

# Pentaquarks

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with results from the Belle Collaboration



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The excitation spectrum of a  $[Q\bar{Q}]$  state is well described by a semi-relativistic phenomenological potential (effective Cornell potential)

$$V(r) = -\frac{4}{3} \frac{\alpha_s(r)}{r} + \sigma r + \delta(1/r^2)$$

- A short-distance **colour potential**
- A long-distance **confinement** term
- **Spin-spin** and **spin-orbit** corrections

Developed in the 70's, particularly accurate to describe and predict the spectrum of  $[c\bar{c}]$  and  $[b\bar{b}]$  states.



All the unpredicted states are labelled as **exotic** states.

- They must contain a  $c\bar{c}$  pair as they all decay into a final state with a charmonium
- They do not present the same properties expected from a pure  $c\bar{c}$  state

As an example, look at X(3872):

- The first exotic state ever observed (Belle, 2003)
- Extremely narrow to be above the open charm threshold
- Radiative decay rates do not match prediction for a  $c\bar{c}$  state
- Decays into two different final states with different isospin (maximal violation)

Furthermore, the  $Z$  states are charged and this implies a minimal quark content of  $[c\bar{c}d\bar{u}]$

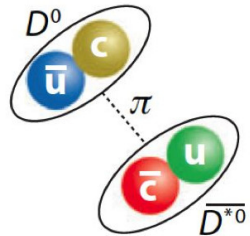
# Models for multiquark states

Several models have been proposed to describe the exotic states.

Main interpretations:

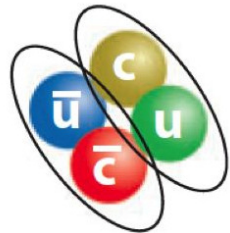
Mesonic (baryonic) molecule

- Low binding energy, narrow states
- Only S-wave, few states predicted
- Independently decaying components



Compact multiquark

- Tightly bound states
- Large widths in principle
- Many states expected



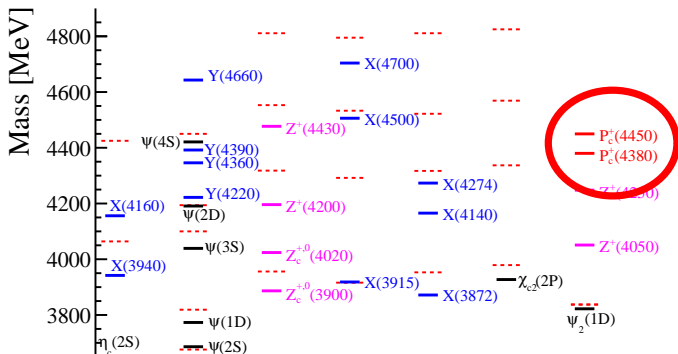
Other models are in principle allowed, as well as mixture of different models

# Pentaquarks

The charmonium spectrum is the ideal place to look for unexpected states

- Large mass difference between states wrt light  $[q\bar{q}]$  states
- Clean environment
- Wide range of detailed studies (better than bottomonium spectrum)

This presentation will focus on measurements and searches for states with 5 constituent quarks  $[qqqq\bar{q}]$ , in particular  $[qqqc\bar{c}]$

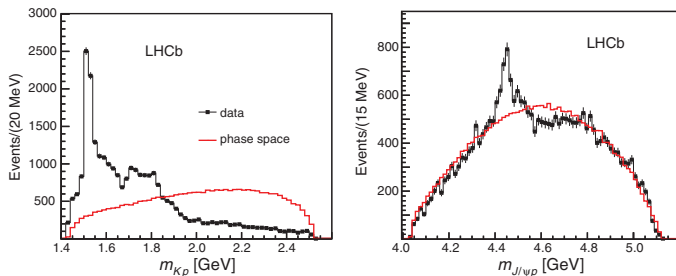




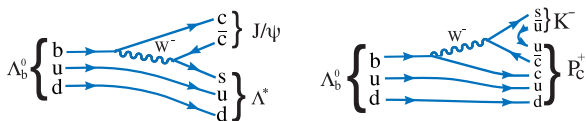
OBSERVATION  
OF  
PENTAQUARKS  
IN  
 $\Lambda_b^0 \rightarrow J/\psi K^- p$   
DECAYS  
(RUN 1)

