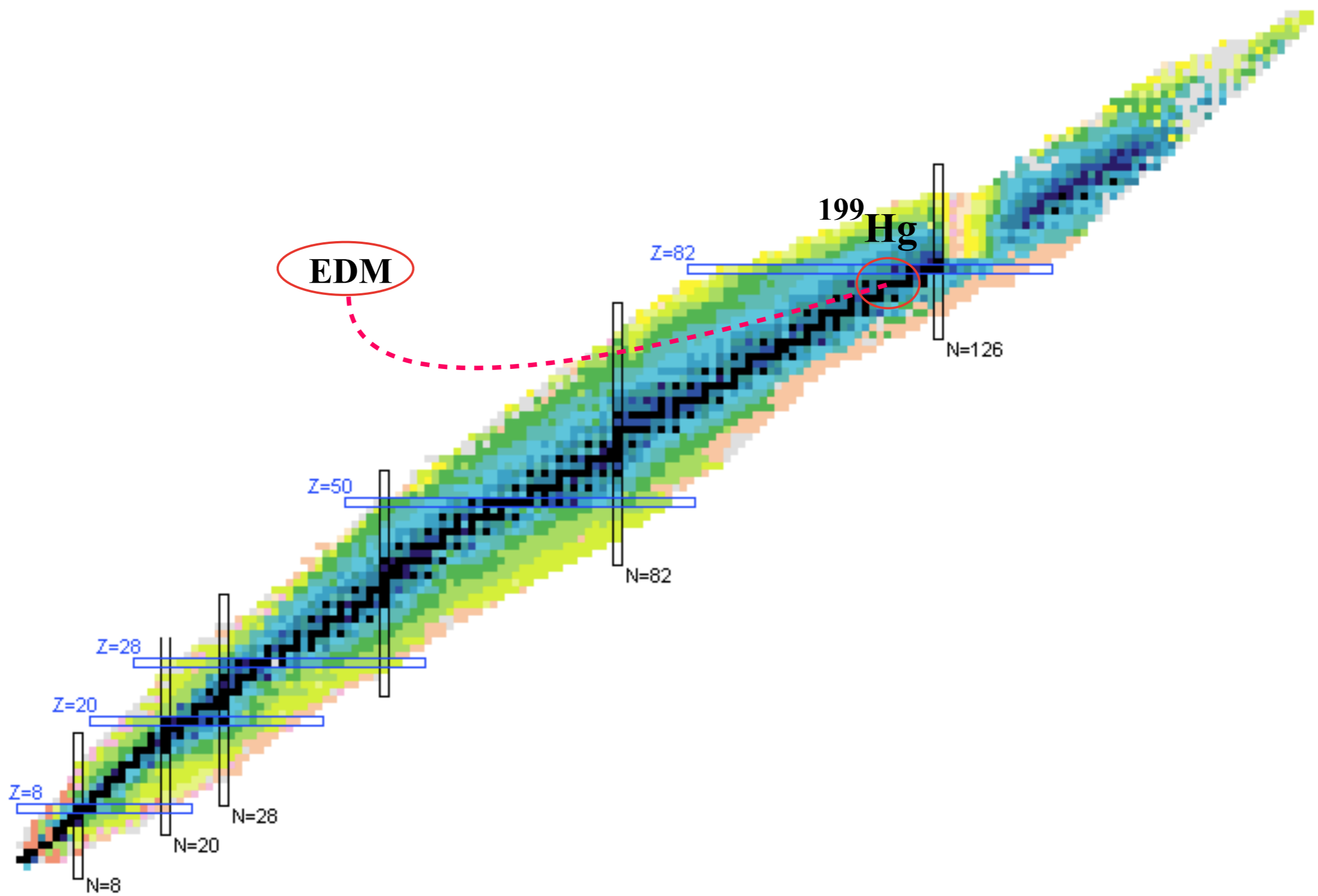


Study of the  $^{198}\text{Hg}(d,d')$  Inelastic  
Scattering Reaction

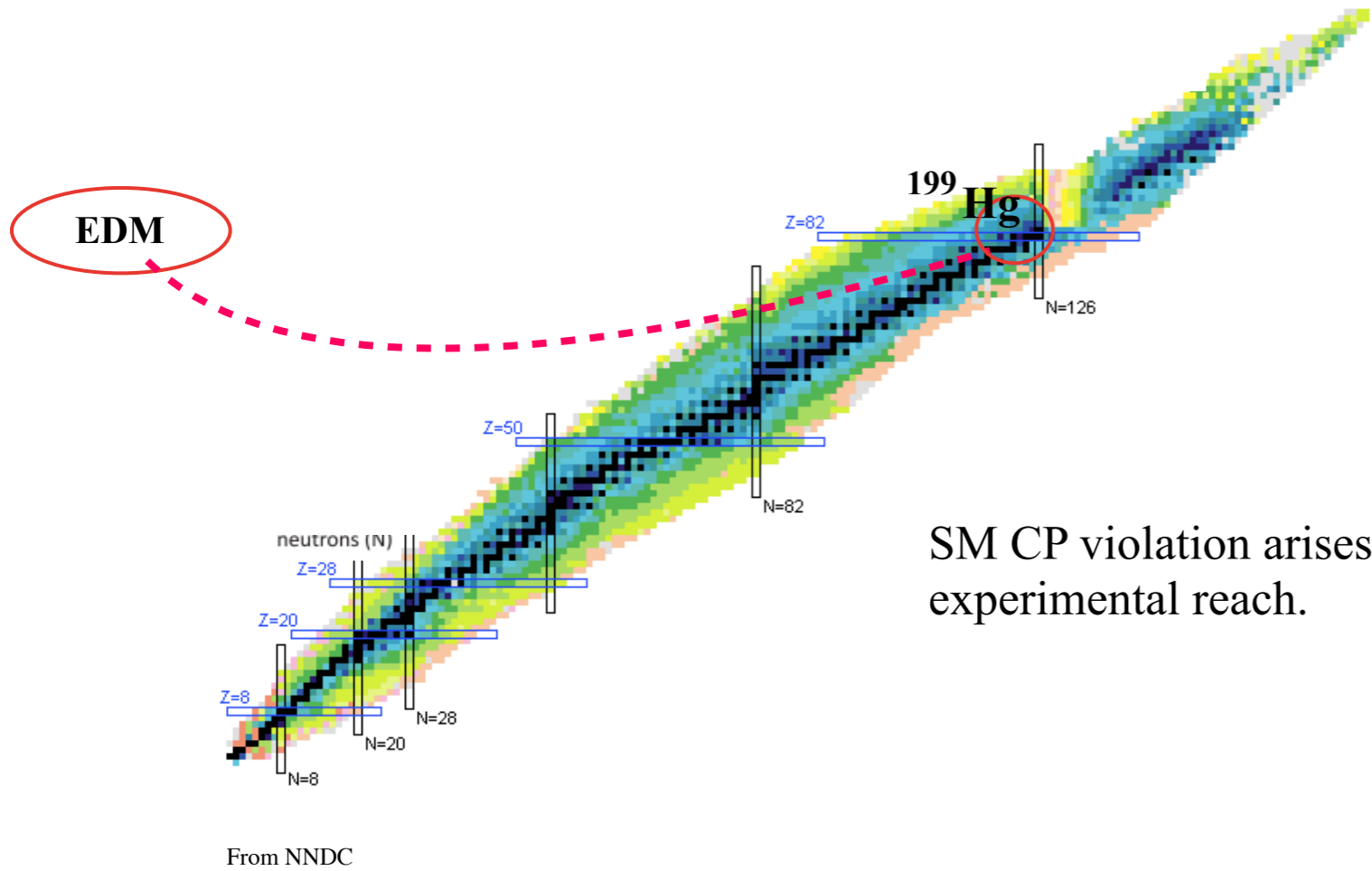
**Sally Valbuena**

Department of Physics  
University of Guelph, Ontario

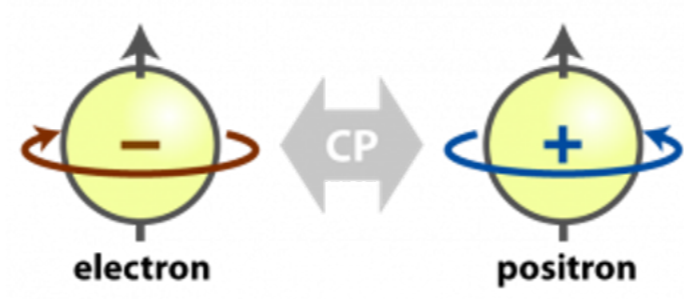
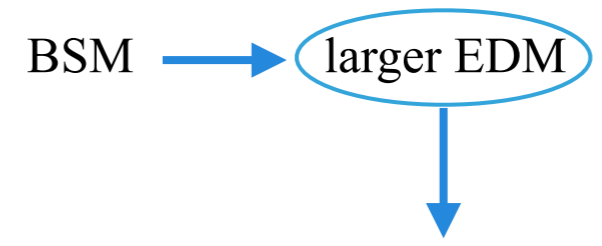
WNPPC February 2024

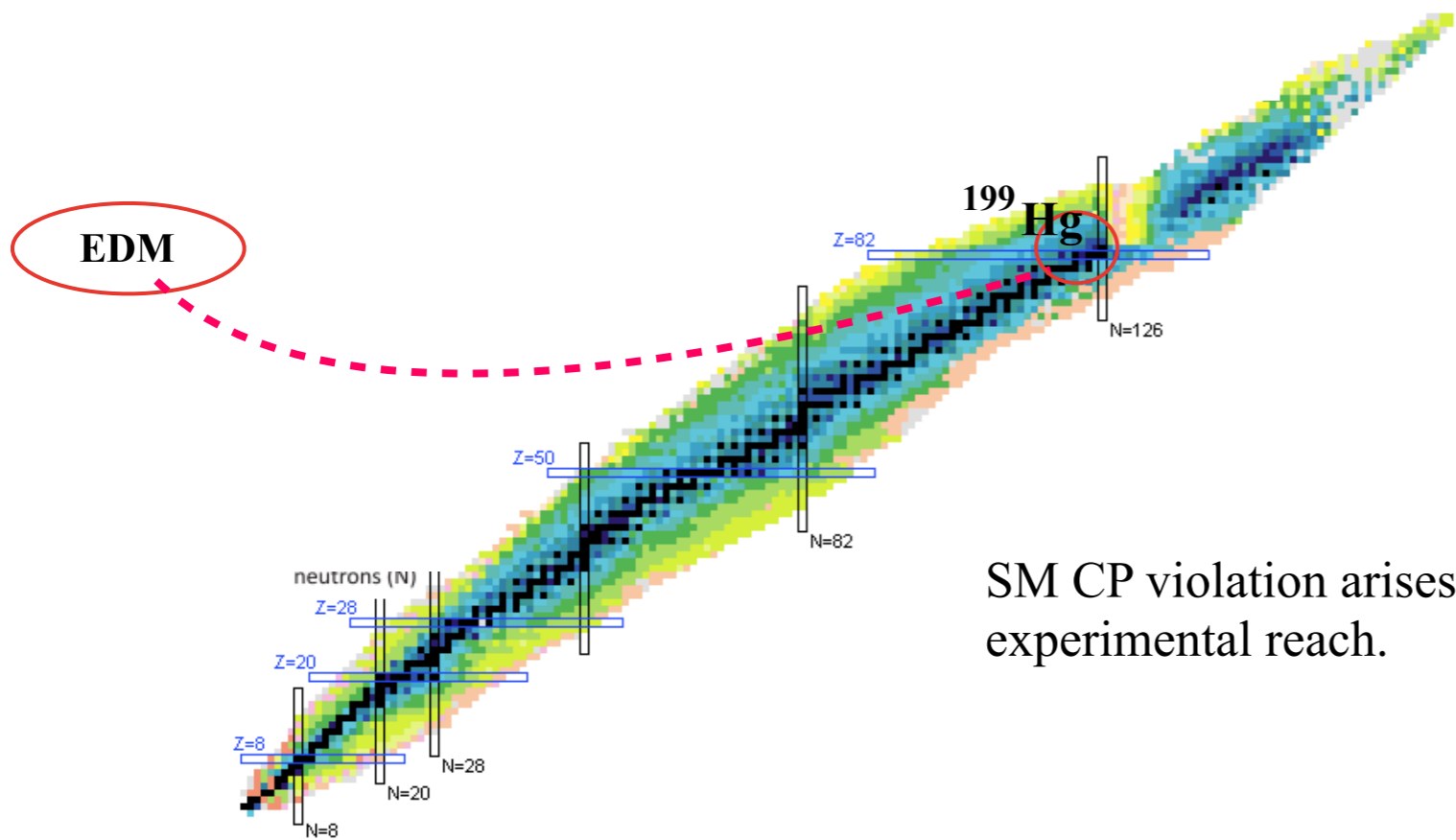


From NNDC

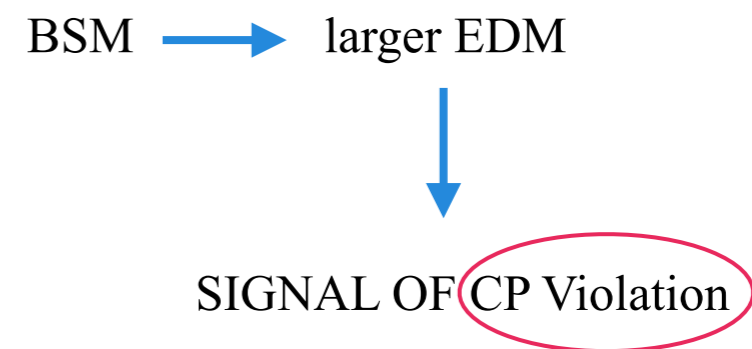
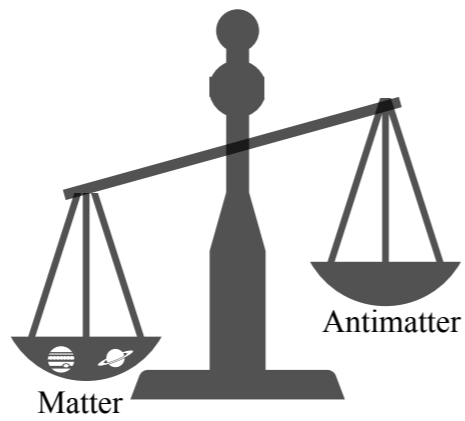


