## Improved medium-mass nuclear structure and responses with the IMSRG









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### Matthias Heinz

with Tom Plies, Jan Hoppe, Frederic Noël, Takayuki Miyagi, Alexander Tichai, Kai Hebeler, Martin Hoferichter, Ragnar Stroberg, Achim Schwenk





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in-medium similarity renormalization group

 IMSRG generates unitary transformation of Hamiltonian

$$\frac{dH}{ds} = [\eta, H]$$

- Normal order with respect to  $|\Phi\rangle$  approximately handles 3N forces and induced many-body forces



Hergert et al., Phys. Rep. 621 (2016)





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- Obtain  $\overline{H} = UHU^{\dagger}$  with  $|\Psi\rangle = U|\Phi\rangle = e^{\Omega}|\Phi\rangle$ and  $\Omega = \Omega_1 + \Omega_2 + \dots$



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Hergert et al., Phys. Rep. **621** (2016)

### **Truncation necessary!** Standard = IMSRG(2) More refined = **IMSRG(3)**

**MH** et al., PRC **103** (2021)









- 1. Input Hamiltonian H
- 2. Solve for mean field (Hartree-Fock, NAT)
  - Input dependence: H,  $e_{max}$ ,  $E_{3max}$ ,  $\hbar\omega$
  - Output: reference state  $|\Phi\rangle$ , basis  $\{\phi_p\}$
- 3. Solve for many-body correlations [IMSRG(2)/(3)]
  - Input dependence: H,  $|\Phi\rangle$ ,  $\{\phi_p\}$ , other ops ...

normal ordering Hebeler, **MH**, et al., PRC **107** (2023)

• Output:  $|\Psi\rangle$ , *E*, expectation values of ops ...



## The commutator core of the IMSRG $[A^{(K)}, B^{(L)}] = \sum_{M} C^{(M)}$ IMSRG(2)

- Normal-ordered commutator induces many-body operators
  See talk by Ragnar Stroberg next
- Fundamental commutator  $[A^{(K)}, B^{(L)}]^{(M)} = C^{(M)}$ with cost  $\mathcal{O}(N^{K+L+M})$
- IMSRG(3): 7  $\rightarrow$  17 terms
- Organize based on computational cost and perturbative importance

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MH et al., PRC 103 (2021)



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- IMSRG(3) = next order in many-body expansion
- Systematic improvement towards exact results
- Benefit greatest for very nonperturbative problems
- Excellent precision on energies







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## IMSRG(3) corrections for 2<sup>+</sup> of <sup>48</sup>Ca

- IMSRG(2) predictions for 2<sup>+</sup> energy in Ca follow experimental trends...
- ... except at <sup>48</sup>Ca
- In CC, similar overprediction resolved by 3-body contributions



Simonis et al., PRC **96** (2017)



## IMSRG(3) corrections for 2<sup>+</sup> of <sup>48</sup>Ca



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- Truncations on 3B operators necessary for realistic calculations:  $e_{max,3b}$ ,  $E_{3max}$
- Convergence in calcium challenging
- Substantial corrections to  $2^+$  energy consistent with CC and experiment
- Revamped numerical implementation to reach convergence Novario et al., PRC 102 (2020)

### Precision IMSRG calculations in medium-mass nuclei possible soon!



## Nuclear responses for elastic electron scattering

- Electron scattering data in <sup>27</sup>Al
- Longitudinal/Coulomb contributions:  $M^{J}, \Phi^{''J} (J = 0, 2, 4)$
- Transverse contributions:  $\Delta^{J}, \Sigma^{'J} (J = 1, 3, 5)$

Hagen et al., Nat. Phys. 12 (2015), Gazda et al., PRD 95 (2016), Hoferichter et al., PRD 102 (2020), Hu et al., PRL 128 (2022)

- Center-of-mass corrections for responses
- Fourier transform to obtain densities  $\rightarrow$  expectation values

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Hagen et al., PRL **103** (2009)



## Nuclear responses for elastic electron scattering

- Electron scattering data in <sup>2</sup> Al
- Longitudinal/Coulomb contributions: **PSA (Martin Hoferichter)** Error in  $r_{so}^2$  from Ong et al., PRC 82 (2010) Fix:  $r_{so}^2 = \frac{1}{Z} \sum_{i} \frac{\mu_i - Q_i}{M^2} (\kappa_i + 1) \rightarrow \frac{1}{Z} \sum_{i} \frac{\mu_i - Q_i/2}{M^2} (\kappa_i + 1)$
- Center-of-mass corrections for responses
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Hagen et al., Nat. Phys. **12** (2015)

Hagen et al., PRL **103** (2009)



- Compute  $F_{ch}$ ,  $F_W$  for set of Hamiltonians (J = 0)



### Weak scattering in nuclei strongly constrained by ab initio nuclear structure!

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### SVD for NN and 3N forces singular value decomposition

$$V = L \cdot \Sigma \cdot R^{\dagger}$$

Largest singular values are most important

$$\Sigma = \operatorname{diag}(s_i)$$

Low-rank approximation via truncation (keeping largest singular values)

$$\tilde{V}$$
 =  $\tilde{L}$  ·  $\tilde{\Sigma}$  ·  $\tilde{R}$ 

Tichai, **MH**, et al., PLB **821** (2021) Tichai, **MH**, et al., arxiv:2307.15572

NN EMN 500 (a)  $^{1}S_{0}$  ${}^{3}S_{1} - {}^{3}D_{1}$  ${}^{3}P_{0}$  $\log s_i$  ${}^{1}D_{2}$ -8 -10 -12 2080 60 40 U rank



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### **Operator basis for low-resolution potentials** Tom Plies @ TU Darmstadt

- SVD to recover new operator basis



- Treat singular values s; as free parameters (LECs)

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Low-resolution potentials lack linear operator structure of chiral EFT

is: 
$$V = \sum_{i} s_i O_i = \sum_{i} s_i |l_i\rangle\langle r_i|$$

Constrain based on chiral EFT uncertainties and propagate to predictions

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## Impact of singular values



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## Matching to low-energy phase shifts

### Single partial wave

- Vary  $s_i$  within reasonable range:  $\rightarrow \vec{s}$
- Constrain 10k samples based on likelihood  $\mathscr{L}(\vec{s}) \sim \prod_{E} \mathscr{N}(\delta(\vec{s}, E) - \delta(\vec{s}_{ref}, E), \sigma_{EKM}^2)$
- Resample to 100 samples based on likelihood
- Multiple partial waves  $\begin{pmatrix} 1 S_0, {}^3S_1, {}^1P_1, {}^3P_0, {}^3P_1, {}^3P_2 \end{pmatrix}$  with charge indep.
- Product space of  $\vec{s}$  in different partial waves
- Reduce to 64 samples





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### PPDs for ground-state energies posterior predictive distributions **NN-only for now**



- EFT uncertainties for low-resolution potentials in nuclei



Reproduction of EKM uncertainties good (improved by more samples)

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## **Conclusion and outlook**

 Establishing IMSRG(3) for high-precision description of medium-mass nuclei and uncertainty quantification

- Exploit correlated uncertainties to constrain difficult-to-measure and nonobservable quantities
- New operator basis from SVD for uncertainty quantification with low-resolution Hamiltonians



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### **Thank you for your attention!**