# Analysis of $\Lambda_b^0 \rightarrow J/\psi K^- p$ decays

Structures are visible, over a non-resonant distribution, in the  $m_{J/\psi p}$  spectrum from  $\Lambda_b^0 \rightarrow J/\psi K^- p$  decays using the full LHCb Run 1 statistics ( $3 \text{ fb}^{-1}$ )



The resonant contributions are expected to be dominated by  $\Lambda^* \rightarrow K^- p$  decays, need to check if structures in  $m_{J/\psi p}$  are reflections in Dalitz plot



[Phys. Rev. Lett. 115, 072001 (2015)]

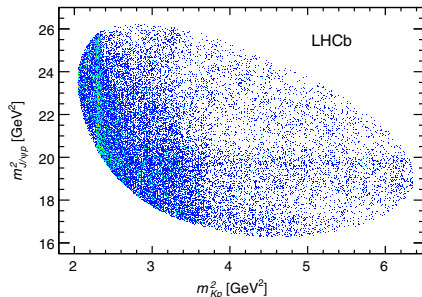
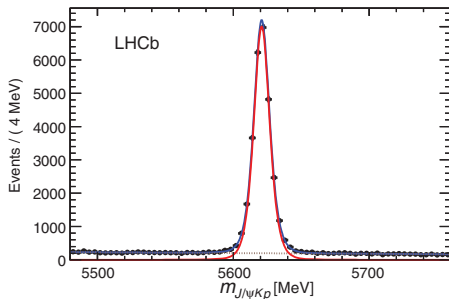


# Analysis strategy



- 14 well established  $\Lambda^* \rightarrow pK^-$  resonances to take into account
- 5 decay angles +  $m_{Kp}$  (6D fit)
- Helicity formalism
- Background-subtracted data

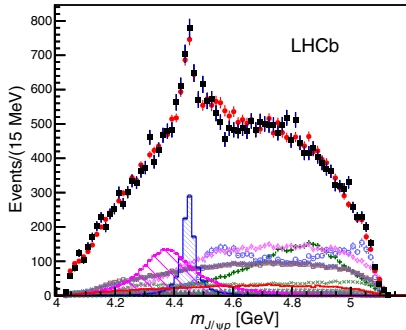
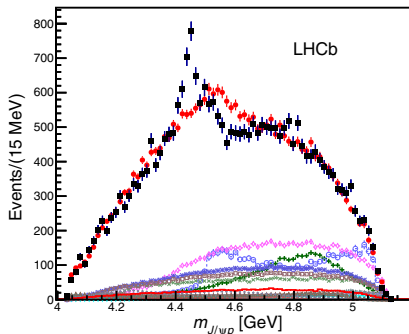
State	$J^P$	$M_0$ (MeV)	$\Gamma_0$ (MeV)	# Reduced	# Extended
$\Lambda(1405)$	$1/2^-$	$1405.1^{+1.3}_{-1.0}$	$50.5 \pm 2.0$	3	4
$\Lambda(1520)$	$3/2^-$	$1519.5 \pm 1.0$	$15.6 \pm 1.0$	5	6
$\Lambda(1600)$	$1/2^+$	1600	150	3	4
$\Lambda(1670)$	$1/2^-$	1670	35	3	4
$\Lambda(1690)$	$3/2^-$	1690	60	5	6
$\Lambda(1800)$	$1/2^-$	1800	300	4	4
$\Lambda(1810)$	$1/2^+$	1810	150	3	4
$\Lambda(1820)$	$5/2^+$	1820	80	1	6
$\Lambda(1830)$	$5/2^-$	1830	95	1	6
$\Lambda(1890)$	$3/2^+$	1890	100	3	6
$\Lambda(2100)$	$7/2^-$	2100	200	1	6
$\Lambda(2110)$	$5/2^+$	2110	200	1	6
$\Lambda(2350)$	$9/2^+$	2350	150	0	6
$\Lambda(2585)$	?	$\approx 2585$	200	0	6



[Phys. Rev. Lett. 115, 072001 (2015)]

