Ab initio description of monopole resonances in light- and medium-mass nuclei

PAINT2024 – Workshop on Progress in Ab Initio Nuclear Theory TRIUMF, Vancouver

February 27th, 2024

Andrea Porro Technische Universität Darmstadt





- Physical introduction
- Existing ab initio theoretical tools

Giant Resonances

- Physical introduction
- Existing ab initio theoretical tools

2 Ab initio PGCM

- Formalisms
- Uncertainty quantification

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Chosen results

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Selected applications

- Shape coexistence
- Deformation

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- PAV and VAP strategies
- Rotation-vibration coupling

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From finite nuclei to Astrophysics

• Preliminary incompressibility results

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[Pictures from Dytrych et al., PRL, 2020]





[[]Pictures from Dytrych et al., PRL, 2020]





EOM and VS extensions

- IMSRG and CC
- Suited for **weakly-collective** excitations only

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CC-LIT Lorenz integral transform (sperical)
 SA-NCSM Application to deformed systems (²⁰Ne)

[Bacca, Barnea, Hagen, Orlandini, Papenbrock, PRL, 2013] [Dytrych, Launey, Draayer, Maris, Vary et al., PRL, 2013]



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(Q)RPA

• Spherical (Q)RPA, 2nd RPA, CC-RPA, IMSRG-RPA, IMSRG-2nd RPA

[R. Trippel, PhD Thesis, 2016]





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- Spherical (Q)RPA, 2nd RPA, CC-RPA, IMSRG-RPA, IMSRG-2nd RPA
- SCGF, RPA with dressed propagators

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 $\sigma_{v}(\omega)/4\pi^{2}\alpha\omega \, [mb/MeV]$

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(Q)RPA

- Spherical (Q)RPA, 2nd RPA, CC-RPA, IMSRG-RPA, IMSRG-2nd RPA
- SCGF, RPA with dressed propagators
- (Q)RPA for axially- and triaxally-deformed systems

[R. Trippel, PhD Thesis, 2016]

[Barbieri, Raimondi, PRC, 2019]

[Beaujeault-Taudière, Frosini, Ebran, Duguet, Roth, Somà, PRC, 2023]



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Schrödinger equation $H |\Psi_n\rangle = E_n |\Psi_n\rangle$

Schrödinger equation

$$H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$$

Open-shell systems

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Open-shell systems

Symmetry-breaking reference states



Schrödinger equation

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Open-shell systems

Symmetry-breaking reference states



1 Constrained HFB solutions $|\Phi(q)
angle$



Schrödinger equation

 $H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$

Open-shell systems



1 Constrained HFB solutions

 $|\Phi(q)\rangle$ Generator coordinates (q can be any coordinate)

Symmetry-breaking reference states



Schrödinger equation

 $H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$

Open-shell systems

Strong static correlations

 $|\Phi(q)\rangle$

1 Constrained HFB solutions

2 PGCM Ansatz

$$|\Psi_n\rangle = \int \mathrm{d}q \, f_n(q) \, |\Phi(q)\rangle$$

Generator coordinates (q can be any coordinate)

Symmetry-breaking reference states

E_{HFB} (q)^h q

Schrödinger equation

 $H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$

Symmetry-breaking reference states

Generator coordinates

(q can be any coordinate)

Open-shell systems

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Linear coefficients

Europhie Participants (q) A q

Schrödinger equation

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Open-shell systems

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 $|\Phi(q)\rangle$

Strong static correlations

1 Constrained HFB solutions

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 $|\Psi_n\rangle = \int \mathrm{d}q f_n(q) \Phi(q)\rangle$

Linear coefficients

Generator coordinates (q can be any coordinate)

Symmetry-breaking reference states

3 HWG Equation

Variational method

$$\delta \frac{\langle \Psi_n | H | \Psi_n \rangle}{\langle \Psi_n | \Psi_n \rangle} = 0$$



Schrödinger equation

 $H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$

Symmetry-breaking reference states

Generator coordinates

(q can be any coordinate)

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Linear coefficients

3 HWG Equation

Variational method $\delta rac{\langle \Psi_n | H | \Psi_n
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angle} = 0$

Schrödinger-like equation $\int \left[\mathcal{H}(p,q) - E_n \, \mathcal{N}(p,q) \right] f_n(q) \, \mathrm{d}q = 0$ Kernels evaluation

