

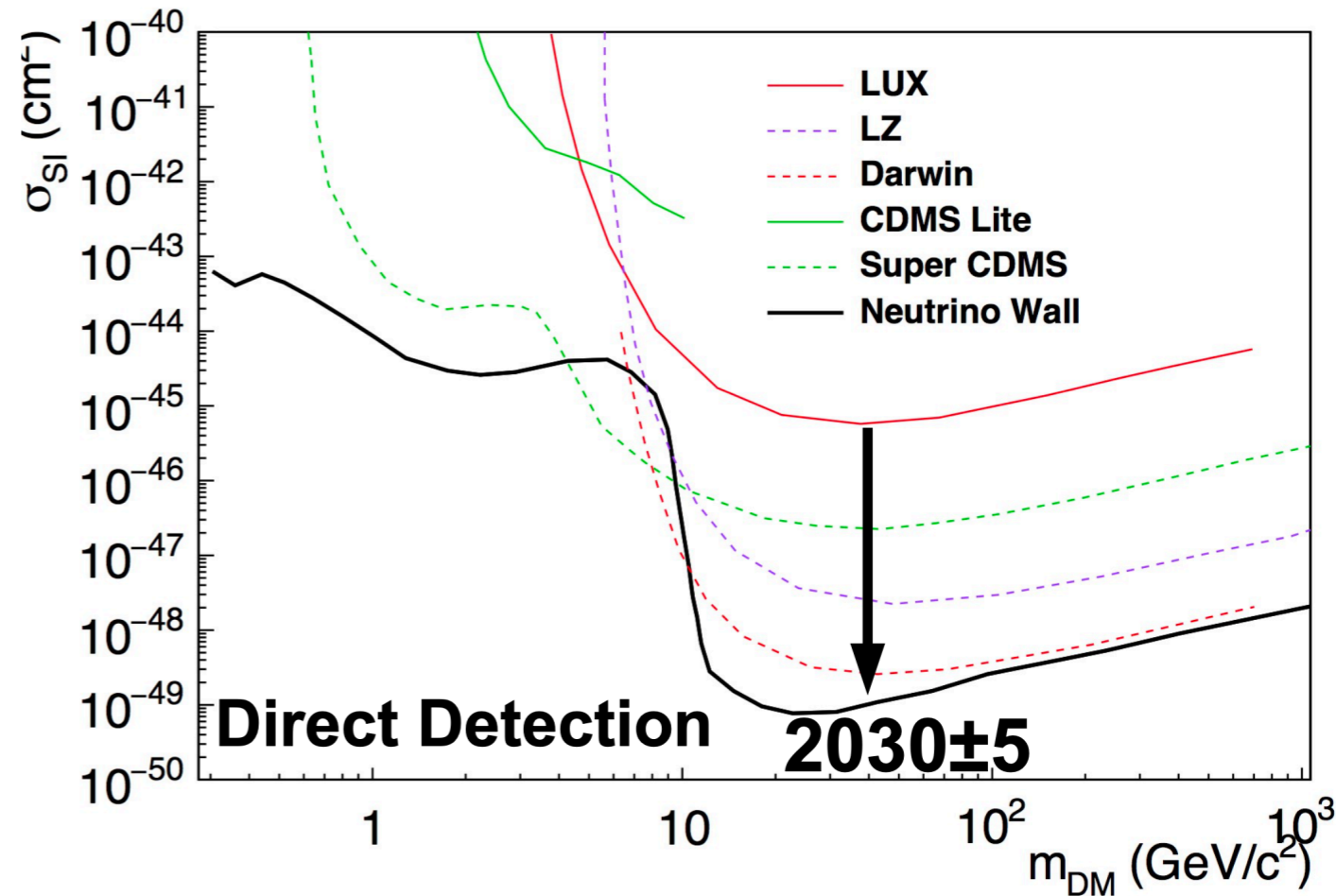
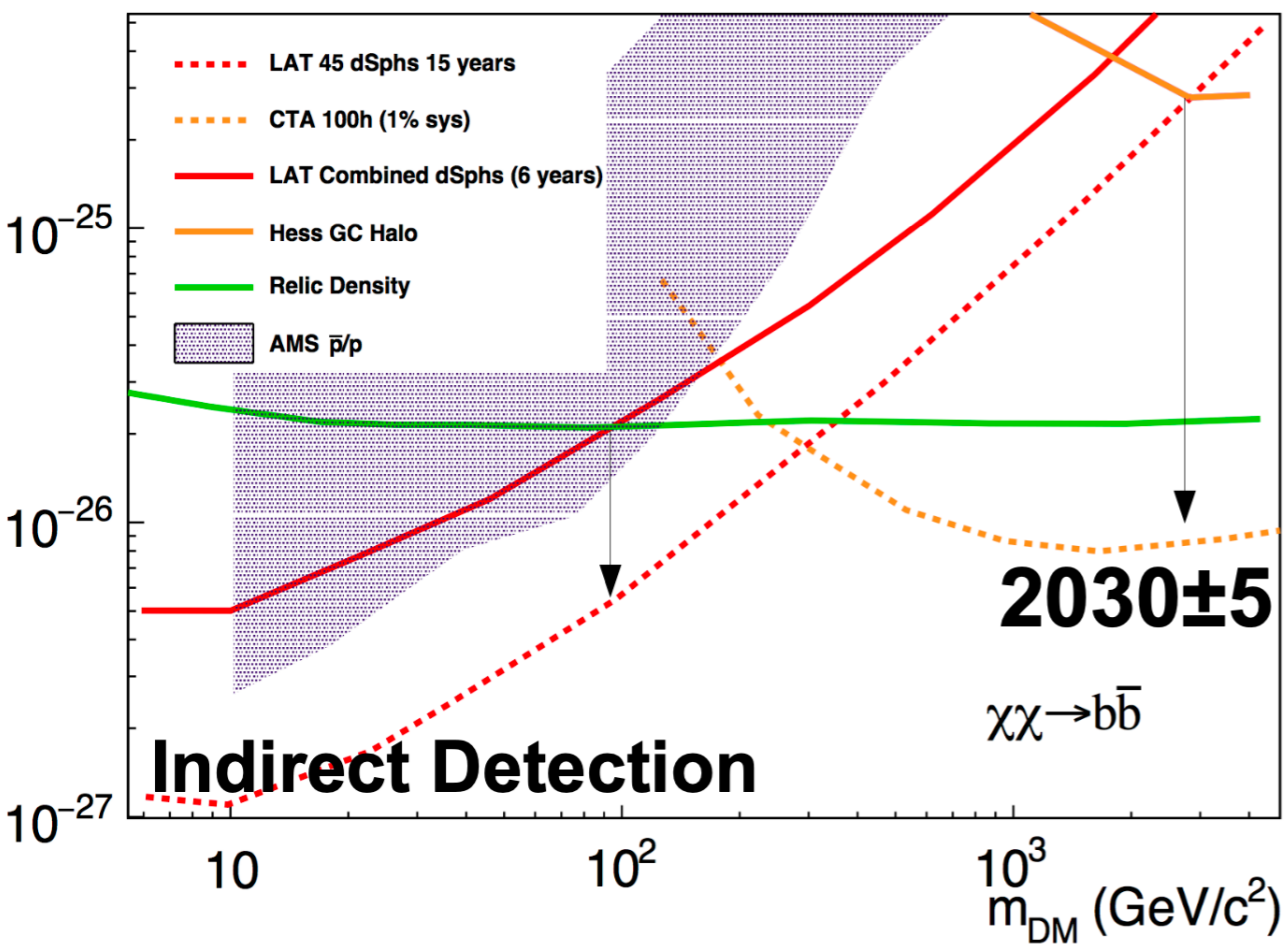


Prospects of DM at Colliders

P. Harris w/ Contributions from Many



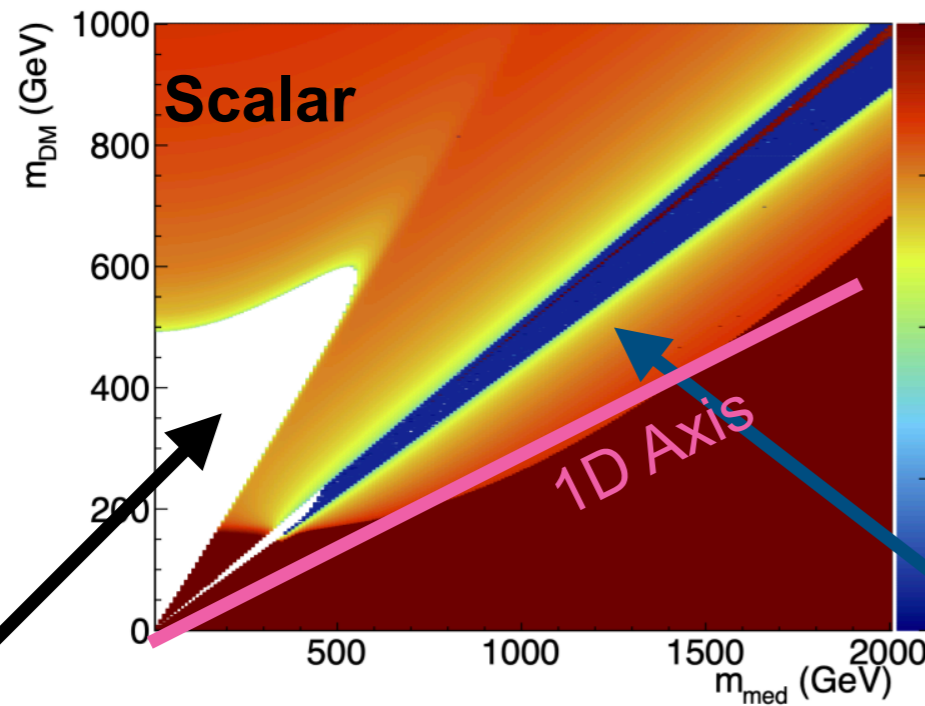
Other Exp: Driving Bounds



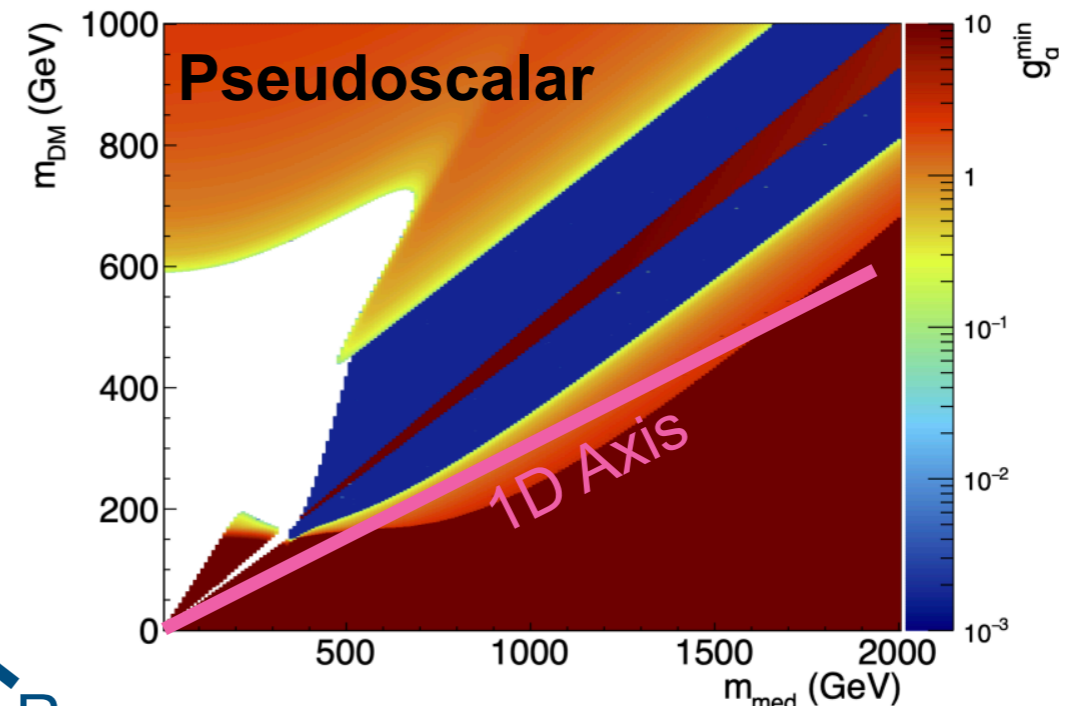
Future Experiments in other domains have clear benchmarks

With current DM searches, we aim to **complement** these

Relic: Driving Bounds

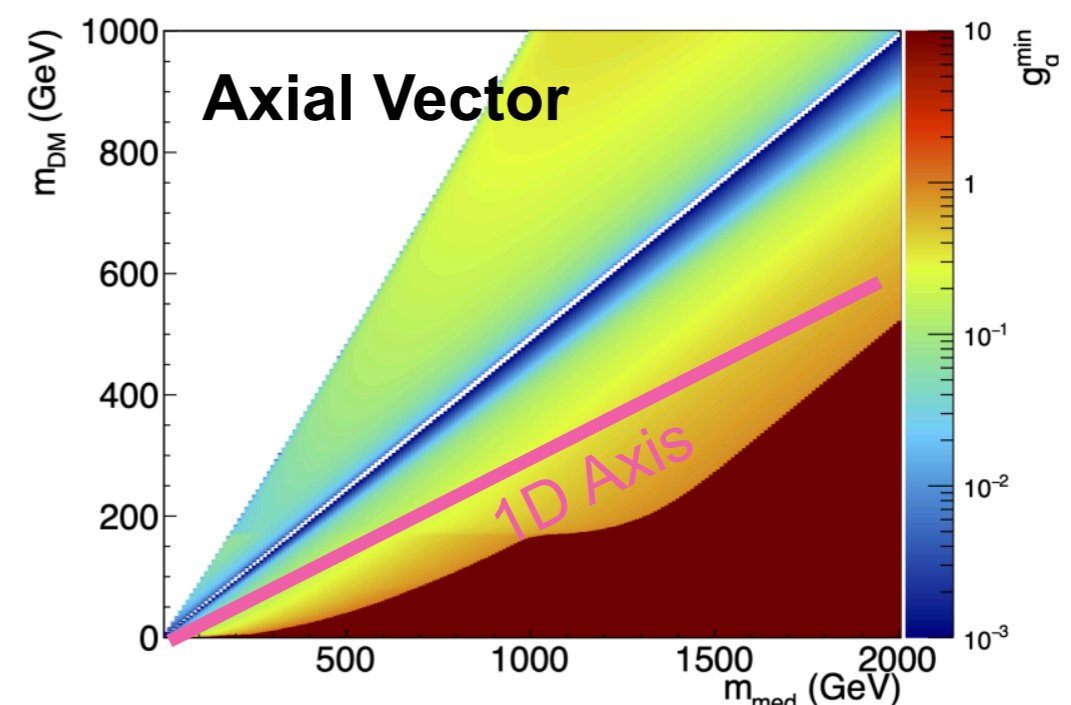
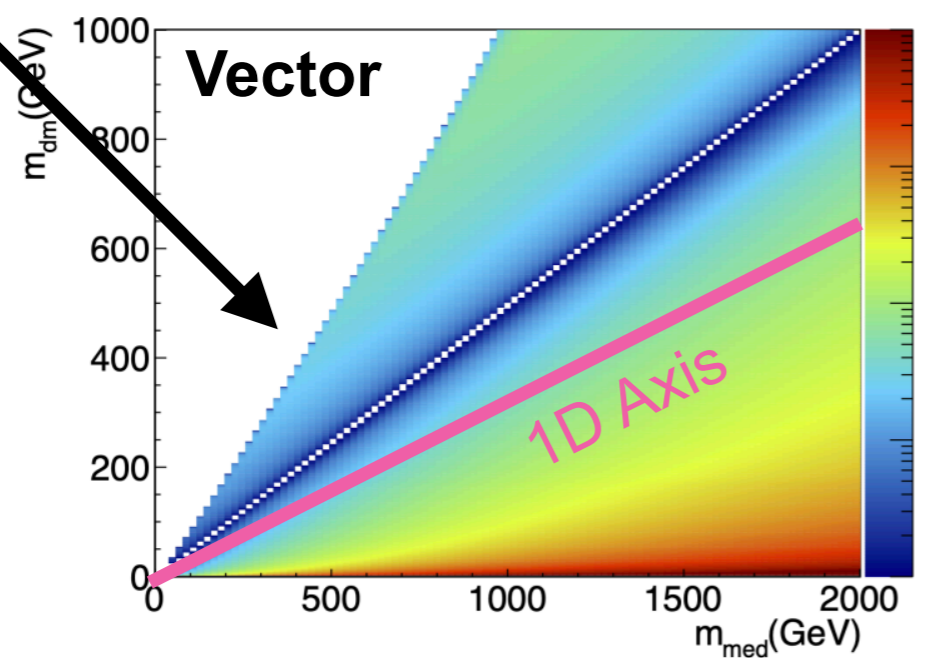


(a)

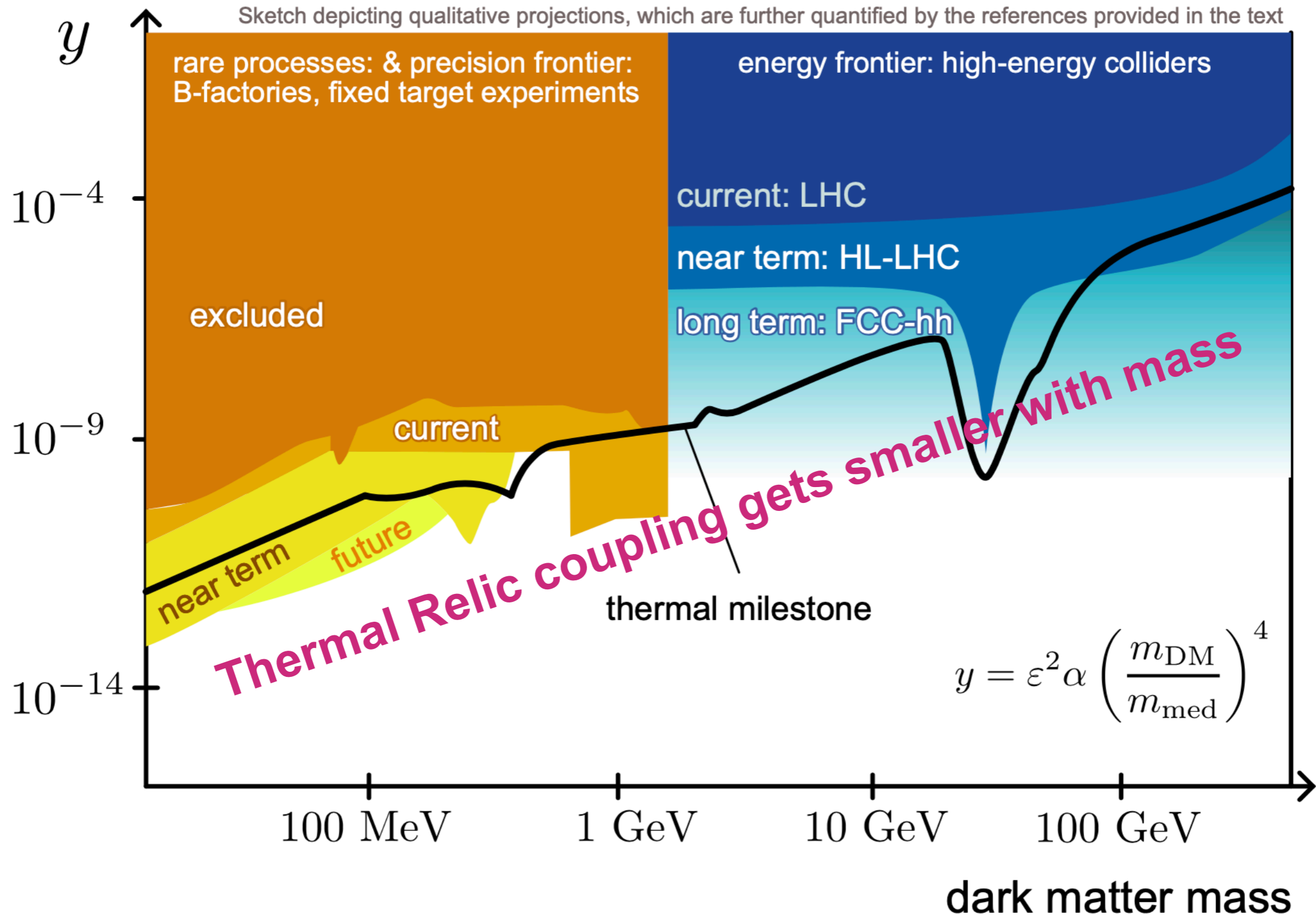


(b)

Resonance
(Loose Constraints)

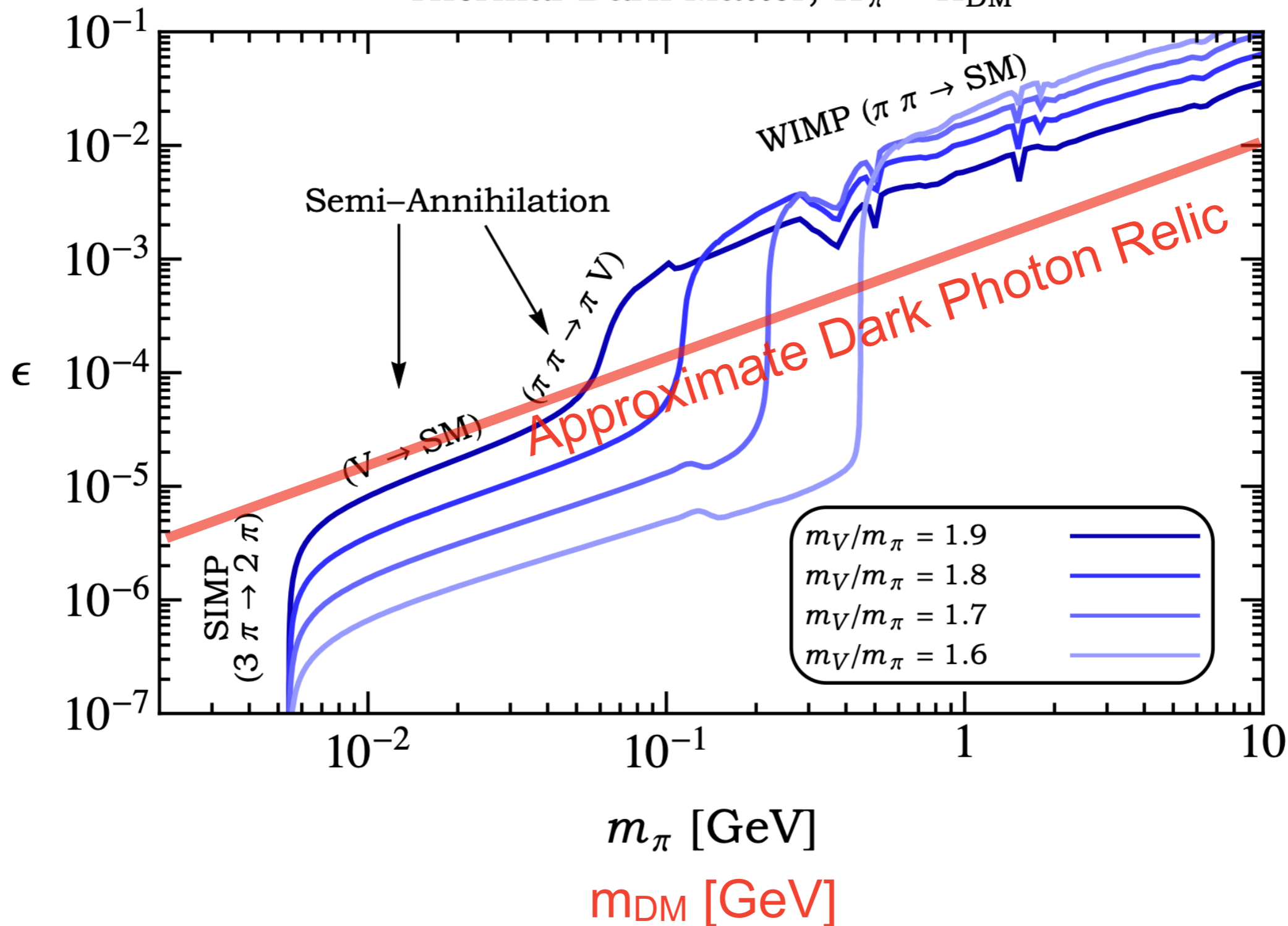


Driving Factors

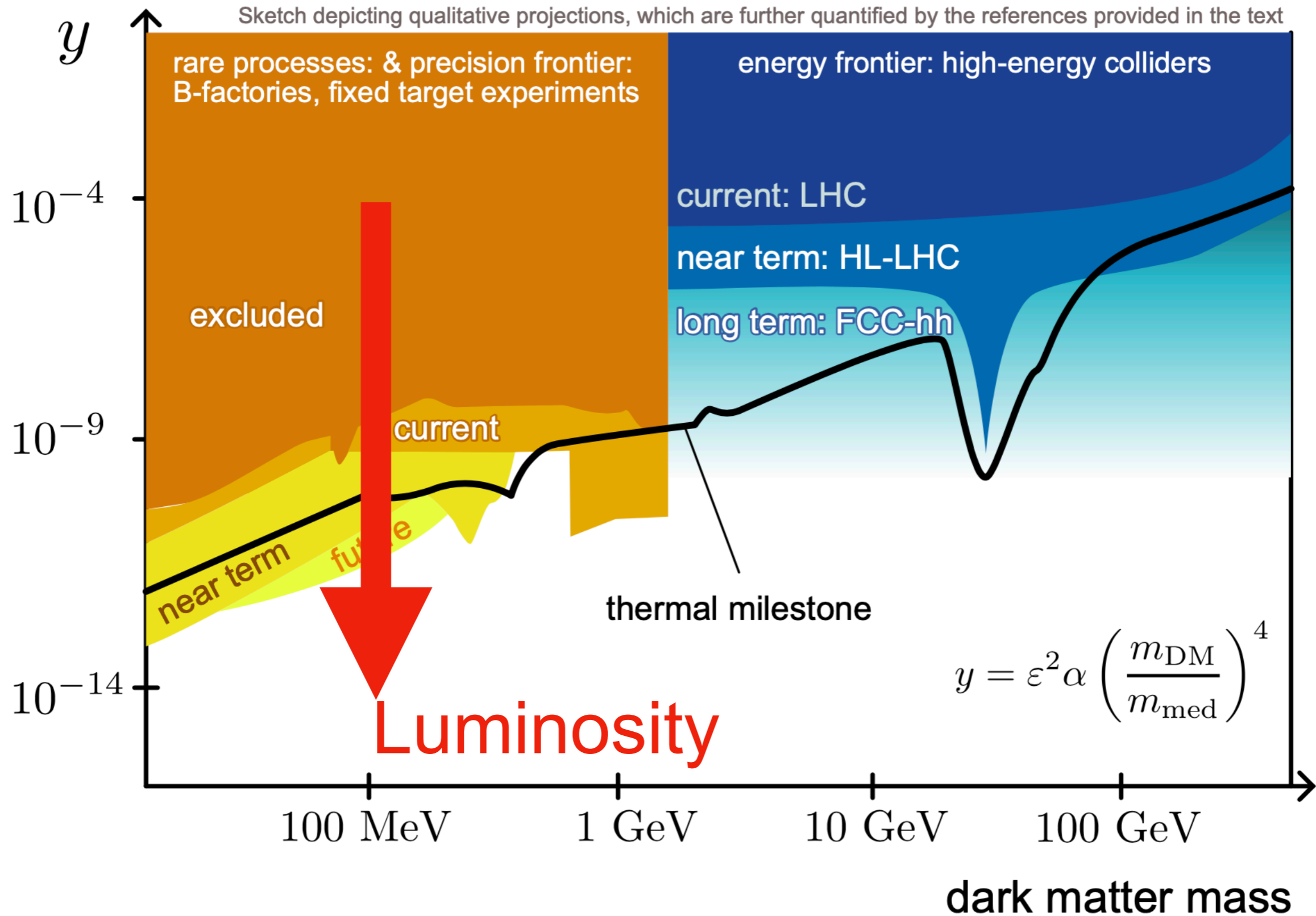


Complicating Models⁵ Complicates Relic

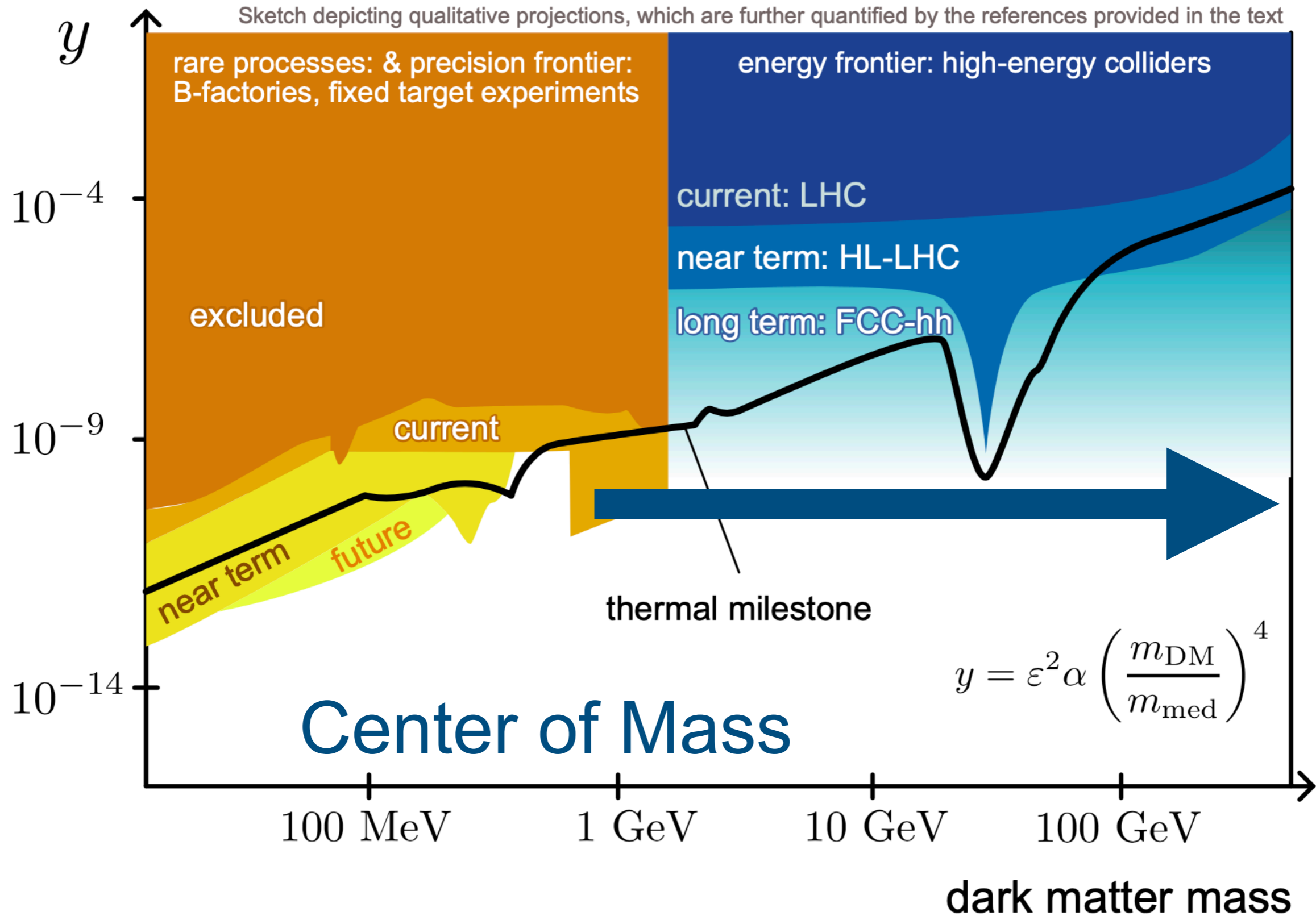
Thermal Dark Matter, $\Omega_\pi = \Omega_{\text{DM}}$



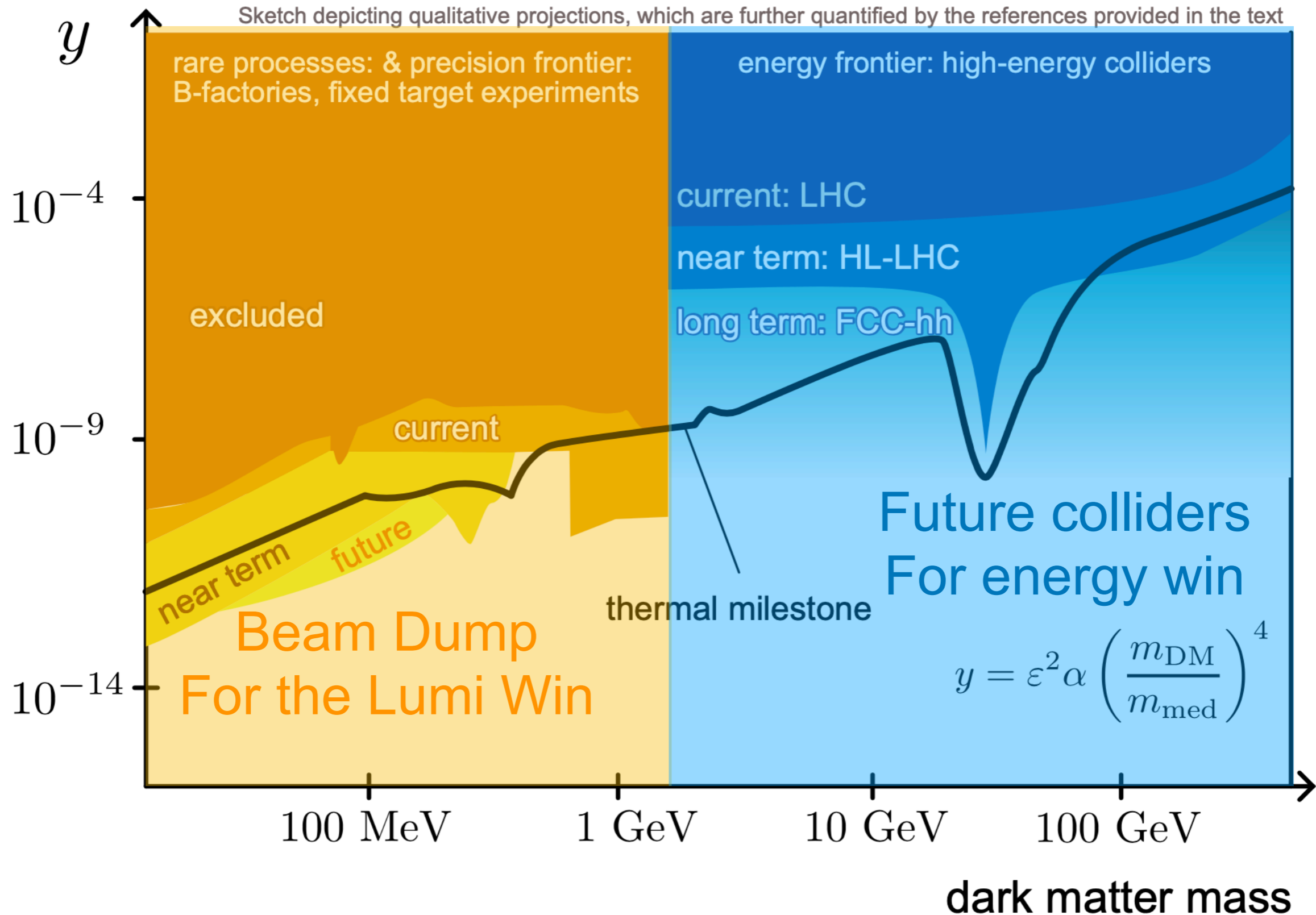
Driving Factors



Driving Factors



Driving Factors

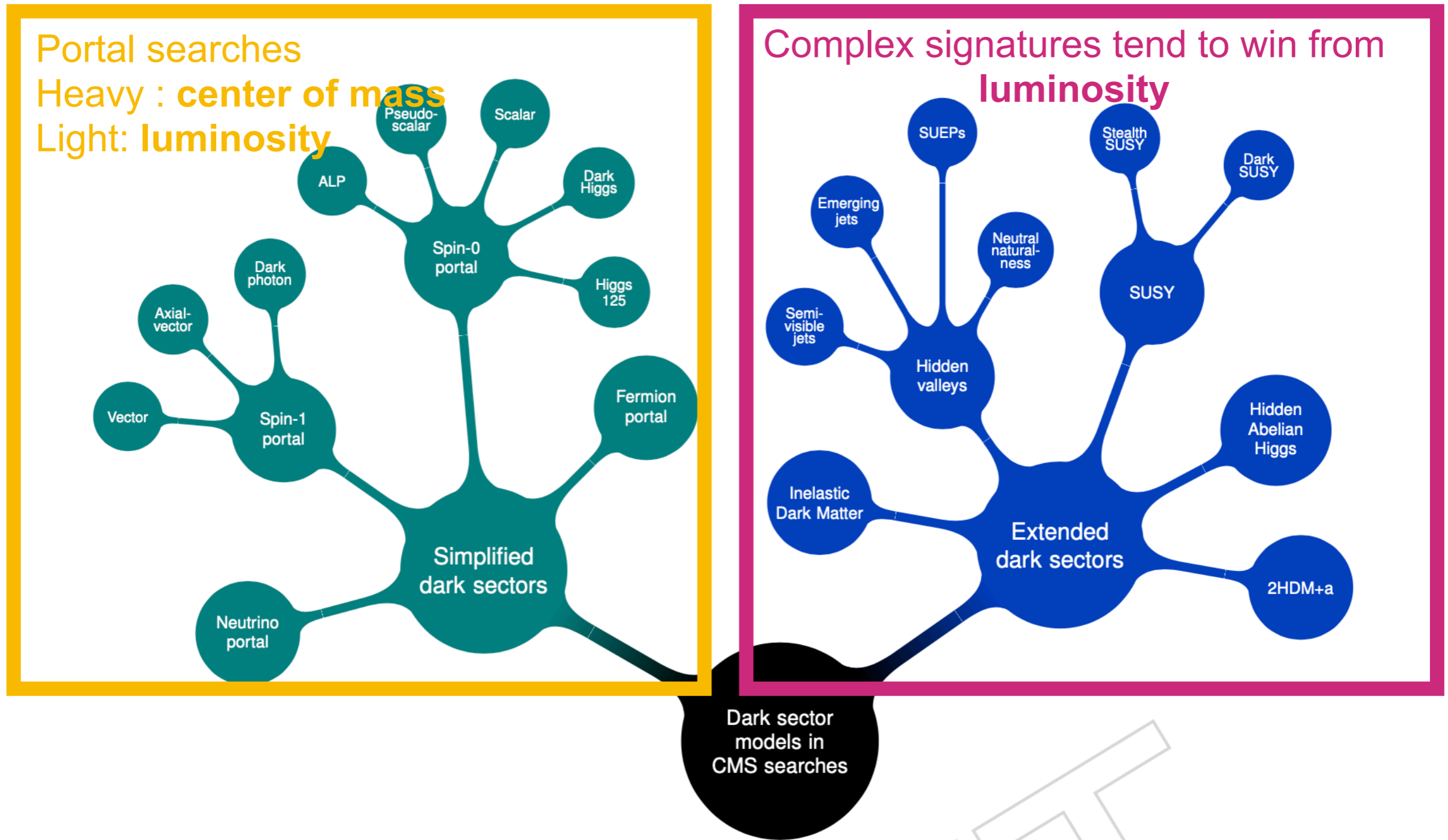


Whats on the Market

BeamDump+

LHC	0.2 x 10 ¹² Z bosons 2 x 10 ¹² W bosons 0.2 x 10 ⁹ H bosons Up to COM of 8 TeV	<10 years
e⁺/e⁻	3 x 10¹² Z bosons 2 x 10 ⁸ W bosons 1 x 10 ⁶ H bosons Up to COM of 500 GeV	2055-2060
<i>μ⁺/μ⁻</i>	Muons are cool Up to COM of 10 TeV	By 2060?
PbPb	50x Current bounds? 0.45x10⁸ γ-enhancement	<10 years
100 TeV-pp	1 x 10 ¹³ Z bosons 1 x 10 ¹⁴ W bosons 1 x 10 ¹⁰ H bosons Up to COM of 8 TeV	The Future

