



# $b \rightarrow s\ell\ell$ decays and $B \rightarrow K^*\ell\ell$ angular distributions

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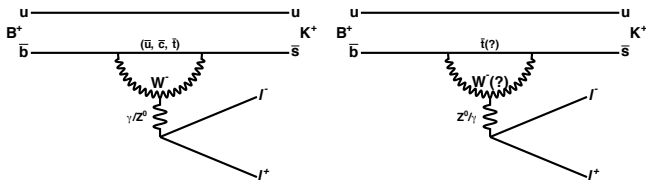


Angular Analysis of  $B \rightarrow K^* \ell \ell$  [Belle, PRL 118, 111801 (2017)]

Test of LFU in  $B \rightarrow K^* \ell \ell$  decays [Belle, arXiv: 1904.02440]  
Test of LFU in  $B \rightarrow K \ell \ell$  decays [LHCb, arXiv: 1903.09252]

Search for LFV  $B^0 \rightarrow K^{*0} \mu e$  decays [Belle, PRD 98, 071101 (2018)]

- $B \rightarrow K\ell\ell$  and  $B \rightarrow K^*\ell\ell$  involve quark transition from  $b \rightarrow s$  which are FCNCs. These processes occur through penguin loop and box diagrams in SM.



- Global analysis of  $B$  decays hint at lepton flavor non universality.
- These decays are highly suppressed and very small BR ( $\mathcal{O}(10^{-6})$ ).
- These decays are very sensitive to NP.
- Rare b-hadron decays place strong constraints on many BSM models by probing energy scales higher than direct searches.

## New physics can contribute by:

- enhancing or suppressing decay rates.
- modifying the angular distribution of the final state particles.

- The amplitude of a hadron decay process is described as:

$$A(M \rightarrow F) = \langle F | \mathcal{H}_{eff} | M \rangle = \frac{G_F}{\sqrt{2}} \sum_i V_{CKM}^i C_i(\mu) \langle F | O_i(\mu) | M \rangle$$

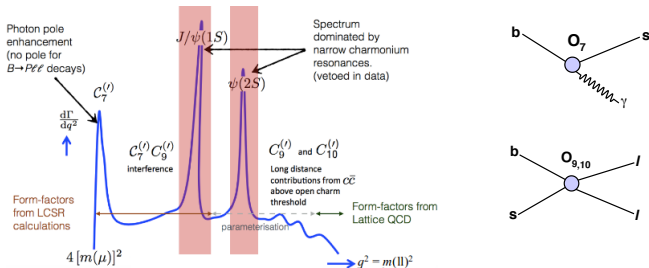
Wilson coefficients  $C_i$  = Perturbative short distance effects

Operators  $O_i$  = non-perturbative long distance effects.

$i = 7$  : Photon penguin

$i = 9, 10$  : Electroweak penguin

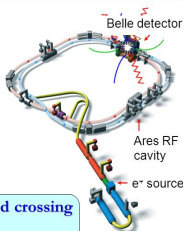
- NP can affect SM operator contributions (Wilson coefficients) and/or enter through new operators.



- Contribution of  $C_7$ ,  $C_9$  and  $C_{10}$  depends on  $q^2$  (invariant mass square of two leptons).

- The angular analysis for  $B \rightarrow K^* \ell \ell$  has complex angular distribution that provides many observables sensitive to different types of BSM physics.
- Each observable depends on different Wilson coefficients and form-factors. [S. Descotes-Genon et al. JHEP 01(2013) 048]
- In the SM, couplings of the gauge bosons to leptons are independent of lepton flavour. [G. Hiller and M. Schmaltz JHEP02(2015) 055]
- Branching fractions of  $e$ ,  $\mu$  and  $\tau$  differ only by phase space and helicity-suppressed contributions.
- Any sign of lepton non-universal interaction would be a direct sign of new physics.
- NP models accommodating LFU violation, will also show LFV [S.L. Glashow et.al PRL 114, 091801 (2015)].