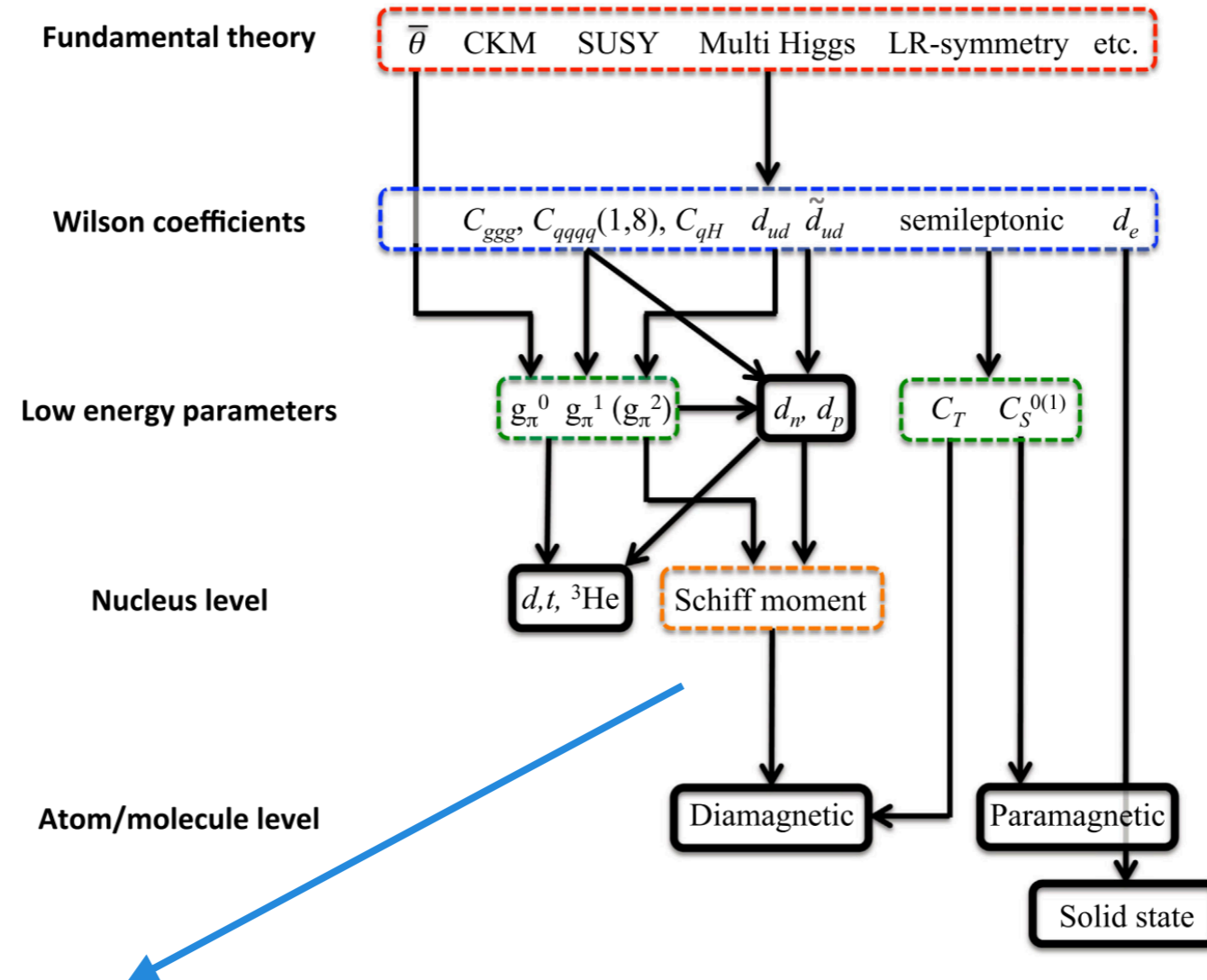
SM CP violation arises from a complex phase in CKM  $\ll$  Current experimental reach.





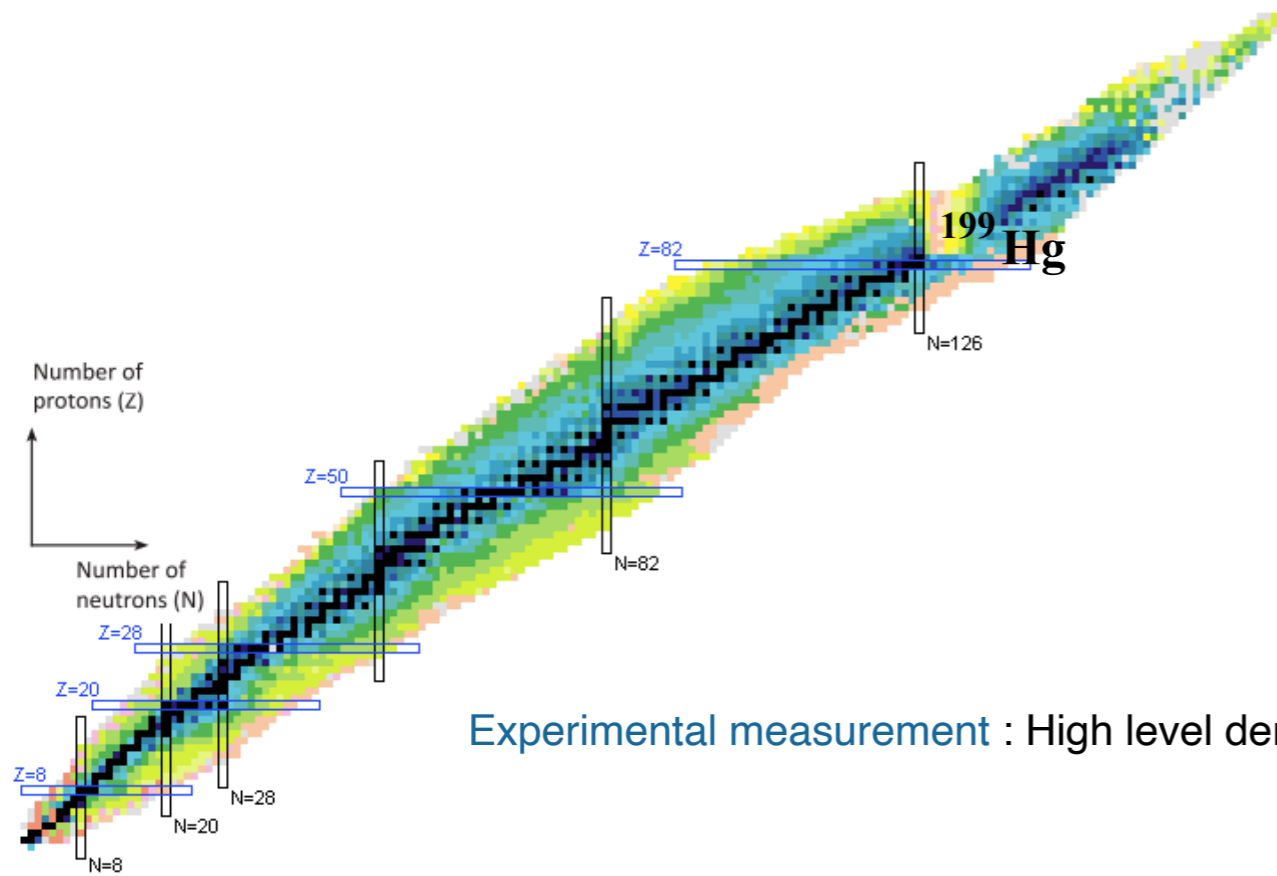
SM CP violation arises from a complex phase in CKM  $\ll$  Current experimental reach.





Lowest-order observable EDM for a neutral atom in its ground state

$$\mathbf{S} = \frac{1}{10} \left( \int e \rho(\mathbf{r}) r^2 \mathbf{r} d^3\mathbf{r} - \frac{5}{3} \mathbf{d} \frac{1}{Z} \int \rho(\mathbf{r}) r^2 d^3\mathbf{r} \right), \quad \mathbf{d} = \frac{1}{10} \left( \int e \mathbf{r} \rho(\mathbf{r}) d^3\mathbf{r} \right)$$

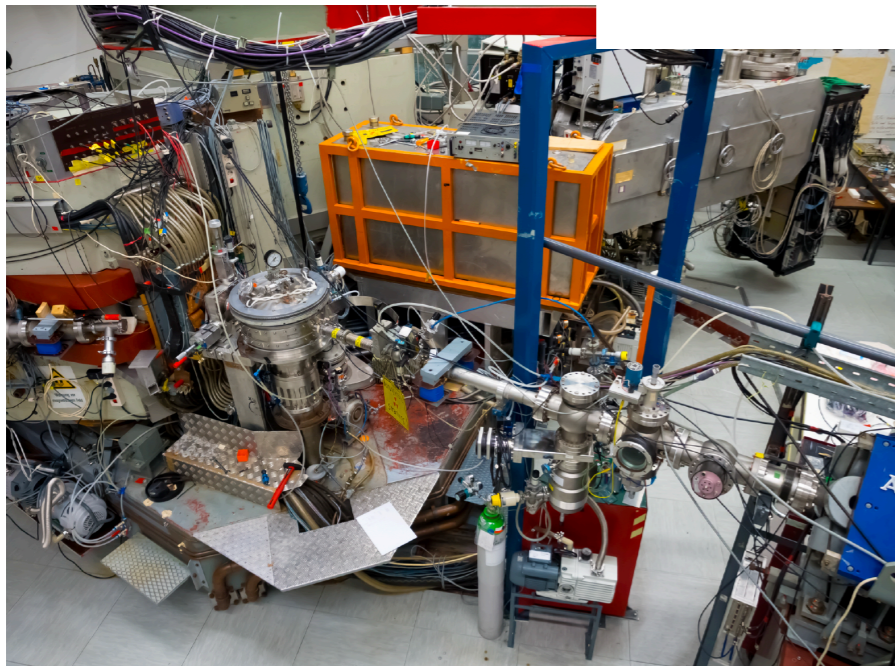
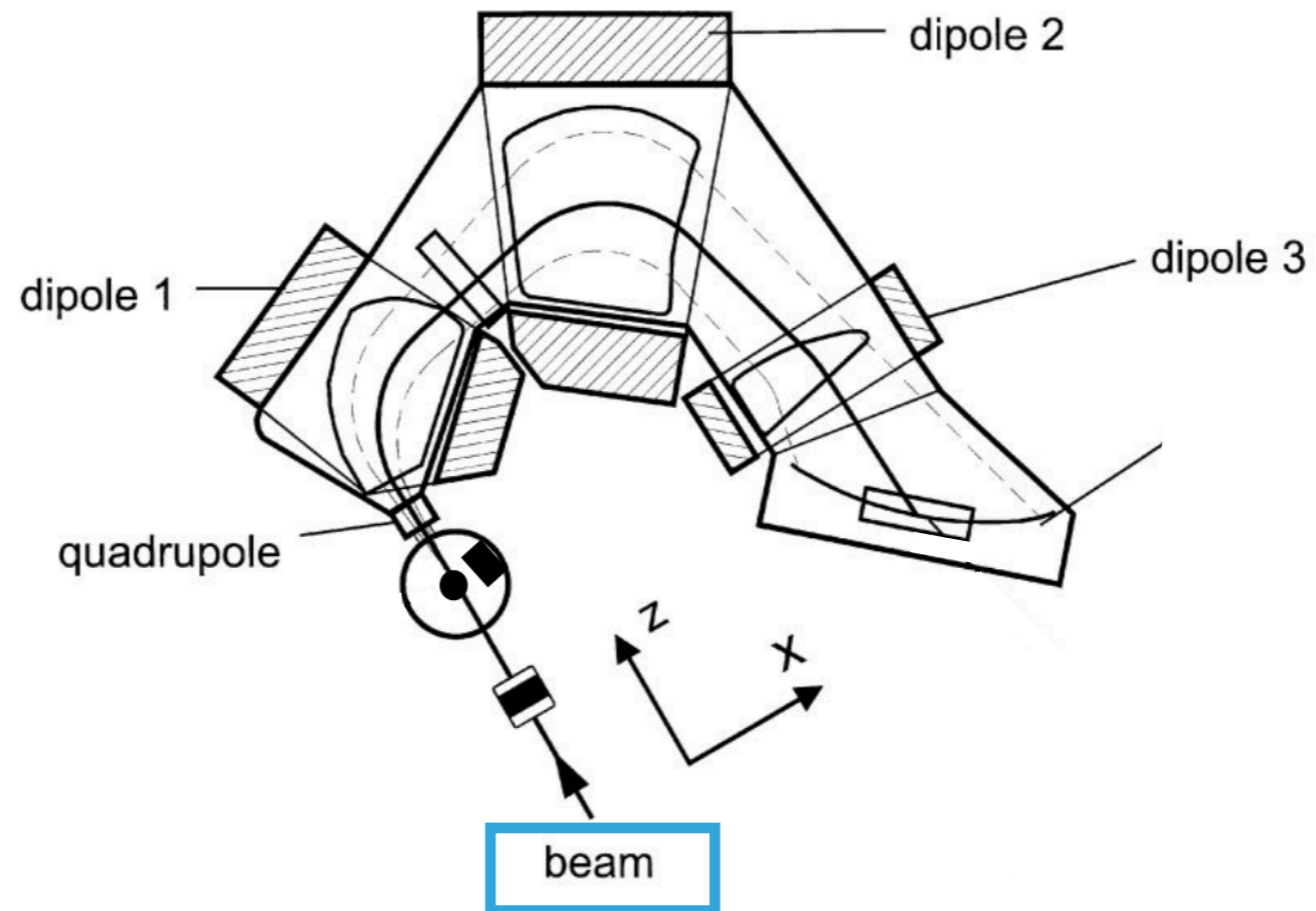


Experimental measurement : High level density of a heavy odd- A nucleus.

