

Neutrino: their origin and their role in matter-antimatter asymmetry

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Outline

● Evidence of mass of neutrinos

↳ ν oscillation

↳ CP? Mass Hierarchy? Nature?

● Mass Models

↳ Majorana mass models

↳ Dirac mass models

● Explanation of Matter Asymmetry

↳ Lepton number violation via Majorana mass

↳ Dirac Leptogenesis

● Correlation & its constraints from Flavor Physics

Evidence of ν -mass

Discovery of ν -oscillation:

- The flavor eigenstate of ν 's are admixtures of mass eigenstates

$$\nu_\alpha = \sum_i U_{\alpha i} \nu_i$$

$$U_{\alpha i} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{atmospheric angle}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}}_{\text{Reactor angle}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar angle}} \underbrace{\begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{only if the particle is Majorana}}$$

$U_{\alpha i} \Rightarrow$ Pontecorvo-Maki-Nakagawa-Satake (PMNS) matrix

Evidence of ν -mass

Discovery of ν -oscillation:

- The flavor eigenstate of ν 's are admixtures of mass eigenstates

$$\nu_\alpha = \sum_i U_{\alpha i} \nu_i$$

$$P_{\alpha \rightarrow \beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2\left(\frac{\Delta m_{ij}^2 L}{4E}\right) + 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin\left(\frac{\Delta m_{ij}^2 L}{2E}\right)$$

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

$$\frac{\Delta m^2 L}{4E} \approx 1.27 \frac{\Delta m^2}{\text{eV}^2} \frac{L}{\text{km}} \frac{\text{GeV}}{E}$$

Evidence of ν -mass

Global Fit to neutrino Oscillation data

| parameter | best fit $\pm 1\sigma$ | 2σ range | 3σ range |
|--|---------------------------|-----------------|-----------------|
| $\Delta m_{21}^2 [10^{-5} \text{eV}^2]$ | $7.50^{+0.22}_{-0.20}$ | 7.12–7.93 | 6.94–8.14 |
| $ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$ (NO) | $2.55^{+0.02}_{-0.03}$ | 2.49–2.60 | 2.47–2.63 |
| $ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$ (IO) | $2.45^{+0.02}_{-0.03}$ | 2.39–2.50 | 2.37–2.53 |
| $\sin^2 \theta_{12} / 10^{-1}$ | 3.18 ± 0.16 | 2.86–3.52 | 2.71–3.69 |
| $\theta_{12} / ^\circ$ | 34.3 ± 1.0 | 32.3–36.4 | 31.4–37.4 |
| $\sin^2 \theta_{23} / 10^{-1}$ (NO) | 5.74 ± 0.14 | 5.41–5.99 | 4.34–6.10 |
| $\theta_{23} / ^\circ$ (NO) | 49.26 ± 0.79 | 47.37–50.71 | 41.20–51.33 |
| $\sin^2 \theta_{23} / 10^{-1}$ (IO) | $5.78^{+0.10}_{-0.17}$ | 5.41–5.98 | 4.33–6.08 |
| $\theta_{23} / ^\circ$ (IO) | $49.46^{+0.60}_{-0.97}$ | 47.35–50.67 | 41.16–51.25 |
| $\sin^2 \theta_{13} / 10^{-2}$ (NO) | $2.200^{+0.069}_{-0.062}$ | 2.069–2.337 | 2.000–2.405 |
| $\theta_{13} / ^\circ$ (NO) | $8.53^{+0.13}_{-0.12}$ | 8.27–8.79 | 8.13–8.92 |
| $\sin^2 \theta_{13} / 10^{-2}$ (IO) | $2.225^{+0.064}_{-0.070}$ | 2.086–2.356 | 2.018–2.424 |
| $\theta_{13} / ^\circ$ (IO) | $8.58^{+0.12}_{-0.14}$ | 8.30–8.83 | 8.17–8.96 |
| δ / π (NO) | $1.08^{+0.13}_{-0.12}$ | 0.84–1.42 | 0.71–1.99 |
| $\delta / ^\circ$ (NO) | 194^{+24}_{-22} | 152–255 | 128–359 |
| δ / π (IO) | $1.58^{+0.15}_{-0.16}$ | 1.26–1.85 | 1.11–1.96 |
| $\delta / ^\circ$ (IO) | 284^{+26}_{-28} | 226–332 | 200–353 |

Evidence of ν -mass

Global Fit to neutrino Oscillation data

Mass Hierarchy? = [

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CP? }

$m_1 ?$ (NO)
 $m_3 ?$ (IO)

Evidence of ν -mass

Nature of ν , Majorana? Dirac?

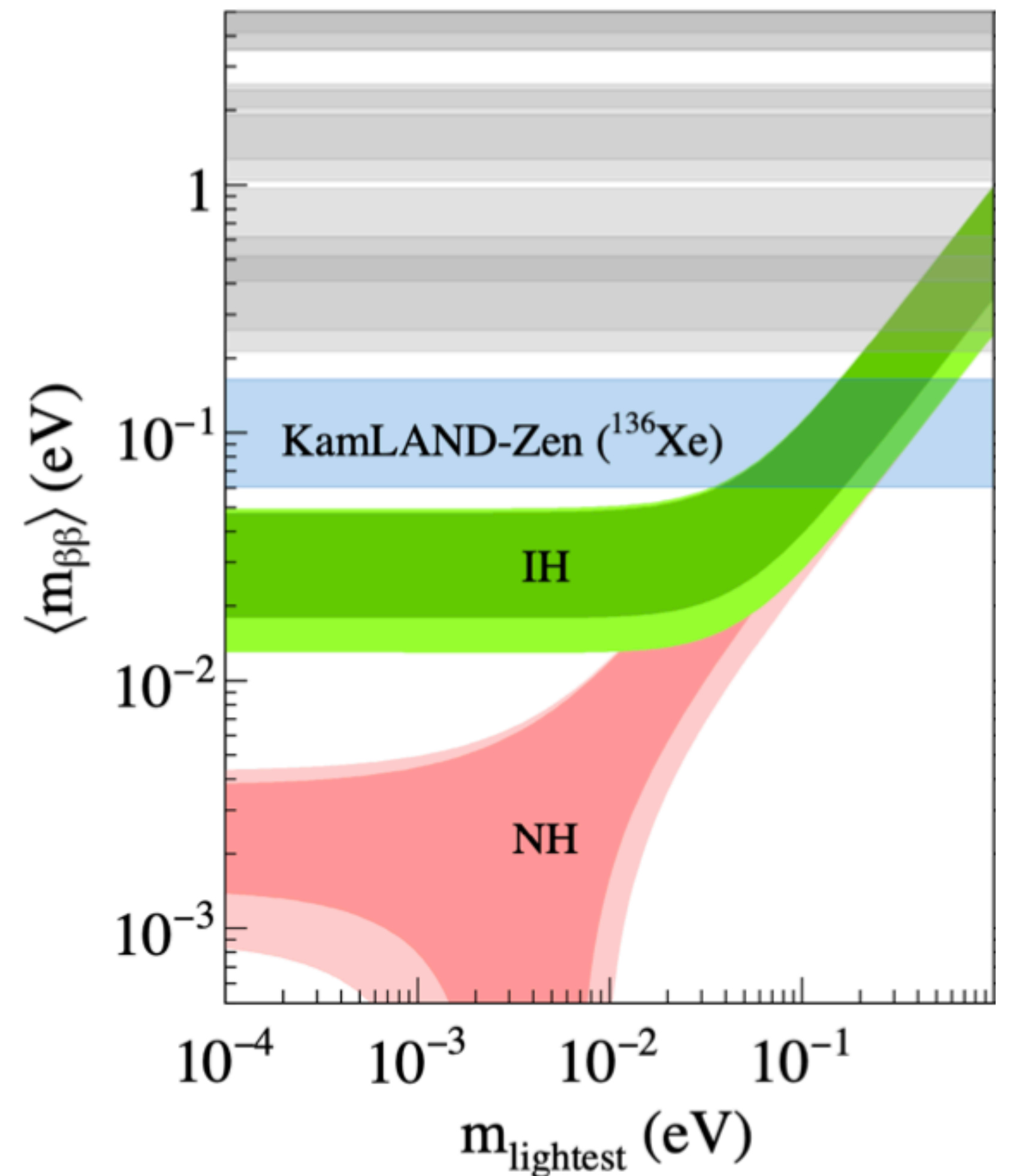
$0\nu\beta\beta$ Decay Experiment

If ν s are majorana the effective mass is given as

$$m_{\beta\beta} = \sum_i (U_{ei})^2 m_i$$

$$m_{\beta\beta} = c_{12}^2 c_{13}^2 e^{i\alpha_1} m_1 + c_{13}^2 s_{12}^2 e^{i\alpha_2} m_2 + s_{13}^2 m_3$$

[arXiv:2108.09364]

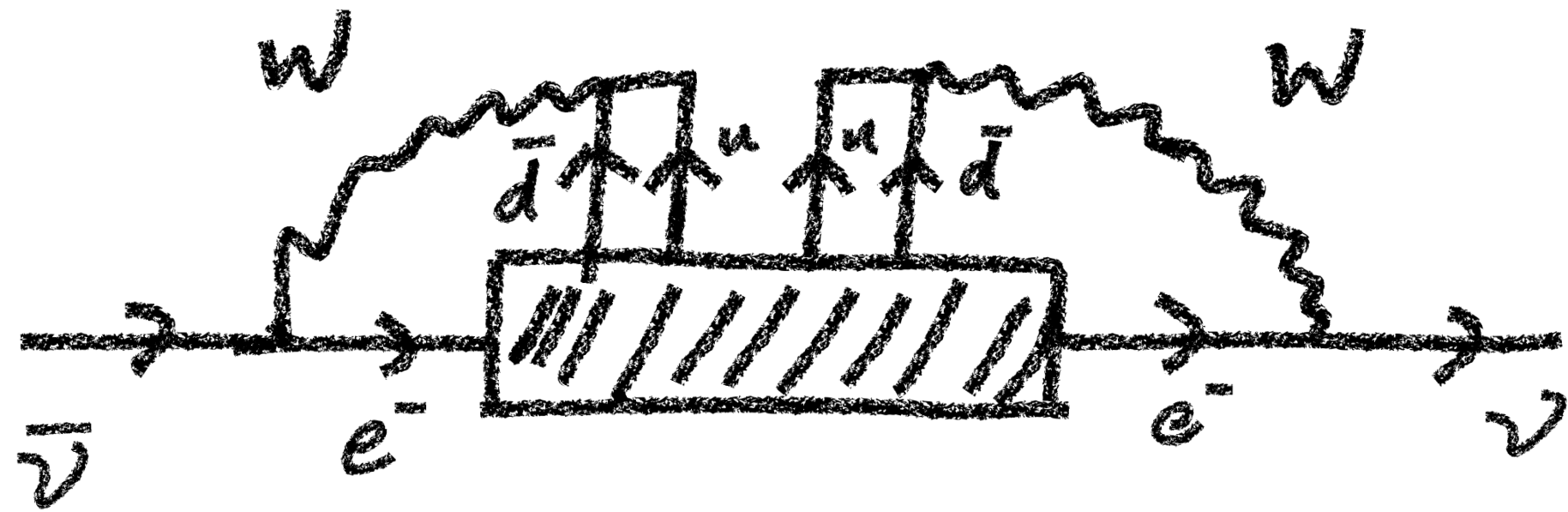


Evidence of ν -mass

Nature of ν , Majorana? Dirac?

If $0\nu\beta\beta$ is discovered \Rightarrow Lepton number is violated

$\Rightarrow \nu_s$ are Majorana!!



Schechter Valle theorem [PRD 25 (1982)
2951]

Mass Models

● Majorana Neutrinos

Mass Models

● Majorana Neutrinos

■ Type I, III

Mass Models

● Majorana Neutrinos

- Type I, III
- Type II

Mass Models

● Majorana Neutrinos

- Type I, III
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- Inverse See-Saw

Mass Models

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- Radiative neutrino mass model

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⇒ Many admixtures

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① Majorana Neutrinos

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- Type II
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⇒ Many admixtures

② Dirac Neutrinos

- Type-I
- Inverse See-Saw
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& Tree $\Rightarrow y_D \bar{L} H \nu_R$

Mass Models (Majorana Scenario)

Type I, III:

| | |
|-------|--|
| | $SU(3)_c \times SU(2)_L \times U(1)_Y$ |
| N_R | $(1, 1, 0)$ |

► Type I: $\mathcal{L} \supset y_\nu \bar{L} H N_R + M_R \bar{N}_R^c N_R + h.c$



Breaking B-L: New scale and new symmetry beyond SM.

