Cosmology of Dirac Neutrino Mass Models

Michael Shamma (He/They) mshamma@triumf.ca





- What we don't know about neutrinos
- It's really tough to produce and detect Dirac ν
- Example Dirac u models which contribute detectably to $\Delta N_{
 m eff}$
- Terrestrial and cosmological ν probes
 - Correlating terrestrial and cosmological ν measurements
- Neutrinos in the bulk

Outline



Michael Shamma (He/They) mshamma@triumf.ca



- EW symmetry permits introduction of "sterile" neutrino ν_R
- Dirac and Majorana masses permitted: $\mathscr{L}_{\nu} \supset -y_{\nu} \bar{L} \tilde{H} \nu_R - i M \bar{\nu}_R^c \nu_R + h \cdot c \,.$



- EW symmetry permits introduction of "sterile" neutrino ν_R
- Dirac and Majorana masses permitted: $\mathscr{L}_{\nu} \supset -y_{\nu} \bar{L} \tilde{H} \nu_R - i M \bar{\nu}_R^c \nu_R + h \cdot c \,.$
 - Majorana Neutrino Hypothesis (MNH) easier to probe: $\Delta L = 2$ permits $0\nu\beta\beta$









- EW symmetry permits introduction of "sterile" neutrino ν_R
- Dirac and Majorana masses permitted: $\mathscr{L}_{\nu} \supset -y_{\nu} \overline{L} \overline{H} \nu_R - i M \overline{\nu}_R^c \nu_R + h.c.$
 - Majorana Neutrino Hypothesis (MNH) easier to probe: $\Delta L = 2$ permits $0\nu\beta\beta$
 - Dirac Neutrino Hypothesis (DNH) more difficult: ν_R interactions with SM are suppressed,

 $y_{\nu} = m_{\nu}/v_{ew} = \mathcal{O}(10^{-12})(m_{\nu}/0.1 \text{ eV})$

Michael Shamma (He/They) mshamma@triumf.ca









- EW symmetry permits introduction of "sterile" neutrino ν_R
- Dirac and Majorana masses permitted: $\mathscr{L}_{\nu} \supset -y_{\nu} \overline{L} \overline{H} \nu_R - i M \overline{\nu}_R^c \nu_R + h.c.$
 - Majorana Neutrino Hypothesis (MNH) easier to probe: $\Delta L = 2$ permits $0\nu\beta\beta$
 - Dirac Neutrino Hypothesis (DNH) more difficult: ν_R interactions with SM are suppressed, $y_{\nu} = m_{\nu}/v_{ew} = \mathcal{O}(10^{-12})(m_{\nu}/0.1 \text{ eV})$

 \implies Negligible lab and cosmological production

Michael Shamma (He/They) mshamma@triumf.ca









- EW symmetry permits introduction of "sterile" neutrino ν_R
- Dirac and Majorana masses permitted: $\mathscr{L}_{\nu} \supset -y_{\nu} \bar{L} \bar{H} \nu_R - i M \bar{\nu}_R^c \nu_R + h.c.$
 - Majorana Neutrino Hypothesis (MNH) easier to probe: $\Delta L = 2$ permits $0\nu\beta\beta$
 - Dirac Neutrino Hypothesis (DNH) more difficult: ν_R interactions with SM are suppressed, $y_{\nu} = m_{\nu}/v_{ew} = \mathcal{O}(10^{-12})(m_{\nu}/0.1 \text{ eV})$

 \implies Negligible lab and cosmological production ... unless new physics accompanies $\nu_R!$

Michael Shamma (He/They) mshamma@triumf.ca









Michael Shamma (He/They) mshamma@triumf.ca





Simplest implementation of Dirac ν_R is $\mathcal{L} \supset -Y_{\nu}^{ij} \overline{L} \widetilde{H} \nu_R + h \cdot c$.

Michael Shamma (He/They) mshamma@triumf.ca





Simplest implementation of Dirac ν_R is $\mathcal{L} \supset -Y_{\nu}^{ij}\overline{L}\widetilde{H}\nu_R + h.c.$

Small ν_R population can freeze-in via e.g.

