

Gravitational collapse to extremal black holes and the third law of black hole thermodynamics

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joint work with Ryan Unger (Princeton)

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The Four Laws of Black Hole Mechanics

J. M. Bardeen★

Department of Physics, Yale University, New Haven, Connecticut, USA

B. Carter and S. W. Hawking

Institute of Astronomy, University of Cambridge, England

Received January 24, 1973

Black hole thermodynamics is a proposed close mathematical **analogy** between black hole dynamics and classical thermodynamics.

BLACK HOLE THERMODYNAMICS

Law	Classical thermodynamics	Black holes

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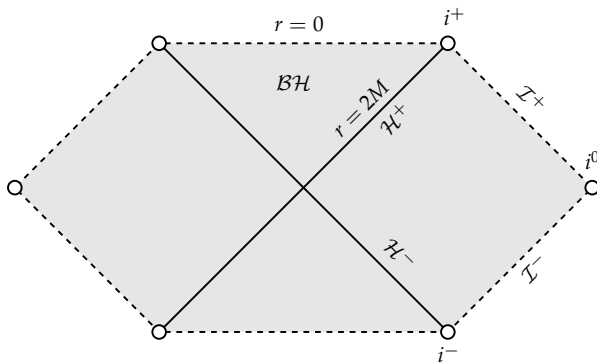
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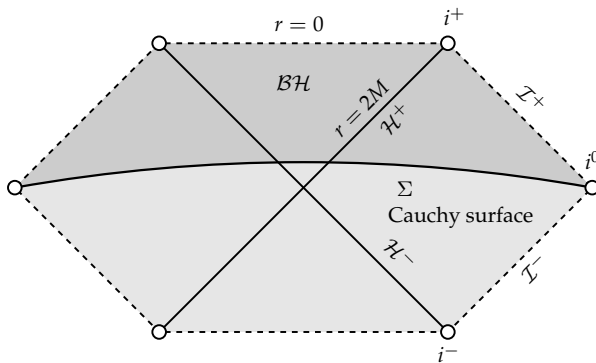
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- Proposed by Bardeen–Carter–Hawking '73 as statements within **classical GR**.
- Hawking radiation—interpret $\kappa \propto T$
- Bekenstein entropy—interpret $A \propto S$
- $\kappa \propto T$, $A \propto S$ further justified once quantum effects are considered, but they are **formal analogies**.
- Laws 0, 1, and 2 proved by Hawking, Carter, Bardeen–Carter–Hawking, Wald, ...

REFRESHER ON SCHWARZSCHILD

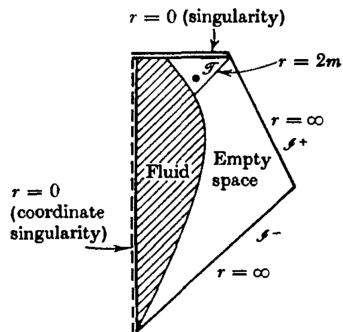


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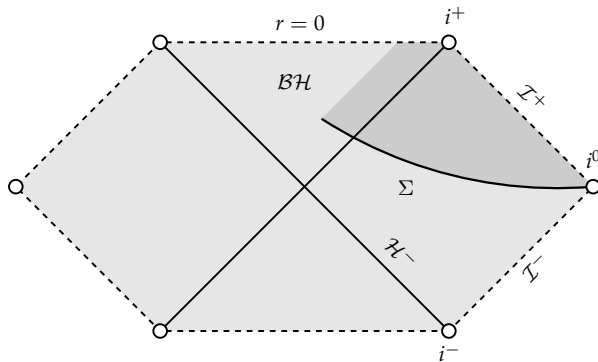
Maximally extended Schwarzschild is the unique maximal Cauchy development of the data induced on a spacelike hypersurface $\Sigma \cong \mathbb{R} \times S^2$.

REFRESHER ON GRAVITATIONAL COLLAPSE

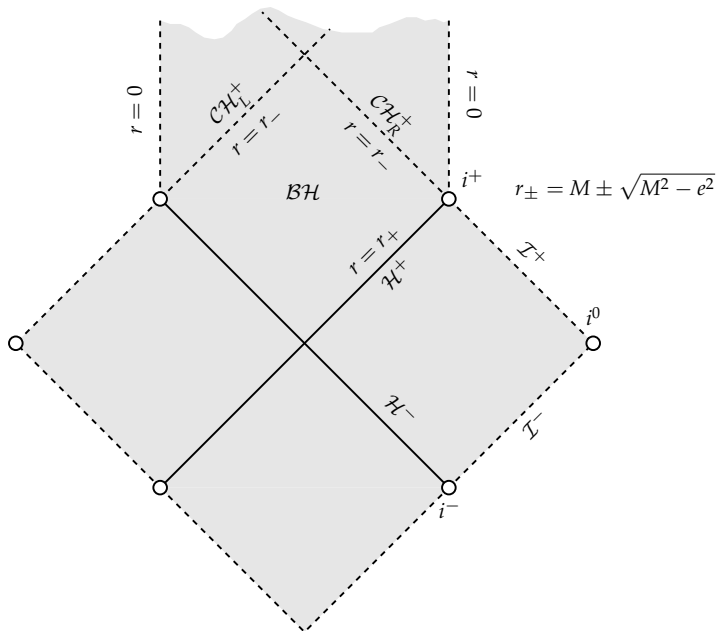


Penrose diagram of gravitational collapse. One-ended Cauchy data!