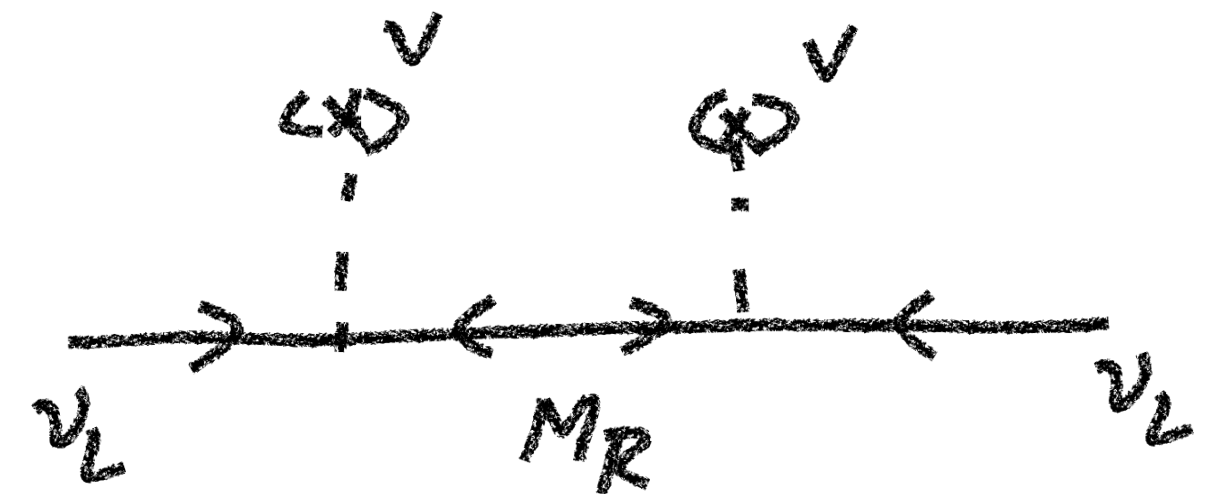
► Type III: $\mathcal{L} \supset y_\nu \bar{L} \Psi H + M_R \text{Tr}[\bar{\Psi}^c \Psi] + h.c$



| | |
|--------|--|
| | $SU(3)_c \times SU(2)_L \times U(1)_Y$ |
| Ψ | $(1, 3, 0)$ |

$$\Psi = \begin{pmatrix} \Sigma^0/\sqrt{2} & \Sigma^+ \\ \Sigma^- & \Sigma^0/\sqrt{2} \end{pmatrix}$$

$$m_\nu = \begin{pmatrix} 0 & m_D \\ M_D^T & M \end{pmatrix} \approx -y_\nu^2 \frac{v^2}{M_R}$$



Mass Models (Majorana Scenario)

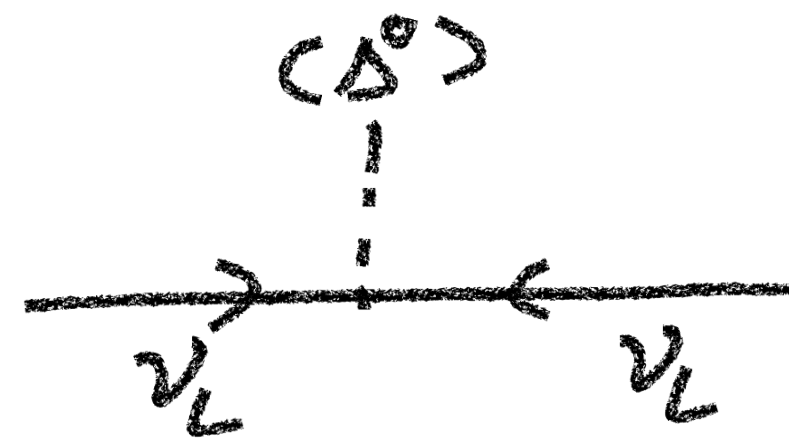
Type II :

$$\mathcal{L} \supset y_\nu \bar{L}^c i\sigma_2 \Delta L + \text{h.c.}$$

$$\Delta = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$$

| | |
|----------|--|
| | $SU(3)_c \times SU(2)_L \times U(1)_Y$ |
| Δ | $(1, 3, 1)$ |

$$m_\nu = y_{\nu\alpha\beta} \langle \Delta^0 \rangle$$



Mass Models (Majorana Scenario)

Inverse See-Saw

$$\mathcal{L} \supset y_D \bar{L} H N_R + y_\nu \phi \bar{S}_R^c N_R + \mu \bar{S}_R^c S_R$$

$$m_\nu = \begin{pmatrix} 0 & y_D v & 0 \\ y_D v & 0 & y_\nu v \phi \\ 0 & y_\nu v \phi & \mu \end{pmatrix} \approx - y_D^T y_\nu^{-1} \mu y_\nu^{-1} y_D \left(\frac{v \phi}{v} \right)^2$$

| | $SU(3)_c \times SU(2)_L \times U(1)_Y$ |
|-------|--|
| S_R | (1, 1, 0) |
| N_R | (1, 1, 0) |

Mass Models (Majorana Scenario)

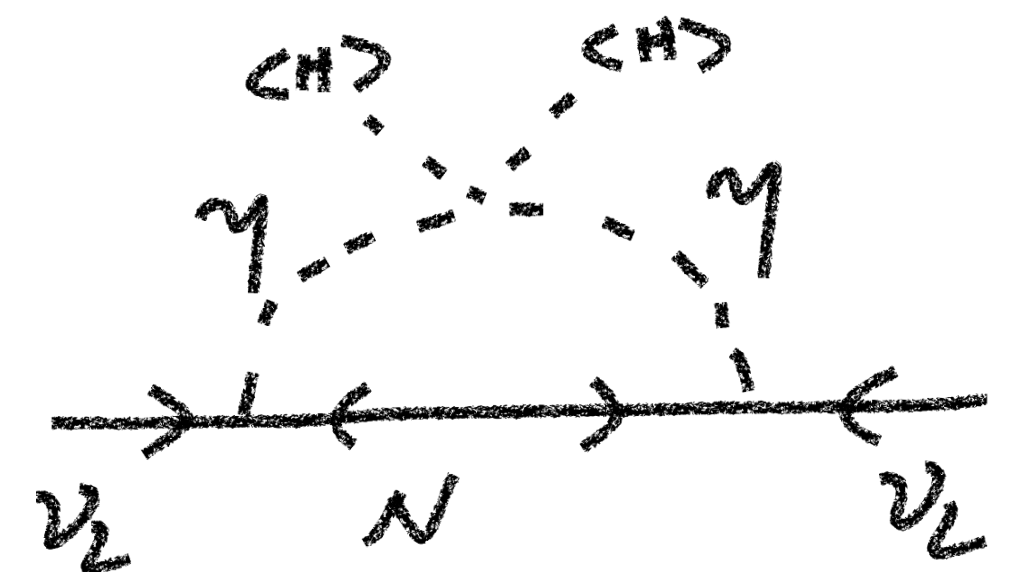
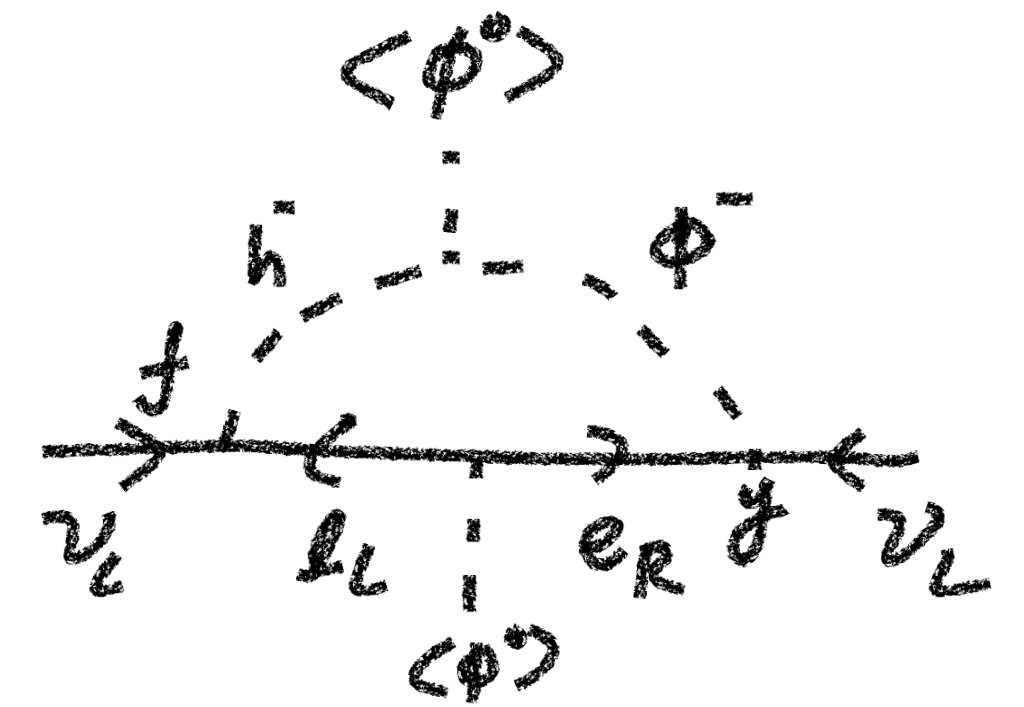
Radiative Neutrino Mass Model (1-Loop only)

Zee-Babu model

$$\mathcal{L} \supset y_D \bar{L} H e_R + f \bar{L} L h^\dagger + h.c$$

Scotogenic model

$$\mathcal{L} \supset y_D \bar{L} \gamma N_R + M \bar{N}_R^c N_R + h.c$$



Mass Models (Majorana Scenario)

