

# Precision theory for charge radii of light nuclei

In collaboration with: [Arseniy Filin](#), Daniel Möller, Vadim Baru, Christopher Körber, Hermann Krebs, Andreas Nogga and Patrick Reinert



In memory of  
Ruprecht Machleidt

- **Fundamental approaches were either not fundamental or not quantitative.**
- **In short: The more quantitative, the less fundamental. The more fundamental, the less quantitative.**

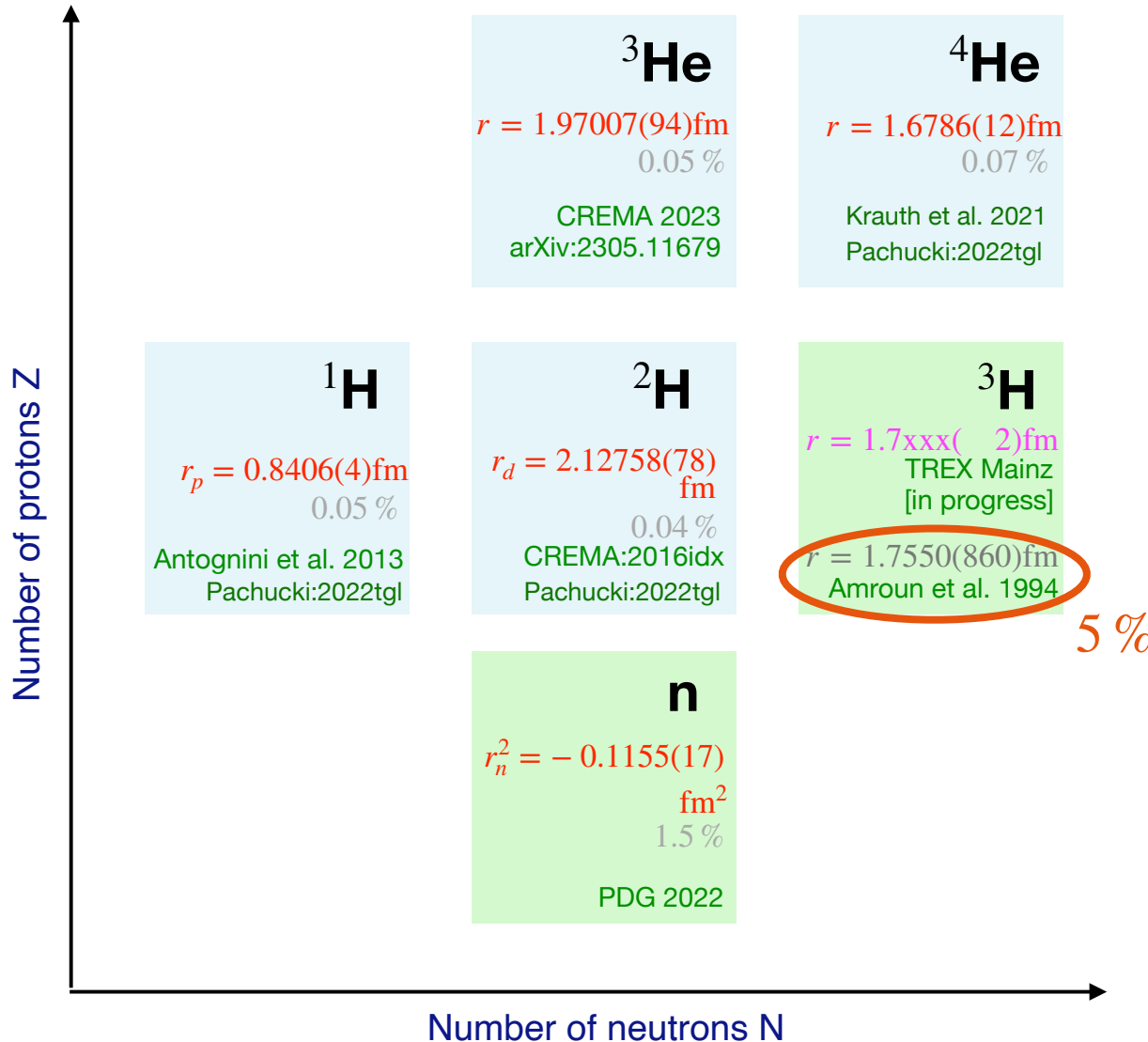
We need ... **a fundamental approach that produces a quantitative description of the nuclear force.**

R. Machleidt

Review: NN and Many-Body Int.  
SURA Workshop, DC, 16-Oct-06

12

# Experimental data



Pohl:2013yb,  
CREMA:2016idx,  
Pohl:2016glp,  
Schmidt:2018kjc,  
Krauth:2021foz  
...

Data from isotope shift

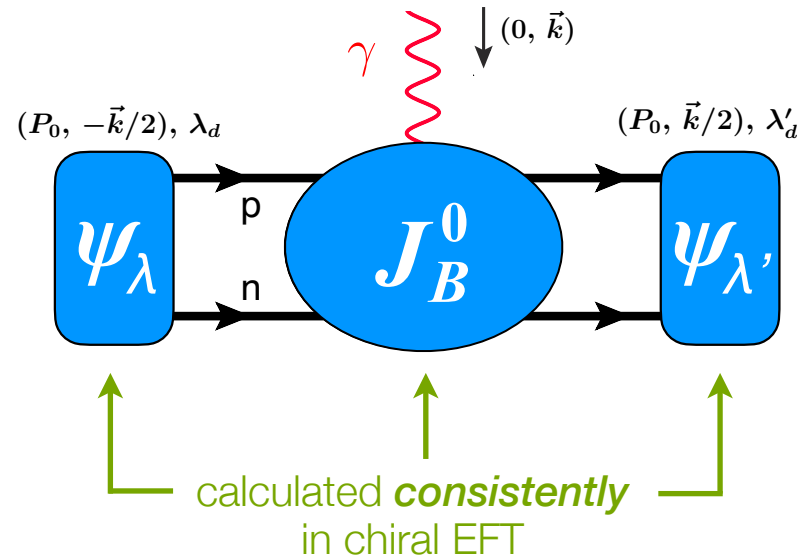
$$(r_d^2 - r_p^2) = 3.82070(31)\text{fm}^2$$

0.01 %
















Pachucki et al. 2018  
Jentschura et al. 2011

# Theory in a nutshell

- **Chiral EFT** for the nuclear Hamiltonian  $H$  and  $J^\mu$
- Use **1N FFs** to avoid reliance on  $\chi$ EFT for  $J_{1N}^\mu$  (i.e., re-summations)
- **Error analysis** (statistical uncertainties of few-N LECs,  $\pi$ N LECs, EFT truncation, parametrizations of the 1N FFs)
- **Regularization and symmetries (3NF, MECs)**
  - *gradient flow method (talk by Hermann)*
- Starting from N<sup>3</sup>LO (Q<sup>4</sup>), one has to worry about:
  - relativistic corrections
  - isospin breaking effects (including neutron-neutron interaction)
  - electromagnetic interactions beyond Coulomb
- **Semi-analytical results for convolution integrals** (to minimize numerical errors)



# The Hamiltonian

	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO ( $Q^0$ )			
NLO ( $Q^2$ )			
N <sup>2</sup> LO ( $Q^3$ )			
N <sup>3</sup> LO ( $Q^4$ )			
N <sup>4</sup> LO ( $Q^5$ )			

