

PERSPECTIVES OF COLLINEAR LASER SPECTROSCOPY

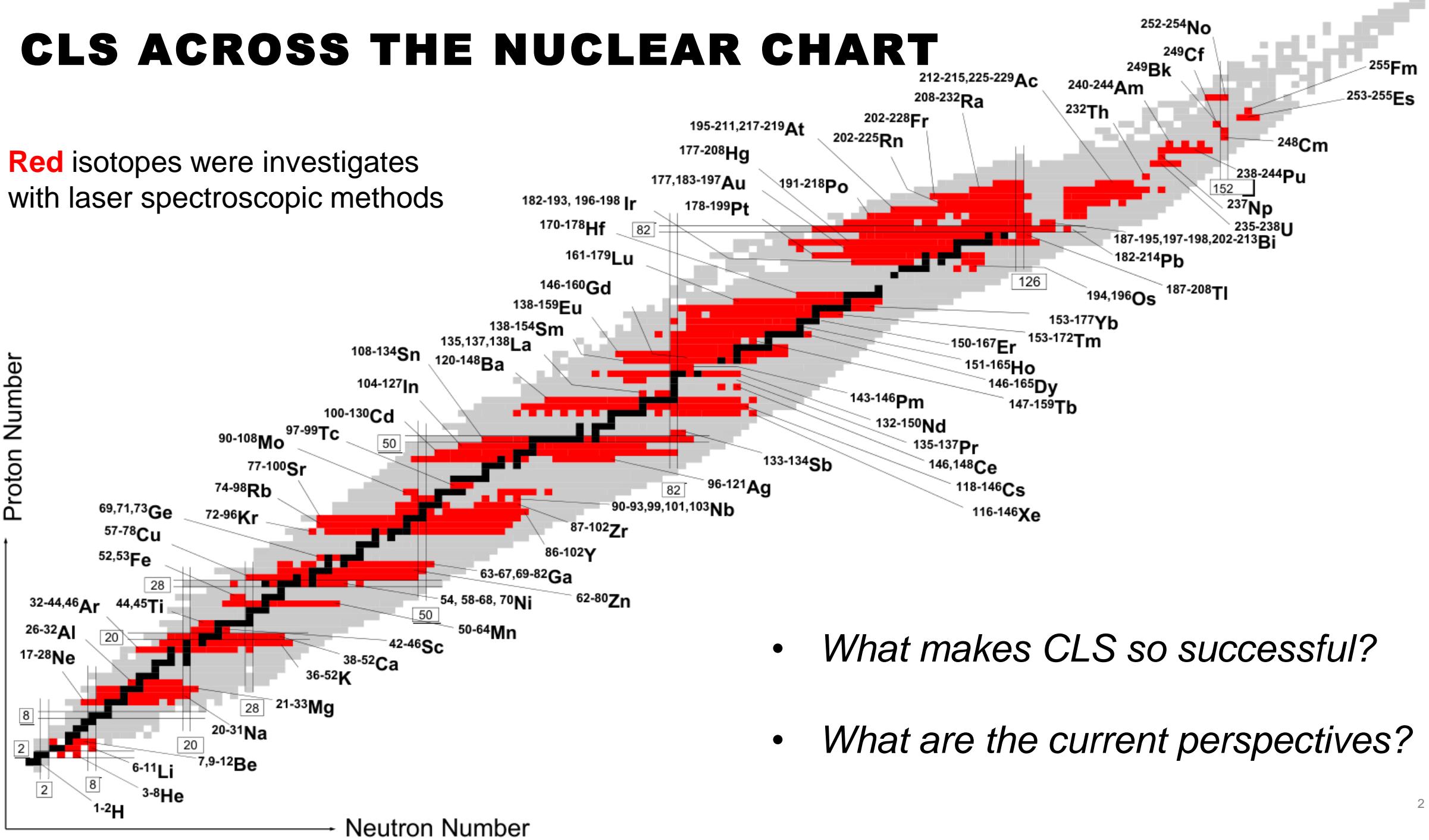


<https://www.triumf.ca/laser-spectroscopy>

CLS ACROSS THE NUCLEAR CHART

Red isotopes were investigated with laser spectroscopic methods

Proton Number ↑

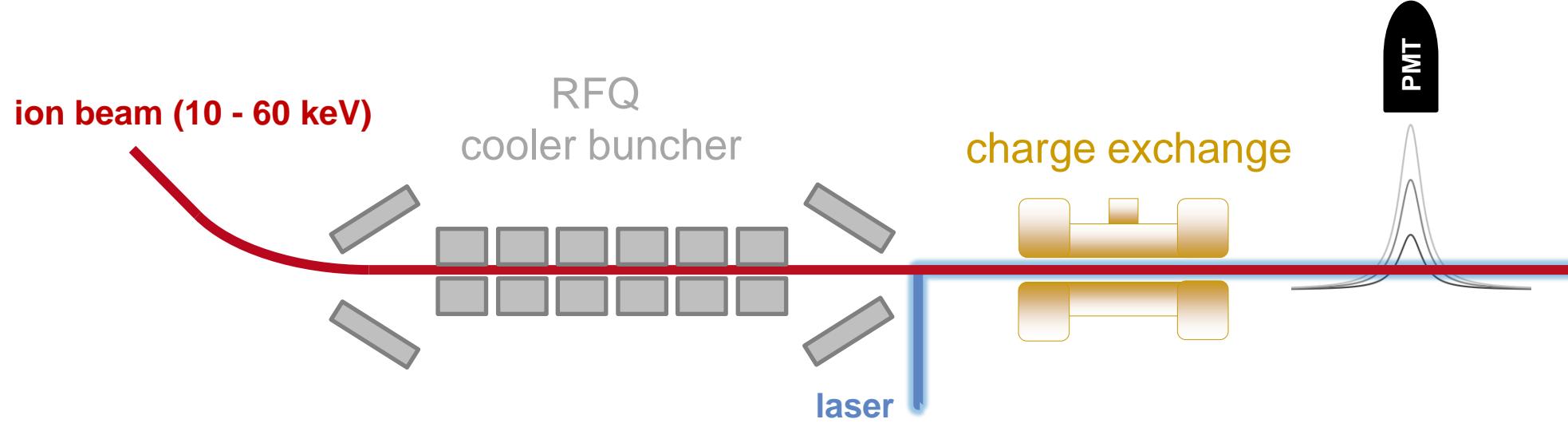


- What makes CLS so successful?
- What are the current perspectives?

COLLINEAR LASER SPECTROSCOPY



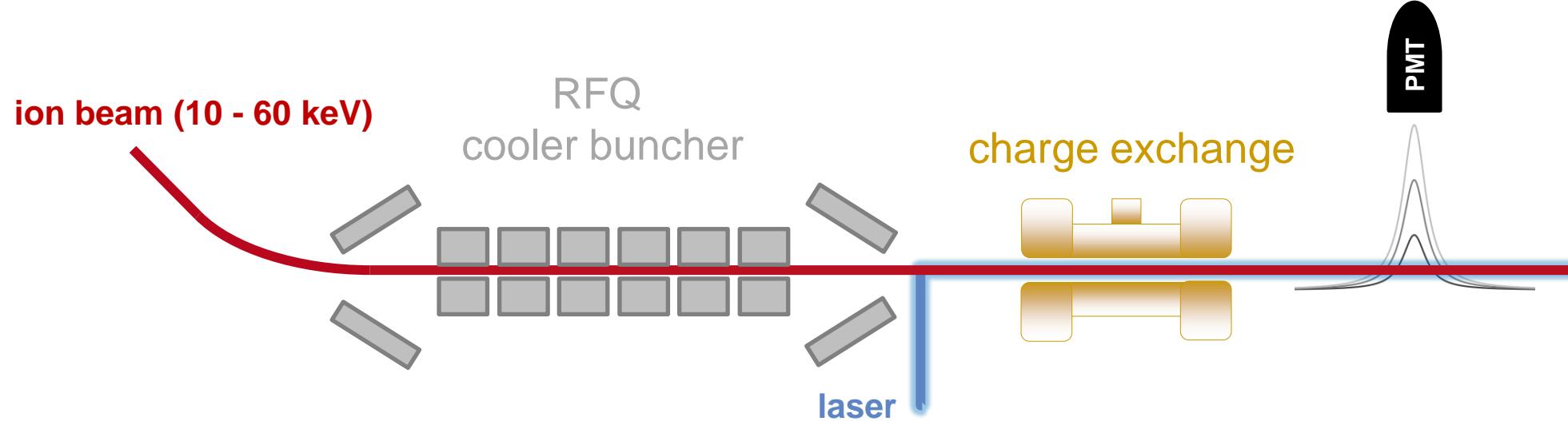
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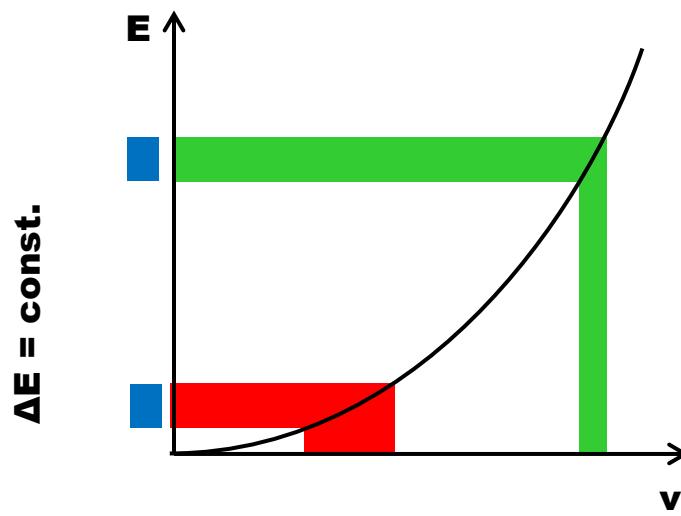
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Doppler compression



$$\frac{\delta v_U}{\delta v_0} = \frac{1}{2} \sqrt{\frac{k_B T}{eU}}$$

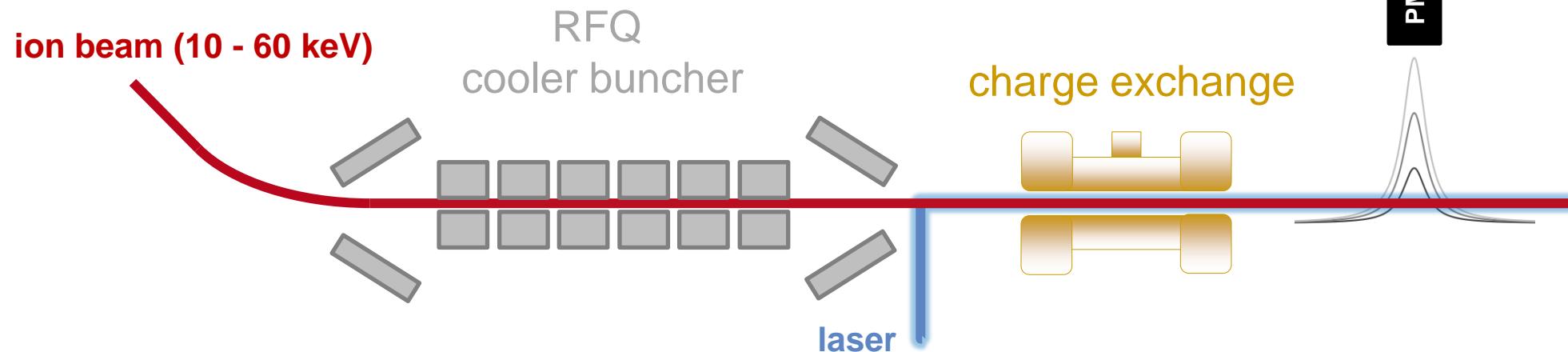
$$\nu_{c/a} = \nu_0 \gamma (1 \pm \beta)$$

