



Centre Canadien de Recherche en
Physique des Astroparticules
Arthur B. McDonald
Canadian Astroparticle Physics Research Institute

Dark matter vs high-energy neutrinos

Aaron Vincent

Dark Interactions | SFU | Vancouver | October 17 2024



Featuring



Carlos Argüelles
Harvard



Ali Kheirandish
UNV Las Vegas



Henry White
Queen's

Qinrui Liu
Queen's



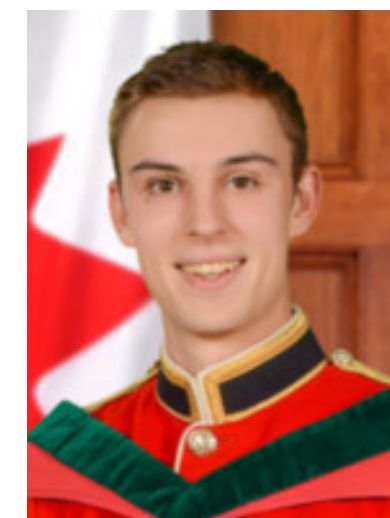
Ibrahim Safa
UW Madison



Andrés Olivares-del-Campo
IPPP Durham



Ningqiang Song
CN Yang ITP



Adam
McMullen
Queen's



Alejandro Diaz
MIT

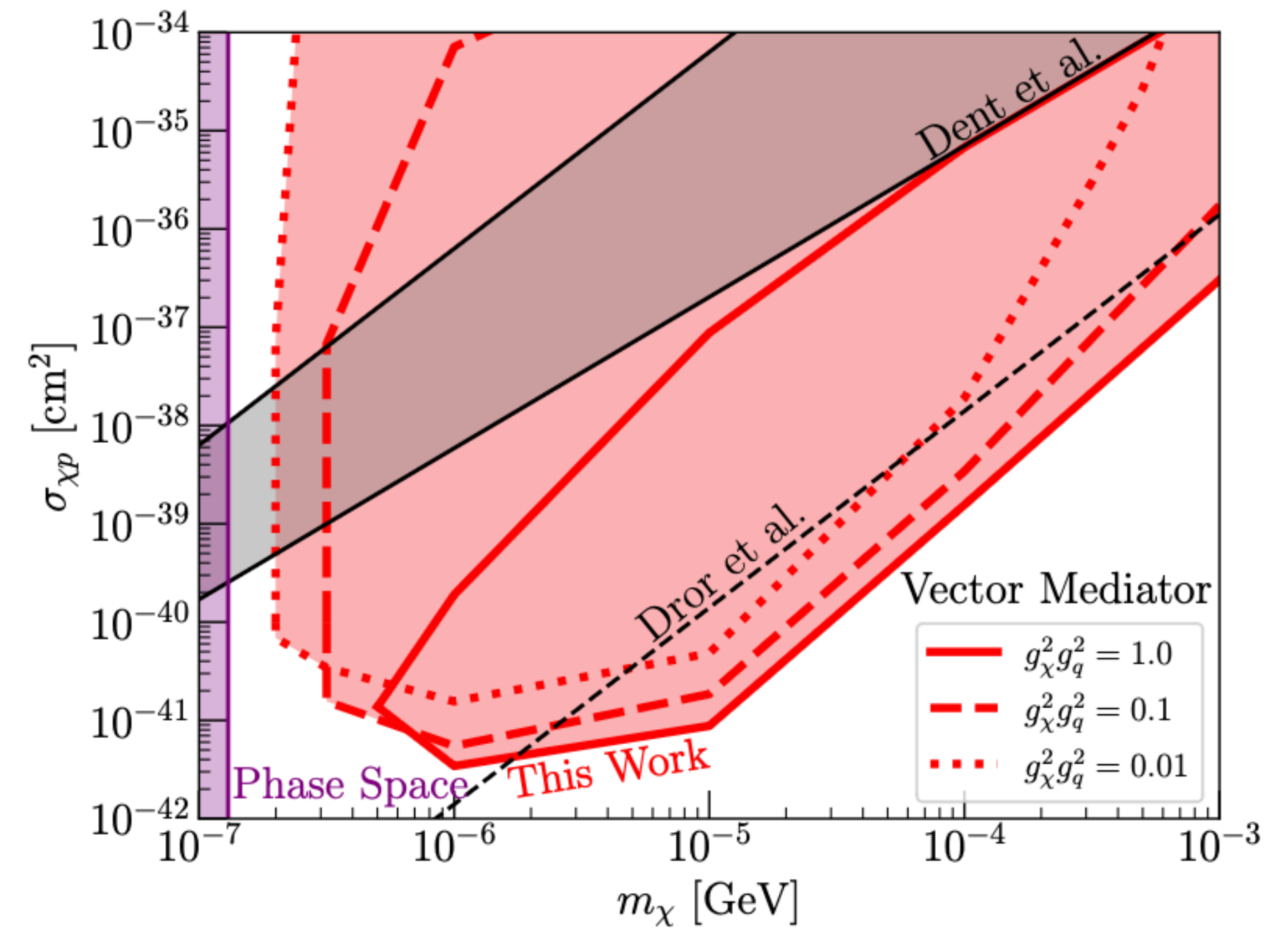
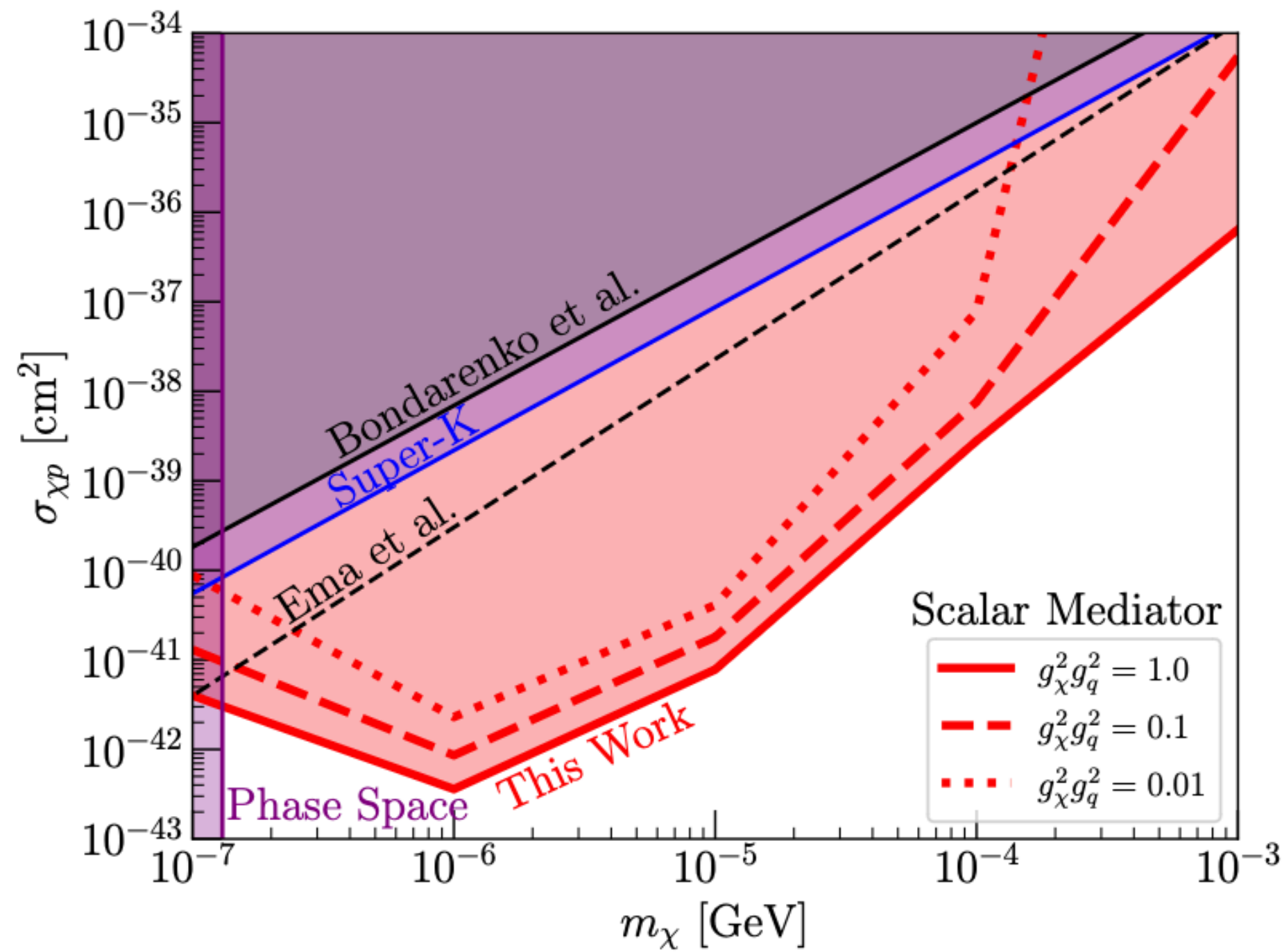


Diyaselis
Delgado
Harvard

I won't talk about

Detecting dark matter at neutrino telescopes

- Cosmic Ray-Boosted Dark Matter at IceCube

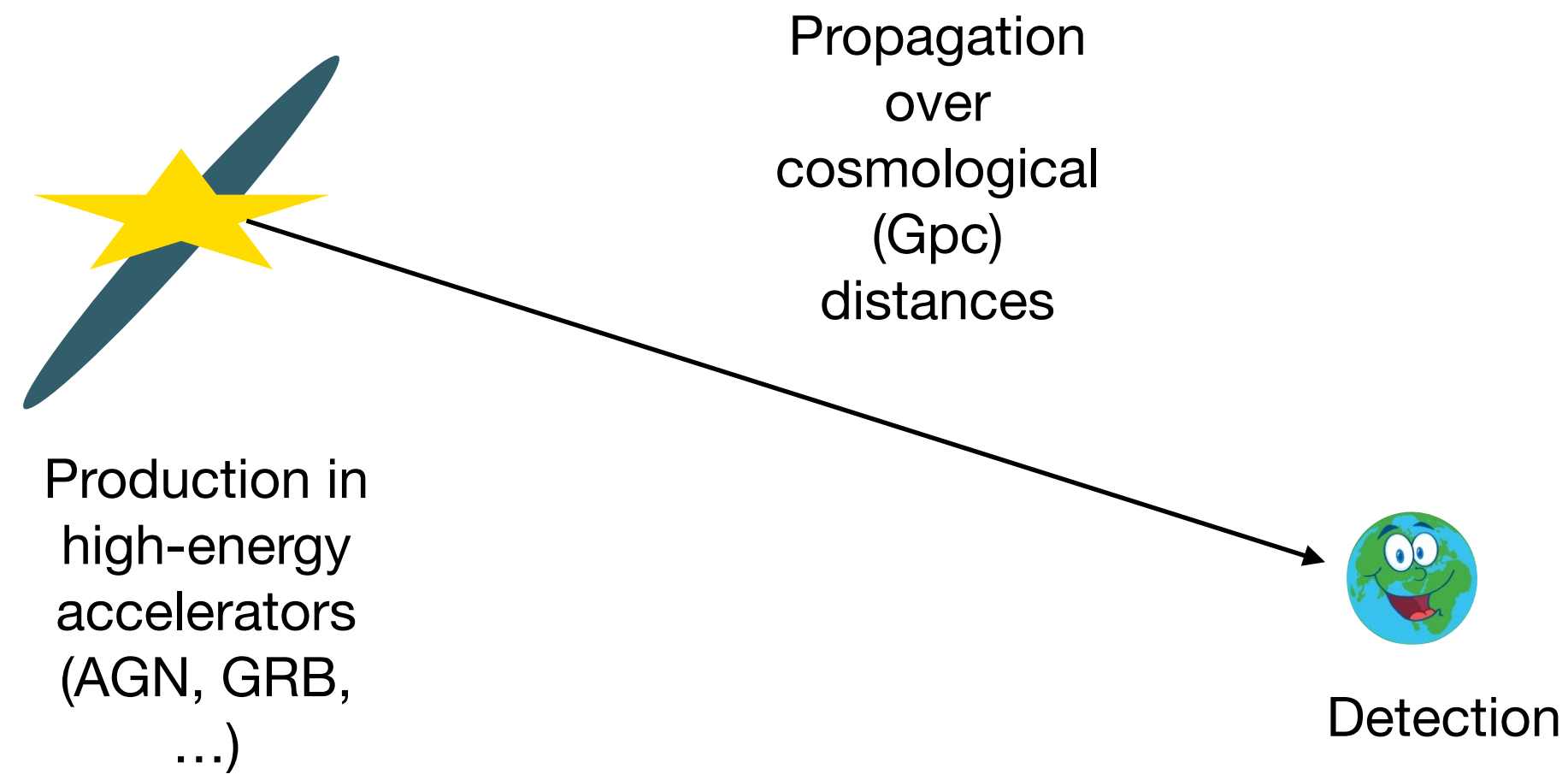


I will talk about

Dark matter-neutrino interactions: how do you see the most invisible particle?



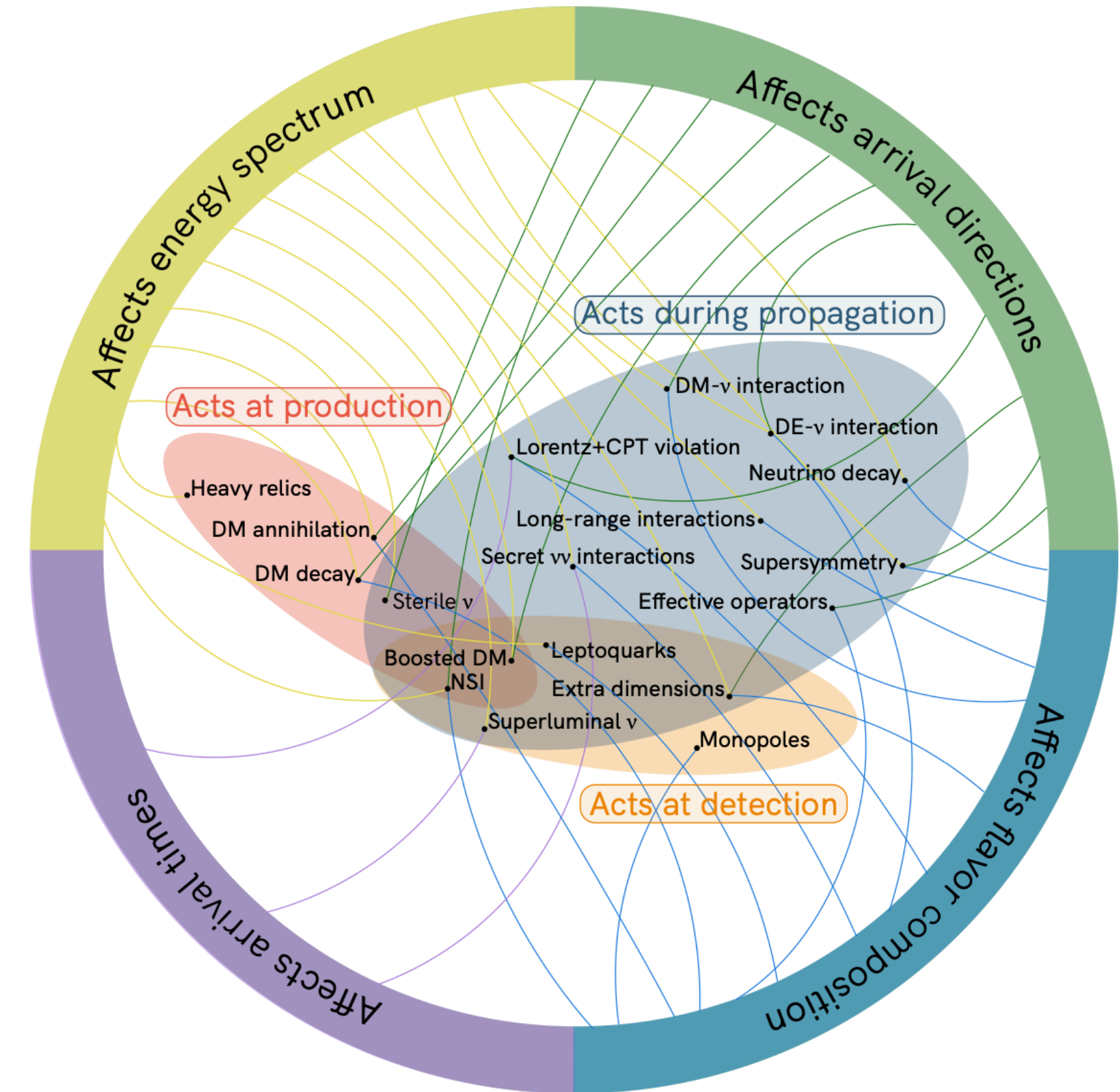
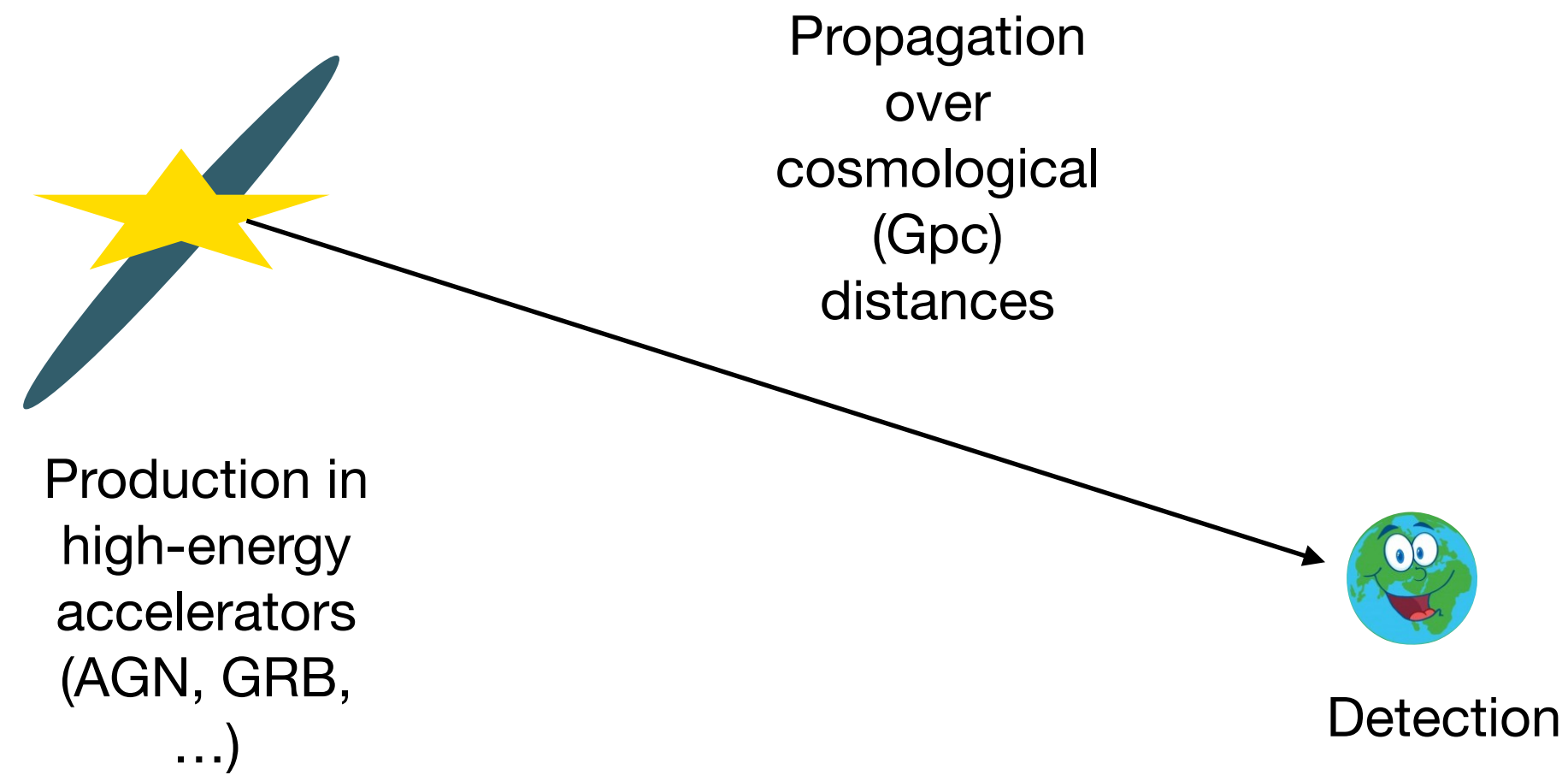
High-energy neutrinos



Neutrinos can tell us about “standard model” physics:

- Nature of these accelerators
- Oscillation, interaction with intergalactic medium
- Detection: high-energy neutrino-nucleus cross sections

High-energy neutrinos



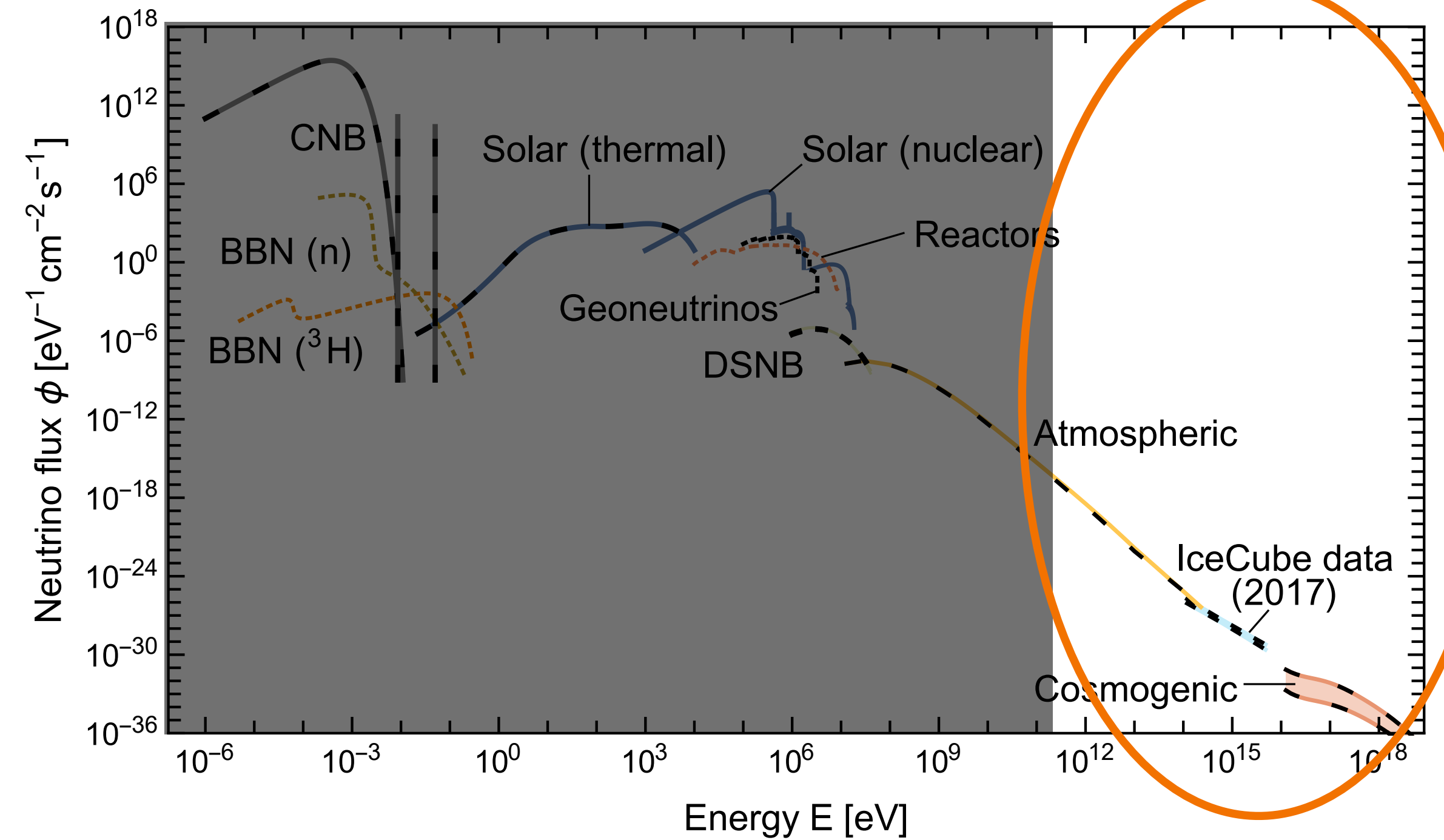
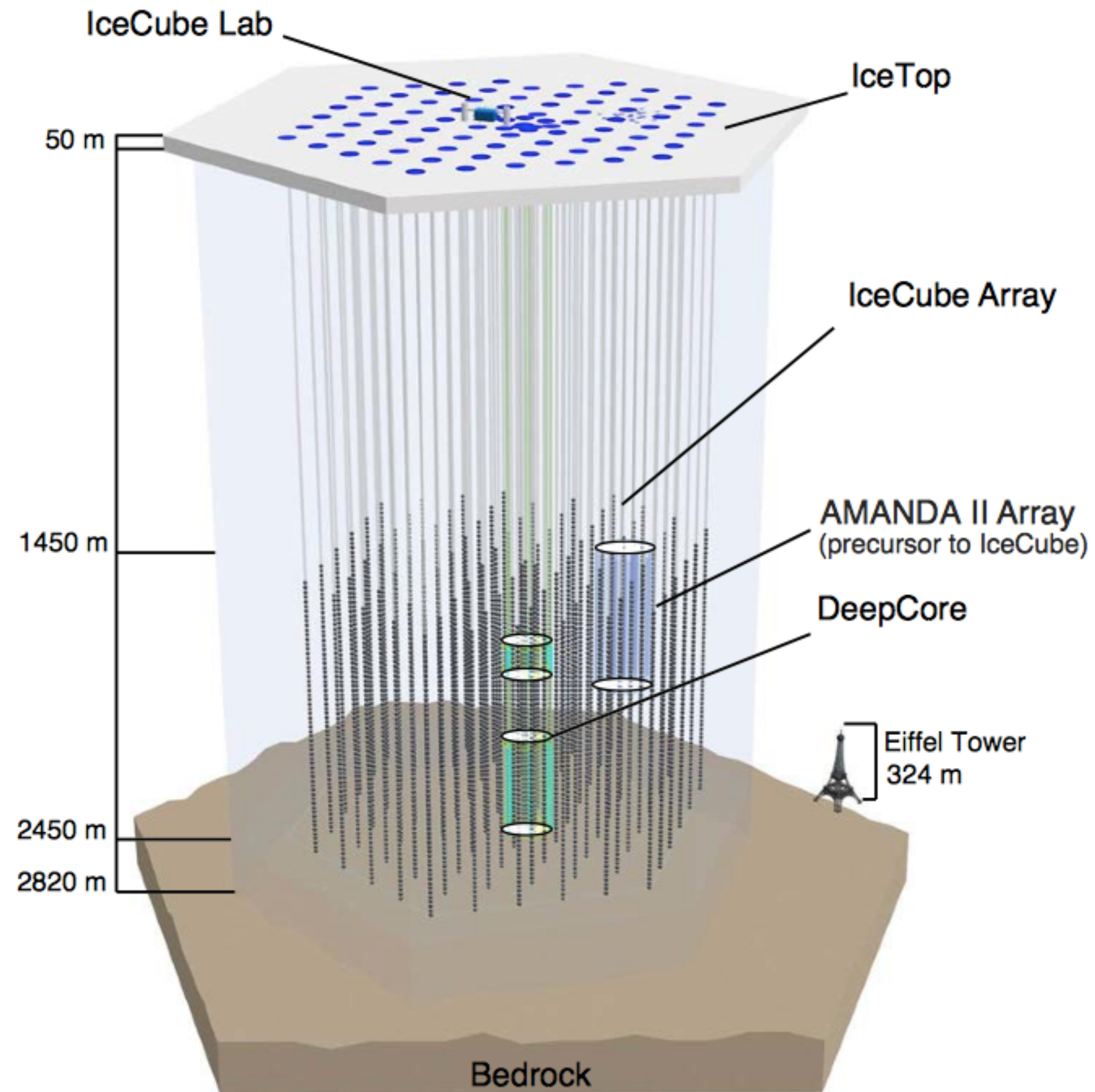
Neutrinos can tell us about “standard model” physics:

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New Physics?

Current observations: IceCube (south pole)

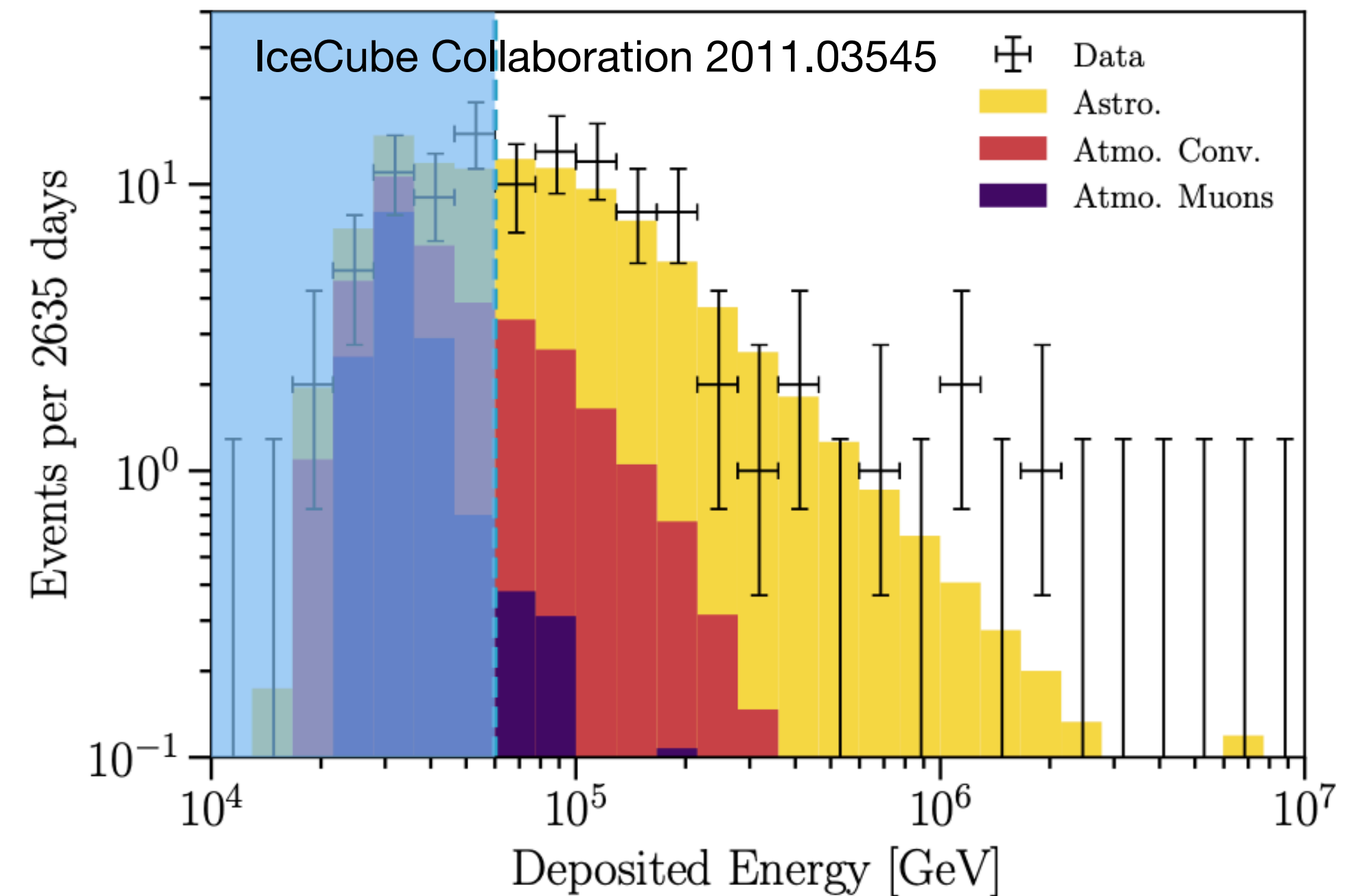
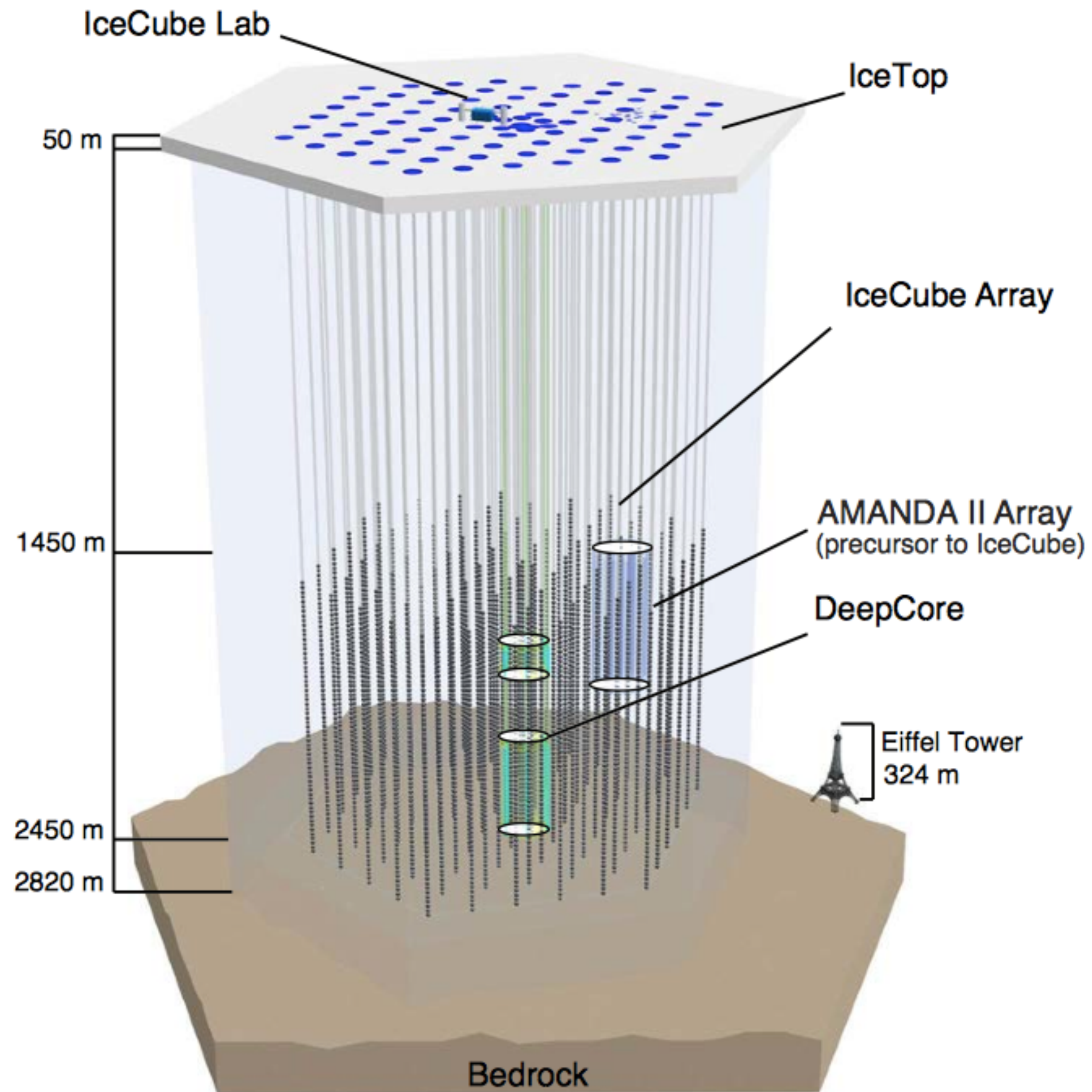
Effective volume $\sim 1 \text{ km}^3$



large exposures necessary due to low fluxes

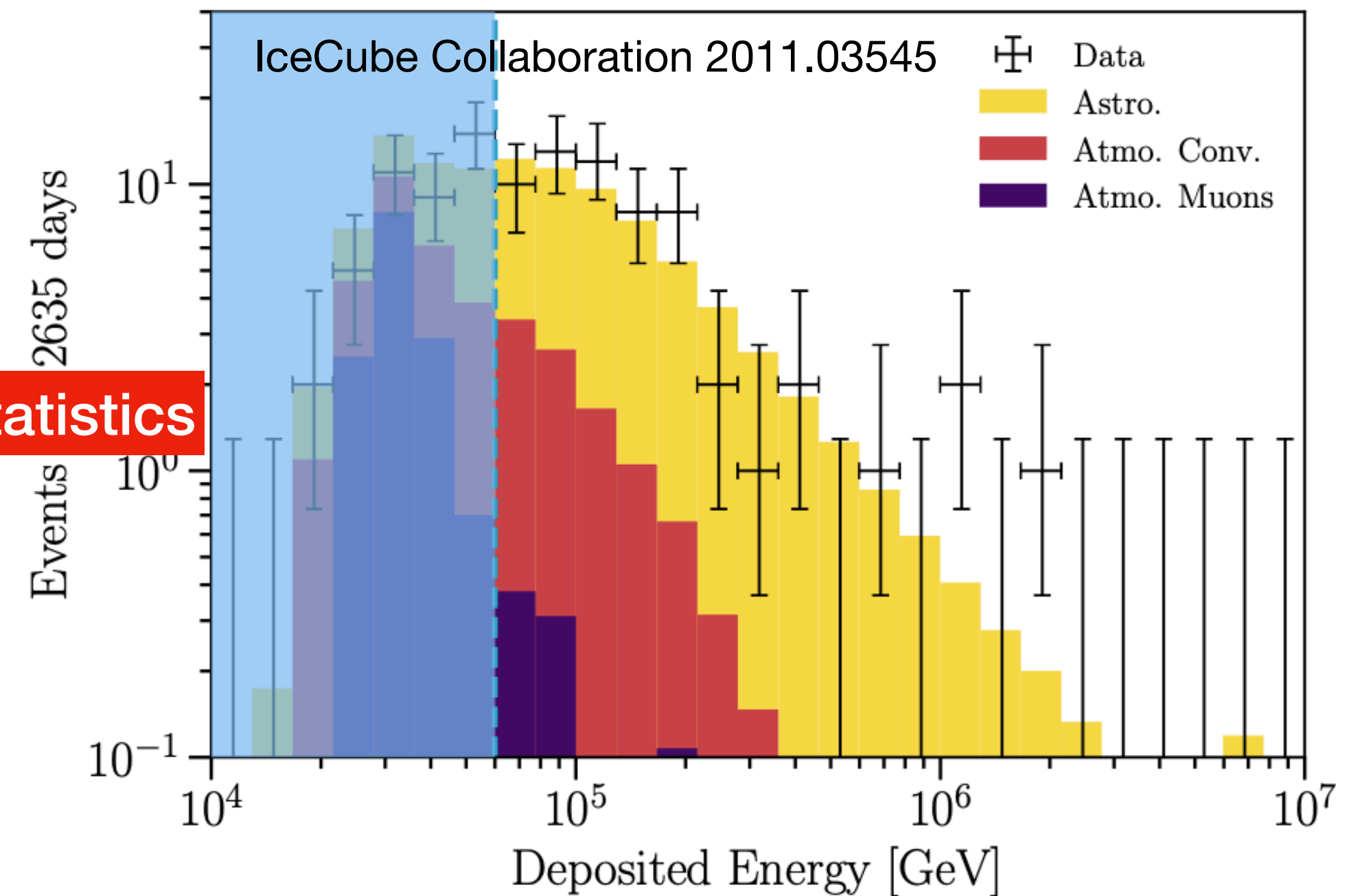
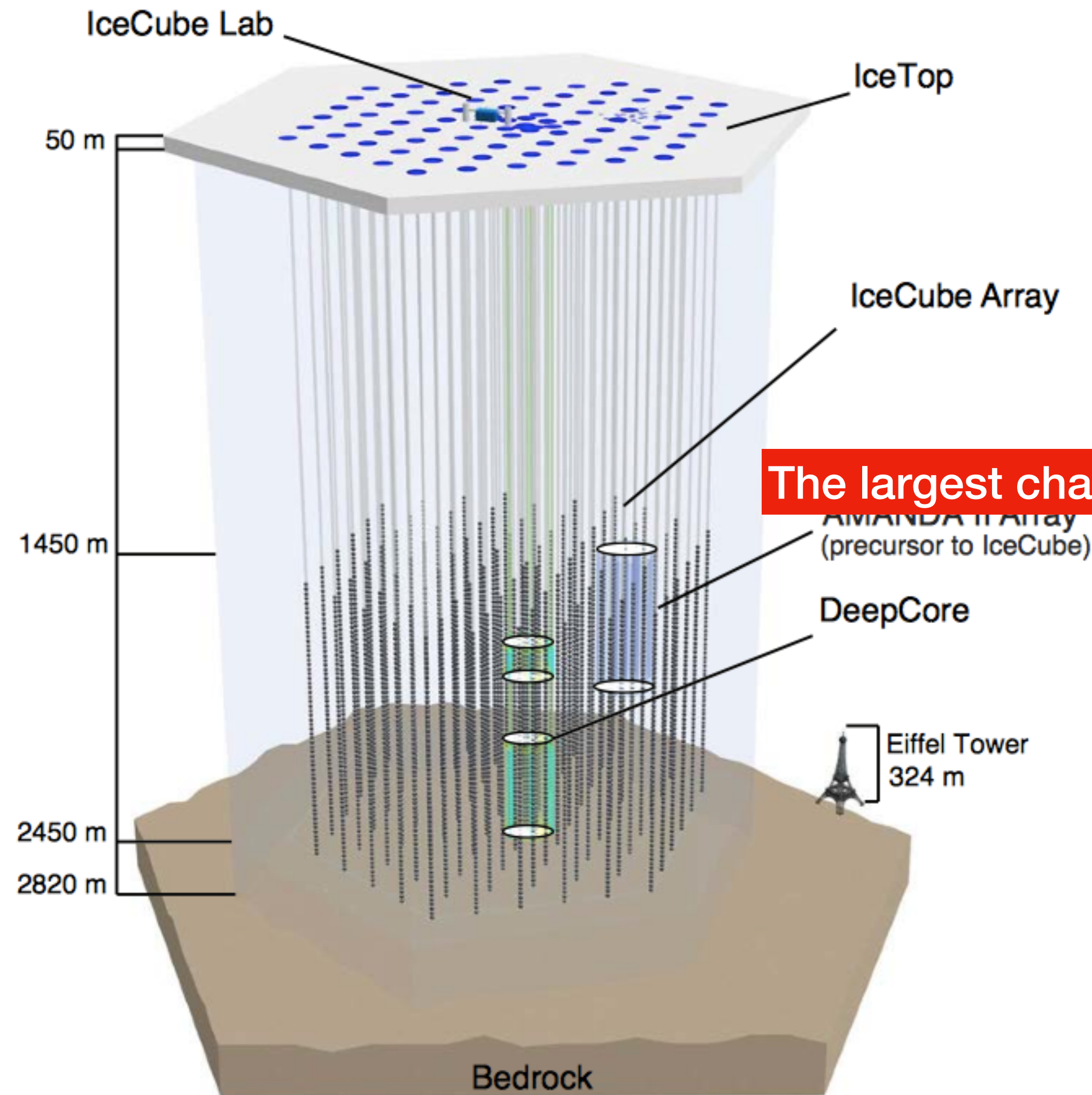
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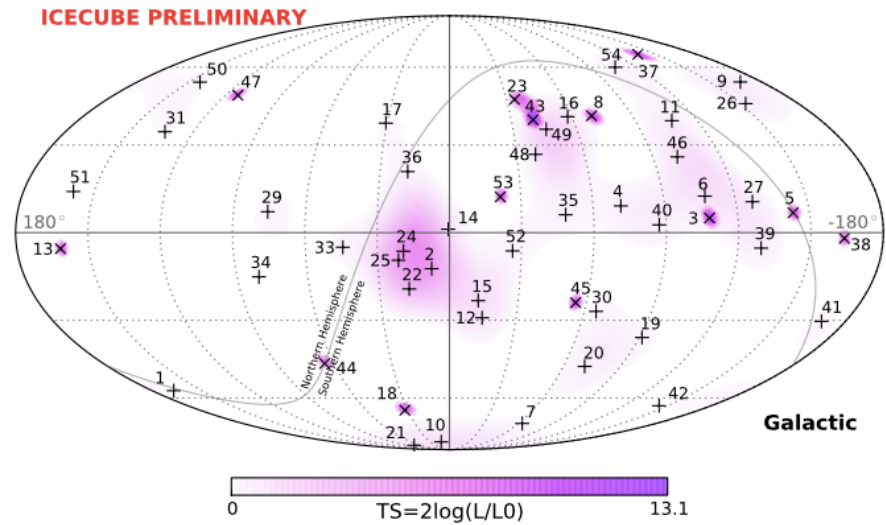
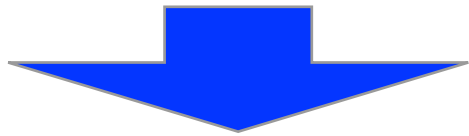


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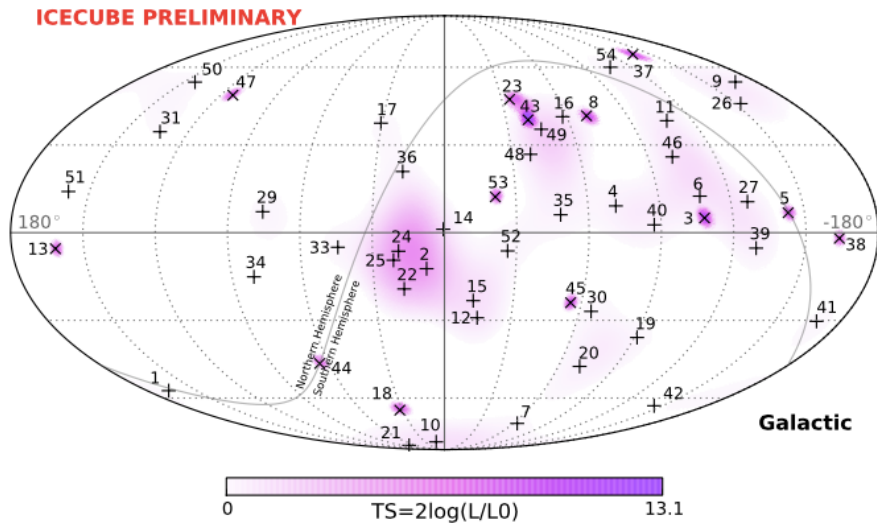
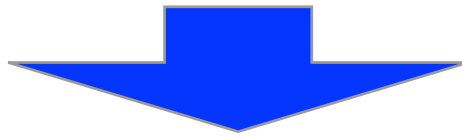
Effective volume $\sim 1 \text{ km}^3$



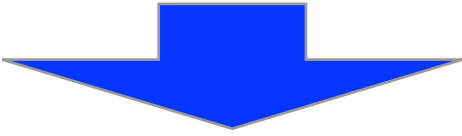
Arrival direction



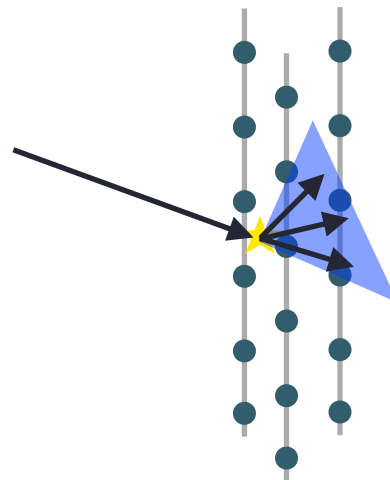
Arrival direction



Energy



Deposited
EM-equivalent



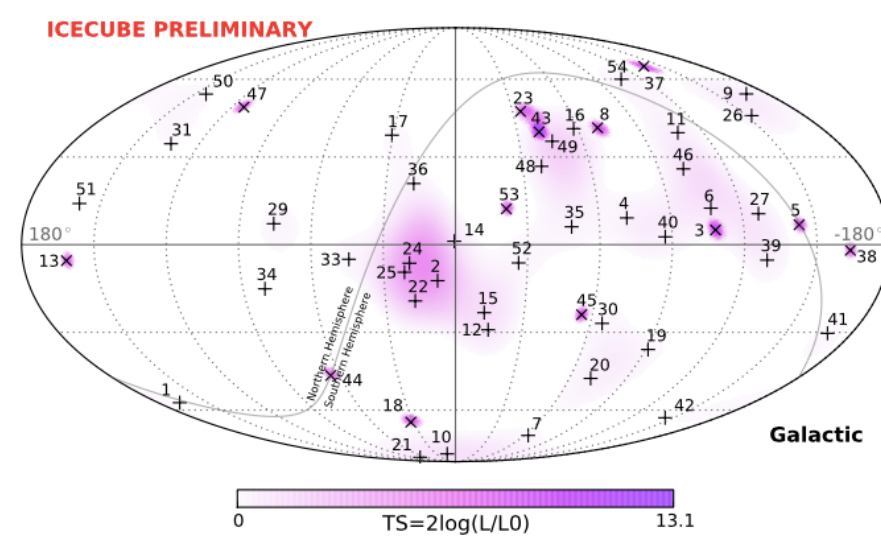
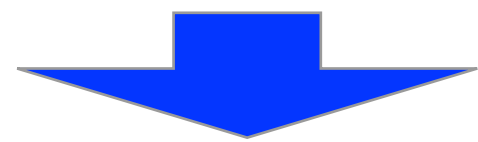
Arrival direction