# Fit projections and results

To have an acceptable fit two new  $P_c^+$  states need to be included



- Black points: data
- Red points: amplitude fit
- $P_c(4380)^+$ ,  $J^P = 3/2^-$ ,  $\Gamma = 205 \pm 18$  MeV, significance  $9\sigma$
- $P_c(4450)^+$ ,  $J^P = 5/2^+$ ,  $\Gamma = 39 \pm 5$  MeV, significance  $12\sigma$

[Phys. Rev. Lett. 115, 072001 (2015)]

# Model-independent confirmation

To confirm the previous result, the analysis is repeated using a different, model-independent approach.

- Minimal assumptions on the excited  $\Lambda^*$  spin and shapes
- Can include also nonresonant  $K^-p$  and  $\Sigma^*$  contributions

The strategy is to describe the 2D plane ( $m_{Kp, \cos \theta_{\Lambda^*}}$ ) expanding the helicity angle  $\theta_{\Lambda^*}$  in Legendre polynomials:

$$dN/d(\cos \theta_{\Lambda^*}) = \sum_{l=0}^{l_{max}} \langle P_l^U \rangle P_l(\cos \theta_{\Lambda^*})$$

where

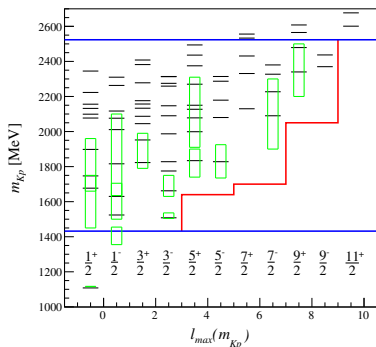
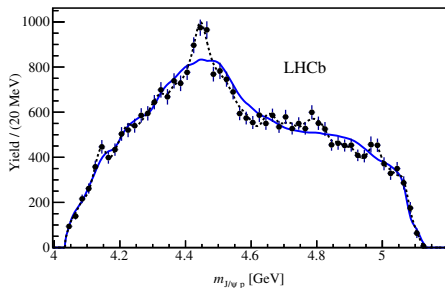
$$\langle P_l^U \rangle = \int_{-1}^{+1} d \cos \theta_{\Lambda^*} P_l(\cos \theta_{\Lambda^*}) dN/d(\cos \theta_{\Lambda^*})$$

and it is extracted from the  $m_{Kp}$  distribution in data.

If no exotic contribution is present and the structures in  $m_{J/\psi p}$  are due to reflections, then this expansion will be enough to describe the  $m_{J/\psi p}$  spectrum

# Model-independent confirmation

In practise, the Legendre moments include all contributions in  $K^-p$  with spin  $2J_{max}$  or less, depending on the given  $m_{Kp}$  range, up to  $J_{max} = 9/2$ .



By looking at  $m_{J/\psi p}$  it is clear that the distribution cannot be explained using only reflections. The discrepancy is more than  $9\sigma$ .

[Phys. Rev. Lett. 117, 082002 (2016)]

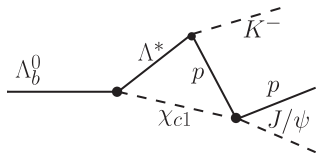


# EXOTIC RESONANCES AND RESCATTERING EFFECTS

# Rescattering effects

The narrow structure at  $4450 \text{ MeV}/c^2$  observed by LHCb happens to be located exactly at the  $\chi_{c1}p$  mass threshold. This can be a signal of a kinematic enhancement due to rescattering effects.

- All intermediate particles must be on shell to have a threshold enhancement
- The  $\Lambda^*$  mass must lie within a kinematically allowed mass range
- One happens to exist:  $\Lambda(1890)$



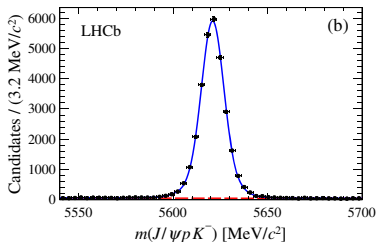
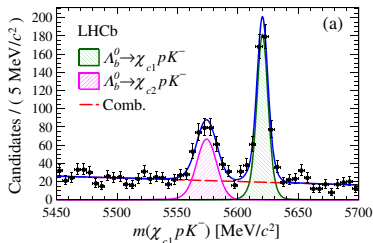
$\Rightarrow$  An observation of  $P_c(4450)^+$  decaying in the  $\chi_{c1}p$  final state (and not  $\chi_{c0,c2}p$ ) would confirm the exotic nature of the resonance

