

QCD Axion-mediated Dark Matter

arXiv:2306.03145 (JHEP 2023)

Collaborators : Jeff Dror² Stefania Gori¹

Speaker : Pankaj Munbodh¹

¹Department of Physics
University of California Santa Cruz

²Department of Physics
University of Florida

October 18, 2024

UC SANTA CRUZ

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term
 $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term
 $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Introduce global $U(1)_{\text{PQ}}$ symmetry which gets broken at high scale f_a
generating axion a as a Goldstone boson. **Weinberg, Wilczek, Peccei, Quinn**

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term
 $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Introduce global $U(1)_{\text{PQ}}$ symmetry which gets broken at high scale f_a
generating axion a as a Goldstone boson. **Weinberg, Wilczek, Peccei, Quinn**

Below QCD confinement scale Λ_{QCD} , instantons generate a potential for
 a , which now acquires a small mass $m_a \approx \frac{\Lambda_{\text{QCD}}^2}{f_a}$.

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Introduce global $U(1)_{\text{PQ}}$ symmetry which gets broken at high scale f_a generating axion a as a Goldstone boson. **Weinberg, Wilczek, Peccei, Quinn**

Below QCD confinement scale Λ_{QCD} , instantons generate a potential for a , which now acquires a small mass $m_a \approx \frac{\Lambda_{\text{QCD}}^2}{f_a}$.

Minimization of the generated potential causes a to get a VEV that cancels the theta term.

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Introduce global $U(1)_{\text{PQ}}$ symmetry which gets broken at high scale f_a generating axion a as a Goldstone boson. **Weinberg, Wilczek, Peccei, Quinn**

Below QCD confinement scale Λ_{QCD} , instantons generate a potential for a , which now acquires a small mass $m_a \approx \frac{\Lambda_{\text{QCD}}^2}{f_a}$.

Minimization of the generated potential causes a to get a VEV that cancels the theta term.

Option 1: Axions can behave as DM, for $f_a \gtrsim 10^{11}$ GeV.

Motivation

Neutron EDM is tightly constrained \rightarrow QCD CP-violating θ term $-\theta \frac{\alpha_s}{8\pi} G\tilde{G}$ is small (or zero) $\bar{\theta} \lesssim 10^{-10}$ (**strong CP problem**).

Introduce global $U(1)_{\text{PQ}}$ symmetry which gets broken at high scale f_a generating axion a as a Goldstone boson. **Weinberg, Wilczek, Peccei, Quinn**

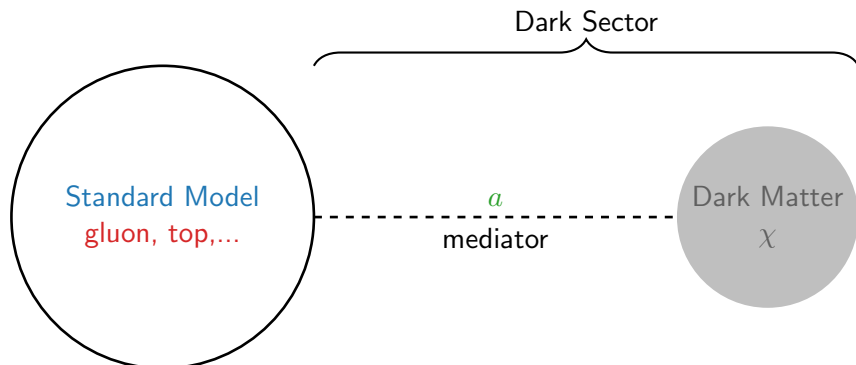
Below QCD confinement scale Λ_{QCD} , instantons generate a potential for a , which now acquires a small mass $m_a \approx \frac{\Lambda_{\text{QCD}}^2}{f_a}$.

Minimization of the generated potential causes a to get a VEV that cancels the theta term.

Option 1: Axions can behave as DM, for $f_a \gtrsim 10^{11}$ GeV.

Option 2: QCD axion can be the mediator between DM and SM for smaller 10^9 GeV $\lesssim f_a \lesssim 10^{11}$ GeV.