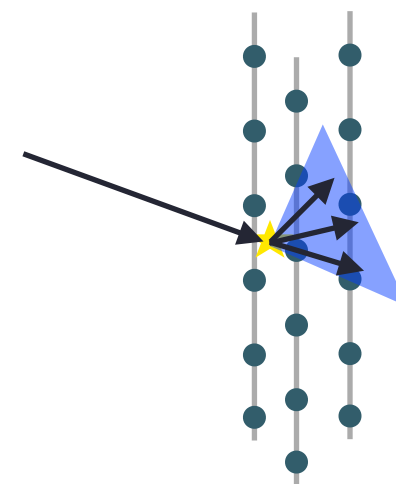
Energy



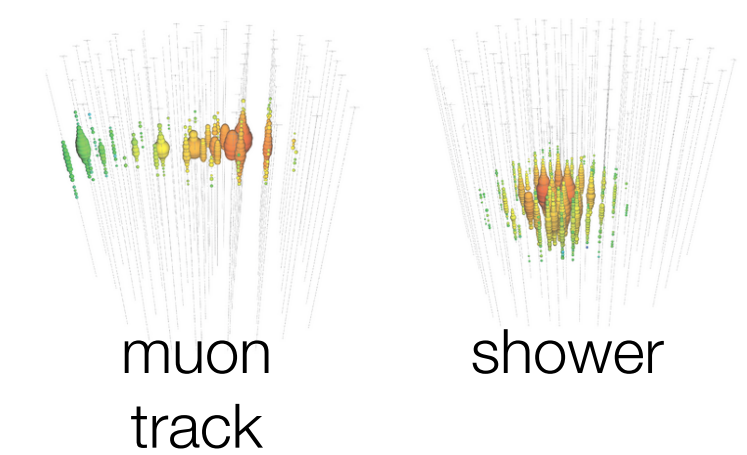
Flavour (e, μ, τ)



Deposited
EM-equivalent

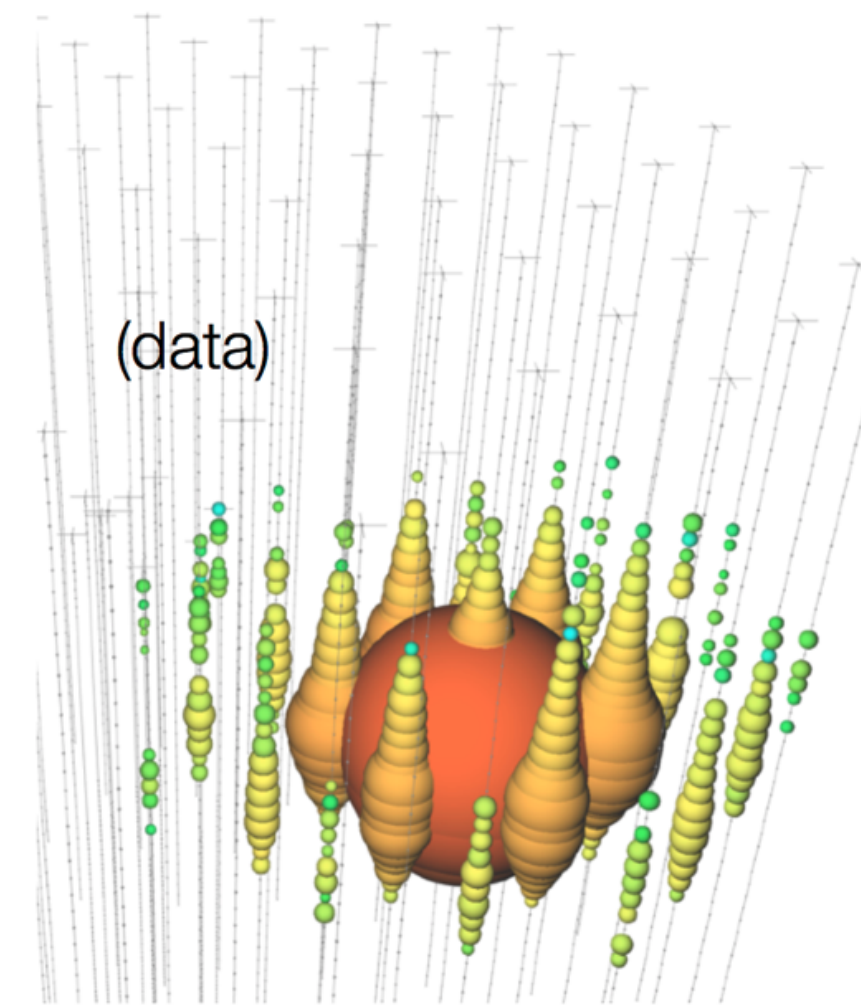


Morphology



Flavour: event morphology

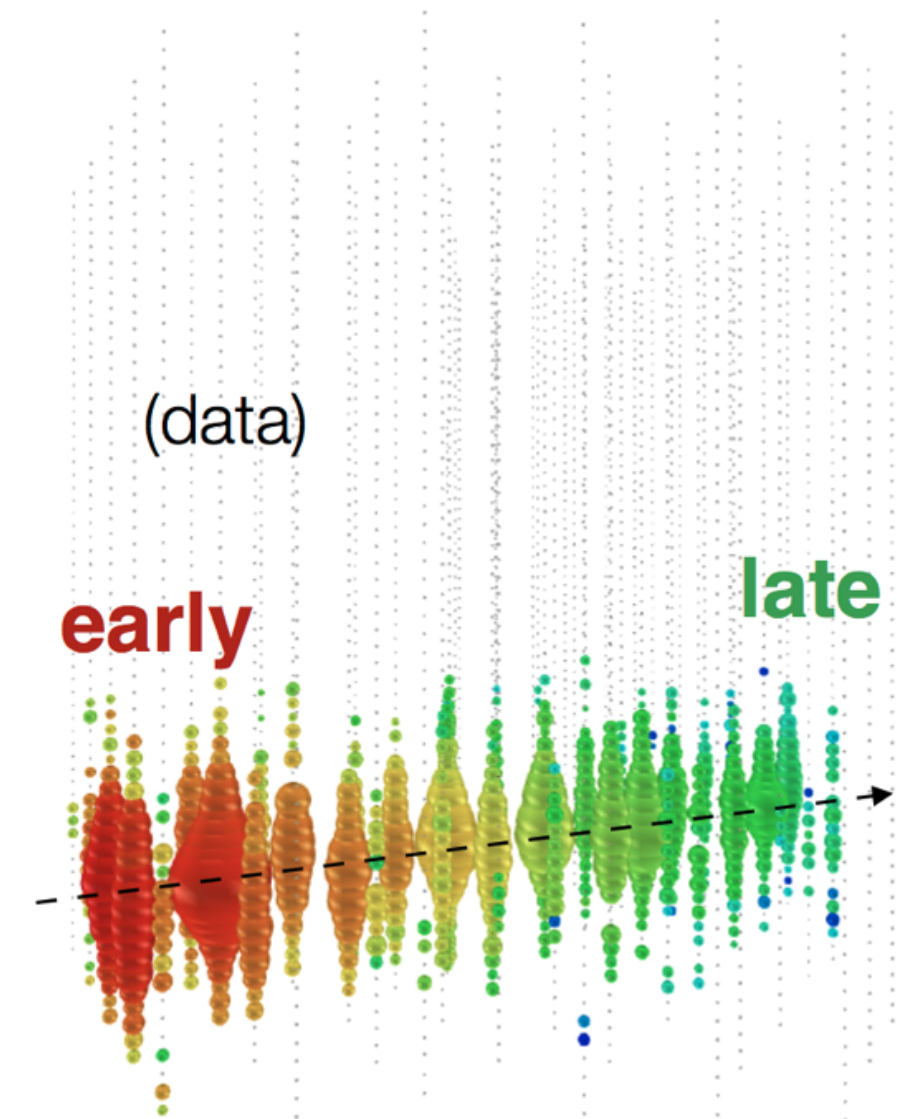
Neutral-current / ν_e



Isolated energy deposition (cascade) with no track

crap

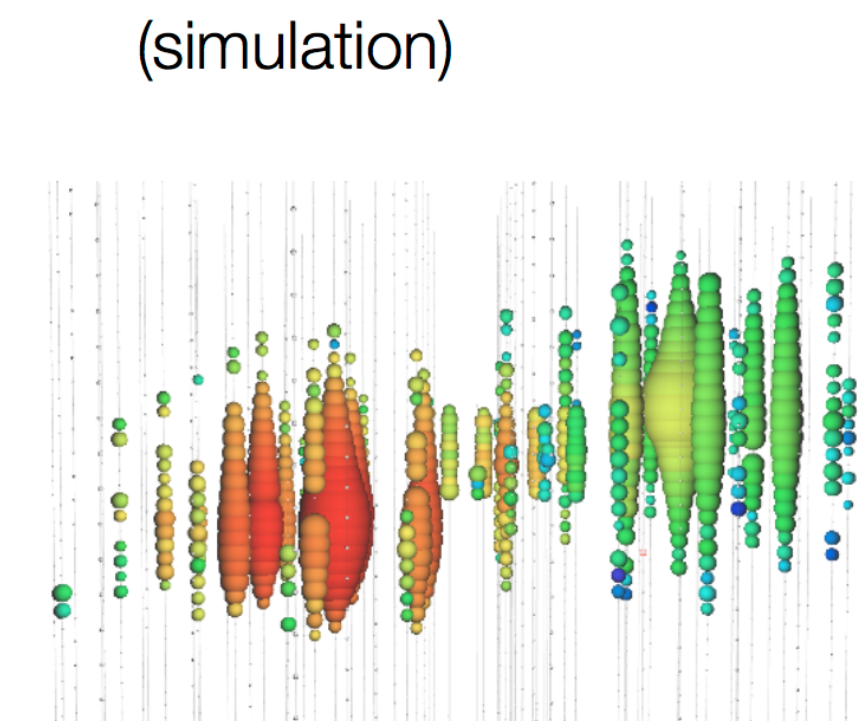
Charged-current ν_μ



Up-going track

ok

Charged-current ν_τ



Double cascade

depends

Angular resolution:

Flavour composition in astrophysical sources

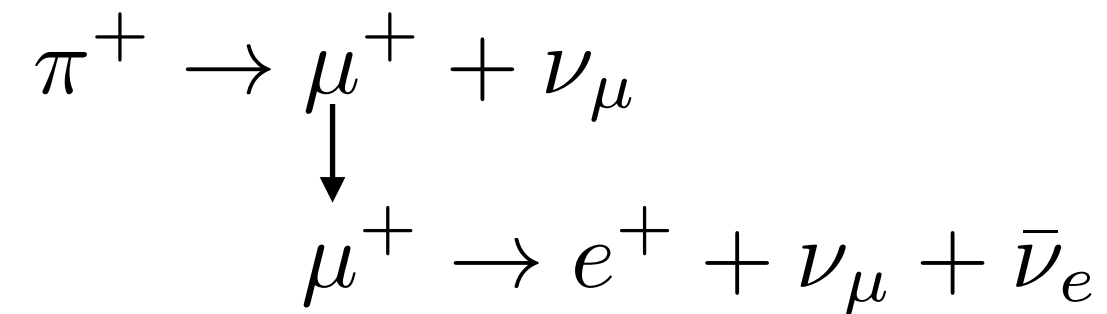
(GRBs, AGNs, blazars, pulsars...)

$(\alpha_e : \alpha_\mu : \alpha_\tau)$



Pion sources

(c.c. for π^-)



(1 : 2 : 0)



“muon-damped”

(c.c. for π^-)

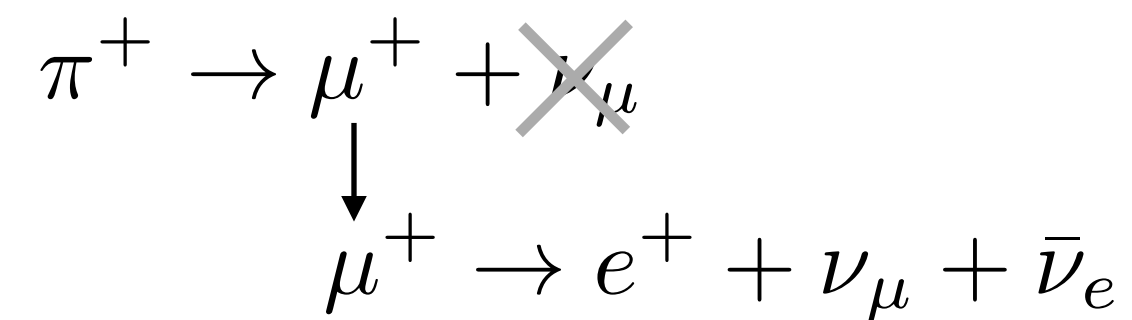


(0 : 1 : 0)



“muon source”

(c.c. for π^-)



(1 : 1 : 0)



Neutron source

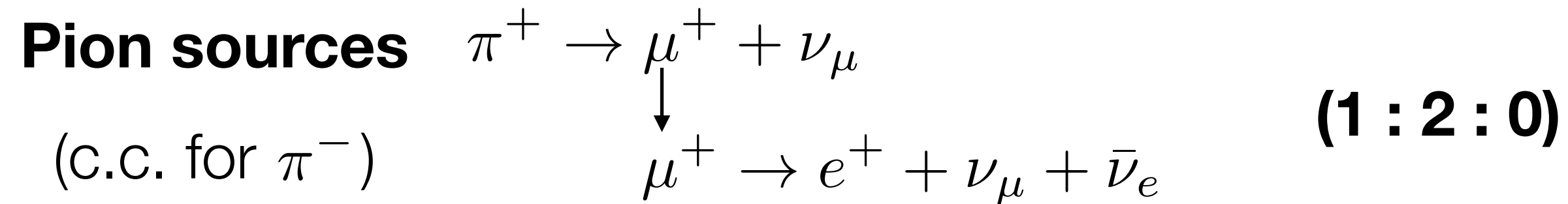


(1 : 0 : 0)

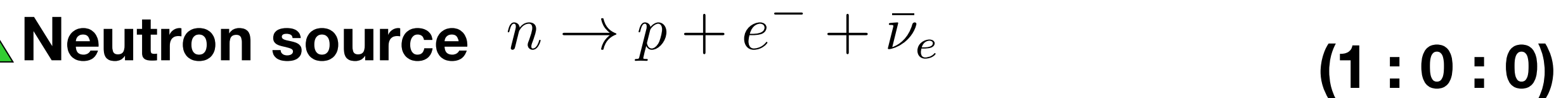
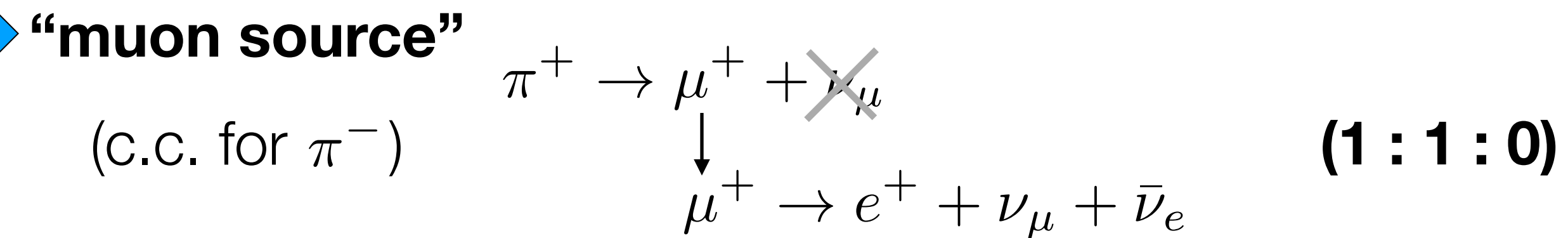
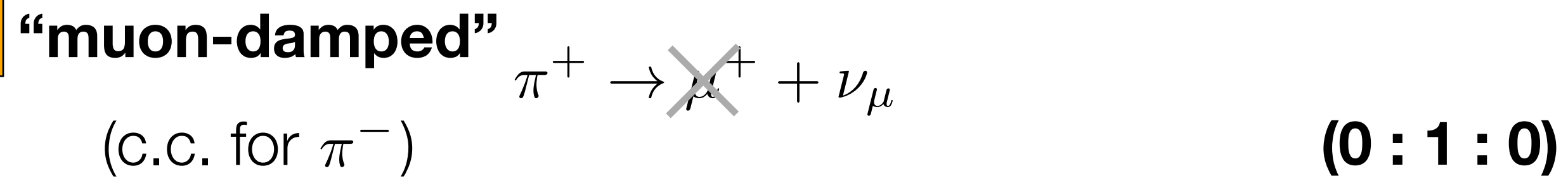
Flavour composition in astrophysical sources

(GRBs, AGNs, blazars, pulsars...)

$(\alpha_e : \alpha_\mu : \alpha_\tau)$



Different scenarios: different production environments

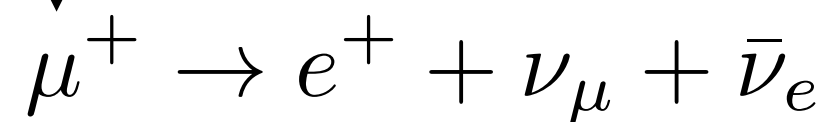
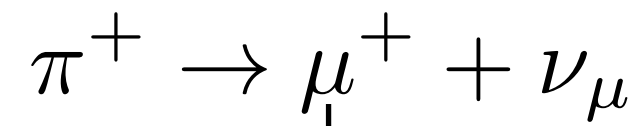


Flavour composition in astrophysical sources

(GRBs, AGNs, blazars, pulsars...)

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Pion sources



(1 : 2 : 0)

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Different scenarios: different production environments

“muon-damped”

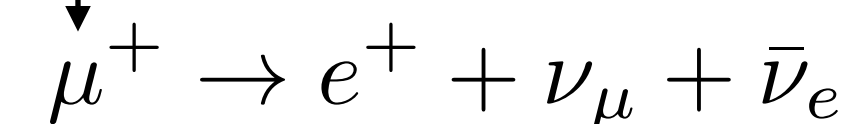
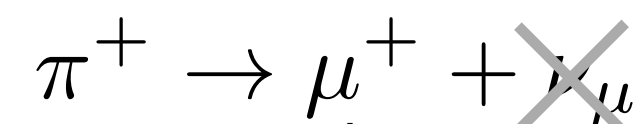


(0 : 1 : 0)

(c.c. for π^-)

Flavour can be distinguished **statistically** in neutrino detectors: different charged-current interactions lead to different event **morphologies** (there is some degeneracy)

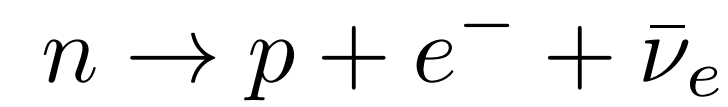
“muon source”



(1 : 1 : 0)

(c.c. for π^-)

Neutron source

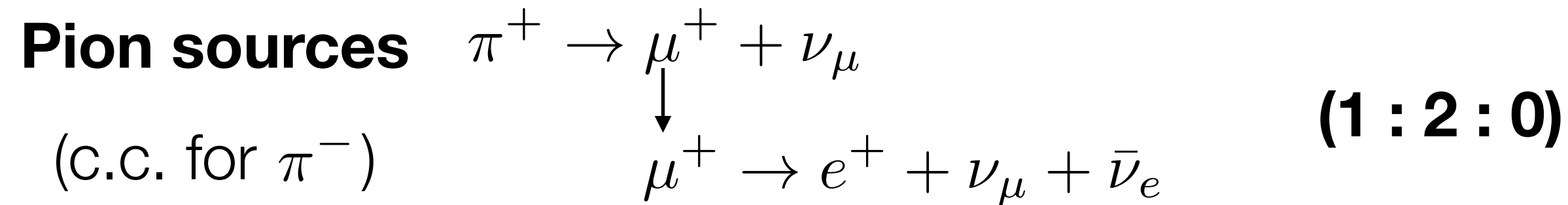


(1 : 0 : 0)

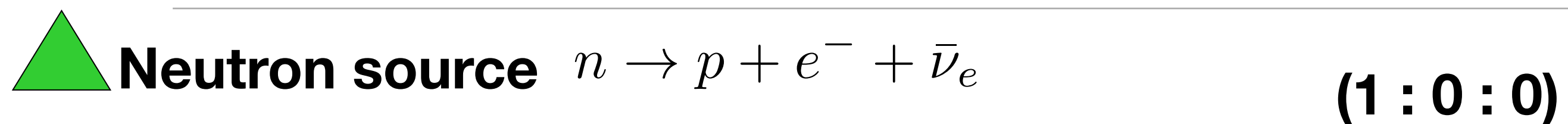
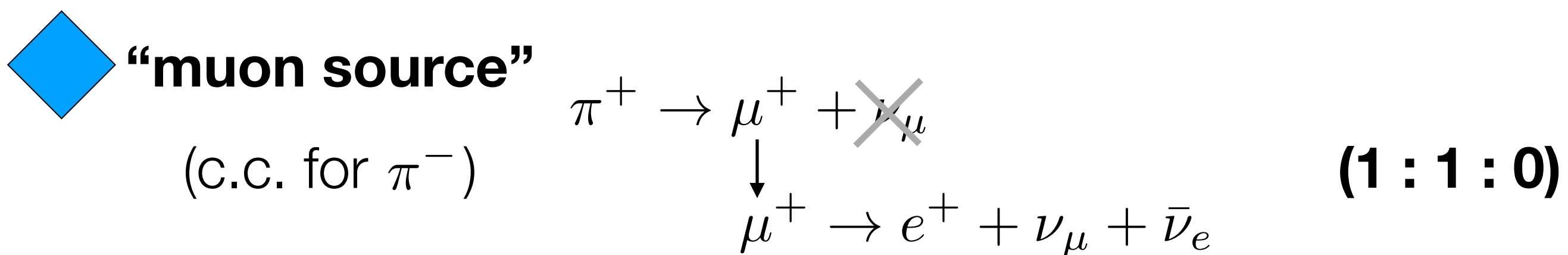
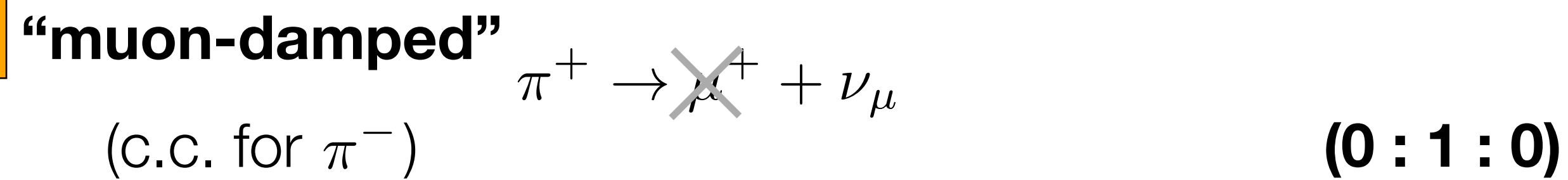
Flavour composition in astrophysical sources

(GRBs, AGNs, blazars, pulsars...)

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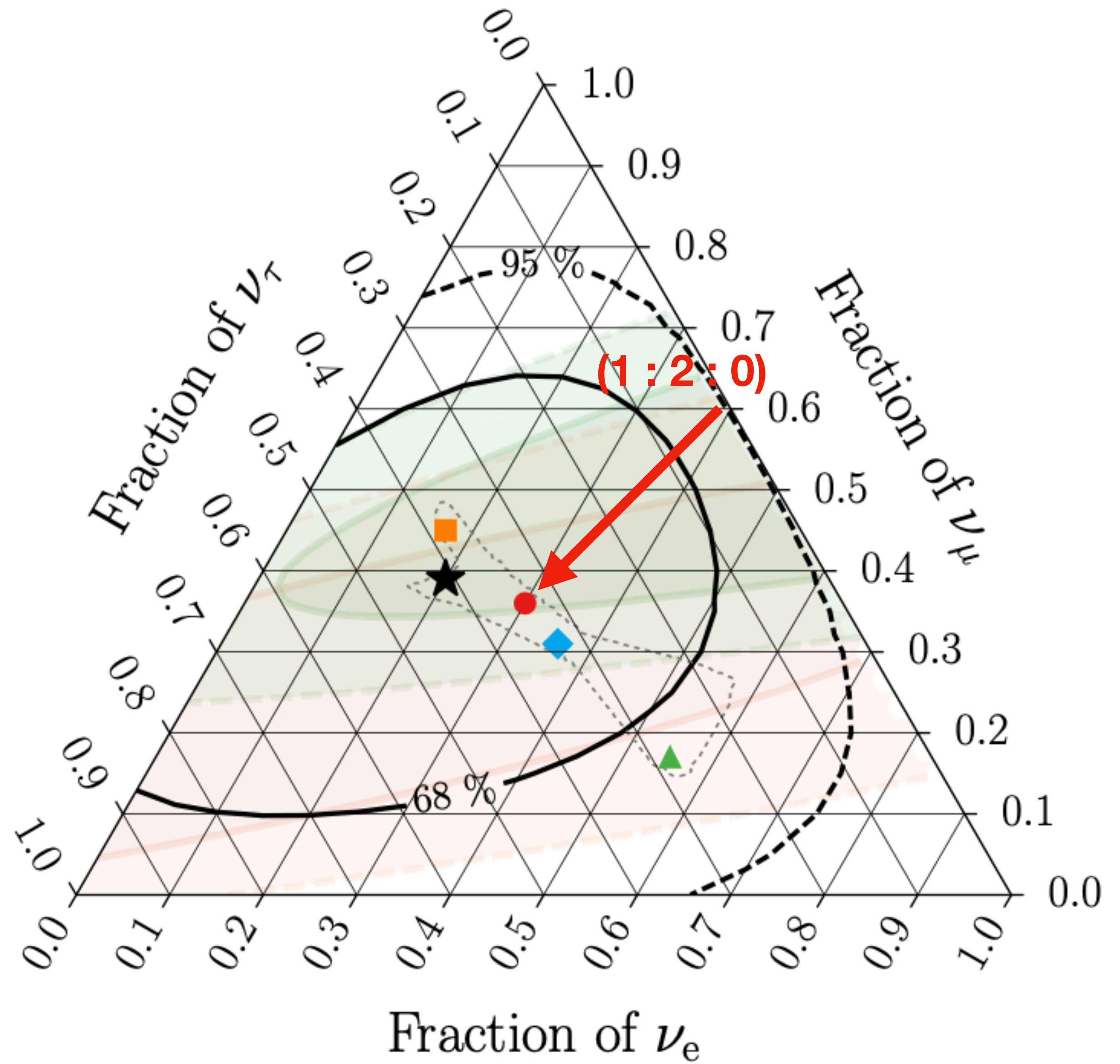


Different scenarios: different production environments



Flavour can be distinguished **statistically** in neutrino detectors: different charged-current interactions lead to different event **morphologies** (there is some degeneracy)

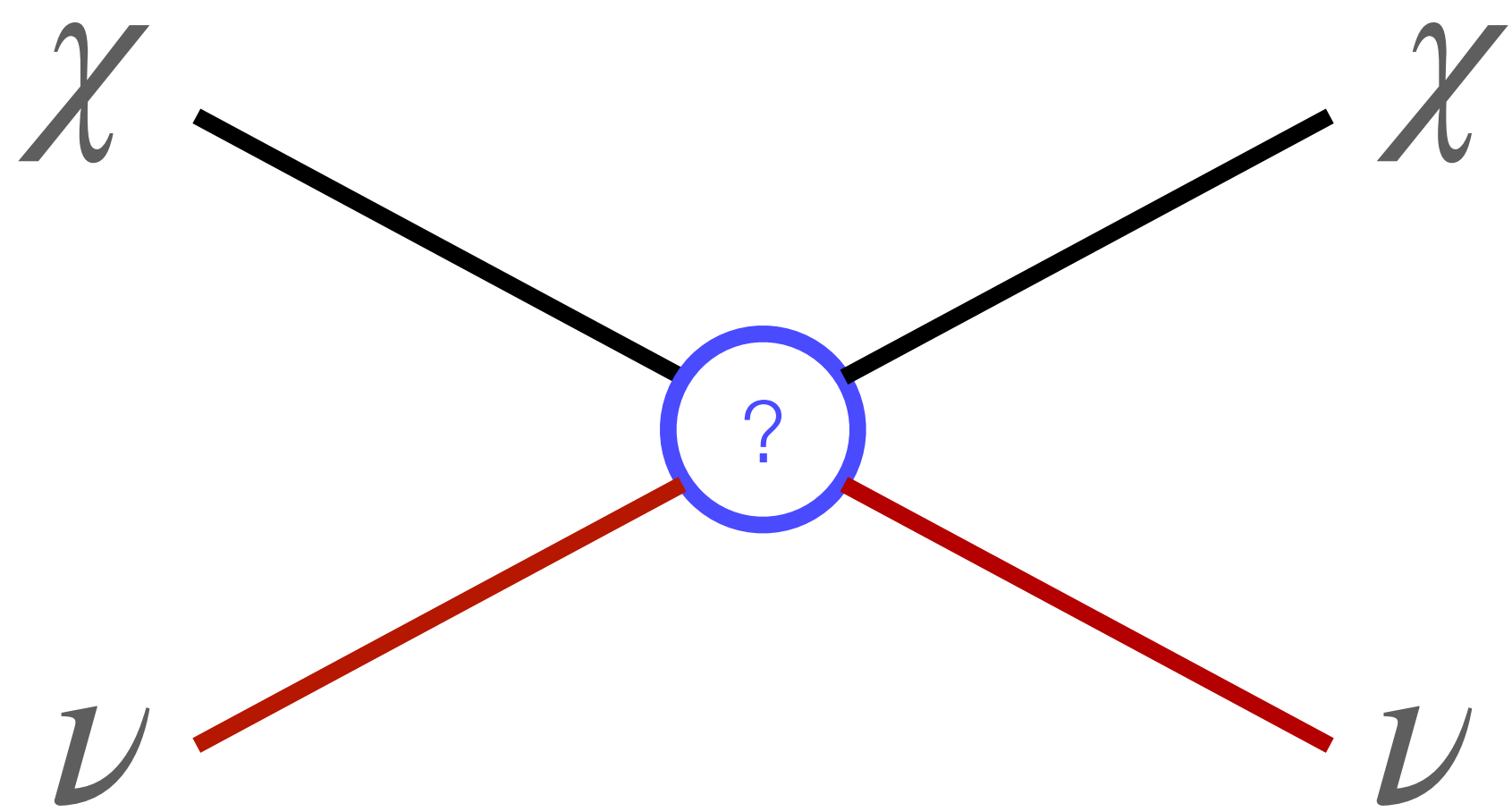
Everything oscillates to $\sim 1:1:1$



Dark Matter?



Elastic scattering



Elastic scattering

We've just heard about limits on dark matter-neutrinos scattering at low (~recombination) energies.

Generically, cross sections grow as

$$s = E_{CM}^2$$

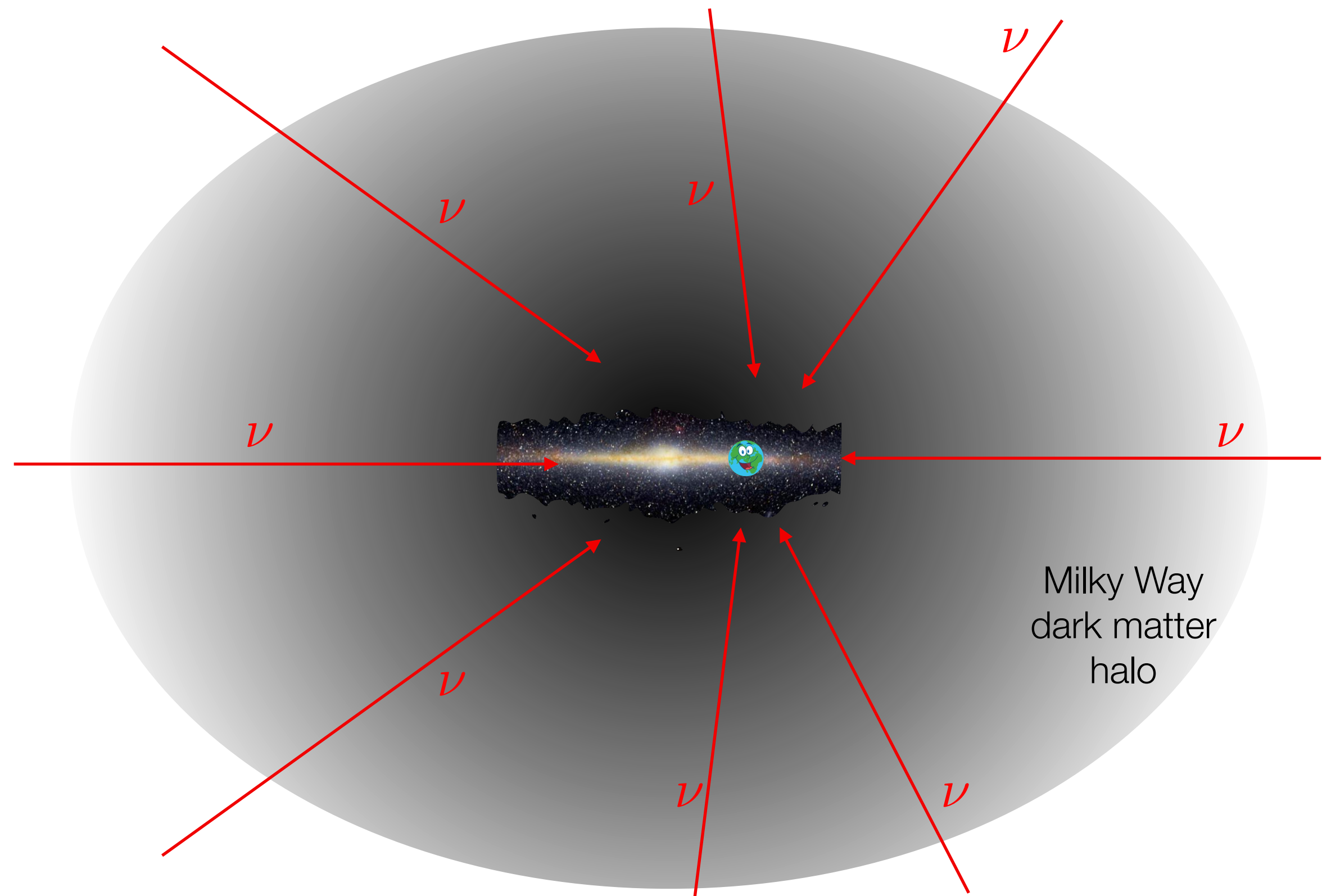
$$\sigma_{DM-\nu} \propto E_\nu^2$$

IceCube has seen events above a PeV....

$$\left(\frac{\text{PeV}}{T_{\nu, recomb.}} \right)^2 \sim 10^{30}$$

Let's look there!

Isotropic extragalactic neutrino flux



b, l : galactic latitude, longitude

column density: $\tau(b, l) = \int_{l.o.s} n_{\chi}(x; b, l) dx.$

$$\frac{d\Phi(E, \tau)}{d\tau} = -\sigma(E)\Phi(E, \tau) + \int_E^{\infty} d\tilde{E} \frac{d\sigma(\tilde{E}, E)}{dE} \Phi(\tilde{E}, \tau)$$



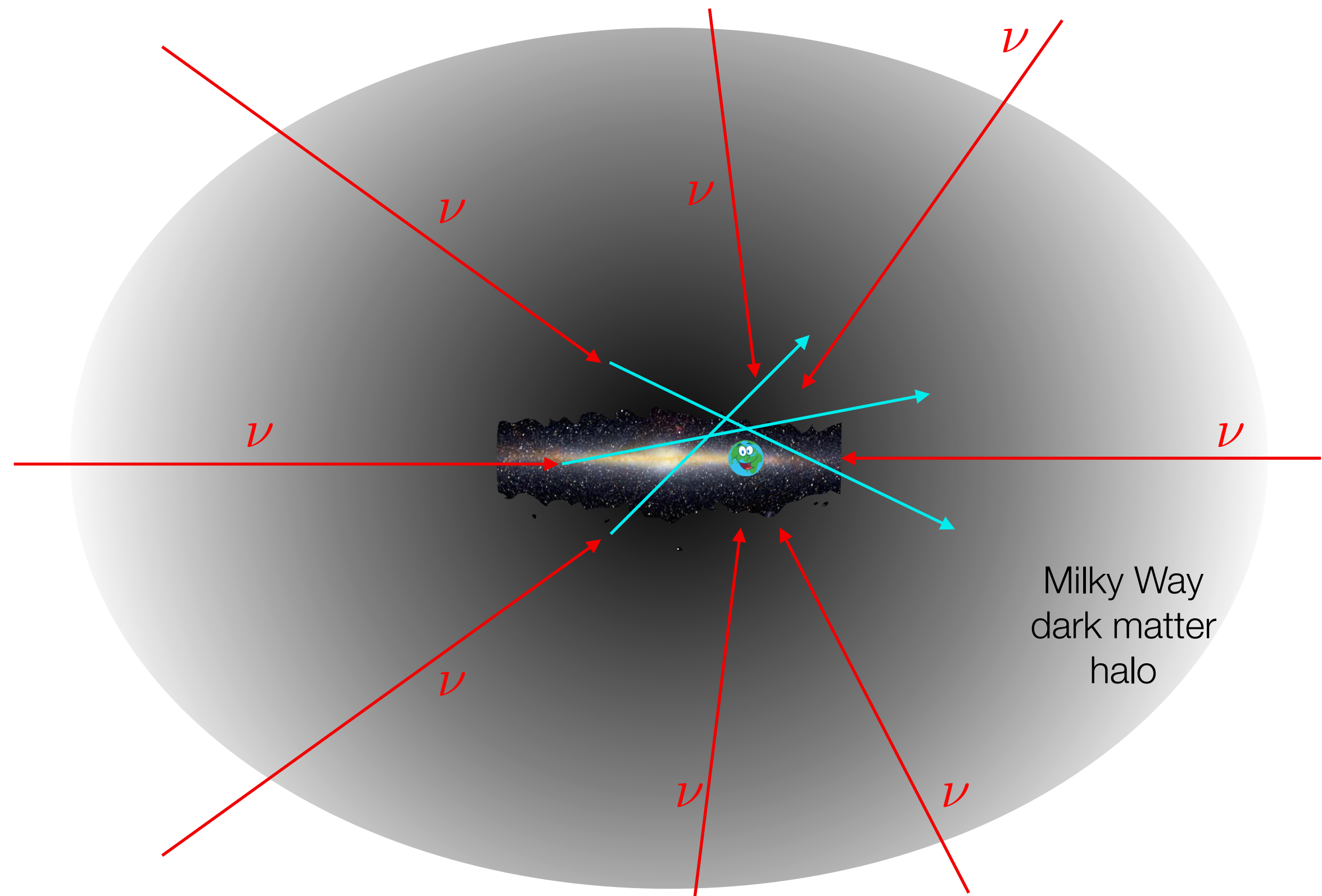
scattering **from** E
to any energy



scattering **to** E from
any energy \tilde{E}

Solve to find flux at earth at energy E and direction (b, l)

Isotropic extragalactic neutrino flux



Anisotropic deflection/energy loss

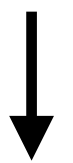
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scattering **from** E
to any energy



scattering **to** E from
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Solve to find flux at earth at energy E and direction (b, l)

Solve by being clever

$$\frac{d\Phi(E, \tau)}{d\tau} = -\sigma(E)\Phi(E, \tau) + \int_E^\infty d\tilde{E} \frac{d\sigma(\tilde{E}, E)}{dE} \Phi(\tilde{E}, \tau)$$

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$$E \rightarrow \vec{E} \quad \Phi \rightarrow \vec{\Phi} \quad C_{ij} = d\tilde{E}_i \frac{d\sigma}{dE}(\tilde{E}_i, E_j)$$

Solve by being clever

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$$\vec{\Phi}'(\tau) = -(\text{diag}(\vec{\sigma}) + C)\vec{\Phi}(\tau) \quad \begin{array}{l} \lambda_i \text{ eigenvalues} \\ \hat{\phi}_i \text{ eigenvectors} \end{array}$$

$$\vec{\Phi} = \sum c_i \hat{\phi}_i e^{\lambda_i \tau}$$

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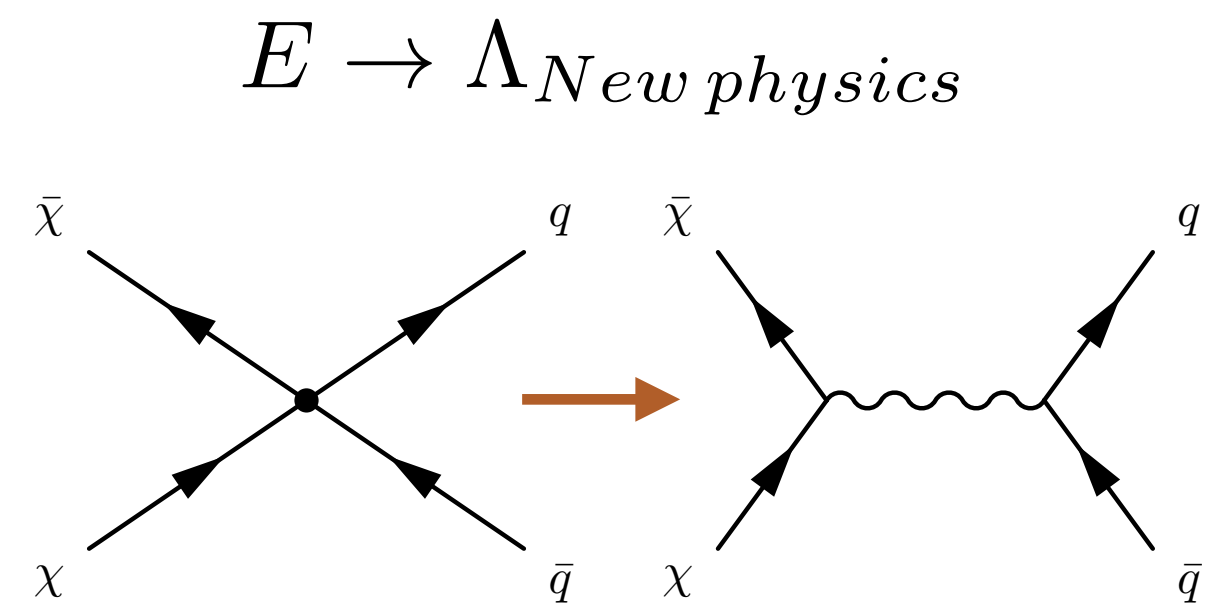
$$\vec{\Phi} = \sum c_i \hat{\phi}_i e^{\lambda_i \tau}$$

c_i 's determined by initial condition of isotropic power law flux

$$\sigma_{DM-\nu} \propto E_\nu^2 \xrightarrow{??} \left(\frac{\text{PeV}}{T_{\nu, \text{recomb.}}} \right)^2 \sim 10^{30}$$

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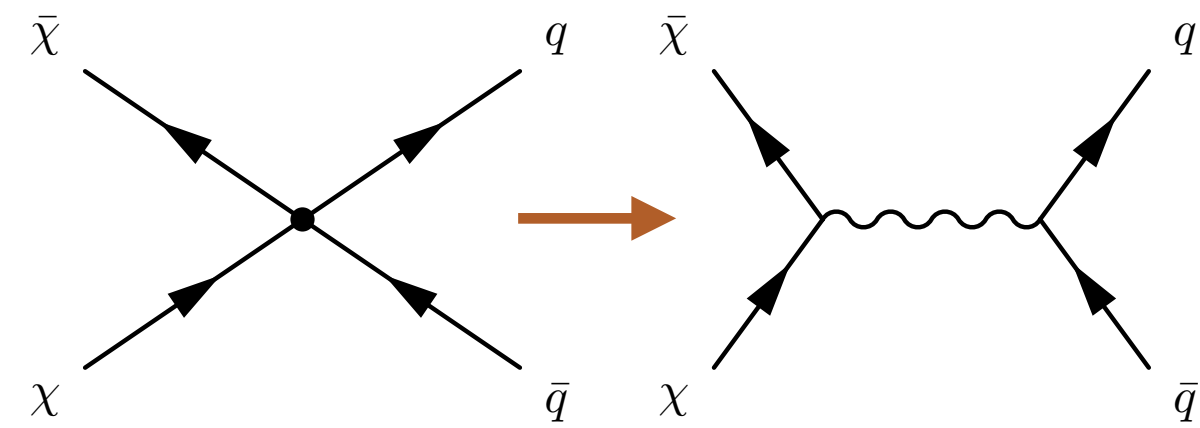


The low energy approximation does not work at a PeV!!