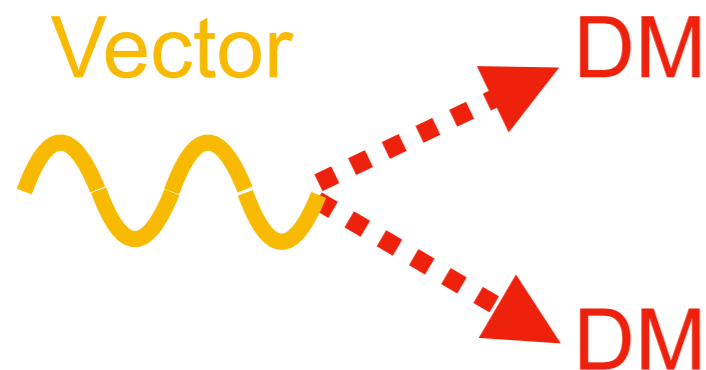
Scope of DM at Colliders



Searches @ colliders

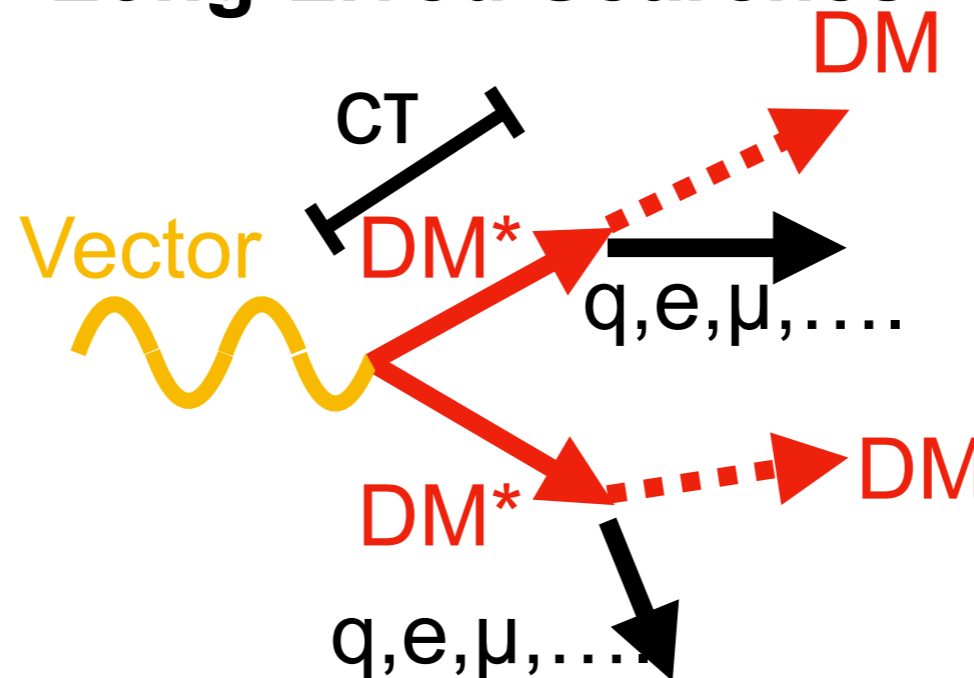
- Light light dark matter there are 3 styles of searches

Invisible Searches



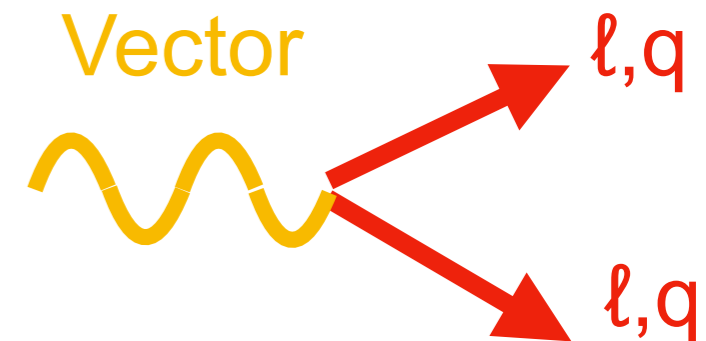
4 π detectors @
LHC (ATLAS/CMS)
use conservation
of transverse
energy for
signatures

Long-Lived Searches



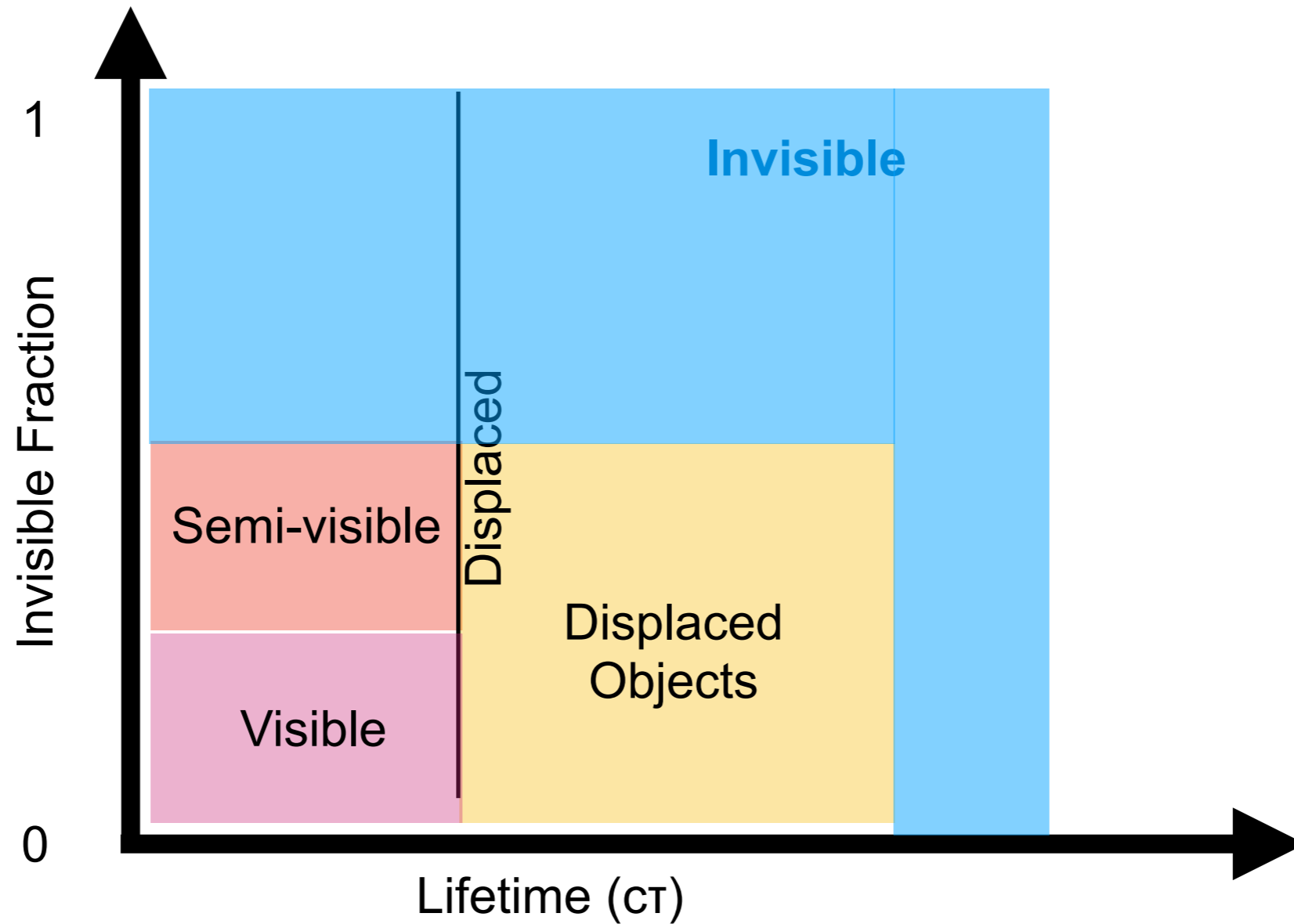
Extensive detector
capabilities make
displaced signatures
up to 5m possible

Visible Searches

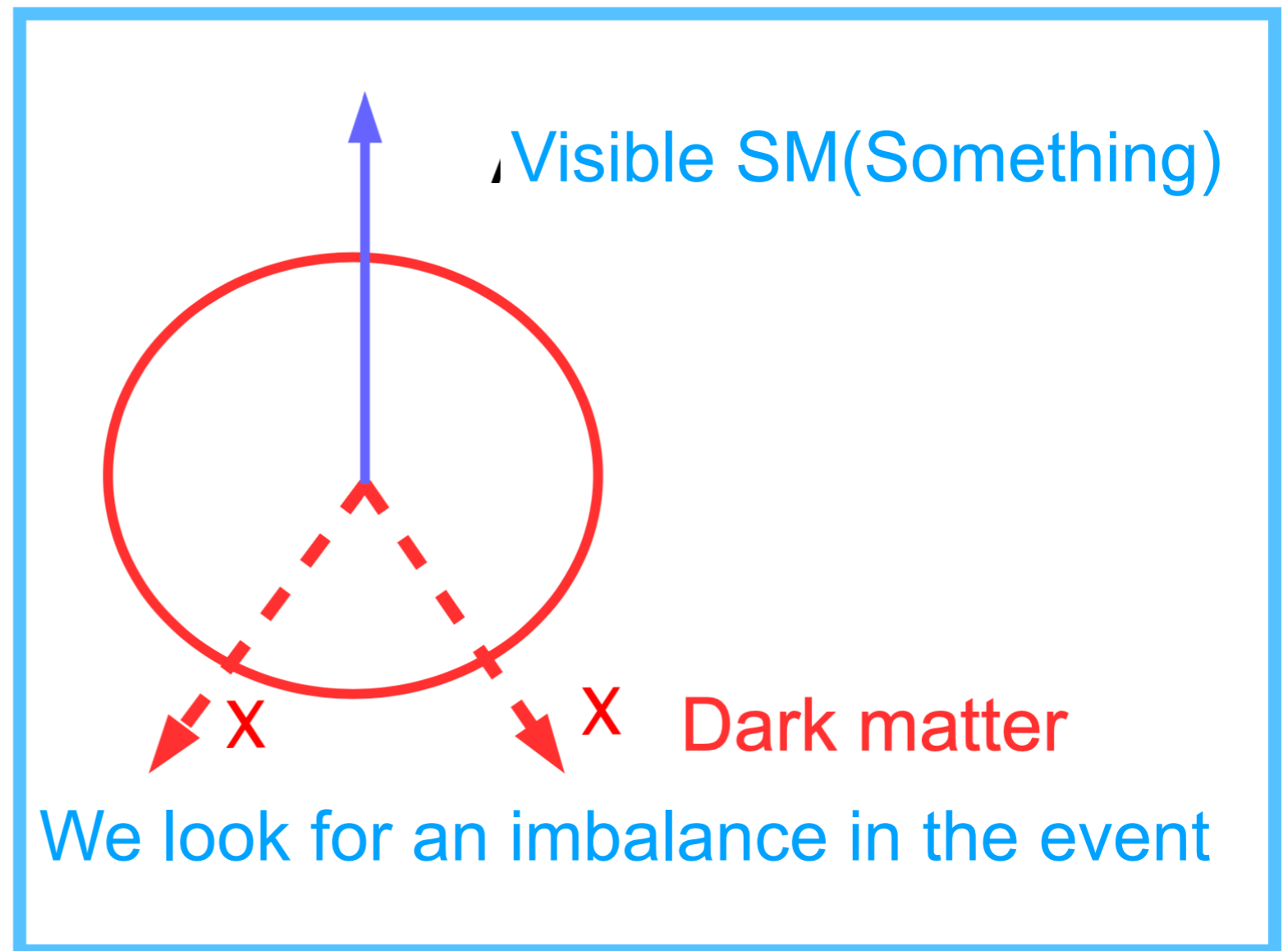
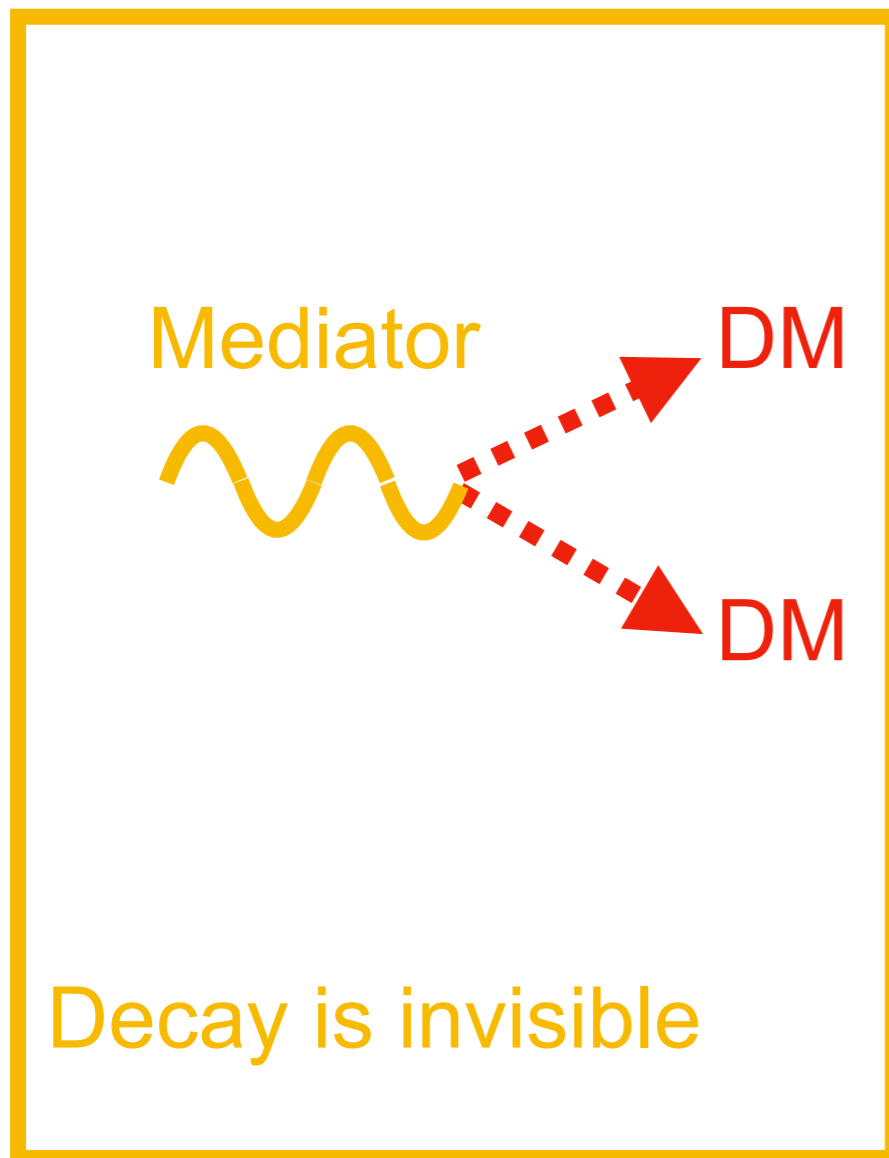


LHC can tag
every type of
object

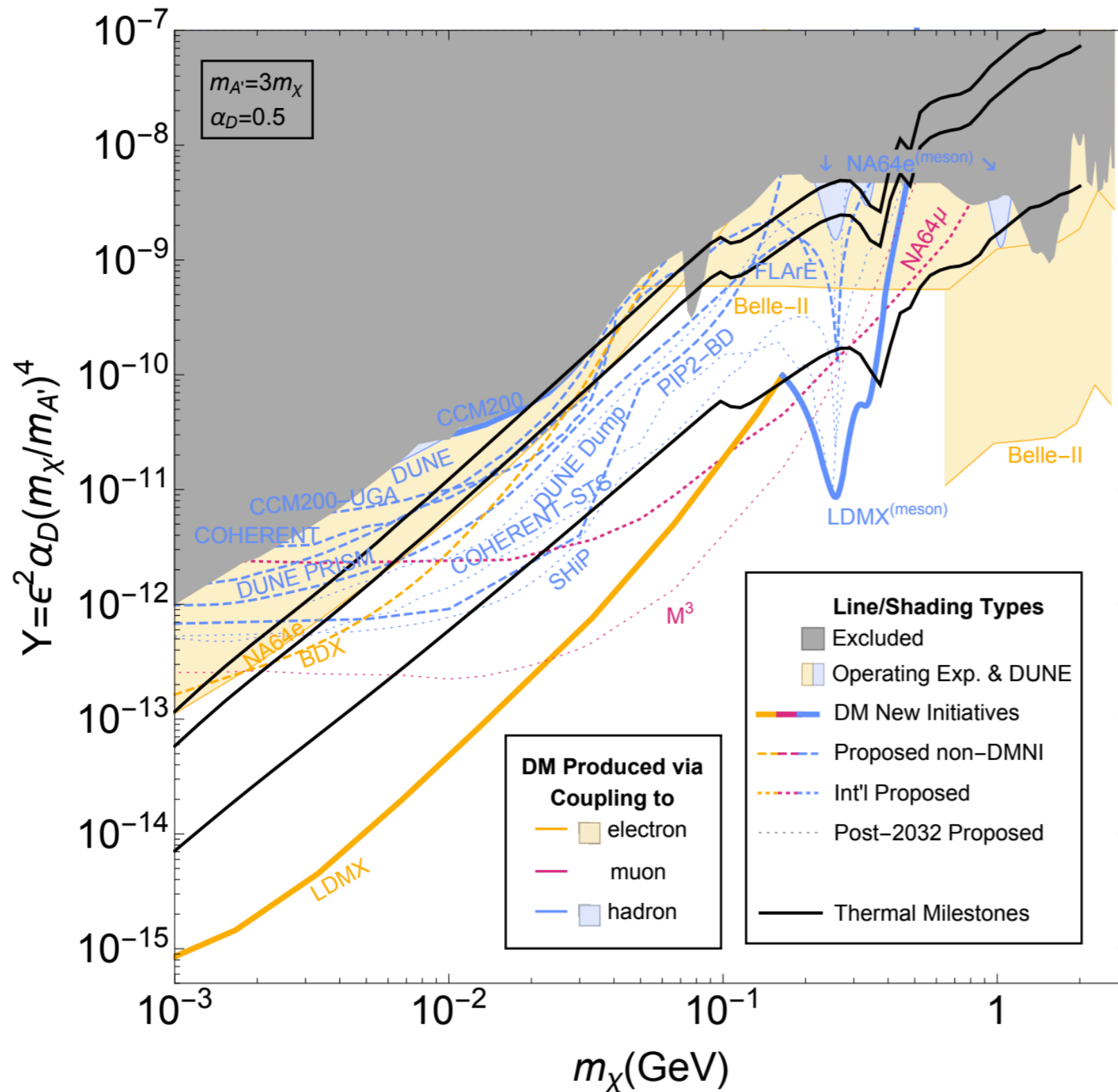
Visualizing the landscape



Invisible Searches



Invisible Searches

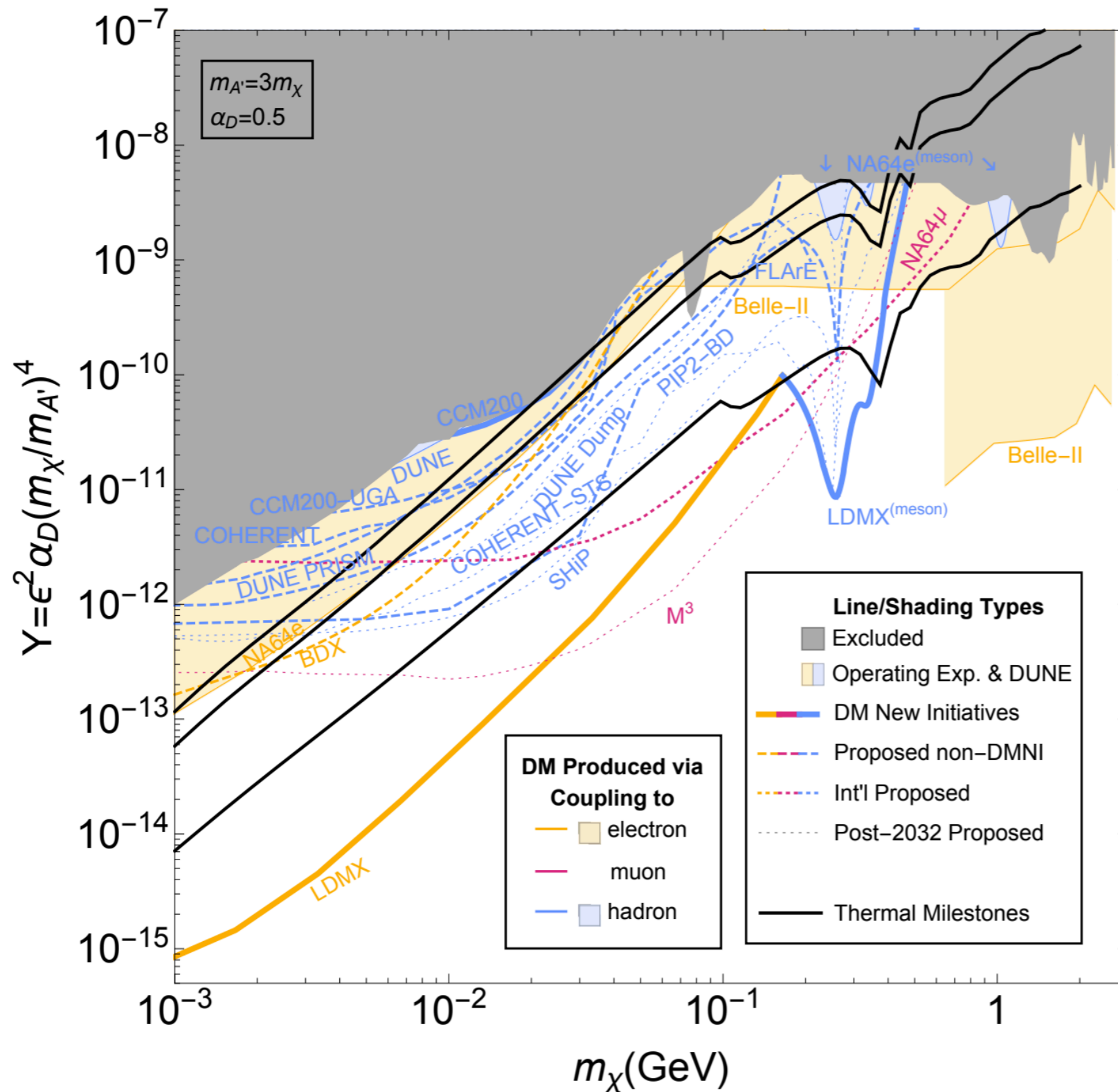


For Light Invisible DM

Next decade bodes well

Experiments like LDMX
Can probe many
conventional invisible
decays

Invisible Searches



What can you do over here?

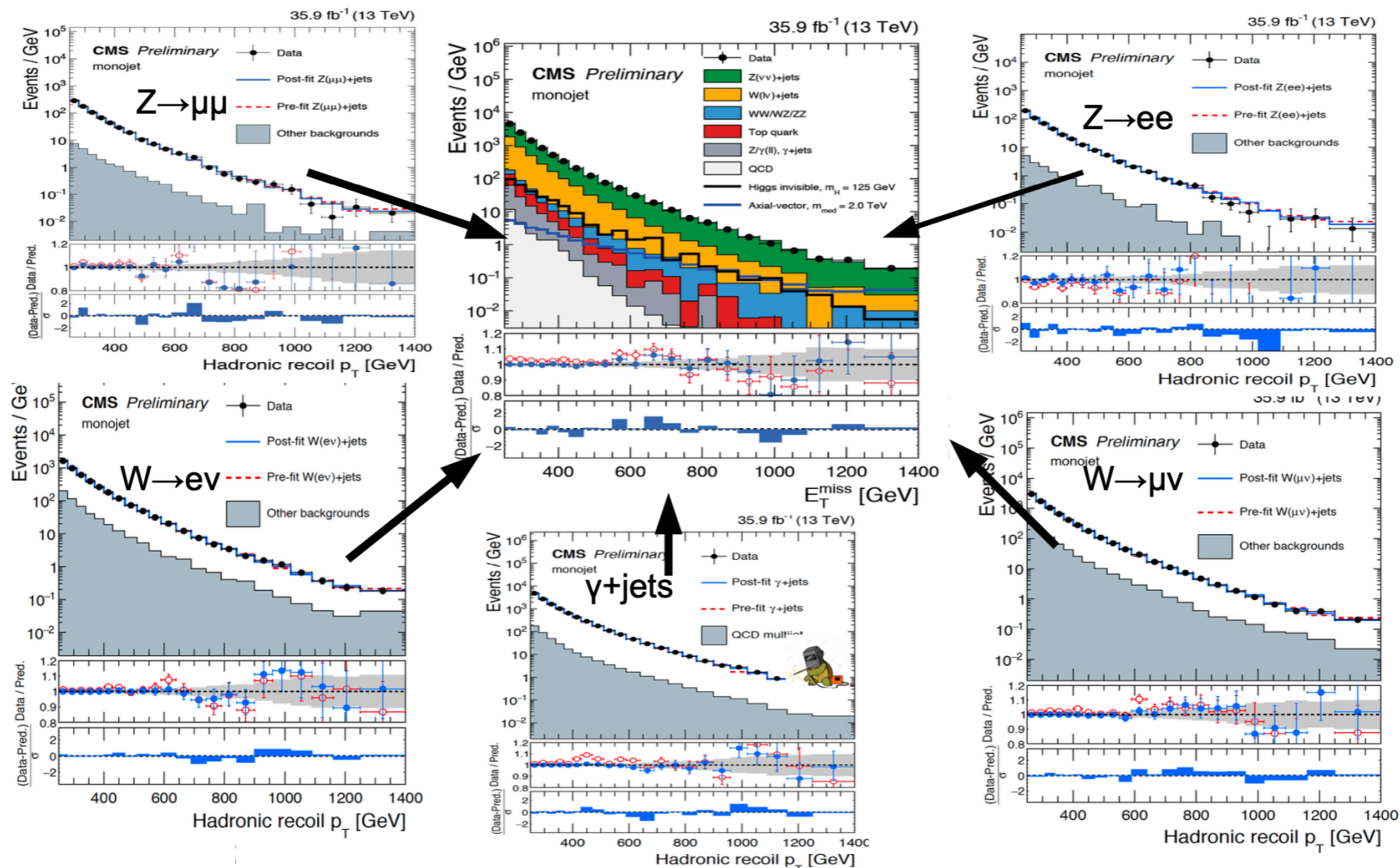
$M_\chi > 3.5 \text{ GeV}$

$M_{\text{med}} > 10 \text{ GeV}$

The LHC is the only game in town

LHC Invisible Search

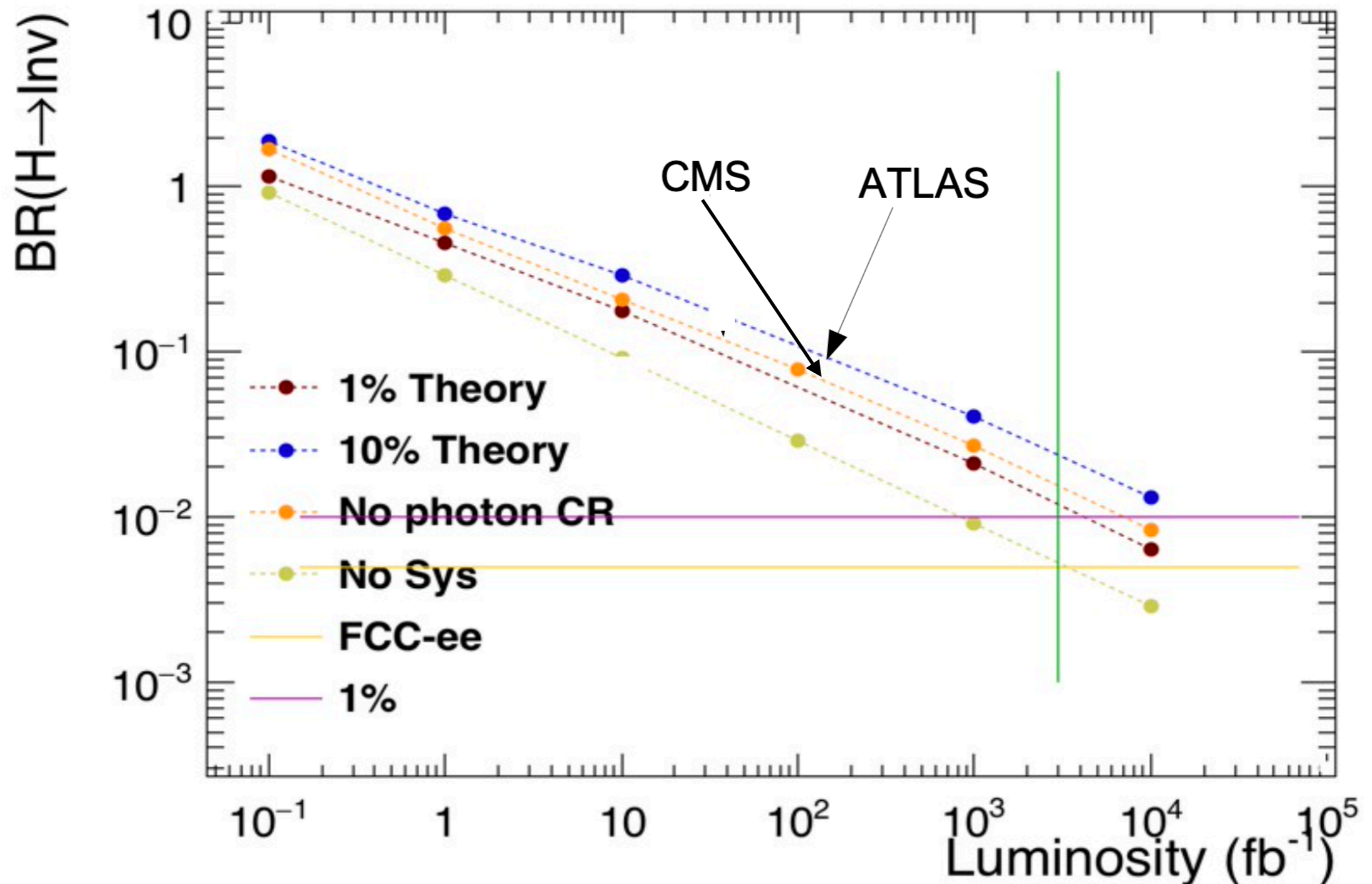
- This analysis tends to be dominated by the monojet search
 - MET+X searches, such as $tt+X$ are also powerful
 - We focus on the monojet to start with



Monojet analysis (MET+X) are built out of multi control regions extrapolations

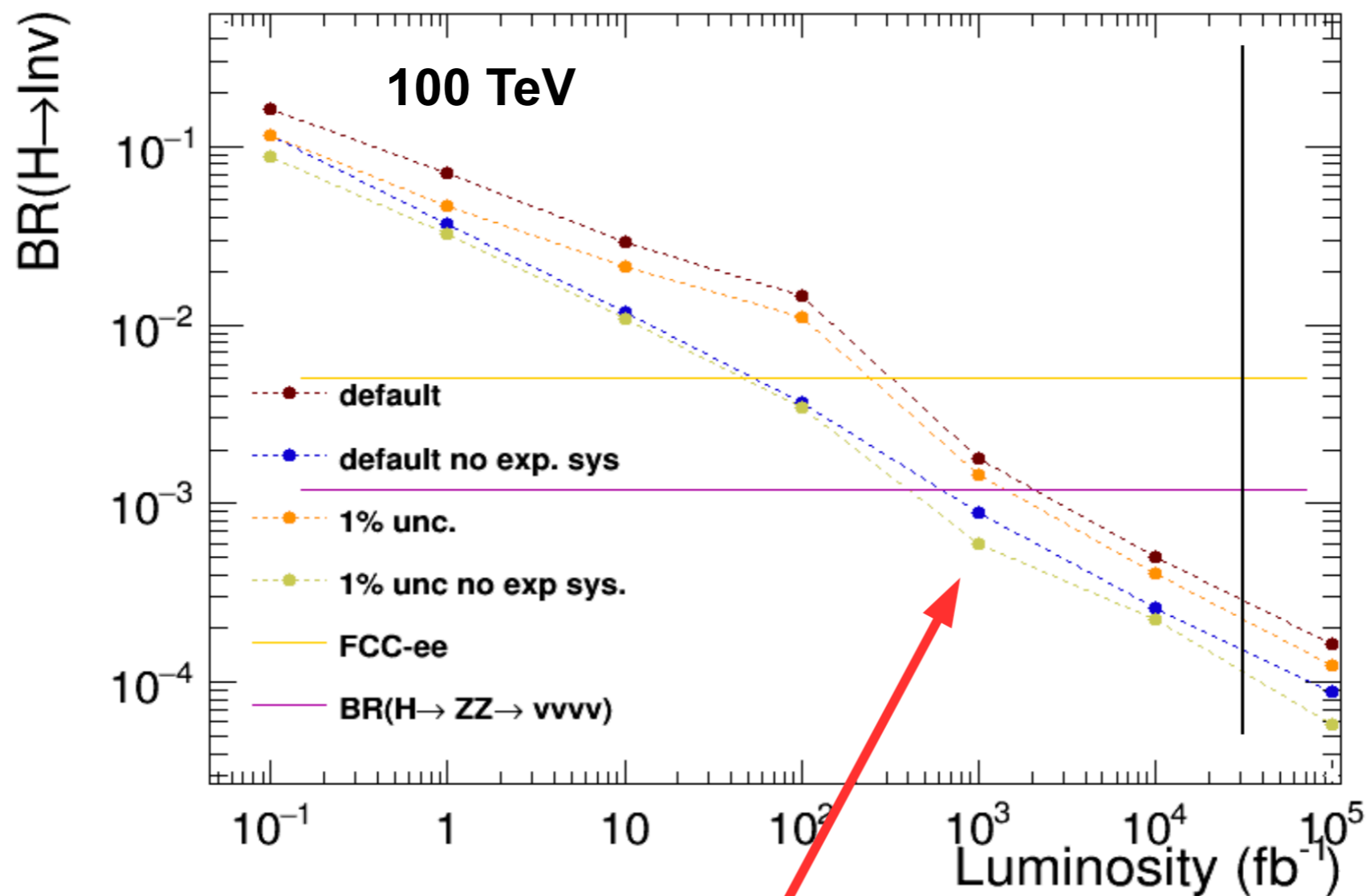
LHC Invisible Search

- Analysis scale with $\sqrt{\text{luminosity}}$ forever
 - Provided we keep experimental uncertainties under control



FCC-hh Invisible Search

- Analysis scale with $\sqrt{\text{luminosity}}$ forever
 - Provided we keep experimental uncertainties under control



