## Electric dipole polarizability in medium-mass nuclei

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## Extreme matter in neutron stars


A. Watts et al, RMP 88, 021001 (2016).

## Neutron matter equation of state (EOS)

Infinite nuclear matter EOS:

$$
\begin{gathered}
\mathcal{E}(\rho, \alpha)=\mathcal{E}_{\mathrm{SNM}}(\rho)+\alpha^{2} \mathcal{S}(\rho)+\mathcal{O}\left(\alpha^{4}\right) \\
\rho=\left(\rho_{n}+\rho_{p}\right) \quad \alpha=\left(\rho_{n}-\rho_{p}\right) / \rho
\end{gathered}
$$

where the symmetry energy is:

$$
\mathcal{S}(\rho)=J+L \frac{\left(\rho-\rho_{0}\right)}{3 \rho_{0}}+\ldots .
$$

symmetry energy
at saturation density
slope parameter, related to pressure of pure neutron matter at saturation density
S. Huth et al, PRC 103, 025803 (2021).


## How to constrain L?

Neutron-skin thickness

B. Hu et al, Nat. Phys. 18, 1196-1200 (2022).

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Electric dipole polarizability


Data from X. Roca-Maza et al, PRC 88, 024316 (2013), B. Hu et al, Nat. Phys. 18, 1196-1200 (2022).

## Nuclear response functions

$$
\left.R(\omega)=\sum_{f}\left|\left\langle\Psi_{f}\right| \Theta\right| \Psi_{0}\right\rangle\left.\right|^{2} \delta\left(E_{f}-E_{0}-\omega\right)
$$



## Coupled-cluster theory

$\square$ Starting point: Hartree-Fock reference state on the HO basis $\left|\Phi_{0}\right\rangle$
$\square$ Add correlations via:

$$
e^{-T} H e^{T}\left|\Phi_{0}\right\rangle=\bar{H}\left|\Phi_{0}\right\rangle=E_{0}\left|\Phi_{0}\right\rangle
$$

with

$$
T=\sum t_{i}^{a} a_{a}^{\dagger} a_{i}+\sum t_{i j}^{a b} a_{a}^{\dagger} a_{b}^{\dagger} a_{j} a_{i}+\sum t_{i j k}^{a b c} a_{a}^{\dagger} a_{b}^{\dagger} a_{c}^{\dagger} a_{k} a_{j} a_{i}+\ldots
$$

$\rightarrow$ coefficients from coupled-cluster equations

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Similarity-transformed
Hamiltonian (non-Hermitian)
with

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## From bound to dipole-excited states <br> $$
\left.R(\omega)=\mathcal{f}_{f}\left|\left\langle\Psi_{f}\right| \Theta\right| \Psi_{0}\right\rangle\left.\right|^{2} \delta\left(E_{f}-E_{0}-\omega\right)
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\left.R(\omega)=\mathcal{F}_{f}\left|\left\langle\Psi_{f}\right| \theta\right| \Psi_{0}\right\rangle\left.\right|^{2} \delta\left(E_{f}-E_{0}-\omega\right) \\
L(\sigma, \Gamma)=\frac{\Gamma}{\pi} \int d \omega \frac{R(\omega)}{(\omega-\sigma)^{2}+\Gamma^{2}}=\frac{\Gamma}{\pi}\left\langle\Psi_{L} \mid \Psi_{R}\right\rangle
\end{gathered}
$$



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\text { Lorentz Integra } \\
\text { Transform (IIT) }
\end{array}} \\
\quad\left(\bar{H}-E_{0}-\sigma-i \Gamma\right)\left|\Psi_{R}\right\rangle=\bar{\Theta}\left|\Phi_{0}\right\rangle \quad \begin{array}{c}
\text { cc equation of motion } \\
\text { with a source }
\end{array}
\end{gathered}
$$

