

Intruder structure, shape coexistence,
and configuration mixing
from an *ab initio* perspective

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Progress in *Ab Initio* Nuclear Theory
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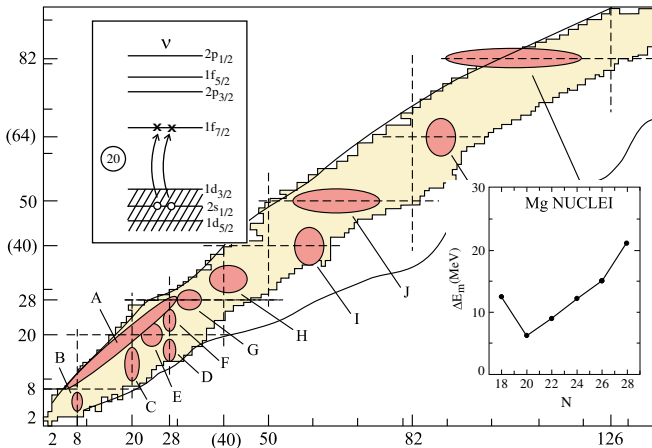


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Collaborators: Calvin W. Johnson (SDSU), Anna E. McCoy (ANL)

Intruder structure (and shape coexistence)

“[T]he intruder configuration ... corresponds to a more correlated state compared to the $0\hbar\omega$ states. Thus, low-lying $2p$ - $2h$ intruder configurations are favored only at and near to the ... shell closure.” *Normal ($0\hbar\omega$) vs. intruder ($2\hbar\omega$)*



In ab initio no-core configuration interaction (NCCI) calculations...

How do “normal” and “intruder” states converge? ${}^9\text{Be}$, ${}^{10}\text{Be}$

What do we find for intruder structure at $N = 8$? ${}^{11}\text{Li}$, ${}^{14}\text{C}$

Can we describe mixing of normal & intruder configurations?

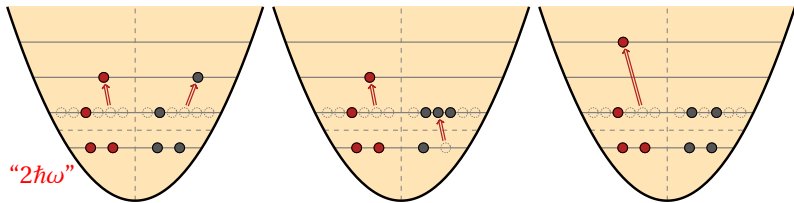
Postscript: Can we see the $4\hbar\omega$ intruder in ${}^{16}\text{O}$?

C 6	${}^9\text{C}^{(3/2-)}$	${}^{10}\text{C}^{0+}$	${}^{11}\text{C}^{3/2-}$	${}^{12}\text{C}^{0+}$	${}^{13}\text{C}^{1/2-}$	${}^{14}\text{C}^{0+}$
B 5	${}^8\text{B}^{2+}$	$[{}^9\text{B}]^{3/2-}$	${}^{10}\text{B}^{3+}$	${}^{11}\text{B}^{3/2-}$	${}^{12}\text{B}^{1+}$	${}^{13}\text{B}^{3/2-}$
Be 4	${}^7\text{Be}^{3/2-}$	$[{}^8\text{Be}]^{0+}$	${}^9\text{Be}^{3/2-}$	${}^{10}\text{Be}^{0+}$	${}^{11}\text{Be}^{1/2+}$	${}^{12}\text{Be}^{0+}$
Li 3	${}^6\text{Li}^{1+}$	${}^7\text{Li}^{3/2-}$	${}^8\text{Li}^{2+}$	${}^9\text{Li}^{3/2-}$		${}^{11}\text{Li}^{3/2-}$
	3	4	5	6	7	8

Many-body problem in an oscillator basis

No-core configuration interaction (NCCI) approach

a.k.a. no-core shell model (NCSM)



Antisymmetrized product basis *Slater determinants*

Distribute nucleons over oscillator shells

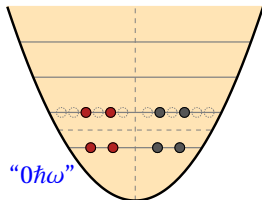
Organize basis by # oscillator excitations N_{ex}

relative to lowest Pauli-allowed filling

$N_{\text{ex}} = 0, 2, \dots$ (i.e., “ $0\hbar\omega$ ”, “ $2\hbar\omega$ ”, ...)

Basis must be truncated: $N_{\text{ex}} \leq N_{\text{max}}$

Convergence towards exact result with increasing N_{max} ...