Cross the SM neutrino wall at FCC with $< 1 \text{ ab}^{-1}$

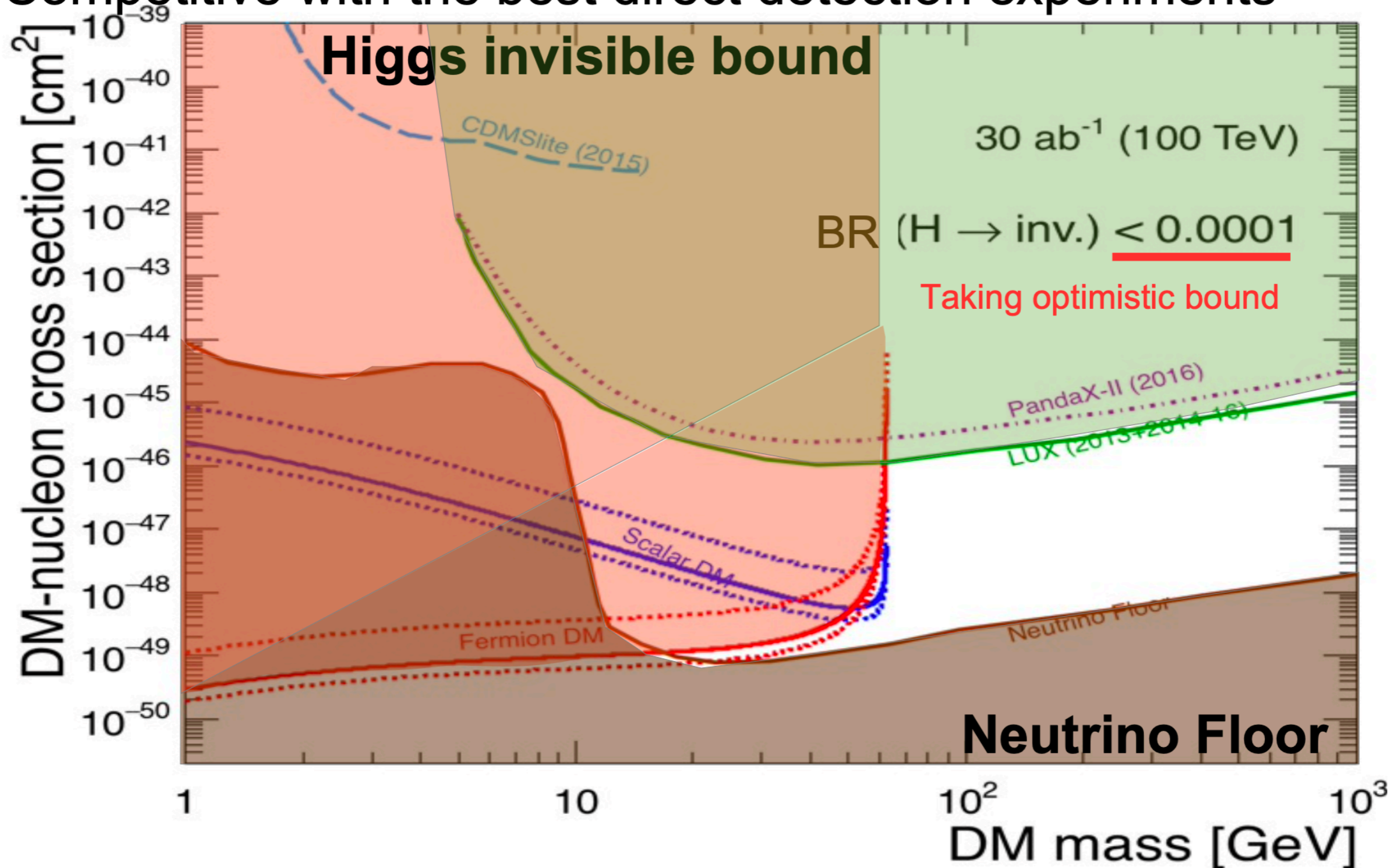
There is no systematics wall

Sensitivity

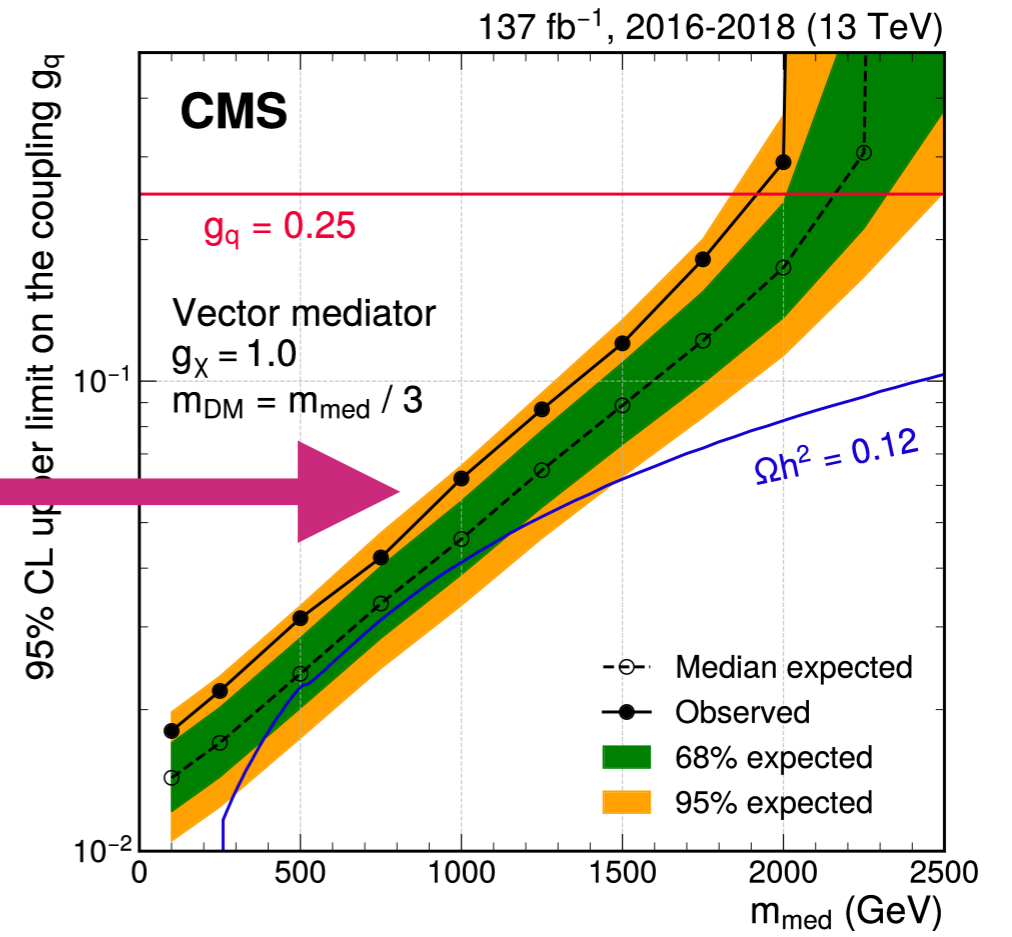
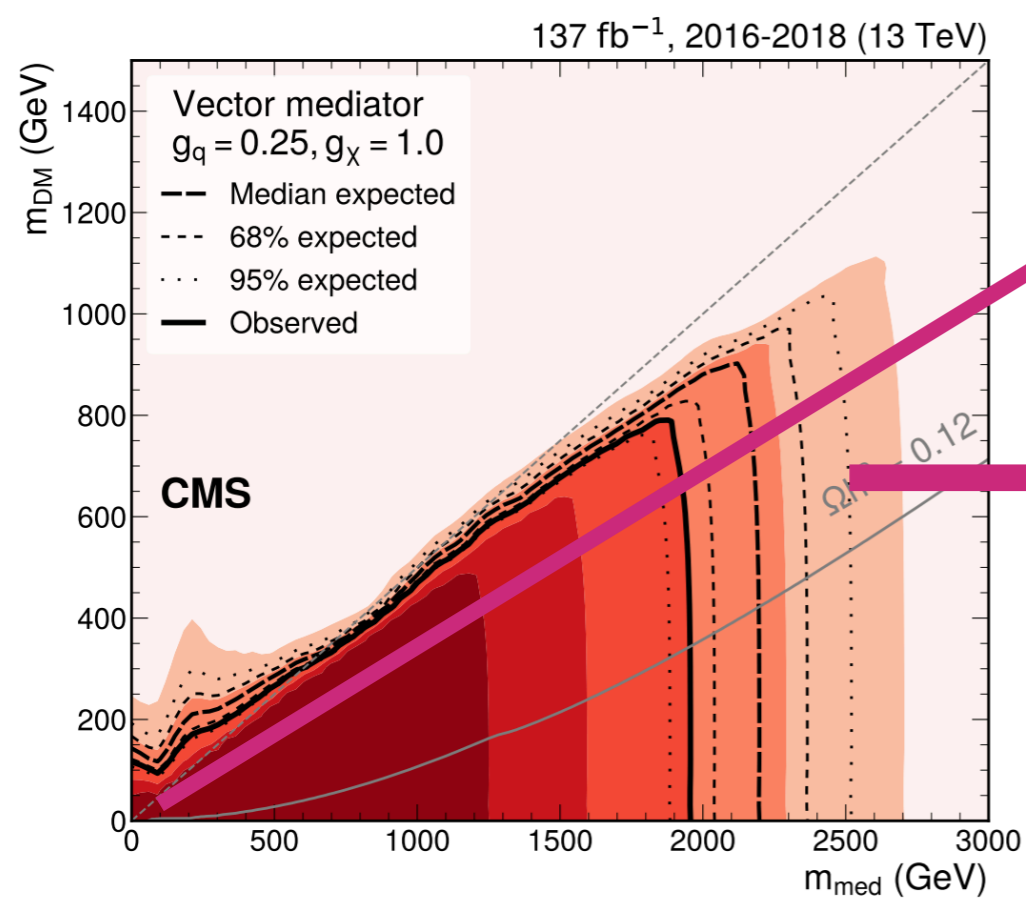
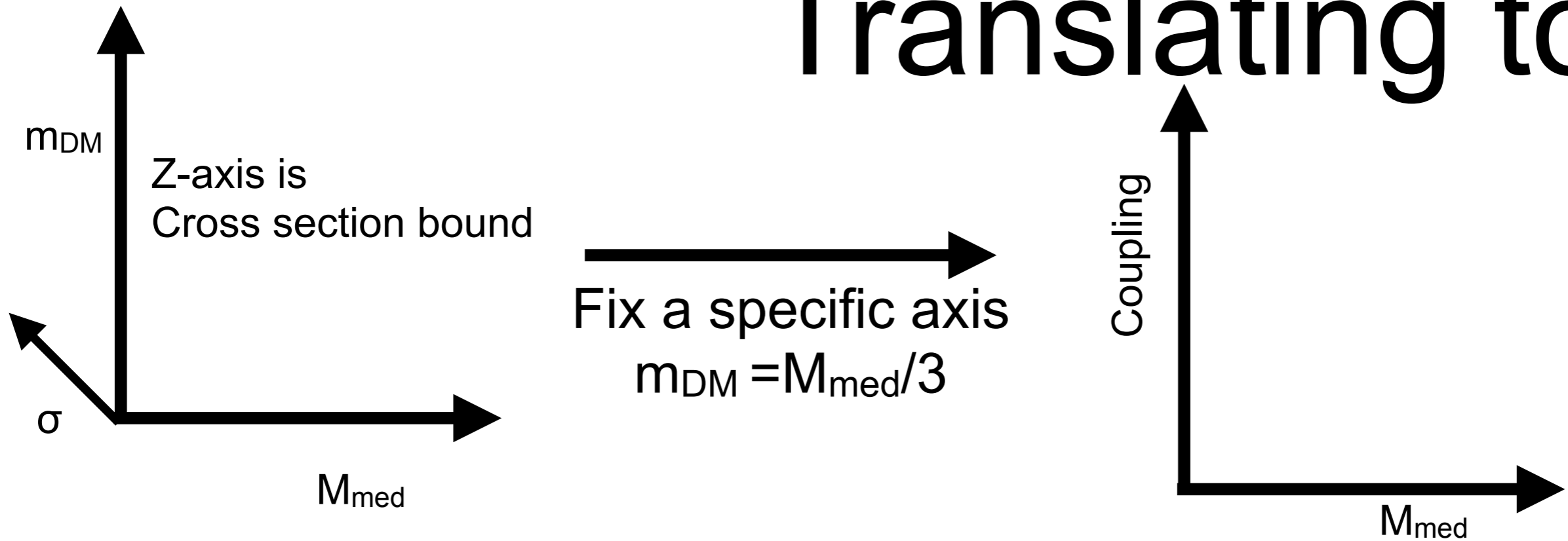
- Benchmarking on Higgs to invisible we find
 - Current LHC $H(\text{inv}) > 0.08$
 - Future LHC $H(\text{inv}) > 0.02$ (Lumi projected is 0.018)
 - FCC-ee $H(\text{inv}) > 0.005$ (Less Higgs and Less Bkg)
 - FCC-hh $H(\text{inv}) > 0.0002$ (Lumi+ σ projected is 0.0002)
 - ▶ Note : $\frac{\sigma(100 \text{ TeV})}{\sigma(14 \text{ TeV})} \approx 20$
- Analytic lumi and cross section scaling works quite well
 - Event with systematics embedded in fit
 - This is because of the all the control regions

Significant Impact

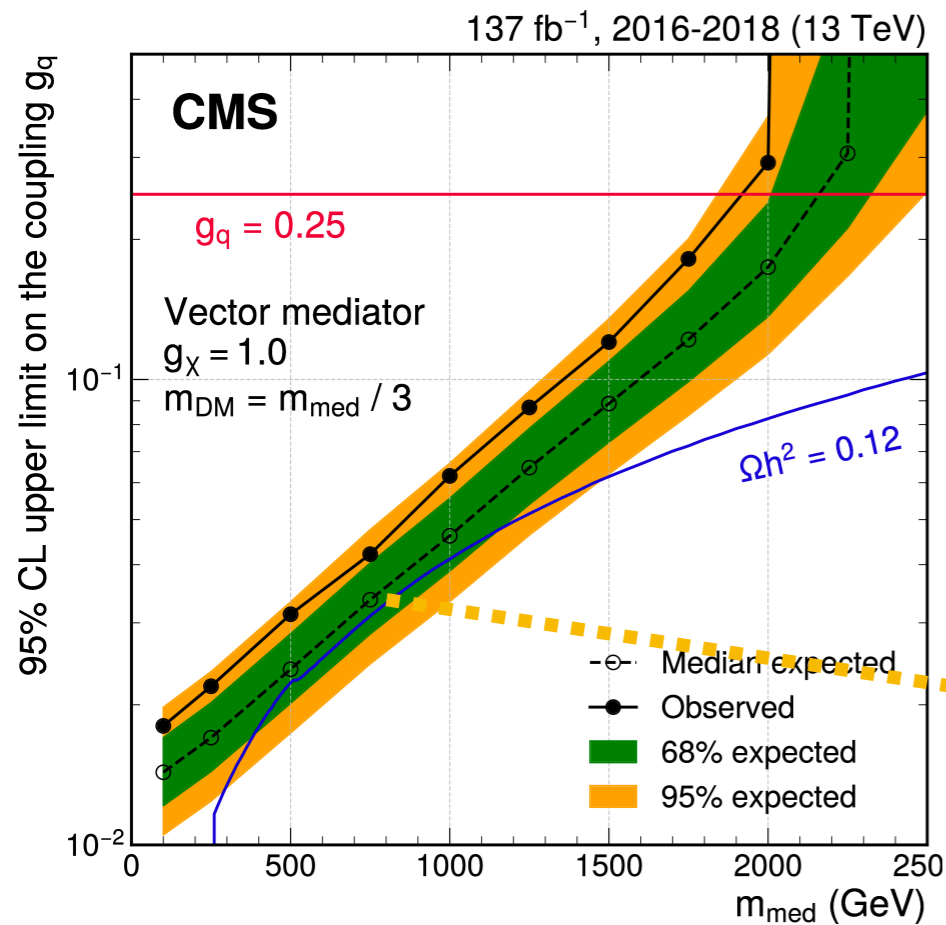
Competitive with the best direct detection experiments



Translating to y

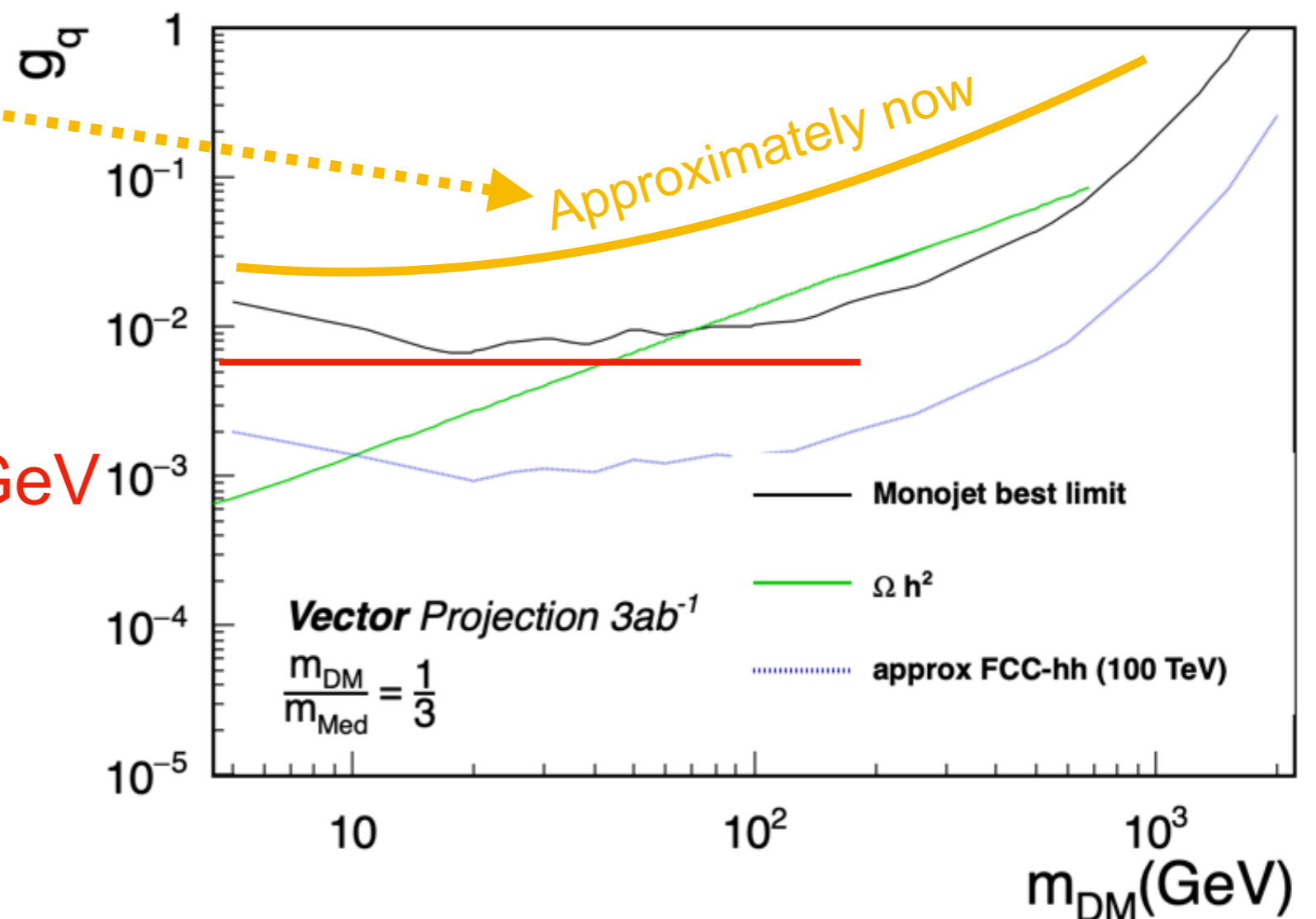


Projecting To the Future



For Heavy Vector and Axial vector DM we exclude with the LHC

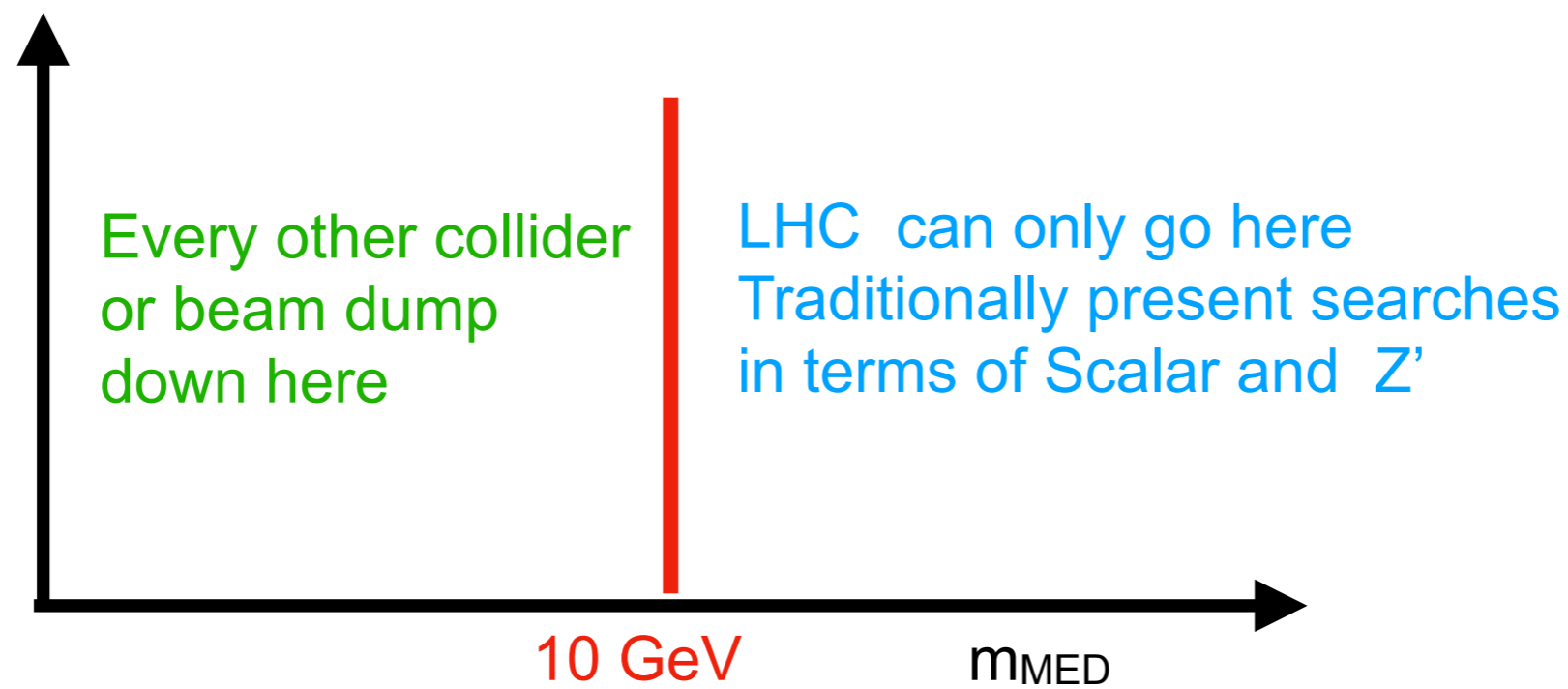
With 100 TeV, we exclude everything



With e⁺/e⁻ polarized 500 GeV
 Linear collider Beam
 FCC is not competitive

Lost in Translation

- LHC has used a different model/notation to present results
 - These models basically the same to FIP/PBC benchmarks
- In the past year we have made an attempt to consolidate
 - General we can recast many analysis towards dark sectors
 - This talk will cover FIP based results at the LHC
 - The full range of LHC goes beyond FIP based models



Mapping Analogies

LHC DM Model	FIP Models	Portal	Coupling
Vector Model	Dark Photon, A_μ		$-\frac{\epsilon}{2\cos\theta_W} F'_{\mu\nu} B^{\mu\nu}$
Scalar Model	Dark Higgs, S		$(\mu S + \lambda S^2) H^\dagger H$
Pseudoscalar Model	Axion, a		$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$
HNL (same)	Sterile Neutrino, N		$y_N L H N$

To get to the FIP Models

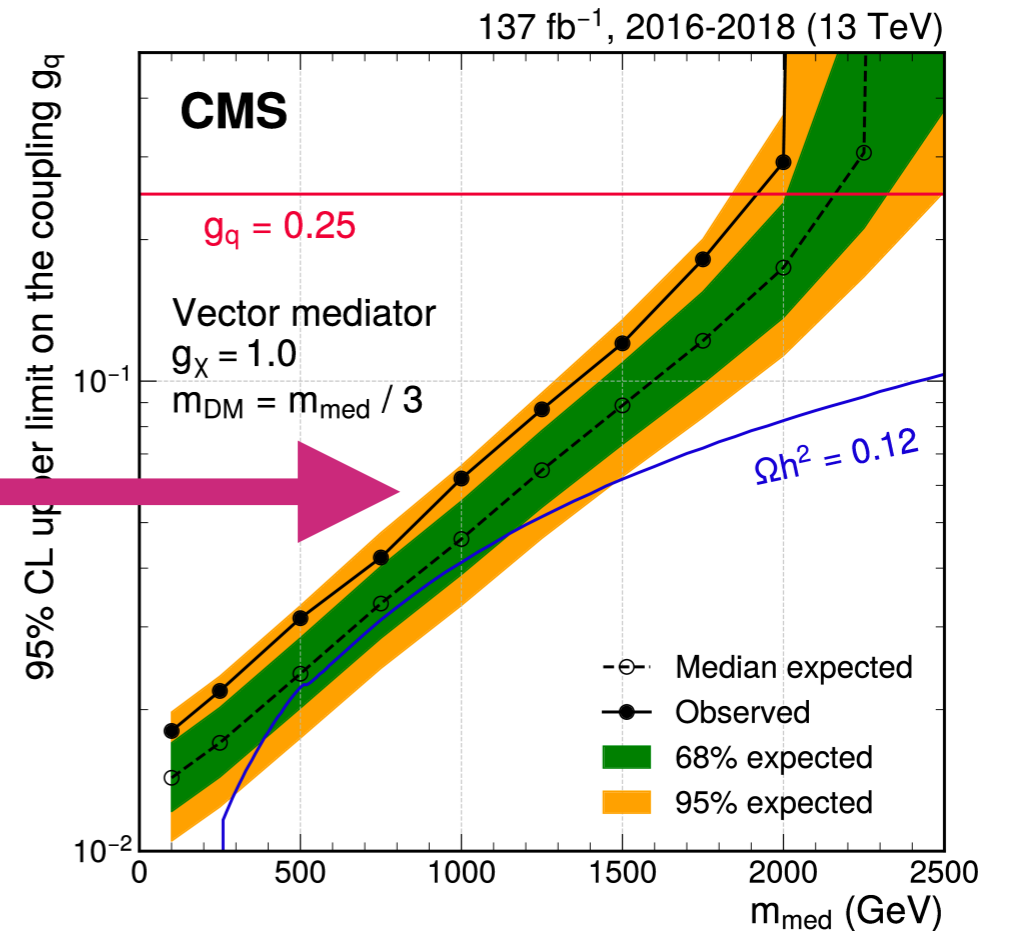
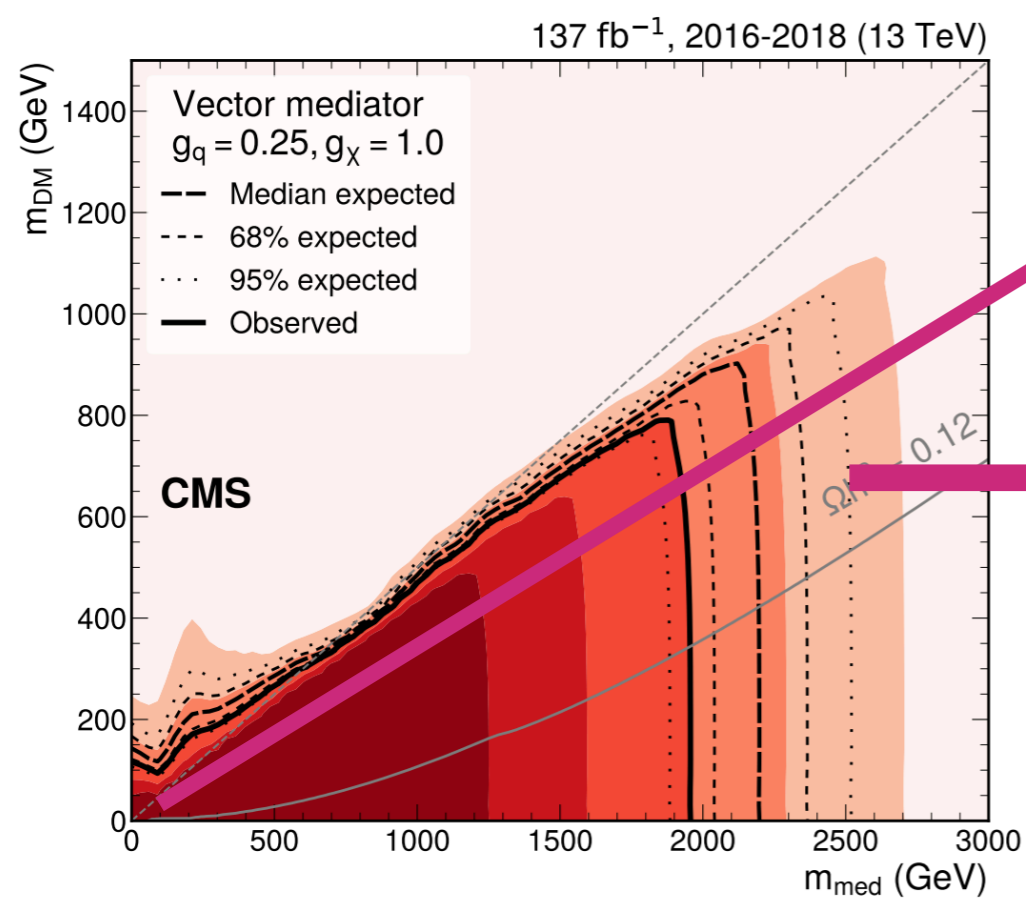
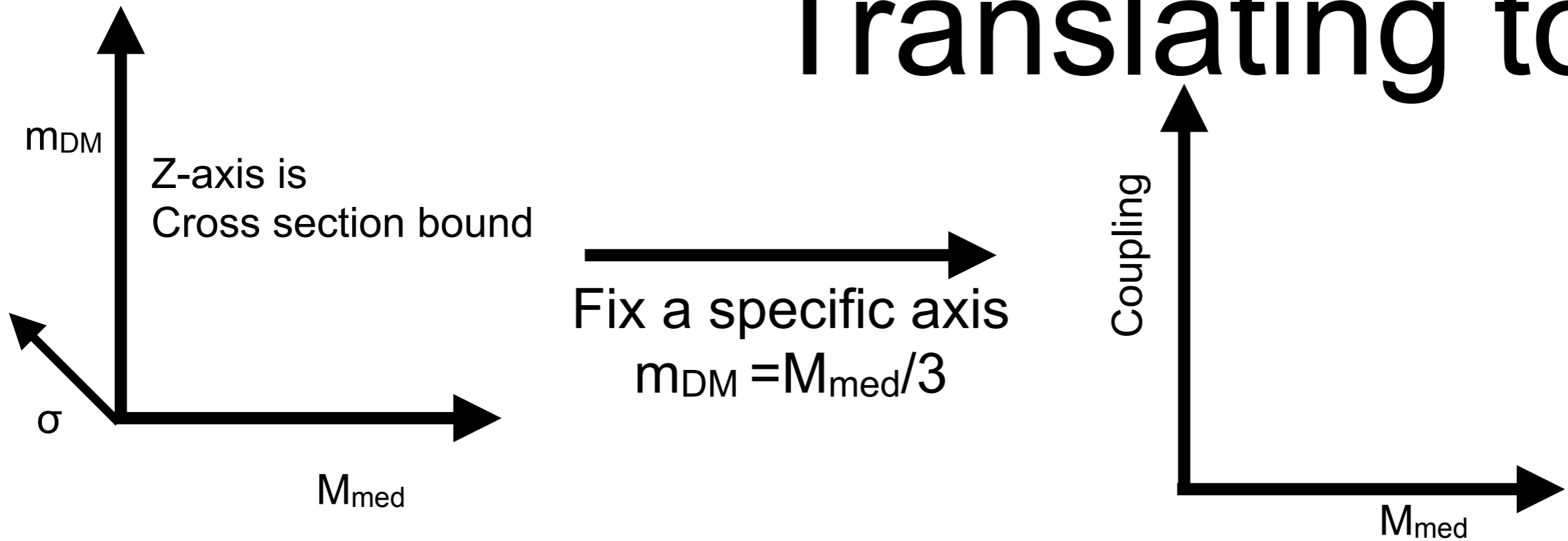
Vector Model : Translate the Couplings and add Z interference

Scalar Model : Add the mixing with the Higgs boson

Pseudoscalar : Translate the Couplings

HNL : Nothing

Translating to y



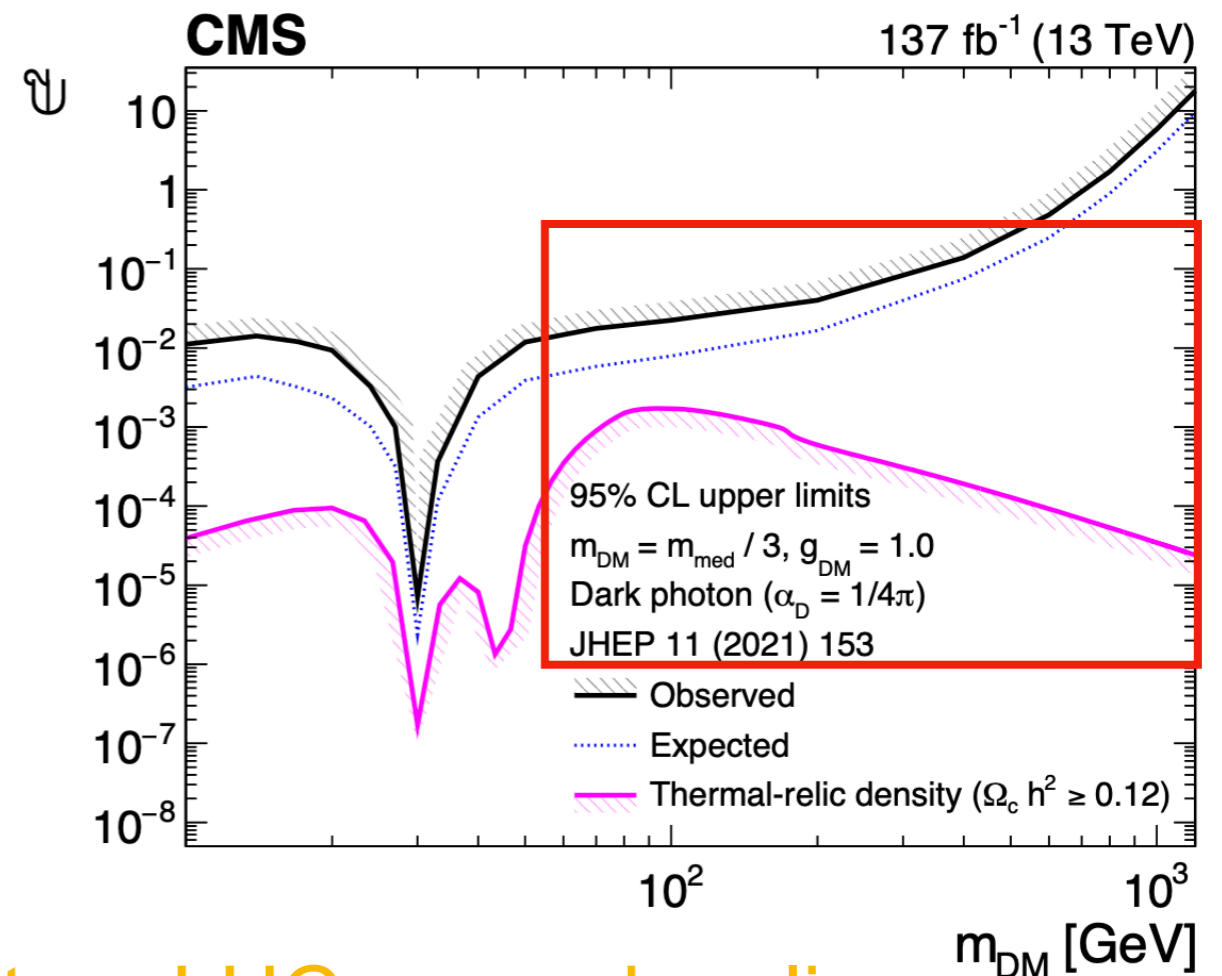
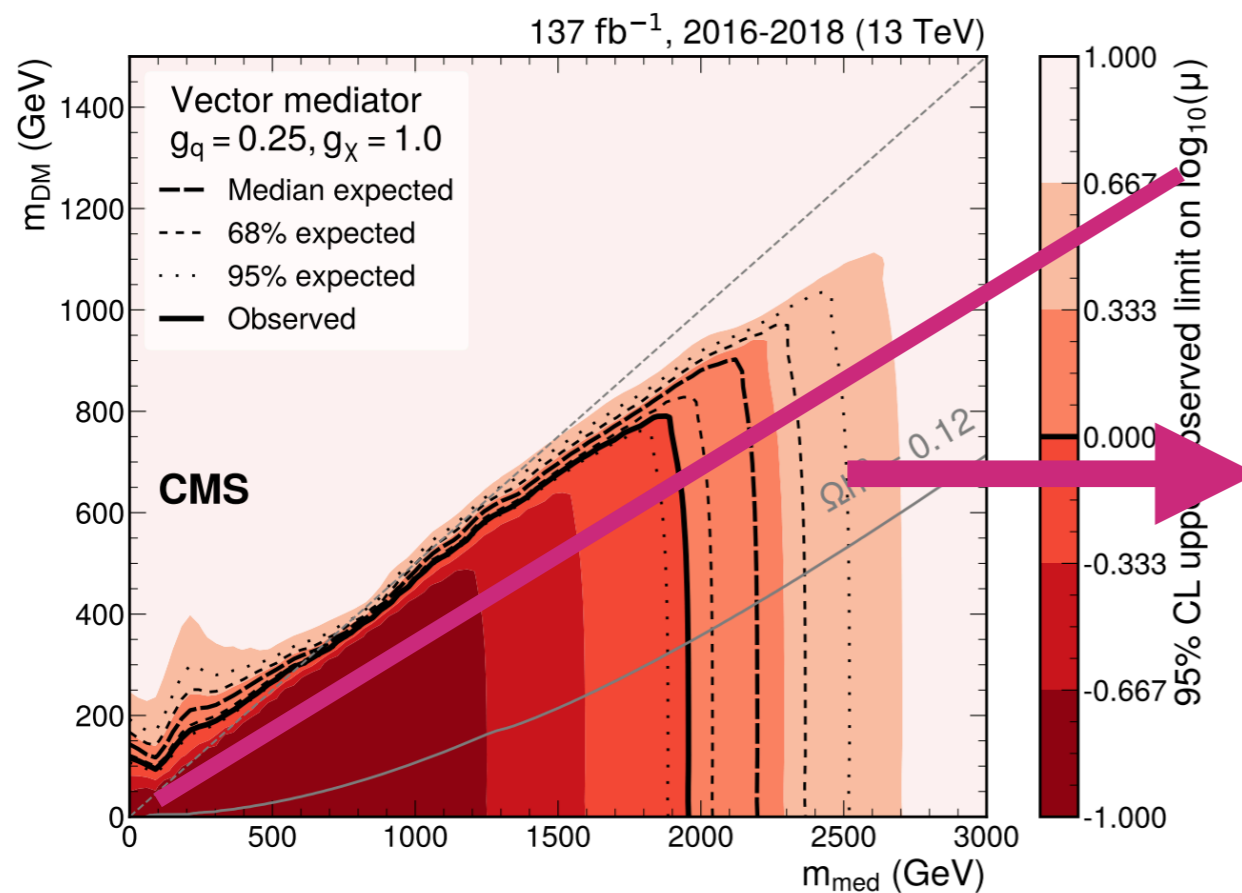
Translating to Dark Photon

$$g_q \simeq \epsilon e \frac{1}{\cos^2 \theta_W} \frac{1}{1 - \Delta_Z} \left(\sqrt{\langle Q^2 \rangle} \cos^2 \theta_W + \Delta_Z \sqrt{\langle Y^2 \rangle} \right)$$

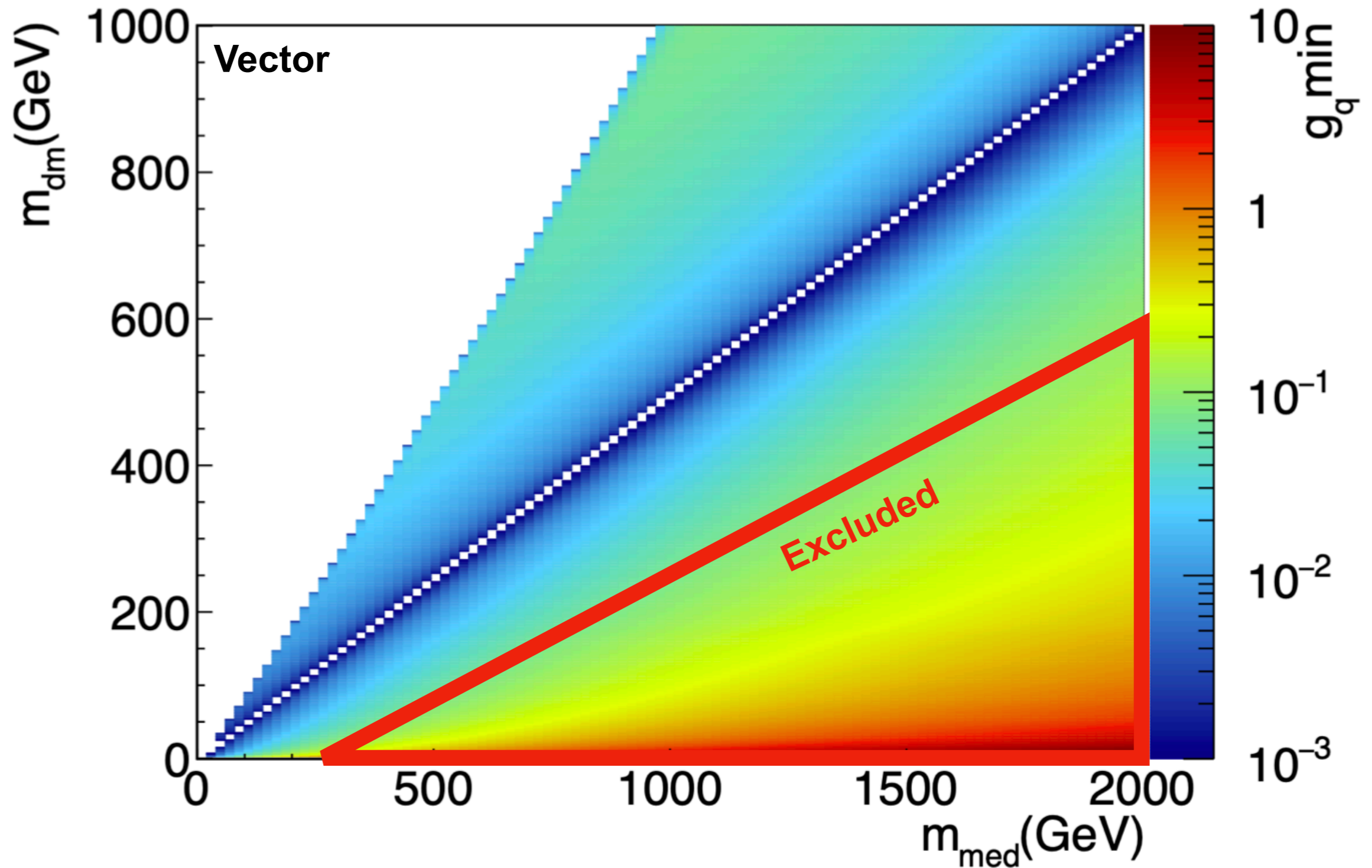
$$\langle Q^2 \rangle = 0.3 \text{ and } \langle Y^2 \rangle \approx 0.7$$

