

Multiphonon excitations in  
dark matter direct detection:  
beyond the spherical ~~cow~~  
guinea pig

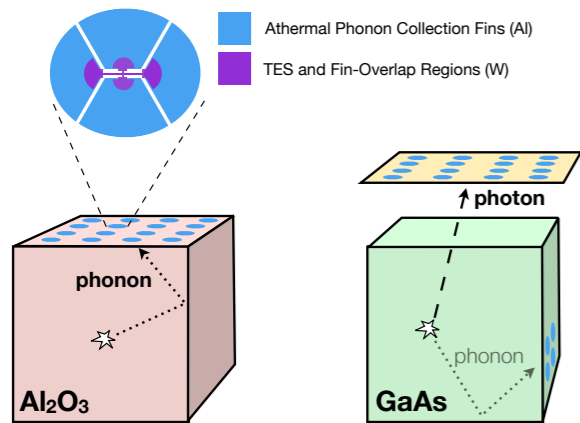
Tongyan Lin

UCSD

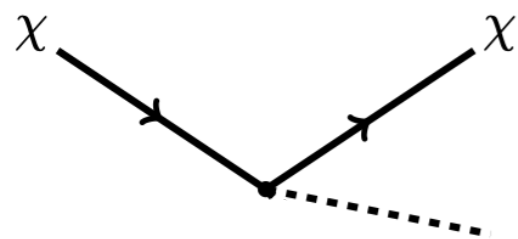
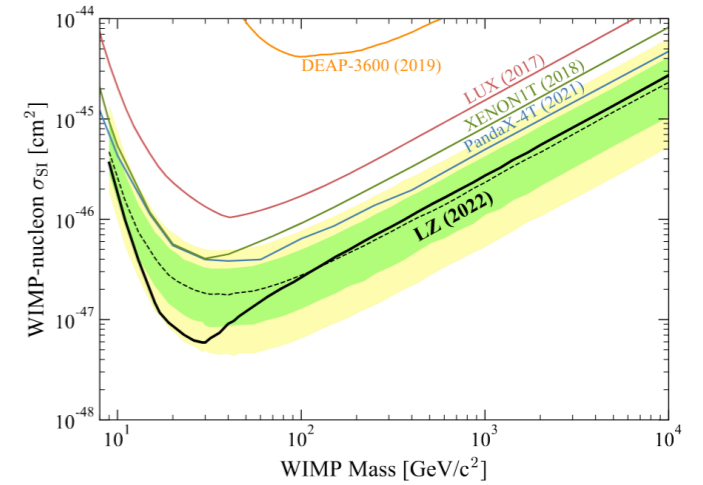
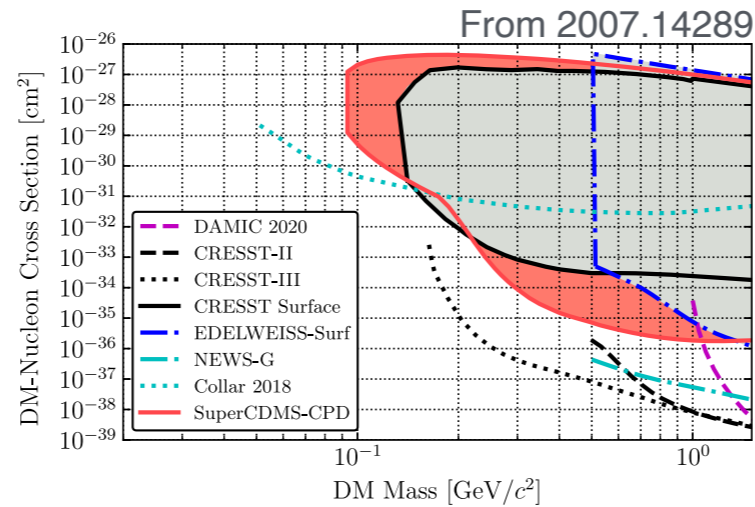
GUINEAPIG Workshop

August 20, 2024

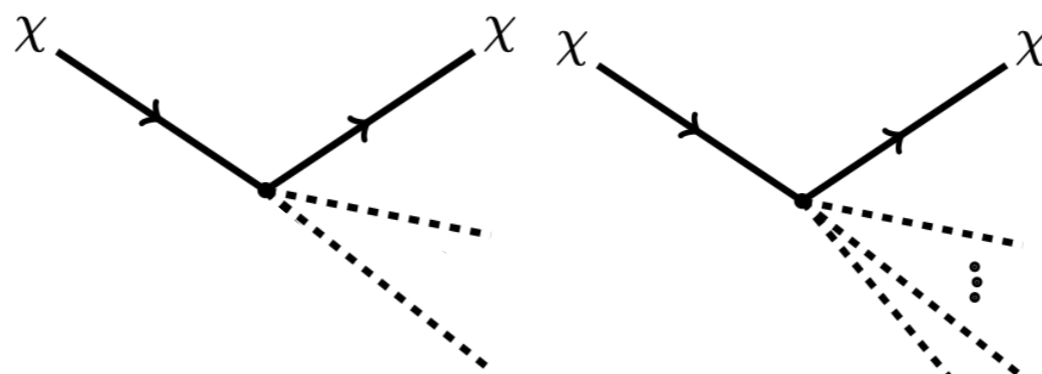
# DM-nucleus scattering in crystals



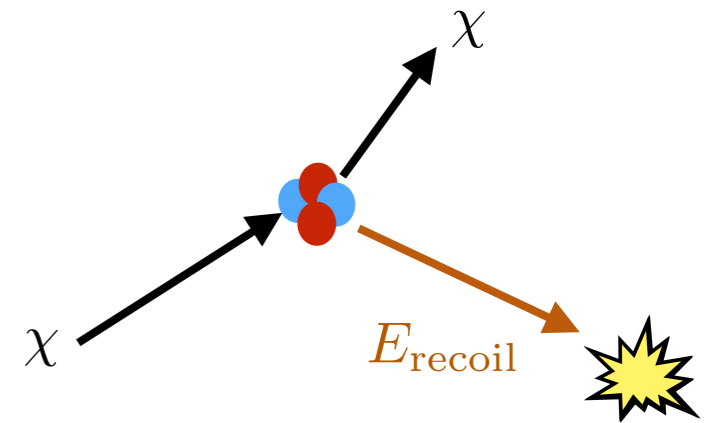
From TESSERACT white paper



Single phonon excitation



Multiphonons



Nuclear recoils

keV

MeV

GeV

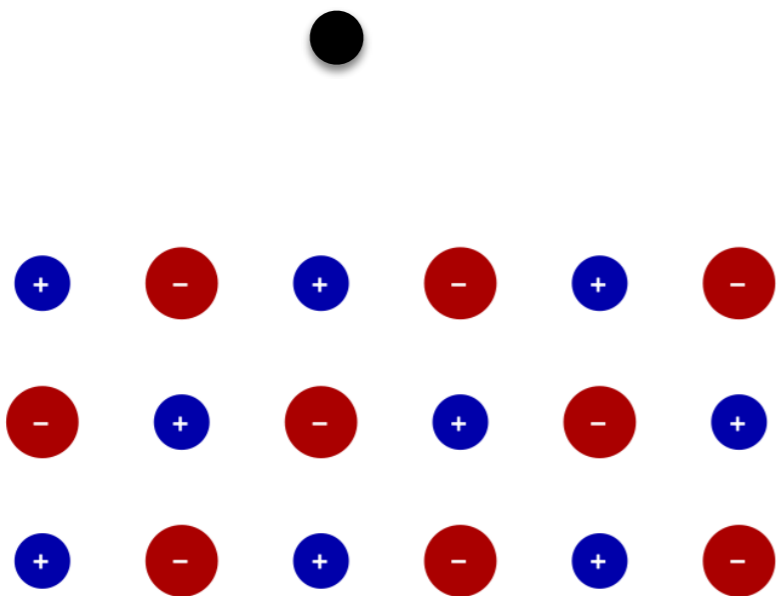
TeV

Dark matter mass

# DM-nucleus scattering in crystals: spherical guinea pig style



# DM-nucleus scattering in crystals: spherical guinea pig style



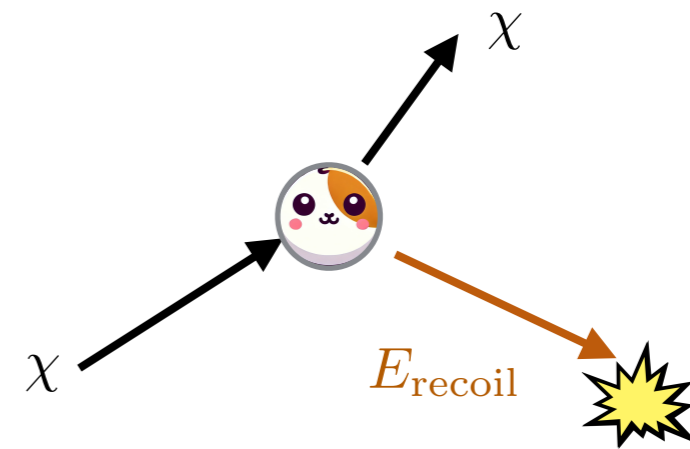
Single phonon excitation:  
scatter against the whole lattice