 $\nu_L \bar{\nu}_L \rightarrow \nu_R \bar{\nu}_R, \ e^- \nu_L \rightarrow e^- \nu_R, \ h \rightarrow \nu_L \nu_R$

Michael Shamma (He/They) mshamma@triumf.ca

Minimal SM Extension: $y_{1}LH\nu_{R}$





Simplest implementation of Dirac ν_R is $\mathcal{L} \supset -Y_{\nu}^{ij} \overline{L} \widetilde{H} \nu_R + h \cdot c$. Small ν_R population can freeze-in via e.g. $\nu_L \bar{\nu}_L \rightarrow \nu_R \bar{\nu}_R, \ e^- \nu_L \rightarrow e^- \nu_R, \ h \rightarrow \nu_L \nu_R$ After EWSB, ν_L , ν_R combine $\implies m_{\nu} \simeq y_{\nu} v_{ew}$

Minimal SM Extension: $y_{\mu}LH\nu_{R}$





Simplest implementation of Dirac ν_R is $\mathcal{L} \supset -Y_{\nu}^{ij} \overline{L} \overline{H} \nu_R + h \cdot c$. Small ν_R population can freeze-in via e.g. $\nu_L \bar{\nu}_L \rightarrow \nu_R \bar{\nu}_R, \ e^- \nu_L \rightarrow e^- \nu_R, \ h \rightarrow \nu_L \nu_R$ After EWSB, ν_I , ν_R combine $\implies m_{\nu} \simeq y_{\nu} v_{ew}$ Taking $m_{\nu} \simeq 0.1 \text{ eV} \implies y_{\nu} \simeq 10^{-12}$

Minimal SM Extension: $y_{1}LH\nu_{R}$





Simplest implementation of Dirac ν_R is $\mathcal{L} \supset -Y_{\nu}^{ij}L\tilde{H}\nu_R + h.c.$ Small ν_R population can freeze-in via e.g. $\nu_L \bar{\nu}_L \rightarrow \nu_R \bar{\nu}_R, \ e^- \nu_L \rightarrow e^- \nu_R, \ h \rightarrow \nu_L \nu_R$ After EWSB, ν_L , ν_R combine $\implies m_{\nu} \simeq y_{\nu} v_{ew}$ Taking $m_{\nu} \simeq 0.1 \text{ eV} \implies y_{\nu} \simeq 10^{-12}$



[P. Adshead, Y. Cui, A. Long, **MS**, hep-ph/2009.07852 X. Luo, W. Rodejohann, X. Xu hep-ph/2011.13059]

Michael Shamma (He/They) mshamma@triumf.ca

Minimal SM Extension: $y_{\mu}LH\nu_{R}$







Michael Shamma (He/They) mshamma@triumf.ca



- Dirac Leptogenesis [K. Dick, et. al. Phys. Rev. Lett., 84:4039–4042, 2000...]
 - Model dependent $\Delta N_{eff} \sim 0.05 0.14$

Michael Shamma (He/They) mshamma@triumf.ca







- Dirac Leptogenesis [K. Dick, et. al. Phys. Rev. Lett., 84:4039–4042, 2000...] - Model dependent $\Delta N_{eff} \sim 0.05 - 0.14$
- Gauged $U(1)_{B-L}$ [V. Barger, et. al. Phys. Rev. D, 67:075009, 2003...]
 - $T_{\text{dec}} \lesssim (m_{Z'}/g'M_{\text{pl}})^{4/3}M_{\text{pl}} \implies \Delta N_{\text{eff}} \simeq 0.13 0.23$





- Dirac Leptogenesis [K. Dick, et. al. Phys. Rev. Lett., 84:4039–4042, 2000...] - Model dependent $\Delta N_{eff} \sim 0.05 - 0.14$
- Gauged $U(1)_{B-L}$ [V. Barger, et. al. Phys. Rev. D, 67:075009, 2003...] $-T_{\rm dec} \lesssim (m_{Z'}/g'M_{\rm pl})^{4/3}M_{\rm pl} \implies \Delta N_{\rm eff} \simeq 0.13 - 0.23$
- Neutrino mass models have a home in an extra dimension [N. Arkani-Hamed, et. al. hep-ph/9811448, K.R. Dienes, et. al. hep-ph/9811428].
 - Compactification generates Dirac ν mass

 $\implies \mathscr{L} \supset -\frac{\lambda v}{\sqrt{2\pi R_{\text{ED}}M_*}} \bar{\nu}_L \nu_R^{(0)}$

- ...also generates mixing of the ν_k modes with ν_L
- And substantial number of relics, strong constraints [K. Abazajian, G. Fuller, M. Patel hep-ph/0011048]

Michael Shamma (He/They) mshamma@triumf.ca





Michael Shamma (He/They) mshamma@triumf.ca



Parametrize absolute neutrino mass scale with

 $m_{\nu_e} = \left[\sum_{i} m_i^2 |U_{ei}|^2\right]^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]

Michael Shamma (He/They) mshamma@triumf.ca





- Parametrize absolute neutrino mass scale with
- $m_{\nu_e} = \left[\sum_{i} m_i^2 |U_{ei}|^2\right]^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]
- $E_{\bar{\nu}_e} = \mathcal{O}(m_{\nu}) \implies \text{downward shift in } E_e$





