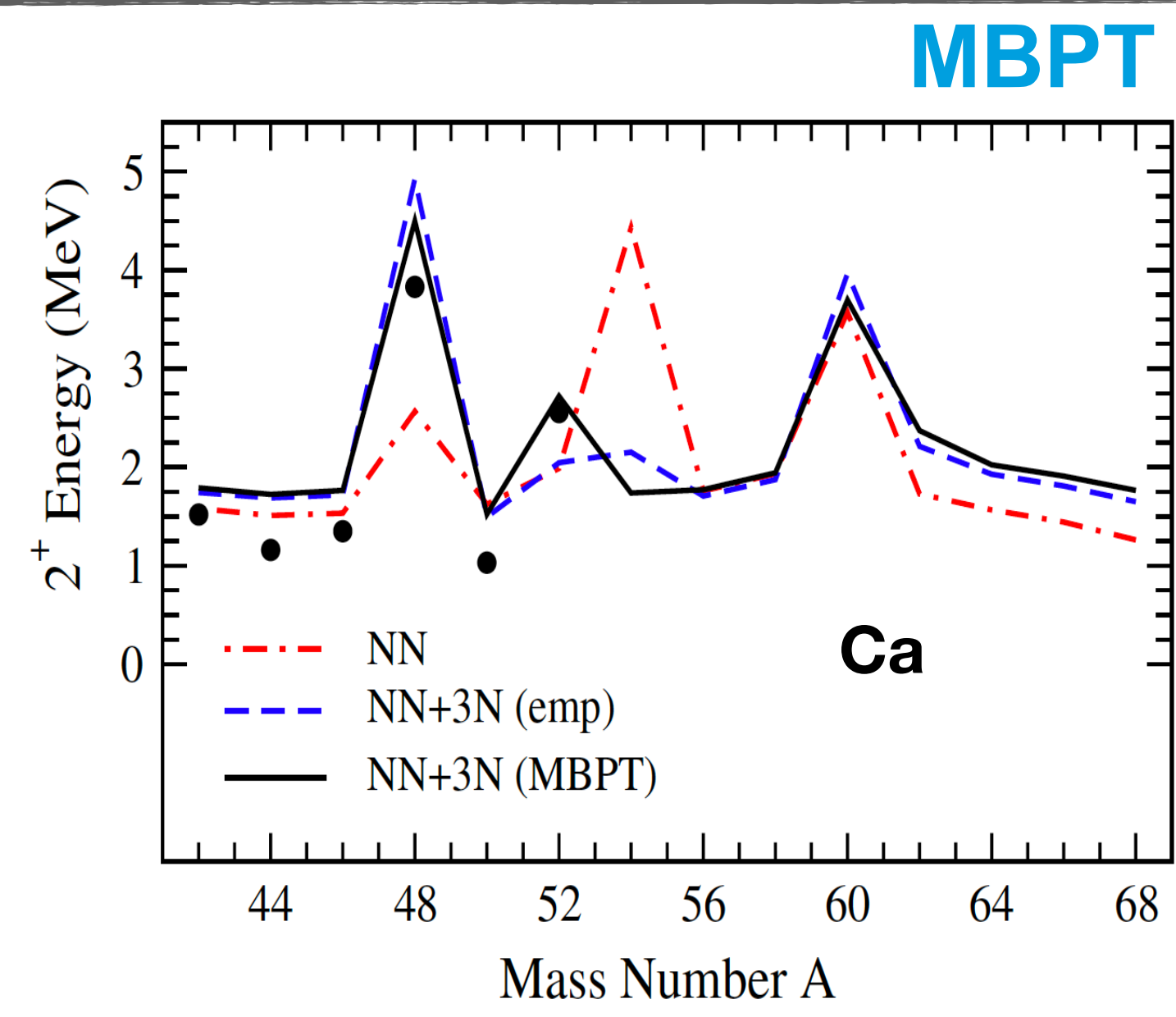


M. Vorabbi, A. Calci, P. Navrátil, *et al.