Begin to resolve microphysics: **need more concrete model**

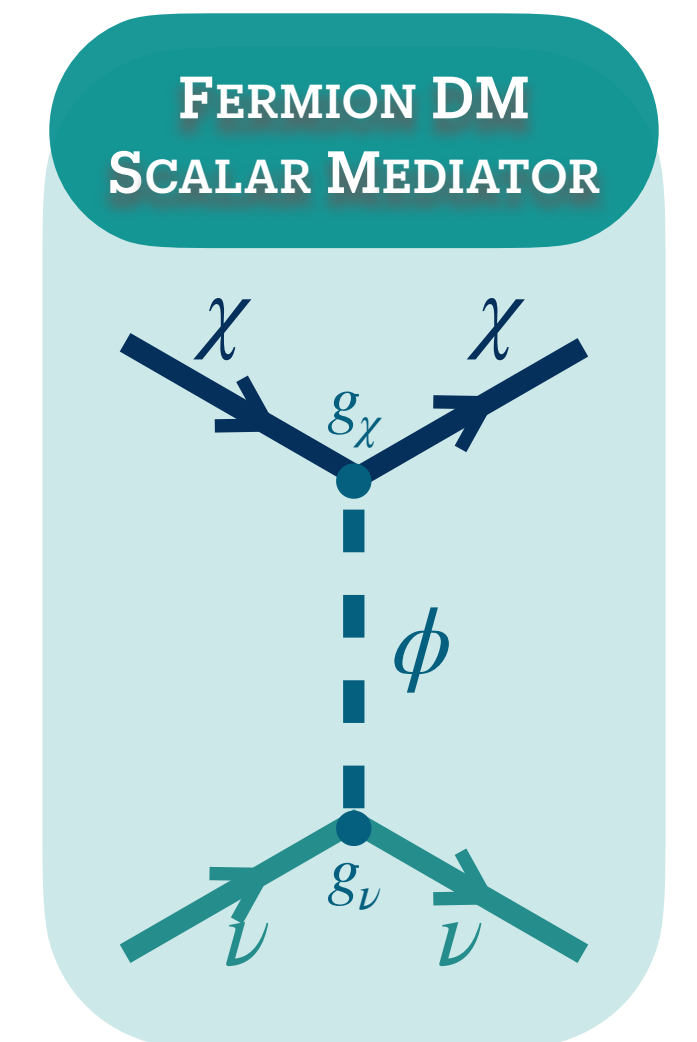
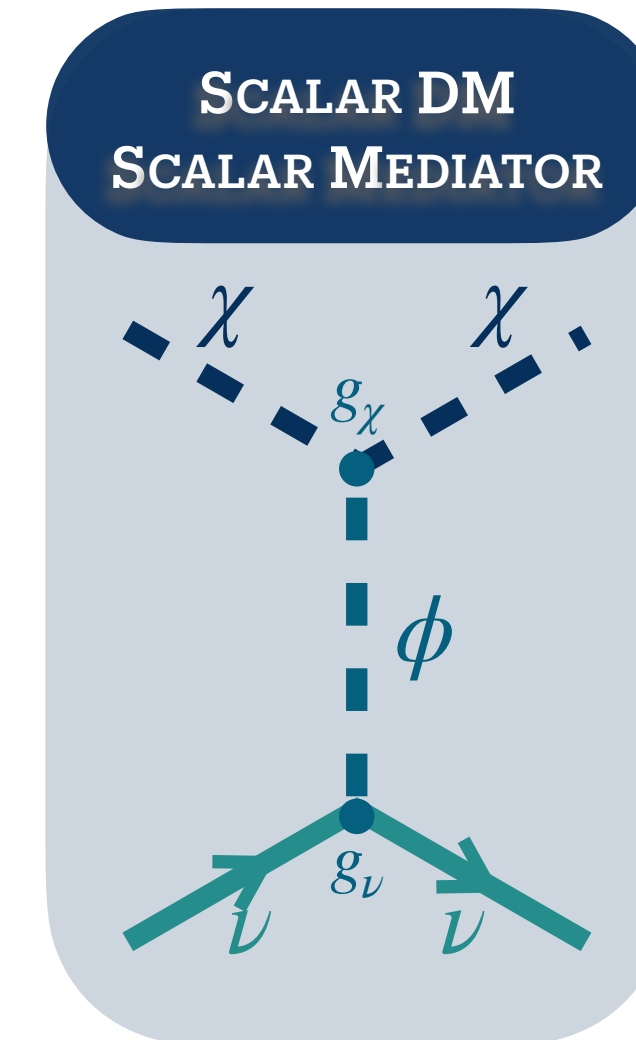
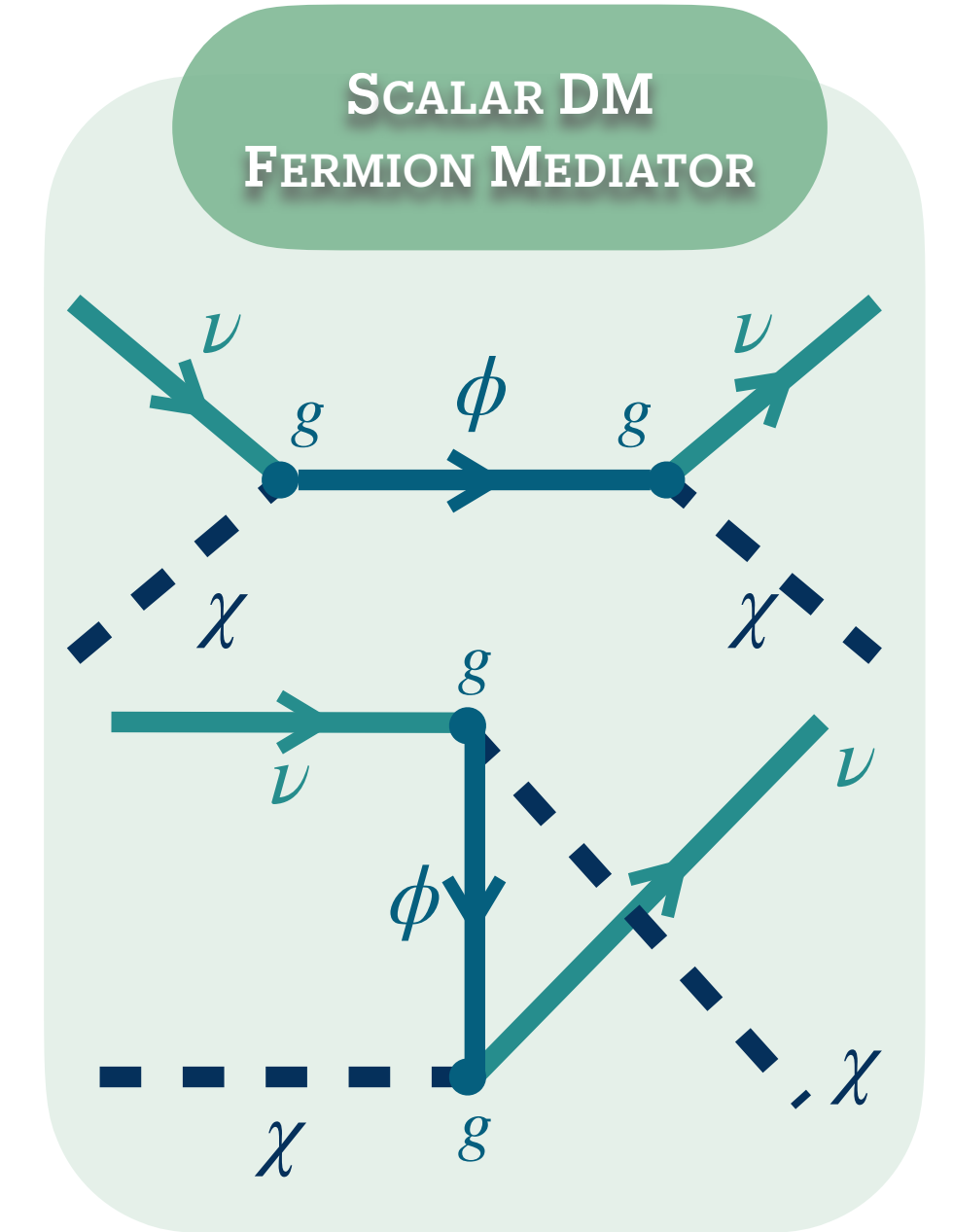
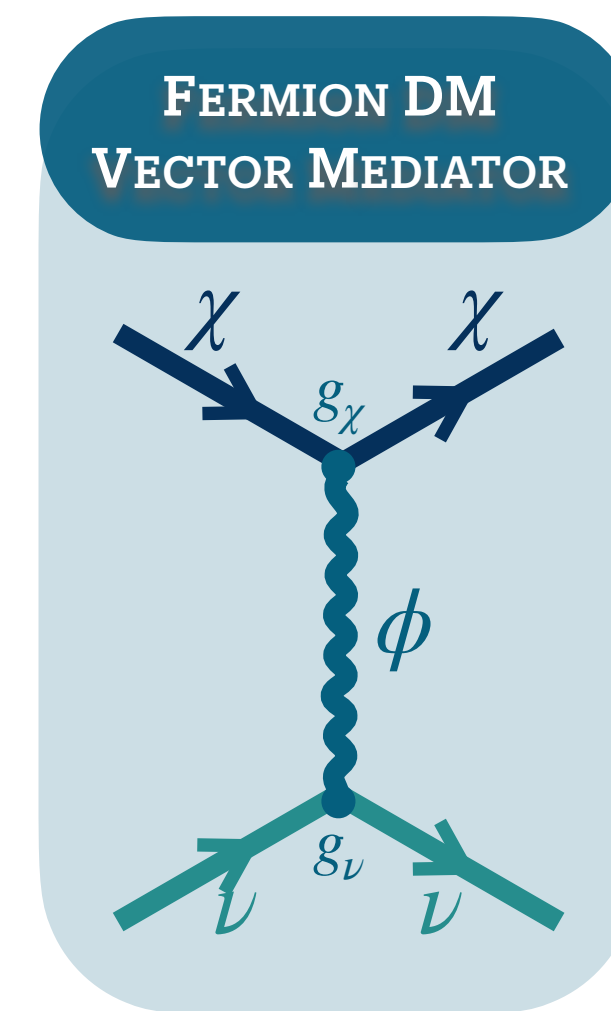
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$E \rightarrow \Lambda_{New\ physics}$

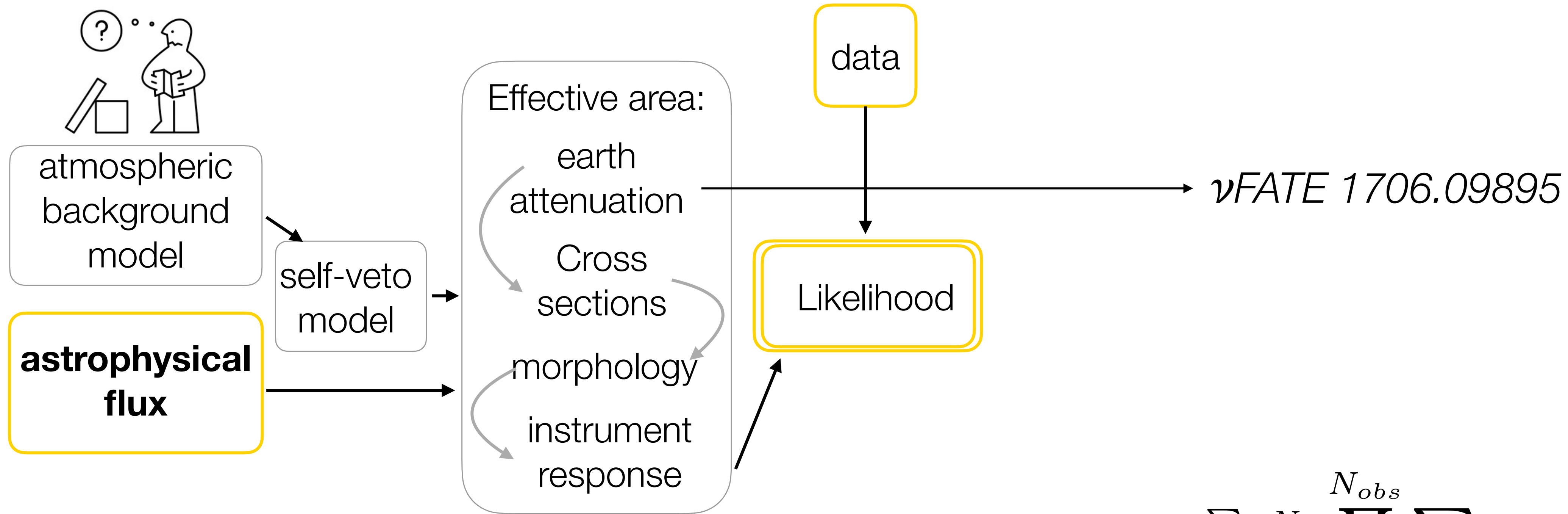


The low energy approximation does not work at a PeV!!

Begin to resolve microphysics: **need more concrete model**



Some simplified models: specify
 -Dark matter spin
 -Mediator spin

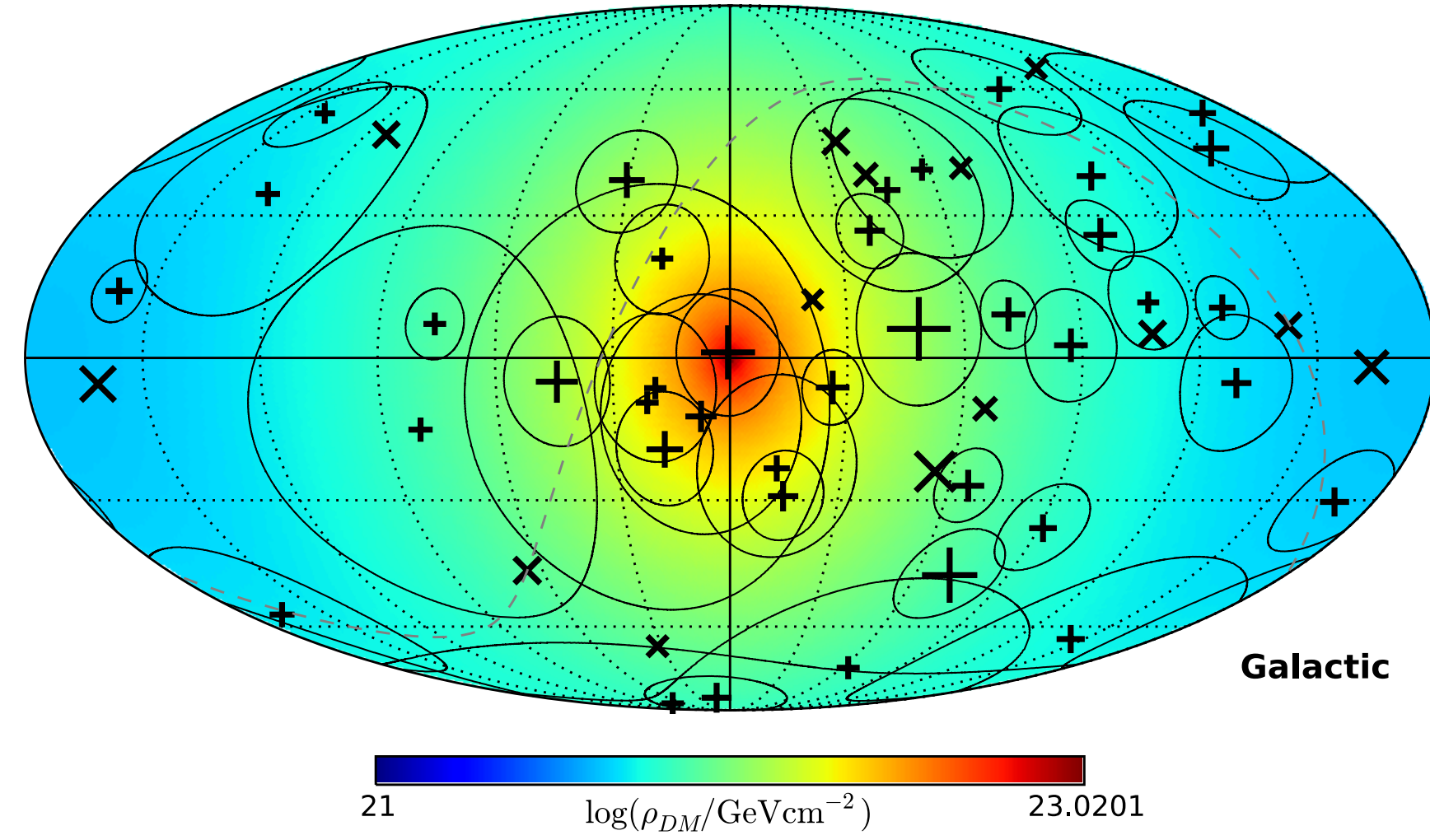


Classical recipe: 1502.02649

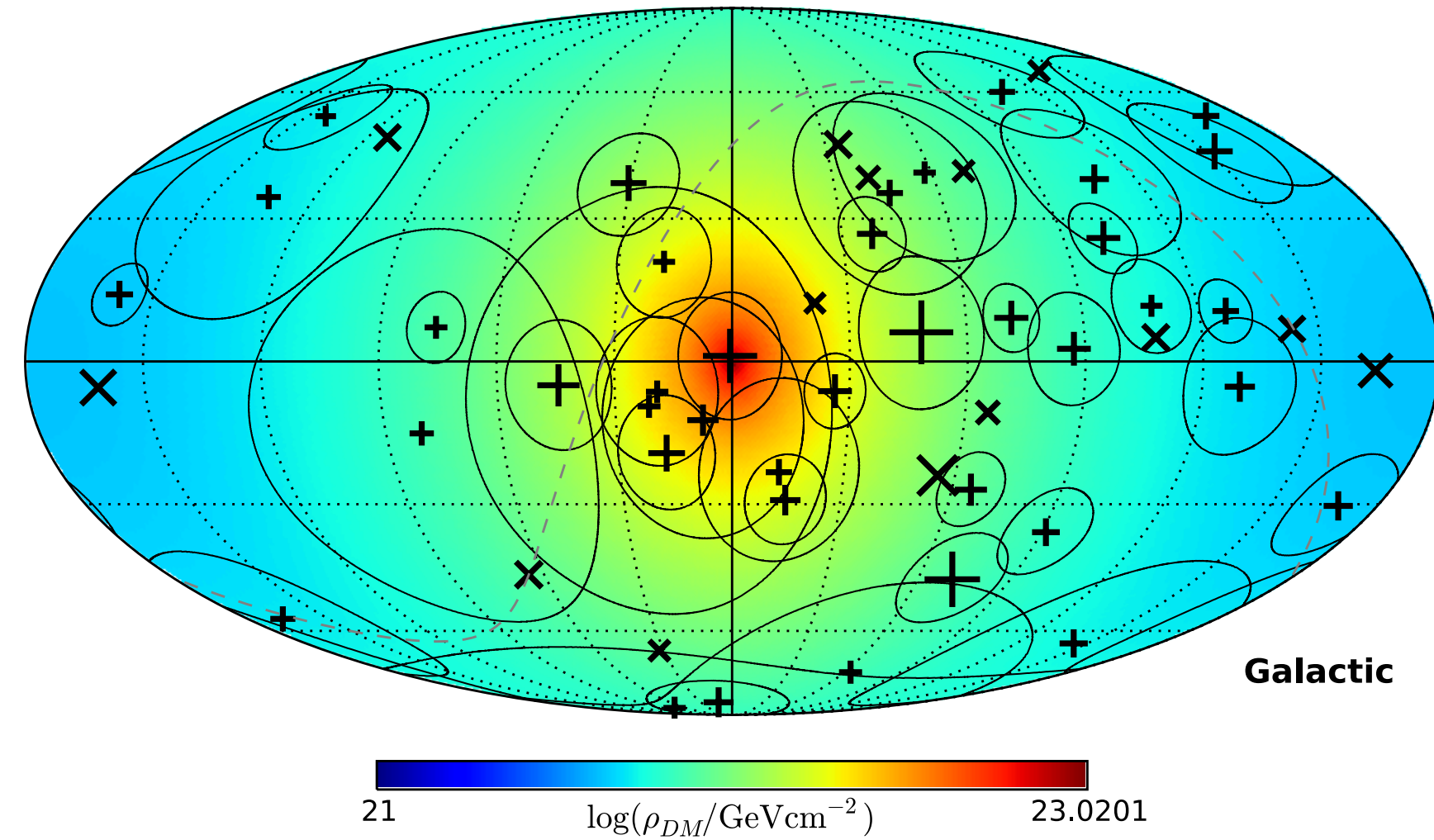
$$\mathcal{L}(\{t, E, \vec{x}\}|\vartheta) = e^{-\sum_a N_a} \prod_{i=1}^{N_{obs}} \sum_a N_a P_a(t_i, E_i, \vec{x}_i|\vartheta),$$

- i : observed (or simulated) event
- a : source (atmospheric muon, atm. neutrino, astro. neutrino)
- t : morphology (track, shower)
- E, x : Energy, arrival direction
- $P_a(t_i, E_i, \vec{x}_i|\vartheta)$ Probability of given event properties, given model parameters ϑ

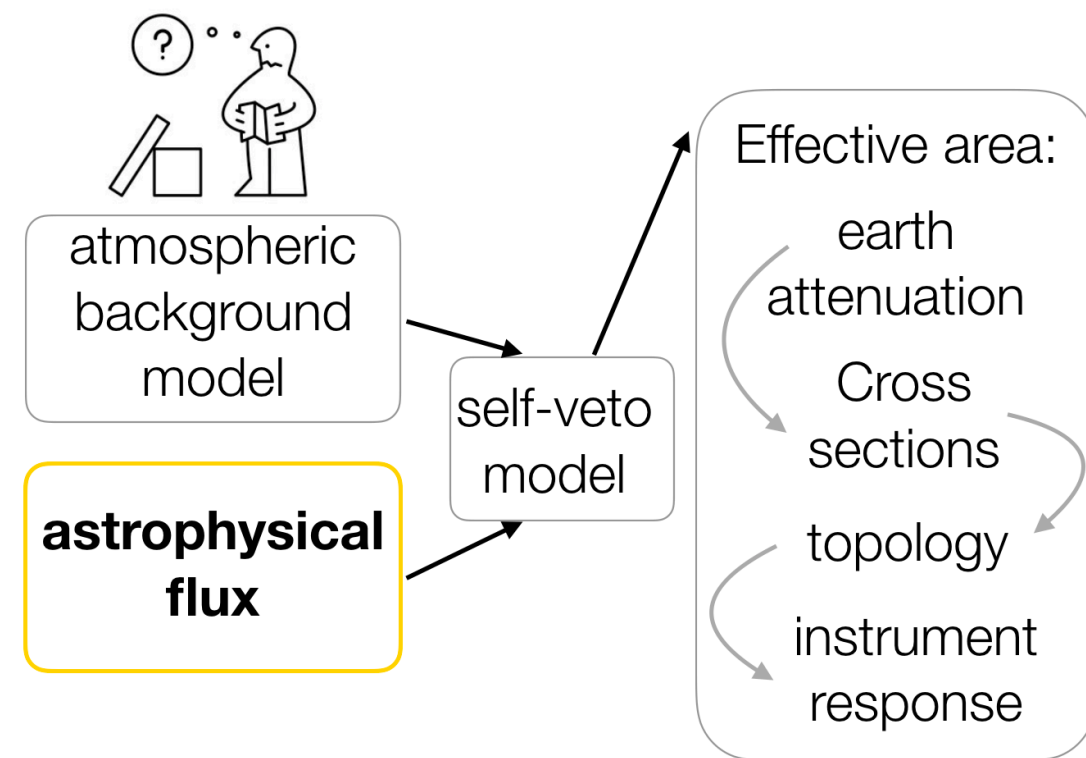
Dark matter column density seen from Earth



Dark matter column density seen from Earth

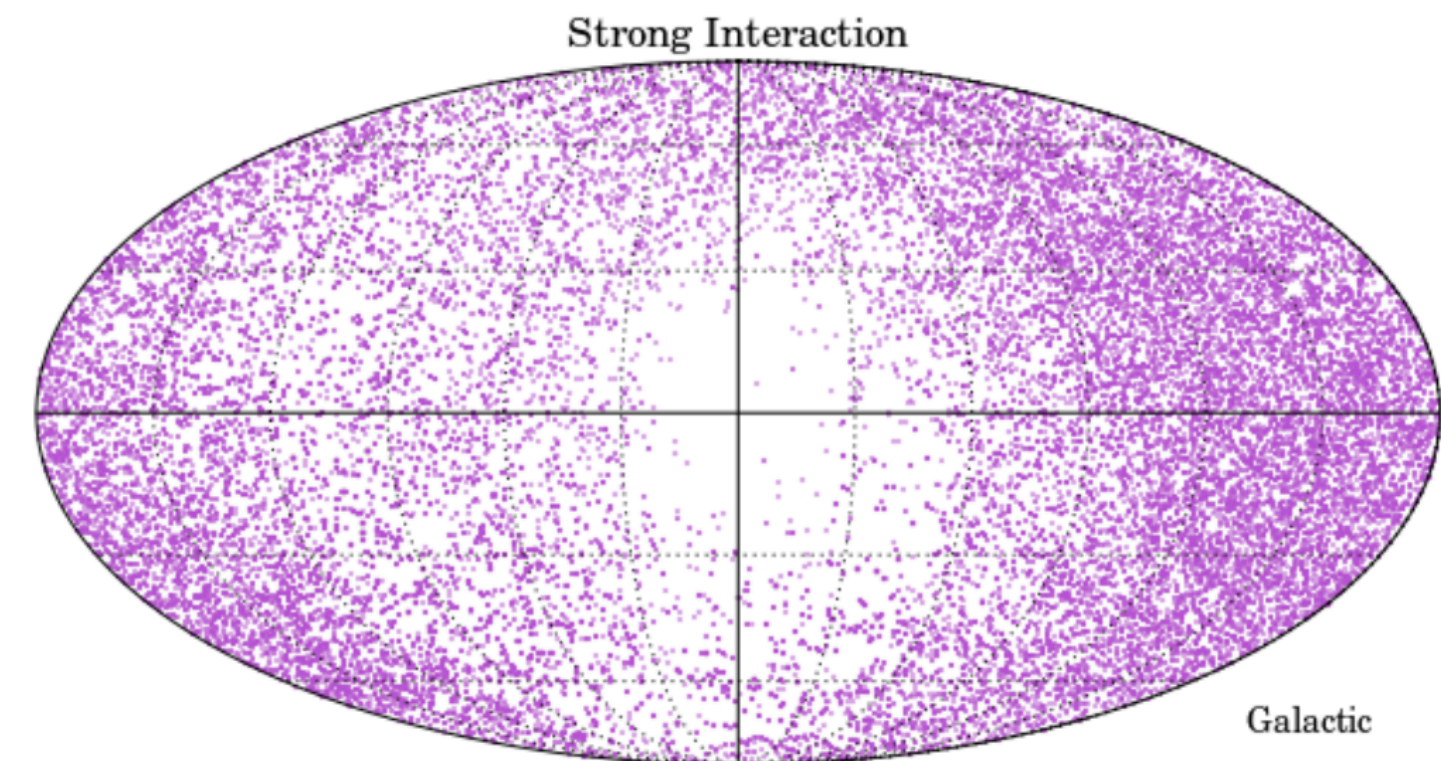
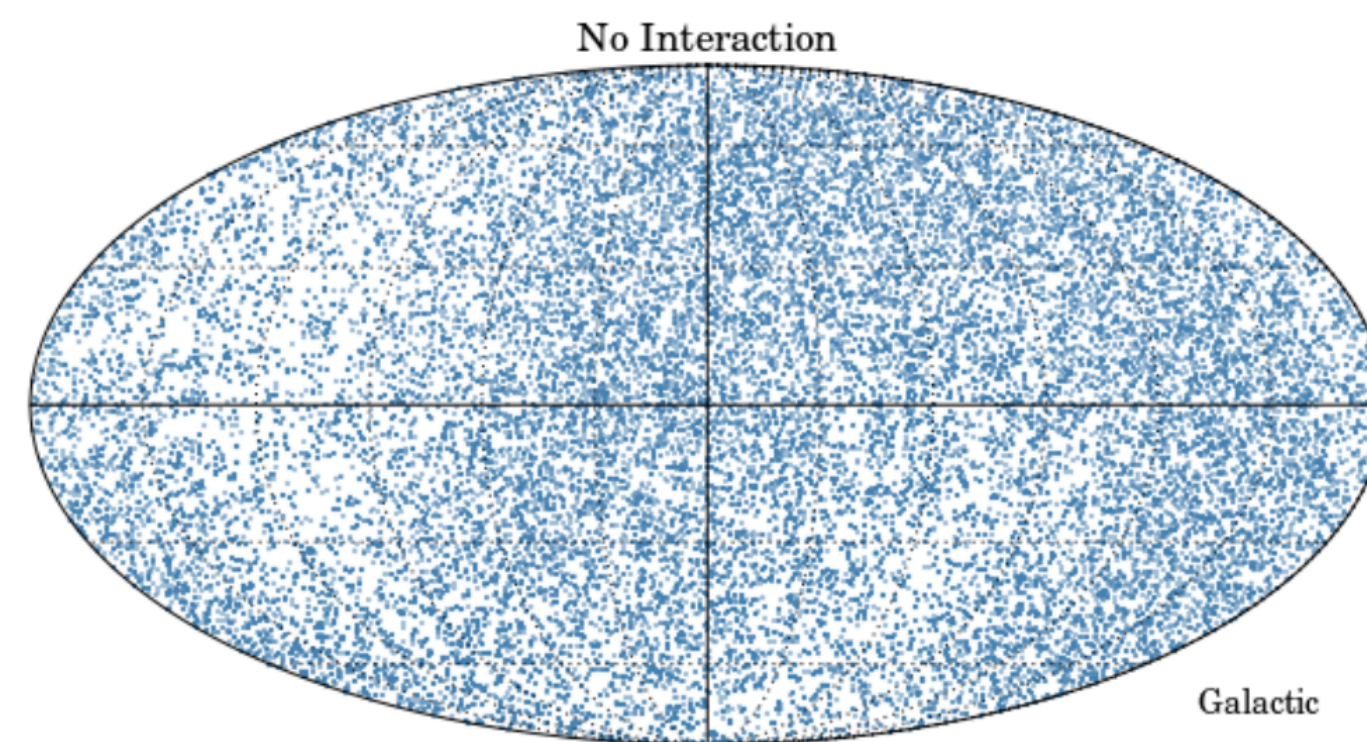