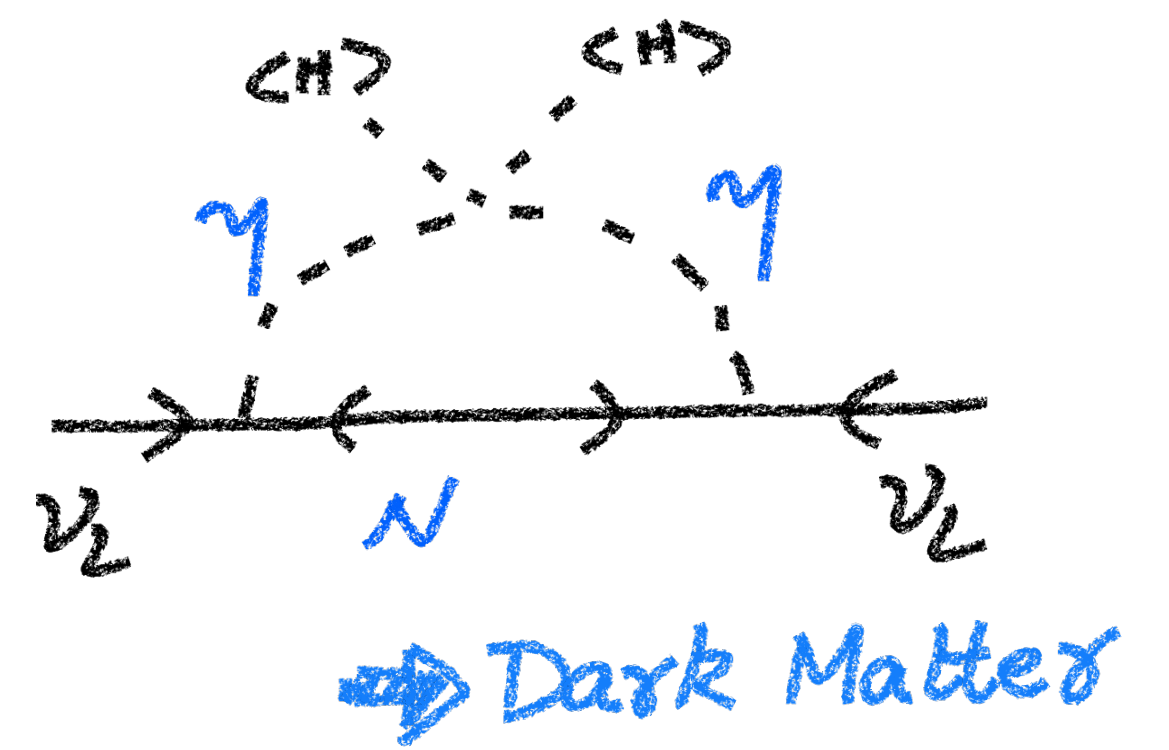
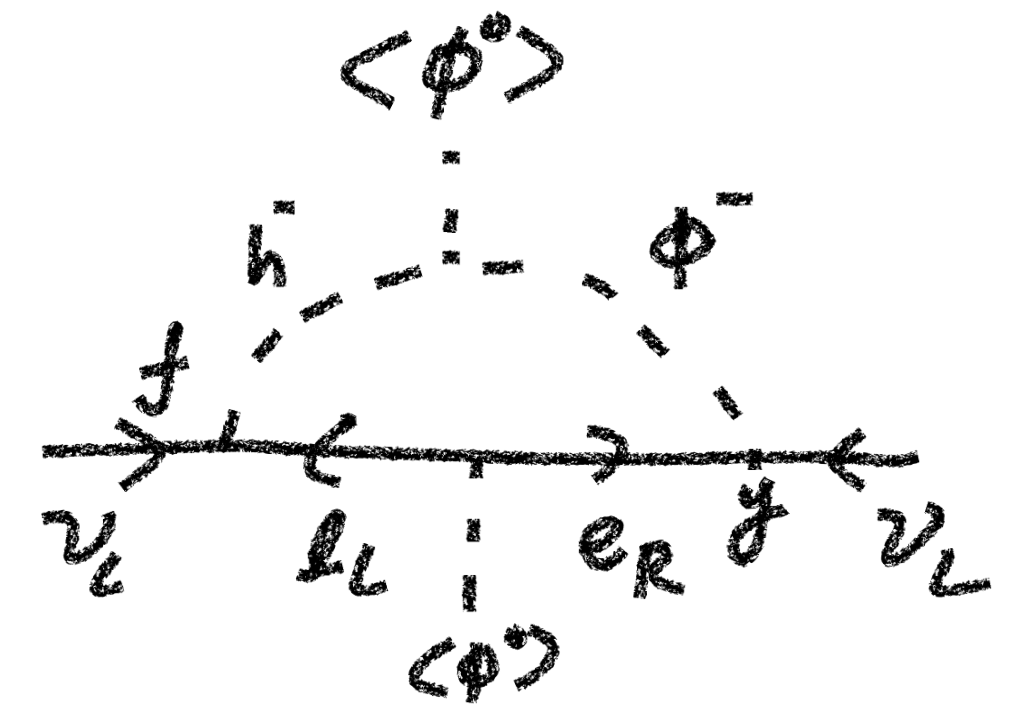
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$$\mathcal{L} \supset y_D \bar{L} H e_R + f \bar{L} L h^\dagger + h.c$$

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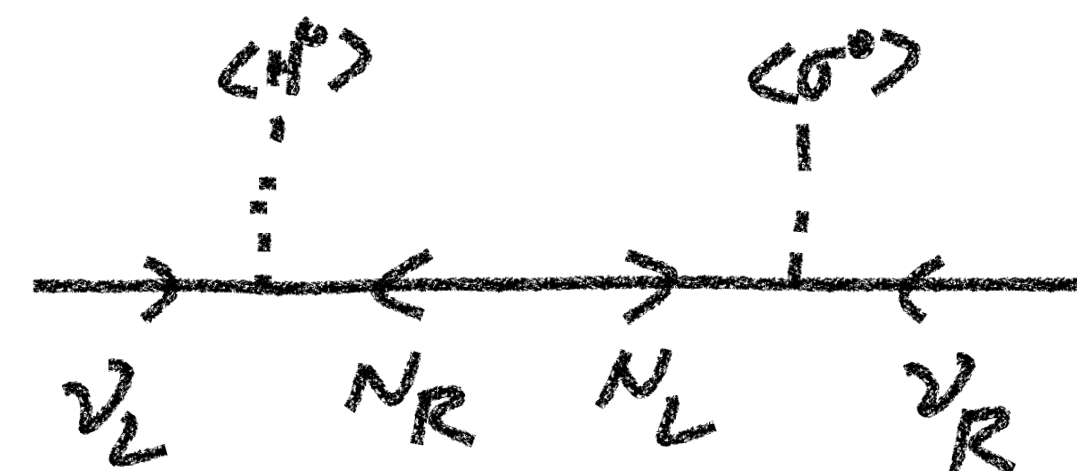


Mass Models (Dirac Scenario)

Type I:

$$\mathcal{L} = y_D \bar{L} H N_R + y_D \bar{N}_L \nu_R \sigma + M_D \bar{N}_L N_R$$

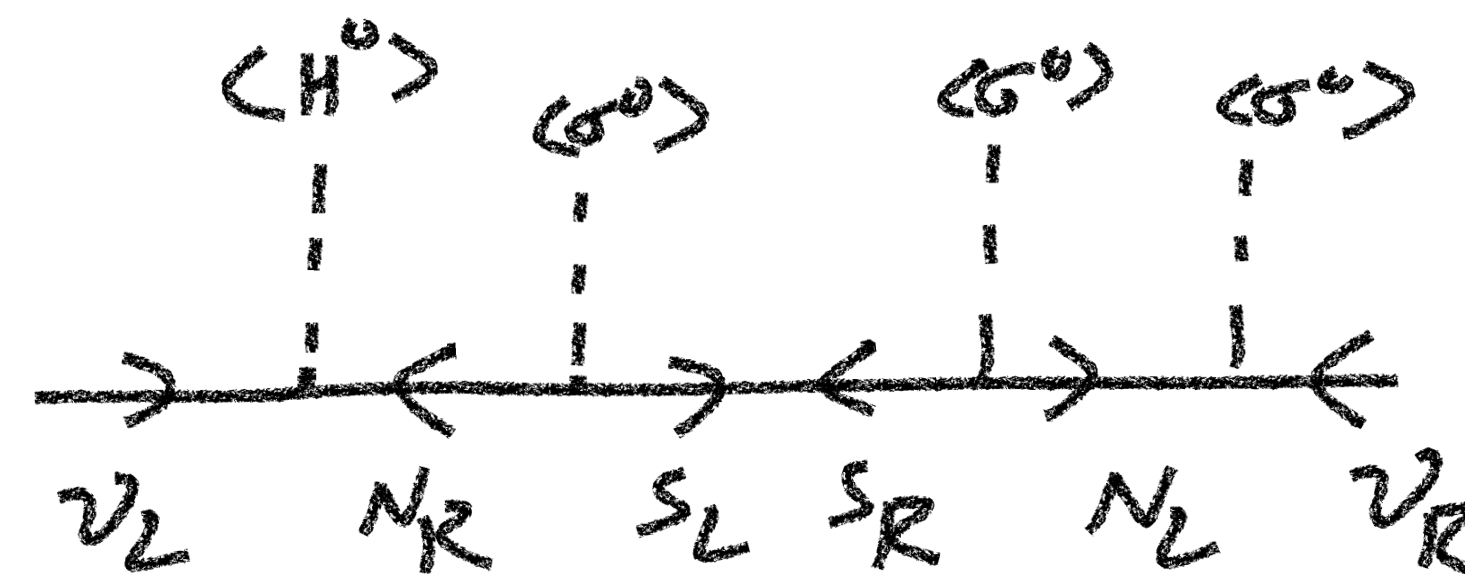
$$m_\nu = \begin{pmatrix} 0 & y_D v_H \\ y_D v_H & M_D \end{pmatrix}$$



Inverse See-Saw:

$$\mathcal{L} = y_D \bar{L} H N_R + y_S \bar{N}_L \nu_R \sigma + y_S \bar{N} S \sigma + M_D \bar{S} S + h.c$$

$$m_\nu = \begin{pmatrix} 0 & y_D v_H & 0 \\ y_D v_H & 0 & y_S v_\sigma \\ 0 & y_S v_\sigma & M_D \end{pmatrix}$$

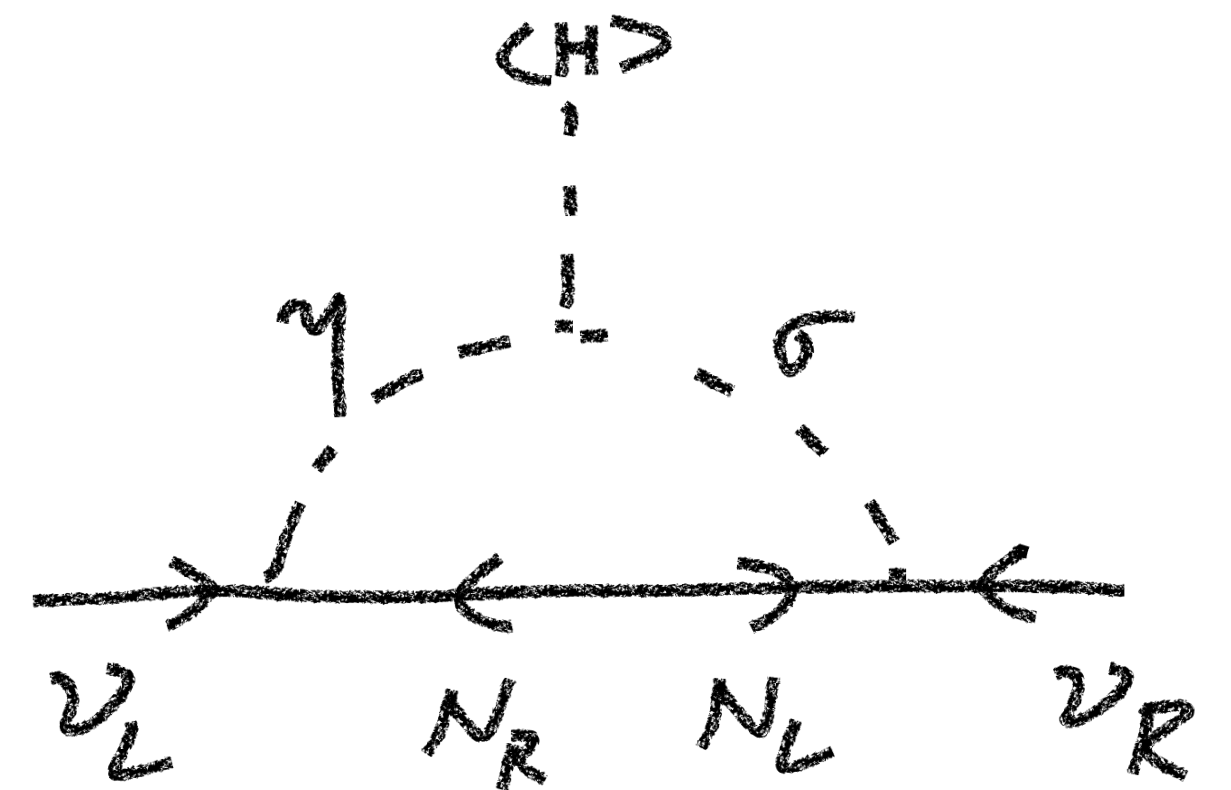


Mass Models (Dirac Scenario)