- Fast – Lifetimes of > 10 ms
- Doppler broadening free
- Only dipole-allowed transitions
- Doppler-shifted resonance

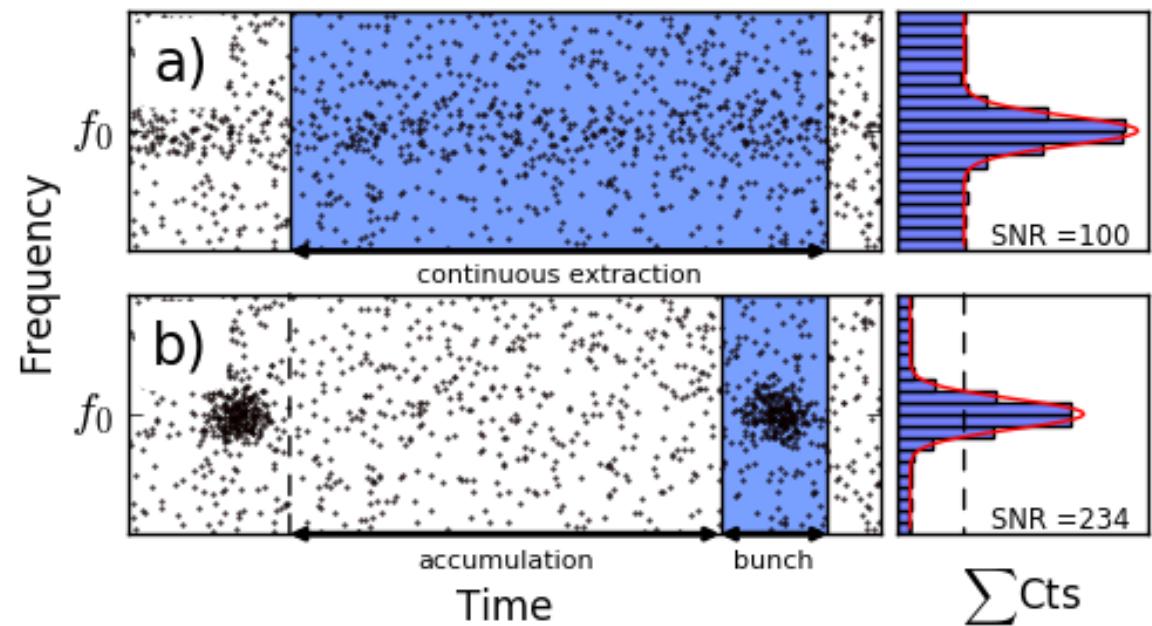
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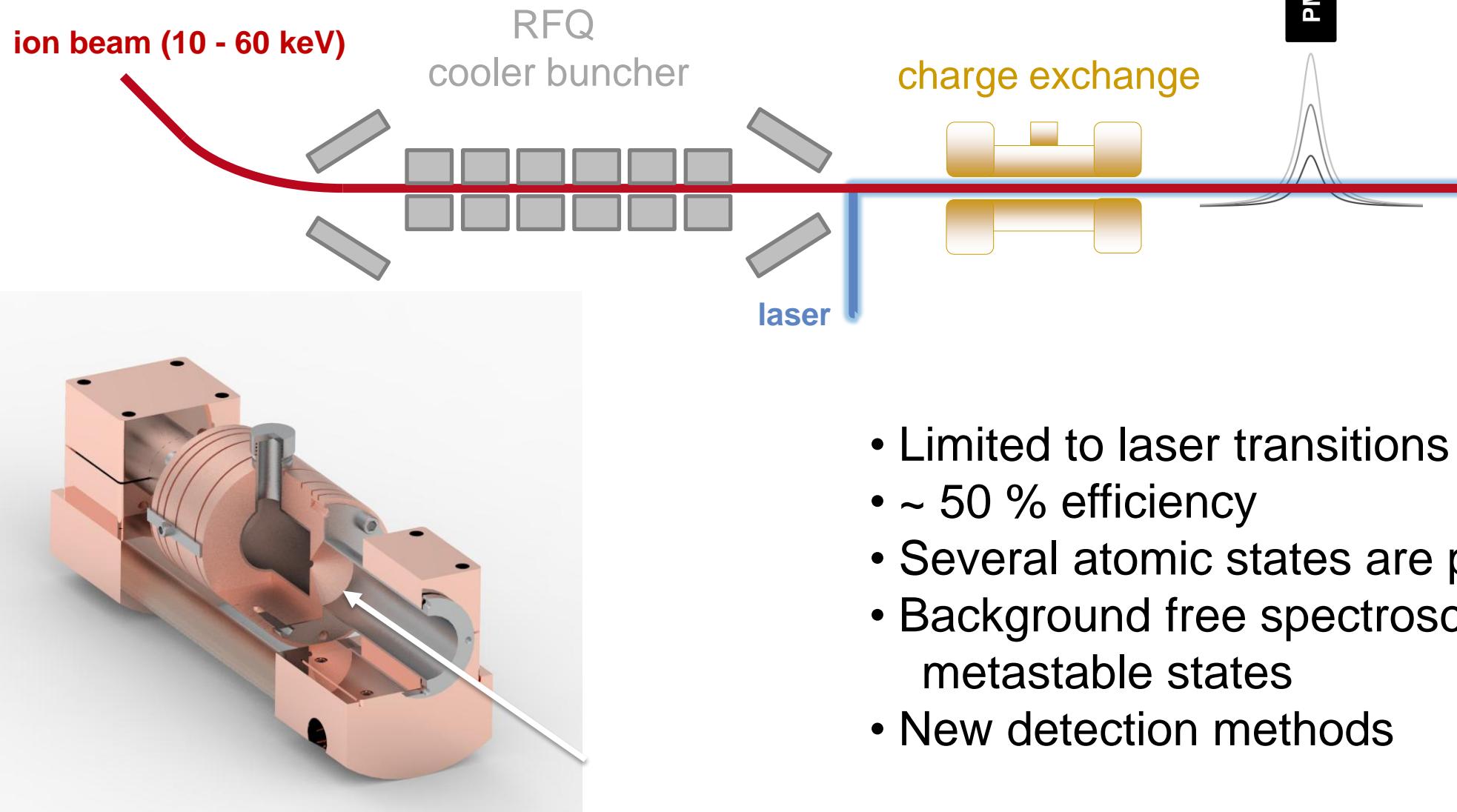
- Cooling of the ion beam
- 10^6 reduction of laser background
- Usage of pulsed spectroscopy lasers
- Optical population transfer



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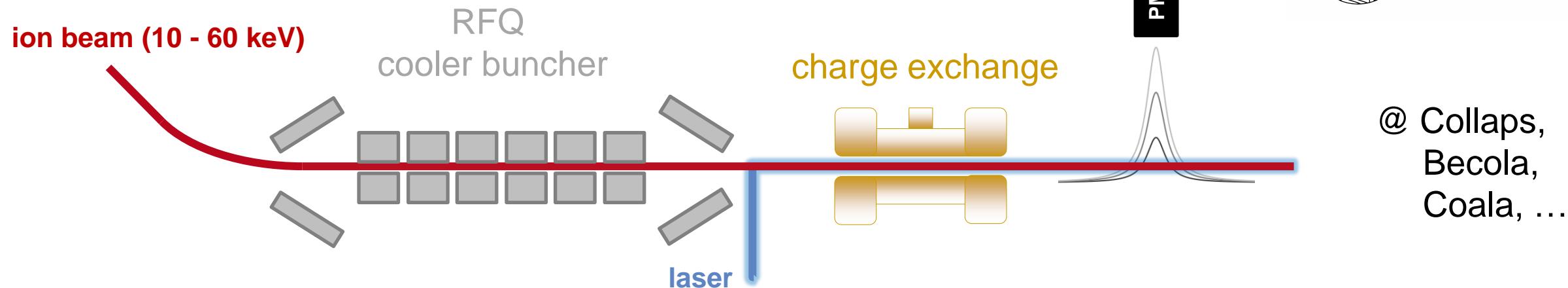


- Limited to laser transitions > 200 nm
- $\sim 50\%$ efficiency
- Several atomic states are populated
- Background free spectroscopy from metastable states
- New detection methods