Simulation including effects of detector, Earth

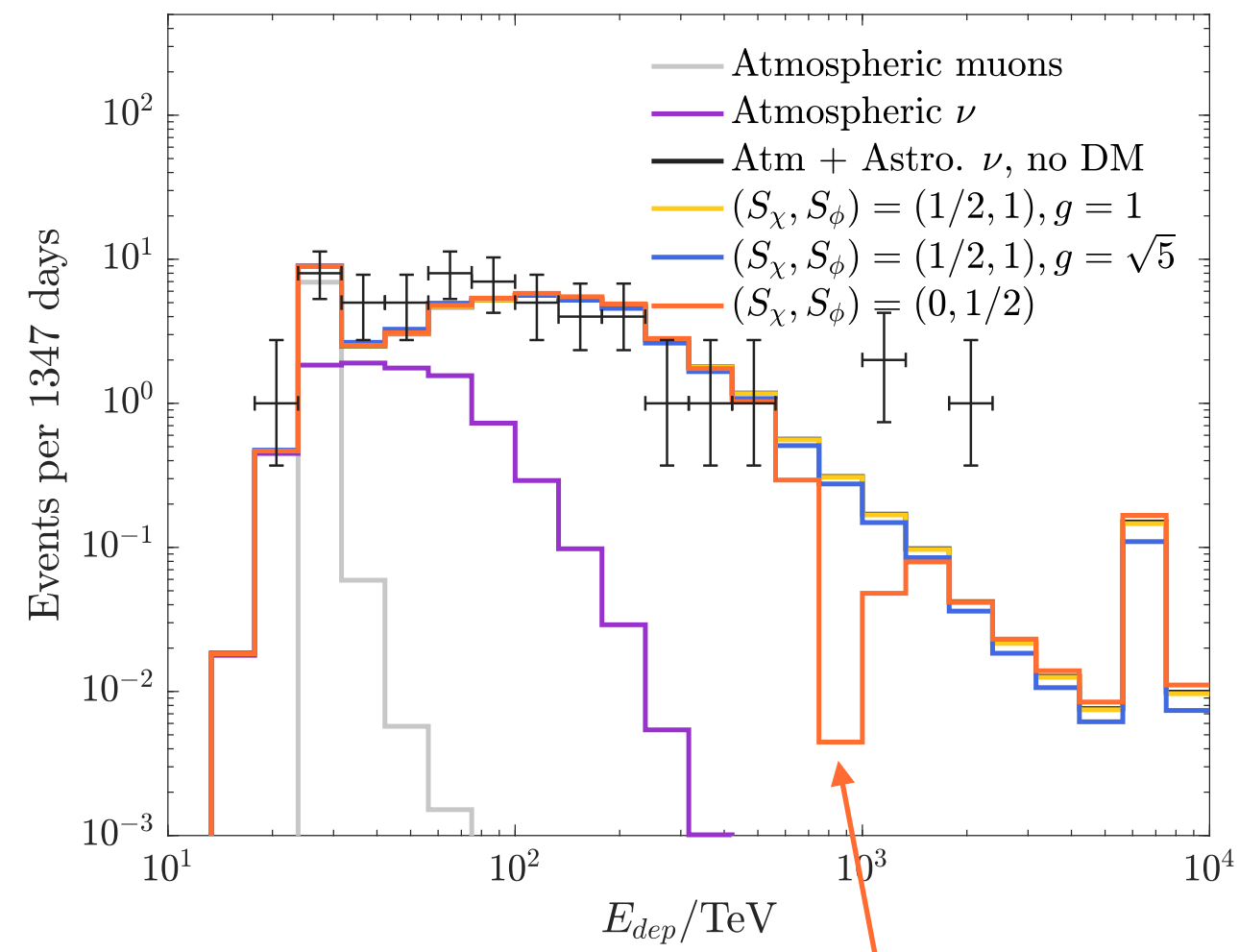


no interaction

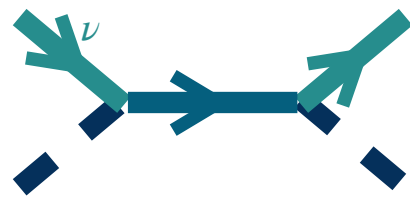
strong interaction



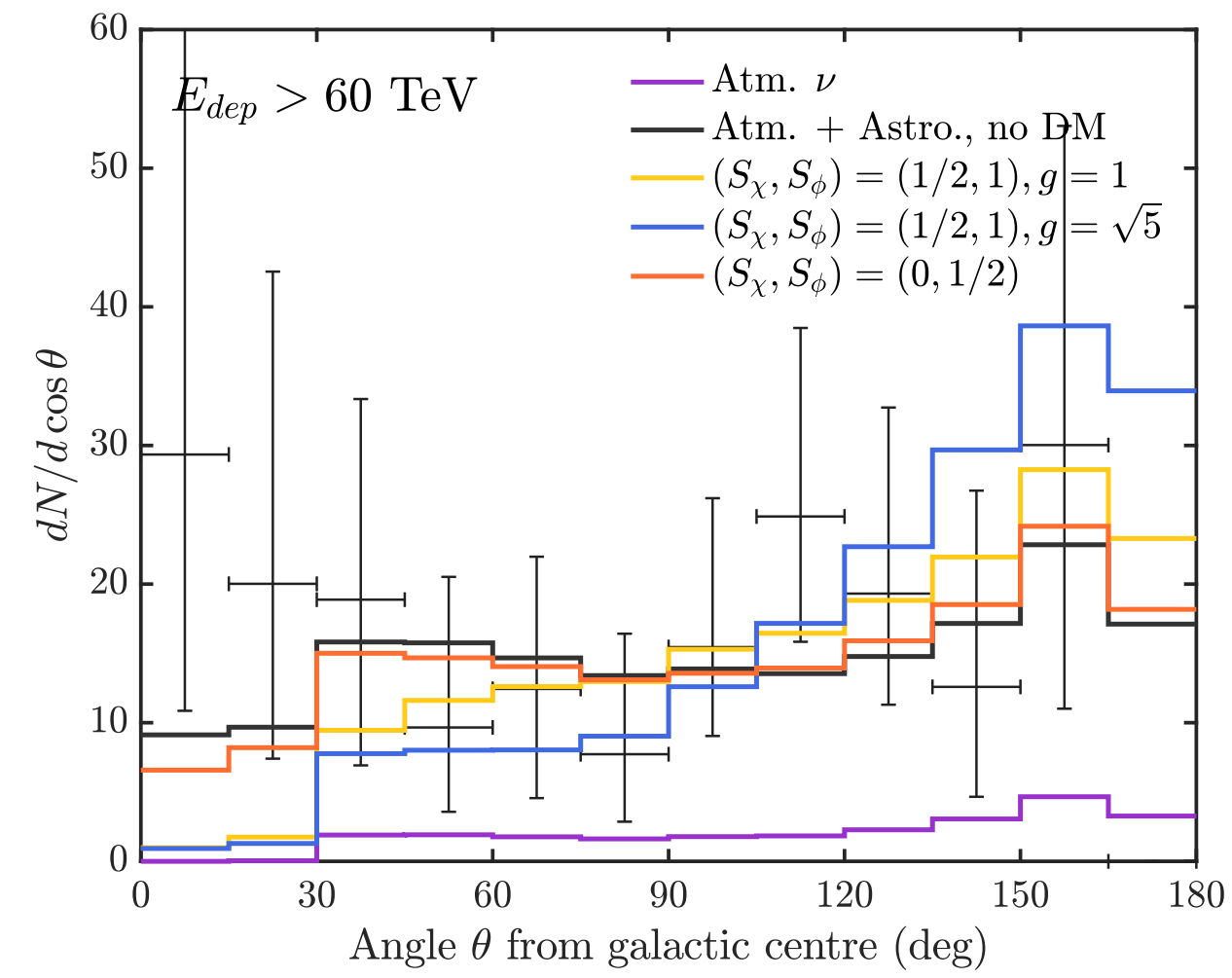
Energy



Resonance @ 810 TeV

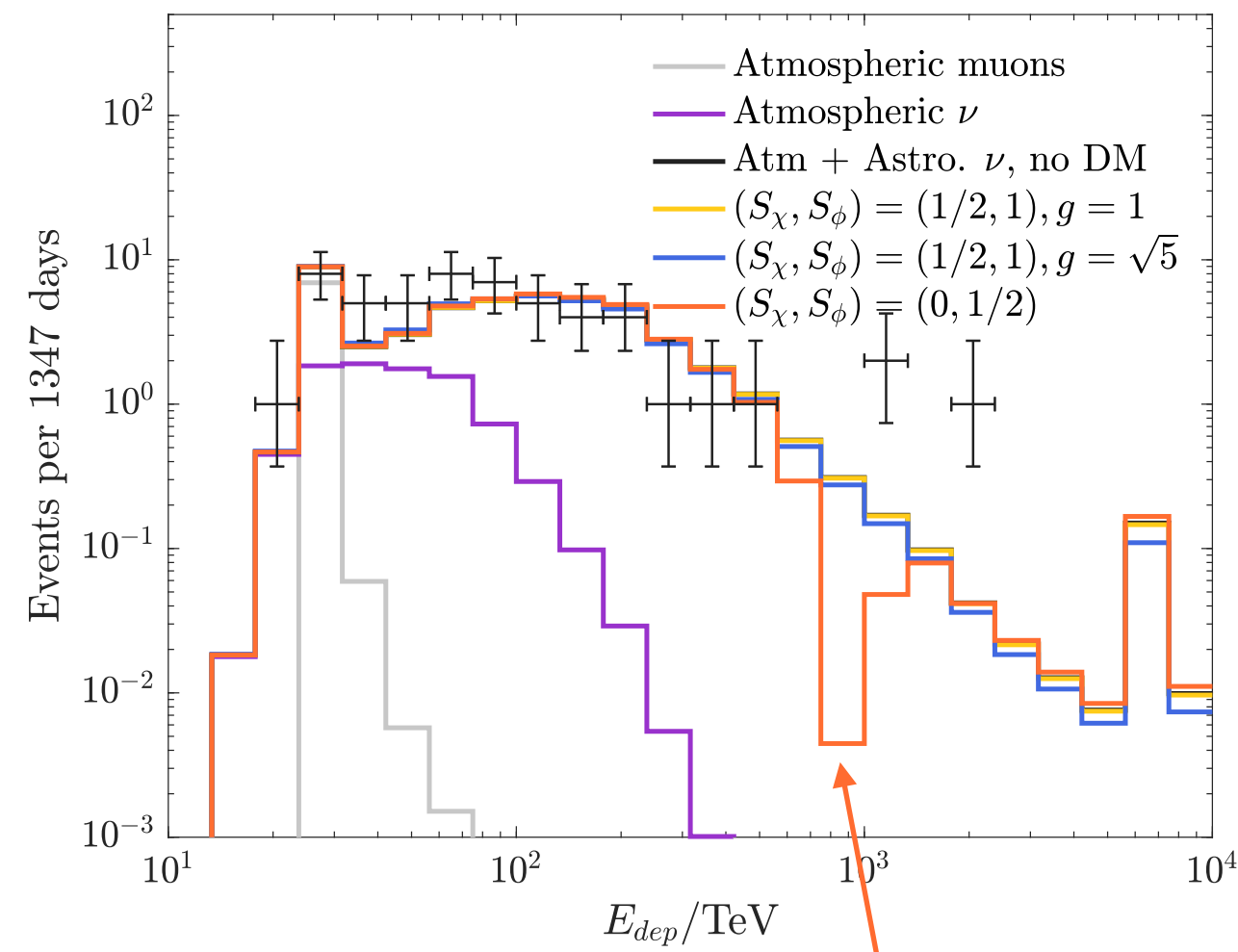


Angle from galactic centre

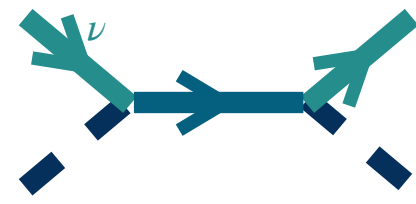


IceCube HESE events

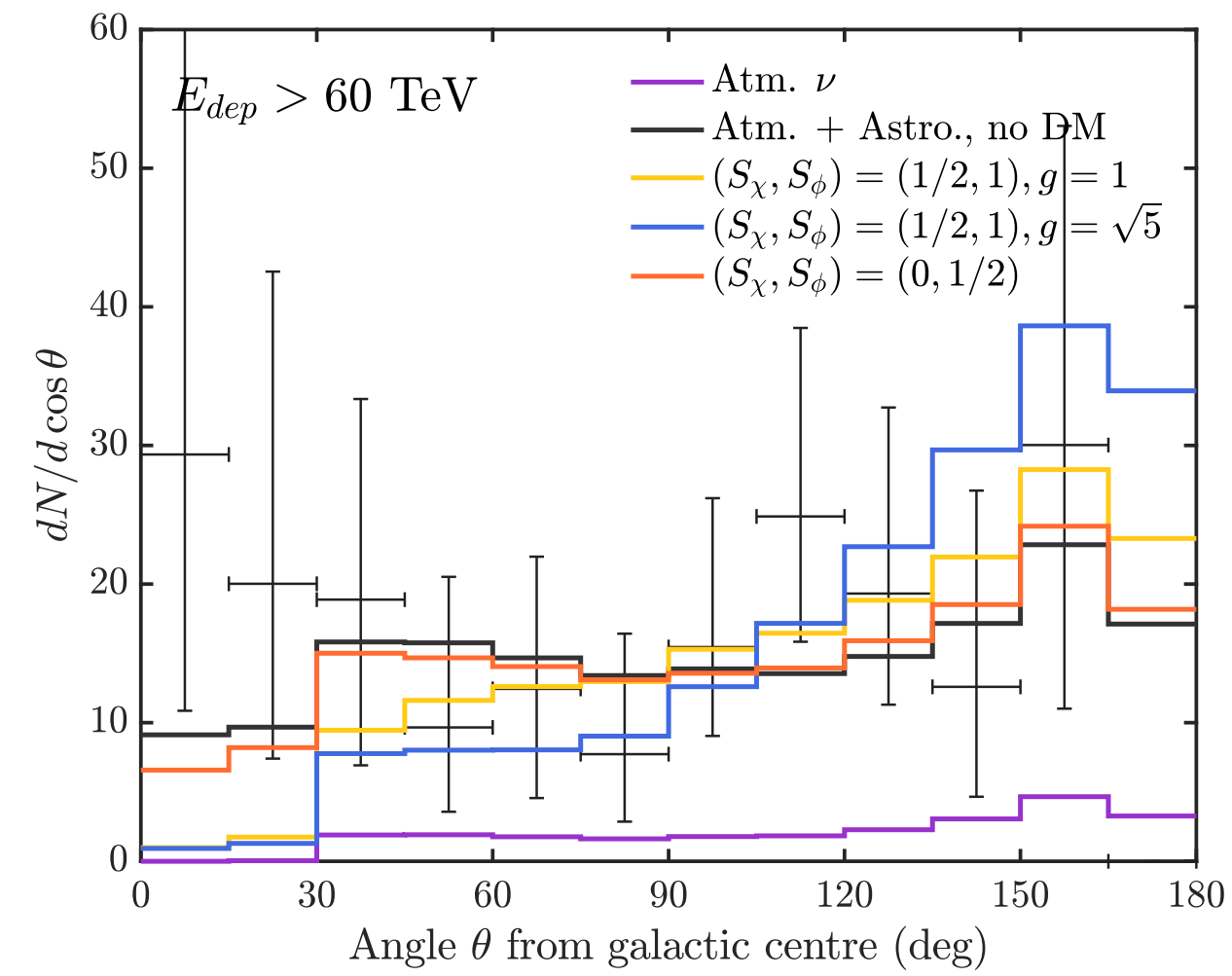
Energy



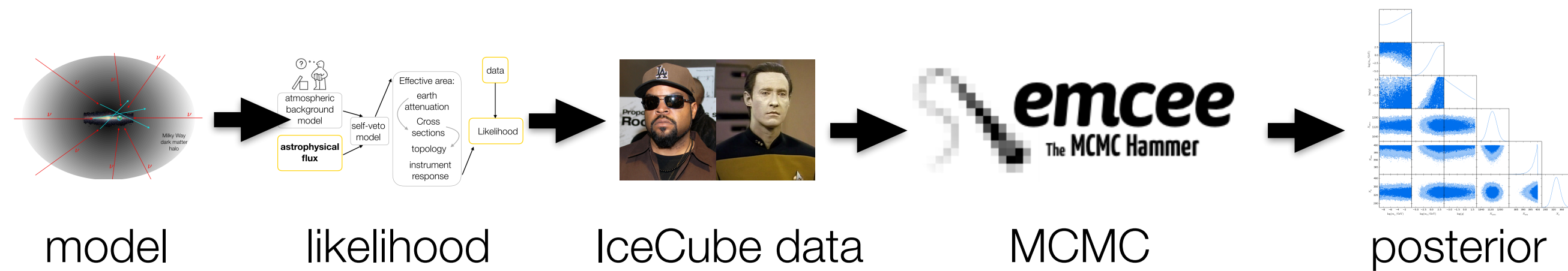
Resonance @ 810 TeV



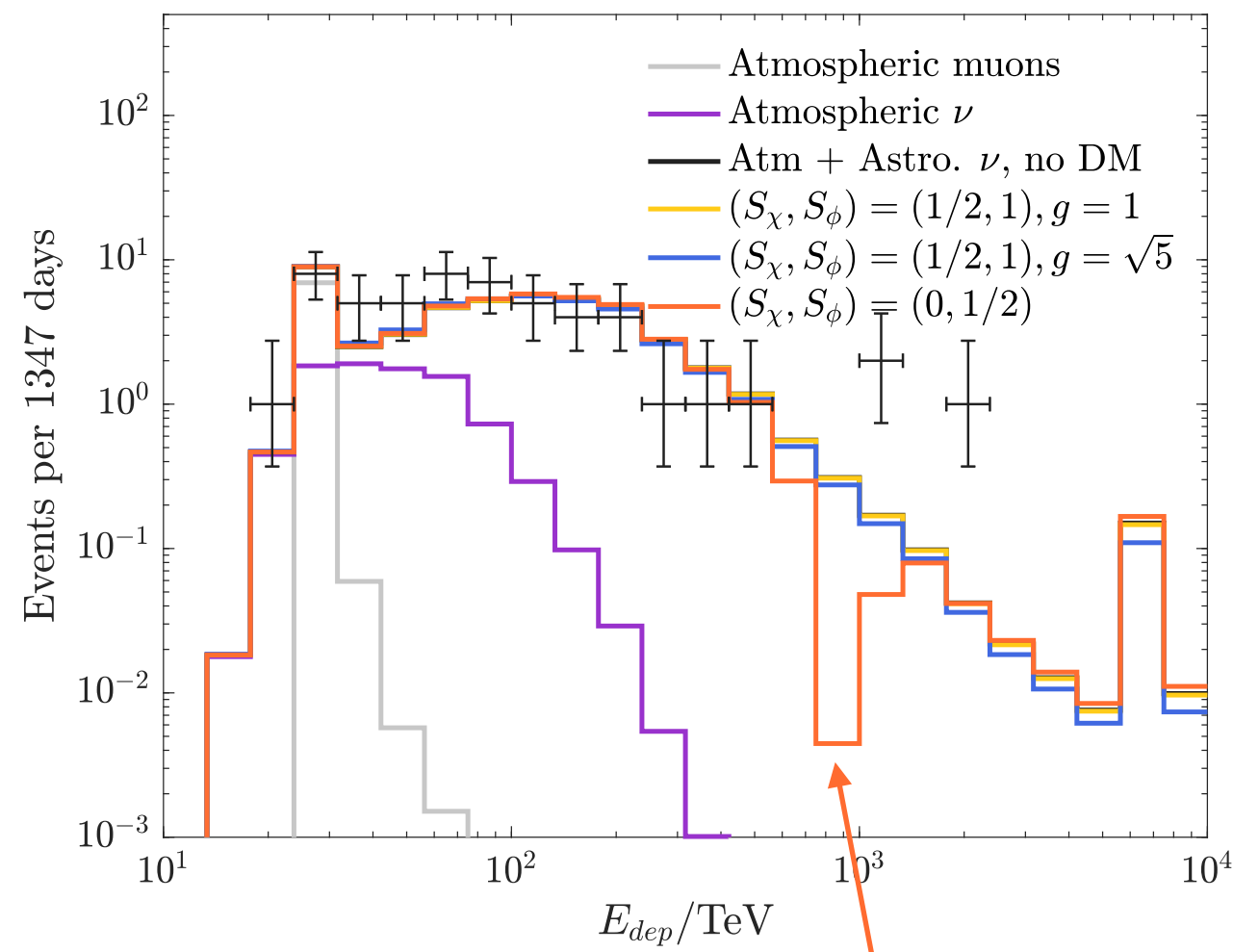
Angle from galactic centre



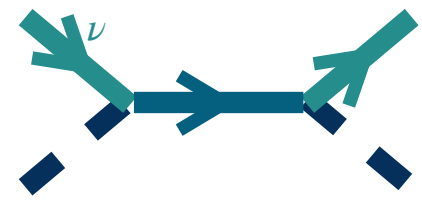
IceCube HESE events



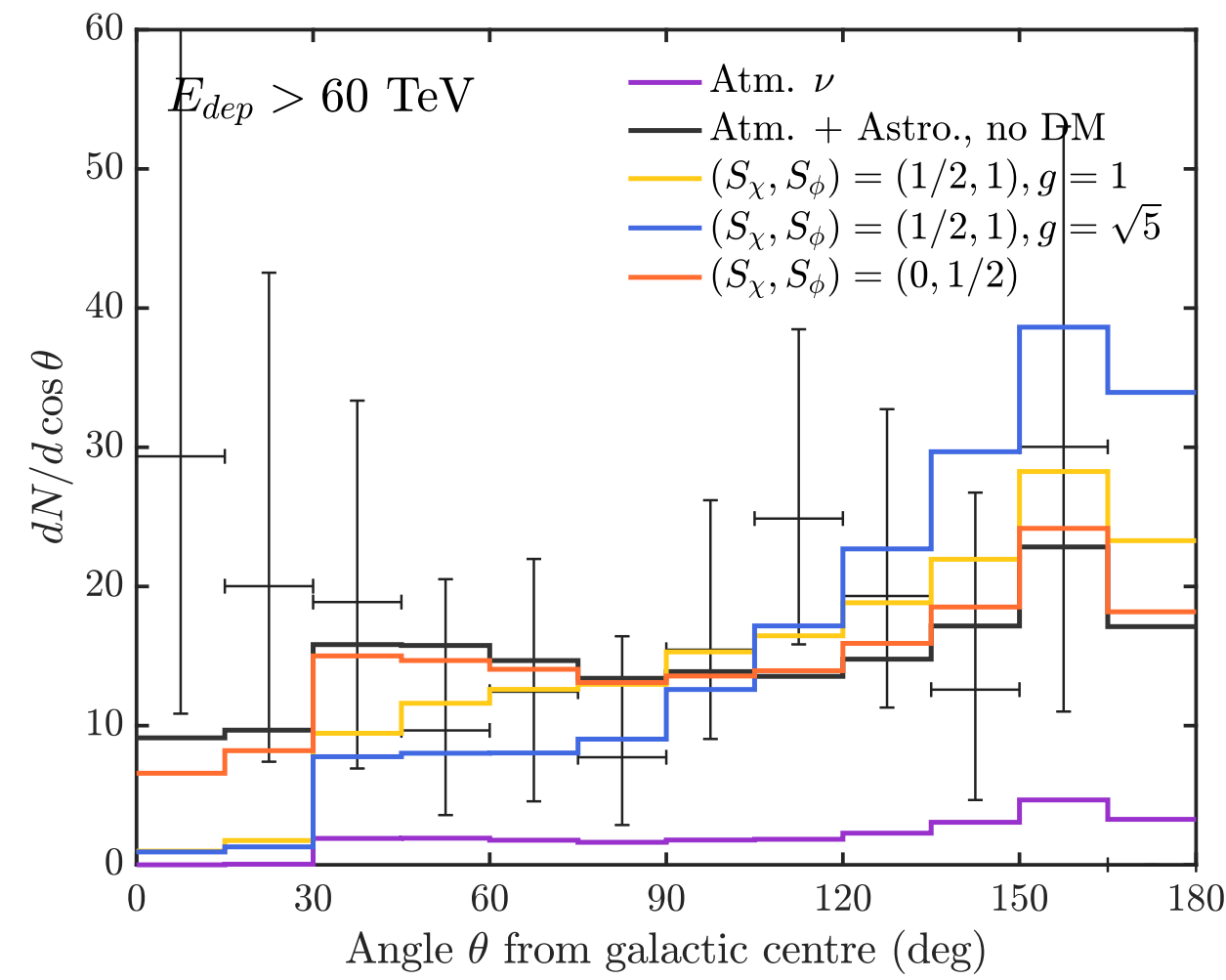
Energy



Resonance @ 810 TeV

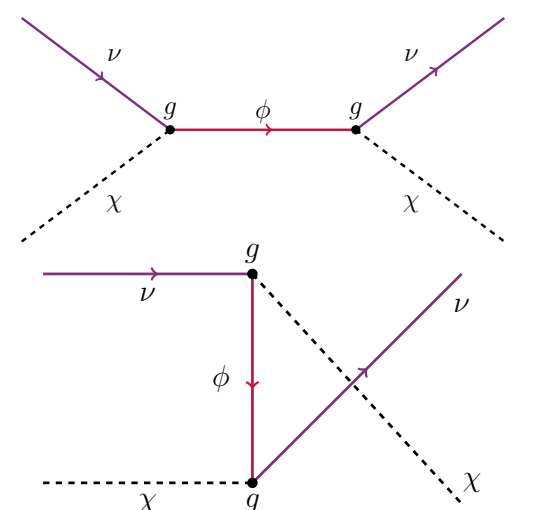
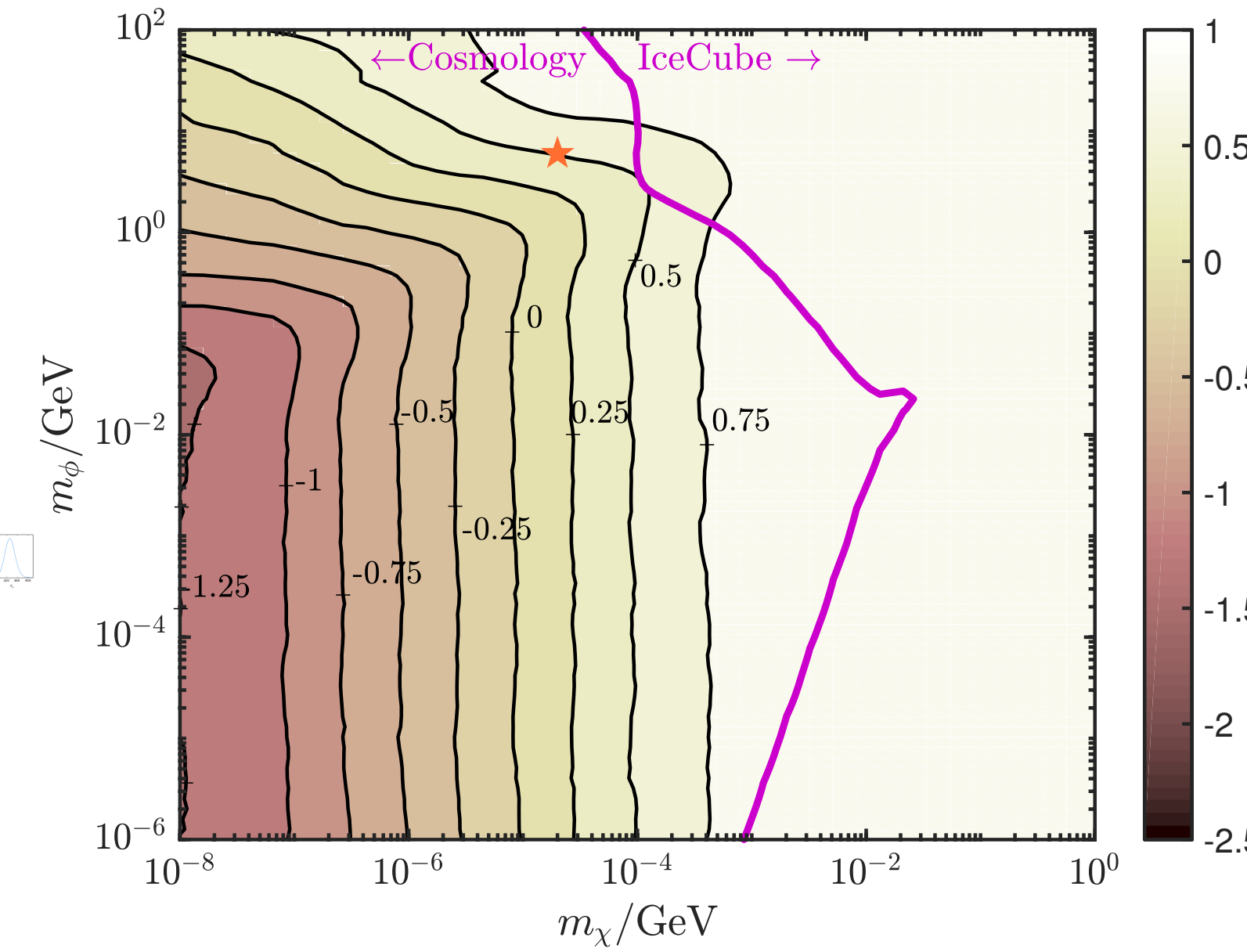
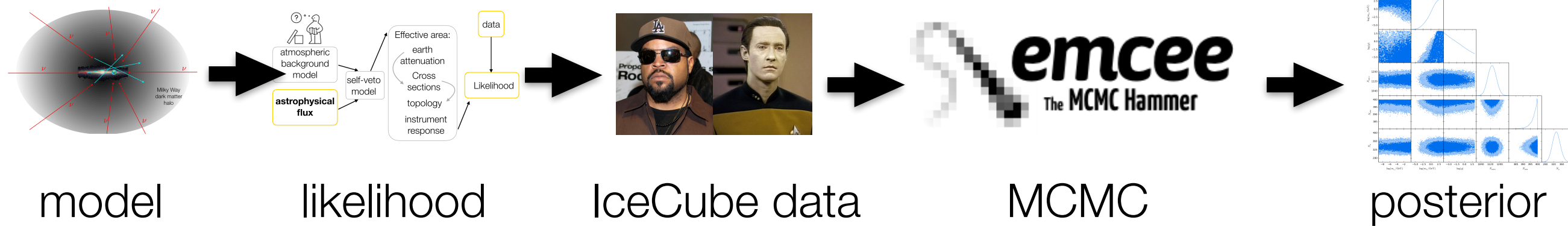
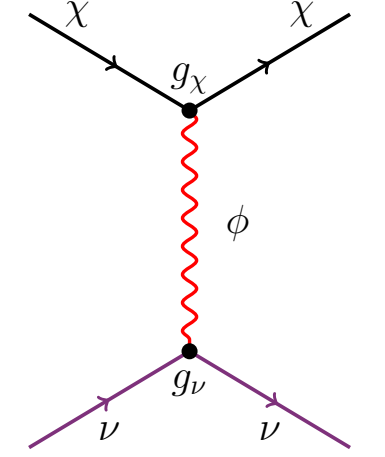
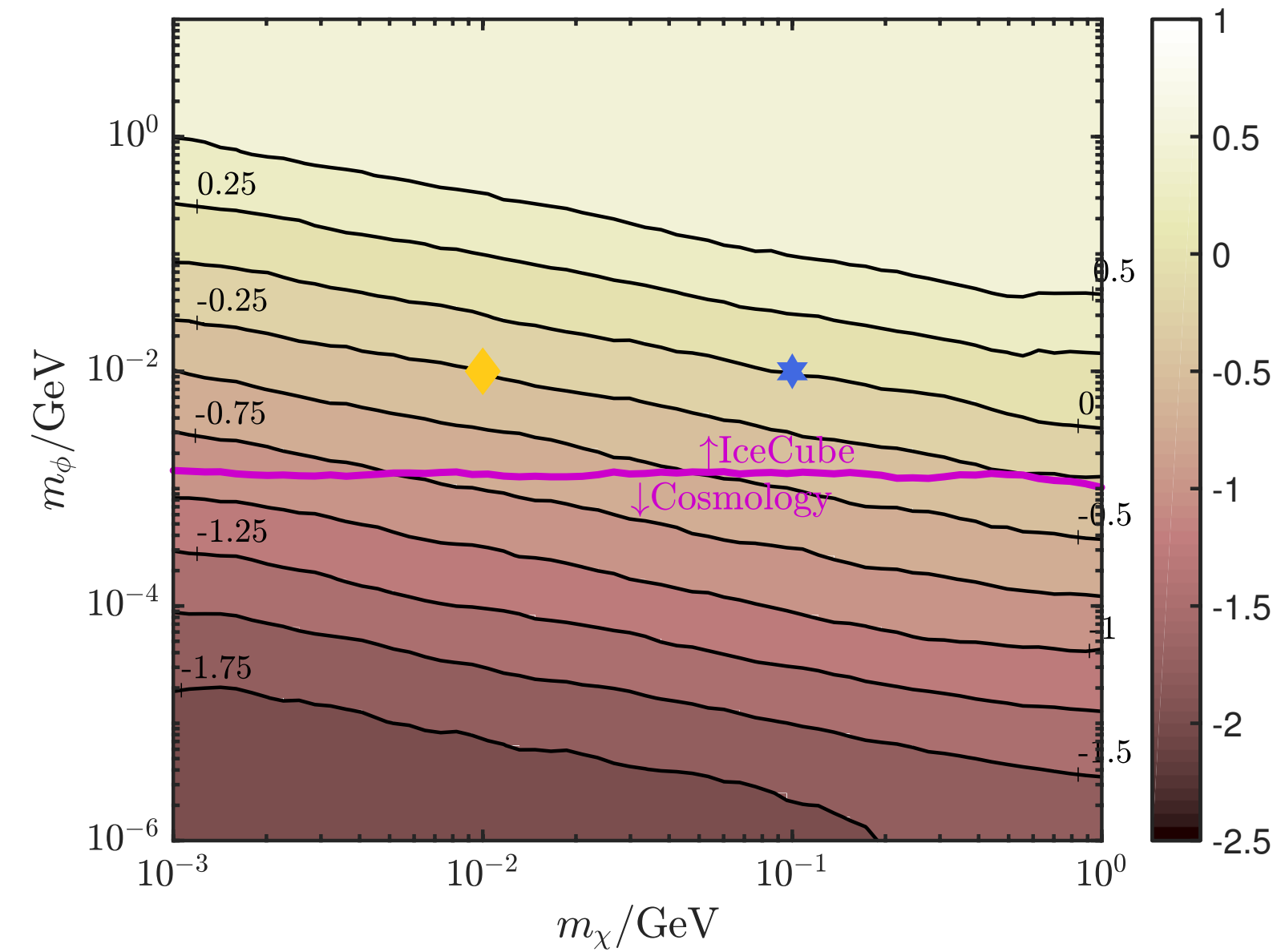


Angle from galactic centre

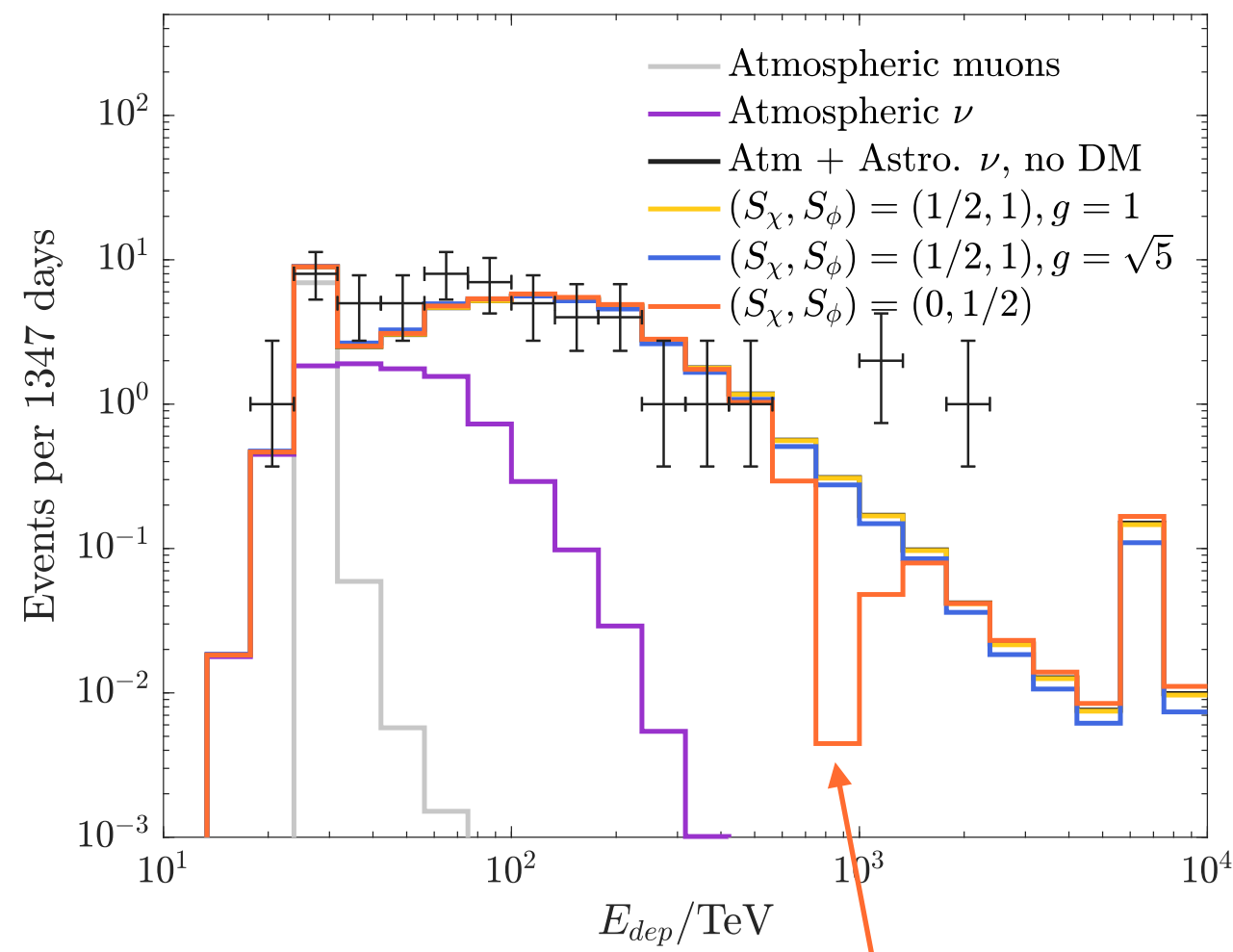


IceCube HESE events

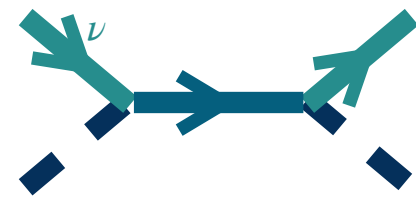
MeV



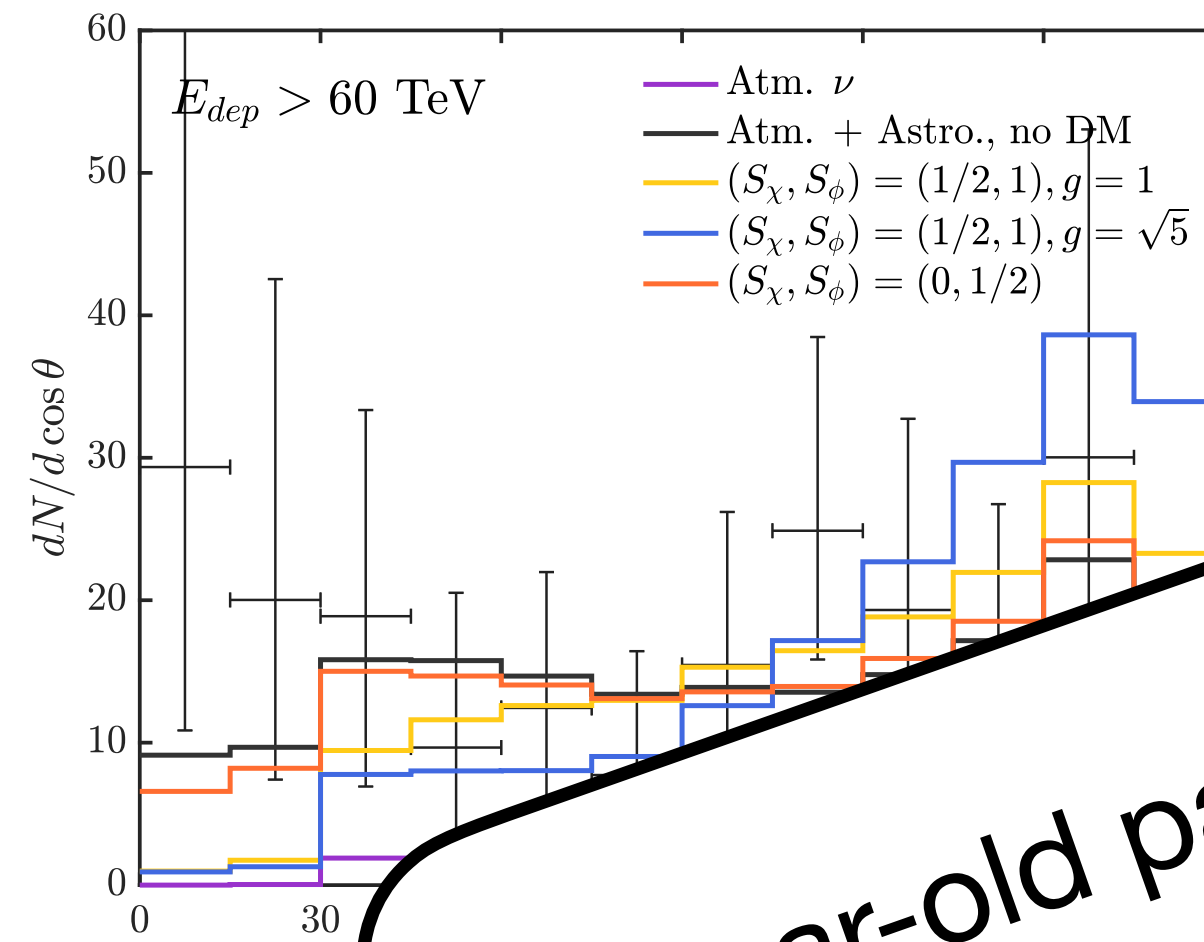
Energy



Resonance @ 810 TeV



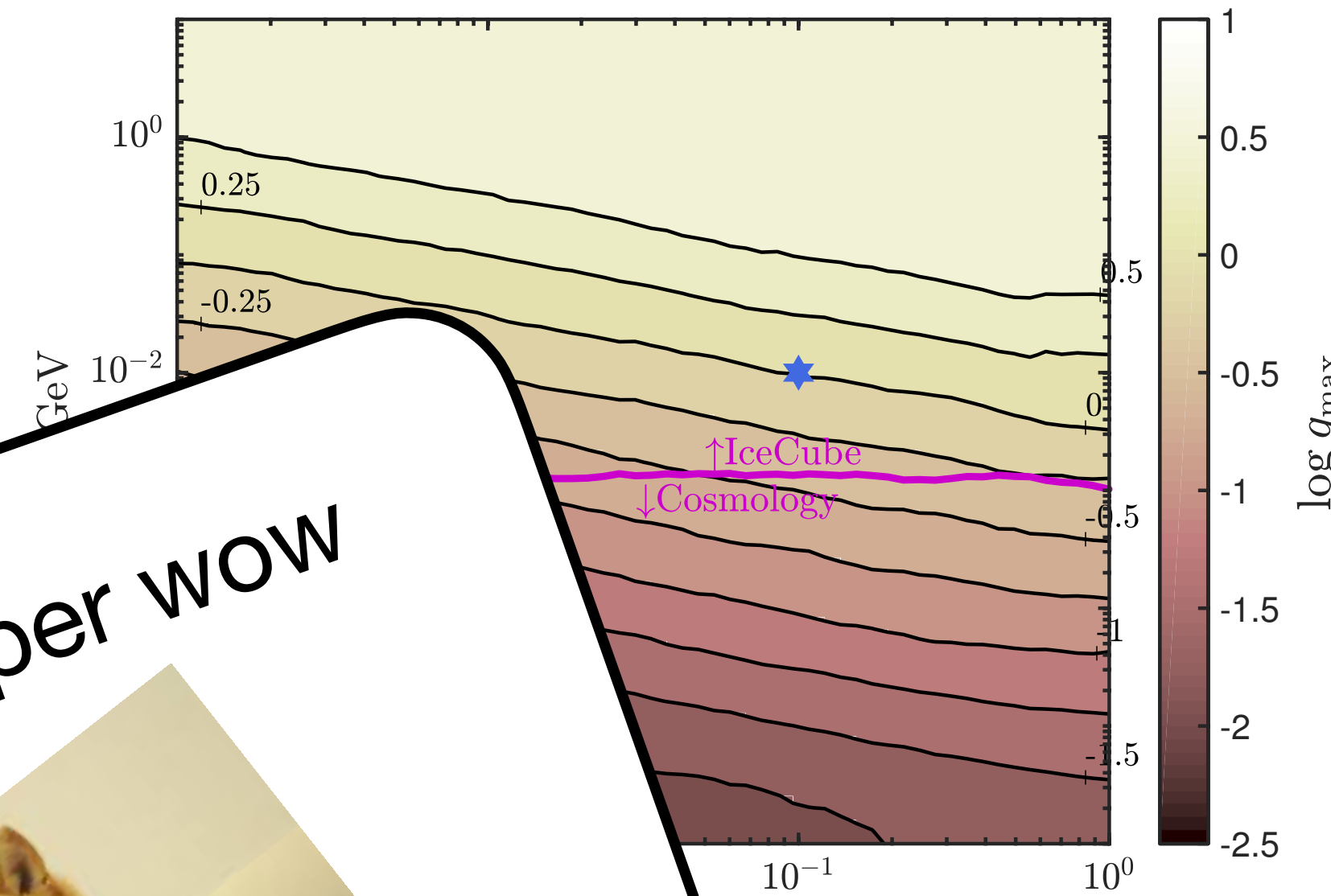
Angle from galactic centre



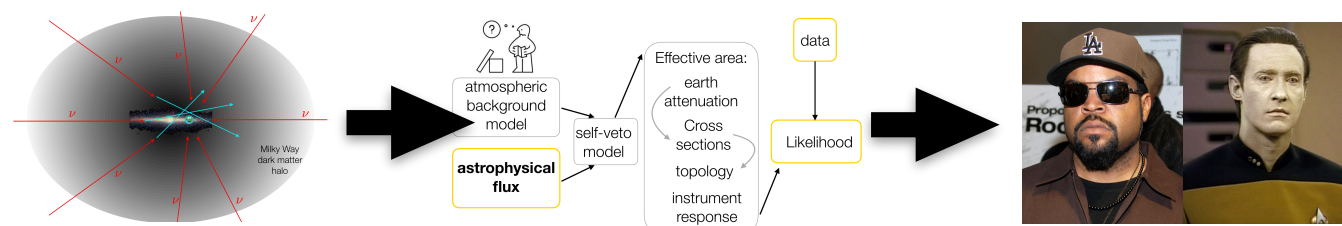
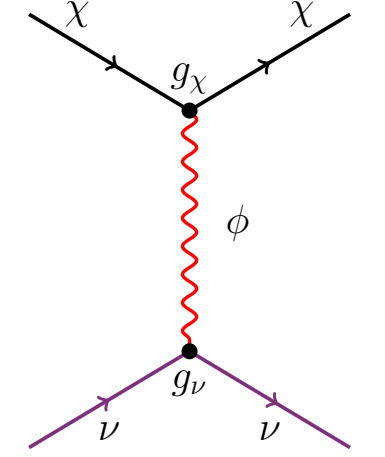
IceCube

MeV

GeV



7 year-old paper wow



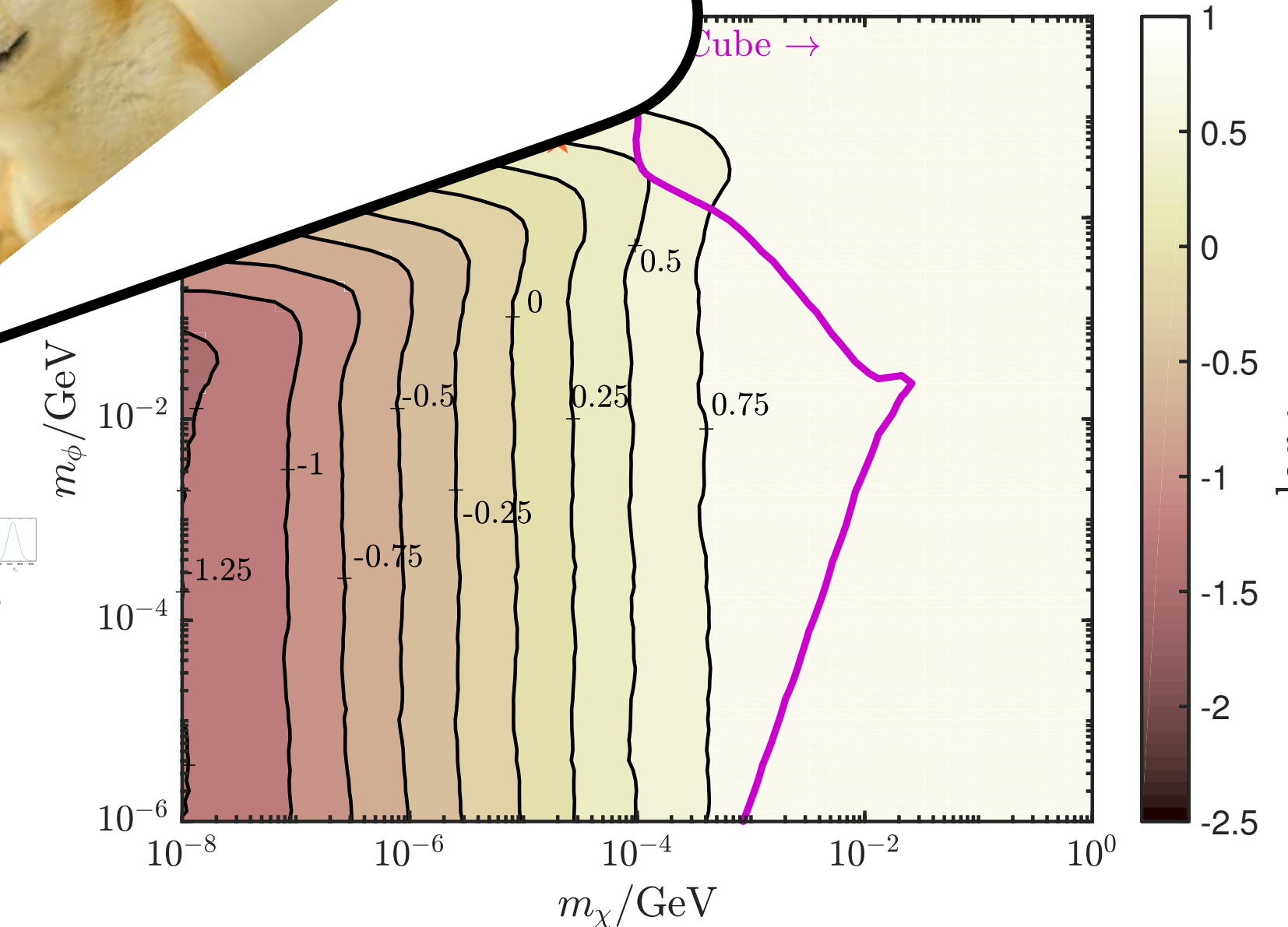
model

likelihood

IceCube data

MCMC

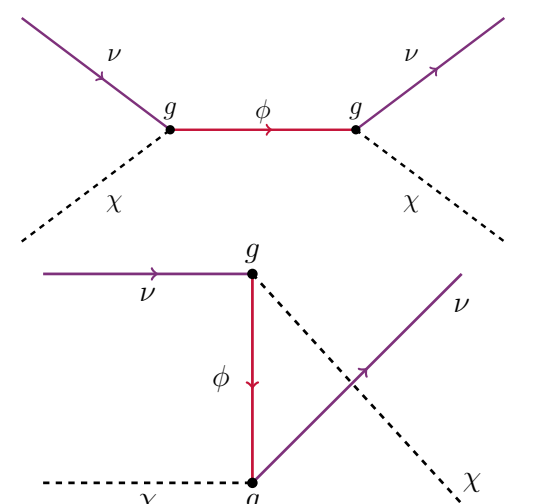
posterior



keV

MeV

GeV



Rēprobate

NOW A MAJOR

IceCube analysis

the lost world of the film and tv tie-in novel

λ

∇

Diyaselis Delgado @ Harvard

also Qinrui

Profile Binned Likelihood: Signal Subtraction Approach

$$\mathcal{L} = \sum_i^{\text{bins}} \log \text{Pois}(\vec{x}_i, \mu_i(\vec{\theta}))$$
$$\text{TS} = -2\Delta\text{LLH}$$

$$\vec{x} = \langle E, \text{RA}, \text{dec} \rangle \quad \vec{\theta} = \langle g, m_\phi, m_\chi \rangle$$

$$\mu(\vec{\theta}) = \mu_{\text{astro}\nu}^{\text{signal}}(\vec{\theta}) + \mu^{\text{bkg}}$$

EXTRAGALACTIC ASTROPHYSICAL NEUTRINOS

Follows a power law spectrum

$$\phi(\vec{x}, \vec{\theta}) = E^{-\gamma} \cdot \mathcal{P}_{\text{surv}} \propto \sum c_i \hat{\phi}_i e^{\lambda_i \tau}$$

BACKGROUND: ATMOSPHERIC NEUTRINOS AND MUONS

$$\mu^{\text{bkg}} = \mu_{\text{atm}\nu} + \mu_{\text{muon}}$$

Scrambling over RA

$$\mu^{\text{bkg}} = \tilde{\mu}_{\text{data}} - \tilde{\mu}_{\text{astro}\nu}^{\text{signal}}(\vec{\theta})$$

Profile Binned Likelihood: Signal Subtraction Approach

$$\mathcal{L} = \sum_i^{\text{bins}} \log \text{Pois}(\vec{x}_i, \mu_i(\vec{\theta}))$$

$$\text{TS} = -2\Delta\text{LLH}$$

$$\vec{x} = \langle E, \text{RA}, \text{dec} \rangle \quad \vec{\theta} = \langle g, m_\phi, m_\chi \rangle$$

$$\mu(\vec{\theta}) = \mu_{\text{astro}\nu}^{\text{ex-gc}}(\vec{\theta}) + \mu_{\text{astro}\nu}^{\text{gc}}(\vec{\theta}) + \mu^{\text{bkg}}$$

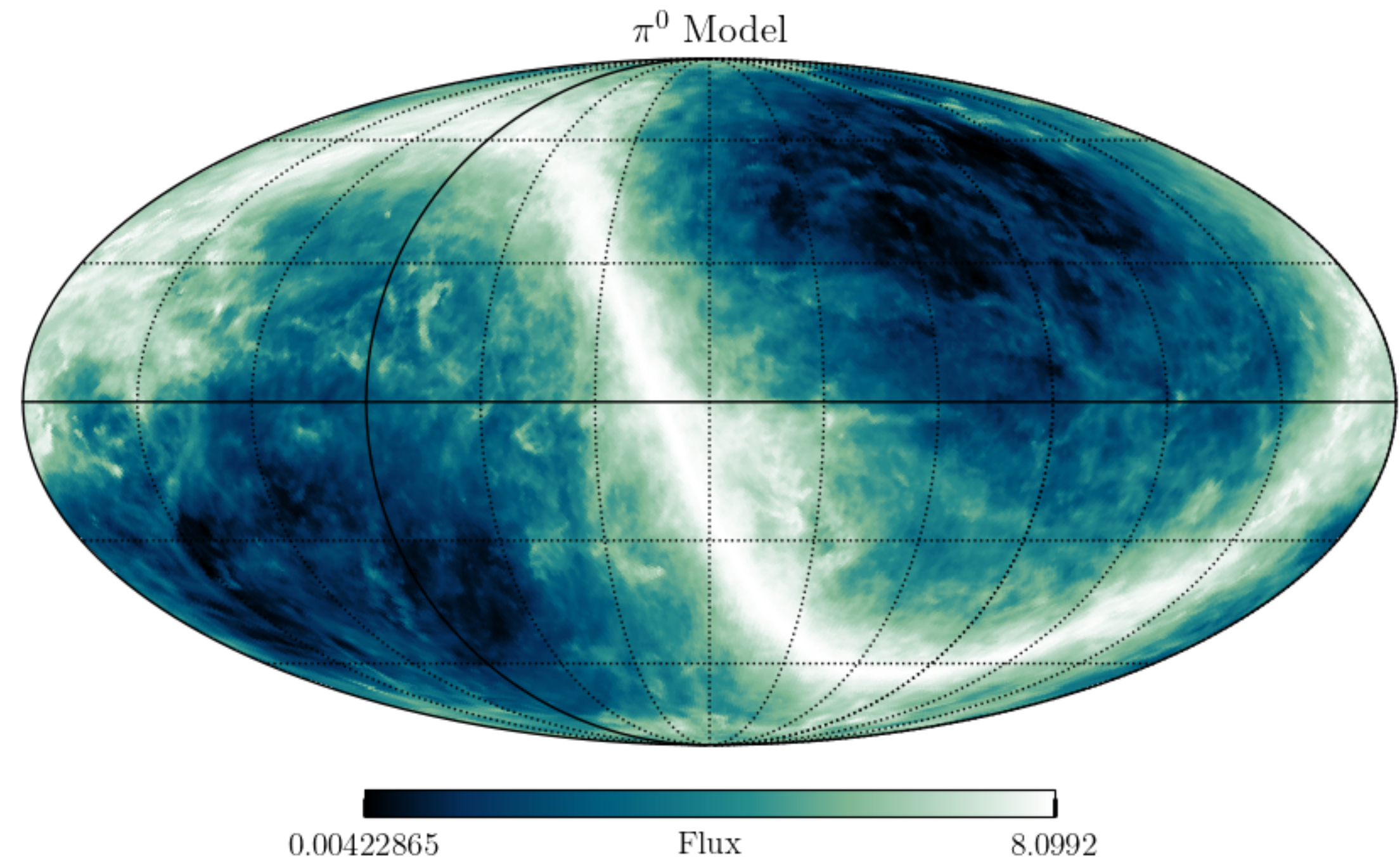
GALACTIC ASTROPHYSICAL NEUTRINOS

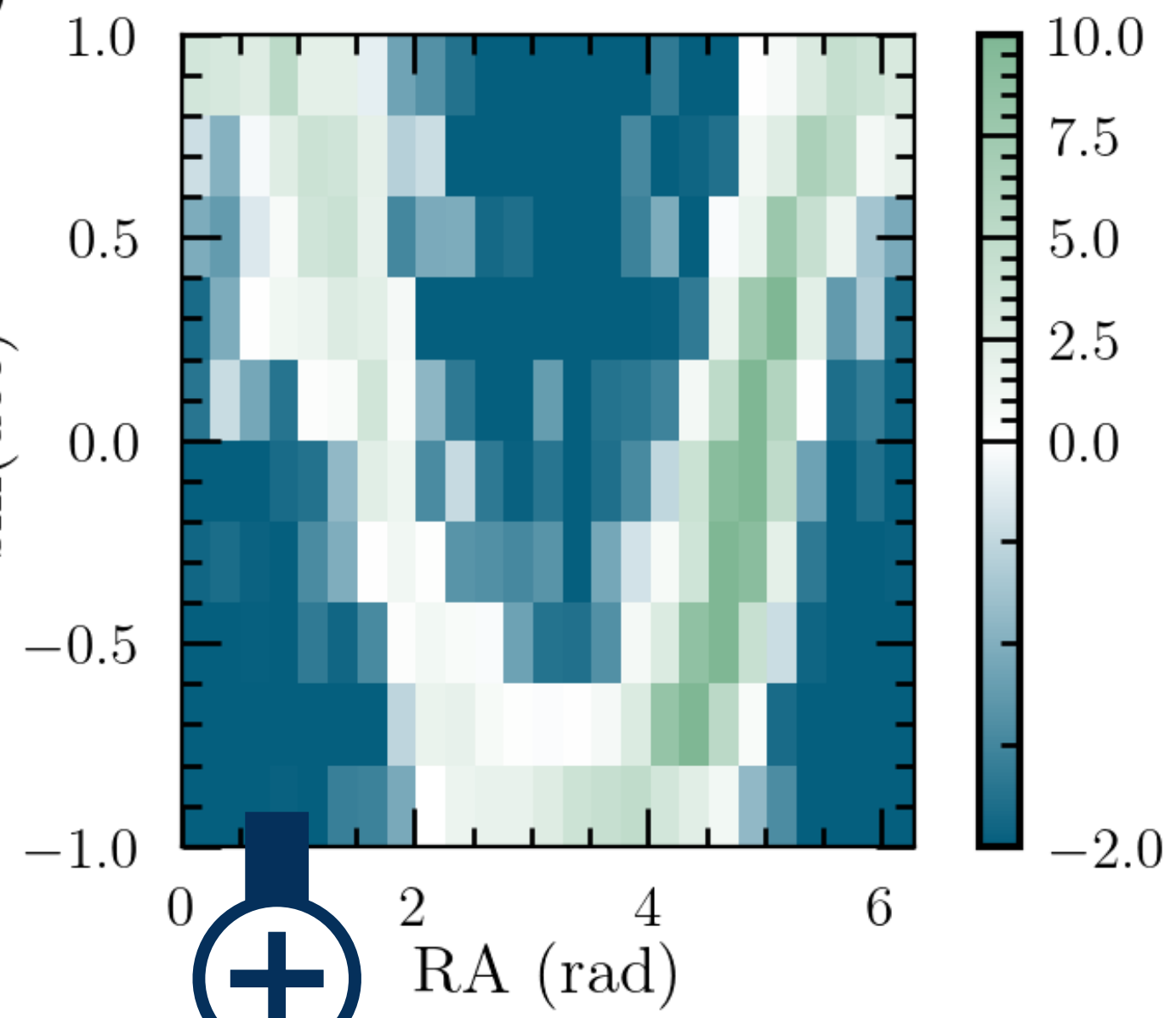
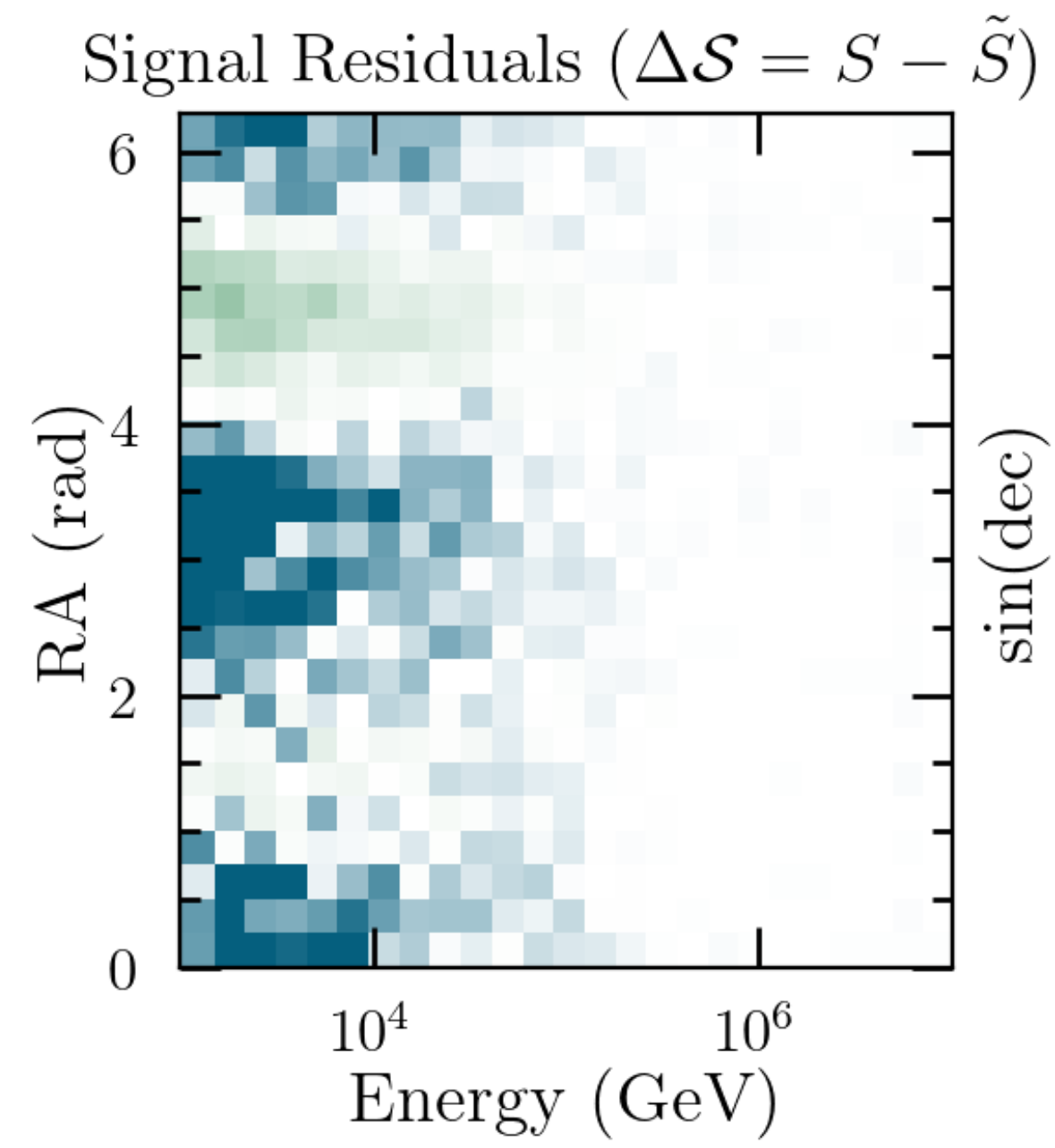
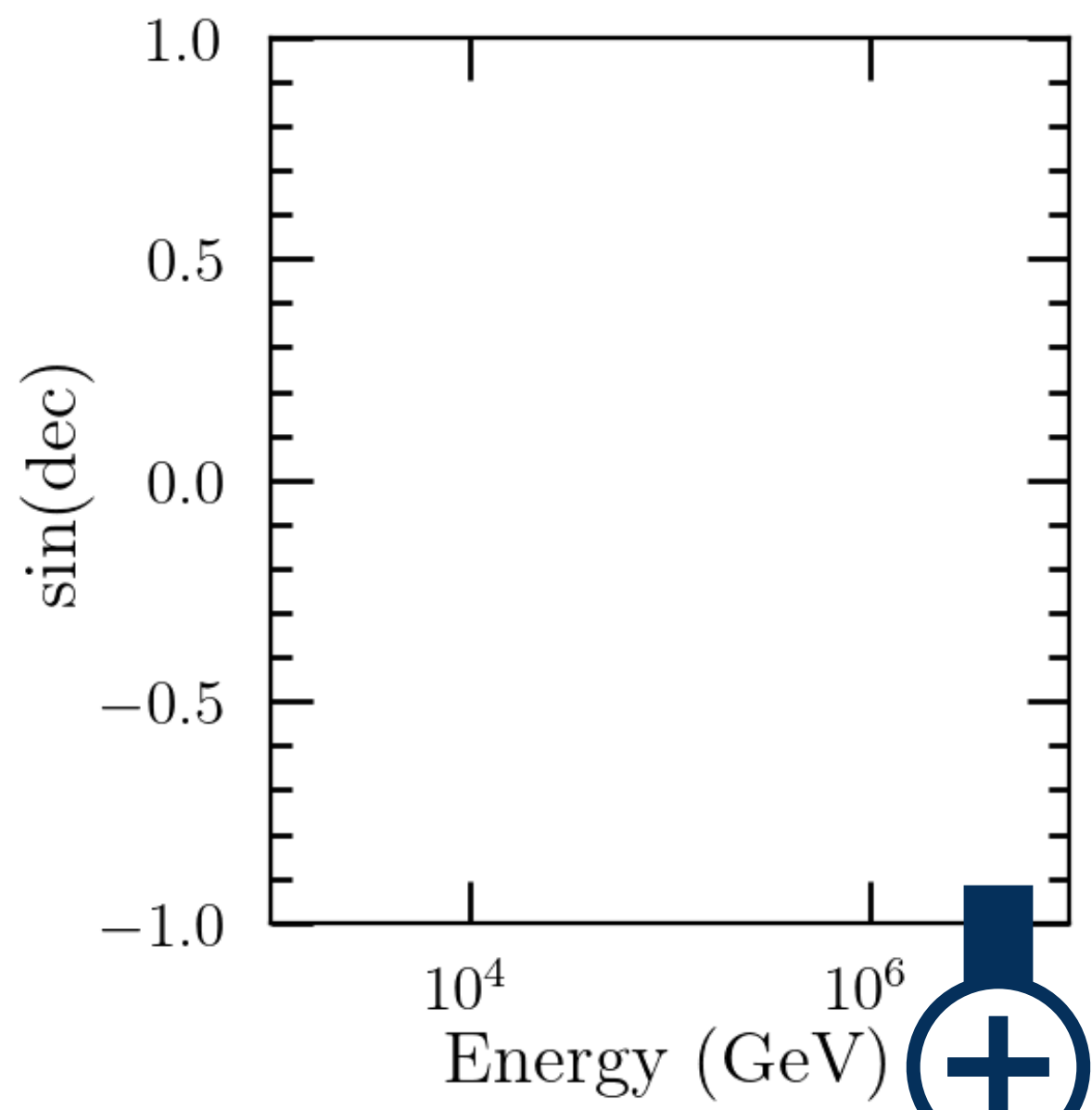
From model template of GC ν emission