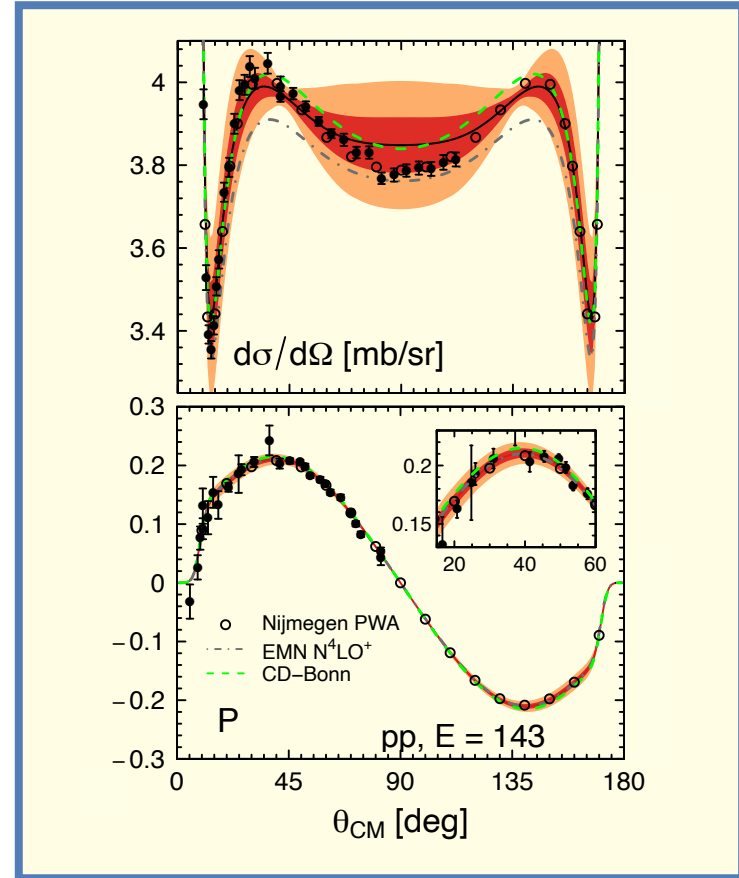
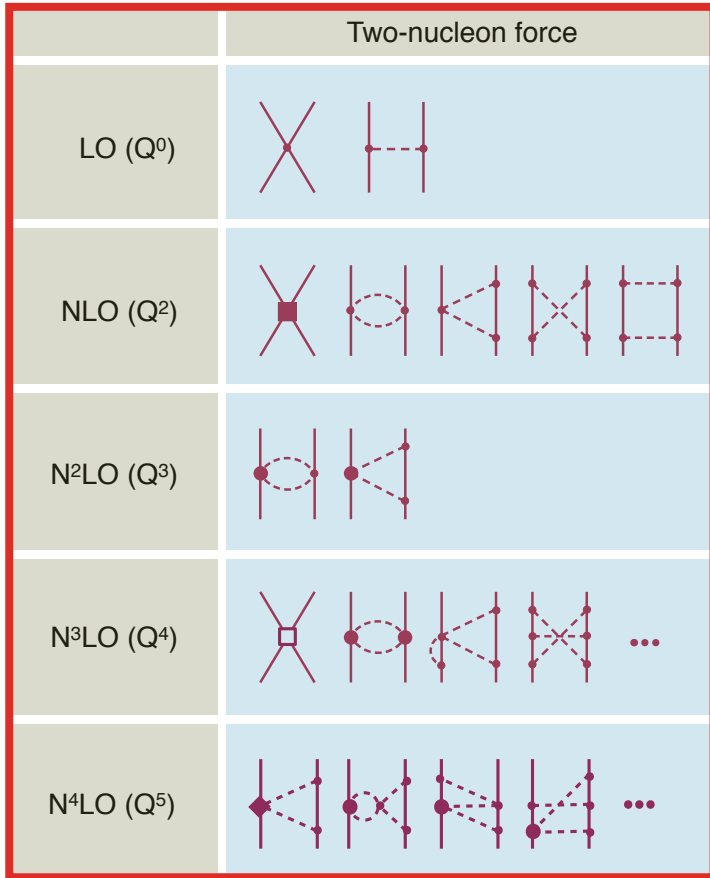
# The Hamiltonian

	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO ( $Q^0$ )		—	—
NLO ( $Q^2$ )		—	—
N <sup>2</sup> LO ( $Q^3$ )			—
N <sup>3</sup> LO ( $Q^4$ )			
N <sup>4</sup> LO ( $Q^5$ )			—

mixing DimReg with Cutoff regularization violates  $\chi$ -symmetry  
 $\Rightarrow$  re-derive using Gradient Flow regulator

— talk by Hermann —

# The Hamiltonian



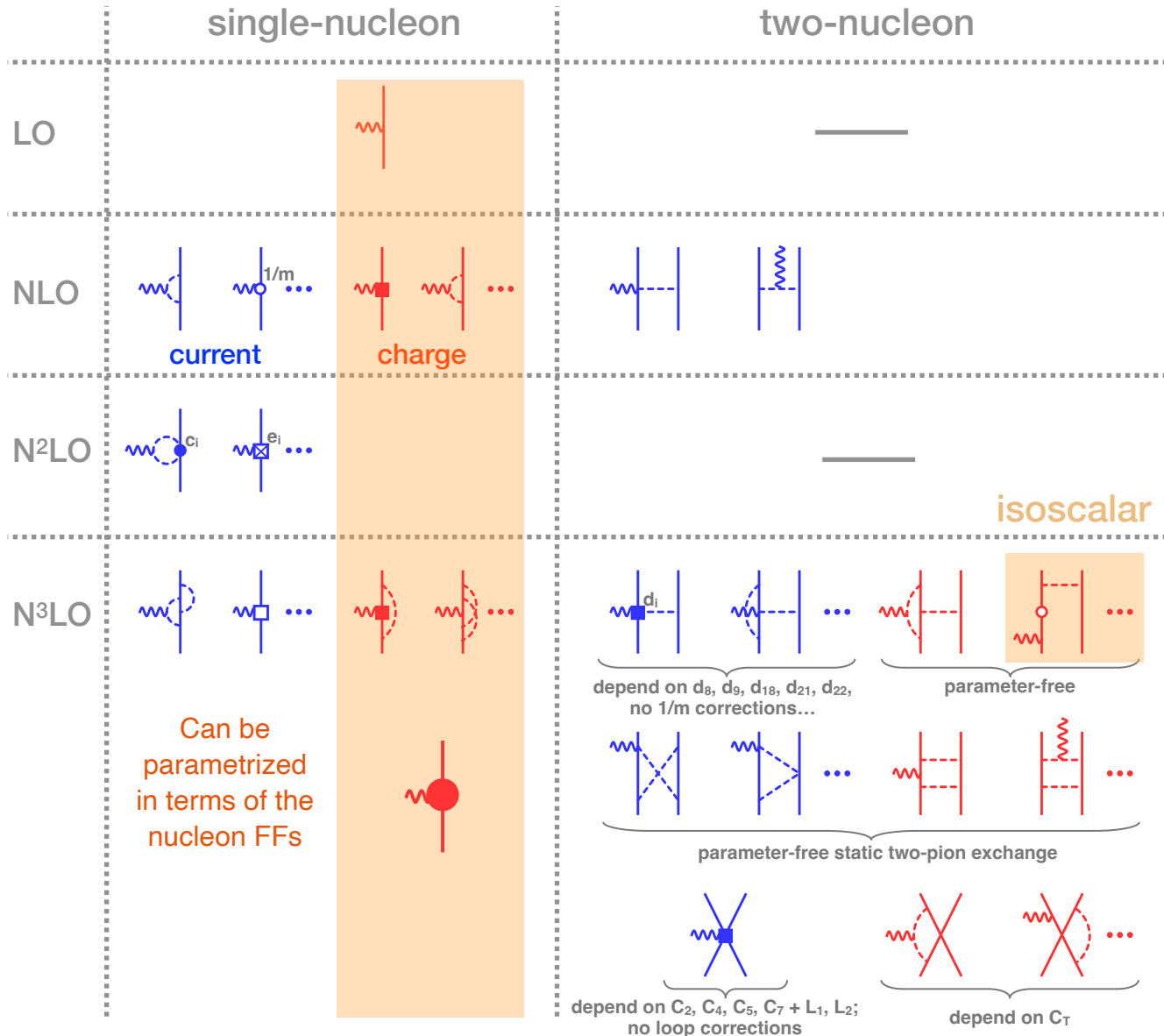
The newest Bochum NN interactions [Reinert, Krebs, EE, EPJA 54 \(2018\) 86](#); [PRL 126 \(2021\) 092501](#)