Radiative Mass Model:

$$d \supset y_D \bar{L} \gamma N_R + M \bar{N}_L N_R + y_\sigma \bar{N}_L \nu_R + h.c$$

$$V \supset \mu_\sigma \gamma^\dagger H \sigma + h.c$$

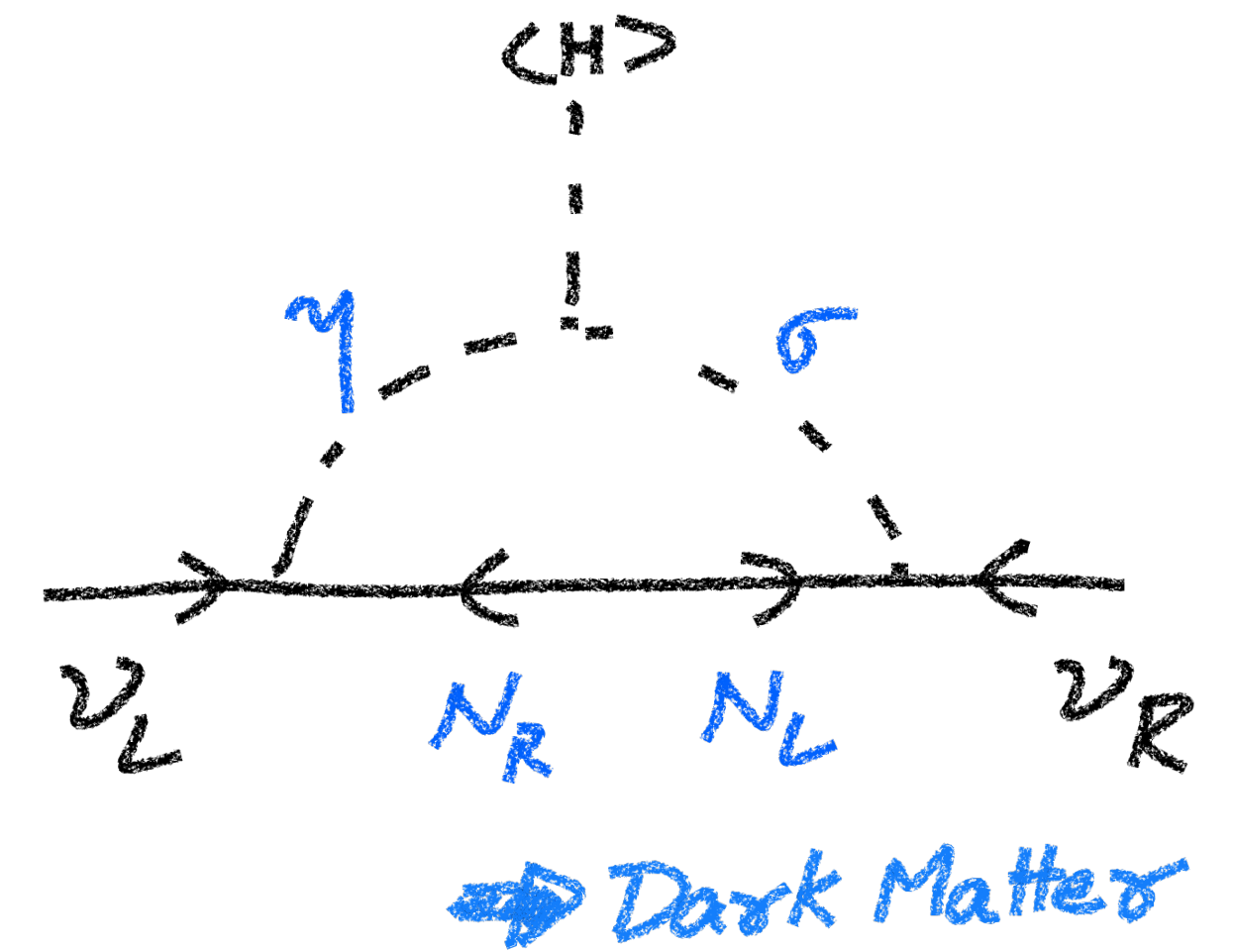


Mass Models (Dirac Scenario)

Radiative Mass Model:

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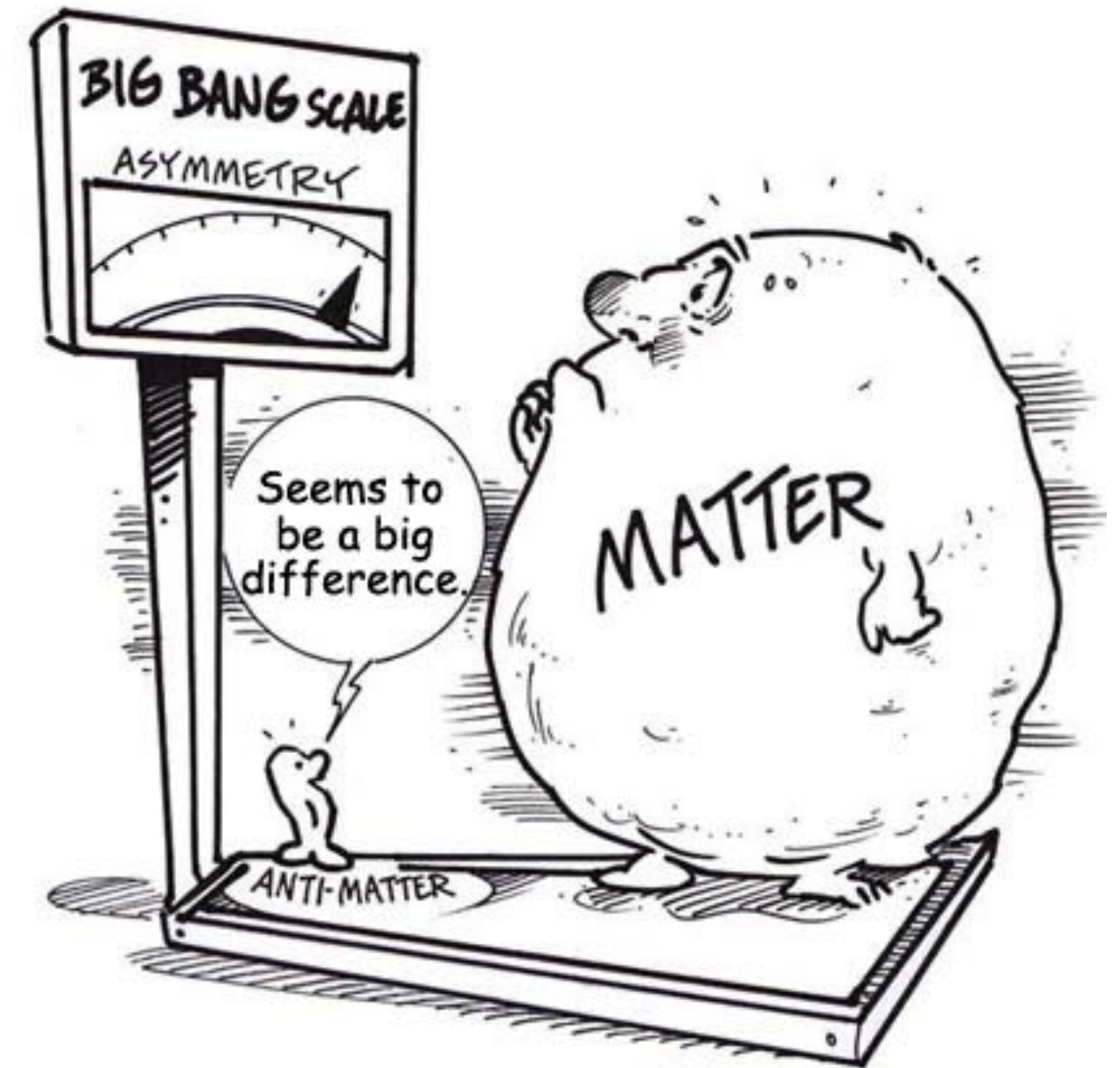
$$V \supset \mu_\sigma \gamma^\dagger H \sigma + \text{h.c.}$$



Matter Asymmetry

① The Standard Model does not explain the present asymmetry

1. The CP violation coming from Jarlskog invariant is of the order 10^{-20} .
2. The experimental lower bound on the Higgs mass implies the transition is not strongly first order.



Matter Asymmetry

- The observed BAU is often quoted in terms of baryon to photon ratio

$$\eta_B = \frac{n_B - \bar{n}_B}{n_\gamma} = 6.04 \pm 0.08 \times 10^{-10} \Rightarrow \eta_B^0 = \eta_B / 7.04 \Rightarrow \Omega_B h^2 = \frac{\rho_0}{\rho_{\text{crit}}} \eta_B^0 \sim 2.26\%$$

- The prediction for this ratio from Big-Bang Nucleosynthesis (BBN) agrees well with the observed value from Cosmic Microwave Background Radiation (CMBR) measurements.

Sakharov's Conditions

• The three basic ingredients necessary to generate a net baryon asymmetry from an initially baryons symmetric Universe (1967)

1. Baryon Number (B) violation $X \rightarrow Y + B$

2. C and CP violation

$$\Gamma(X \rightarrow Y + B) \neq \Gamma(\bar{X} \rightarrow \bar{Y} + \bar{B})$$

3. Departure from thermal Equilibrium

Recipe for Asymmetry

- Now, in order to get a non-zero CP violation we start need at-least two distinct amplitude for a particular process.
- In order to understand the above claim we start with the amplitude of a B-violating process ($X \rightarrow b$)

$$i\mathcal{M} = (C_1 \mathcal{A}_1 + C_2 \mathcal{A}_2) f$$

Similarly, for anti-particle

$$i\bar{\mathcal{M}} = (C_1^* \mathcal{A}_1 + C_2^* \mathcal{A}_2) f^*$$

spinor wave function

Recipe for Asymmetry

- The difference between the two processes comes out to be

$$S = 4 \operatorname{Im}[c_1^* c_2] \operatorname{Im}[A_1^* A_2] |f|^2$$

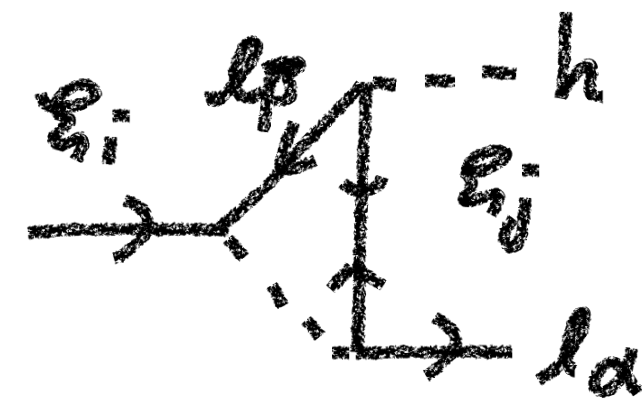
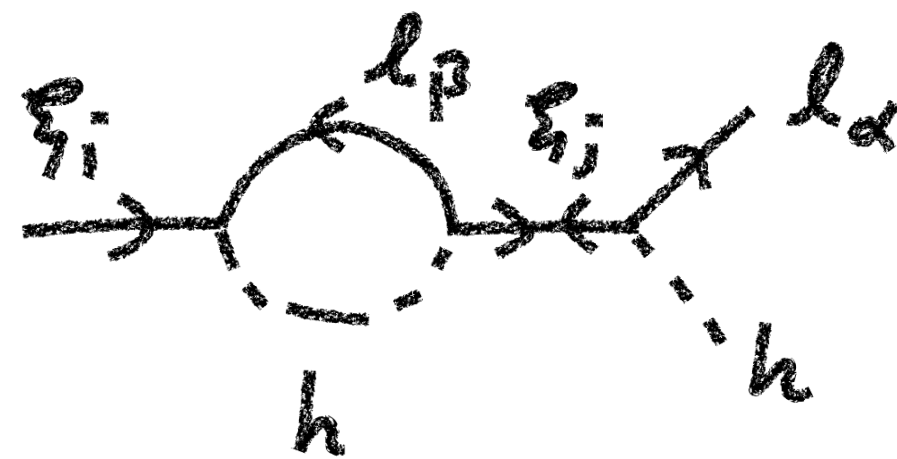
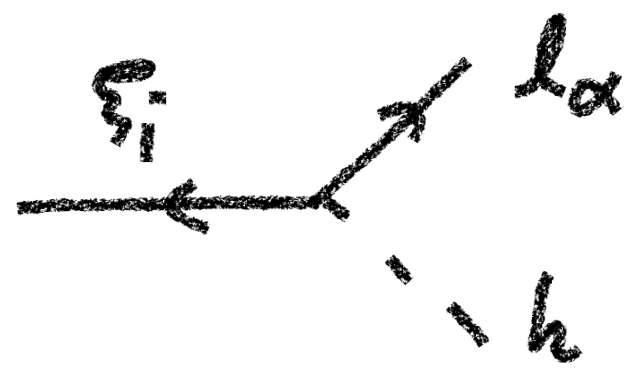
└─ Purely from coupling

└─ Purely from Amplitudes

Leptogenesis from ν -Mass model (Majorana case)