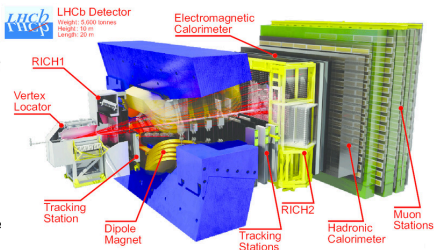
## KEKB Accelerator



8 GeV e<sup>-</sup> x 3.5 GeV e<sup>+</sup>; 22mrad crossing  
 $L_{\text{peak}} = 2.11 \times 10^{34}$   
 Integ. Lum. ~1040 fb<sup>-1</sup>

$$e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$$

- The Belle experiment is located at the KEKB accelerator in Tsukuba, Japan.
- Data taking from 1999 to 2010.
- It is designed as a B-factory.
- Data collected: 772 million  $B\bar{B}$  pairs.



- Located at the CERN LHC proton-proton collider
- Forward spectrometer with Vertex, Tracking, PID and Calorimetry
- Run 1, Collision energies 7,8 TeV (2011,2012) 3 fb<sup>-1</sup>
- Run 2, Collision energy 13 TeV (2015,2016) 3 fb<sup>-1</sup>

The differential decay rate is given by:

$$\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d \cos \theta_\ell d \cos \theta_K d\phi dq^2} =$$

$$\frac{9}{32\pi} \left[ \frac{3}{4} (1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4} (1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + \right.$$

$$S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_6 \sin^2 \theta_K \cos \theta_\ell +$$

$$\left. S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

- In the lepton massless limit there are eight independent observables:

$F_L$ : Fraction of the longitudinal polarization of the  $K^*$

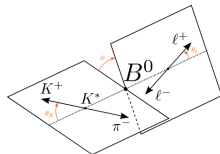
$S_6$ : The forward-backward asymmetry of the dimuon system

$S_{3,4,5,7,8,9}$ : The remaining CP-averaged observables

- $F_L$  and  $S_i$  are function of  $q^2$ .
- Observable  $P'_i$  and  $Q_i$

$$P'_{i=4,5,6,8} = \frac{S_{j=4,5,7,8}}{\sqrt{F_L(1 - F_L)}} \quad \text{JHEP 05(2013) 137}$$

$$Q_i = P'_i - P_i^e, \quad i = 4, 5 \quad \text{JHEP 10(2016) 075}$$



- $P'_i$  are free of form-factor uncertainties.
- Any deviation from zero for  $Q_i$ , would be a hint for NP.

- The channels reconstructed for analysis are

$$\begin{aligned}
 B^0 &\rightarrow K^{*0} \mu^+ \mu^-, & B^+ &\rightarrow K^{*+} \mu^+ \mu^- \\
 B^0 &\rightarrow K^{*0} e^+ e^-, & B^+ &\rightarrow K^{*+} e^+ e^-
 \end{aligned}$$

- $K^*$  is reconstructed from:

$$\begin{aligned}
 K^{*0} &\rightarrow K^+ \pi^- \\
 K^{*+} &\rightarrow K^+ \pi^0 \\
 K^{*+} &\rightarrow K_S^0 \pi^+
 \end{aligned}$$

- Multivariate analysis technique (NN) is used to identify each particle type in the decay chain.
- Kinematic variables which distinguish signal from background are

$$\begin{aligned}
 M_{bc} &= \sqrt{E_{beam}^2/c^4 - |p_B|^2/c^4} \\
 \Delta E &= E_B - E_{beam}
 \end{aligned}$$

- Requirement on kinematic variables:

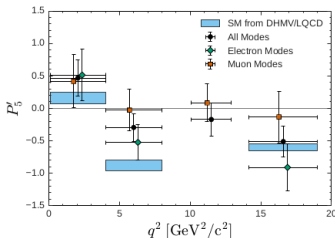
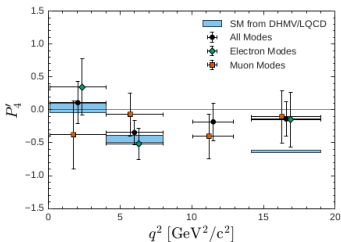
$$5.22 < M_{bc} < 5.20 \text{ GeV}/c^2 \text{ and } -0.10 \text{ } (-0.05) < \Delta E < 0.05 \text{ GeV for } ee(\mu\mu)$$