[1]

**SOLUTION** → Information from states in the neighbouring even-even isotopes of  $^{198}\text{Hg}$  and  $^{200}\text{Hg}$

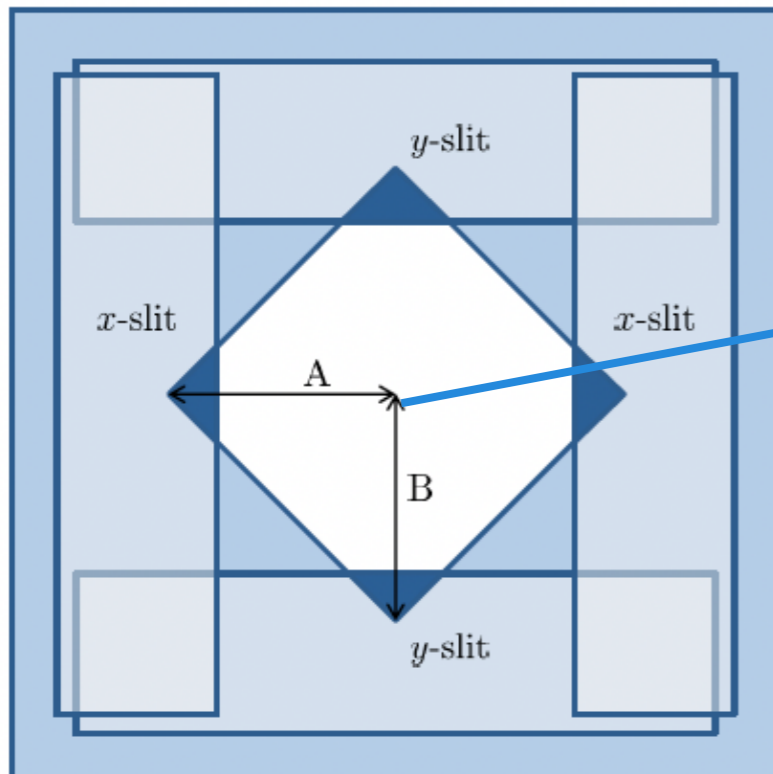
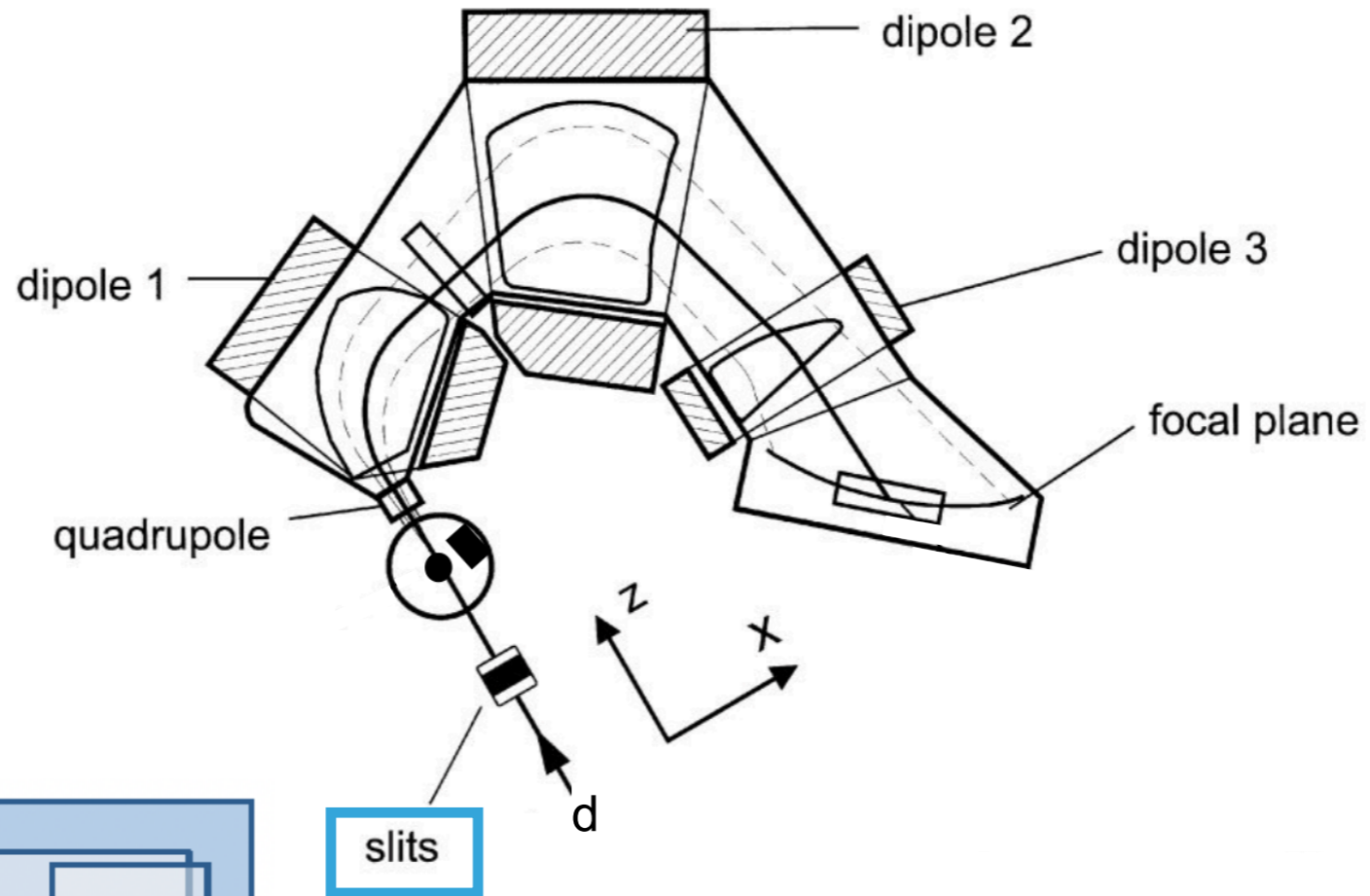
E2 and E3 strengths → Constrains present models.



Q3D magnetic spectrograph at MLL

- ▶ Beam: d
- ▶ Beam energy: 22 MeV
- ▶ Beam current: 1  $\mu\text{A}$

From E.T. Rand.

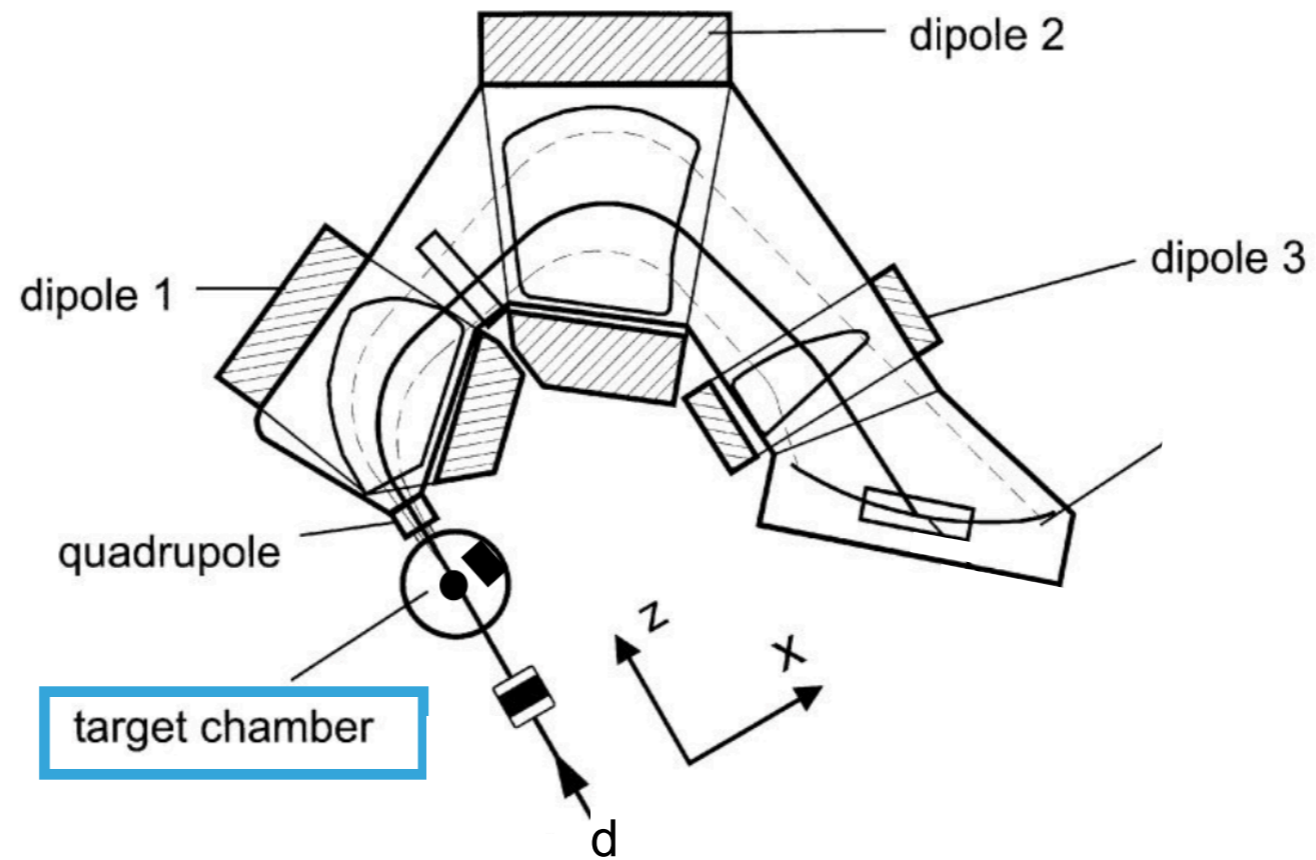


Q3D acceptance solid angle

$$\begin{matrix} A=31.5 \text{ mm} \\ B=36.75 \text{ mm} \end{matrix}$$

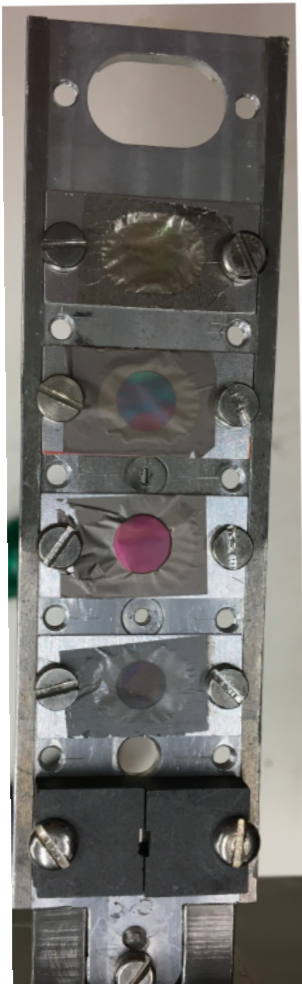
$$\Omega = \frac{(F(x, y) - \Delta F(x, y))}{D^2},$$





Q3D magnetic spectrograph at MLL

- ▶ Target:  $^{198}\text{Hg}^{32}\text{S}$  (thickness  $\sim 40 - 95 \mu\text{g}/\text{cm}^2$ )
- ▶ Mounted to a target ladder.
- ▶ MLL's target ladder hold up five target.