COLLINEAR LASER SPECTROSCOPY



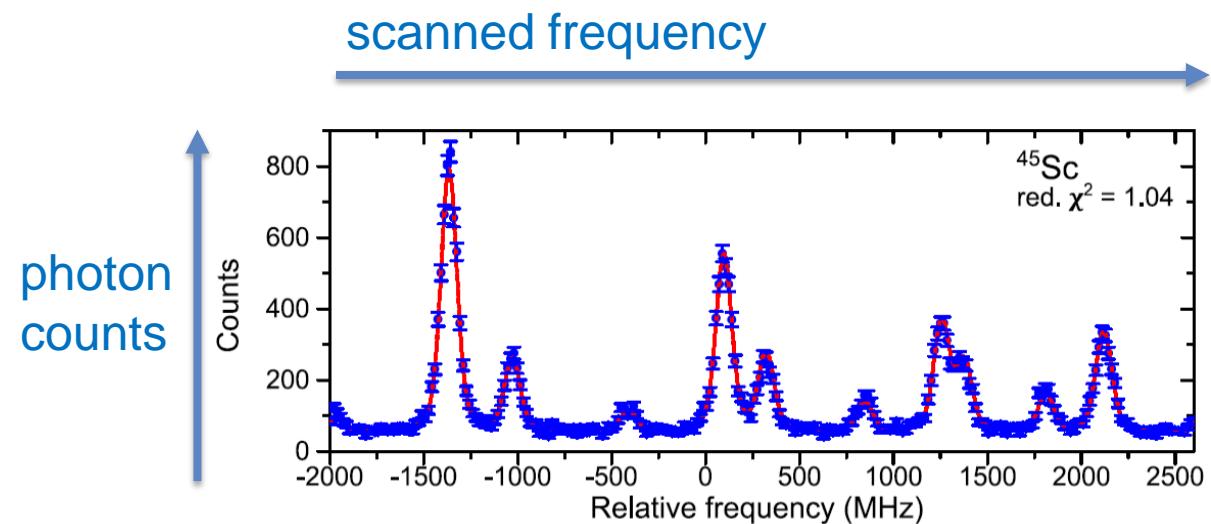
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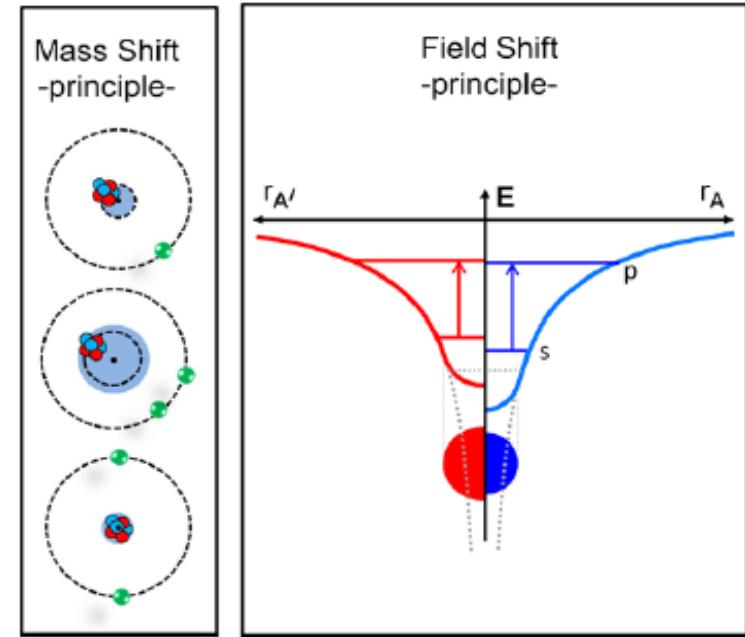
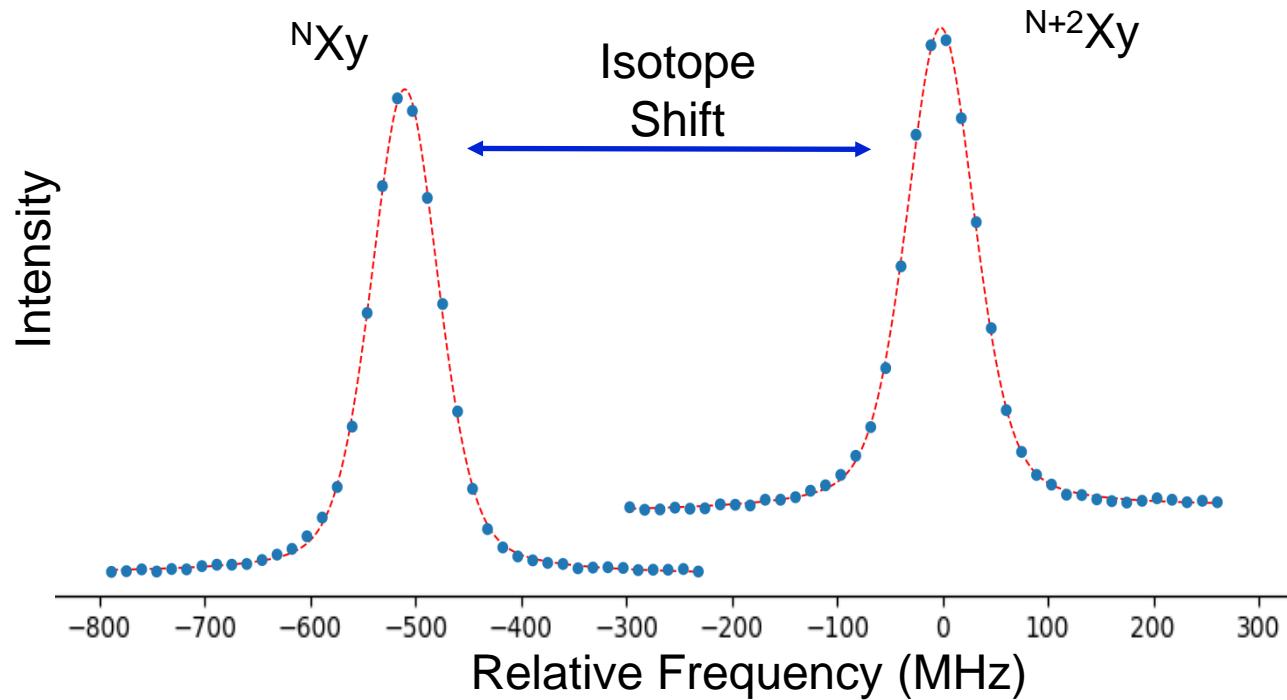
@ Collaps,
Becola,
Coala, ...

$$\Delta E_{\text{HFS}} = A \cdot \frac{C}{2} + B \cdot \frac{\frac{3}{4}C(C+1) - I(I+1)J(J+1)}{2I(I-1)J(2J-1)}$$

$$C = F(F+1) - I(I+1) - J(J+1)$$



NUCLEAR CHARGE RADII



Nörterhäuser and Ch. Geppert, Lecture Notes in Physics 879 (2014)

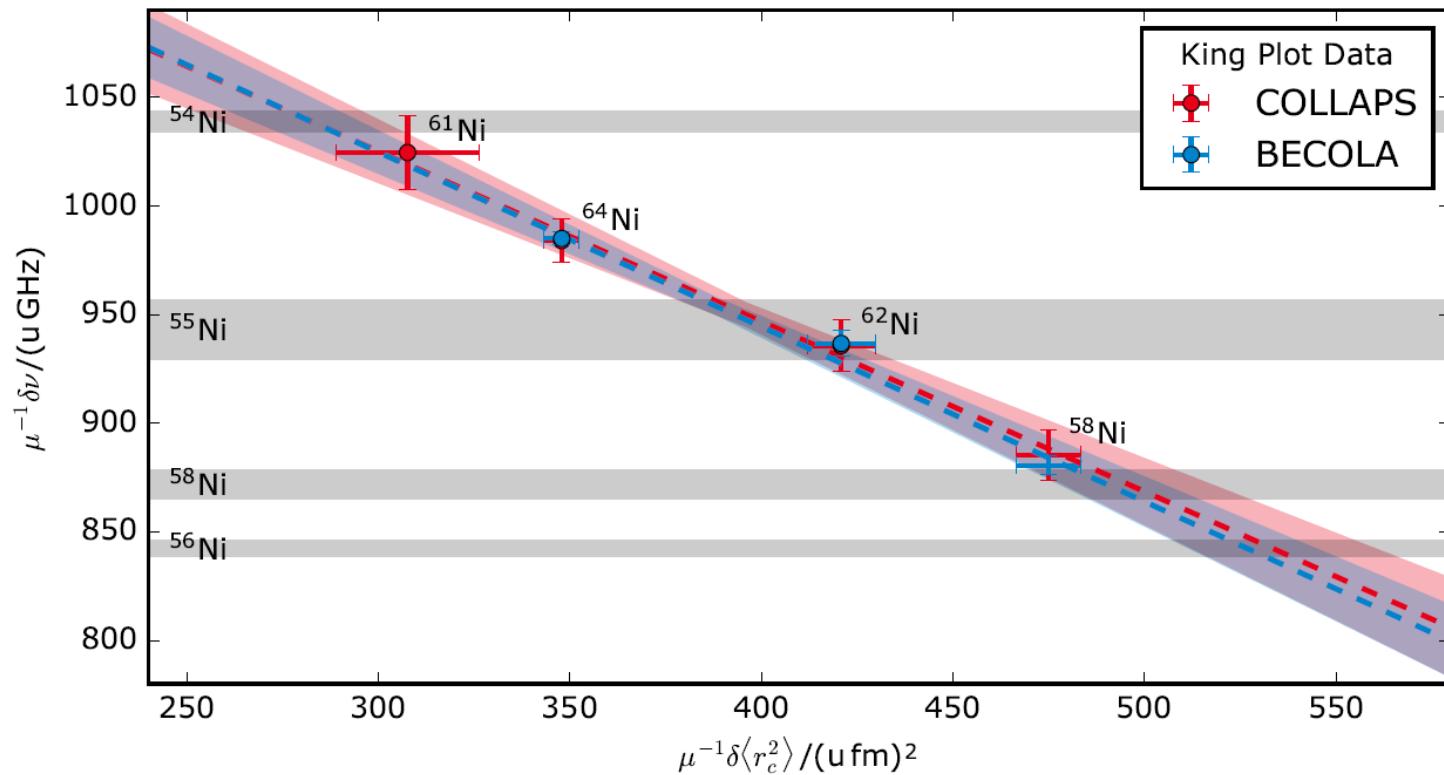
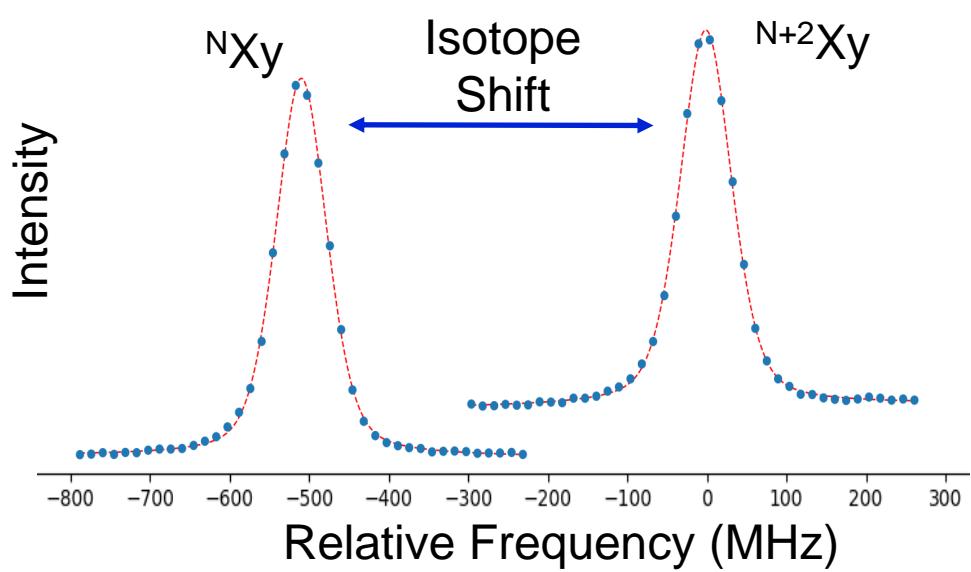
$$\delta\nu_i^{A,A'} = \nu_i^{A'} - \nu_i^A = F_i \delta\langle r^2 \rangle^{A,A'} + M_i \frac{m'_A - m_A}{m'_A m_A}$$

M_i and F_i from theory or from experiment if ≥ 3 stable isotopes (King plot)

