

New Opportunities for Nuclear Structure Calculations for BSM Physics

Ayala Glick-Magid



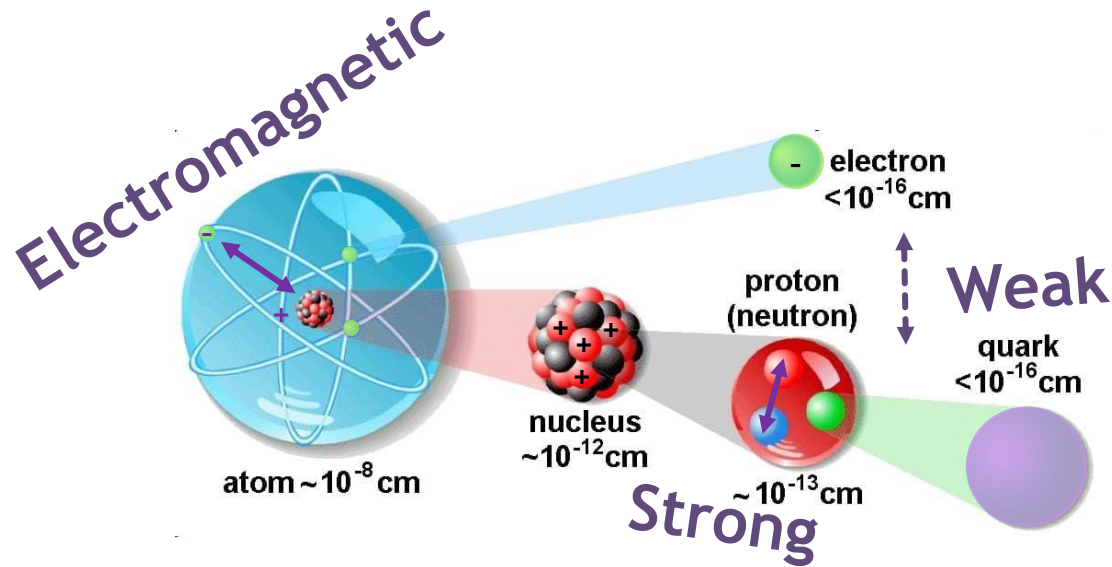
INSTITUTE for
NUCLEAR THEORY

PAINT 2024

Standard Model (SM)

Fundamental Forces

Elementary Particles



three generations of matter (fermions)

| | I | II | III | | |
|---------|--|---|--|---|--|
| LEPTONS | $\approx 0.511 \text{ MeV}/c^2$ -1 1/2 e electron | $\approx 105.67 \text{ MeV}/c^2$ -1 1/2 μ muon | $\approx 1.7768 \text{ GeV}/c^2$ -1 1/2 τ tau | 0 0 1 g gluon | $\approx 125.09 \text{ GeV}/c^2$ 0 0 0 H Higgs |
| | $< 2.2 \text{ eV}/c^2$ 0 1/2 ν_e electron neutrino | $< 1.7 \text{ MeV}/c^2$ 0 1/2 ν_μ muon neutrino | $< 15.5 \text{ MeV}/c^2$ 0 1/2 ν_τ tau neutrino | 0 0 1 γ photon | |
| | $\approx 2.4 \text{ MeV}/c^2$ 2/3 1/2 u up | $\approx 1.275 \text{ GeV}/c^2$ 2/3 1/2 c charm | $\approx 172.44 \text{ GeV}/c^2$ 2/3 1/2 t top | $\approx 91.19 \text{ GeV}/c^2$ 0 1 Z Z boson | |
| QUARKS | $\approx 4.8 \text{ MeV}/c^2$ -1/3 1/2 d down | $\approx 95 \text{ MeV}/c^2$ -1/3 1/2 s strange | $\approx 4.18 \text{ GeV}/c^2$ -1/3 1/2 b bottom | GAUGE BOSONS $\approx 80.39 \text{ GeV}/c^2$ 1 1 W W boson | SCALAR BOSONS |

Beyond Standard Model (BSM)

NOBEL PRIZE IN PHYSICS 2015

The Nobel Prize in Physics 2015 was awarded to Takaaki Kajita and Arthur B. McDonald for discovery of neutrino oscillations, which shows neutrinos have mass.

WHAT IS A NEUTRINO? Neutrinos are tiny subatomic particles, produced by nuclear reactions that take place in stars, including our sun, as well as in radioactive decay processes. They come in three 'flavours'.

V_e

ELECTRON NEUTRINO

V_μ

MUON NEUTRINO

V_τ

TAU NEUTRINO

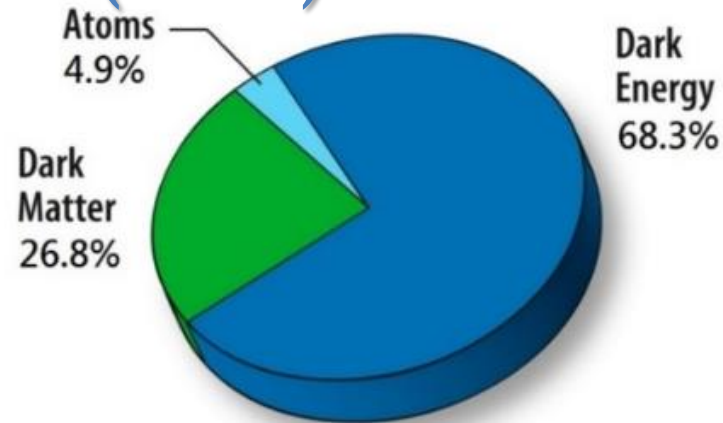
The nuclear reactions in the sun produce neutrinos, which we can detect.

The number of neutrinos detected was only a third of the expected value.

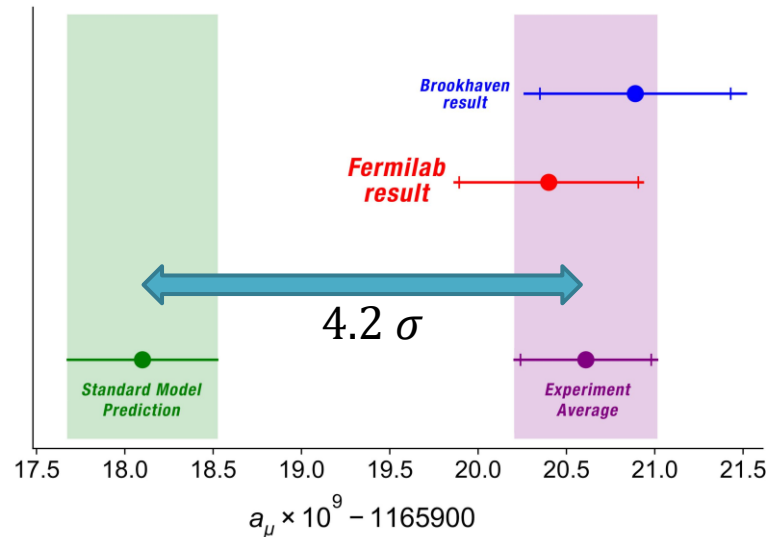
Neutrinos 'flip' between the three flavours, and only one type was being detected.

WHY DOES IT MATTER? If neutrinos oscillate between types, they must have mass, even if this mass is incredibly small. This contradicts the standard model of particle physics, which states they are massless.

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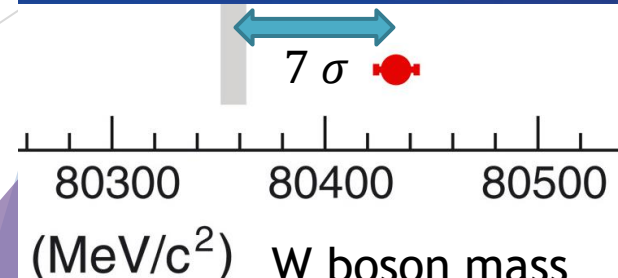


Dark Matter & Energy



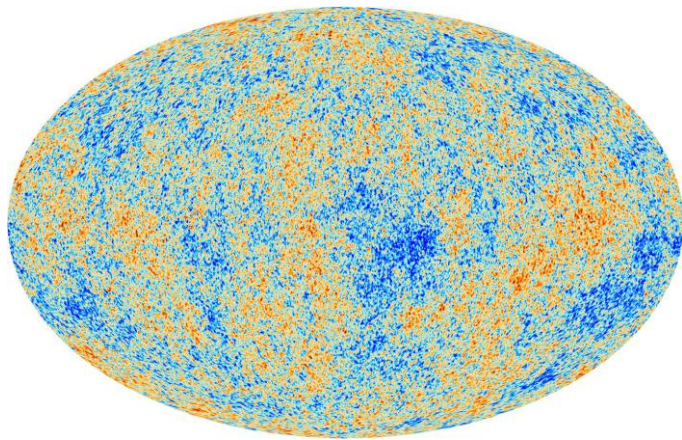
The Neutrino has mass, even though according to the SM it should not

Deviations from the SM at high precision: muon $g-2$, W mass



Searches for BSM physics

Astronomical Frontier
Astronomy

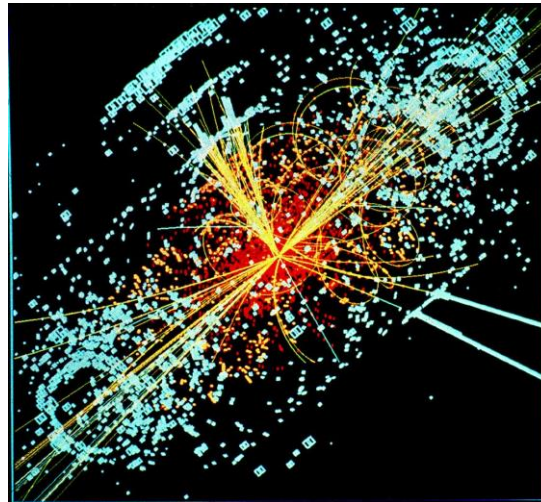


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E.g.

➤ Dark Matter

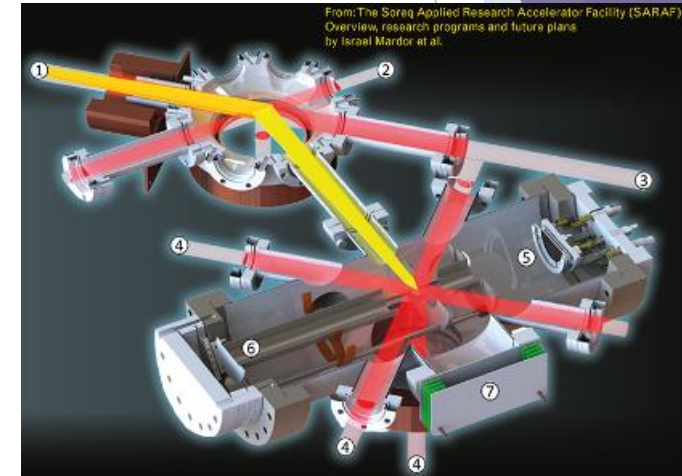
High Energy Frontier
Particles Physics



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➤ Lepton Flavor Violation

Precision Frontier
Nuclear Physics

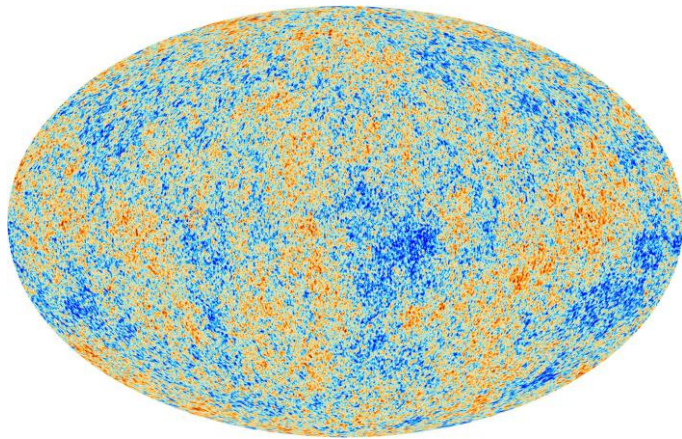


From: The Soreq Applied Research Accelerator Facility (SARAF):
Overview, research programs and future plans
by Israel Mardor et al.
Mardor et al., *Eur. Phys. J. A* 54, 91 (2018)

➤ New Weak Interactions

Searches for BSM physics

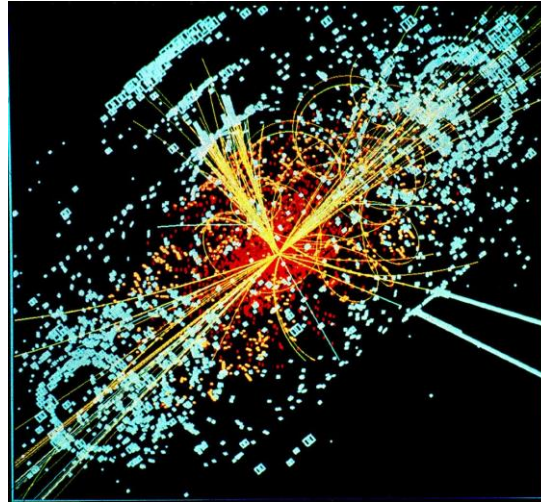
Astronomy



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➤ Dark Matter

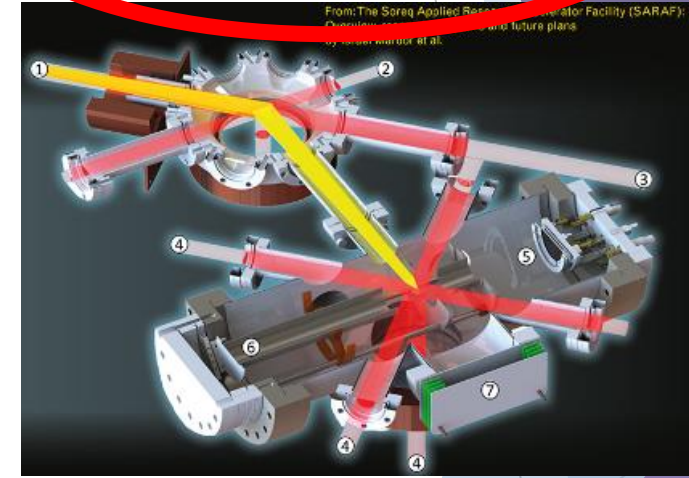
Particles Physics



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➤ Lepton Flavor Violation

Precision Frontier Nuclear Physics



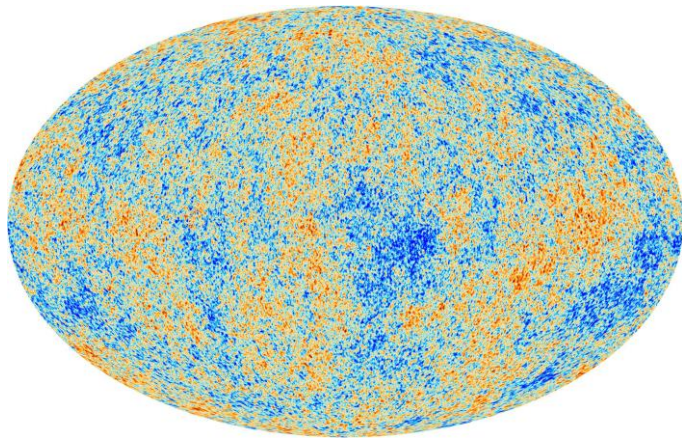
Mardor *et al.*, [Eur. Phys. J. A 54, 91](https://doi.org/10.1007/s00034-018-0618-1) (2018)

➤ New Weak Interactions

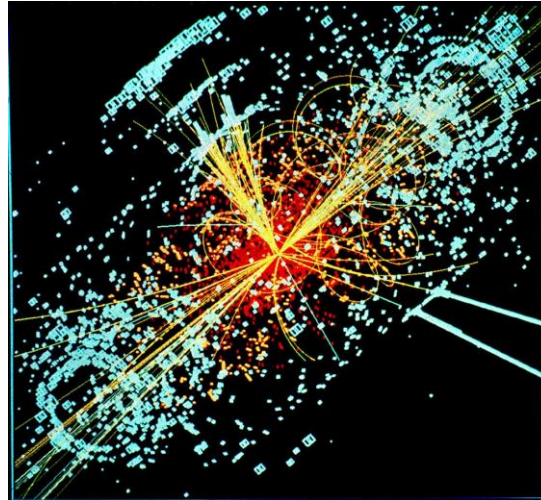
with nuclei...

Searches for BSM physics

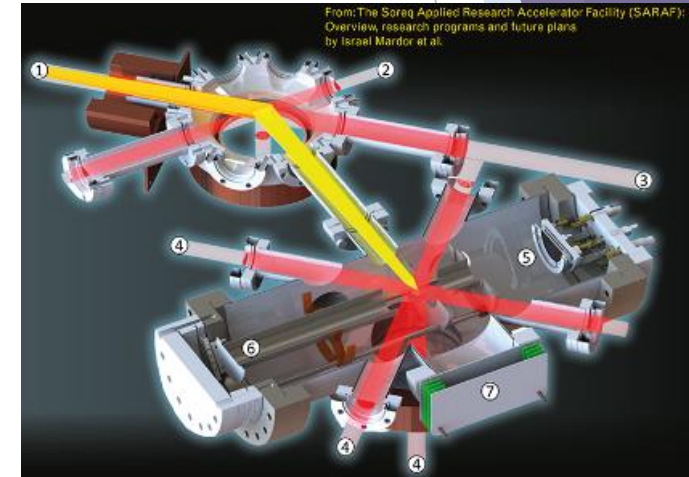
✓ Introduction



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➤ Dark Matter

➤ Lepton Flavor Violation

➤ New Weak Interactions

with nuclei...