- Parametrize absolute neutrino mass scale with
- $m_{\nu_e} = \left[\sum_{i} m_i^2 |U_{ei}|^2\right]^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]
- $E_{\bar{\nu}_e} = \mathcal{O}(m_{\nu}) \implies \text{downward shift in } E_e$
- Larger $m_{\nu_e} \implies$ larger shift





- Parametrize absolute neutrino mass scale with
- $m_{\nu_e} = \left| \sum_{i} m_i^2 |U_{ei}|^2 \right|^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]
- $E_{\bar{\nu}_e} = \mathcal{O}(m_{\nu}) \implies \text{downward shift in } E_e$
- Larger $m_{\nu_o} \implies$ larger shift
- Δm_{ii}^2 , θ_{ij} , mass hierarchy, and $m_{\nu_e} \implies \Sigma m_{\nu}$





- Parametrize absolute neutrino mass scale with
- $m_{\nu_e} = \sum_{i} m_i^2 |U_{ei}|^2 |^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]
- $E_{\bar{\nu}_e} = \mathcal{O}(m_{\nu}) \implies \text{downward shift in } E_e$
- Larger $m_{\nu_o} \implies$ larger shift
- Δm_{ii}^2 , θ_{ij} , mass hierarchy, and $m_{\nu_o} \implies \Sigma m_{\nu}$
- KATRIN gives best (current) limit: $m_{\nu_o} < 0.8 \text{ eV}$ $\implies \Sigma m_{\nu} \simeq 3 \text{ eV}$ [J.Angrik, et. al. 2004.22005]

Michael Shamma (He/They) mshamma@triumf.ca





- Parametrize absolute neutrino mass scale with
- $m_{\nu_e} = \sum_{i} m_i^2 |U_{ei}|^2 |^{1/2}$ [M. Tanabashi, Phys. Rev. D, 98(3):030001, 2018.]
- $E_{\bar{\nu}_e} = \mathcal{O}(m_{\nu}) \implies \text{downward shift in } E_e$
- Larger $m_{\nu_o} \implies$ larger shift
- Δm_{ii}^2 , θ_{ij} , mass hierarchy, and $m_{\nu_o} \implies \Sigma m_{\nu}$
- KATRIN gives best (current) limit: $m_{\nu_o} < 0.8 \text{ eV}$ $\implies \Sigma m_{\nu} \simeq 3 \text{ eV}$ [J.Angrik, et. al. 2004.22005]
- Project 8 projected sensitivity: $m_{\nu_o} < 0.04 \text{ eV}$ $\implies \Sigma m_{\nu} \simeq 0.14 \text{ eV}$ (NO) and $\Sigma m_{\nu} \simeq 0.099 \text{ eV}$ (IO) [A.A. Esfahani et al. J. Phys. G, 44(5):054004]

Michael Shamma (He/They) mshamma@triumf.ca





Michael Shamma (He/They) mshamma@triumf.ca



• ν_L relativistic during BBN and contribute to radiation energy density

Michael Shamma (He/They) mshamma@triumf.ca



- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\text{eff}} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$

Michael Shamma (He/They) mshamma@triumf.ca



- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\text{eff}} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$
- If ν_R present, contribute to $N_{eff} > N_{eff}^{SM} \simeq 3.044$



- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\text{eff}} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$
- If ν_R present, contribute to $N_{eff} > N_{eff}^{SM} \simeq 3.044$
- Increase in radiation energy density increases H at CMB



- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\rm eff} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$
- If ν_R present, contribute to $N_{eff} > N_{eff}^{SM} \simeq 3.044$
- Increase in radiation energy density increases H at CMB
- These free-streaming ν lead to phase-shift in CMB's acoustic peaks





- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\text{eff}} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$
- If ν_R present, contribute to $N_{eff} > N_{eff}^{SM} \simeq 3.044$
- Increase in radiation energy density increases H at CMB
- These free-streaming ν lead to phase-shift in CMB's acoustic peaks
- Planck restricts $\Delta N_{\text{eff}} \equiv N_{\text{eff}} N_{\text{eff}}^{\text{SM}} < 0.3$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]





- ν_L relativistic during BBN and contribute to radiation energy density
- Generally measured through $N_{\text{eff}} = (8/7)(11/4)^{4/3} \rho_{\nu} / \rho_{\gamma}$
- If ν_R present, contribute to $N_{eff} > N_{aff}^{SM} \simeq 3.044$
- Increase in radiation energy density increases H at CMB
- These free-streaming ν lead to phase-shift in CMB's acoustic peaks
- Planck restricts $\Delta N_{\text{eff}} \equiv N_{\text{eff}} N_{\text{off}}^{\text{SM}} < 0.3$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]
- CMB Stage-IV will constrain $\Delta N_{eff} < 0.06$