NUCLEAR CHARGE RADII

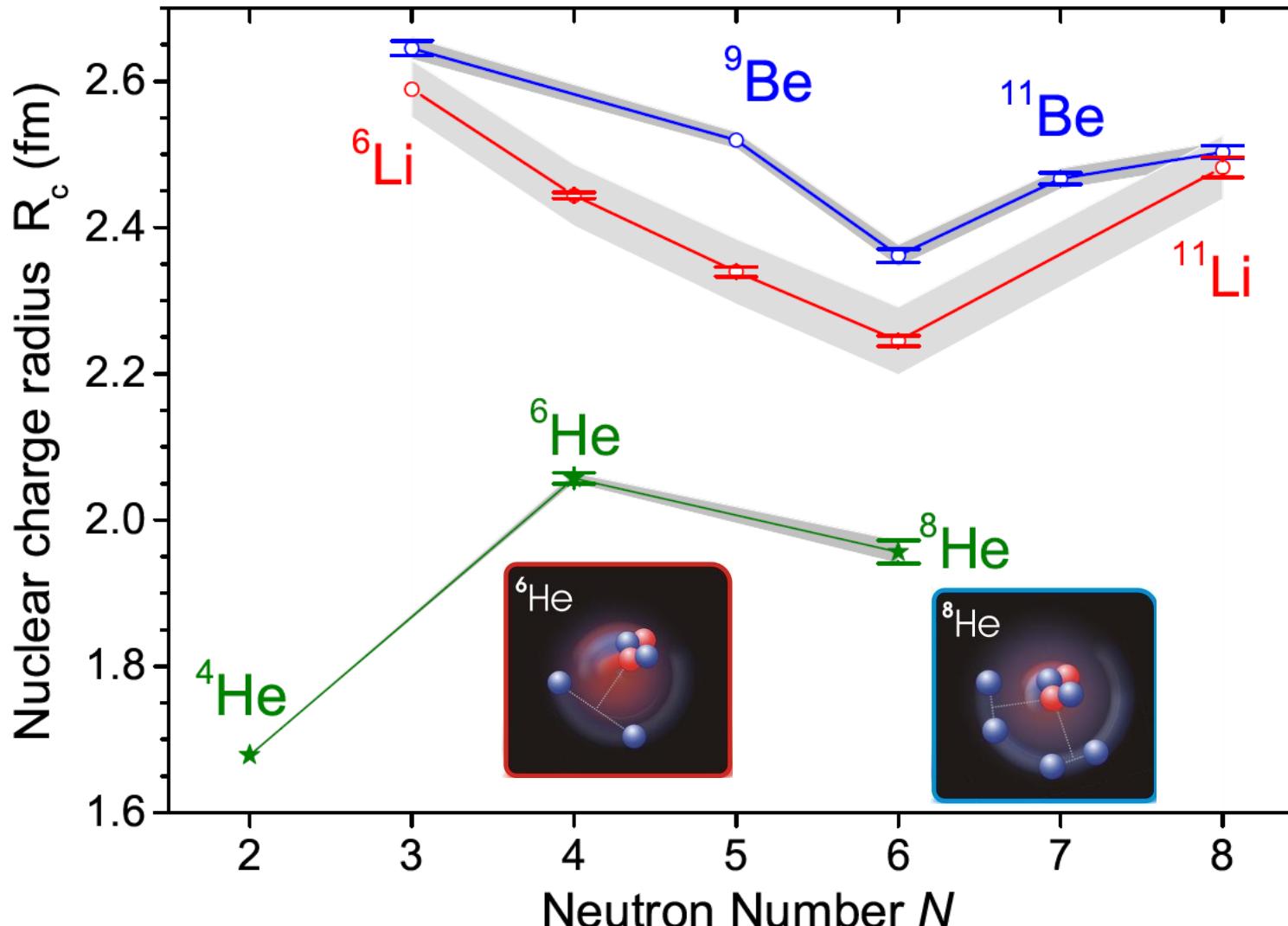


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$$\delta\nu_i^{A,A'} = \nu_i^{A'} - \nu_i^A = F_i \delta\langle r^2 \rangle^{A,A'} + M_i \frac{m'_A - m_A}{m'_A m_A}$$

NUCLEAR CHARGE RADII LIGHT NUCLEI



$R_c(^9\text{Be}) = 2.519(12) \text{ fm}$
Jansen et al., Nucl. Phys. A 188, 337 (1972)

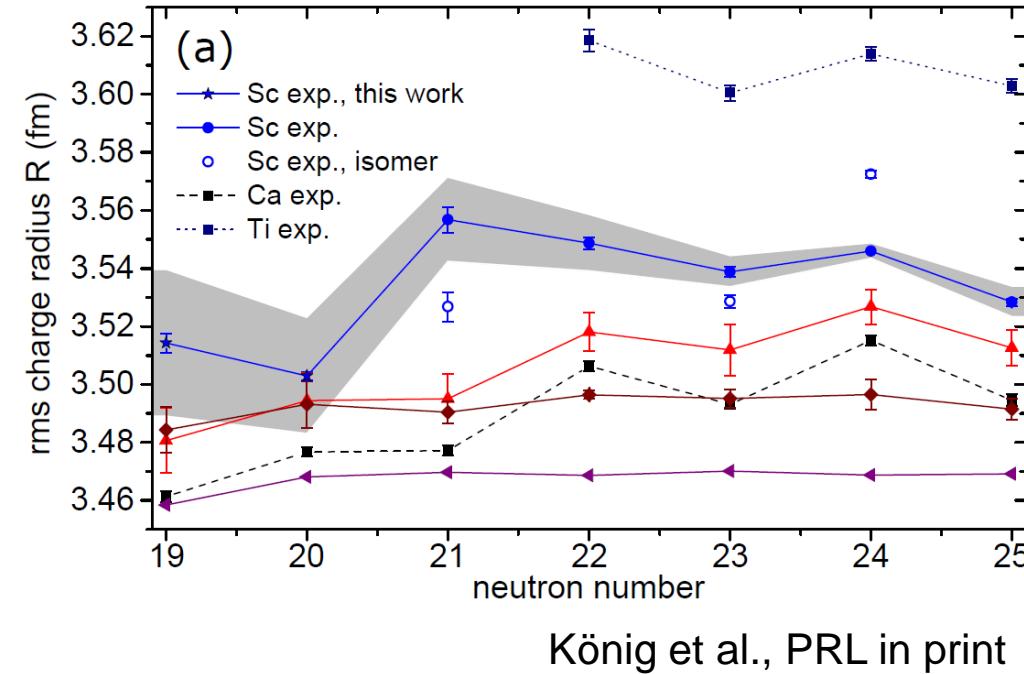
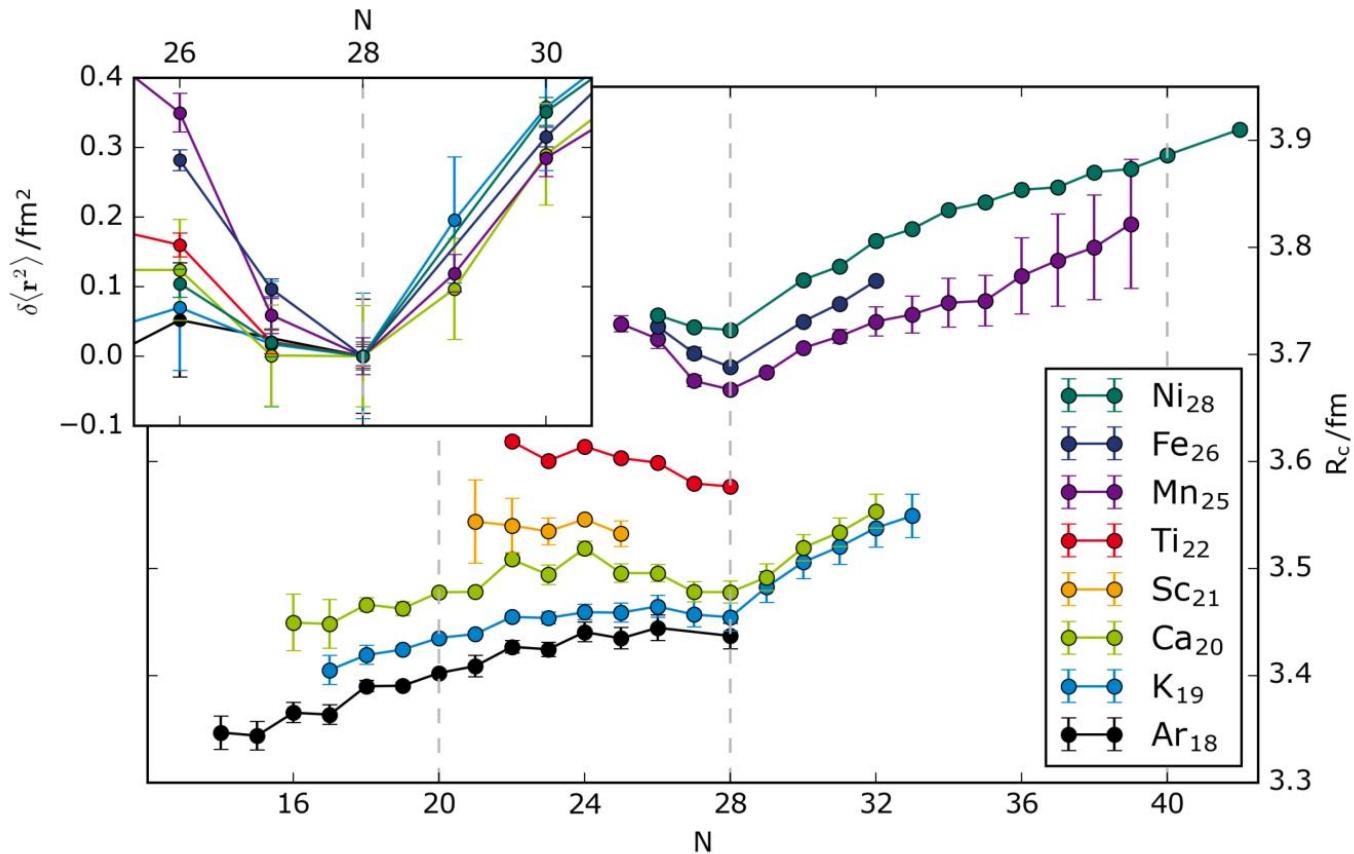
$R_c(^6\text{Li}) = 2.589(39) \text{ fm}$
Nörtershäuser et al., Phys. Rev. C 84, 024307 (2011)

$r_\alpha = 1.678\,24(83) \text{ fm}$
Krauth et al., Nature 589, 527 (2021)
from μHe^+ (one-electron system)
“all-optical charge radius”