*, PRC97, 034314 (2018)

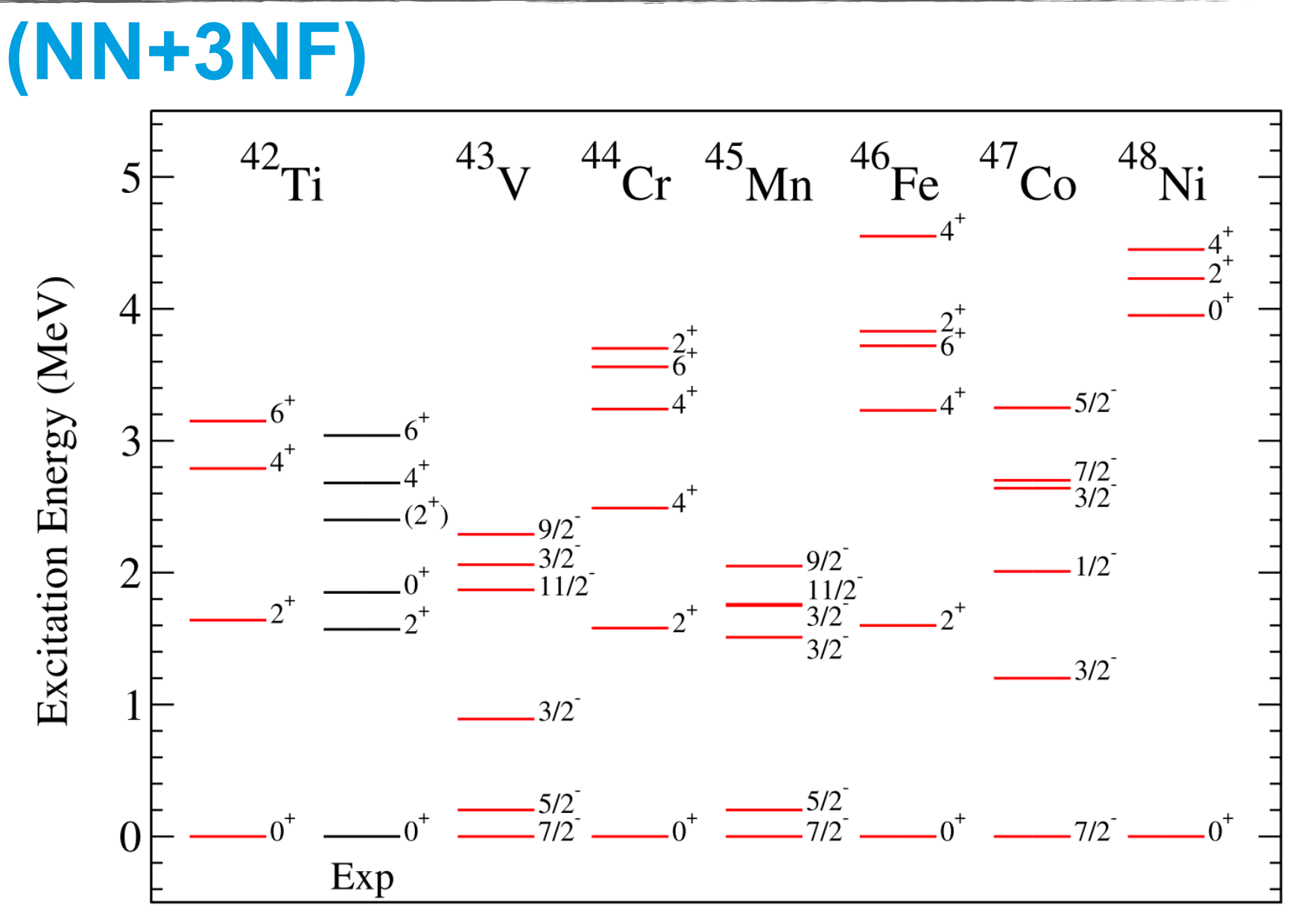




Baishan Hu - TRIUMF (2022/8/9)

