



Chiral Symmetry in Nuclear Medium Observed in Pionic Atoms

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nature physics

T. Nishi, K.I. et al., Nat. Phys. (2023)

Article

<https://doi.org/10.1038/s41567-023-02001-x>

**Chiral symmetry restoration at high matter
density observed in pionic atoms**

- T.Nishi, KI et al., N. Phys. **19**, 788 (2023)
Article DOI: 10.1038/s41567-023-02001-x
- Nature Physics (2023/3/23)
News and Views "Modified in Medium"

Chiral Symmetry in Nuclear Medium Observed in Pionic Atoms

- **Dominant symmetry of the vacuum in low-energy region.**
- **Spontaneous breakdown due to non-perturbative strong interaction.**
- **Non-trivial structure of the QCD vacuum.**

nature physics

T. Nishi, K.I. et al., Nat. Phys. (2023)

Article

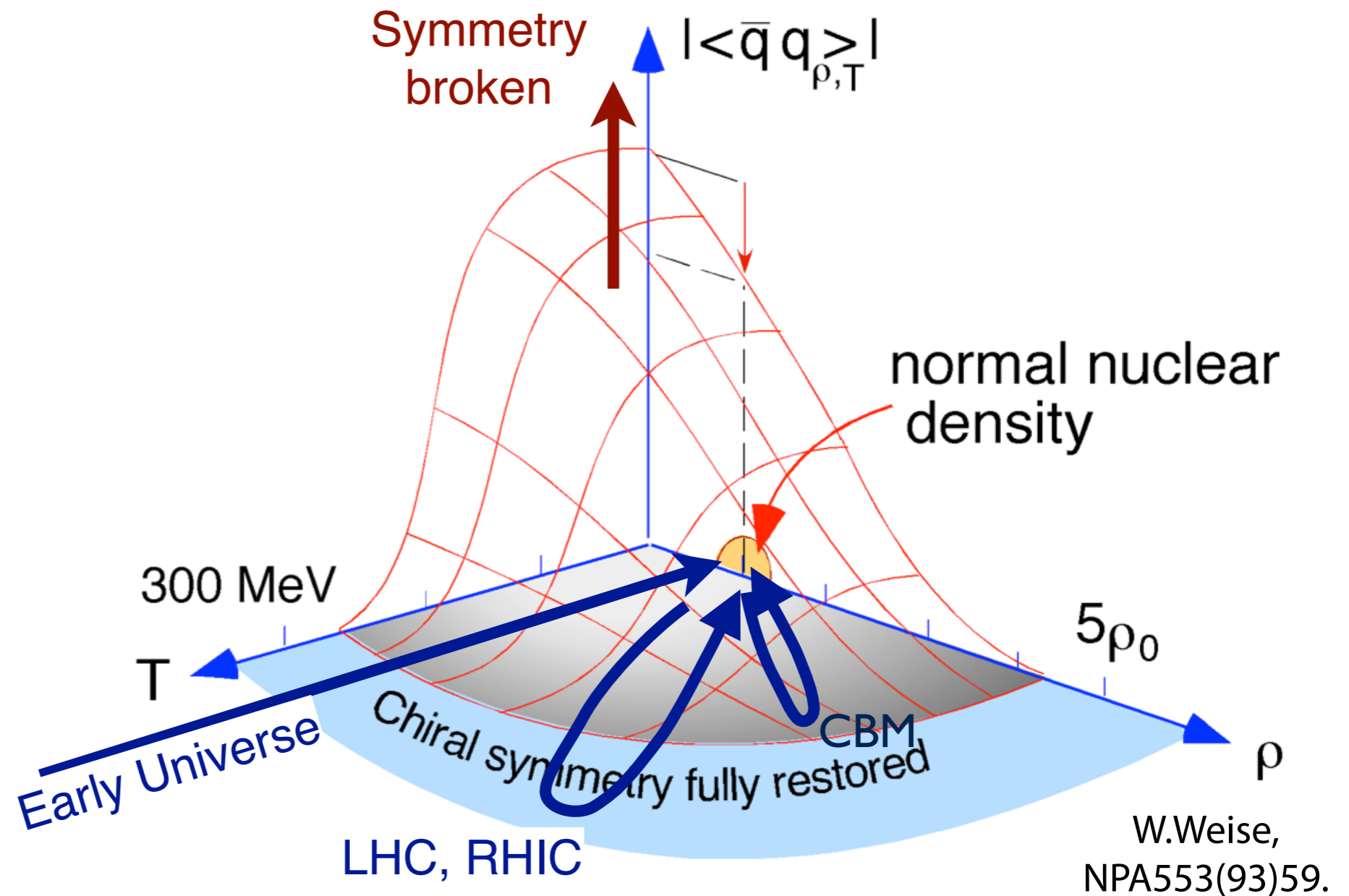
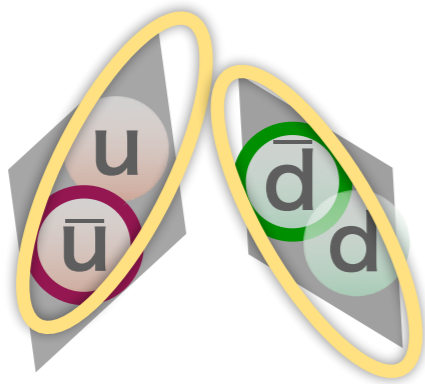
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Chiral symmetry restoration at high matter density observed in pionic atoms

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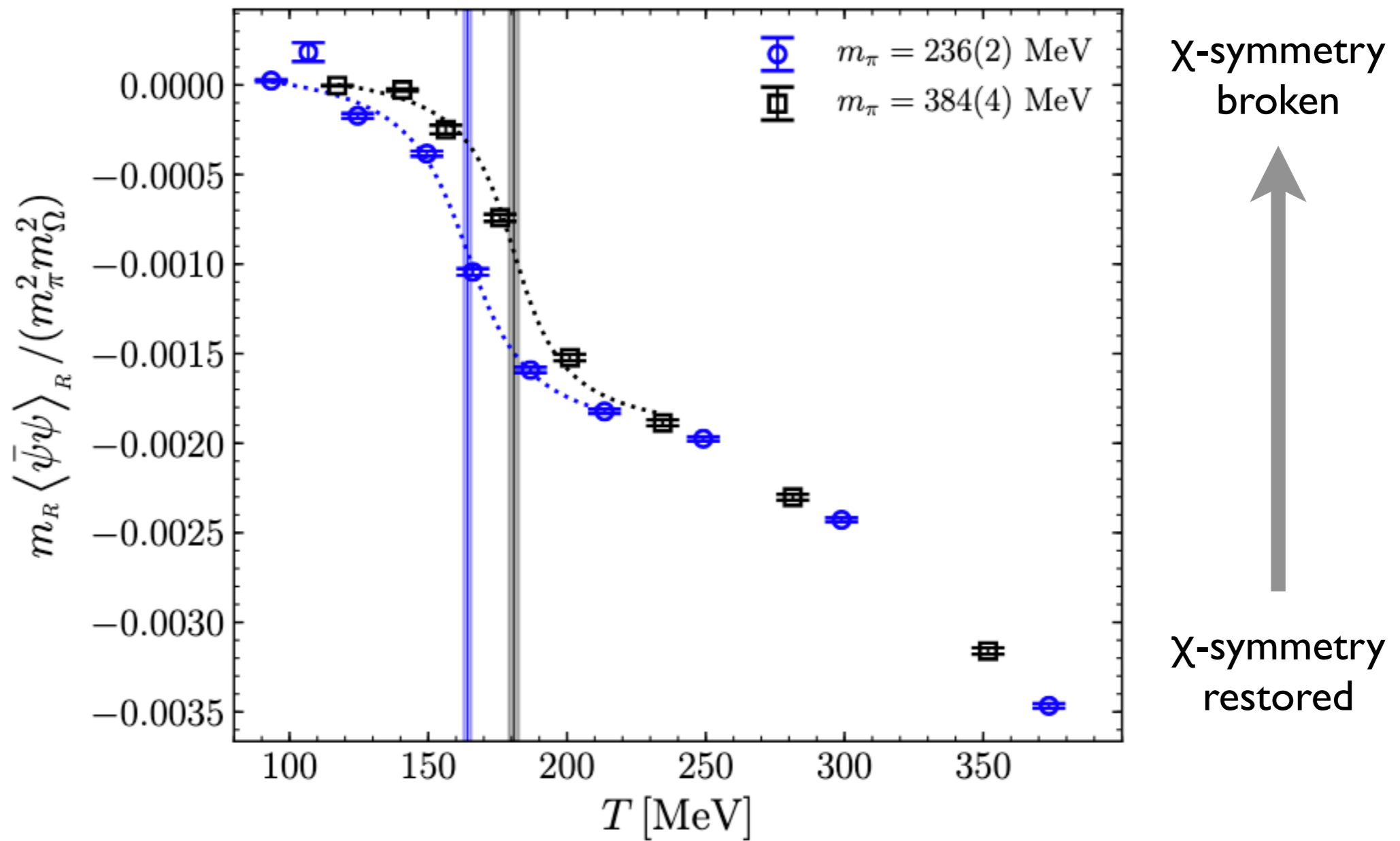
Chiral condensate $\bar{q}q$ on $T\rho$ plane

An order parameters of χ -symmetry



Material properties
of QCD vacuum

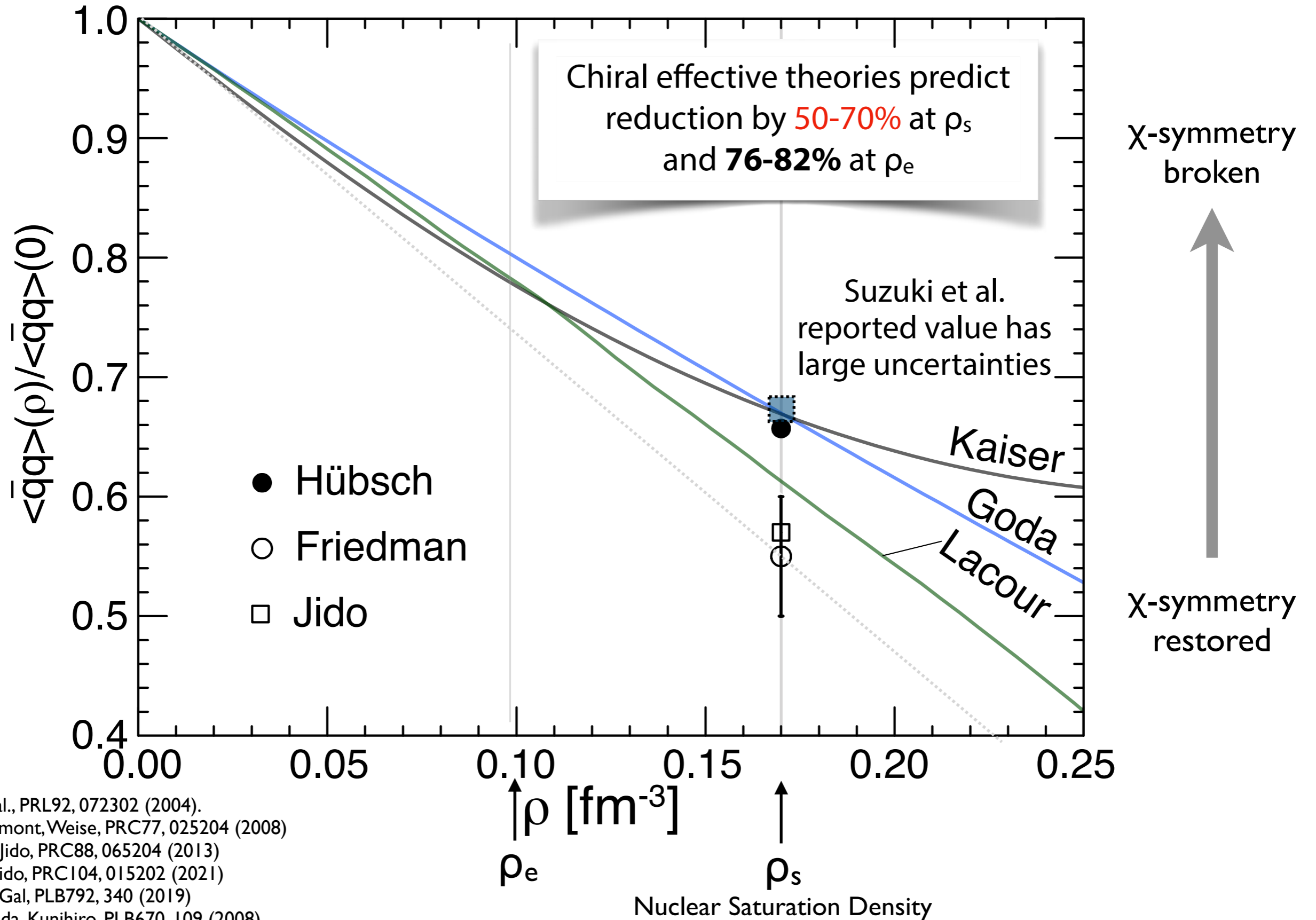
Lattice QCD calculated T dependence of $\langle \bar{q}q \rangle$



Remark: sign problem makes it difficult for lattice to approach non-zero ρ region

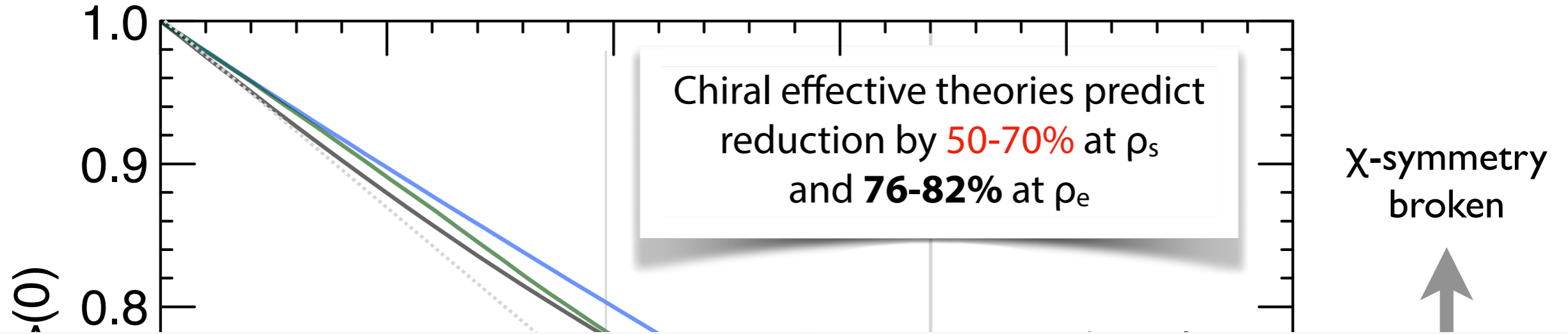
Jon-Ivar Skullerud
PRD105(2022)034504

ρ dependence of $\langle \bar{q}q \rangle$ known so far

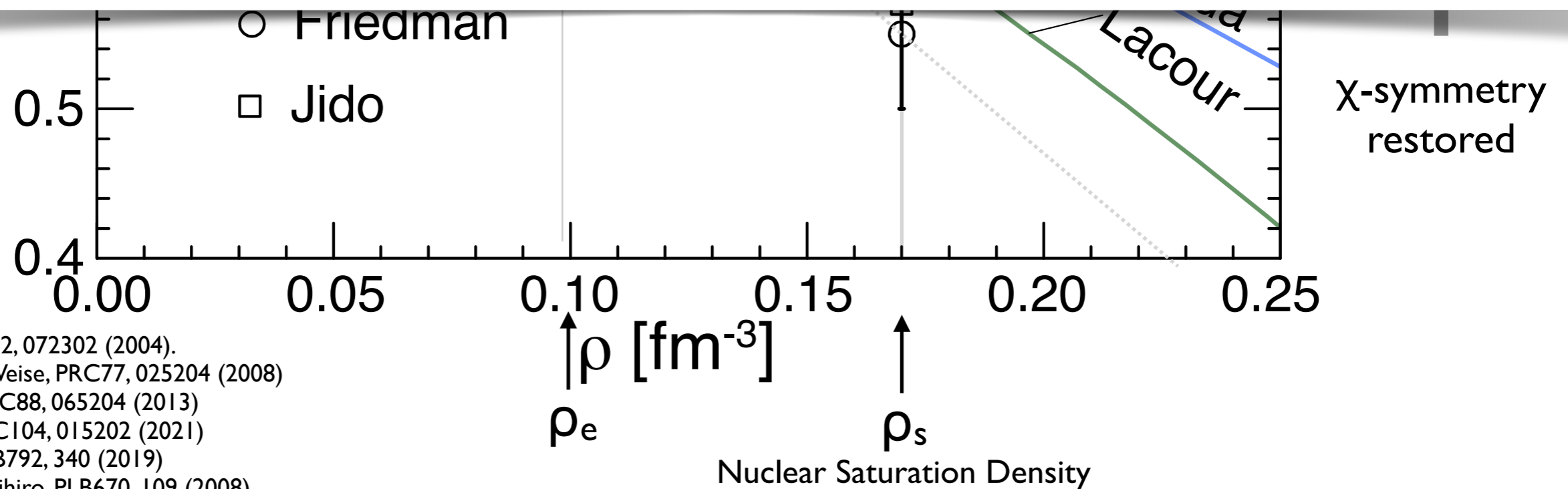


Suzuki et al., PRL92, 072302 (2004).
 Kaiser, Homont, Weise, PRC77, 025204 (2008)
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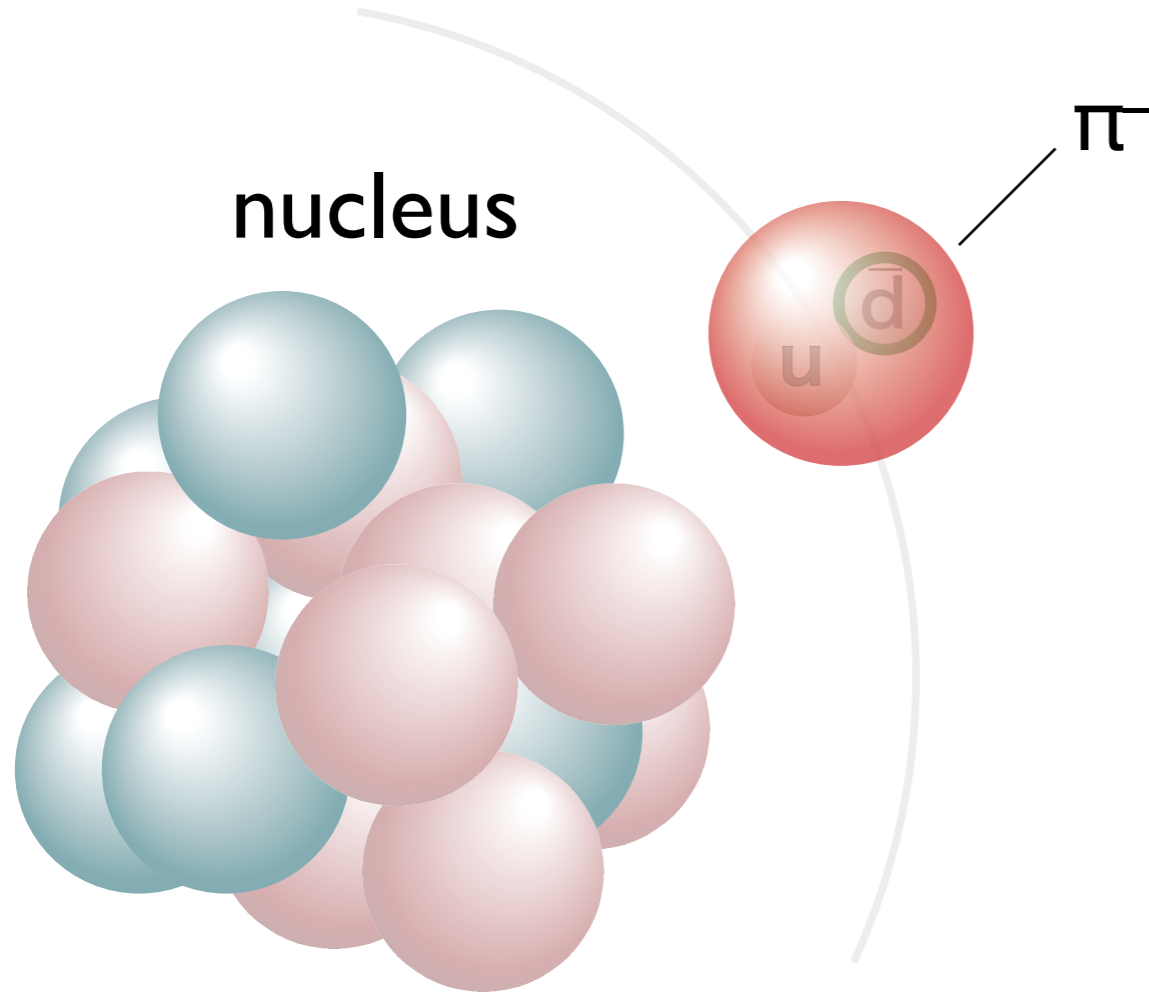