NUCLEAR CHARGE RADII THE CALCIUM – NICKEL REGION



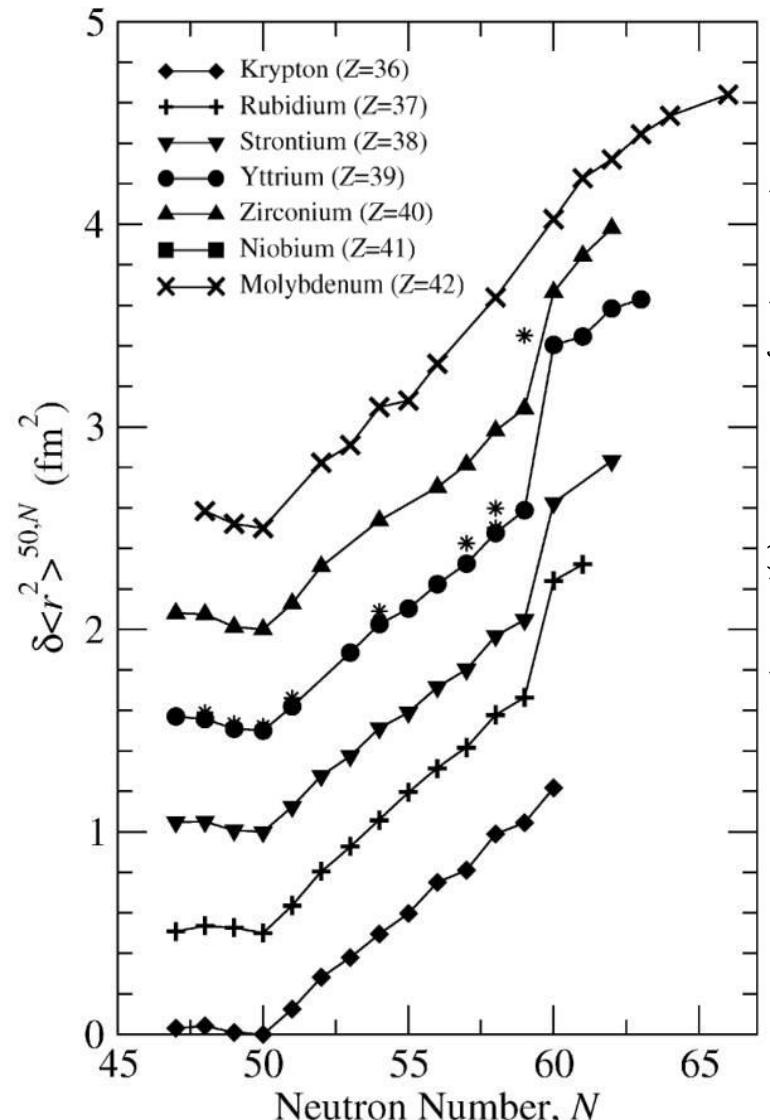
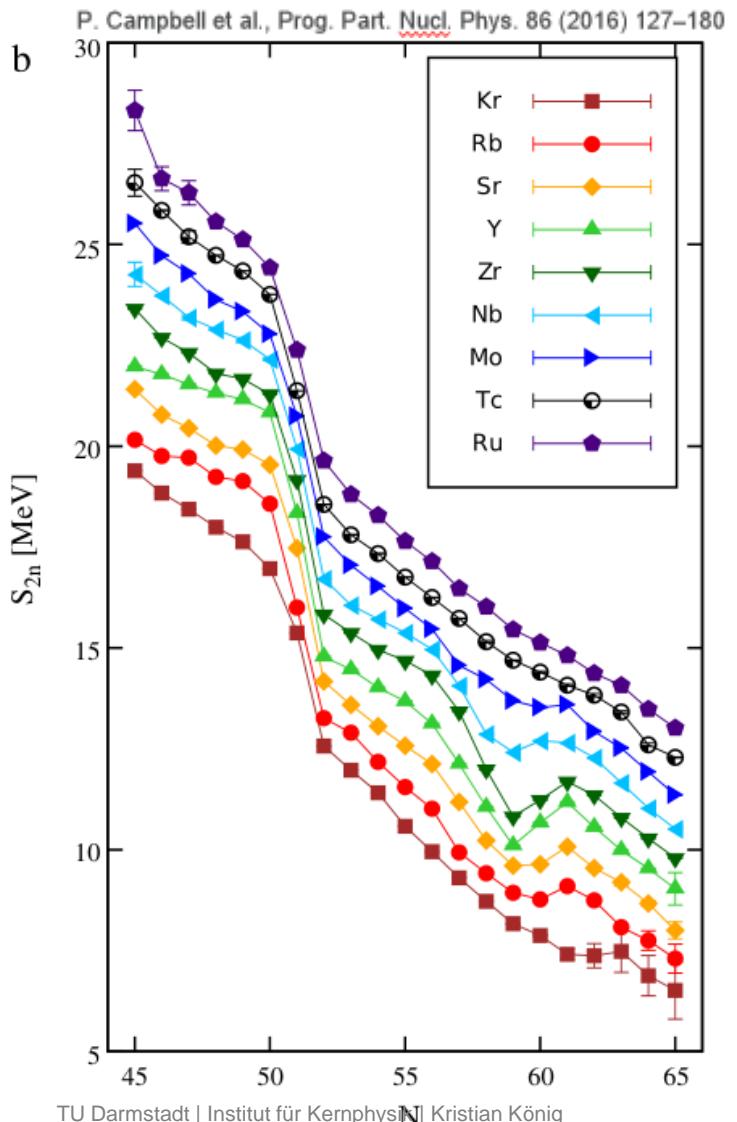
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König et al., PRL in print

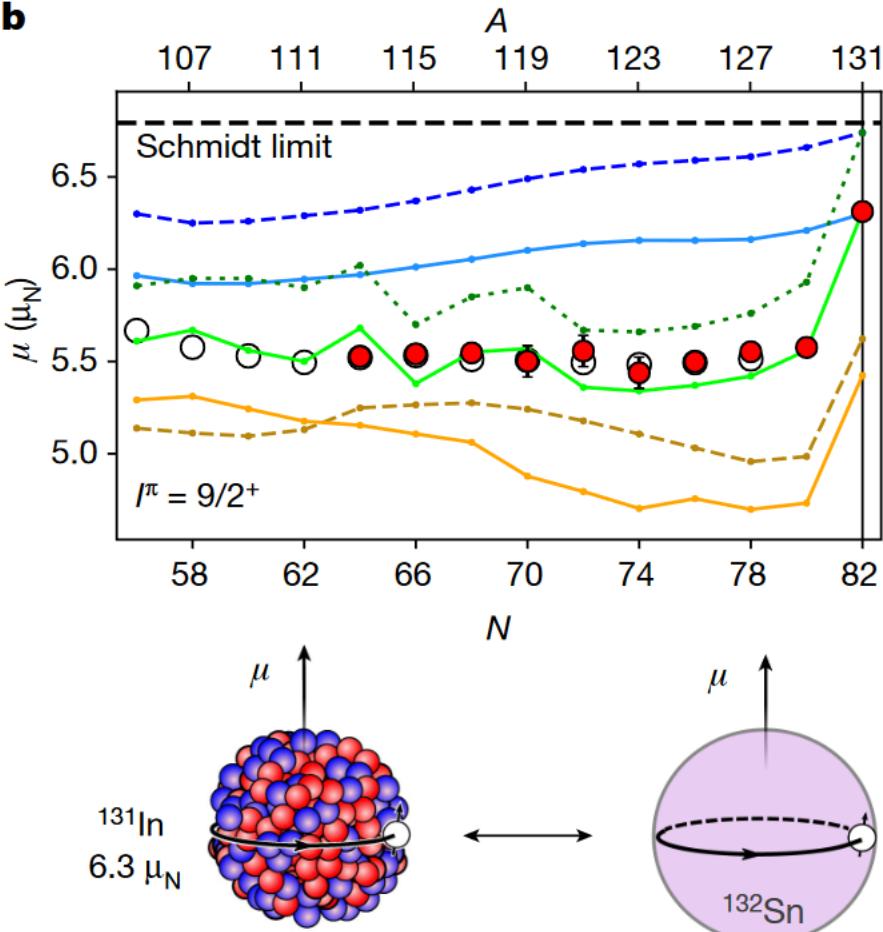
NUCLEAR CHARGE RADII THE STRONTIUM REGION

- Evolution of nuclear charge radii along a chain of isotopes
- Correlation with S_{2n} - two-neutron separation energy
- Shows structural effects
- Onset of deformation at $N=60$
- Is this effect still visible for Mo, Tc, Ru, Rh?
- Observed „change in slope“ for $N>60$

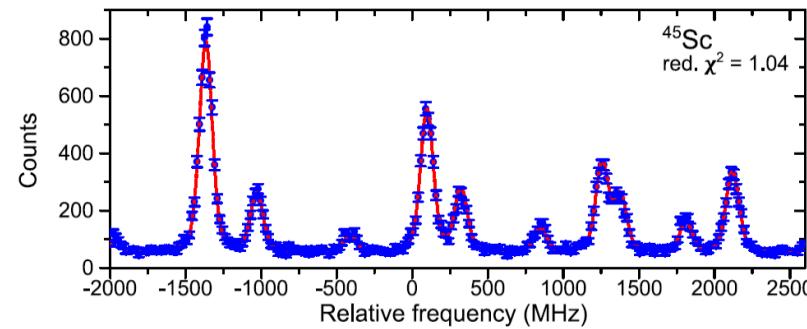


NUCLEAR MOMENTS IN INDIUM ISOTOPES

b



A. Vernon et al., Nature 607, 260 (2022)



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$$\Delta E_{\text{HFS}} = A \cdot \frac{C}{2} + B \cdot \frac{\frac{3}{4}C(C+1) - I(I+1)J(J+1)}{2I(I-1)J(2J-1)}$$