$$\phi(\vec{x}, \theta) = \int_{l.o.s.} ds n_\nu(s(\vec{x})) \phi(\vec{x}, \vec{\theta}; s(\vec{x}))$$

n_ν SPATIAL DISTRIBUTION

ϕ ENERGY DISTRIBUTION





Scalar χ - Fermion ϕ

$g = 1$

$m_\phi = 1 \text{ GeV}$

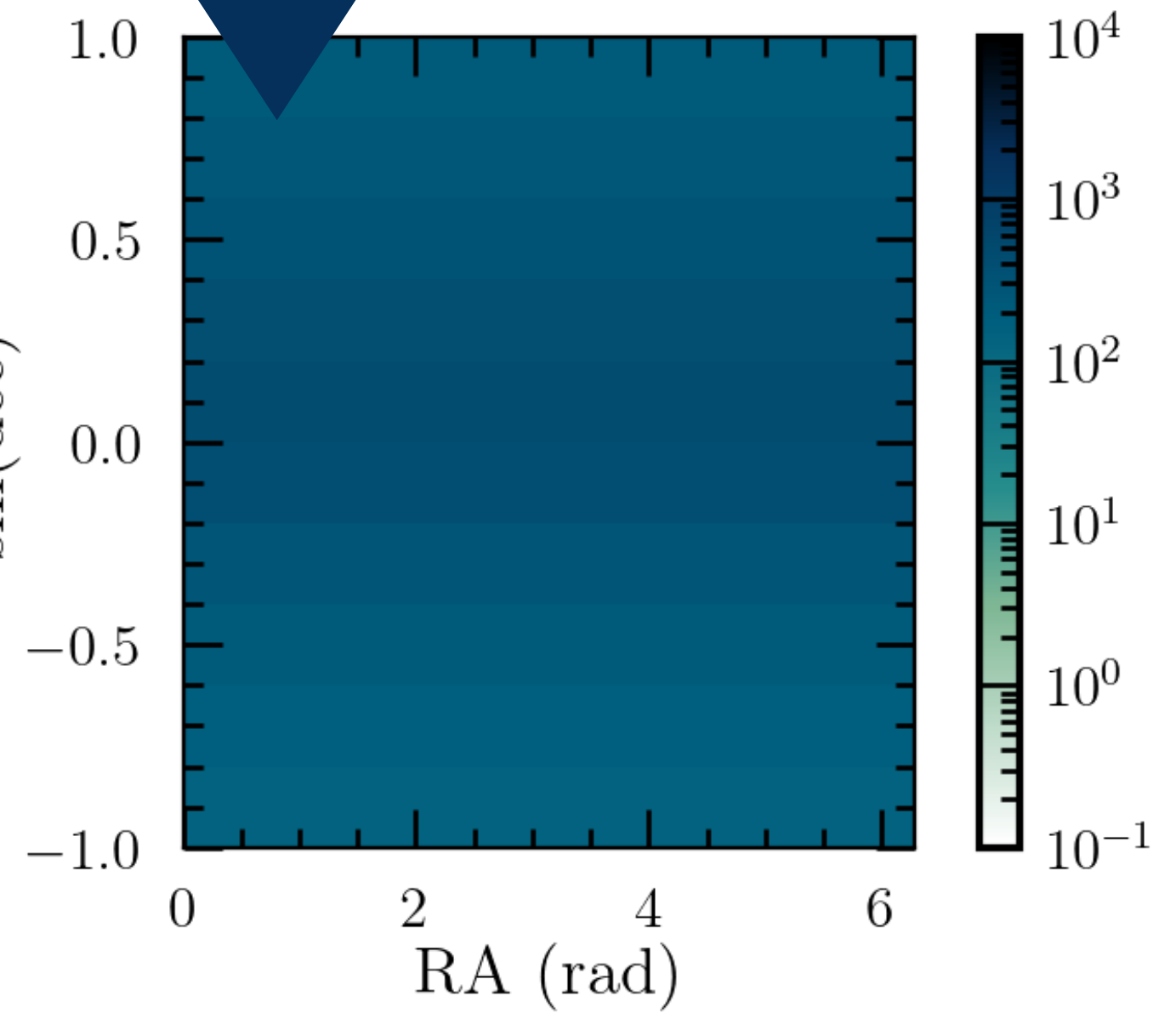
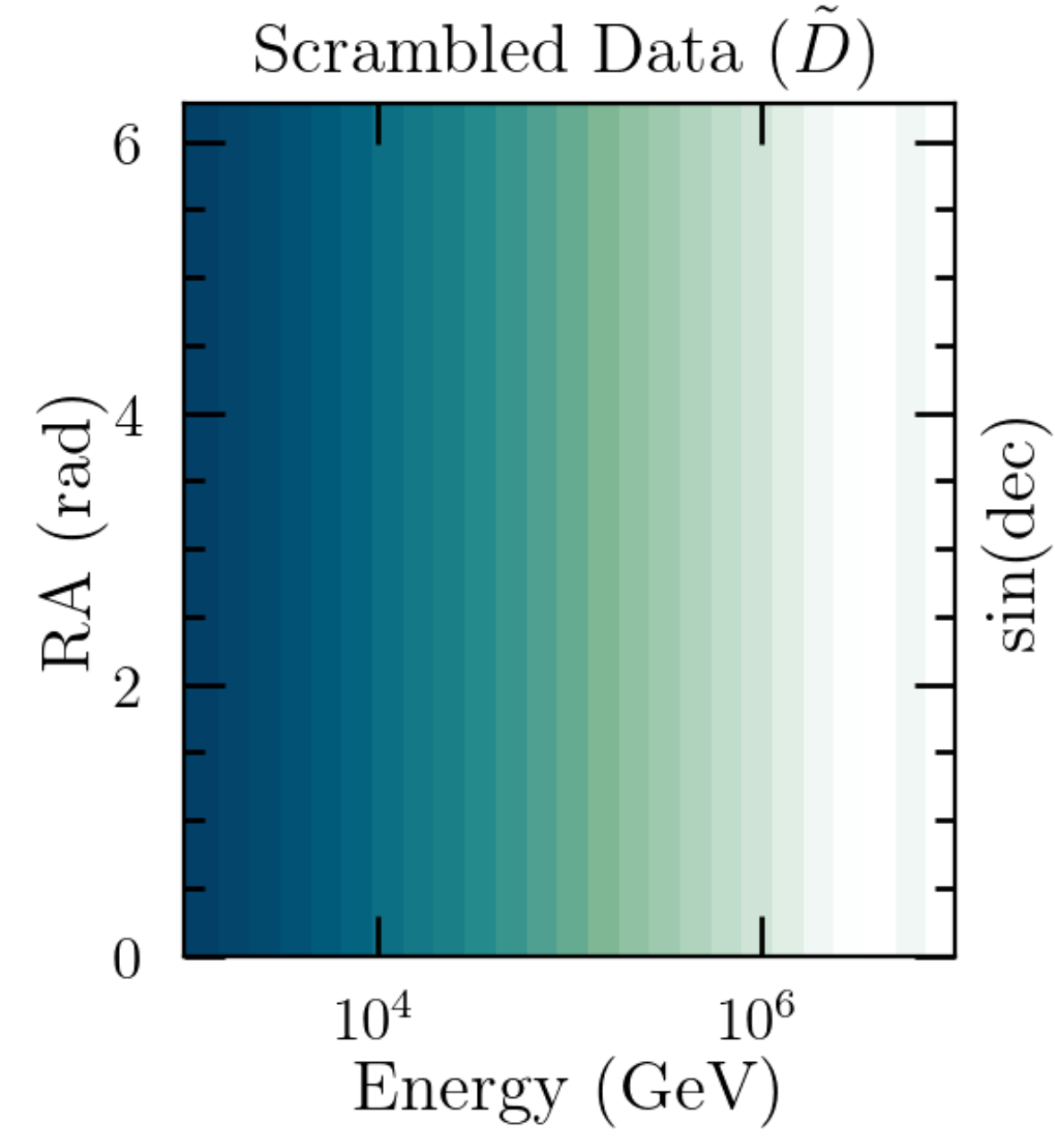
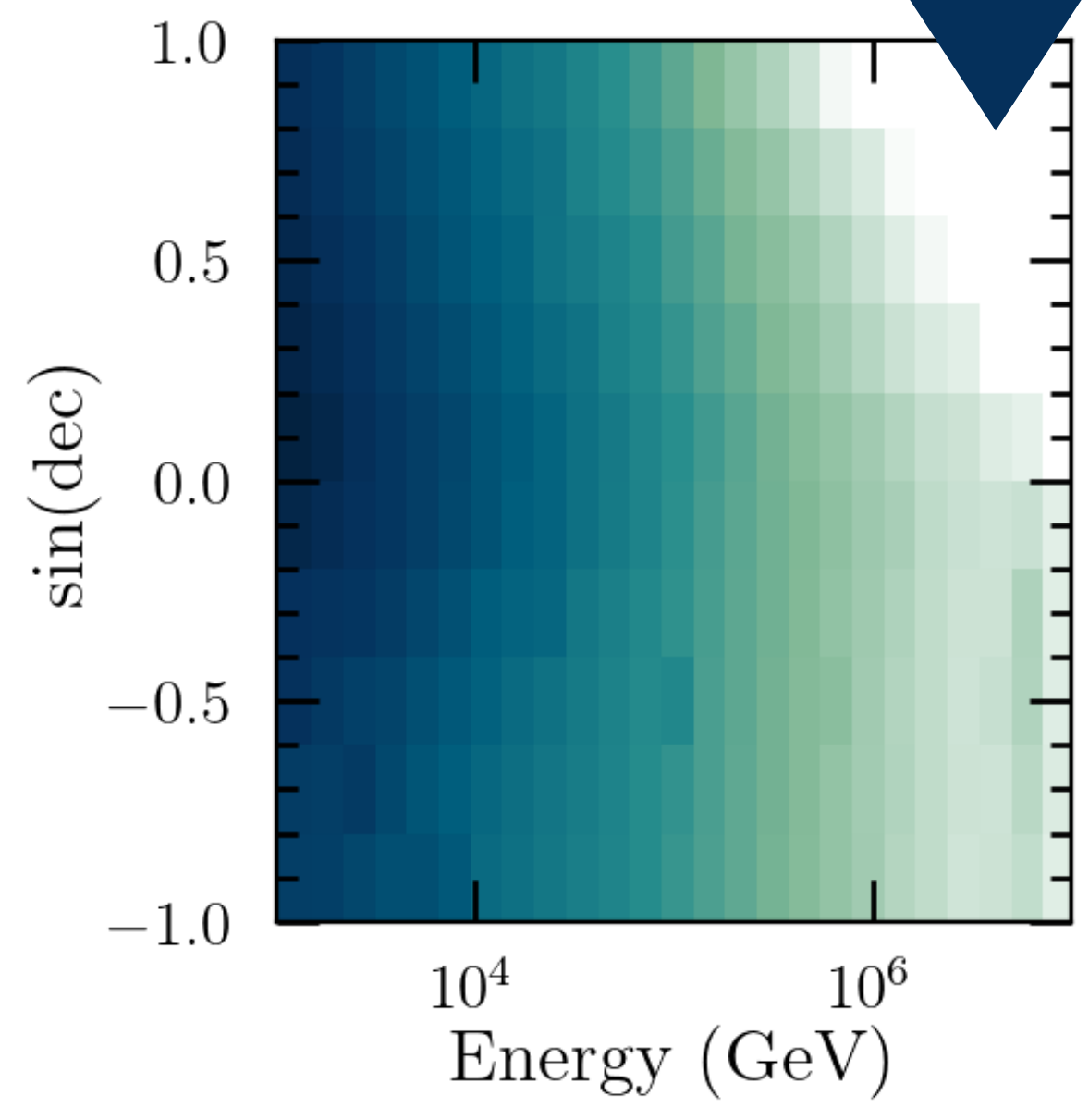
$m_\chi = 1\text{e-}06 \text{ GeV}$

$N_{\text{astro}} = 1.66$

$\gamma_{\text{astro}} = 2.53$

$N_{\text{galactic}} = 21.8$

Event Distribution =



Signal (S) -

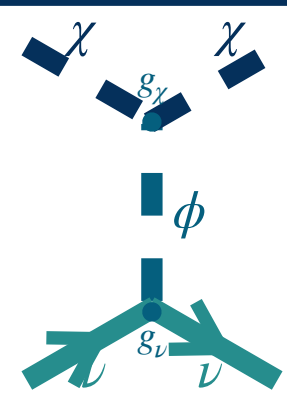
Scrambled

Signal (\tilde{S}) +

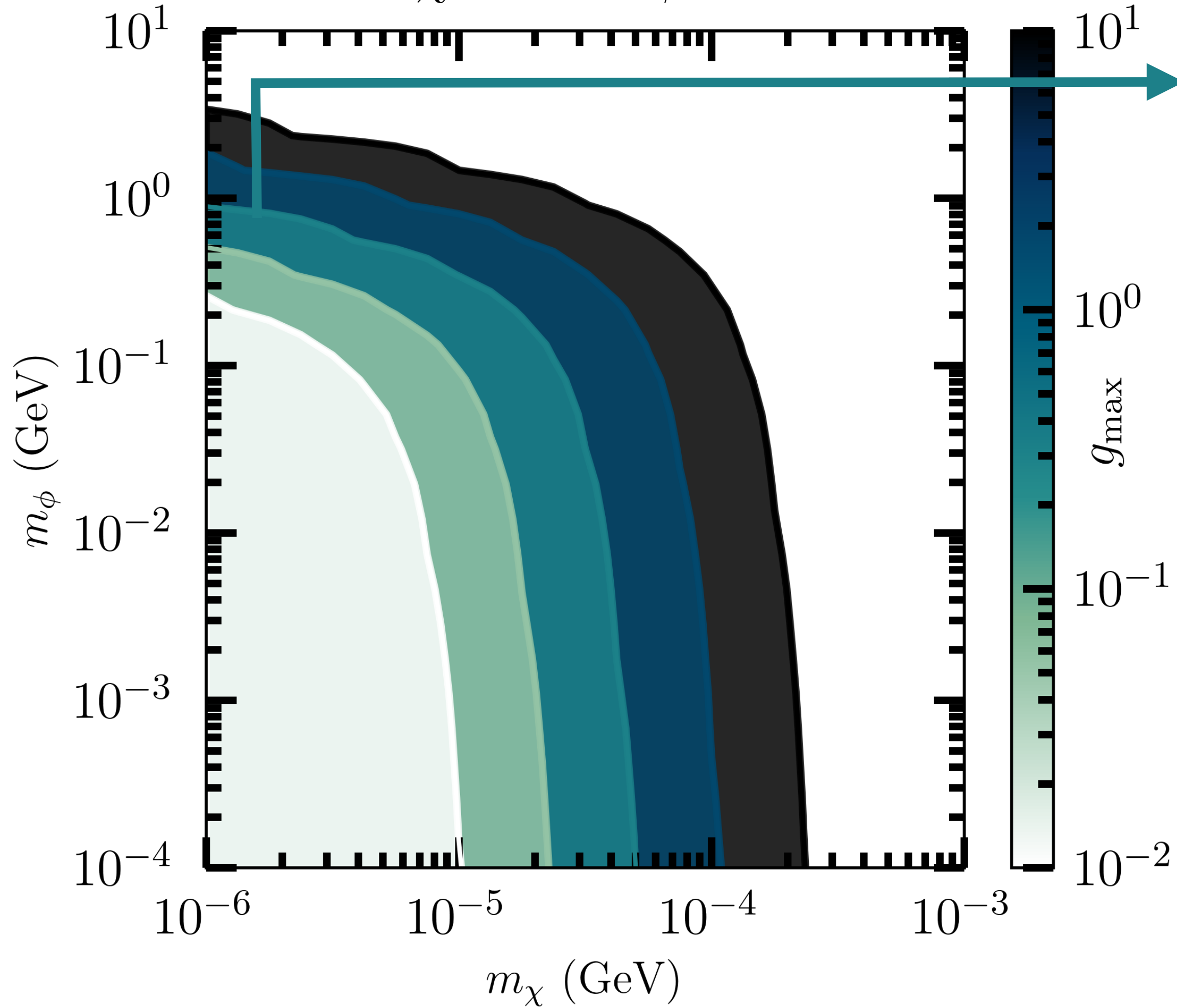
Scrambled

Data (\tilde{D})

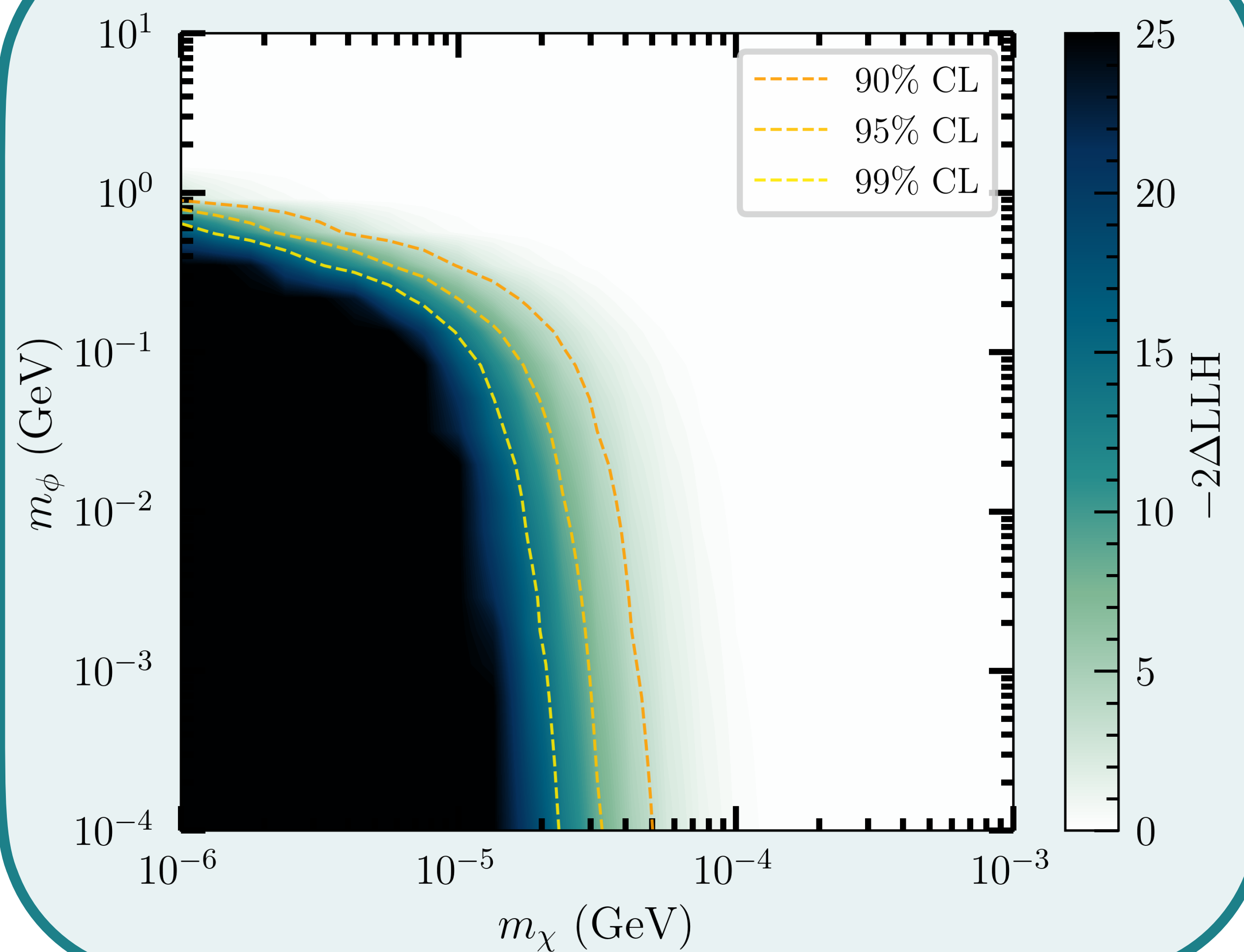
Asimov Sensitivities (One Example)



Scalar χ - Scalar ϕ : 90% CL

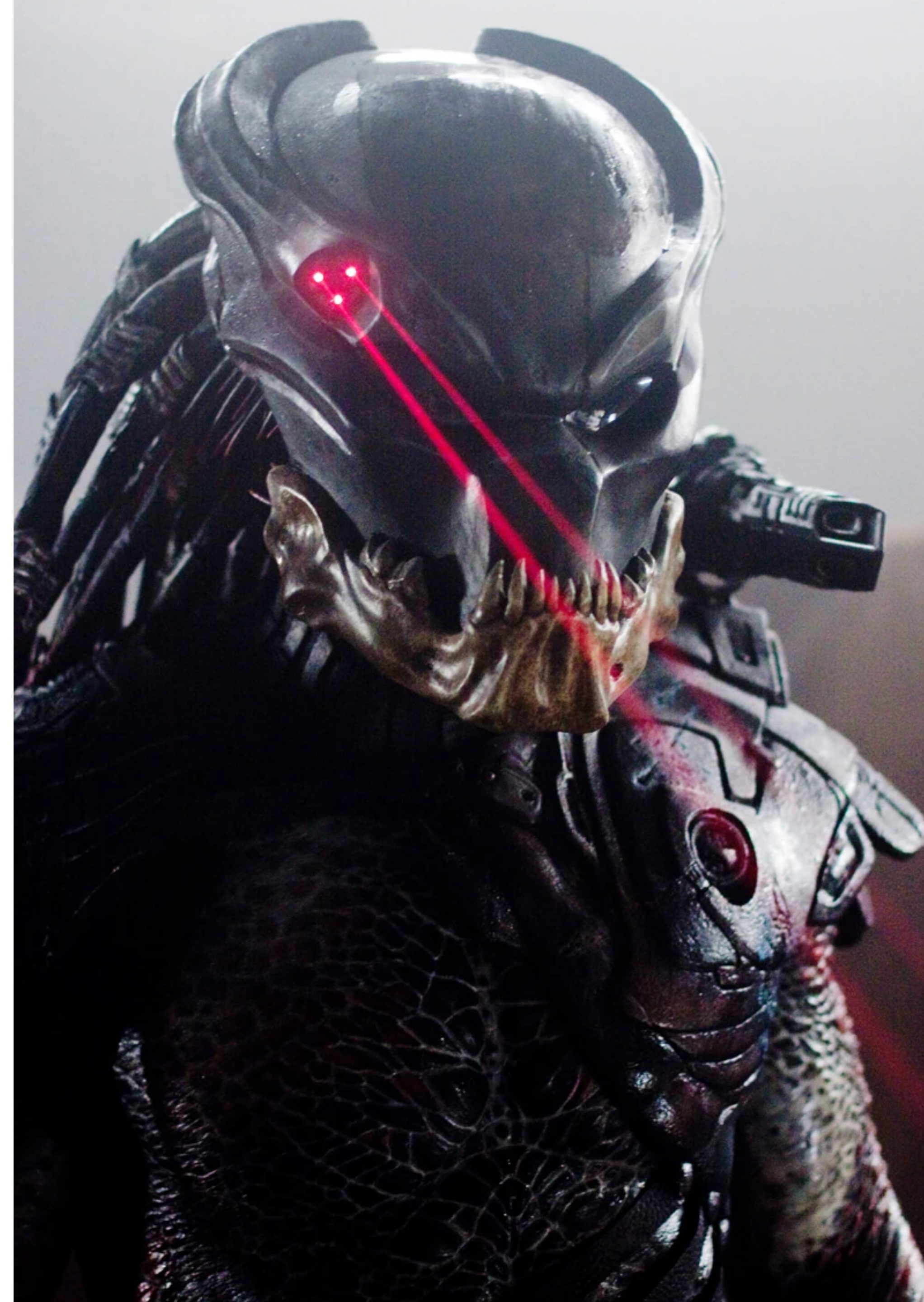
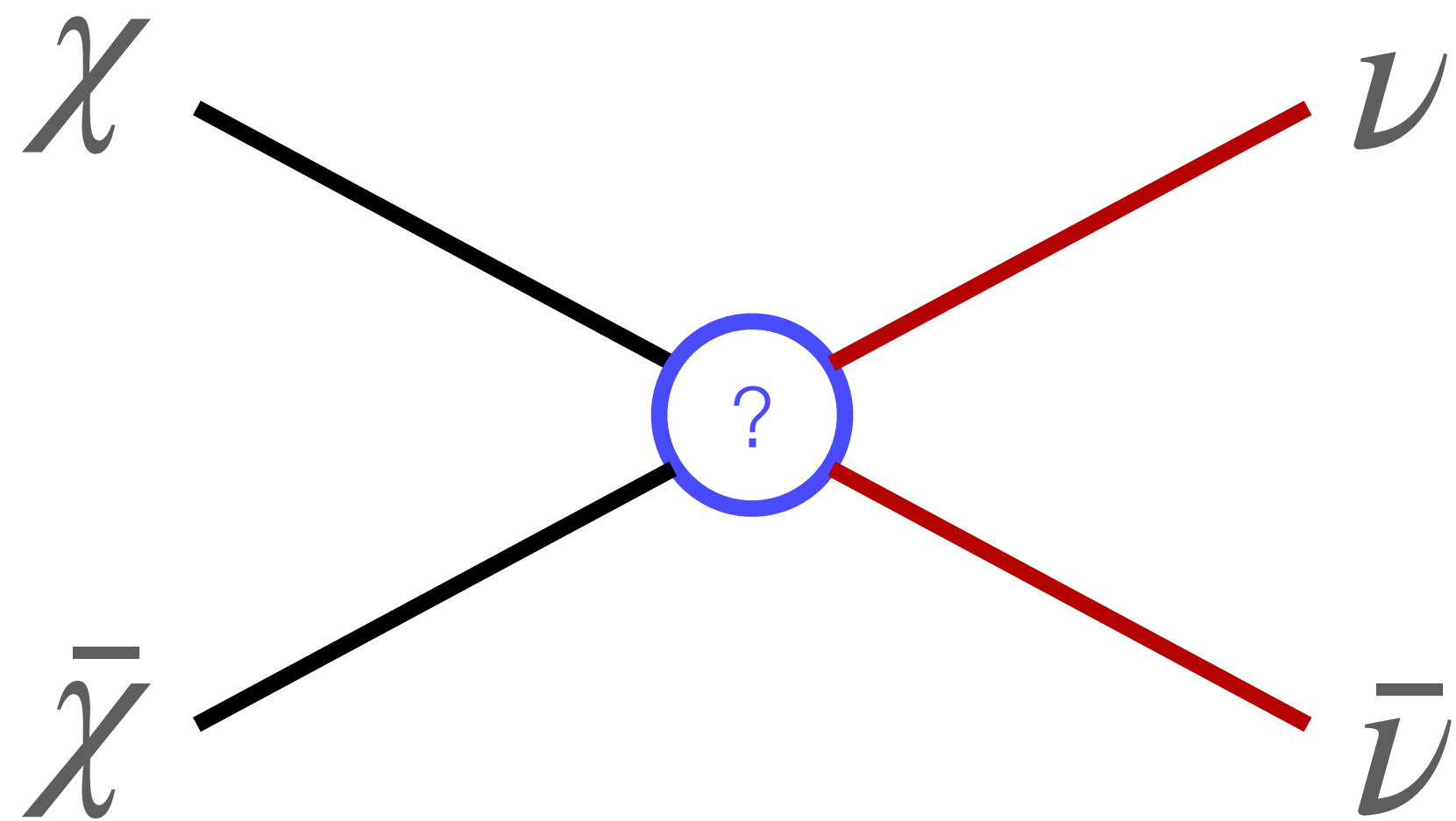


Scalar χ - Scalar ϕ : $g = 1.00e+00$ GeV



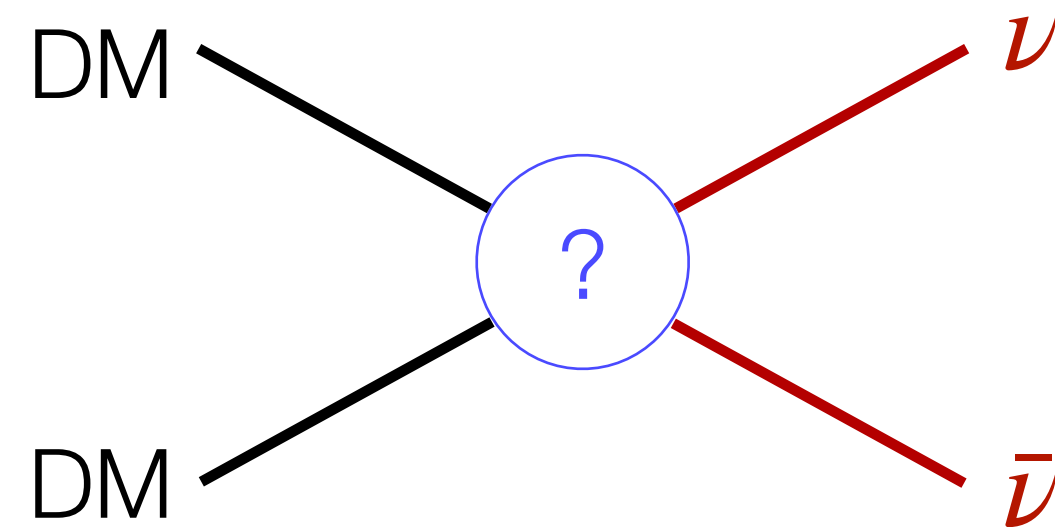
Unblinding This Week!

Annihilation



Annihilation

Indirect searches $\chi\chi \rightarrow SM, SM$: gammas dominate, *except* if neutrinos are the only product

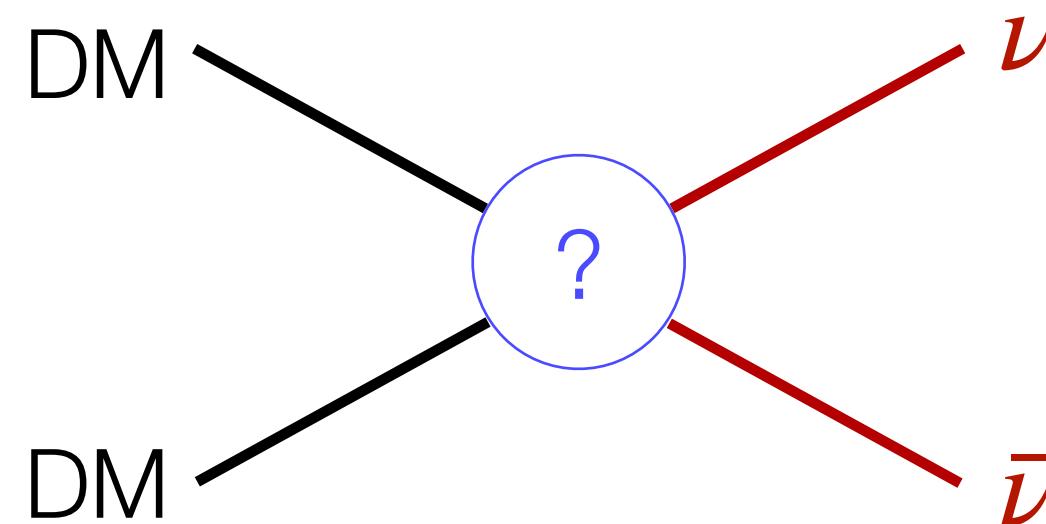


Annihilation

Indirect searches $\chi\chi \rightarrow SM, SM$: gammas dominate, *except* if neutrinos are the only product

$$\frac{d\Phi_{\nu+\bar{\nu}}}{dE_\nu} = \frac{1}{4\pi} \frac{\langle\sigma v\rangle}{\kappa m_\chi^2} \frac{1}{3} \frac{dN_\nu}{dE_\nu} J(\Omega)$$

$$J \equiv \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_\chi^2(x) dx$$



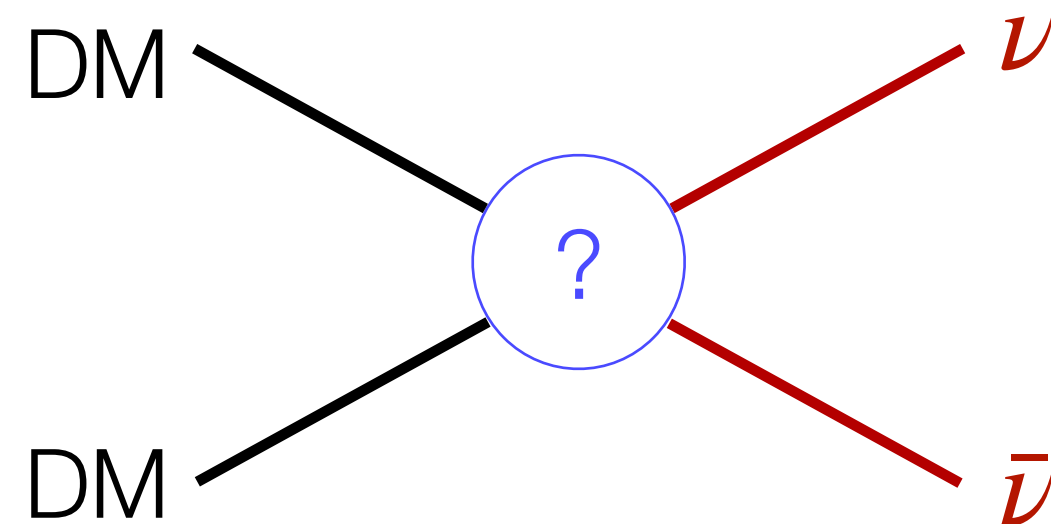
$$\frac{dN_\nu}{dE_\nu} = 2\delta(1 - E/m_\chi) m_\chi / E^2 .$$

Annihilation

Indirect searches $\chi\chi \rightarrow SM, SM$: gammas dominate, *except* if neutrinos are the only product

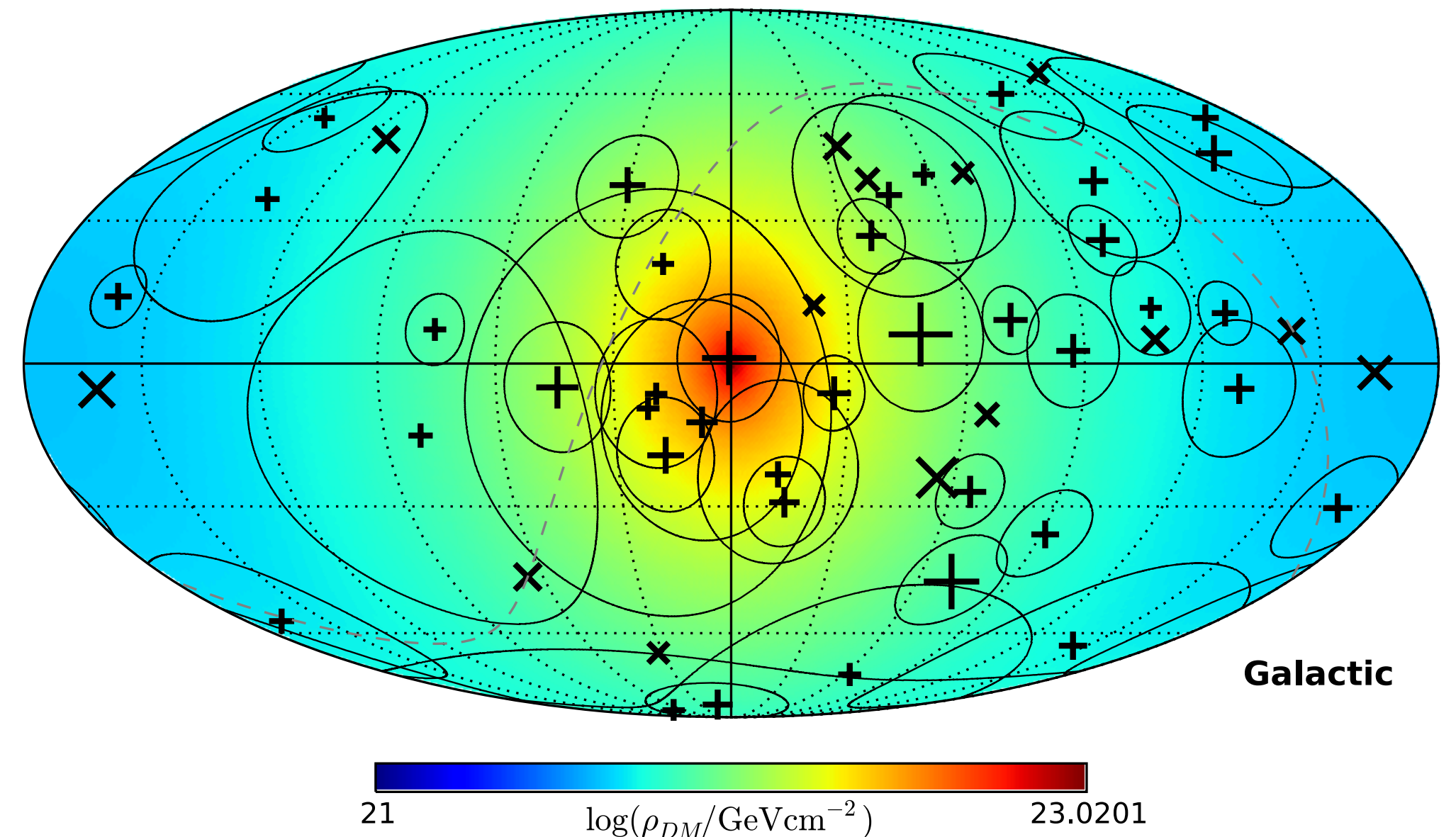
$$\frac{d\Phi_{\nu+\bar{\nu}}}{dE_\nu} = \frac{1}{4\pi} \frac{\langle\sigma v\rangle}{\kappa m_\chi^2} \frac{1}{3} \frac{dN_\nu}{dE_\nu} J(\Omega)$$

$$J \equiv \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_\chi^2(x) dx$$



$$\frac{dN_\nu}{dE_\nu} = 2\delta(1 - E/m_\chi) m_\chi / E^2.$$

Dark matter column density

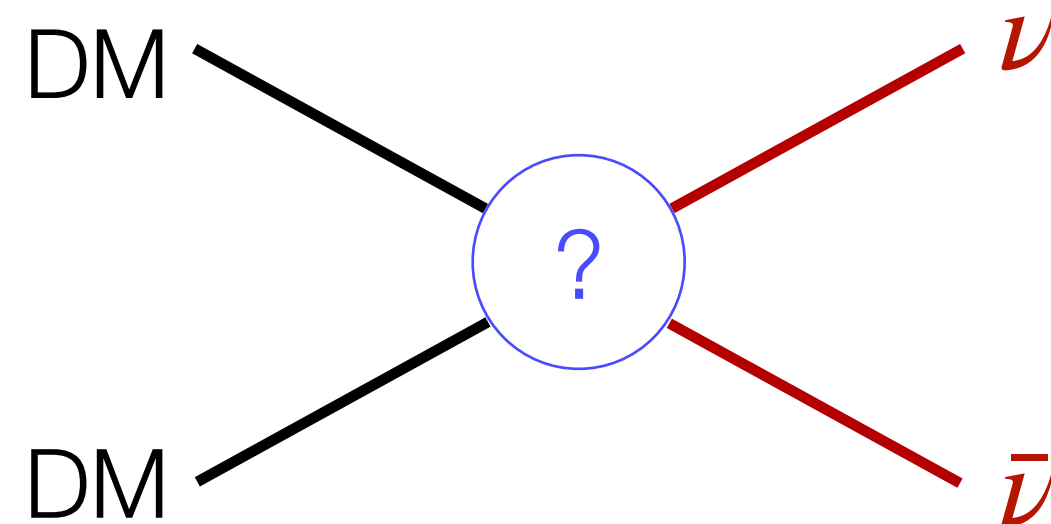


Annihilation

Indirect searches $\chi\chi \rightarrow SM, SM$: gammas dominate, *except* if neutrinos are the only product

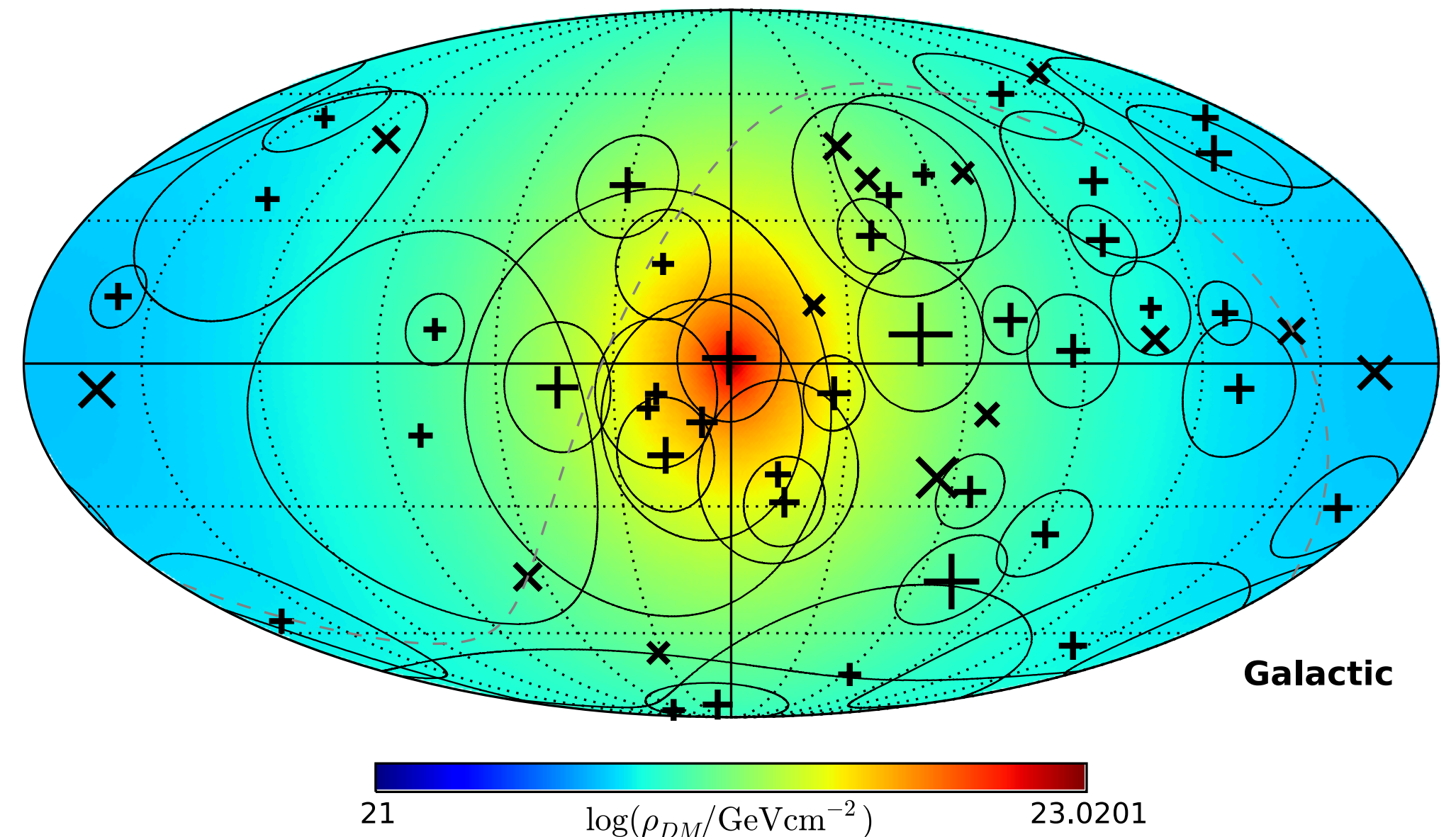
$$\frac{d\Phi_{\nu+\bar{\nu}}}{dE_\nu} = \frac{1}{4\pi} \frac{\langle\sigma v\rangle}{\kappa m_\chi^2} \frac{1}{3} \frac{dN_\nu}{dE_\nu} J(\Omega)$$

$$J \equiv \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_\chi^2(x) dx$$



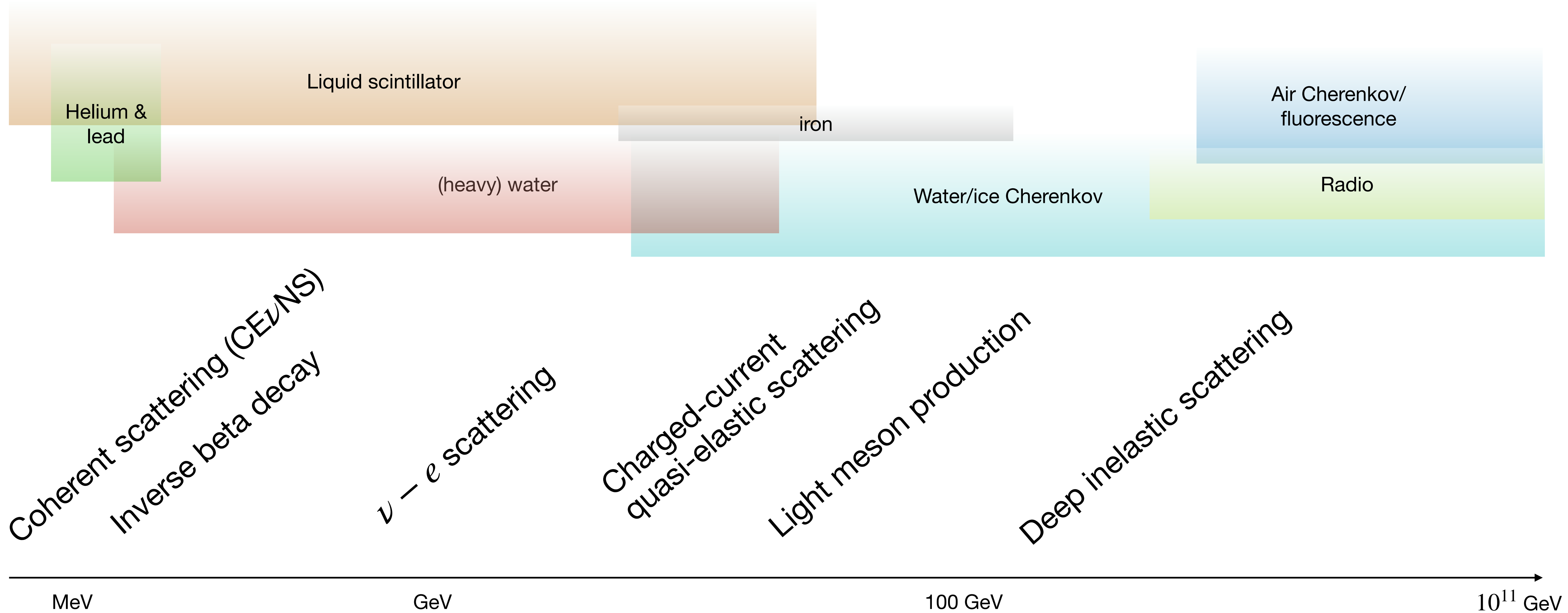
$$\frac{dN_\nu}{dE_\nu} = 2\delta(1 - E/m_\chi) m_\chi / E^2.$$