Need high-quality experimental information to quantify $\langle \bar{q}q \rangle$ reduction and confirm theoretical scenario of vacuum evolution



Suzuki et al., PRL92, 072302 (2004).
 Kaiser, Homont, Weise, PRC77, 025204 (2008)
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Precision Spectroscopy of Pionic Atoms

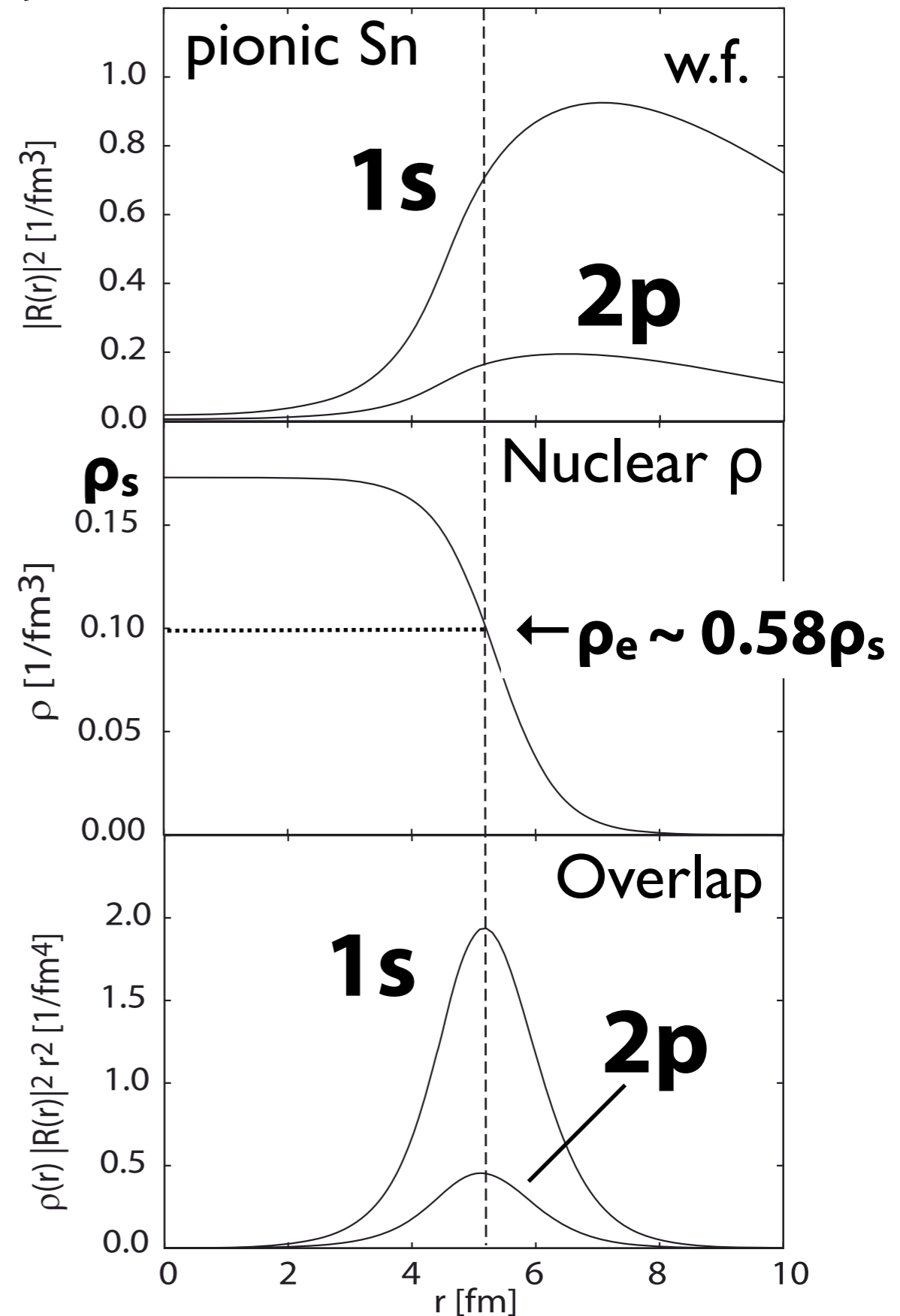


Ericson-Ericson potential

$$U_{\text{opt}}(r) = U_s(r) + U_p(r),$$

$$U_s(r) = b_0 \rho + \mathbf{b}_1 (\rho_n - \rho_p) + B_0 \rho^2$$

$$U_p(r) = \frac{2\pi}{\mu} \vec{\nabla} \cdot [c(r) + \varepsilon_2^{-1} C_0 \rho^2(r)] L(r) \vec{\nabla}$$



Pion-nucleus strong interaction

Overlap between
pion w.f. and nucleus
→ π works as a probe
at $\rho_e \sim 0.58\rho_s$



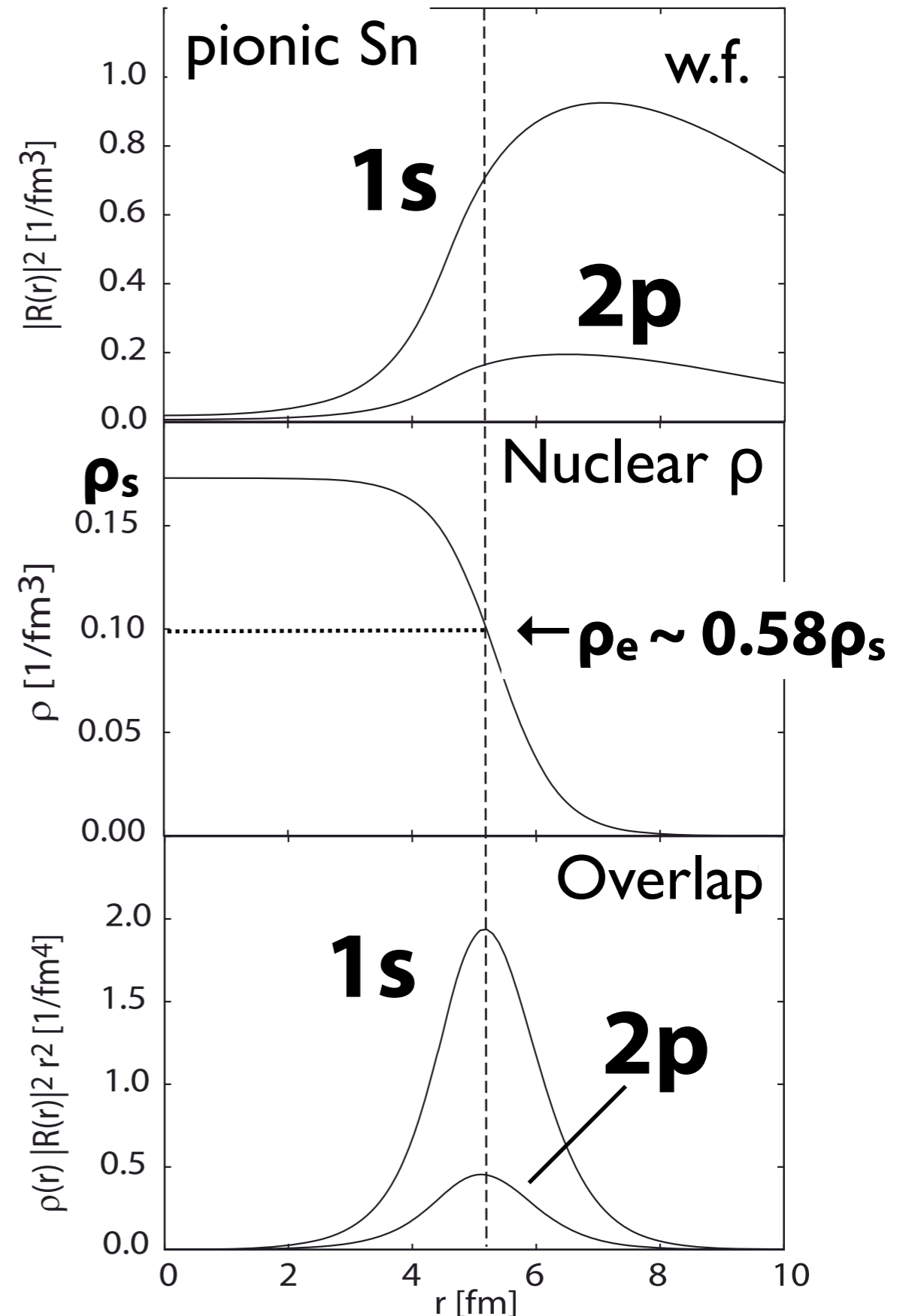
π -nucleus **interaction is changed** in
nuclear medium for wavefunction
renormalization effect

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Strong interaction and chiral condensate

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In-medium Glashow-Weinberg relation

$$\frac{\langle \bar{q}q \rangle^*}{\langle \bar{q}q \rangle^v} \simeq \left(\frac{b_1^v}{b_1} \right)^{1/2} \left(1 - \gamma \frac{\rho}{\rho_0} \right)$$

$$\gamma = 0.184 \pm 0.003$$

Jido, Hatsuda, Kunihiro, PLB670, 109 (2008)

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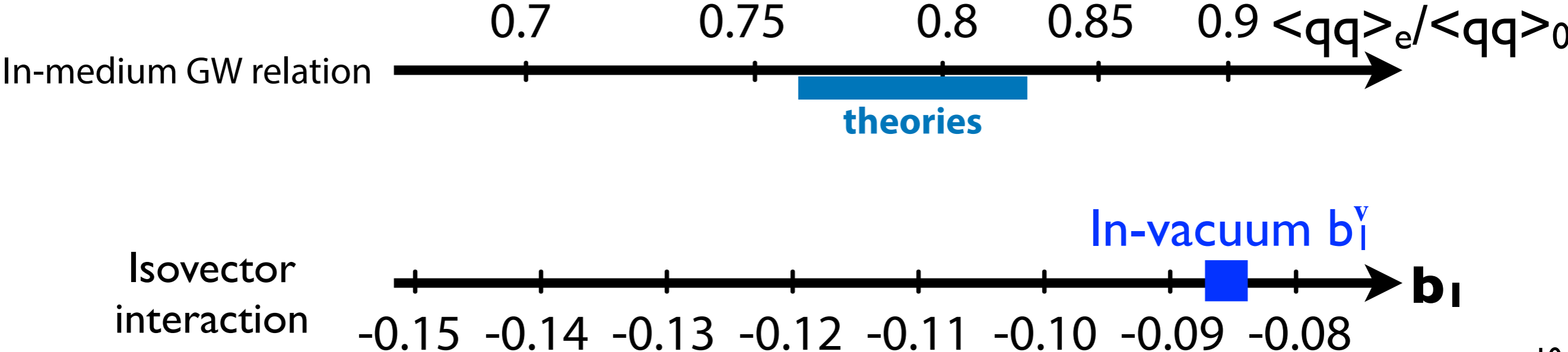
$$\gamma = 0.184 \pm 0.003$$

Jido, Hatsuda, Kunihiro, PLB670, 109 (2008)

Pionic hydrogen and deuterium

$$b_1^v = 0.0866 \pm 0.0010$$

Hirtl et al., EPJA57, 70 (2021)

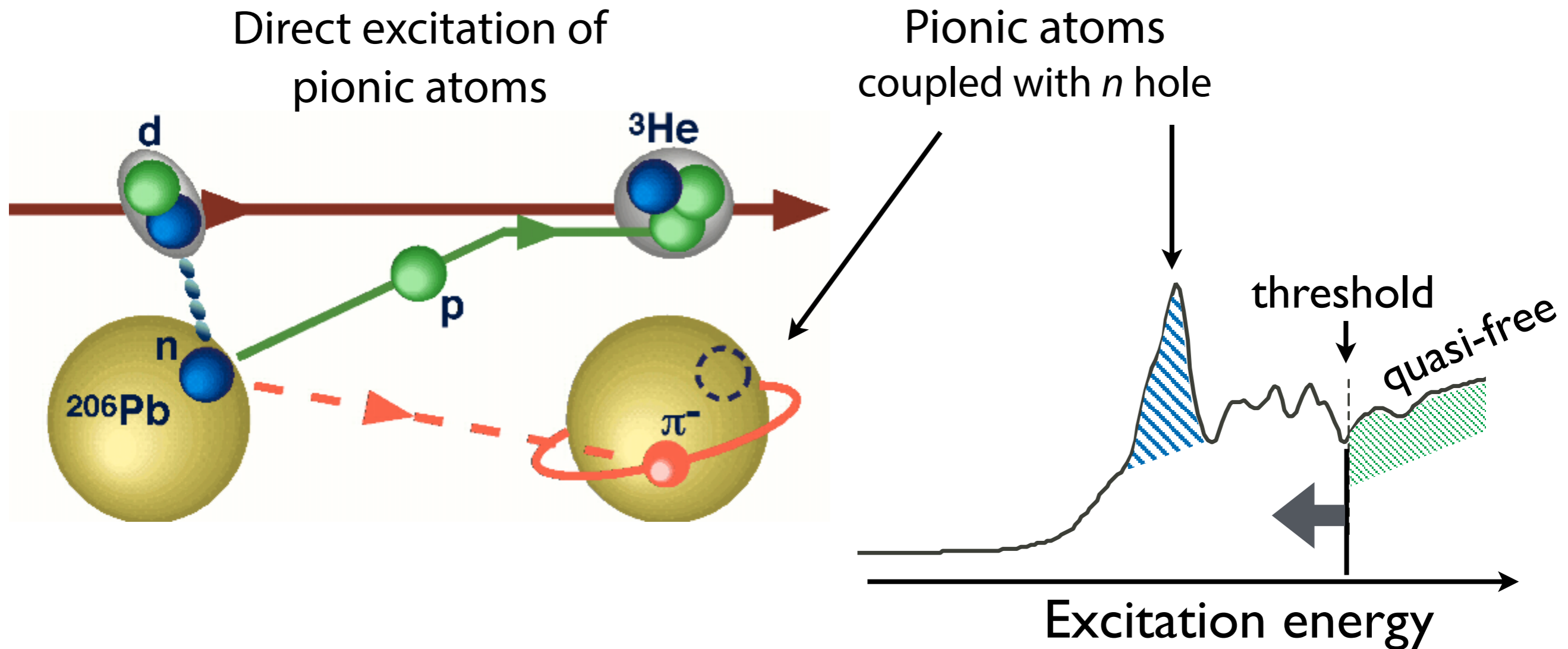


Deduction of pion-nucleus interaction in medium by

Spectroscopy of pionic atoms in $(d, {}^3\text{He})$ reactions

Based on energy-momentum conservation law:

$$\text{Excitation energy} \sim T_d - T_{{}^3\text{He}}$$

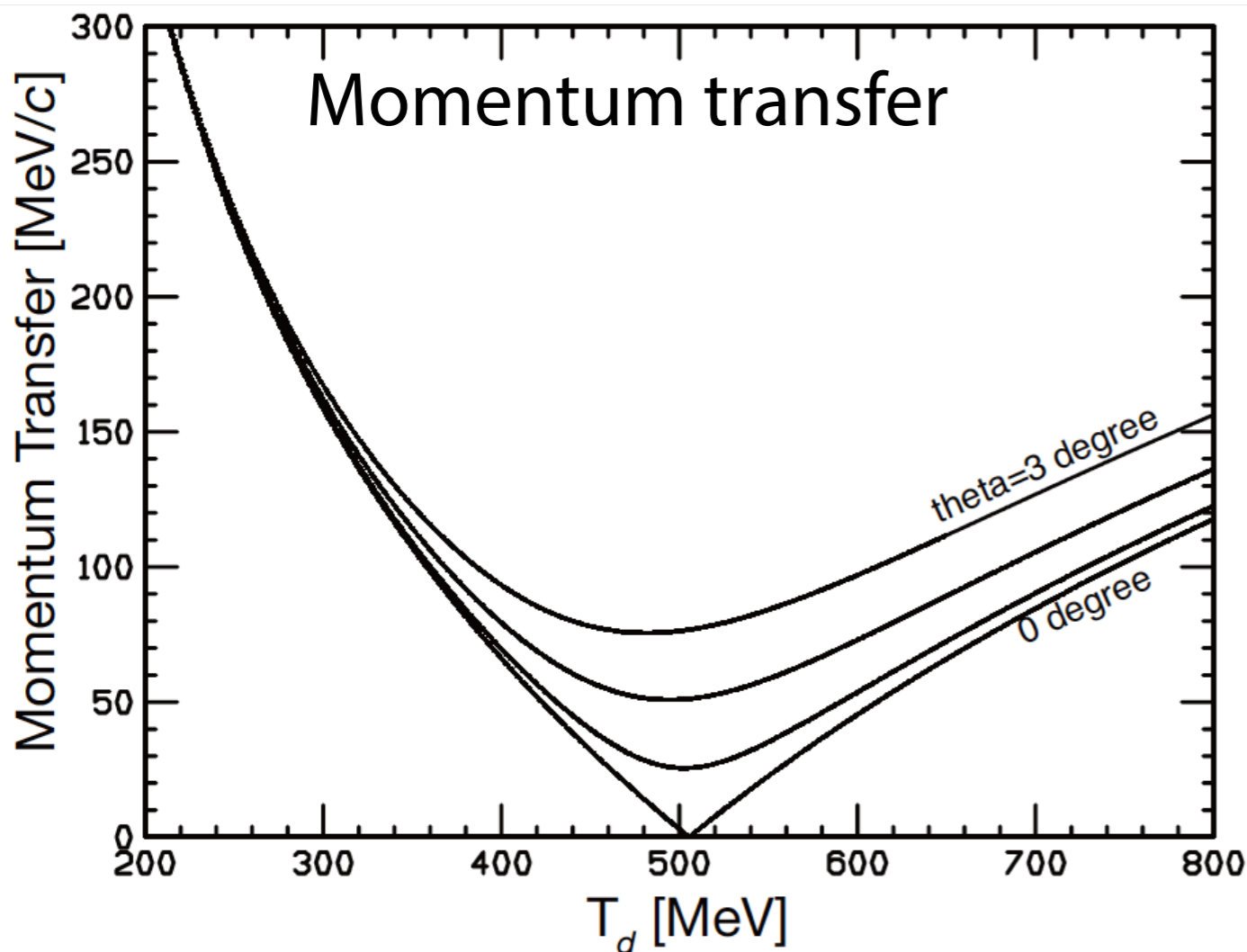


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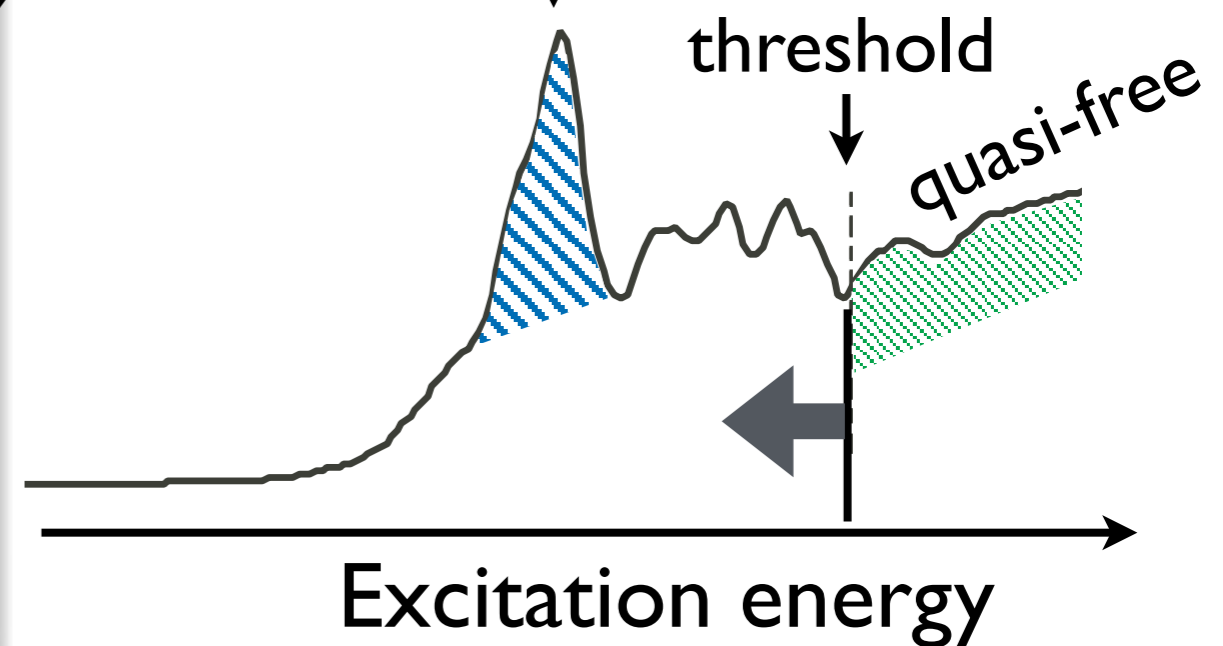
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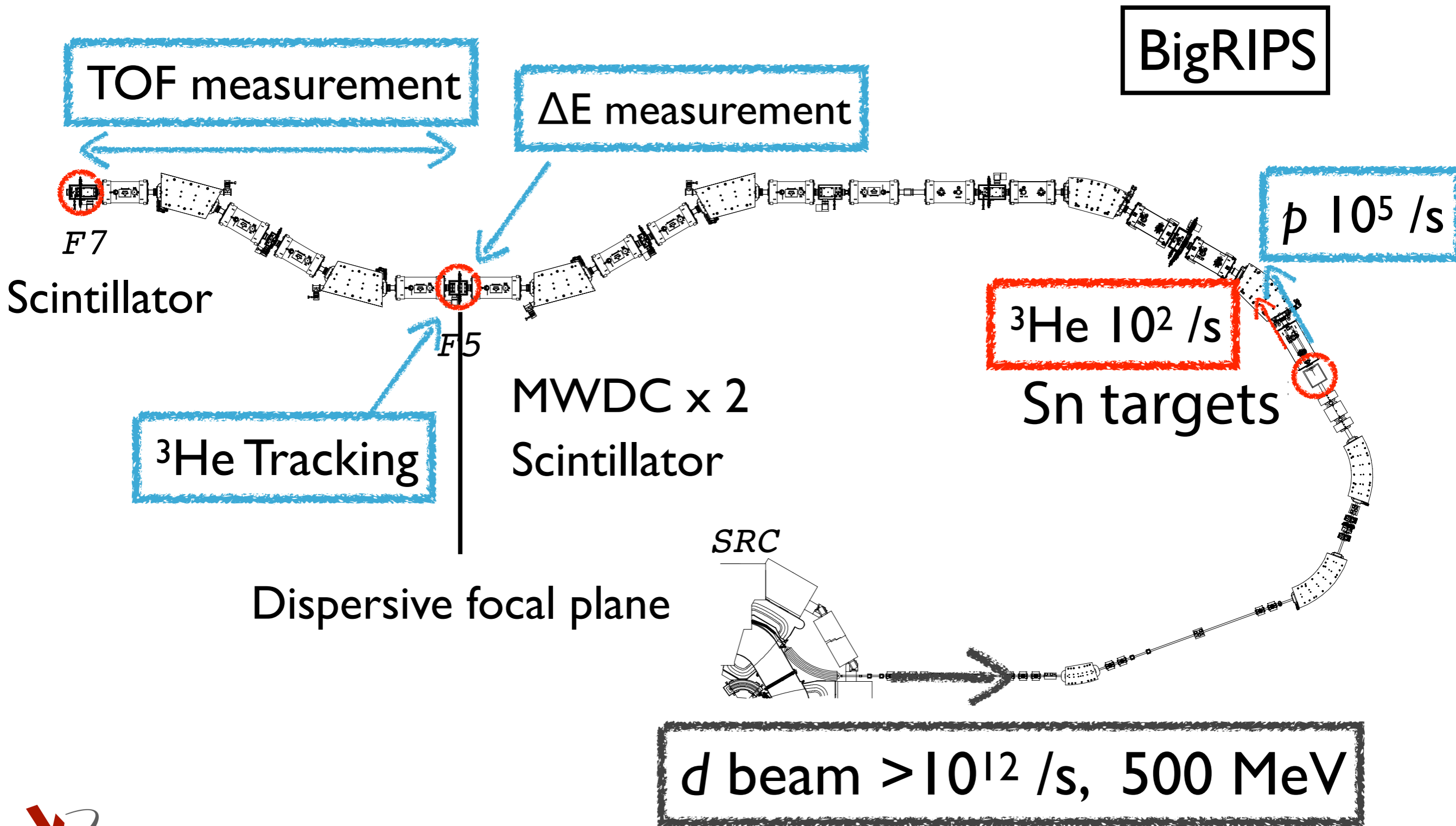
$$\text{Excitation energy} \sim T_d - T_{3\text{He}}$$



Pionic atoms
coupled with n hole



(d,³He) Reaction Spectroscopy in RIBF



First pionic atom in RIBF (2010)

Pionic ^{121}Sn atom

First simultaneous 1s and 2p observation

$$B_{1s} = 3.828 \pm 0.013(\text{stat})_{-0.033}^{+0.036}(\text{syst}) \text{ MeV}$$

$$\Gamma_{1s} = 0.252 \pm 0.054(\text{stat})_{-0.070}^{+0.053}(\text{syst}) \text{ MeV}$$

$$B_{2p} = 2.238 \pm 0.015(\text{stat})_{-0.043}^{+0.046}(\text{syst}) \text{ MeV}$$

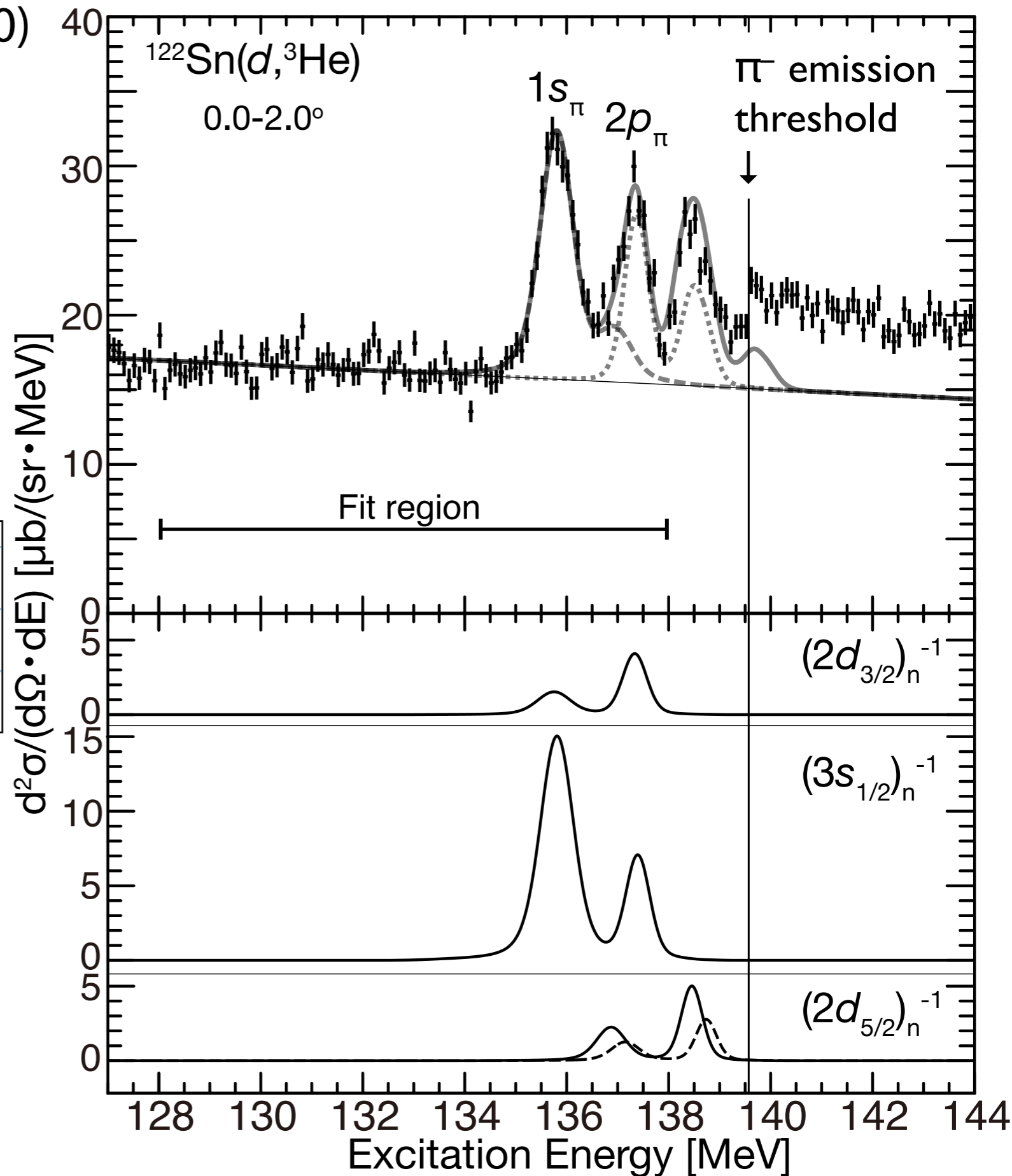
Resolution 394 keV (FWHM)

Theories

$$B_{1s} = 3.787\text{--}3.850 \text{ MeV}$$

$$\Gamma_{1s} = 0.306\text{--}0.324 \text{ MeV}$$

$$B_{2p} = 2.257\text{--}2.276 \text{ MeV}$$



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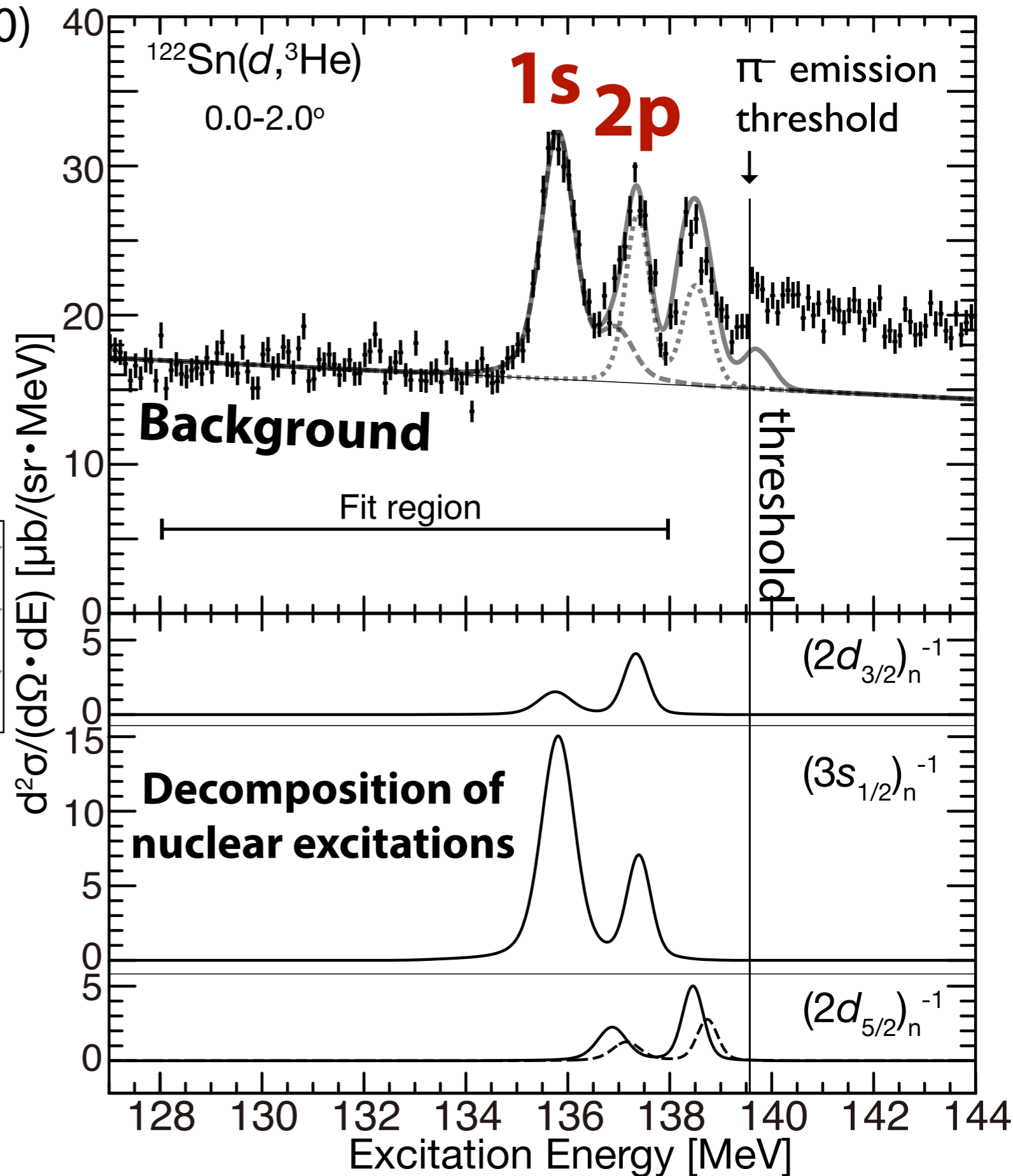
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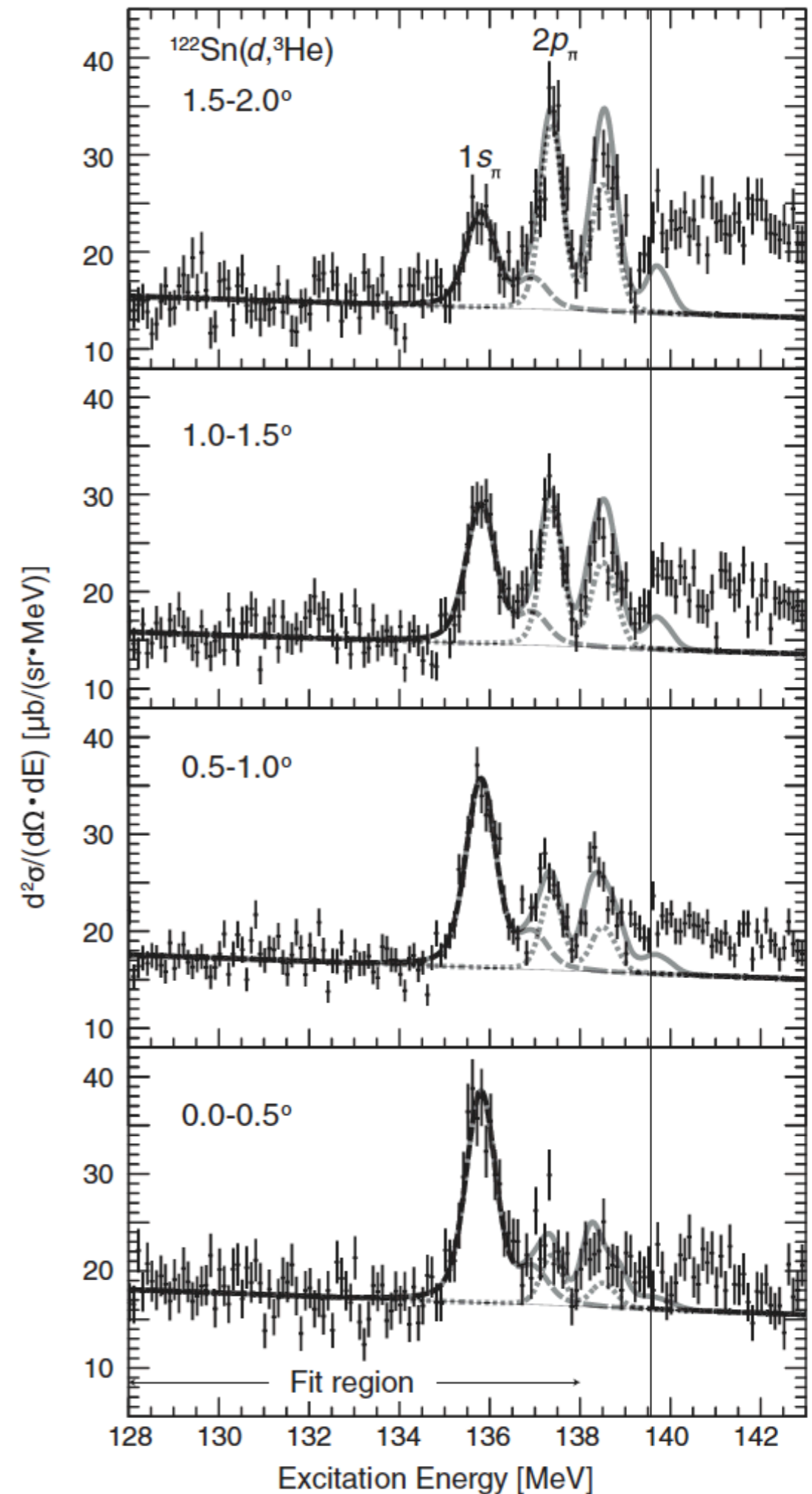
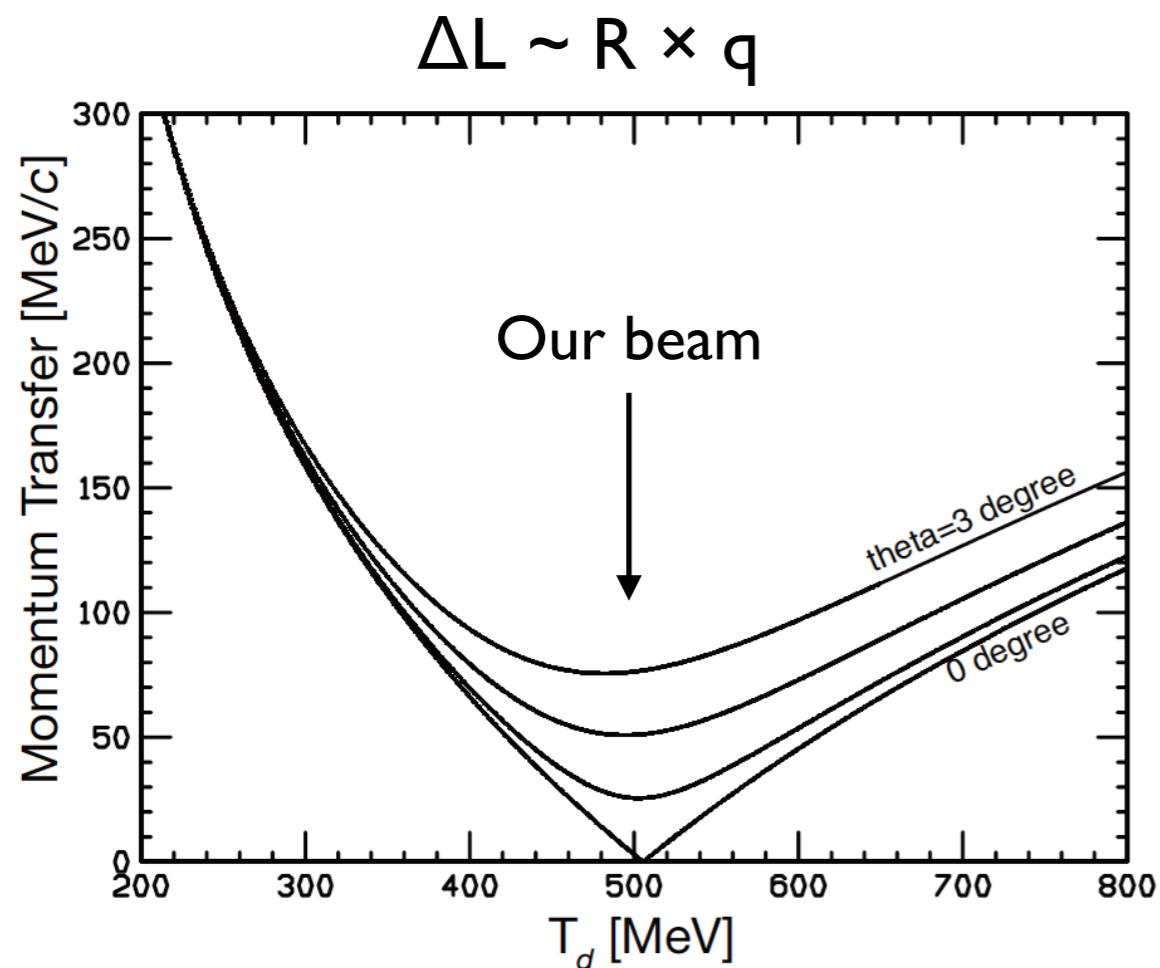
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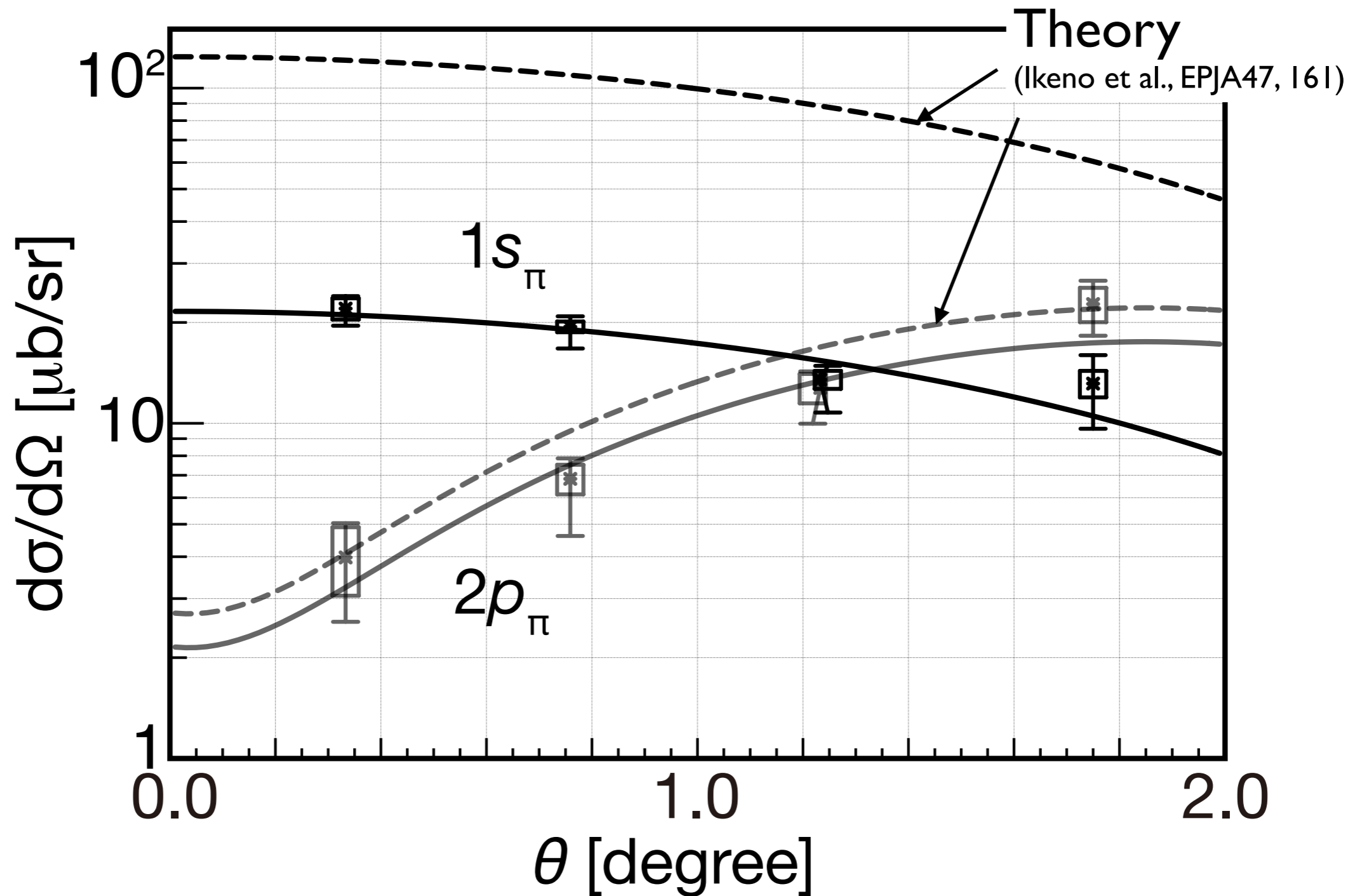
First pionic atom in RIBF (2010)