Lagrangian

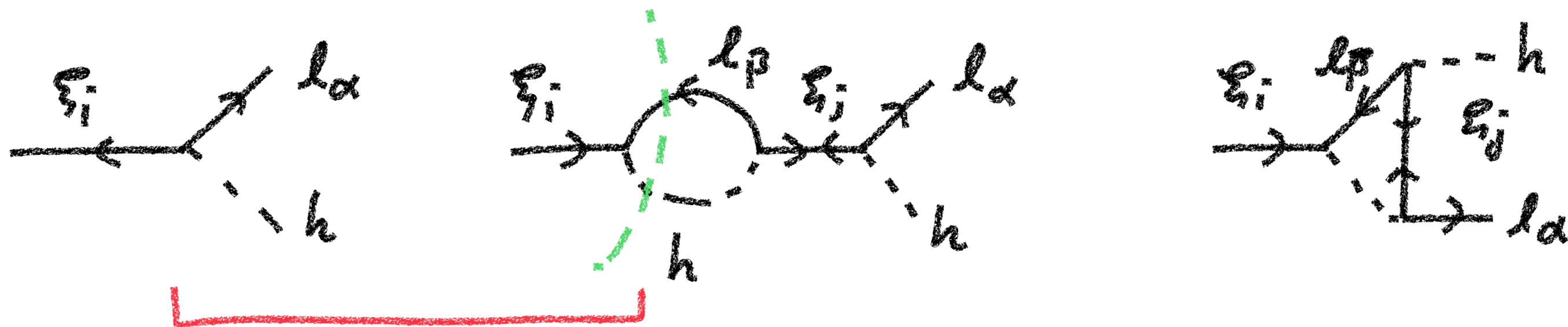
$$\mathcal{L} \supset y_{\alpha i} \bar{l}_{\alpha} h \xi_i + M_{ij} \xi_i \xi_j + h.c.$$



Leptogenesis from ν -Mass model (Majorana case)

• Lagrangian

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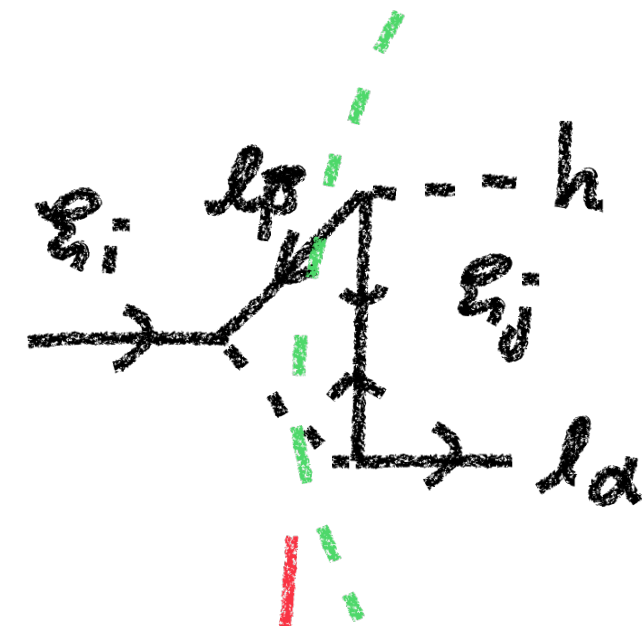
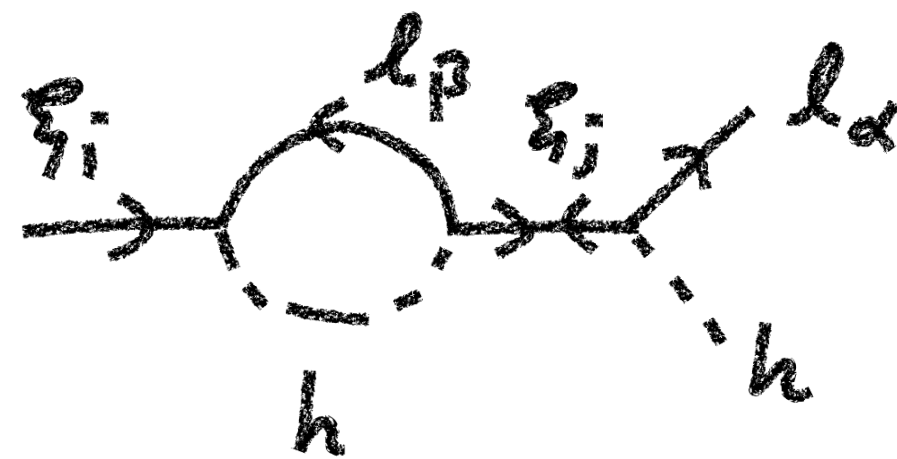
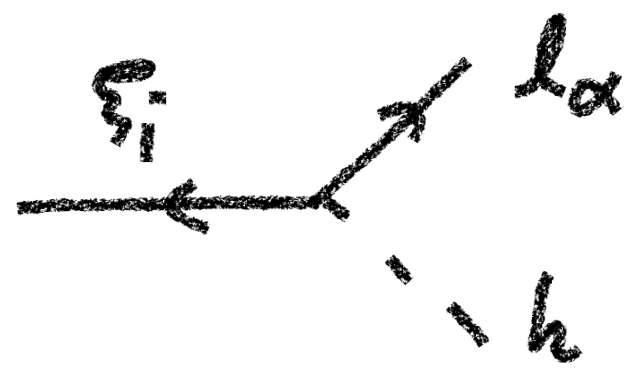


(Liu & Segre '93)

Leptogenesis from ν -Mass model (Majorana case)

• Lagrangian

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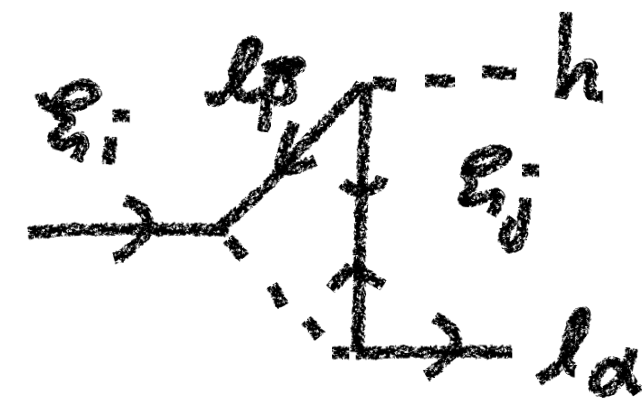
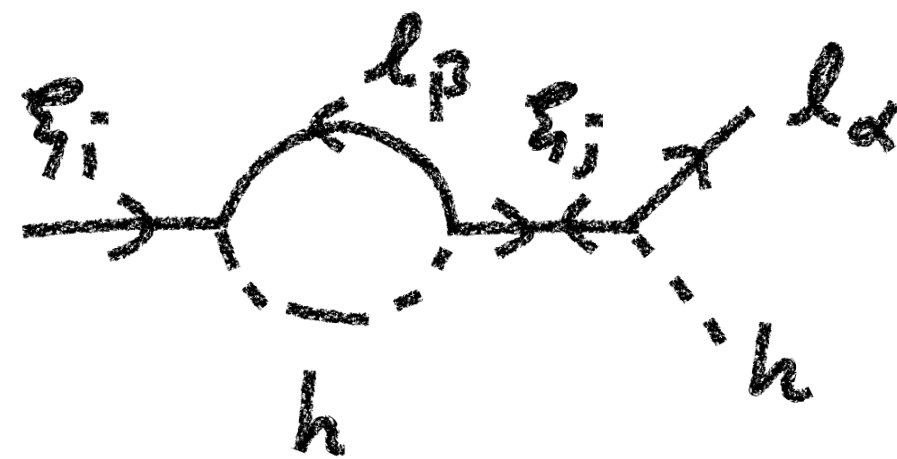
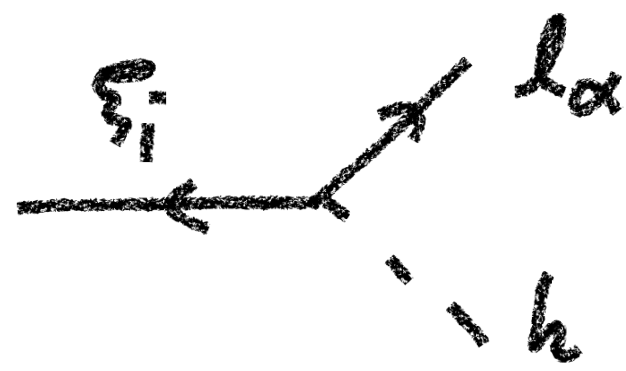


(Fukugida & Yanagida '86)

Leptogenesis from ν -Mass model (Majorana case)

Lagrangian

$$\mathcal{L} \supset y_{\alpha i} \bar{l}_\alpha h \xi_i + M_{ij} \xi_i \xi_j + h.c$$



$$y_{\alpha i} = \frac{1}{\sqrt{v}} \left(U D_{\sqrt{M_\nu}} R^\dagger D_{\sqrt{\Lambda}} \right)_{\alpha i},$$

$\hookrightarrow RR^\dagger = \mathbb{1}$

$$M_\nu = y^* \Lambda^{-1} y^\dagger$$

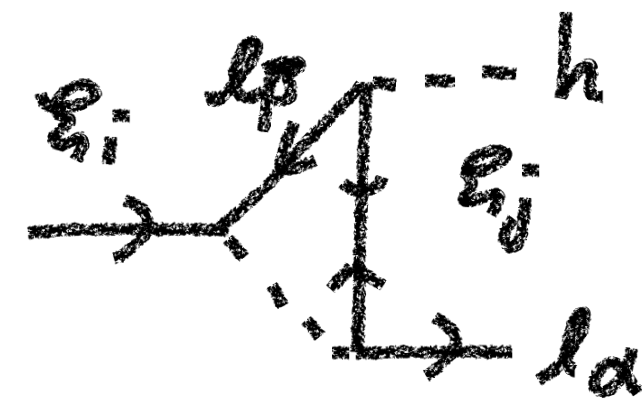
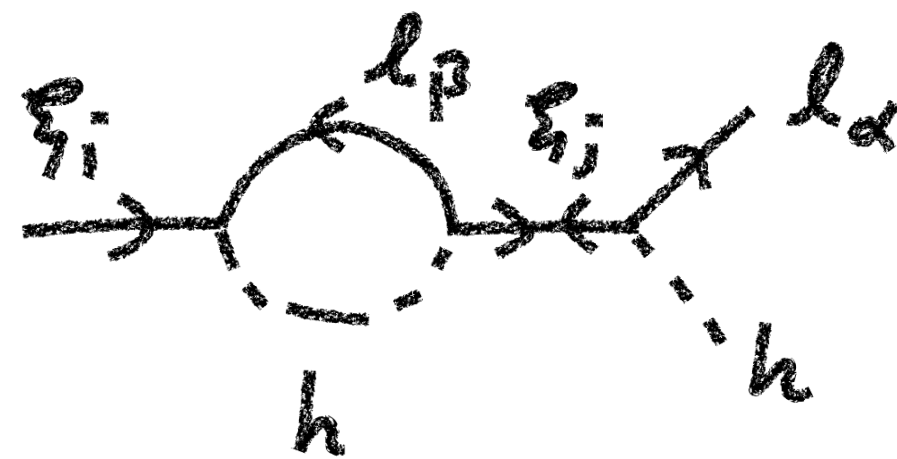
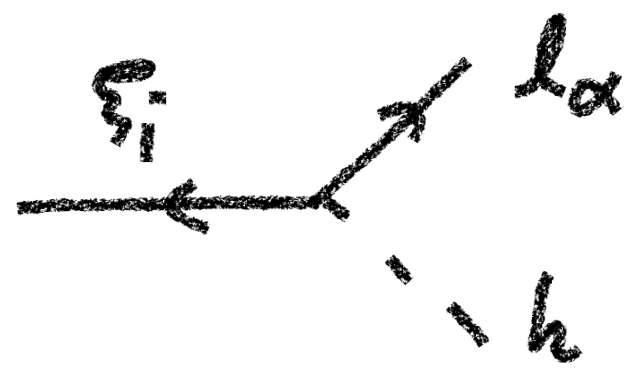
$\hookrightarrow \nu\text{-mass}$

$$D_{\sqrt{M_\nu}} = \begin{pmatrix} \sqrt{M_{\nu_1}} & 0 & 0 \\ 0 & \sqrt{M_{\nu_2}} & 0 \\ 0 & 0 & \sqrt{M_{\nu_3}} \end{pmatrix}; \quad D_{\sqrt{\Lambda}} = \begin{pmatrix} \sqrt{\Lambda_1} & 0 & 0 \\ 0 & \sqrt{\Lambda_2} & 0 \\ 0 & 0 & \sqrt{\Lambda_3} \end{pmatrix}$$