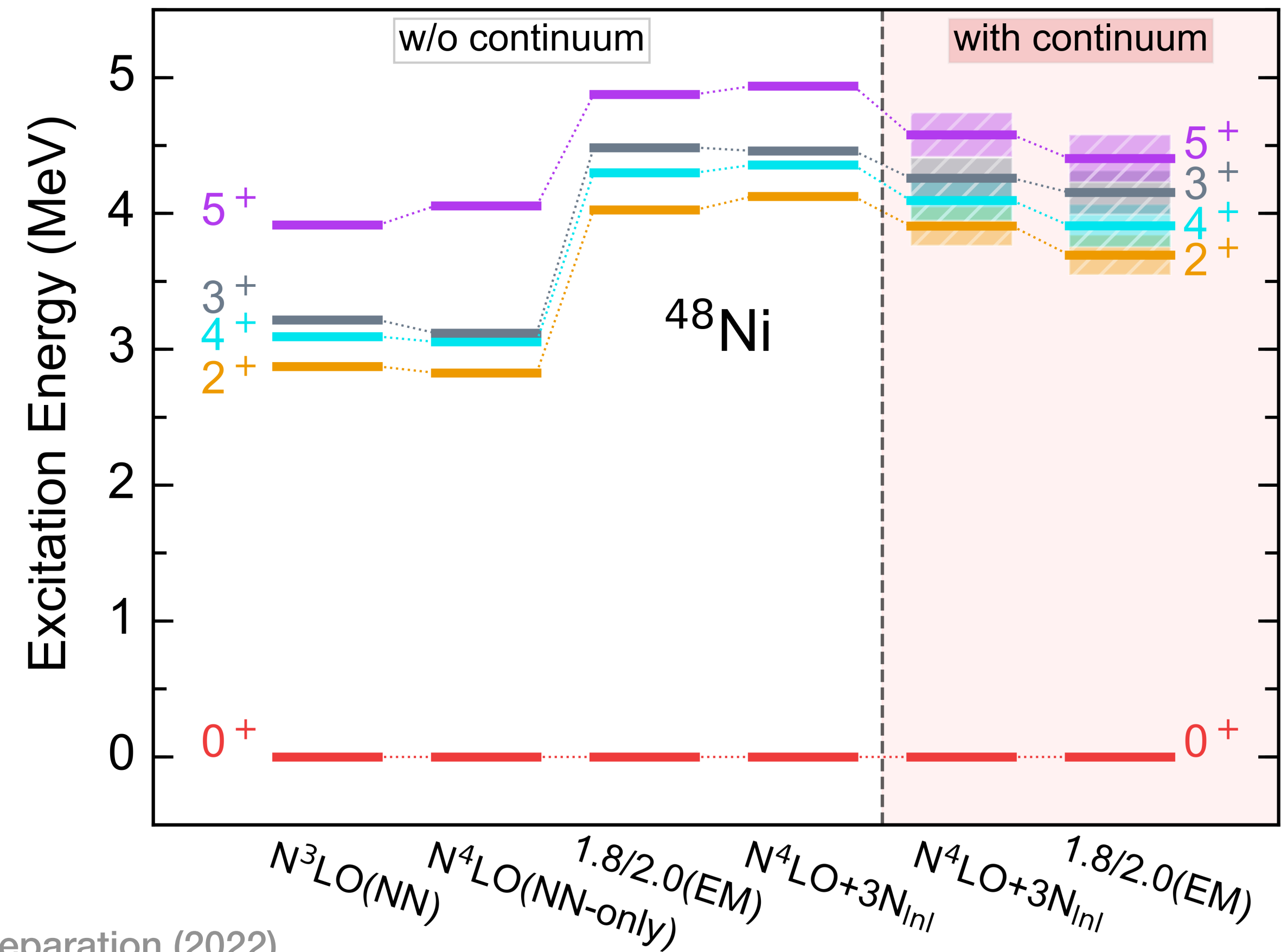