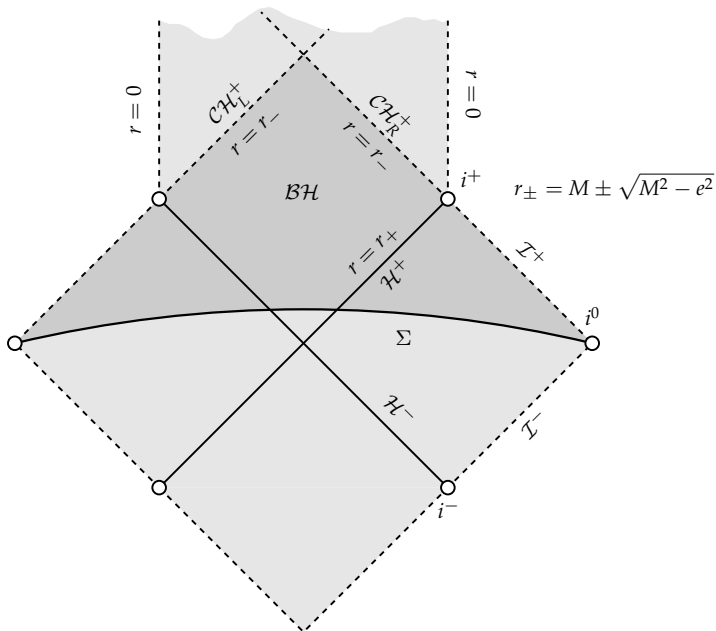
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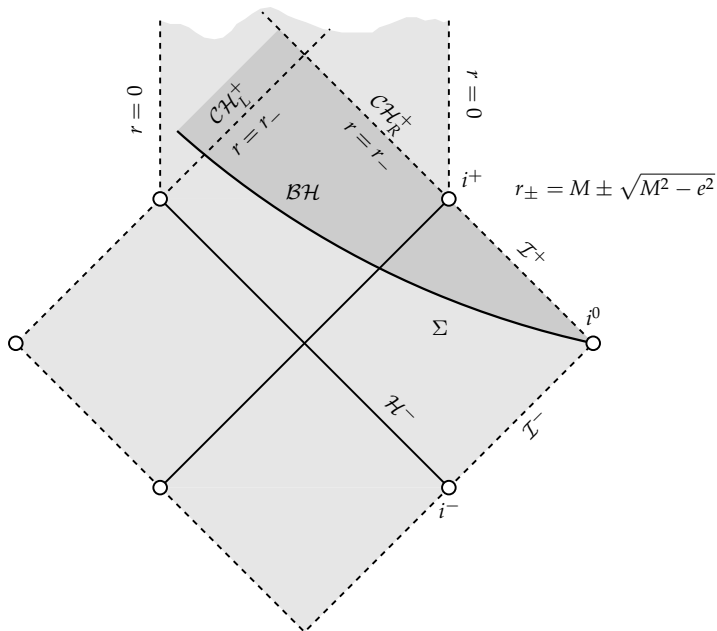
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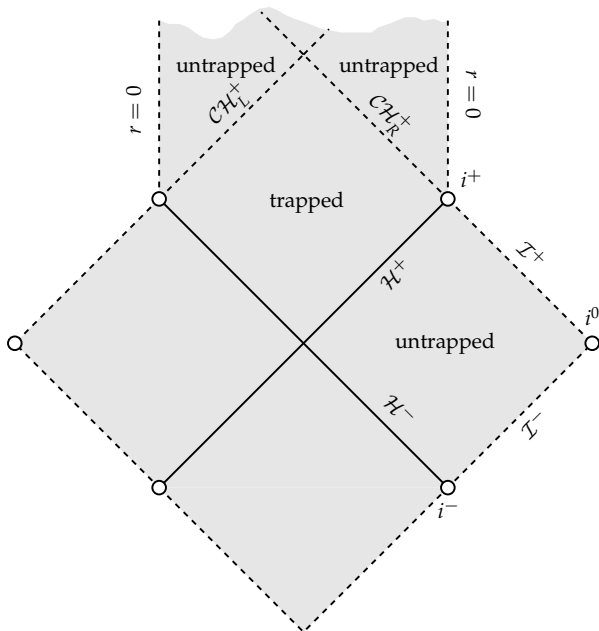
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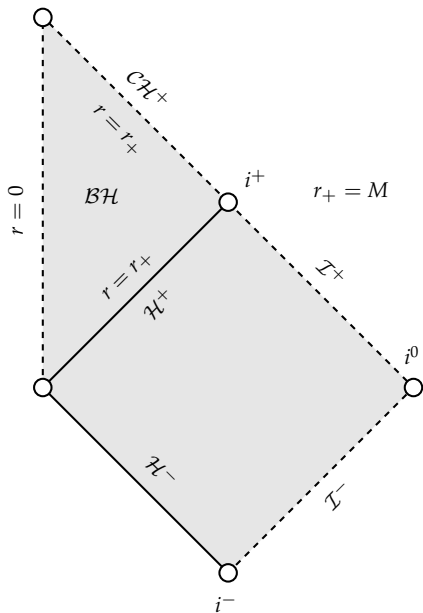
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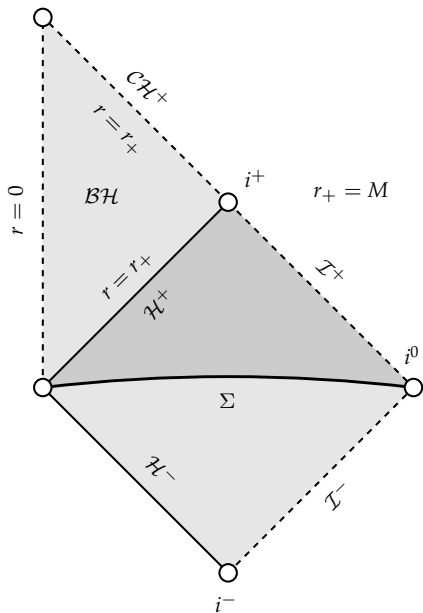
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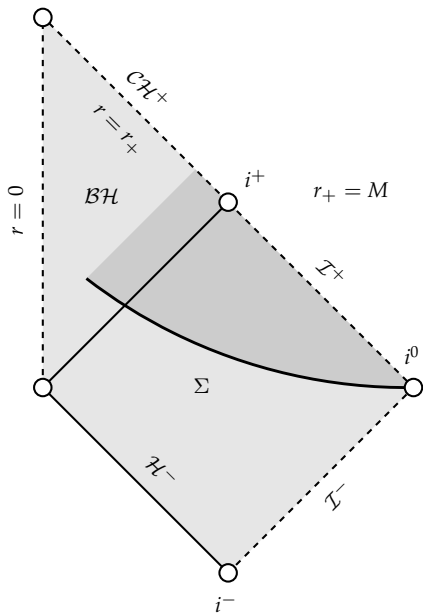
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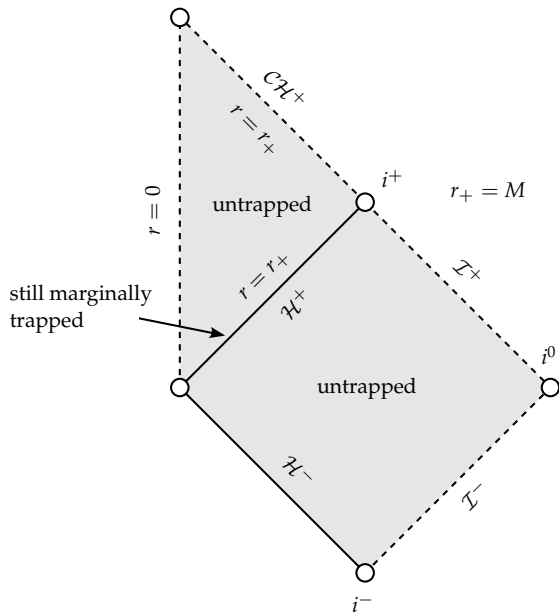
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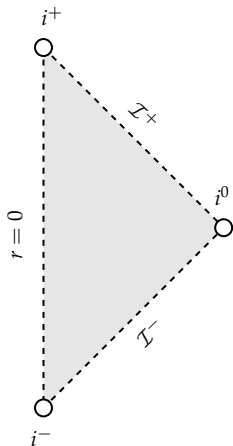
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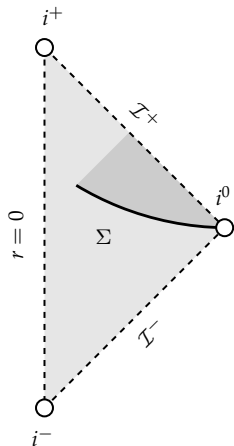
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REFRESHER ON SUPEREXTREMAL REISSNER–NORDSTRÖM: $0 < M < |e|$