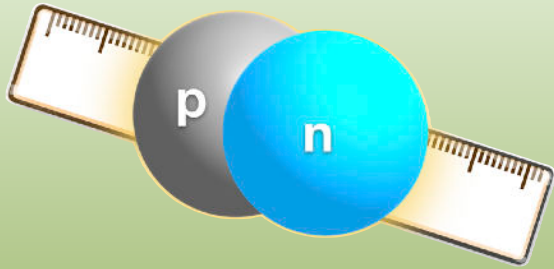
$$V_{1\pi}(q) = \frac{\alpha}{\vec{q}^2 + M_\pi^2} e^{-\frac{\vec{q}^2 + M_\pi^2}{\Lambda^2}} + \text{subtraction}, \quad V_{2\pi}(q) = \frac{2}{\pi} \int_{2M_\pi}^{\infty} d\mu \mu \frac{\rho(\mu)}{\vec{q}^2 + \mu^2} e^{-\frac{\vec{q}^2 + \mu^2}{2\Lambda^2}} + \text{subtractions}$$

+ nonlocal (Gaussian) cutoff for contacts

# Electromagnetic currents

Kölling, EE, Krebs, Meißner, PRC 80 (09) 045502; PRC 86 (12) 047001; Krebs, EE, Meißner, FBS 60 (2019) 31





# The deuteron ( $A = 2$ )

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Patrick Reinert, Phys. Rev. Lett. 124 (2020) 082501;

Phys. Rev. C103 (2021) 024313

$$\rho_{1N}^{\text{DF}} = -e \frac{\mathbf{k}^2}{8m_N^2} G_E(\mathbf{k}^2)$$

$$G(Q^2) = G^{\text{Main}}(Q^2) + G^{\text{DF}}(Q^2) + G^{\text{SO}}(Q^2) + G^{\text{Boost}}(Q^2) + G^{1\pi}(Q^2) + G^{\text{Cont}}(Q^2)$$

$$\rho_{1N}^{\text{Main}} = e G_E(\mathbf{k}^2)$$

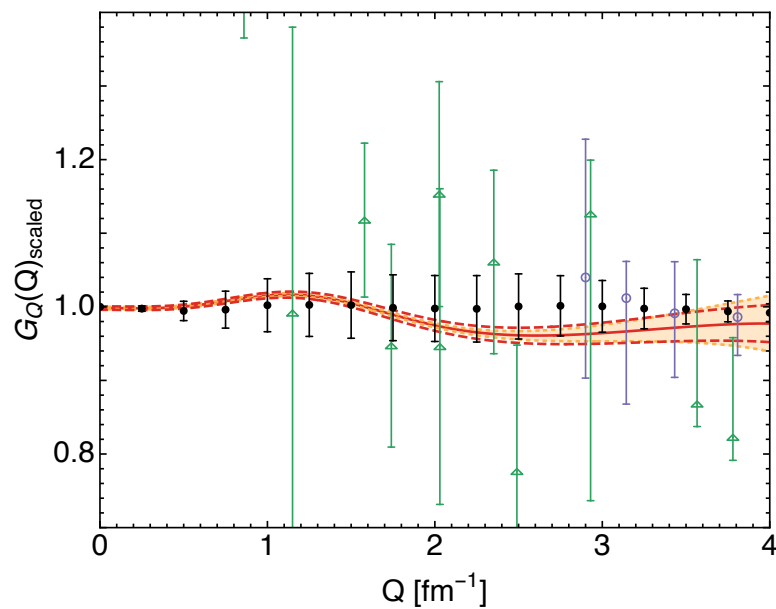
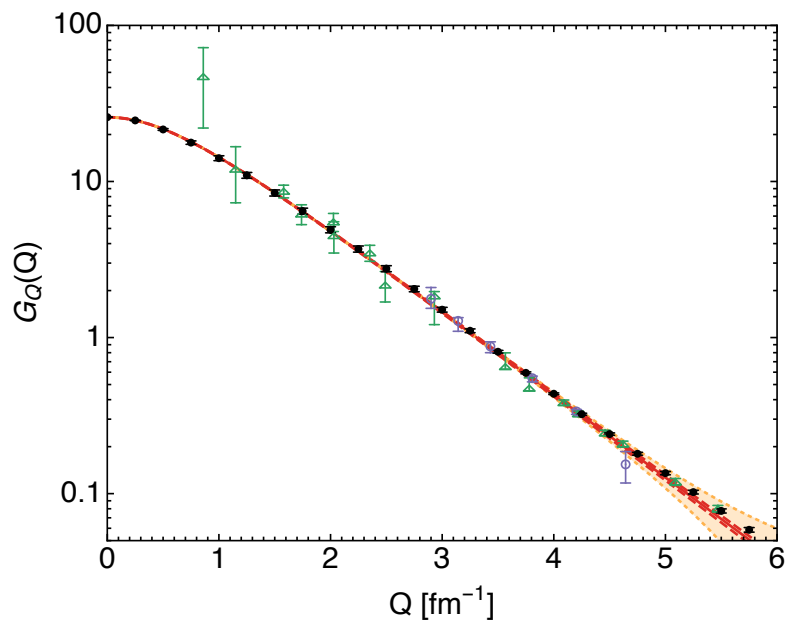
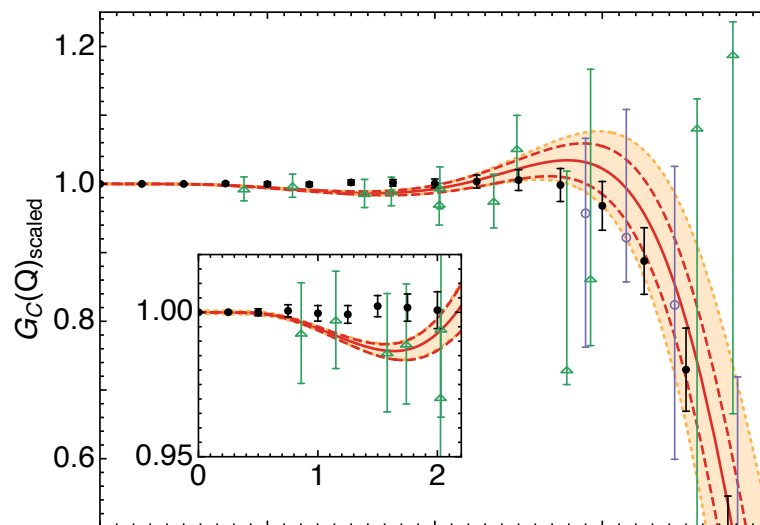
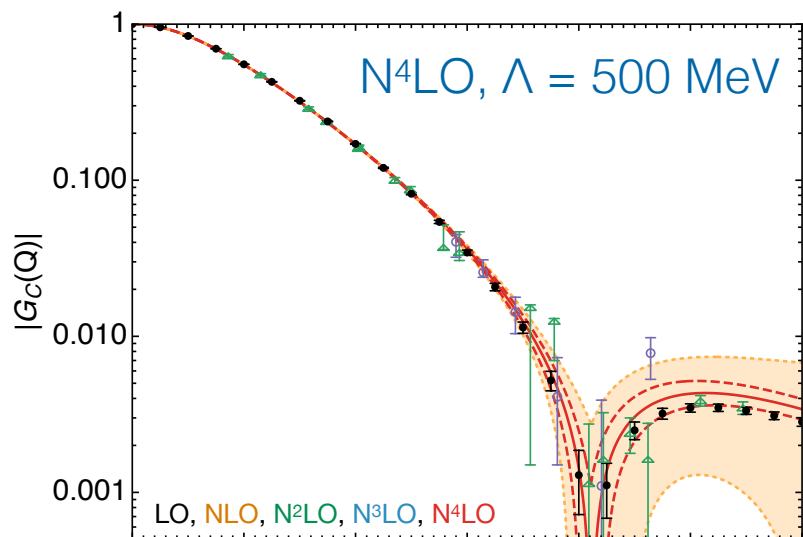
$$\rho_{1N}^{\text{SO}} = ie \frac{2G_M(\mathbf{k}^2) - G_E(\mathbf{k}^2)}{4m_N^2} \boldsymbol{\sigma} \cdot \mathbf{k} \times \mathbf{p}$$