$$\Delta_Z = (m_{Z'} / m_Z)^2$$

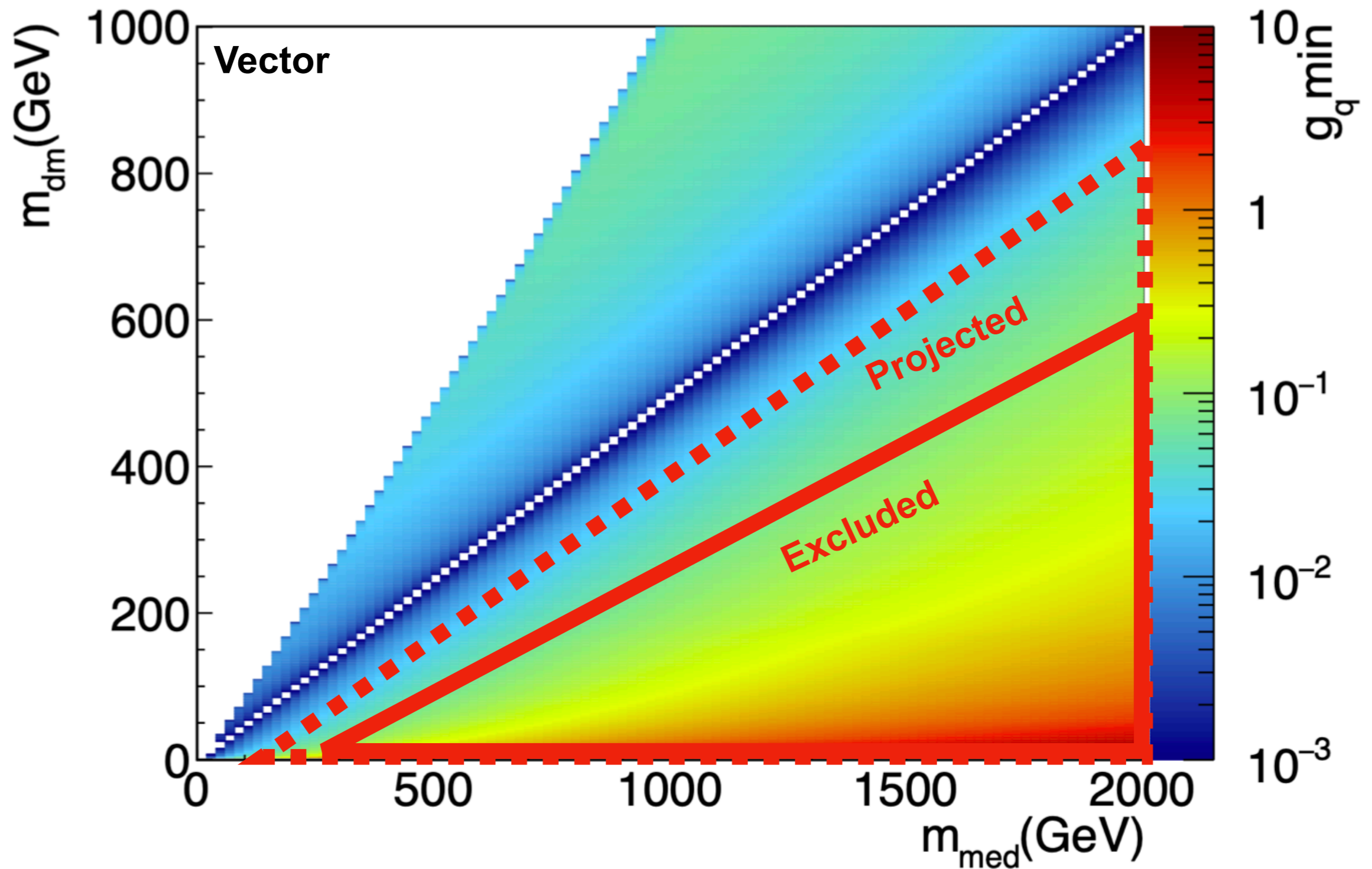
Heavy Dark Photons



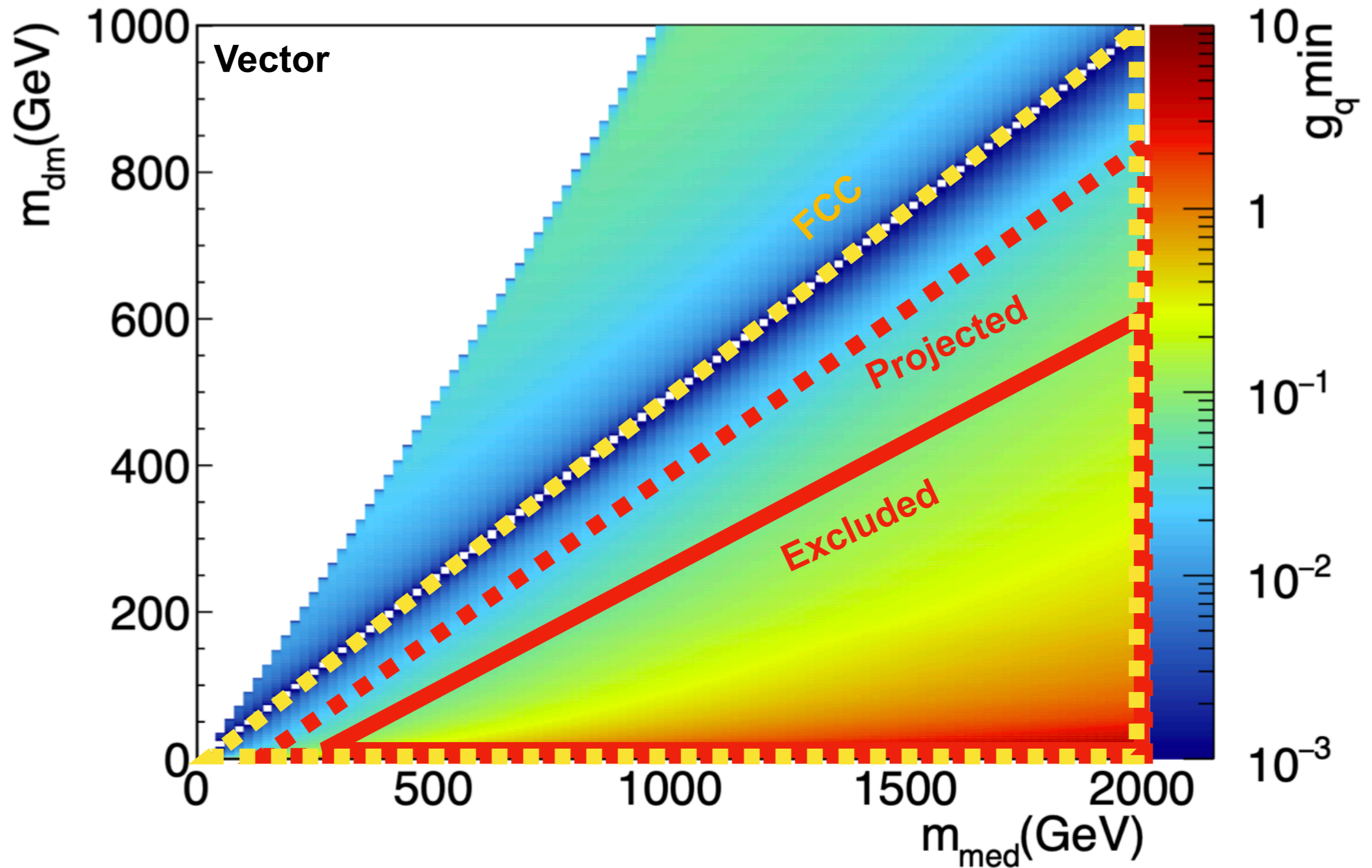
Min Coupling For Relic



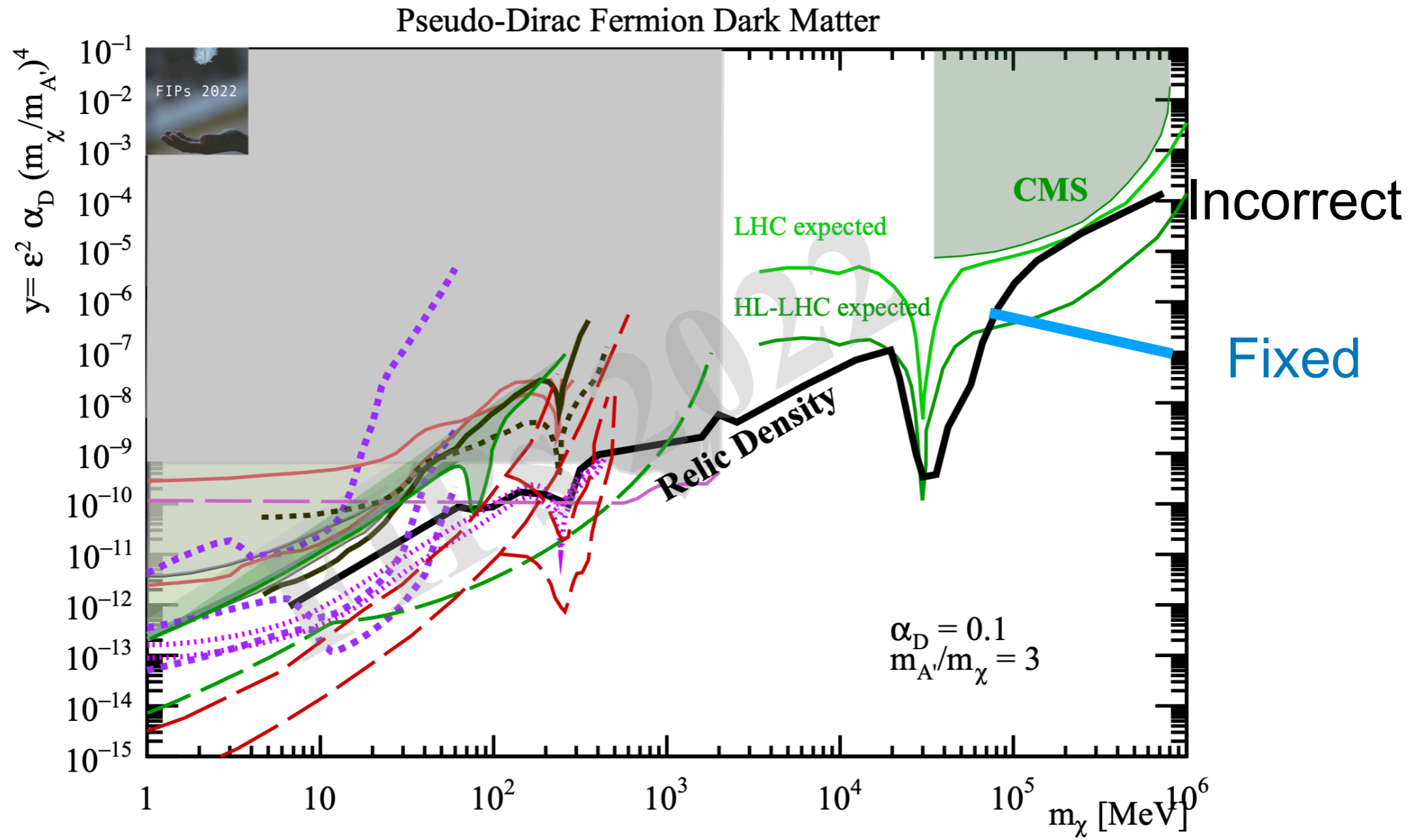
Min Coupling For Relic



Min Coupling For Relic



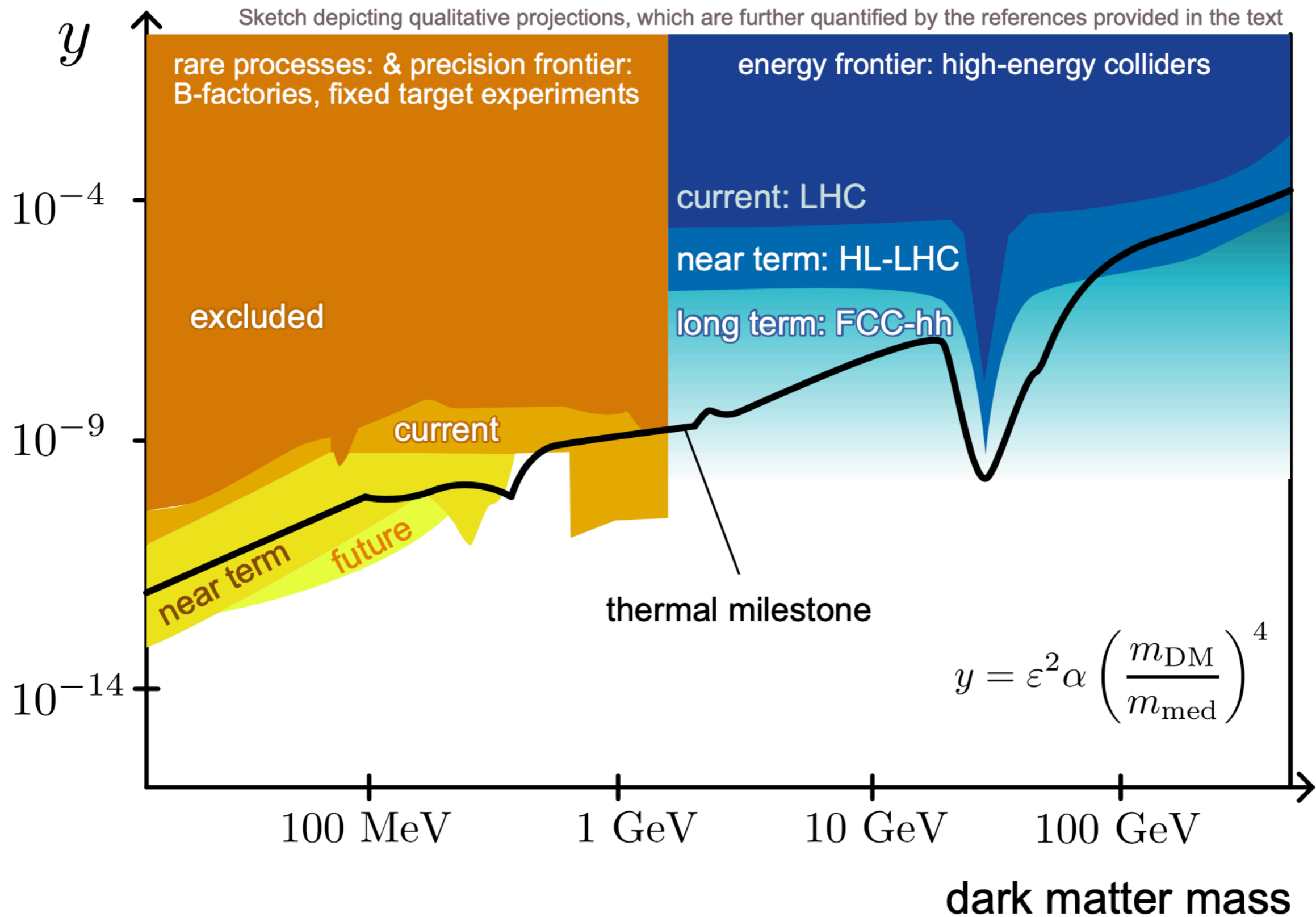
Connecting with low mass



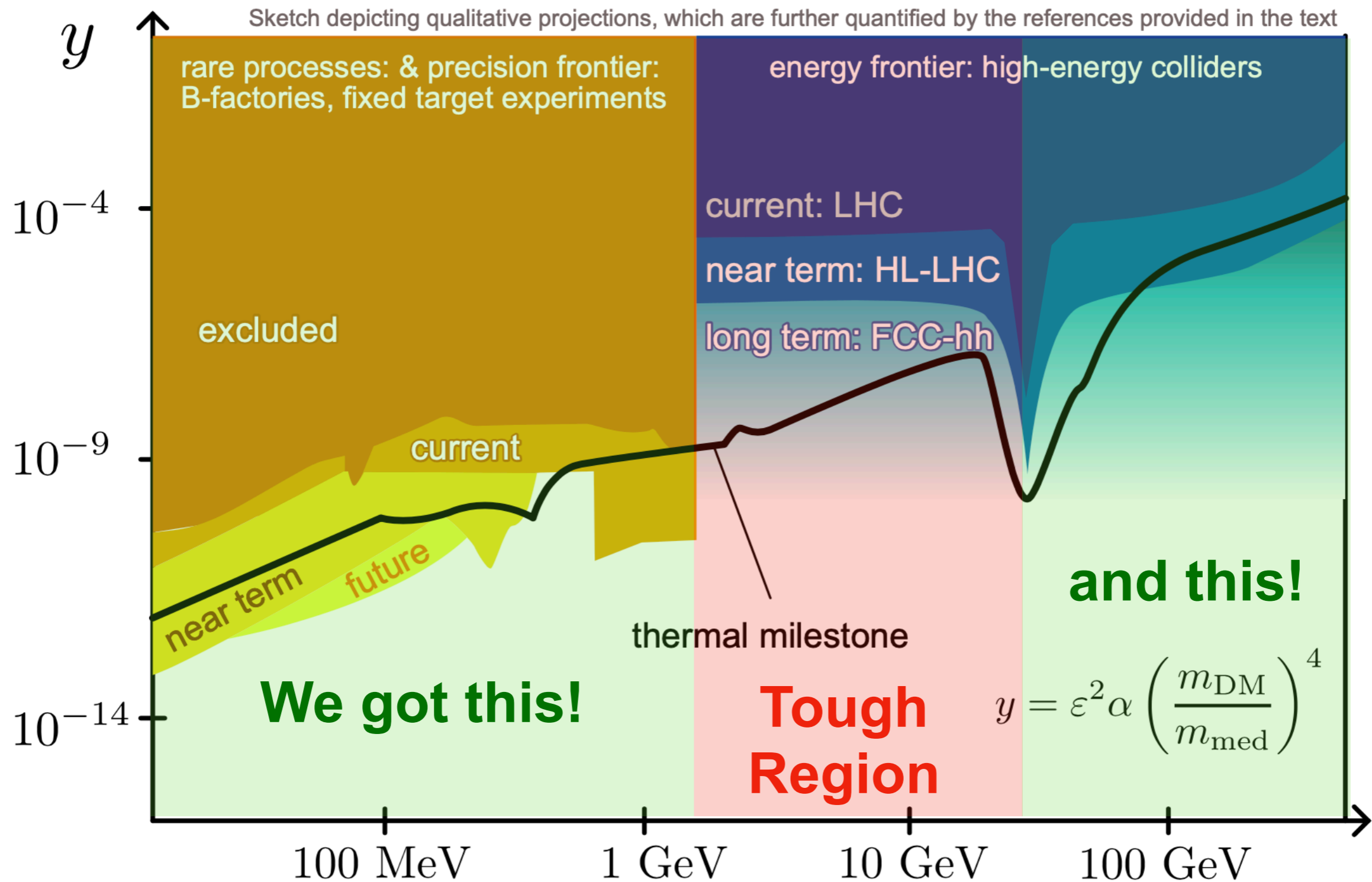
LHC goes left when mediator mass gets larrger

Note Also that as $\frac{m_{DM}}{m_{med}}$ gets larger LHC DM searches are the only game in town

A Nicer Plot



A Better Connection



**Its going to be hard to probe invisible dark Photon
in the intermediate range**

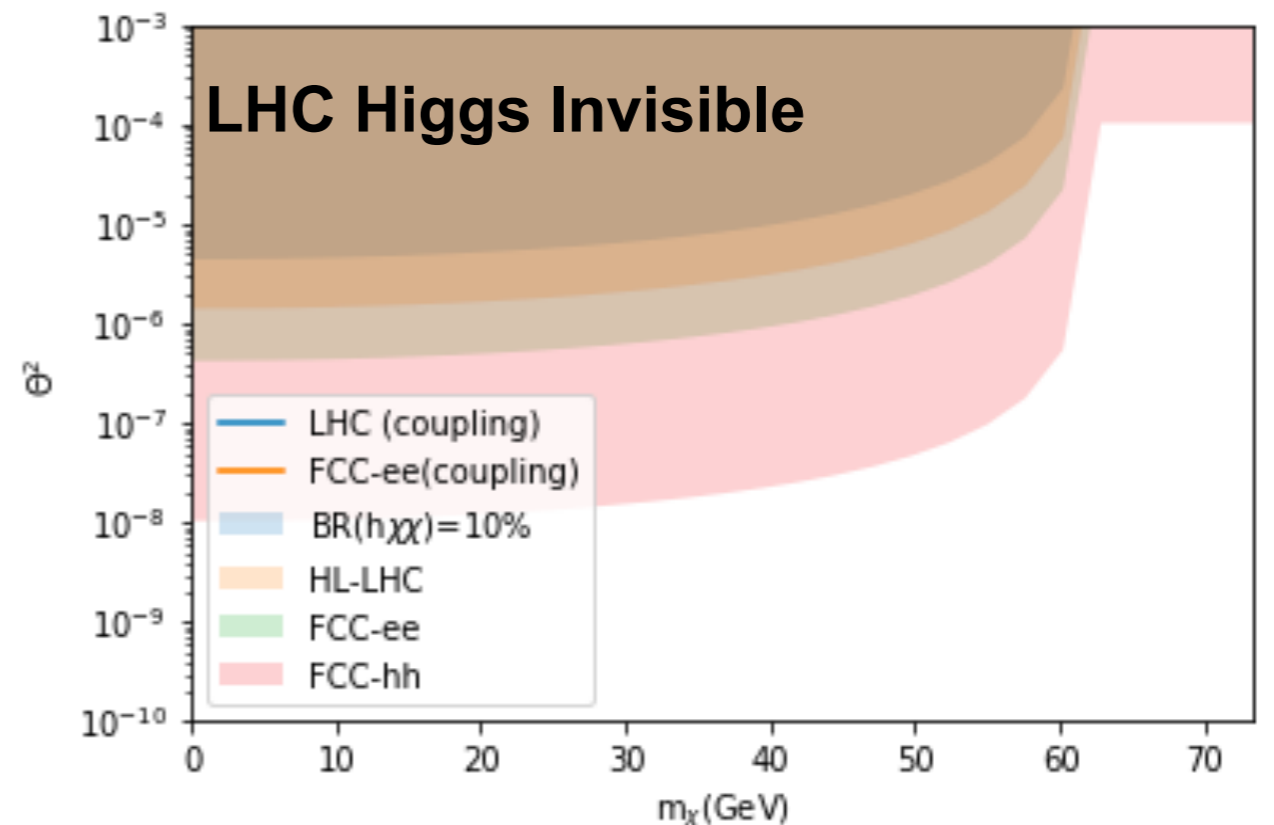
dark matter mass

Scalar Portal(Dark Higgs)

- In a similar vein we can recast Dark Higgs model to LHC

$$\Gamma(h_1 \rightarrow \chi\bar{\chi}) = \frac{y_{DM}^2 \sin^2 \theta m_{h_1}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{h_1}^2}\right)^{3/2}$$

- Higgs to invisible Bounds
 - Current LHC $H(\text{inv}) > 0.1$
 - Future LHC $H(\text{inv}) > 0.02$
 - FCC-ee $H(\text{inv}) > 0.005$
 - FCC-hh $H(\text{inv}) > 0.0001$

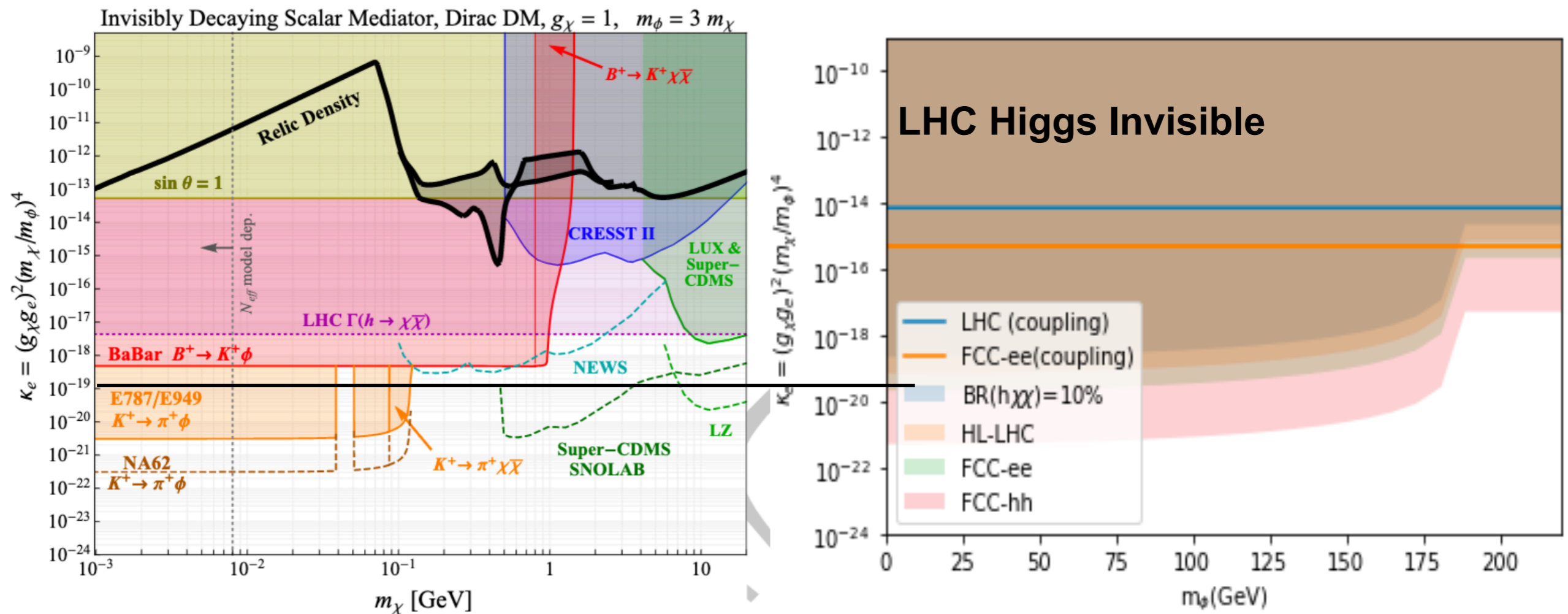


- It is hard to have the Dark Higgs model explain DM (Natalia's talk)
 - Requires a very large coupling to satisfy relic

Scalar Portal (Dark Higgs)

- In a similar vein we can recast Dark Higgs model to LHC

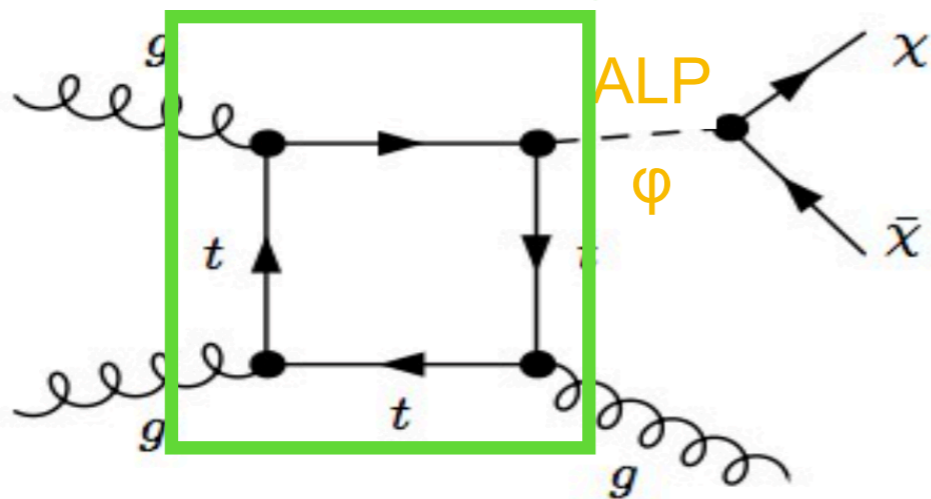
$$\Gamma(h_1 \rightarrow \chi\bar{\chi}) = \frac{y_{DM}^2 \sin^2 \theta m_{h_1}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{h_1}^2}\right)^{3/2}$$



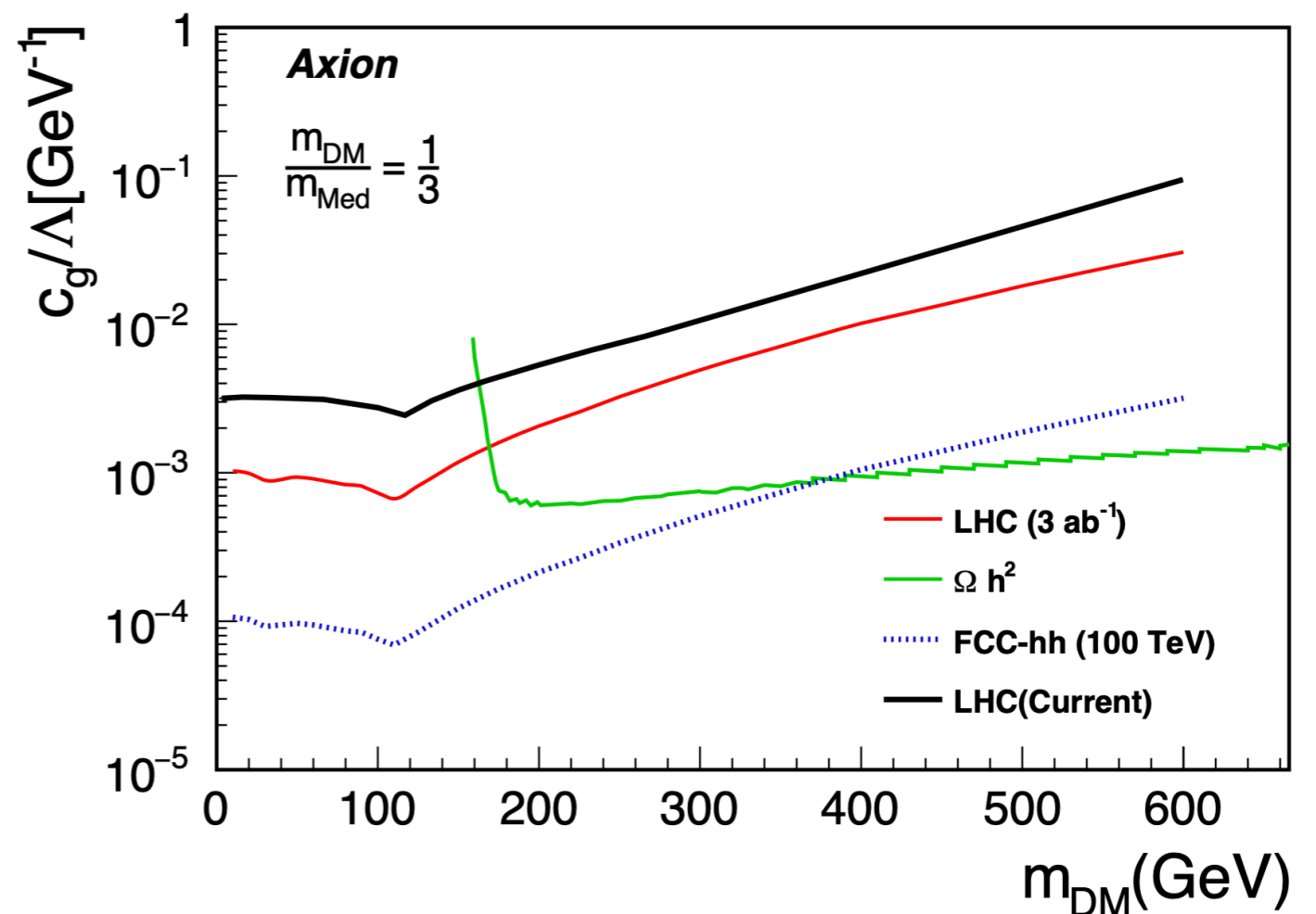
Axion Portal

- The Axion portal decaying to invisible particles
 - Generally less explored
 - We can again recast existing LHC bounds to this
- The simplest model requires a heavy mediator to reconcile relic

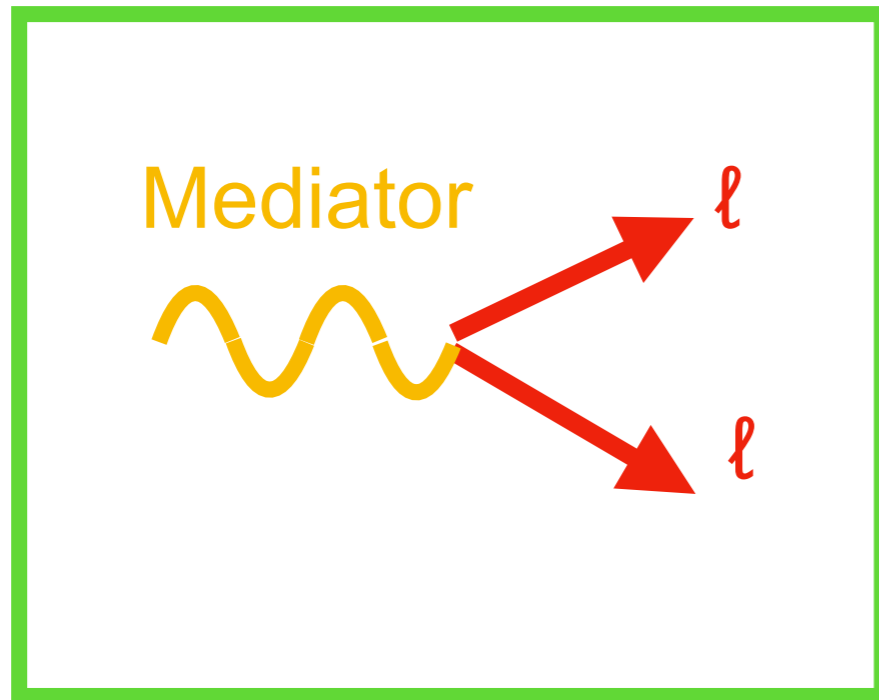
Just Gluon Coupling



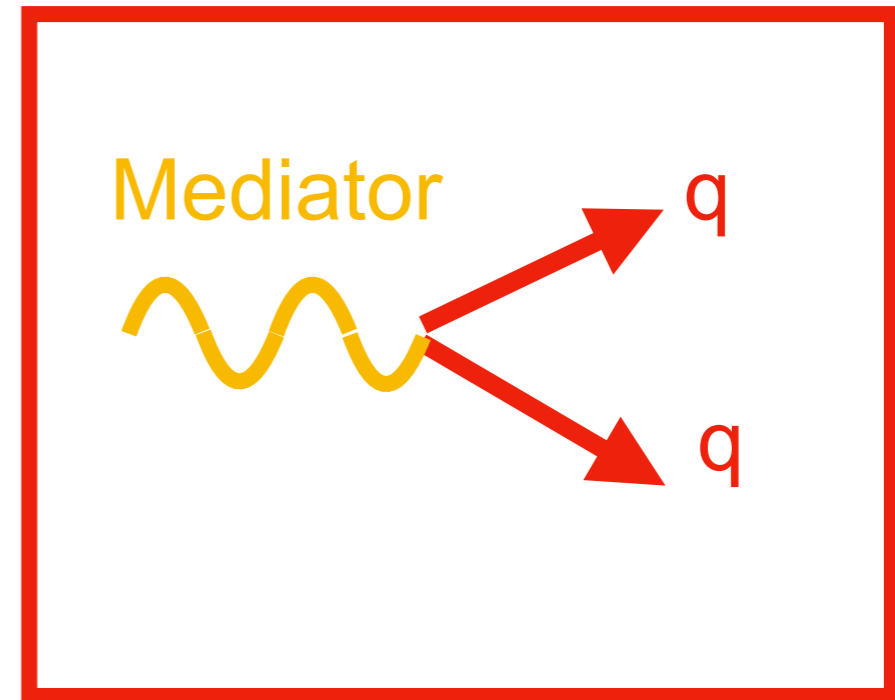
Axion, a $\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}$, $\frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}$ $\frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$



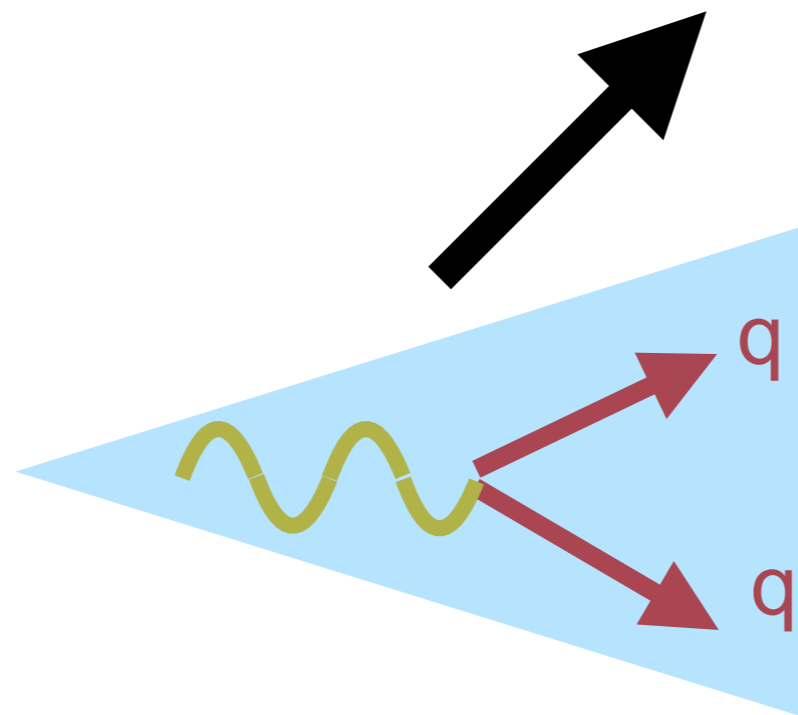
Visible Searches



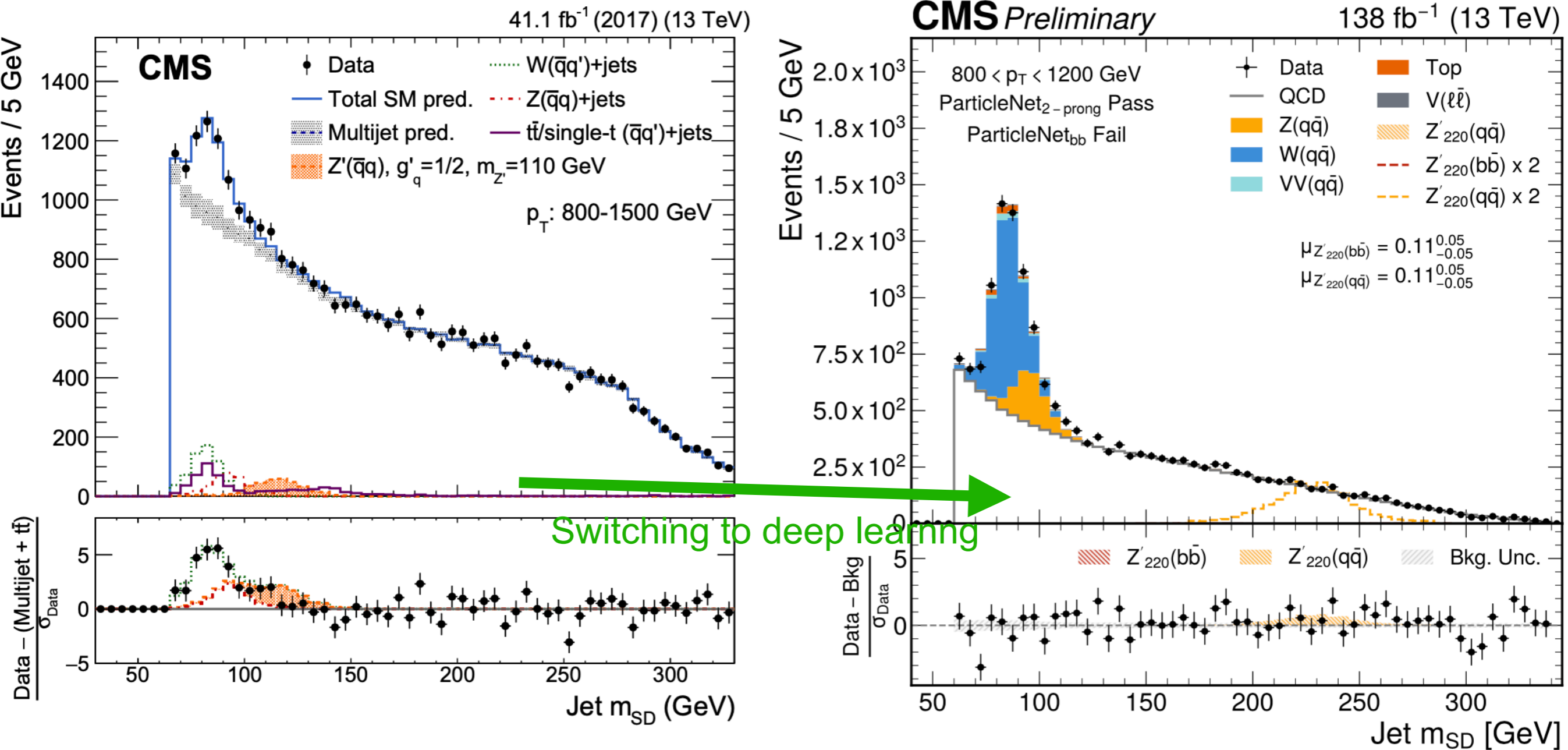
LHC Produces a huge number of leptons



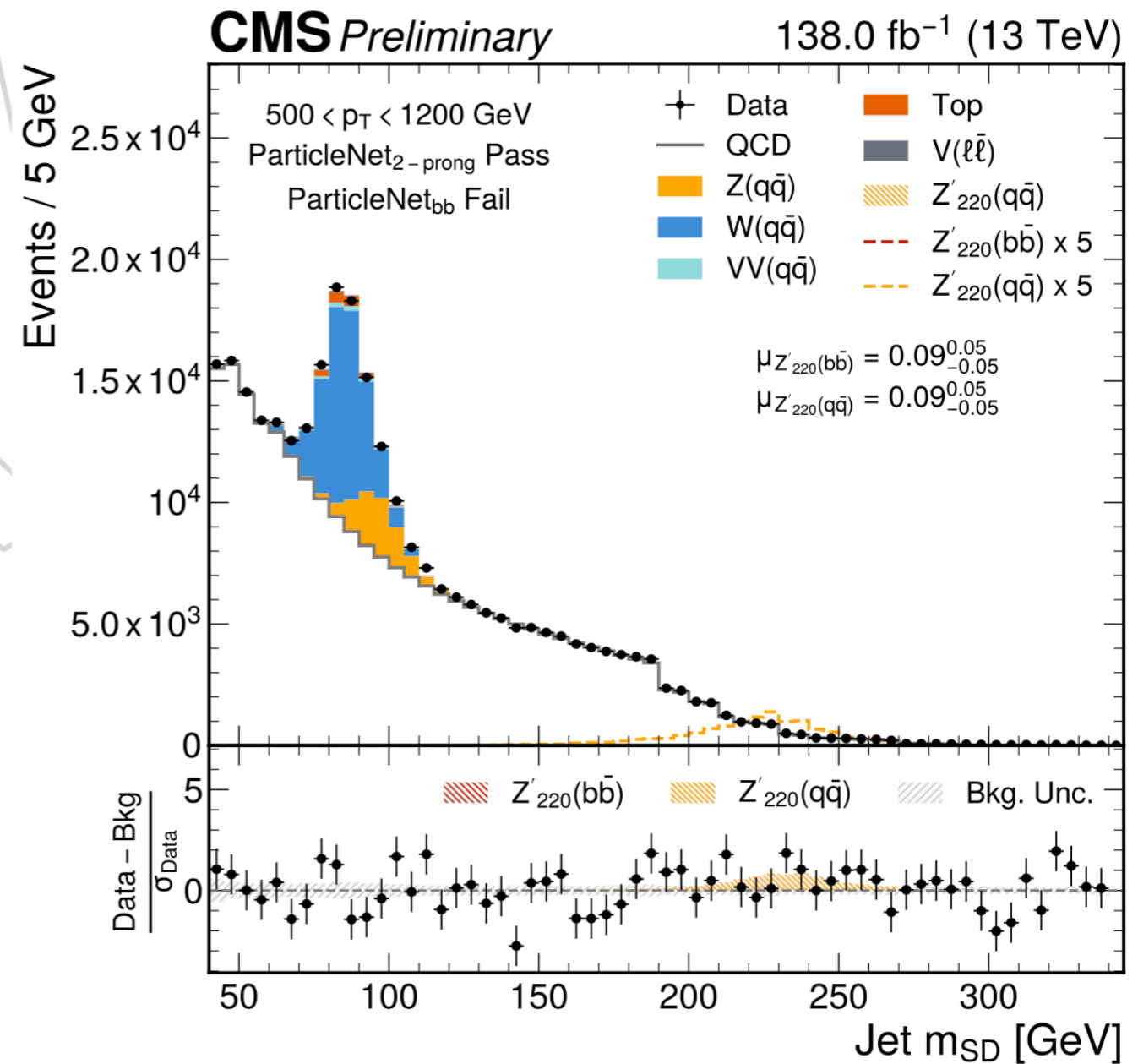
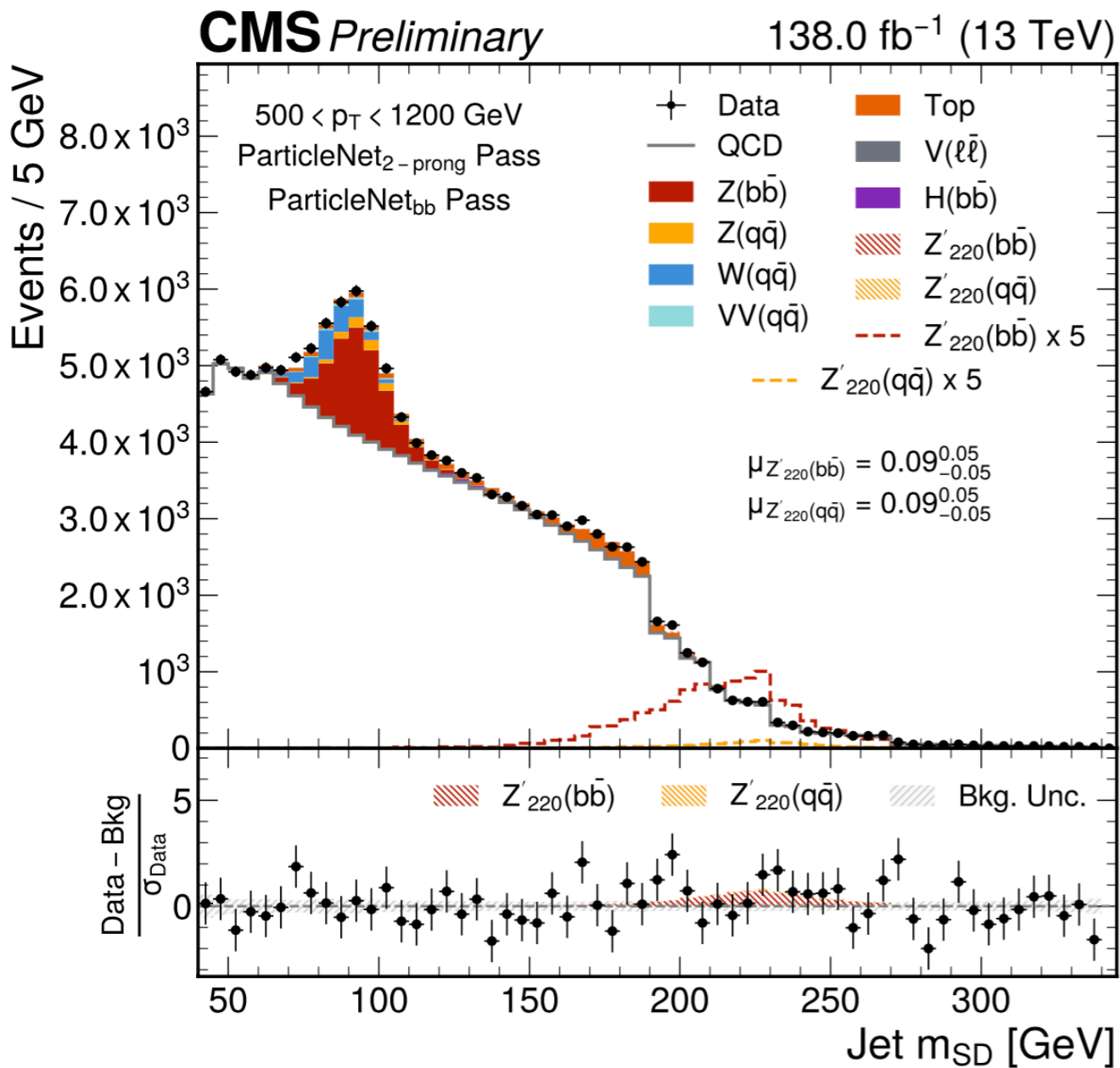
This is very hard but we still do it