➤ Summary

Dark Matter

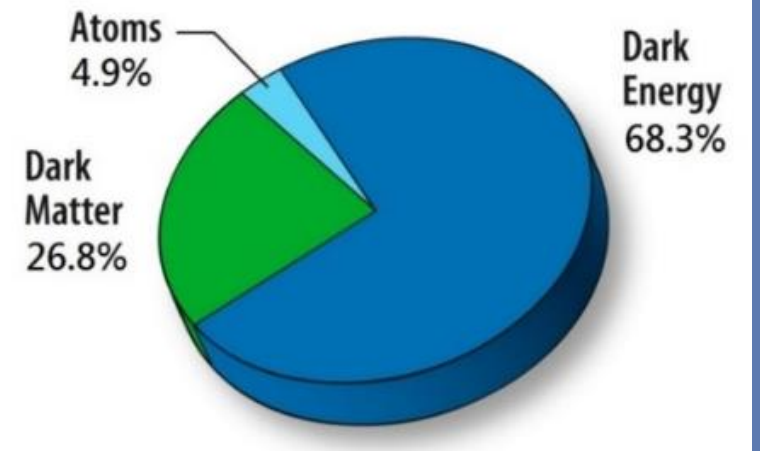
direct detection

Dark Matter Direct Detection

Promising candidates - WIMPs:
Weakly-Interacting Massive Particles

▶ Challenge - Direct detection:

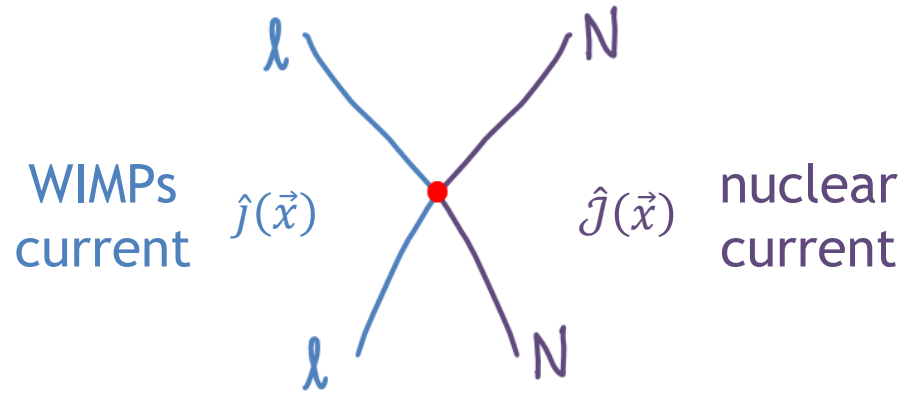
- ▶ Measuring WIMP scattering off nuclei on detectors
 - ▶ Nuclear matrix elements & structure factors
- ▶ Detection capabilities: $q \sim 100 \text{ MeV}/c$
- ▶ The structure of the coupling is determined only by symmetry considerations



q - momentum transfer

WIMPs scattering off nuclei

Low energy reaction of
WIMPs with nucleons



Non-Relativistic
Nuclear Reduction:
contact interaction
between
WIMP's & Nucleon's
currents

$$\mathcal{L}_{int} \sim \bar{\chi} O_{\chi} \chi \bar{N} O_N N$$

WIMPs scattering off nuclei

Non-Relativistic Nuclear Reduction

$$\mathcal{L}_{int} \sim \bar{\chi} O_{\chi} \chi \bar{N} O_N N$$

Scalar $\langle p(p_p) | \bar{u} d | n(p_n) \rangle = g_S(q^2) \bar{u}_p(p_p) u_n(p_n)$

$$2 \times 2 = 4$$

$$\bar{N} N$$

$$\bar{N} \gamma^5 N$$

Pseudoscalar $\langle p(p_p) | \bar{u} \gamma_5 d | n(p_n) \rangle = g_P(q^2) \bar{u}_p(p_p) \gamma_5 u_n(p_n)$

Vector $\langle p(p_p) | \bar{u} \gamma_{\mu} d | n(p_n) \rangle = \bar{u}_p(p_p) \left[g_V(q^2) \gamma_{\mu} - i \frac{\tilde{g}_{T(V)}(q^2)}{2M_N} \sigma_{\mu\nu} q^{\nu} + \frac{\tilde{g}_S(q^2)}{2M_N} q_{\mu} \right] u_n(p_n)$

Axial Vector $\langle p(p_p) | \bar{u} \gamma_{\mu} \gamma_5 d | n(p_n) \rangle = \bar{u}_p(p_p) \left[g_A(q^2) \gamma_{\mu} - i \frac{\tilde{g}_{T(A)}(q^2)}{2M_N} \sigma_{\mu\nu} q^{\nu} + \frac{\tilde{g}_P(q^2)}{2M_N} q_{\mu} \right] \gamma_5 u_n(p_n)$

$$4 \times 4 = 16$$

$$\bar{N} \frac{p^{\mu}}{m_N} N$$

$$\bar{N} \sigma^{\mu\nu} \frac{q_{\nu}}{m_N} N$$

$$\bar{N} \frac{p^{\mu}}{m_N} \gamma^5 N$$

$$\bar{N} \gamma^{\mu} \gamma^5 N$$

And similar terms for the WIMPs

WIMPs scattering off nuclei

Non-Relativistic Nuclear Reduction

$$\mathcal{L}_{int} \sim \bar{\chi} O_{\chi} \chi \bar{N} O_N N \approx \sum_{i=1}^{16} c_i O_i \bar{\chi} \chi \bar{N} N$$

$\{O_i\}_{i=1}^{16}$ - 16 non-relativistic operators

built of 4 three-vectors:

- ▶ $\frac{i\vec{q}}{m_N}$
- ▶ $\vec{v}^{\perp} \equiv \frac{\vec{P}}{2m_{\chi}} - \frac{\vec{K}}{2m_N}$
- ▶ $\vec{S}_{\chi}, \vec{S}_N$

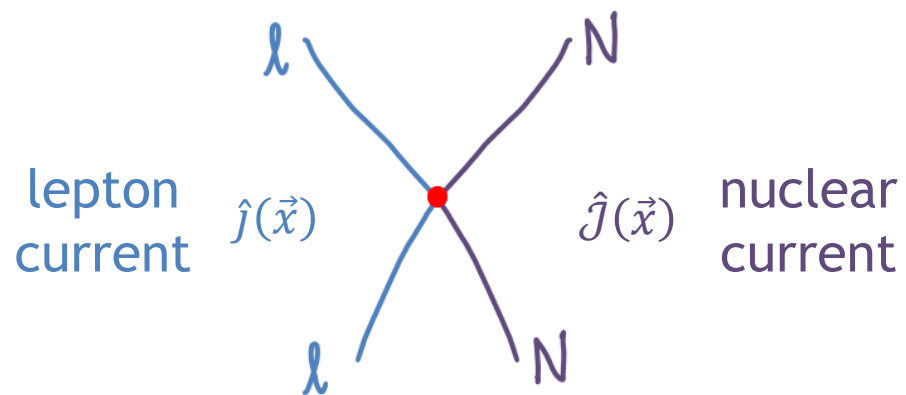
Missing tensor couplings

| j | \mathcal{L}_{int}^j | Nonrelativistic Reduction | $\sum_i c_i \mathcal{O}_i$ |
|-----|---|--|---|
| 1 | $\bar{\chi} \chi \bar{N} N$ | $1_{\chi} 1_N$ | \mathcal{O}_1 |
| 2 | $i \bar{\chi} \chi \bar{N} \gamma^5 N$ | $i \frac{\vec{q}}{m_N} \cdot \vec{S}_N$ | \mathcal{O}_{10} |
| 3 | $i \bar{\chi} \gamma^5 \chi \bar{N} N$ | $-i \frac{\vec{q}}{m_{\chi}} \cdot \vec{S}_{\chi}$ | $-\frac{m_N}{m_{\chi}} \mathcal{O}_{11}$ |
| 4 | $\bar{\chi} \gamma^5 \chi \bar{N} \gamma^5 N$ | $-\frac{\vec{q}}{m_{\chi}} \cdot \vec{S}_{\chi} \frac{\vec{q}}{m_N} \cdot \vec{S}_N$ | $-\frac{m_N}{m_{\chi}} \mathcal{O}_6$ |
| 5 | $\frac{P^{\mu}}{m_M} \bar{\chi} \chi \frac{K_{\mu}}{m_M} \bar{N} N$ | $4 \frac{m_{\chi} m_N}{m_M^2} 1_{\chi} 1_N$ | $4 \frac{m_{\chi} m_N}{m_M^2} \mathcal{O}_1$ |
| 6 | $\frac{P^{\mu}}{m_M} \bar{\chi} \chi \bar{N} i \sigma_{\mu\alpha} \frac{q^{\alpha}}{m_M} N$ | $-\frac{m_{\chi}}{m_N} \frac{\vec{q}^2}{m_M^2} 1_{\chi} 1_N - 4i \frac{m_{\chi}}{m_M} \vec{v}^{\perp} \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_N \right)$ | $-\frac{m_{\chi}}{m_N} \frac{\vec{q}^2}{m_M^2} \mathcal{O}_1 + 4 \frac{m_{\chi} m_N}{m_M^2} \mathcal{O}_3$ |
| 7 | $\frac{P^{\mu}}{m_M} \bar{\chi} \chi \bar{N} \gamma_{\mu} \gamma^5 N$ | $-4 \frac{m_{\chi}}{m_M} \vec{v}^{\perp} \cdot \vec{S}_N$ | $-4 \frac{m_{\chi}}{m_M} \mathcal{O}_7$ |
| 8 | $i \frac{P^{\mu}}{m_M} \bar{\chi} \chi \frac{K_{\mu}}{m_M} \bar{N} \gamma^5 N$ | $4i \frac{m_{\chi}}{m_M} \frac{\vec{q}}{m_M} \cdot \vec{S}_N$ | $4 \frac{m_{\chi} m_N}{m_M^2} \mathcal{O}_{10}$ |
| 9 | $\bar{\chi} i \sigma^{\mu\nu} \frac{q_{\nu}}{m_M} \chi \frac{K_{\mu}}{m_M} \bar{N} N$ | $\frac{m_N}{m_{\chi}} \frac{\vec{q}^2}{m_M^2} 1_{\chi} 1_N + 4i \frac{m_N}{m_M} \vec{v}^{\perp} \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_{\chi} \right)$ | $\frac{m_N}{m_{\chi}} \frac{\vec{q}^2}{m_M^2} \mathcal{O}_1 - 4 \frac{m_N^2}{m_M^2} \mathcal{O}_5$ |
| 10 | $\bar{\chi} i \sigma^{\mu\nu} \frac{q_{\nu}}{m_M} \chi \bar{N} i \sigma_{\mu\alpha} \frac{q^{\alpha}}{m_M} N$ | $4 \left(\frac{\vec{q}}{m_M} \times \vec{S}_{\chi} \right) \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_N \right)$ | $4 \left(\frac{\vec{q}^2}{m_M^2} \mathcal{O}_4 - \frac{m_N^2}{m_M^2} \mathcal{O}_6 \right)$ |
| 11 | $\bar{\chi} i \sigma^{\mu\nu} \frac{q_{\nu}}{m_M} \chi \bar{N} \gamma^{\mu} \gamma^5 N$ | $-4i \left(\frac{\vec{q}}{m_M} \times \vec{S}_{\chi} \right) \cdot \vec{S}_N$ | $-4 \frac{m_N}{m_M} \mathcal{O}_9$ |
| 12 | $i \bar{\chi} i \sigma^{\mu\nu} \frac{q_{\nu}}{m_M} \chi \frac{K_{\mu}}{m_M} \bar{N} \gamma^5 N$ | $\left[i \frac{\vec{q}^2}{m_{\chi} m_M} - 4 \vec{v}^{\perp} \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_{\chi} \right) \right] \frac{\vec{q}}{m_M} \cdot \vec{S}_N$ | $\frac{m_N}{m_{\chi}} \frac{\vec{q}^2}{m_M^2} \mathcal{O}_{10} + 4 \frac{\vec{q}^2}{m_M^2} \mathcal{O}_{12} + 4 \frac{m_N^2}{m_M^2} \mathcal{O}_{15}$ |
| 13 | $\bar{\chi} \gamma^{\mu} \gamma^5 \chi \frac{K_{\mu}}{m_M} \bar{N} N$ | $4 \frac{m_N}{m_M} \vec{v}^{\perp} \cdot \vec{S}_{\chi}$ | $4 \frac{m_N}{m_M} \mathcal{O}_8$ |
| 14 | $\bar{\chi} \gamma^{\mu} \gamma^5 \chi \bar{N} i \sigma_{\mu\alpha} \frac{q^{\alpha}}{m_M} N$ | $-4i \vec{S}_{\chi} \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_N \right)$ | $4 \frac{m_N}{m_M} \mathcal{O}_9$ |
| 15 | $\bar{\chi} \gamma^{\mu} \gamma^5 \chi \bar{N} \gamma^{\mu} \gamma^5 N$ | $-4 \vec{S}_{\chi} \cdot \vec{S}_N$ | $-4 \mathcal{O}_4$ |
| 16 | $i \bar{\chi} \gamma^{\mu} \gamma^5 \chi \frac{K_{\mu}}{m_M} \bar{N} \gamma^5 N$ | $4i \vec{v}^{\perp} \cdot \vec{S}_{\chi} \frac{\vec{q}}{m_M} \cdot \vec{S}_N$ | $4 \frac{m_N}{m_M} \mathcal{O}_{13}$ |
| 17 | $i \frac{P^{\mu}}{m_M} \bar{\chi} \gamma^5 \chi \frac{K_{\mu}}{m_M} \bar{N} N$ | $-4i \frac{m_N}{m_M} \frac{\vec{q}}{m_M} \cdot \vec{S}_{\chi}$ | $-4 \frac{m_N^2}{m_M^2} \mathcal{O}_{11}$ |
| 18 | $i \frac{P^{\mu}}{m_M} \bar{\chi} \gamma^5 \chi \bar{N} i \sigma_{\mu\alpha} \frac{q^{\alpha}}{m_M} N$ | $\frac{\vec{q}}{m_M} \cdot \vec{S}_{\chi} \left[i \frac{\vec{q}^2}{m_N m_M} - 4 \vec{v}^{\perp} \cdot \left(\frac{\vec{q}}{m_M} \times \vec{S}_N \right) \right]$ | $\frac{\vec{q}^2}{m_M^2} \mathcal{O}_{11} + 4 \frac{m_N^2}{m_M^2} \mathcal{O}_{15}$ |
| 19 | $i \frac{P^{\mu}}{m_M} \bar{\chi} \gamma^5 \chi \bar{N} \gamma_{\mu} \gamma^5 N$ | $4i \frac{\vec{q}}{m_M} \cdot \vec{S}_{\chi} \vec{v}^{\perp} \cdot \vec{S}_N$ | $4 \frac{m_N}{m_M} \mathcal{O}_{14}$ |
| 20 | $\frac{P^{\mu}}{m_M} \bar{\chi} \gamma^5 \chi \frac{K_{\mu}}{m_M} \bar{N} \gamma^5 N$ | $-4 \frac{\vec{q}}{m_M} \cdot \vec{S}_{\chi} \frac{\vec{q}}{m_M} \cdot \vec{S}_N$ | $-4 \frac{m_N^2}{m_M^2} \mathcal{O}_6$ |

Why do we need the tensor?

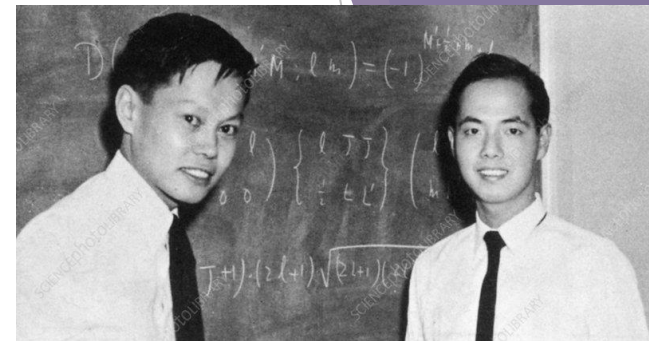
We already have 16 operator basis

Weak interaction: **Low energy reaction** of leptons with nucleons



$$\hat{H}_W \sim C \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

A-priori: {
Scalar (C_S)
PseudoScalar (C_P)
Vector (C_V)
Axial vector (C_A)
Tensor (C_T)

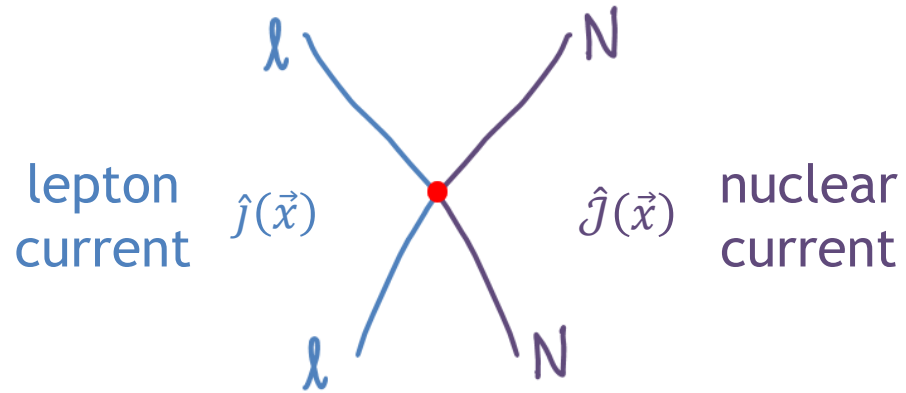


Theory: C.N. Yang and T.D. Lee (Nobel 1957)

Why do we need the tensor?