- Both the nuclear force and the 2N charge density are available to N<sup>4</sup>LO
- Simple numerics



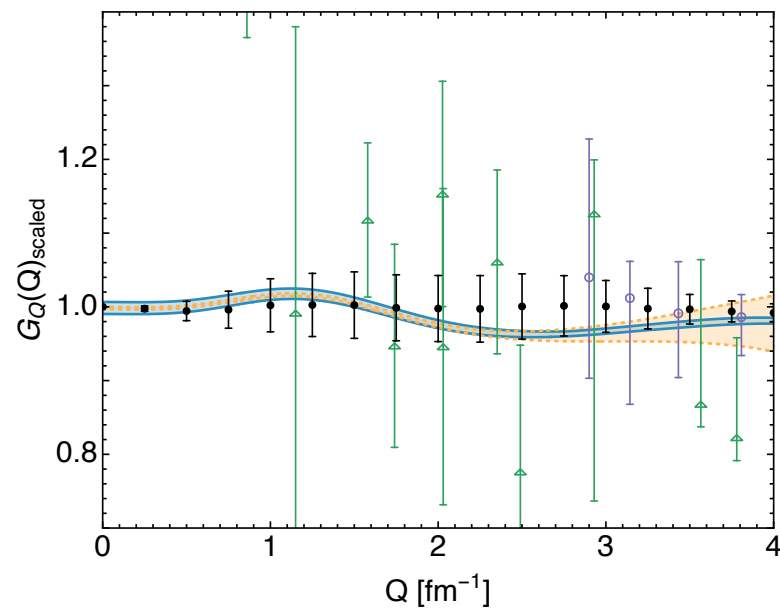
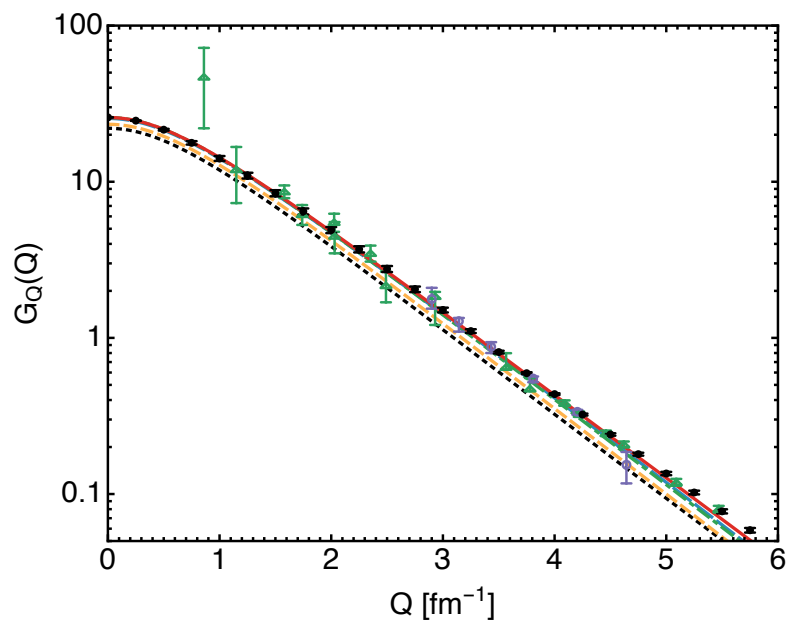
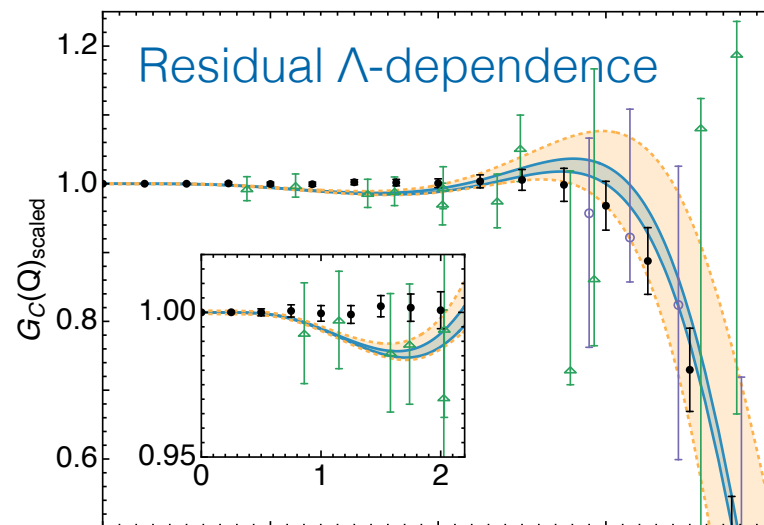
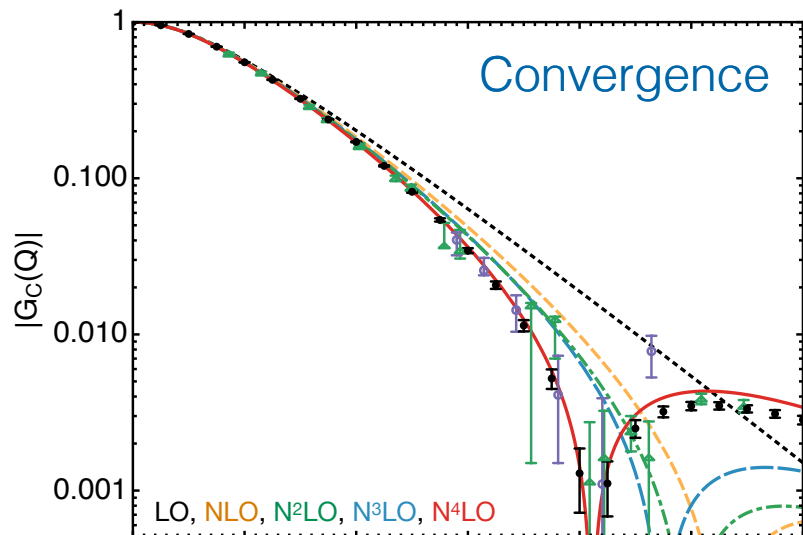
# The charge and quadrupole FFs of ${}^2\text{H}$