$\Rightarrow$  An observation of  $P_c(4450)^+ \rightarrow J/\psi p$  from  $\Lambda_b^0 \rightarrow J/\psi p \pi^-$  decays would be harder to accommodate in this picture (dominated by  $N^*$ )

# Search for $P_c^+ \rightarrow \chi_{c1}p$

First observation of the decays  $\Lambda_b^0 \rightarrow \chi_{c1}pK^-$  and  $\Lambda_b^0 \rightarrow \chi_{c2}pK^-$

- First investigation, with limited statistics ( $3 \text{ fb}^{-1}$ , full LHCb Run 1)
- $N(\Lambda_b^0 \rightarrow \chi_{c1}pK^-) = 453 \pm 25$
- Not enough to analyse the  $\chi_{c1}p$  mass spectrum, will be updated with Run 2 data
- First measurement of the branching fractions relative to  $\Lambda_b^0 \rightarrow J/\psi pK^-$
- $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \chi_{c1}pK^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi pK^-)} = 0.242 \pm 0.014 \pm 0.013 \pm 0.009$
- $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \chi_{c2}pK^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi pK^-)} = 0.248 \pm 0.020 \pm 0.014 \pm 0.009$

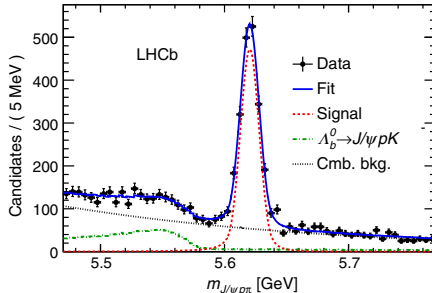


[Phys. Rev. Lett. 119, 062001 (2017)]

# Analysis of the $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ channel

- Data:  $3 \text{ fb}^{-1}$ , full LHCb Run 1
- Thanks to the  $\Delta I = 1/2$  rule the  $\Lambda^*$  contributions are suppressed
- 14 well established  $N^* \rightarrow p\pi^-$  resonances to take into account
- 5 decay angles +  $m_{Kp}$
- Helicity formalism
- Background-subtracted data

State	$J^P$	Mass (MeV)	Width (MeV)	RM	EM
$NR \ p\pi$	$1/2^-$	...	...	4	4
$N(1440)$	$1/2^+$	1430	350	3	4
$N(1520)$	$3/2^-$	1515	115	3	3
$N(1535)$	$1/2^-$	1535	150	4	4
$N(1650)$	$1/2^-$	1655	140	1	4
$N(1675)$	$5/2^-$	1675	150	3	5
$N(1680)$	$5/2^+$	1685	130	...	3
$N(1700)$	$3/2^-$	1700	150	...	3
$N(1710)$	$1/2^+$	1710	100	...	4
$N(1720)$	$3/2^+$	1720	250	3	5
$N(1875)$	$3/2^-$	1875	250	...	3
$N(1900)$	$3/2^+$	1900	200	...	3
$N(2190)$	$7/2^-$	2190	500	...	3
$N(2300)$	$1/2^+$	2300	340	...	3
$N(2570)$	$5/2^-$	2570	250	...	3
Free parameters				40	106

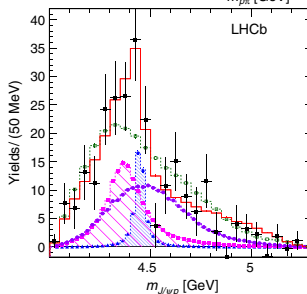
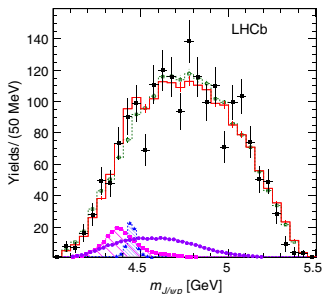
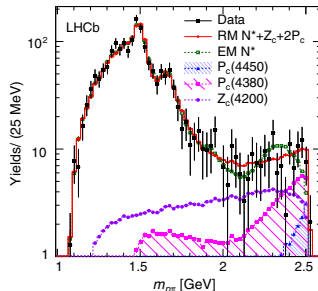