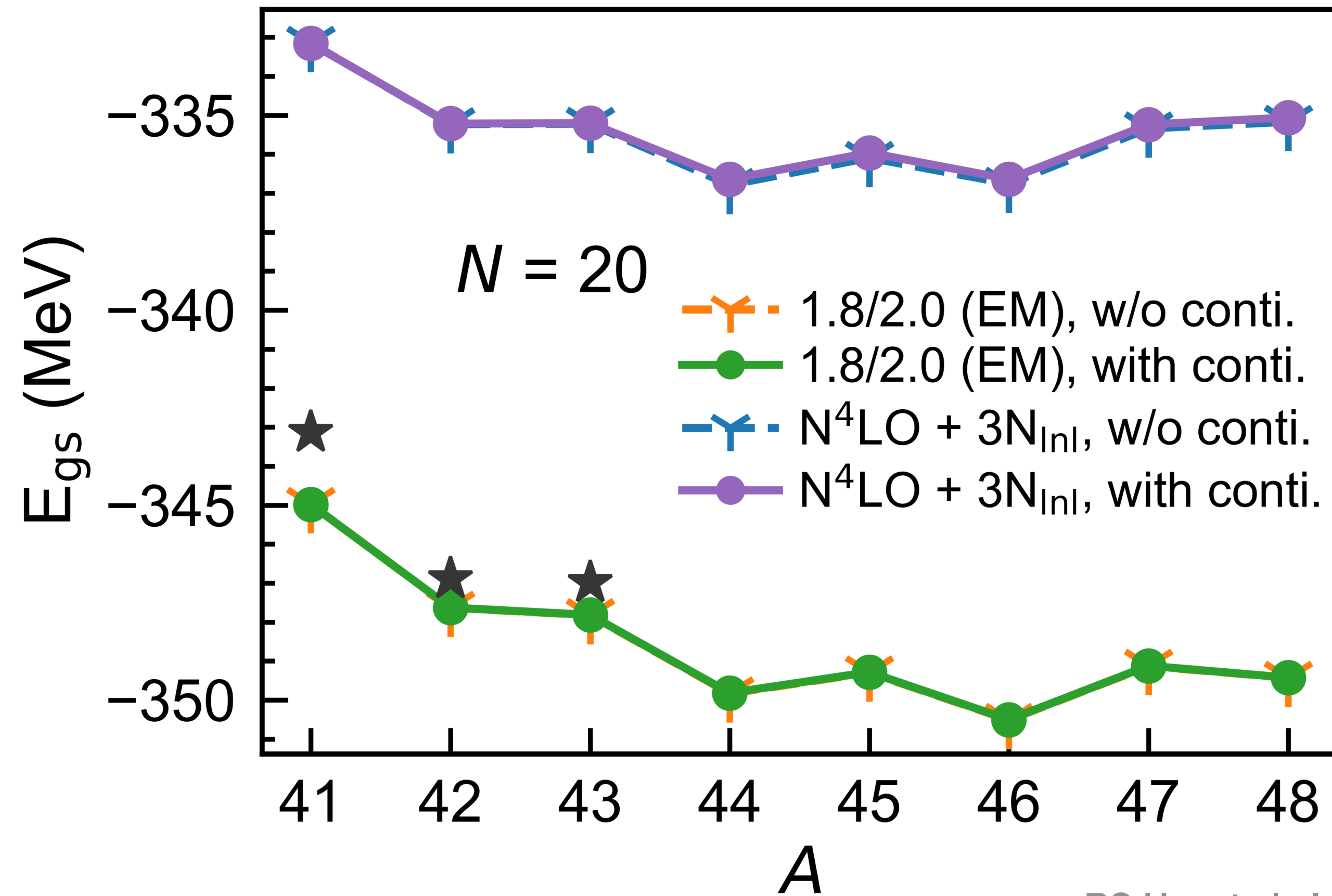


