Muon Cascade Calculations

Philip D. Jones

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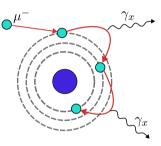
Engineering and Physical Sciences Research Council

ISIS Neutron and Muon Source

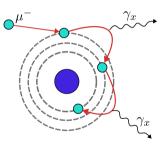
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- Radiative Transitions
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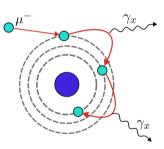
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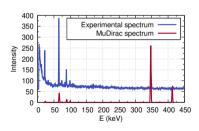
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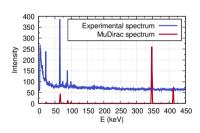
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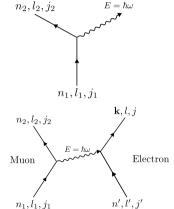


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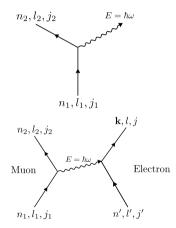


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- Energies are currently predicted well when compared to experiment, but not intensities

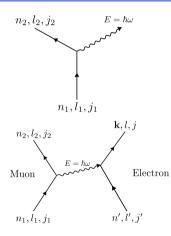




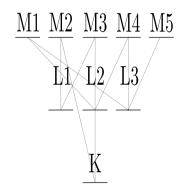
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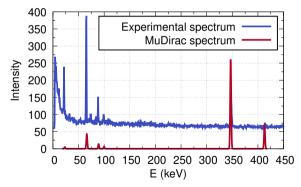


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- With each transition, the muon carries some of its population
- The final measurable intensities are therefore a function of all possible previous transitions

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- This is useful for validation of new relativistic results

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- Energies are calculated accurately through the use of Dirac spinors, finite nuclear models, and QED corrections
- However, intensities are not calculated correctly due to missing physics
- Radiative dipoles are calculated, but Auger transitions, electron refilling, and angular momentum distributions are not considered

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- Radiative Transitions

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- The full relativistic, radiative transition rate between an initial and final muon state is found to be

$$\Gamma_R = \frac{4E_{if}}{3c} |\langle \psi_f | (\alpha \cdot \hat{\mathbf{e}}_{\lambda}) e^{-i\mathbf{k}_{\lambda} \cdot \mathbf{r}} |\psi_i \rangle|^2$$
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 The exponential can be expanded out in Taylor series, and each term in the expansion corresponds to a term in the multipole expansion

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Relativistic Radiative Transitions

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- Comparison to the Akylas-Vogel code is valid for systems in the non-relativistic limit i.e low Z atoms
- Individual quadrupole rates were compared for a given set of transitions:

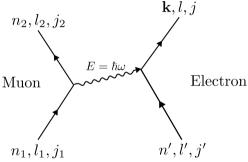
| (n,l) pairs | Akylas-Vogel | Fine pairs | MuDirac |
|---------------------------|----------------------|------------|----------------------|
| $(3,2) \rightarrow (1,0)$ | 5.48×10^{9} | K1-M4 | 3.70×10^{9} |
| | | K1-M5 | $1.61 	imes 10^9$ |
| | | Sum | 5.31×10^{9} |
| $(3,2) \rightarrow (2,0)$ | 4.72×10^{8} | L1-M4 | 3.18×10^{8} |
| | | L1-M5 | 1.39×10^{8} |
| | | Sum | 4.57×10^{8} |

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- Four wavefunctions are needed to calculate this: the initial muon state, final muon state, a bound electron, and a continuum electron

$$\Gamma_A = \left| \frac{e^2}{\hbar} \int \int \chi_f^*(\mathbf{r}_2) \psi_f^{\dagger}(\mathbf{r}_1) \frac{1}{||\mathbf{r}_1 - \mathbf{r}_2||} \chi_i(\mathbf{r}_2) \psi_i(\mathbf{r}_1) d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 \right|^2$$

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- Coulomb operator can be expanded out in spherical harmonics using the Laplace expansion, and expressions for Auger multipole rates were obtained
- Validation against Akylas-Vogel is yet to be performed

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- Density Functional Theory (DFT) provides a more accurate description of the electronic structure of a system
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- TCDFT has been used for electron-positive muon systems, including a parameterisation of the correlation functional
- Being able to solve this problem for the negative muon would give a more accurate description of the system

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- Ultimately, the ability to solve for muonic states in a bulk system may allow for more accurate capture probabilities, potentially leading to better intensities

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- MuDirac does not get intensities correct due to missing physics
- Relativistic radiative and Auger transitions are the first steps taken to getting correct intensities, with the radiative rates being successfully validated against the Akylas-Vogel code

References

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