REFRESHER ON SUPEREXTREMAL REISSNER–NORDSTRÖM: $0 < M < |e|$



SURFACE GRAVITY κ OF REISSNER–NORDSTRÖM

- RN with mass M and charge e , $|e| \leq M$, has

$$\kappa = 2\pi T = \frac{\sqrt{M^2 - e^2}}{(M + \sqrt{M^2 - e^2})^2}$$

- **Subextremal:** $\kappa > 0$
- **Extremal:** $\kappa = 0$

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Original formulation of Bardeen–Carter–Hawking:

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It is impossible by any procedure, no matter how idealized, to reduce κ to zero by a finite sequence of operations.

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A nonextremal black hole cannot become extremal (i.e., lose its trapped surfaces) at a finite advanced time in any continuous process in which the stress-energy tensor of accreted matter stays bounded and satisfies the weak energy condition in a neighborhood of the outer apparent horizon.

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 - ▶ If singularities allowed, counterexample using massive dust shell. (Farrugia-Hajicek '79)

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4. **Weak energy condition** must be enforced.
 - ▶ Otherwise: counterexample using charged null dust. (Sullivan-Israel '80)

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Conjecture (The third law, BCH '73, Israel '86).

*A subextremal black hole cannot become extremal in finite time by any continuous process, no matter how idealized, in which the spacetime and matter fields remain *regular* and obey the *weak energy condition*.*

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Theorem (K.–Unger '22).

Subextremal black holes can become extremal in finite time, evolving from regular Cauchy data for the Einstein–Maxwell-charged scalar field system.

In particular, the “third law of black hole thermodynamics” is false.