Minimal Setup



Other studies : [2209.03932](#) Bharucha et al. (JHEP 2022)
[2306.03128](#) Fitzpatrick et al. (PRD 2023)
[2207.02221](#) Coffey et al. (PRD 2022)

Model

$$\mathcal{L} \supset \frac{c_\chi}{2f_a} \partial_\mu a \bar{\chi} \gamma^\mu \gamma^5 \chi + \frac{c_{\psi_i}}{2f_a} \partial_\mu a \bar{\psi}_i \gamma^\mu \gamma^5 \psi_i + \frac{c_\gamma}{4f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

ψ_i = SM leptons (e) and up/down-type quarks (u/d).

c_χ : axion-DM coupling.

$c_{\psi_i} = c_e, c_u, c_d$ (axion-matter couplings).

c_γ : axion-photon coupling.

$g_{a\chi} \equiv \frac{c_\chi m_\chi}{f_a}$, m_χ and $f_a \rightarrow$ parameters of the model.

DFSZ : SM matter-axion couplings depend on an angle β .

M. Dine et al. (1981), Zhitnitsky (1980)

Model

$$\mathcal{L} \supset \frac{c_\chi}{2f_a} \partial_\mu a \bar{\chi} \gamma^\mu \gamma^5 \chi + \frac{c_{\psi_i}}{2f_a} \partial_\mu a \bar{\psi}_i \gamma^\mu \gamma^5 \psi_i + \frac{c_\gamma}{4f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

ψ_i = SM leptons (e) and up/down-type quarks (u/d).

c_χ : axion-DM coupling.

$c_{\psi_i} = c_e, c_u, c_d \xrightarrow{\alpha_s} c_p, c_n$ (axion-matter couplings).

c_γ : axion-photon coupling.

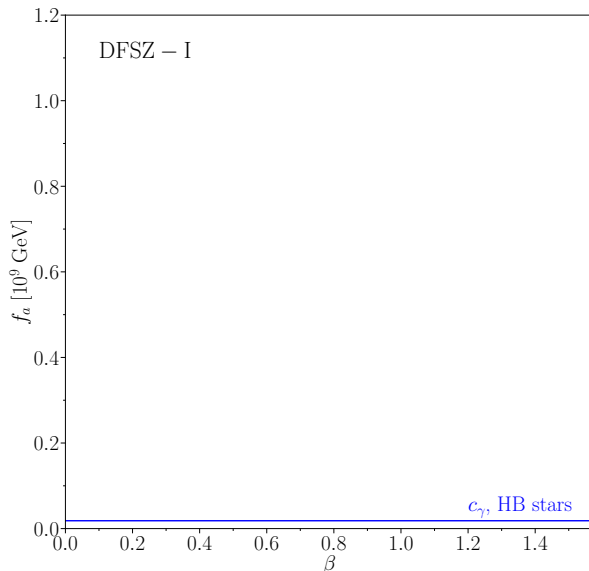
$g_{a\chi} \equiv \frac{c_\chi m_\chi}{f_a}$, m_χ and $f_a \rightarrow$ parameters of the model.

DFSZ : SM matter-axion couplings depend on an angle β .

M. Dine et al. (1981), Zhitnitsky (1980)

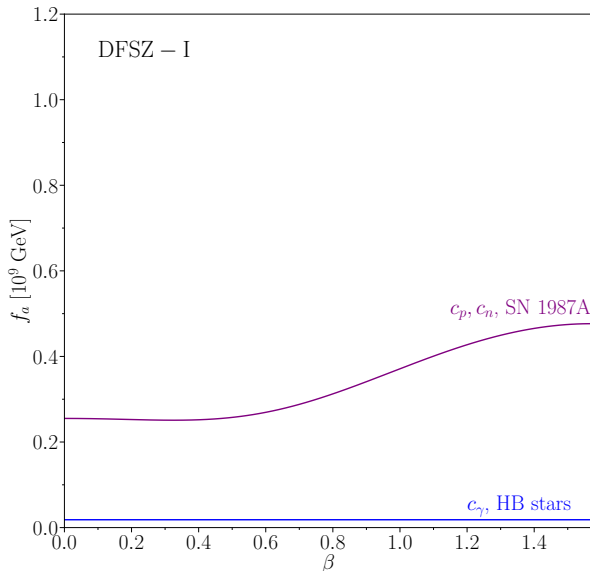
Experimental Constraints I (Stellar Cooling)

Axions \rightarrow additional cooling of stars \rightarrow constrain axion couplings.