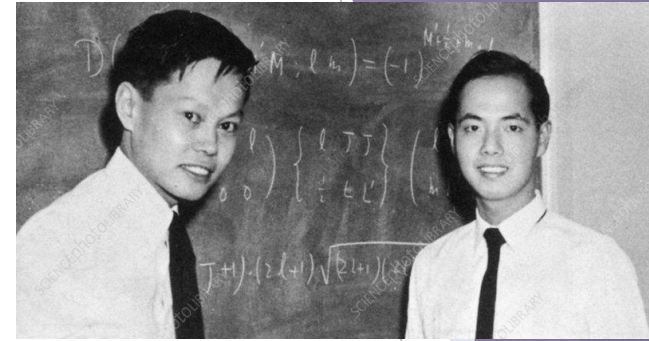
We already have 16 operator basis

Weak interaction: **Low energy reaction** of leptons with nucleons



$$\hat{H}_W \sim C \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

- A-priori:
- Scalar (C_S)
 - PseudoScalar (C_P)
 - Vector (C_V)**
 - Axial vector (C_A)**
 - Tensor (C_T)



Theory: C.N. Yang and T.D. Lee (Nobel 1957)

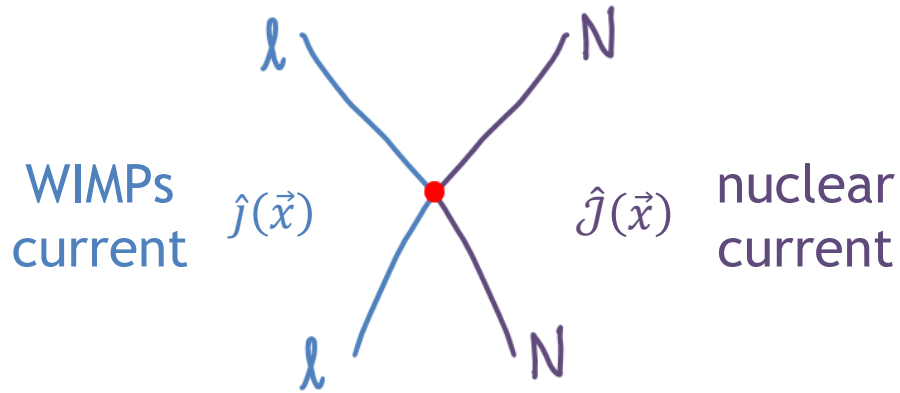


Experiment: C.S. Wu:
Parity violation in nuclear β -decays
 \Rightarrow Weak SM structure: "**V - A**"

To identify the interaction's nature, we need to know the operators & symmetries involved in each of S, P, V, A, T

How do we find the tensor NR EFT?

Low energy reaction of
WIMPs with nucleons



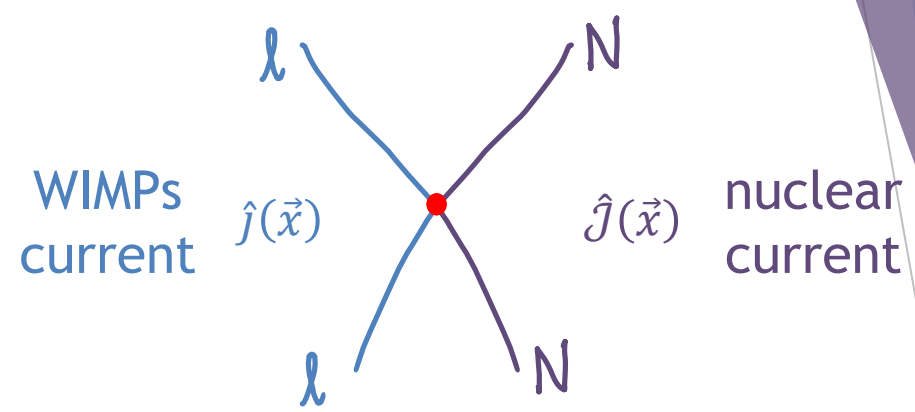
$$\hat{\mathcal{H}} \sim C \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

$$\langle k_f | \mathcal{J}_{\mu\nu}^a | k_i \rangle = \bar{u}(k_f) \frac{1}{2} \left[g_T(q^2) \sigma_{\mu\nu} + \tilde{g}_T^{(1)}(q^2) \left(\frac{q_\mu}{m_M} \gamma_\nu - \frac{q_\nu}{m_M} \gamma_\mu \right) + \tilde{g}_T^{(2)}(q^2) \left(\frac{q^\mu}{m_M} \frac{K^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{K^\mu}{m_M} \right) + \tilde{g}_T^{(3)}(q^2) \left(\gamma_\mu \frac{\not{q}}{m_M} \gamma_\nu - \gamma_\nu \frac{\not{q}}{m_M} \gamma_\mu \right) \right] \tau^a u(k_i)$$

$$4 \times 4 = 16$$

$$\begin{aligned} & \bar{N} \sigma_{\mu\nu} N \\ & \bar{N} \frac{q_\mu}{m_N} \gamma_\nu N \\ & \bar{N} \frac{q_\mu}{m_N} \frac{K_\nu}{m_N} N \\ & \bar{N} \gamma_\mu \frac{\not{q}}{m_N} \gamma_\nu N \end{aligned}$$

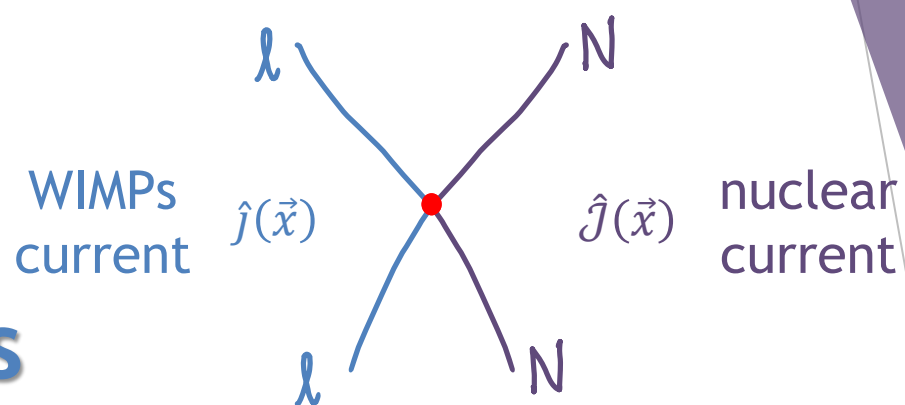
Tensor



$$\mathcal{L}_{int} \sim \hat{j}^{\mu\nu}(\vec{x}) \hat{J}_{\mu\nu}(\vec{x})$$

Tensor

→ **vector-like objects**



Tensor interactions

$$\mathcal{L}_{int} \sim \hat{j}^{\mu\nu}(\vec{x}) \hat{J}_{\mu\nu}(\vec{x})$$

▶ Symmetric:

- ▶ A space-time-metric and the stress-energy tensor

▶ Antisymmetric

- ▶ Fermionic probes

$$\Rightarrow l_{00} = 0$$

$$\Rightarrow l_{.0} = -l_0.$$

$$\Rightarrow l_{ij} \rightarrow [l_{ij}]^{(1)}$$

$$l_{\mu\nu} = \begin{pmatrix} \cancel{l_{00}} & (\leftarrow \vec{l}_0 \rightarrow) \\ \begin{pmatrix} \uparrow \\ \textcircled{-\vec{l}_0} \\ \downarrow \end{pmatrix} & \begin{pmatrix} \textcircled{\vec{l}^{(1)}} \end{pmatrix} \end{pmatrix}$$

DM Tensor Interactions

| j | $\mathcal{L}_{\text{int}}^j$ | Nonrelativistic Reduction | $\Sigma_i c_i \mathcal{O}_i$ |
|-----|---|---|--|
| 21 | $\bar{\chi} \sigma^{\mu\nu} \chi \bar{N} \sigma_{\mu\nu} N$ | $8 \frac{\vec{\sigma}_\chi \cdot \vec{\sigma}_N}{2} + o\left(\frac{1}{m^2}\right)$ | $8 \mathcal{O}_4$ |
| 22 | $\bar{\chi} \sigma^{\mu\nu} \chi \bar{N} \left(\frac{q^\mu}{m_M} \gamma_\nu - \frac{q_\nu}{m_M} \gamma^\mu \right) N$ | $-\frac{iq^2}{m_\chi m_M} 1\chi 1_N - \frac{4iq^2}{m_N m_M} \left(\frac{\vec{\sigma}_\chi \cdot \vec{\sigma}_N}{2} \right) - \frac{4}{m_M} \frac{\vec{\sigma}_\chi}{2} \cdot (\vec{q} \times \vec{v}^\perp) + \frac{4i}{m_N m_M} \left(\frac{\vec{\sigma}_N \cdot \vec{q}}{2} \right) \left(\frac{\vec{\sigma}_\chi \cdot \vec{q}}{2} \right) + o\left(\frac{1}{m^3}\right)$ | $-i \frac{q^2}{m_M m_\chi} \mathcal{O}_1 - 4i \frac{q^2}{m_M m_N} \mathcal{O}_4 + 4i \frac{m_N}{m_M} \mathcal{O}_5 + 4i \frac{m_N}{m_M} \mathcal{O}_6$ |
| 23 | $\bar{\chi} \sigma^{\mu\nu} \chi \bar{N} \left(\frac{q^\mu}{m_M} \frac{K^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{K^\mu}{m_M} \right) N$ | $-2i \frac{m_N}{m_\chi} \frac{q^2}{m_M^2} 1\chi 1_N - 8 \frac{m_N}{m_M^2} \frac{\vec{\sigma}_\chi}{2} \cdot (\vec{q} \times \vec{v}^\perp) + o\left(\frac{1}{m^4}\right)$ | $-2i \frac{m_N}{m_\chi} \frac{q^2}{m_M^2} \mathcal{O}_1 + 8i \frac{m_N}{m_M^2} \mathcal{O}_5$ |
| 24 | $\bar{\chi} \sigma^{\mu\nu} \chi \bar{N} \left(\gamma_\mu \frac{q^\nu}{m_M} \gamma_\nu - \gamma_\nu \frac{q^\mu}{m_M} \gamma_\mu \right) N$ | $8i \left(\frac{\vec{\sigma}_\chi \cdot \vec{q}}{m_M} \right) \left(\frac{\vec{\sigma}_N \cdot \vec{v}^\perp}{2} \right) + o\left(\frac{1}{m^3}\right)$ | $8 \frac{m_N}{m_M} \mathcal{O}_{14}$ |
| 25 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \gamma^\nu - \frac{q^\nu}{m_M} \gamma^\mu \right) \chi \bar{N} \sigma_{\mu\nu} N$ | $\frac{iq^2}{m_N m_M} 1\chi 1_N + \frac{4}{m_M} \frac{\vec{\sigma}_N}{2} \cdot (\vec{q} \times \vec{v}^\perp) + \frac{4i}{m_\chi m_M} q^2 \left(\frac{\vec{\sigma}_\chi \cdot \vec{\sigma}_N}{2} \right) - \frac{4i}{m_\chi m_M} (\vec{q} \cdot \frac{\vec{\sigma}_\chi}{2}) (\vec{q} \cdot \frac{\vec{\sigma}_N}{2}) + o\left(\frac{1}{m^4}\right)$ | $i \frac{q^2}{m_N m_M} \mathcal{O}_1 - 4i \frac{m_N}{m_M} \mathcal{O}_3 + 4i \frac{q^2}{m_\chi m_M} \mathcal{O}_4 - 4i \frac{m_N}{m_\chi m_M} \mathcal{O}_6$ |
| 26 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \gamma^\nu - \frac{q^\nu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \gamma_\nu - \frac{q_\nu}{m_M} \gamma_\mu \right) N$ | $-\frac{iq^2}{m_\chi m_M} 1\chi 1_N - \frac{4iq^2}{m_N m_M} \left(\frac{\vec{\sigma}_\chi \cdot \vec{\sigma}_N}{2} \right) - \frac{4}{m_M} \frac{\vec{\sigma}_\chi}{2} \cdot (\vec{q} \times \vec{v}^\perp) + \frac{4i}{m_N m_M} \left(\frac{\vec{\sigma}_N \cdot \vec{q}}{2} \right) \left(\frac{\vec{\sigma}_\chi \cdot \vec{q}}{2} \right) + o\left(\frac{1}{m^4}\right)$ | $-i \frac{q^2}{m_\chi m_M} \mathcal{O}_1 - 4i \frac{q^2}{m_N m_M} \mathcal{O}_4 + 4i \frac{m_N}{m_M} \mathcal{O}_5 + 4i \frac{m_N}{m_M} \mathcal{O}_6$ |
| 27 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \gamma^\nu - \frac{q^\nu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \frac{K^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{K^\mu}{m_M} \right) N$ | $-4 \frac{m_N}{m_M} \frac{q^2}{m_M^2} 1\chi 1_N + o\left(\frac{1}{m^4}\right)$ | $-4 \frac{m_N}{m_M} \frac{q^2}{m_M^2} \mathcal{O}_1$ |
| 28 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \gamma^\nu - \frac{q^\nu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\gamma_\mu \frac{q^\nu}{m_M} \gamma_\nu - \gamma_\nu \frac{q^\mu}{m_M} \gamma_\mu \right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 29 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \frac{P^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{P^\mu}{m_M} \right) \chi \bar{N} \sigma_{\mu\nu} N$ | $2i \frac{m_\chi}{m_N} \frac{q^2}{m_M^2} 1\chi 1_N + 8 \frac{m_\chi}{m_M^2} \frac{\vec{\sigma}_N}{2} \cdot (\vec{q} \times \vec{v}^\perp) + o\left(\frac{1}{m^4}\right)$ | $2i \frac{m_\chi}{m_N} \frac{q^2}{m_M^2} \mathcal{O}_1 - 8i \frac{m_\chi m_N}{m_M^2} \mathcal{O}_3$ |
| 30 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \frac{P^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{P^\mu}{m_M} \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \gamma_\nu - \frac{q_\nu}{m_M} \gamma_\mu \right) N$ | $-4 \frac{m_\chi}{m_M} \frac{q^2}{m_M^2} 1\chi 1_N + o\left(\frac{1}{m^4}\right)$ | $-4 \frac{m_\chi}{m_M} \frac{q^2}{m_M^2} \mathcal{O}_1$ |
| 31 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \frac{P^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{P^\mu}{m_M} \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \frac{K^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{K^\mu}{m_M} \right) N$ | $-8 \frac{m_\chi m_N}{m_M^2} \frac{q^2}{m_M^2} 1\chi 1_N + o\left(\frac{1}{m^4}\right)$ | $-8 \frac{m_\chi m_N}{m_M^2} \frac{q^2}{m_M^2} \mathcal{O}_1$ |
| 32 | $\bar{\chi} \left(\frac{q^\mu}{m_M} \frac{P^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{P^\mu}{m_M} \right) \chi \bar{N} \left(\gamma_\mu \frac{q^\nu}{m_M} \gamma_\nu - \gamma_\nu \frac{q^\mu}{m_M} \gamma_\mu \right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 33 | $\bar{\chi} \left(\gamma^\mu \frac{q^\nu}{m_M} \gamma^\nu - \gamma^\nu \frac{q^\mu}{m_M} \gamma^\mu \right) \chi \bar{N} \sigma_{\mu\nu} N$ | $-8i \left(\frac{\vec{\sigma}_N \cdot \vec{q}}{m_M} \right) \left(\frac{\vec{\sigma}_\chi \cdot \vec{v}^\perp}{2} \right) + o\left(\frac{1}{m^3}\right)$ | $-8 \frac{m_N}{m_M} \mathcal{O}_{13}$ |
| 34 | $\bar{\chi} \left(\gamma^\mu \frac{q^\nu}{m_M} \gamma^\nu - \gamma^\nu \frac{q^\mu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \gamma_\nu - \frac{q_\nu}{m_M} \gamma_\mu \right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 35 | $\bar{\chi} \left(\gamma^\mu \frac{q^\nu}{m_M} \gamma^\nu - \gamma^\nu \frac{q^\mu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\frac{q^\mu}{m_M} \frac{K^\nu}{m_M} - \frac{q^\nu}{m_M} \frac{K^\mu}{m_M} \right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 36 | $\bar{\chi} \left(\gamma^\mu \frac{q^\nu}{m_M} \gamma^\nu - \gamma^\nu \frac{q^\mu}{m_M} \gamma^\mu \right) \chi \bar{N} \left(\gamma_\mu \frac{q^\nu}{m_M} \gamma_\nu - \gamma_\nu \frac{q^\mu}{m_M} \gamma_\mu \right) N$ | $\frac{32}{m_M^2} \left[q^2 \left(\frac{\vec{\sigma}_\chi \cdot \vec{\sigma}_N}{2} \right) - \left(\frac{\vec{\sigma}_N \cdot \vec{q}}{2} \right) \left(\frac{\vec{\sigma}_\chi \cdot \vec{q}}{2} \right) \right] + o\left(\frac{1}{m^4}\right)$ | $32 \frac{q^2}{m_M^2} \mathcal{O}_4 - 32 \frac{m_N^2}{m_M^2} \mathcal{O}_6$ |

To identify the interaction's nature we need to know the operators & symmetries involved in each of S, P, V, A, T

$$\mathcal{O}_1 \equiv 1_\chi 1_N,$$

$$\mathcal{O}_3 \equiv i \vec{S}_N \cdot \left(\frac{\vec{q}}{m_N} \times \vec{v}^\perp \right),$$

$$\mathcal{O}_4 \equiv \vec{S}_\chi \cdot \vec{S}_N,$$

$$\mathcal{O}_5 \equiv i \vec{S}_\chi \cdot \left(\frac{\vec{q}}{m_N} \times \vec{v}^\perp \right),$$

$$\mathcal{O}_6 \equiv \left(\vec{S}_\chi \cdot \frac{\vec{q}}{m_N} \right) \left(\vec{S}_N \cdot \frac{\vec{q}}{m_N} \right),$$

$$\mathcal{O}_{13} \equiv i \left(\vec{S}_\chi \cdot \vec{v}^\perp \right) \left(\vec{S}_N \cdot \frac{\vec{q}}{m_N} \right),$$

$$\mathcal{O}_{14} \equiv i \left(\vec{S}_\chi \cdot \frac{\vec{q}}{m_N} \right) \left(\vec{S}_N \cdot \vec{v}^\perp \right).$$