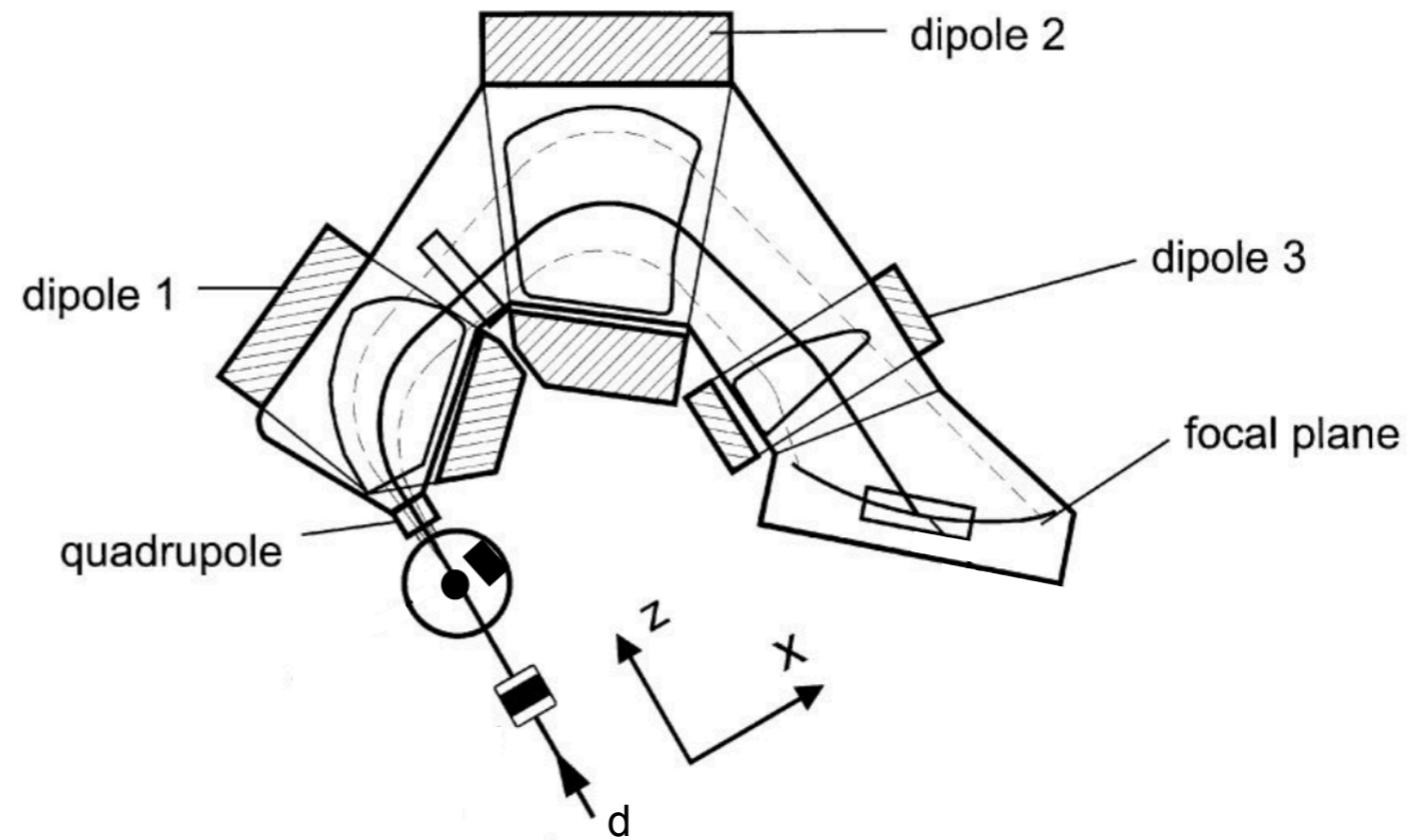


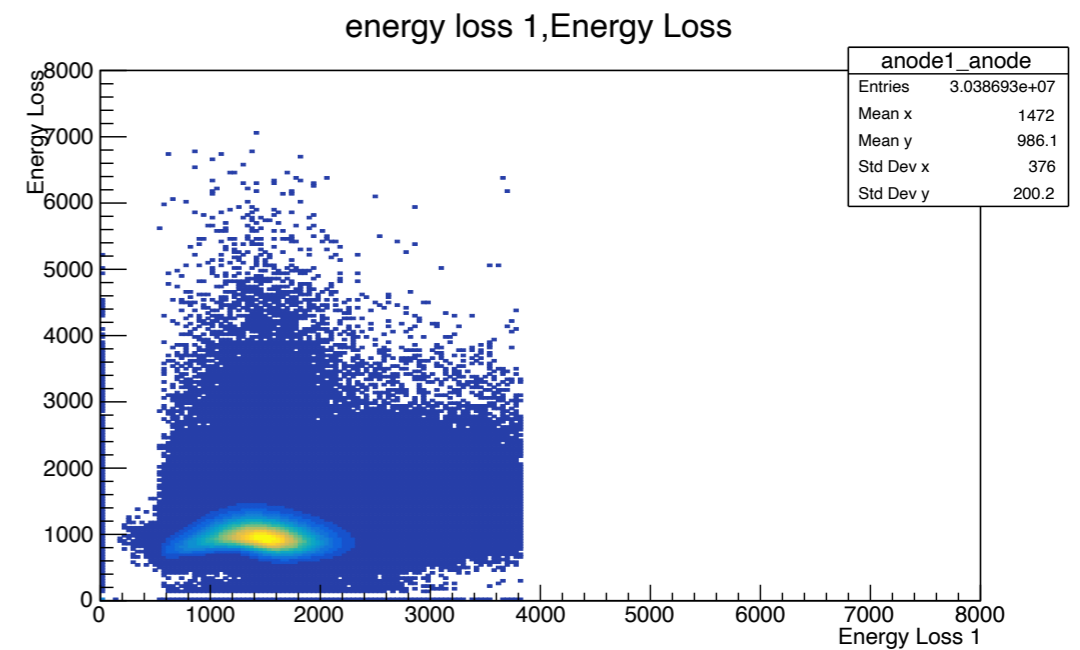
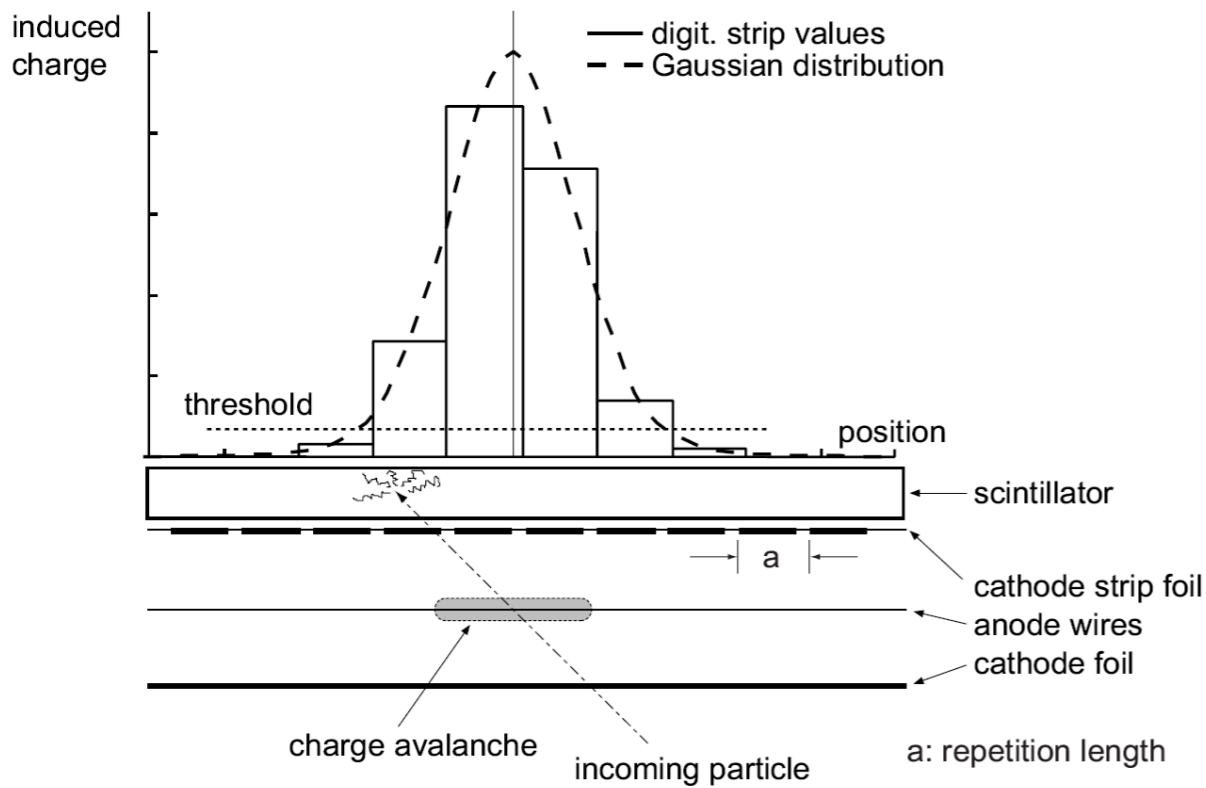
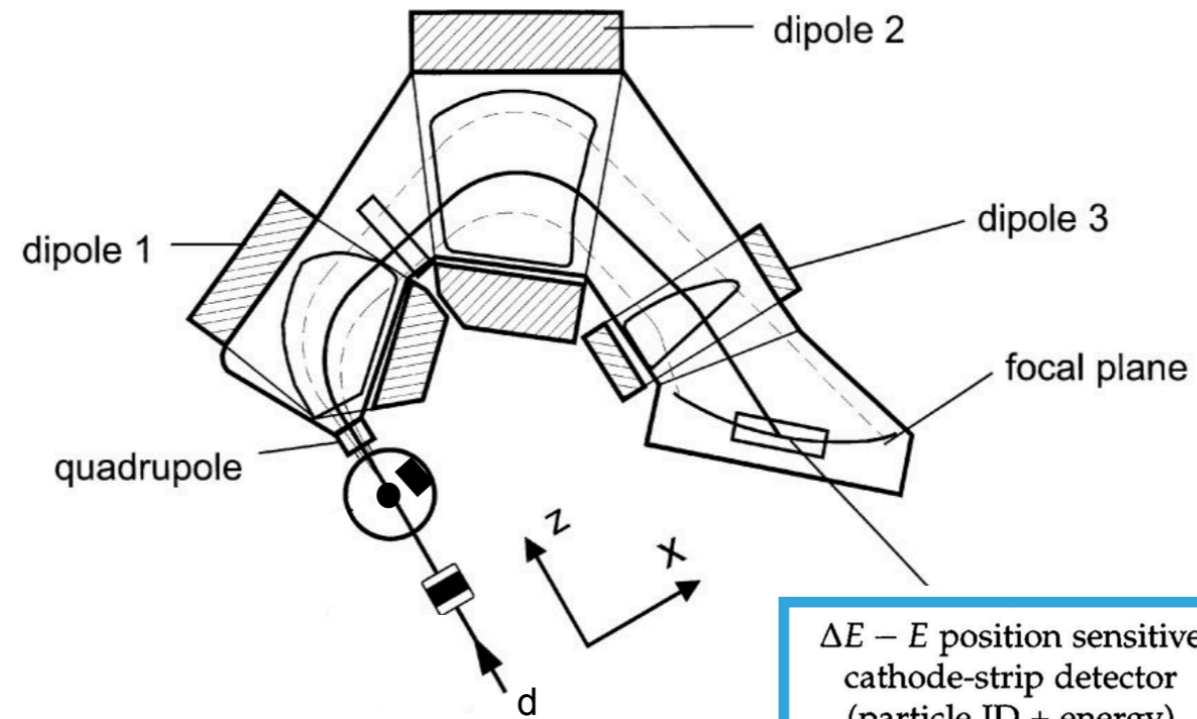
From C. Burbadge.

Deuterons travel in a circular path:

$$F = qvB = m \frac{v^2}{R},$$

$$R = \frac{mv}{qB},$$



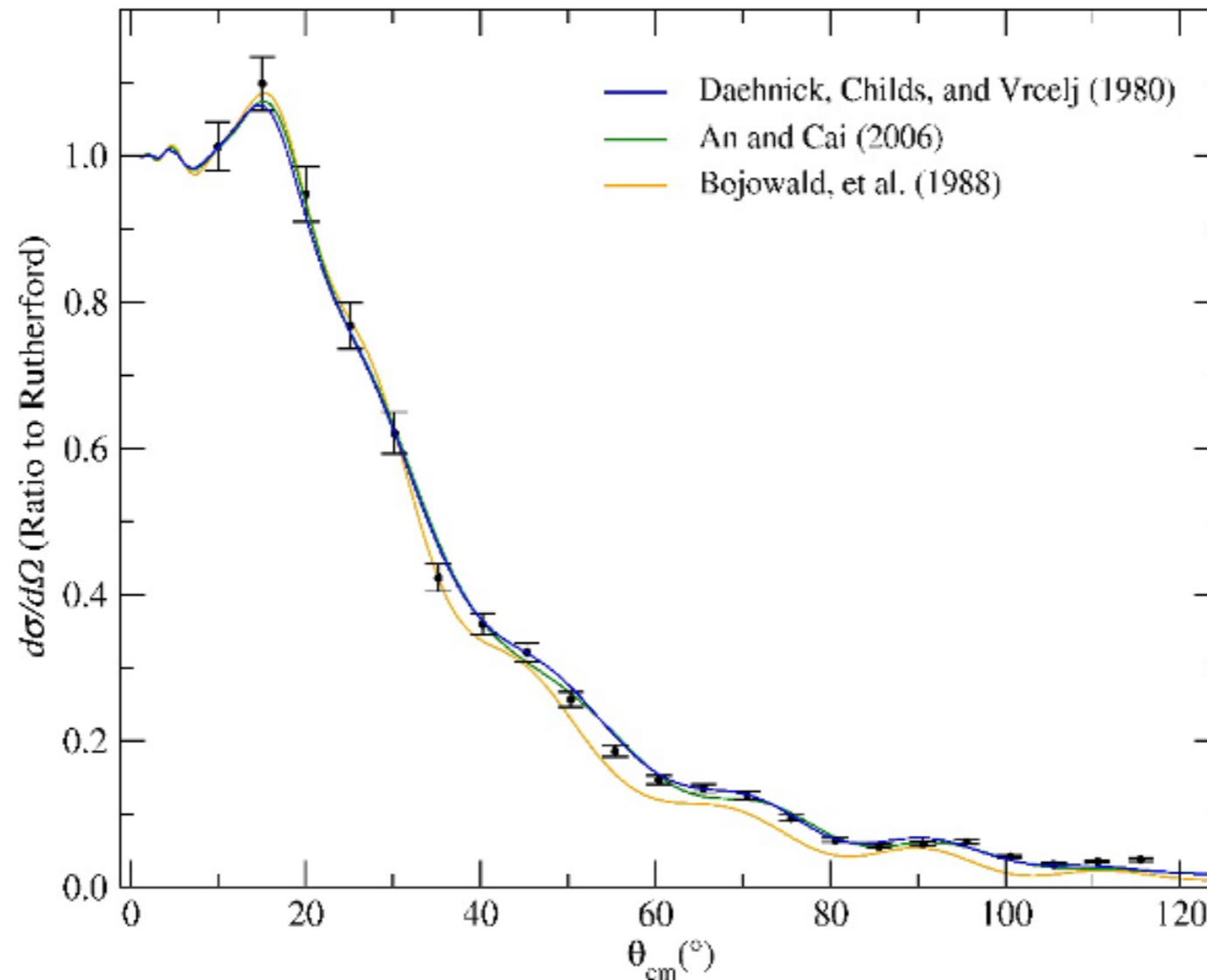


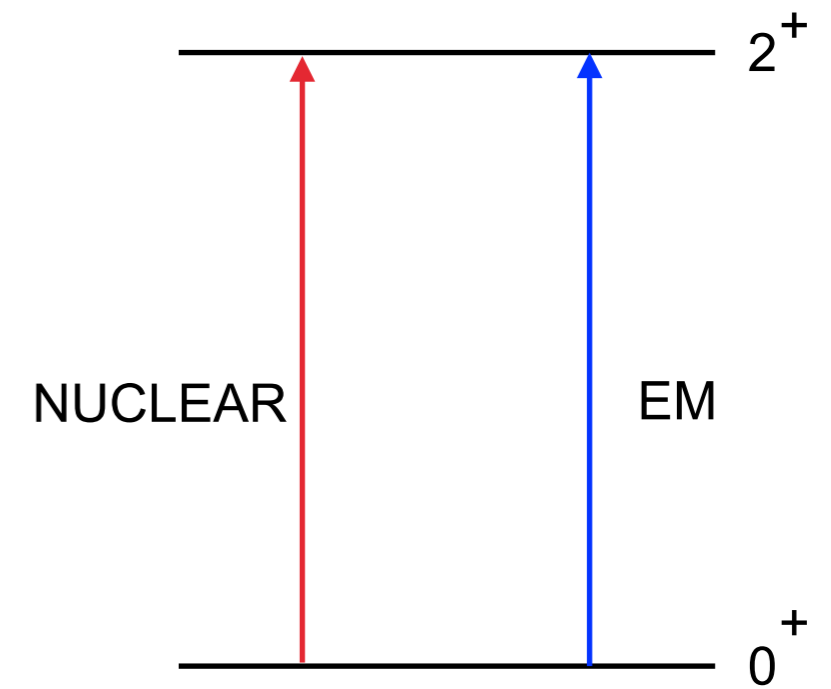
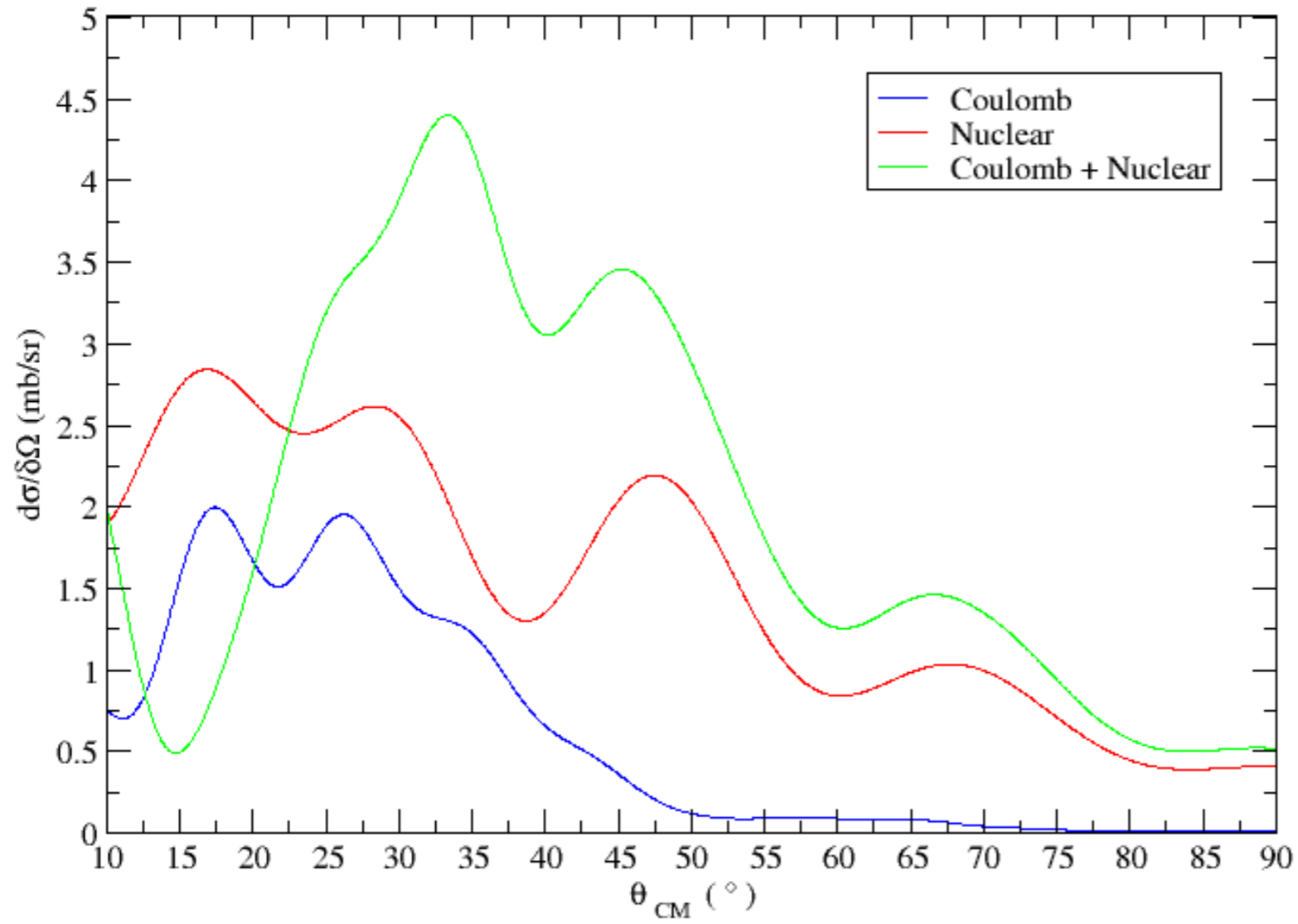
[3]

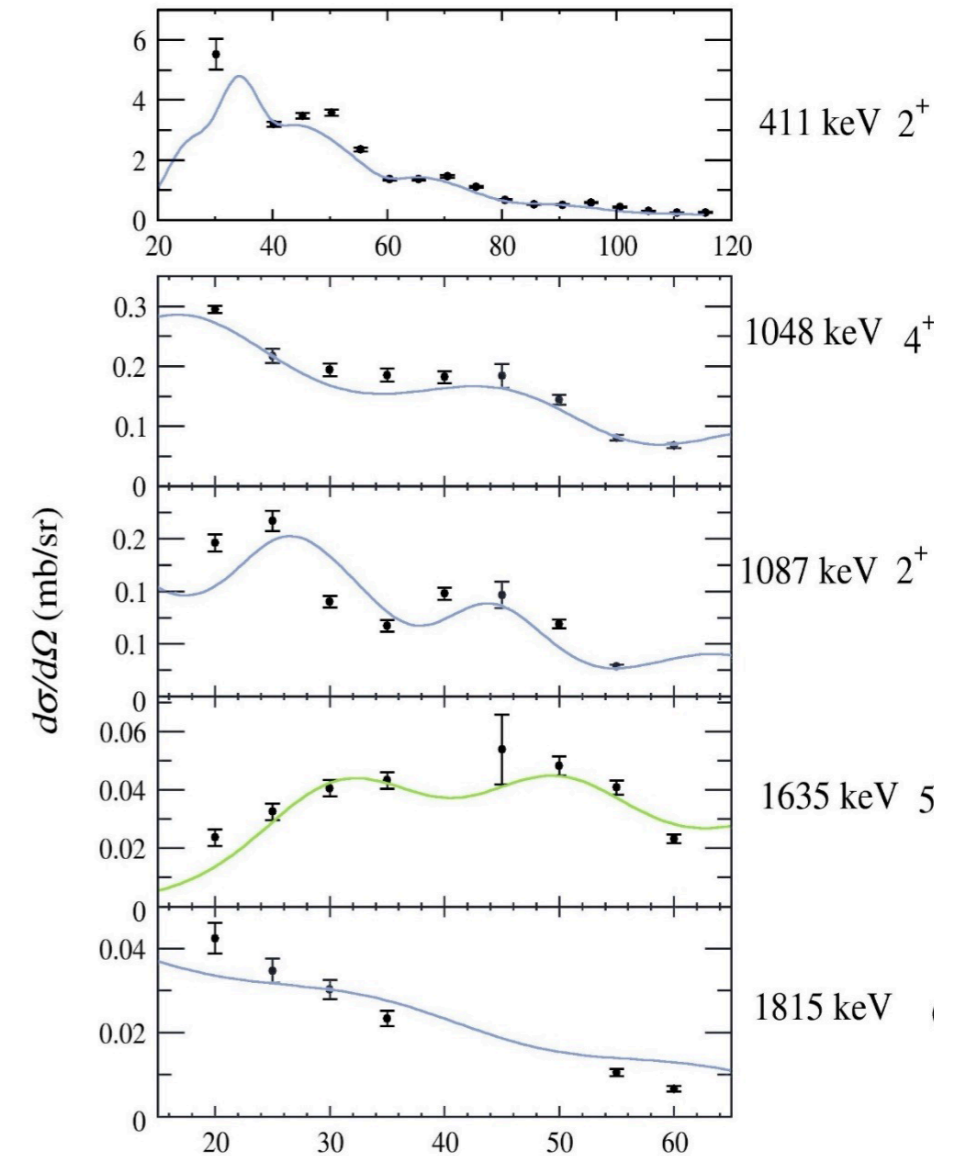
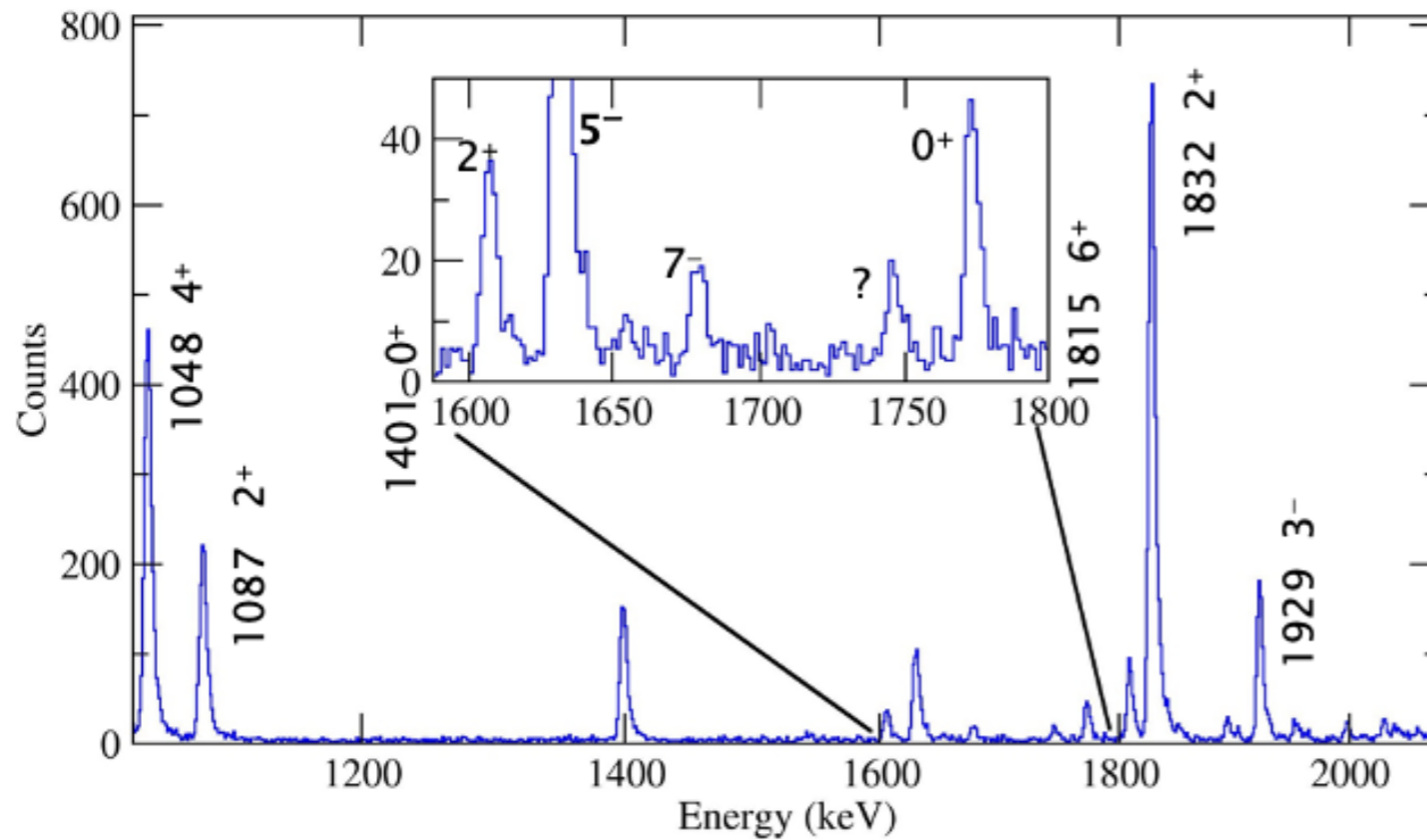
Experimental cross section measurements:  $\left(\frac{d\sigma}{d\Omega}\right)_{\text{lab}} = \frac{N_c}{N_{t,\text{eff}} N_b \Omega_{LT} \text{DAQ} \text{LT}_{\text{DET}}} \times 10^{31} \left[\frac{\text{mb}}{\text{sr}}\right]$

Optical potential by *An and Cai* [7]:

$$V(r) = -V f_r(r) - iW_v f_v(r) + i4a_s W_s \frac{df_s(r)}{dr} + \lambda_{\pi}^2 \frac{V_{\text{so}} + W_{\text{so}}}{r} \frac{df_{\text{so}}(r)}{dr} \vec{\sigma} \cdot \vec{l} + V_C(r)$$







1. National nuclear data center. <https://www.nndc.bnl.gov/ensdf/>. National nuclear data center. <https://www.nndc.bnl.gov/ensdf/>.
2. J.M. Pendlebury and E.A. Hinds. Particle electric dipole moments. 2000. [https://doi.org/10.1016/S0168-9002\(99\)01023-2](https://doi.org/10.1016/S0168-9002(99)01023-2)
3. E.T. Rand. Investigation of the E2 and E3 matrix elements in  $^{200}\text{Hg}$  using direct nuclear reactions. PhD dissertation, University of Guelph, 2015.
4. T. E. Chupp, P. Fierlinger, M.J. Ramsery-Musolf and J.T. Singh. Electric dipole moments of atoms, molecules, nuclei and particles.
5. S. M. Wong. Introductory Nuclear Physics. Wiley-VCH Verlag GmbH & Co. KGaA, 2004.
6. P.E. Garret, E. T. Rand, A. Diaz Varela, G.C. Ball, V. Bildstein, T. Faestermann, B. Hadinia, R. Hertenberg, D. S. Jamieson, B. Jigmeddorj, K. G. Leach, C. E. Svensson, and H.-F. Wirth. Direct reactions for nuclear structure required for fundamental symmetry tests. DOI: 10.1051/epjconf/201612303003, 2016.
7. H. An and C. Cai. Global deuteron optical model potential for the energy range up to 183 MeV. Physical Review C, 73(054605), May 2006.
8. J. Bojowald, H. Machner, H. Nann, M. Oelert, M. Rogge, and P. Turek. Elastic deuteron scattering and optical model parameters at energies up to 100 MeV. Physical Review C, 38(3), September (1988).
9. W. Daehnick, J. Childs, and Z. Vrcelj. Global optical model potential for elastic deuteron scattering from 12 to 90 MeV. Physical Review C, 21(number 6), June 1980.
10. Y. Han, Y. Shi, and Q. Shen. Deuteron global optical model potential for energies up to 200 MeV. Physical Review C, 74(044615), October 2006.

