



Electric dipole polarizability in medium-mass nuclei

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Extreme matter in neutron stars



A. Watts et al, RMP 88, 021001 (2016).

Neutron matter equation of state (EOS)

Infinite nuclear matter EOS:

$$\mathcal{E}(\rho, \alpha) = \mathcal{E}_{\text{SNM}}(\rho) + \alpha^2 \mathcal{S}(\rho) + \mathcal{O}(\alpha^4)$$
$$\rho = (\rho_n + \rho_p) \qquad \alpha = (\rho_n - \rho_p)/\rho$$

where the **symmetry energy** is:

$$\mathcal{S}(\rho) = \frac{J}{J} + \frac{L}{3\rho_0} \frac{(\rho - \rho_0)}{3\rho_0} + \dots$$

symmetry energy at saturation density slope parameter, related to pressure of pure neutron matter at saturation density



How to constrain L?

Neutron-skin thickness



B. Hu et al, Nat. Phys. 18, 1196–1200 (2022).

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Electric dipole polarizability

$$\alpha_D = 2\alpha \int d\omega \; \frac{R(\omega)}{\omega}$$



Data from X. Roca-Maza et al, PRC 88, 024316 (2013), B. Hu et al, Nat. Phys. 18, 1196-1200 (2022).

Nuclear response functions

$$R(\omega) = \sum_{f} |\langle \Psi_{f} | \Theta | \Psi_{0} \rangle |^{2} \delta(E_{f} - E_{0} - \omega)$$



Coupled-cluster theory

 \Box Starting point: Hartree-Fock reference state on the HO basis $|\Phi_0
angle$

Add correlations via:

$$e^{-T}He^{T}|\Phi_{0}\rangle = \overline{H}|\Phi_{0}\rangle = E_{0}|\Phi_{0}\rangle$$

with

$$T = \sum t_i^a a_a^{\dagger} a_i + \sum t_{ij}^{ab} a_a^{\dagger} a_b^{\dagger} a_j a_i + \sum t_{ijk}^{abc} a_a^{\dagger} a_b^{\dagger} a_c^{\dagger} a_k a_j a_i + \dots$$

→ coefficients from coupled-cluster equations

Similarity-transformed

Hamiltonian

(non-Hermitian)

Coupled-cluster theory

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angle$

□ Add correlations via:

(CCSD)

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singles and doubles
$$\xrightarrow{\rightarrow} \text{ coefficients from coupled-cluster equations}}$$

G. Hagen, T. Papenbrock, M. Hjorth-Jensen, D. J. Dean, RPP 77, 096302 (2014).

Similarity-transformed

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G. Hagen, T. Papenbrock, M. Hjorth-Jensen, D. J. Dean, RPP 77, 096302 (2014).

Hamiltonian

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From bound to dipole-excited states $R(\omega) = \sum_{f} |\langle \Psi_{f} | \Theta | \Psi_{0} \rangle |^{2} \delta(E_{f} - E_{0} - \omega)$

From bound to dipole-excited states $R(\omega) = \oint_{f} |\langle \Psi_{f} | \Theta | \Psi_{0} \rangle |^{2} \delta(E_{f} - E_{0} - \omega)$ $L(\sigma, \Gamma) = \frac{\Gamma}{\pi} \int d\omega \, \frac{R(\omega)}{(\omega - \sigma)^{2} + \Gamma^{2}} = \frac{\Gamma}{\pi} \langle \Psi_{L} | \Psi_{R} \rangle$ Lorentz Integral Transform (LIT)

$$(\overline{H} - E_0 - \sigma - i\Gamma) |\Psi_R\rangle = \overline{\Theta} |\Phi_0\rangle \qquad \begin{array}{c} \text{CC equation of motion} \\ \text{with a source} \end{array}$$

From bound to dipole-excited states $R(\omega) = \sum_{f} |\langle \Psi_{f} | \Theta | \Psi_{0} \rangle |^{2} \delta(E_{f} - E_{0} - \omega)$ Transform (ĽIT) where $\left(\overline{H} - E_0 - \sigma - i\Gamma\right) \left|\Psi_R\right\rangle = \overline{\Theta} \left|\Phi_0\right\rangle$ CC equation of motion with a source

LIT-CC ansatz for closed-shell nuclei:

$$|\Psi_R\rangle = \mathcal{R} |\Phi_0\rangle \qquad \mathcal{R} = r_0 + \sum r_i^a a_a^{\dagger} a_i + \sum r_{ij}^{ab} a_a^{\dagger} a_b^{\dagger} a_j a_i + \dots$$

S. Bacca, N. Barnea, G. Hagen, G. Orlandini, T. Papenbrock, PRL 111, 122502 (2013).

The case of ⁴⁰Ca

R. Fearick, P. von Neumann-Cosel, S. Bacca, FB et al, Phys. Rev. Research 5, L022044 (2023).



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Constraints on **slope parameter** from chEFT interactions of this work: L = 41 - 49 MeV Global analysis of all Ca+Pb skin data [J. Lattimer, Particles (2023)]: $L = 40 \pm 8$ MeV









We need to extend the LIT-CC method beyond closed-shell nuclei!

LIT-CC for open-shell nuclei: the 2-Particle-Attached (2PA) case



Many-body truncations in LIT-CC

To solve

$$\left(\overline{H} - E_0 - \sigma - i\Gamma\right) \left|\Psi_R\right\rangle = \overline{\Theta} \left|\Phi_0\right\rangle$$

we have two CC expansions, for ground and excited states.

In the closed-shell case we consider excited states @CCSD and vary ground-state scheme.

We then estimate the **many-body uncertainty** as:

$$\delta_{\alpha_D}^{\rm CC} \approx \frac{\alpha_D^{\rm CCSD} - \alpha_D^{\rm CCSDT-1}}{2}$$



What about the 2PA case?

In this case, for the moment we do not have access to orders higher than 3p1h...

If we can't go higher, let's go lower!

We use the **3p1h scheme for excited states**, to keep as much correlations as possible in the LIT calculation, and we look at:

$$\delta_{\alpha_D}^{\rm CC} \approx \frac{\alpha_D^{\rm 3p1h/3p1h} - \alpha_D^{\rm 2p0h/3p1h}}{2}$$













α_D along oxygen isotopes



... and heavier: α_D in calcium isotopes



FB et al., in preparation.

... and heavier: α_D in calcium isotopes



... and heavier: α_D in calcium isotopes



Conclusions

□ Dipole polarizabilities provide a way to cast light on the collective excitations of the nucleus as well as to put constraints on the nuclear symmetry energy.

□ We extended the reach of ab initio calculations of these electromagnetic observables to **nuclei in the vicinity of closed shells**, and investigated the **effect of many-body truncations**.

□ Soon new experimental benchmarks for 2PA nuclei!

Thank you for your attention!