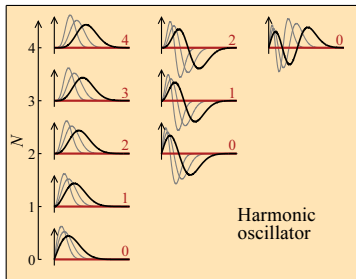
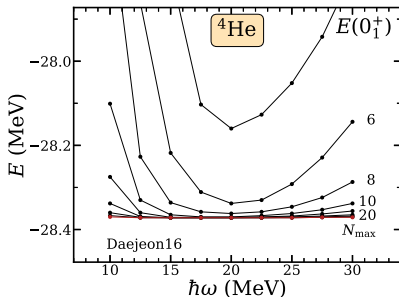
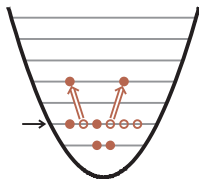
Convergence of NCCI calculations

Results in finite space depend upon:

- Many-body truncation N_{\max}
- Oscillator length b (or $\hbar\omega$)

$$b = \frac{(\hbar c)}{[(m_N c^2)(\hbar\omega)]^{1/2}}$$

Convergence of results signaled
by independence of N_{\max} & $\hbar\omega$



Harmonic oscillator

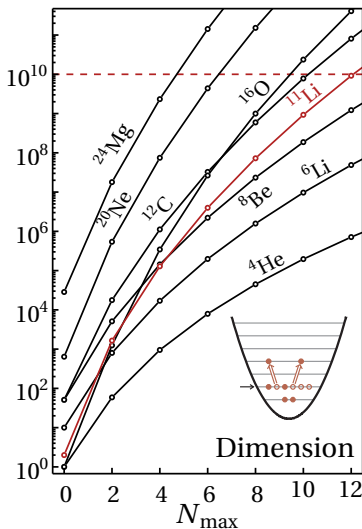
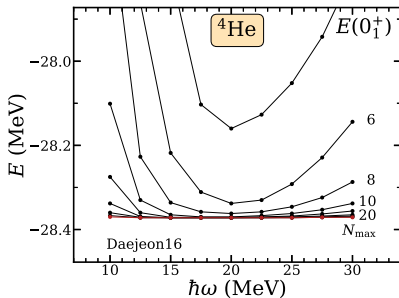
Convergence of NCCI calculations

Results in finite space depend upon:

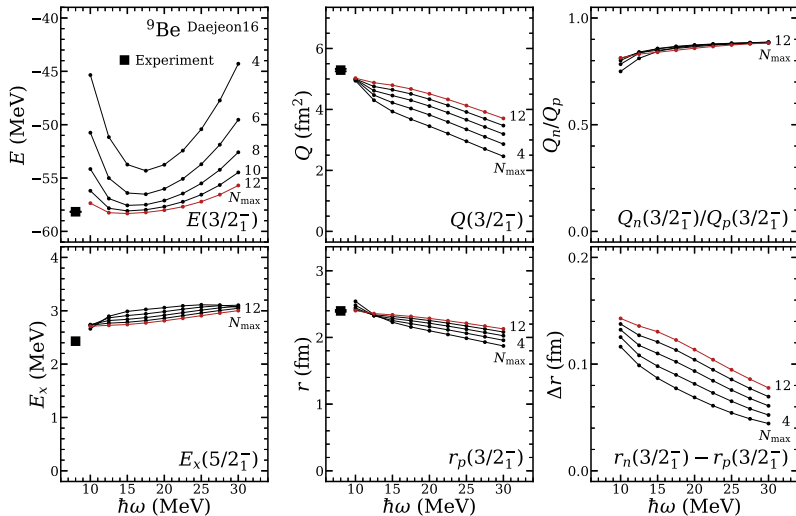
- Many-body truncation N_{\max}
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$$b = \frac{(\hbar c)}{[(m_N c^2)(\hbar\omega)]^{1/2}}$$

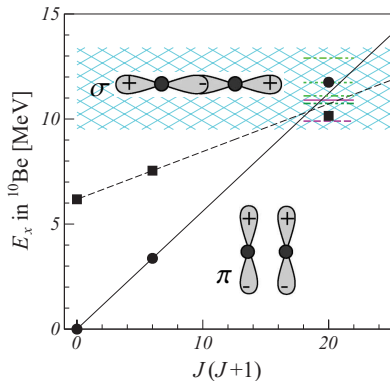
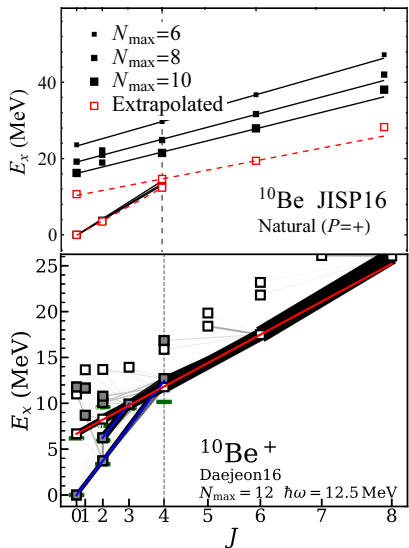
Convergence of results signaled
by independence of N_{\max} & $\hbar\omega$