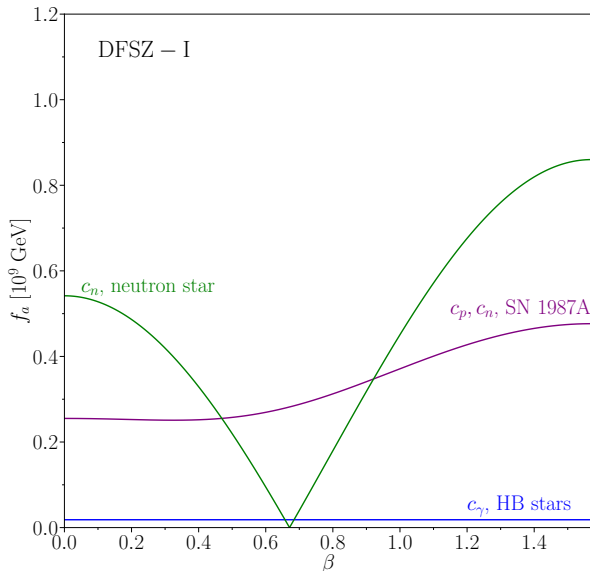
Experimental Constraints I (Stellar Cooling)

Axions \rightarrow additional cooling of stars \rightarrow constrain axion couplings.



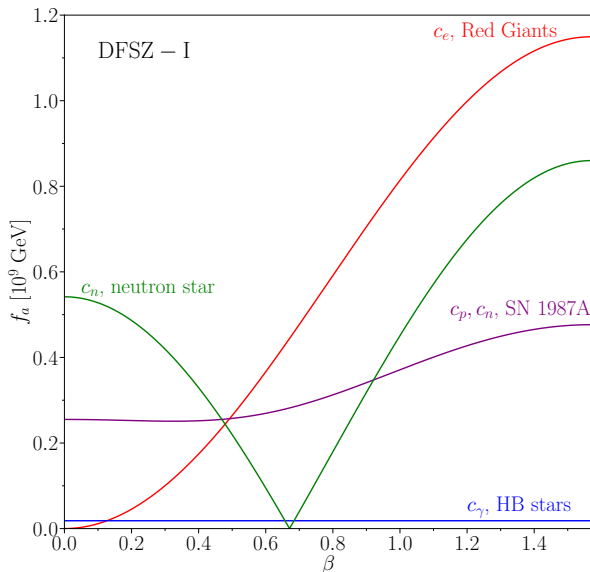
Experimental Constraints I (Stellar Cooling)

Axions \rightarrow additional cooling of stars \rightarrow constrain axion couplings.



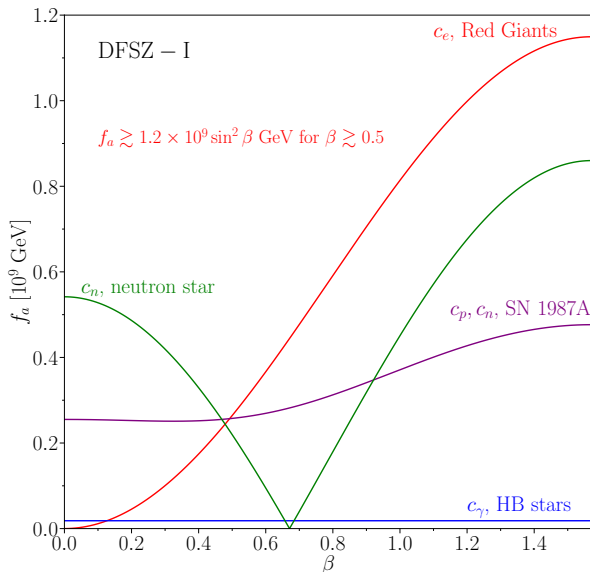
Experimental Constraints I (Stellar Cooling)

Axions \rightarrow additional cooling of stars \rightarrow constrain axion couplings.