Pionic ^{121}Sn atom

First observation of
 θ dependence of
 π atom cross section



1s and 2p pionic atom cross sections in (d,³He)

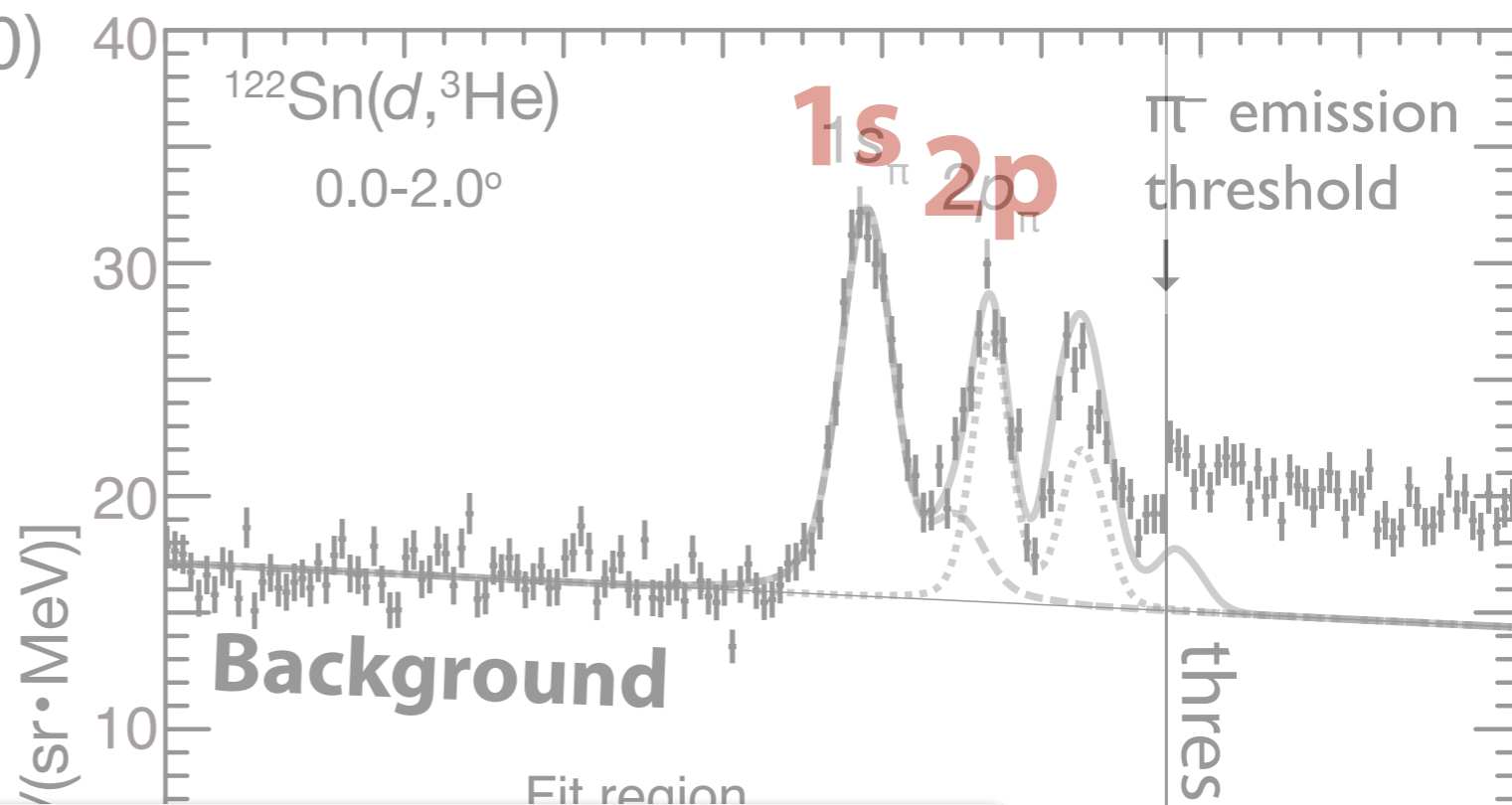


θ dependence is well reproduced.
Theory calculates 5x larger cross section for 1s

First pionic atom in RIBF (2010)

Pionic ^{121}Sn atom

First simultaneous 1s and 2p observation



The precision was not enough...

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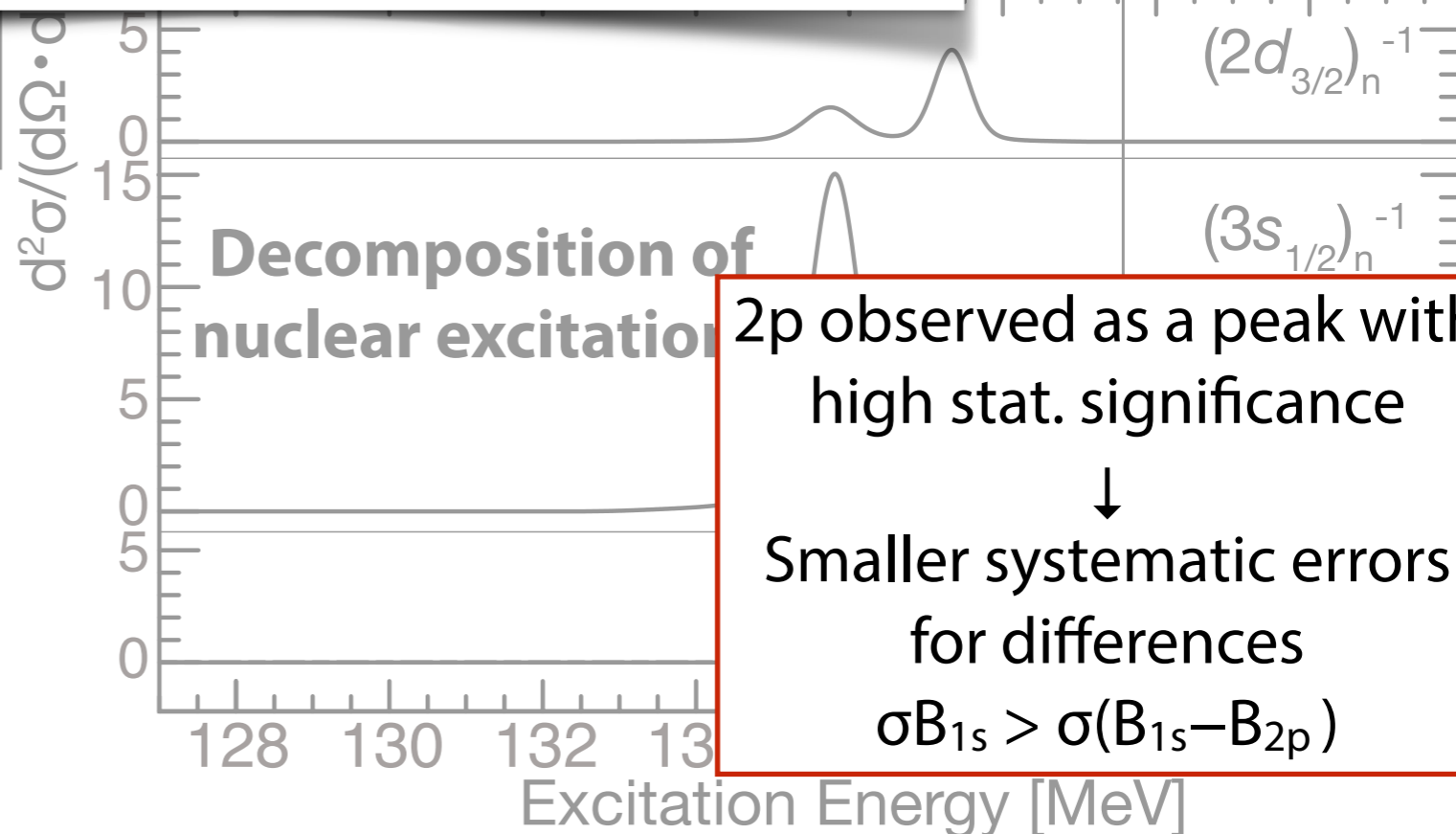
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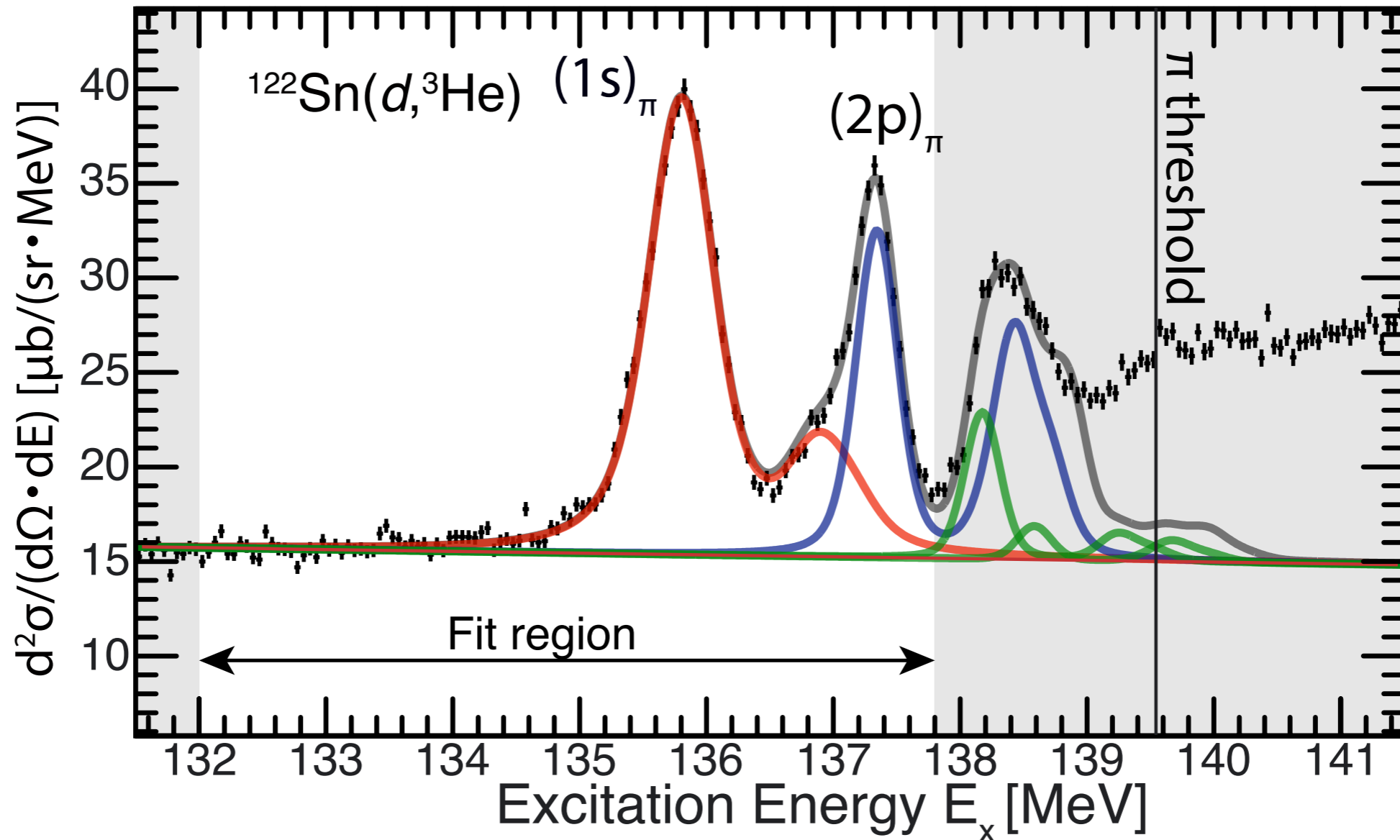


2p observed as a peak with high stat. significance

↓
Smaller systematic errors for differences

$$\sigma B_{1s} > \sigma(B_{1s} - B_{2p})$$

High Precision Spectrum of $^{122}\text{Sn}(d,^3\text{He})$ in RIBF-54

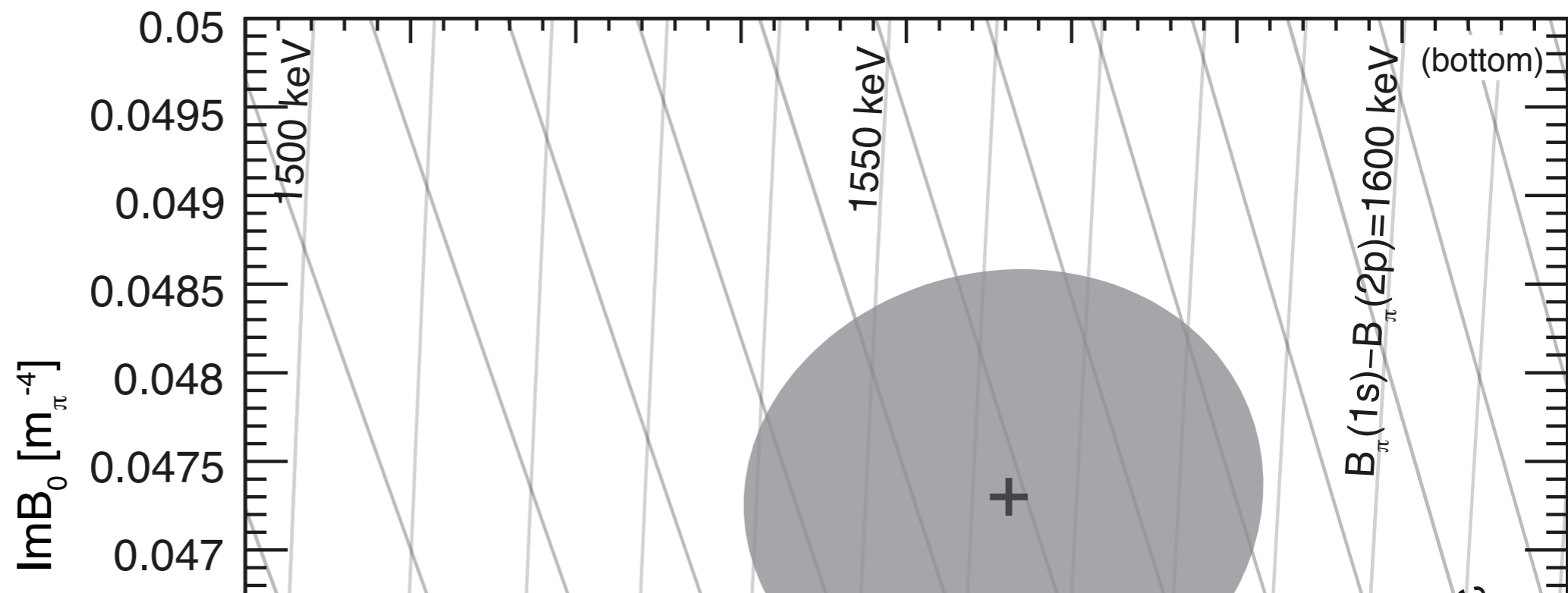
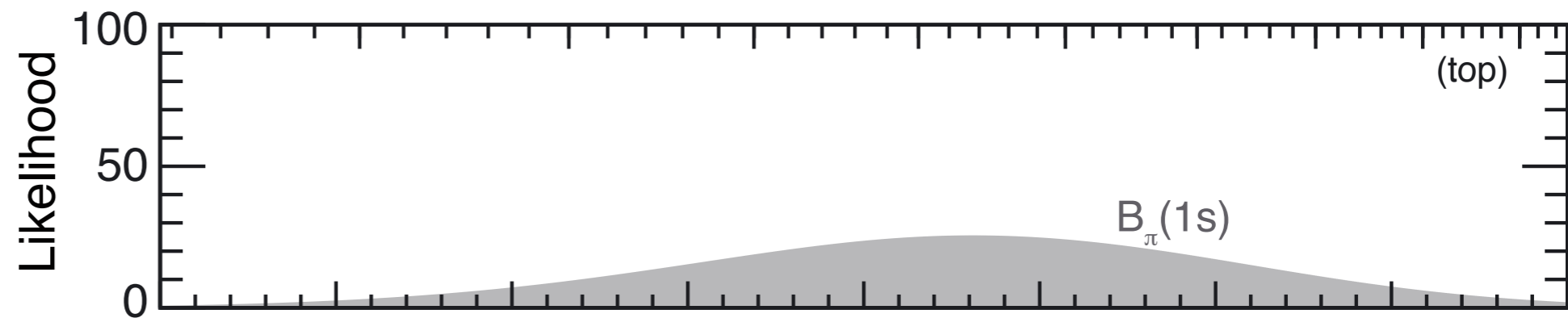