Convergence for “normal” states ${}^9\text{Be}$

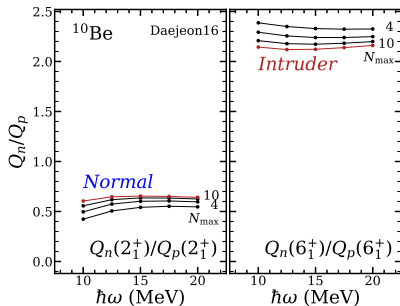
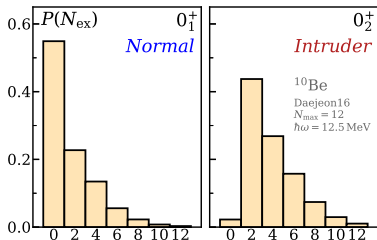
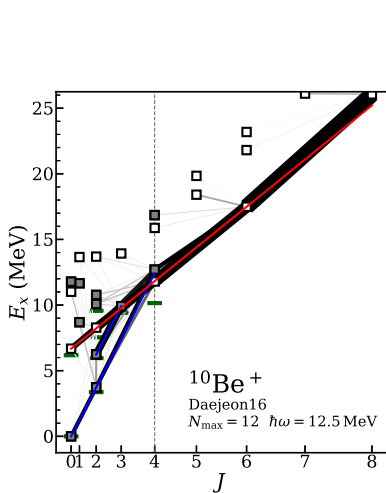


Convergence for “intruder” band ^{10}Be



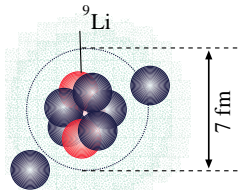
From D. Suzuki *et al.*, Phys. Rev. C **87**, 054301 (2013).
 Orbital schematics from Y. Kanada-En'yo, H. Horiuchi,
 and A. Doté, Phys. Rev. C **60**, 064304 (1999).

Structure of “intruder” band ^{10}Be



Structure of the ^{11}Li ground state?

Li 3			6Li^{1+}	$7\text{Li}^{3/2-}$	8Li^{2+}	$9\text{Li}^{3/2-}$	$^{11}\text{Li}^{3/2-}$
He 2	$^3\text{He}^{1/2+}$	$^4\text{He}^{0+}$		$^6\text{He}^{0+}$		$^8\text{He}^{0+}$	
H 1	$^2\text{H}^{1+}$	$^3\text{H}^{1/2+}$					
	1	2	3	4	5	6	7
	N						



Shell model: Doubly-magic plus one proton

Closed shell neutron

Interaction σ & p scattering: Enhanced matter radius

Neutron halo

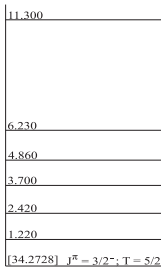
Neutron knockout: Neutron $(0p_{3/2})^2$ and $(1s_{1/2})^2$ contribute about equally to ground state

Intruder configurations

Excitation spectrum: Not particularly illuminating!

No J^P assignments. Unbound, but relatively narrow.

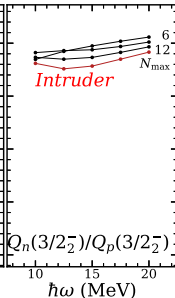
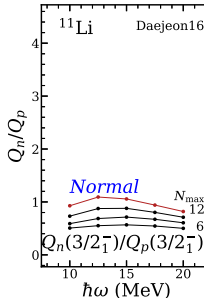
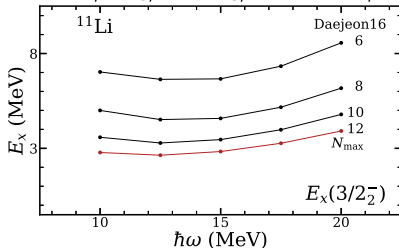
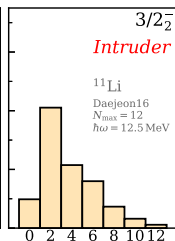
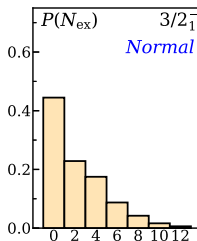
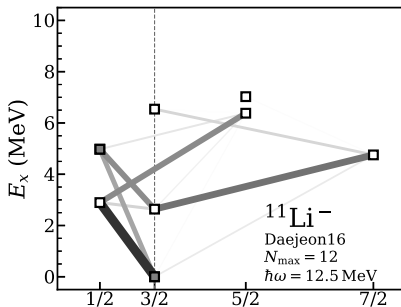
^{11}Li B. Jonson