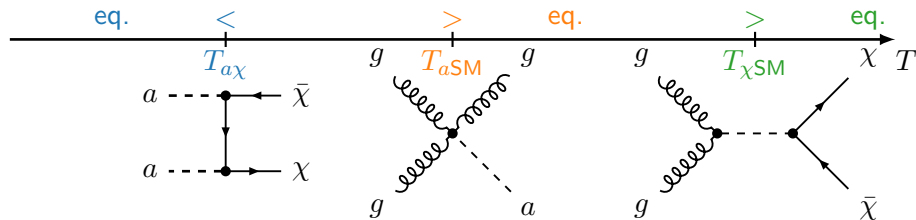


Experimental Constraints I (Stellar Cooling)

Axions \rightarrow additional cooling of stars \rightarrow constrain axion couplings.



Thermalization I

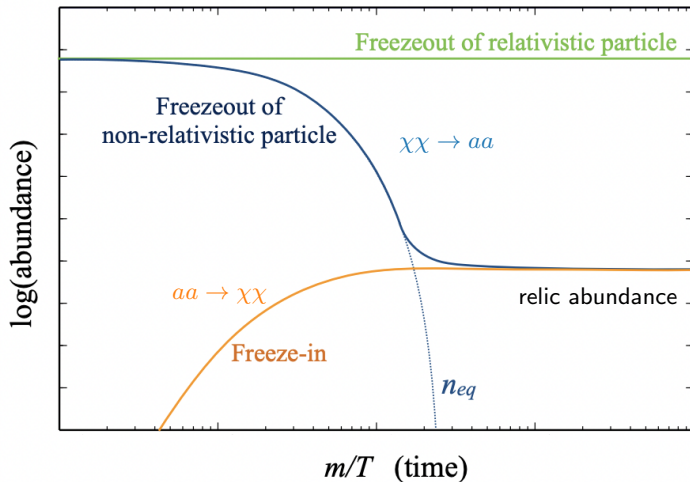


three possible hierarchies: $T_{a\chi} \ll T_{aSM} \ll T_{\chi SM}$ (as shown),
 $T_{aSM} \ll T_{a\chi} \ll T_{\chi SM}$, and $T_{aSM} \ll T_{\chi SM} \ll T_{a\chi}$ depending on the size of
 $g_{a\chi}$ and f_a .

+ reheat temperature $T_{RH} \rightarrow$ cosmological history \rightarrow dominant
 production mechanism.