 $\mathcal{H}(p,q) \equiv \langle \Phi(p) | H | \Phi(q) \rangle$ $\mathcal{N}(p,q) \equiv \langle \Phi(p) | \Phi(q) \rangle$

- · A

Schrödinger equation





 $|\Phi(q)\rangle$

1 Constrained HFB solutions





Linear coefficients

Generator coordinates (q can be any coordinate)

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EHFB (q)

Kernels evaluation

Emin

 $\mathcal{H}(p,q) \equiv \langle \Phi(p) | H | \Phi(q) \rangle$ $\mathcal{N}(p,q) \equiv \langle \Phi(p) | \Phi(q) \rangle$

Schrödinger equation





 $|\Phi(q)\rangle$

1 Constrained HFB solutions

2 PGCM Ansatz



Linear coefficients

Generator coordinates (q can be any coordinate)

3 HWG Equation

Variational method $\delta \frac{\langle \Psi_n | H | \Psi_n \rangle}{\langle \Psi_n | \Psi_n \rangle} = 0$

+ Projection

Schrödinger-like equation $\int \left[\mathcal{H}(p,q) - E_n \,\mathcal{N}(p,q) \right] f_n(q) \,\mathrm{d}q = 0$ Kernels evaluation

 $\mathcal{H}(p,q) \equiv \langle \Phi(p) | H | \Phi(q) \rangle$ $\mathcal{N}(p,q) \equiv \langle \Phi(p) | \Phi(q) \rangle$

$$E_{HFB}(q)^{h}$$

 q
 E_{min}

Setting

Studied quantity: monopole strength

- Transition amplitudes: height of peaks
- Energy difference: position of peaks

$$S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$$


Studied quantity: monopole strength

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Studied quantity: monopole strength

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 ω [MeV]

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Related moments

٠

$$\begin{split} m_k &\equiv \int_0^\infty S_{00}(\omega) \,\omega^k \,d\omega \\ &= \sum_\nu (E_\nu - E_0)^k |\langle \Psi_\nu | r^2 | \Psi_0 \rangle|^2 \end{split}$$



Studied quantity: monopole strength

- Transition amplitudes: height of peaks
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Related moments $m_k \equiv$

$$S_{00}(\omega) \omega^{k} d\omega$$
$$= \sum_{\nu} (E_{\nu} - E_{0})^{k} |\langle \Psi_{\nu} | r^{2} | \Psi_{0} \rangle|^{2}$$

 $S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$

Quantify the most relevant features of the strength $\bar{E}_1 = \frac{m_1}{m_0}$ $\sigma^2 = \frac{m_2}{m_0} - \left(\frac{m_1}{m_0}\right)^2 \ge 0$

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Studied quantity: monopole strength

- Transition amplitudes: height of peaks ٠
- Energy difference: position of peaks

Related moments $m_k \equiv \int_{-\infty}^{\infty} S_{00}(\omega) \omega^k d\omega$

$$= \sum_{\nu} (E_{\nu} - E_0)^k |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2$$

Quantify the most relevant features of the strength $\bar{E}_1 = \frac{m_1}{m_0}$ $\sigma^2 = \frac{m_2}{m_0} - \left(\frac{m_1}{m_0}\right)^2 \ge 0$



 $S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$

200

150

100

30

40

50

Ab-initio PGCM and QRPA consistent numerical settings (systematic study in ⁴⁶Ti)

Quantities expanded on harmonic oscillator basis (characterised by $\hbar\omega$, e_{max} , e_{3max})

Studied quantity: monopole strength

- Transition amplitudes: height of peaks
- Energy difference: position of peaks

Related moments $m_k \equiv \int_0^\infty S_{00}(\omega) \, \omega^k \, d\omega$

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Ab-initio PGCM and QRPA **consistent numerical settings** (systematic study in ⁴⁶Ti)

- Quantities expanded on harmonic oscillator basis (characterised by $\hbar\omega$, e_{max} , e_{3max})
- Family of chiral NN + in-medium 3N interactions (NLO, N2LO and N3LO)
 - T. Hüther, K. Vobig, K. Hebeler, R. Machleidt and R. Roth, "Family of chiral two-plus three-nucleon interactions for accurate nuclear structure studies", *Phys. Lett. B*, 808, 2020
 - In-vacuum SRG evolution (α =0.04 fm⁴, α =0.08 fm⁴)
 - M. Frosini, T. Duguet, B. Bally, Y. Beaujeault-Taudière, J.-P. Ebran and V. Somà, "In-medium k-body reduction of n-body operators", *The European Physical Journal A*, *57*(4), 2021