- Final selection requirement on the top-level NN is optimized by maximizing a figure of merit

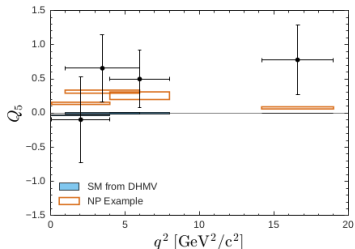
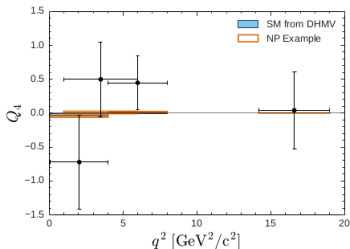
$$\text{FOM} = \frac{n_s}{\sqrt{n_s + n_b}}$$

- Extended maximum likelihood fit to extract signal.





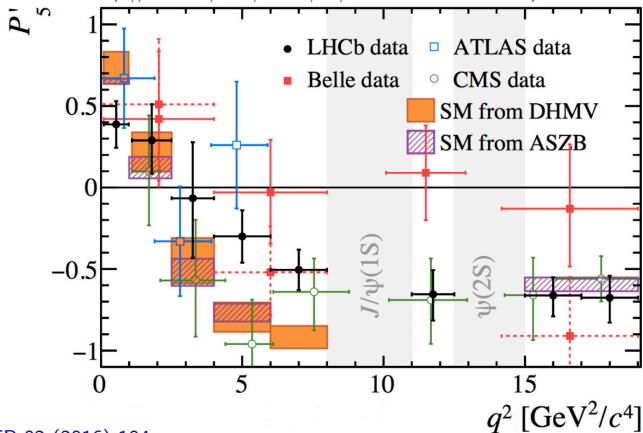
- All measurements are compatible with SM predictions.
- The strongest tension of  $2.6\sigma$  is observed in  $P'_5$  of the muon modes for the region  $4 < q^2 < 8 \text{ GeV}^2/c^4$ .
- For  $4 < q^2 < 8 \text{ GeV}^2/c^4$  bin,  $1.3\sigma$  deviation is found in electron mode.
- Combining muon and electron modes, deviation is  $2.5\sigma$ .



- $Q_{4,5}$  observables show no significant deviation from zero.

$q^2$ GeV <sup>2</sup> /c <sup>4</sup>	$Q_4$	$Q_5$
[1.00, 6.00]	$0.498 \pm 0.527 \pm 0.166$	$0.656 \pm 0.485 \pm 0.103$
[0.10, 4.00]	$-0.723 \pm 0.676 \pm 0.163$	$-0.097 \pm 0.601 \pm 0.164$
[4.00, 8.00]	$0.448 \pm 0.392 \pm 0.076$	$0.498 \pm 0.410 \pm 0.095$
[14.18, 19.00]	$0.041 \pm 0.565 \pm 0.082$	$0.778 \pm 0.502 \pm 0.065$

P. Cartelle, Dark Matter @ LHC Heidelberg, April 2018

<https://cds.cern.ch/record/2311960/files/FlavourAnomaliesLHCbAlvarez.pdf>

- LHCb [JHEP 02 \(2016\) 104](#)
- Belle: [PRL 118 \(2017\) 111801](#)
- ATLAS: [JHEP 10 \(2018\) 517](#)
- CMS: [Phys. Lett. B 781\(2018\) 517](#)

$3.4\sigma$  deviation for LHCb  $4 < q^2 < 6 \text{ GeV}^2/c^4$

# Test of LFU ( $R_{K^*}$ ) for $B \rightarrow K^* \ell \ell$ Prior to Moriond, 2019

- LHCb measurement of

$$R_{K^*} = \frac{BR(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{BR(B^0 \rightarrow K^{*0} e^+ e^-)}$$