Third Law of Black-Hole Dynamics: A Formulation and Proof

W. Israel^(a)

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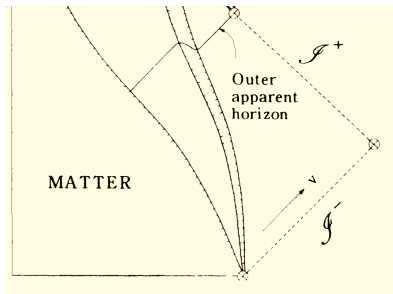
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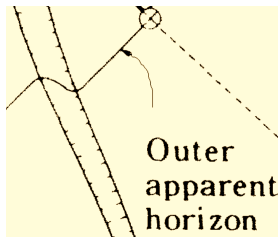
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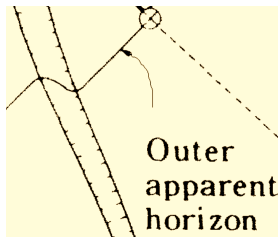
- ▶ First incoming matter flux creates (dynamical) subextremal apparent horizon.
- ▶ Second matter flux pushes to extremal horizon.

ISRAEL'S ARGUMENT II



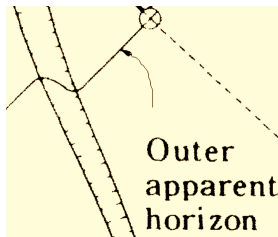
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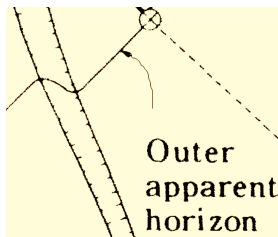


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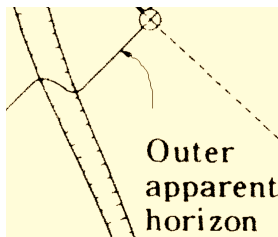
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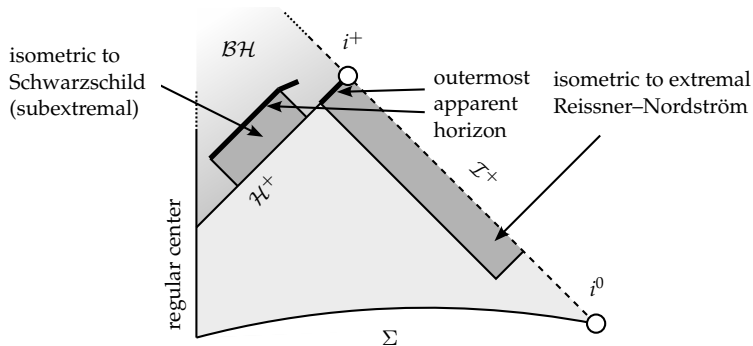
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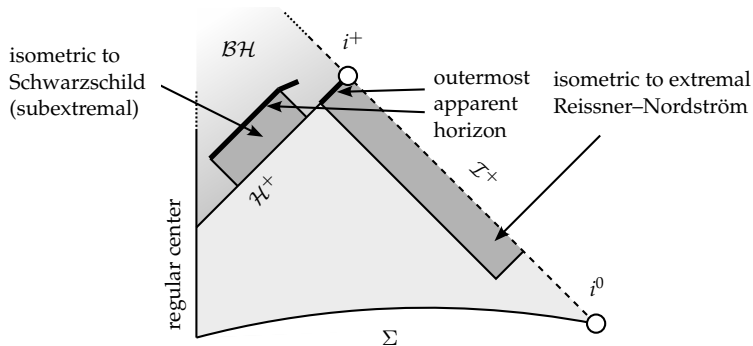
However, outer apparent horizon can jump in smooth spacetimes.