J. Holt, et al, JPG:NPP 40 (2013) 075105



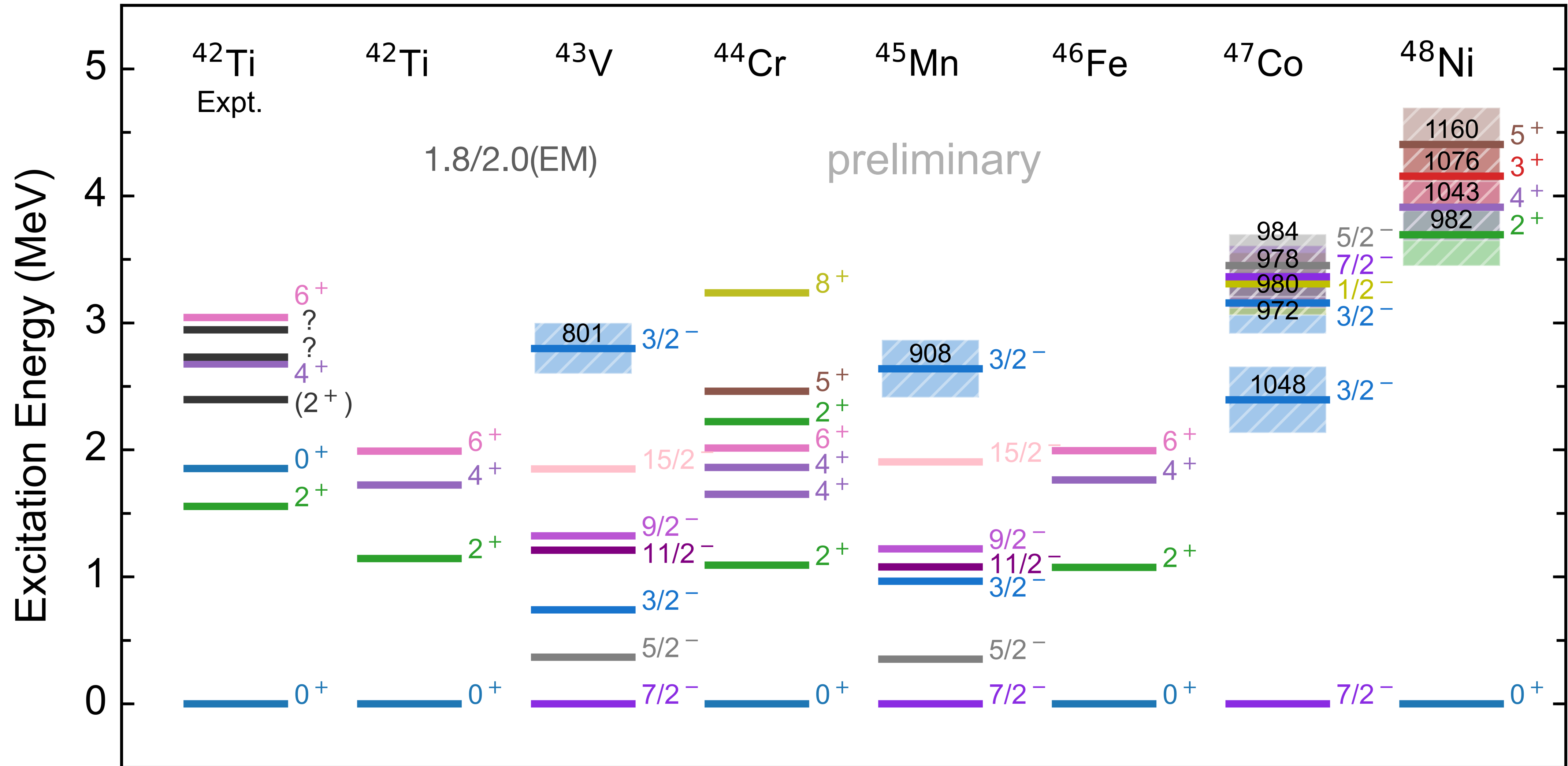
J. Holt, J. Menéndez and A. Schwenk, PRL110 (2013) 022502