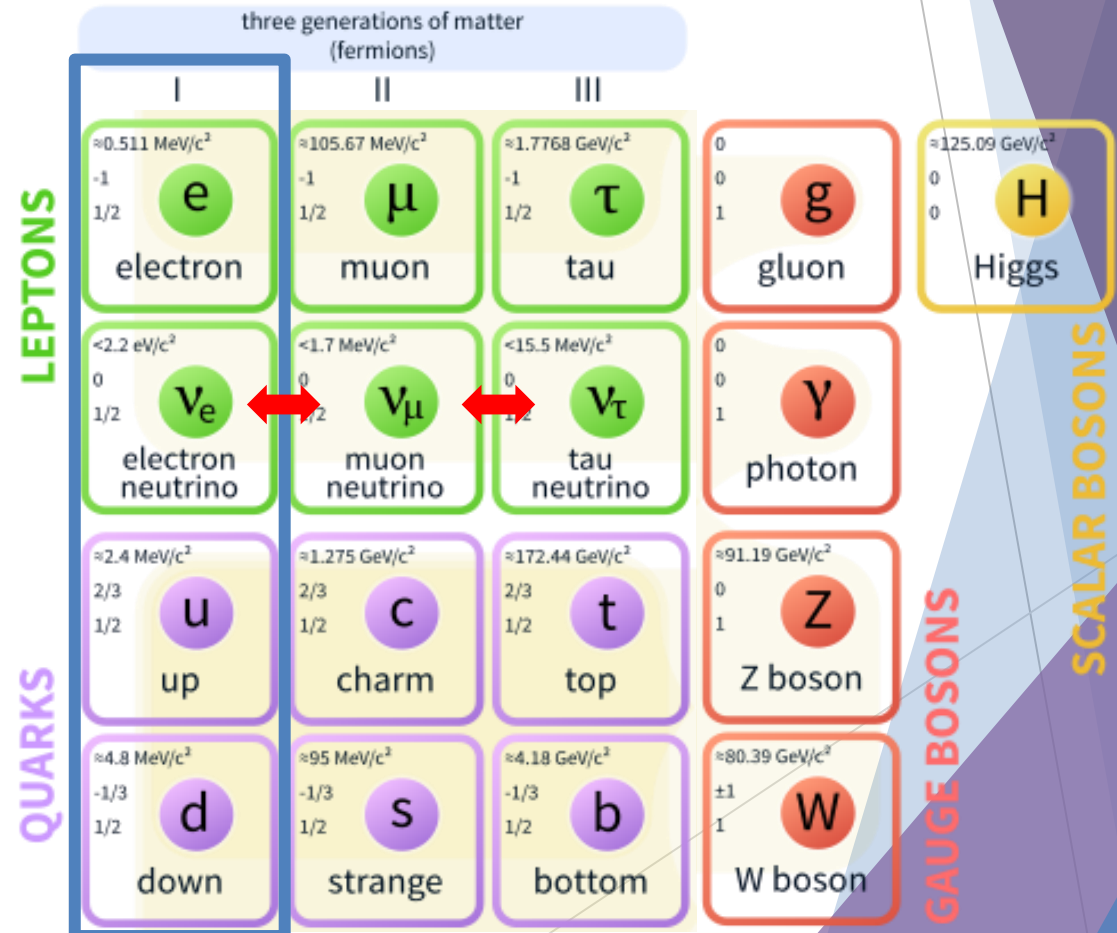


Lepton Flavor Violation

$\mu \rightarrow e$ conversion

Beyond Standard Model (BSM)

Elementary Particles



NOBEL PRIZE IN PHYSICS 2015

The Nobel Prize in Physics 2015 was awarded to **Takaaki Kajita** and **Arthur B. McDonald** for discovery of neutrino oscillations, which shows neutrinos have mass.

WHAT IS A NEUTRINO?

Neutrinos are tiny subatomic particles, produced by nuclear reactions that take place in stars, including our sun, as well as in radioactive decay processes. They come in three 'flavours'.



ELECTRON NEUTRINO



MUON NEUTRINO



TAU NEUTRINO



NOBEL PRIZE



The nuclear reactions in the sun produce neutrinos, which we can detect.

The number of neutrinos detected was only a third of the expected value.

Neutrinos 'flip' between the three flavours, and only one type was being detected.

WHY DOES IT MATTER?

If neutrinos oscillate between types, they must have mass, even if this mass is incredibly small. This contradicts the standard model of particle physics, which states they are massless.



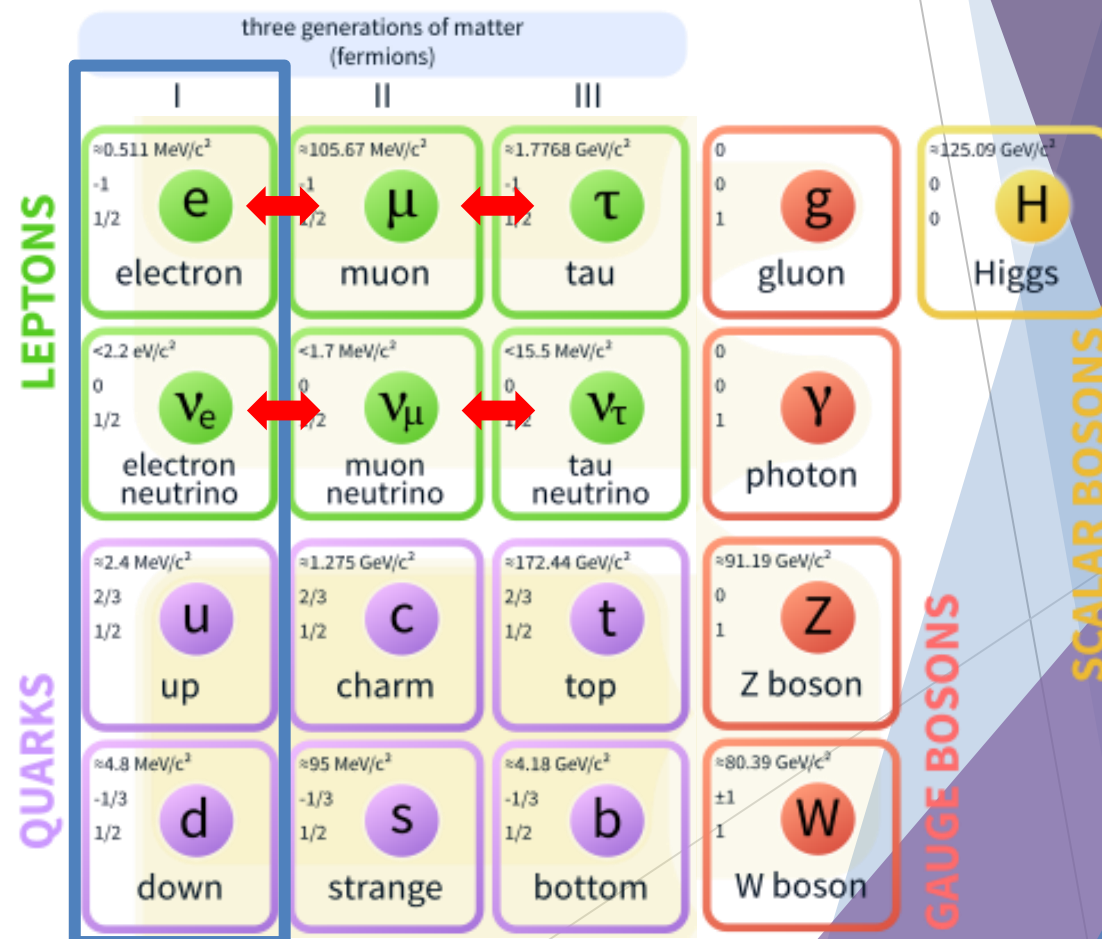
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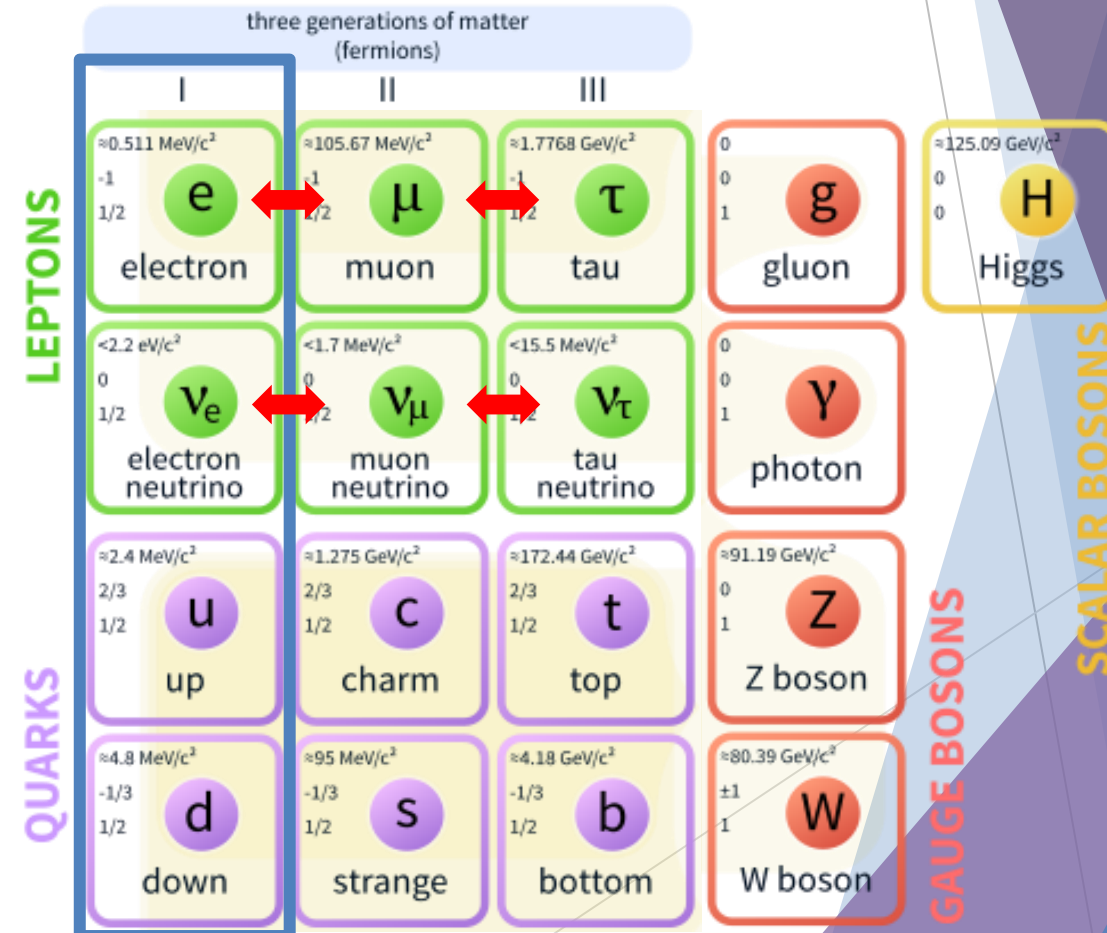
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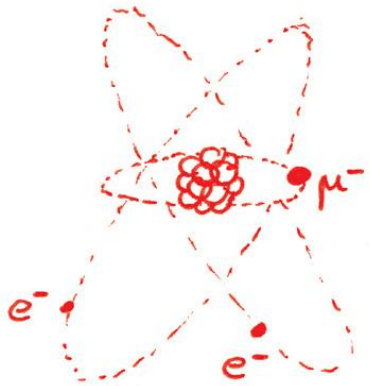
Charged Lepton Flavor Violation

Beyond Standard Model (BSM) with nuclei...

Elementary Particles



This is what we start with.



(Credit: symmetry magazine)

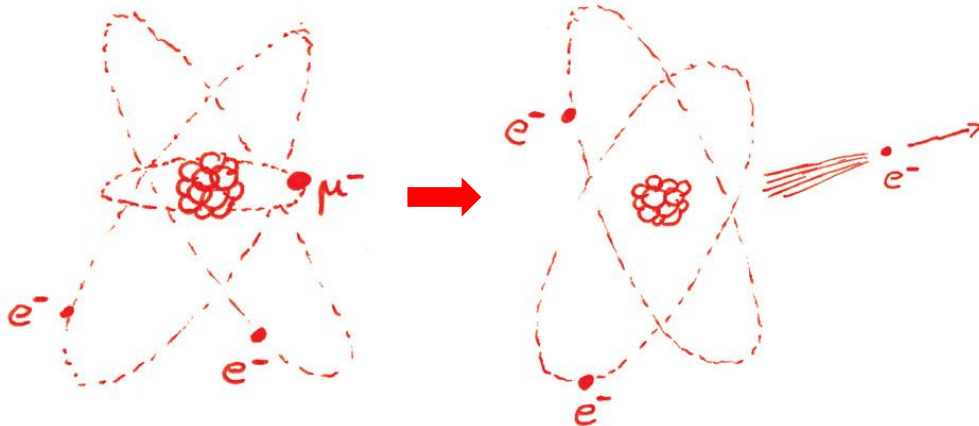
Charged Lepton Flavor Violation

Beyond Standard Model (BSM) with nuclei...

$\mu \rightarrow e$ conversion

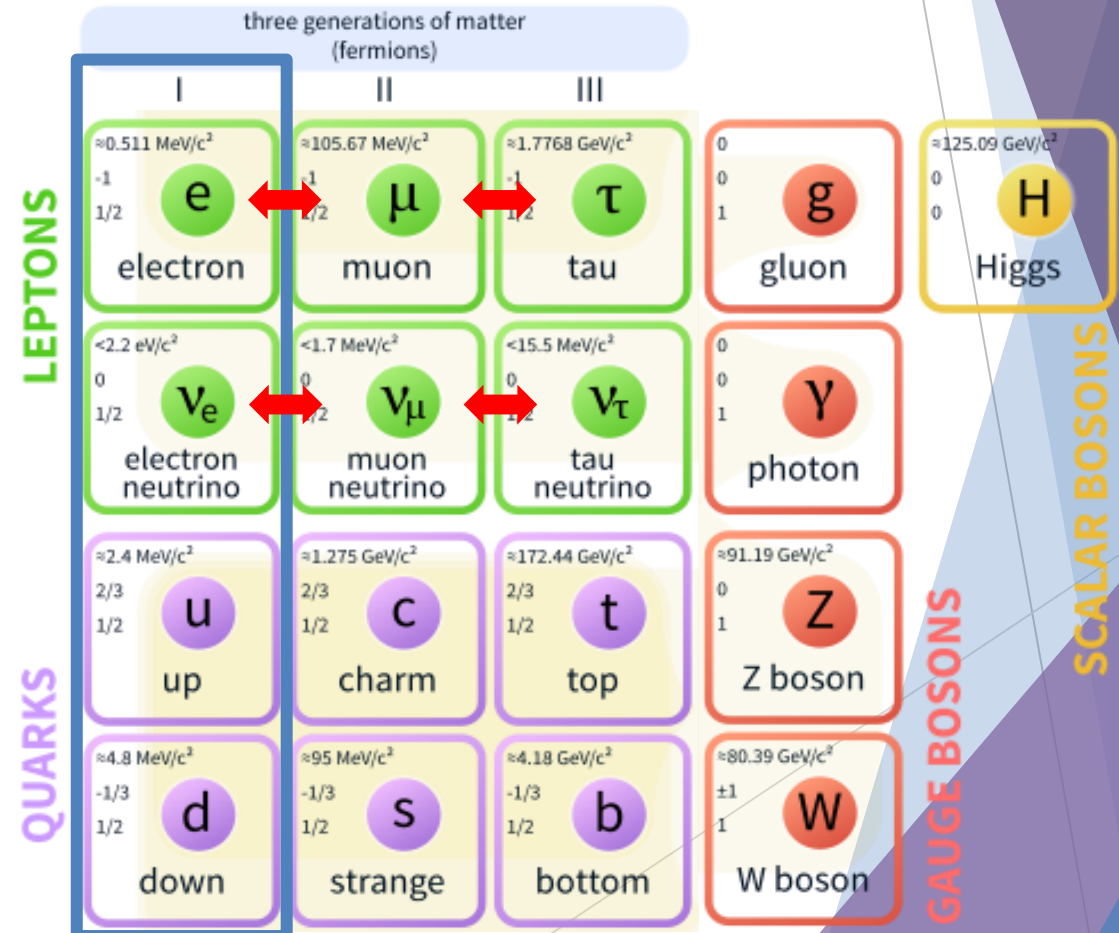
This is what we start with.

This is the process we are looking for.



(Credit: symmetry magazine)

Elementary Particles

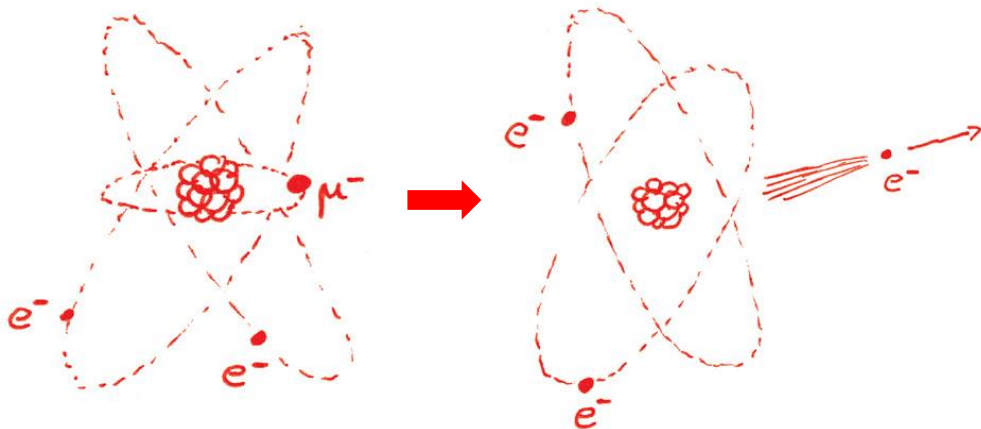


Charged Lepton Flavor Violation

$\mu \rightarrow e$ conversion

This is what we start with.