COUNTEREXAMPLE TO THE THIRD LAW



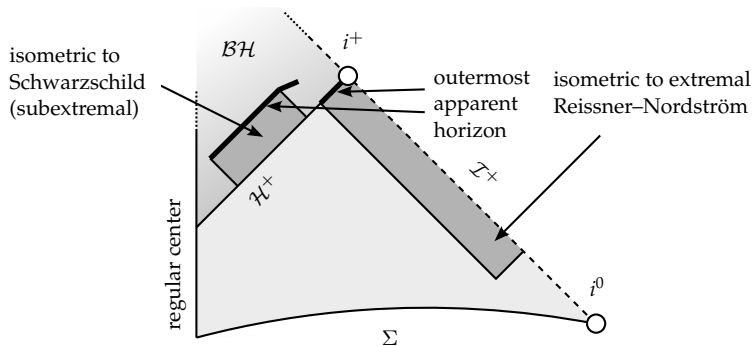
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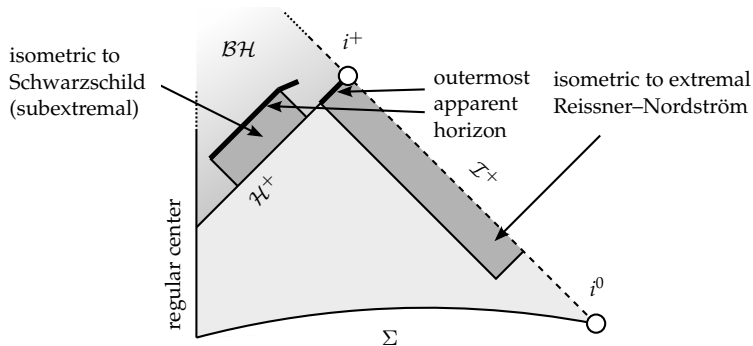
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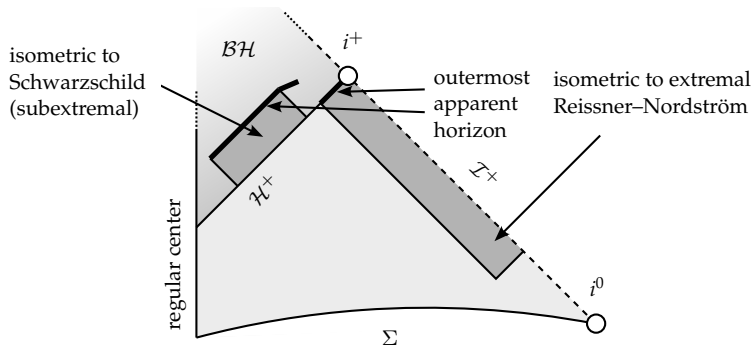
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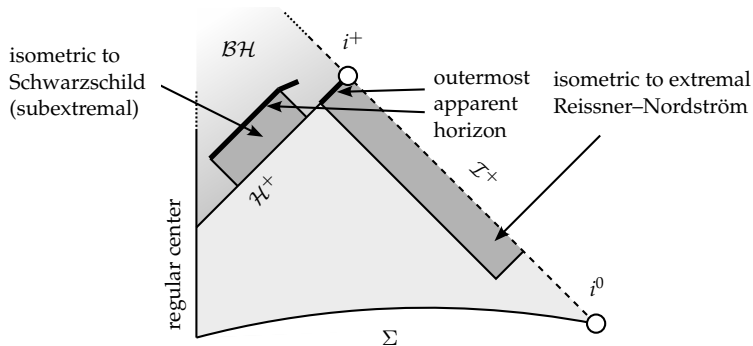
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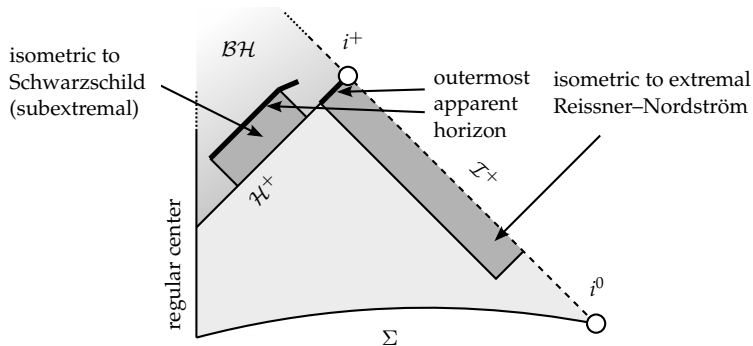
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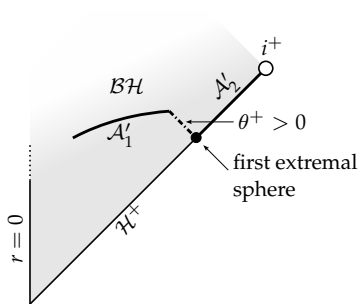
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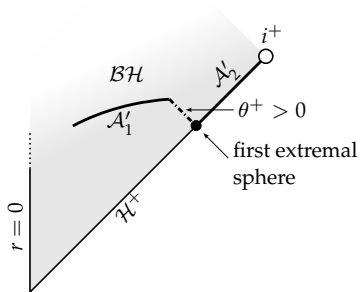
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- Proof directly extends to massive fields for $m \ll |\epsilon|$.

ISRAEL'S PAPER REINTERPRETED



Outermost apparent horizon becomes disconnected
the instant the black hole becomes extremal!

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Outermost apparent horizon becomes disconnected
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This is a feature, not a bug!

EINSTEIN-MAXWELL-CHARGED SCALAR FIELD SYSTEM

- ▶ Lorentzian manifold (\mathcal{M}^{3+1}, g)
- ▶ 2-form $F = dA$ (electromagnetism)
- ▶ Charged (complex) scalar field ϕ

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$$\nabla^\mu F_{\mu\nu} = 2\mathfrak{e} \operatorname{Im}(\phi \overline{D_\nu \phi})$$

$$g^{\mu\nu} D_\mu D_\nu \phi = 0$$

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$$T_{\mu\nu}^{\text{CSF}} = \operatorname{Re}(D_\mu \phi \overline{D_\nu \phi}) - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} D_\alpha \phi \overline{D_\beta \phi}$$

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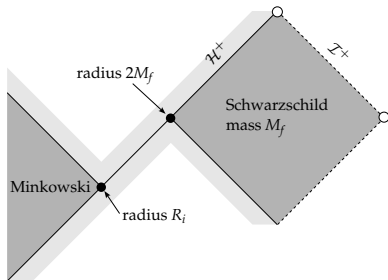
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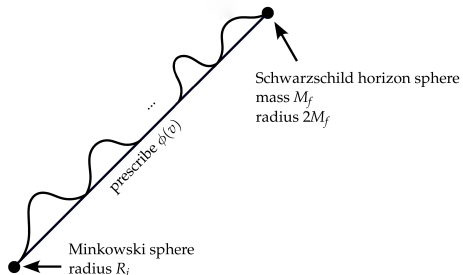
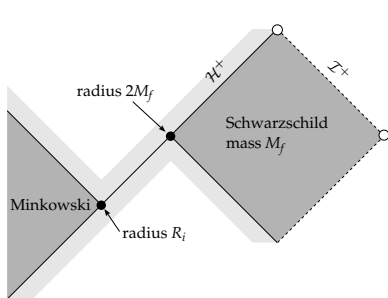
- Spherical symmetry!