[K. Abazajian, et. al. astro-ph.IM/1907.04473]

Michael Shamma (He/They) mshamma@triumf.ca





Michael Shamma (He/They) mshamma@triumf.ca



• At late times, decoupled ν_L become non-relativistic

Michael Shamma (He/They) mshamma@triumf.ca



• At late times, decoupled ν_L become non-relativistic

- Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB

Michael Shamma (He/They) mshamma@triumf.ca


- At late times, decoupled ν_L become non-relativistic
 - Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB
 - Gravitational lensing and clustering measurements are sensitive to Σm_{ν}



- At late times, decoupled ν_L become non-relativistic
 - Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB
 - Gravitational lensing and clustering measurements are sensitive to $\Sigma m_{
 u}$
 - In 6-parameter ACDM extended with Σm_{ν} , Planck finds (model-dependent) $\Sigma m_{\nu} \lesssim 0.1 \text{ eV}$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]



- At late times, decoupled ν_L become non-relativistic
 - Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB
 - Gravitational lensing and clustering measurements are sensitive to Σm_{ν}
 - In 6-parameter ACDM extended with Σm_{μ} , Planck finds (model-dependent) $\Sigma m_{\mu} \leq 0.1 \text{ eV}$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]
 - In 12-parameter eACDM extended bound relaxes to $\sum m_{\nu} \leq 0.4 \text{ eV}$ [E. Di Valentino, A. Melichiorri, J. Silk, astro-ph.CO/1507.06646]



- At late times, decoupled ν_L become non-relativistic
 - Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB
 - Gravitational lensing and clustering measurements are sensitive to Σm_{ν}
 - In 6-parameter ACDM extended with Σm_{μ} , Planck finds (model-dependent) $\Sigma m_{\mu} \leq 0.1 \text{ eV}$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]
 - [E. Di Valentino, A. Melichiorri, J. Silk, astro-ph.CO/1507.06646]
- In 12-parameter eACDM extended bound relaxes to $\sum m_{\nu} \leq 0.4 \text{ eV}$ • Contribution of ν_R also constrained by Planck $\Omega_{\nu,\text{sterile}}h^2 \equiv m_{\nu,\text{sterile}}^{\text{eff}}/94 \text{ eV}$



- At late times, decoupled ν_L become non-relativistic
 - Free-streaming ν_L suppresses growth of structure \implies less lensing in CMB
 - Gravitational lensing and clustering measurements are sensitive to $\Sigma m_{
 u}$
 - In 6-parameter ACDM extended with Σm_{μ} , Planck finds (model-dependent) $\Sigma m_{\mu} \leq 0.1 \text{ eV}$ [N. Aghanim, et. al. astro-ph.CO/1807.06209]
 - [E. Di Valentino, A. Melichiorri, J. Silk, astro-ph.CO/1507.06646]
- In 12-parameter eACDM extended bound relaxes to $\sum m_{\mu} \leq 0.4 \text{ eV}$ • Contribution of ν_R also constrained by Planck $\Omega_{\nu,\text{sterile}}h^2 \equiv m_{\nu,\text{sterile}}^{\text{eff}}/94 \text{ eV}$ • Placnk+BAO gives joint bound of $m_{\nu,\text{sterile}}^{\text{eff}} < 0.23 \text{ eV}, N_{\text{eff}} < 3.34$



Can use conservation of entropy for thermalized ν_R

Michael Shamma (He/They) mshamma@triumf.ca





Can use conservation of entropy for thermalized ν_R

$$\Delta N_{\rm eff} = \frac{\rho}{\rho_{\nu}} = \frac{g}{(7/8)2} \left(\frac{T}{T_{\nu}}\right)^4 = \frac{4g}{7} \left[\frac{g'_{\star}}{g_{\star}(T_d)}\right]^4$$

Michael Shamma (He/They) mshamma@triumf.ca



4/3



Thermalized Dirac ν_R

Can use conservation of entropy for thermalized ν_R



Michael Shamma (He/They) mshamma@triumf.ca

¬ 4/3



Thermalized Dirac ν_R

Can use conservation of entropy for thermalized ν_R



$$\Delta N_{\rm eff} = \frac{4g}{7} \left[\frac{43}{4g_{\star}(T_d)} \right]^{4/3} = 0.027g \left[\frac{10}{g_{\star}} \right]^{4/3}$$

Michael Shamma (He/They) mshamma@triumf.ca

¬ 4/3

4/3 6.75



Can use conservation of entropy for thermalized ν_R



$$\Delta N_{\rm eff} = \frac{4g}{7} \left[\frac{43}{4g_{\star}(T_d)} \right]^{4/3} = 0.027g \left[\frac{10}{g_{\star}} \right]^{4/3}$$