**THANK YOU FOR THE ATTENTION**



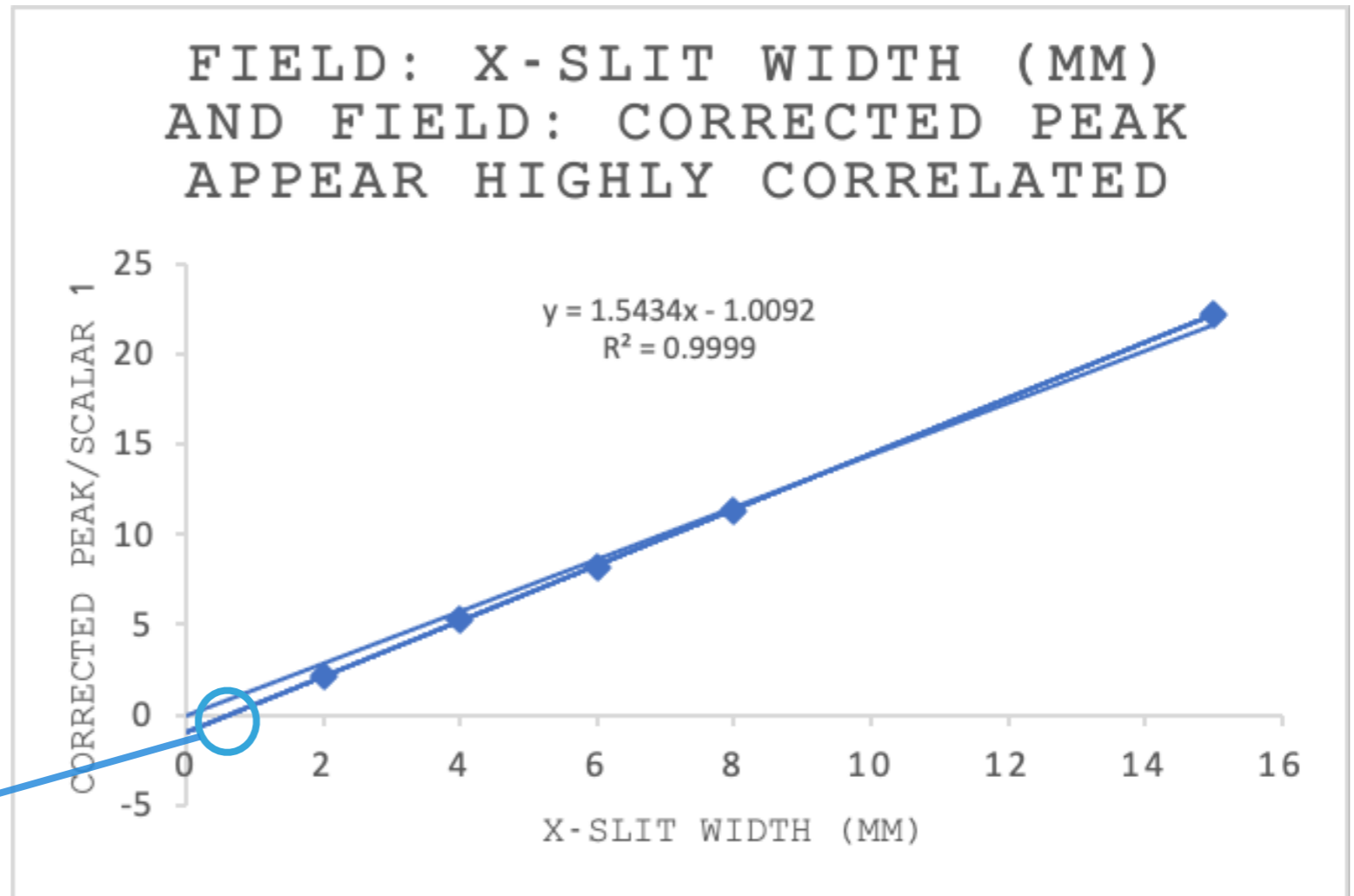
QUESTIONS?

# BACKUP SLIDES

Beginning date of Experiment	Momentum Bite	Number of Runs
January 23, 2014	730, 750	46
	1730, 1760	11
	3130, 3160	15
December 17, 2016	700	14
December 17, 2016 $^{138}\text{Ba}(d, d')$ data-Solid angle	500	6
July 14, 2017	750	8
July 24, 2017 $^{136}\text{Ba}(p, P)$ data-Solid angle	400	5

Angle 30°	
$x$ -slit $\times$ $y$ -slit (mm)	corrected peak/Scalar 1
2 $\times$ 24.5	2.088784711
4 $\times$ 24.5	5.264463509
6 $\times$ 24.5	8.114882187
8 $\times$ 24.5	11.33647109
15 $\times$ 24.5	22.17004426
-21.5 $\times$ 24.5	

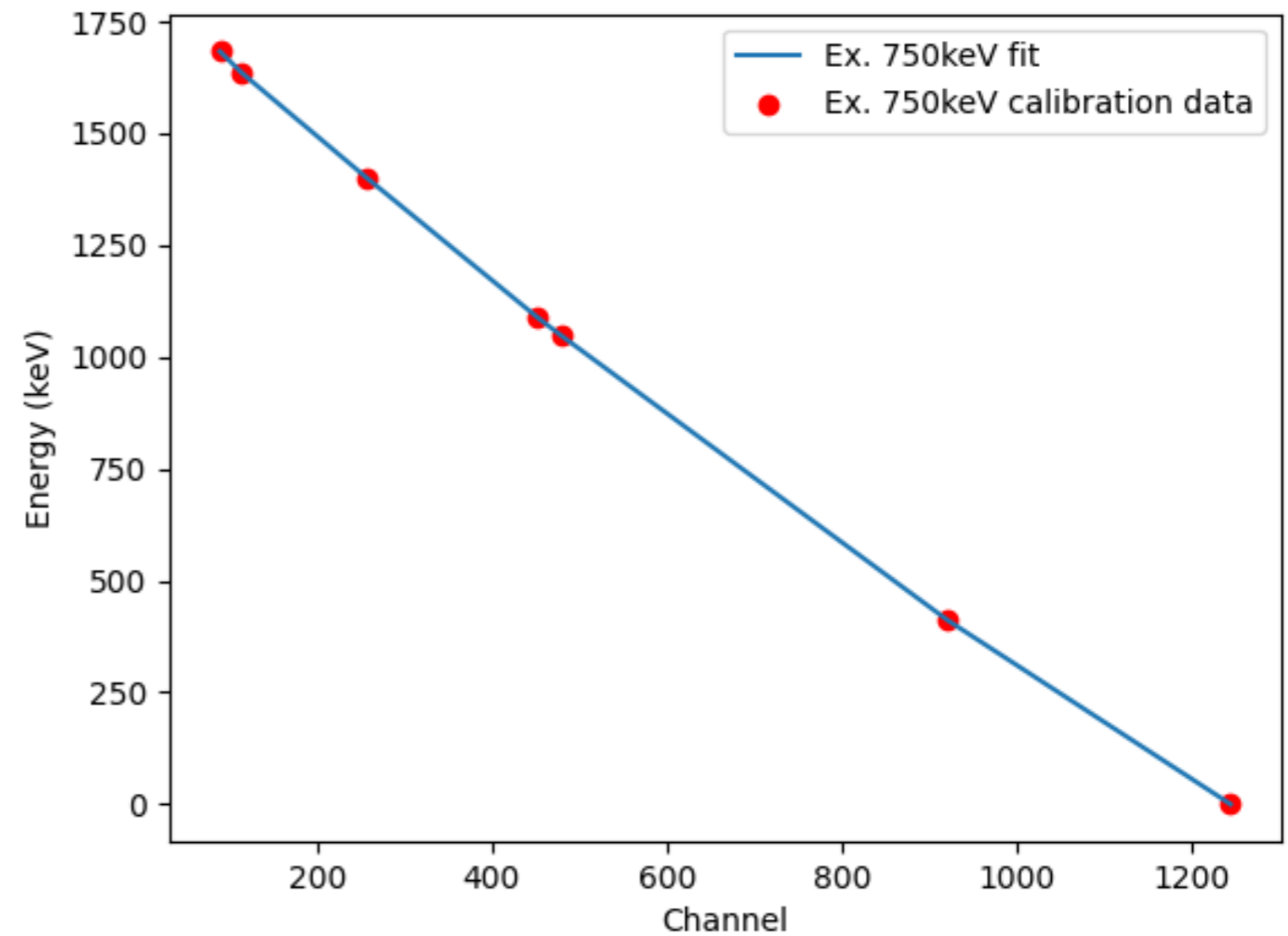
$A=31.5$  mm  
 $B=36.75$  mm  
 $D=354.8$  mm  
 $x_{slit}=0.654$  mm  
 $x - x_{off}=1.346$  mm  
 $y_{slit}=24.5$



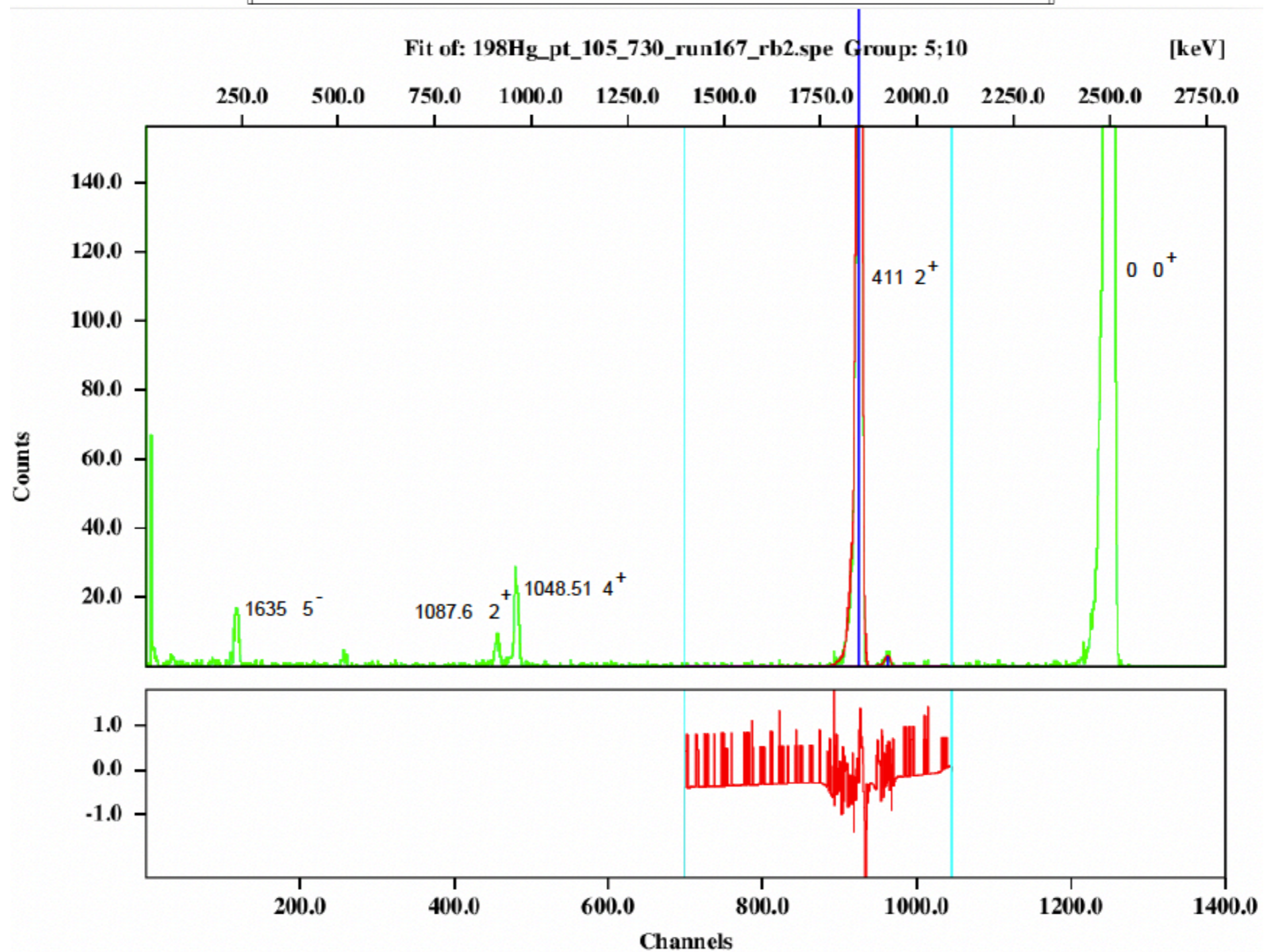
Beginning date of Experiment	Momentum Bite	Number of Runs
January 23, 2014	730, 750	46