Resonant Searches

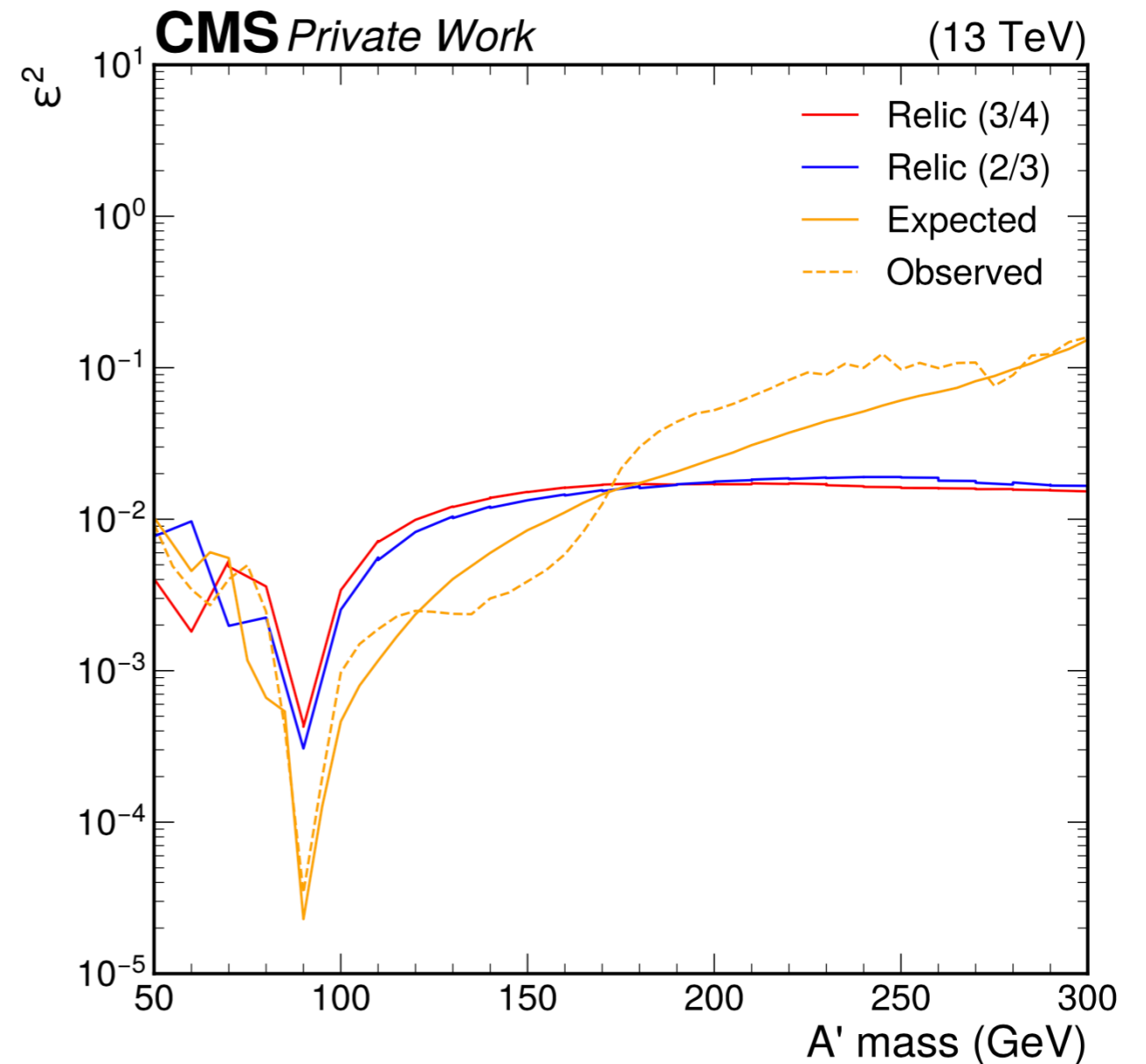
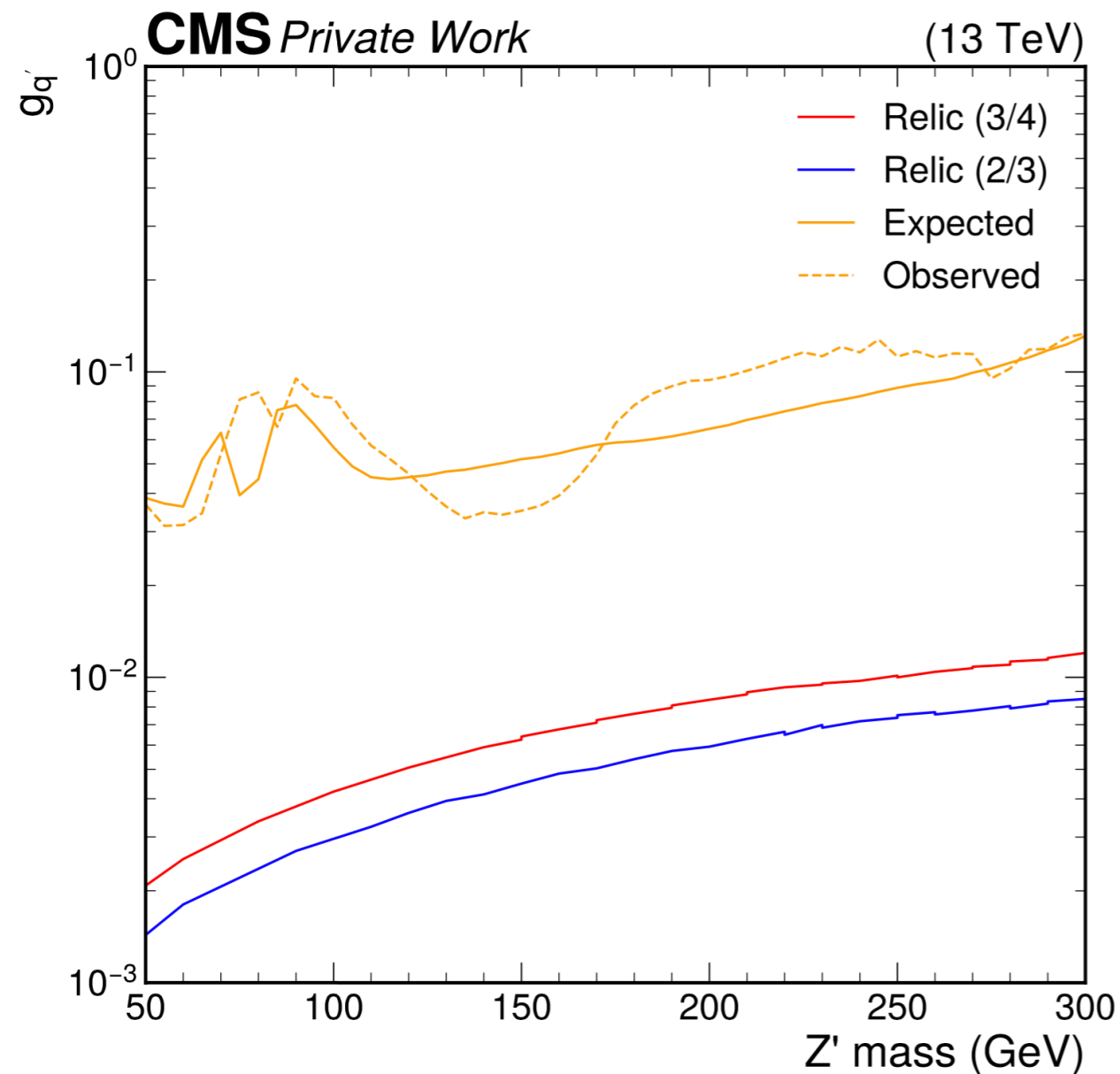


The Search



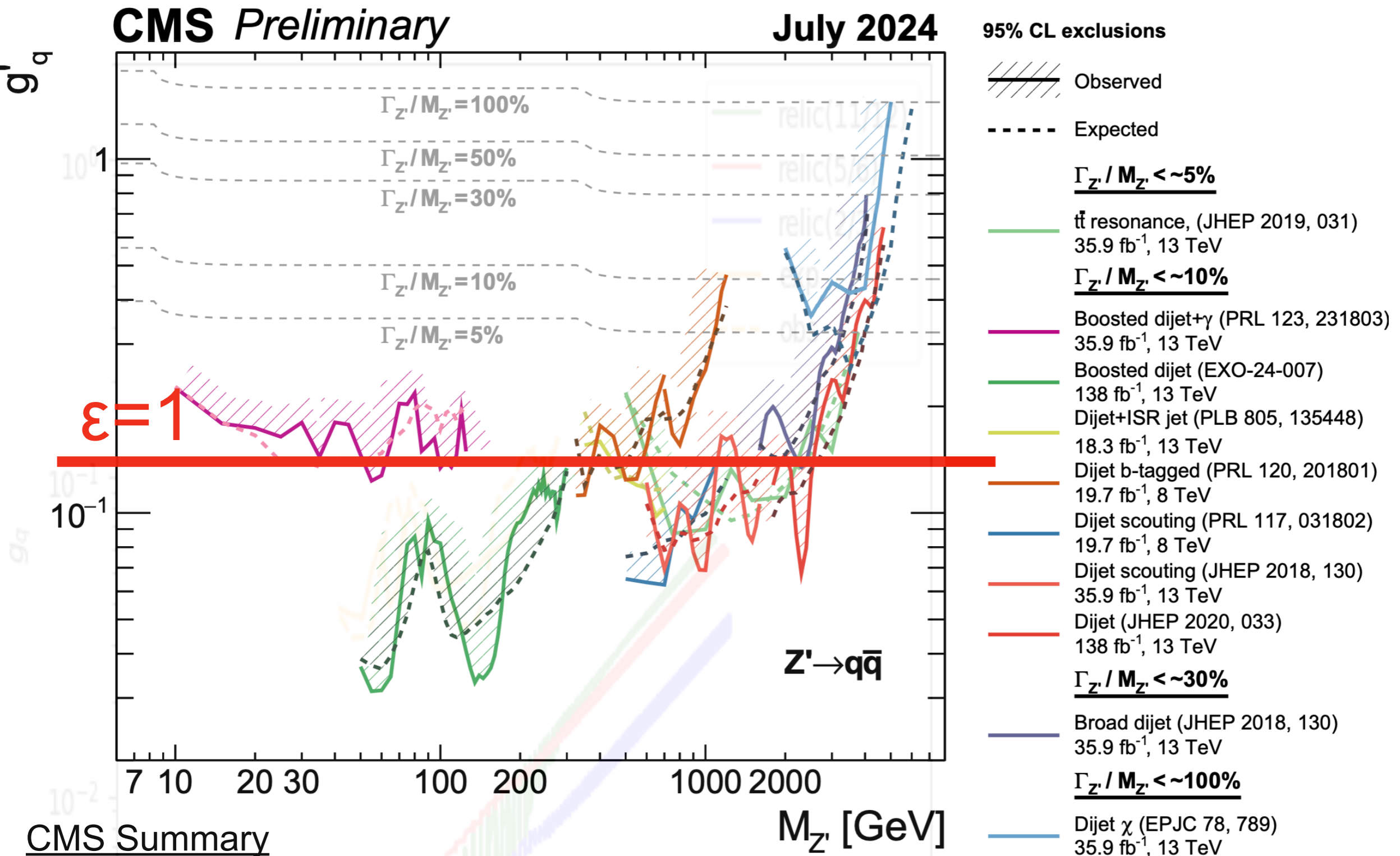
- More Hadronic W bosons here than at LEP
 - These are very highly boosted

Comparison w/dark sector

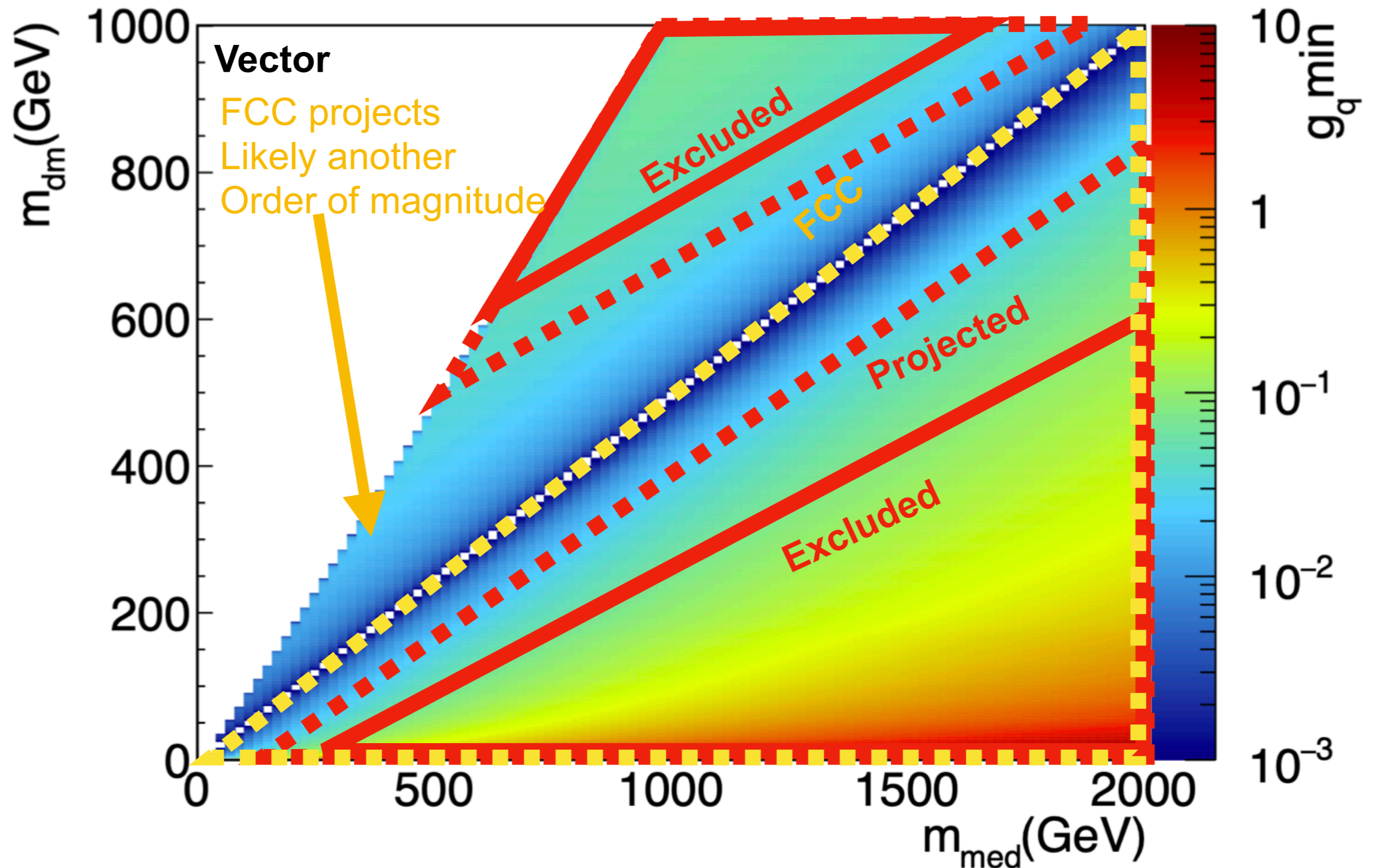


- Significant constraints that start to approach interesting regions
 - Relic density is just a motivator here, its not a true benchmark

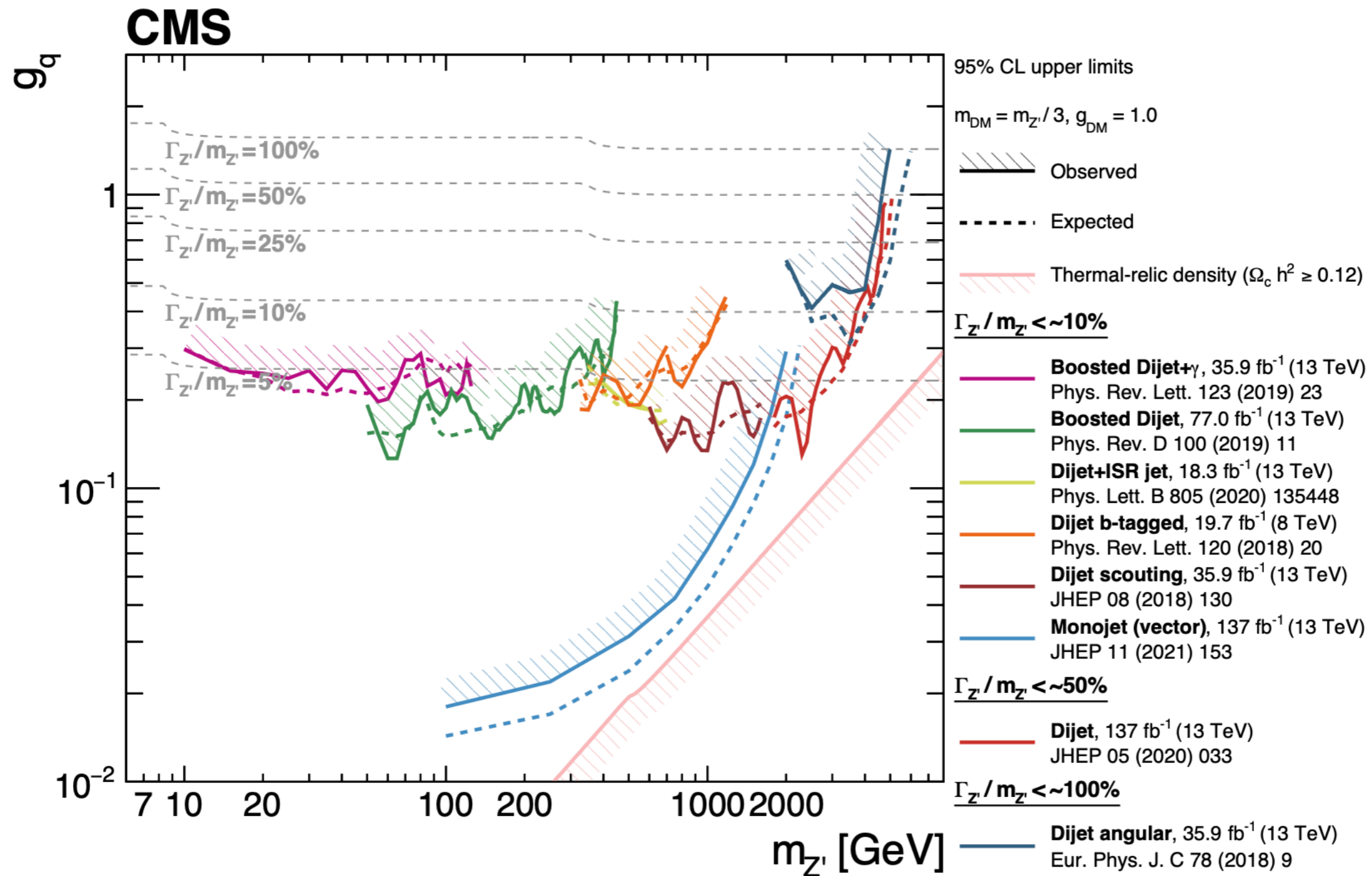
Just Quark Couplings



Min Coupling For Relic

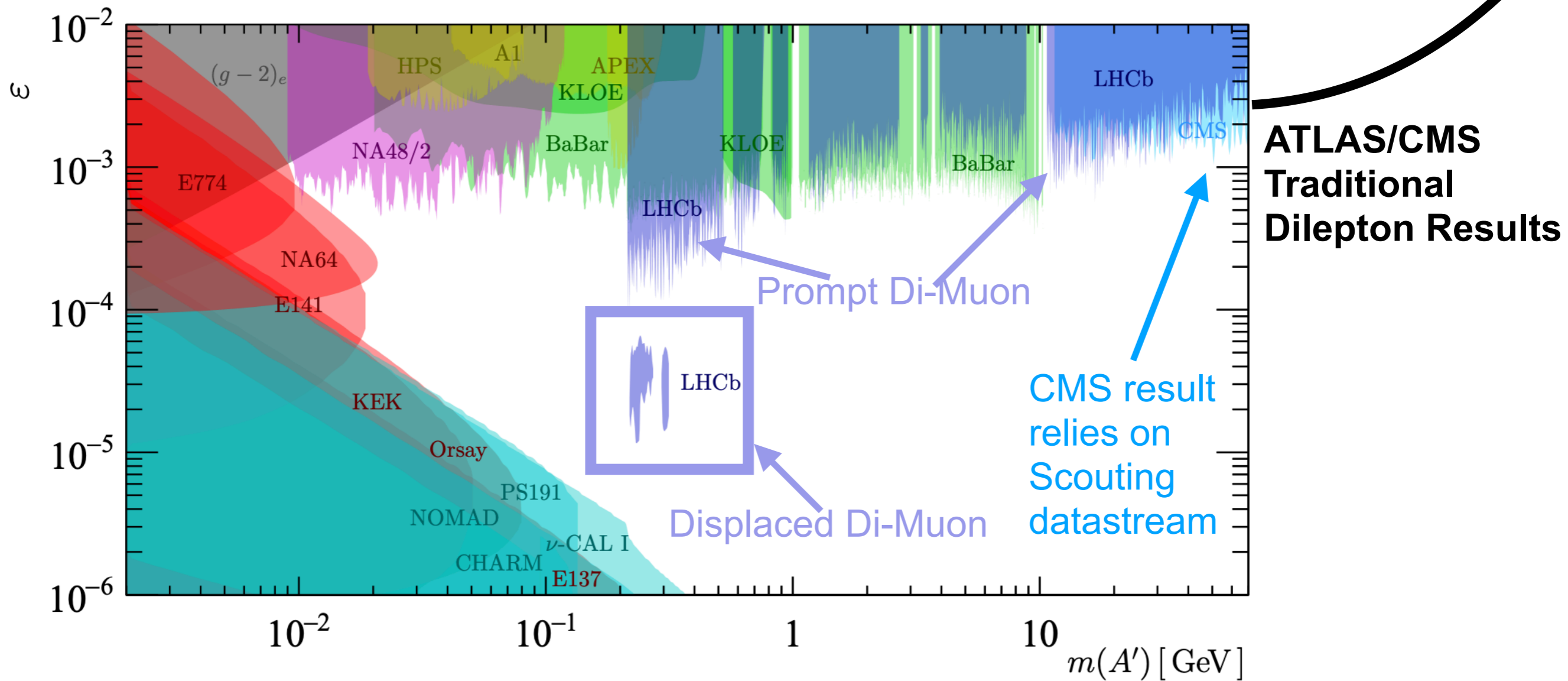


Now considering Invisible



Dark Photons

Visible Dark Photons at High Mass dominated by LHCb+CMS

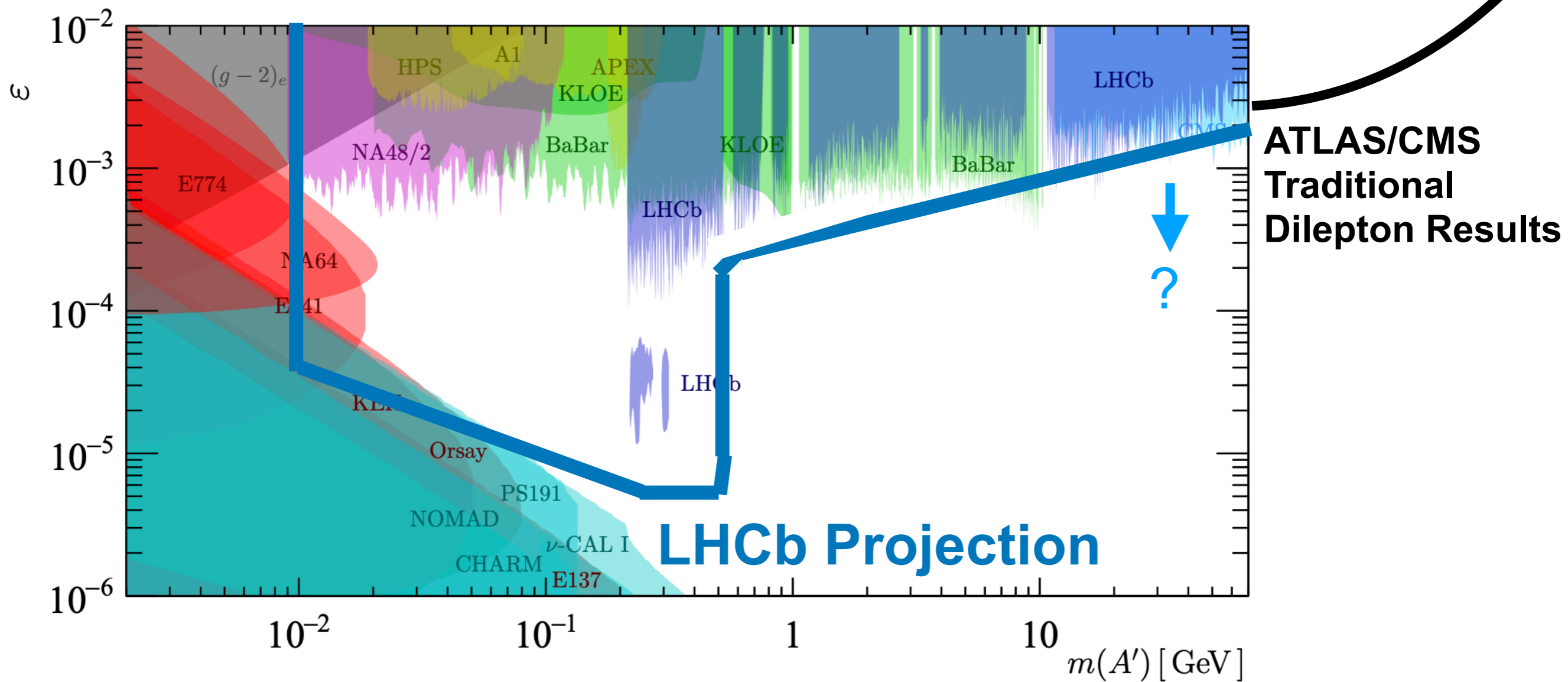


LHCb capable of triggering all di-muons

CMS result relies on innovative Scouting Stream

Future Dark Photons

Addition of di-electron channel makes LHCb very sensitive

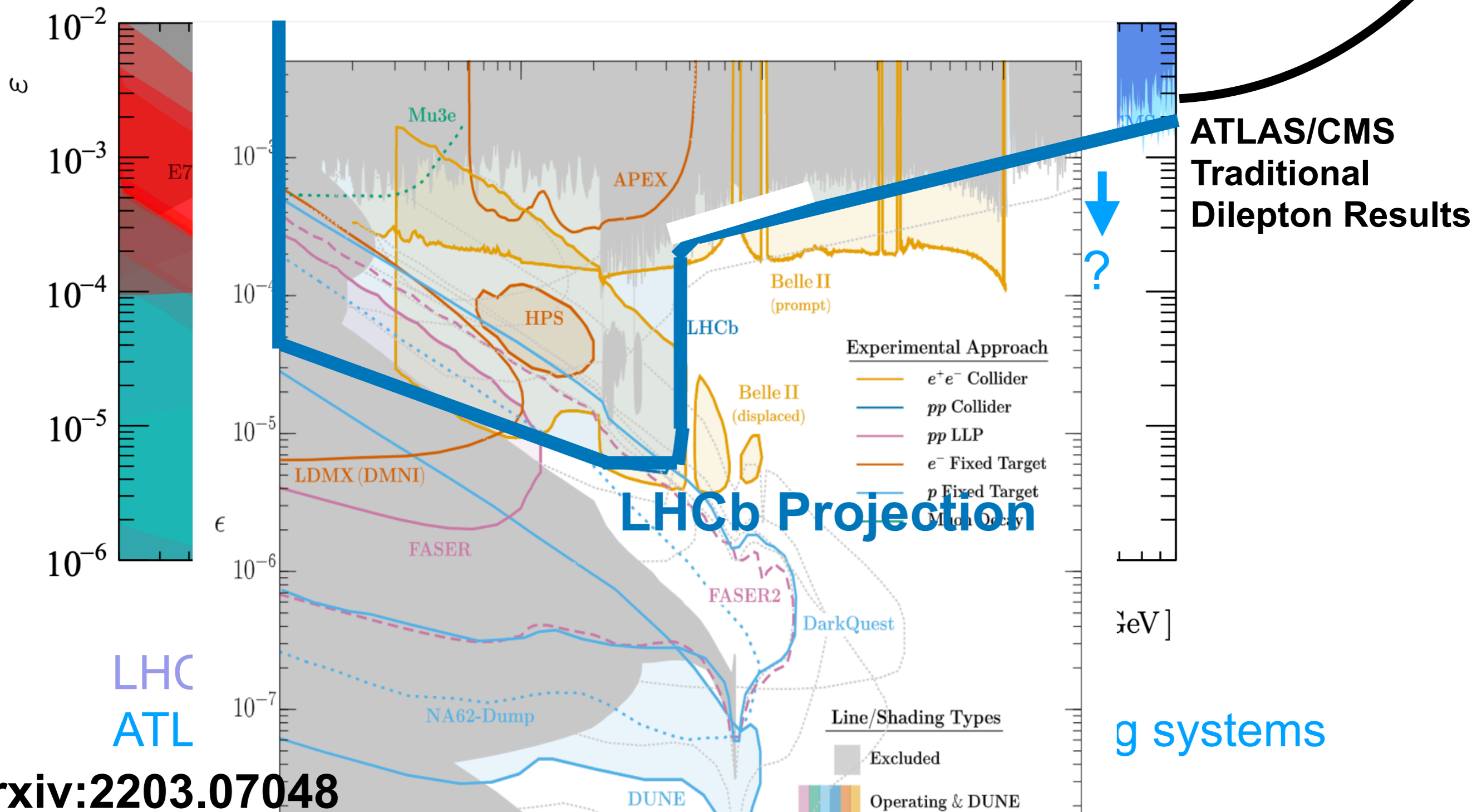


LHCb will add $D^* \rightarrow A + \text{SM Decays}$

ATLAS and CMS will have better trigger/scouting systems

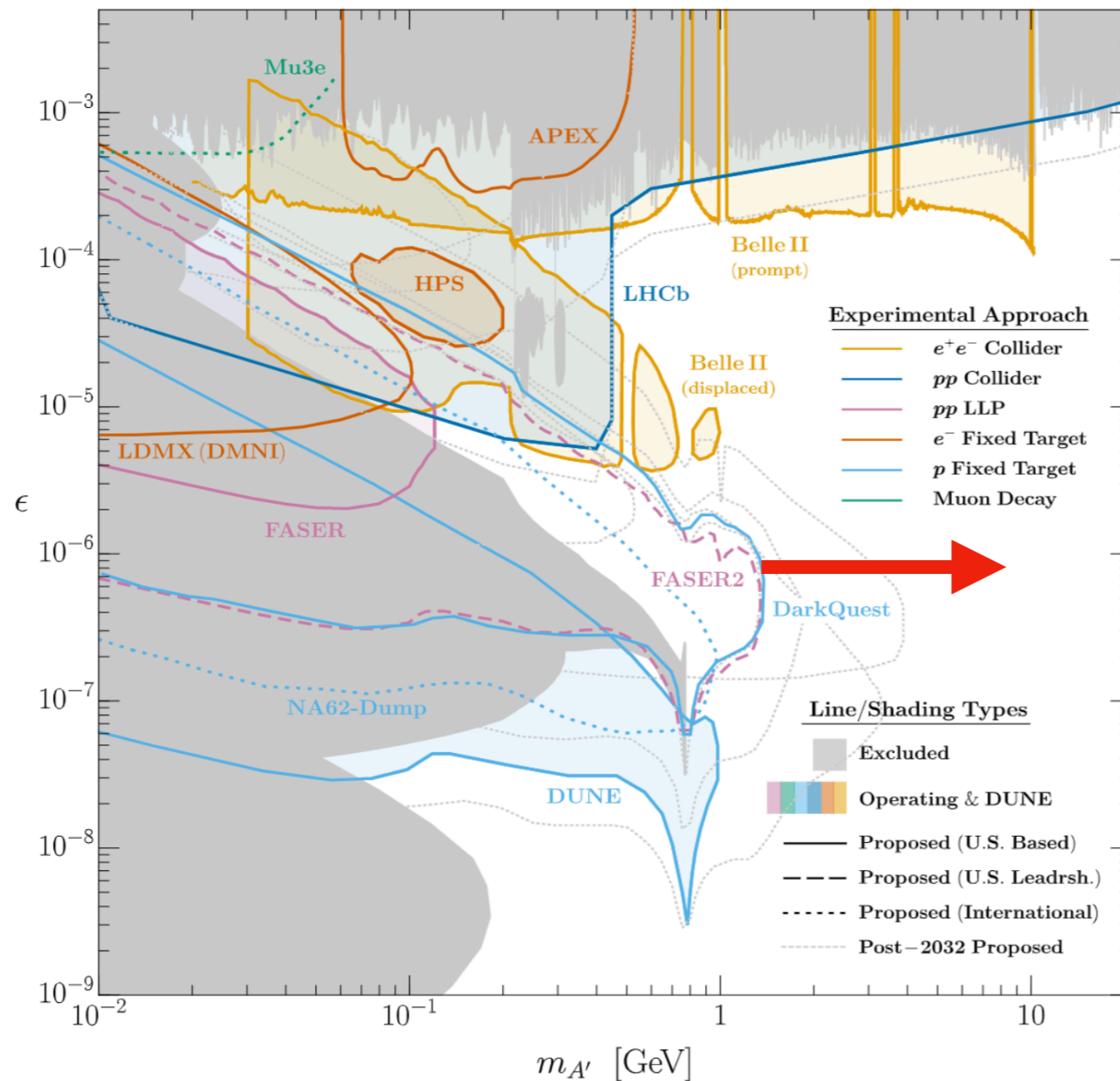
Future Dark Photons

Addition of di-electron channel makes LHCb very sensitive



Future Dark Photons

Addition of di-electron channel makes LHCb very sensitive

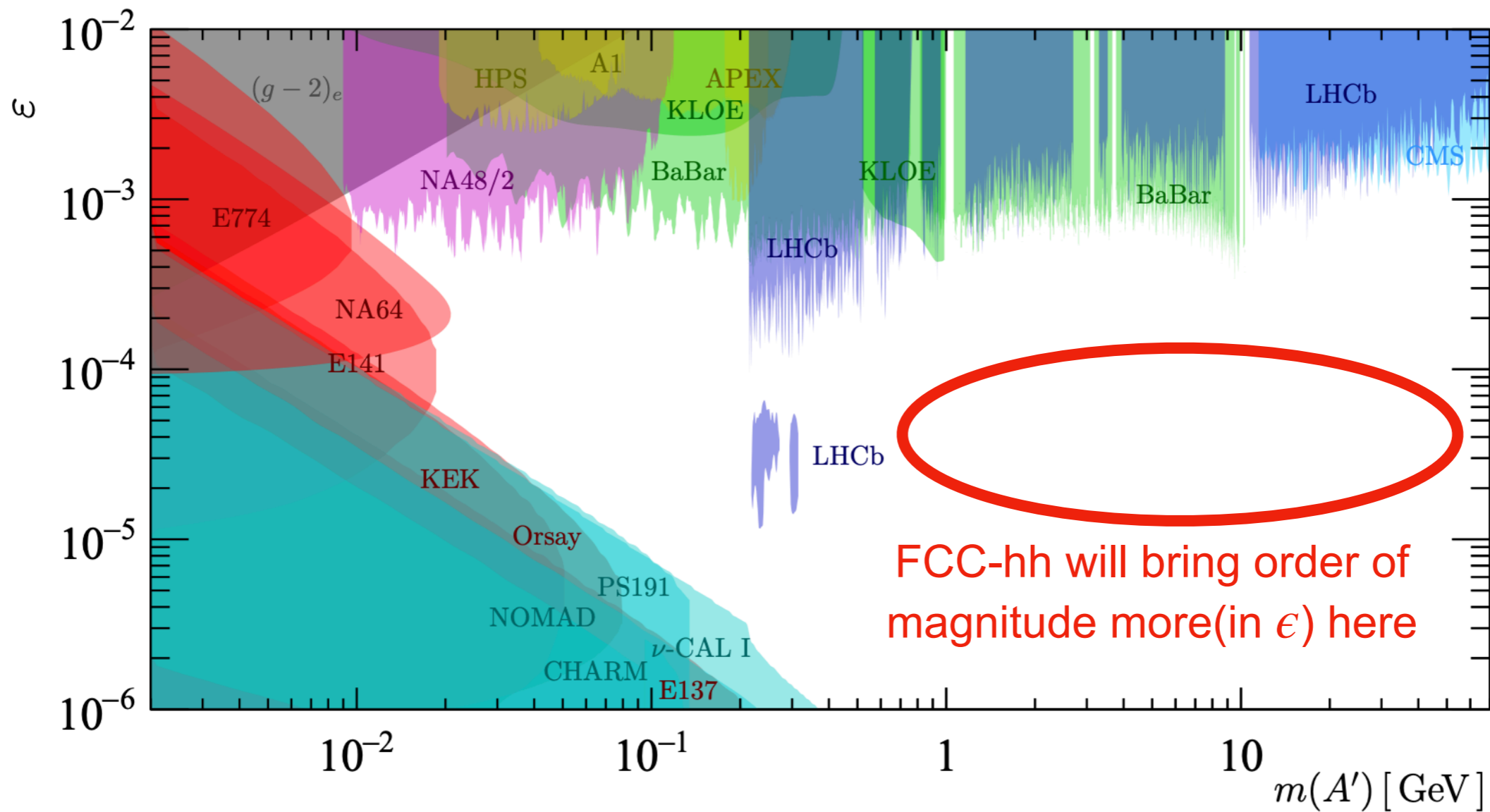


A small coupling is still motivated for heavy Dark Matter and a light mediator

Higher COM beam dumps help here?