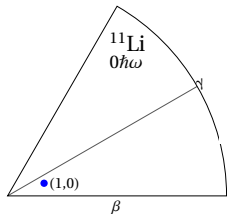
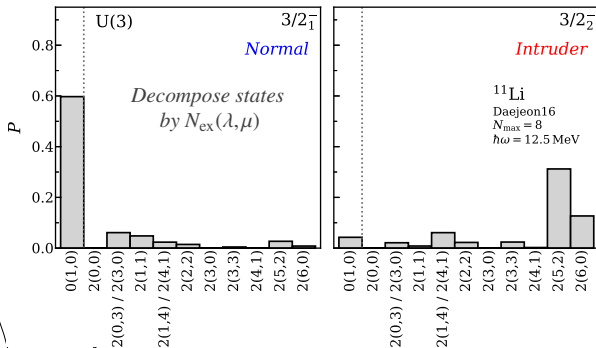
^{11}Li

TUNL (2012)

Low-lying intruder structure in ^{11}Li



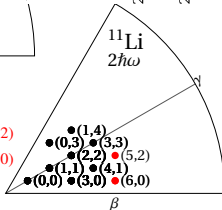
Elliott SU(3) symmetry for ^{11}Li



Proton-neutron structure

$$\pi(1,0) \times \nu(4,2) \Rightarrow (5,2)$$

$$\pi(1,0) \times \nu(5,0) \Rightarrow (6,0)$$



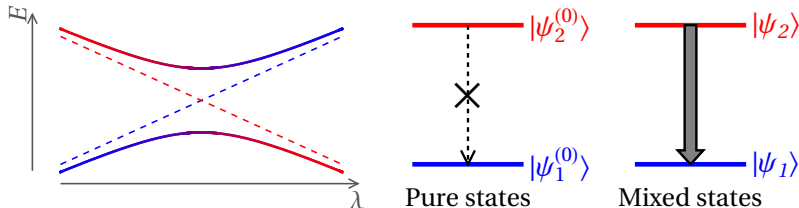
Intrinsic deformation for irrep (λ, μ)

$$\beta \propto (Q \cdot Q)^{1/2}$$

$$\propto (\lambda^2 + \lambda\mu + \mu^2 + 3\lambda + 3\mu + 3)^{1/2}$$

$$\gamma = \tan^{-1}\left(\frac{\sqrt{3}(\mu+3)}{2\lambda+\mu+3}\right)$$

Transition as measure of intruder mixing



$$H = \begin{pmatrix} E_1(\lambda) & V \\ V & E_2(\lambda) \end{pmatrix} \quad \begin{pmatrix} |\psi_1\rangle \\ |\psi_2\rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} |\psi_1^{(0)}\rangle \\ |\psi_2^{(0)}\rangle \end{pmatrix}$$

Mixing depends on *energy difference* $E_2 - E_1$ and *mixing matrix element* V .

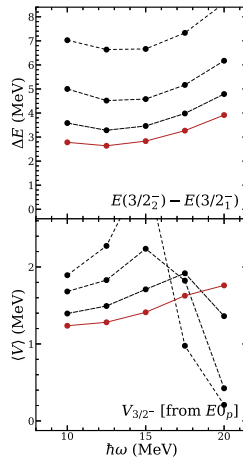
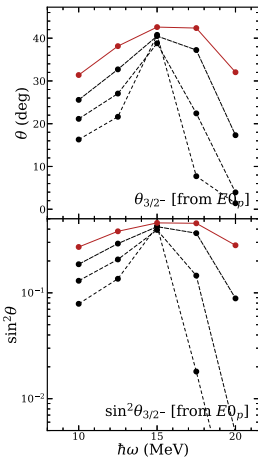
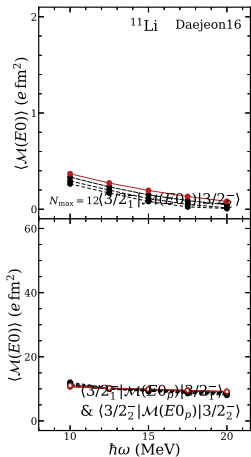
If transition operator \mathcal{M} does not connect “pure” (unmixed) states, transition matrix element for “mixed” states measures: (1) their *mixing* and (2) the difference in diagonal matrix elements, *i.e.*, moments $M_2 - M_1$:

$$\langle \psi_1 | \mathcal{M} | \psi_2 \rangle = \cos\theta \sin\theta \left[\langle \psi_2^{(0)} | \mathcal{M} | \psi_2^{(0)} \rangle - \langle \psi_1^{(0)} | \mathcal{M} | \psi_1^{(0)} \rangle \right]$$

Mixing analysis of *ab initio* calculations for ^{11}Li

Assume $\langle 0\hbar\omega | \mathcal{M}(E0) | 2\hbar\omega \rangle$ vanishes for “pure” (unmixed) $3/2^-$ states.