PROTOTYPE: MINKOWSKI TO SCHWARZSCHILD GLUING



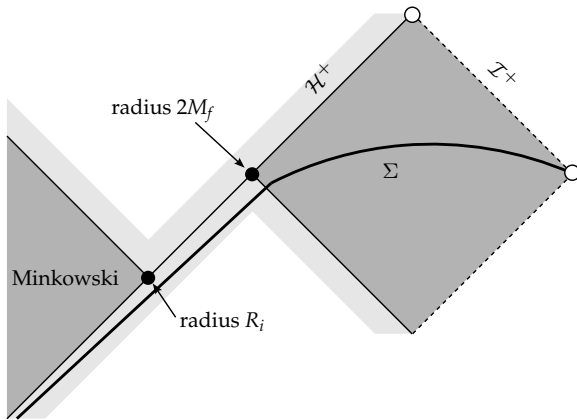
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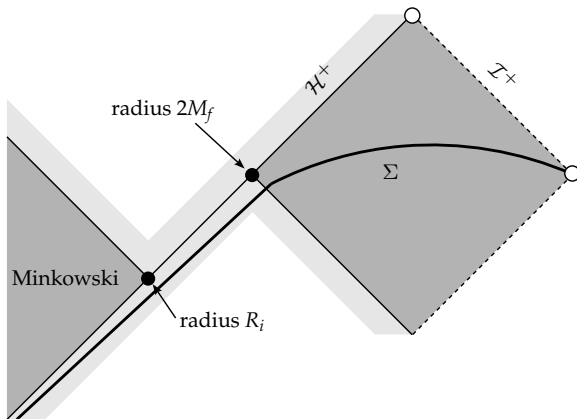
- **Goal:** Create spacetime from grav. collapse containing the above Minkowski and Schwarzschild patches.
- **Enemy:** Decoherence of waves.
- **Solution:** Characteristic gluing makes superposition of waves **purely ingoing** along a single outgoing null hypersurface, e.g. \mathcal{H}^+

CHARACTERISTIC DATA TO CAUCHY DATA



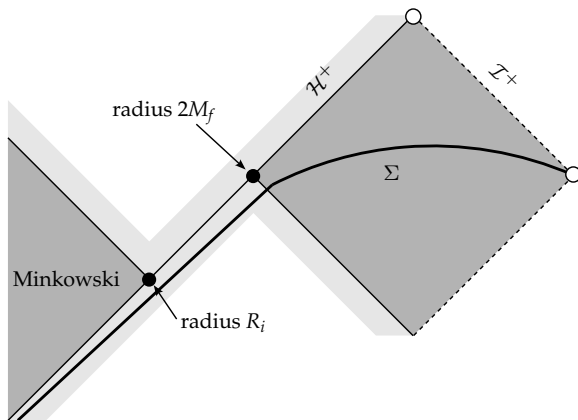
Obtain **Cauchy data** by solving backwards.

PROTOTYPE: MINKOWSKI TO SCHWARZSCHILD GLUING



Characteristic gluing for the vacuum Einstein equations recently initiated in fundamental works of Aretakis–Czimek–Rodnianski and refined by Czimek–Rodnianski, but in a **perturbative** regime around Minkowski space which is inapplicable here.

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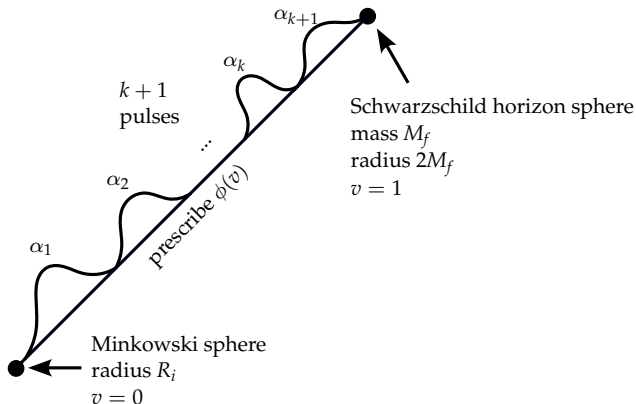
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The present problem is fundamentally **nonlinear** and **nonperturbative**, but is tractable because of the symmetry assumption.

MINKOWSKI TO SCHWARZSCHILD GLUING

Theorem (K–Unger '22).

For any $k \in \mathbb{N}$ and $0 < R_i < 2M_f$, the Minkowski sphere of radius R_i can be characteristically glued to the Schwarzschild event horizon sphere with mass M_f to order C^k within the Einstein-scalar field model in spherical symmetry.