Michael Shamma (He/They) mshamma@triumf.ca

Thermalized Dirac ν_R



Michael Shamma (He/They) mshamma@triumf.ca





Michael Shamma (He/They) mshamma@triumf.ca



$$\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB} , \frac{\Sigma m_\nu^{\text{tot}}}{\Sigma m_\nu} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB}$$

 ν_L relativistic at t_{CMB} and non-relativistic today Since ν_R colder, also non-relativistic today **Assume** ν_R relativistic at t_{CMB} Entropy conservation between $t_{\nu_I, \text{ dec}}, t_{\nu_R, \text{ dec}}$

 $+ \rho_R / \rho_L) |_{t=t_0},$



$$\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB} , \frac{\Sigma m_\nu^{\text{tot}}}{\Sigma m_\nu} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB}$$

 ν_L relativistic at t_{CMB} and non-relativistic today Since ν_R colder, also non-relativistic today **Assume** ν_R relativistic at t_{CMB} Entropy conservation between $t_{\nu_I, \text{ dec}}, t_{\nu_R, \text{ dec}}$

$$\implies \frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = 1 + \left(\frac{a_{\text{dec},R}T_{\text{dec},R}}{a_{\text{dec},L}T_{\text{dec},L}}\right)^4, \frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}} = 1 + \frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}}$$

Michael Shamma (He/They) mshamma@triumf.ca

 $+ \rho_R / \rho_L) |_{t=t_0},$

 $\frac{m_{\nu_R}}{m_{\nu_I}} \left(\frac{a_{\text{dec},R} T_{\text{dec},R}}{a_{\text{dec},L} T_{\text{dec},L}} \right)^3$



$$\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB} , \frac{\Sigma m_\nu^{\text{tot}}}{\Sigma m_\nu} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB}$$

 ν_L relativistic at t_{CMB} and non-relativistic today Since ν_R colder, also non-relativistic today **Assume** ν_R relativistic at t_{CMB} Entropy conservation between $t_{\nu_I, \text{ dec}}, t_{\nu_R, \text{ dec}}$

$$\implies \frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = 1 + \left(\frac{a_{\text{dec},R}T_{\text{dec},R}}{a_{\text{dec},L}T_{\text{dec},L}}\right)^4, \frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}} = 1 + \frac{\sum m_{\nu}^{T}}{\sum m_{\nu}}$$

Michael Shamma (He/They) mshamma@triumf.ca

 $+ \rho_R / \rho_L |_{t=t_0}$,

 $\frac{m_{\nu_R}}{m_{\nu_L}} \left(\frac{a_{\text{dec},R} T_{\text{dec},R}}{a_{\text{dec},L} T_{\text{dec},L}} \right)_3$



$$\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB} , \frac{\Sigma m_\nu^{\text{tot}}}{\Sigma m_\nu} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB}$$

 ν_L relativistic at t_{CMB} and non-relativistic today Since ν_R colder, also non-relativistic today **Assume** ν_R relativistic at t_{CMB} Entropy conservation between $t_{\nu_I, \text{ dec}}, t_{\nu_R, \text{ dec}}$

$$\implies \frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = 1 + \left(\frac{a_{\text{dec},R}T_{\text{dec},R}}{a_{\text{dec},L}T_{\text{dec},L}}\right)^4, \frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}} = 1 + \frac{\sum m_{\nu_R}}{\sum m_{\nu_L}} \left(\frac{a_{\text{dec},R}T_{\text{dec},R}}{a_{\text{dec},L}T_{\text{dec},L}}\right)^3$$
$$\left(\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} - 1\right)^{1/4} \left(\frac{\Delta N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}}\right)^{1/4} = \left(\frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}} - 1\right)^{1/3} = \left(\frac{m_{\nu,\text{sterile}}}{\sum m_{\nu}}\right)^{1/3}$$

Michael Shamma (He/They) mshamma@triumf.ca

 $+ \rho_R / \rho_L) |_{t=t_0},$

[κ. Abazajian, J. Heeck, hep-ph/ 1908.03286: P. Adshead, Y. Cui, A. Long, MS, hep-ph/ 2009.07852]





$$\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB} , \frac{\Sigma m_\nu^{\text{tot}}}{\Sigma m_\nu} = (1 + \rho_R / \rho_L) |_{t=t} \text{CMB}$$