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(Credit: symmetry magazine)

TABLE IX. Existing limits on branching ratios for $\mu \rightarrow e$ conversion, taken from the tabulation of [75].

| Process | Limit | Lab/Reference |
|---|-----------------------|---------------|
| $\mu^- + {}^{32}\text{S} \rightarrow e^- + {}^{32}\text{S}$ | 7×10^{-11} | SIN [76] |
| $\mu^- + \text{Ti} \rightarrow e^- + \text{Ti}$ | 1.6×10^{-11} | TRIUMF [77] |
| $\mu^- + \text{Ti} \rightarrow e^- + \text{Ti}$ | 4.6×10^{-12} | TRIUMF [78] |
| $\mu^- + \text{Ti} \rightarrow e^- + \text{Ti}$ | 4.3×10^{-12} | PSI [79] |
| $\mu^- + \text{Ti} \rightarrow e^- + \text{Ti}$ | 6.1×10^{-13} | PSI [80] |
| $\mu^- + \text{Cu} \rightarrow e^- + \text{Cu}$ | 1.6×10^{-8} | SREL [81] |
| $\mu^- + \text{Au} \rightarrow e^- + \text{Au}$ | 7×10^{-13} | PSI [82] |
| $\mu^- + \text{Pb} \rightarrow e^- + \text{Pb}$ | 4.9×10^{-10} | TRIUMF [78] |
| $\mu^- + \text{Pb} \rightarrow e^- + \text{Pb}$ | 4.6×10^{-11} | PSI [83] |

branching ratio with respect to muon capture in the same nucleus

- ▶ Future experiments: mu2e @ Fermilab, COMET @ J-PARC
 $({}^{27}\text{Al}) \sim 10^{-17}$

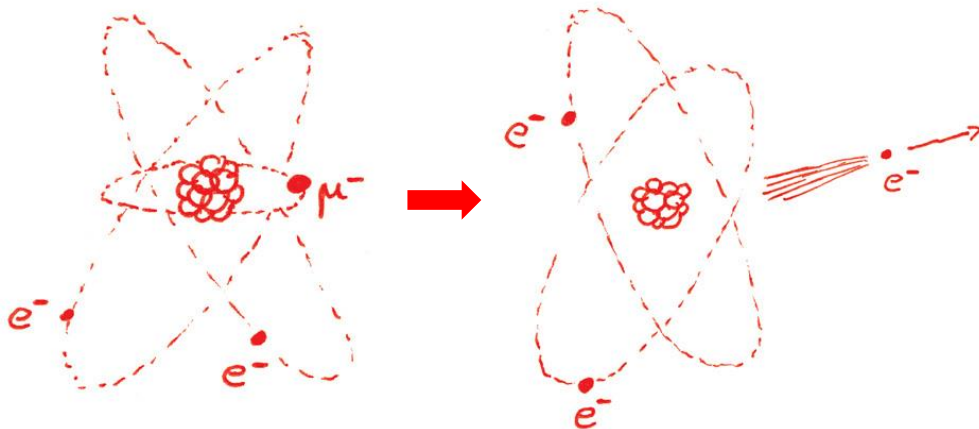
4 orders of magnitude enhancement!

- ▶ $q \sim m_\mu$
- ▶ The electron is “fully relativistic”

$\mu \rightarrow e$ conversion

This is what we start with.

This is the process we are looking for.



(Credit: symmetry magazine)

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branching ratio with respect to muon capture in the same nucleus

- ▶ Future experiments: mu2e @ Fermilab, COMET @ J-PARC
 $({}^{27}\text{Al}) \sim 10^{-17}$

NREFT Missing tensor couplings

4 orders of magnitude enhancement!

$\mu \rightarrow e$ Tensor Interactions

New operators!
Easier for identifying the nature of the CLFV

| j | $\mathcal{L}_{\text{int}}^j$ | Pauli Operator Reduction | $\sum_i c_i \mathcal{O}_i$ |
|-----|---|--|---|
| 21 | $\bar{\chi} e \sigma^{\mu\nu} \chi_\mu \bar{N} \sigma_{\mu\nu} N$ | $-\frac{q}{m_N} 1_L 1_N - 2i 1_L \hat{q} \cdot (\vec{v}_N \times \vec{\sigma}_N) + 2\vec{\sigma}_L \cdot \vec{\sigma}_N + 2\vec{\sigma}_L \cdot [\hat{q} \times (\vec{v}_N \times \vec{\sigma}_N)]$ $+ i(\hat{q} \times \vec{v}_\mu) \cdot \vec{\sigma}_N - (\vec{v}_\mu \cdot \vec{\sigma}_L)(\hat{q} \cdot \vec{\sigma}_N) - [\hat{q} \times (\vec{v}_\mu \times \vec{\sigma}_L)] \cdot \vec{\sigma}_N + o\left(\frac{q^2}{m_N^2}\right)$ | $-\frac{q}{m_N} \mathcal{O}_1 - 2\mathcal{O}_3 + 2\mathcal{O}_4 - 2i\mathcal{O}_{13}$ $+ \mathcal{O}_5^f + 2i\mathcal{O}_{14}^f + 2i\mathcal{O}_{13}^{f'}$ |
| 22 | $\bar{\chi} e \sigma^{\mu\nu} \chi_\mu \bar{N} \left(\frac{q\mu}{m_N} \gamma_\nu - \frac{q\nu}{m_N} \gamma_\mu\right) N$ | $-2i\frac{q}{m_N} 1_L 1_N + i\frac{q^2}{m_N^2} (\vec{\sigma}_L \cdot \vec{\sigma}_N) - 2\frac{q}{m_N} \vec{\sigma}_L \cdot (\hat{q} \times \vec{v}_N)$ $+ i\frac{q^2}{m_N^2} (\hat{q} \cdot \vec{\sigma}_N)(\vec{\sigma}_L \cdot \hat{q}) + i\frac{q}{m_N} (\vec{v}_\mu \cdot \hat{q}) 1_N - \frac{q}{m_N} \hat{q} \cdot (\vec{v}_\mu \times \vec{\sigma}_L) 1_N + o\left(\frac{q^3}{m_N^3}\right)$ | $-2i\frac{q}{m_N} \mathcal{O}_1 + i\frac{q^2}{m_N^2} \mathcal{O}_4 + 2i\frac{q}{m_N} \mathcal{O}_5$ $-i\frac{q^2}{m_N^2} \mathcal{O}_6 + 2i\frac{q}{m_N} \mathcal{O}_2^{f'} + 2i\frac{q}{m_N} \mathcal{O}_3^f$ |
| 23 | $\bar{\chi} e \sigma^{\mu\nu} \chi_\mu \bar{N}' \left(\frac{q\mu}{m_N} v_{N\nu} - \frac{q\nu}{m_N} v_{N\mu}\right) N$ | $\frac{q}{m_N} [-2i 1_L \cdot 1_N + 2\vec{\sigma}_L \cdot (\hat{q} \times \vec{v}_N) + i(\hat{q} \cdot \vec{v}_\mu) 1_N - \hat{q} \cdot (\vec{v}_\mu \times \vec{\sigma}_L) 1_N] + o\left(\frac{q^3}{m_N^3}\right)$ | $2\frac{q}{m_N} (-i\mathcal{O}_1 - i\mathcal{O}_5 + \mathcal{O}_2^{f'} + i\mathcal{O}_3^f)$ |
| 24 | $\bar{\chi} e \sigma^{\mu\nu} \chi_\mu \bar{N} \left(\gamma_\mu \frac{q}{m_N} \gamma_\nu - \gamma_\nu \frac{q}{m_N} \gamma_\mu\right) N$ | $-4i\frac{q}{m_N} [(\vec{\sigma}_L \cdot \vec{\sigma}_N) + 4i\frac{q}{m_N} (\vec{\sigma}_L \cdot \hat{q})(\vec{\sigma}_N \cdot \hat{q}) + i(\hat{q} \cdot \vec{\sigma}_L)(\vec{v}_N \cdot \vec{\sigma}_N)]$ $-4i\frac{q}{m_N} \{i(\hat{q} \times \frac{\vec{v}_\mu}{2}) \cdot \vec{\sigma}_N - [\hat{q} \times (\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L)] \cdot \vec{\sigma}_N\} + o\left(\frac{q^3}{m_N^3}\right)$ | $-4i\frac{q}{m_N} (\mathcal{O}_4 + \mathcal{O}_6 - i\mathcal{O}_{14})$ $-4i\frac{q}{m_N} (\mathcal{O}_5^f + i\mathcal{O}_{13}^{f'})$ |
| 25 | $\bar{\chi} e \left(\frac{q\mu}{m_L} \gamma^\nu - \frac{q\nu}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \sigma_{\mu\nu} N$ | $\frac{q}{m_L} \{2i(\vec{\sigma}_L \cdot \vec{\sigma}_N) - 2i(\hat{q} \cdot \vec{\sigma}_L)(\hat{q} \cdot \vec{\sigma}_N) + \frac{i}{2}\frac{q}{m_N} 1_L 1_N - 1_L \hat{q} \cdot (\vec{v}_N \times \vec{\sigma}_N)\}$ $+ \frac{q}{m_L} \{(\hat{q} \times \vec{v}_\mu) \cdot \vec{\sigma}_N + i[\hat{q} \times (\vec{v}_\mu \times \vec{\sigma}_L)] \cdot \vec{\sigma}_N\} + o\left(\frac{q}{m_L} \frac{q^2}{m_N^2}\right)$ | $i\frac{q}{m_L} (2\mathcal{O}_4 + 2\mathcal{O}_6 + \frac{1}{2}\frac{q}{m_N} \mathcal{O}_1 + \mathcal{O}_3)$ $+ 2\frac{q}{m_L} (-i\mathcal{O}_5^f + \mathcal{O}_{13}^{f'})$ |
| 26 | $\bar{\chi} e \left(\frac{q\mu}{m_L} \gamma^\nu - \frac{q\nu}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \left(\frac{q\mu}{m_N} \gamma_\nu - \frac{q\nu}{m_N} \gamma_\mu\right) N$ | $\frac{q^2}{m_L m_N} [-1_L 1_N + \frac{q}{m_N} (\vec{\sigma}_L \cdot \vec{\sigma}_N) - \frac{q}{m_N} (\hat{q} \cdot \vec{\sigma}_N)(\hat{q} \cdot \vec{\sigma}_L) + 2i\vec{\sigma}_L \cdot (\hat{q} \times \vec{v}_N)]$ $+ \frac{q^2}{m_L m_N} [- (\hat{q} \cdot \vec{v}_\mu) 1_N - i\hat{q} \cdot (\vec{v}_\mu \times \vec{\sigma}_L) 1_N] + o\left(\frac{q}{m_L} \frac{q^3}{m_N^3}\right)$ | $\frac{q^2}{m_L m_N} (-\mathcal{O}_1 + \frac{q}{m_N} \mathcal{O}_4 + \frac{q}{m_N} \mathcal{O}_6 + 2\mathcal{O}_5)$ $+ \frac{q^2}{m_L m_N} (2i\mathcal{O}_2^{f'} - 2\mathcal{O}_3^f)$ |
| 27 | $\bar{\chi} e \left(\frac{q\mu}{m_L} \gamma^\nu - \frac{q\nu}{m_L} \gamma^\mu\right) \chi_\mu \bar{N}' \left(\frac{q\mu}{m_N} v_{N\nu} - \frac{q\nu}{m_N} v_{N\mu}\right) N$ | $\frac{q}{m_L} \frac{q}{m_N} \{-1_L 1_N + 2i\vec{\sigma}_L \cdot (\hat{q} \times \vec{v}_N)\}$ $- \frac{q}{m_L} \frac{q}{m_N} \{(\hat{q} \cdot \vec{v}_\mu) 1_N + i\hat{q} \cdot (\vec{v}_\mu \times \vec{\sigma}_L) 1_N\} + o\left(\frac{q}{m_L} \frac{q^3}{m_N^3}\right)$ | $\frac{q^2}{m_L m_N} (-\mathcal{O}_1 + 2\mathcal{O}_5)$ $+ \frac{q^2}{m_L m_N} (2i\mathcal{O}_2^{f'} - 2\mathcal{O}_3^f)$ |
| 28 | $\bar{\chi} e \left(\frac{q\mu}{m_L} \gamma^\nu - \frac{q\nu}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \left(\gamma_\mu \frac{q}{m_N} \gamma_\nu - \gamma_\nu \frac{q}{m_N} \gamma_\mu\right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 29 | $\bar{\chi} e \left(\frac{q\alpha}{m_L} v_\mu^\nu - \frac{q\nu}{m_L} v_\mu^\alpha\right) \chi_\mu \bar{N} \sigma_{\alpha\nu} N$ | $2\frac{q}{m_L} \left\{\frac{i}{2}\frac{q}{m_N} 1_L 1_N - 1_L \hat{q} \cdot (\vec{v}_N \times \vec{\sigma}_N) + (\hat{q} \times \vec{v}_\mu) \cdot \vec{\sigma}_N\right\} + o\left(\frac{q}{m_L} \frac{q^2}{m_N^2}\right)$ | $i\frac{q}{m_L} \left(\frac{q}{m_N} \mathcal{O}_1 + 2\mathcal{O}_3 - 4\mathcal{O}_5^f\right)$ |
| 30 | $\bar{\chi} e \left(\frac{q\alpha}{m_L} v_\mu^\nu - \frac{q\nu}{m_L} v_\mu^\alpha\right) \chi_\mu \bar{N} \left(\frac{q\alpha}{m_N} \gamma_\nu - \frac{q\nu}{m_N} \gamma_\alpha\right) N$ | $2\frac{q^2}{m_L m_N} \left\{-1_L 1_N + (\hat{q} \cdot \frac{\vec{v}_\mu}{2}) 1_N + i\hat{q} \cdot \left(\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L\right) 1_N\right\} + o\left(\frac{q}{m_L} \frac{q^3}{m_N^3}\right)$ | $2\frac{q^2}{m_L m_N} (-\mathcal{O}_1 - i\mathcal{O}_2^{f'} + \mathcal{O}_3^f)$ |
| 31 | $\bar{\chi} e \left(\frac{q\alpha}{m_L} v_\mu^\nu - \frac{q\nu}{m_L} v_\mu^\alpha\right) \chi_\mu \bar{N}' \left(\frac{q\alpha}{m_N} v_{N\nu} - \frac{q\nu}{m_N} v_{N\alpha}\right) N$ | $2\frac{q^2}{m_L m_N} \left\{-1_L 1_N + (\hat{q} \cdot \frac{\vec{v}_\mu}{2}) 1_N + i\hat{q} \cdot \left(\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L\right) 1_N\right\} + o\left(\frac{q}{m_L} \frac{q^3}{m_N^3}\right)$ | $2\frac{q^2}{m_L m_N} (-\mathcal{O}_1 - i\mathcal{O}_2^{f'} + \mathcal{O}_3^f)$ |
| 32 | $\bar{\chi} e \left(\frac{q\alpha}{m_L} v_\mu^\nu - \frac{q\nu}{m_L} v_\mu^\alpha\right) \chi_\mu \bar{N} \left(\gamma_\alpha \frac{q}{m_N} \gamma_\nu - \gamma_\nu \frac{q}{m_N} \gamma_\alpha\right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 33 | $\bar{\chi} e \left(\gamma^\mu \frac{q}{m_L} \gamma^\nu - \gamma^\nu \frac{q}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \sigma_{\mu\nu} N$ | $-4i\frac{q}{m_L} \{(\hat{q} \cdot \vec{\sigma}_L)(\hat{q} \cdot \vec{\sigma}_N) - \vec{\sigma}_L \cdot [\hat{q} \times (\vec{v}_N \times \vec{\sigma}_N)]\} + o\left(\frac{q}{m_L} \frac{q^2}{m_N^2}\right)$ $-4i\frac{q}{m_L} \left(\frac{\vec{v}_\mu}{2} \cdot \vec{\sigma}_L\right) (\hat{q} \cdot \vec{\sigma}_N) + o\left(\frac{q}{m_L} \frac{q^2}{m_N^2}\right)$ | $4\frac{q}{m_L} (i\mathcal{O}_6 + \mathcal{O}'_{13})$ $-4\frac{q}{m_L} \mathcal{O}_{14}^f$ |
| 34 | $\bar{\chi} e \left(\gamma^\mu \frac{q}{m_L} \gamma^\nu - \gamma^\nu \frac{q}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \left(\frac{q\mu}{m_N} \gamma_\nu - \frac{q\nu}{m_N} \gamma_\mu\right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 35 | $\bar{\chi} e \left(\gamma^\mu \frac{q}{m_L} \gamma^\nu - \gamma^\nu \frac{q}{m_L} \gamma^\mu\right) \chi_\mu \bar{N}' \left(\frac{q\mu}{m_N} v_{N\nu} - \frac{q\nu}{m_N} v_{N\mu}\right) N$ | $o\left(\frac{1}{m^6}\right)$ | |
| 36 | $\bar{\chi} e \left(\gamma^\mu \frac{q}{m_L} \gamma^\nu - \gamma^\nu \frac{q}{m_L} \gamma^\mu\right) \chi_\mu \bar{N} \left(\gamma_\mu \frac{q}{m_N} \gamma_\nu - \gamma_\nu \frac{q}{m_N} \gamma_\mu\right) N$ | $8\frac{q^2}{m_L m_N} \{(\vec{\sigma}_L \cdot \vec{\sigma}_N) - (\hat{q} \cdot \vec{\sigma}_L)(\hat{q} \cdot \vec{\sigma}_N) - (\hat{q} \cdot \vec{\sigma}_L)(\vec{v}_N \cdot \vec{\sigma}_N)\}$ $+ 8\frac{q^2}{m_L m_N} \{-i(\hat{q} \times \frac{\vec{v}_\mu}{2}) \cdot \vec{\sigma}_N + [\hat{q} \times (\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L)] \cdot \vec{\sigma}_N\} + o\left(\frac{q}{m_L} \frac{q^3}{m_N^3}\right)$ | $8\frac{q^2}{m_L m_N} (\mathcal{O}_4 + \mathcal{O}_6 + i\mathcal{O}_{14})$ $-8\frac{q^2}{m_L m_N} (\mathcal{O}_5^f + i\mathcal{O}_{13}^{f'})$ |