Momentum transfer

$$q \ll q_{\text{BZ}} = \frac{2\pi}{a} \sim \text{keV}$$

and energy deposition

$$\omega \sim \bar{\omega}_{\text{phonon}} < 0.2 \text{ meV}$$



Nuclear recoils: scatter  
against a single atom

Momentum transfer

$$q \gg \sqrt{2m\bar{\omega}_{\text{phonon}}} \sim 50 \text{ keV}$$

and energy deposition

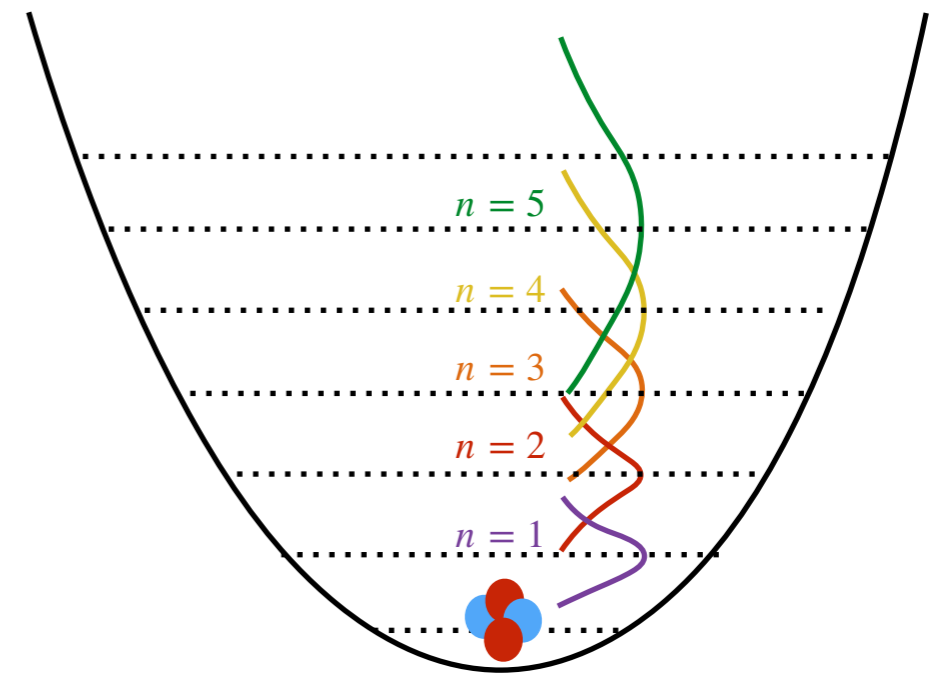
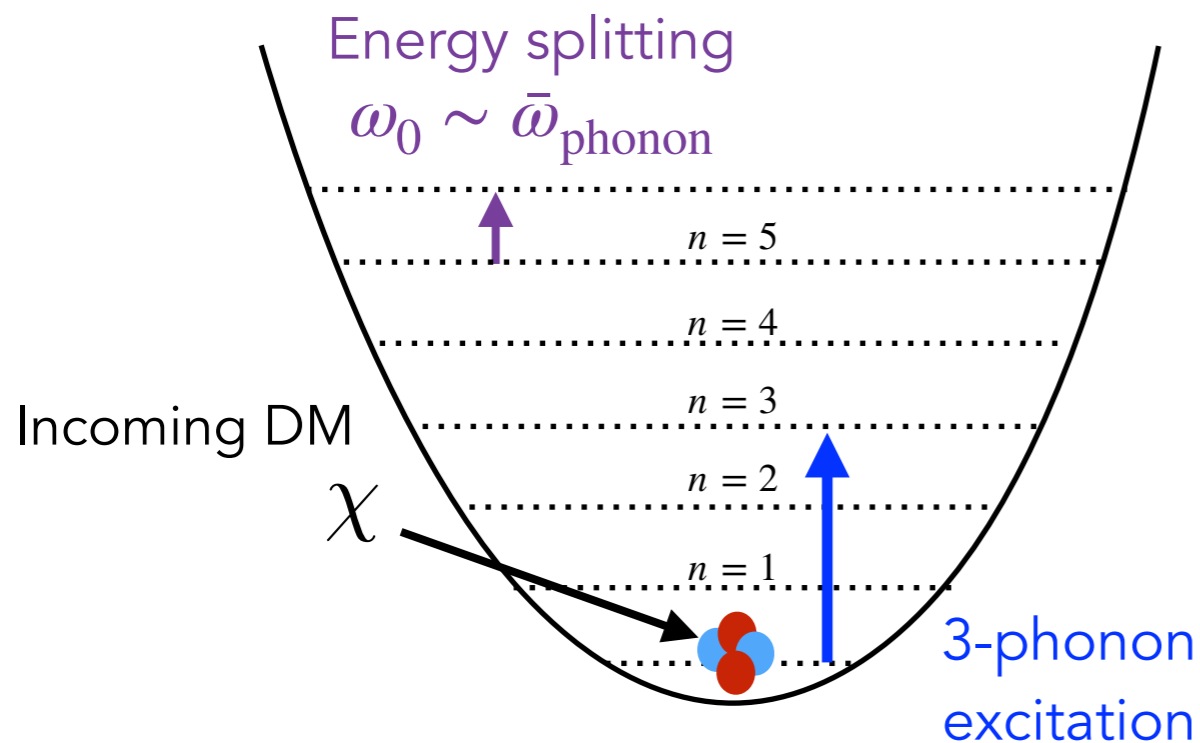
$$\omega \gg \bar{\omega}_{\text{phonon}}$$

# Multiphonon scattering



A first approximation: scattering off a nucleus in a bound harmonic oscillator potential

Phonons in a crystal exist in a continuum! An excellent second approximation:



Account for density of states  $D(\omega)$ :

$$\frac{\delta(\omega - \omega_0)}{\omega_0} \rightarrow \frac{D(\omega)}{\omega}$$

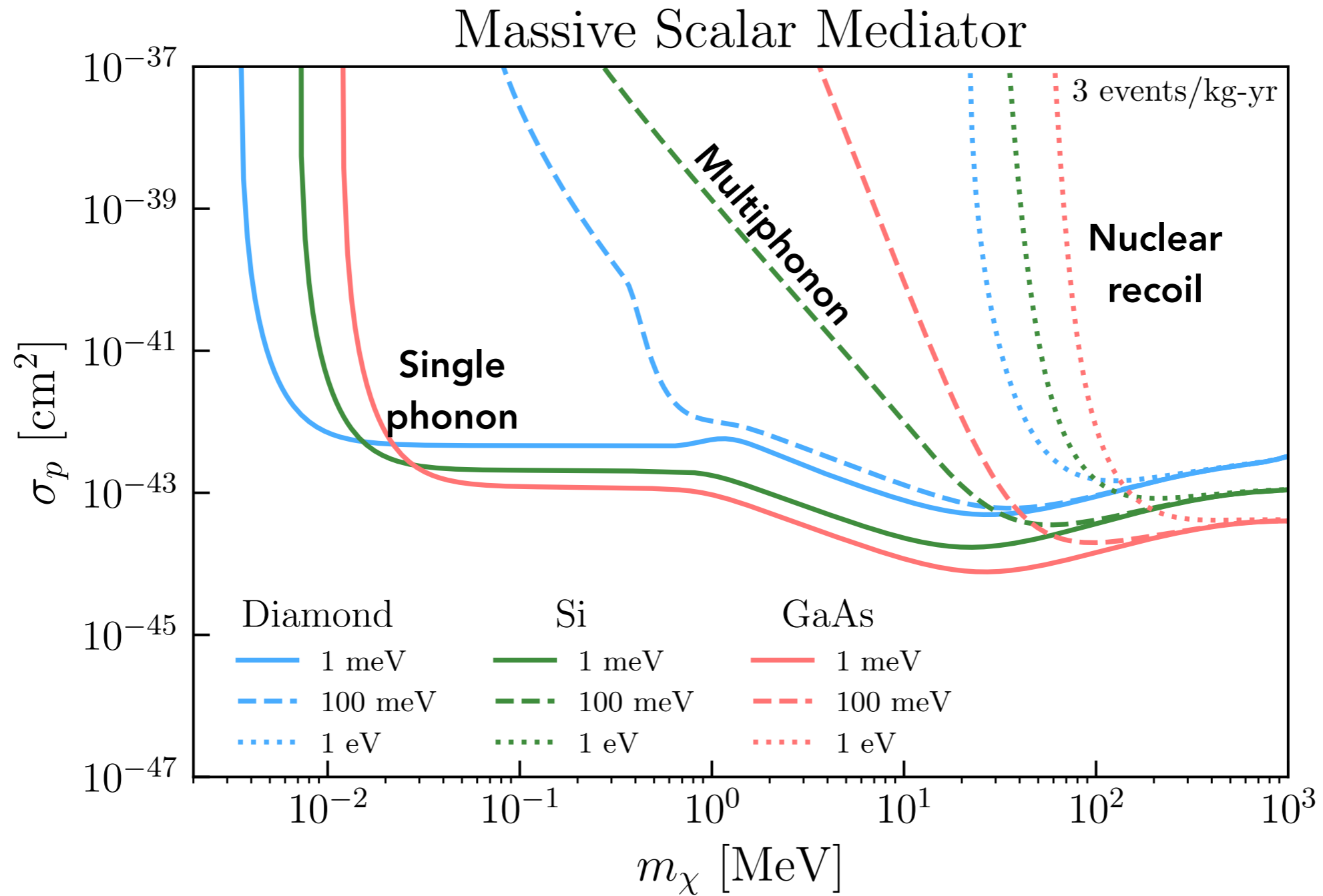
Knapen, Kozaczuk, TL 2011.09496

Kahn, Krnjaic, Mandava 2011.09477

Campbell-Deem, Knapen, TL, Villarama 2205.02250

TL, Shen, Sholapurkar, Villarama 2309.10839

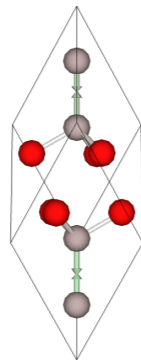
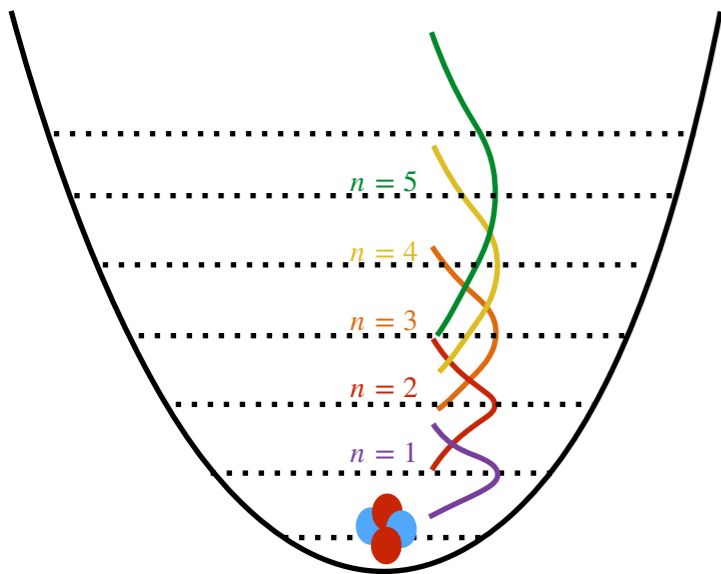
# A unified picture of low-energy nuclear recoils in crystals



# Multiphonon scattering



**Spherical Guinea Pig Approach**

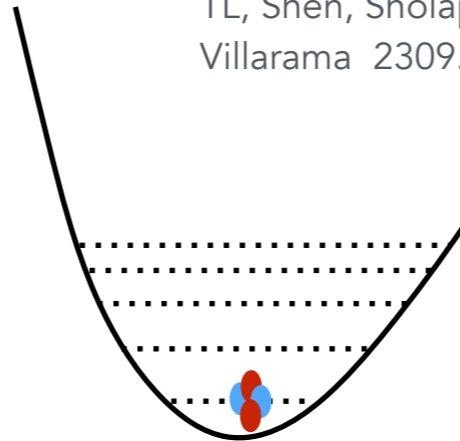


**Anisotropic crystals & daily modulation**

Stratman, TL (to appear)

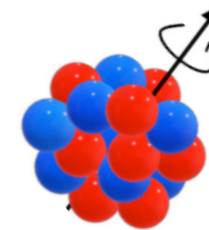
**Anharmonic effects**

TL, Shen, Sholapurkar,  
Villarama 2309.10839



**Spin-dependent scattering**

In progress with Knapen, Gori,  
Munbodh, Suter



**Incoherent approximation**

In progress with Fang, Lin



# DM scattering rate

$$\frac{d\sigma}{d^3\mathbf{q} d\omega} \propto \sigma_{\chi p} \overbrace{|\tilde{F}_{\text{med}}(q)|^2}^{\text{DM-mediator form factor}} \underbrace{S(\mathbf{q}, \omega)}_{\text{Dynamic structure factor}} \delta\left(\omega - \mathbf{q} \cdot \mathbf{v} + \frac{q^2}{2m_\chi}\right)$$

Dynamic structure factor  
captures response of target

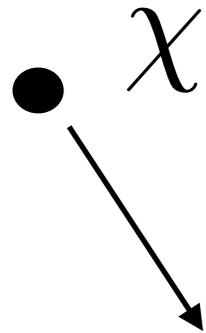
For free nuclei and spin-independent interactions:

$$S(\mathbf{q}, \omega) \propto A_N^2 \delta\left(\omega - \frac{q^2}{2m_N}\right)$$

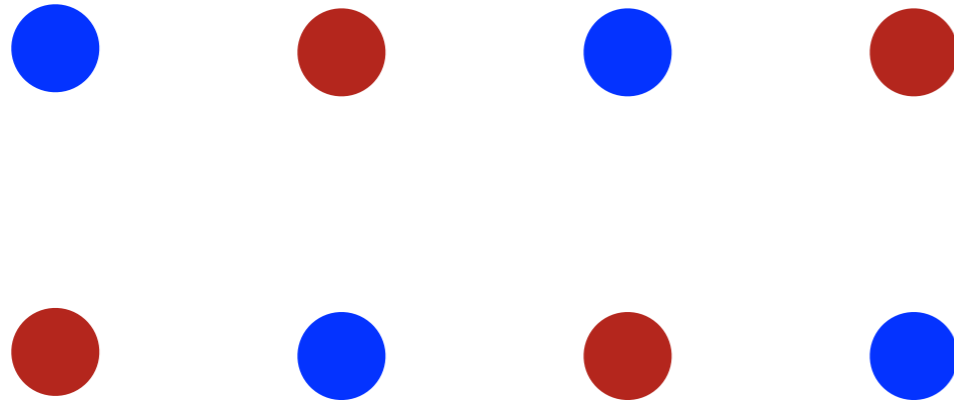
Goal: understand response function  $S(\mathbf{q}, \omega)$  from the single phonon to the nuclear recoil regime