 ν_L relativistic at t_{CMB} and non-relativistic today Since ν_R colder, also non-relativistic today **Assume** ν_R relativistic at t_{CMB} Entropy conservation between $t_{\nu_L, \text{ dec}}, t_{\nu_R, \text{ dec}}$

$$\implies \frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} = 1 + \left(\frac{a_{\text{dec},R}T_{\text{dec},R}}{a_{\text{dec},L}T_{\text{dec},L}}\right)^4, \frac{\sum m_{\nu}^{\text{tot}}}{\sum m_{\nu}} = 1 + \frac{\sum m_{\nu}^{N}}{\sum m_{\nu}}$$
$$\left(\frac{N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}} - 1\right)^{1/4} \left(\frac{\Delta N_{\text{eff}}}{N_{\text{eff}}^{\text{SM}}}\right)^{1/4} = \left(\frac{\sum m_{\nu}^{1}}{\sum m_{\nu}}\right)^{1/4}$$

Michael Shamma (He/They) <u>mshamma@triumf.ca</u>



BSM Dirac Neutrino Masses

- Dirac Leptogenesis [K. Dick, et. al. Phys. Rev. Lett., 84:4039–4042, 2000...] - Model dependent $\Delta N_{eff} \sim 0.05 - 0.14$
- Gauged $U(1)_{B-L}$ [V. Barger, et. al. Phys. Rev. D, 67:075009, 2003...] $-T_{\rm dec} \lesssim (m_{Z'}/g'M_{\rm pl})^{4/3}M_{\rm pl} \implies \Delta N_{\rm eff} \simeq 0.13 - 0.23$

 Neutrino mass models have a home in an extra dimension [N. Arkani-Hamed, et. al. hep-ph/9811448, K.R. Dienes, et. al. hep-ph/9811428]

- Compactification generates Dirac ν mass $\implies \mathscr{L} \supset -\frac{\lambda v}{\sqrt{2\pi R_{\text{ED}}M_*}} \bar{\nu}_L \nu_R^{(0)}$

– ...also generates mixing of the ν_k modes with ν_L - And substantial number of relics, strong constraints [K. Abazajian, G. Fuller, M. Patel hep-ph/0011048]

Michael Shamma (He/They) mshamma@triumf.ca





Michael Shamma (He/They) mshamma@triumf.ca



Bulk sterile neutrinos mix with active states:

 $\sin^2 2\theta_k \simeq \left(\frac{8mR}{k}\right)^2 = \left(\frac{8m}{M_k}\right)^2$

Michael Shamma (He/They) mshamma@triumf.ca



Bulk sterile neutrinos mix with active states: $\sin^2 2\theta_k \simeq \left(\frac{8mR}{k}\right)^2 = \left(\frac{8m}{M_k}\right)^2$

Produce a substantial number through freeze-in



Michael Shamma (He/They) mshamma@triumf.ca



Bulk sterile neutrinos mix with active states: $\sin^2 2\theta_k \simeq \left(\frac{8mR}{k}\right)^2 = \left(\frac{8m}{M_1}\right)^2$

Produce a substantial number through freeze-in

 $\frac{dY_k}{dT} = -\frac{1}{4} \frac{\sin^2 2\theta_k}{\left[1 + (T/T_V^k)^6\right]^2} \frac{M_{\text{Pl}}\Gamma_\nu Y_\nu}{\gamma T^3}$

For $T < T_{\rm ew} \implies$ Fermi interaction $\Gamma_{\nu} \simeq G_F^2 T^5$ Effective temperature $T_V^k \approx (T_{ew}^2 M_k)^{1/3}$

Michael Shamma (He/They) mshamma@triumf.ca



Bulk sterile neutrinos mix with active states: $\sin^2 2\theta_k \simeq \left(\frac{8mR}{k}\right)^2 = \left(\frac{8m}{M_k}\right)^2$

Produce a substantial number through freeze-in

 $\frac{dY_k}{dT} = -\frac{1}{4} \frac{\sin^2 2\theta_k}{\left[1 + (T/T_V^k)^6\right]^2} \frac{M_{\text{Pl}}\Gamma_\nu Y_\nu}{\gamma T^3}$

For $T < T_{ew} \implies$ Fermi interaction $\Gamma_{\nu} \simeq G_F^2 T^5$ Effective temperature $T_V^k \approx (T_{ew}^2 M_k)^{1/3}$

Max production occurs with $T \sim \min(T_V^k, T_{ew}) \implies M_k \leq \min(T_V^k, T_{ew})$

Michael Shamma (He/They) mshamma@triumf.ca



Bulk sterile neutrinos mix with active states: $\sin^2 2\theta_k \simeq \left(\frac{8mR}{k}\right)^2 = \left(\frac{8m}{M}\right)^2$