- $\mathcal{O}_1 \equiv 1_L 1_N,$
- $\mathcal{O}_3 \equiv 1_L i\hat{q} \cdot (\vec{v}_N \times \vec{\sigma}_N)$
- $\mathcal{O}_4 \equiv \vec{\sigma}_L \cdot \vec{\sigma}_N,$
- $\mathcal{O}_5 \equiv \vec{\sigma}_L \cdot (i\hat{q} \times \vec{v}_N),$
- $\mathcal{O}_6 \equiv (i\hat{q} \cdot \vec{\sigma}_L)(i\hat{q} \cdot \vec{\sigma}_N),$
- $\mathcal{O}'_{13} \equiv \vec{\sigma}_L \cdot [i\hat{q} \times (\vec{v}_N \times \vec{\sigma}_N)]$
- $\mathcal{O}_{14} \equiv (i\hat{q} \cdot \vec{\sigma}_L)(\vec{v}_N \cdot \vec{\sigma}_N).$

- $\mathcal{O}_2^{f'} \equiv i\hat{q} \cdot \frac{\vec{v}_\mu}{2} 1_N,$
- $\mathcal{O}_3^f \equiv i\hat{q} \cdot \left(\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L\right) 1_N,$
- $\mathcal{O}_5^f \equiv \left(i\hat{q} \times \frac{\vec{v}_\mu}{2}\right) \cdot \vec{\sigma}_N,$
- $\mathcal{O}_{13}^{f'} \equiv \left[i\hat{q} \times \left(\frac{\vec{v}_\mu}{2} \times \vec{\sigma}_L\right)\right] \cdot \vec{\sigma}_N,$
- $\mathcal{O}_{14}^f \equiv \left(\frac{\vec{v}_\mu}{2} \cdot \vec{\sigma}_L\right)(i\hat{q} \cdot \vec{\sigma}_N).$

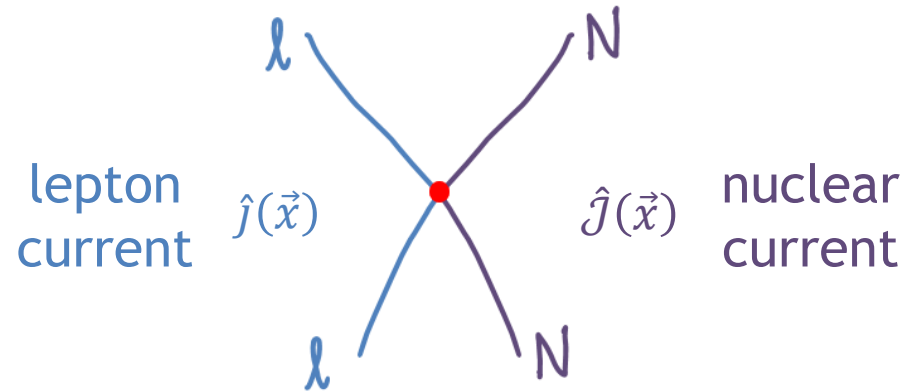
Matching data
 \Rightarrow Must be Tensor

New weak interactions

Nuclear β -decay

Weak interaction

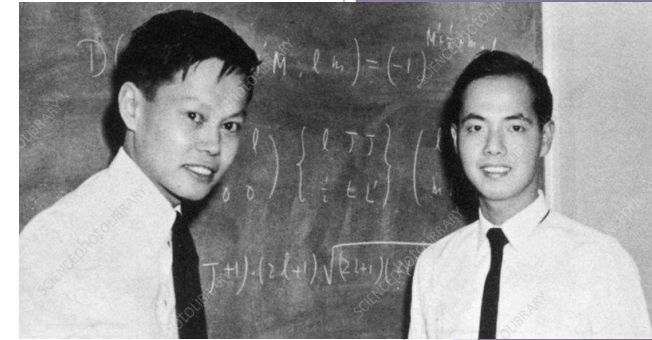
Low energy reaction of leptons with nucleons



$$\hat{\mathcal{H}}_W \sim C \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

A-priori: {

- Scalar (C_S)
- PseudoScalar (C_P)
- Vector (C_V)**
- Axial vector (C_A)**
- Tensor (C_T)



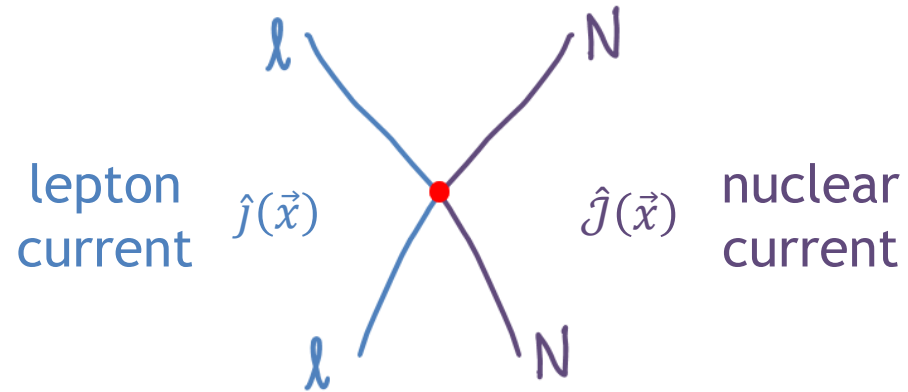
Theory: C.N. Yang and T.D. Lee (Nobel 1957)



Experiment: C.S. Wu:
Parity violation in *nuclear* β -decays
 \Rightarrow Weak SM structure: “**V – A**”

Weak interaction

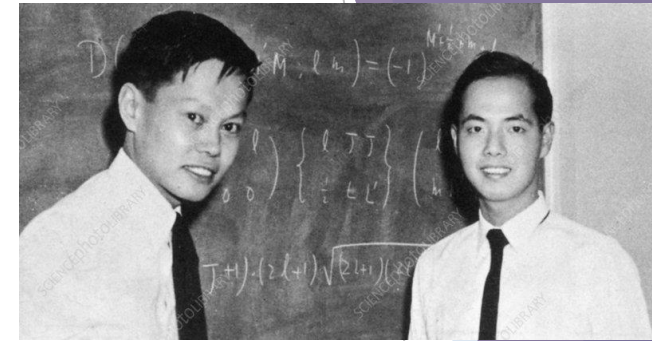
Low energy reaction of leptons with nucleons



$$\hat{\mathcal{H}}_W \sim C \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

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Theory: C.N. Yang and T.D. Lee (Nobel 1957)



Experiment: C.S. Wu:

Parity violation in *nuclear β -decays*

\Rightarrow Weak SM structure: “**V – A**”

The SM is incomplete

\gg Ongoing searches for C_S, C_P, C_T in precision *nuclear β -decay* experiments

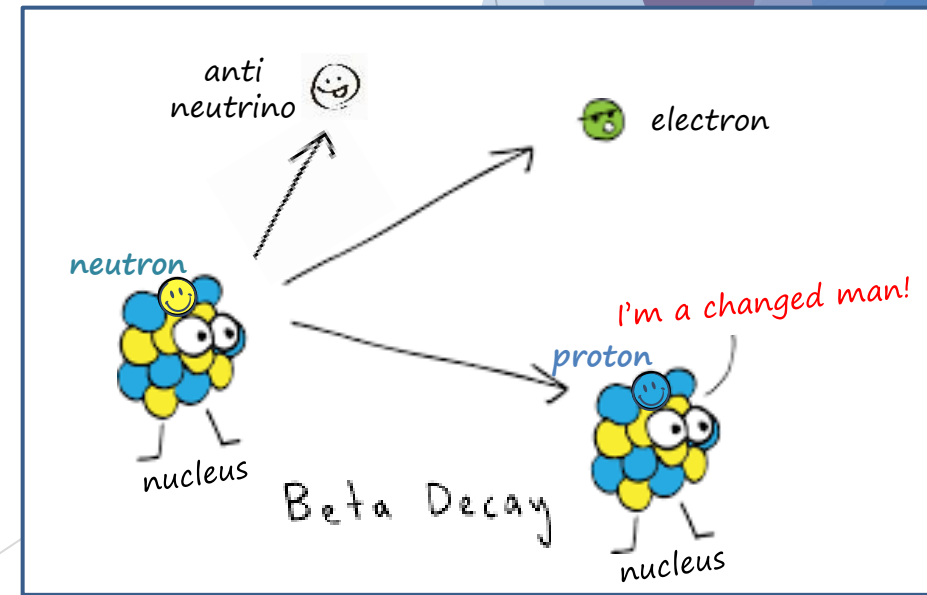
Nuclear β -decay

Low momentum transfer: $q \sim 0 - 10 \text{ MeV}/c$

angular
momentum \swarrow \nearrow parity

Transitions $J^{\Delta\pi}$:

Beta decay, Khan Academy, cdn.kastatic.org/ka-perseus-images/8d978444f15f9bbc3bcadb0549816bc7e264b977.svg



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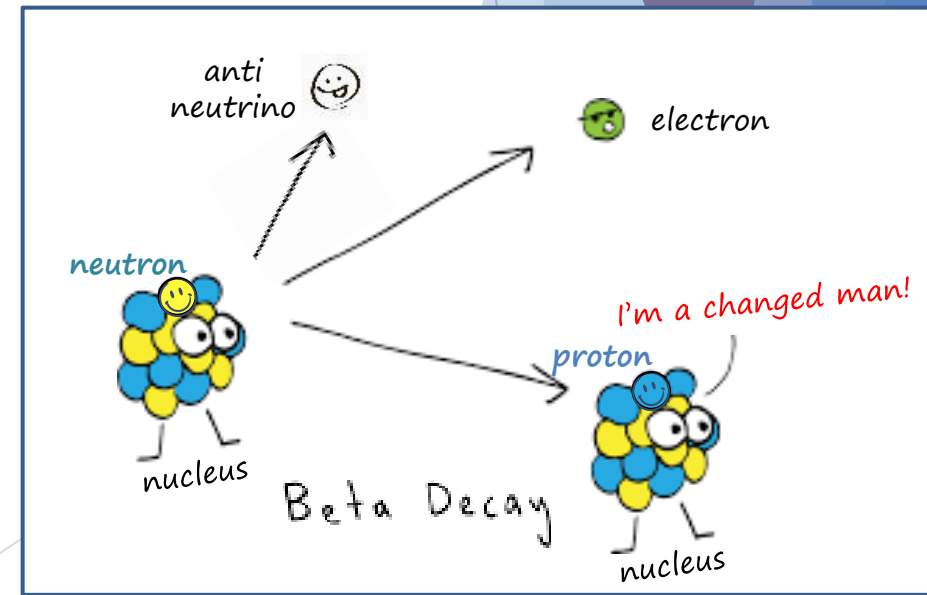
“Allowed”
(when $q \rightarrow 0$)

- 0^+ : Fermi
- 1^+ : Gamow-Teller

“Forbidden”
(vanish for $q \rightarrow 0$)

- All the rest ($J^{\Delta\pi}$)

Beta decay, Khan Academy, cdn.kastatic.org/ka-perseus-images/8d978444f15f9bbc3bcadb0549816bc7e264b977.svg

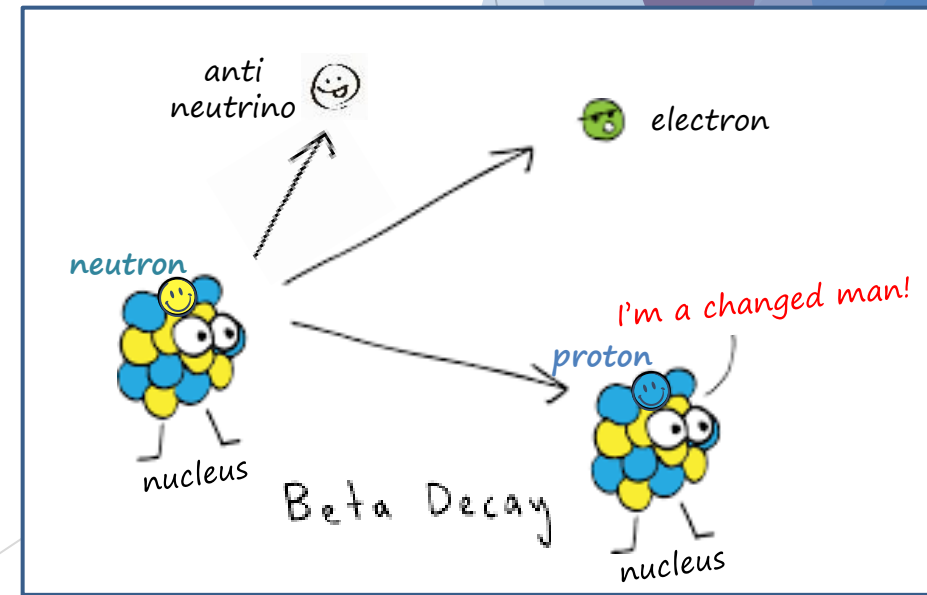


Nuclear β -decay

- ▶ β -decay rate:

$$d\omega \propto |\langle \psi_f || \hat{H}_W || \psi_i \rangle|^2 \underset{\text{allowed}}{\propto} 1 + \underbrace{a_{\beta\nu} \vec{\beta} \cdot \hat{v}}_{\text{Observables}} + \underbrace{b_F}_{\text{Observables}} \frac{m_e}{E}$$

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Nuclear β -decay

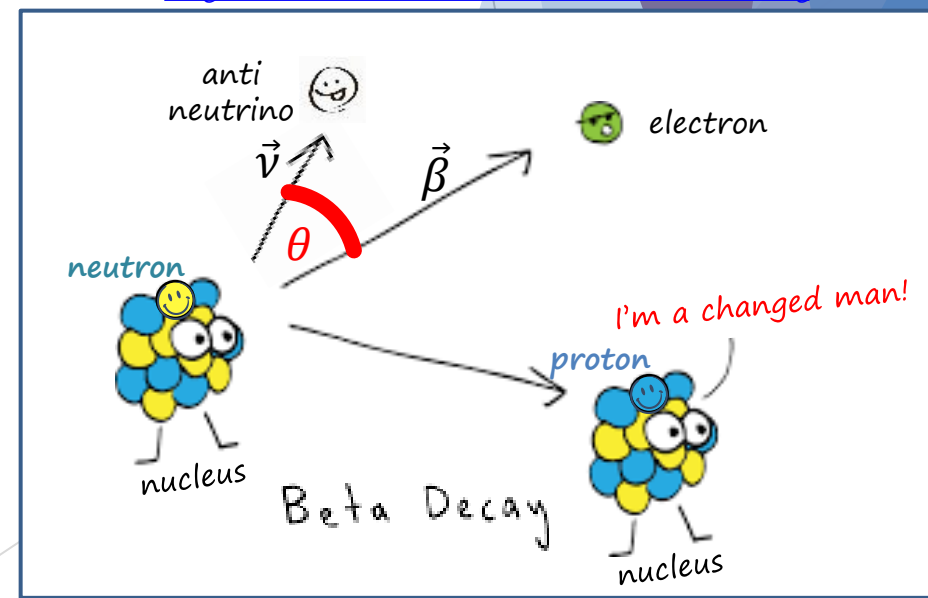
- ▶ β -decay rate:

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Measurements (e.g., Gamow-Teller):

- ▶ **Angular correlation:** $a_{\beta\nu} = -\frac{1}{3} \left(1 - \frac{|C_T^+|^2 + |C_T^-|^2}{4|C_A|^2} \right)$
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ϵ_T
 $C_A = 1.27$ Axial vector coupling constant (SM)

C_T^+ (C_T^-) $\lesssim 10^{-3}$ Tensor left (right) coupling constants (BSM), unknown

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electron's mass, energy

Observables

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SM (1) BSM (terms in fraction)

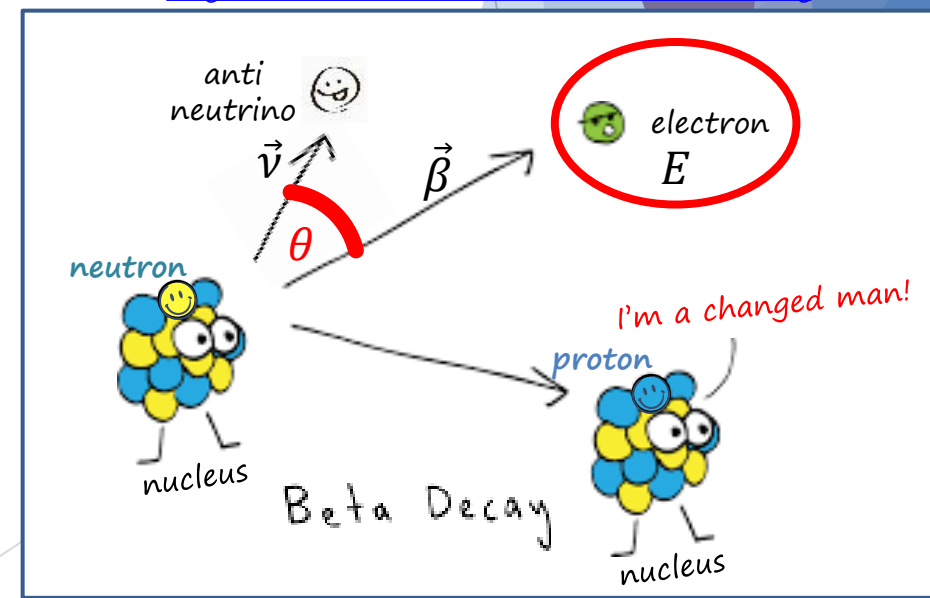
- ▶ Quadratic in C_T^+ , C_T^-

- ▶ **Energy spectrum:** Fierz term $b_F^{\beta\bar{F}} = 0 \pm \frac{C_T^+}{C_A}$