# DM-nucleus interaction



$f_J$  - effective coupling strength between DM and ion  $J$



Short range SI interaction

$$\sigma_{\chi p} = 4\pi b_p^2$$

Scattering potential in Fourier space

$$V(\mathbf{q}) \propto b_p \sum_J f_J e^{i\mathbf{q}\cdot\mathbf{r}_J}$$

$$S(\mathbf{q}, \omega) \equiv \frac{2\pi}{V} \sum_f \left| \sum_J \langle \Phi_f | f_J e^{i\mathbf{q}\cdot\mathbf{r}_J} | 0 \rangle \right|^2 \delta(E_f - \omega)$$

$$= \frac{1}{V} \sum_{J, J'} f_J f_{J'}^* \int_{-\infty}^{\infty} dt \langle e^{-i\mathbf{q}\cdot\mathbf{r}_{J'}(0)} e^{i\mathbf{q}\cdot\mathbf{r}_J(t)} \rangle e^{-i\omega t}$$

Contains correlations between different atoms  $\rightarrow$  single phonon excitations

# Dynamic structure factor

Phonons appear through positions of ions:

$$\mathbf{r}_J(t) = \mathbf{r}_J^0 + \mathbf{u}_J(t)$$

↑  
Quantized phonon field given in terms of  
phonon dispersions  $\omega_{\mathbf{q}}$  and eigenvectors  $\mathbf{e}_{\mathbf{q}}$

Single phonon contribution has been studied extensively in literature,  
with  $\omega_{\mathbf{q}}$ ,  $\mathbf{e}_{\mathbf{q}}$  calculated from first principles approaches

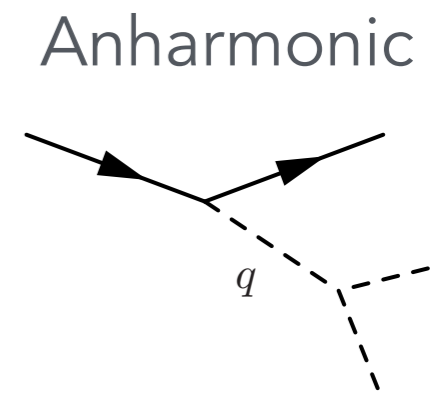
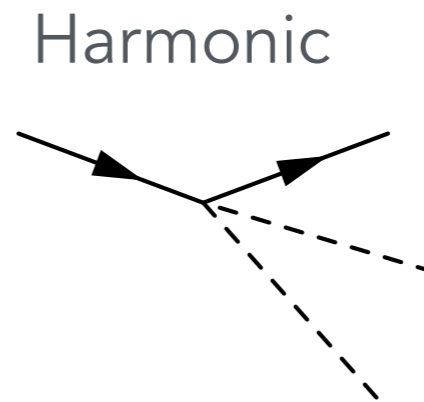
$$S^{n=1}(\mathbf{q}, \omega) \sim \sum_{J, J'} f_J f_{J'} \int dt \langle \mathbf{q} \cdot \mathbf{u}_J(0) \mathbf{q} \cdot \mathbf{u}_{J'}(t) \rangle e^{-i\omega t}$$

Griffin, Knapen, TL, Zurek 1807.10291; Griffin, Inzani, Trickle, Zhang, Zurek 1910.10716  
Griffin, Hochberg, Inzani, Kurinsky, TL, Yu 2020; Coskuner, Tickle, Zhang, Zurek 2102.09567

# Dynamic structure factor

Expansion in  $q^2/(M_N\omega)$  (and anharmonic interactions):

$$S(\mathbf{q}, \omega) = \begin{aligned} & \text{(0-phonon)} \\ & + \text{(1-phonon)} \\ & + \text{(2-phonon)} + \dots \end{aligned}$$



Quickly becomes more complicated to evaluate for more than 1 phonon

Analytic approximations are valuable to make progress

# The spherical guinea pig

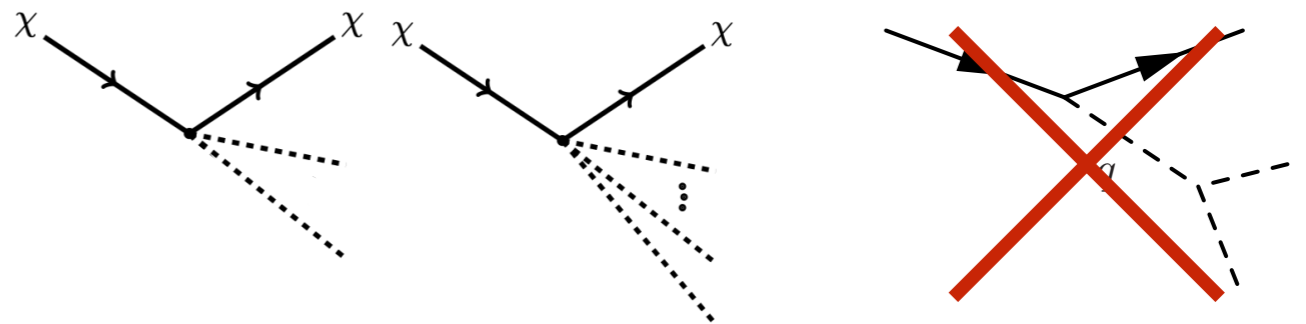
Neglect interference terms in structure factor:

Incoherent approximation

$$S(\mathbf{q}, \omega) \approx \frac{1}{V} \sum_J^N (f_J)^2 \int_{-\infty}^{\infty} dt \langle e^{-i\mathbf{q} \cdot \mathbf{u}_J(0)} e^{i\mathbf{q} \cdot \mathbf{u}_J(t)} \rangle e^{-i\omega t}$$

Good approx for  $q \gg q_{BZ} \sim \text{keV}$

Harmonic crystal



Isotropic crystal

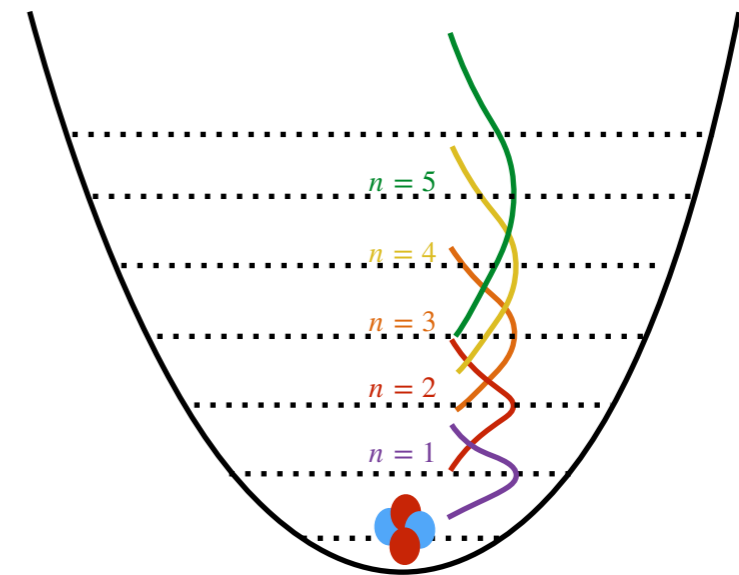
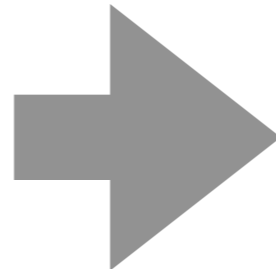
Phonon density of states  $D(\omega)$  does not depend on direction of phonon mode

# The spherical guinea pig

Incoherent approximation

Harmonic crystal

Isotropic crystal



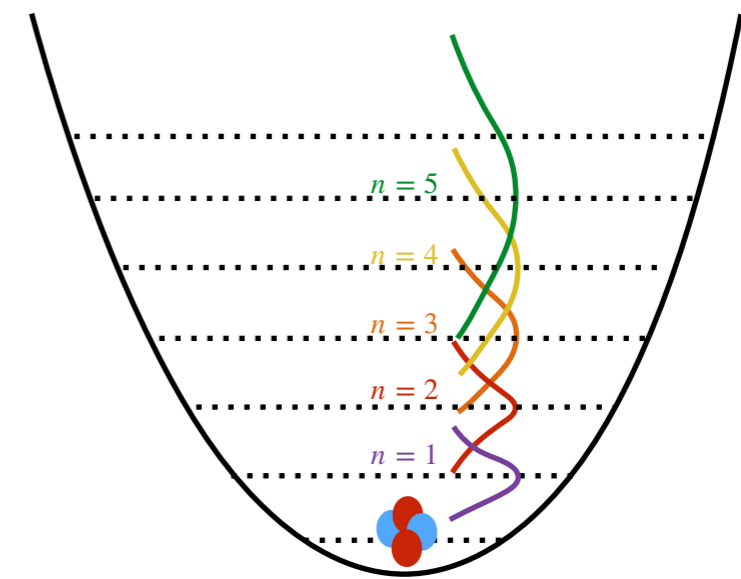
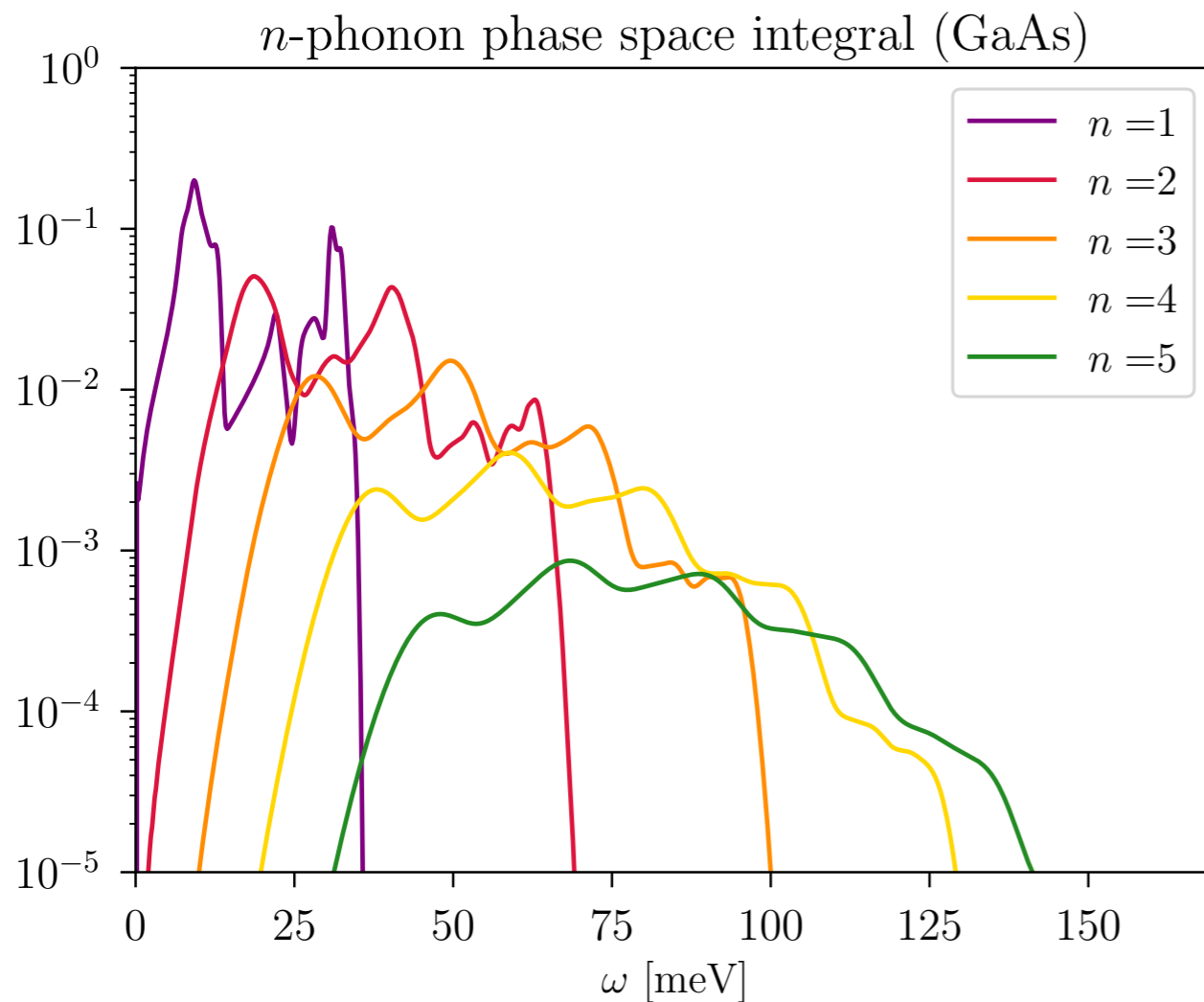
Integrate over n-phonon phase space:

$$\sum_n \frac{1}{n!} \left( \frac{q^2}{2m_N} \right)^n \left( \prod_{i=1}^n \int d\omega_i \frac{D(\omega_i)}{\omega_i} \right) \delta \left( \sum_j \omega_j - \omega \right)$$

Public code: **DarkELF**

<https://github.com/tongylin/DarkELF>

# The spherical guinea pig



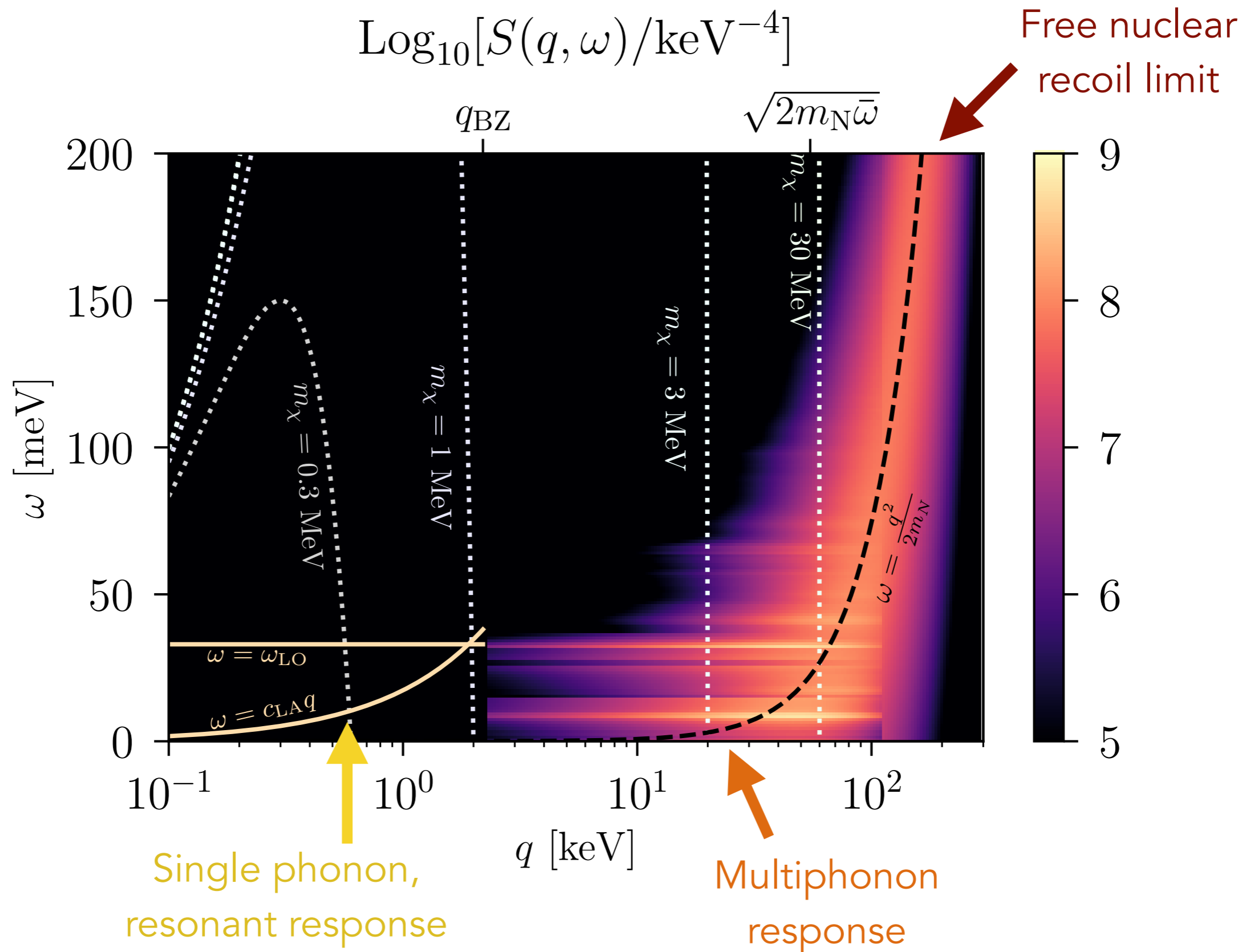
Integrate over *n*-phonon phase space:

$$\sum_n \frac{1}{n!} \left( \frac{q^2}{2m_N} \right)^n \left( \prod_{i=1}^n \int d\omega_i \frac{D(\omega_i)}{\omega_i} \right) \delta \left( \sum_j \omega_j - \omega \right)$$

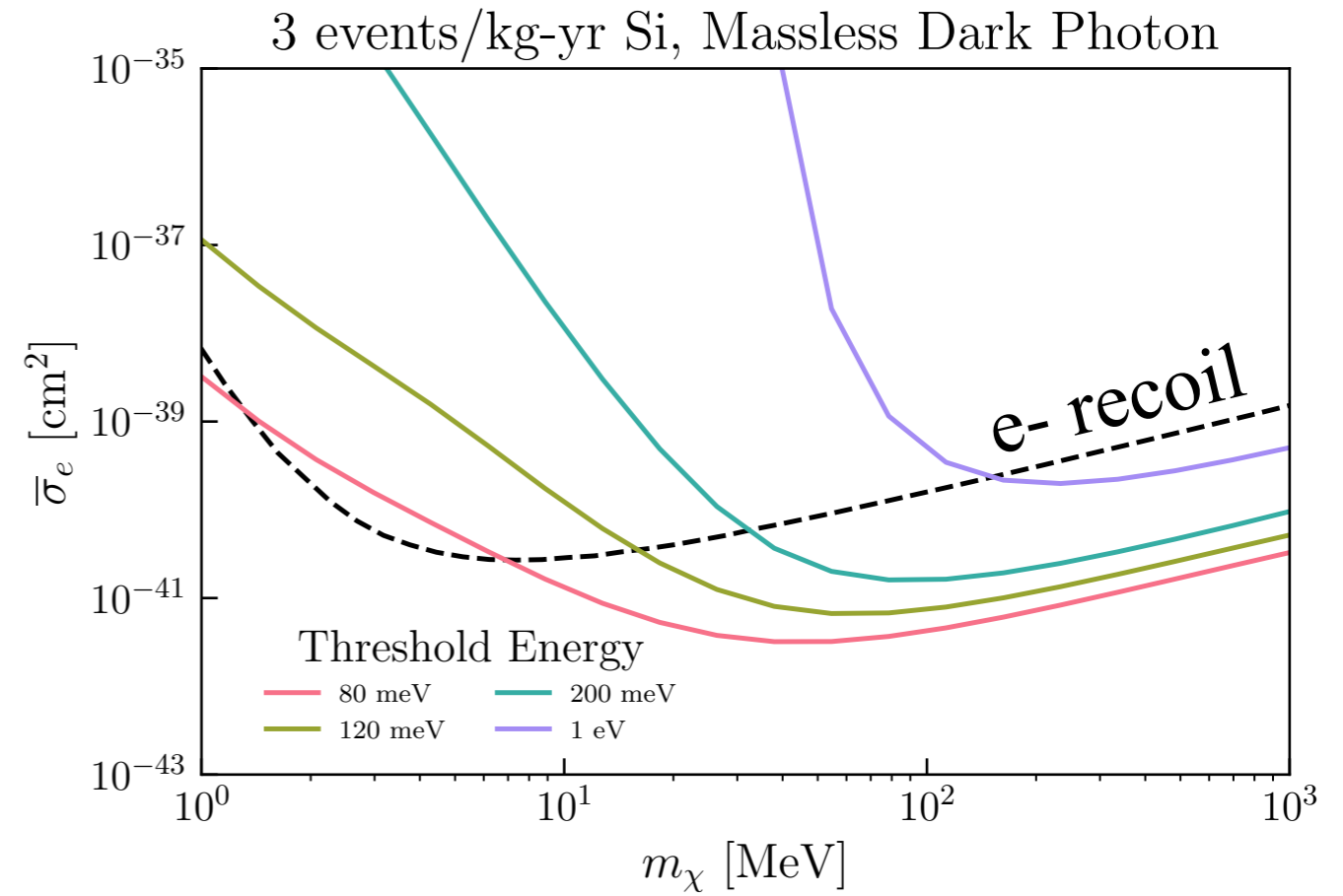
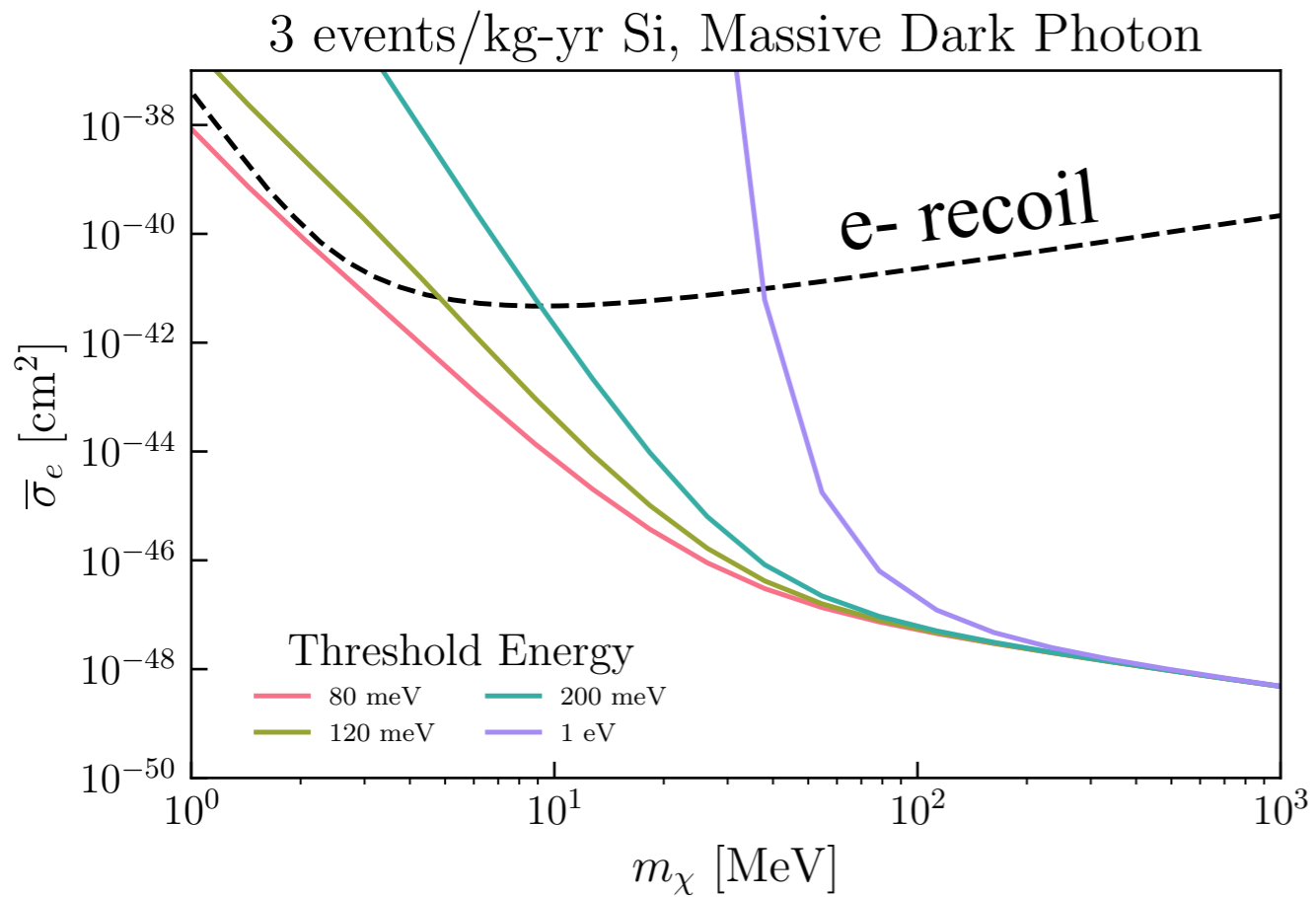
Public code: **DarkELF**