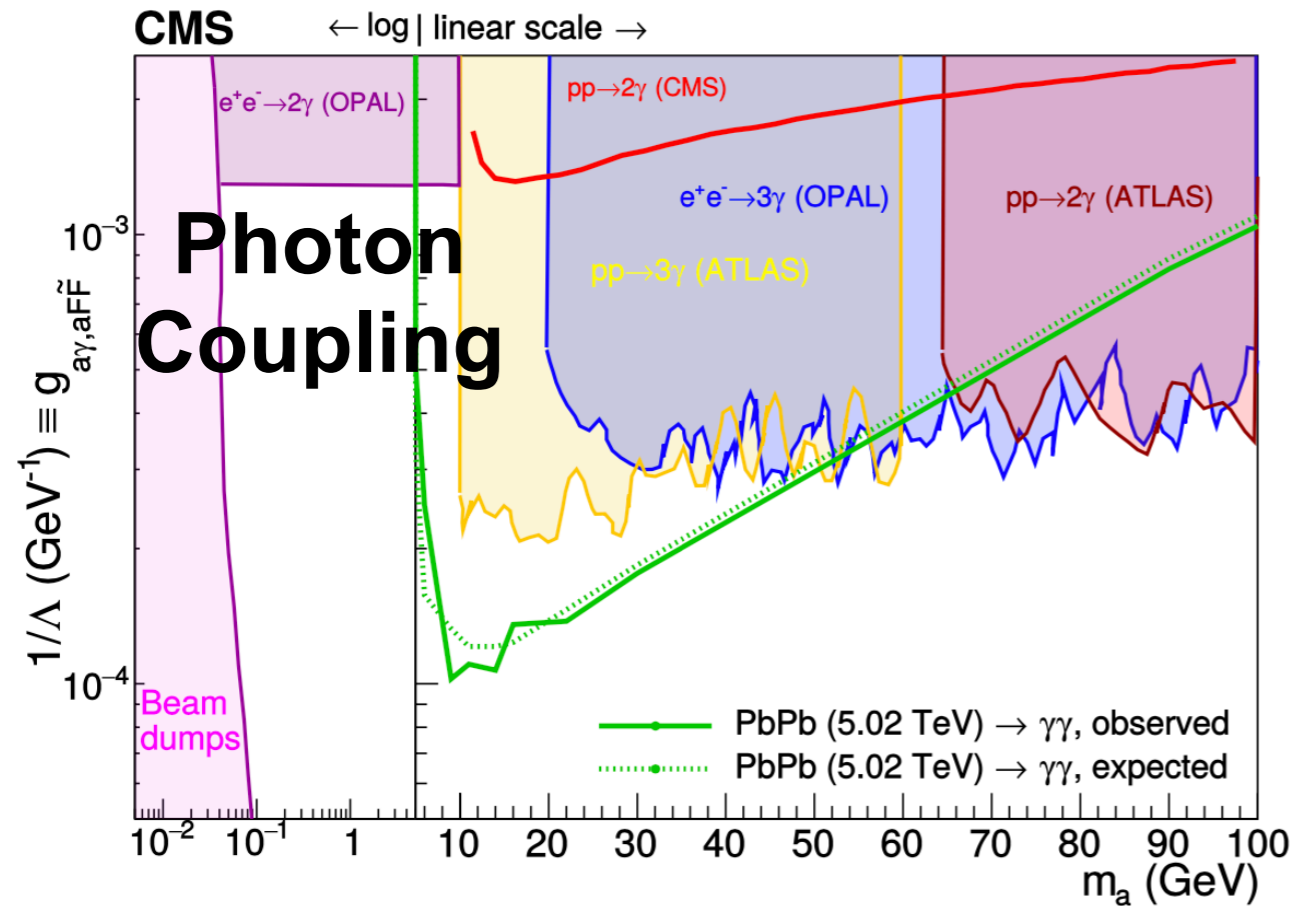
Beyond LHC?

FCC-hh or a higher energy beam dump would help
Not clear that others would

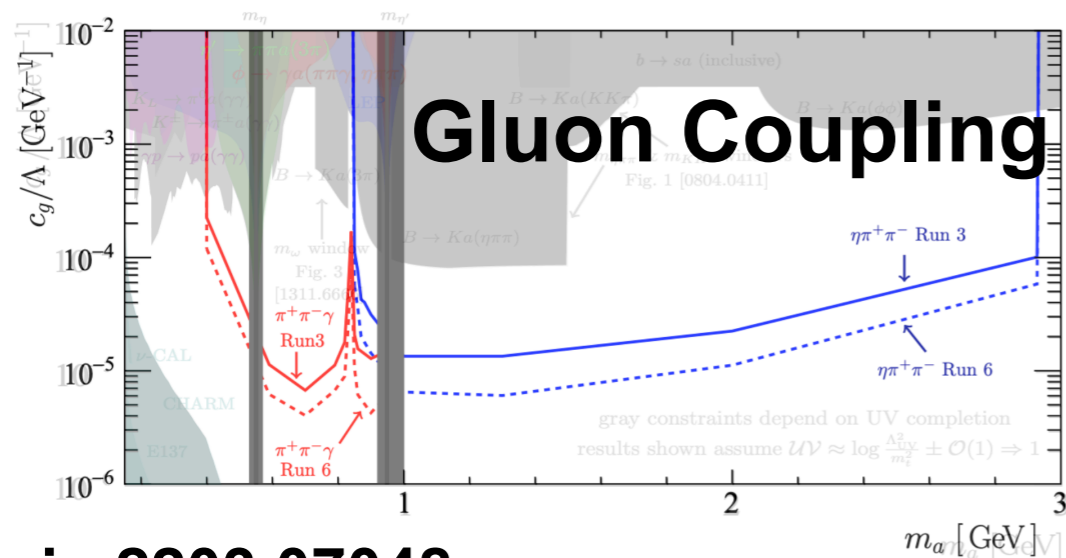
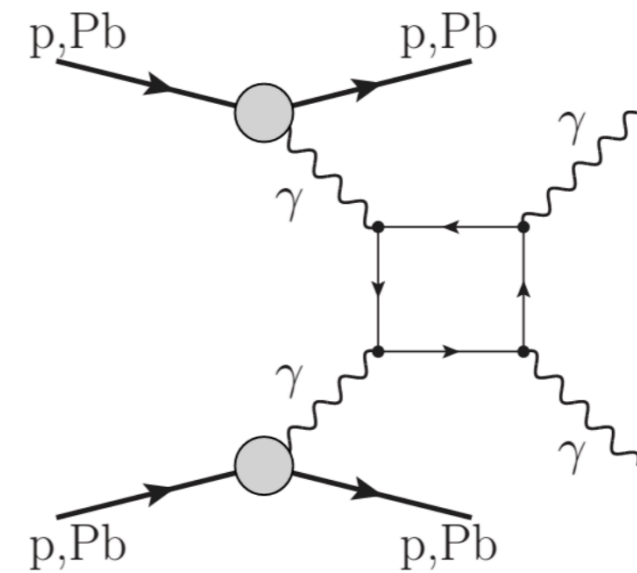


Axion Portal

- Visible bounds dominated by light-light scatter



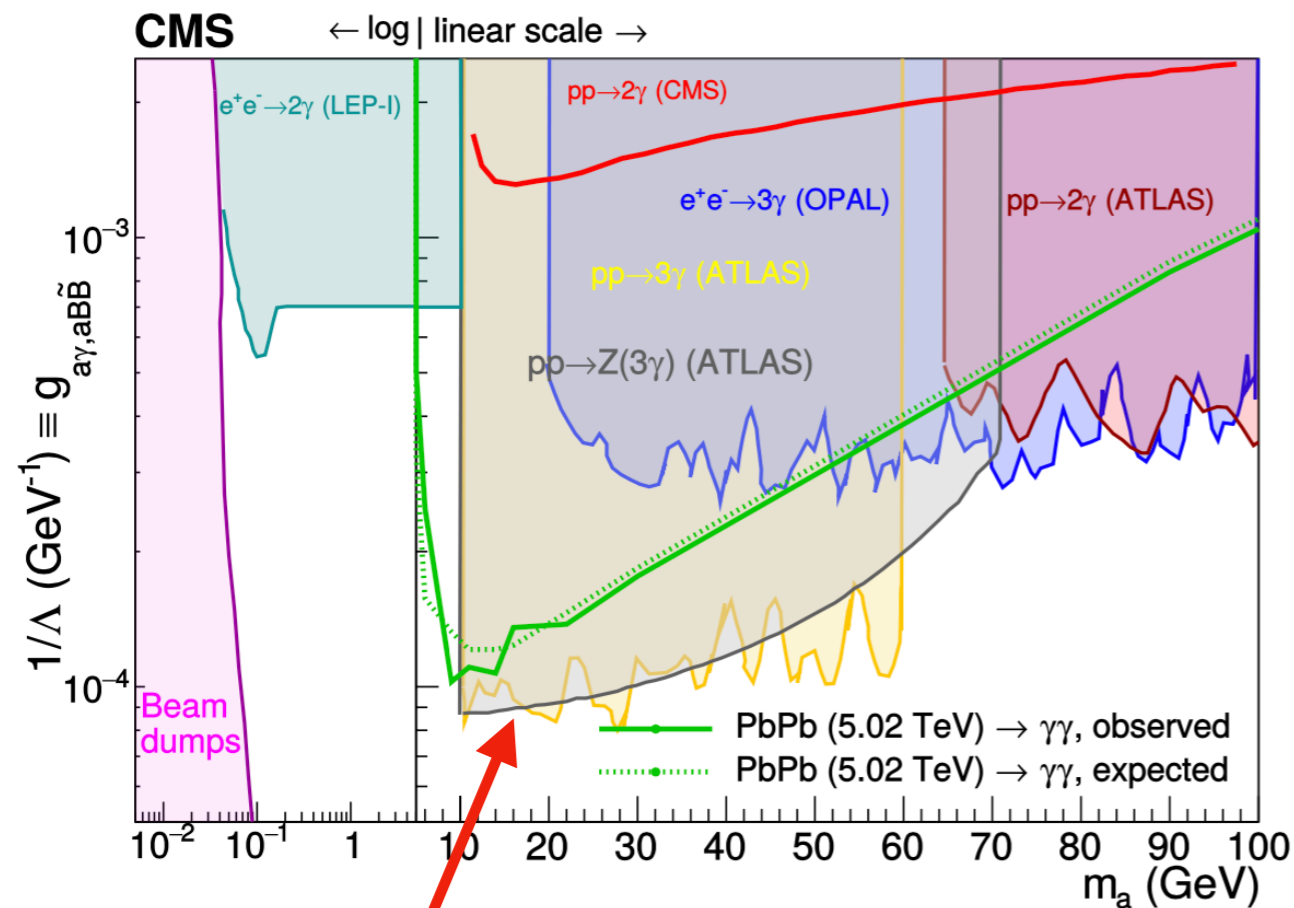
Heavy Ions very sensitive
Z⁴ enhancement



Gluon Coupling is much harder
LHCb has proposed searches in $\eta\pi\pi\pi$

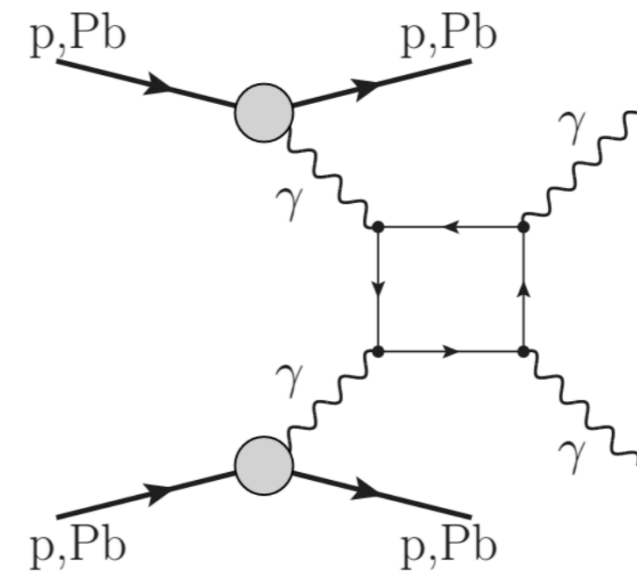
Axion Portal

- Visible bounds dominated by light-light scatter



Extending Model to have a Z coupling

Heavy Ions very sensitive
Z⁴ enhancement

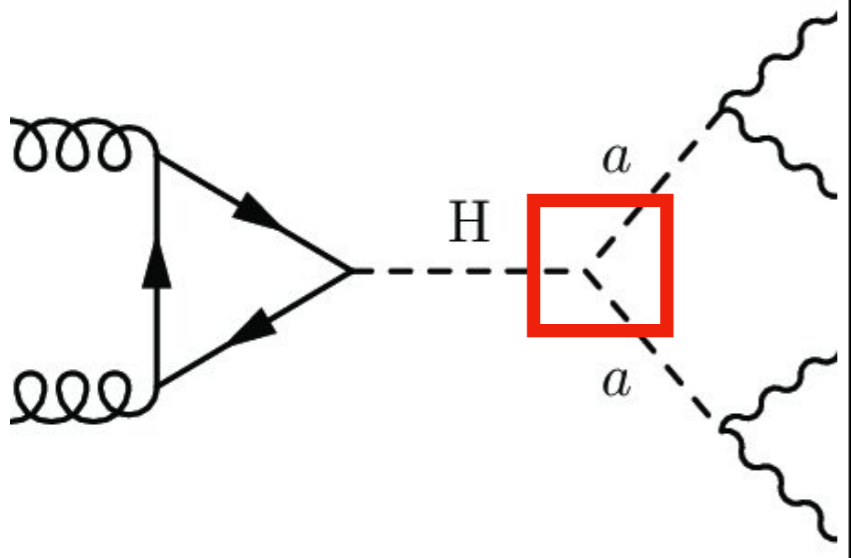


Adding Z decays (Fcc-ee) may be competitive
However a big heavy ion run may help much more

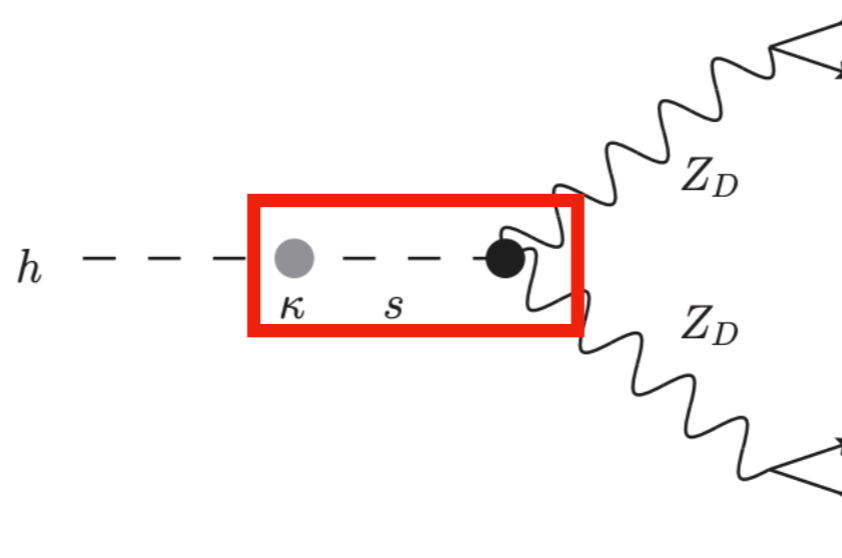
Unique Models

- Unlike other experiments, LHC has access to the Higgs
 - The Higgs enables the possibility of new final states
- The LHC also has **more Z bosons than anywhere on earth**
 - Z boson decays gives rise to additional final states

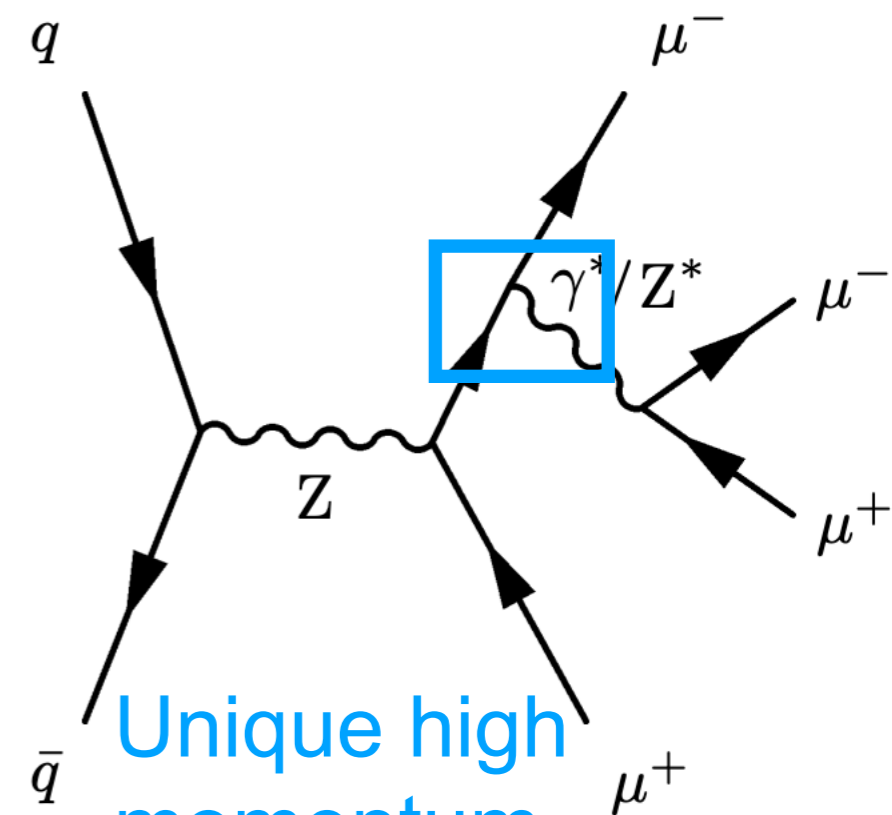
ALPs



Dark Photons



mu-Philic Photons



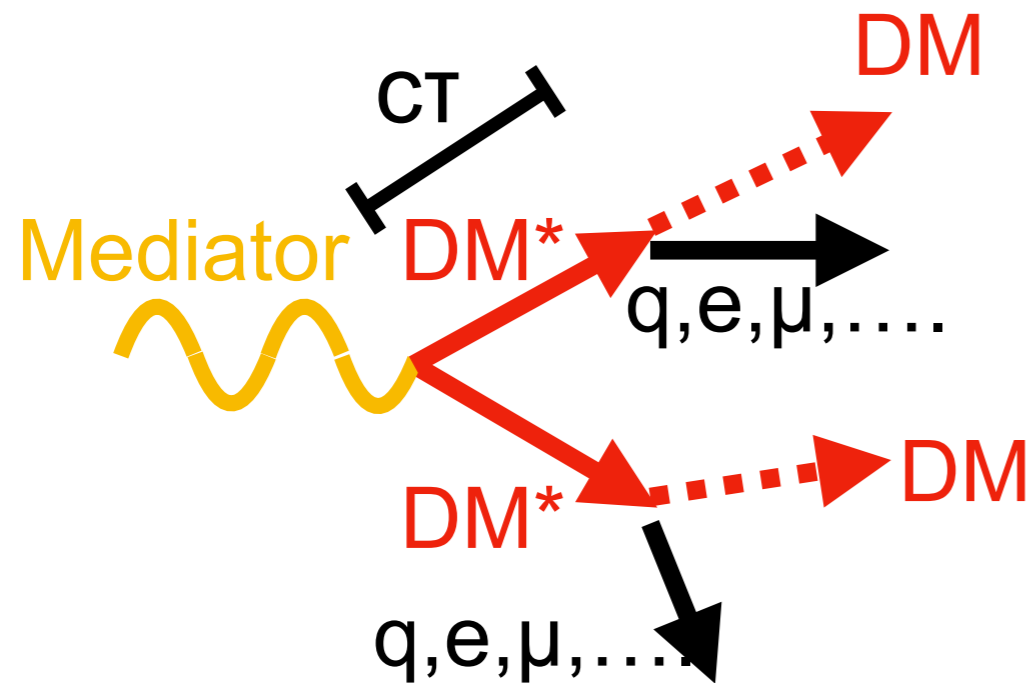
Requires a Higgs Coupling to the Dark Sector
 Large additional coupling can make LHC sensitive

Unique high momentum Muon Beam

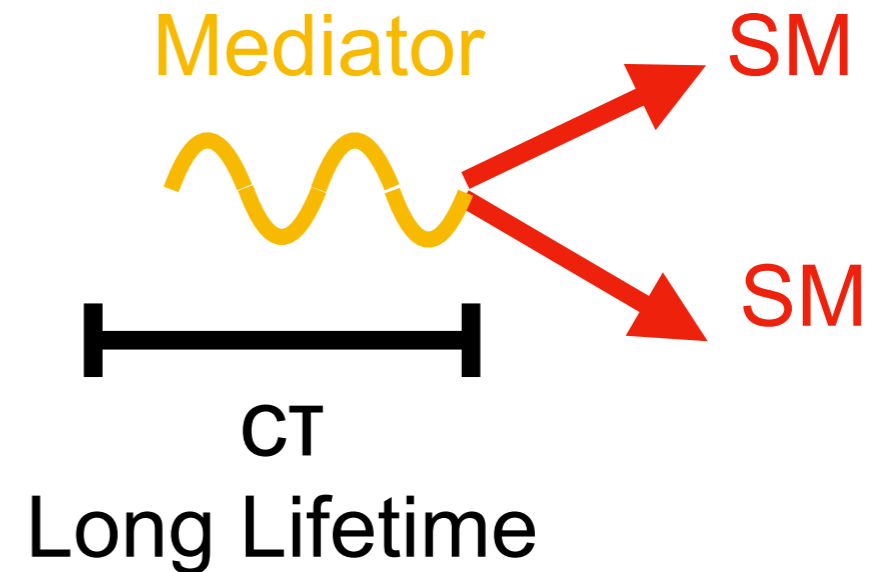
Luminosity @ LHC

- There will be an enormous amount of data at the LHC
 - Strong bounds provided we can tap this dataset
- Higgs or Z boson couplings to Dark Sector yield enhancements
 - Provided we have the right model
 - ▶ Higgs and Z boson couplings are needed for this

Long-Lived Searches



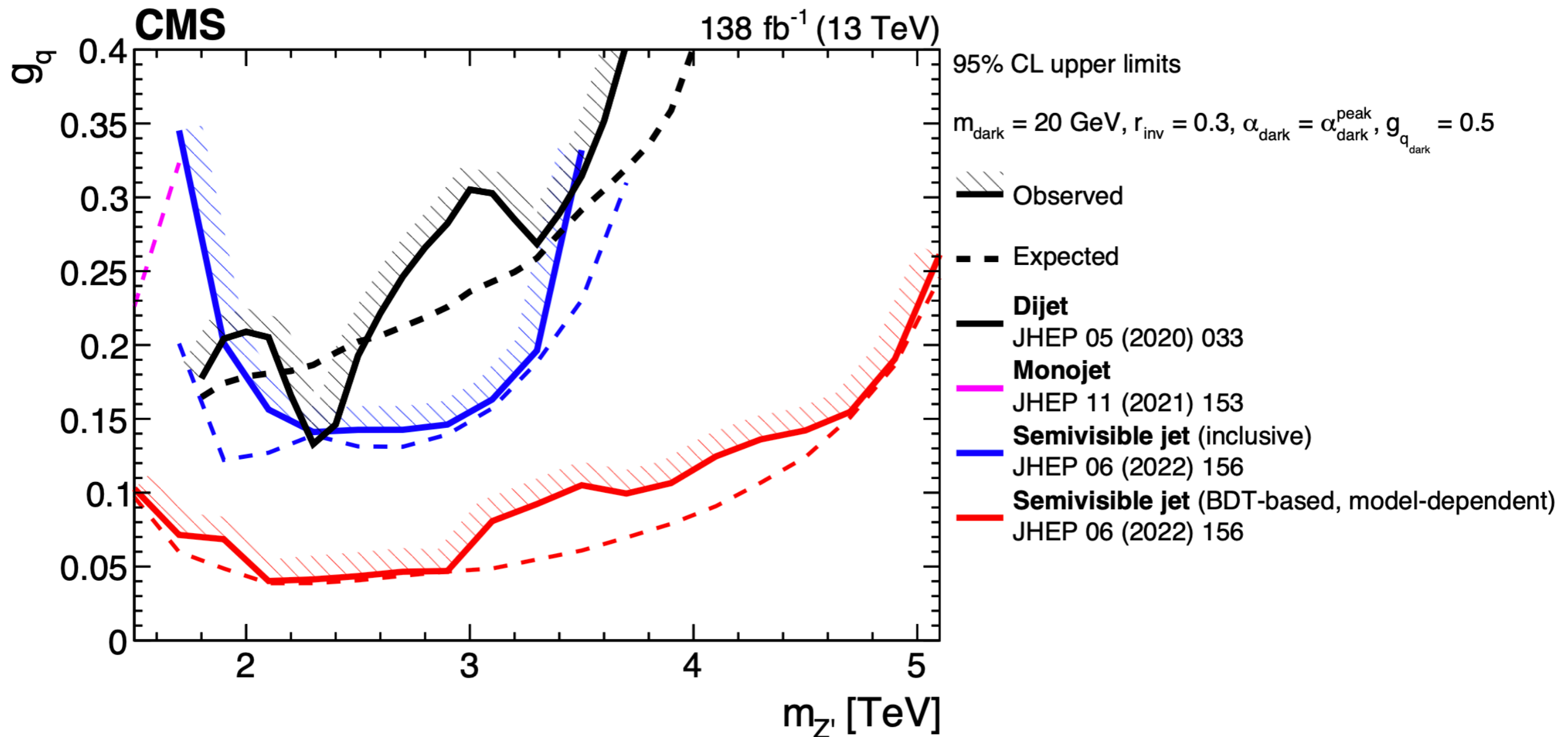
Dark Matter is Unstable



Mediator is Weakly Coupled

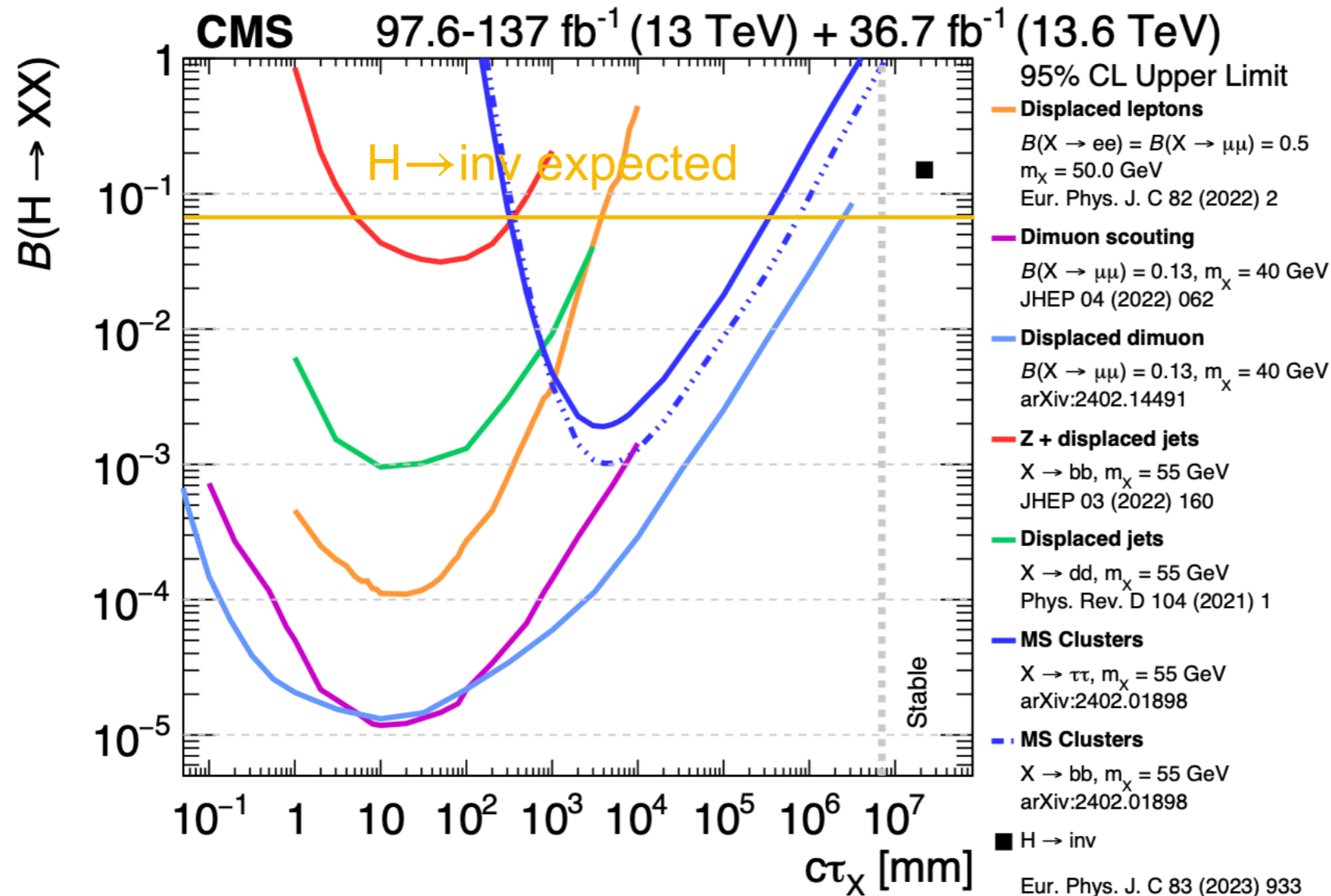
Semi-visible decays help

- Adding complexity to the model enhances overall sensitivity



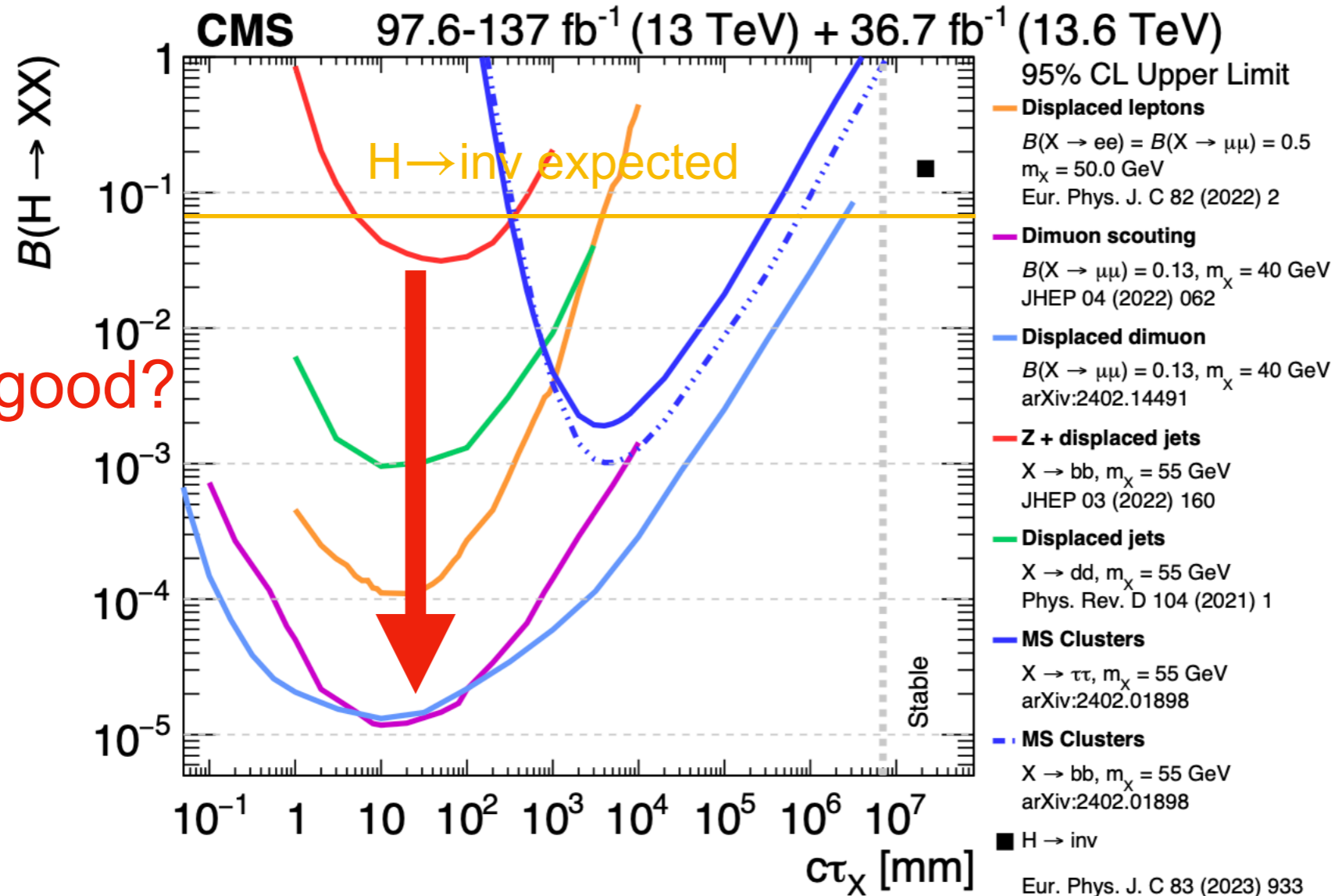
Long-lived decays help

- Adding complexity to the model enhances overall sensitivity



Long-lived decays help

- Adding complexity to the model enhances overall sensitivity

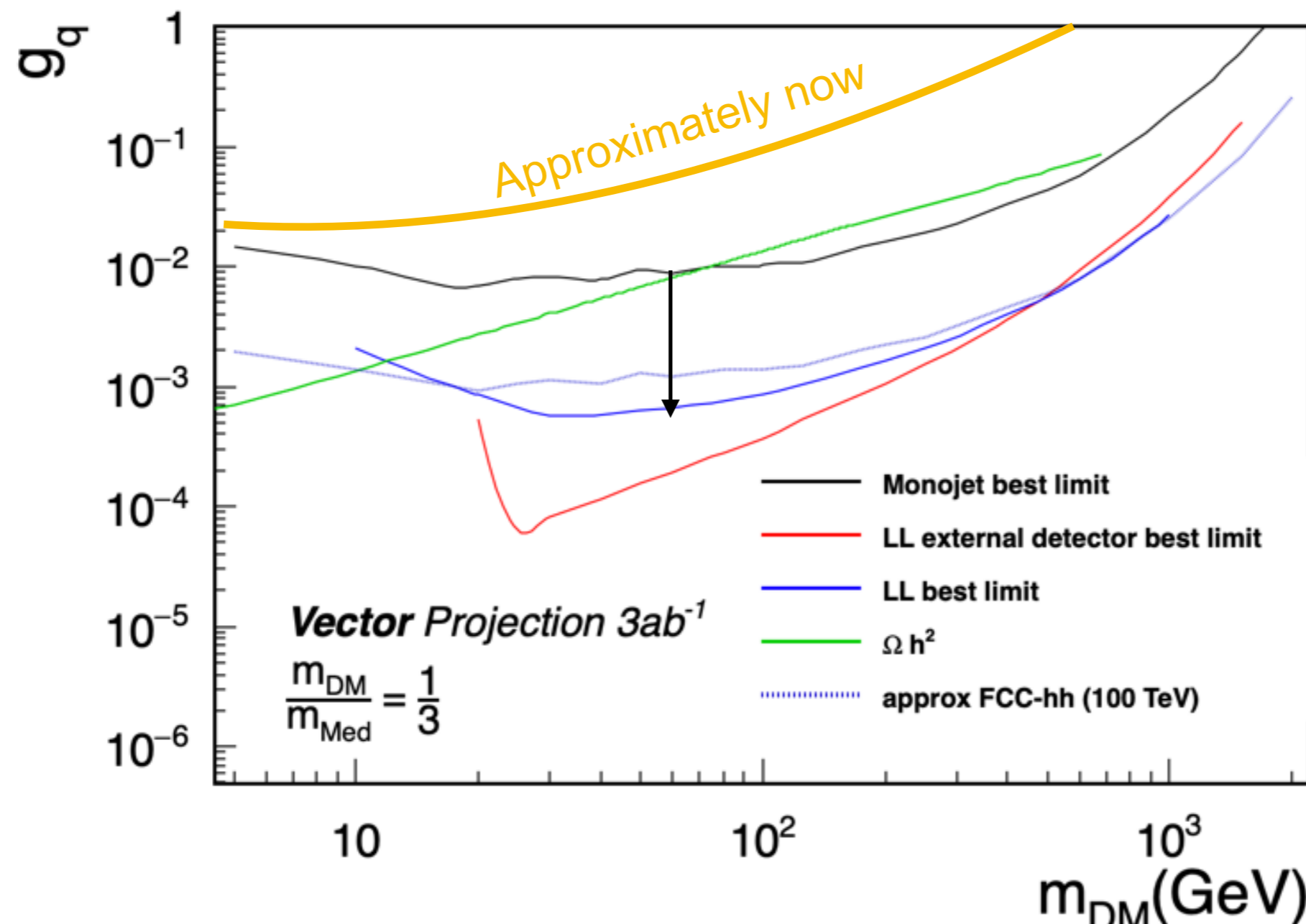


Why so good?