$$C = F(F+1) - I(I+1) - J(J+1)$$

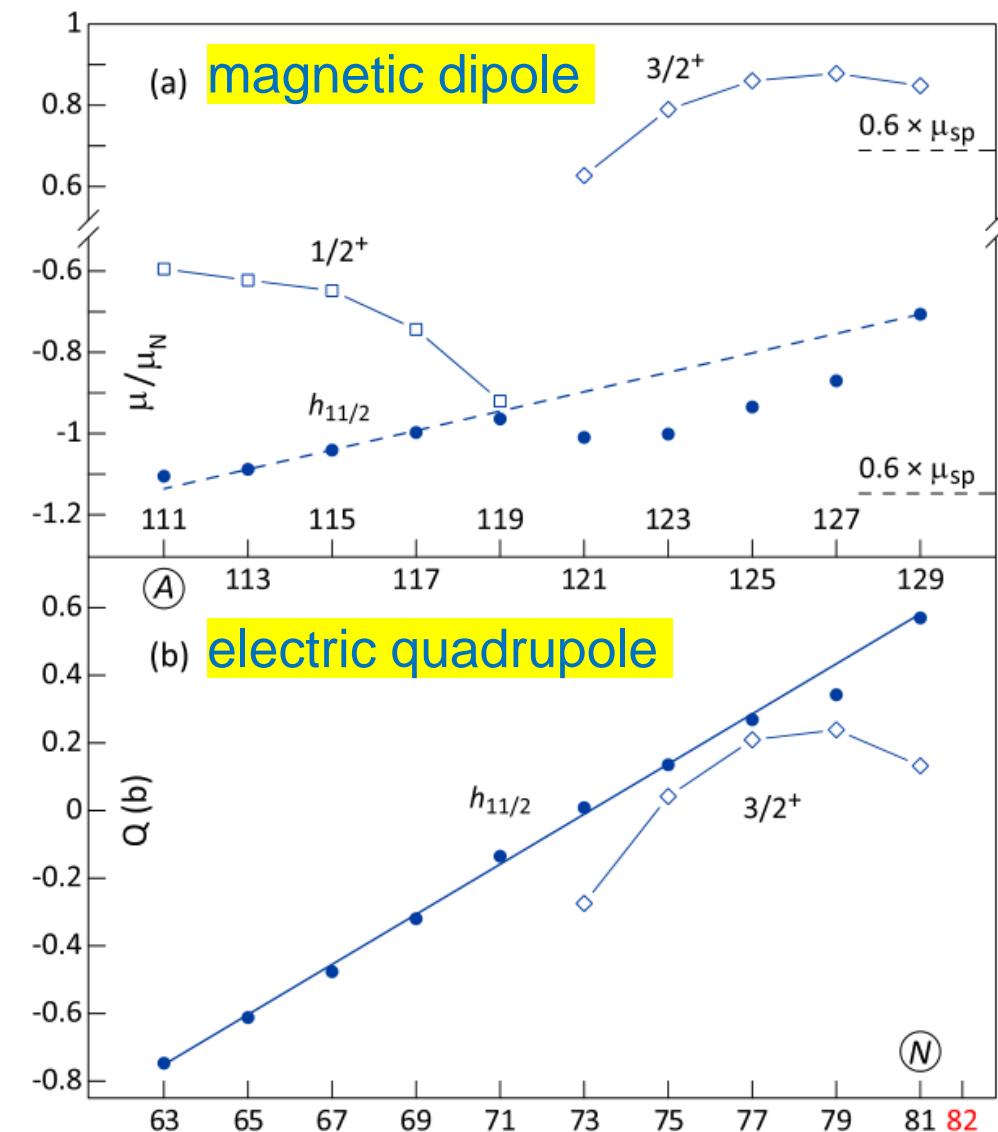
- Spin I , magnetic dipole moment μ and electric quadrupole moment Q can be deduced from hyperfine splitting

$$\mu_{\text{meas}} = \mu_{\text{ref}} \frac{A_{\text{meas}}^{\text{hf}}}{A_{\text{ref}}^{\text{hf}}} \frac{I_{\text{meas}}}{I_{\text{ref}}}$$

- μ probes the single-particle nature of the valence nucleon

NUCLEAR MOMENTS IN CADMIUM ISOTOPES

- Linear trend in magnetic dipole and electrical quadrupole moments for cadmium $I=11/2$ states
- Large effective charge of the one neutron (Core polarizability for the Z-2 magic nucleus)
- Trend continues across 10 isotopes, more than the volume of the occupied $h_{11/2}$ shell
- Paired neutrons occupy neighbouring states, valence neutron is always in $h_{11/2}$
- Follows remarkably simple trend in a complex nuclei
- For Sn (magic) the QP moment is not linear – the core is not as polarizable

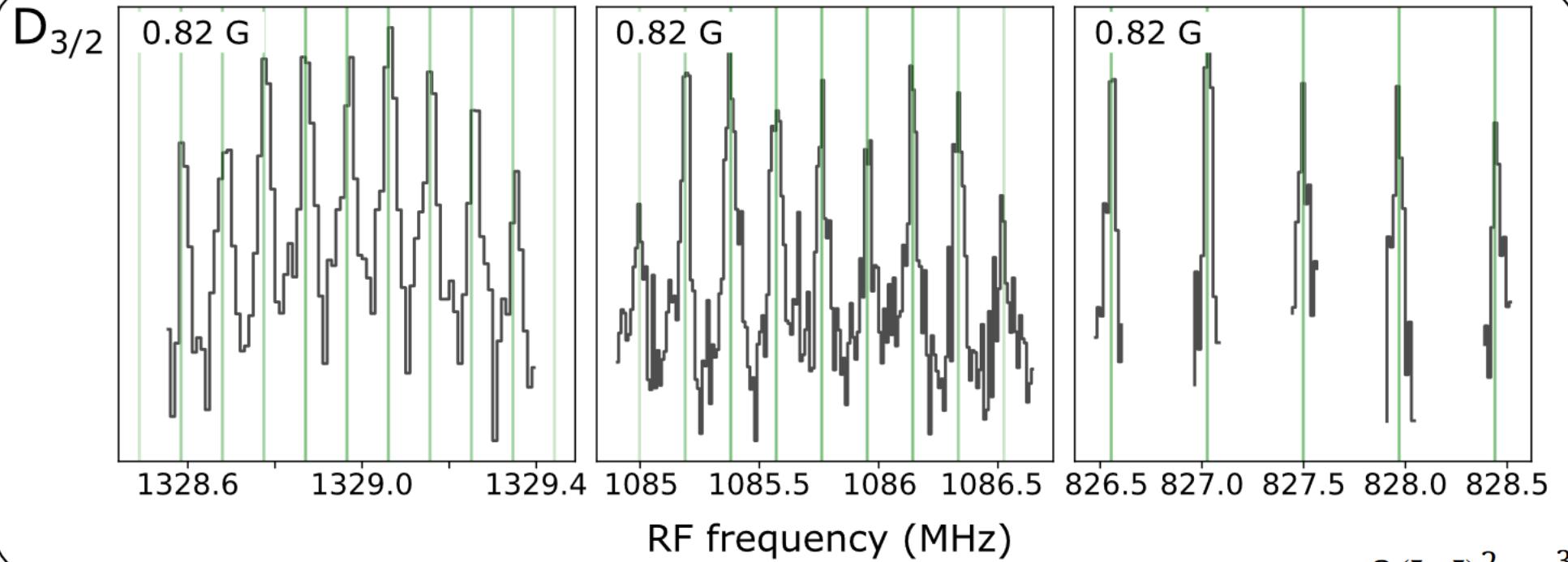


Yordanov *et al.*, Phys. Rev. Lett 110, 192501 (2013)

OCTUPOLE MOMENTS



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R. de Groote, Physics Letters B 827 (2022) 136930

$$\begin{aligned} \mathcal{H}_{\text{hyp}} = & A \mathbf{I} \cdot \mathbf{J} + B \frac{3(\mathbf{I} \cdot \mathbf{J})^2 + \frac{3}{2}(\mathbf{I} \cdot \mathbf{J}) - I(I+1)J(J+1)}{2I(2I-1)J(2J-1)} \\ & + C \left[\frac{10(\mathbf{I} \cdot \mathbf{J})^3 + 20(\mathbf{I} \cdot \mathbf{J})^2}{I(I-1)(2I-1)J(J-1)(2J-1)} \right. \\ & \left. + \frac{2\mathbf{I} \cdot \mathbf{J}\{I(I+1) + J(J+1) - 3N + 3\} - 5N}{I(I-1)(2I-1)J(J-1)(2J-1)} \right], \end{aligned}$$

IMPLICATIONS FOR NUCLEAR EOS



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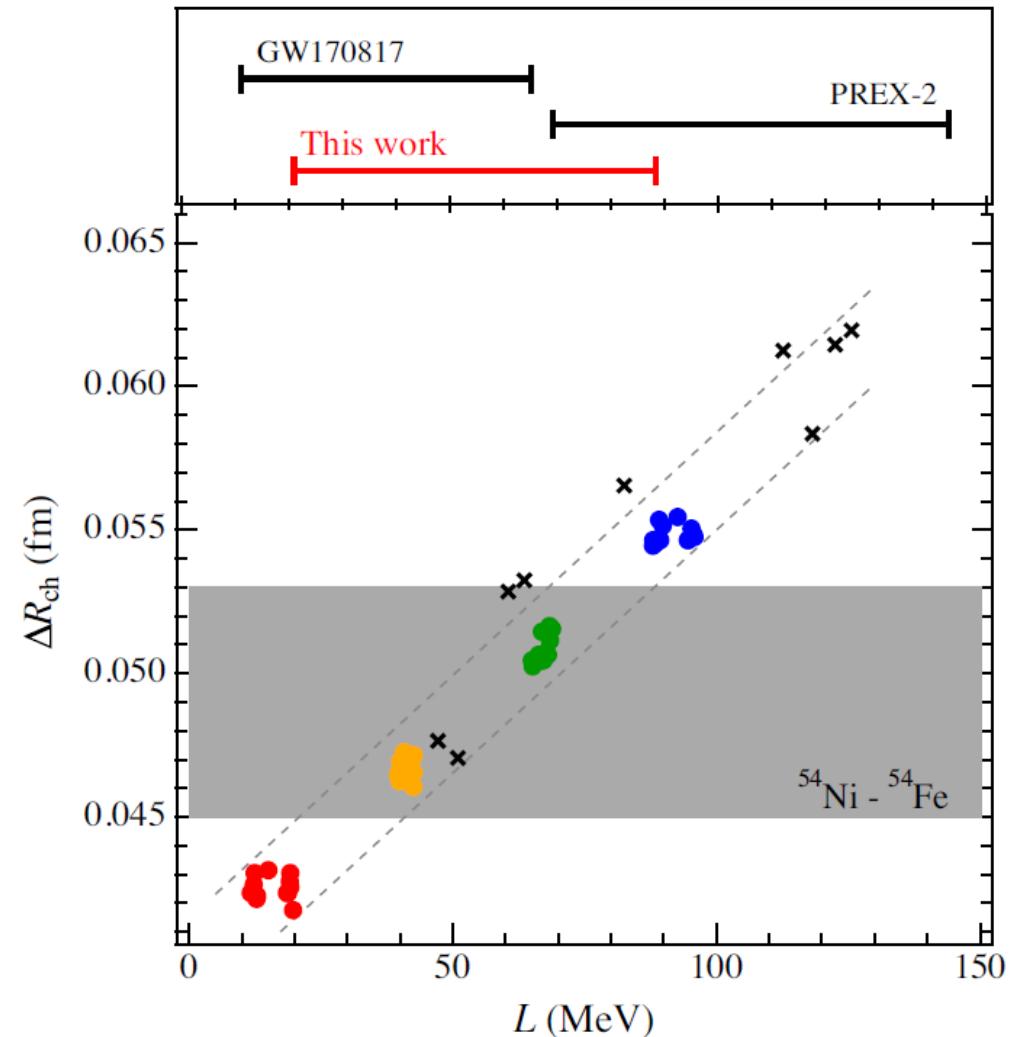
Constrain the slope of the symmetry energy at nuclear saturation density (L)

Important to predict the properties of both super-heavy nuclei and neutron stars

$$L = 3\rho_0 \left[\frac{\partial E_{sym}(\rho)}{\partial \rho} \right]_{\rho=\rho_0}$$

$$R_{skin}(N, Z) = R_p(N, Z) - R_p(Z, N) \equiv \Delta R_{murr}(Z, N)$$

- ⇒ Favoring a soft neutron matter EOS
- ⇒ Good agreement with new CREX and most theoretical results