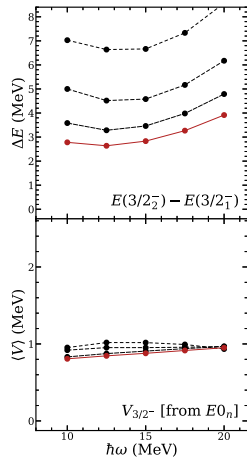
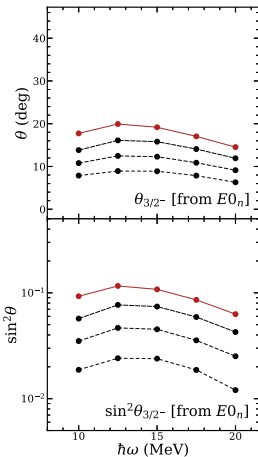
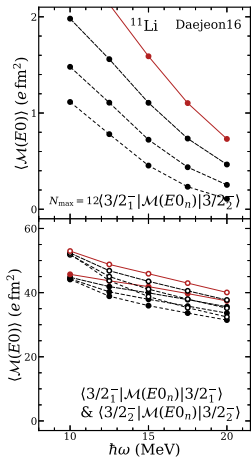
Deduce mixing from matrix elements for NCCI calculated (mixed) states.



Mixing analysis of *ab initio* calculations for ^{11}Li

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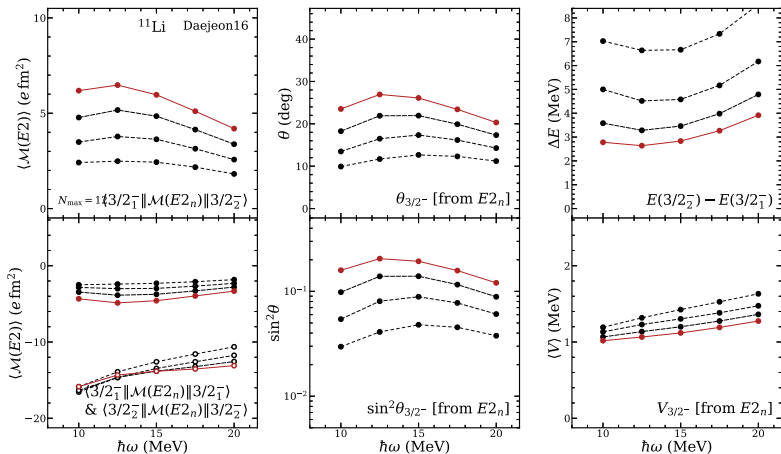
Deduce mixing from matrix elements for NCCI calculated (mixed) states.



Mixing analysis of *ab initio* calculations for ^{11}Li

Assume $\langle 0\hbar\omega | \mathcal{M}(E2) | 2\hbar\omega \rangle$ vanishes for “pure” (unmixed) $3/2^-$ states.

Deduce mixing from matrix elements for NCCI calculated (mixed) states.



The $E2$ strength to the first 2^+ state(s) in ^{14}C ?

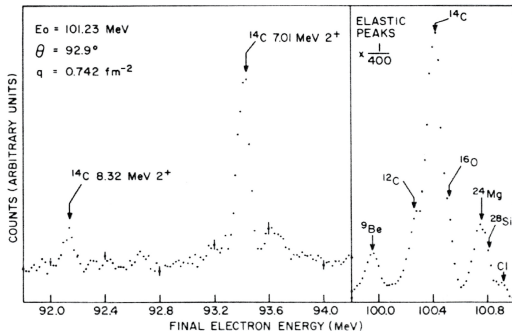
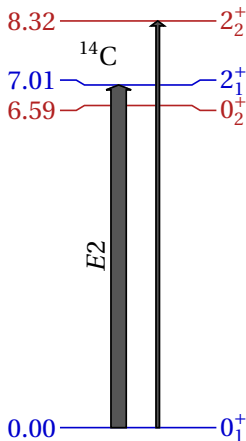
Electron Scattering from Low Lying 2^+ States in $^{14}\text{C}^*$

Hall Crannell, P.L. Hallowell, J.T. O'Brien,
J.M. Finn and F.J. Kline[†]

The Catholic University of America, Washington, D.C.

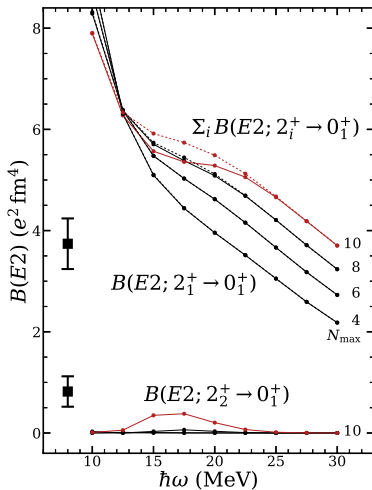
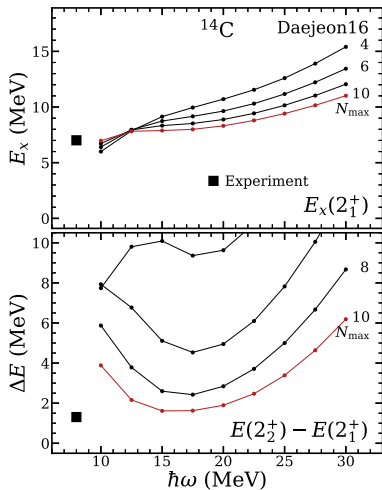
and