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Patrick Reinert, PRL 124 (2020) 082501; PRC 103 (2021) 024313



# The charge and quadrupole FFs of ${}^2\text{H}$

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# Charge radius and quadrupole moment

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Patrick Reinert, PRL 124 (2020) 082501; PRC 103 (2021) 024313

Deuteron charge *and structure* radii:  $r_d^2 = r_{str}^2 + r_p^2 + r_n^2 + \frac{3}{4m_p^2}$

EFT truncation, choice of fitting range, NN,  $\pi$ N and  $\gamma$ NN LECs

Our results:  $r_{str} = 1.9729^{+0.0015}_{-0.0012}$  fm,  $Q_d = 0.2854^{+0.0038}_{-0.0017}$  fm<sup>2</sup>

$Q_d^{exp} = 0.285\,699(15)(18)$  fm<sup>2</sup> Puchalski et al., PRL 125 (2020)

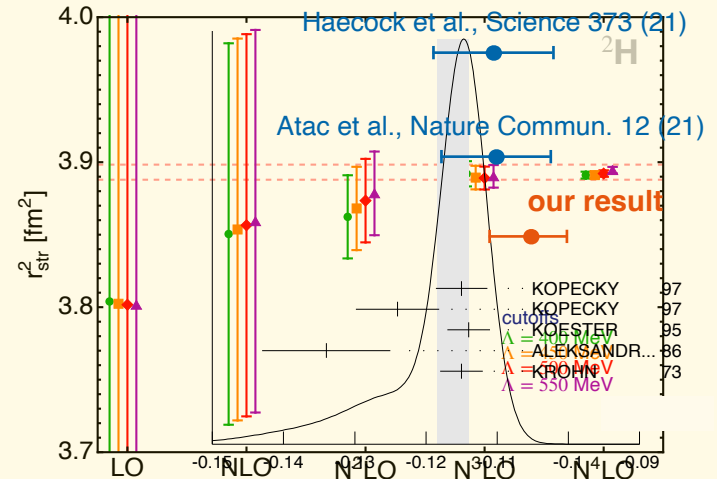
Error budget:

	central	truncation	$\rho_{Cont}^{reg}$	$\pi$ N LECs RSA	2N and $\pi$ N LECs	Q-range	total
$r_{str}^2$ [fm <sup>2</sup> ]	3.8925	$\pm 0.0030$	$\pm 0.0024$	$\pm 0.0003$	$\pm 0.0025$	$+0.0035$ $-0.0005$	$+0.0058$ $-0.0046$
$Q_d$ [fm <sup>2</sup> ]	0.2854	$\pm 0.0005$	$\pm 0.0007$	$\pm 0.0003$	$\pm 0.0016$	$+0.0035$ $-0.0005$	$+0.0038$ $-0.0017$

Combining our result for  $r_{str}^2$  with the <sup>1</sup>H-<sup>2</sup>H isotope shift datum  $r_d^2 - r_p^2 = 3.82070(31)$  fm<sup>2</sup> Jentschura et al., PRA 83 (2011)

leads to the prediction for the neutron radius:  
(residual cutoff dependence):

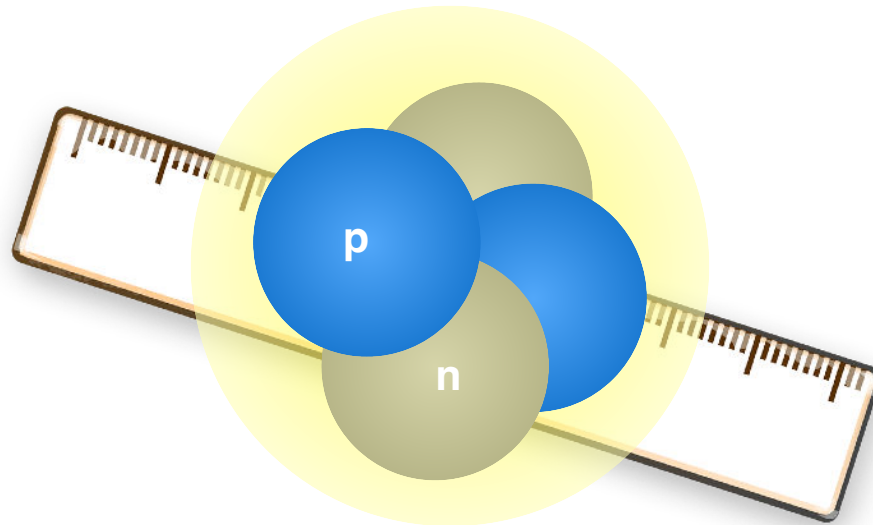
$r_n^2 = -0.105^{+0.005}_{-0.006}$  fm<sup>2</sup>



# The charge FF of ${}^4\text{He}$ ( $A = 4$ )

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

3NF beyond  $N^2\text{LO}$  not yet available (However, no sensitivity to 3NF once the BEs are reproduced...)