Dark matter column density

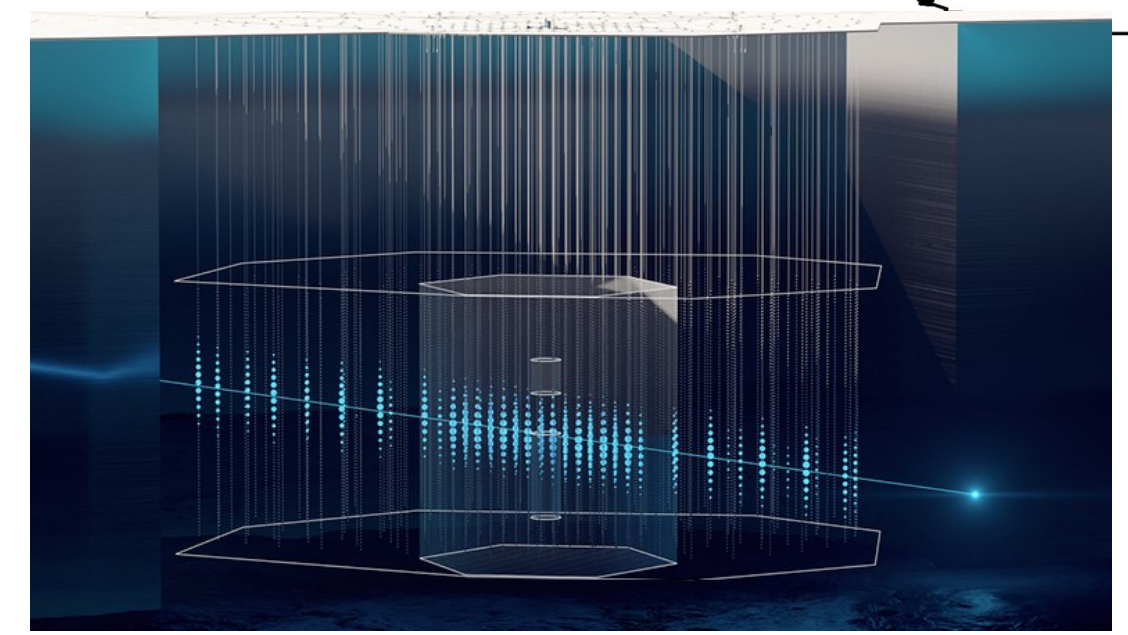
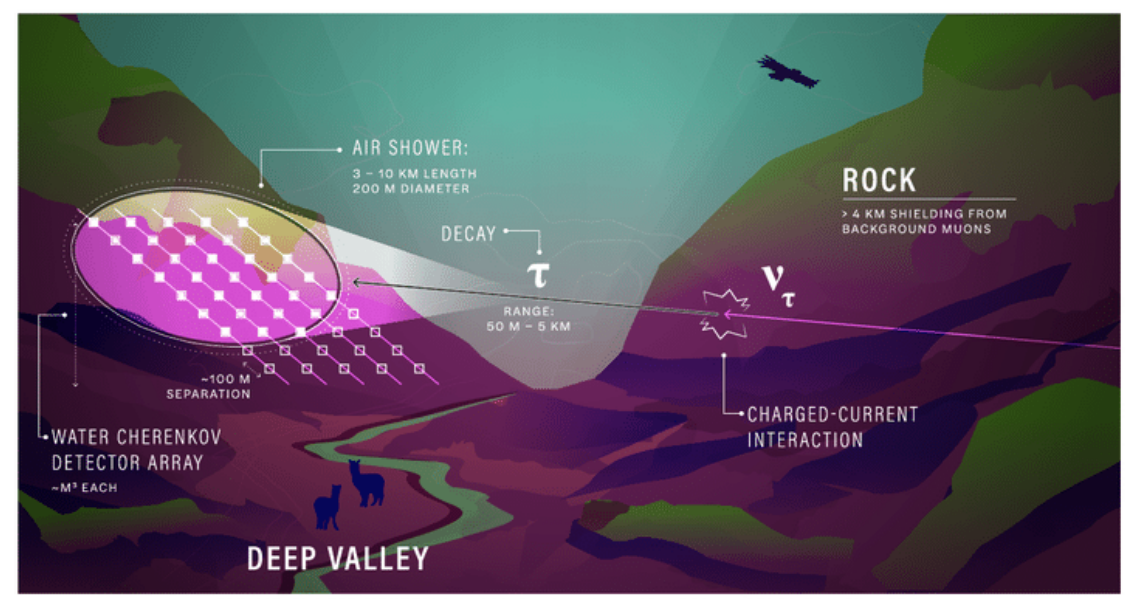
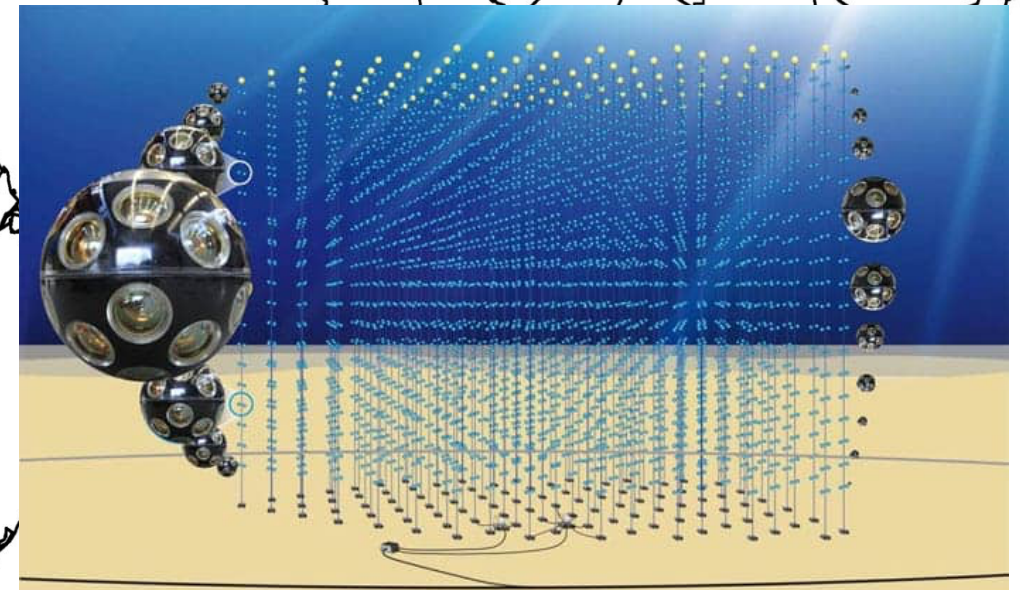
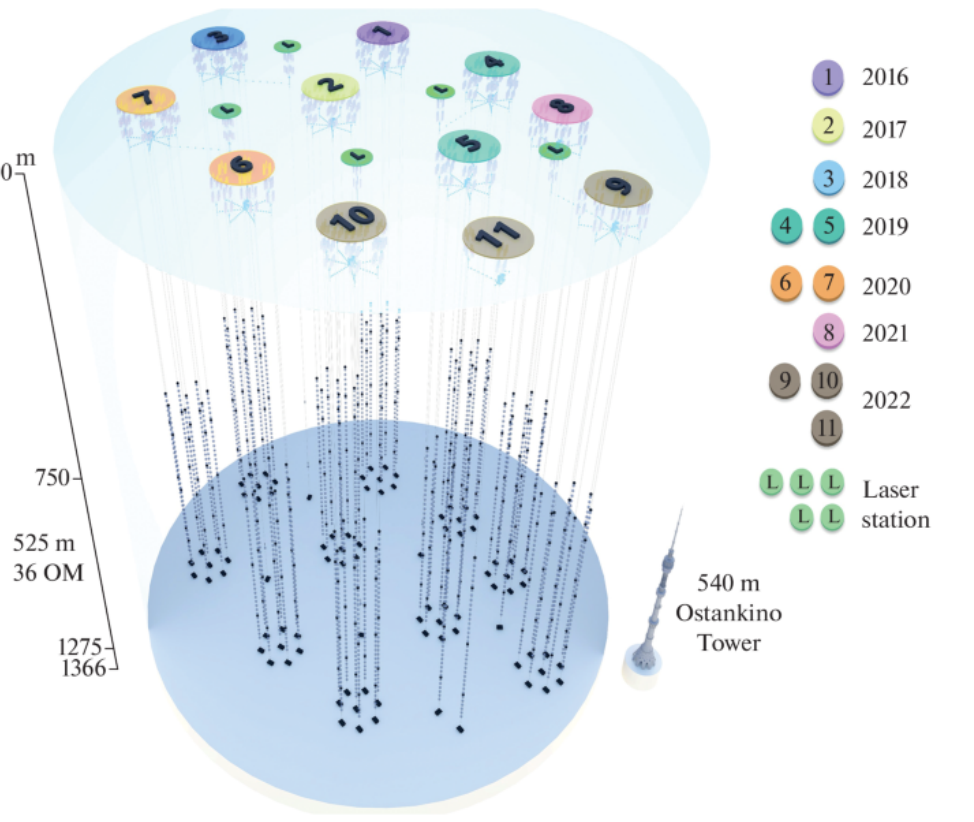
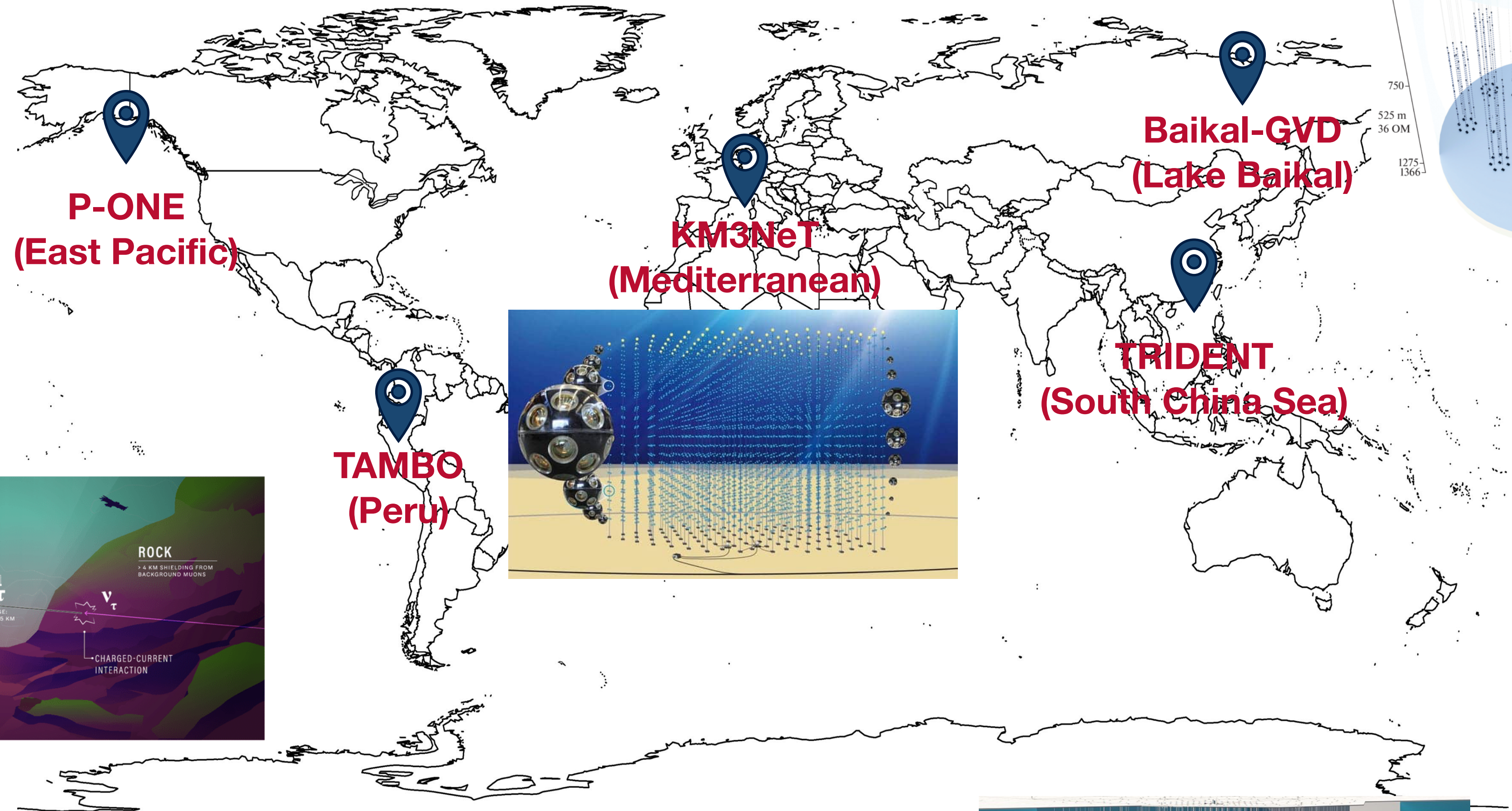
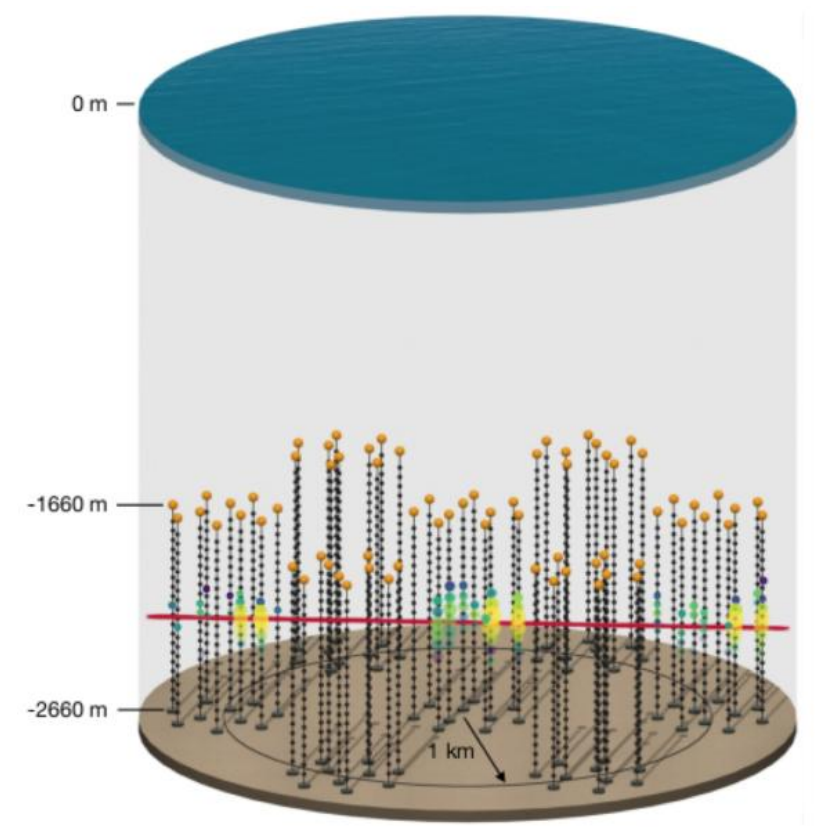


*We also spent a long time calculating extragalactic constraints. They are subdominant though

What if we looked at *every neutrino telescope in the world*?

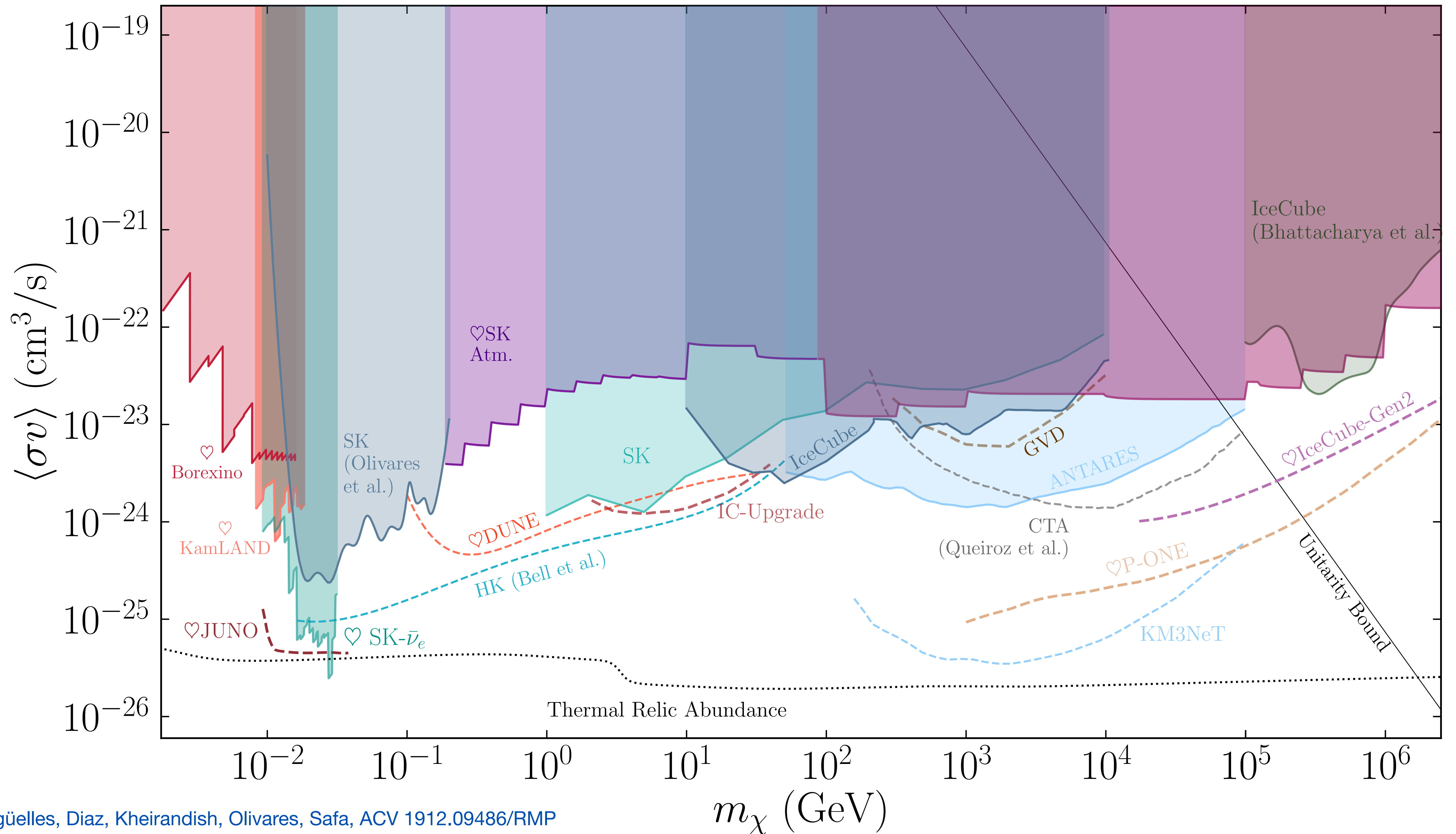


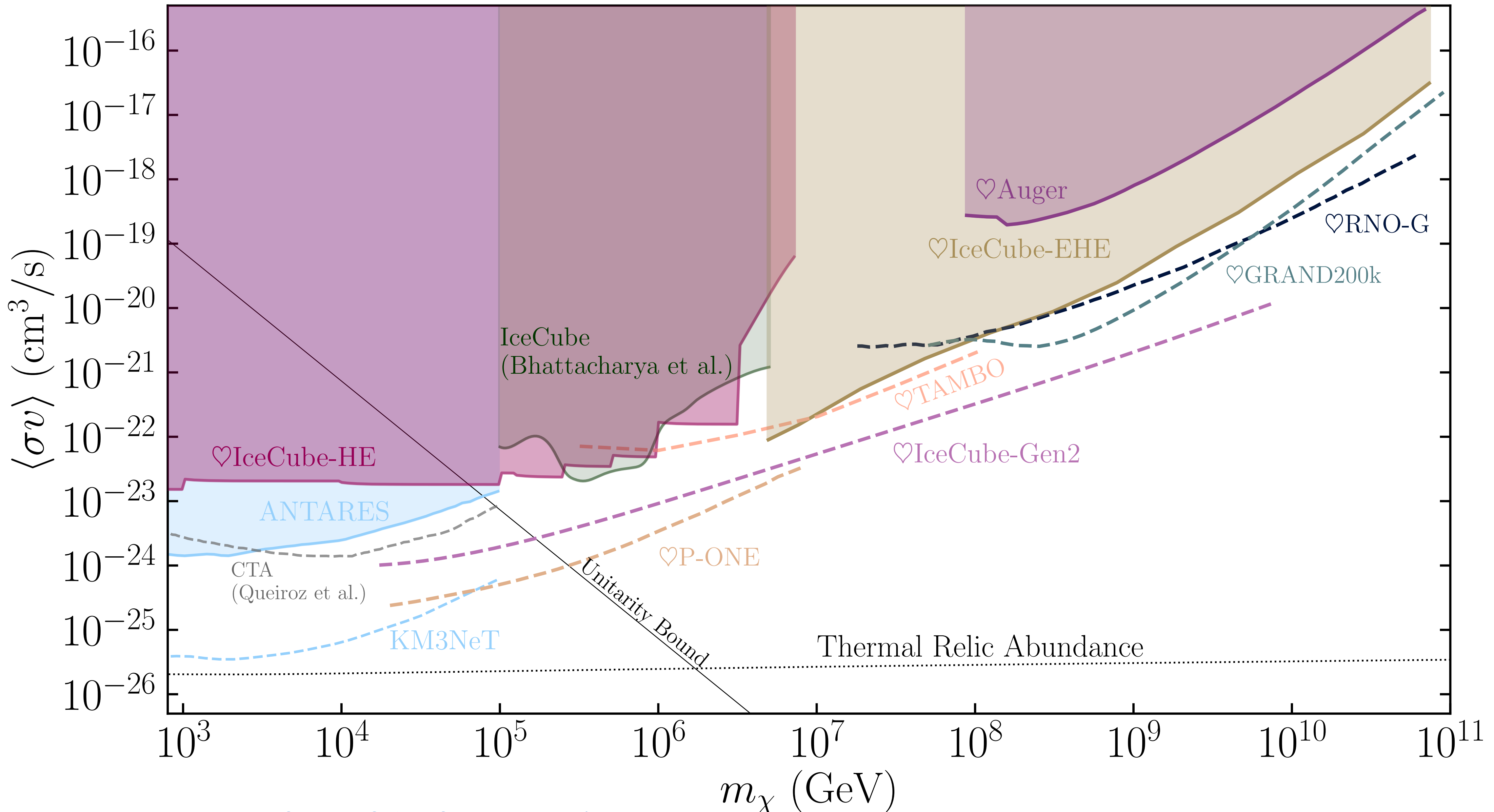
Next-Generation High-Energy Neutrino Telescopes



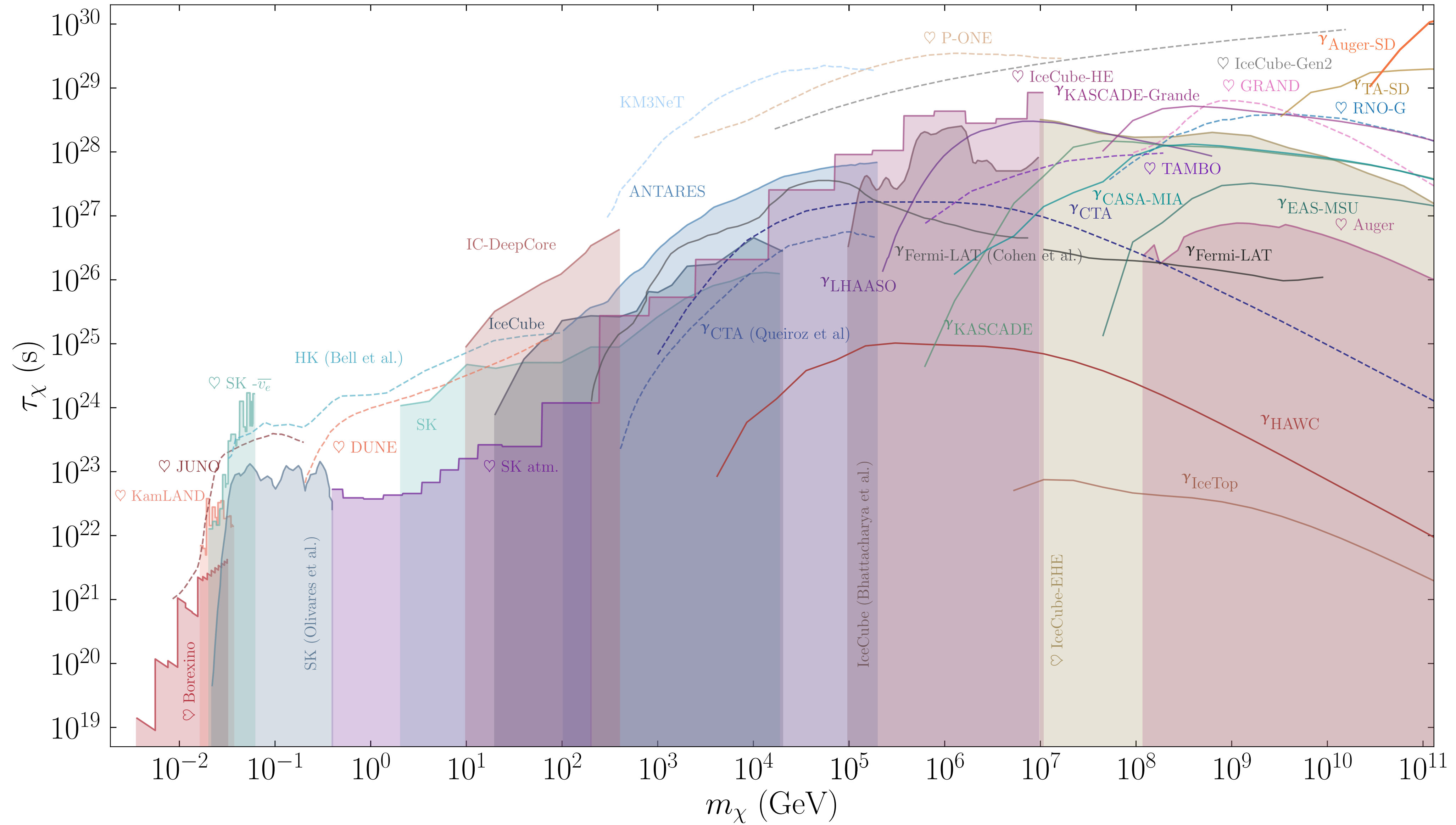
**IceCube-Gen2
(South Pole)**

More telescopes with larger exposure!



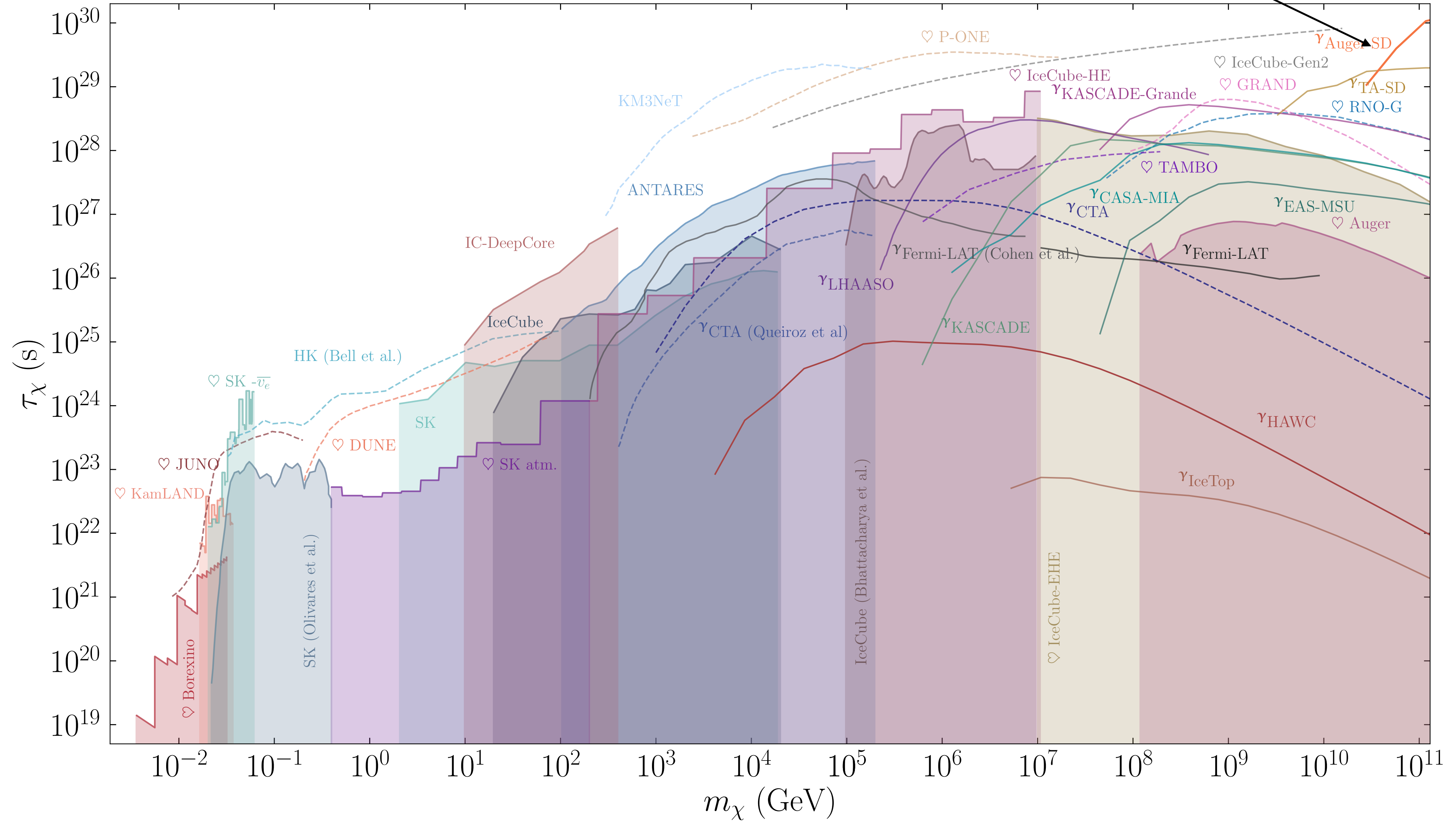


Decay to $\chi \rightarrow \bar{\nu}\nu$

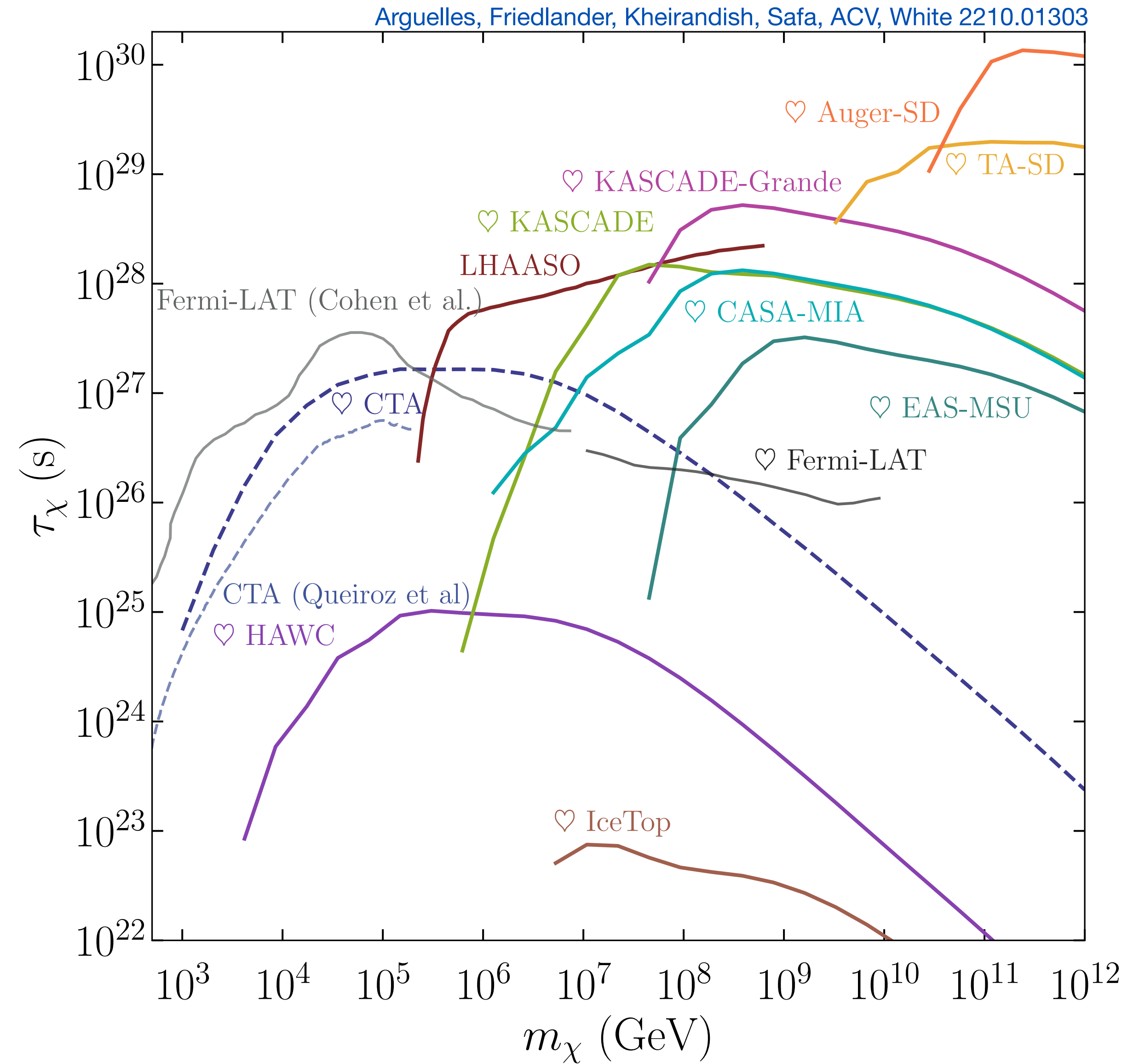
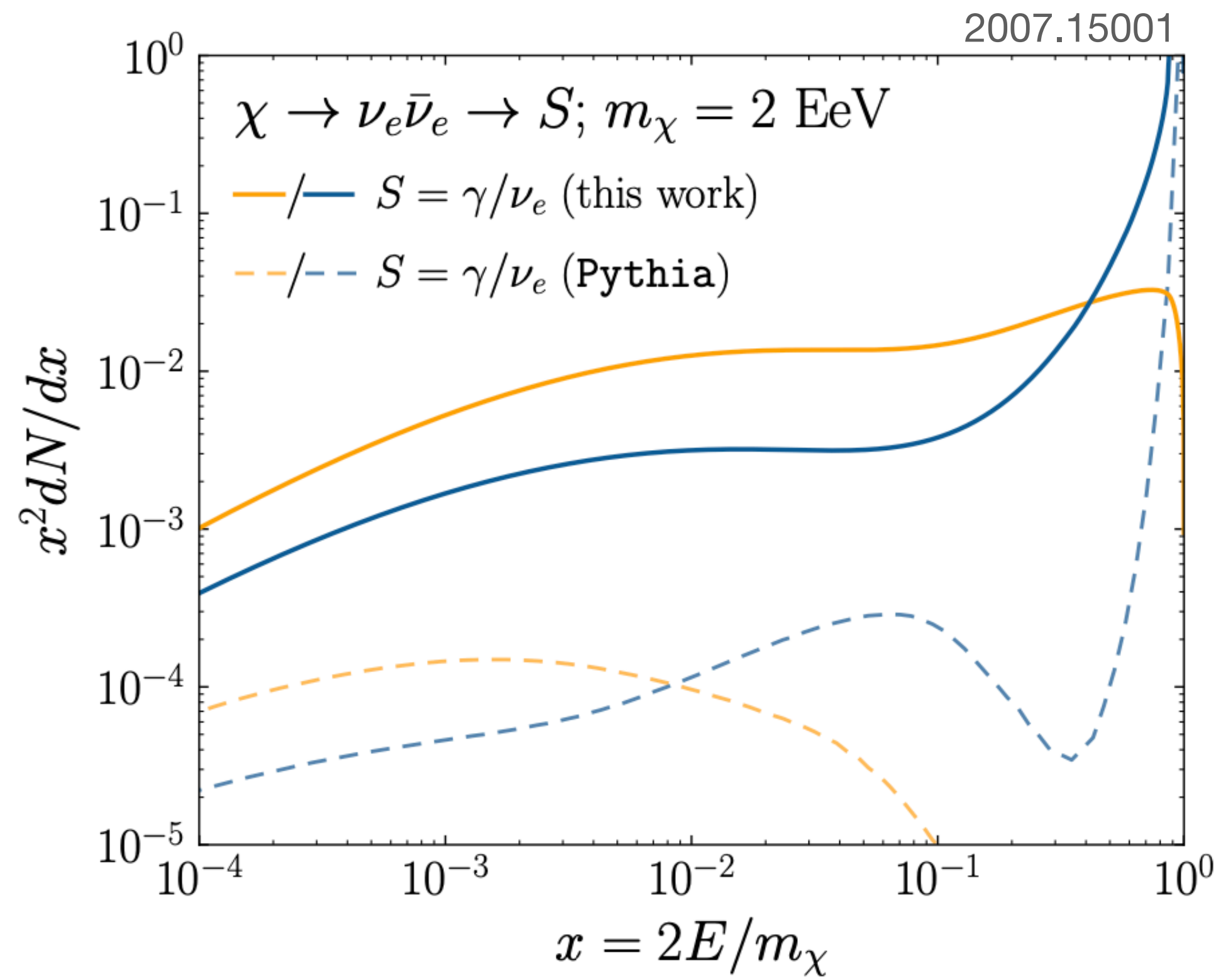


Decay to $\chi \rightarrow \bar{\nu}\nu$

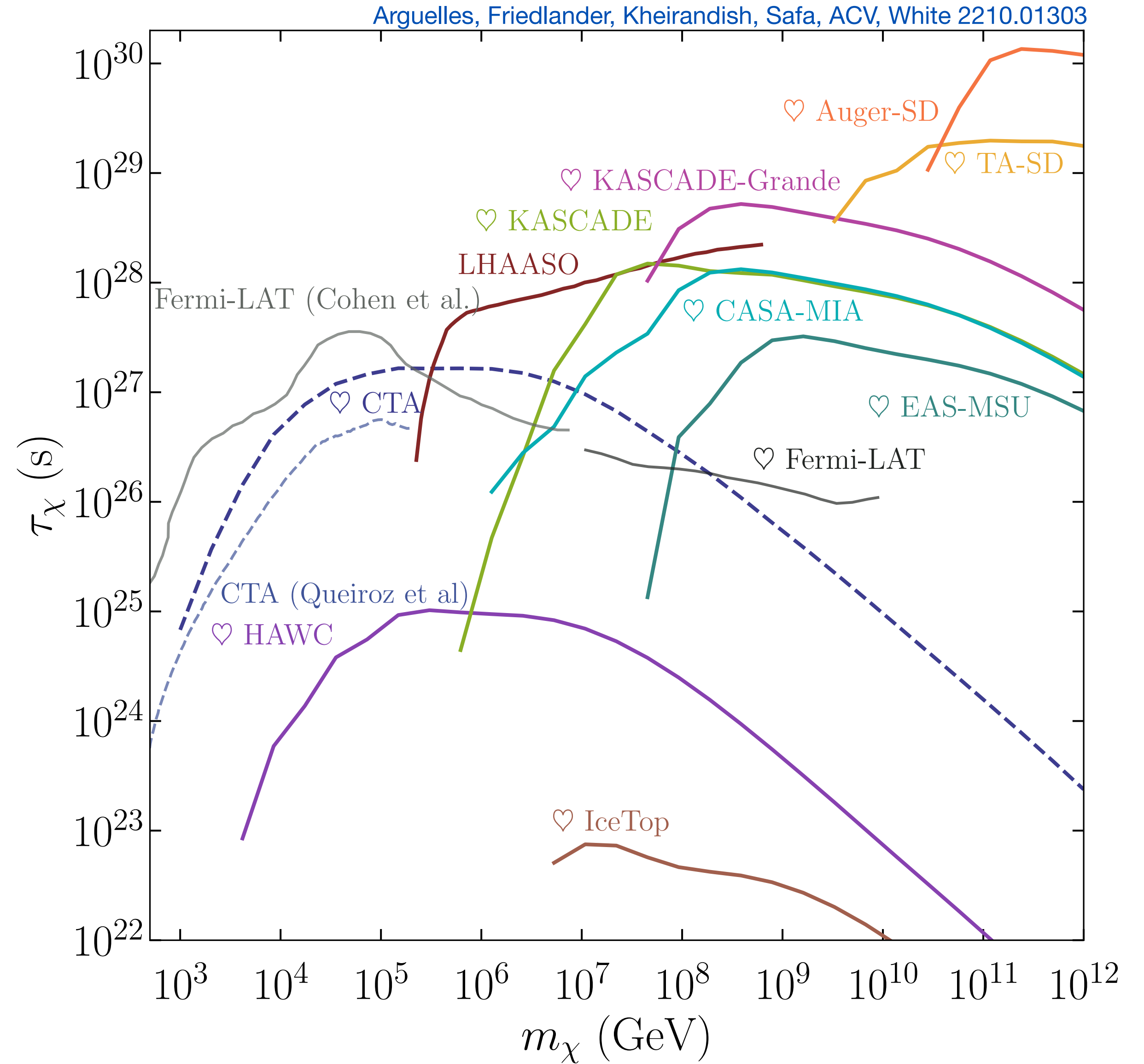
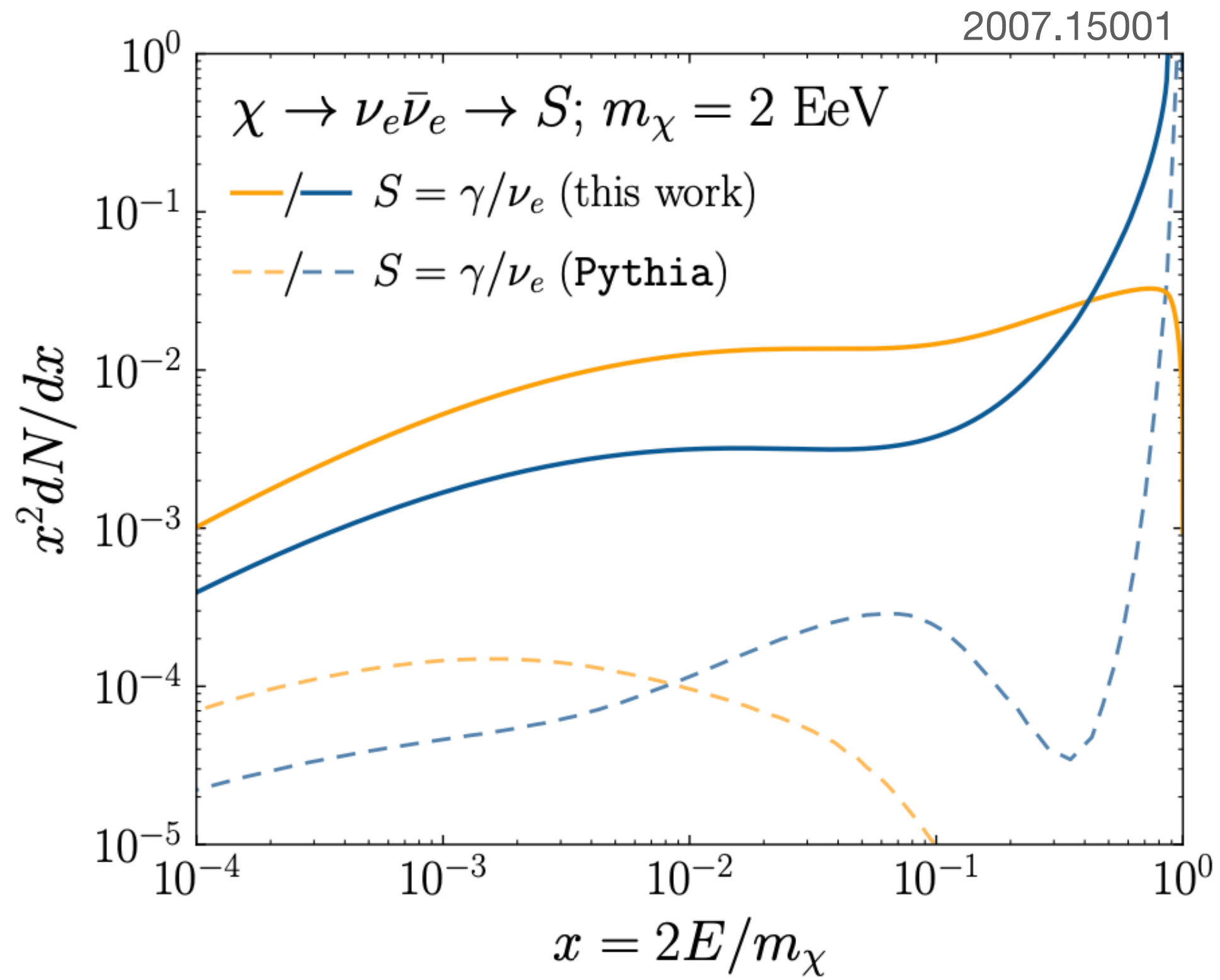
these are gamma-ray telescopes



Electroweak corrections



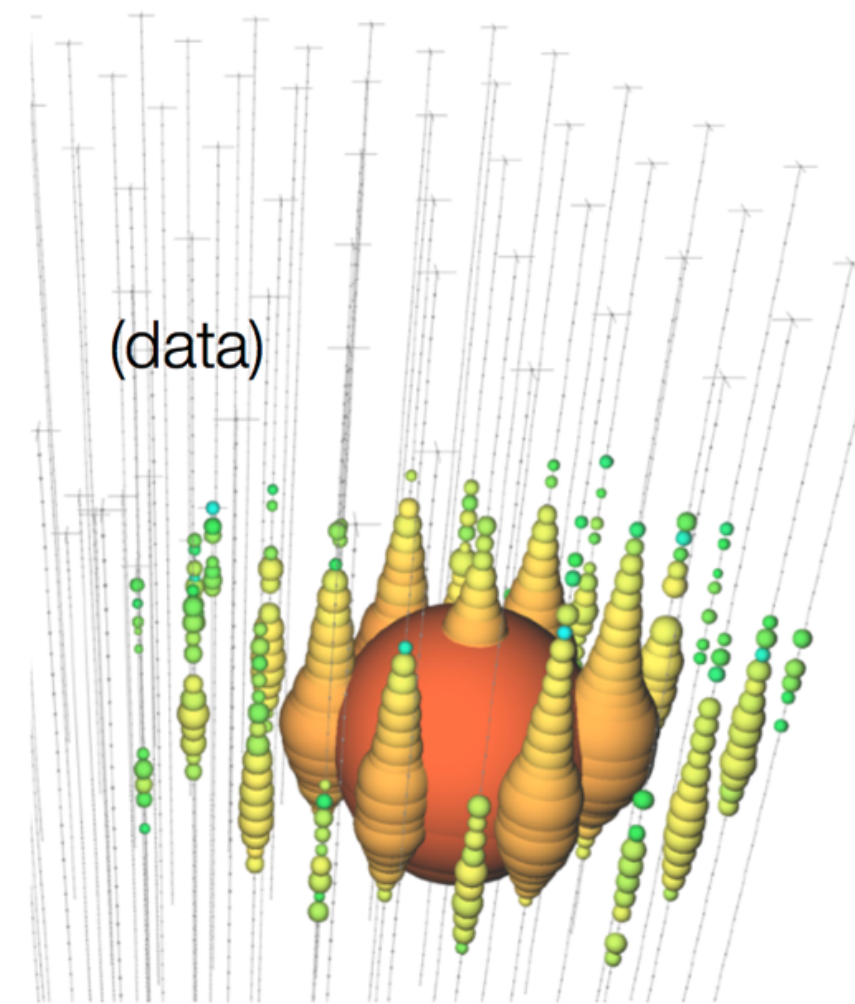
Electroweak corrections



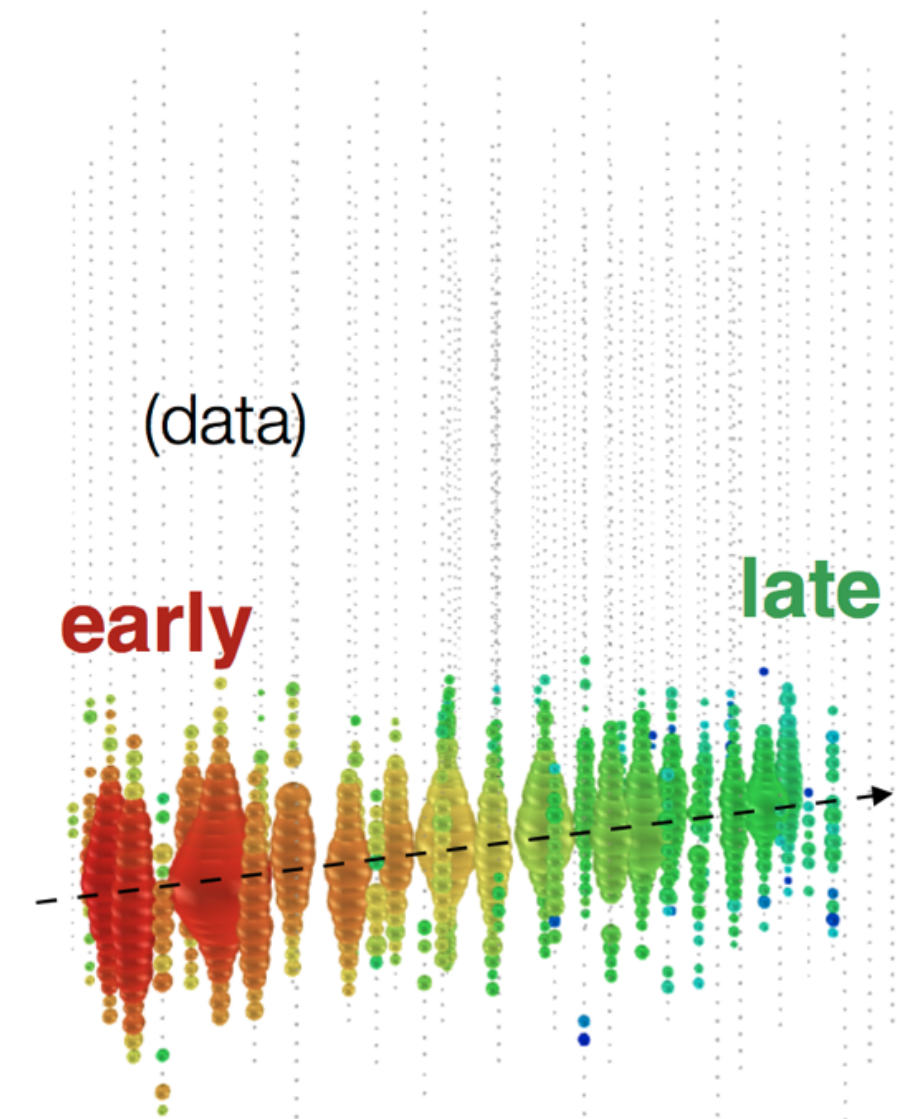
Looks like the gamma rays do better than neutrinos... but not quite the end of the story!