S. Penner, J.W. Lightbody, Jr., and S.P. Pivovinsky
National Bureau of Standards, Washington, D.C.

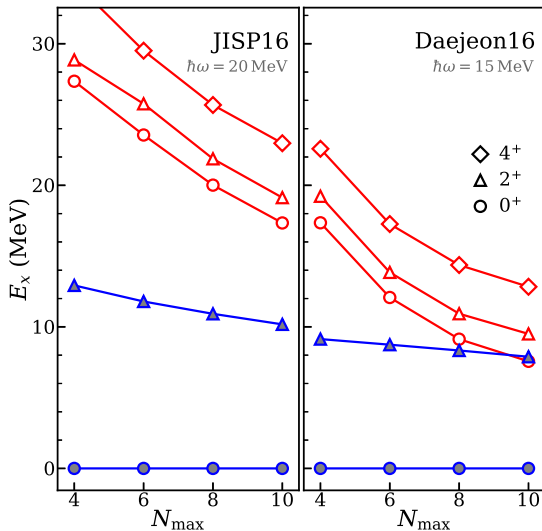
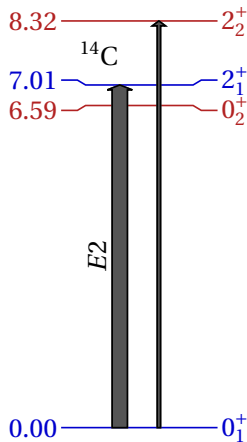


H. Crannell *et al.*, Proc. Int. Conf. Nucl. Struct. Studies Using Electron Scattering and Photoreaction, Sendai, Japan (1972).

The $E2$ strength to the first 2^+ state(s) in ^{14}C ?



Convergence of intruder state energies in ^{14}C

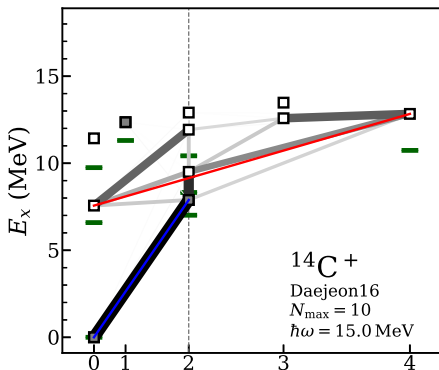
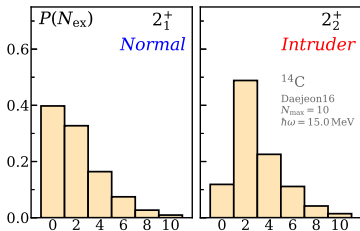
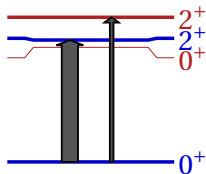


Low-lying intruder structure in ^{14}C

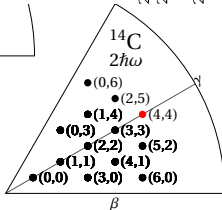
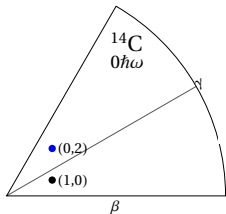
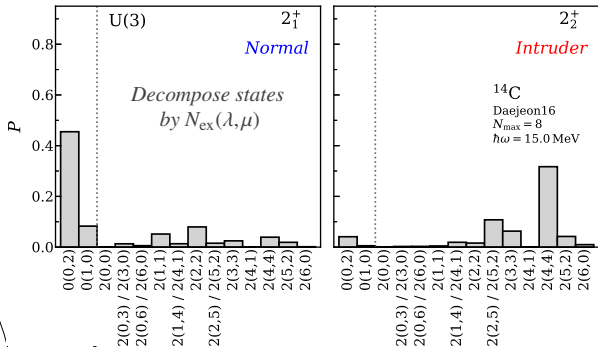
Coexisting $0^+ - 2^+$ sequences: $0\hbar\omega$ and $2\hbar\omega$