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# Dynamic structure factor

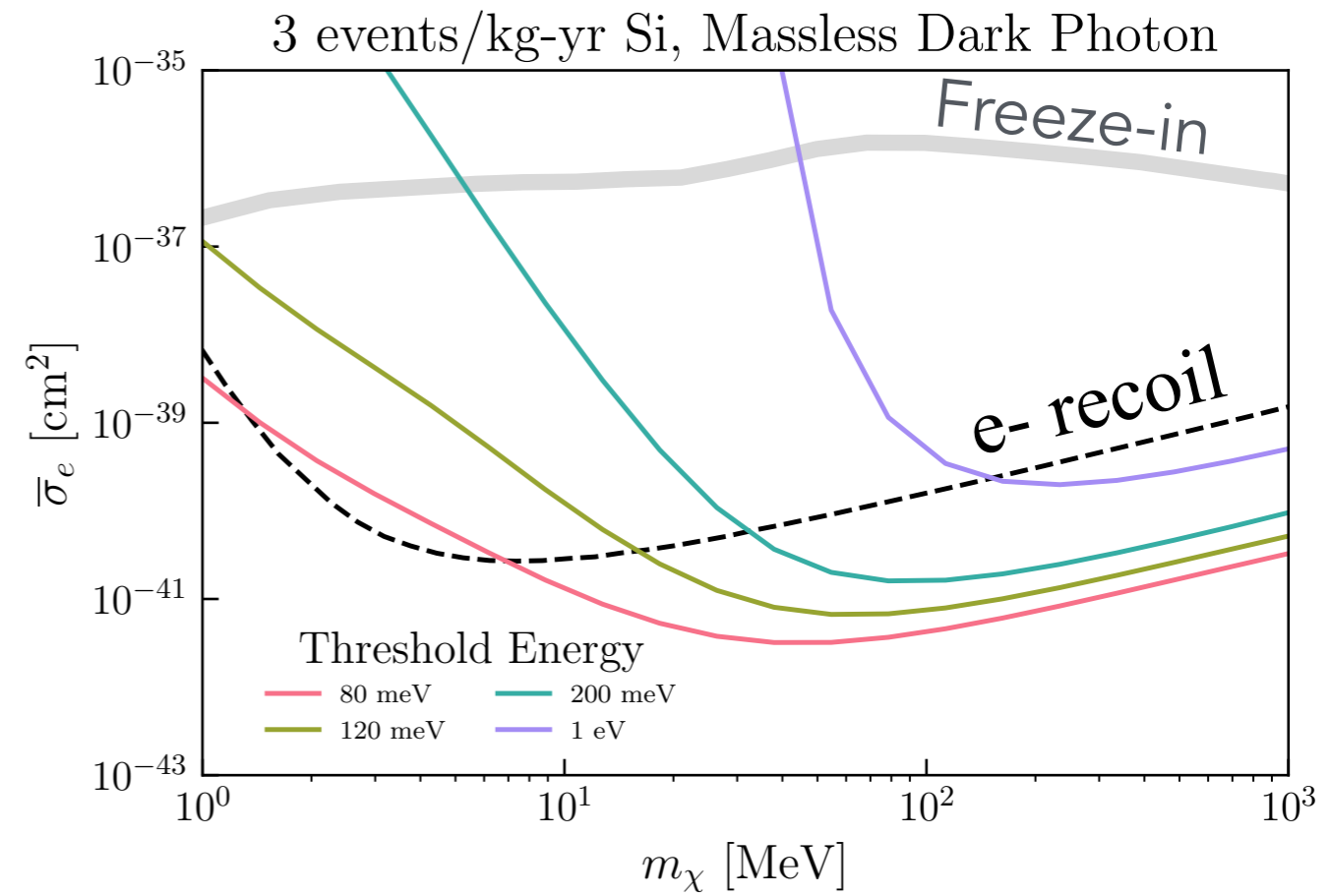
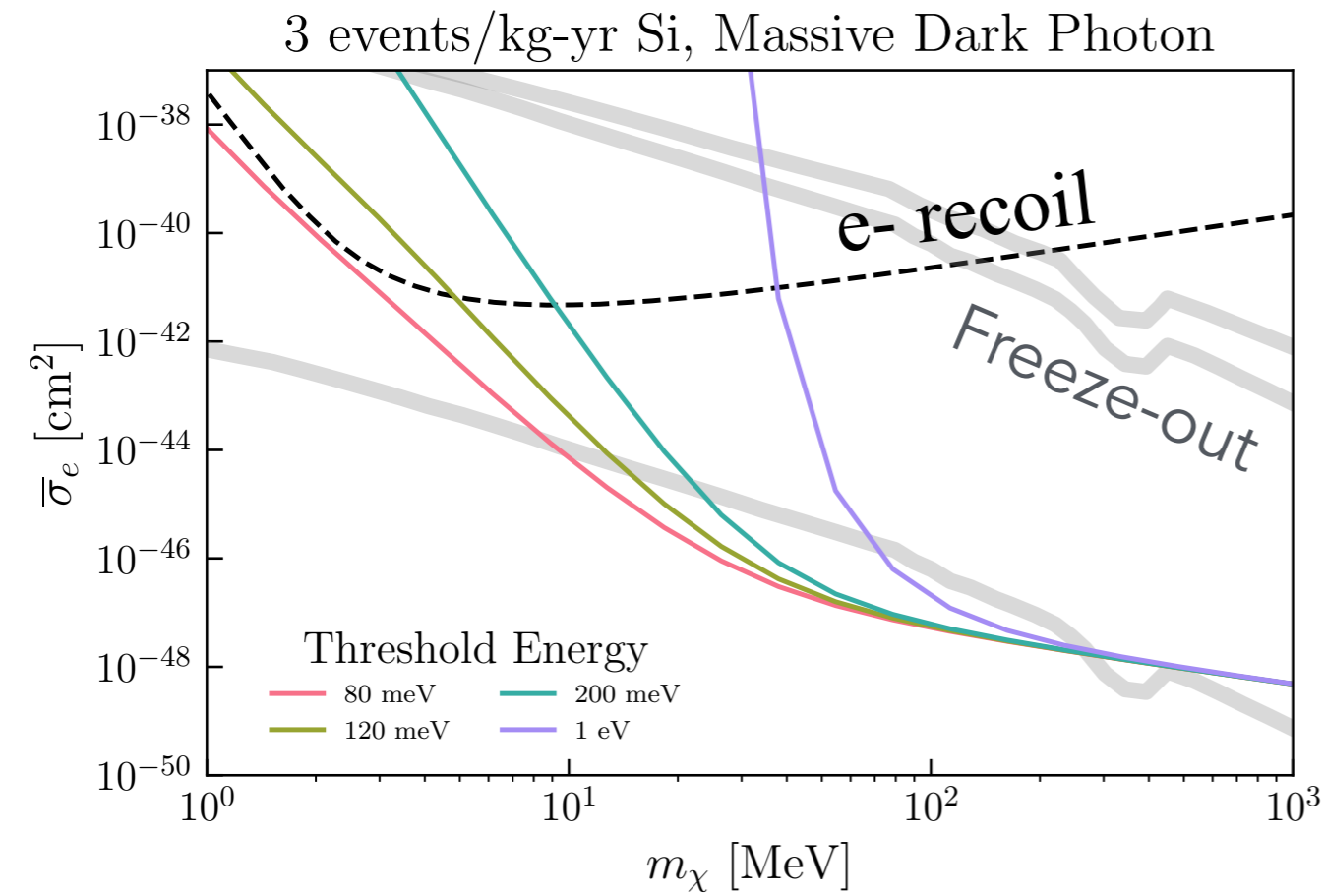


# Dark photon mediator





# Dark photon mediator





# Anharmonic crystals

$$V_{\text{lattice}} = \frac{1}{2} \sum_{\ell, d, \ell', d'} \sum_{\alpha, \beta} k_{\alpha\beta}^{(2)}(\ell d, \ell' d') u_{\alpha}(\ell d) u_{\beta}(\ell' d')$$

Harmonic  
(2nd order  
force constants)

$$+ \frac{1}{3!} \sum_{\ell, d, \ell', d', \ell'', d''} \sum_{\alpha, \beta, \gamma} k_{\alpha\beta\gamma}^{(3)}(\ell d, \ell' d', \ell'' d'') \\ \times u_{\alpha}(\ell d) u_{\beta}(\ell' d') u_{\gamma}(\ell'' d'')$$

3rd order  
force constants

+ ...

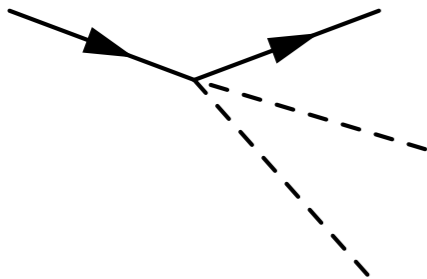
Anharmonicity related to phonon decay, thermal conductivity and expansion observed in crystals.