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\text { where } \\
\left(\bar{H}-E_{0}-\sigma-i \Gamma\right)\left|\Psi_{R}\right\rangle=\bar{\Theta}\left|\Phi_{0}\right\rangle \quad \begin{array}{c}
\sigma \\
\text { Lorentz Integra } \\
\text { Transform (LIT) } \\
\text { cc equation of motion } \\
\text { with a source }
\end{array}
\end{gathered}
$$

LIT-CC ansatz for closed-shell nuclei:

$$
\left|\Psi_{R}\right\rangle=\mathcal{R}\left|\Phi_{0}\right\rangle \quad \mathcal{R}=r_{0}+\sum r_{i}^{a} a_{a}^{\dagger} a_{i}+\sum r_{i j}^{a b} a_{a}^{\dagger} a_{b}^{\dagger} a_{j} a_{i}+\ldots
$$

## The case of ${ }^{40} \mathrm{Ca}$

R. Fearick, P. von Neumann-Cosel, S. Bacca, FB et al, Phys. Rev. Research 5, L022044 (2023).


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Constraints on slope parameter from chEFT interactions of this work: $\mathrm{L}=41-49 \mathrm{MeV}$ Global analysis of all Ca+Pb skin data [J. Lattimer, Particles (2023)]: L $=40 \pm 8 \mathrm{MeV}$

## Happy ending for ${ }^{40,48} \mathrm{Ca}$... but what's next?



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## LIT-CC for open-shell nuclei: the 2-Particle-Attached (2PA) case

$$
\begin{aligned}
& \mathcal{R}=\frac{1}{2} \sum r^{a b} a_{a}^{\dagger} a_{b}^{\dagger}+\frac{1}{6} \sum r_{i}^{a b c} a_{a}^{\dagger} a_{b}^{\dagger} a_{c}^{\dagger} a_{i}+\ldots \\
& \left|\Psi_{R}\right\rangle=\mathcal{R}\left|\Phi_{0}\right\rangle \\
& \text { 3p1h }
\end{aligned}
$$

## Many-body truncations in LIT-CC

To solve

$$
\left(\bar{H}-E_{0}-\sigma-i \Gamma\right)\left|\Psi_{R}\right\rangle=\bar{\Theta}\left|\Phi_{0}\right\rangle
$$

we have two CC expansions, for ground and excited states.

In the closed-shell case we consider excited states @CCSD and vary ground-state scheme.

We then estimate the many-body uncertainty as:

$$
\delta_{\alpha_{D}}^{\mathrm{CC}} \approx \frac{\alpha_{D}^{\mathrm{CCSD}}-\alpha_{D}^{\mathrm{CCSDT}-1}}{2}
$$



## What about the 2PA case?

In this case, for the moment we do not have access to orders higher than 3p1h...

If we can't go higher, let's go lower!

We use the 3p1h scheme for excited states, to keep as much correlations as possible in the LIT calculation, and we look at:

$$
\delta_{\alpha_{D}}^{\mathrm{CC}} \approx \frac{\alpha_{D}^{3 \mathrm{p} 1 \mathrm{~h} / 3 \mathrm{p} 1 \mathrm{~h}}-\alpha_{D}^{2 \mathrm{p} 0 \mathrm{~h} / 3 \mathrm{p} 1 \mathrm{~h}}}{2}
$$



## Electric dipole polarizability of ${ }^{24} \mathrm{O}$



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CCSDT-1


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## Electric dipole polarizability of ${ }^{24} \mathrm{O}$



## $\alpha_{D}$ along oxygen isotopes



## $\ldots$ and heavier: $\alpha_{D}$ in calcium isotopes



## $\ldots$ and heavier: $\alpha_{D}$ in calcium isotopes


$\ldots$ and heavier: $\alpha_{D}$ in calcium isotopes


## Conclusions

$\square$ Dipole polarizabilities provide a way to cast light on the collective excitations of the nucleus as well as to put constraints on the nuclear symmetry energy.
$\square$ We extended the reach of ab initio calculations of these electromagnetic observables to nuclei in the vicinity of closed shells, and investigated the effect of many-body truncations.Soon new experimental benchmarks for 2PA nuclei!