ASPECTS OF THE PROOF: SCHWARZSCHILD

- Null constraint system: coupled nonlinear system of ODEs sourced by $\phi(v)$

ASPECTS OF THE PROOF: SCHWARZSCHILD

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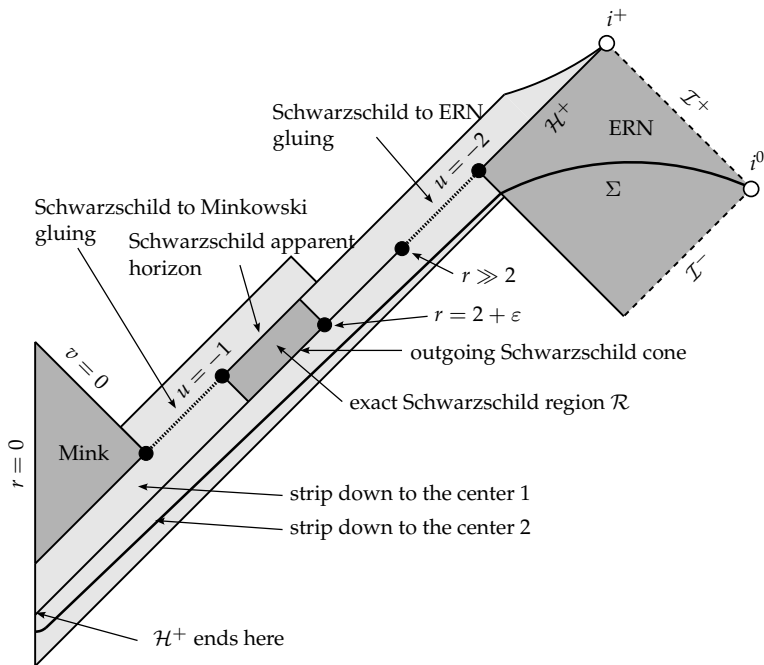
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- ▶ Charged scalar field: sphere of α 's is deformed topological sphere ensuring the “charge” condition.

DISPROOF OF THE THIRD LAW



THE THIRD LAW IN THE VACUUM CASE

Conjecture.

There exist Cauchy data for the Einstein vacuum equations

$$R_{\mu\nu} = 0$$

*which form an exactly Schwarzschild apparent horizon, only for the spacetime to form an **exactly extremal Kerr event horizon** at a later advanced time. In particular, **already in vacuum**, the “third law of black hole thermodynamics” is **false**.*

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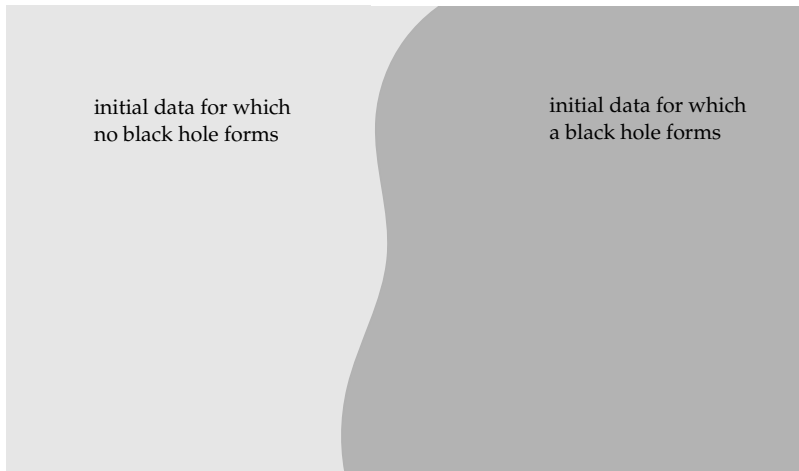
Theorem (K.–Unger, '23).

For any $0 \leq |a| \ll M$, there exist Cauchy data for the Einstein vacuum equations