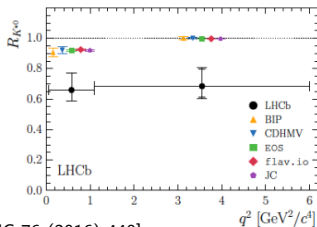
shows deviations from SM expectation.

$$R_{K^*}(0.045 < q^2 < 1.1 \text{ GeV}^2/c^4) = 0.66^{+0.11}_{-0.07} \pm 0.03$$

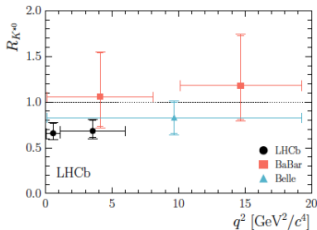
$$R_{K^*}(1.1 < q^2 < 6 \text{ GeV}^2/c^4) = 0.69^{+0.11}_{-0.07} \pm 0.05$$

- Compatibility with the SM estimated to be at the level of  $2.1 - 2.3\sigma$  for low  $q^2$  and  $2.4 - 2.5\sigma$  at central  $q^2$  for a data sample of  $3\text{fb}^{-1}$ .
- Belle measurement for whole  $q^2$  region,  $R_{K^*} = 0.83 \pm 0.17 \pm 0.08$ , is consistent with SM prediction.
- BaBar measured for low and high  $q^2$  bins and are consistent with SM with high uncertainty.

P. Cartelle, Dark Matter @ LHC Heidelberg, April 2018  
<https://cds.cern.ch/record/2311960/files/FlavourAnomaliesLHCbAlvarez.pdf>



- ▲ BIP [EPJC 76 (2016) 440]
- ▼ CDHMV [JHEP 04 (2017) 016]
- EOS [PRD 95 (2017) 035029]
- ◆ flav.io [EPJC 77 (2017) 377]
- JC [PRD 93 (2016) 014028]



- LHCb [JHEP 08(2017) 055]
- BaBar [PRD 86 (2012) 032012]
- ▲ Belle [PRL 103 (2009) 171801]

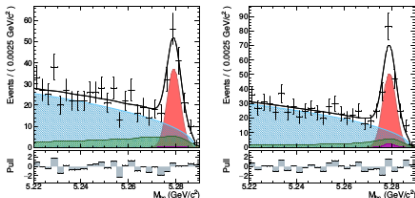
- Similar particle selection and fitting procedure as that of  $B \rightarrow K^* \ell \ell$  angular analysis.

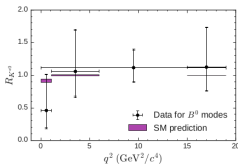
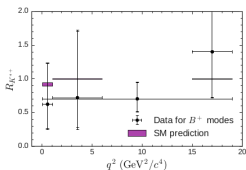
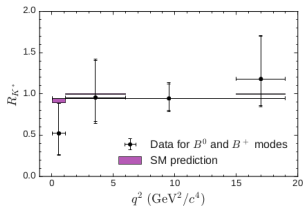
- The  $B \rightarrow K^* \ell \ell$  mode is reconstructed by hierarchical NN.
- Background is suppressed by multivariate analysis technique which uses event topology and NN outputs.
- $B \rightarrow KJ/\psi$  is used as a control sample.
- Signal is extracted using extended maximum likelihood fit.
- $R_{K^*}$  is calculated.

$$R_{K^*} = \frac{B \rightarrow K^* \mu \mu}{B \rightarrow K^* e e}$$

- Performed extended maximum likelihood fit.
- PDF parameterization:
  - Signal : CB with the shape parameters determined from  $B \rightarrow K^* J/\psi (\rightarrow \ell\ell)$  candidates in  $J/\psi$  region.
  - Combinatorial background: Argus shape.
  - Events from charmonium decay: The events which pass the charmonium veto because of misreconstruction are modeled with kernel density function.
  - Peaking background: Peaking from double flavor misidentification fitted with KDF.
- The normalization from peaking and charmonium are derived from MC and fixed in the fit for the signal yield.
- Example fit for  $q^2 > 0.045 \text{ GeV}^2/c^4$ .
- $103.0^{+13.4}_{-12.7}$  and  $139.9^{+16.0}_{-15.4}$  events for electron and muon modes, respectively.