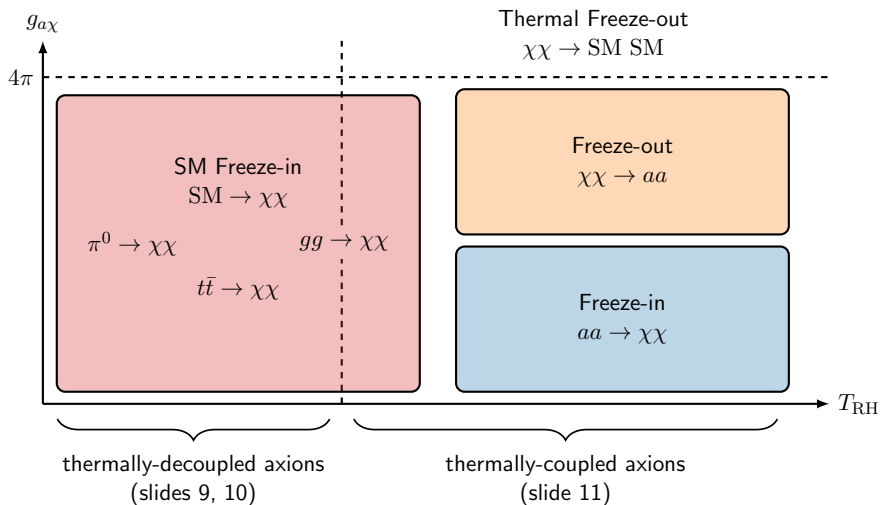
Pedestrian's guide to DM production

T. Lin arXiv 1904.07915



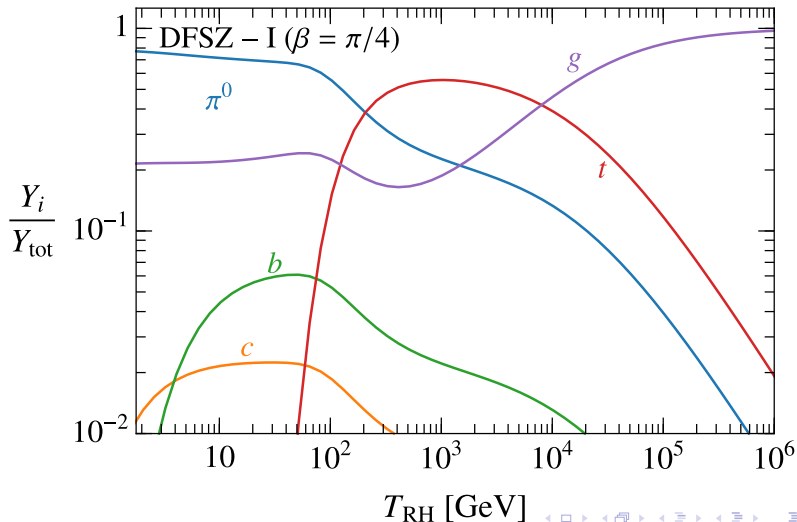
A Bird's Eye view

arXiv:2306.03145 (Dror, Gori and Munbodh)



Thermally decoupled axions I

$T_{RH} < T_{aSM}$. Freeze-in : $\pi^0 \rightarrow \chi\chi$, $gg \rightarrow \chi\chi$, $t\bar{t} \rightarrow \chi\chi \dots$



Thermally decoupled axions II

$$T_{\text{RH}} < T_{a\text{SM}}$$

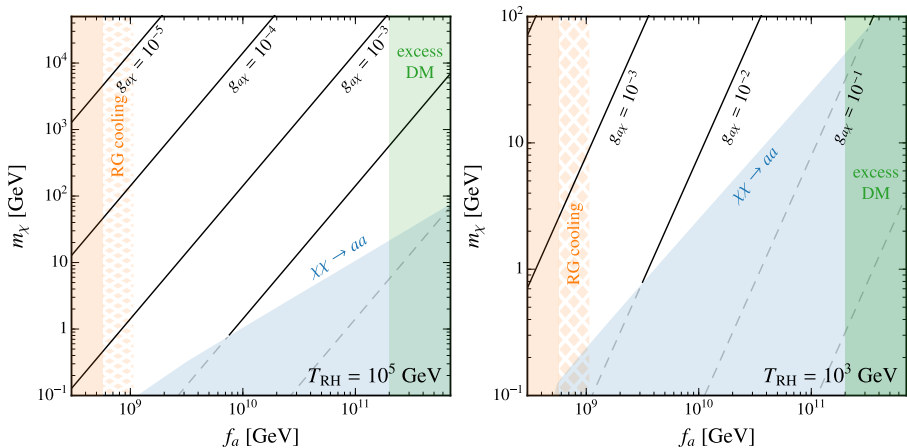
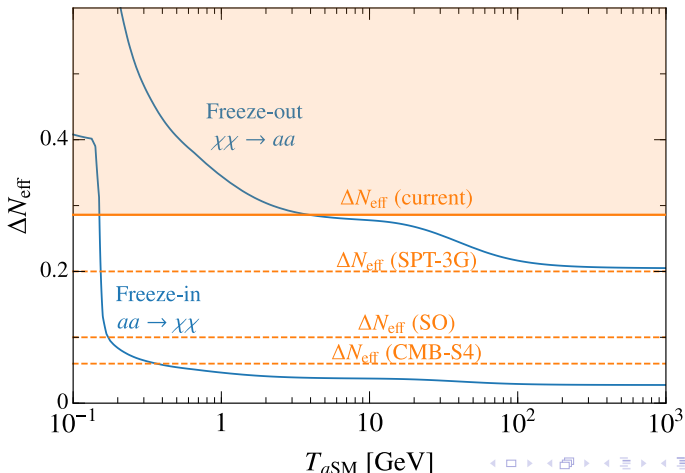


Figure: LEFT : $gg \rightarrow \chi\chi$, RIGHT: $gg \rightarrow \chi\chi$, $t\bar{t} \rightarrow \chi\chi$

Thermally coupled axions I

$$T_{\text{RH}} > T_{a\text{SM}}$$

Dark Sector decouples from SM at $T_{a\text{SM}} \rightarrow$ Dark radiation \rightarrow change from SM prediction $N_{\text{eff}} = 3.044$.



Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Parameters $g_{a\chi}, f_a, m_\chi \rightarrow$ Temperature hierarchy.

$T_{RH} \rightarrow$ production mechanism of dark matter.

Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Parameters $g_{a\chi}, f_a, m_\chi \rightarrow$ Temperature hierarchy.

$T_{RH} \rightarrow$ production mechanism of dark matter.

Vast collection of production mechanisms to be understood and probed fully.

Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Parameters $g_{a\chi}, f_a, m_\chi \rightarrow$ Temperature hierarchy.

$T_{RH} \rightarrow$ production mechanism of dark matter.

Vast collection of production mechanisms to be understood and probed fully.

Future studies of out-of-equilibrium dynamics :

Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Parameters $g_{a\chi}, f_a, m_\chi \rightarrow$ Temperature hierarchy.

$T_{RH} \rightarrow$ production mechanism of dark matter.

Vast collection of production mechanisms to be understood and probed fully.

Future studies of out-of-equilibrium dynamics :

- interplay of $\chi\chi \rightarrow aa$ with SM $\rightarrow \chi\chi$ (blue region on slide 10).

Summary

QCD axion can play a crucial role as the mediator between the DM and SM.

Parameters $g_{a\chi}, f_a, m_\chi \rightarrow$ Temperature hierarchy.

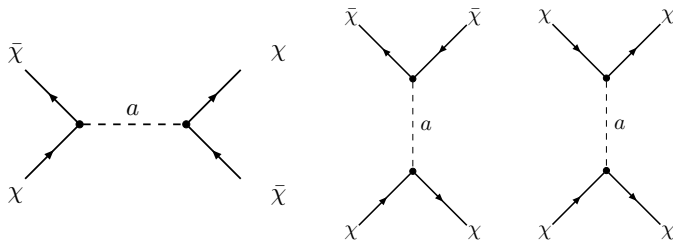
$T_{RH} \rightarrow$ production mechanism of dark matter.

Vast collection of production mechanisms to be understood and probed fully.

Future studies of out-of-equilibrium dynamics :

- interplay of $\chi\chi \rightarrow aa$ with SM $\rightarrow \chi\chi$ (blue region on slide 10).
- Out-of-equilibrium collisions of the axions $aa \rightarrow \chi\chi$ frozen-in from SM \rightarrow SM a (Sequential Freeze-in).

Backup Slide I : Experimental Constraints II (SIDM)



Self-interacting DM

Transfer

$$\sigma_T = \int d\Omega \frac{d\sigma_{\text{SIDM}}}{d\Omega} (1 - \cos \theta) \lesssim 1 \frac{\text{cm}^2}{\text{g}},$$

$$\Rightarrow g_{a\chi} \lesssim 0.21 \left(\frac{m_\chi}{1 \text{ MeV}} \right)^{\frac{3}{4}}.$$

Backup slide II : Thermally coupled axions I

Freeze-out (secluded) $\chi\bar{\chi} \rightarrow aa$.

Hierarchy : $T_{\text{RH}} \gtrsim T_{\chi\text{SM}}$ or $T_{a\chi} \gtrsim T_{\text{RH}} \gtrsim T_{a\text{SM}}$.

For $m_\chi \sim 10$ GeV, $g_{a\chi} \sim 0.1$.

Freeze-in $aa \rightarrow \bar{\chi}\chi$.

Hierarchy : $T_{a\text{SM}} \lesssim T_{\text{RH}} \lesssim T_{\chi\text{SM}}$ and $T_{\text{RH}} \gtrsim T_{a\chi}$

$g_{a\chi}$ can be in its natural regime m_χ/f_a for weak-scale m_χ .