# Relativistic effects

- Relativistic corrections start appearing in  $V_{2N}$ ,  $V_{3N}$  and  $\rho_{2N}$  at  $N^3\text{LO}$
- Boosting  ${}^2\text{H}$ ,  ${}^3\text{H}(e)$ ,  ${}^4\text{He}$  to the Breit frame straightforward (effect decreases with  $A$ )
- Subtleties related to the form of the Schrödinger equation

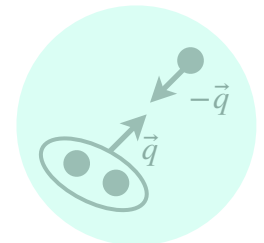
For 2N:  $\underbrace{\left(2\sqrt{\hat{p}^2 + m^2} + \tilde{V}\right) \Psi = \underbrace{2\sqrt{k^2 + m^2}}_{E_{\text{CMS}}} \Psi}_{\text{equation to be used for } A > 2} \Rightarrow \underbrace{\left(\frac{\hat{p}^2}{m} + \tilde{V}\right) \Psi = \frac{k^2}{m} \Psi}_{\text{equation we solve for } A = 2}$

$\swarrow$   $\chi\text{EFT NN potential}$

Relationship between  $\tilde{V}$  and  $V$ :  $\underbrace{2\left\{\sqrt{\hat{p}^2 + m^2}, \tilde{V}\right\} + \tilde{V}^2}_{\text{can be solved by iterations to obtain } \tilde{V}} = 4mV$

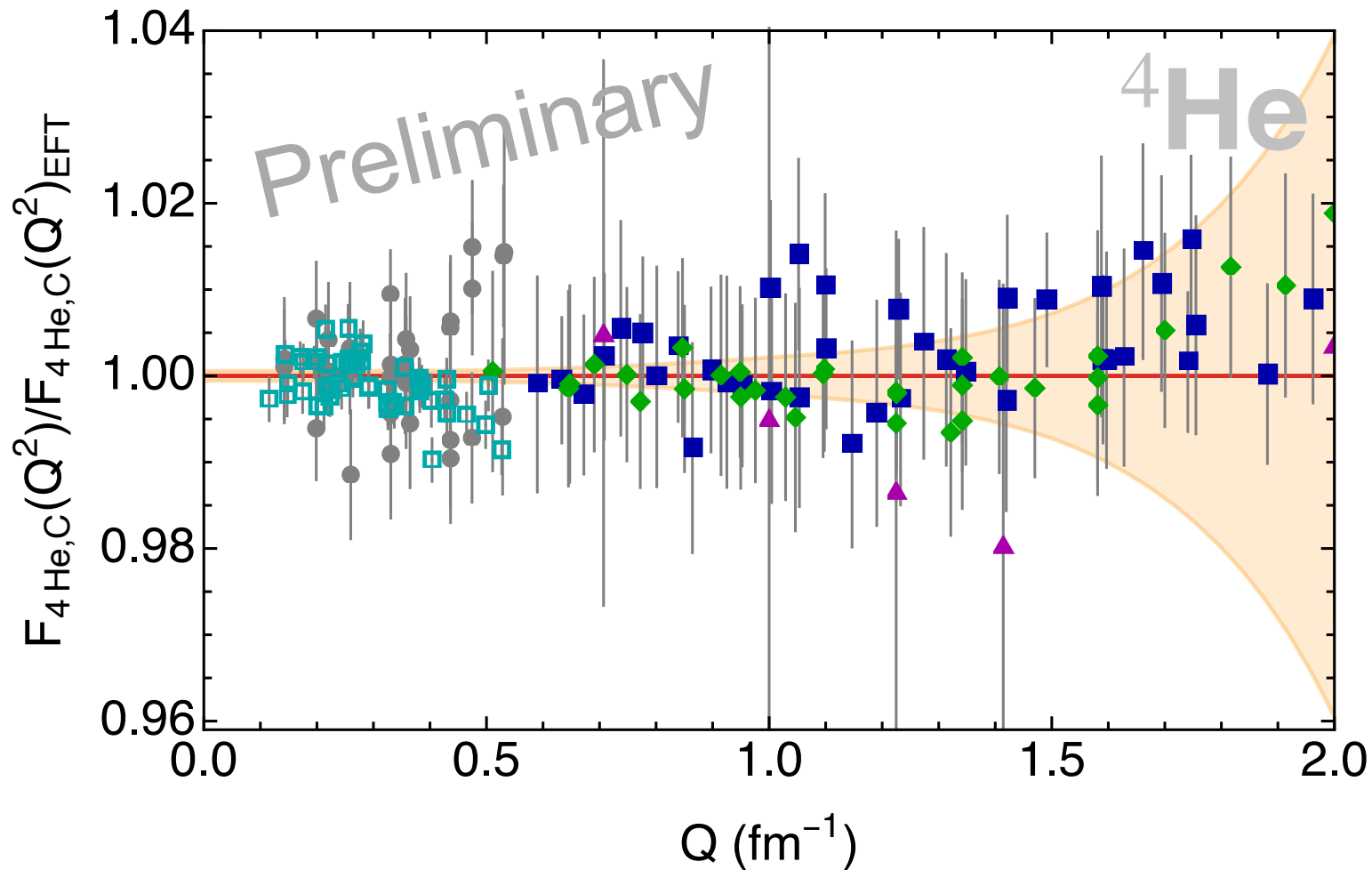
- Boosted potentials from the dynamical mass operator Polyzou et al., Few Body Syst. 49 (2011)

$$\tilde{V}_q := \sqrt{\underbrace{\left(2\sqrt{\hat{p}^2 + m^2} + \tilde{V}\right)^2}_{\hat{M}_{12}(\tilde{V})} + q^2} - \sqrt{\underbrace{\left(2\sqrt{\hat{p}^2 + m^2}\right)^2}_{\hat{M}_{12}^0} + q^2}$$



- Relativistic Faddeev/FY equations Witala et al. '09; Kamada et al. '19; Kamada '20; Hadizadeh et al. '20