# Anharmonic crystals

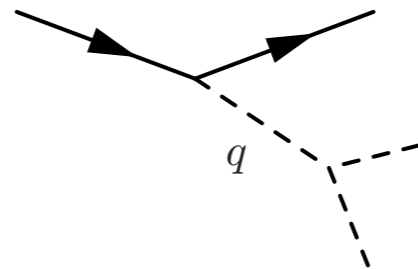
Momentum transfer

$$q \ll q_{\text{BZ}} = \frac{2\pi}{a} \sim \text{keV}$$

Harmonic



Anharmonic



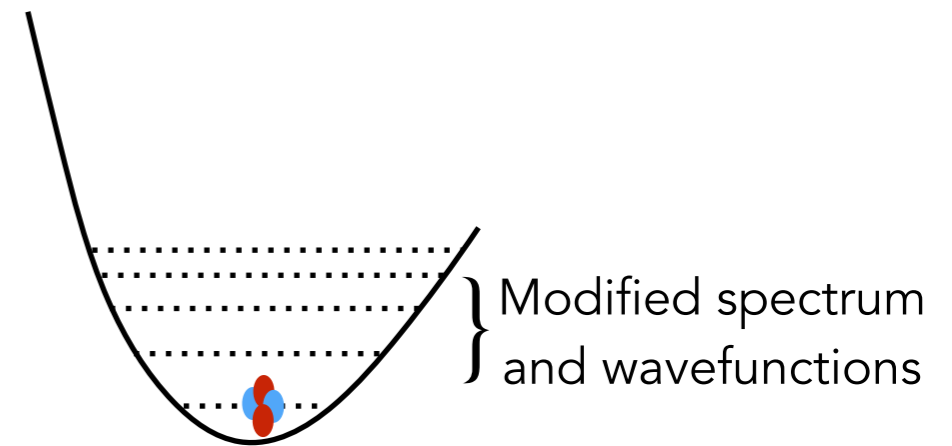
Use EFT of elastic waves

$$\delta H_{anh} \sim \int d^3\mathbf{r} \alpha (\nabla_i \mathbf{u}_j)^3$$

to calculate for acoustic phonons

Momentum transfer

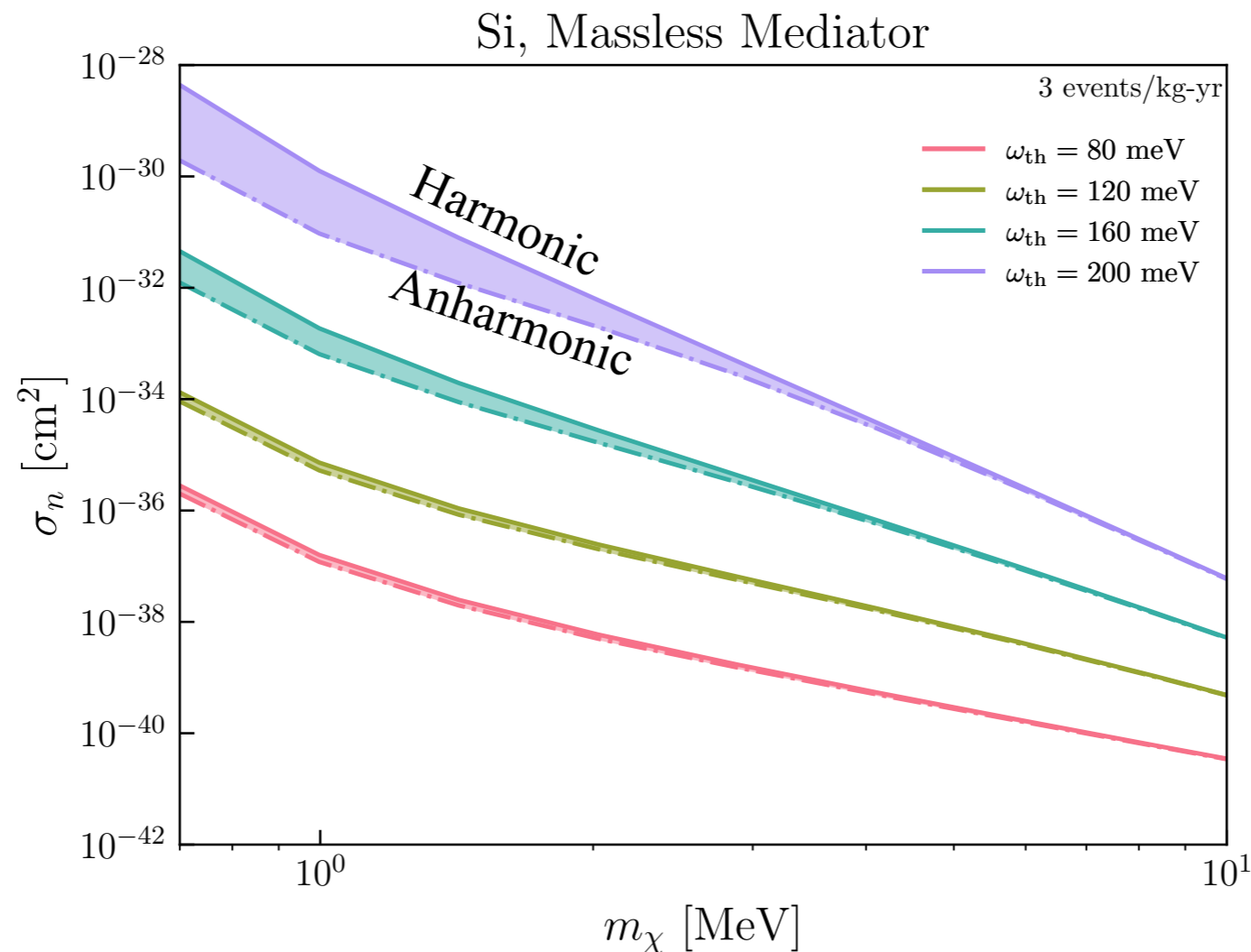
$$q \gtrsim q_{\text{BZ}} = \frac{2\pi}{a} \sim \text{keV}$$



Toy model: scattering off a particle in a 1D anharmonic potential  $V(x)$

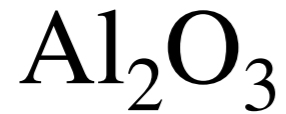


# Anharmonic interactions

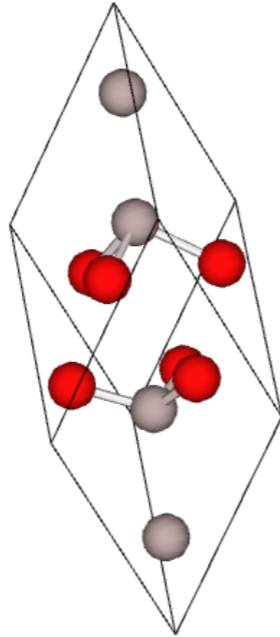


Takeaway: anharmonic effects are **limited** to DM masses  $\lesssim$  MeV and high cross sections (typically not viable in realistic models).