Leptogenesis from ν -Mass model (Majorana Case)

Lagrangian

$$\mathcal{L} \supset y_{\alpha i} \bar{l}_{\alpha} h \xi_i + M_{ij} \xi_i \xi_j + h.c$$



$$y_{\alpha i} = \frac{1}{\sqrt{v}} \left(U D_{\alpha} R^T D_{\alpha} \right)_{\alpha i},$$

$\hookrightarrow RR^T = \mathbb{1}$

$$M_{\nu} = y^* \Lambda^{-1} y^T$$

$\hookrightarrow \nu\text{-mass}$

$$D_{\sqrt{M_{\nu}}} = \begin{pmatrix} \sqrt{M_{\nu_1}} & 0 & 0 \\ 0 & \sqrt{M_{\nu_2}} & 0 \\ 0 & 0 & \sqrt{M_{\nu_3}} \end{pmatrix}; \quad D_{\sqrt{\Lambda}} = \begin{pmatrix} \sqrt{\Lambda_1} & 0 & 0 \\ 0 & \sqrt{\Lambda_2} & 0 \\ 0 & 0 & \sqrt{\Lambda_3} \end{pmatrix}$$

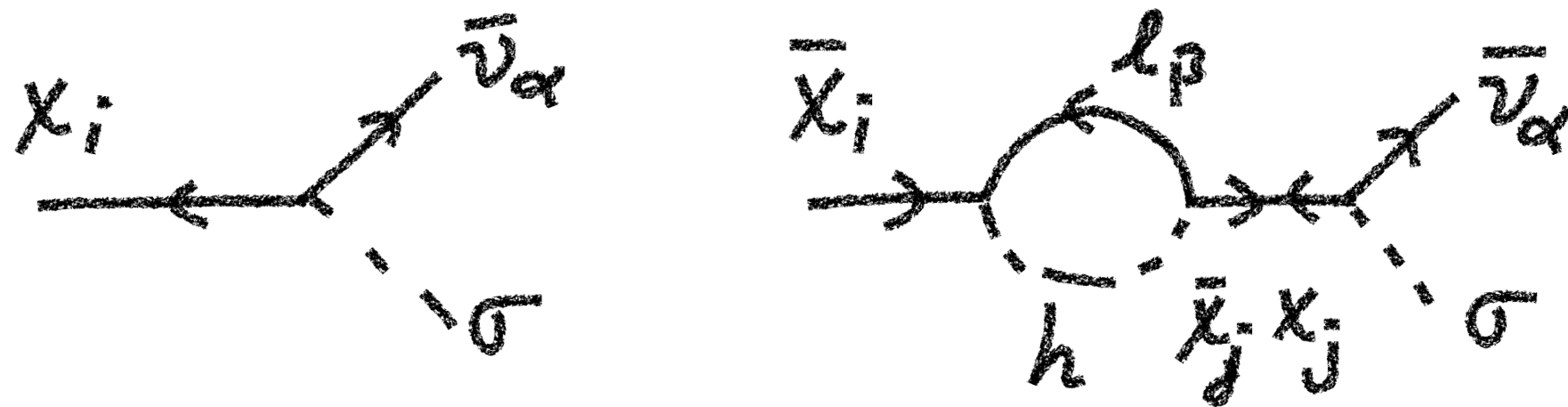
Casas-Ibarra Parameterization

Leptogenesis from ν -Mass model (Dirac case)

Lagrangian

$$\mathcal{L} \supset y_{\alpha i} \bar{l}_\alpha h \bar{\chi}_i + m_{ij} \bar{\chi}_i \chi_j + \tilde{y}_{i\alpha} \chi_i \sigma \bar{\nu}_\alpha + h.c$$

(Cerdeño, Dedes
& Underwood '06)



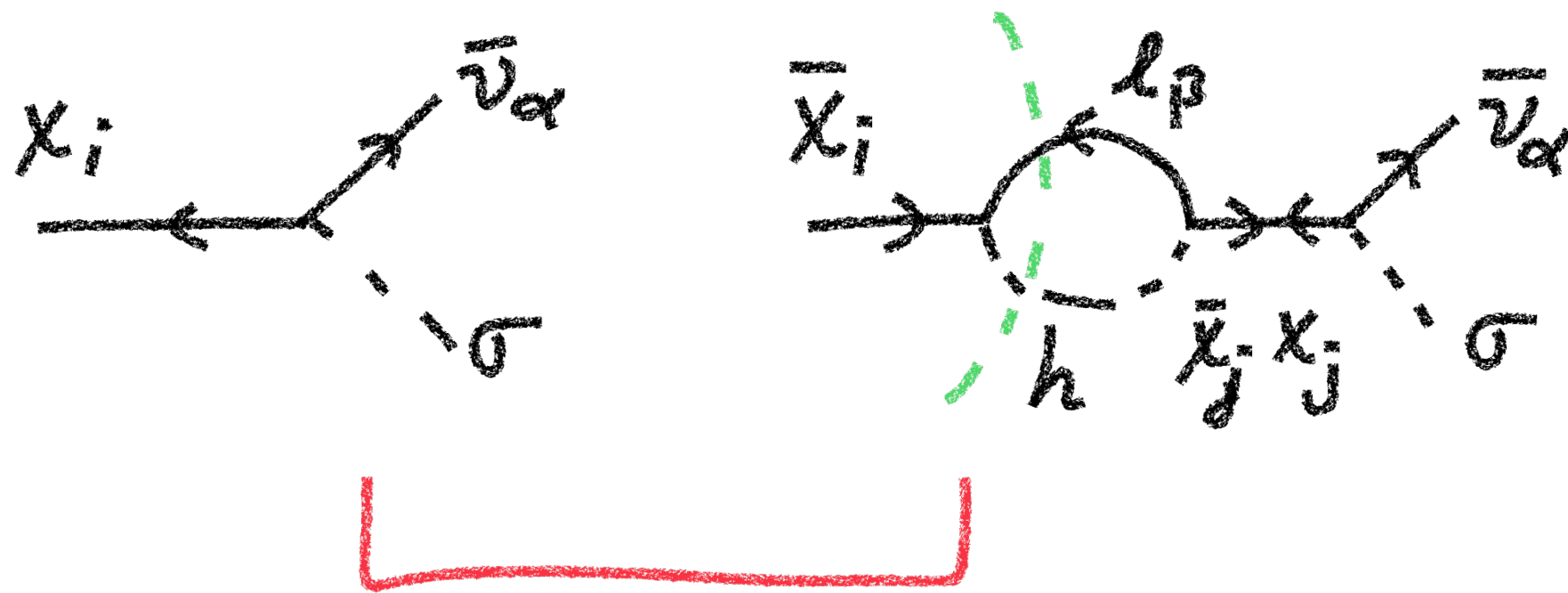
$$y_{\alpha i} = \frac{1}{v} \left(A D_{\sqrt{M_\nu}} W D_{\sqrt{M}} \right)_{\alpha i}; \quad \tilde{y}_{i\alpha} = \frac{1}{v_\sigma} \left(D_{\sqrt{M}} X^\dagger D_{\sqrt{M_\nu}} B^\dagger \right)_{i\alpha} \Rightarrow W X^\dagger = 1$$

$$m_{\nu_{\alpha\beta}} = v v_\sigma \left(y M^{-1} \tilde{y} \right)_{\alpha\beta}$$

Leptogenesis from ν -Mass model (Dirac case)

Lagrangian

$$\mathcal{L} \supset y_{\alpha i} \bar{l}_\alpha h \bar{\chi}_i + m_{ij} \bar{\chi}_i \chi_j + \tilde{y}_{i\alpha} \chi_i \sigma \bar{\nu}_\alpha + h.c.$$



Dirac Leptogenesis

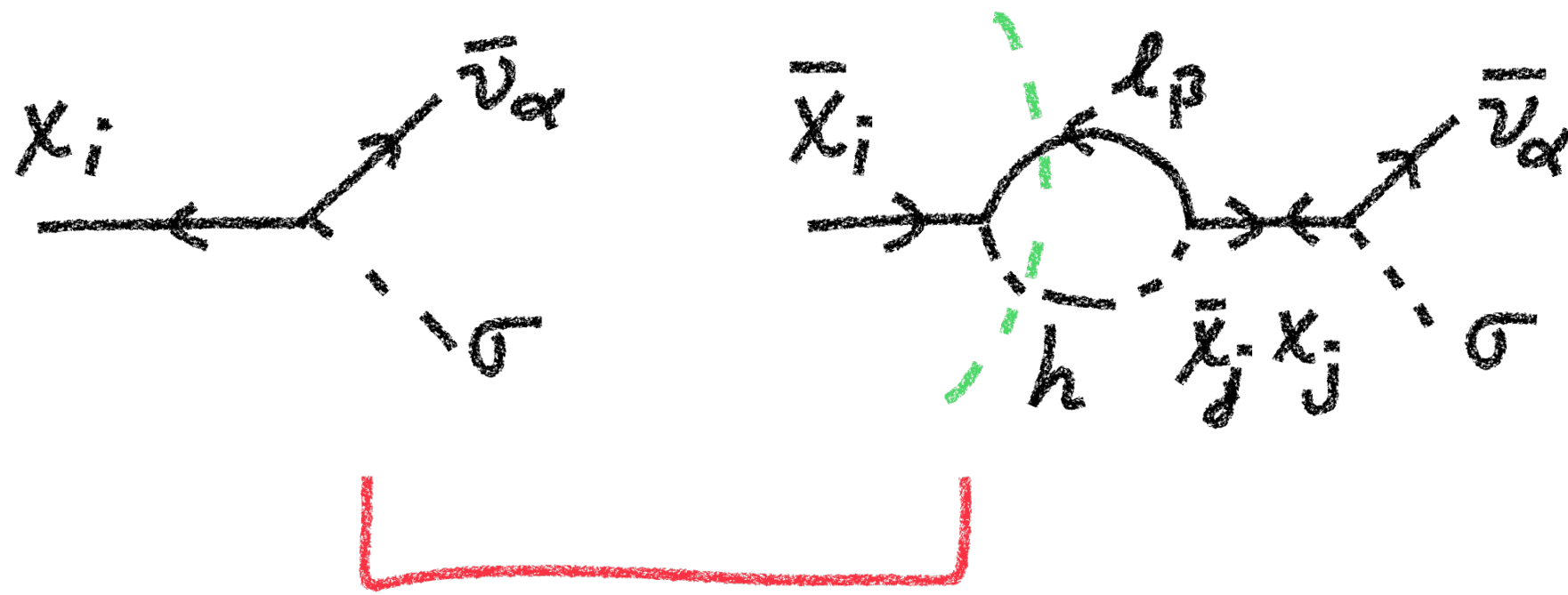
$$Y_B = \frac{28}{79} (Y_B - Y_{LSM}) \Rightarrow Y_B = -\frac{28}{79} Y_{LVR}$$

$$Y_B - Y_{LSM} - Y_{LVR} = 0$$

Leptogenesis from ν -Mass model (Dirac case)

Lagrangian

$$\mathcal{L} \supset y_{\alpha i} \bar{l}_\alpha h \bar{\chi}_i + m_{ij} \bar{\chi}_i \chi_j + \tilde{y}_{i\alpha} \chi_i \sigma \bar{\nu}_\alpha + h.c.$$



Dirac Leptogenesis

$$Y_B = \frac{28}{79} (Y_B - Y_{LSM}) \Rightarrow Y_B = -\frac{28}{79} Y_{LVR}$$

$$Y_B - Y_{LSM} - Y_{LVR} = 0$$

Avoid left-right equilibration

$$LH \leftrightarrow \sigma \nu_R \quad ; \quad \Gamma_{L \leftrightarrow R} \sim \frac{|y|^2 |\tilde{y}|^2}{m_1} T^5$$

$$H(T) = \sqrt{\frac{8\pi^3 g_*}{90}} \frac{T^2}{M_{Pl}}$$

$$\left. \frac{\Gamma_{L \leftrightarrow R}}{H} \right|_{T=m_1} < 1 \Rightarrow \frac{|y|^2 |\tilde{y}|^2}{m_1} \lesssim \frac{1}{M_{Pl}} \sqrt{\frac{8\pi^3 g_*}{90}}$$