 $S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$

Uncertainty budget

SRG dependence Many-body truncation Strong centroid dependence ~ 10 % • Comparison to PGCM-PT Dispersion relative error ~ 20 % • Only tested for **low-lying exc** Truncates both H and many-body Correlated to SRG and generator coords Generator coordinates choice **Empirical knowledge**, two coords **r** and β_2 More systematic choice needed Three-body treatment NO2B approximation 1-2 % uncertainty in low-lying exc Not tested for giant resonances • Hamiltonian parameters C 3 me C WAX • **LEC** dependence of χ forces Few interactions compared • Correlated to SRG Need for emulators (EC)

Chiral Order

- Good overall convergence ٠
- Centroid relative error ~ 1,6 % ٠
- Dispersion relative error ~ 9,8 % .

Harmonic Oscillator width

- Good overall convergence
- Centroid relative error ~ 1,6 %
- Dispersion relative error ~ 6 %

- Good overall convergence •
- Centroid relative error ~ 0.6 %
- Dispersion relative error ~ 1,7 %
- e_{3max} not studied (14 safe for GS)



Good overall convergence



- Good overall convergence
- Centroid relative error ~ 0,6 %





- Good overall convergence
- Centroid relative error ~ 0,6 %
- Dispersion relative error ~ 1,7 %





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• Preliminary incompressibility results







Shape coexistence [Jenkins et al., 2012]

• Oblate GS and prolate-shape isomer



Shape coexistence [Jenkins et al., 2012]

- Oblate GS and prolate-shape isomer •
- Proper study of shape coexistence in PGCM •



- Oblate GS and prolate-shape isomer
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- Oblate GS and prolate-shape isomer
- Proper study of shape coexistence in PGCM
 - Shape coexistence but weak mixing



- Oblate GS and prolate-shape isomer
- Proper study of shape coexistence in PGCM
 - Shape coexistence but weak mixing
 - Nuclei with stronger signature ? 12









40 3.6 • • 30 [MeV] 3.4 [^E] 3.2 یں 20 س^{یت} $E_{_{
m HFB}}$ 3 10 2.8 0 0.5 -0.5 0 1 -1β2

Total Energy Surface $E_{HFB}(\beta_2, r)$

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- Focus on the prolate-shape isomer
- Coupling to GQR generates splitting
- X High peak = shifted "spherical" breathing mode
- **x** Low peak = induced by coupling to GQR (K=0)
- Two-peak GMR on the prolate shape isomer



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Comparison to experimental data



Comparison to experimental data



Comparison to experimental data



Ab initio PGCM nicely reproduces the experimental data

• Better description of the main resonance and fragmentation

Experimental data are useful and promising to test different many-body methods

Data are not unambiguous, i.e. higher resolution would be beneficial

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Projection in GCM and QRPA

(Q) R P A

DIFFERENT FLAVOURS OF SYMMETRY BREAKING AND RESTORATION

Symmetry breaking

GCM

Symmetry conserving

Projection in GCM and QRPA






(2) [Federschmidt and Ring, NucPhysA, 1985]

DIFFERENT FLAVOURS OF SYMMETRY BREAKING AND RESTORATION



(2) [Federschmidt and Ring, NucPhysA, 1985]

DIFFERENT FLAVOURS OF SYMMETRY BREAKING AND RESTORATION



(1) [Erler, PhD Thesis, TUD, 2012]

(2) [Federschmidt and Ring, NucPhysA, 1985]

Symmetry conserving



GCM : symmetry-breaking solutions $|\text{GS}\rangle_{\text{def}}$ $|\omega\rangle_{\text{def}}$ PGCM : symmetry-conserving solutions $|\text{GS}\rangle_{\text{sym}}$ $|\omega\rangle_{\text{sym}}$



GCM : symmetry-breaking solutions

PGCM : symmetry-conserving solutions (

$$\begin{split} |\mathrm{GS}\rangle_{\mathrm{def}} & |\omega\rangle_{\mathrm{def}} \\ |\mathrm{GS}\rangle_{\mathrm{sym}} & |\omega\rangle_{\mathrm{sym}} \end{split}$$

Projection effects

- Not too dissimilar
- Increased **fragmentation** (e.g. ²⁴Mg)
- More quantitative agreement



GCM : symmetry-breaking solutions

PGCM : symmetry-conserving solutions $|GS\rangle_{sym}$

 $|\mathrm{GS}\rangle_{\mathrm{def}}$

 $|\omega\rangle_{\rm def}$

 $|\omega\rangle_{\rm sym}$

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Can we treat projection a posteriori?