	[keV]	Statistical	Systematic
$B_\pi(1s)$	3831	± 3	+78 – 76
$B_\pi(2p)$	2276	± 3	+84 – 83
$B_\pi(1s) - B_\pi(2p)$	1555	± 4	± 12
$\Gamma_\pi(1s)$	316	± 12	+36 – 39
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b_1 parameter Deduction

based on
pionic C, N, O,
and Sn atoms

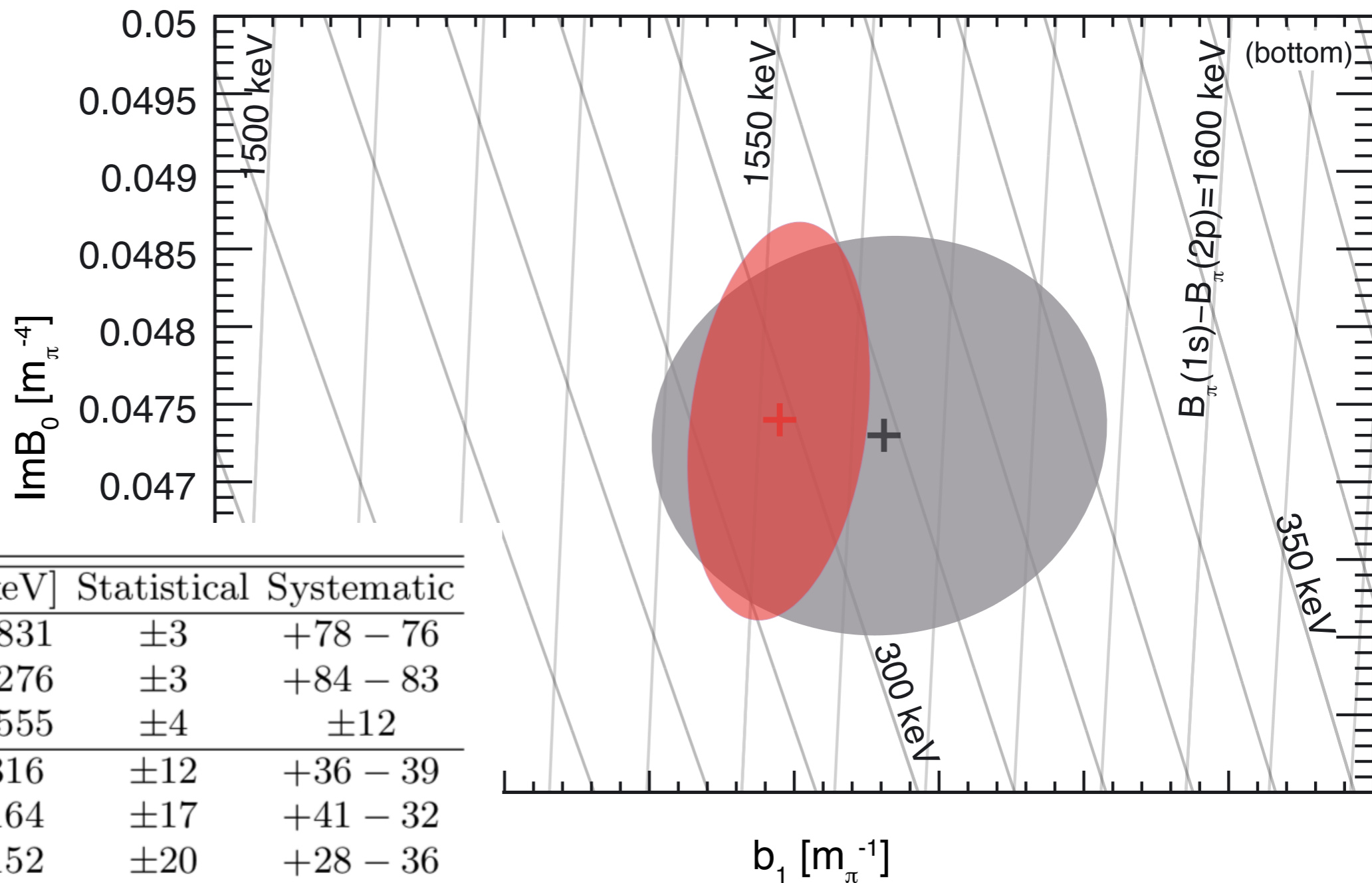
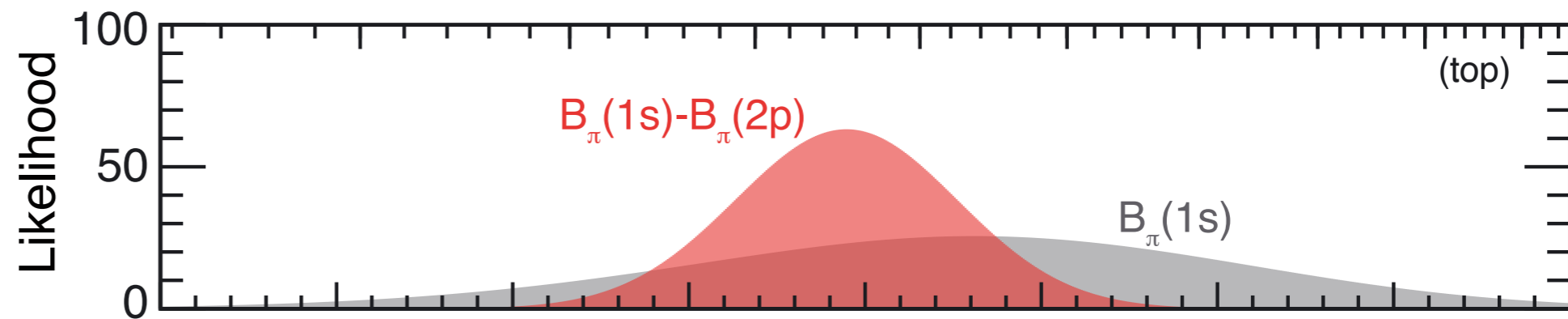


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b_1 [m_π^{-1}]

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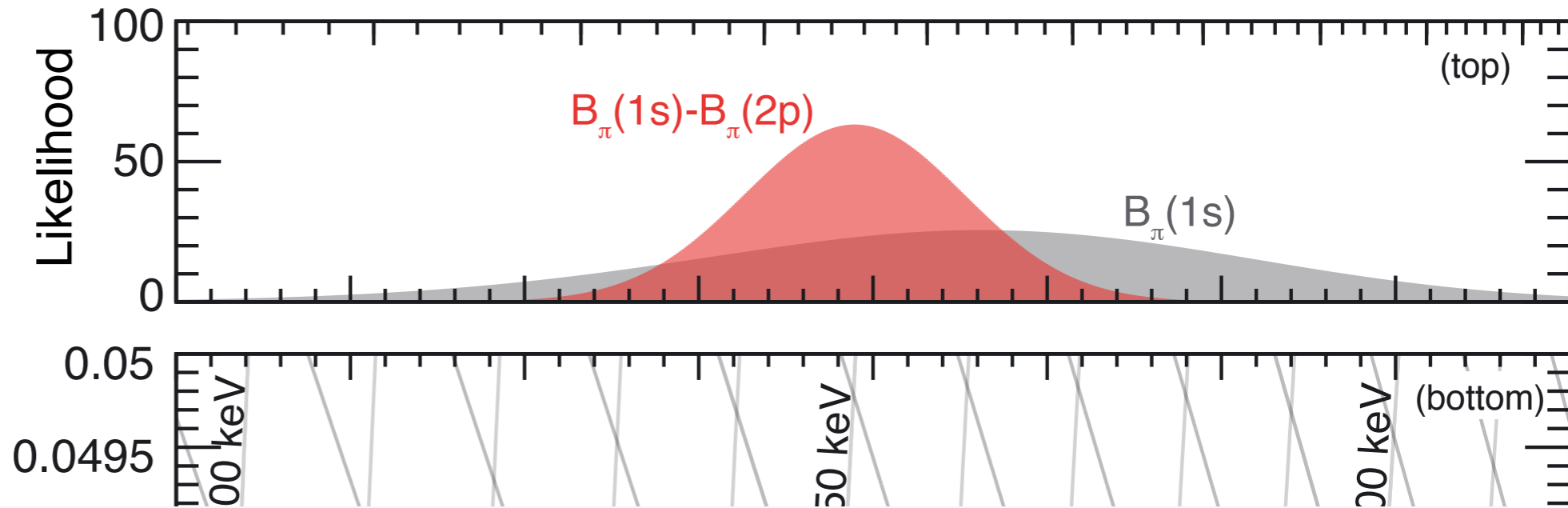


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b_1 parameter

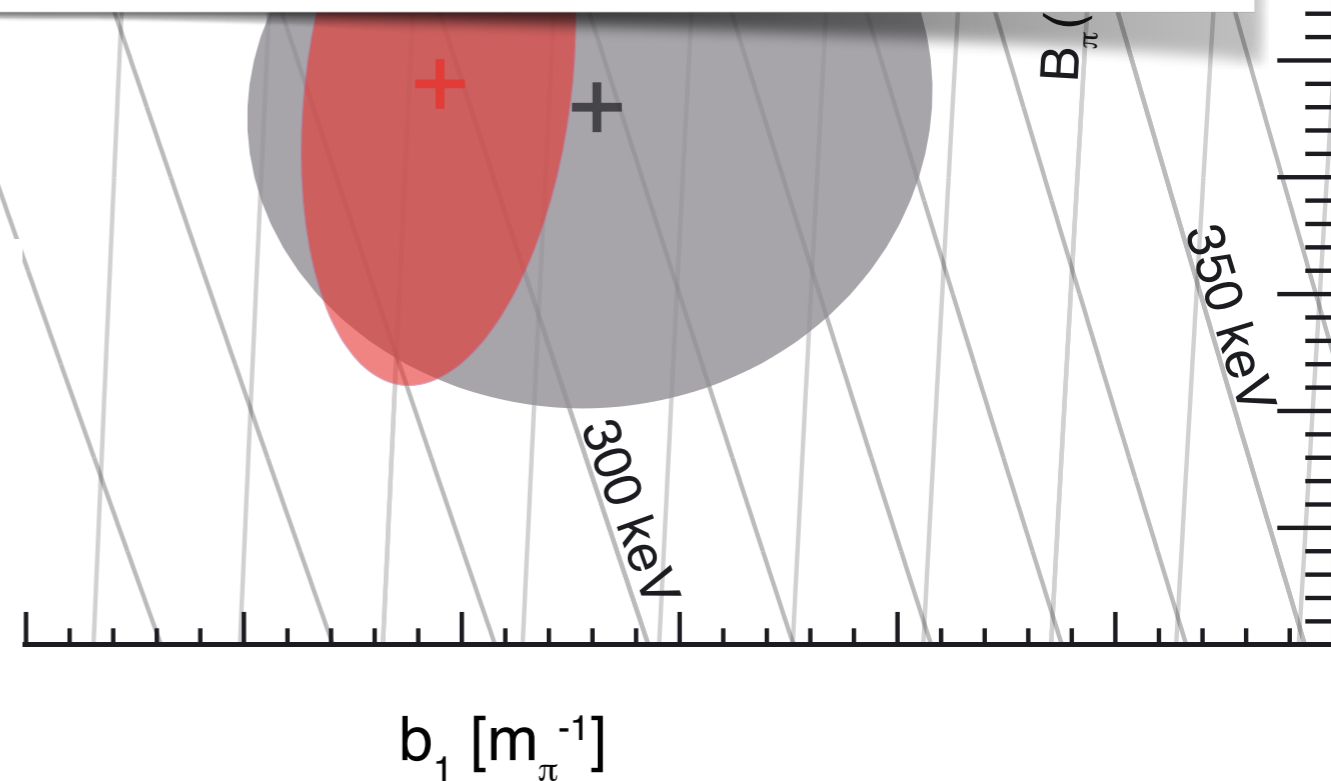
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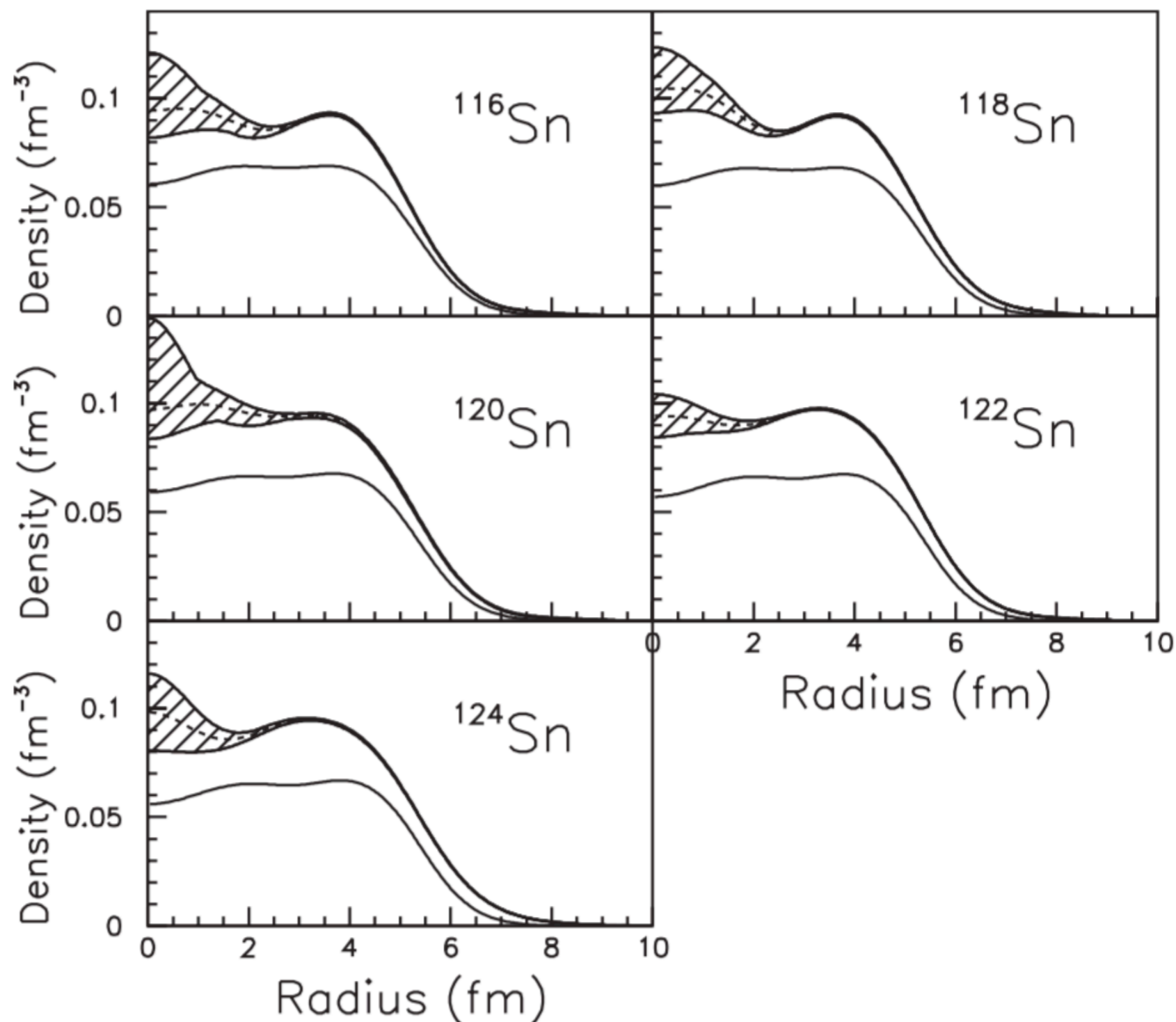
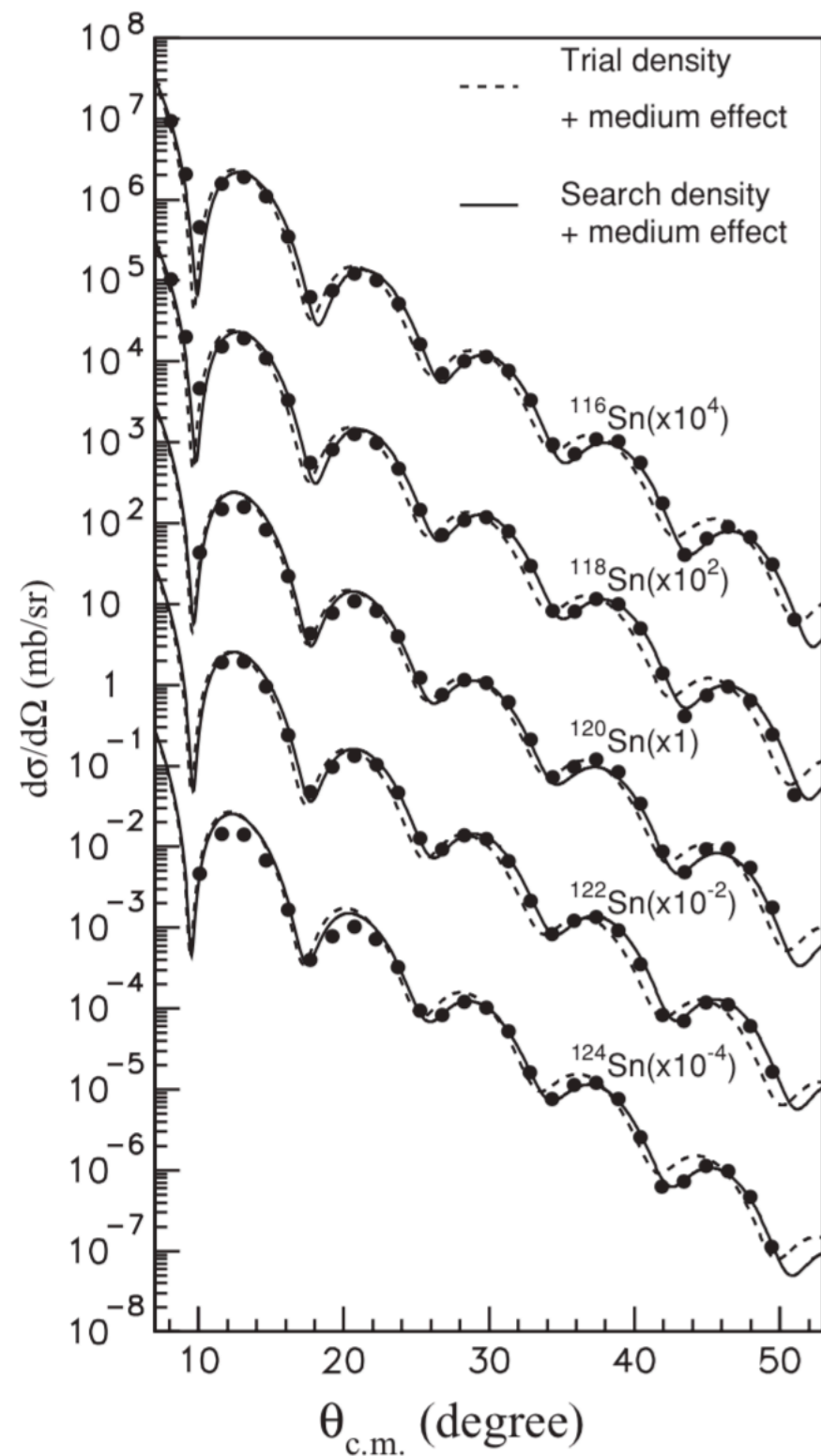
Before determining b_1 , we included updated theories, nuclear parameters etc.