~1.27 keV per channel

Position (centroid)/channel	Energy (keV)	$J^\pi$
1245	0	$0^+$
921	411.8025	$2^+$
478.85	1048.51	$4^+$
452.56	1087.6874	$2^+$
256.1615	1401.52	$0^+$
167.49	1550	$0^+$
114	1635	$5^-$
89.02	1683	$7^-$



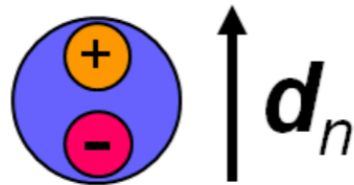
Beginning date of Experiment	Momentum Bite	Number of Runs
January 23, 2014	730, 750	46



$$\begin{aligned}STR_{\text{Coul}} &= M(E\lambda) \\ &= (-1)^{\frac{I-I'+|I-I'|}{2}} \langle I' \| M(E\lambda) \| I \rangle \\ &= (-1)^{\frac{I-I'+|I-I'|}{2}} \sqrt{(2I+1)B(E\lambda; I \rightarrow I')}\end{aligned}$$

$$\begin{aligned}STR_{\text{Nuc}} &= RDEF(\lambda; I \rightarrow I') \\ &= (-1)^{\frac{I-I'+|I-I'|}{2}} \langle I' \| \delta_\lambda \| I \rangle \\ &= (-1)^{\frac{I-I'+|I-I'|}{2}} \sqrt{2I+1} \langle IK\lambda 0 | I'K \rangle \beta_\lambda R_0\end{aligned}$$

What is an EDM?

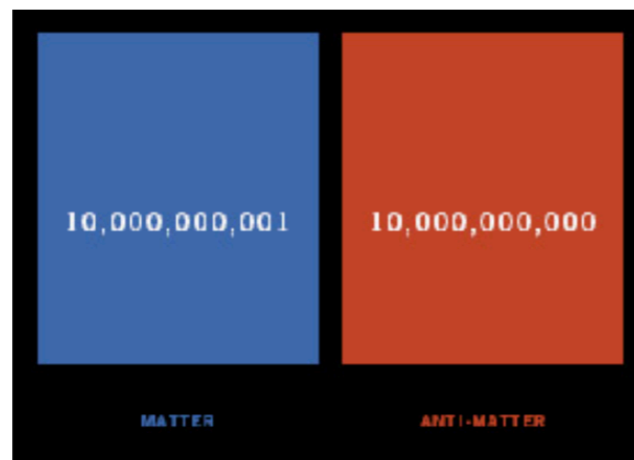


Neutron Electric Dipole  
Moment

A bar magnet is an example of a "magnetic dipole" with a north pole at one end and a south pole at the other. Although the neutron has no net electric charge, it is made of positive and negative quarks. If the positive charges were slightly to one side and the negative charges were slightly to the other, it would create an "electric dipole". The "electric dipole moment", or EDM, is the product of the electric charges and the distance between them. The neutron EDM is usually given in e-cm, the electronic charge times the distance in cm.

[1]

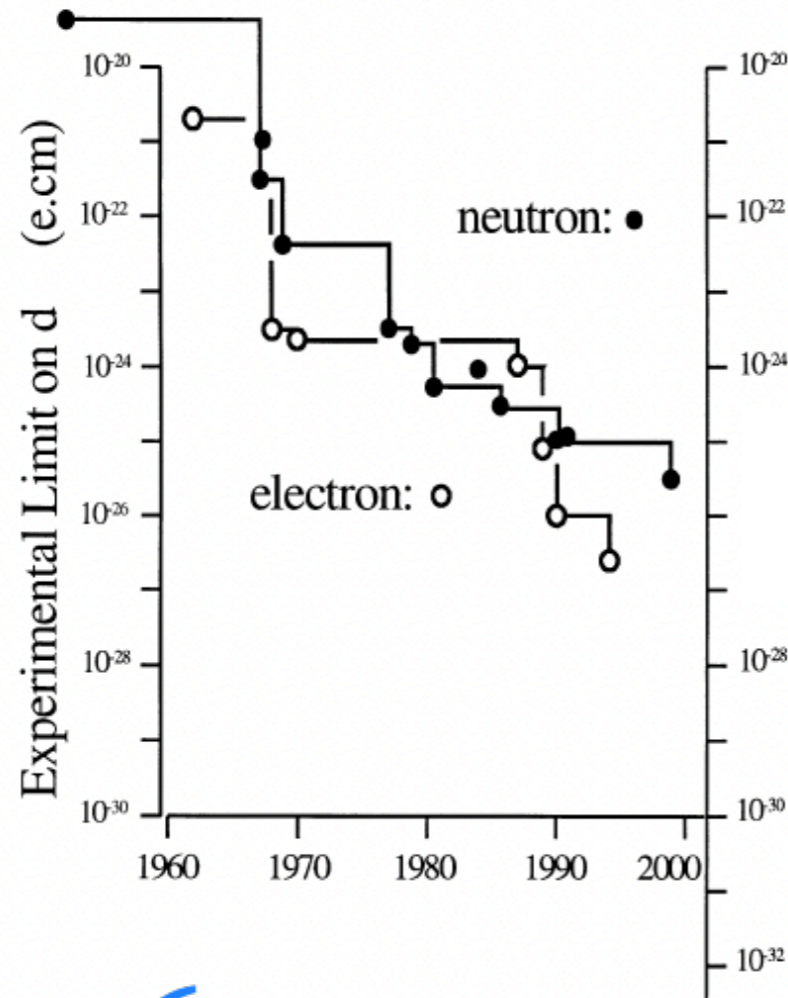
Why measure the Neutron EDM?



Matter and Antimatter in the Early  
Universe



- ▶ Increased electric polarizability: a molecule or an atomic nucleus with octuple collectivity.



SM

Specie	d 95% u.l. (e cm)
Paramagnetic Systems	
$Xe^m$	$3.1 \times 10^{-22}$
Cs	$1.4 \times 10^{-23}$
	$1.2 \times 10^{-25}$
	$2 \times 10^{-5}$
	$2.6 \times 10^{-7} \mu_N RC_s$
Tl	$1.1 \times 10^{-24}$
	$1.9 \times 10^{-27}$
YbF	$1.2 \times 10^{-27}$
ThO	$9.7 \times 10^{-29}$
	$6.4 \times 10^{-9}$
$HfF^+$	$1.6 \times 10^{-28}$
Diamagnetic System	
$^{199}Hg$	$7.4 \times 10^{-30}$
$^{129}Xe$	$6.6 \times 10^{-27}$
$^{225}Ra$	$1.4 \times 10^{-23}$
TlF	$6.5 \times 10^{-23}$
n	$3.6 \times 10^{-26}$
Particle Systems	
$\mu$	$1.8 \times 10^{-19}$
$\tau$	$3.9 \times 10^{-17}$
$\Lambda$	$1.6 \times 10^{-16}$

[2], [3] [4]

Optical potential by *An and Cai* [7]:

$$V(r) = -V f_r(r) - iW_v f_v(r) + i4a_s W_s \frac{df_s(r)}{dr} + \lambda_{\pi}^2 \frac{V_{so} + W_{so}}{r} \frac{df_{so}(r)}{dr} \vec{\sigma} \cdot \vec{l} + V_C(r)$$

Saxon-Woods form factor:  $f_i(r) = \left\{ 1 + \exp\left[\frac{(r - r_i A^{1/3})}{a_i}\right] \right\}^{-1}$

Parameter	An and Cai [1]	Bojowald et al. [2]	Deehnick, Childs and Vrcelj [3]	Han, Shi and Shen [4]
$r_C$ (fm)	1.303000	1.300000	1.3000	1.6980
$V_r$ (MeV)	95.240100	95.667800	94.8586	86.336100
$r_r$ (fm)	1.150670	1.180000	1.170000	1.174000
$a_r$ (fm)	0.792439	0.839997	0.746400	0.809000
$W_V$ (MeV)	2.472400	0.000000	0.603400	3.043564
$r_V$ (fm)	1.322100	-	1.325000	1.563000
$a_V$ (fm)	0.272200	-	0.937993	0.962281
$W_S$ (MeV)	10.156800	13.861600	12.168600	10.892290
$r_S$ (fm)	1.360080	1.270000	1.325000	1.328000
$a_S$ (fm)	0.892366	0.890398	0.937993	0.727281
$V_{SO}$ (MeV)	1.778500	3.000000	3.346000	1.851500
$r_{SO}$ (fm)	0.972000	1.001480	1.070000	1.234000
$a_{SO}$ (fm)	1.011000	1.001480	0.660000	0.813000
$W_{SO}$ (MeV)	0.000000	0.000000	0.000000	-0.103000

appropriate to express the combination of contributions to a measured atomic EDM for paramagnetic systems, diamagnetic systems, and nucleons as

$$d_i = \sum_j \alpha_{ij} C_j, \quad (9)$$

where  $i$  labels the system, and  $j$  labels the specific low-energy parameter (*e.g.*,  $d_e$ ,  $C_S$ , *etc.*). The  $\alpha_{ij} = \partial d_i / \partial C_j$

imental results. In the global analysis, paramagnetic systems are used to set limits on the electron EDM  $d_e$  and the nuclear spin-independent electron-nucleus coupling  $C_S$ . Diamagnetic systems set limits on four dominant parameters: two pion-nucleon couplings ( $\bar{g}_\pi^{(0)}$ ,  $\bar{g}_\pi^{(1)}$ ), a specific isospin combination of nuclear spin-dependent couplings, and the “short distance” contribution to the

- ▶ Search for EDM in Atoms with octupole-deformed nuclei

- ▶ EDM = 0 if there is invariance under P and T

- ▶ **d same orientation of I angular momentum,** d must change sign the same way under P and T

$$\mathbf{d} = q\mathbf{r}$$

$$\mathbf{I} = \mathbf{r} \times \mathbf{p}$$

$$\left\{ \begin{array}{l} \mathbf{d} \xrightarrow{P} -\mathbf{d} \\ \mathbf{I} \xrightarrow{P} \mathbf{I} \\ \mathbf{d} \xrightarrow{T} \mathbf{d} \\ \mathbf{I} \xrightarrow{T} -\mathbf{I} \end{array} \right.$$

[1]

All particle and atomic EDMs are odd under P and T symmetry

- The observable EDM in the electron cloud is induced by the Schiff Moment of the nucleus.

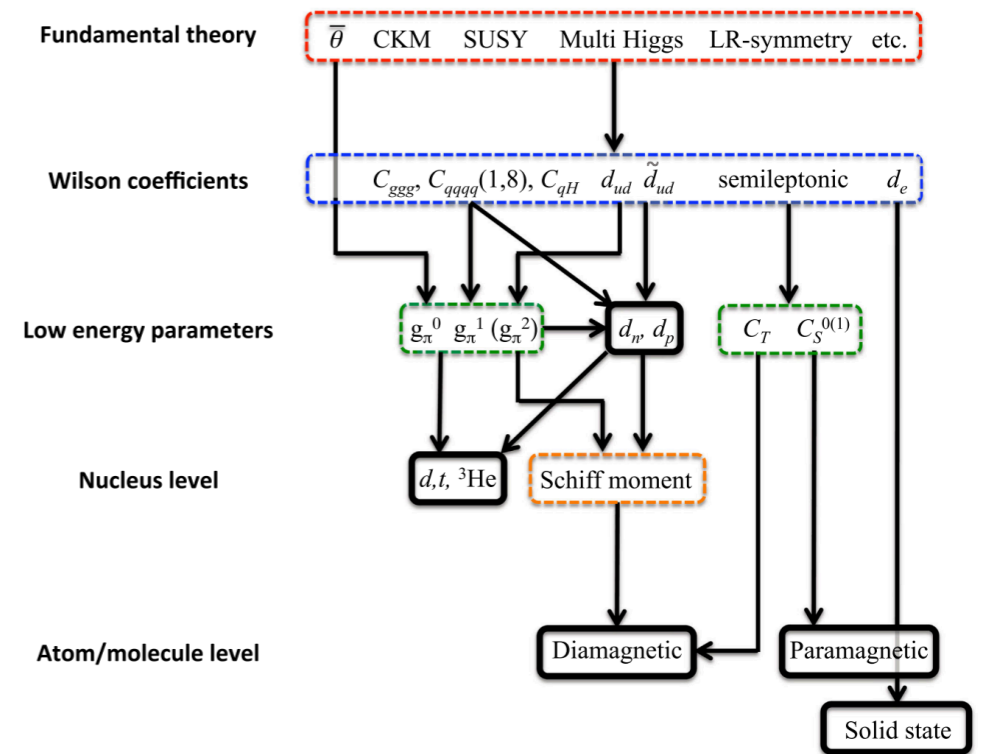
Schiff Moment  $\rightarrow \vec{S}_{int} = \frac{1}{10}[\vec{O}_0 - \frac{5}{3}d_0 \langle r^2 \rangle_{ch}]$

- Large nuclear gs EDMs (large Schiff moments)  $\rightarrow$  The action of a P- and T- violating  $\hat{V}_{PT}$

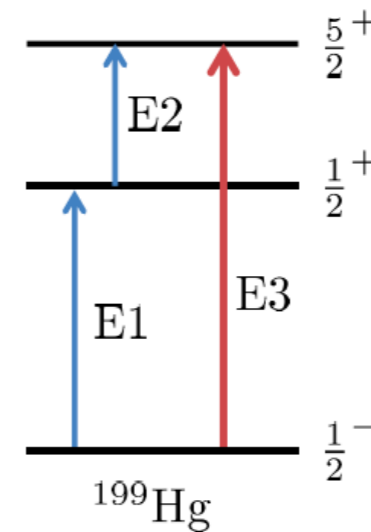
$$S \equiv \langle S_z \rangle = \sum_{i \neq 0} \frac{\langle \Psi_0 | S_z | \Psi_i \rangle \langle \Psi_0 | \hat{V}_{PT} | \Psi_i \rangle}{\Delta E} + \dots$$

In the laboratory frame:  $S \approx -2 \frac{I}{I+1} \hat{S}_{int} \frac{\langle \Psi^+ | \hat{V}_{PT} | \Psi^- \rangle}{\Delta E}$

- The most direct ways to measure E2 and E3 matrix  $\rightarrow$  Inelastic scattering.