[Phys. Rev. Lett. 117, 082003 (2016)]



# Fit projections and results

- Adding  $P_c(4380)^+$ ,  $P_c(4450)^+$  and a  $Z_c(4200)^- \rightarrow J/\psi\pi^-$  contribution significantly improves the fit
- $P_c^+$  production rates as expected from previous observation (including Cabibbo suppression)
- Combined significance:  $3.1\sigma$



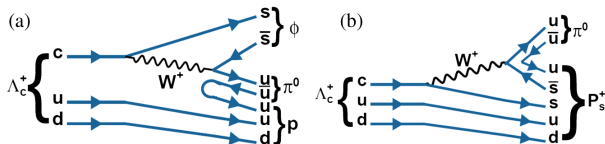
[Phys. Rev. Lett. 117, 082003 (2016)]



# RECENT SEARCHES FOR STRANGE AND BEAUTY PENTAQUARKS

# Search for $s$ -flavoured pentaquarks

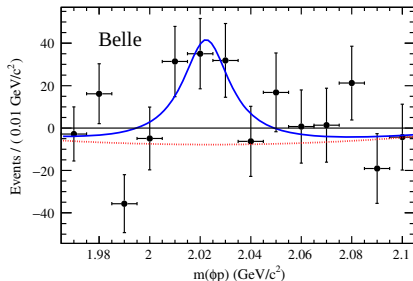
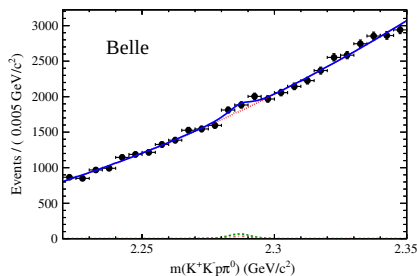
- Strange-flavour analogue of the  $P_c^+$  discovery channel:  $\Lambda_c^+ \rightarrow \phi p \pi^0$
- This channel has never been studied before
- Dataset:  $915 \text{ fb}^{-1}$  at  $\Upsilon(4S)$  and  $\Upsilon(5S)$  collected by the Belle experiment
- $P_s^+$  can be observed as peak in the  $\phi p$  mass spectrum if the same production mechanism holds, and if  $m_{P_s^+} < m_{\Lambda_c^+} - m_{\pi^0}$



[Phys. Rev. D 96, 051102 (2017)]

# Search for $s$ -flavoured pentaquarks

No signal is observed in a mass window of 20 MeV/ $c^2$  around the  $\phi$  peak, upper limits at 90%CL are set on the branching fraction product, normalised using  $\Lambda_c^+ \rightarrow pK^- \pi^+$  decays



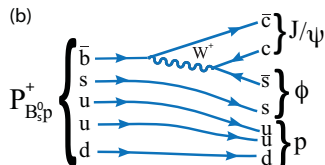
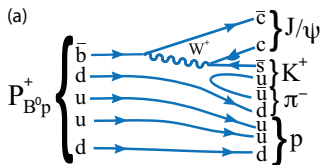
$$\mathcal{B}(\Lambda_c^+ \rightarrow P_s^+ \pi^0) \times \mathcal{B}(P_s^+ \rightarrow \phi p) < 8.3 \times 10^{-5}$$

(as a reference)

$$\mathcal{B}(\Lambda_b^0 \rightarrow P_c(4450)^+ K^-) \times \mathcal{B}(P_c(4450)^+ \rightarrow J/\psi p) = (1.3 \pm 0.4) \times 10^{-5}$$