Pineda et al., Phys. Rev. Lett. 127, 182503 (2021)

INPUT FOR CKM MATRIX



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CKM matrix differs $\approx 2 \sigma$ from unitarity

$$V_{ud} = G_V/G_F$$

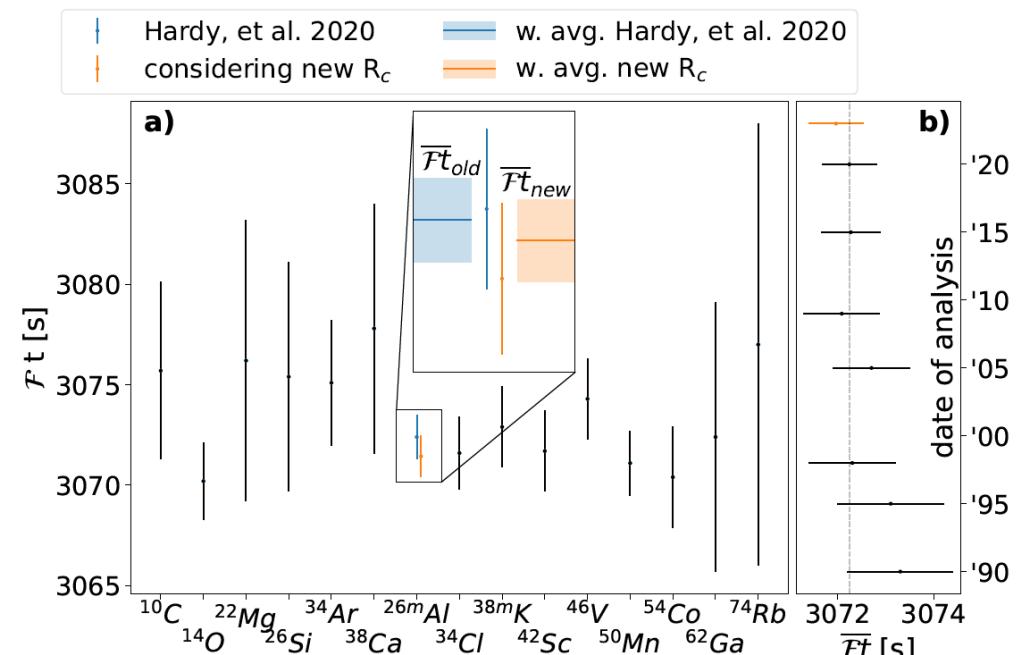
At present, superallowed $0+ \rightarrow 0+$ nuclear β decays remain the most precise way to access V_{ud}

$$\mathcal{F}t \equiv ft(1 + \delta'_R)(1 + \delta_{NS} - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R^V)}$$

Charge radius essential input parameter for isospin-symmetry breaking (ISB) corrections

^{74}Rb : Mane et al., PRL 107, 212502 (2011)

^{26m}Al : Plattner et al., submitted to PRL (2023)