# Gamow VS-IMSRG results

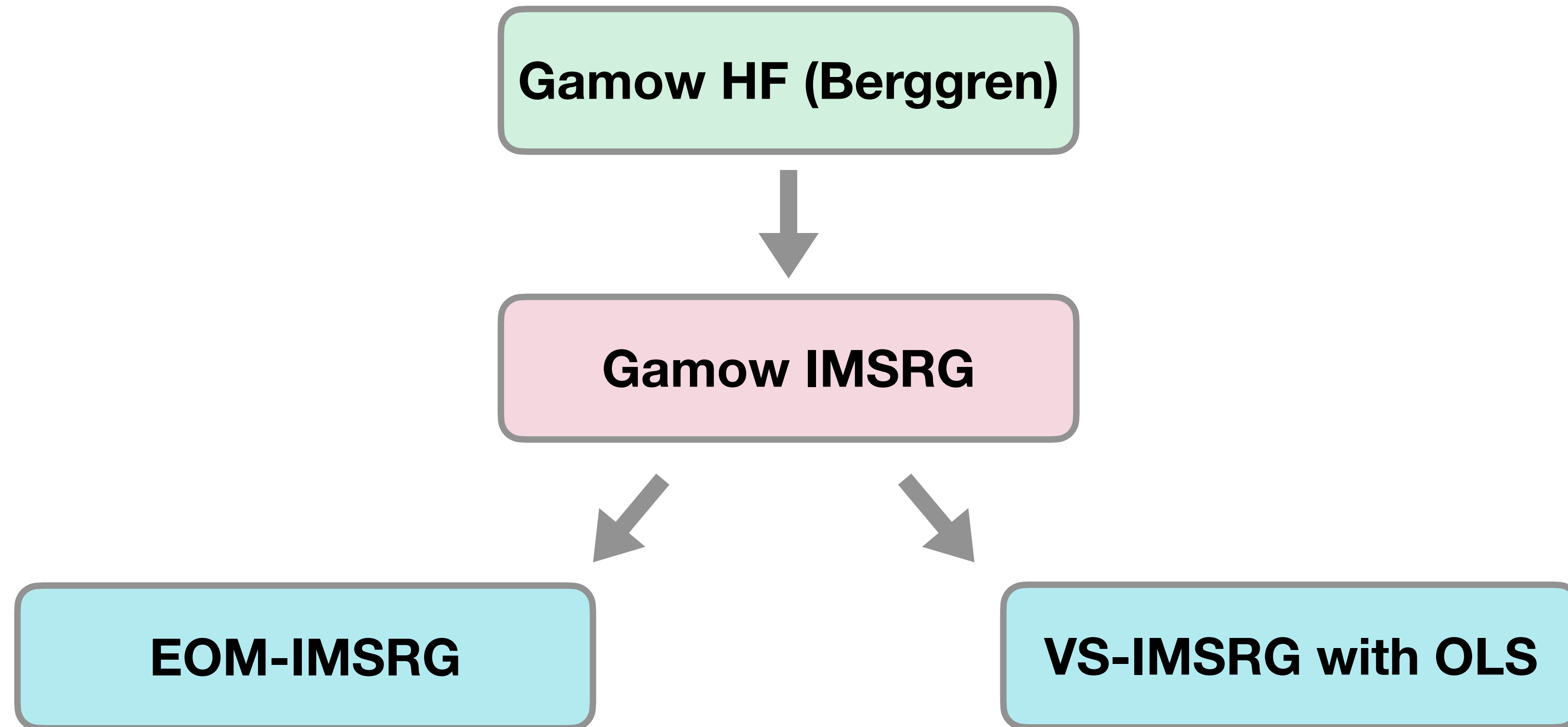


BS Hu, et al., In preparation (2022)

# Gamow VS-IMSRG results



# Summary



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Thank you  
Merci