Very different “moments of inertia” $\Rightarrow 2^+$ states approach and mix

Excited structure as triaxial rotor? *Elliott SU(3)*



Elliott SU(3) symmetry for ^{14}C



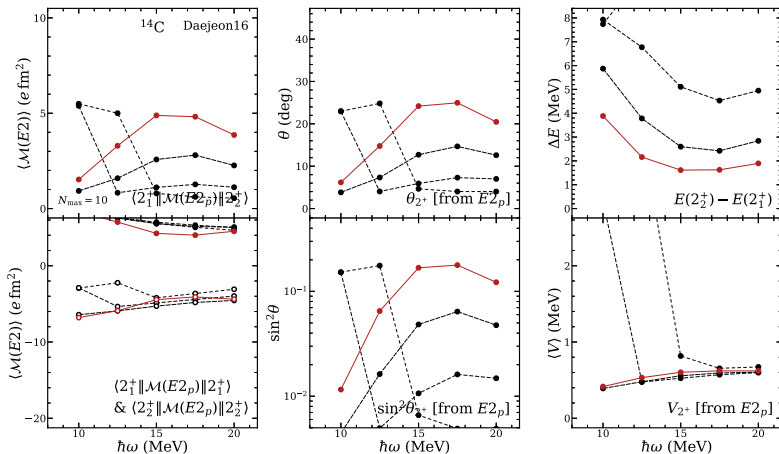
Proton-neutron triaxiality

$$\underbrace{\pi(0,2)}_{\text{oblate}} \times \underbrace{\nu(4,2)}_{\text{prolate}} \Rightarrow (4,4)$$

Mixing analysis of *ab initio* calculations for ^{14}C

Assume $\langle 0\hbar\omega | \mathcal{M}(E2) | 2\hbar\omega \rangle$ vanishes for “pure” (unmixed) 2^+ states.

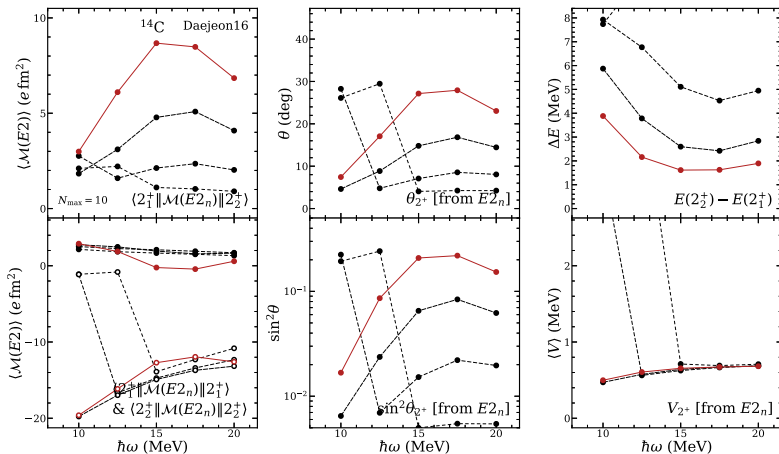
Deduce mixing from matrix elements for NCCI calculated (mixed) states.



Mixing analysis of *ab initio* calculations for ^{14}C

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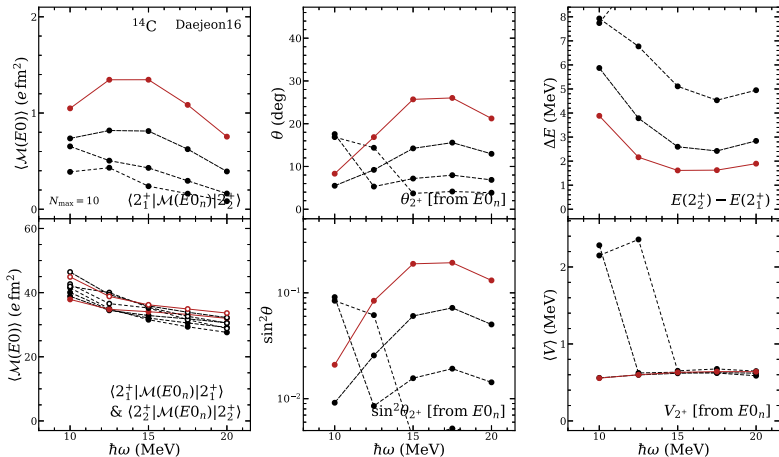
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The $E2$ strength to the first 2^+ state(s) in ^{14}C ?

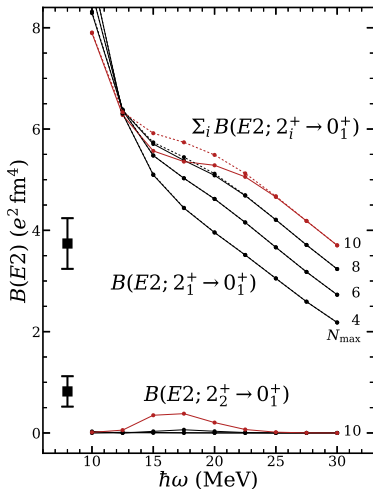
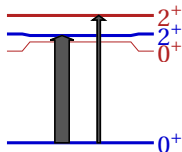
Two-state mixing estimate