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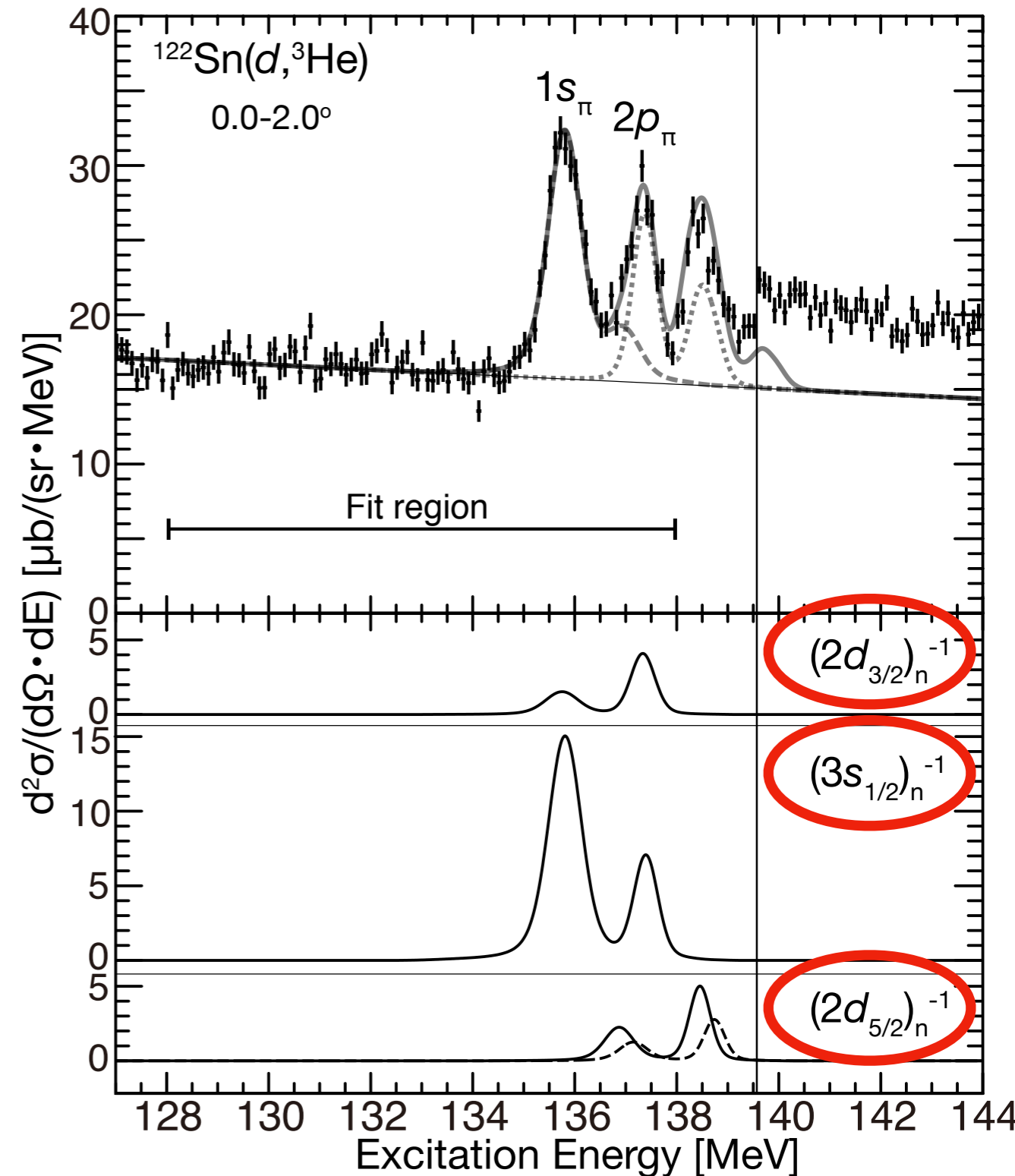
Nuclear ρ distribution measured in $\text{Sn}(p,p')$

at RCNP, Osaka



Neutron spectroscopic factors in Sn isotopes

$(d, ^3\text{He})$ requires n -spectroscopic factor information

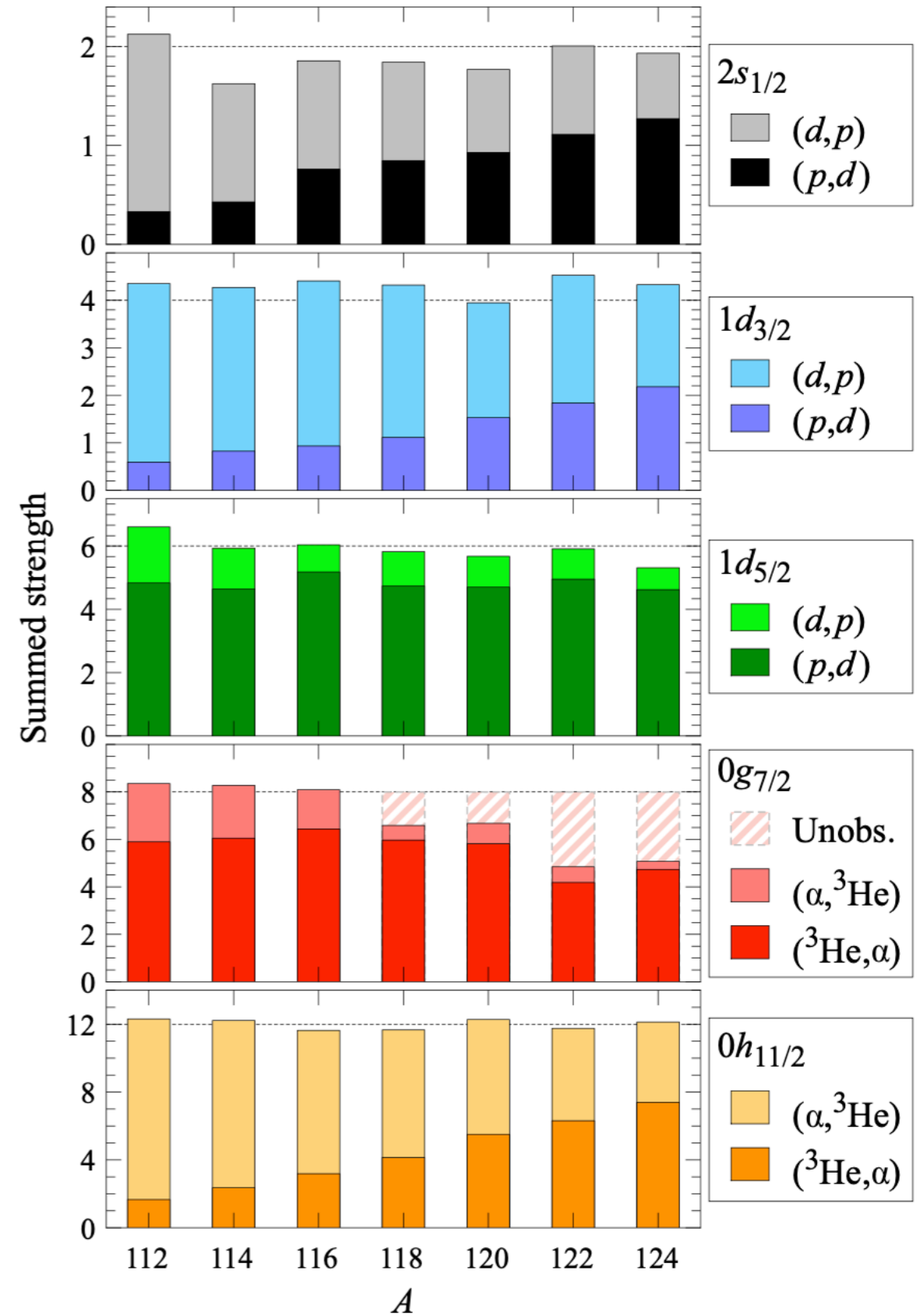
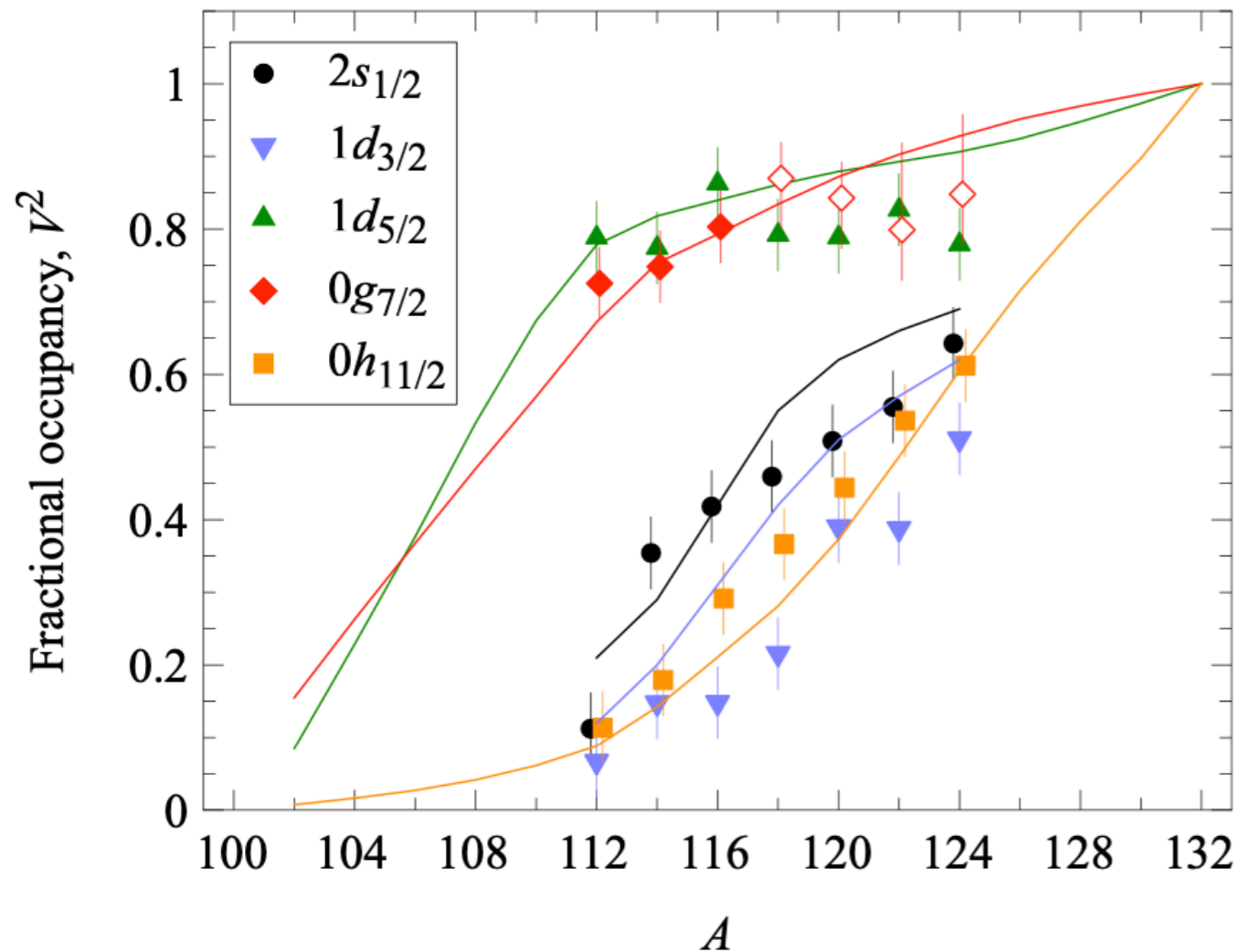


Neutron spectroscopic factors in Sn isotopes

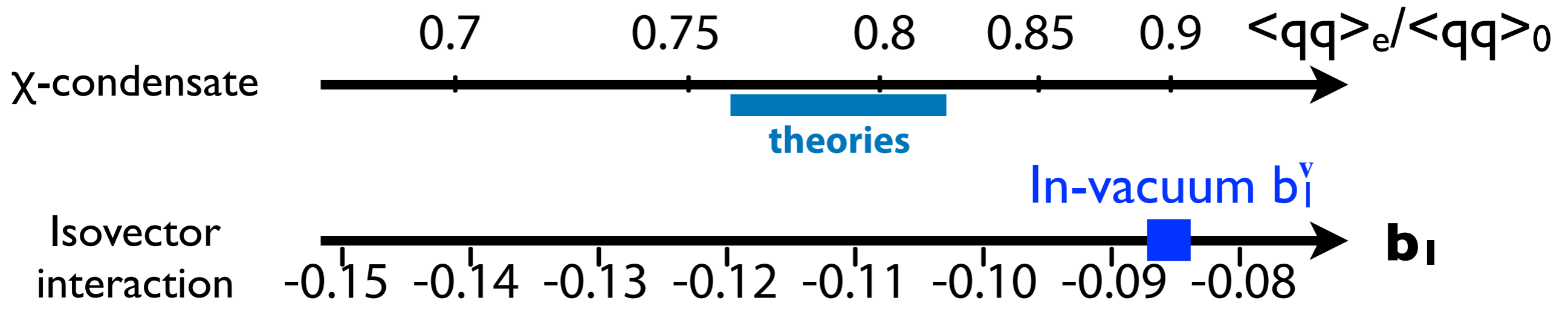
$(d, ^3\text{He})$ requires n -spectroscopic factor information

Spectroscopic factors are measured in (p,d) , (d,p) ... nuclear reactions confirming vacancy+occupancy $\sim 2j+1$

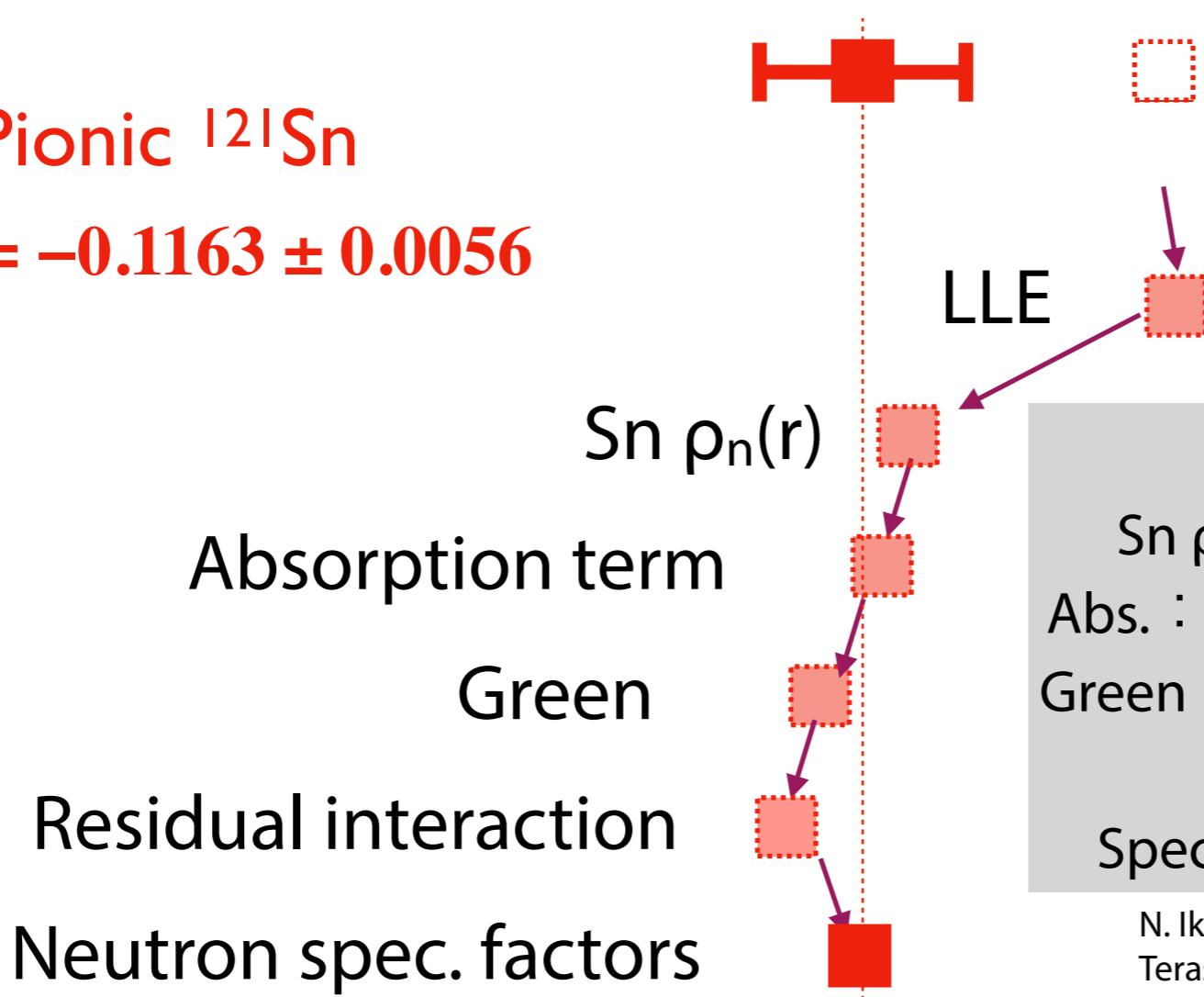
S.V. Szwece et al., PRC104, 054308 (2021)



Deduced b_1 with updated parameters



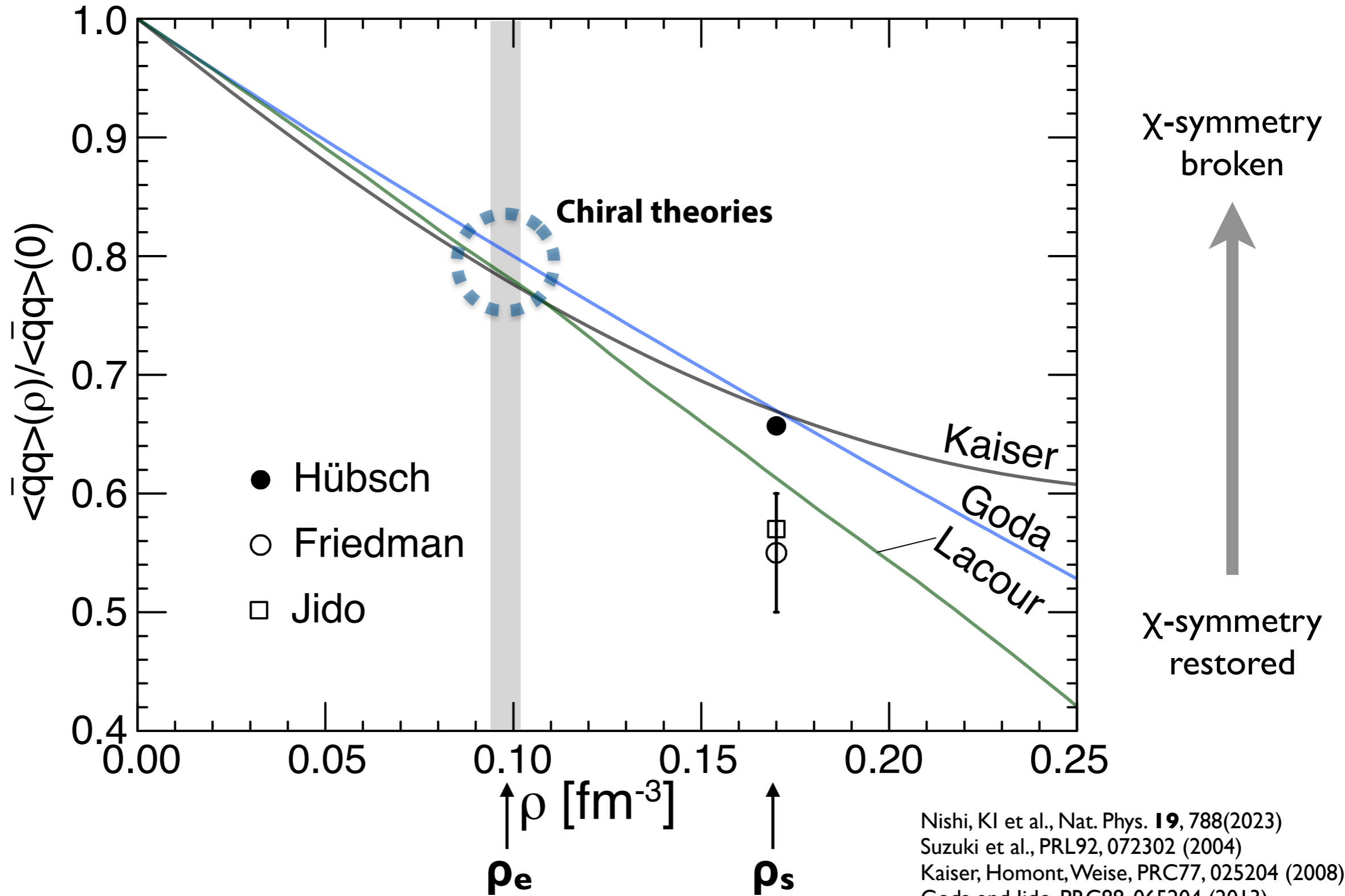
Pionic ^{121}Sn
 $b_1 = -0.1163 \pm 0.0056$



LLE : short-range effect
 Sn ρ : neutron density distribution
 Abs. : representation of absorption term
 Green : cross section calculation method
 Res. : Residual interaction
 Spec. : neutron spectroscopic factors