SM (0) BSM (term in fraction)

- ▶ Vanishes for right-handed neutrinos ($C_T^+ = 0$)

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electron's mass, energy

Observables

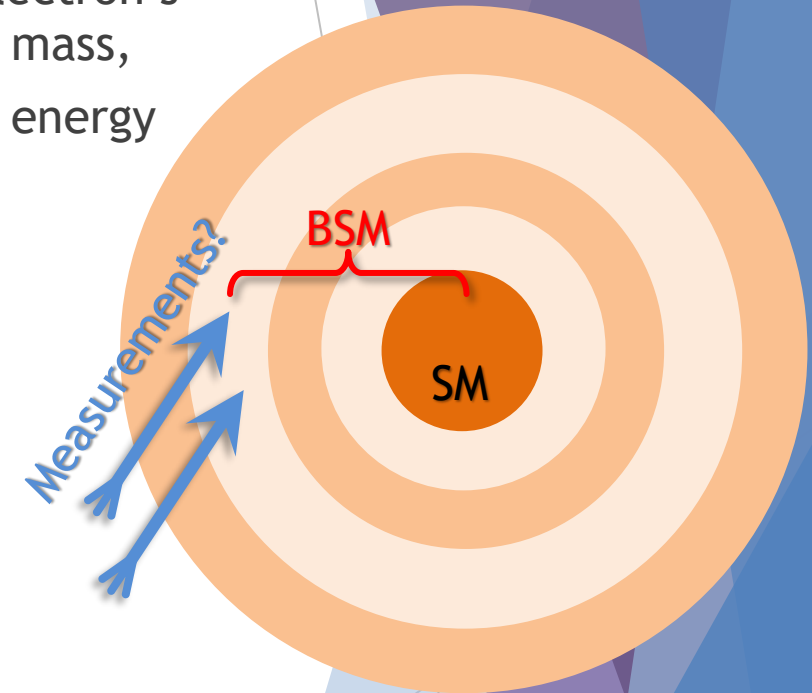
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Searches for deviations from the SM “V-A” structure

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electron's mass, energy

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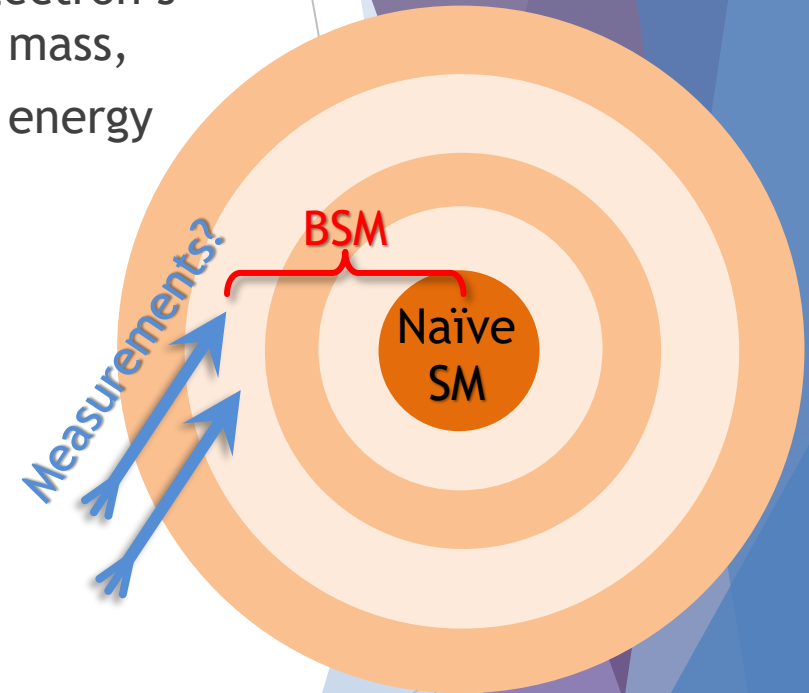
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electron's mass, energy
 →
 →

Observables

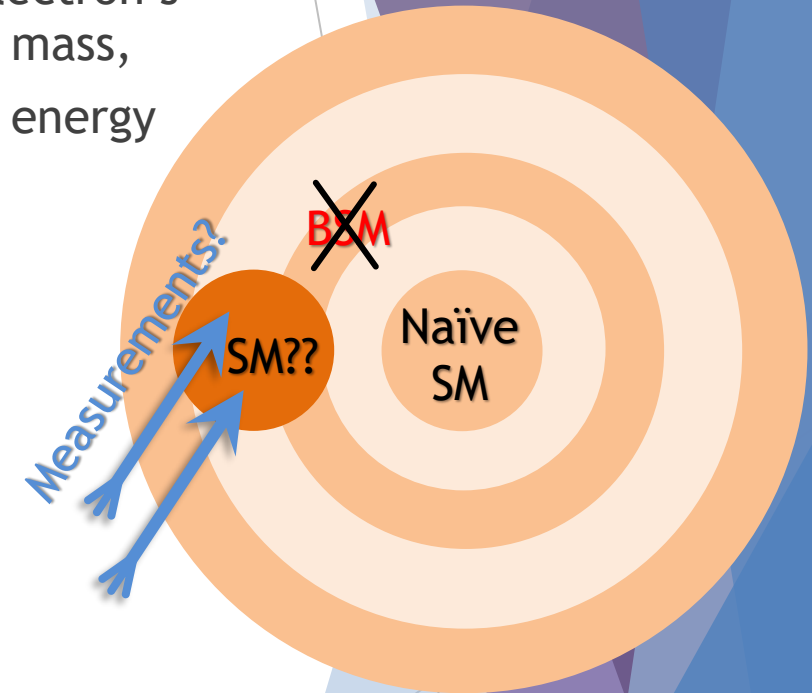
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Observables

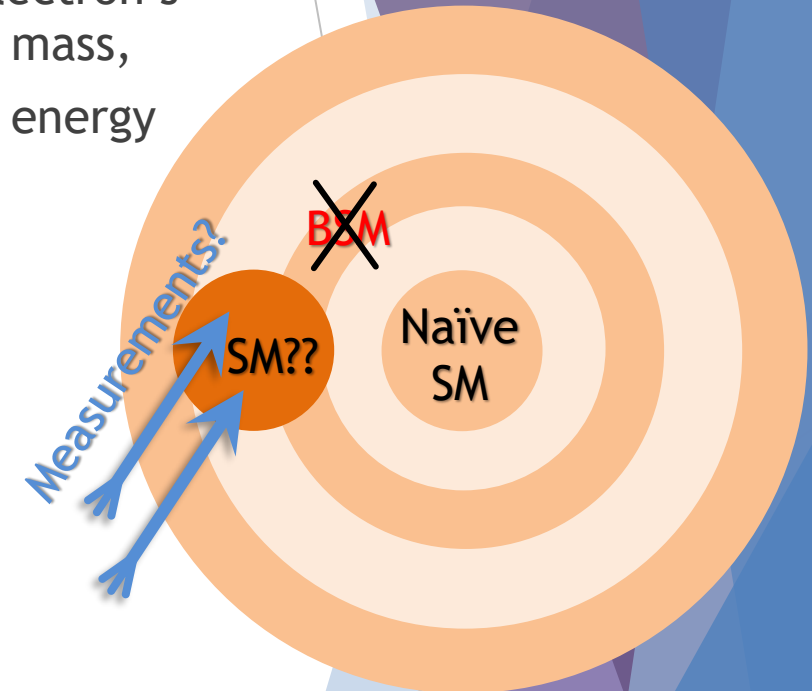
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Searches for deviations from the SM “V-A” structure

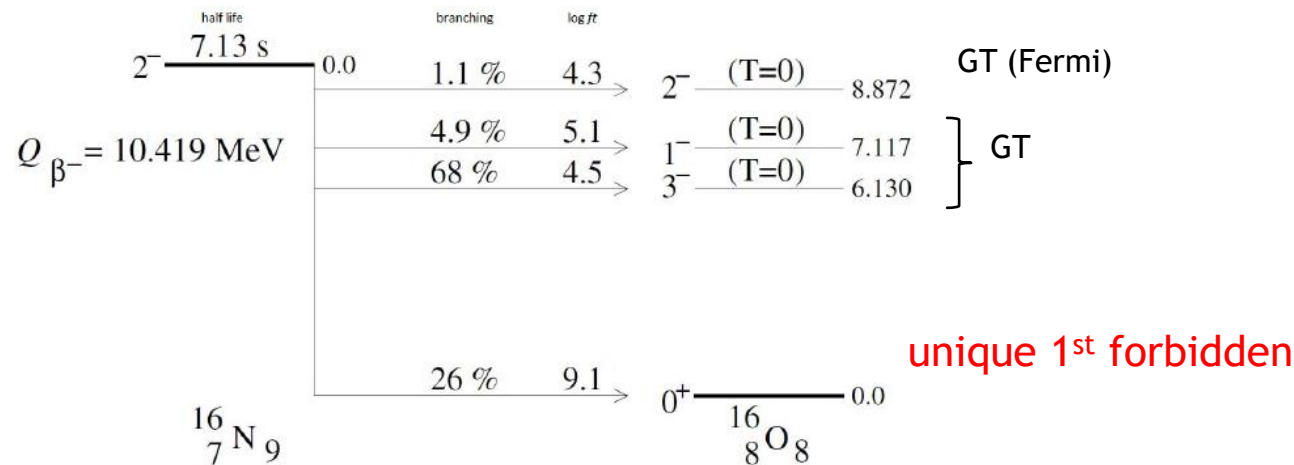
>> More accurate theory is needed

Unique 1st-forbidden decays

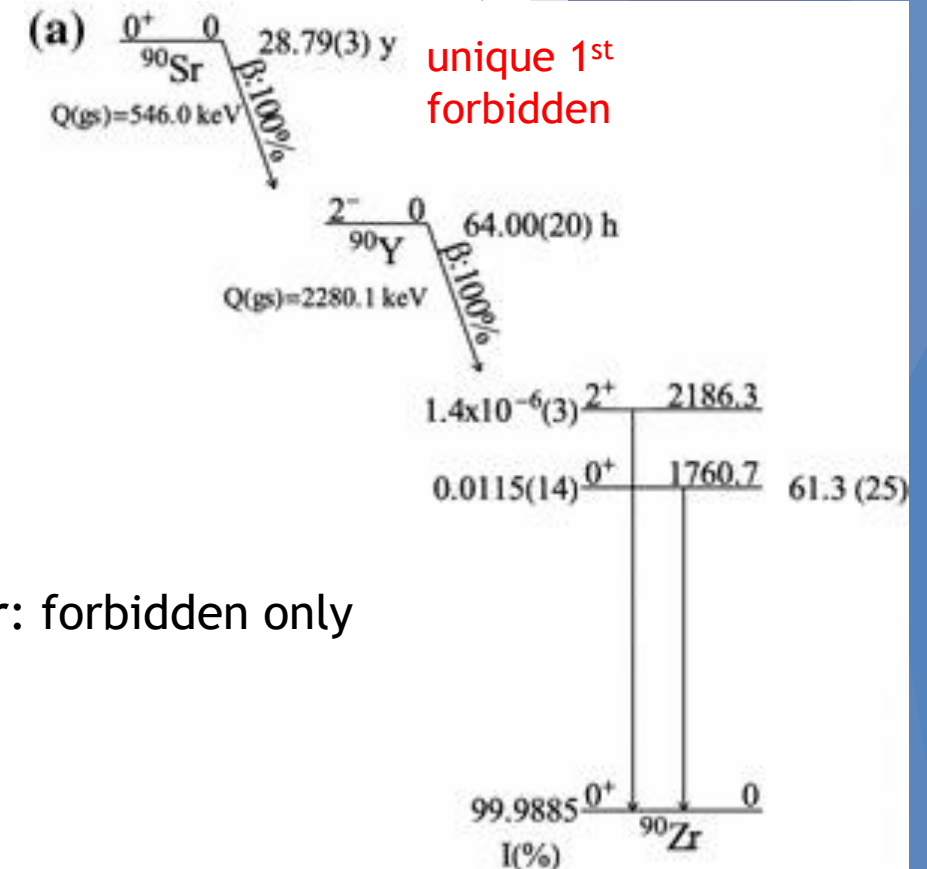
$$d\omega \propto 1 + a_{\beta\nu} \left[1 - (\hat{\beta} \cdot \hat{\nu})^2 \right] + b_F \frac{m_e}{\epsilon}$$

The β -energy spectrum is sensitive to both $a_{\beta\nu}$ & b_F

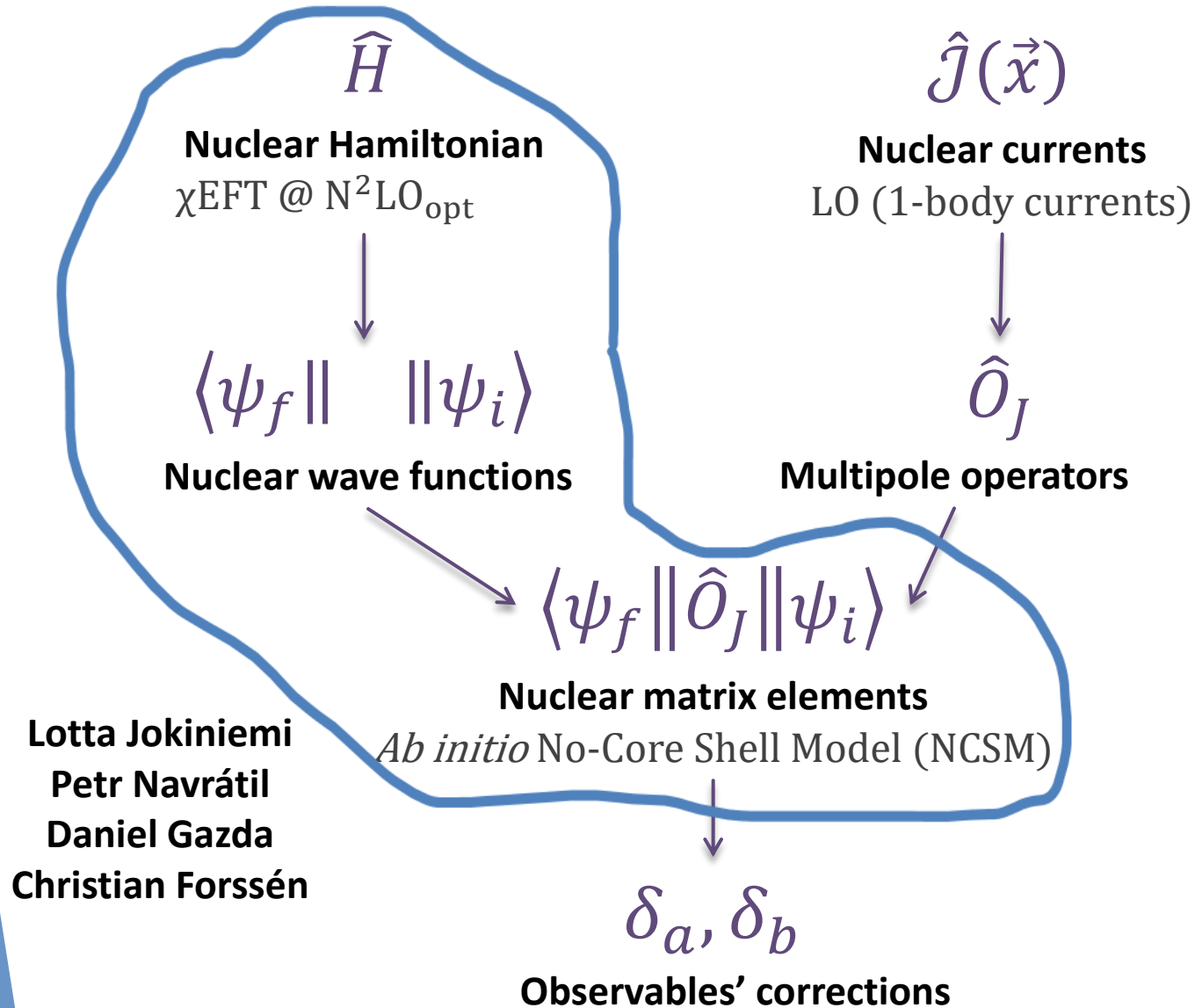
- ▶ Allows simultaneous extraction of C_T^+ and C_T^-
- ▶ Increases the accuracy level