# Search for $b$ -flavoured pentaquarks

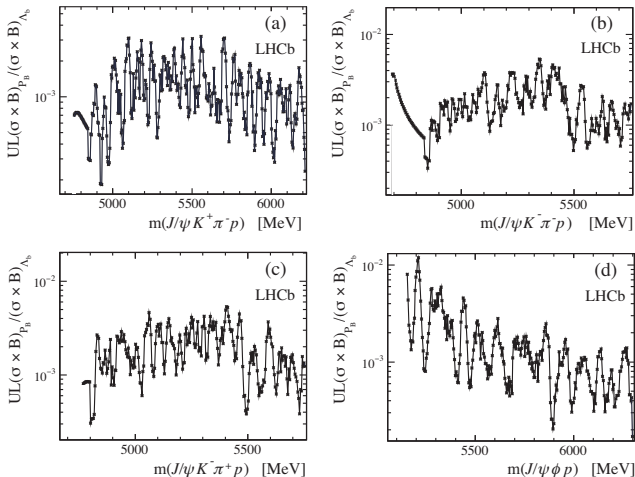
- According to the Skyrme model, the heavier the constituent quarks are, the more tightly bound the state is
- No searches for  $b$ -flavoured pentaquarks have ever been published
- Full LHCb Run 1 integrated luminosity ( $3 \text{ fb}^{-1}$ )
- Four different states considered:
  - $P_{B^0 p}^+ \rightarrow J/\psi K^+ p \pi^-$
  - $P_{\Lambda_b^0 \pi^+}^+ \rightarrow J/\psi K^+ p \pi^+$
  - $P_{\Lambda_b^0 \pi^-}^- \rightarrow J/\psi K^+ p \pi^-$
  - $P_{B_s^0 p}^+ \rightarrow J/\psi \phi p$
- Mass ranges chosen to be below the strong decay threshold



[RSPA 260, 1300 (1961)], [Phys. Rev. D 97, 032010 (2018)]

# Search for $b$ -flavoured pentaquarks

No signal is observed, upper limits at 90%CL are set on the production cross sections times the BR, normalised using  $\Lambda_b^0 \rightarrow J/\psi K^- p$  decays



[Phys. Rev. D 97, 032010 (2018)]

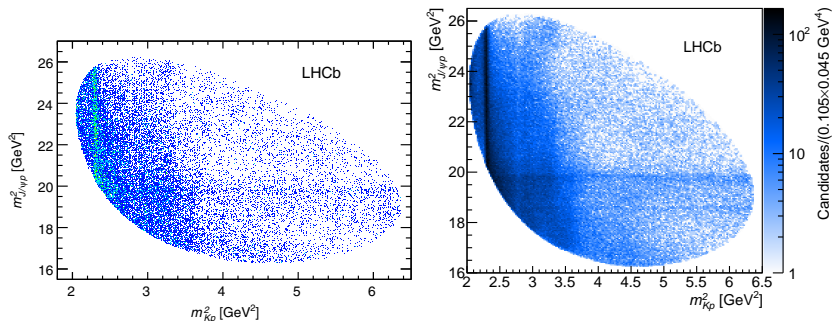


OBSERVATION  
OF  
PENTAQUARKS  
IN  
 $\Lambda_b^0 \rightarrow J/\psi K^- p$   
DECAYS  
(RUN 1 + RUN 2)

# Update with full Run 1 and Run 2 statistics



- Latest LHCb result on pentaquark searches: update of 2015 analysis
- Integrated luminosity  $9 \text{ fb}^{-1}$ , better data selection, increase in production cross-section (13 TeV instead of 7 and 8 TeV)
- 9 times more statistics  $\implies$  improved resolution on mass spectra

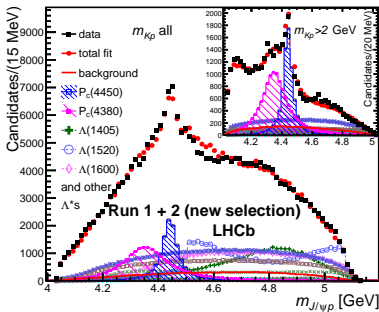
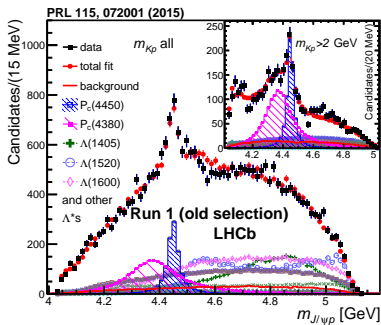