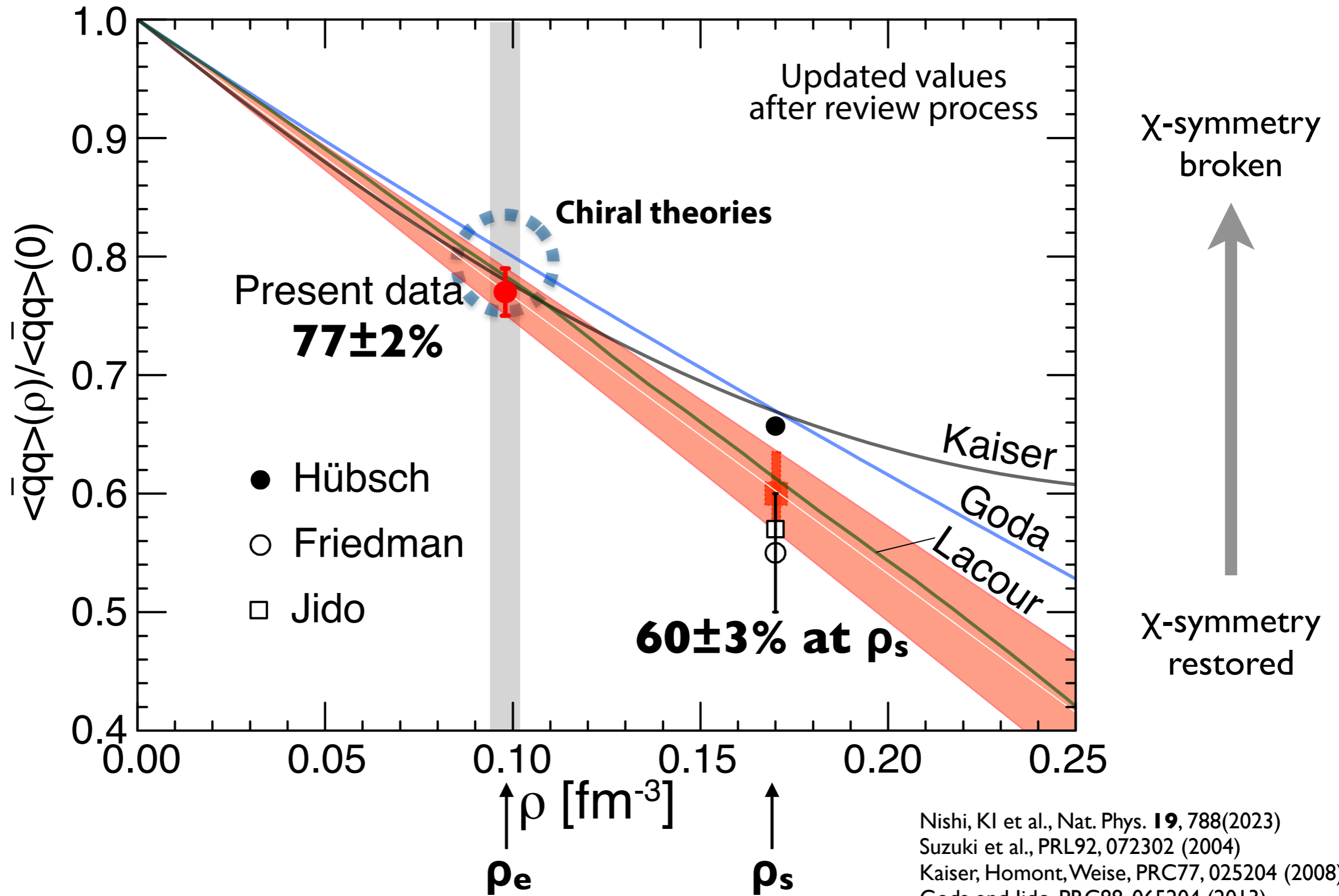
N. Ikeno et al., PTEP 2015, 033D01 (2015)
 Terashima et al., PHYSICAL REVIEW C 77, 024317 (2008)
 Nose-Togawa et al., PRC71, 061601(R) (2005)
 Szwec et al., PRC104,054308 (2021)

Result: deduced chiral condensate



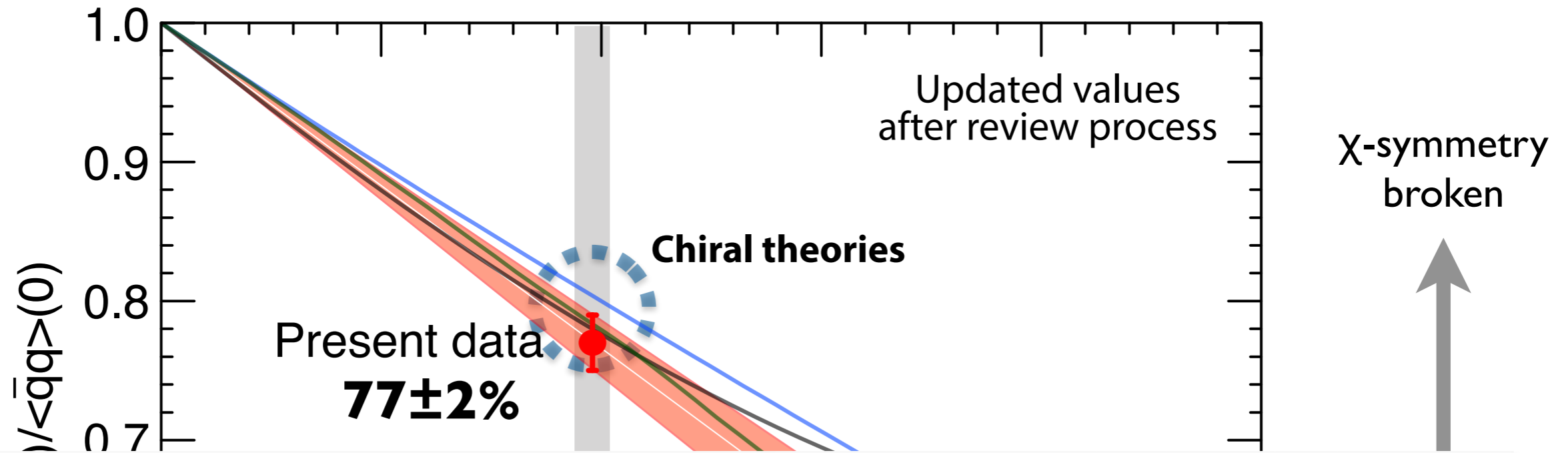
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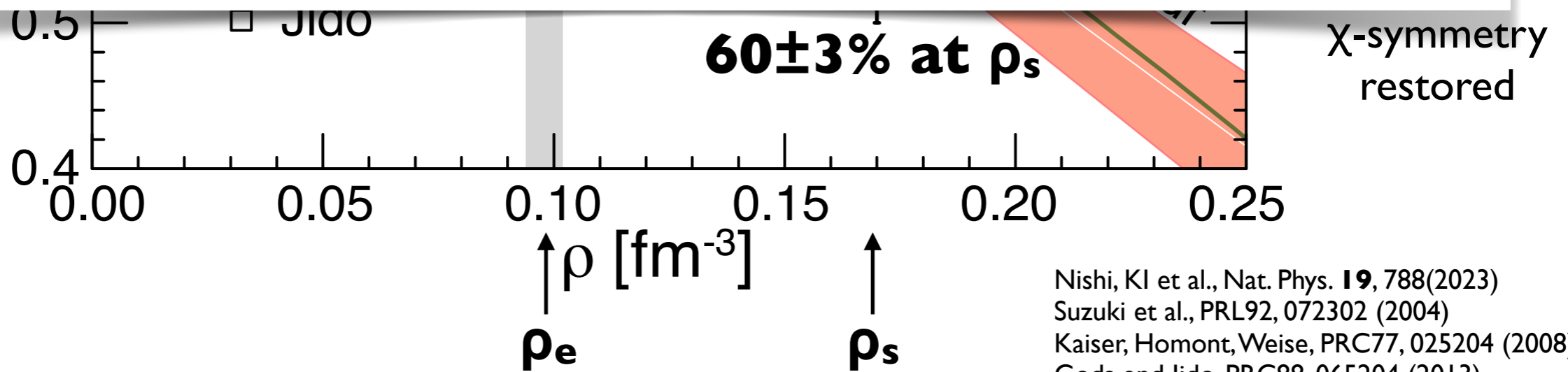


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Result: deduced chiral condensate



Support existence of non-trivial structure in the vacuum



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**To step further forward,
RIBF-135 for systematic study of Sn isotopes**

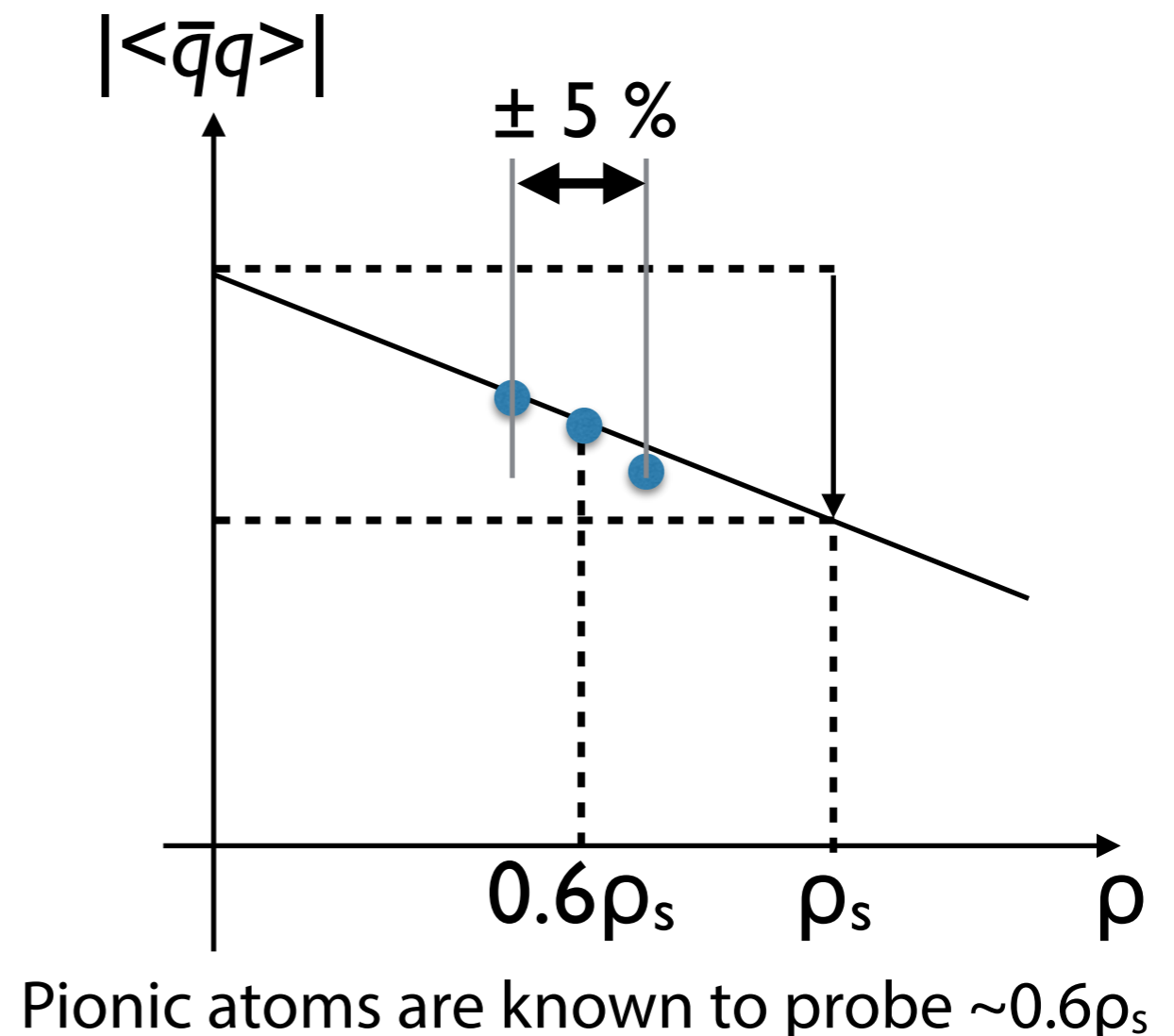
Density Dependence of Chiral Condensate

*Q. what can be achieved by measuring isotopes?
why not single isotope? How far can we discuss?*

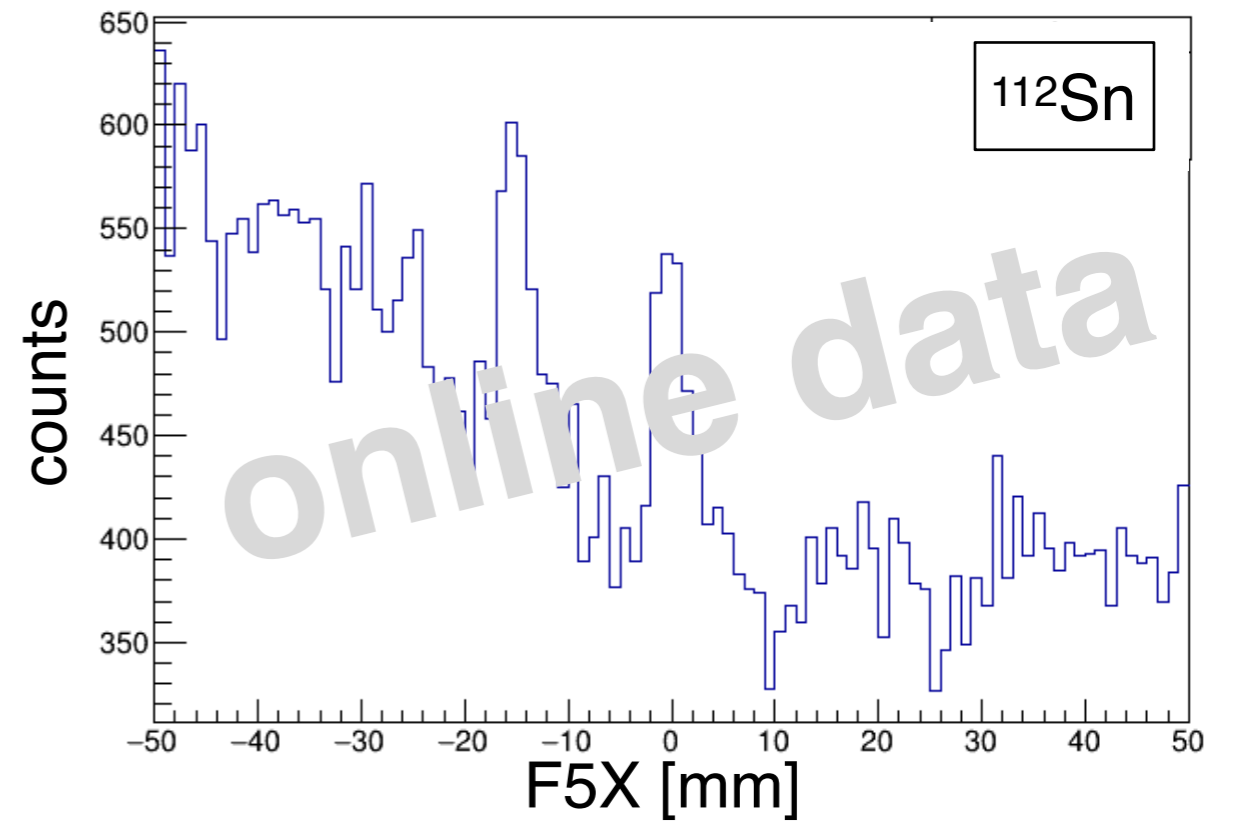
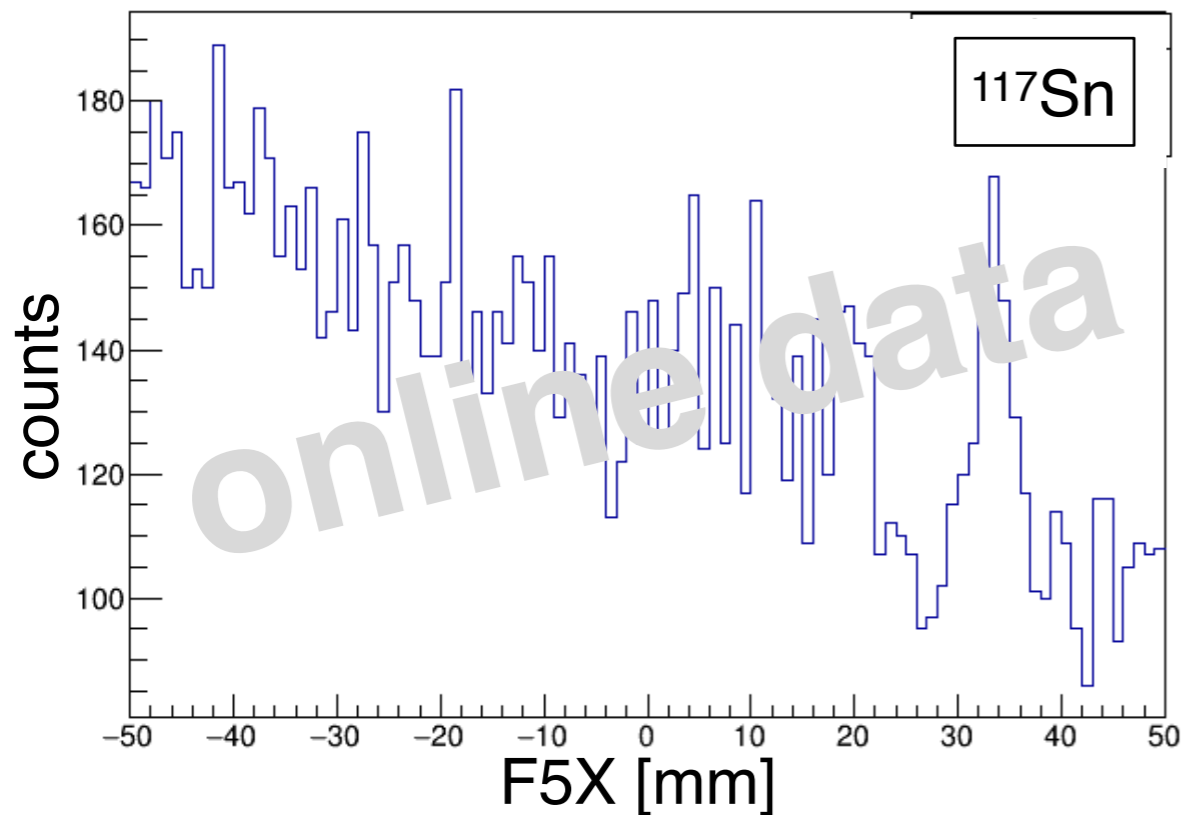
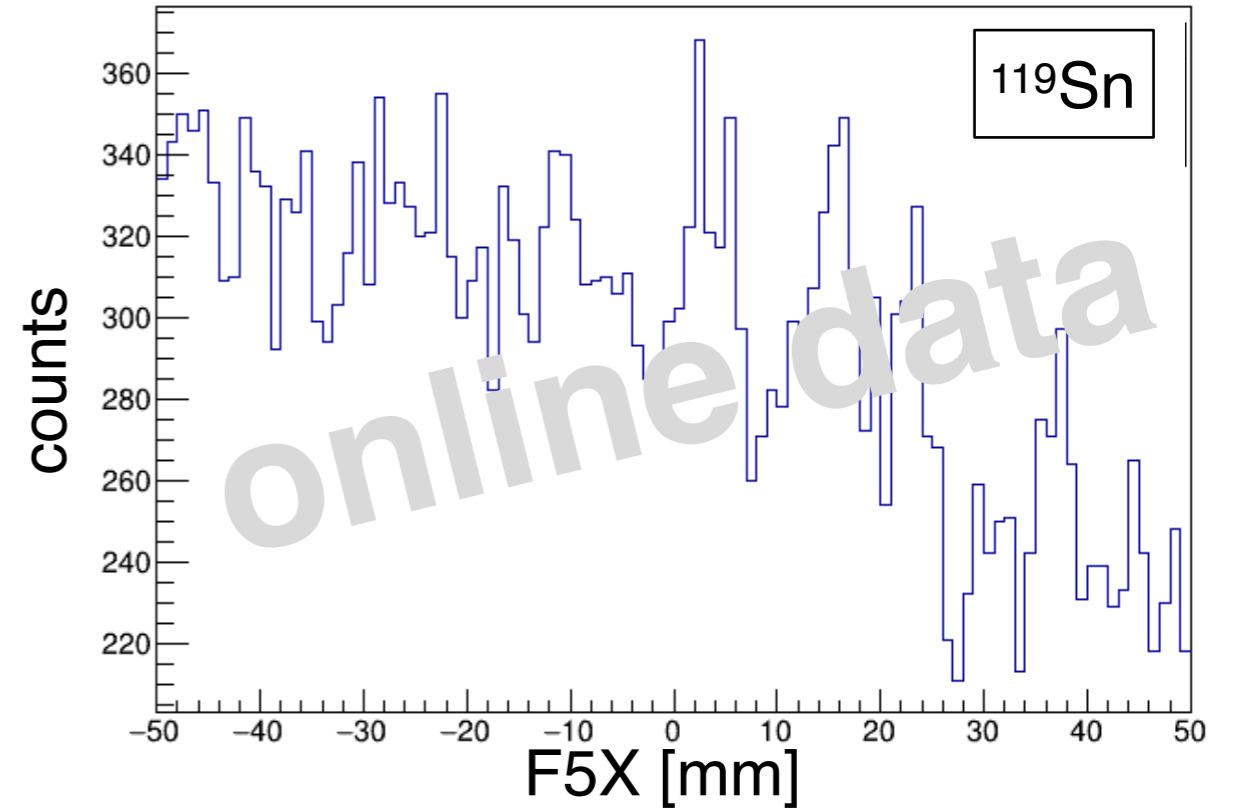
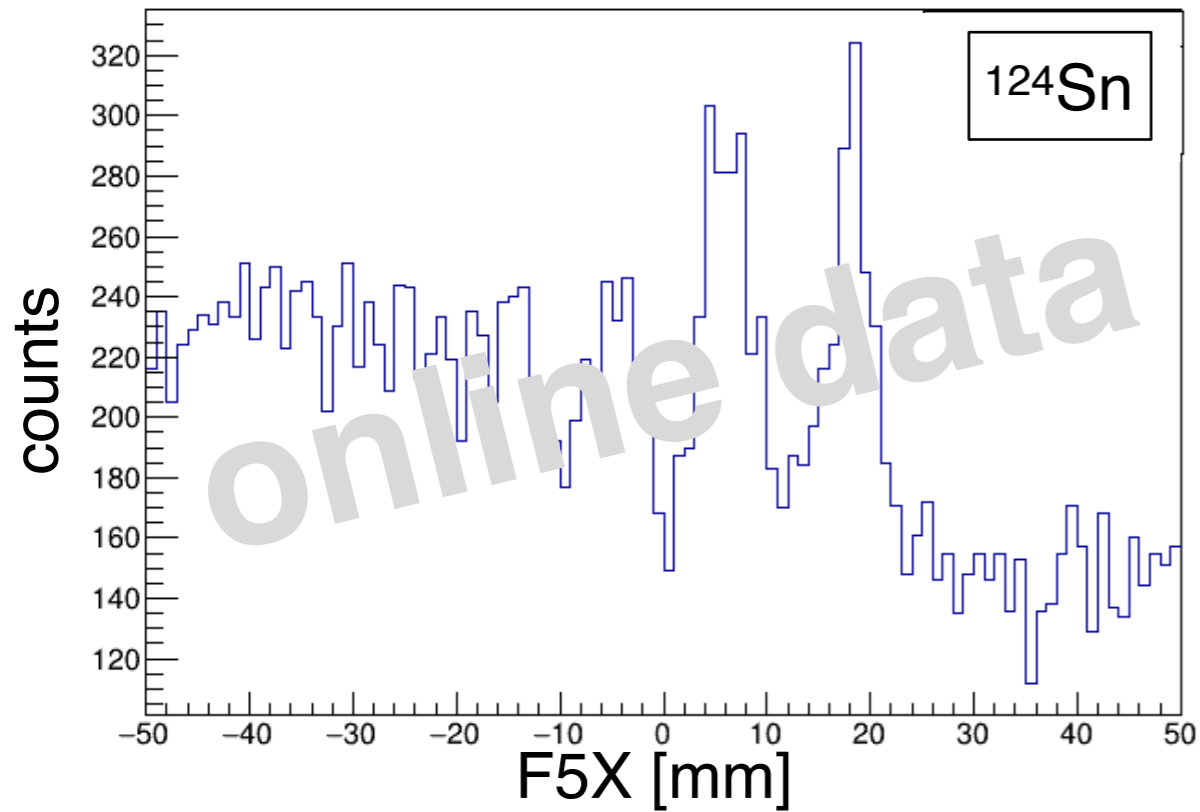
Ans.: ρ derivative of $\langle qq \rangle = d\langle qq \rangle/d\rho$ can be studied based on pionic Sn isotopes