GCM : symmetry-breaking solutions $|\text{GS}\rangle_{\text{def}}$ $|\omega\rangle_{\text{def}}$ PGCM : symmetry-conserving solutions $|\text{GS}\rangle_{\text{sym}}$ $|\omega\rangle_{\text{sym}}$

PAV GCM: projection of symmetry-breaking solution

- Anomalous spectrum
- Zero-frequency rotations (Goldstone modes)
- Born-Oppenheimer-like approximation



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Rotational state $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\omega\rangle_{\rm sym}$ $|\mathrm{GS}\rangle_{\mathrm{sym}}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum ٠ Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$ Rotational state $\langle \mathrm{ROT} | \omega \rangle_{\mathrm{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\mathrm{GS}\rangle_\mathrm{sym}$ $|\omega\rangle_{\rm sym}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum • Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$ Rotational state $\langle \mathrm{ROT} | \omega \rangle_{\mathrm{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$

$$|\omega\rangle_{\rm def} = N_{\rm rot} |{
m ROT}\rangle + N_{\rm vib} |{
m VIB}\rangle$$



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $\ket{\mathrm{GS}}_{\mathrm{sym}}$ $|\omega\rangle_{\rm sym}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$ Rotational state $\langle \mathrm{ROT} | \omega \rangle_{\mathrm{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$ $|\omega\rangle_{\rm def} = N_{\rm rot} |{\rm ROT}\rangle + N_{\rm vib} |{\rm VIB}\rangle$

Can be subtracted!



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\omega\rangle_{
m sym}$ $|\mathrm{GS}\rangle_{\mathrm{sym}}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum • Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$ Rotational state $\langle \text{ROT} | \omega \rangle_{\text{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$ $|\omega\rangle_{\rm def} = N_{\rm rot} |{\rm ROT}\rangle + N_{\rm vib} |{\rm VIB}\rangle$

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 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\omega\rangle_{\rm sym}$ $|\mathrm{GS}\rangle_{\mathrm{sym}}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum • Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\mathrm{ROT}\rangle = \hat{R}(\Omega) |\mathrm{GS}\rangle_{\mathrm{def}}$ Rotational state $\langle \mathrm{ROT} | \omega \rangle_{\mathrm{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$ $|\omega\rangle_{\rm def} = N_{\rm rot} |{\rm ROT}\rangle + N_{\rm vib} |{\rm VIB}\rangle$ Can be subtracted!

Observed both in GCM and RPA⁽¹⁾

- Does not depend on the many-body method
- Consequence of deformed ground state



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\omega\rangle_{\rm sym}$ $|\mathrm{GS}\rangle_{\mathrm{svm}}$ PGCM : symmetry-conserving solutions PAV GCM: projection of symmetry-breaking solution Anomalous spectrum • Zero-frequency rotations (Goldstone modes) • Born-Oppenheimer-like approximation ٠ $|\mathrm{ROT}\rangle = \hat{R}(\Omega) |\mathrm{GS}\rangle_{\mathrm{def}}$ Rotational state $\langle \mathrm{ROT} | \omega \rangle_{\mathrm{sym}} = 0$ Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}}$ $|\omega\rangle_{\rm def} = N_{\rm rot} |{\rm ROT}\rangle + N_{\rm vib} |{\rm VIB}\rangle$

Rotations must be treated variationally

- PGCM already does
- Projected QRPA needed

Observed both in GCM and RPA⁽¹⁾

- Does not depend on the many-body method
- Consequence of deformed ground state

(1) INFN collaboration, G. Colò and D. Gambacurta

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From finite nuclei to Astrophysics