[4]



[3]

▶ WEAK INTERACTION VIOLATES C, P and T transformations.

▶ **Charge Conjugation**  $\longrightarrow$

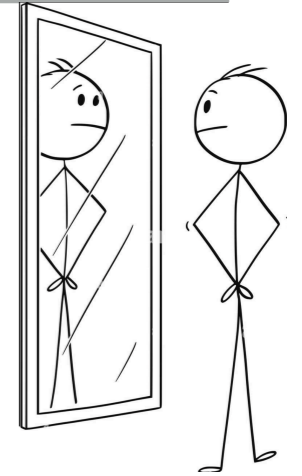
$$|p\rangle \xrightarrow{C} (-1)^{1/2+1/2} |\bar{p}\rangle = - |\bar{p}\rangle$$
$$|n\rangle \xrightarrow{C} (-1)^{1/2+1/2} |\bar{n}\rangle = + |\bar{n}\rangle$$

▶ **CP violation**  $\longrightarrow$  First observed in the decay of  $K^0$  meson.

$$K_L^0 \rightarrow \begin{cases} \pi^+ + e^- + \bar{\nu}_e \\ \pi^- + e^+ + \nu_e \end{cases}$$

$$CP |K^0\rangle = |\bar{K}^0\rangle$$

CP violation - The standard model of elementary particle suggests that when the universe was less than 10-12 sec old, the condition was ripe for the production of more matter than antimatter with CP violation to provide the mechanism for different reaction rate (to produce matter and antimatter).

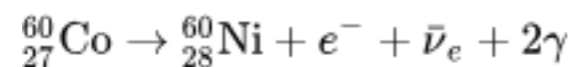
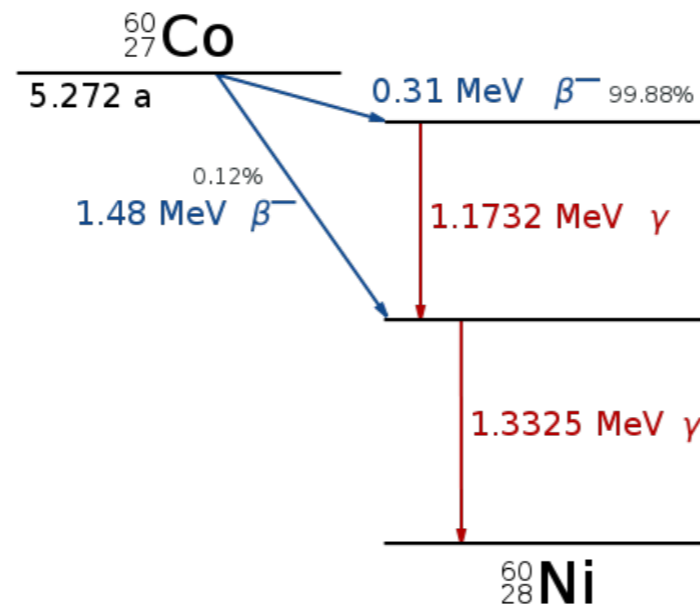


▶ WEAK INTERACTION VIOLATES C, P and T transformations.

▶ Parity  $\longrightarrow (x, y, z) \xrightarrow{P} (-x, -y, -z)$

Violation of Parity  $\longrightarrow K^+ \rightarrow \begin{cases} \pi^+ + \pi^0 \\ \pi^+ + \pi^- + \pi^+ \end{cases}$

$\beta^-$  decay of  $^{60}\text{Co}$   $\longrightarrow$  First evidence that confirmed parity nonconservation.



From [commons.wikimedia.org](https://commons.wikimedia.org)

## Reduced transition probabilities in terms of intrinsic moments

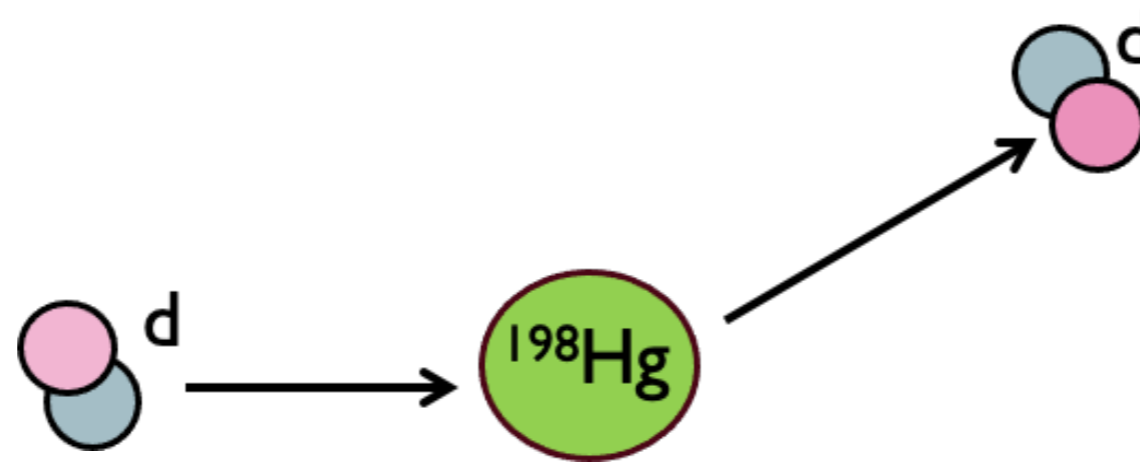
▶ E1 transition →  $B(E1; I_i \rightarrow I_f) = \frac{3}{4\pi} Q_1^2 (I_i K_i 10 | I_f K_f)^2$

▶  $\gamma$  -ray studies

▶ Higher order transitions →  $B(E\lambda; I_i \rightarrow I_f) = \frac{1}{2I_i + 1} \langle I_i || M(E\lambda) || I_f K_f \rangle$   
 $= \frac{2\lambda + 1}{16\pi} Q_\lambda^2 (I_i K_i \lambda 0 | I_f K_f)^2, \quad Q_\lambda = \frac{3}{\sqrt{2\lambda + 1} \pi Z R_0^\lambda \beta_\lambda}$

▶ E2 and E3 transitions → Inelastic scattering, Coulomb excitation





$$\cos \theta_{cm} = \cos \theta_{lab} [x \cos \theta_{lab} + \sqrt{1 - x^2 \sin^2 \theta_{lab}}] - x$$

Elastic scattering

$$x = \frac{m_p}{m_t}$$

Inelastic scattering

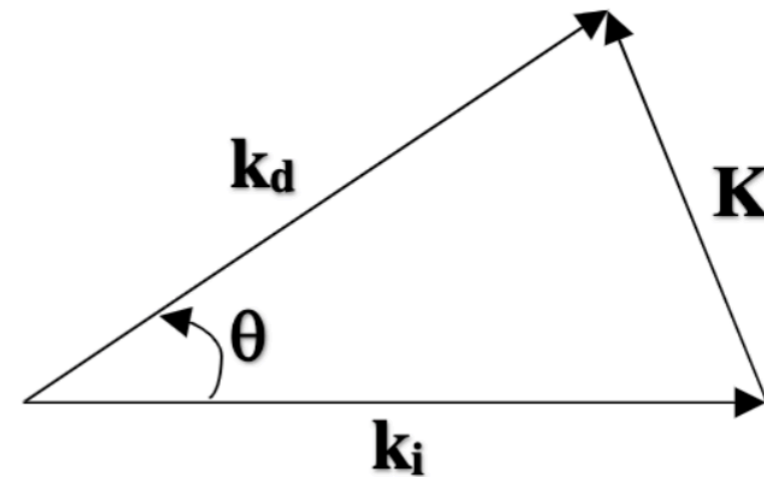
$$x = \left[ \frac{m_a m_b}{m_A m_B} \frac{1}{1 + \frac{Q}{E_{cm}}} \right]^{1/2}$$

For spherically symmetric potential:

$$f(\theta) \cong -\frac{2m}{\hbar^2 K} \int_0^\infty dr r_o V(r_o) \sin(Kr_o)$$

↙

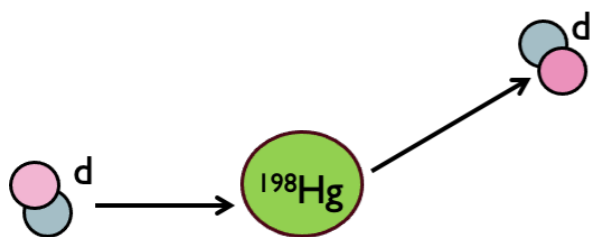
$$f(\theta) = -\frac{2mV_o}{\hbar^2 \mu(\mu^2 + K^2)}$$



In the limit  $\mu \rightarrow 0$ , the Yukawa potential tends to the Coulomb potential  $V(r) \sim \frac{Z_1 Z_2 e^2}{r}$

$$\frac{d\sigma}{d\Omega} = \frac{Z_1 Z_2 e^4}{16E^2 \sin^4 \frac{\theta}{2}}$$

In MKSA units:



$$\begin{aligned} \left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}} &= \frac{d\sigma}{d\Omega} = \frac{1}{4\pi\epsilon_o} \frac{Z_1 Z_2 e^4}{16E^2 \sin^4 \frac{\theta}{2}} \\ &= 17.408 \frac{1}{\sin^4 \frac{\theta}{2}} \text{ mb/sr.} \end{aligned}$$

# TARGET THICKNESS CORRECTION

Beginning date of Experiment	Momentum Bite	Number of Runs
December 17, 2016	700	14

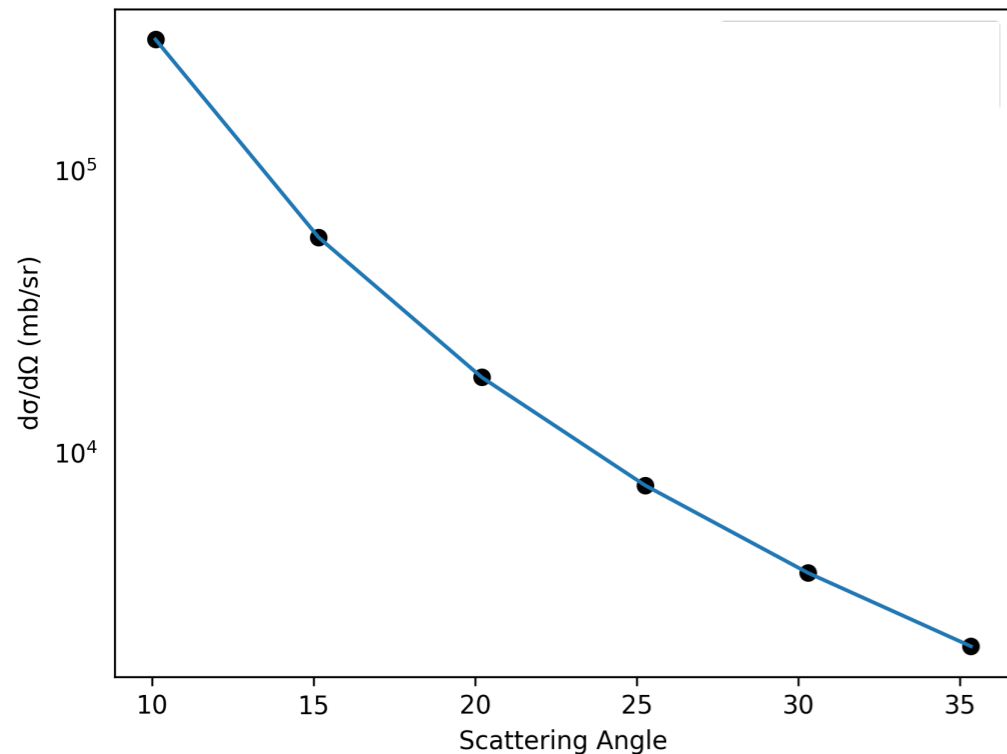
$$N_{t, \text{eff}} = \frac{N_t}{\cos \theta_{\text{target}}} [\text{mb}^{-1}]$$

$$N_t = \frac{\rho t}{M} N_A$$

Beginning date of Experiment	Nominal Target Thickness $\rho t$ ( $\mu\text{g}/\text{cm}^2$ )
January 23, 2014	95
December 17, 2016	70
July 24, 2017	40

Beginning date of Experiment	Momentum Bite	Number of Runs
December 17, 2016	700	14

$$N_{t,eff} = 0.980639890 \left( \frac{N_c}{\left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}} N_b \Omega_{LT} \text{DAQ}_{LT} \text{DET} \left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}}} \right)$$



## RUTHERFORD CROSS SECTION-FRESCO:

$$\frac{\left(\frac{d\sigma}{d\Omega}\right)_{\text{DWBA}}}{\left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}}} = 1.065.$$

# DETECTOR