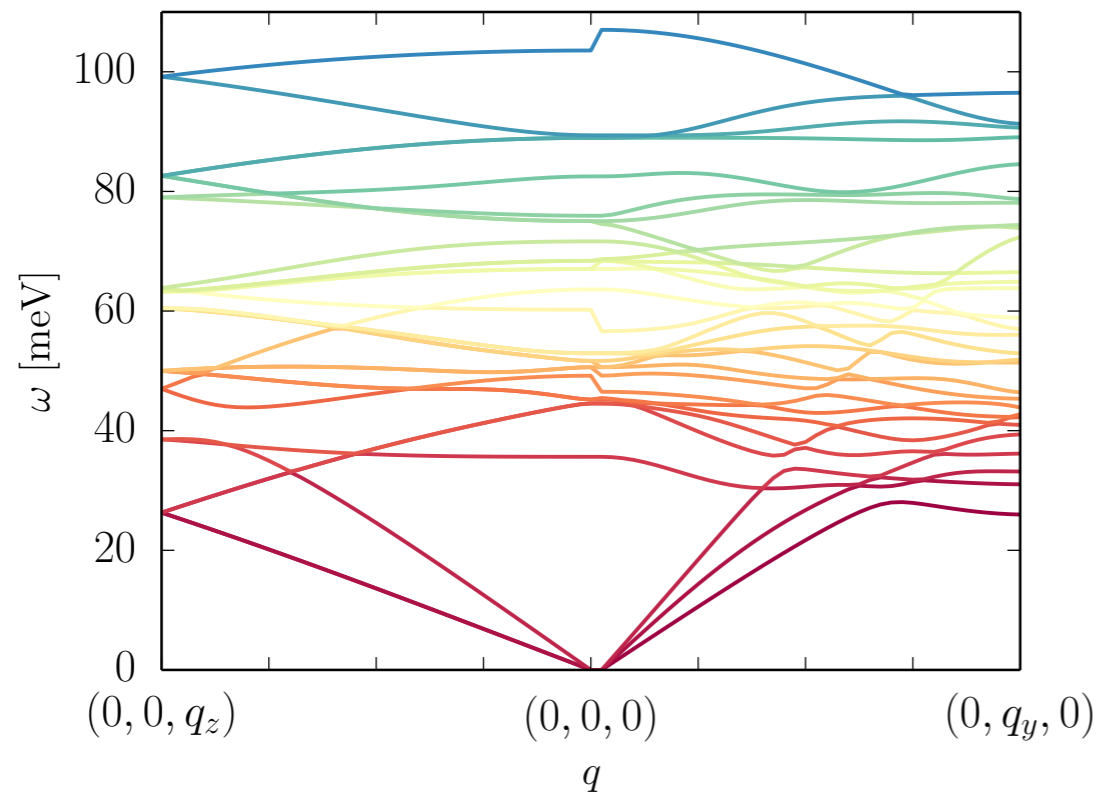
# Anisotropic crystals



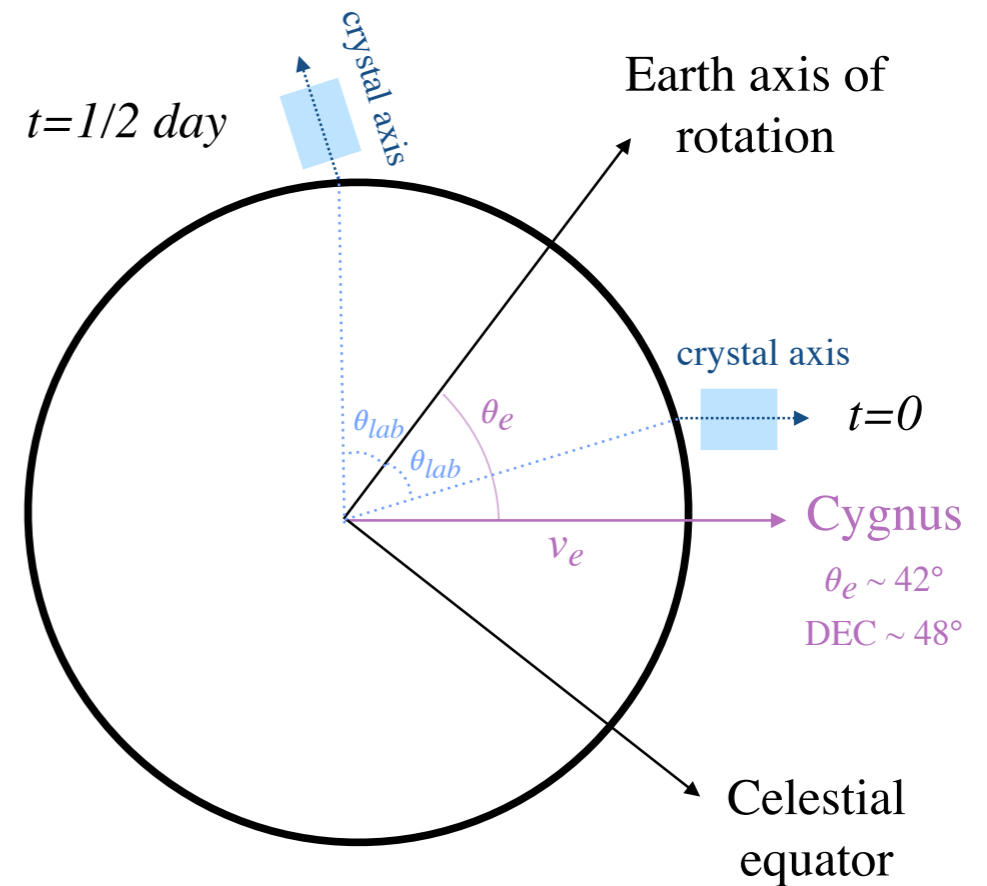
Sapphire; to be used in SPICE



Band structure

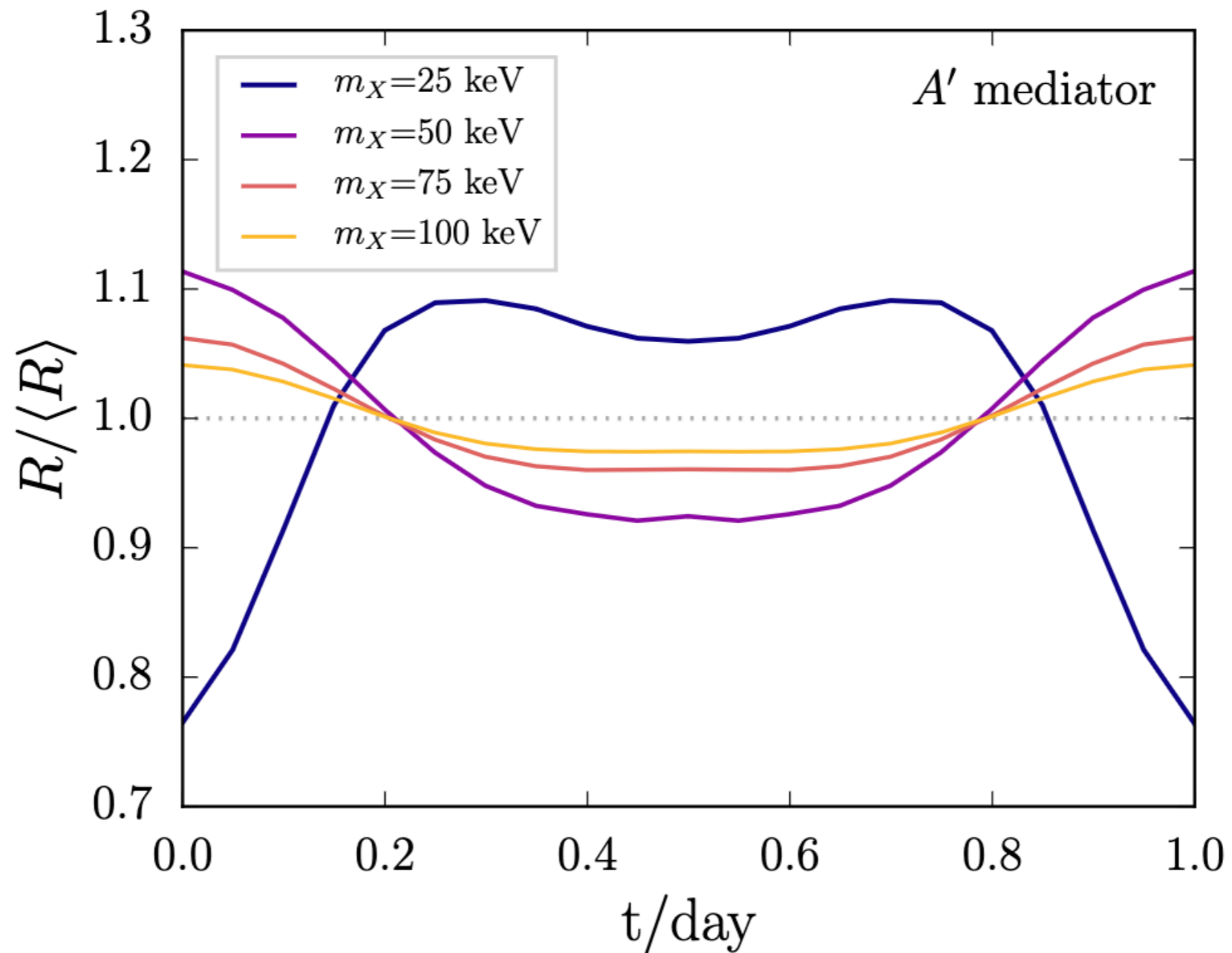


Phonon modes (DoS) depend on direction of momentum



Daily modulation in rate as Earth rotates

# Daily modulation with **single** phonons

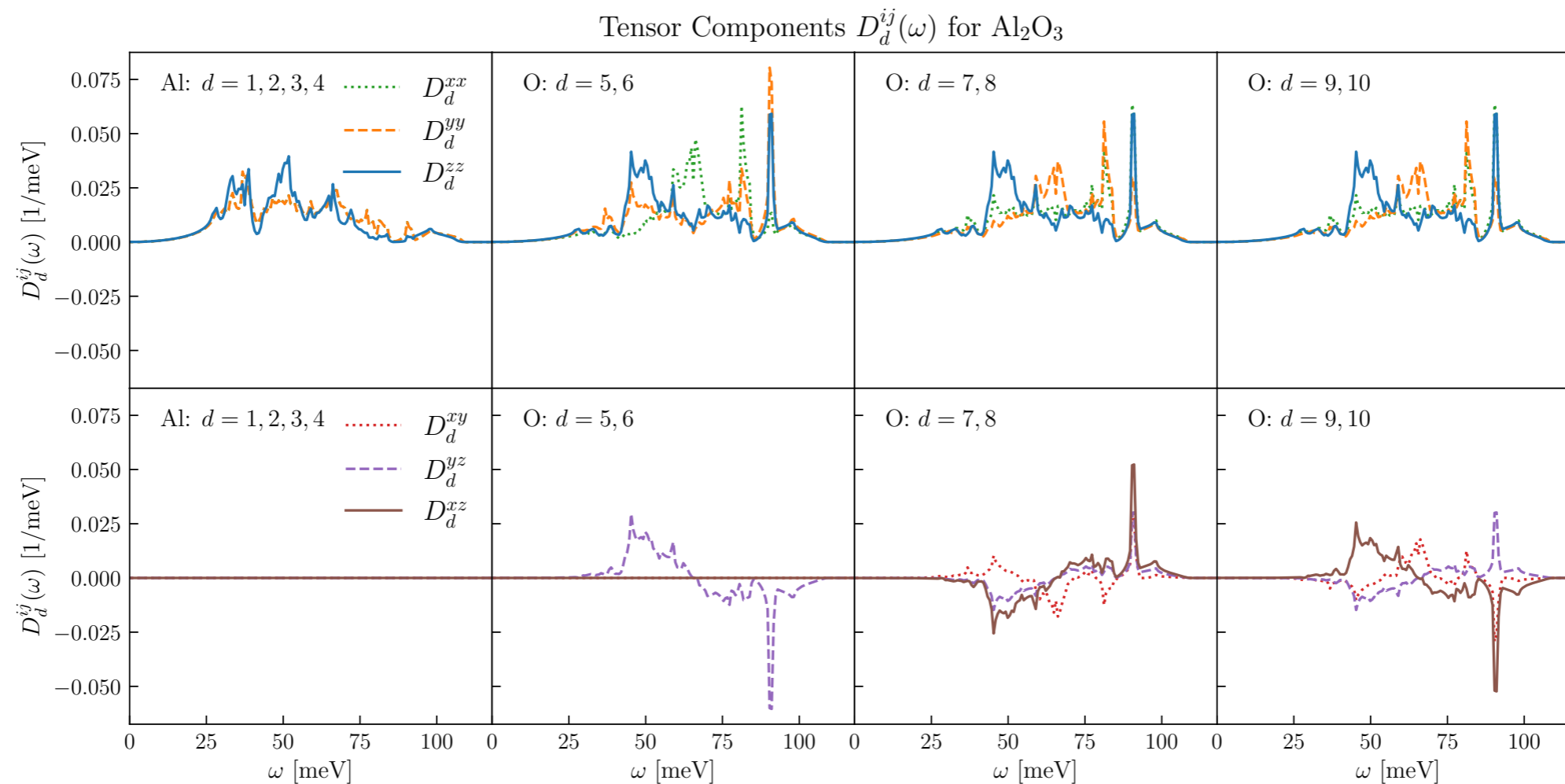


# Directional detection with multi phonons



With Connor Stratman, to appear

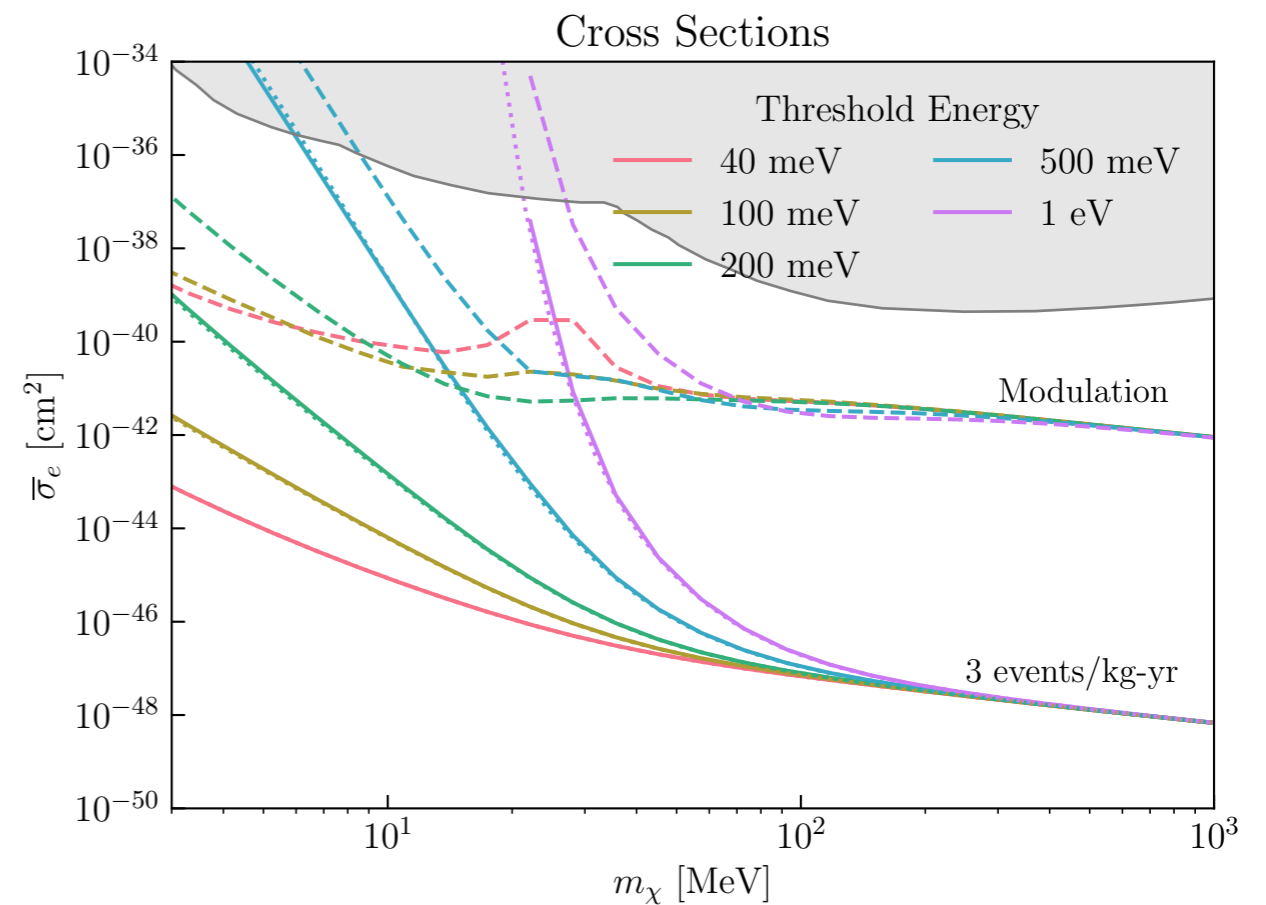
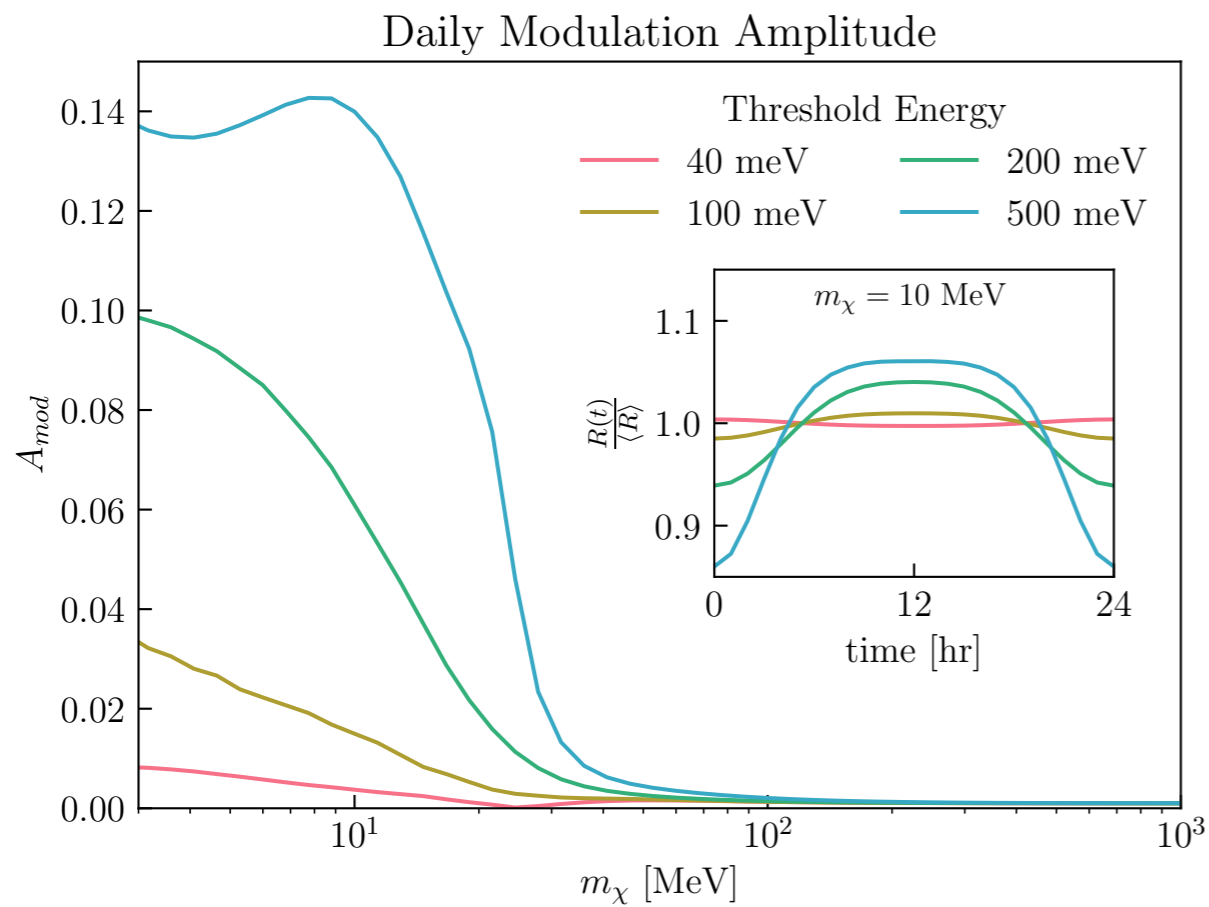
Density of states is now a tensor:  $D(\omega) \rightarrow D^{ij}(\omega)$   
 and scattering now depends on  $\hat{q}_i D^{ij}(\omega) \hat{q}_j$