Many new ideas
For triggering
Are helping a lot

Projecting To the Future

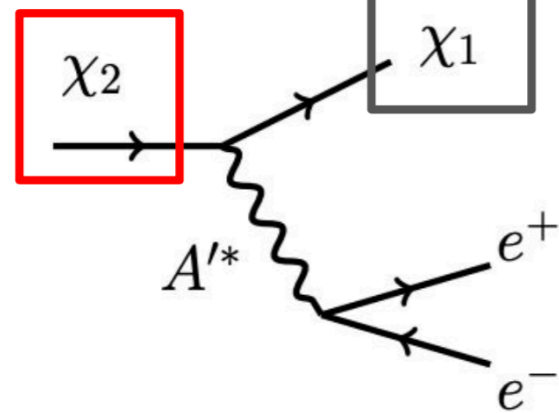
With a naive recast of the monojet analysis for displaced vertices



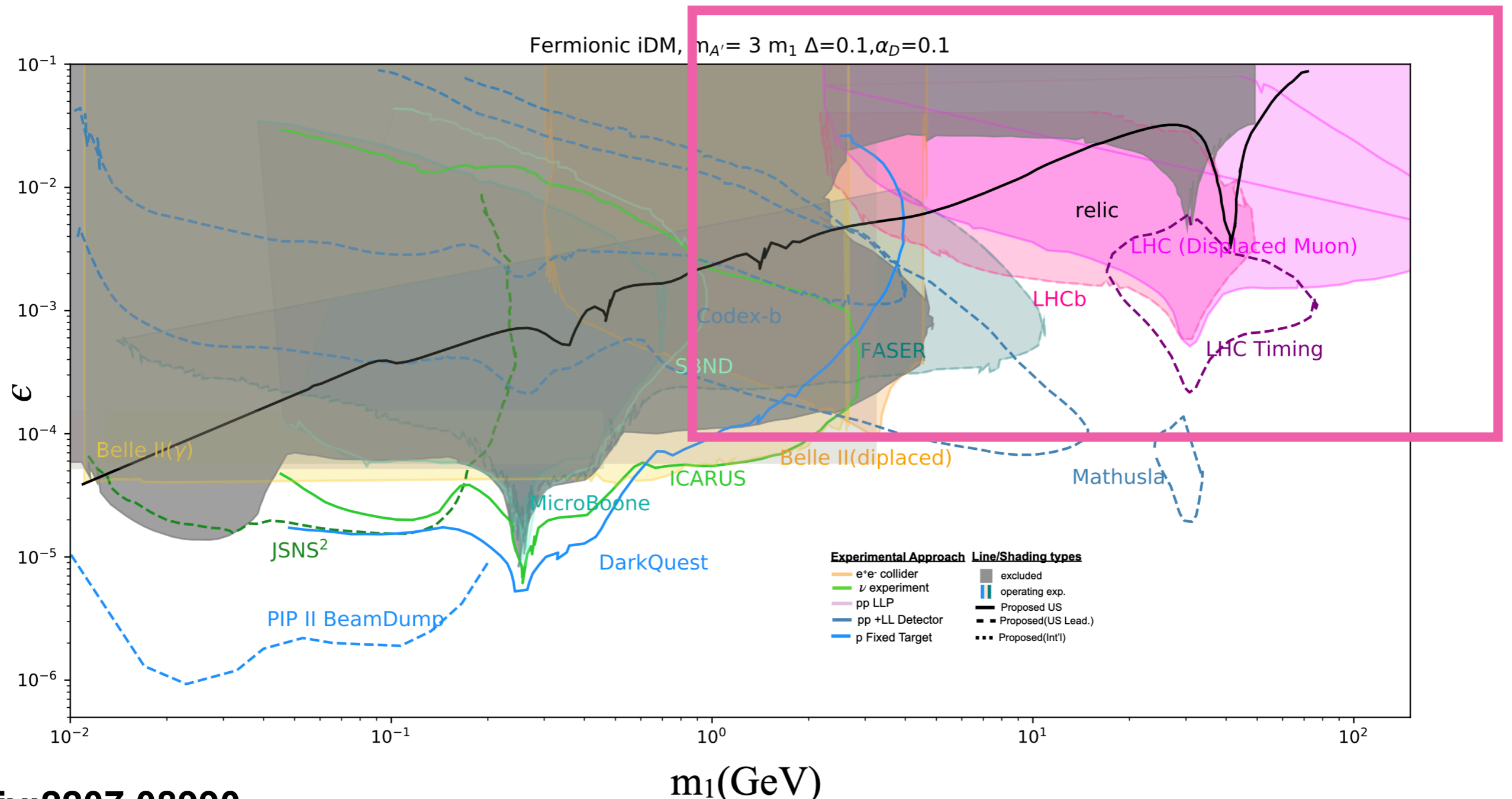
Can get another order of magnitude in coupling pretty generically

Many Others

Additional
unstable DM

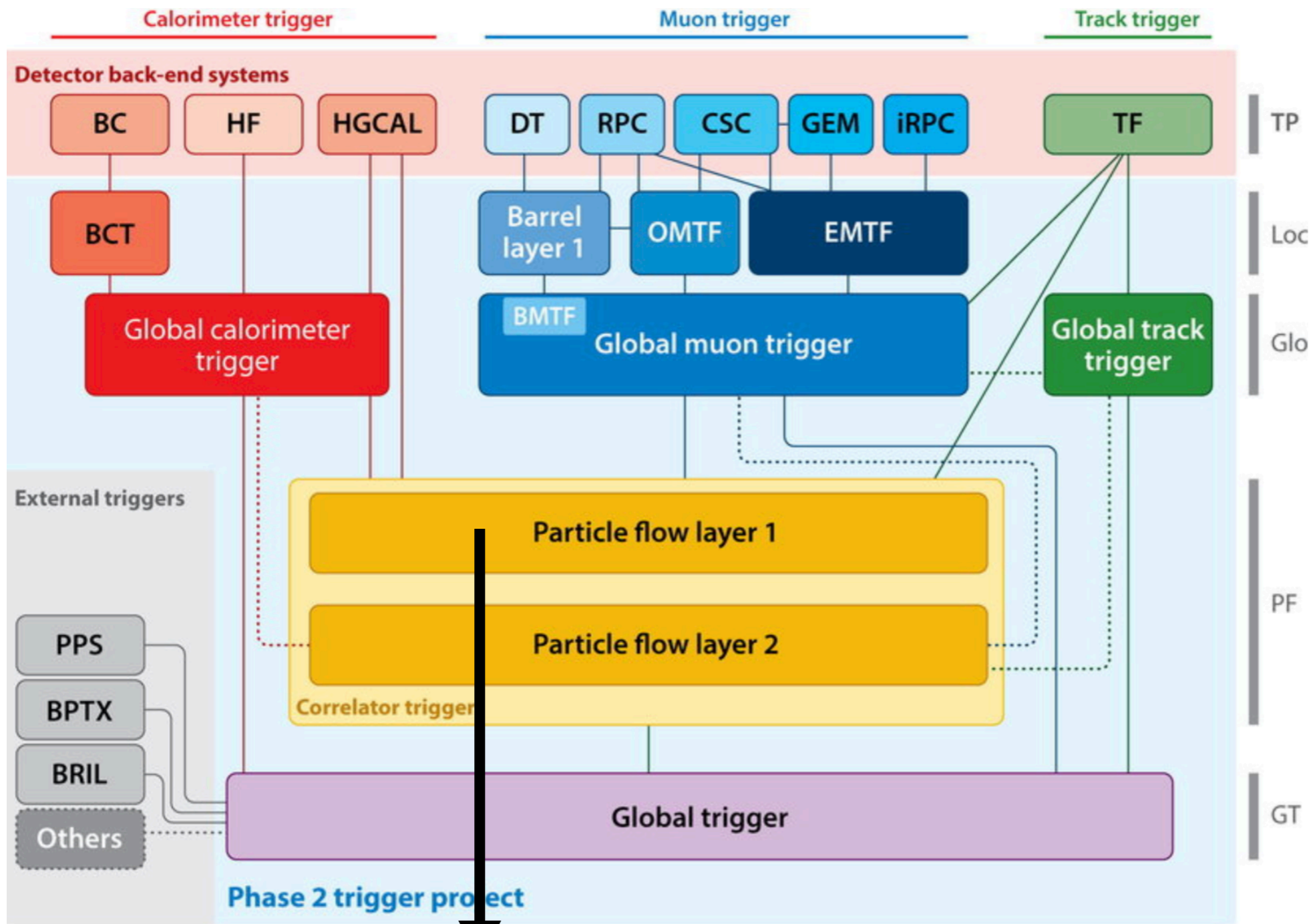


By Making the Dark Matter Unstable
Invisible searches become Long lived LHC



We are developing a system that hopefully can
 Do any analysis on all events
 It can't do every analysis on all events

HL-LHC



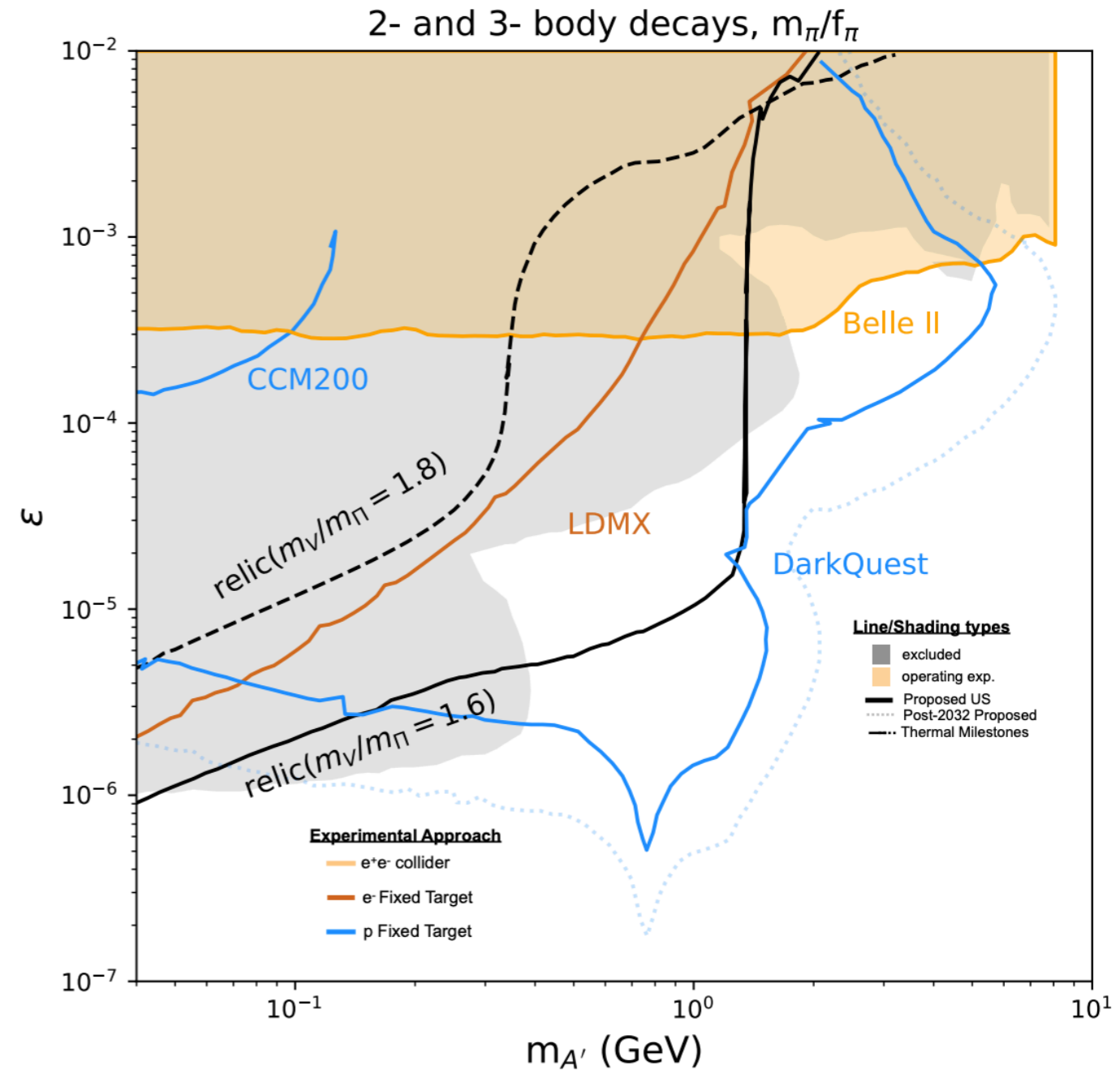
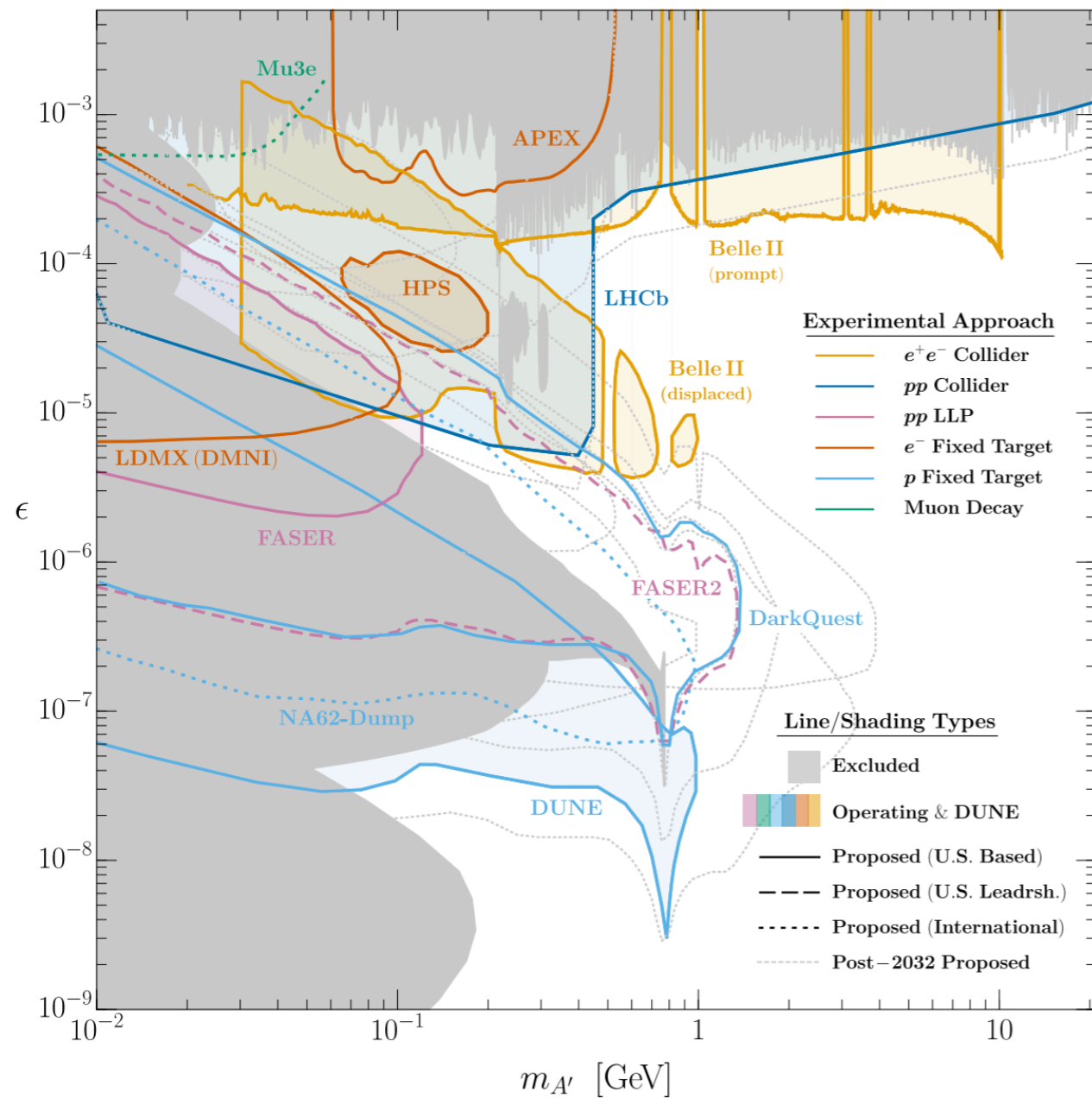
40 MHz Scouting system

This system will start form L1 Puppi candidates but will possibly add more

Conclusions

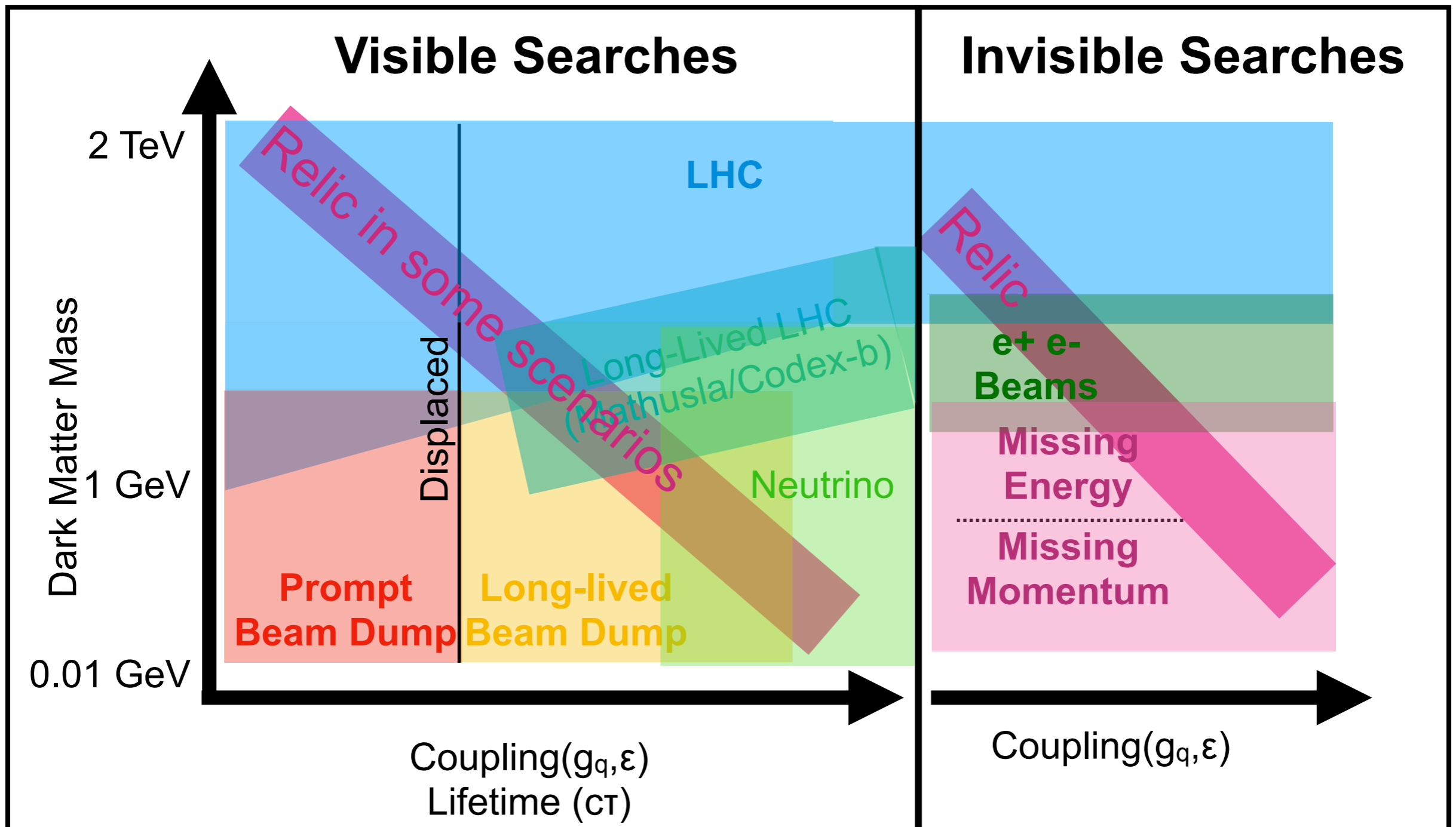
- Invisible searches
 - γ_D benchmark : will cover $m < 3$ GeV and $m > 45$ GeV
 - It is **not clear that we can cover the middle region**
- Visible searches
 - LHC will cover broad range of prompt di-lepton signatures
 - ▶ Will be the **leading constraint for γ_D benchmark at large ϵ**
 - Higgs and Z decays can enhance things
- Long-lived searches : further enhancements
- Future colliders:
 - My view: high mass is not the most exciting region
 - FCC-hh will help with lumi, not clear others will help
 - **A high COM beam-dump would be something to think about**

DarkQuest

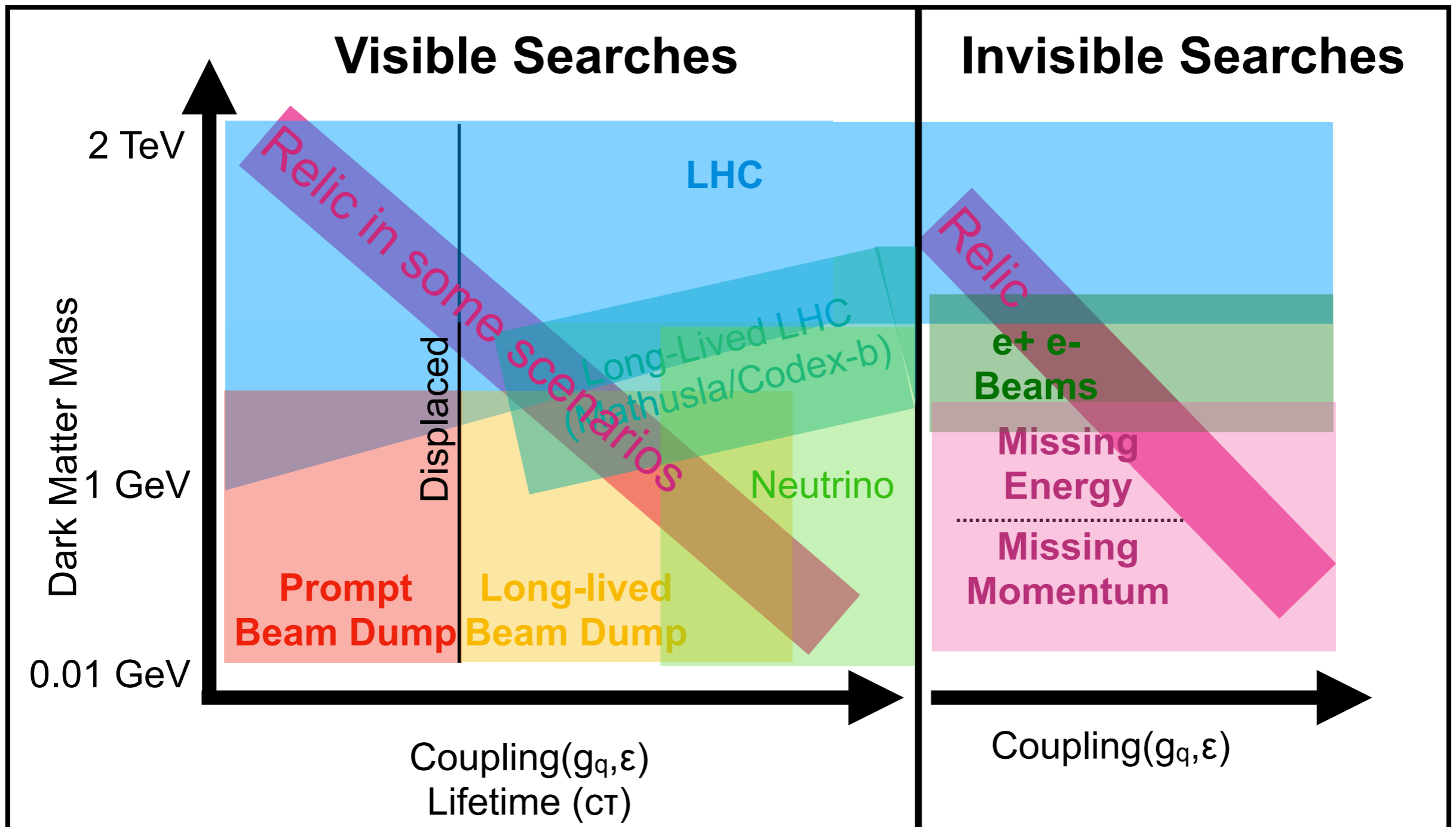


DarkQuest is a funded experiment
Come and join if you are interested!

Thanks!

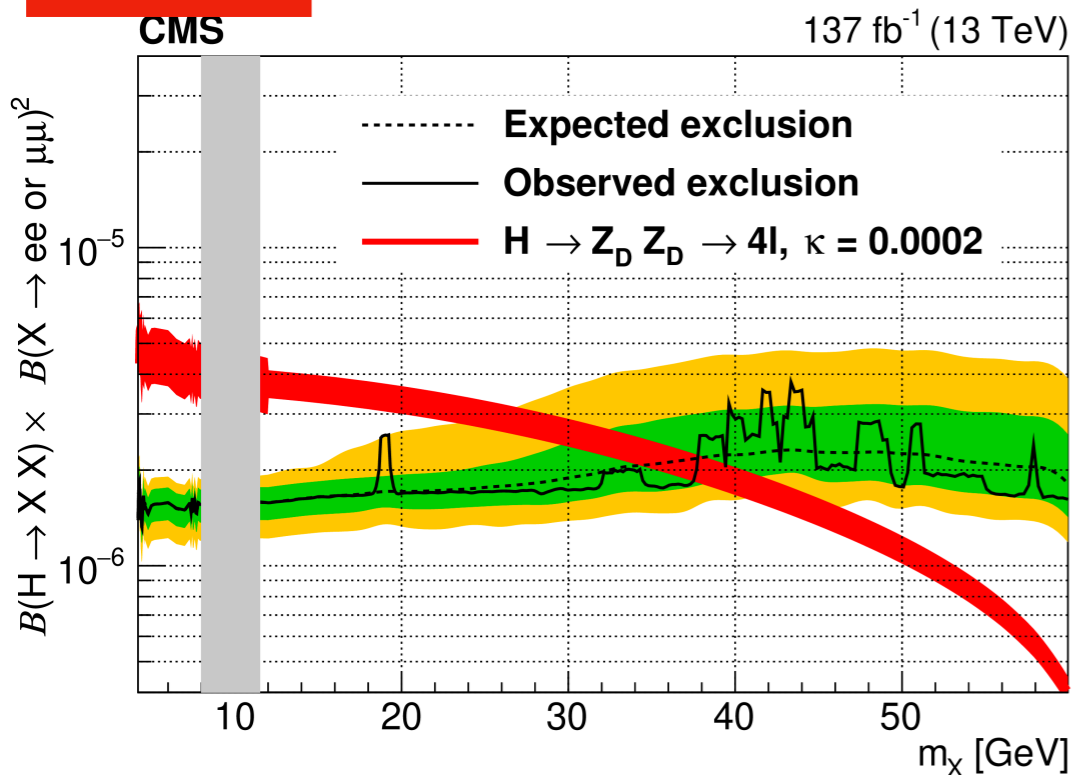
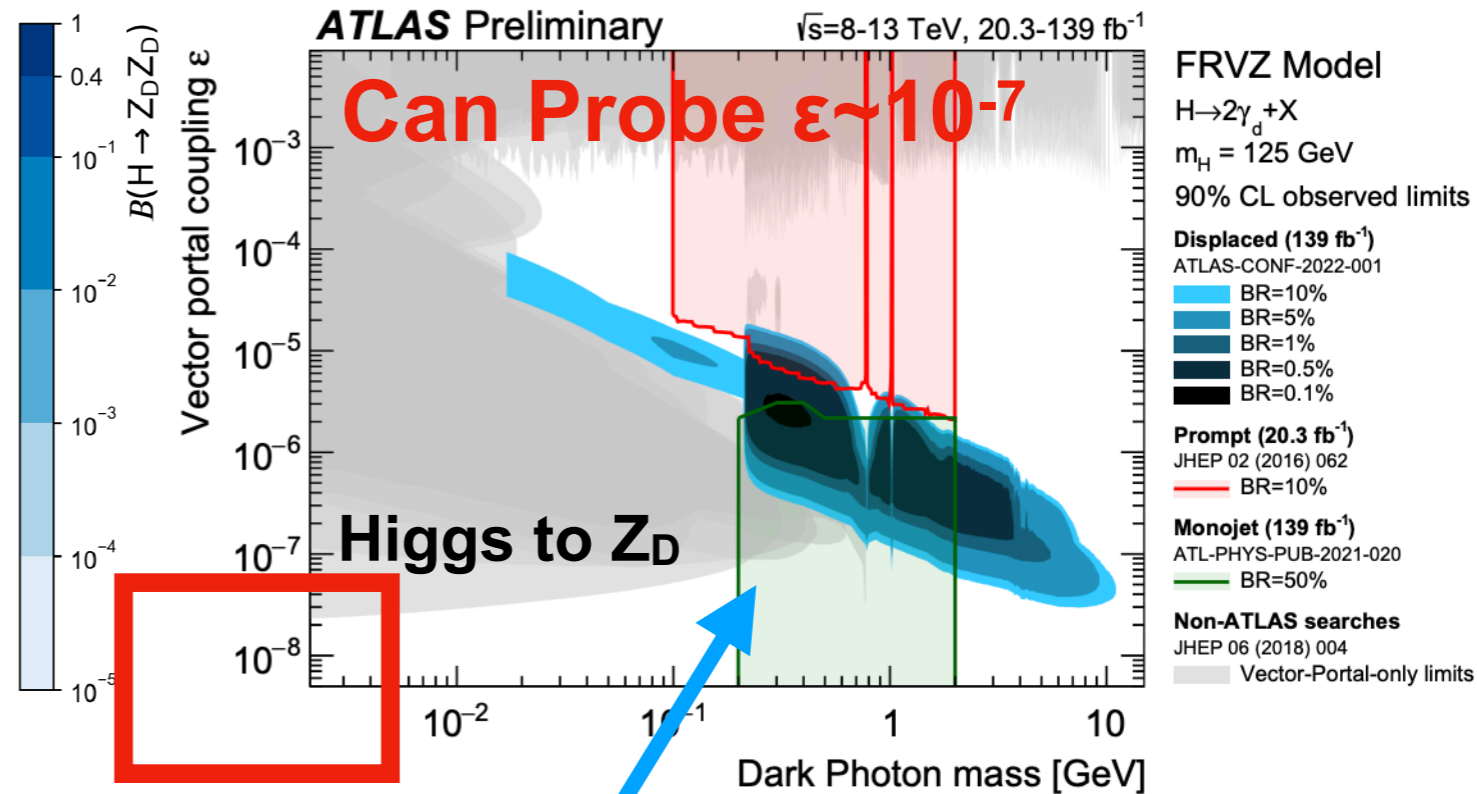
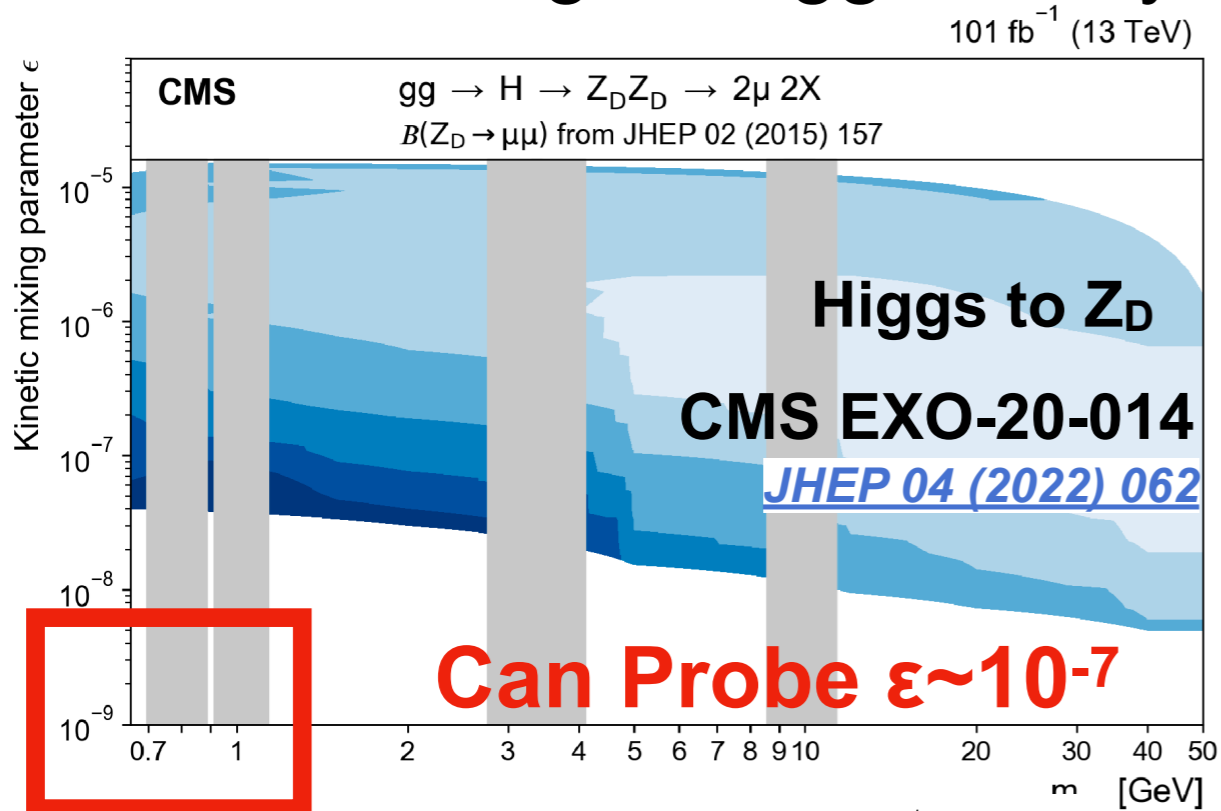


Visualizing the landscape



What Results are there?

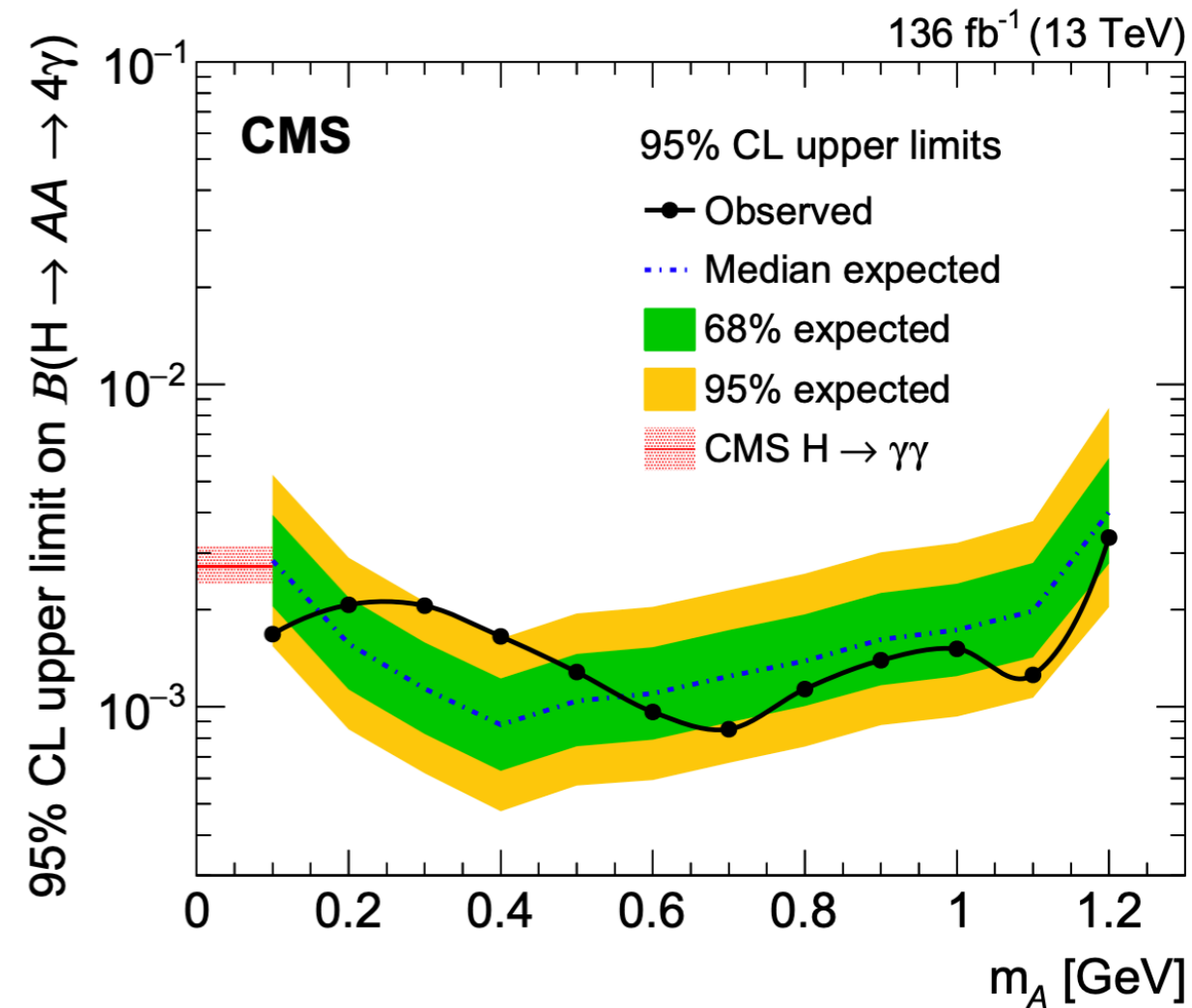
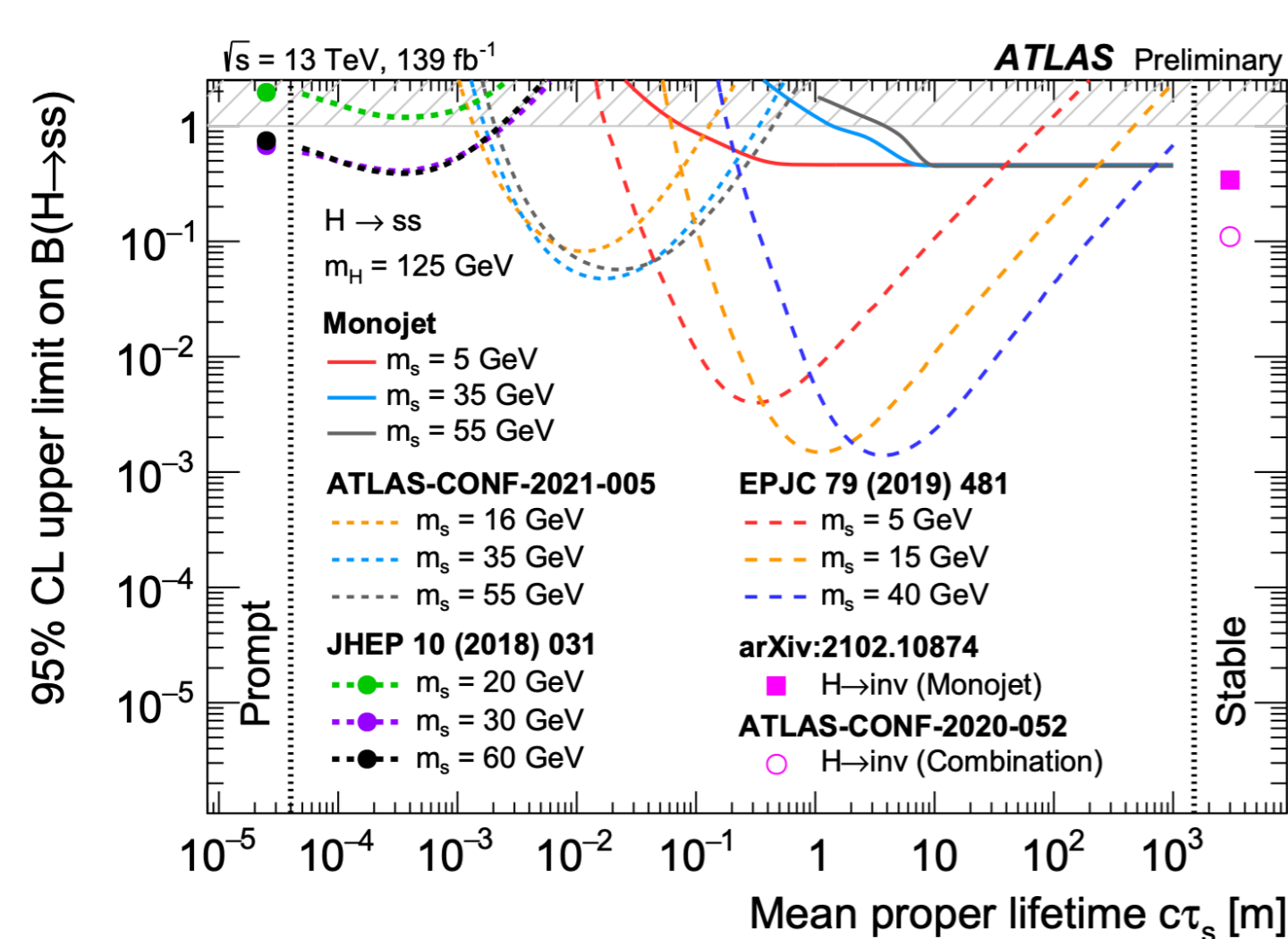
- Looking at Higgs decays to dark photons



Given right choice of Higgs Dark Coupling
 Can probe allowed range of sensitivity
 Long Lived particle signatures emerge
 Invisible and visible also contribute