$$V \approx 0.6 \text{ MeV} \quad \textit{Ab initio}$$

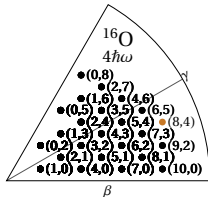
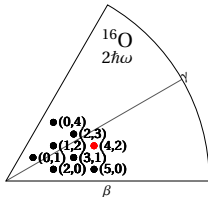
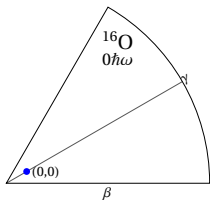
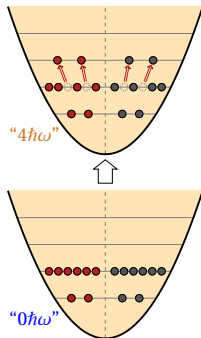
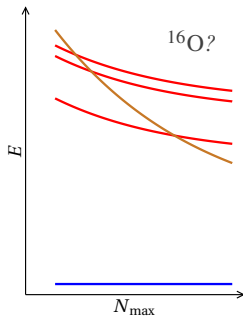
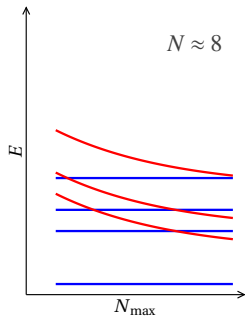
$$\begin{aligned} \Delta &\equiv \frac{1}{2} [E(2_2^+) - E(2_1^+)] \\ &= \frac{1}{2} [(8.32 \text{ MeV}) - (7.01 \text{ MeV})] \\ &= 0.65 \text{ MeV} \quad \textit{Experiment} \end{aligned}$$

$$\sin 2\theta = -\frac{V}{\Delta} \Rightarrow |\theta| \approx 34^\circ \quad \textit{Delicate!}$$

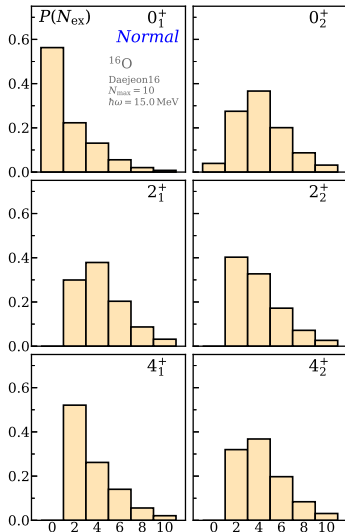
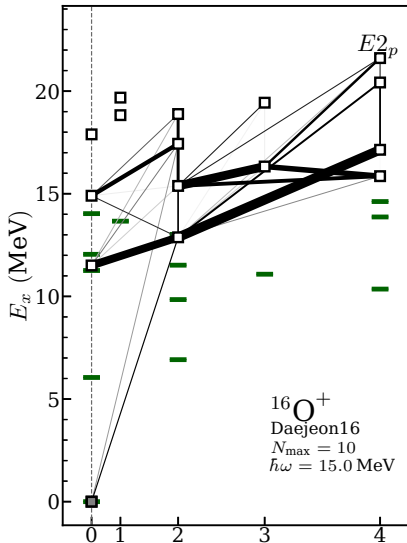
$$\frac{B(E2; 0_1^+ \rightarrow 2_2^+)}{B(E2; 0_1^+ \rightarrow 2_1^+)} = \frac{\sin^2 \theta}{\cos^2 \theta} \approx \boxed{0.44}$$



Convergence of normal, intruder, and “super-intruder” states?



Low-lying (*super-*)intruder structure in ^{16}O ?



Summary

Different states in low-lying spectrum have different...

- Rotational moments of inertia *Energy spacing within band*
- Shell model character *Normal* ($0\hbar\omega$) vs. *intruder* ($2\hbar\omega$)
- Proton/neutron asymmetric deformation Q_n/Q_p
- Elliott SU(3) symmetry (\approx “shape”)

Intruders hard to converge, but tractable with soft interaction *Daejeon16*

Mixing in *ab initio* results... *Emergent two-state mixing?*

- Strong mixing as same- J states approach *Within a few MeV*
- Mixing can be transient as energies cross $^{10}\text{Be } 4_1^+ \& 4_2^+$
- Mixing can be physical $^{11}\text{Li } \textit{ground state} / ^{14}\text{C } 2_1^+ \& 2_2^+$
- Transition matrix element provides handle on mixing angle θ
- Calculated “energy denominator” may be unconverged or inexact
- *But...* Can robustly extract emergent mixing matrix element V
 \Rightarrow *Estimate expected mixing at “physical” energy difference*

Beware! Ignore imminent mixing with an intruder at your own risk!

For $N = Z...$ Elusive $4\hbar\omega$ intruder states within reach? $^{16}\text{O } 0_2^+$