

Nuclear Incompressibility, *K*∞, and the Asymmetry Term, *K*τ, from Measurements of the Giant Monopole Resonance in Neutron-Rich Nuclei

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The Compressional Mode Giant Resonances

GMR

I = 0



ISGDR



"Breathing Mode"

 $å r_i^2$

"Squeezing Mode" $a r_i^3 Y_1$

 $2 \square W$

3 WThe energies of both these resonances are directly related to

Nuclear Incompressibility.

$$E_{GMR} = \Box \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

 $E_{ISGDR} = \Box \sqrt{\frac{7}{3} \frac{K_A + \frac{27}{25} e_F}{m \langle r^2 \rangle}}$









240 MeV (α , α ') [Texas A & M]







From GMR data on ²⁰⁸Pb and ⁹⁰Zr $K_{\infty} = 240 \pm 20 \text{ MeV}$

This number is consistent with both GMR and ISGDR data and with non-relativistic and relativistic calculations



We know K_A from E_{GMR}:

$$E_{GMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

In an approximate way, K_A may be expressed as:

 $K_A \sim K_\infty (1 + cA^{-1/3}) + K_\tau ((N - Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$

c ~ -1

Kcoul is, basically, model independent

K_{τ} ??

Measurements over a series of isotopes gives K_{τ}

 $K_{\tau} = K_{\rm sym} - 6L - Q_0 L/K_{\infty}$



$$\begin{split} \mathsf{K}_{\mathsf{A}} &\sim \mathsf{K}_{\mathsf{vol}} \left(1 + \mathsf{c}\mathsf{A}^{-1/3}\right) + \mathsf{K}_{\tau} \left((\mathsf{N} - \mathsf{Z})/\mathsf{A}\right)^{2} + \mathsf{K}_{\mathsf{Coul}} \, \mathsf{Z}^{2} \mathsf{A}^{-4/3} \\ \mathsf{K}_{\mathsf{A}} &- \mathsf{K}_{\mathsf{Coul}} \, \mathsf{Z}^{2} \mathsf{A}^{-4/3} \sim \mathsf{K}_{\mathsf{vol}} \left(1 + \mathsf{c}\mathsf{A}^{-1/3}\right) + \mathsf{K}_{\tau} \left((\mathsf{N} - \mathsf{Z})/\mathsf{A}\right)^{2} \\ &\sim \mathsf{Constant} + \mathsf{K}_{\tau} \left((\mathsf{N} - \mathsf{Z})/\mathsf{A}\right)^{2} \end{split}$$

We use K_{Coul} = - 5.2 MeV (from Sagawa)







 K_τ = -550 \pm 100 MeV

$$K_{\rm t} = K_{\rm sym} - 6L - Q_0 L/K_{\infty}$$





J. Button et al., Phys. Rev. C 96, 054330 (2017)





K. B. Howard et al., Phys. Lett. B 801, 135185 (2020)



Can we do any better on K_{∞} and K_{τ} ?

-- Not much, I think, on $K_{\!\scriptscriptstyle\infty}$

-- Can hope for improvement in K_{τ}



± 50 MeV, perhaps









RIBF 113

¹³²Sn + ²H 100 MeV/A >50 kHz ¹³²Sn



Preliminary uncorrected energy spectra for 132 Sn, 133 Sb, and 134 Te.

Particle identifications of beam and recoil particles have been done.



Preliminary results from RIKEN d(¹³²Sn, ¹³²Sn')d experiment



132Sn





133Sb



Preliminary results from RIKEN d(132Sn, 132Sn')d experiment





Preliminary results from RIKEN d(132Sn, 132Sn')d experiment









50 MeV/A (α , α '); He with 5% CF₄



⁴He(⁷⁰Ni, ⁷⁰Ni')⁴He @ NSCL



82 MeV/A; ~20k pps





⁴He(⁷⁰Ni, ⁷⁰Ni')⁴He @ NSCL





J. Bradt et al., Nucl. Inst. Meth. Phys. Res. A 875, 65 (2017) Giraud et al., Nucl. Inst. Meth. Phys. Res. A 1051, 168213 (2023)











Preliminary results from J.S. Randhawa et al.





- L=0 : Minima at ~ 2.7 degrees
- L=2 : nearly flat in 1-4 degrees





In ⁷⁰Ni GMR expected at ~17 MeV

Preliminary results from J.S. Randhawa et al.







Ex= 12-16 MeV



Preliminary results from J.S. Randhawa et al.



⁴He(¹³²Sn, ¹³²Sn')⁴He @ FRIB

Aims To investigate:

- i) the asymmetry term in the nuclear incompressibility;
- ii) the phenomenon of "softness" of open-shell nuclei as observed in the ISGMR strength distributions in the Sn and Cd nuclei; and,
- iii) to directly test the impact of the recent results from PREX-II on the observed ISGMR strengths in very neutron-rich nuclei.



⁴He(¹³²Sn, ¹³²Sn')⁴He @ FRIB



Calculations courtesy Prof. J. Piekarewicz.





Rigidities of Sn-isotopes and S800 acceptance are shown. The bars for the Sn isotope points represent their expected momentum spread





Why are tins so "Fluffy"?

F. Li et al., Phys. Rev. Lett. 99, 162503 (2007)





Fig. 3. (Color online.) Systematics of the moment ratio, m_1/m_0 for the ISGMR strength distributions in the Cd isotopes investigated in this work. The experimental results (squares) are compared with relativistic calculations performed using the FSUGold (circles) and NL3 (triangles) effective interactions. Also presented are results from non-relativistic calculations performed using the Sly5 parameter set in the HF-BCS + QRPA formalism with and without the mixed pairing interaction (diamonds and stars, respectively) [36]. The solid lines are to guide the eye.

D. Patel *et al.*, Phys. Lett. B **718**, 447 (2012)





P. Vesely et al., Phys. Rev. C 86, 024303 (2012)





K. B. Howard et al., Phys. Lett. B 807, 135608 (2020)





the connection between the line shape of the monopole strength ISGMR and the deformation-induced coupling between the ISGMR and the K = 0 branch of the ISGQR. The ISGMR is best described by the force SkP^{δ}, having a low incompressibility K_{∞} = 202 MeV.

Gianluca Colò et al., Phys. Lett. B 811, 135940 (2020)





Z. Z. Li, Y. F. Niu, and G. Colo, arXiv:2211.01264v1





FIG. 1. ISGMR in ¹²⁰Sn and ²⁰⁸Pb: RQRPA and RQTBA strength distributions compared to experimental data [35] (¹²⁰Sn) and [1] (²⁰⁸Pb).



Elena Litvinova, Phys. Rev. C 107, L041302 (2023).



We have "experimental" values now for K_{∞} and K_{τ} . There is a lot of room for "improvement" in the latter.

 We have completed a measurement at NSCL on ⁷⁰Ni using pure ⁴He gas and trigger with S800. Very clean spectra have been observed with a "peak" at ~17 MeV consistent with GMR.

 An experiment to measure ISGMR in ¹³²Sn was approved in the "first round" proposals for FRIB.

In the Cd and Sn isotopes, the ISGMR energy was significantly lower than that expected from the accepted value of K_∞.

The "fluffiness" appears in the Mo isotopes, beginning with ⁹²Mo, just two nucleons out of the "doubly-closed" nucleus ⁹⁰Zr.

Until very recently, there had been no satisfactory theoretical explanation of this "fluffiiness" of open-shell nuclei.

MAYBE, Now there is!



Mercí धन्य वा ट Thanks







The Question Kitten