Flavour again

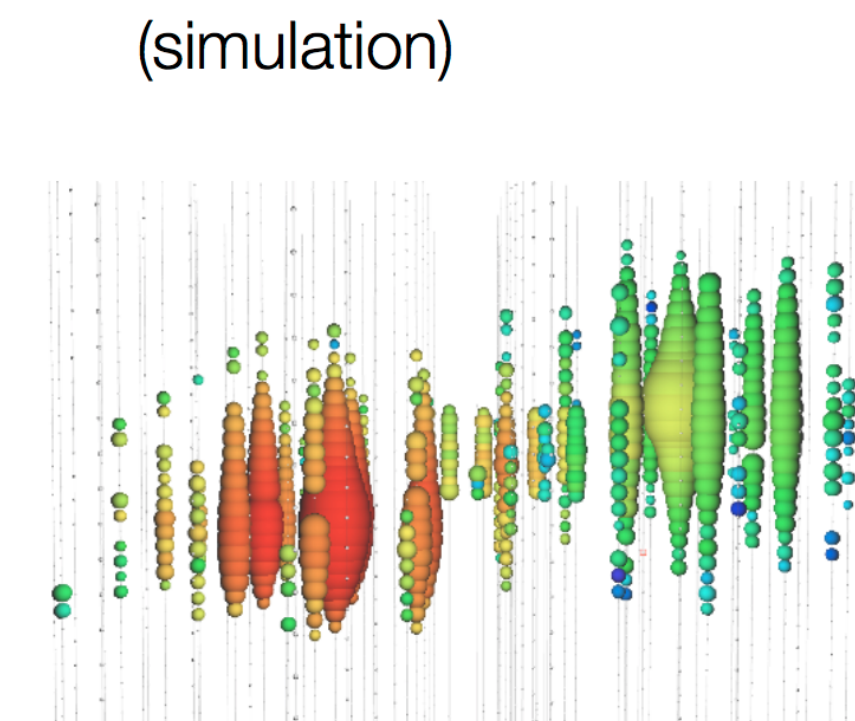
Neutral-current / ν_e



Charged-current ν_μ



Charged-current ν_τ



both ν and $\bar{\nu}$ look the same

What does ν vs $\bar{\nu}$ do for you?

Are neutrinos coming from pp or $p\gamma$ collisions? These give different π^+/π^- ratios

$$\{\nu_e, \bar{\nu}_e\} : \{\nu_\mu, \bar{\nu}_\mu\} : \{\nu_\tau, \bar{\nu}_\tau\}$$

$$p + p \rightarrow n_\pi [\pi^0 + \pi^+ + \pi^-]$$

$$p + \gamma \rightarrow \Delta^+ \rightarrow \pi^+ + n$$

Production	Source flavor ratio	Earth flavor ratio $\nu + \bar{\nu}$	Earth flavor ratio	$f_{\bar{\nu}_e}$
pp	$\{1, 1\} : \{2, 2\} : \{0, 0\}$	$0.33 : 0.34 : 0.33$	$\{0.17, 0.17\} : \{0.17, 0.17\} : \{0.16, 0.16\}$	0.17
$p\gamma$	$\{1, 0\} : \{1, 1\} : \{0, 0\}$	$0.33 : 0.34 : 0.33$	$\{0.26, 0.08\} : \{0.21, 0.13\} : \{0.20, 0.13\}$	0.08

What does ν vs $\bar{\nu}$ do for you?

Are neutrinos coming from pp or $p\gamma$ collisions? These give different π^+/π^- ratios

$$\{\nu_e, \bar{\nu}_e\} : \{\nu_\mu, \bar{\nu}_\mu\} : \{\nu_\tau, \bar{\nu}_\tau\}$$

$$p + p \rightarrow n_\pi [\pi^0 + \pi^+ + \pi^-]$$

$$p + \gamma \rightarrow \Delta^+ \rightarrow \pi^+ + n$$

Production	Source flavor ratio	Earth flavor ratio $\nu + \bar{\nu}$	Earth flavor ratio	$f_{\bar{\nu}_e}$
pp	$\{1, 1\} : \{2, 2\} : \{0, 0\}$	$0.33 : 0.34 : 0.33$	$\{0.17, 0.17\} : \{0.17, 0.17\} : \{0.16, 0.16\}$	0.17
$p\gamma$	$\{1, 0\} : \{1, 1\} : \{0, 0\}$	$0.33 : 0.34 : 0.33$	$\{0.26, 0.08\} : \{0.21, 0.13\} : \{0.20, 0.13\}$	0.08
$pp \mu$ damped	$\{0, 0\} : \{1, 1\} : \{0, 0\}$	$0.23 : 0.39 : 0.38$	$\{0.11, 0.11\} : \{0.20, 0.20\} : \{0.19, 0.19\}$	0.11
$p\gamma \mu$ damped	$\{0, 0\} : \{1, 0\} : \{0, 0\}$	$0.23 : 0.39 : 0.38$	$\{0.23, 0.00\} : \{0.39, 0.00\} : \{0.38, 0.00\}$	0

New physics?

ν vs $\bar{\nu}$

- Symmetric dark matter annihilation/decay: equal parts neutrino/antineutrino
- Asymmetric dark matter

$$\mathcal{O}_{\text{ADM}} = \frac{\mathcal{O}_X \mathcal{O}_{\text{SM}}}{\Lambda^{m+n-4}}$$

Benchmark	B1 (scalar X)	B2 (scalar X)	B3 (fermion X)	B4 (fermion X)
$\mathcal{O}_{X \rightarrow \nu}$	$\frac{1}{\Lambda} X \psi L \Phi$	$\frac{1}{\Lambda^2} X (L \Phi)^2$	$\frac{1}{\Lambda^2} X L \psi^2$	$\frac{1}{\Lambda^2} X L L \nu^c$
Decay	$X \rightarrow \bar{\nu} \psi / \nu \psi$	$X \rightarrow \bar{\nu} \bar{\nu} / \nu \nu$	$X \rightarrow \bar{\nu} \psi \bar{\psi} / \nu \psi \bar{\psi}$	$X \rightarrow \bar{\nu} \nu \bar{\nu} / \nu \nu \bar{\nu}$

(Aside: India-based Neutrino Observatory proposal)



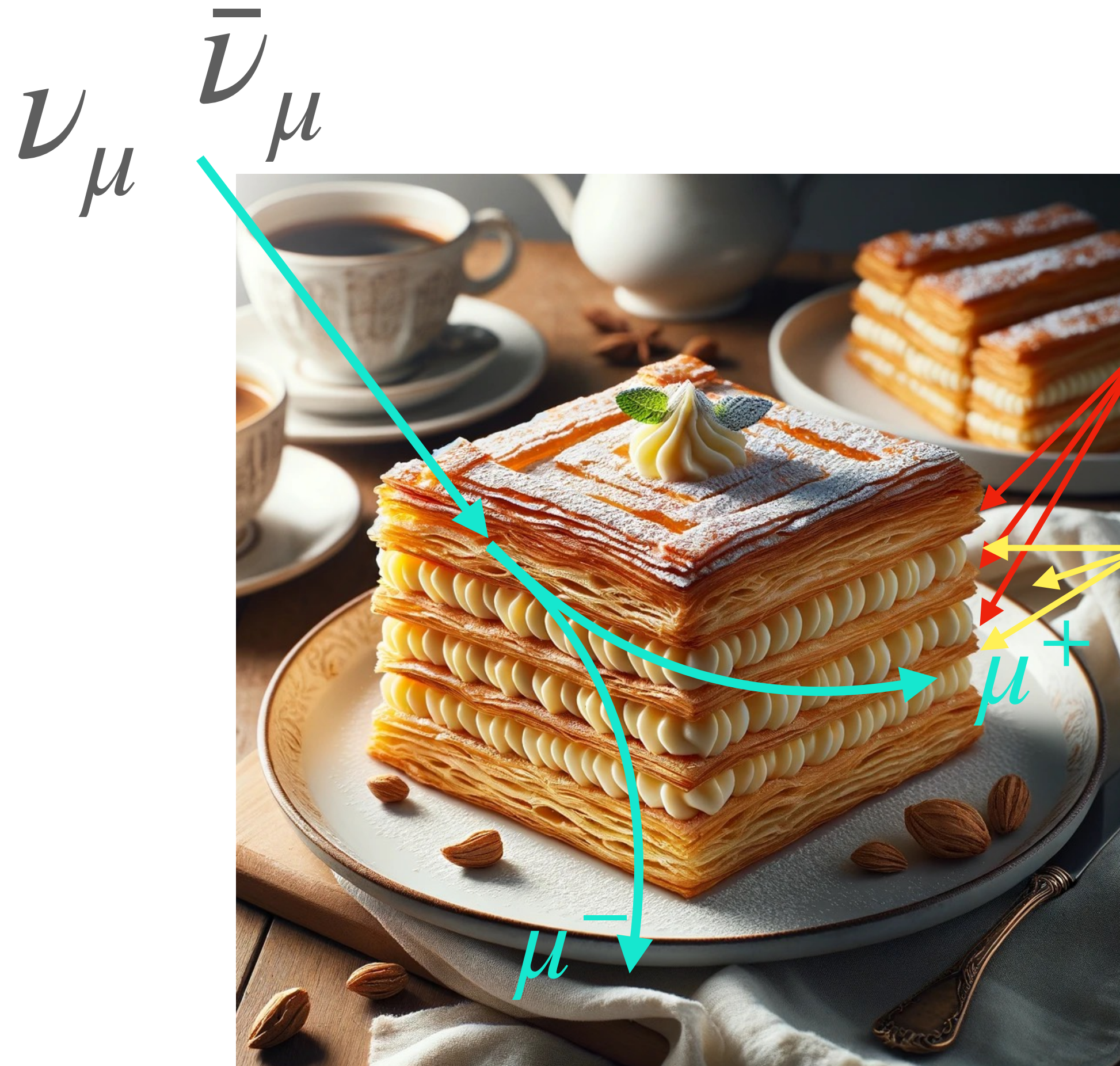
(Aside: India-based Neutrino Observatory proposal)



50,000 tons magnetized iron leaves

Resistive plate chambers

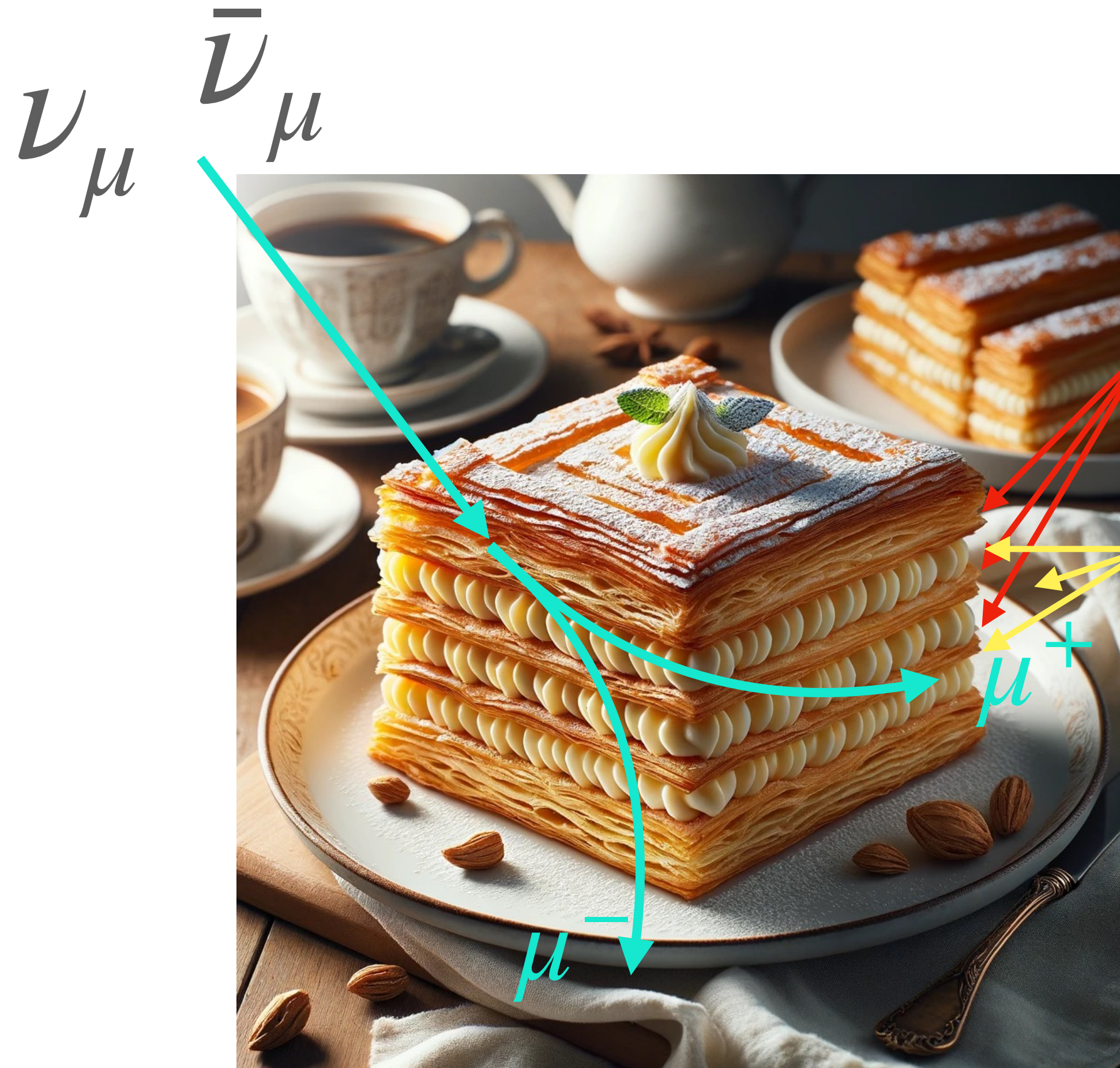
(Aside: India-based Neutrino Observatory proposal)



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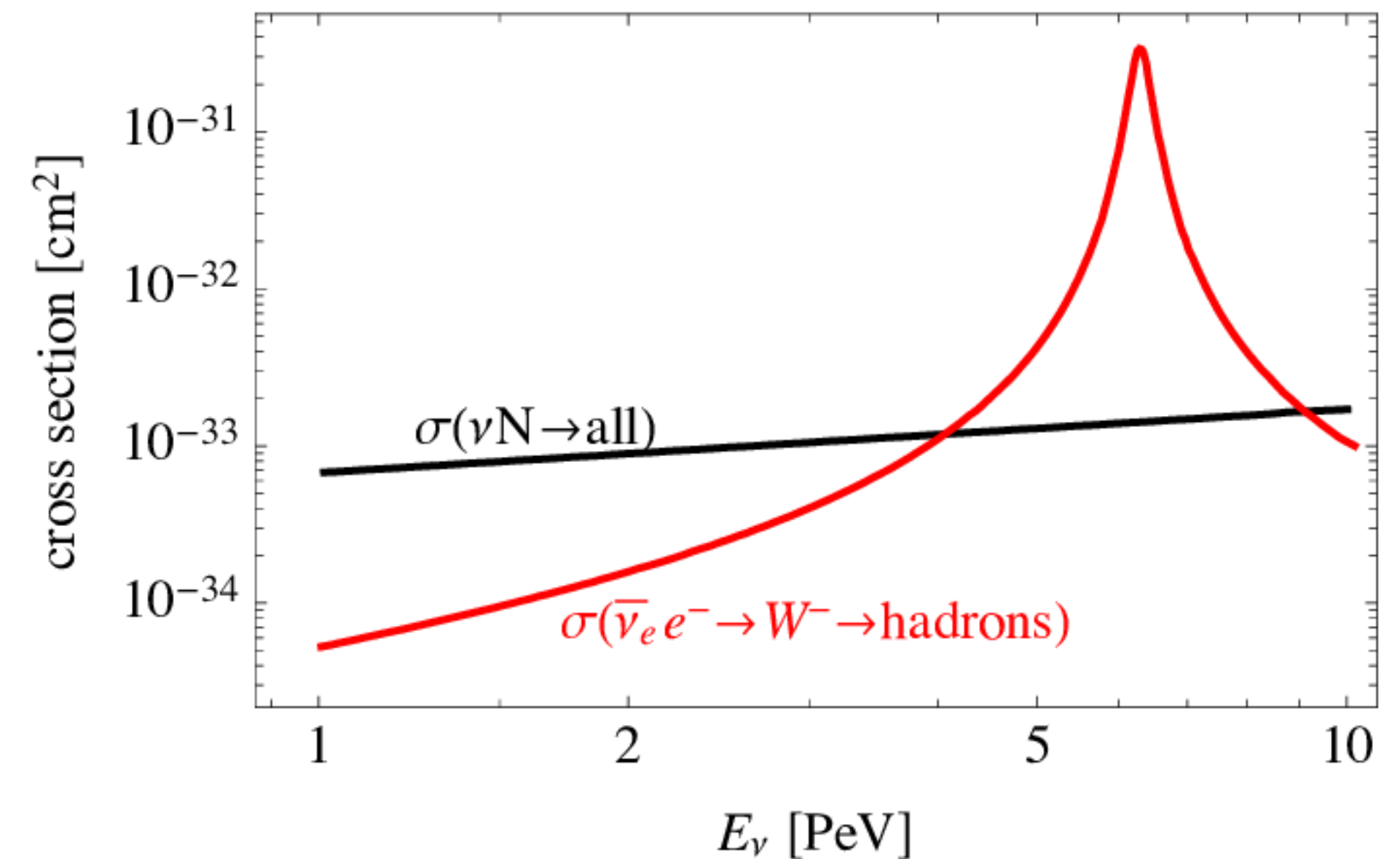
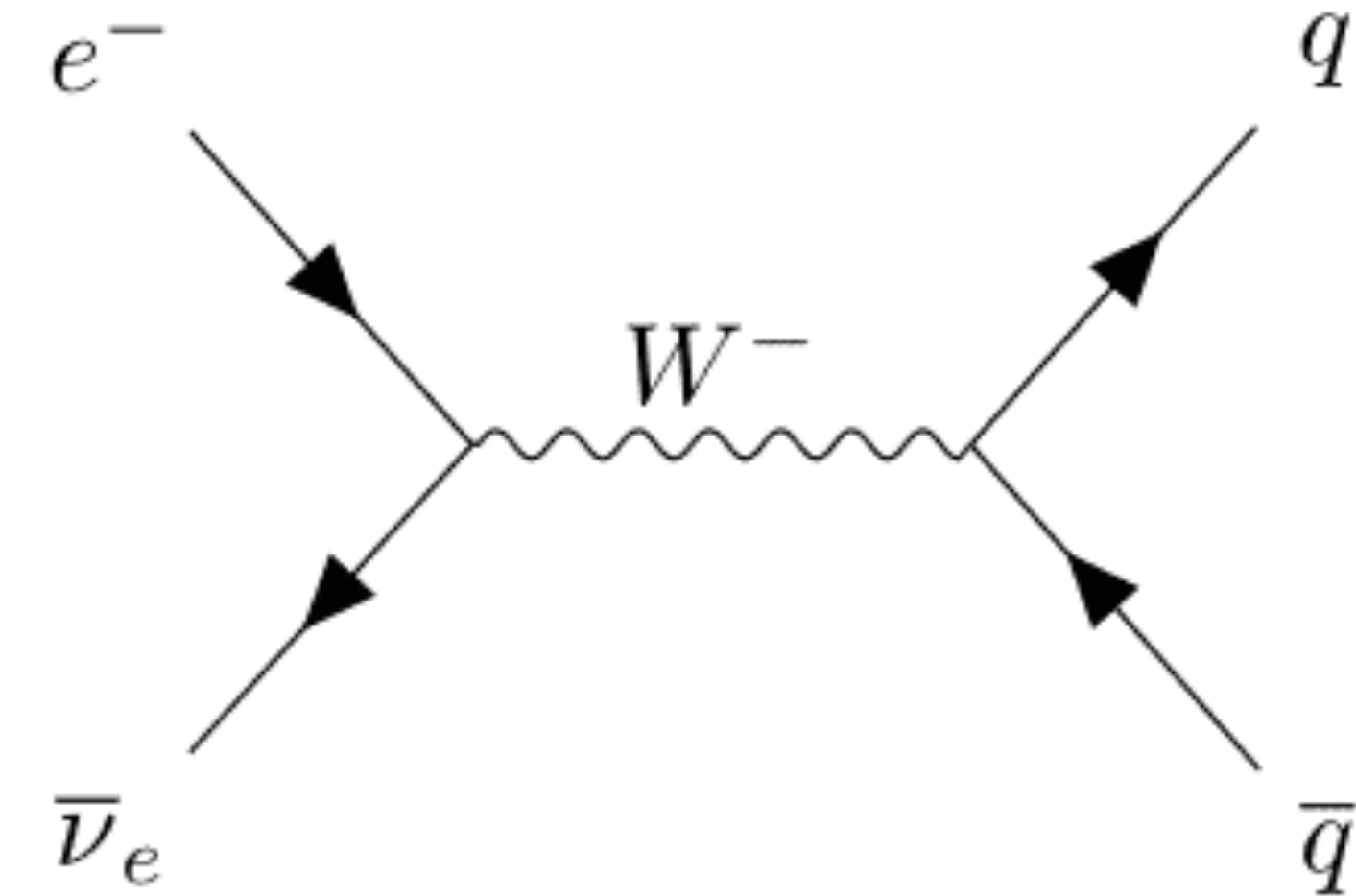
50,000 tons magnetized iron leaves

Resistive plate chambers

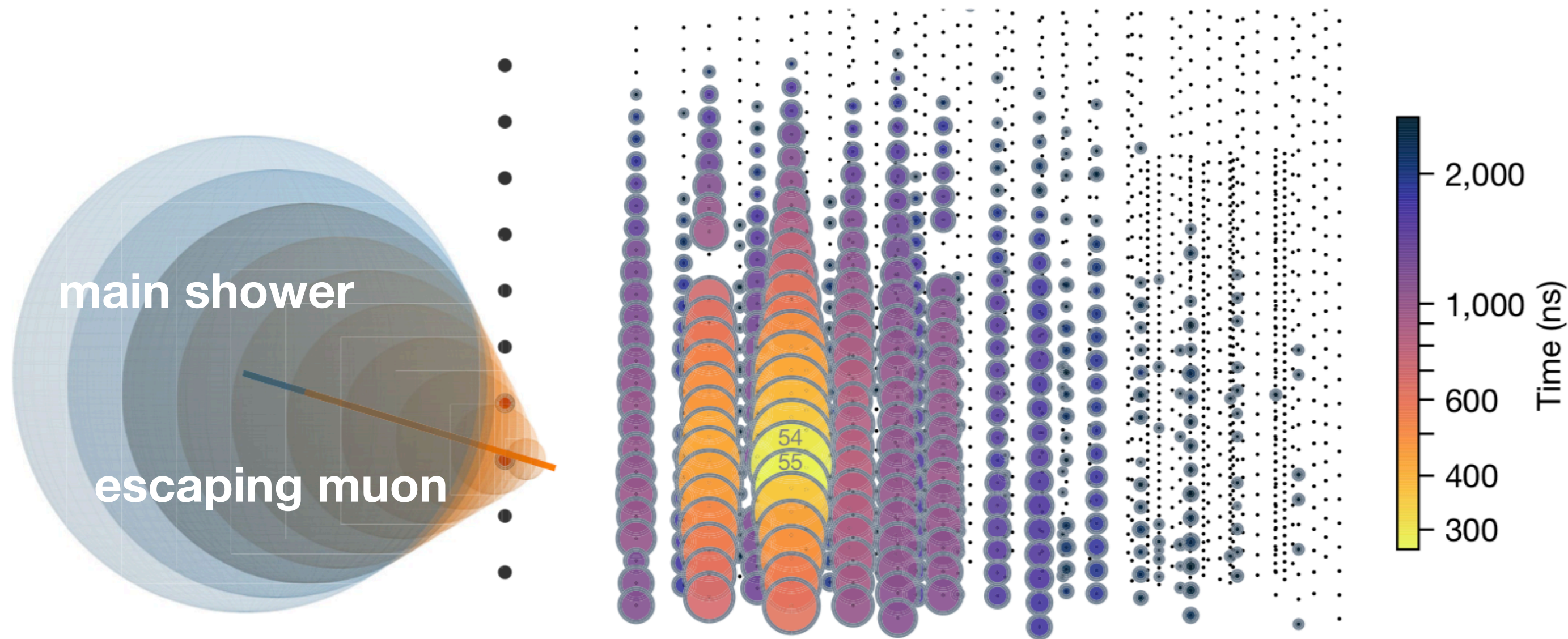
Currently stalled due to ecological concerns (blasting, excavation, ...)

Distinguishing ν vs $\bar{\nu}$?

- At high energies, neutrino/antineutrino separation is almost impossible
- Exception: the Glashow resonance.
At $E_{CM} = M_W$, or $E_\nu = 6.3$ PeV, can produce on-shell W for $\bar{\nu}_e$ only.

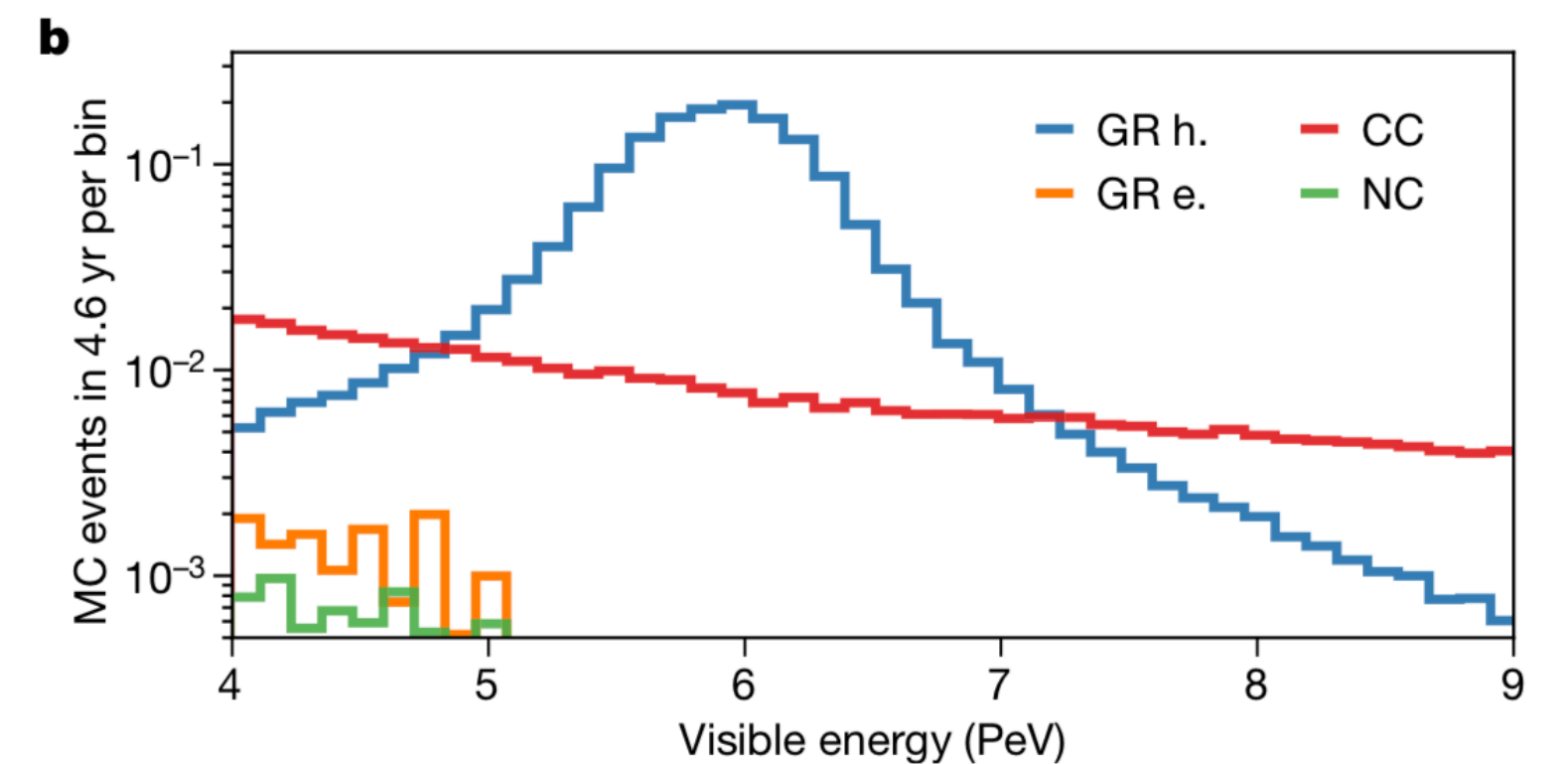
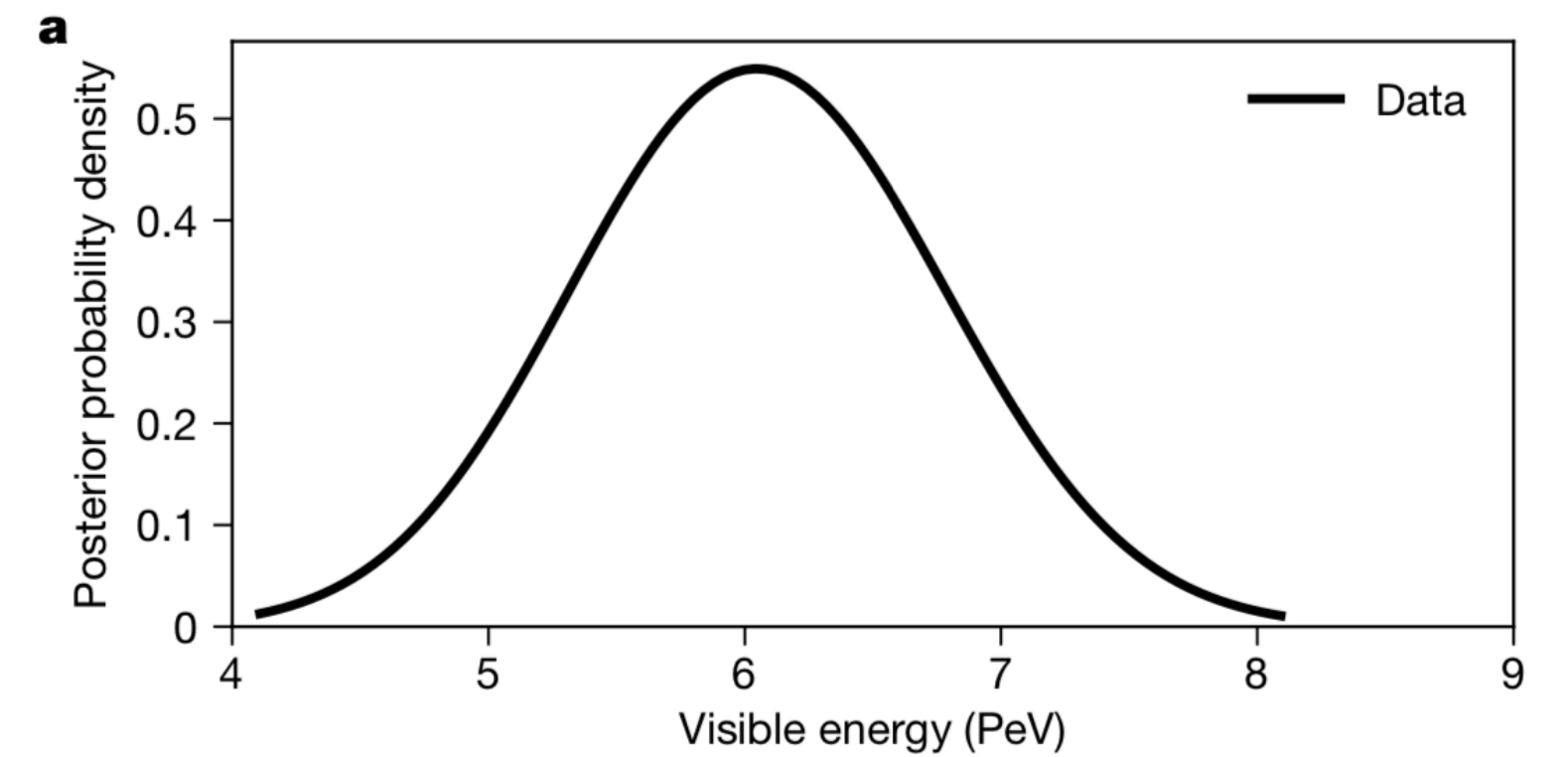
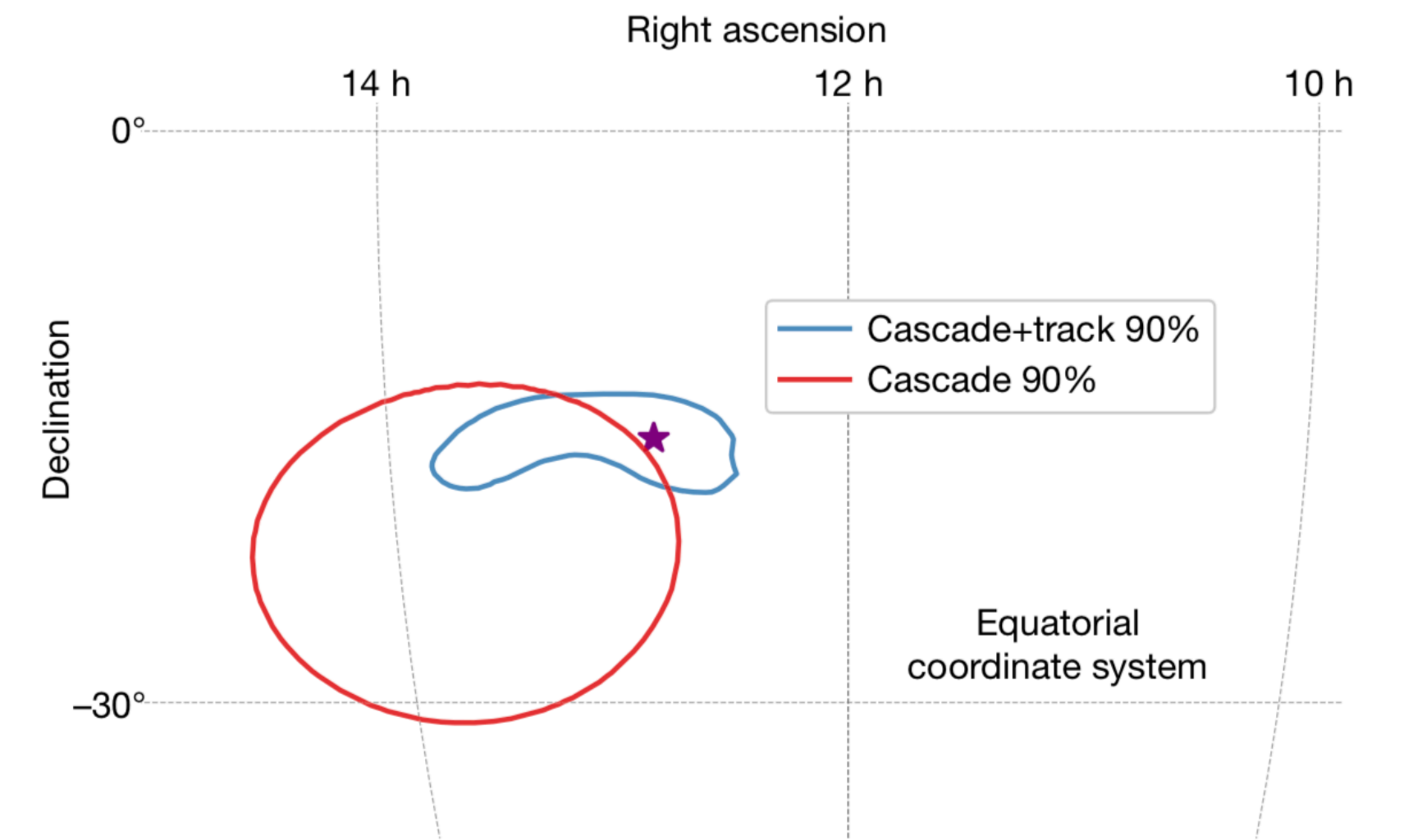


First Detection of Glashow Resonance

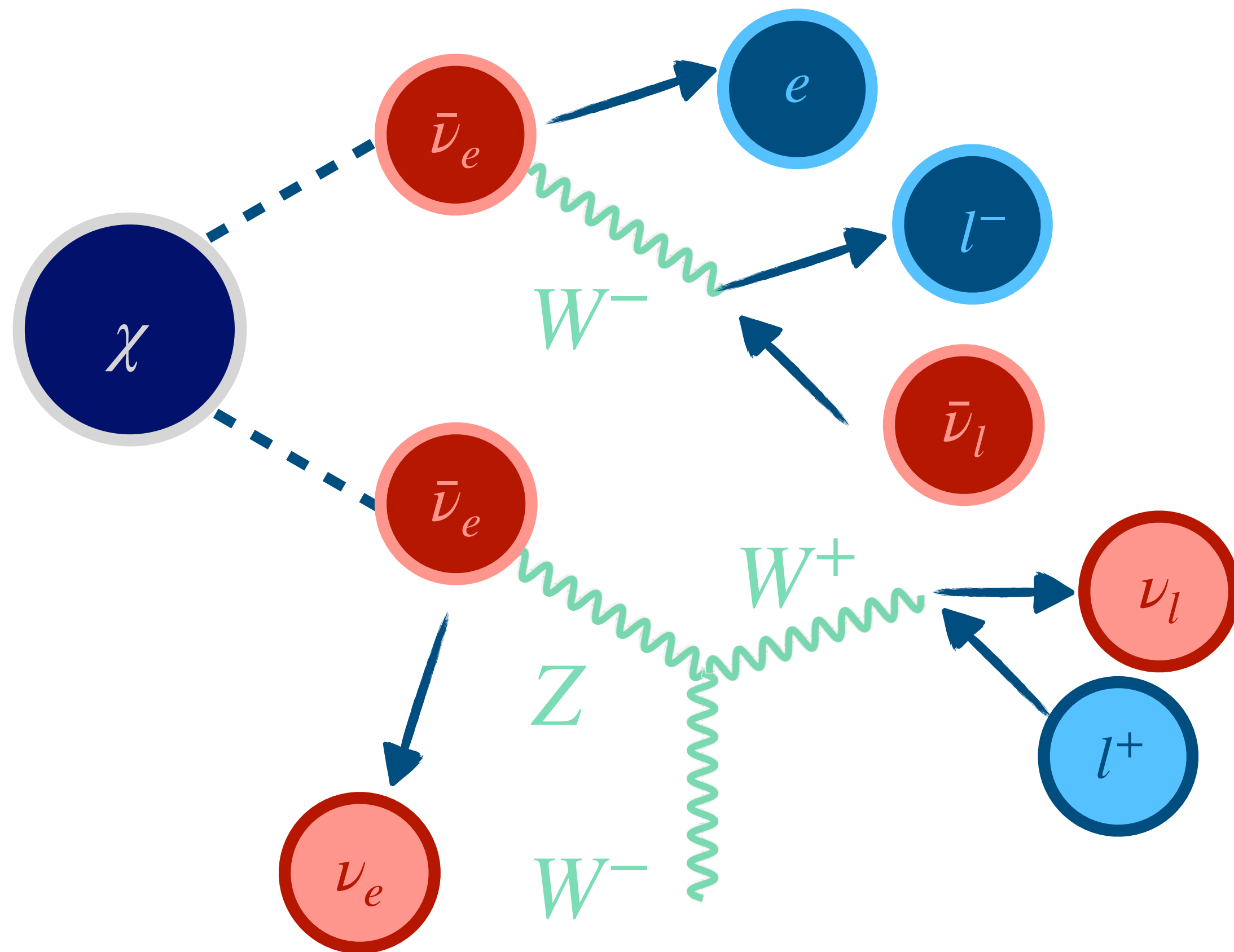


PeV energy partially-contained event selection

- Reconstructed energy of 6.05 ± 0.72 PeV.
- The detectable escaping muon suggests it's a hadronic shower.

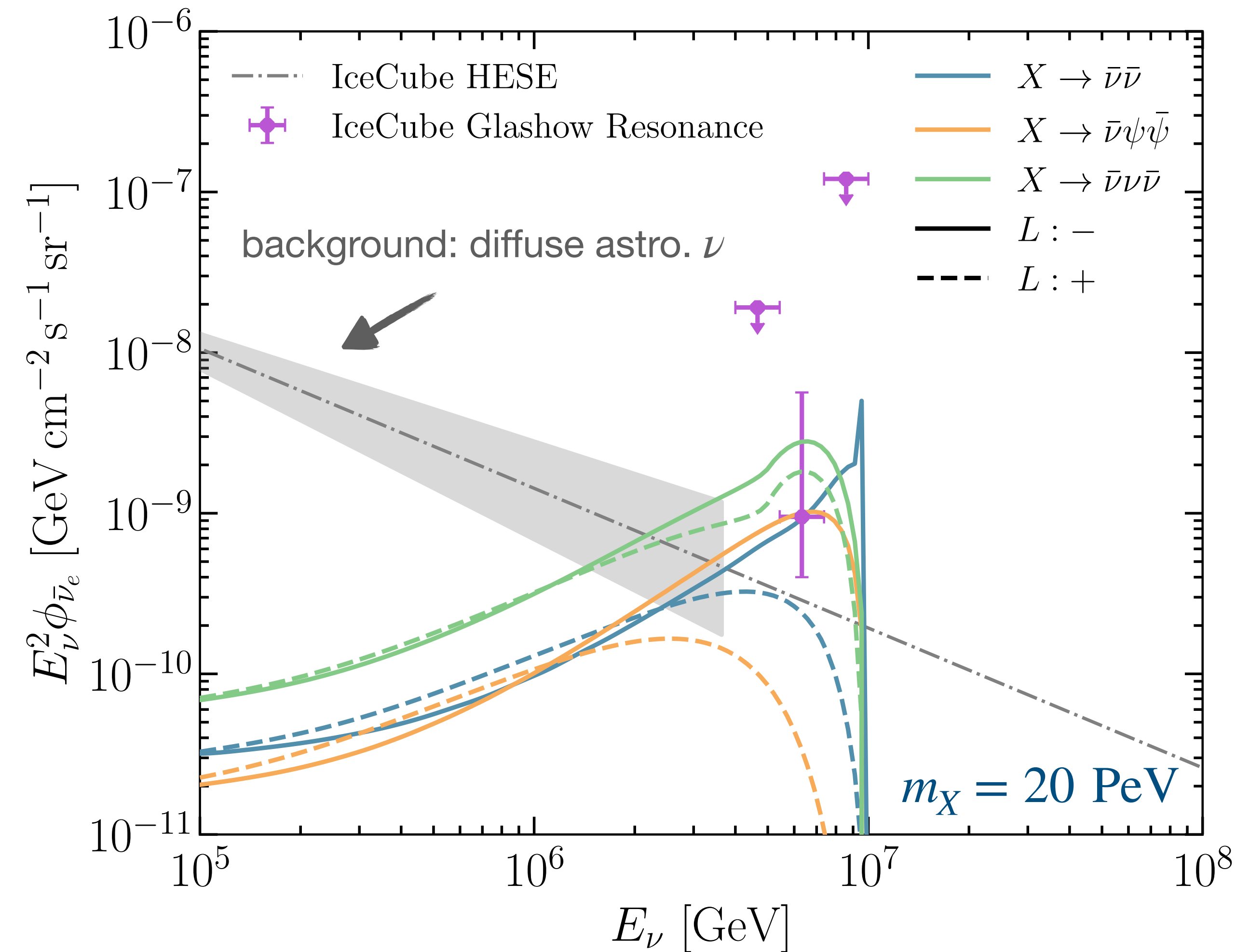


- Both ν and $\bar{\nu}$ can be produced no matter whether the lepton number is positive or negative.
- The spectrum $dN_{\bar{\nu}_i}^{\text{ch}}/dE_\nu$ becomes softer.