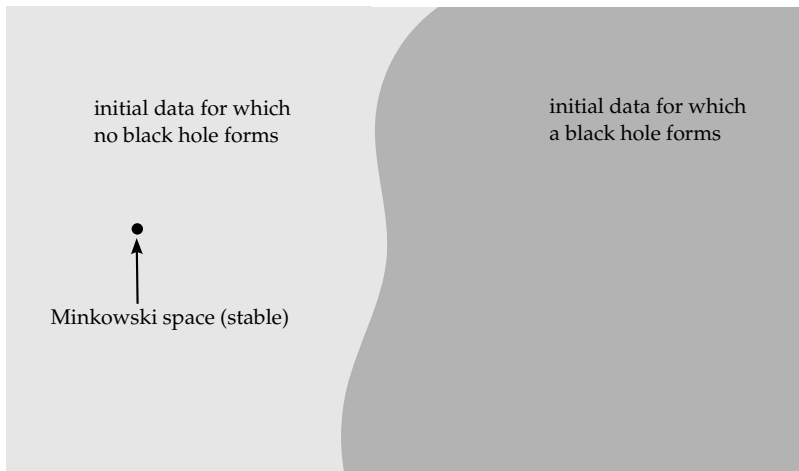
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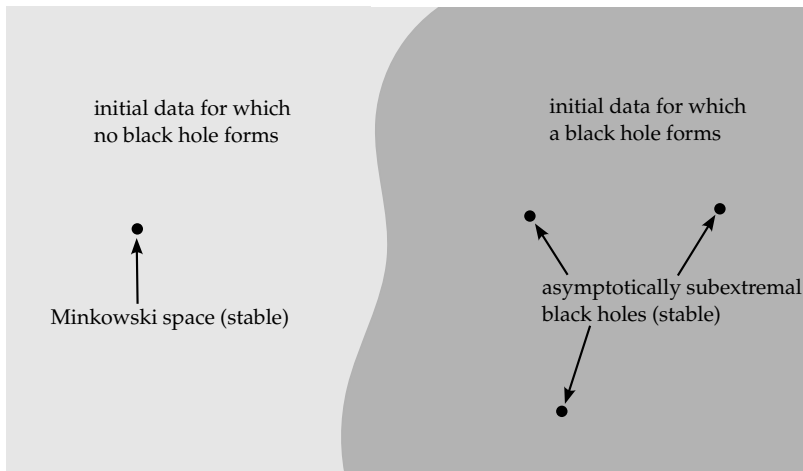
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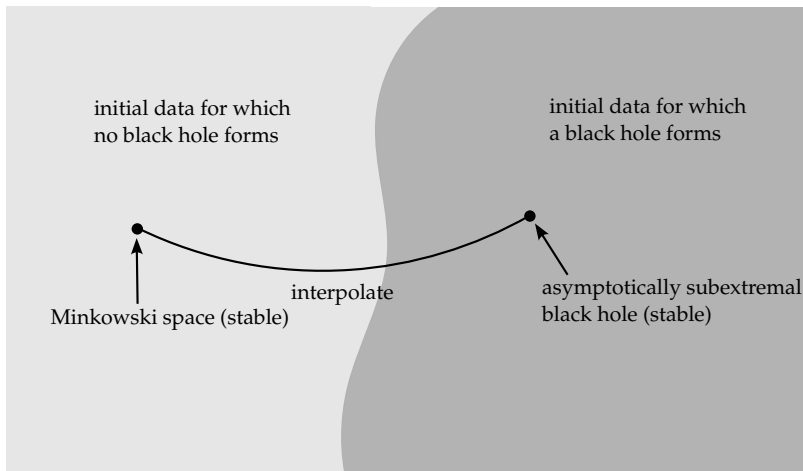
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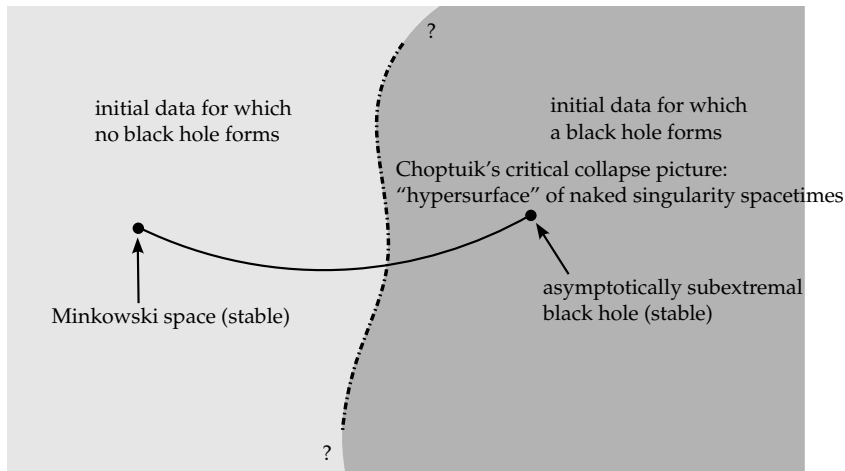
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Additional critical behavior: Some extremal black holes lie at the threshold!

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We propose to first study this phenomenon in the Einstein–Maxwell–massless Vlasov model—inspired by classical work of A. Ori.

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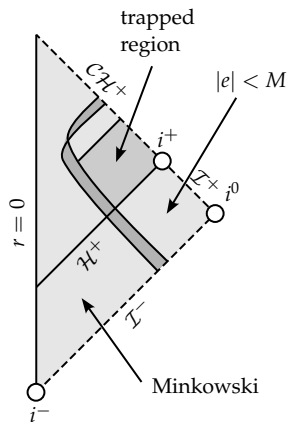
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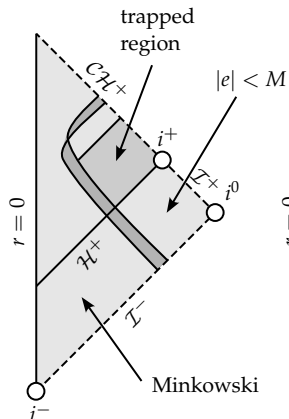


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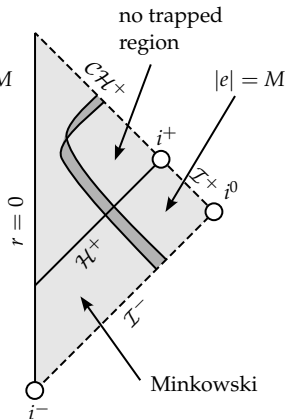
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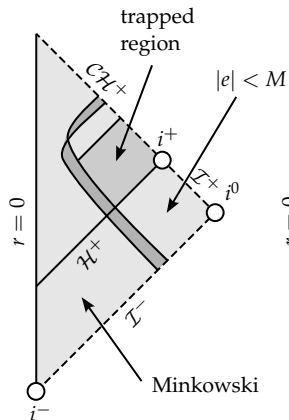


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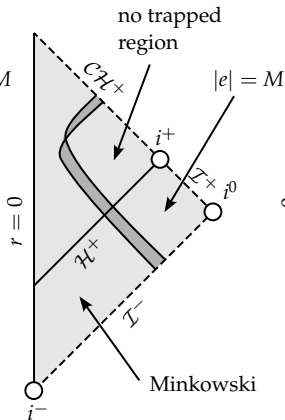
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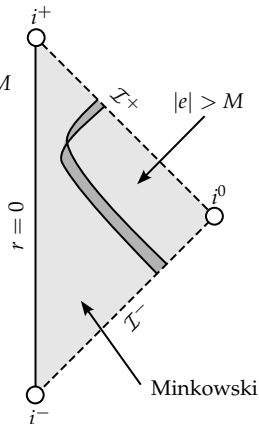
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