Belle [[arXiv:1904.02440](https://arxiv.org/abs/1904.02440)]





First measurement of  $R_{K^{*+}}$

$q^2$ GeV <sup>2</sup> /c <sup>4</sup>	All modes	$B^0$ modes	$B^+$ modes
[0.045, 1.1]	$0.52^{+0.36}_{-0.26} \pm 0.05$	$0.46^{+0.55}_{-0.27} \pm 0.07$	$0.62^{+0.60}_{-0.36} \pm 0.10$
[1.1, 6]	$0.96^{+0.45}_{-0.29} \pm 0.11$	$1.06^{+0.63}_{-0.38} \pm 0.13$	$0.72^{+0.99}_{-0.44} \pm 0.18$
[0.1.8]	$0.90^{+0.27}_{-0.21} \pm 0.10$	$0.86^{+0.33}_{-0.24} \pm 0.08$	$0.96^{+0.56}_{-0.35} \pm 0.14$
[15 - 19]	$1.18^{+0.52}_{-0.32} \pm 0.10$	$1.12^{+0.61}_{-0.36} \pm 0.10$	$1.40^{+1.99}_{-0.68} \pm 0.11$
[0.045, ]	$0.94^{+0.17}_{-0.14} \pm 0.08$	$1.12^{+0.27}_{-0.21} \pm 0.09$	$0.70^{+0.24}_{-0.19} \pm 0.07$

- All results are found to compatible with SM prediction.

# Test of LFU ( $R_K$ ) for $B \rightarrow K\ell\ell$ Prior to Moriond 2019

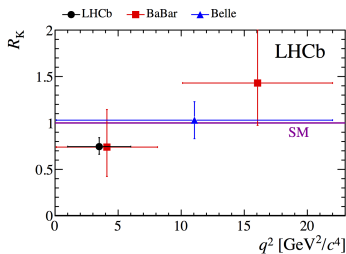
- SM prediction is very accurate.  $R_K^{(SM)} = 1 \pm \mathcal{O}(10^{-2})$

- LHCb (PRL 113, 151601(2014)) shows deviation from SM

$$R_K = \frac{BR(B^+ \rightarrow K^+ \mu^+ \mu^-)}{BR(B^+ \rightarrow K^+ e^+ e^-)} = 0.745_{-0.074}^{+0.090} \pm 0.036$$

in  $q^2 = [1 - 6] \text{ GeV}^2/c^4$  :  $2.6\sigma$  tension for  $3\text{fb}^{-1}$  data sample (2011-12 data).

- This observable is theoretically very clean, as most of the hadronic uncertainties cancel out in the ratio.
- The value of  $R_K$  for Belle was consistent with unity within the uncertainty limit measured for a data sample of  $605\text{fb}^{-1}$ .



Bin	$R_K$	Collaboration
$1 < q^2 < 6$	$0.745_{-0.074}^{+0.090} \pm 0.036$	LHCb (2014)
$1.1 < q^2 < 6$	$0.846_{-0.054}^{+0.060+0.016} - 0.014$	LHCb (2019)
whole $q^2$	$1.03 \pm 0.19 \pm 0.06$	Belle
$0.10 < q^2 < 8.12$	$0.74_{-0.31}^{+0.40} \pm 0.06$	BaBar
$q^2 > 10.11$	$1.43_{-0.44}^{+0.65} \pm 0.12$	BaBar



- The analysis of 2011 and 2012 data is re-optimized and analysis strategy is re-designed.
- 2015 and 2016 LHCb data are added.
- This analysis uses twice as many  $B$ 's as the previous analysis.