Densities probed by pionic Sn with wide range of A

Important for $\sigma_{\pi N}$ for investigation of origin of matter mass



Online spectra in RIBF-135 (2021)



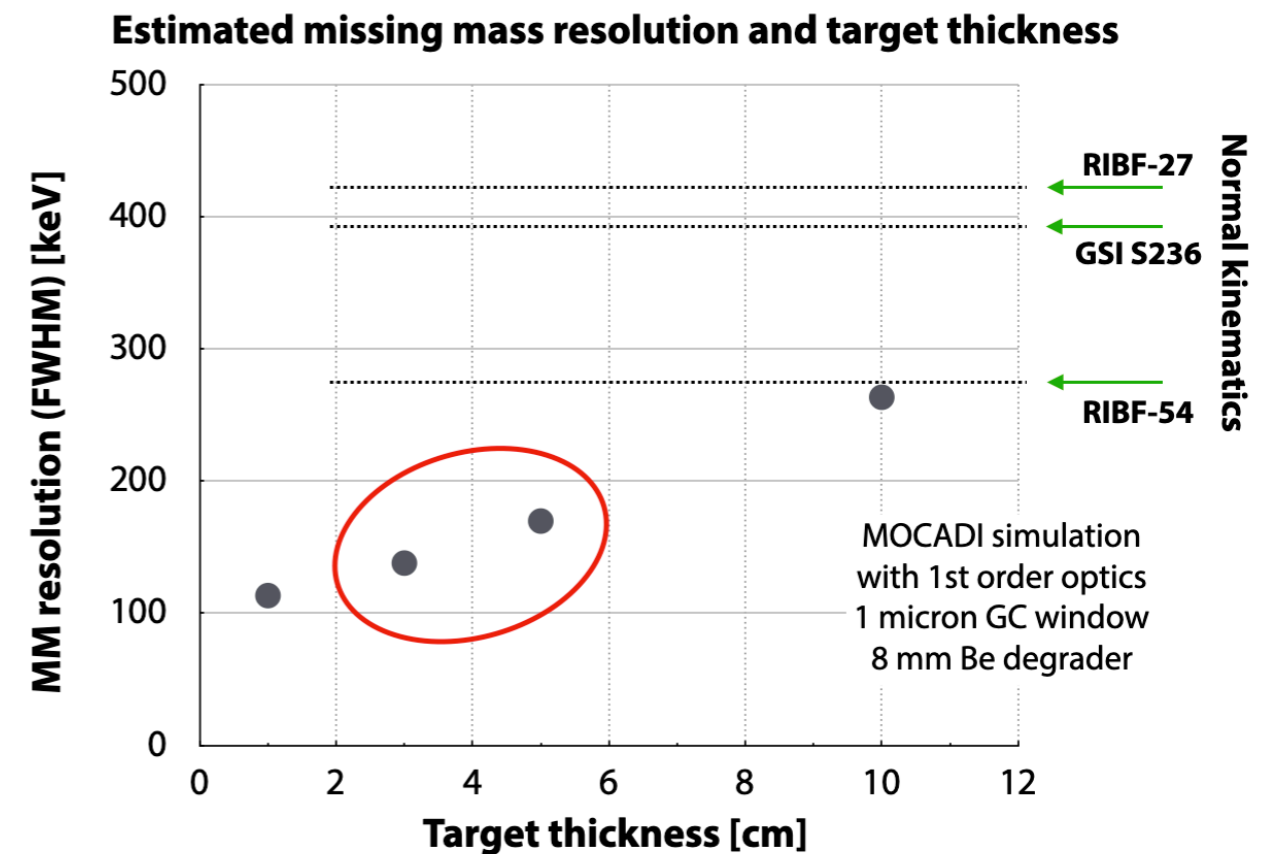
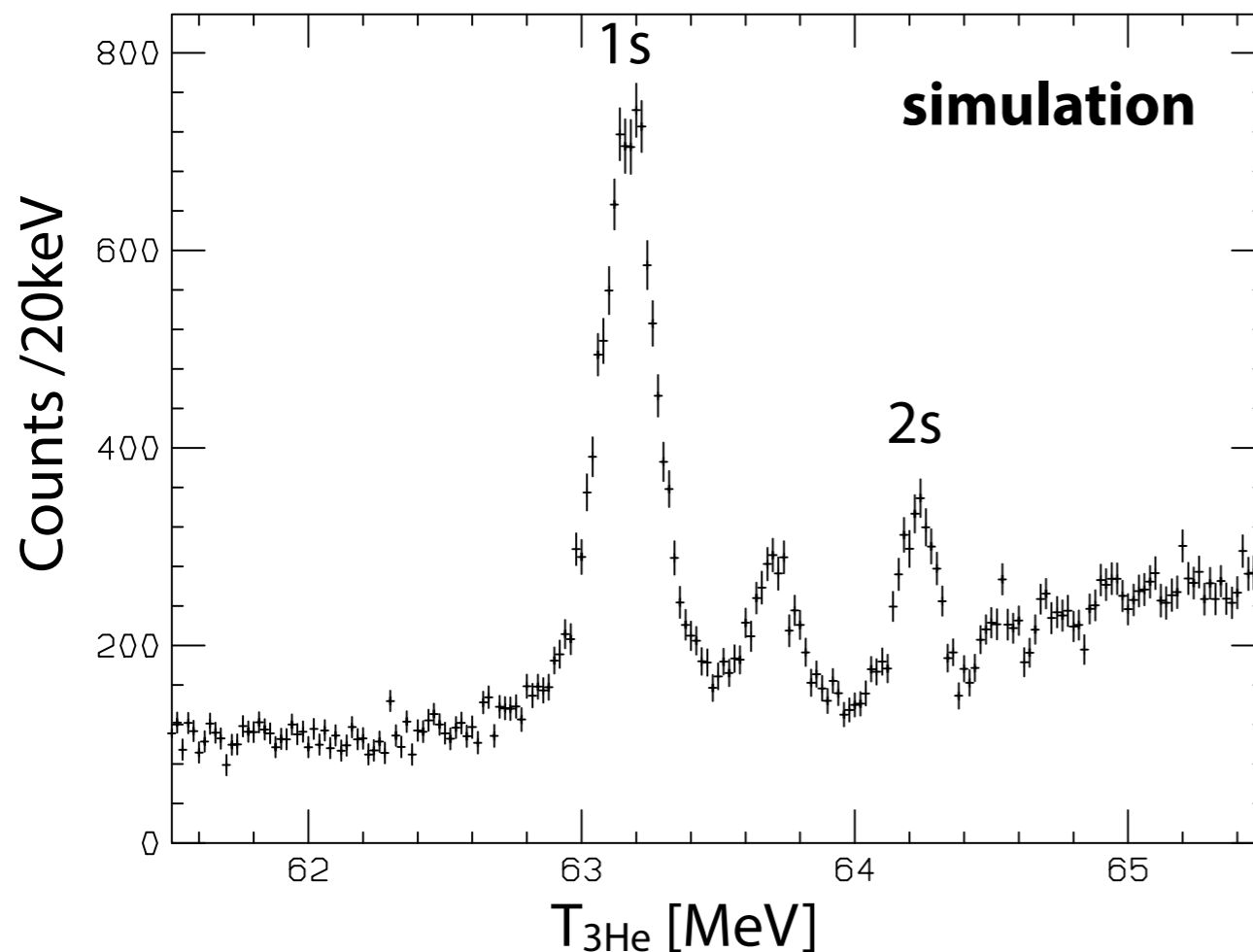
We are preparing for
Inverse kinematics (RIBF-214)

For kinematical reasons, ambiguities in the incident beam energy do not affect the results.

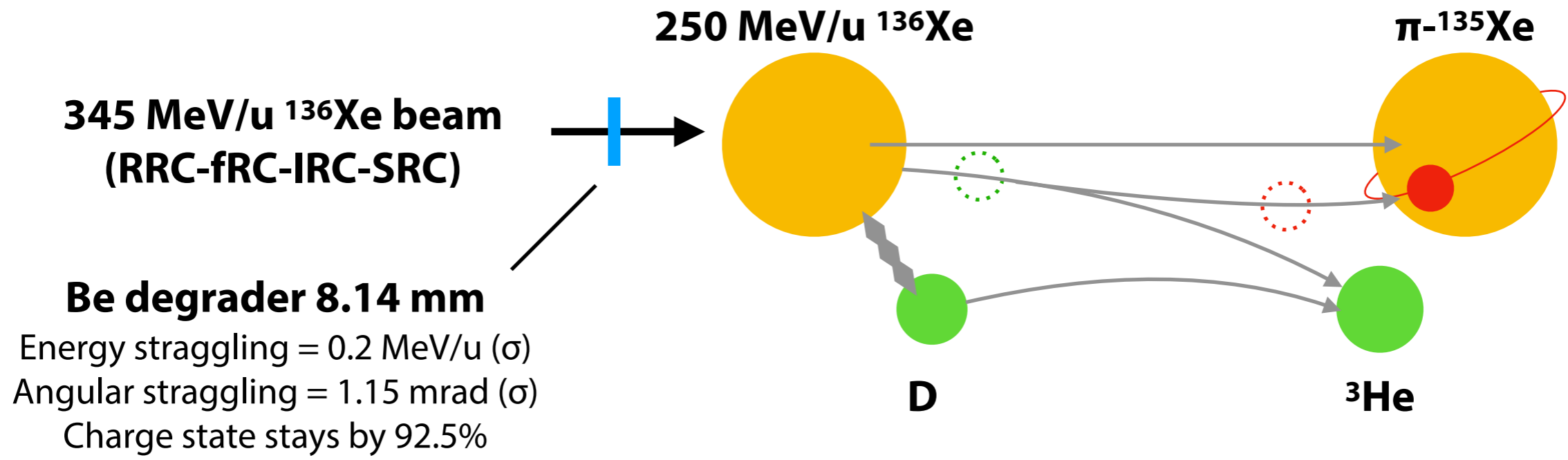
The resolution will be even improved.

Proposing $D(^{136}\text{Xe}, ^3\text{He})$ reaction at $T = 250 \text{ MeV/u}$ at RIBF

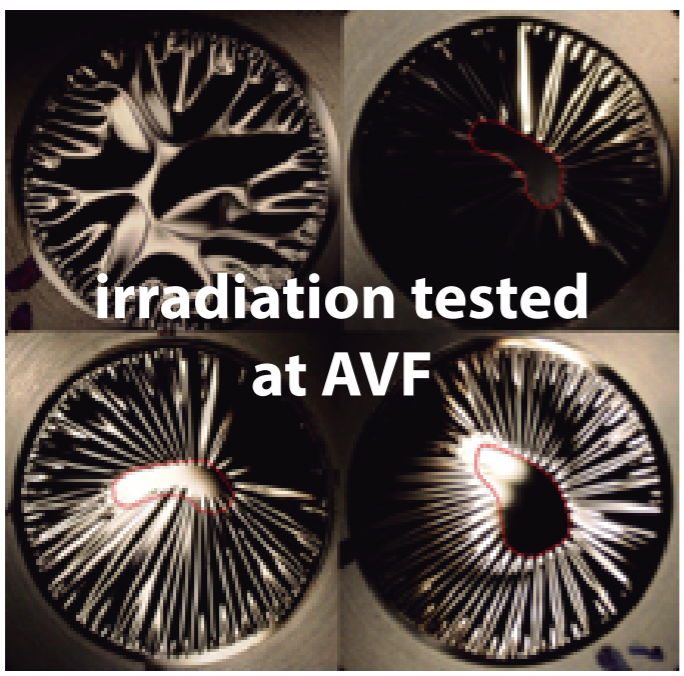
72 hours with $10^{10}/\text{s}$ ^{136}Xe beam



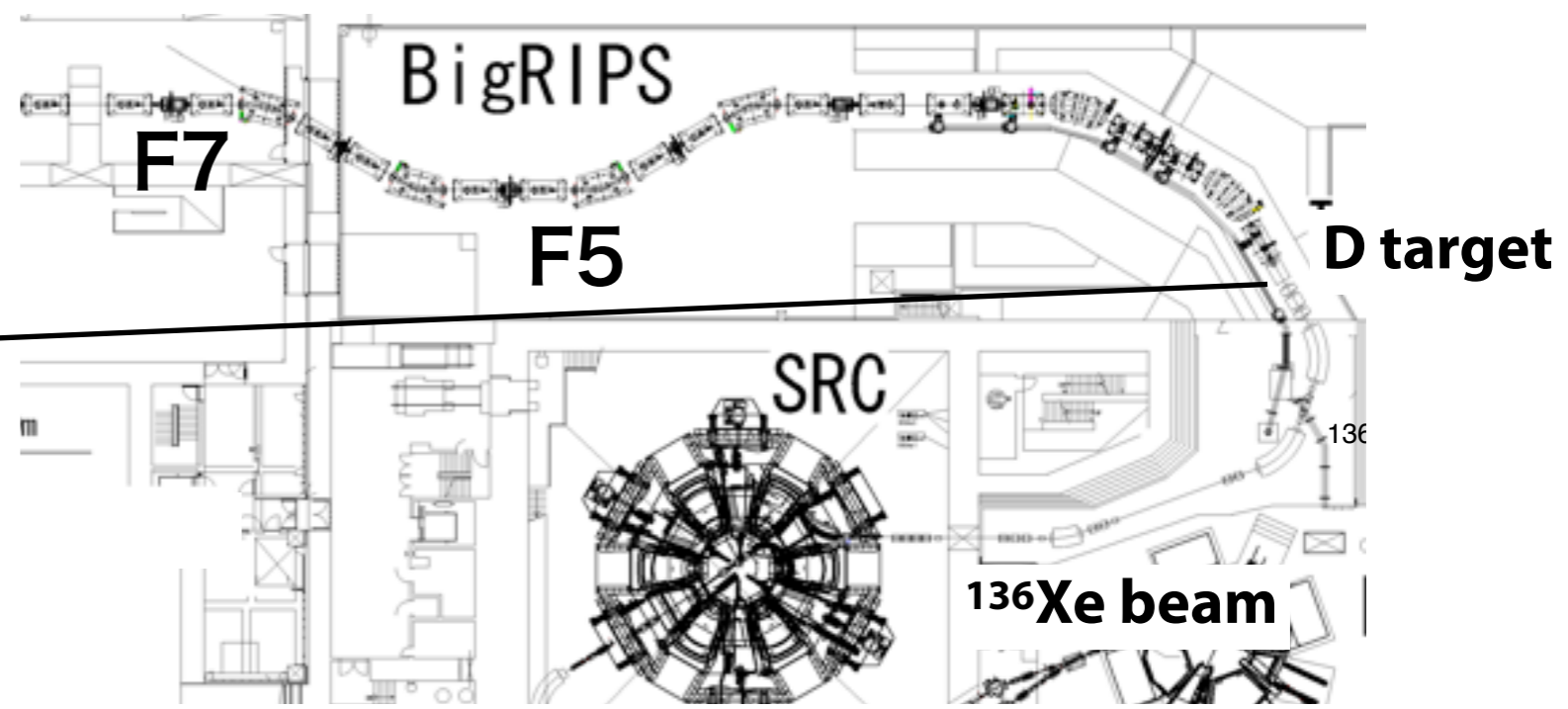
Experimental setup



1-atm deuterium gas target at F0 with **graphenic carbon** windows



BigRIPS as spectrometer to measure $\sim 1 \text{ Tm } ^3\text{He}$ momentum



S. Purushothaman et al., APR **53**, 134 (2019)

Summary

- Chiral condensate at ρ_e is evaluated to be reduced by $77\pm 2\%$, which is linearly extrapolated to $60\pm 3\%$ at the nuclear saturation density [N. Phys. 19 (2023) 788].
- The binding energies and widths of the pionic $1s$ and $2p$ states in $\text{Sn}121$ are determined with very high precision. Difference between the $1s$ and $2p$ values drastically reduces the systematic errors.
- Recent theoretical progress was adopted to the $\langle \bar{q}q \rangle$ deduction, which directly relates the chiral condensate and the pion-nucleus interaction.
- We included various updates for the first time. The updated parameters made substantial effects leading to much higher accuracy.
- For future, we are analyzing data of systematic study of pionic Sn isotopes to deduce ρ derivative of chiral condensate. We also plan measurement with “inverse kinematics” reactions for pionic xenon, which leads to future experiments for pionic unstable nuclei.



Hadron 2025 (Osaka, Mar. 27-31)
Hadron in Nucleus 2025 (Kyoto, Apr. 2-4)

for yielding different $V(0)$ values. If we allow $\text{Re } B_0$ to be varied, we have to change the $V(0)$ value accordingly. However the two parameters, b_0 and $\text{Re } B_0$, are interrelated as in the Seki-Masutani relation obtained by reading from Fig. 1 in Ref. [26],

$$b_0 \rho(0) + 0.50 \times \text{Re } B_0 \rho^2(0) = 0.062 \text{ fm}^{-2}. \quad (4.8)$$

This relation can be derived by asserting that the binding energies are determined essentially by the local potential strength at the nuclear radius ($r=R_0$). Since $\rho(R_0)$