[arXiv:1904.03947]

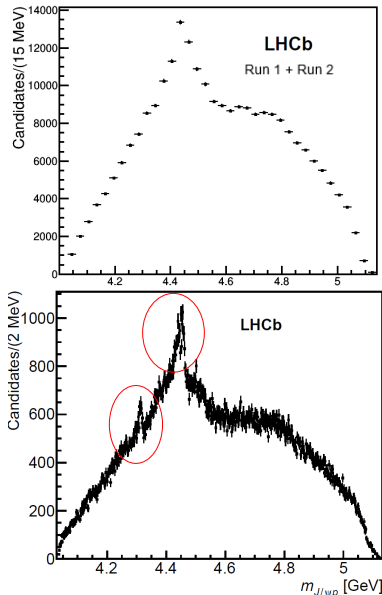


# Consistency check

First check: using the new dataset, the new selection and the same amplitude model we get compatible results



# New features

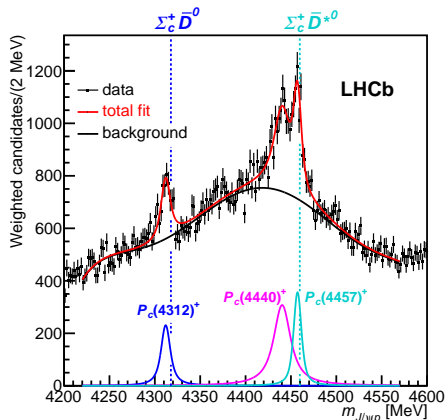


- Increase in mass resolution ( $\approx 2.5$  MeV)
- New narrow structure at 4.3 GeV,  $P_c(4450)^+$  is resolved into two peaks
- Amplitude fit computationally challenging, currently work in progress
- Very narrow states, cannot be artificial reflections
- Cut at  $m_{Kp} > 1.9$  GeV to suppress the dominant  $\Lambda^* \rightarrow pK^+$  contributions
- 1-dimensional fit using different composition of  $\Lambda^*$  reflections to model the background
- This analysis is not sensitive to broad  $J/\psi p$  contribution, like  $P_c(4380)^+$

[arXiv:1904.03947]

# Fit to the $J/\psi p$ invariant mass

- The masses of the narrow peaks are just below the  $\Sigma_c^+ \bar{D}^{(*)0}$  masses
- Although the compact pentaquark model is not ruled out, these features favour the molecular interpretation
- Need to measure quantum numbers and find isospin partners in order to have a definitive answer



State	$M$ [MeV]	$\Gamma$ [MeV]	(95% CL)	$\mathcal{R}$ [%]
$P_c(4312)^+$	$4311.9 \pm 0.7^{+6.8}_{-0.6}$	$9.8 \pm 2.7^{+3.7}_{-4.5}$	(< 27)	$0.30 \pm 0.07^{+0.34}_{-0.09}$
$P_c(4440)^+$	$4440.3 \pm 1.3^{+4.1}_{-4.7}$	$20.6 \pm 4.9^{+8.7}_{-10.1}$	(< 49)	$1.11 \pm 0.33^{+0.22}_{-0.10}$
$P_c(4457)^+$	$4457.3 \pm 0.6^{+4.1}_{-1.7}$	$6.4 \pm 2.0^{+5.7}_{-1.9}$	(< 20)	$0.53 \pm 0.16^{+0.15}_{-0.13}$

[arXiv:1904.03947]



# CONCLUSIONS

- Exotic spectroscopy is an extremely rich and productive field
- Several observations and searches for pentaquark states in the last 4 years
- Quite a recent discovery - this is just the beginning of a new era in both discovery of new states and understanding of QCD binding mechanisms
- We still do not know what the real nature of these new states is
- The LHCb Run 2 update measurement is the strongest evidence so far towards a molecular interpretation of the  $P_c^+$  states
- Amplitude analysis is challenging, but ongoing
- LHCb clearly dominates the scene for now, waiting for Belle II to join



# BACKUP