Correlation & bounds from Flavor Physics

(ν -mass & mixing)

(Baryon Asymmetry)



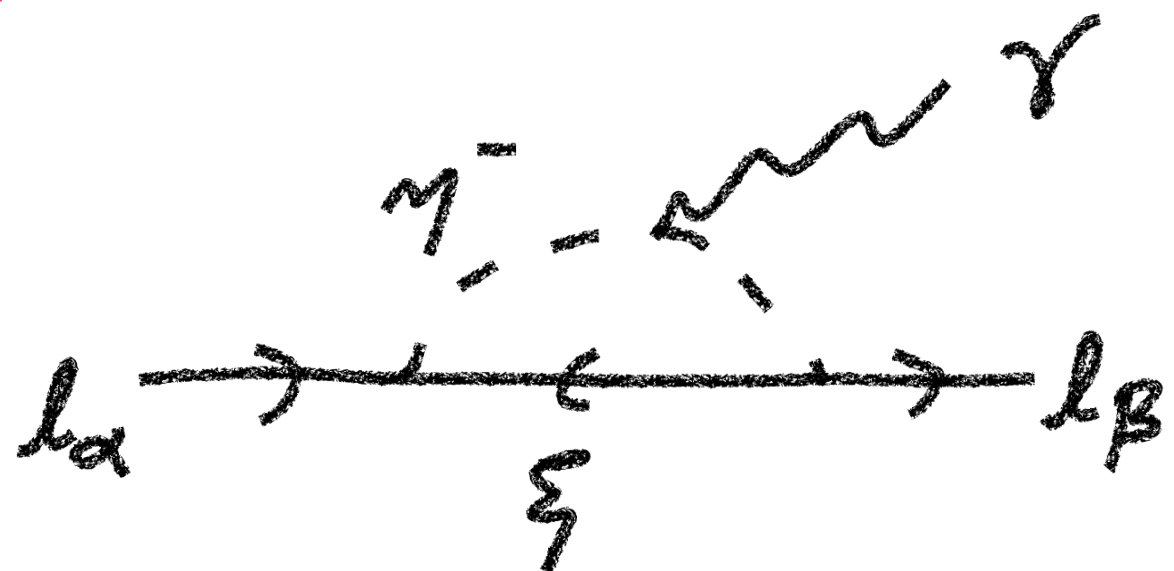
$$\rightarrow y = \frac{1}{v} \left(U D R^\dagger D \sqrt{M} \right) \quad (\text{Majorana case})$$

$$\rightarrow y = \frac{1}{v} \left(A D W D \sqrt{M} \right) \quad (\text{Dirac case})$$

$$\tilde{y} = \frac{1}{v} \left(D \sqrt{M} X^\dagger D \sqrt{M} B^\dagger \right)$$

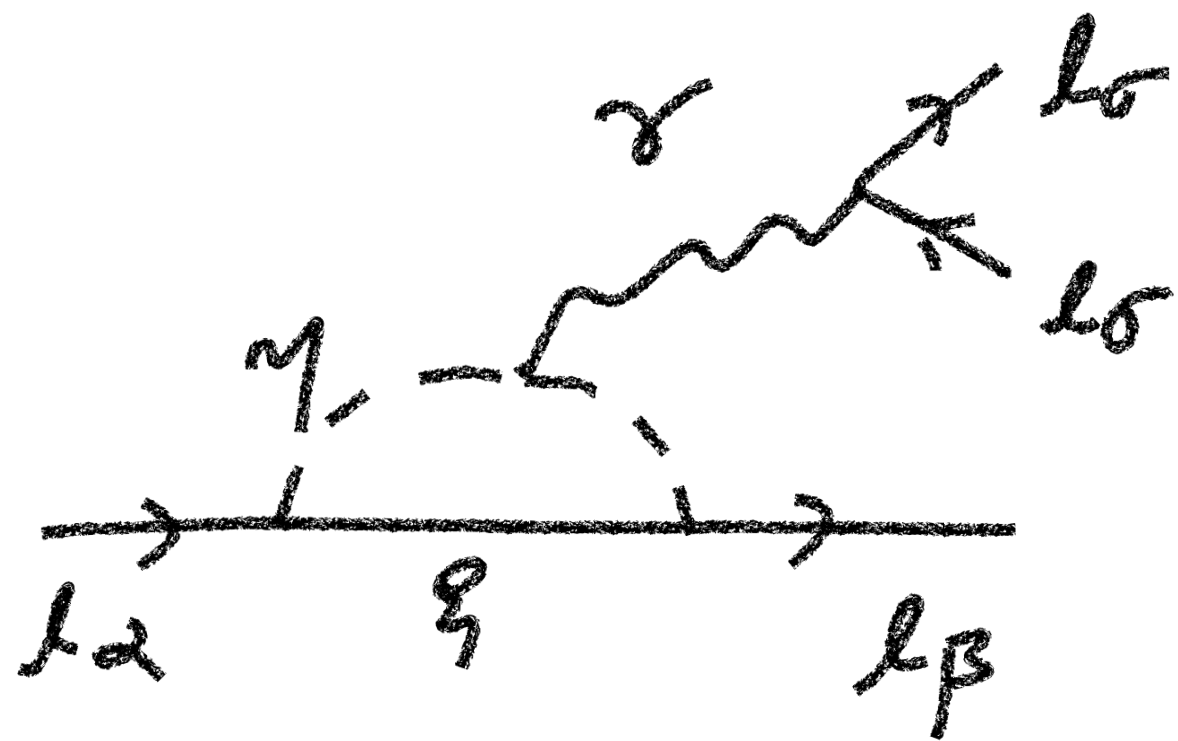
Correlation & bounds from Flavor Physics

● $l_\alpha \rightarrow l_\beta \gamma$



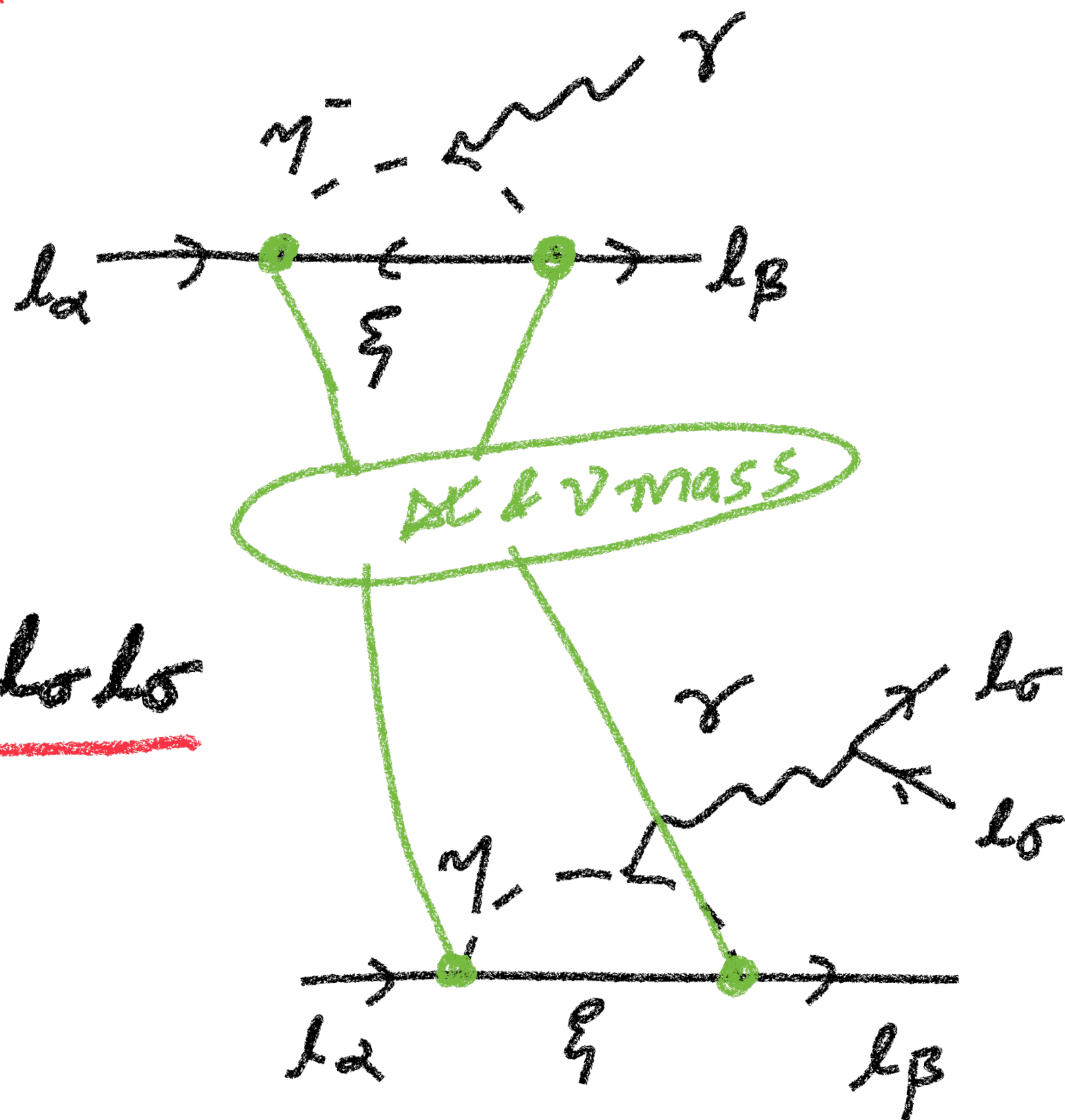
(Scotogenic & it's variant)