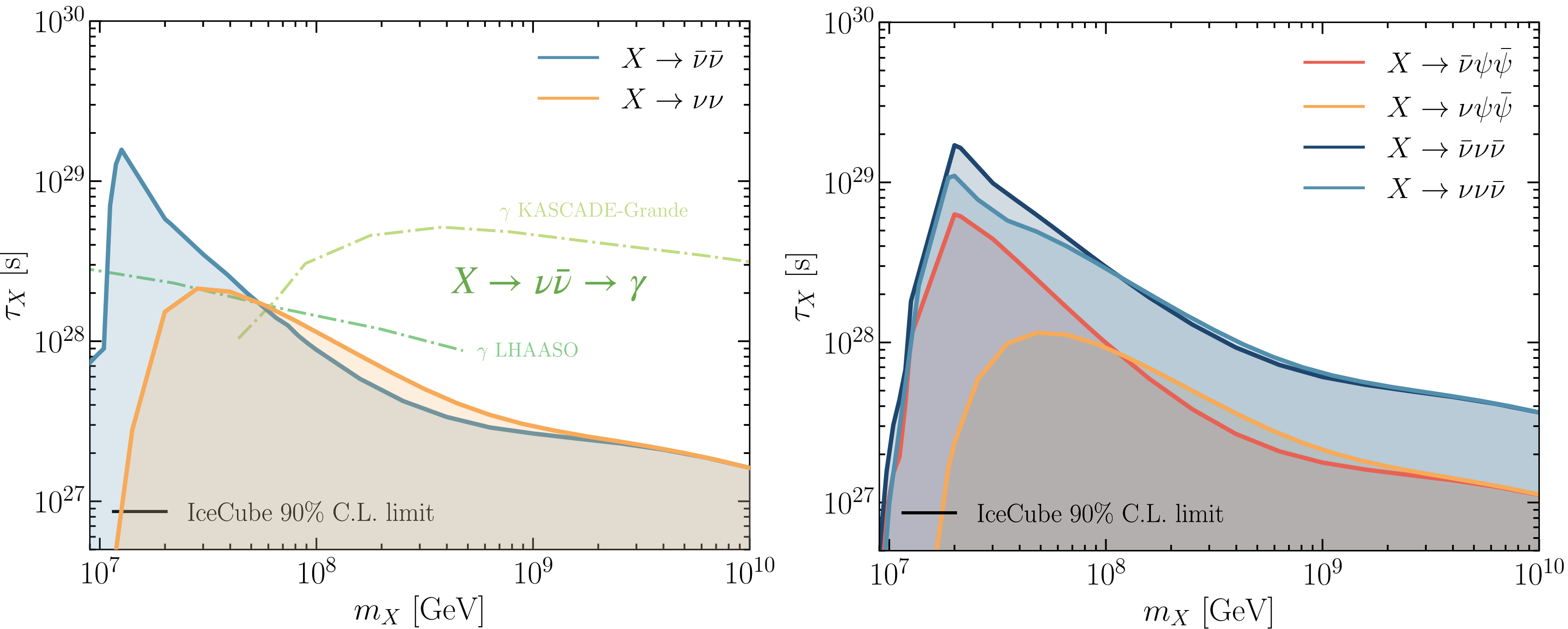


fragmentation functions from [HDMSpectra](#)

Bauer, Rodd, Webber 2007.15001

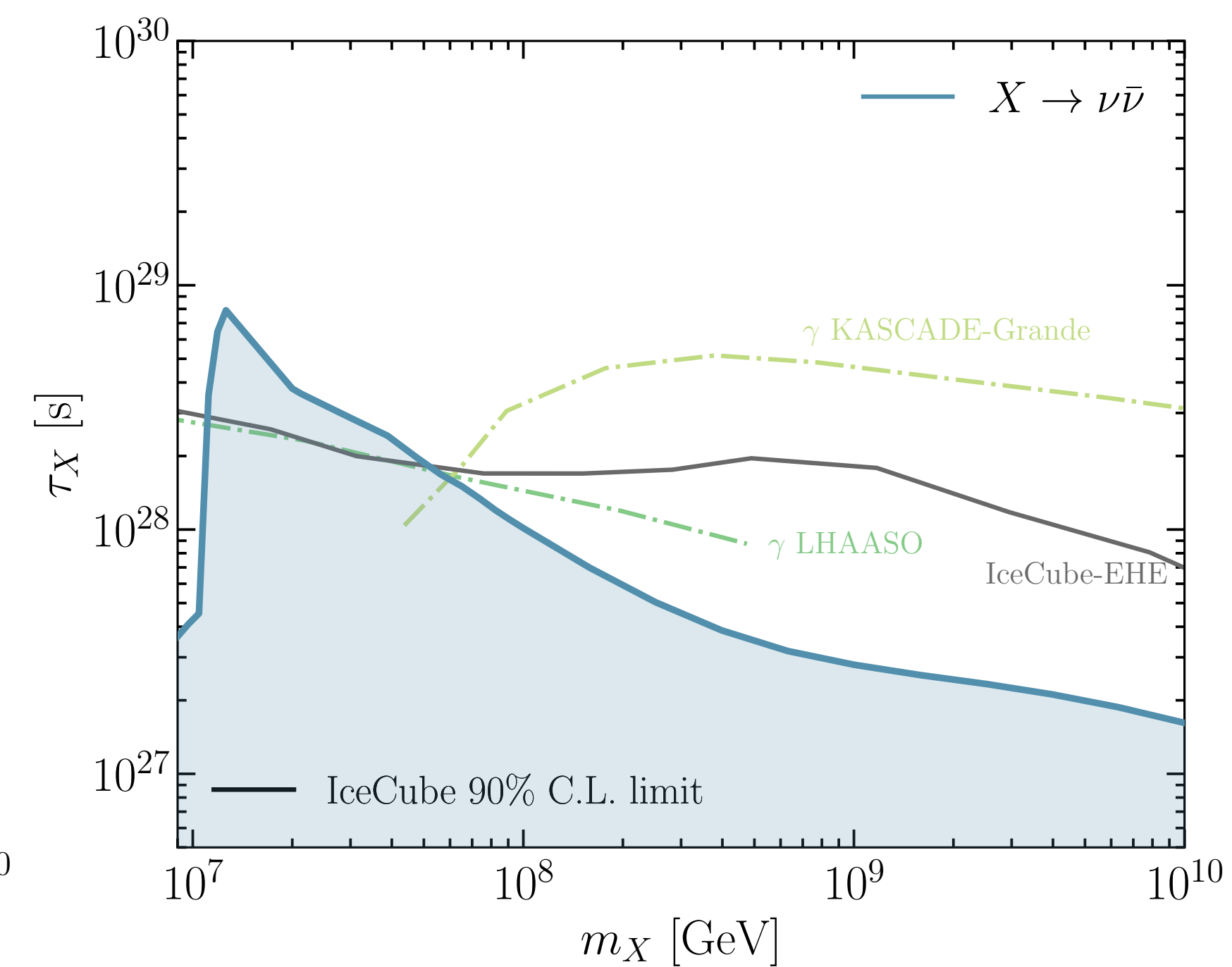
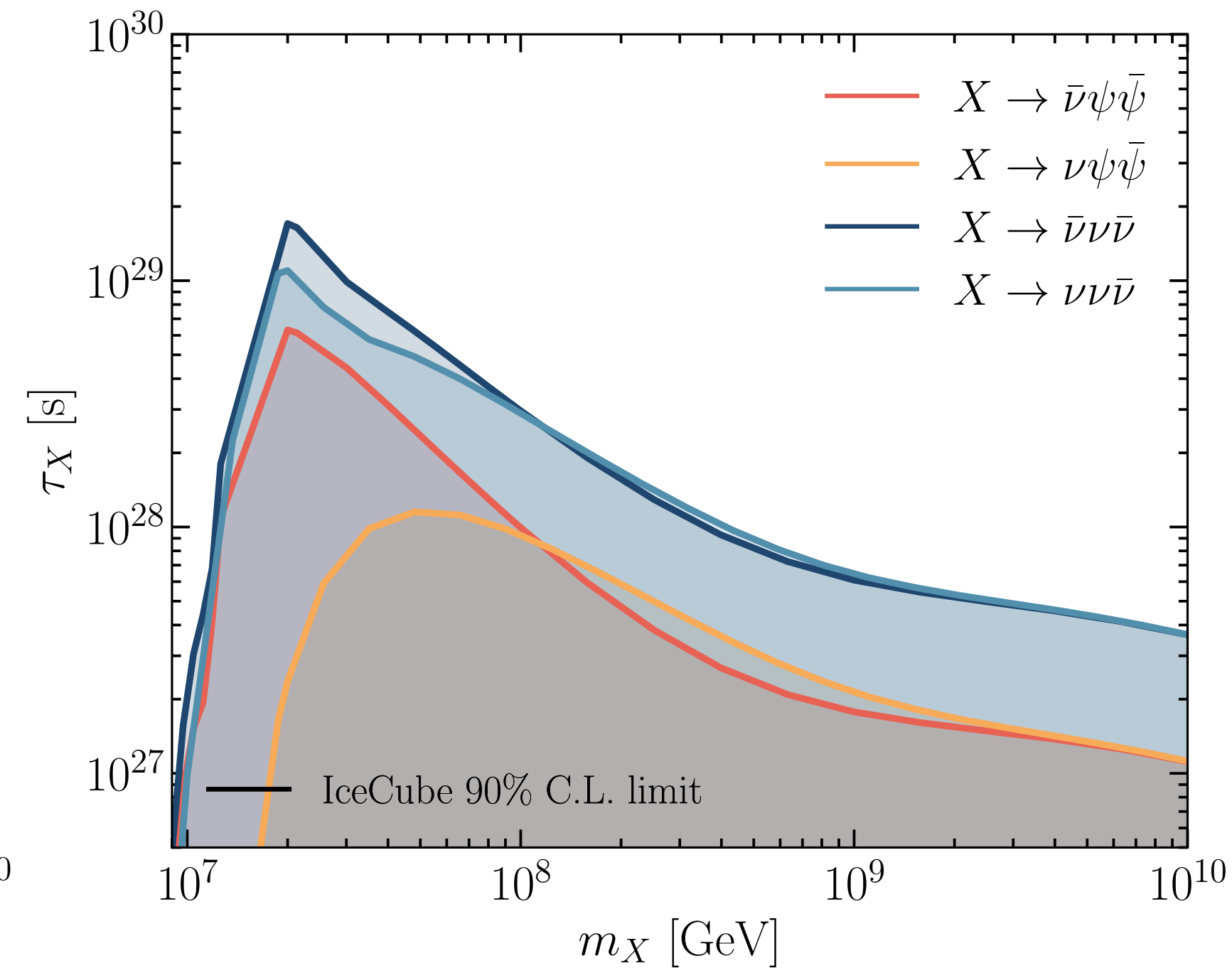
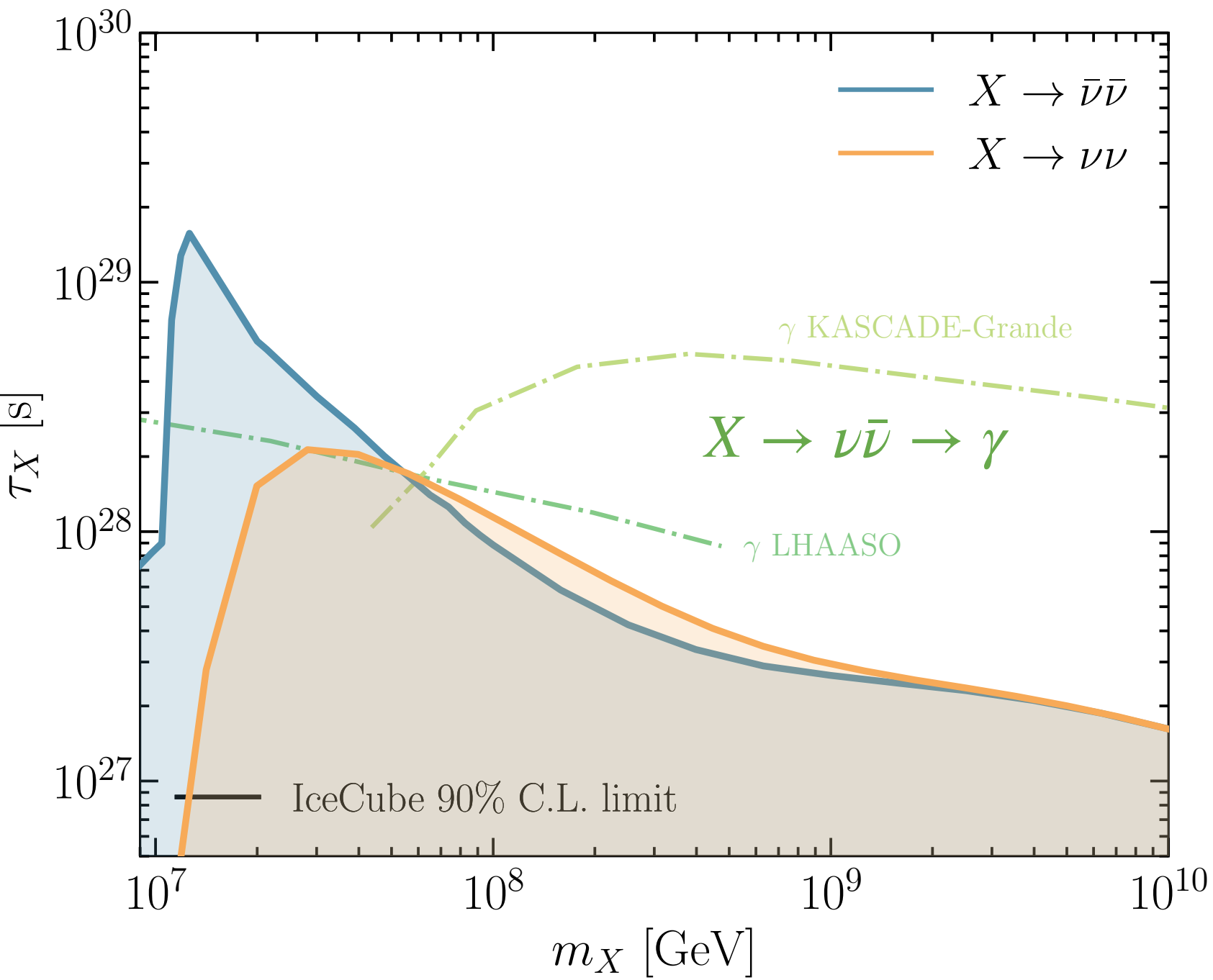


Constraints with Current Observation



- **Scenarios with positive/negative lepton numbers can be constrained respectively for $m_X \sim \text{PeV} - \text{EeV}$.**
- **The sensitivity of Glashow Resonance weakens when the number of decay products increases as $\nu : \bar{\nu} \rightarrow 1 : 1$.**

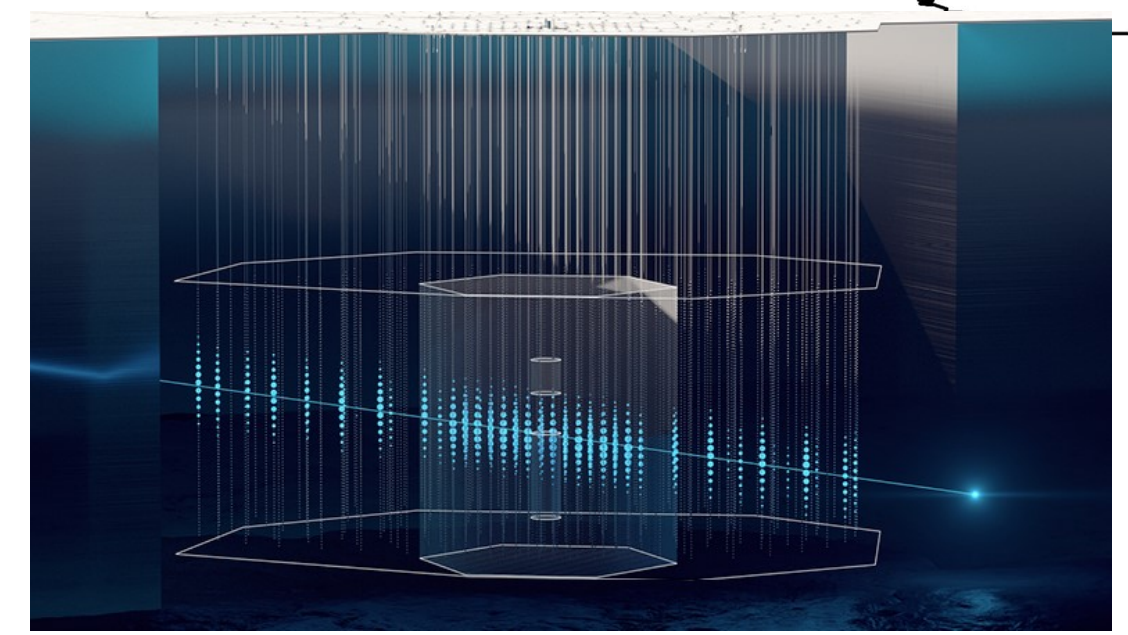
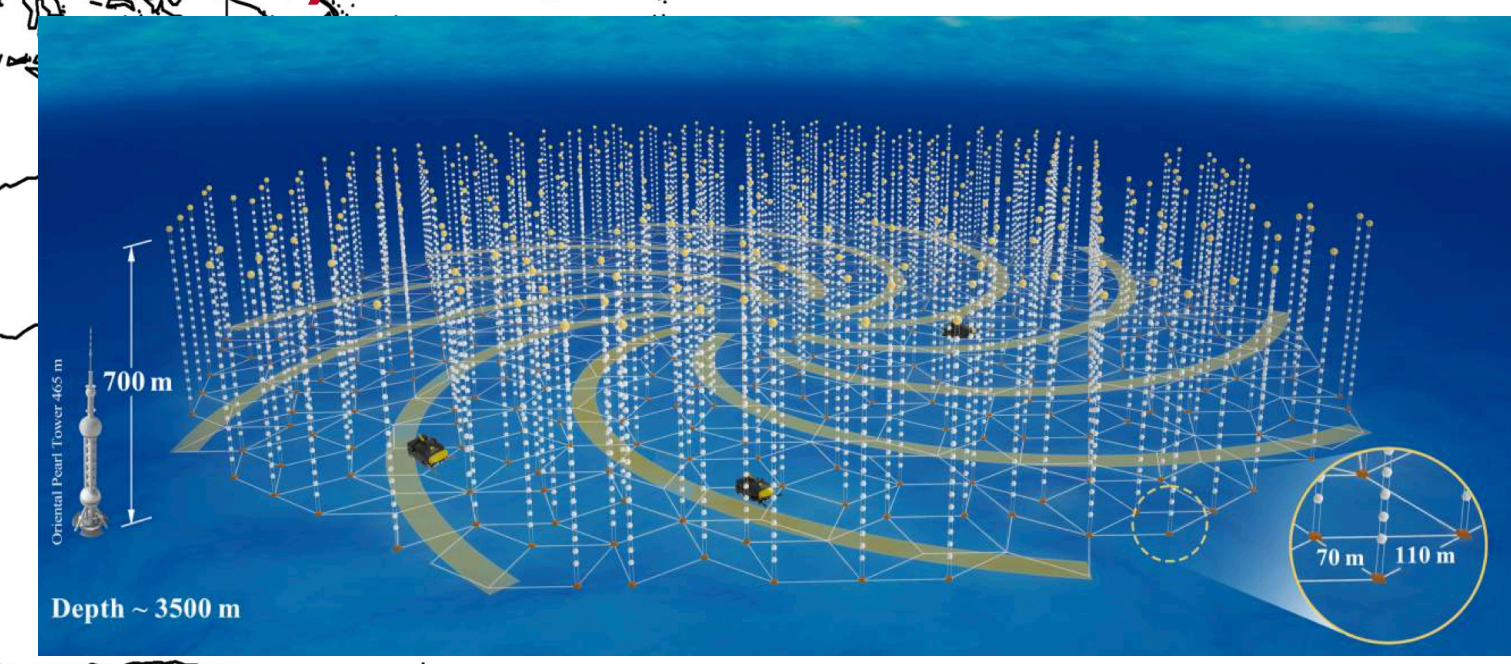
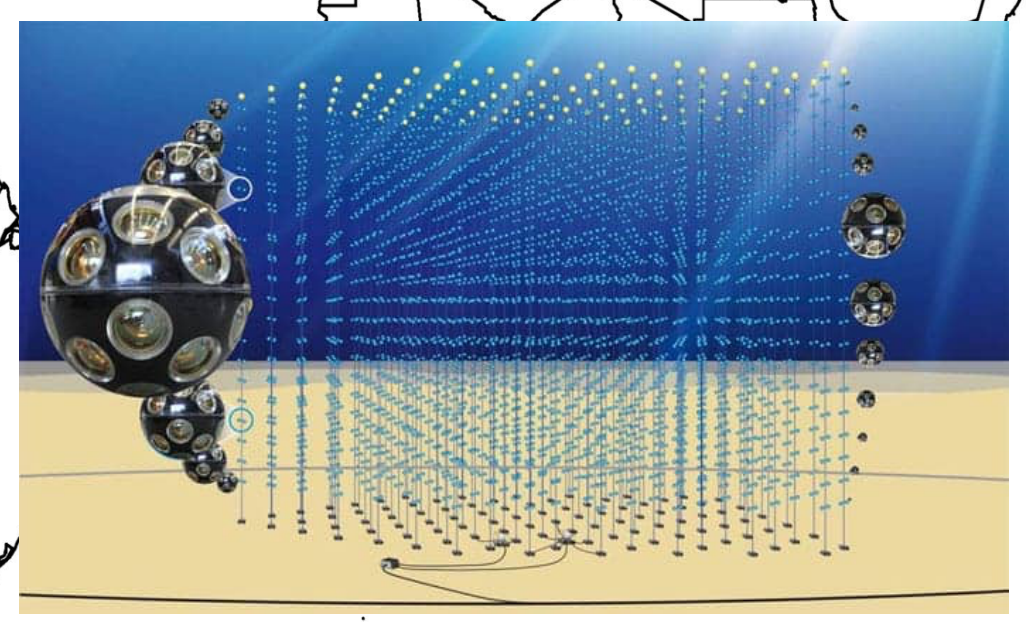
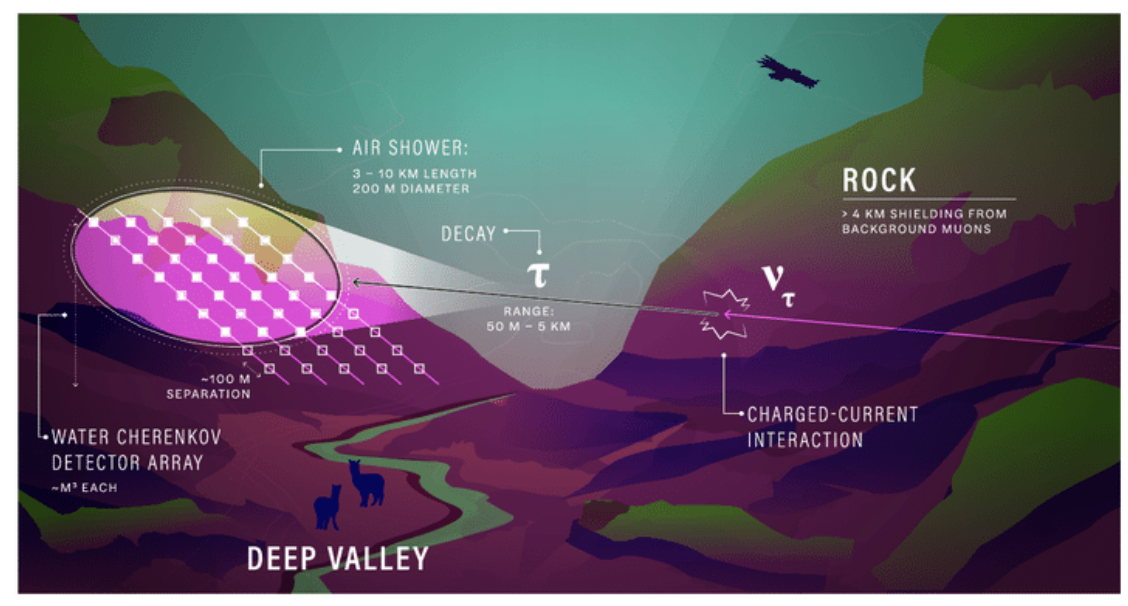
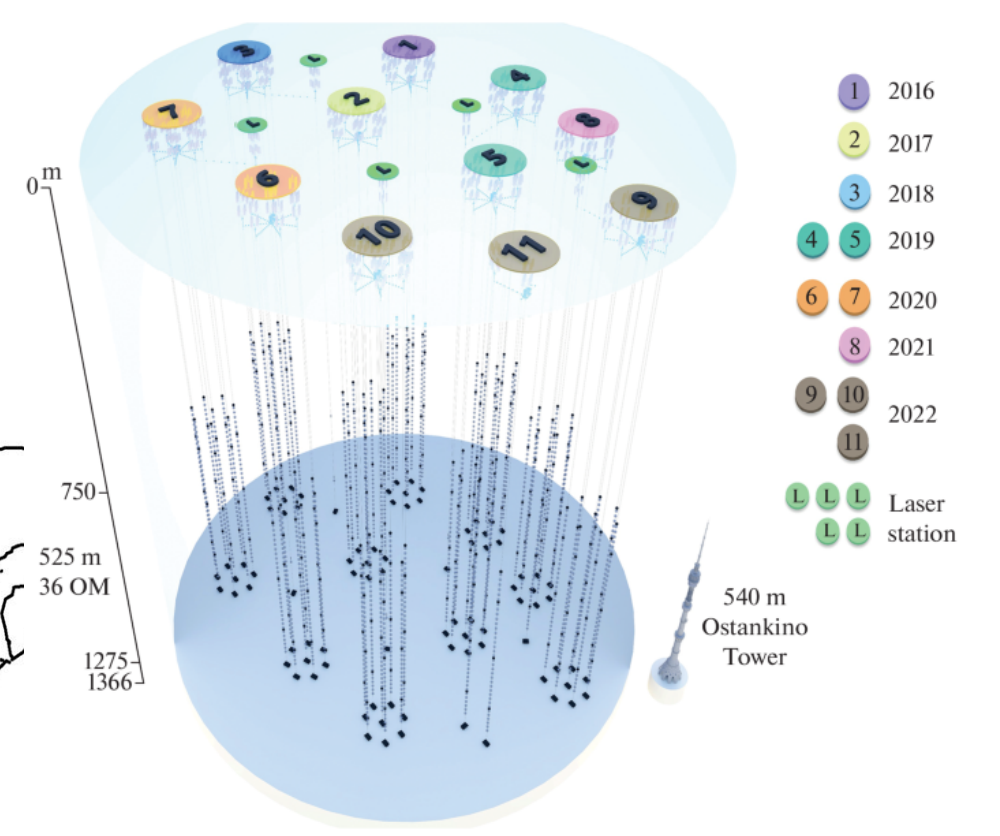
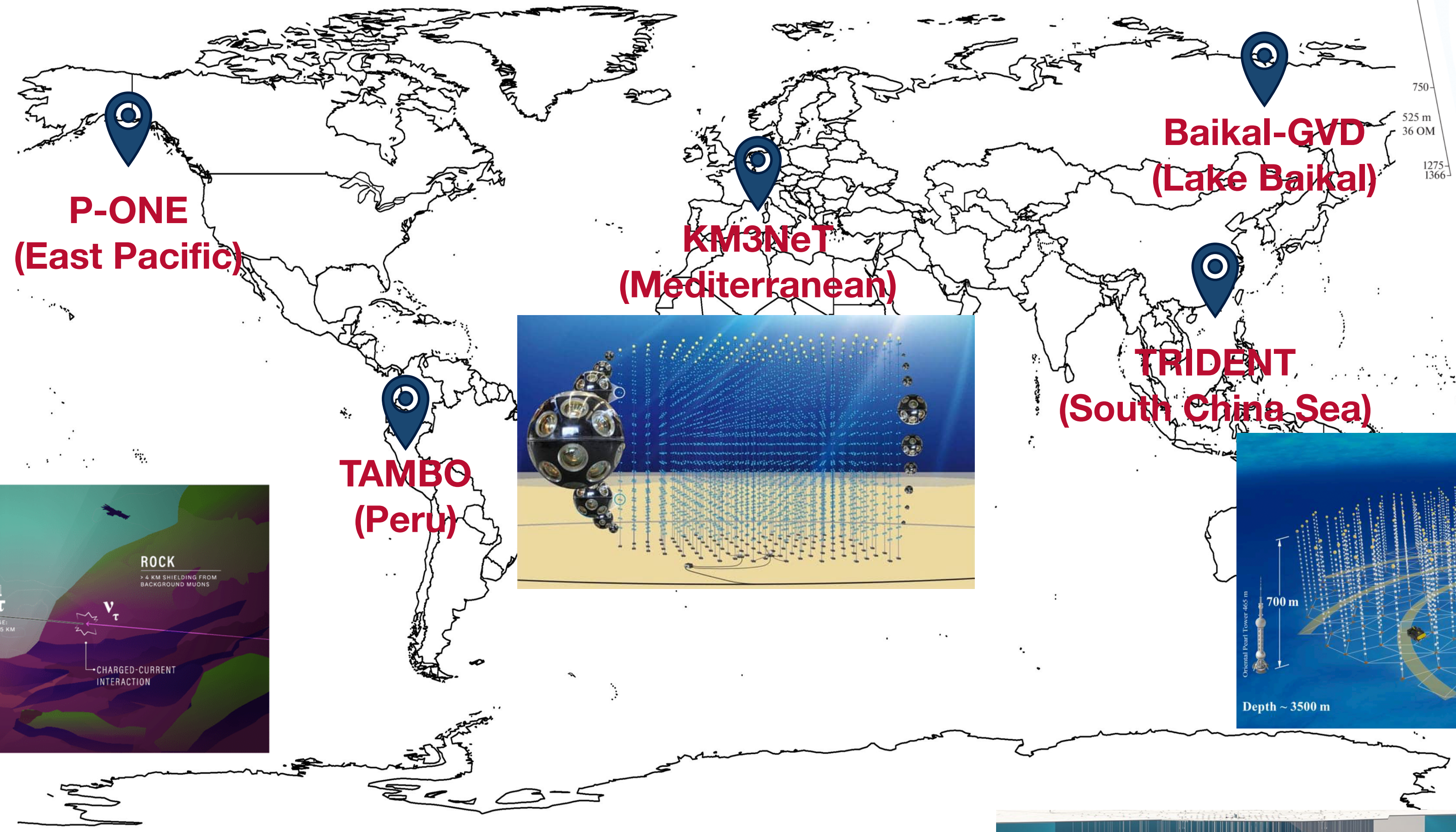
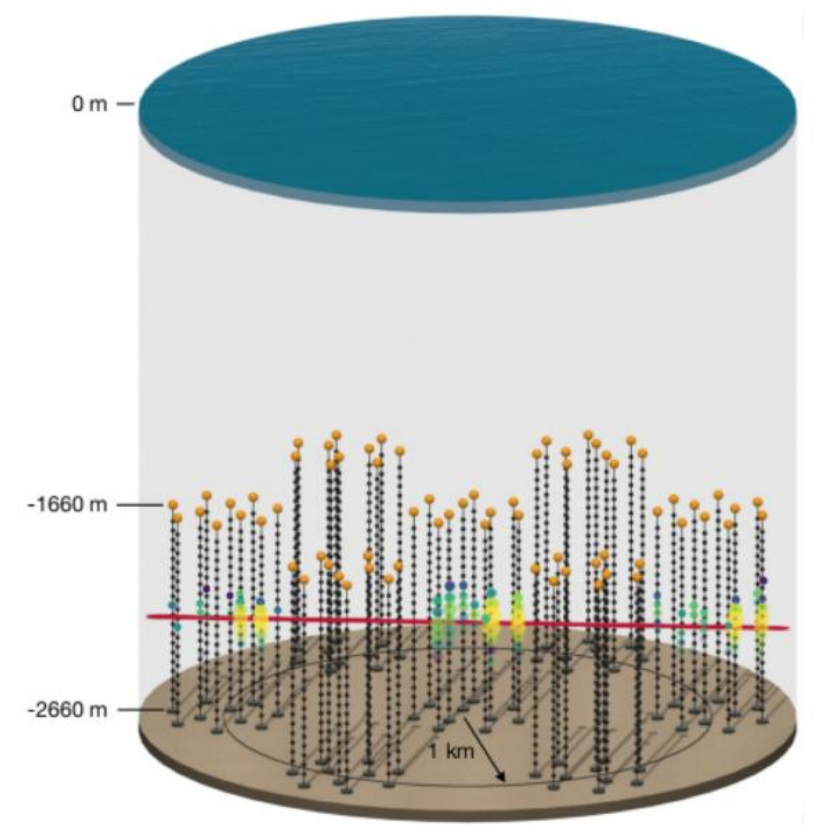
Constraints with Current Observation



- Scenarios with positive/negative lepton numbers can be constrained respectively for $m_X \sim \text{PeV} - \text{EeV}$.
- The sensitivity of Glashow Resonance weakens when the number of decay products increases as $\nu : \bar{\nu} \rightarrow 1 : 1$.

Better constraints on neutrino portal of symmetric DM decay for $m_X \sim 10 - 100 \text{ PeV}$

Next-Generation High-Energy Neutrino Telescopes



IceCube-Gen2
(South Pole)

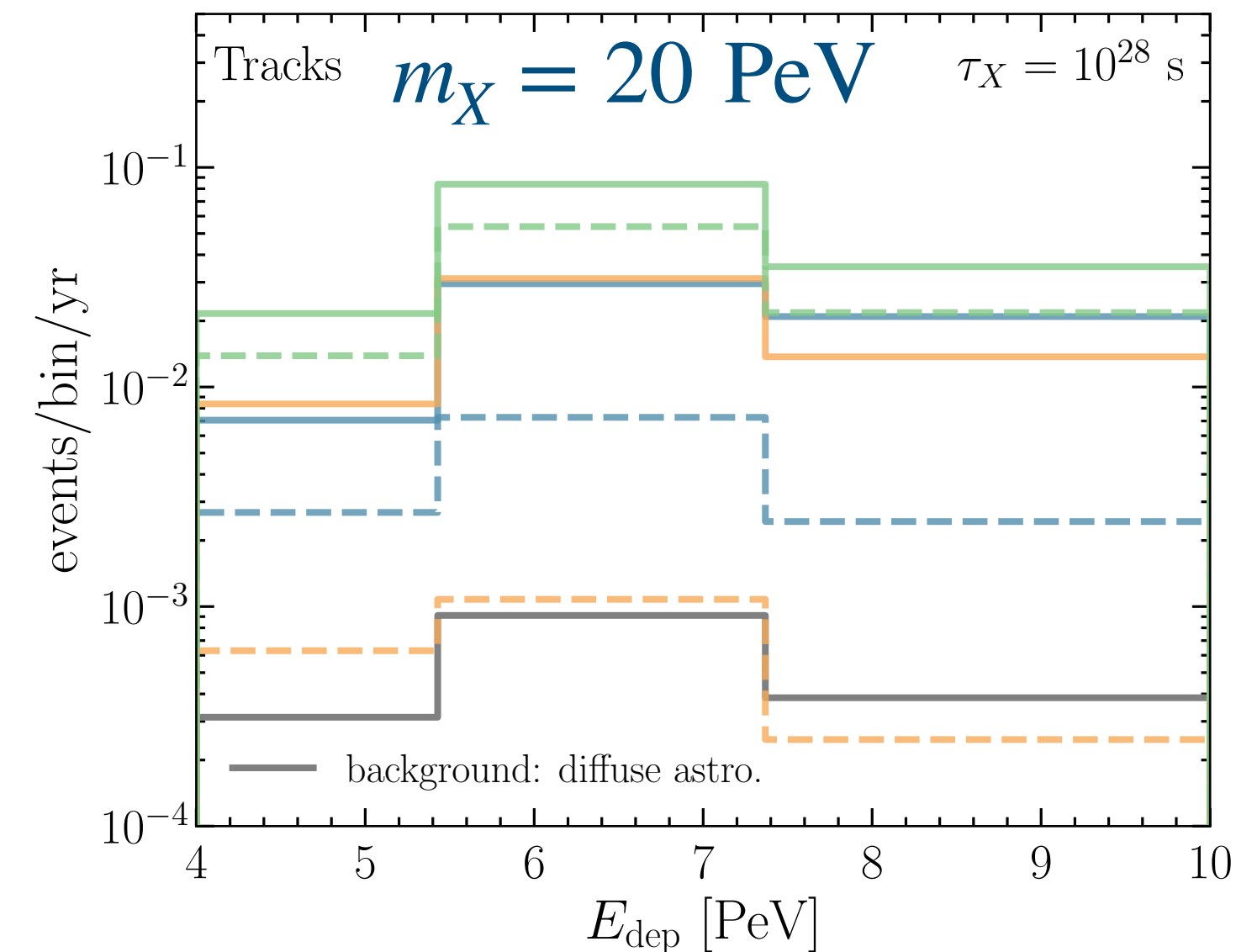
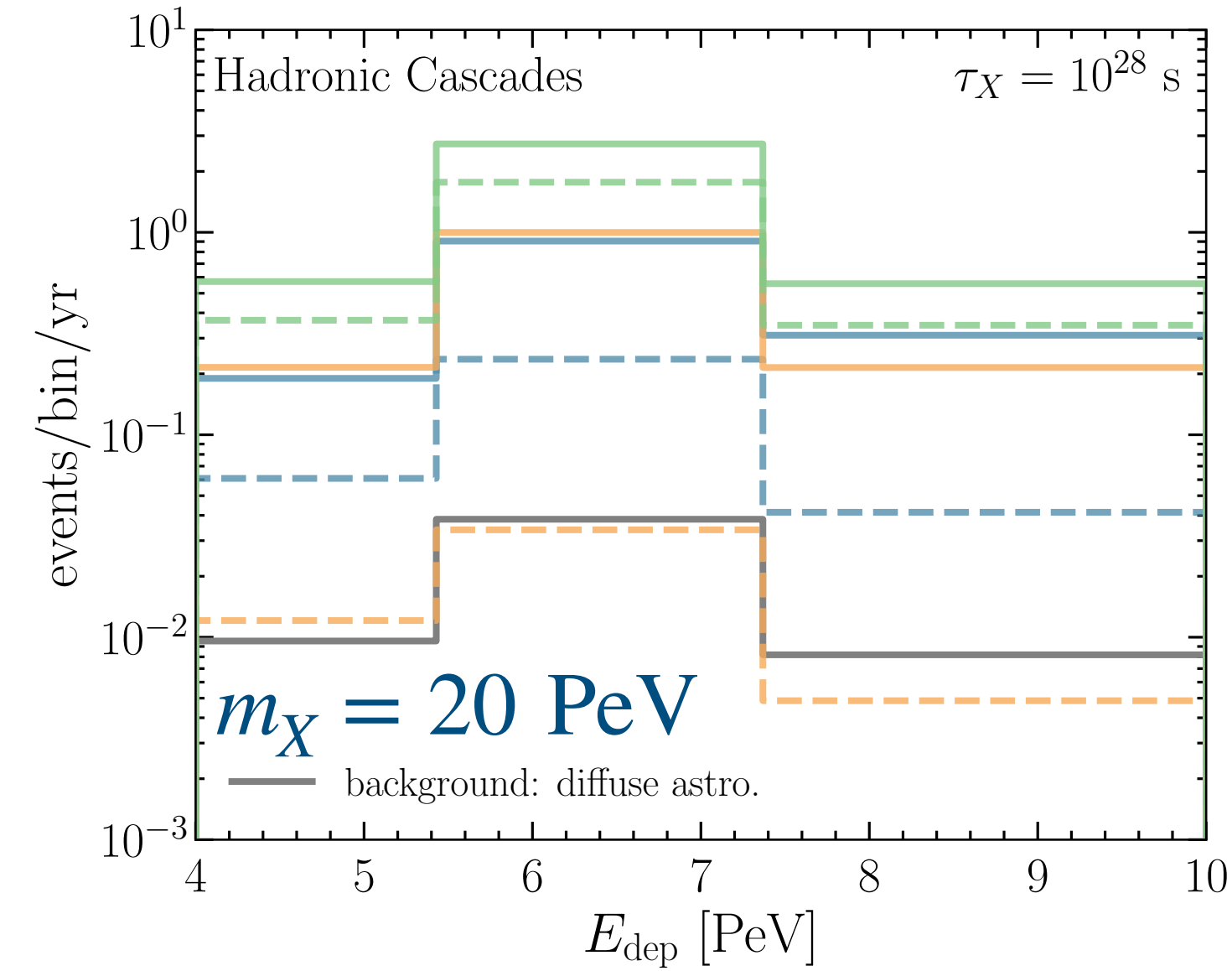
More telescopes with larger exposure!

Glashow Resonance Signal

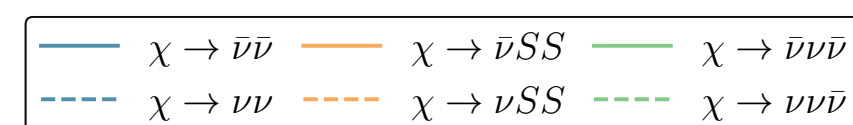
Glashow resonant events can be identified on an event-wise basis in the [4,10] PeV deposited energy window.

- ★ $W^- \rightarrow \text{hadrons}$ BR ~67 %
 - ✓ escaping muons, the only irreducible background is from neutral-current events
- ★ $W^- \rightarrow e^- \bar{\nu}_e / \tau^- \bar{\nu}_\tau$ BR ~11 %
 - ✗ Indistinguishable from a deep-inelastic-scattering cascade
- ★ $W^- \rightarrow \mu^- \bar{\nu}_\mu$ BR ~11 %
 - ✓ track without the initial cascade compared to ν_μ charged-current events

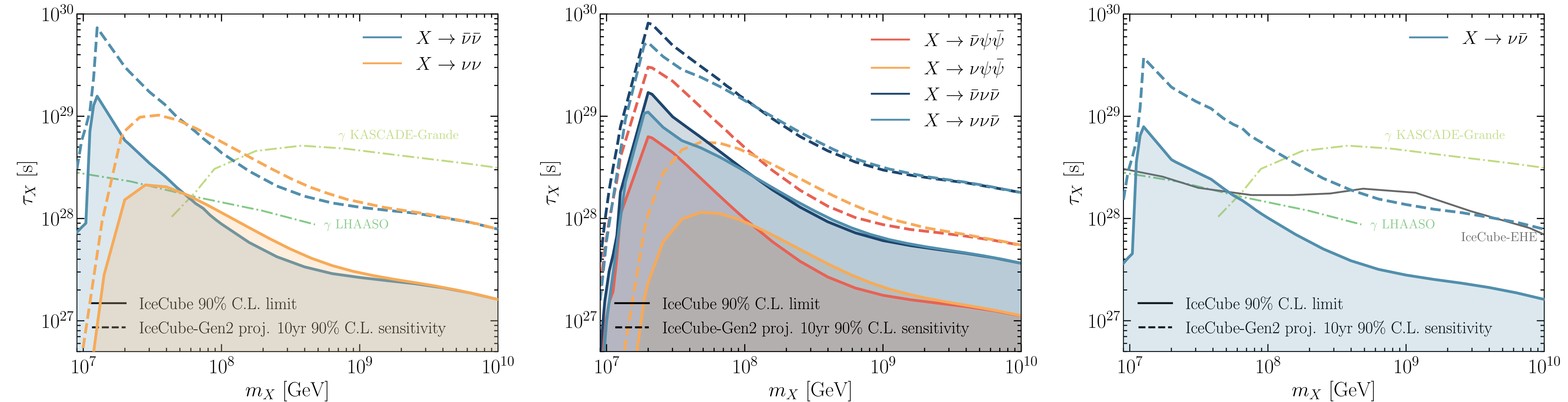
Liu, Song, ACV, 2304.06068



**Event rates of
Glashow
resonance at
IceCube as
partially
contained events**



Projected Sensitivities in the Future



- 90% C.L. sensitivities are estimated.
- Projected 10yr IceCube-Gen2 ($8 \times$ IceCube) sensitivities have lifetimes ~ 5 of current constraints.

Liu, Song, ACV, [2304.06068](#)

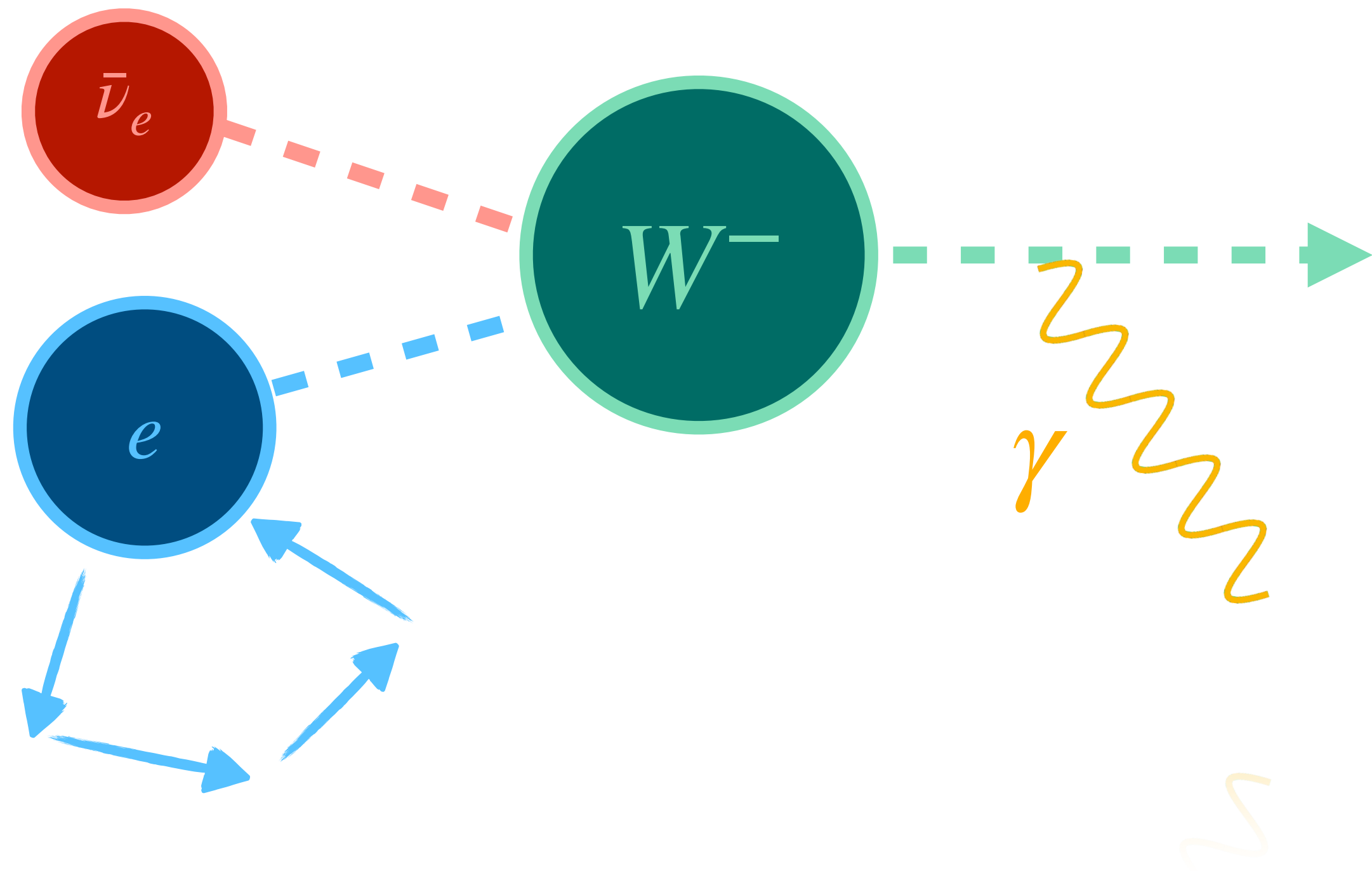
(not really a) Summary

- Even though neutrinos are the most invisible channel, neutrino telescopes cover at least **14 orders of magnitude in energy** & can say all sorts of things about the dark sector & new physics.
- Electroweak cross section rises with energy, but events become rarer (big detectors)
- Electroweak corrections means that even if the dark matter tries to cloak itself with the neutrino... it cannot hide forever.



Corrected Cross Section

subleading effects that affect the cross section



Atomic e motion:
Doppler Broadening

Initial State Radiation