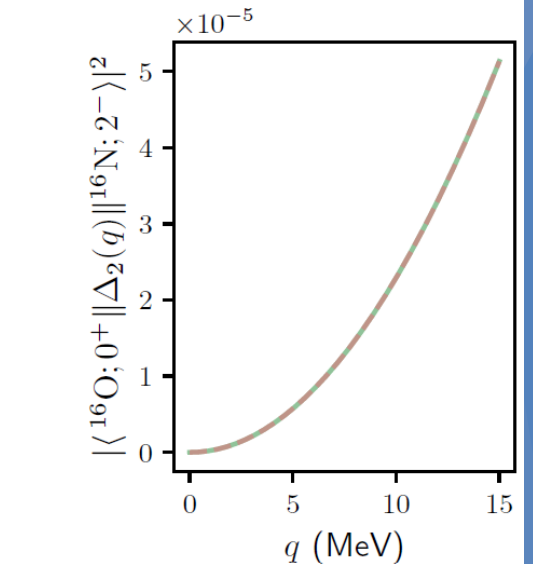
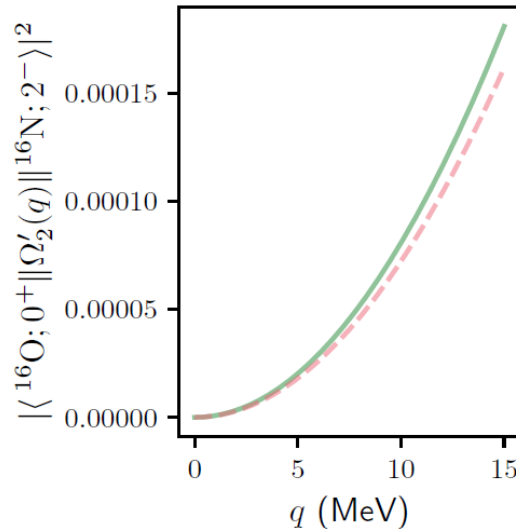
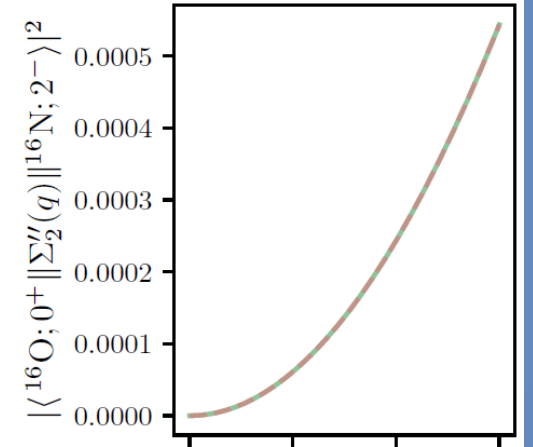
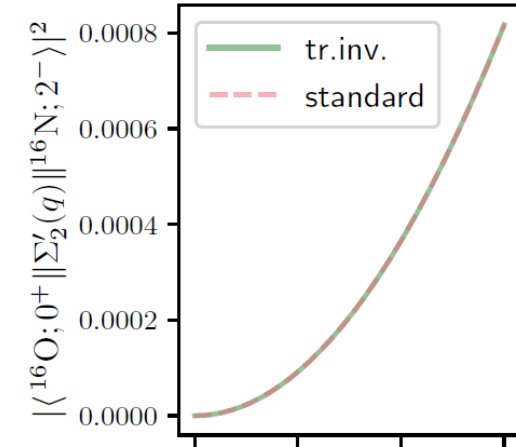
¹⁶N: Large energy separation between the forbidden and allowed branches



Ab initio calculations of $^{16}\text{N} \xrightarrow{\beta^-} ^{16}\text{O}$ forbidden decay



Lotta Jokiniemi
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$^{16}\text{N} \rightarrow ^{16}\text{O}$ forbidden spectrum

- ▶ Experiments are aiming a 10^{-3} accuracy
- ▶ The spectrum can be used to extract b_F & $a_{\beta\nu}$

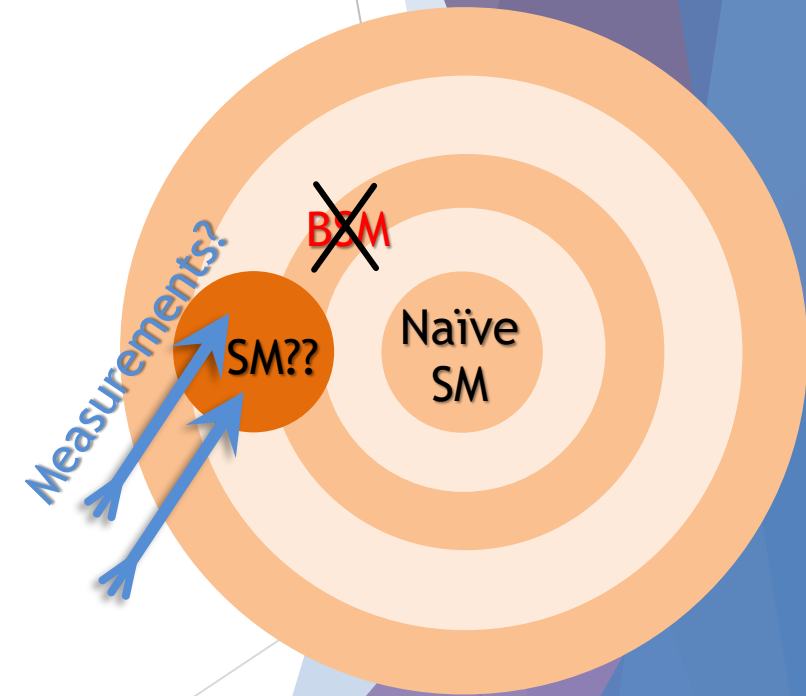
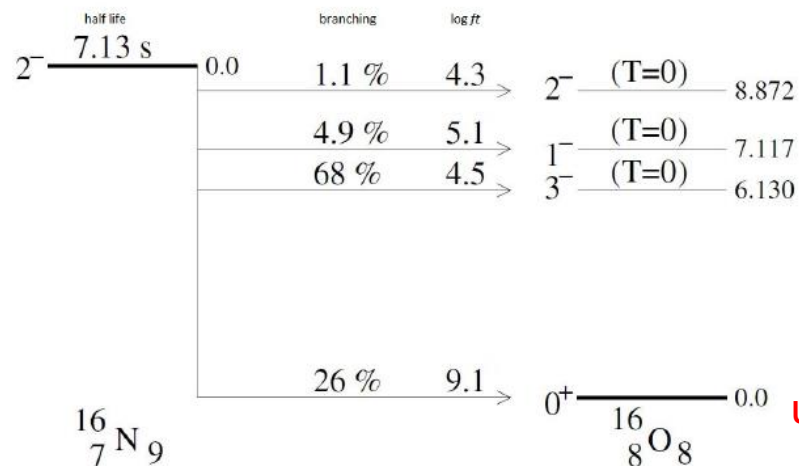
$$b_F = 0 + \delta_b + \frac{C_T^+}{C_A}$$

SM SM correction BSM

▶ Looking for $\frac{C_T^+}{C_A} \sim 10^{-3}$

$$\delta_b = -1.04(13) \cdot 10^{-3}$$

Preliminary



^{16}N : Large energy separation between the forbidden and allowed branches

$^{16}\text{N} \rightarrow ^{16}\text{O}$ forbidden spectrum

- ▶ Experiments are aiming a 10^{-3} accuracy
- ▶ The spectrum can be used to extract b_F & $a_{\beta\nu}$

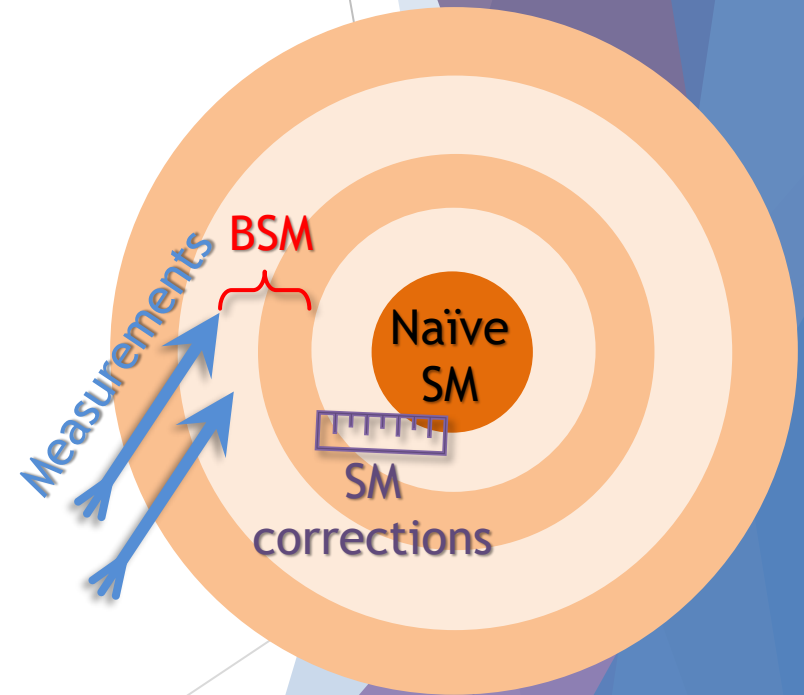
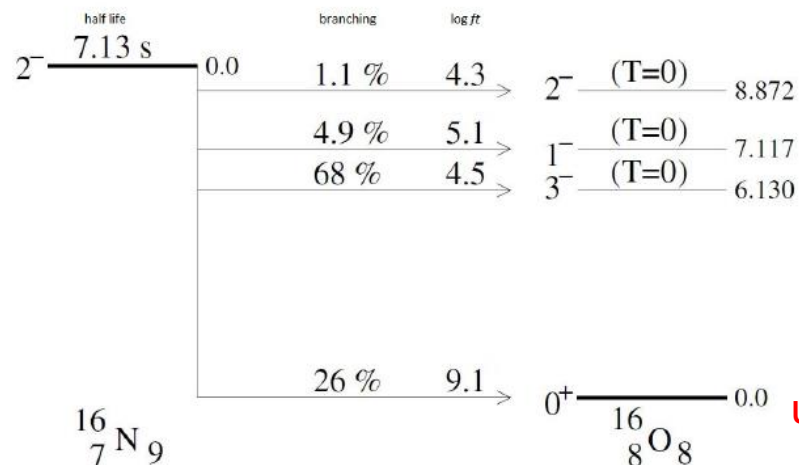
$$b_F = 0 + \overset{\text{SM}}{\delta_b} + \overset{\text{BSM}}{\frac{C_T^+}{C_A}}$$

$$\text{Looking for } \frac{C_T^+}{C_A} \sim 10^{-3}$$

$$\delta_b = -1.07(13) \cdot 10^{-3}$$

$$\text{Uncertainty} \sim 10^{-4}$$

Preliminary



^{16}N : Large energy separation between the forbidden and allowed branches

$^{16}\text{N} \rightarrow ^{16}\text{O}$ forbidden angular correlation

- ▶ Experiments are aiming a 10^{-3} accuracy

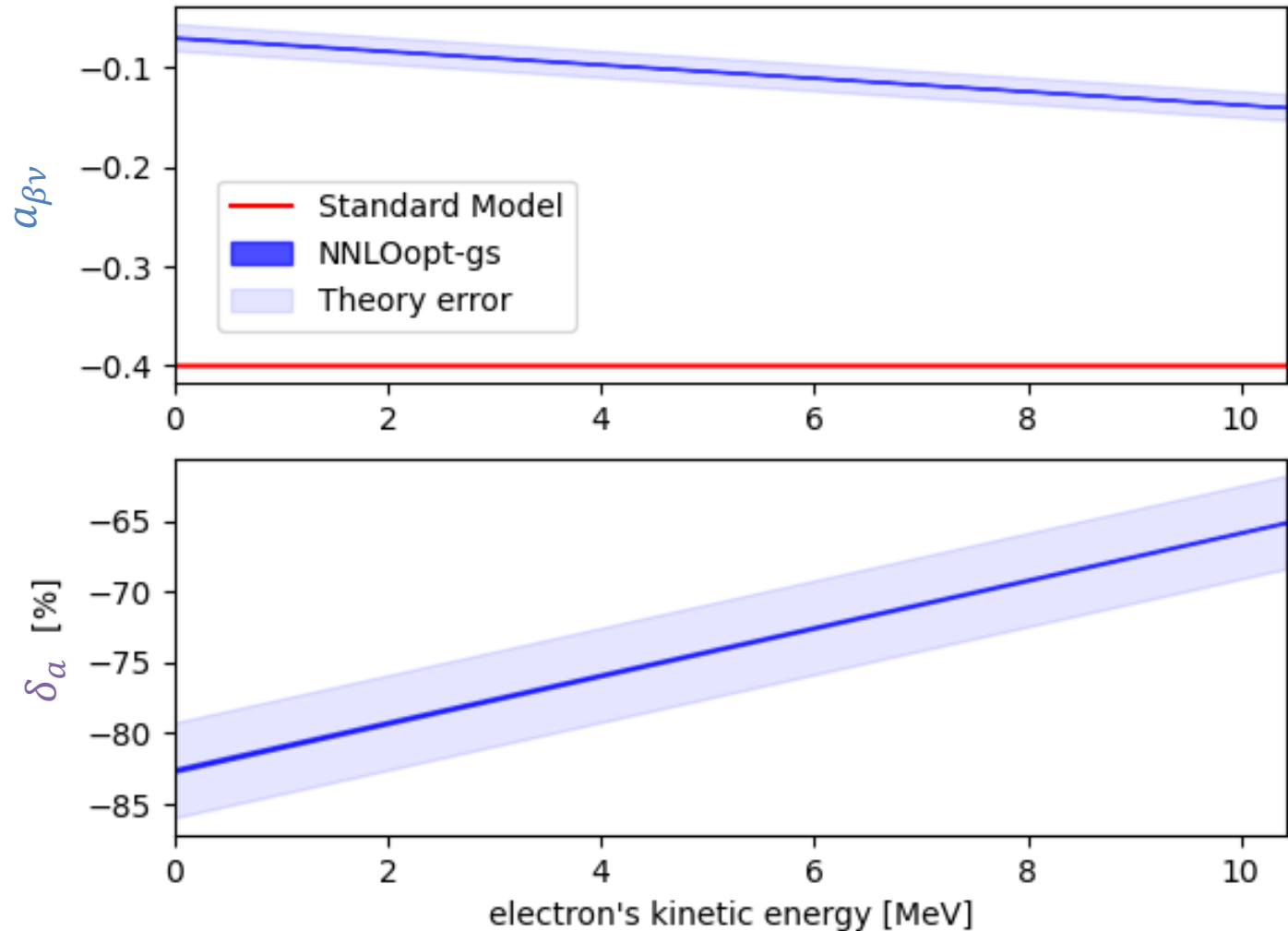
$$C_T^+(C_T^-) \sim 10^{-3}$$

- ▶
$$a_{\beta\nu} = -\frac{2}{5} \left(1 + \overset{\text{SM}}{\delta_a} + \overset{\text{BSM}}{\frac{|C_T^+|^2 + |C_T^-|^2}{4|C_A|^2}} \right)$$

- ▶ $\langle \delta_a \rangle = -0.609(15)$

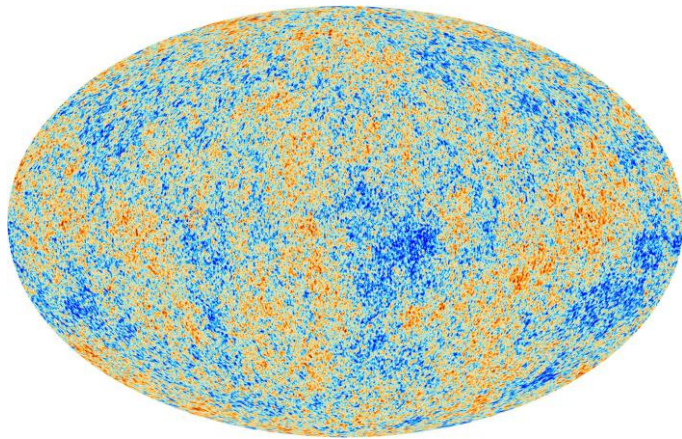
Very Preliminary

16N forbidden beta decay angular correlation



Summary: BSM Searches **with nuclei...**

Astronomy



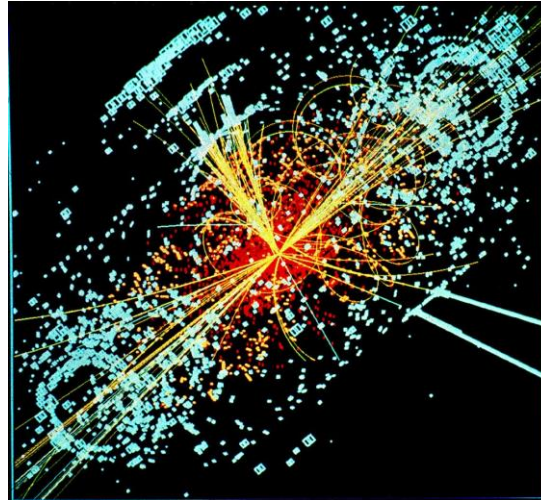
https://www.esa.int/ESA_Multimedia/Images/2013/03/Planck_CMB
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➤ Dark Matter direct detection



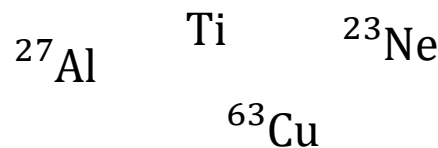
and many more...

Particles Physics

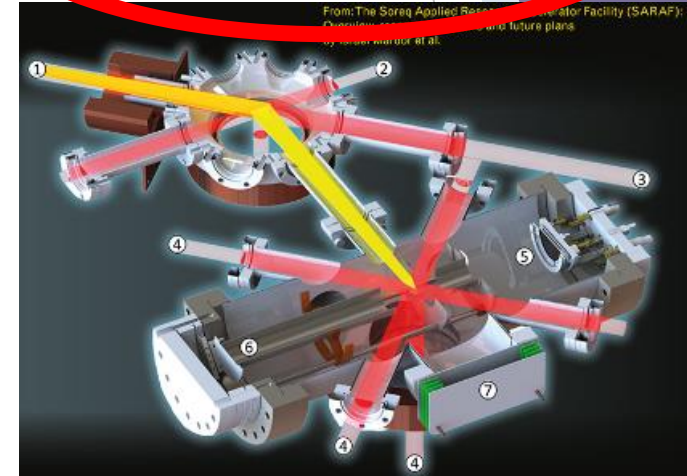


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➤ Lepton Flavor Violation with $\mu \rightarrow e$ conversion

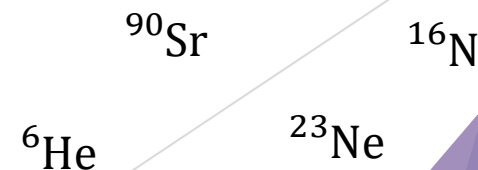


Precision Frontier Nuclear Physics



Mardor et al., *Eur. Phys. J. A* 54, 91 (2018)

➤ New Weak Interactions with β -decays



Thanks!

INT

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