Symmetry energy

- IV GDR
- Dipole polarizability
- Neutron skin

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Nuclear compressibility

• GMR

$$K_{\rm A} = (M/\hbar^2) \langle r^2 \rangle E_{\rm GMR}^2$$
$$\tilde{E}_k = \sqrt{\frac{m_{\rm k}}{m_{k-2}}} \qquad \bar{E}_1 = \frac{m_1}{m_0}$$

From finite nuclei to Astrophysics

Symmetry energy

- IV GDR
- Dipole polarizability
- Neutron skin

Nuclear compressibility

• GMR



Preliminary evaluation of K_{∞}

- Starting from deformed systems
- Extrapolation in **agreement** with commonly accepted values
- Systematic investigation in heavier systems (Sn, Mo isotopic chains, neutron rich)

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Current frontiers







Perspectives



Perspectives



Perspectives



Thanks for the attention



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Alexander Tichai Robert Roth Achim Schwenk



Gianluca Colò Danilo Gambacurta cea

Thomas Duguet Vittorio Somà Mikael Frosini Benjamin Bally Jean-Paul Ebran Alberto Scalesi

Backup slides



• GRs can be interpreted as the first phonon of a collective excitation



- GRs can be interpreted as the first phonon of a collective excitation
- Higher phonons also exist ! Multi-phonon states



- GRs can be interpreted as the first phonon of a collective excitation
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- Not accessible to QRPA



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- Higher phonons also exist ! Multi-phonon states
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• GRs can be interpreted as the first phonon of a collective excitation

 $\log_{10} S_0 + 2$

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One-dimensional PGCM calculation



• PGCM predicts high-lying states





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- PGCM predicts high-lying states
- Close to the harmonic oscillator eigen-solutions

[fm⁴MeV⁻¹]

400 0

200

 $B(E0)[0_{\chi}^{+}$





- GRs can be interpreted as the first phonon of a collective excitation
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- PGCM predicts high-lying states
- Close to the harmonic oscillator eigen-solutions
- Transitions maximised between neighbouring phonons

 $B(E0)[0_{\chi}^{+}$



GRs can be interpreted as the first phonon of a collective excitation •

2

 $\log_{10} S_0 + 2$

- Higher phonons also exist! Multi-phonon states •
- Not accessible to **QRPA** •





- PGCM predicts high-lying states
- Close to the harmonic oscillator eigen-solutions
- Transitions maximised between neighbouring phonons
 - Linear trend in the transition strength Х
• 2-D PGCM in the (r, β_2) plane





- 2-D PGCM in the (r, β_2) plane
- Good agreement with experiment



- 2-D PGCM in the (r, β_2) plane
- Good agreement with experiment
- Multi-phonon states observed





- 2-D PGCM in the (r, β_2) plane
- Good agreement with experiment
- Multi-phonon states observed
- Harmonicity well confirmed







Harmonic Oscillator width











PGCM : multi-reference unperturbed state



PGCM : multi-reference unperturbed state





(2) [Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao and Somà, EPJA 58(64), 2022]

^{(1) [}Frosini, Duguet, Ebran and Somà, EPJA 58(62), 2022]



(2) [Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao and Somà, EPJA 58(64), 2022]



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Ab initio

Exp



(1) [Frosini, Duguet, Ebran and Somà, EPJA 58(62), 2022]

(2) [Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao and Somà, EPJA 58(64), 2022]

SRG dependence



SRG dependence



SRG dependence

A-body Hilbert space \mathcal{H}_{A} $|\Psi_k^{\rm A}\rangle = \Omega \Theta_k^{(0)}\rangle$ Wave operator action $\mathcal{H}_{\mathsf{A}}^{\mathsf{PGCM}}$ SRG ? PGCM subspace













[Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao and Somà, EPJA 58(64), 2022]







One coordinate insufficient (deformed systems)

Two coordinates necessary: empirical knowledge r and β_2





- One coordinate insufficient (deformed systems)
- Two coordinates necessary: empirical knowledge **r** and β_2
- Additional coordinates ?





PGCM alone suited for ab initio?



One coordinate insufficient (deformed systems)

- Two coordinates necessary: empirical knowledge r and β_2
- Additional coordinates ?



[S. Bofos, ongoing] Systematic VS-PGCM study

Many possible directions

MCSM-like calculations (greedy algorithm)

Momentum-like coordinates (DGCM)

PGCM alone suited for ab initio?











Must exhaust the PGCM subspace


Mesh refinement