- Electron and muon tracks are very different in LHCb.
  - Because of bremsstrahlung, electron has worse  $q^2$  resolution and low reconstruction efficiency.
  - Better PID and trigger performances for muons.
- Use double ratio to cancel out most of the systematic uncertainties.

$$R_K = \frac{BR(B^+ \rightarrow K^+ \mu^+ \mu^-)}{BR(B^+ \rightarrow K^+ e^+ e^-)} \bigg/ \frac{BR(B^+ \rightarrow K^+ J/\psi(\rightarrow \mu^+ \mu^-))}{BR(B^+ \rightarrow K^+ J/\psi(\rightarrow e^+ e^-))}$$

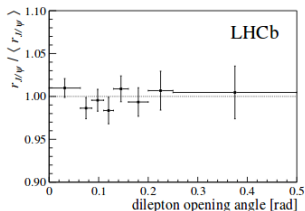
$$\Rightarrow R_K = \frac{N(K^+ \mu\mu)}{N(K^+ J/\psi(\mu\mu))} \cdot \frac{N(K^+ J/\psi(ee))}{N(K^+ ee)} \cdot \frac{\varepsilon(K^+ J/\psi(\mu\mu))}{\varepsilon(K^+ \mu\mu)} \cdot \frac{\varepsilon(K^+ ee)}{\varepsilon(K^+ J/\psi(ee))}$$

- To check efficiencies are correct,

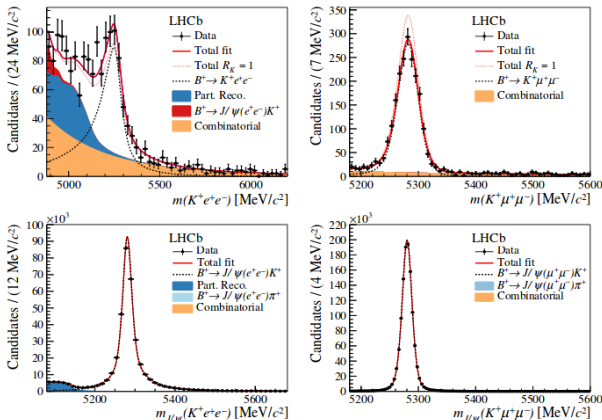
$$r_{J/\psi} = \frac{BR(B^+ \rightarrow K^+ J/\psi(\rightarrow \mu\mu))}{BR(B^+ \rightarrow K^+ J/\psi(\rightarrow ee))} = 1.0$$

$r_{J/\psi}$  is found to be  $1.014 \pm 0.035$  (stat. + syst.)

- Efficiencies should be understandable as a function of any variable.
  - $r_{J/\psi}$  should be flat for all variables examined.



- A single fit to the  $m(K^+ \ell^+ \ell^-)$  distributions is performed to determine  $R_K$  from the entire 2011-2016 dataset.



- The red-dotted line shows the distribution that would be expected from the observed number of  $B^+ \rightarrow K^+ \mu^+ \mu^-$  or  $B^+ \rightarrow K^+ e^+ e^-$  decays and  $R_K = 1$ .

- Using 2011 and 2012 LHCb data,  $R_K$  was:

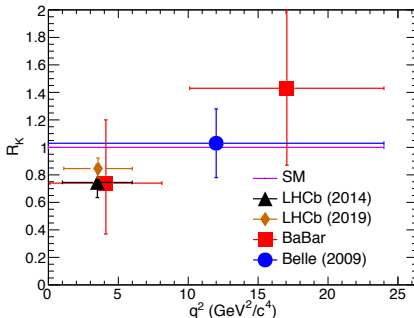
$$R_K[1.0 - 6.0] = 0.745_{-0.074}^{+0.090} \pm 0.036$$

2.6 $\sigma$  from SM prediction.

- Adding 2015 and 2016 data,  $R_K$  become:

$$R_K([1.1 - 6.0]) = 0.846_{-0.054}^{+0.016+0.060}$$

$\sim 2.5\sigma$  from SM.



- The deviation from SM expectation in  $R_K$  and  $R_{K^*}$  from LHCb result possibly show LFU violation.
- LFV can come together with LFU violation [S. L. Glashow et.al PRL 114, 091801 (2015)].
- Belle has published LFV decays  $B^0 \rightarrow K^{*0} \ell \ell'$ , where  $\ell = \mu, e$  [PRD 98.071101(2018)].