● $l_\alpha \rightarrow l_\beta l_\sigma l_\sigma$



Correlation & bounds from Flavor Physics

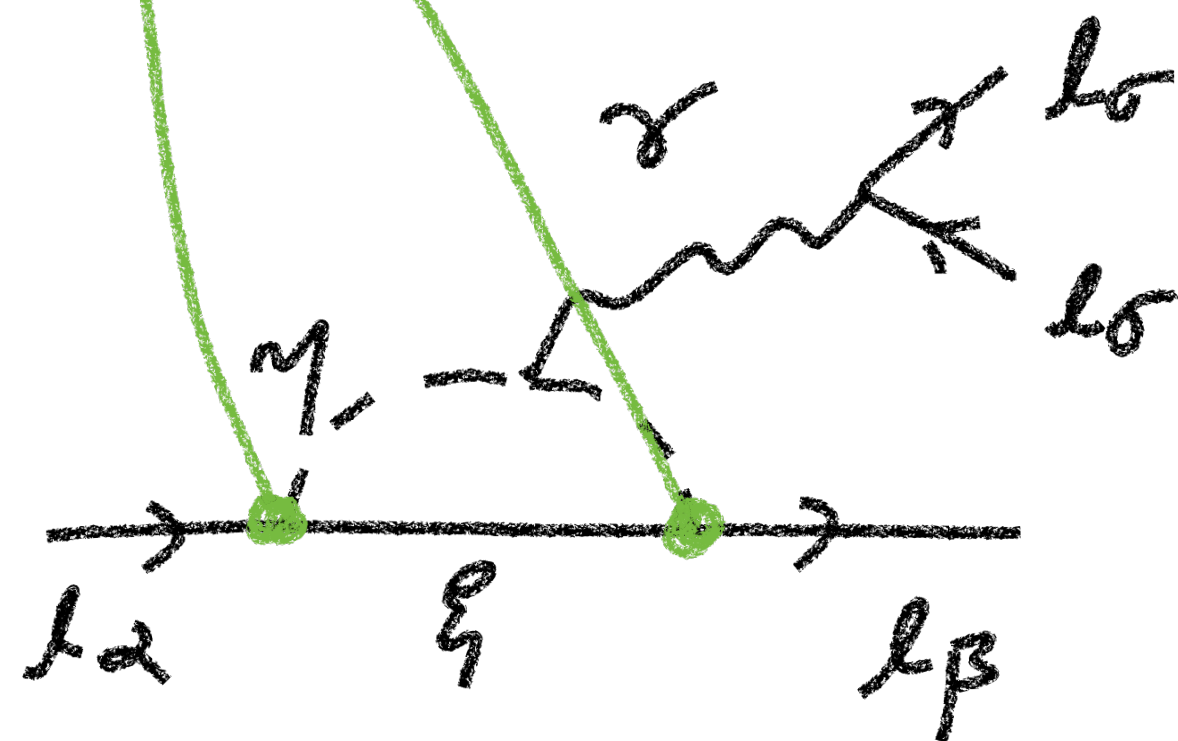
● $l_\alpha \rightarrow l_\beta \gamma$



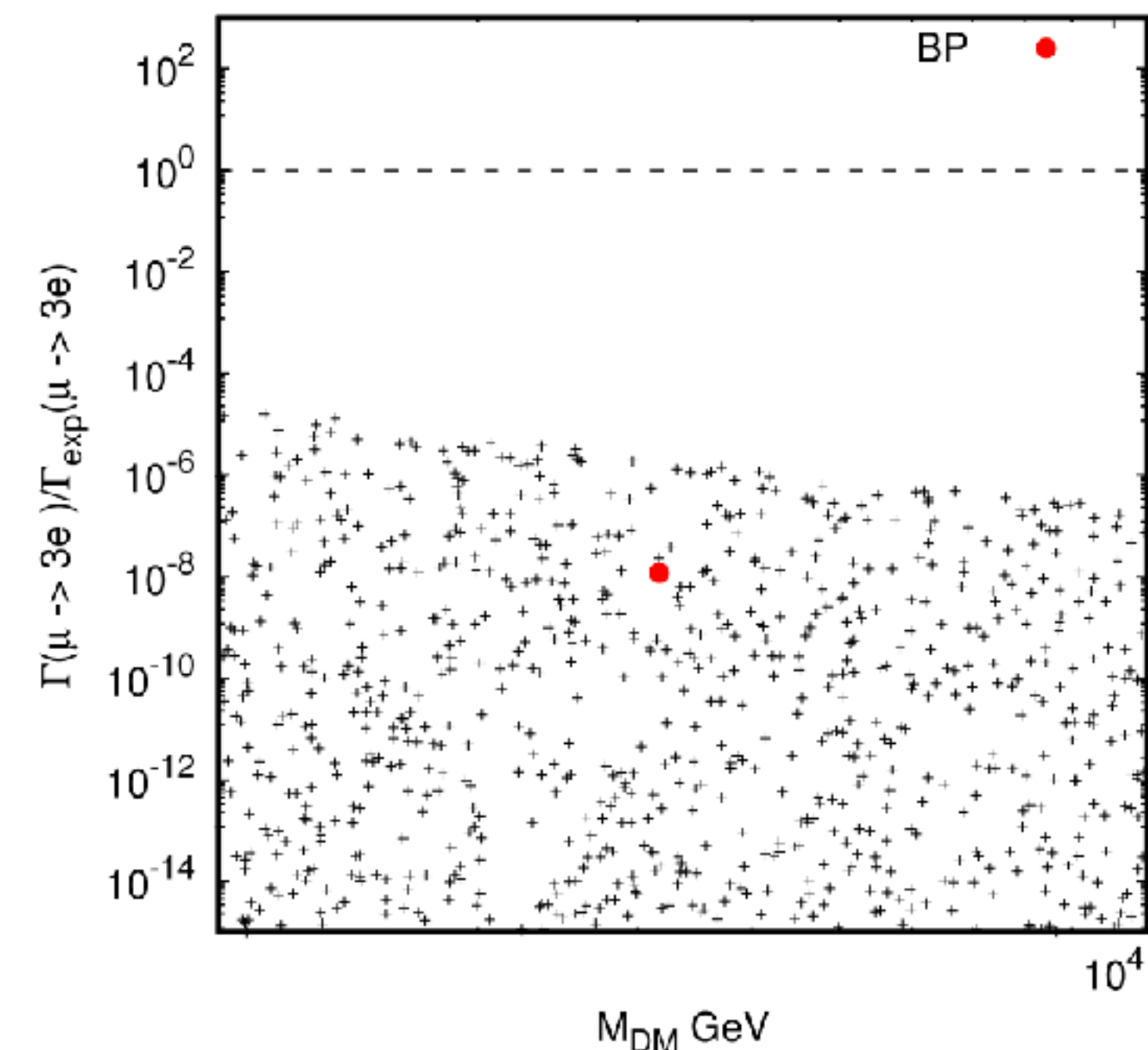
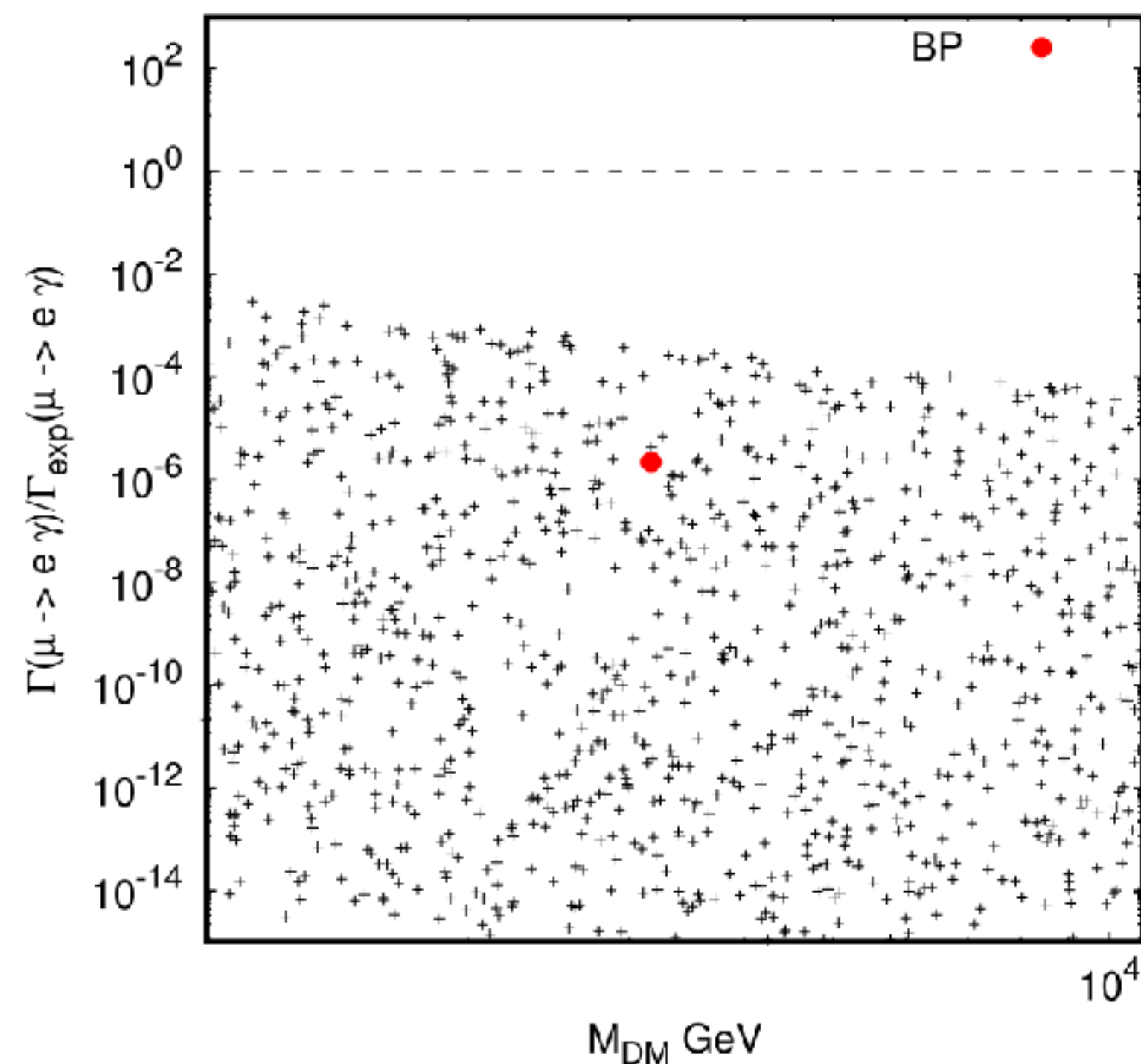
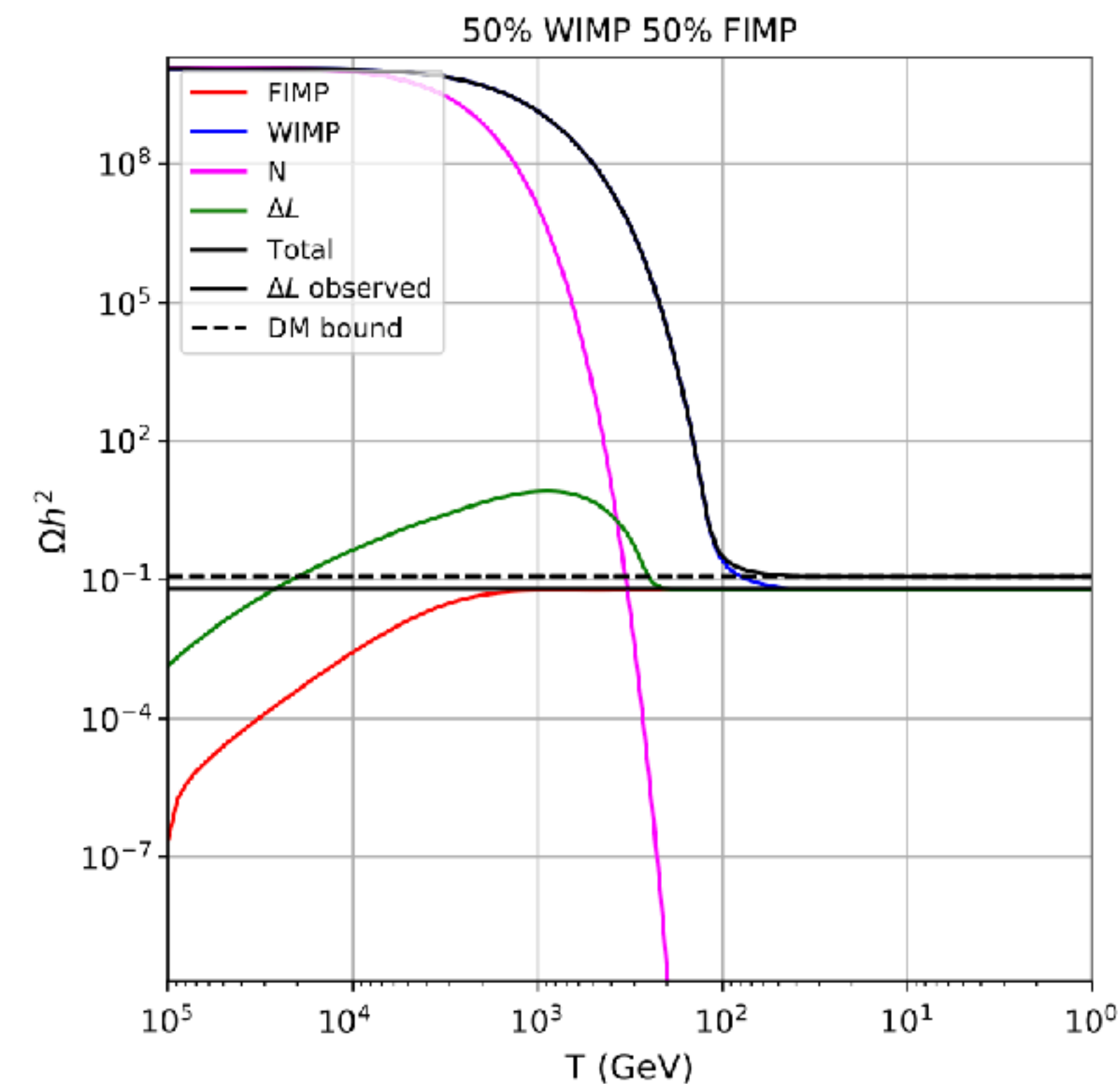
(Scotogenic & it's variant)

An example

● $l_\alpha \rightarrow l_\beta l_\sigma l_\sigma$



Correlation & bounds from Flavor Physics



[arXiv:1903.10516 D. Borah, A. Dasgupta, S.K. Kang]

Thank you !!