Produce a substantial number through freeze-in

 $\frac{dY_k}{dT} = -\frac{1}{4} \frac{\sin^2 2\theta_k}{\left[1 + (T/T_V^k)^6\right]^2} \frac{M_{\text{Pl}}\Gamma_\nu Y_\nu}{\gamma T^3}$

For $T < T_{\rm ew} \implies$ Fermi interaction $\Gamma_{\nu} \simeq G_F^2 T^5$ Effective temperature $T_V^k \approx (T_{ew}^2 M_k)^{1/3}$

Max production occurs with $T \sim \min(T_V^k, T_{ew}) \implies M_k \leq \min(T_V^k, T_{ew})$

Michael Shamma (He/They) mshamma@triumf.ca

Probing Large Extra Dimensions? [D. Mckeen, J. Ng., MS in prep]



All modes $M_k < M_W, M_Z$

$$\tau_k \simeq 1.9 \times 10^{26} \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3} k^{-3}$$

All modes $M_k < M_W, M_Z$

$$\tau_k \simeq 1.9 \times 10^{26} \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3} k^{-3}$$

Heaviest mode has

$$\tau_k^{\rm min} \simeq 1.4 \times 10^3 \, {\rm s} \left(\frac{0.1 \, {\rm eV}}{m} \right)^2$$

All modes $M_k < M_W, M_Z$

$$\tau_k \simeq 1.9 \times 10^{26} \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3} k^{-3}$$

 $\tau_k < \tau_{\rm CMB} \implies \Delta N_{\rm eff}$

Heaviest mode has

$$\tau_k^{\min} \simeq 1.4 \times 10^3 \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2$$

 $\tau_k \gtrsim 14 \text{ Gyr} \implies \text{Stable DM } \Omega_{\text{DM}} h^2 < 0.1$ $\tau_{\rm CMB} < \tau_k < 14 \; {\rm Gyr} \implies {\rm Decaying \; DM}$

All modes $M_k < M_W, M_Z$

$$\tau_k \simeq 1.9 \times 10^{26} \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^3 k^{-3}$$

 $\tau_k < \tau_{\rm CMB} \implies \Delta N_{\rm eff}$

High Reheat Universe, $T_{\rm RH} > T_{\rm ew}$:

$$\frac{n_k}{n_\nu} \simeq 2.32 \times 10^{-2} \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left(\frac{R}{10^{-4} \ \mu m}\right) \frac{1}{k} \left[\frac{10.75}{g_*(T_V^k)}\right]^{11/6}$$
$$\implies k_{\text{max}} = \min\left[6.5 \times 10^7 \left(\frac{R}{10^{-4} \ \mu m}\right), 5.1 \times 10^7 \left(\frac{R}{10^{-4} \ \mu m}\right)\right]$$

Heaviest mode has

$$\tau_k^{\rm min} \simeq 1.4 \times 10^3 \, {\rm s} \left(\frac{0.1 \, {\rm eV}}{m} \right)^2$$

 $\tau_k \gtrsim 14 \text{ Gyr} \implies \text{Stable DM } \Omega_{\text{DM}} h^2 < 0.1$ $\tau_{\rm CMB} < \tau_k < 14 \ {\rm Gyr} \implies {\rm Decaying \ DM}$

All modes $M_k < M_W, M_Z$

$$\tau_k \simeq 1.9 \times 10^{26} \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^3 k^{-3}$$

 $\tau_k < \tau_{\rm CMB} \implies \Delta N_{\rm eff}$

High Reheat Universe, $T_{\rm RH} > T_{\rm ew}$:

$$\frac{n_k}{n_\nu} \simeq 2.32 \times 10^{-2} \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left(\frac{R}{10^{-4} \ \mu m}\right) \frac{1}{k} \left[\frac{10.75}{g_*(T_V^k)}\right]^{11/6}$$
$$\implies k_{\text{max}} = \min\left[6.5 \times 10^7 \left(\frac{R}{10^{-4} \ \mu m}\right), 5.1 \times 10^7 \left(\frac{R}{10^{-4} \ \mu m}\right)\right]$$

Heaviest mode has

$$\tau_k^{\min} \simeq 1.4 \times 10^3 \text{ s} \left(\frac{0.1 \text{ eV}}{m}\right)^2$$

 $\tau_k \gtrsim 14 \text{ Gyr} \implies \text{Stable DM } \Omega_{\text{DM}} h^2 < 0.1$ $\tau_{\rm CMB} < \tau_k < 14 \; {\rm Gyr} \implies {\rm Decaying } \, {\rm DM}$