# The charge FFs of $^4\text{He}$

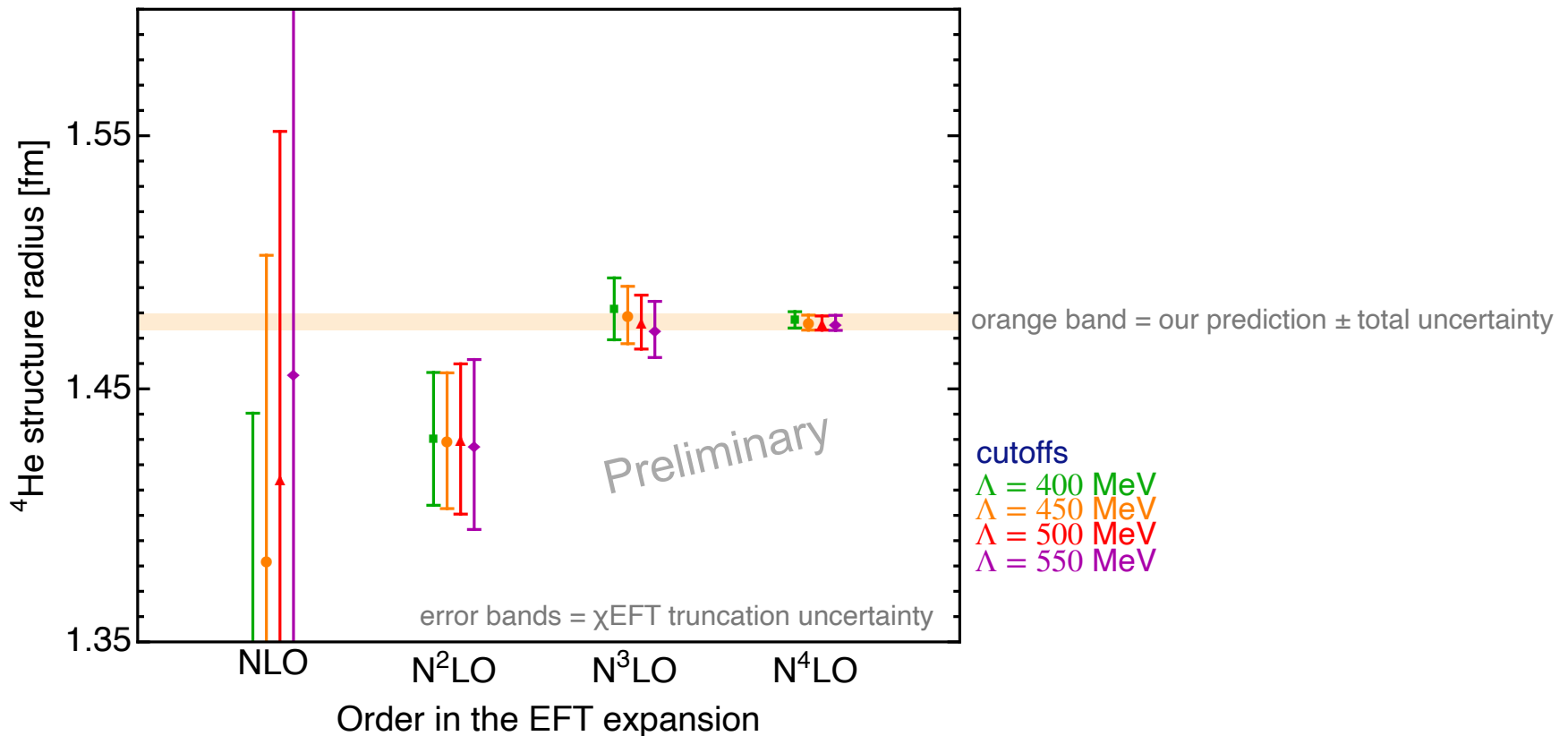


# The structure radius of ${}^4\text{He}$

Preliminary result for the  ${}^4\text{He}$  structure radius:

$$r_{\text{str}}({}^4\text{He}) = 1.47\text{xx} \pm 0.0028_{\text{trunc}} \pm 0.0011_{\text{stat}} \pm 0.0010_{\text{nucl-FF}} \text{ fm (Preliminary)}$$

Consistency check (residual cutoff dependence):



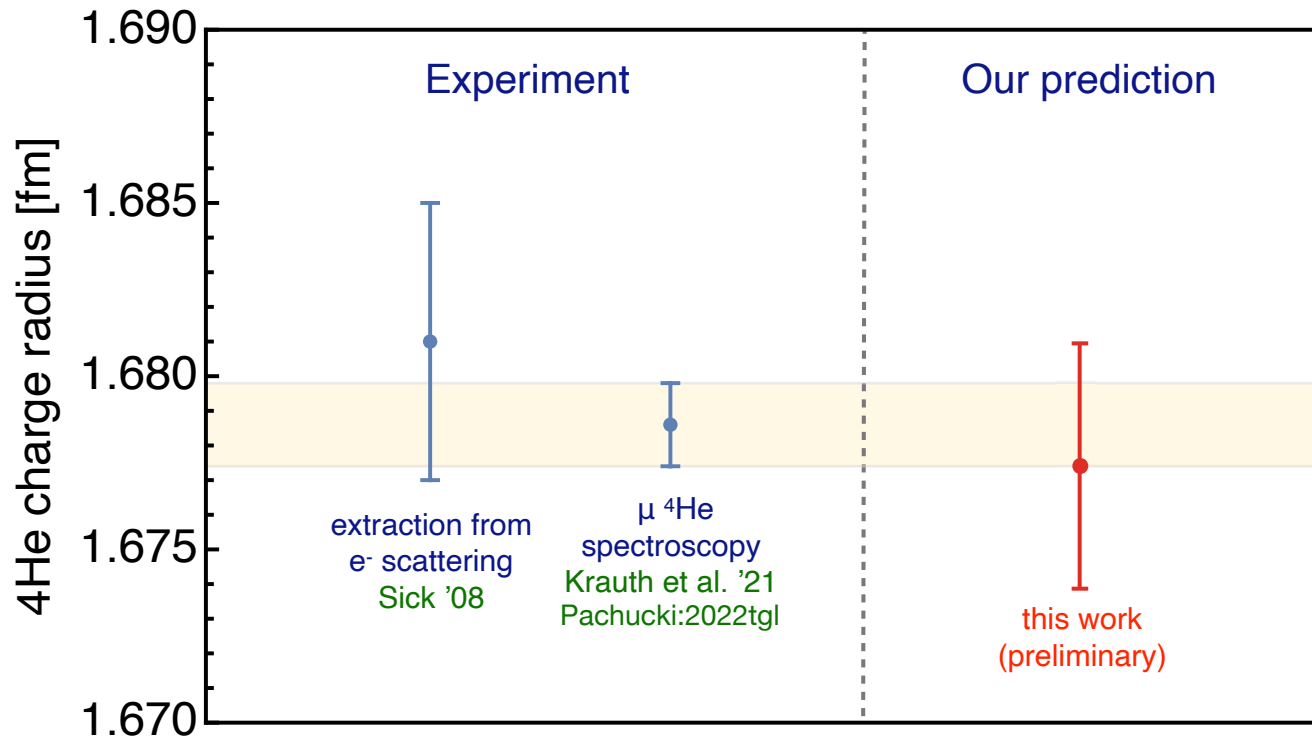
# The charge radius of ${}^4\text{He}$

Preliminary result for the  ${}^4\text{He}$  charge radius:

$$r_C^2({}^4\text{He}) = r_{\text{str}}^2({}^4\text{He}) + \left( r_p^2 + \frac{3}{4m_p^2} \right) + r_n^2 \Rightarrow r_C({}^4\text{He}) = (1.67\text{xx} \pm 0.0035) \text{ fm}$$

preliminary (CODATA 2018  $r_p$  + own determination of  $r_n$ )

Theory versus experiment:





# Nucleon size from the $^4\text{He}$ charge radius

Alternatively: Nucleon size from  $^4\text{He}$  radius

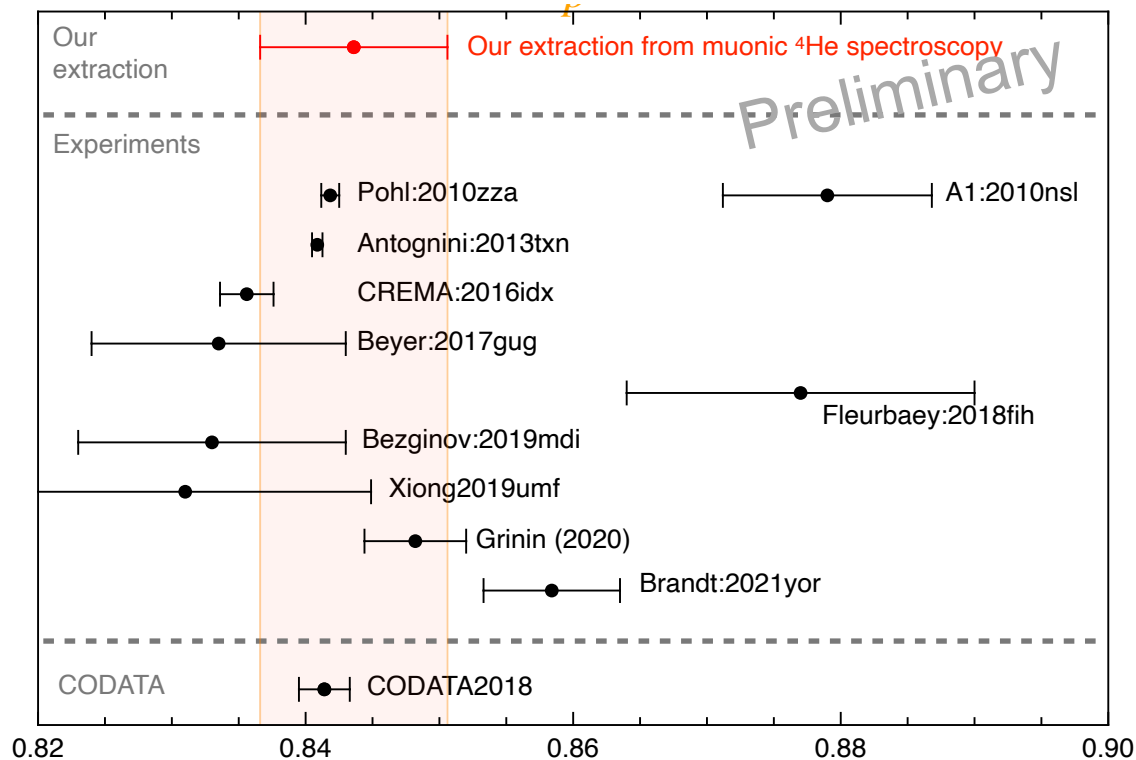
$$r_C^2(^4\text{He}) = r_{\text{str}}^2(^4\text{He}) + \left( r_p^2 + \frac{3}{4m_p^2} \right) + r_n^2 \Rightarrow$$

$$r_p^2 + r_n^2 = (0.6\text{xx} \pm 0.010) \text{ fm}^2$$

$$r_p = (0.8\text{xx} \pm 0.007) \text{ fm}$$

preliminary (own determination of  $r_n$ )

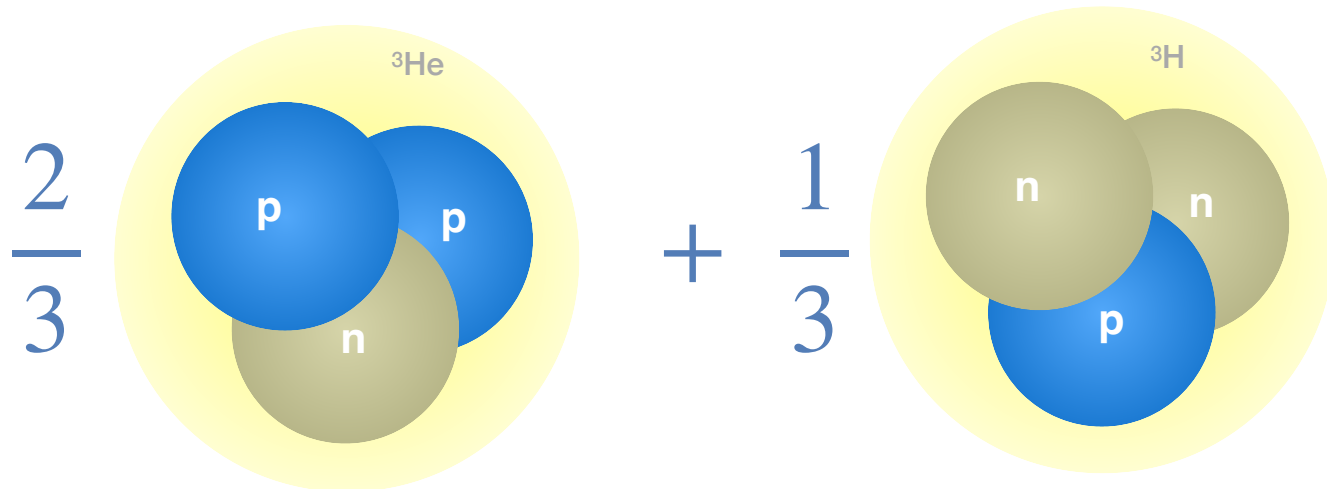
## Proton charge radius



# Isoscalar charge radius of $A = 3$ nuclei

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

Isovector  $\rho_{2N}$  beyond  $N^2LO$  not yet available  $\Rightarrow$  focus on the *isoscalar* radius



# Isoscalar charge radius of $A = 3$ nuclei

Predicted value of the isoscalar 3N charge radius  $r_C^{\text{isoscalar}} = \sqrt{\frac{1}{3}(r_C^{3H})^2 + \frac{2}{3}(r_C^{3He})^2}$

$$r_C^{\text{isoscalar}} = (1.90\text{xx} \pm 0.0026) \text{ fm}$$

preliminary (own determination of  $r_n$ )

Experimental value:  $r_{C, \text{exp}}^{\text{isoscalar}} = (1.9010 \pm 0.0260) \text{ fm}$

error dominated by the  $^3\text{H}$  datum

⇒ our prediction is 10x more precise than the current experimental value

The ongoing T-REX experiment in Mainz [Pohl et al.] aims at measuring the  $^3\text{H}$  charge radius within  $\pm 0.0002 \text{ fm}$  (i.e., 400x more precise) ⇒ the isoscalar radius will be known within  $\pm 0.0009 \text{ fm}$

⇒ precision test of nuclear chiral EFT

$^3\text{He}$

$r = 1.97007(94) \text{ fm}$   
0.05 %

CREMA 2023  
arXiv:2305.11679

$^3\text{H}$

$r = 1.7\text{xxx}( \text{ } 2) \text{ fm}$   
T-REX Mainz  
[in progress]

$r = 1.7550(860) \text{ fm}$   
Amroun et al. 1994

# Summary and outlook

- Charge & quadrupole FFs of  ${}^2\text{H}$  are in good shape (N<sup>4</sup>LO, high-precision)
- Other systems and processes are limited to N<sup>2</sup>LO accuracy due to unavailability of (consistently regularized) many-body forces & exchange currents
  - ⇒ **symmetry-preserving gradient flow regularization** talk by Hermann
- Correlations between BEs and radii can be employed to obtain precise results for the charge FFs of  ${}^4\text{He}$  &  ${}^3\text{H}(\text{e})_{\text{isoscalar}}$  already at this stage Arseniy Filin et al., in progress
- ${}^4\text{He}$ : Nuclear effects under control ⇒ new source of information about 1N radii
- ${}^3\text{He}/{}^3\text{H}$ : prediction for the isoscalar 3N charge radius 10x more precise than exp!

Thank you for your attention



# The 11th International Workshop on Chiral Dynamics

Aug 26 – 30, 2024  
Ruhr University Bochum, Germany  
Europe/Berlin timezone



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Scientific Program

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Travel Information

## Contact

✉ [cd2024@rub.de](mailto:cd2024@rub.de)

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The 11<sup>th</sup> International Workshop on Chiral Dynamics (CD2024) will take place August 26-30, 2024 at the [Ruhr University Bochum](#), Germany. This series of workshops started at MIT in 1994 and brings together theorists and experimentalists every three years to discuss the status, progress and challenges in the physics of low-energy QCD, Goldstone Boson dynamics, meson-baryon Interactions, few-body physics, lattice QCD and ChPT. Previous workshops took place in [Pisa \(2015\)](#), Durham, NC (2018) and [Beijing \(2021\)](#).