## Particle selection and Background suppression

- Charged particles are selected which satisfy PID criteria and originate from a region near the  $e^+e^-$  interaction point.
- Kaon and pion candidates are combined to form  $K^{*0}$ .
- $B$  candidate is reconstructed by combining  $K^{*0}$ ,  $\mu^\pm$  and  $e^\pm$  candidates.
- Constraint on kinematic variables are

$$M_{bc} > 5.2 \text{ GeV}/c^2$$

$$-0.05 < \Delta E < 0.04 \text{ GeV}$$

- Strong contribution from continuum ( $q\bar{q}$ ) and generic  $B$  ( $B\bar{B}$ ) backgrounds.
- Two stage  $NN$  is used to suppress the backgrounds. *i.e.*, Optimization of generic  $B$  background from the optimal cut of continuum background.

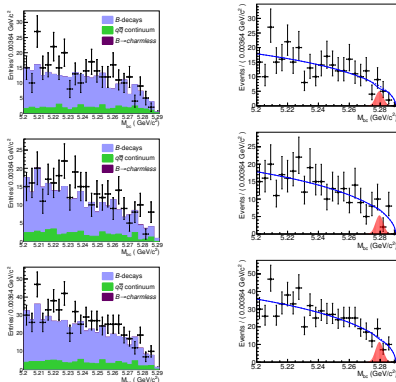
## Peaking Background

Peaking background due  $B^0 \rightarrow K^{*0}(\rightarrow K^+\pi^-)J/\psi(\rightarrow \ell\ell)$ , and PID misidentification

We veto:

- For  $B^0 \rightarrow K^{*0}\mu^+e^-$ 
  - $M(\ell^+\ell^-) \notin [3.04, 3.12] \text{ GeV}/c^2$
  - $M(K^+e^-) \notin [2.90, 3.12] \text{ GeV}/c^2$
  - $M(\pi^-\mu^+) \notin [3.06, 3.12] \text{ GeV}/c^2$
- For  $B^0 \rightarrow K^{*0}\mu^-e^+$ 
  - $M(\ell^+\ell^-) \notin [3.02, 3.12] \text{ GeV}/c^2$
  - $M(\pi^-e^+) \notin [3.02, 3.12] \text{ GeV}/c^2$

- Good agreement between data and MC.
- No evidence of signal observed  $\rightarrow$  upper limit is estimated.



Mode	$\epsilon$ (%)	$N_{sig}$	$N_{sig}^{(UL)}$	$\mathcal{B}^{(UL)} (10^{-7})$
$B^0 \rightarrow K^{*0} \mu^+ e^-$	8.8	$-1.5^{+4.7}_{-4.1}$	5.2	1.2
$B^0 \rightarrow K^{*0} \mu^- e^+$	9.3	$0.4^{+4.8}_{-4.5}$	7.4	1.6
$B^0 \rightarrow K^{*0} \mu^\pm e^\mp$	9.0	$-1.2^{+6.8}_{-6.2}$	8.0	1.8

- The angular analysis variable  $P'_5$  of  $B \rightarrow K^* \ell \ell$  show a deviation of  $2.5\sigma$  from SM prediction for the bin of  $4 < q^2 < 8 \text{ GeV}^2/c^4$ . This deviation is maximum for  $\mu\mu$  mode.
- $R_{K^*}$  measurements are compatible with SM prediction for Belle data.
- Updated  $R_K$  analysis from LHCb has a significantly improved precision. There is  $\sim 2.5\sigma$  tension with SM.
- Most stringent upper limit is found for  $B^0 \rightarrow K^{*0} \mu e$  mode.
- More data from LHCb collected in 2017 and 2018 is being analyzed.
- Upgraded LHCb detector will collect many times more data in the early 2020's.
- The Belle II experiment (50 times more data than Belle) will also provide stringent limits on any deviation from SM predictions.
- Lots more data to come!