Low Reheat Universe, $T_{\rm RH} \sim \mathcal{O}(10 \text{ MeV})$:

$$\frac{n_k}{n_\nu} \simeq 7.1 \times 10^{-7} \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left(\frac{R}{10^{-4} \ \mu m}\right)^2 \frac{1}{k^2} \left[\frac{10.75}{g_*(T_{\text{RH}})}\right]^{1/2}$$
$$\implies k_{\text{max}} = 5.07 \times 10^3 \left(\frac{R}{10^{-4} \ \mu m}\right) \left(\frac{T_{\text{RH}}}{10 \text{ MeV}}\right)$$

Michael Shamma (He/They) mshamma@triumf.ca



$$\Delta N_{\text{eff}} \simeq \sum_{k} 3.25 \times 10^{6} \left(\frac{m}{0.1 \text{ eV}}\right) \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3/2} k^{-3/2}$$
$$\Omega_{DM} \simeq \sum_{k=1}^{k_{tU}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$
$$\Omega_{DM} \simeq \sum_{k=k_{tU}}^{k_{\text{CMB}}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$

Michael Shamma (He/They) <u>mshamma@triumf.ca</u>



$$\Delta N_{\text{eff}} \simeq \sum_{k} 3.25 \times 10^{6} \left(\frac{m}{0.1 \text{ eV}}\right) \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3/2} k^{-3/2}$$
$$\Omega_{DM} \simeq \sum_{k=1}^{k_{tU}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$
$$\Omega_{DM} \simeq \sum_{k=k_{tU}}^{k_{\text{CMB}}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$

Michael Shamma (He/They) mshamma@triumf.ca



$$\Delta N_{\text{eff}} \simeq \sum_{k} 3.25 \times 10^{6} \left(\frac{m}{0.1 \text{ eV}}\right) \left(\frac{R}{10^{-4} \ \mu\text{m}}\right)^{3/2} k^{-3/2}$$
$$\Omega_{DM} \simeq \sum_{k=1}^{k_{tU}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$
$$\Omega_{DM} \simeq \sum_{k=k_{tU}}^{k_{\text{CMB}}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^{2} \left[\frac{10.75}{g_{*}(T_{V}^{k})}\right]^{11/6}$$

Michael Shamma (He/They) mshamma@triumf.ca



March 2024

$$\Delta N_{\text{eff}} \simeq \sum_{k} 3.25 \times 10^{6} \left(\frac{m}{0.1 \text{ eV}}\right) \left(\frac{R}{10^{-4} \ \mu \text{m}}\right)^{3/2} k^{-3/2}$$

$$\Omega_{DM} \simeq \sum_{k=1}^{k_{tU}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left[\frac{10.75}{g_*(T_V^k)}\right]^{11/6}$$

$$\Omega_{DM} \simeq \sum_{k=k_{tU}}^{k_{\text{CMB}}} 0.5 \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left[\frac{10.75}{g_*(T_V^k)}\right]^{11/6}$$

Low Reheat Universe, $T_{\rm RH} > T_{\rm ew}$:

$$\Delta N_{\rm eff} \simeq \sum_{k=1}^{k_{\rm TCMB}} 7.9 \times 10^{-17} \left(\frac{m}{0.1 \text{ eV}}\right)^4 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{T_{\rm RH}}{10 \text{ MeV}}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{R}\right)^3 + \sum_{k=k_{t_{\rm CMB}}}^{k_{\rm max}} 6.7 \times 10^3 \left(\frac{10^{-4} \ \mu m}{R}\right)^2 k^2 \left(\frac{10^{-4} \ \mu m}{$$

$$\Omega_{DM} \simeq \sum_{k=1}^{k_{tU}} 1.5 \times 10^{-5} \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right) k^{-1} \left[\frac{10.75}{g_*(T_{\text{RH}})}\right]$$

$$\Omega_{DM} \simeq \sum_{k=k_{tU}}^{k_{\text{CMB}}} 1.5 \times 10^{-5} \left(\frac{m}{0.1 \text{ eV}}\right)^2 \left(\frac{R}{10^{-4} \ \mu\text{m}}\right) k^{-1} \left[\frac{10.75}{g_*(T_{\text{RH}})}\right]$$

Michael Shamma (He/They) mshamma@triumf.ca



March 2024

- New physics with Dirac ν provide well-motivated explanation of m_{ν}
- New interactions can produce substantial numbers of ν_R which potentially thermalize
 - Upcoming cosmological surveys can probe Dirac ν !
- Future surveys may probe BSM physics through correlated observables - This diagnostic test can distinguish generic eV-scale relics ($m_{relic} \neq m_{\nu_{I}}$) from the DNH ($m_{\nu_{I}} = m_{\nu_{P}}$)
- Upcoming work explores the scenario in which many possibly non-degenerate eV-keV scale relics are produced alongside the degenerate ν_R
 - E.g. will provide constraints on the size of dark dimensions, neutrino masses

Summary






Thank you!