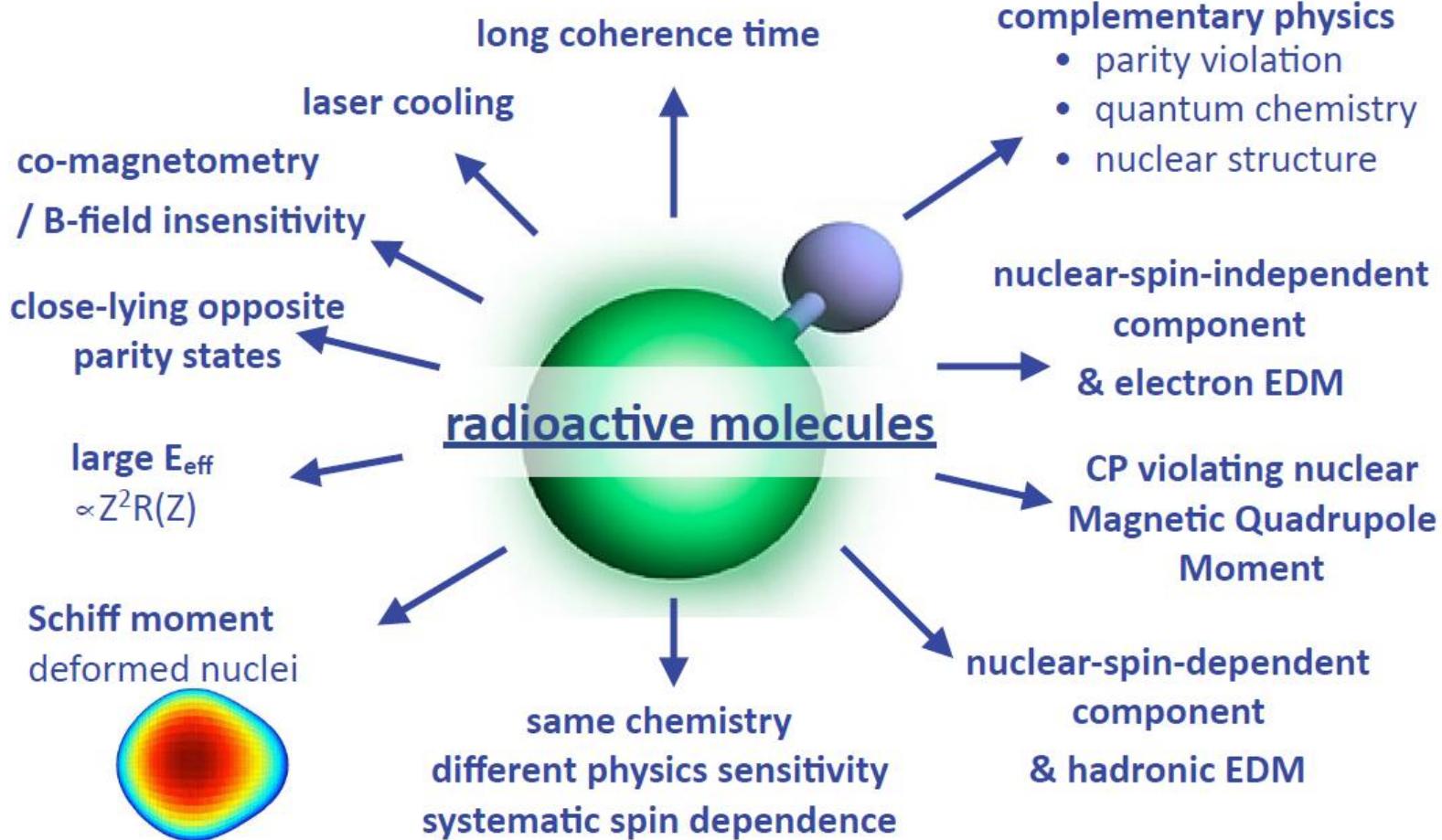


$\approx 1/10 \sigma$ closer towards unitarity

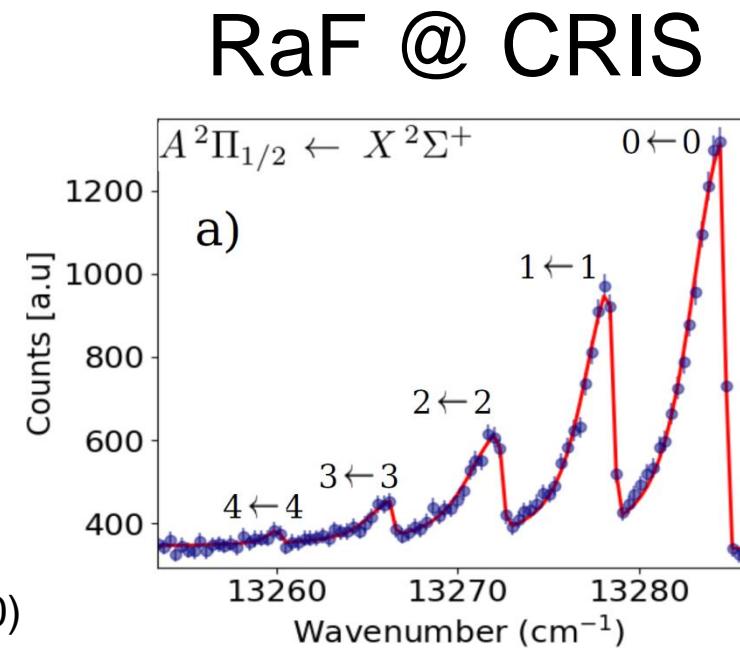
SPECTROSCOPY OF MOLECULES



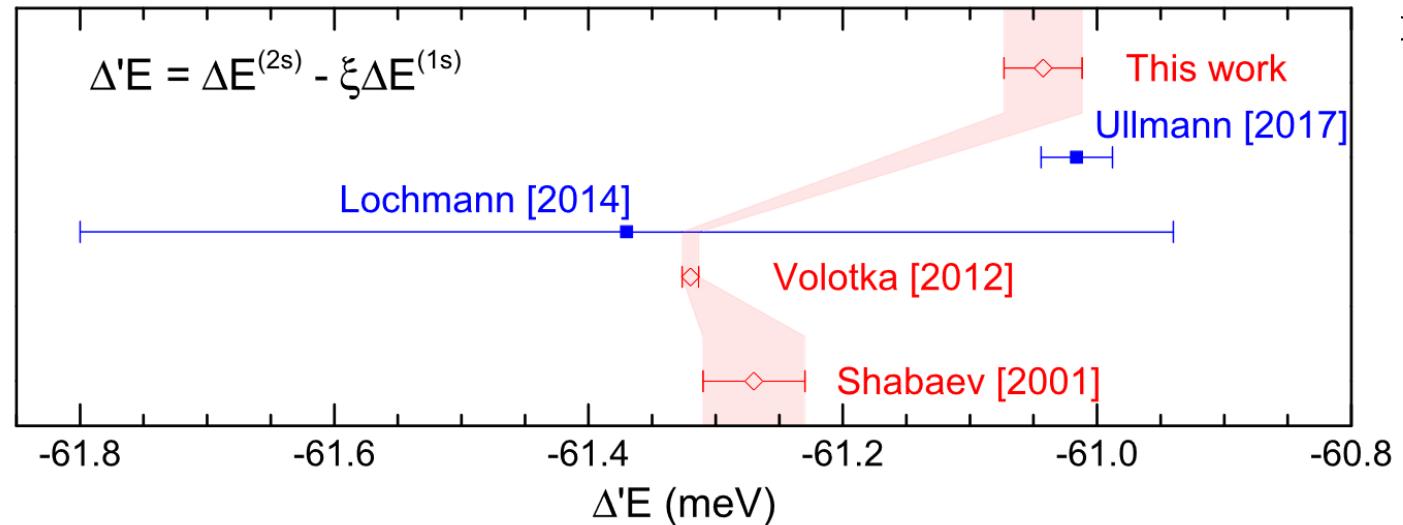
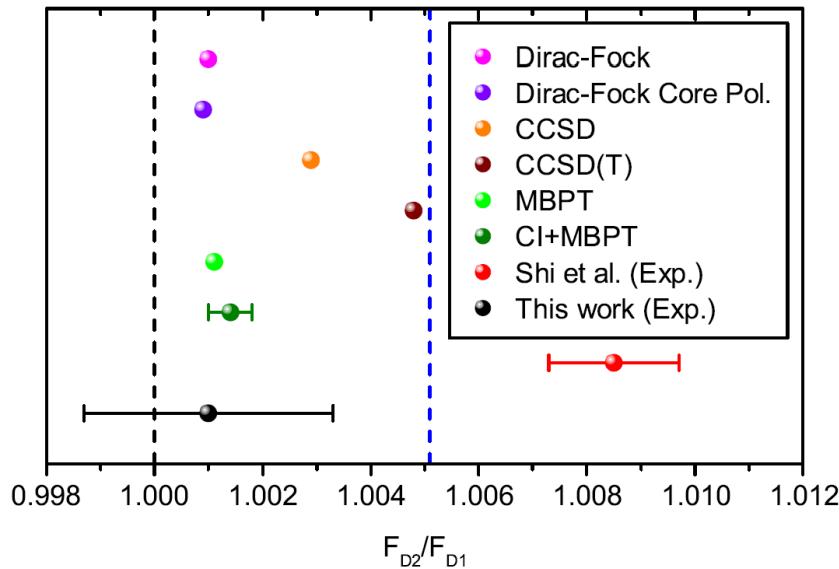
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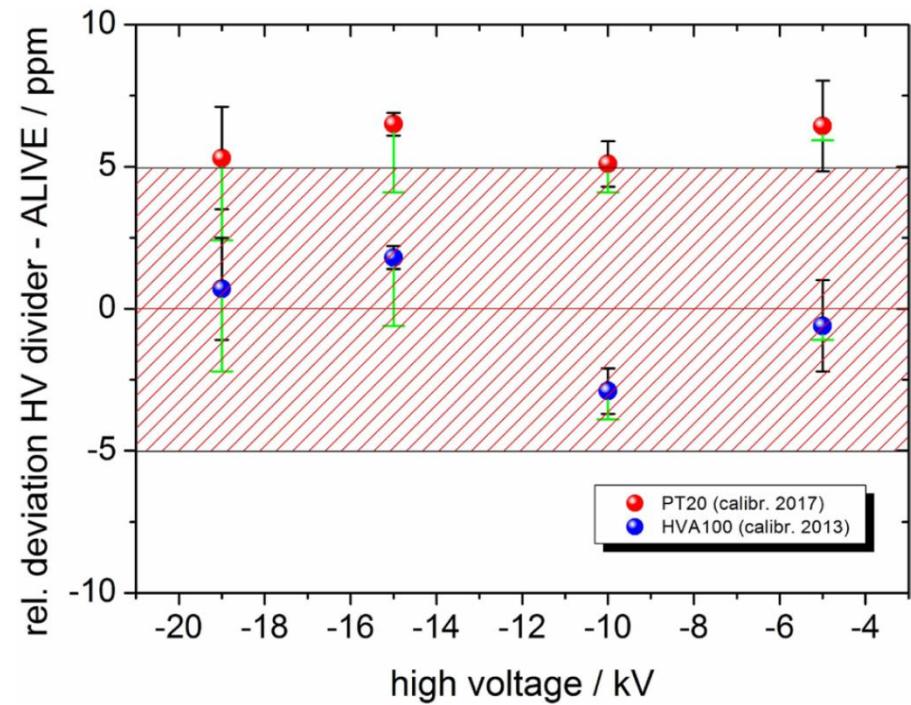
R. F. Garcia Ruiz et al., Nature 581, 396 (2020)



ATOMIC PHYSICS & HV METROLOGY



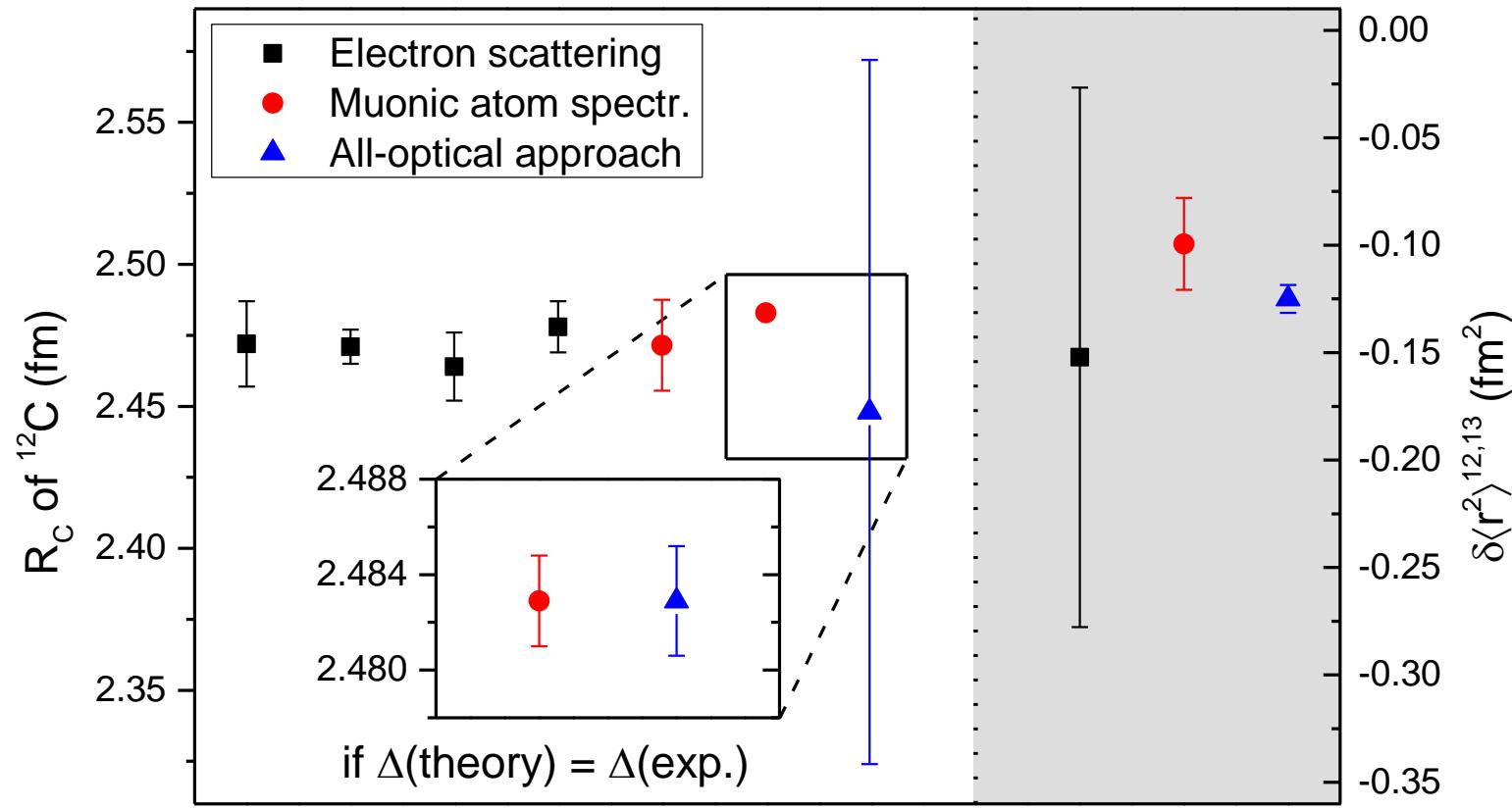
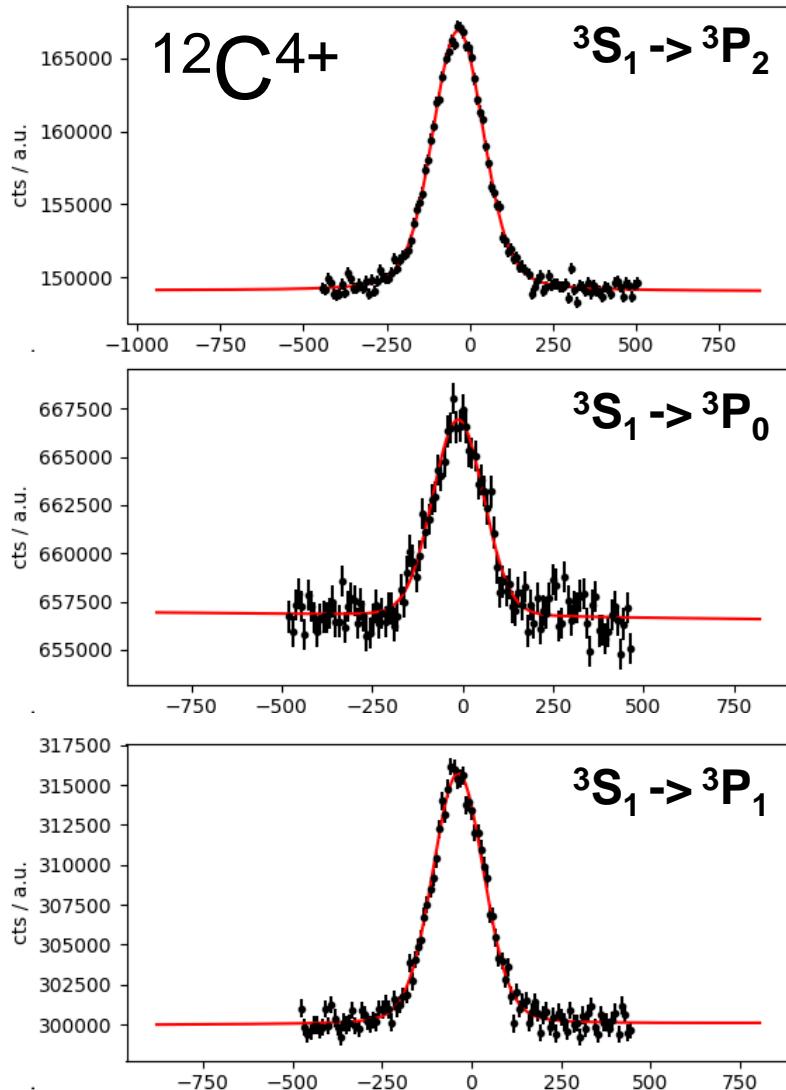
P. Müller et al., Phys. Rev. Research 2, 043351 (2020)
 L. Skripnikov et al., Phys. Rev. Lett. 120, 093001 (2018)
 J. Krämer et al 2018 Metrologia 55 268



HIGHLY CHARGED IONS FROM EBIS



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P. Imgram et al., submitted to PRL
P. Müller et al., in preparation

THANK YOU FOR YOUR ATTENTION



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