# Daily modulation for multi phonons

Scattering in  $\text{Al}_2\text{O}_3$  via massive dark photon mediator



Modulation calculation also to be made public in DarkELF  
(<https://github.com/tongylin/DarkELF>) when paper submitted.

# Spin-dependent scattering



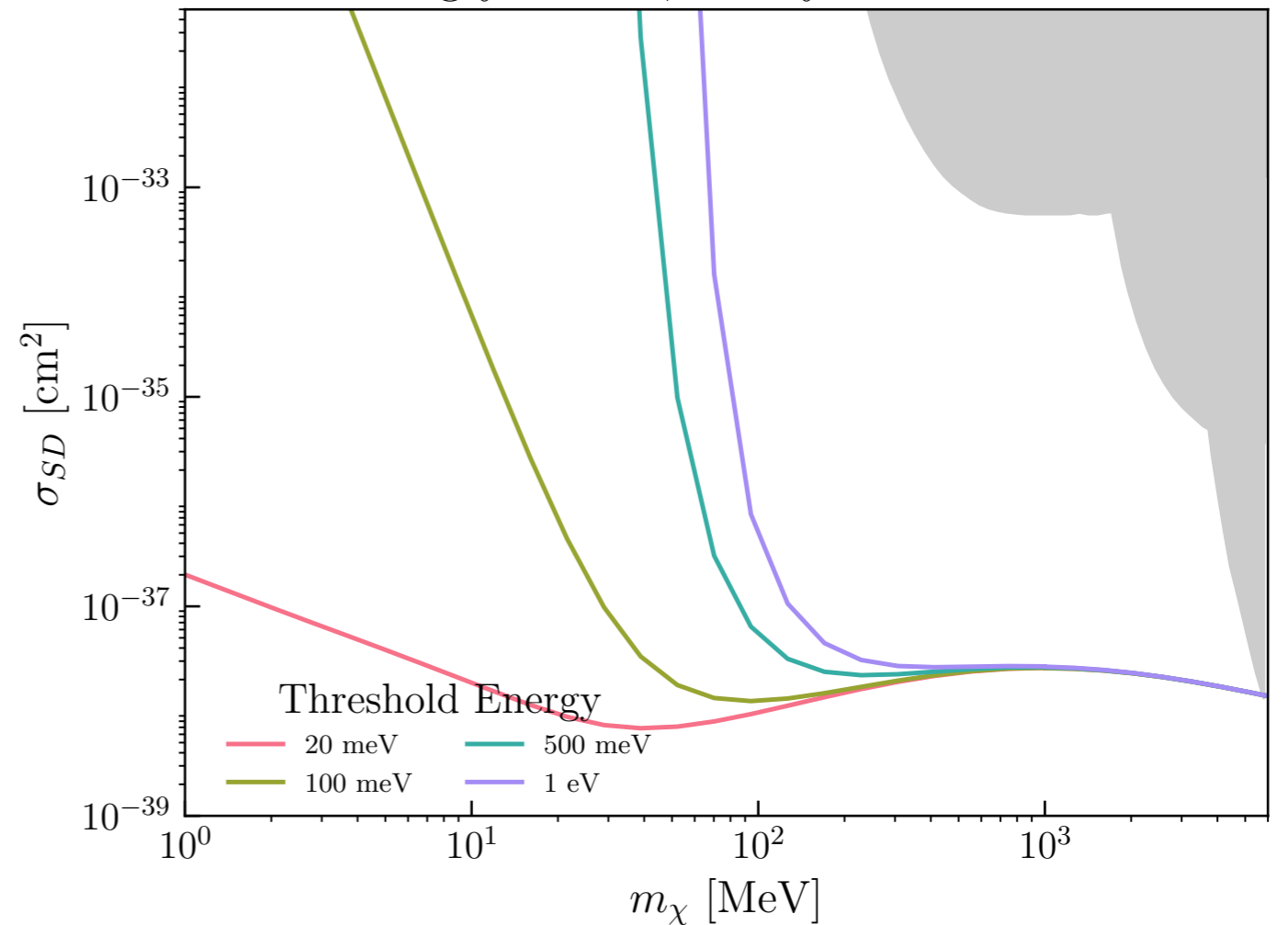
With Pankaj Munbodh, Bethany Suter,  
Stefania Gori, Simon Knapen (in progress)

No long-range spin order in (typical)  
direct detection materials →  
spherical guinea pig approximation.

Also soon to appear in DarkELF!

SD scattering off protons

kg-yr GaAs, Heavy Mediator

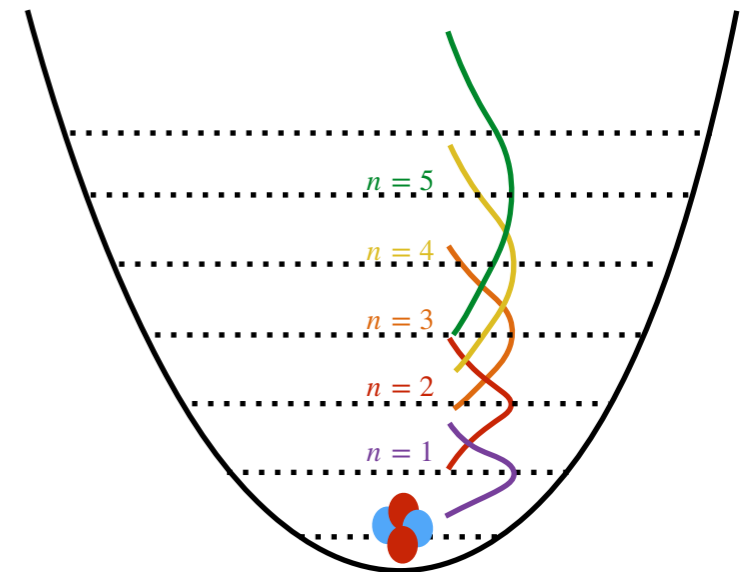
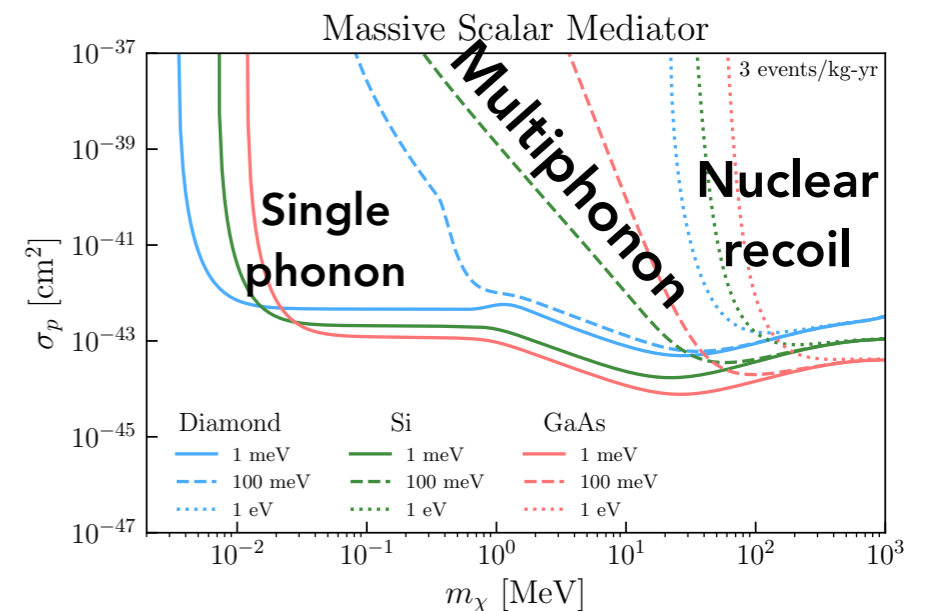


# Key takeaways

Multiphonons are the next step in the road towards single-phonon & MeV DM detection.

Full DFT multiphonon calculations are extremely daunting — we instead rely on a variety of approximations.

Theory progress: calculate anharmonic effects, anisotropies (daily modulation), spin-dependent scattering. Guinea pigs are actually pretty spherical.





**Thanks!**