Higgs to spin-0

- Looking at Higgs decays to ALPs and Scalars



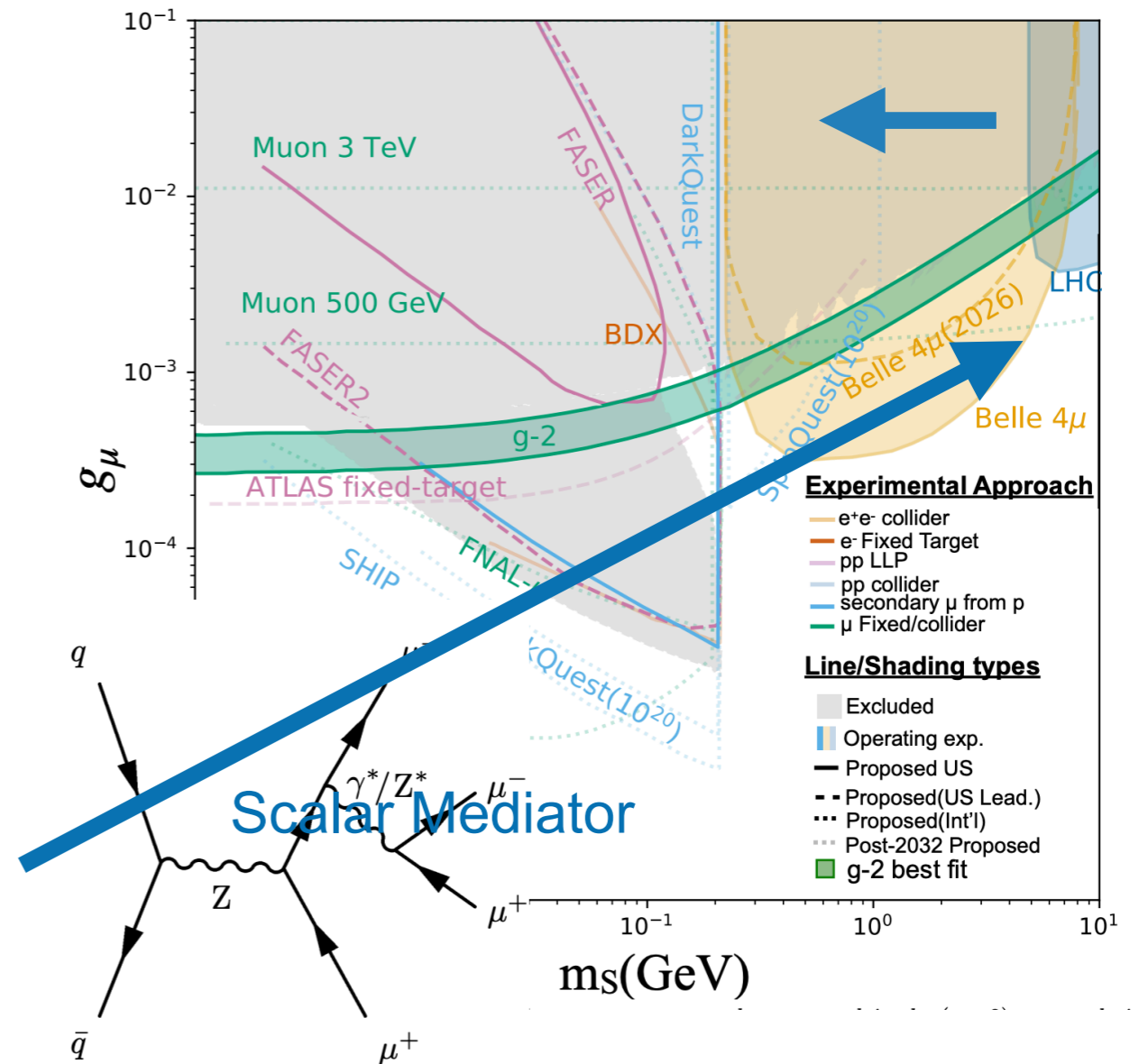
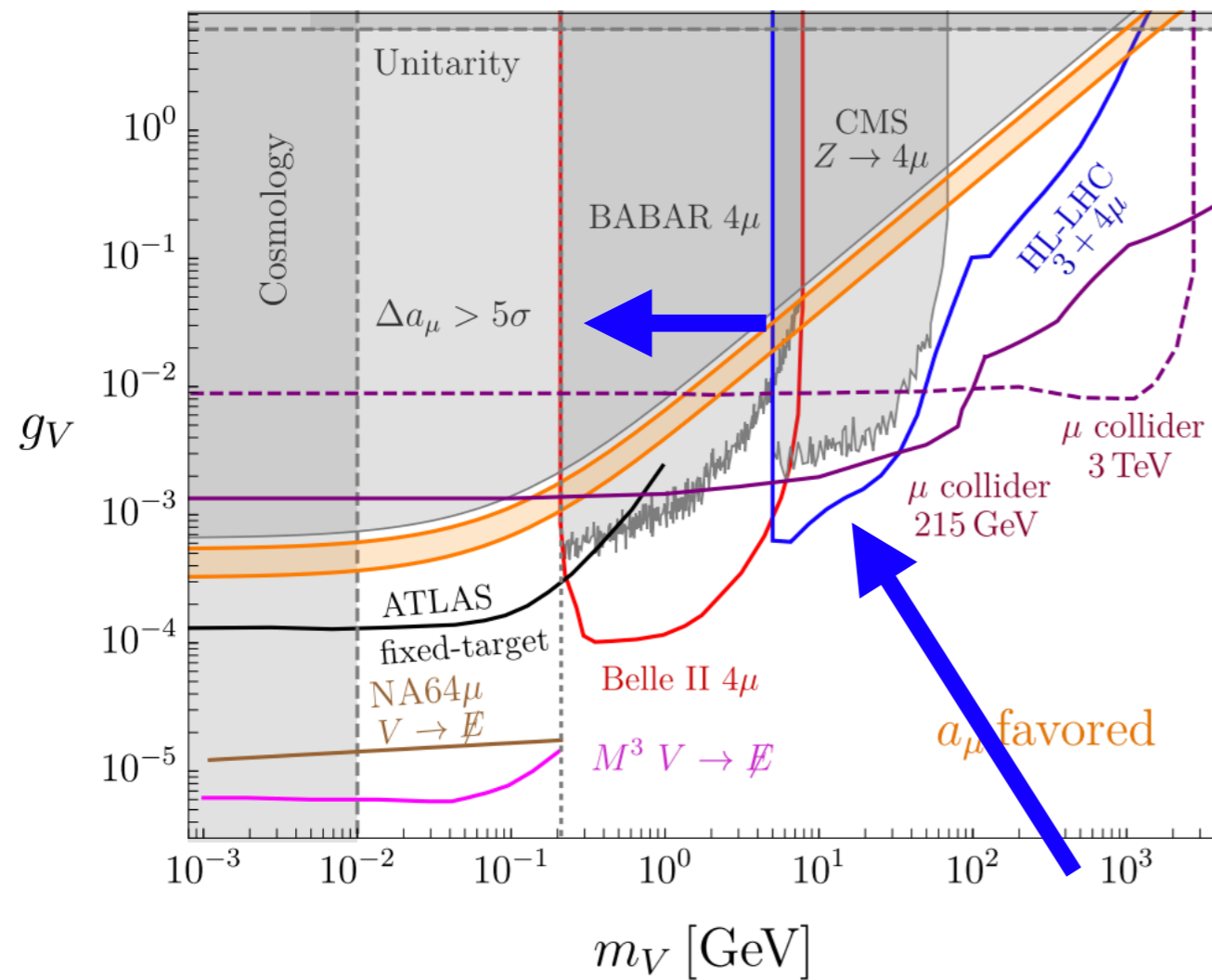
Long lifetimes often lead to very small couplings

Longer Lifetime: Higgs to invisible bounds dominate

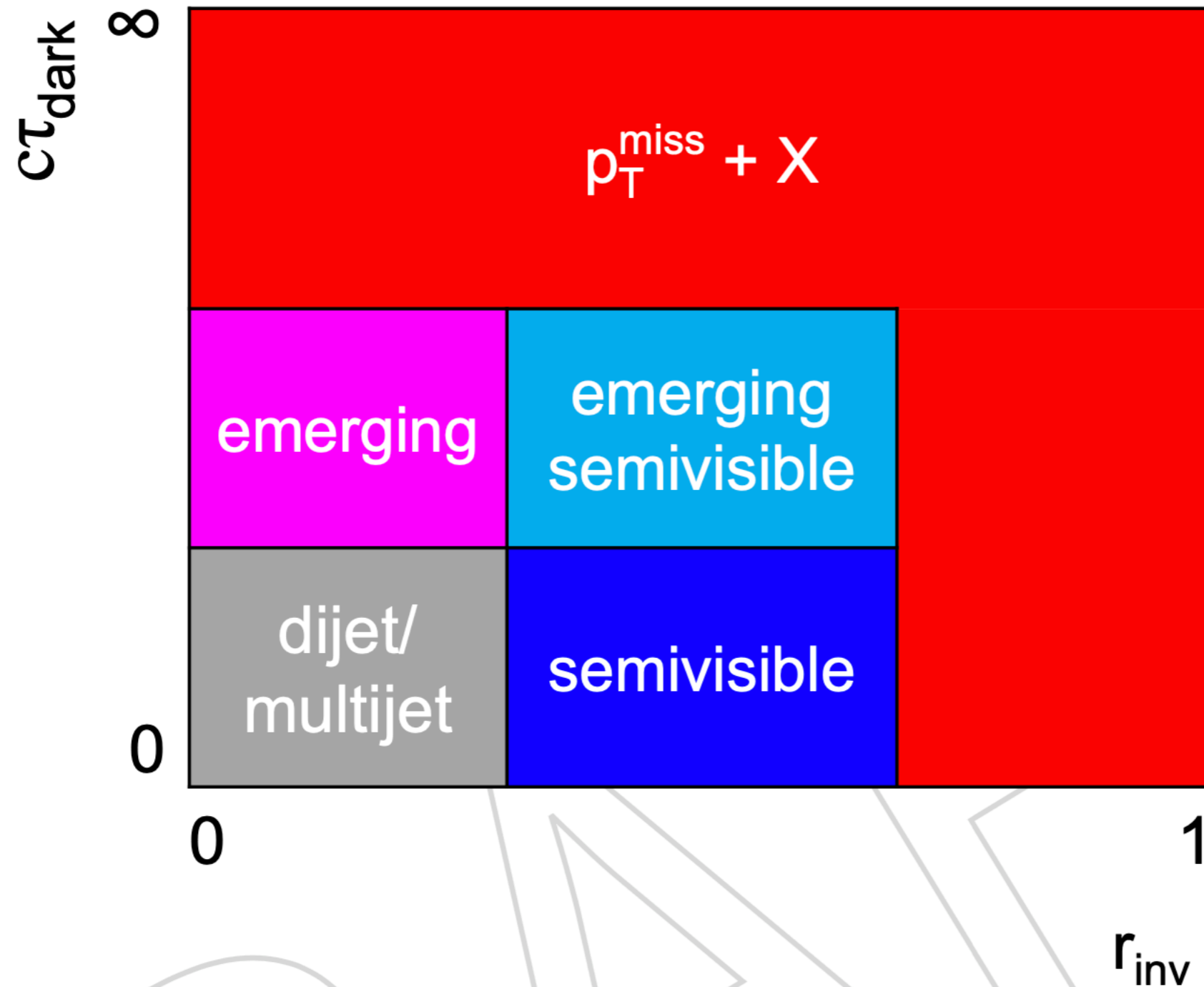
$g-2$ mediator (Just muons)

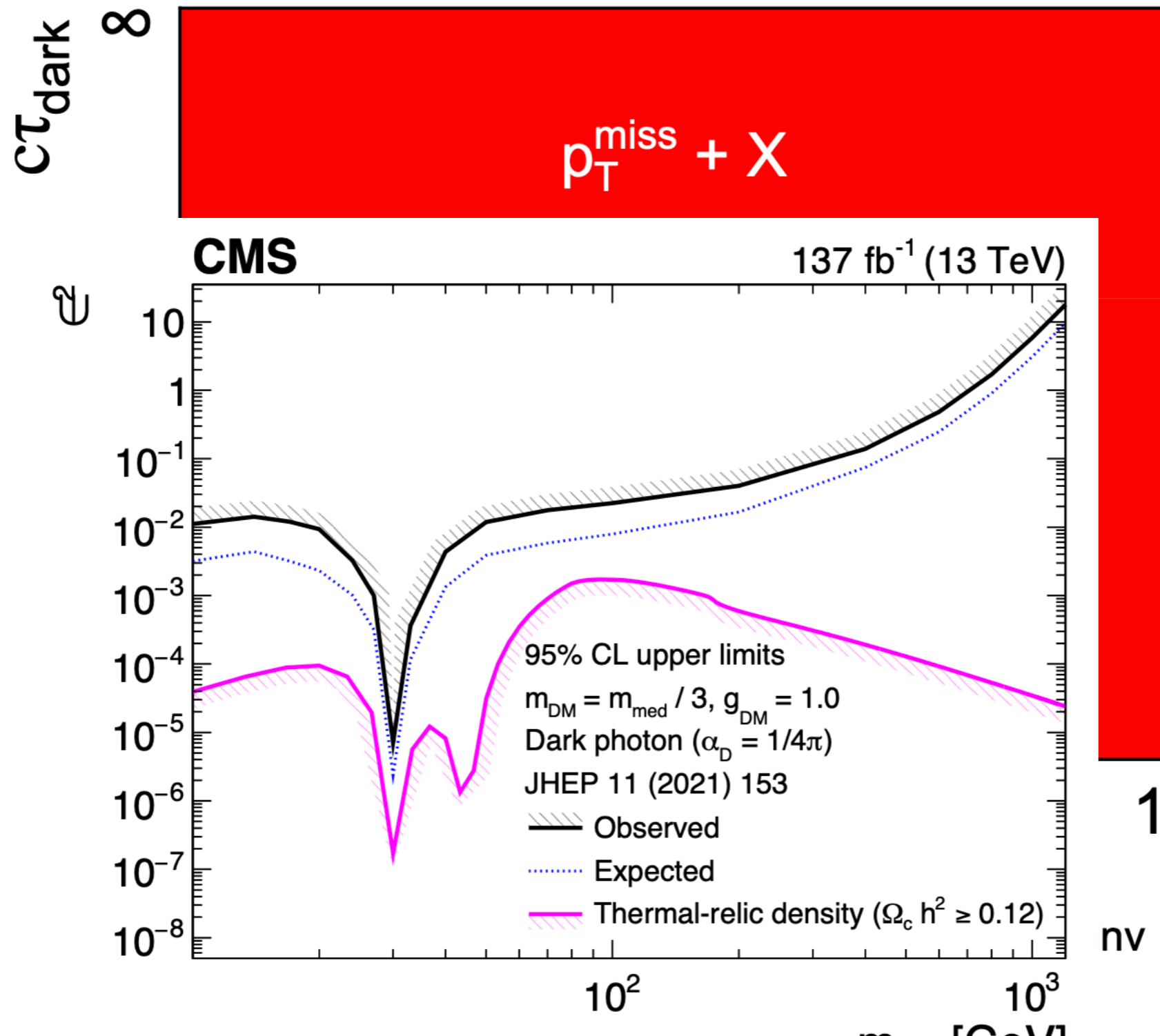
- Z to 4 muon channel provides unique high mass constraints

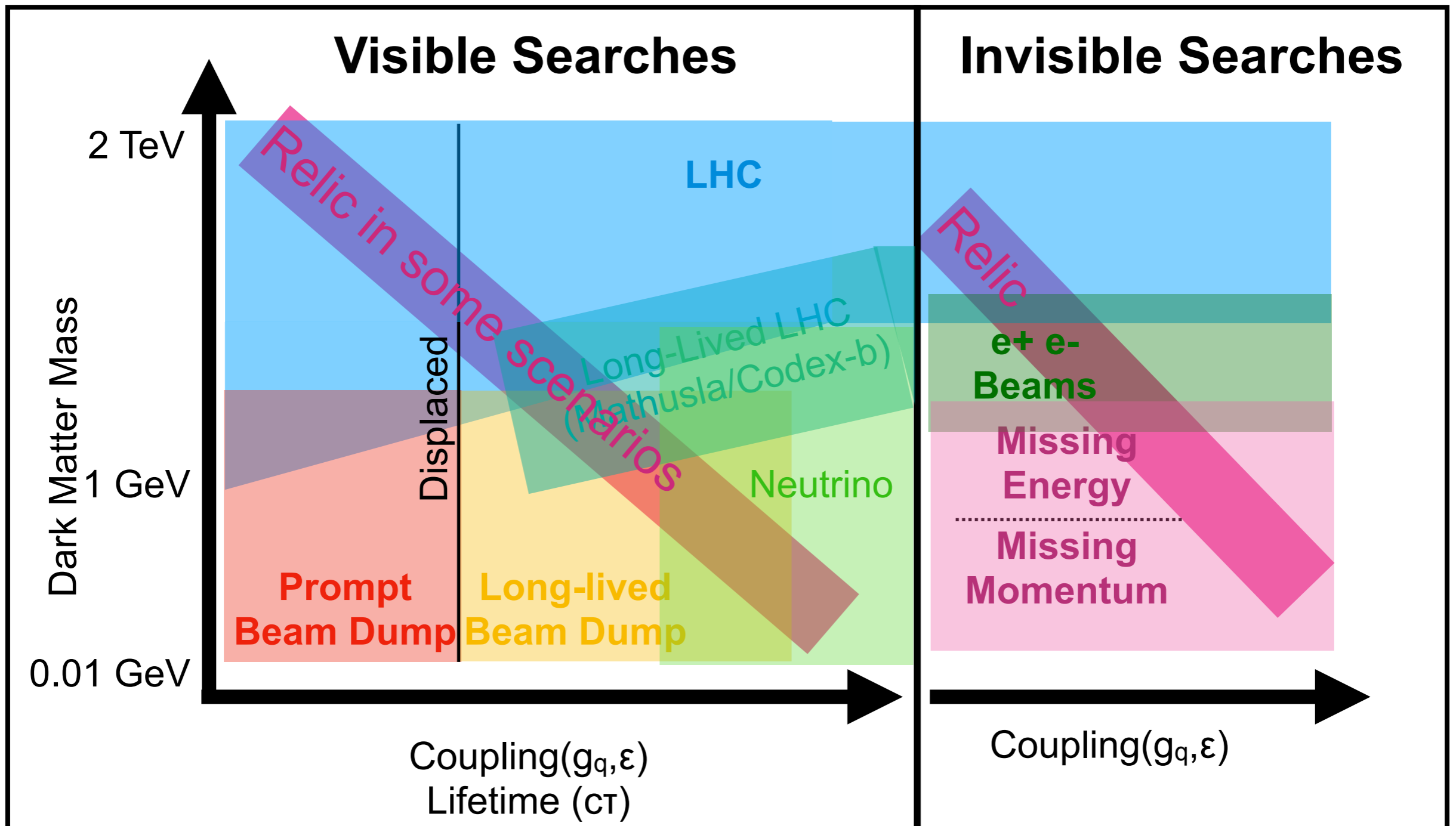
Vector, $\text{BR}(V \rightarrow \mu^+ \mu^-) = 1$ for $m_V > 2m_\mu$



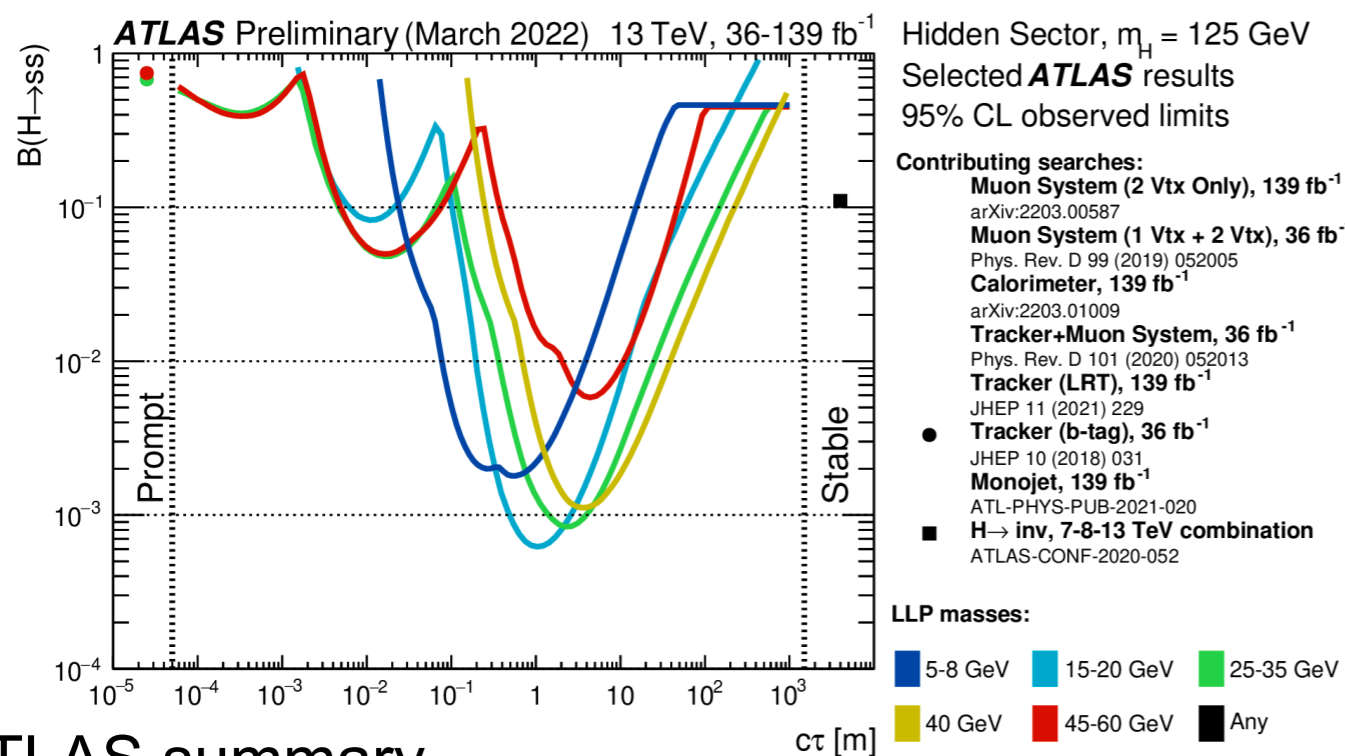
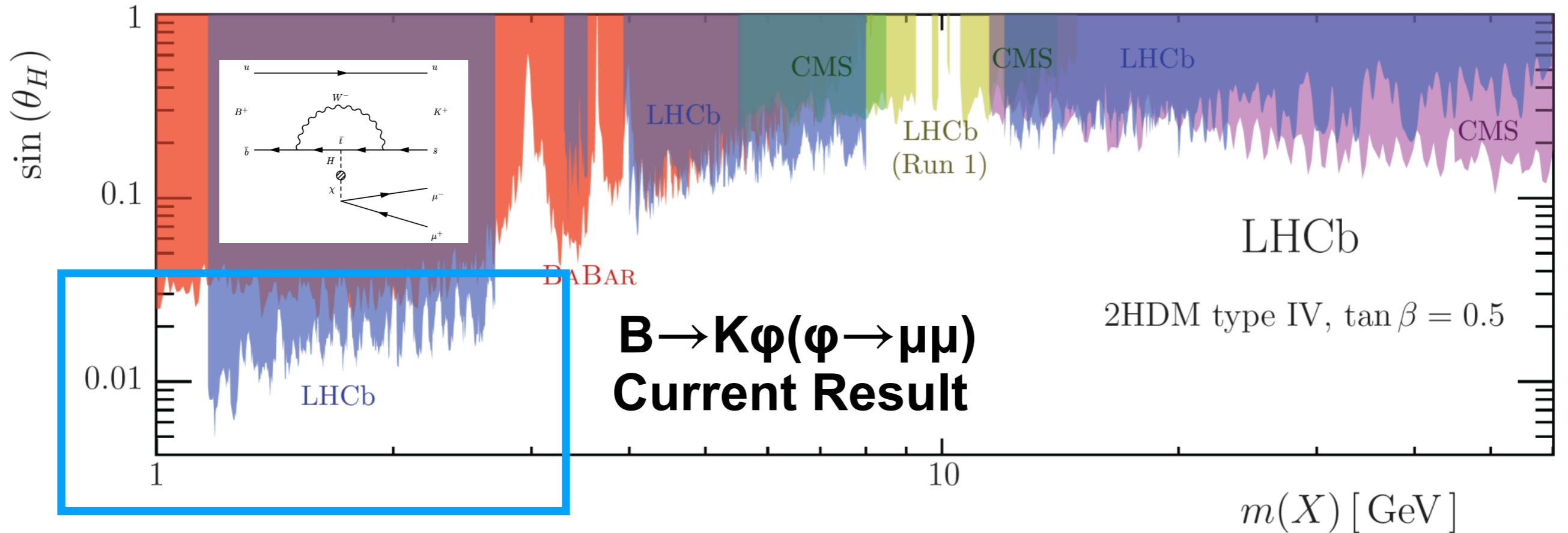
Visualizing the landscape







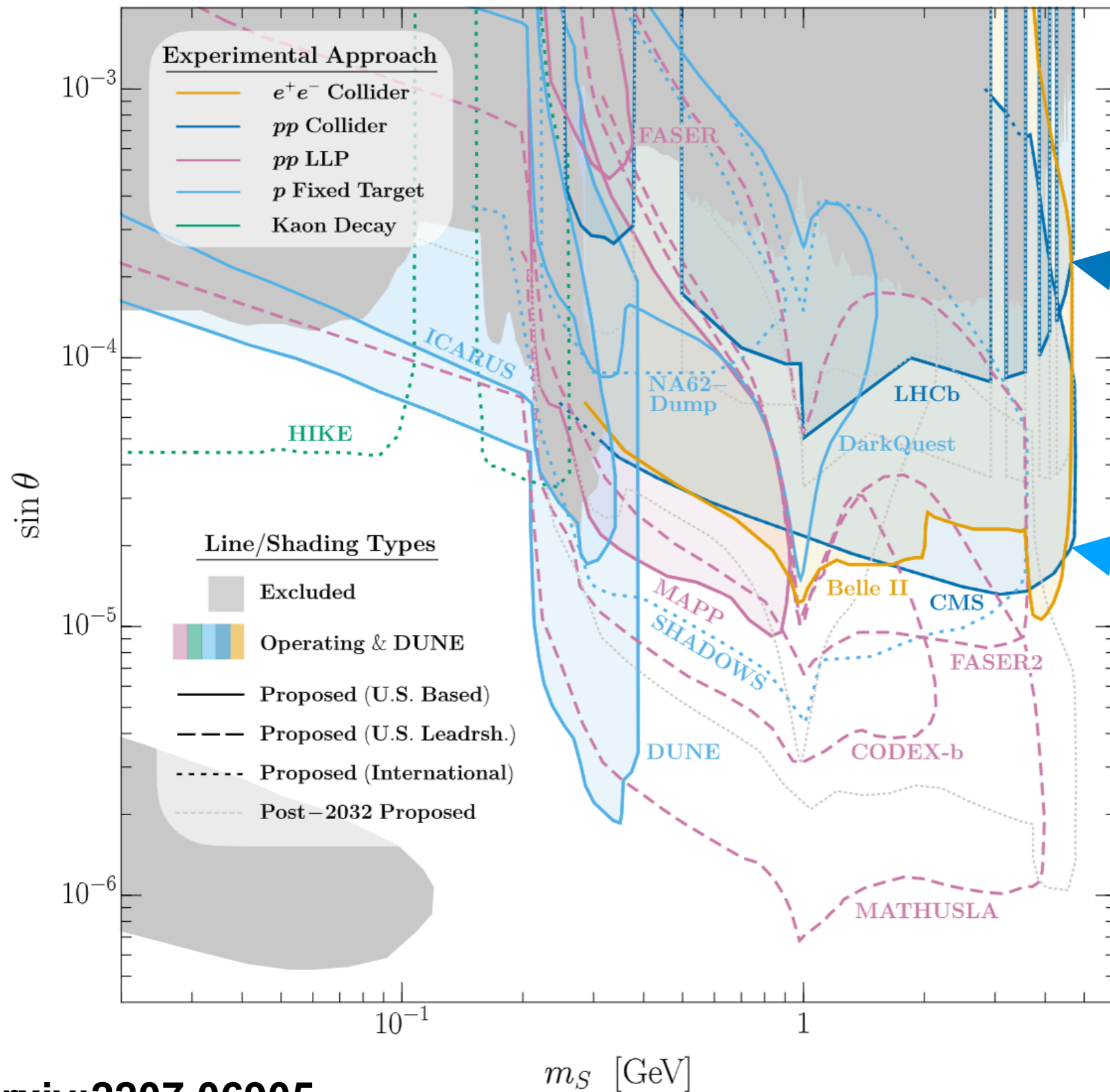
Scalar Portal



Once we get into the small coupling regime

Long-Lived searches start to dominate

Scalar Portal



$B \rightarrow X\phi (\phi \rightarrow \mu\mu)$

LHCb Future Projections

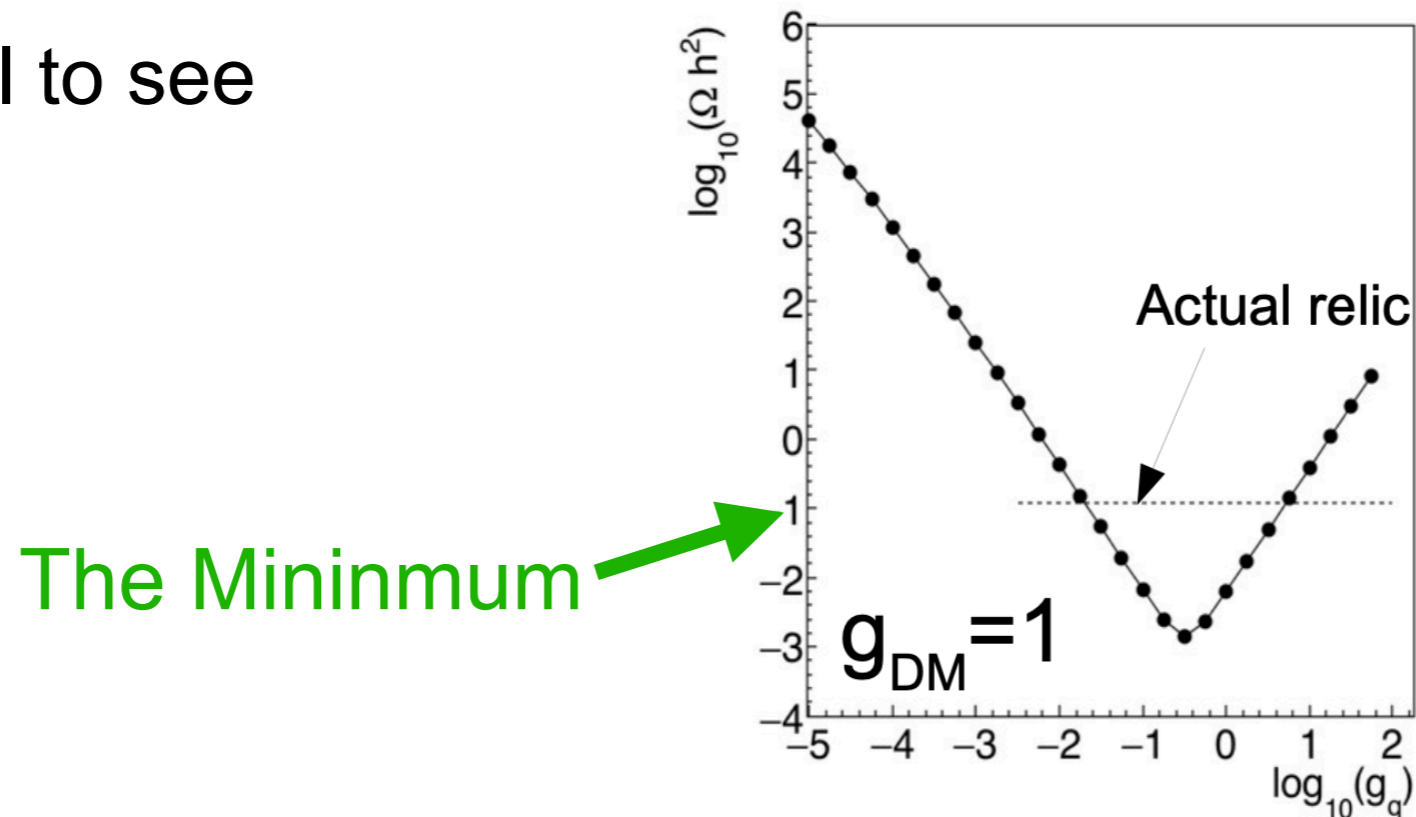
$B \rightarrow X\phi (\phi \rightarrow \mu\mu)$

Future CMS Long Lived
Track Trigger

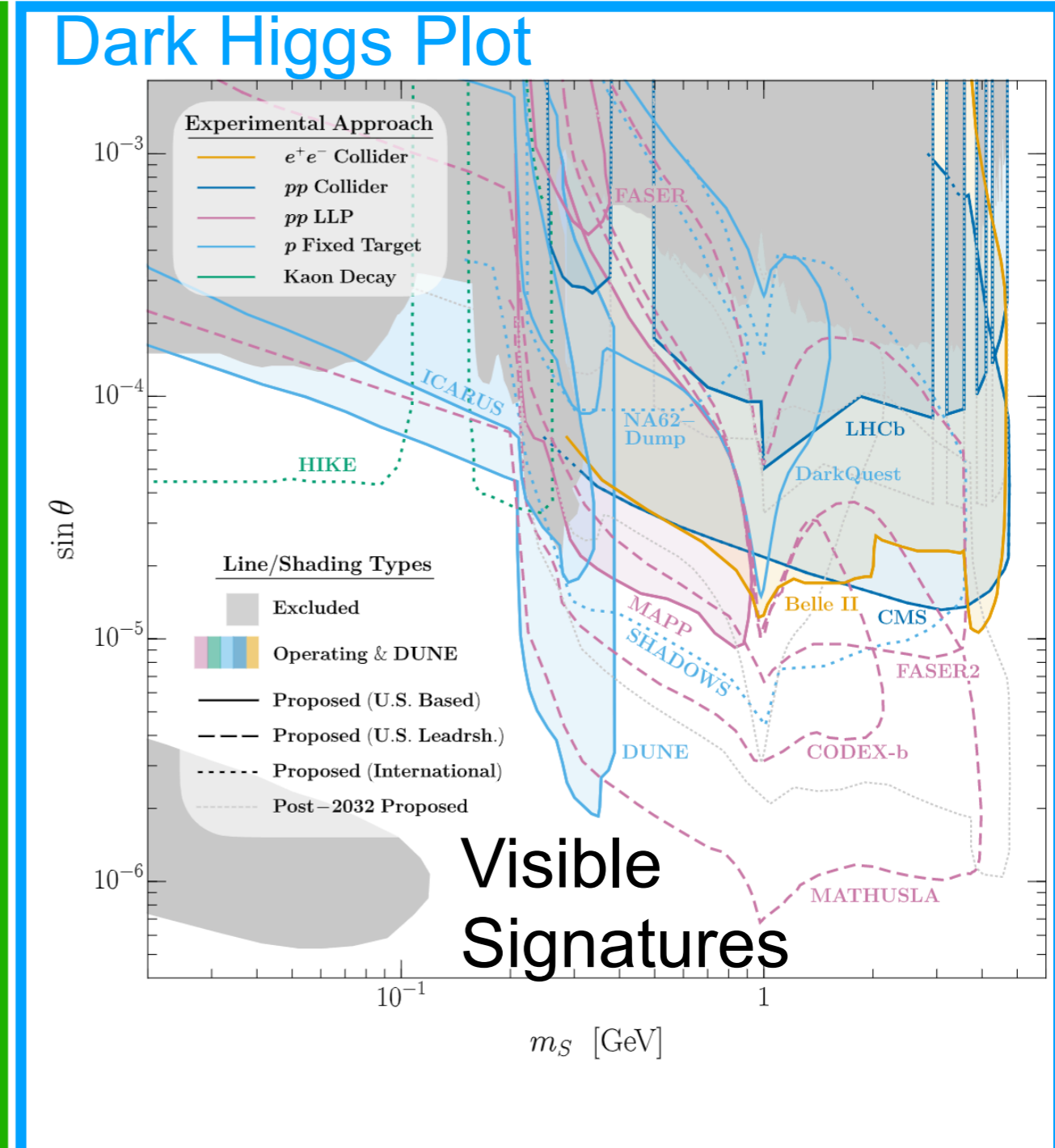
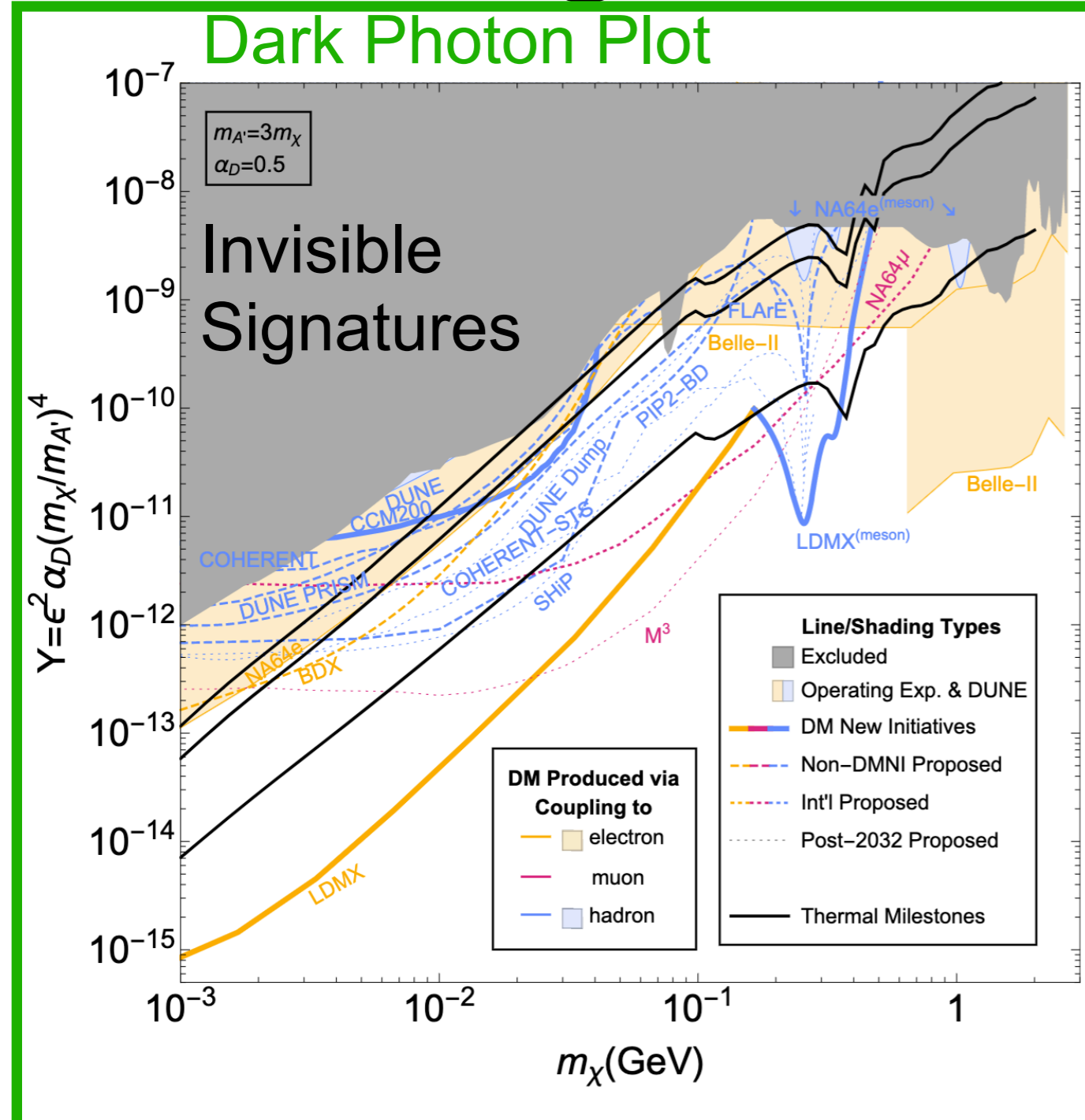
Era of Long Lived
Exploration is just starting
Many more like CMS on
their way

Minimum Coupling Scan

- As w/all simplified DM models there is a minimum coupling
- For the LHC models we can compute the relic density
 - Simplified models, so **relic calculation is simplified**
 - Compute relic density with MadDM
- We scan the full dark matter mass vs mediator mass
 - Useful to see



Light DM at Snowmass

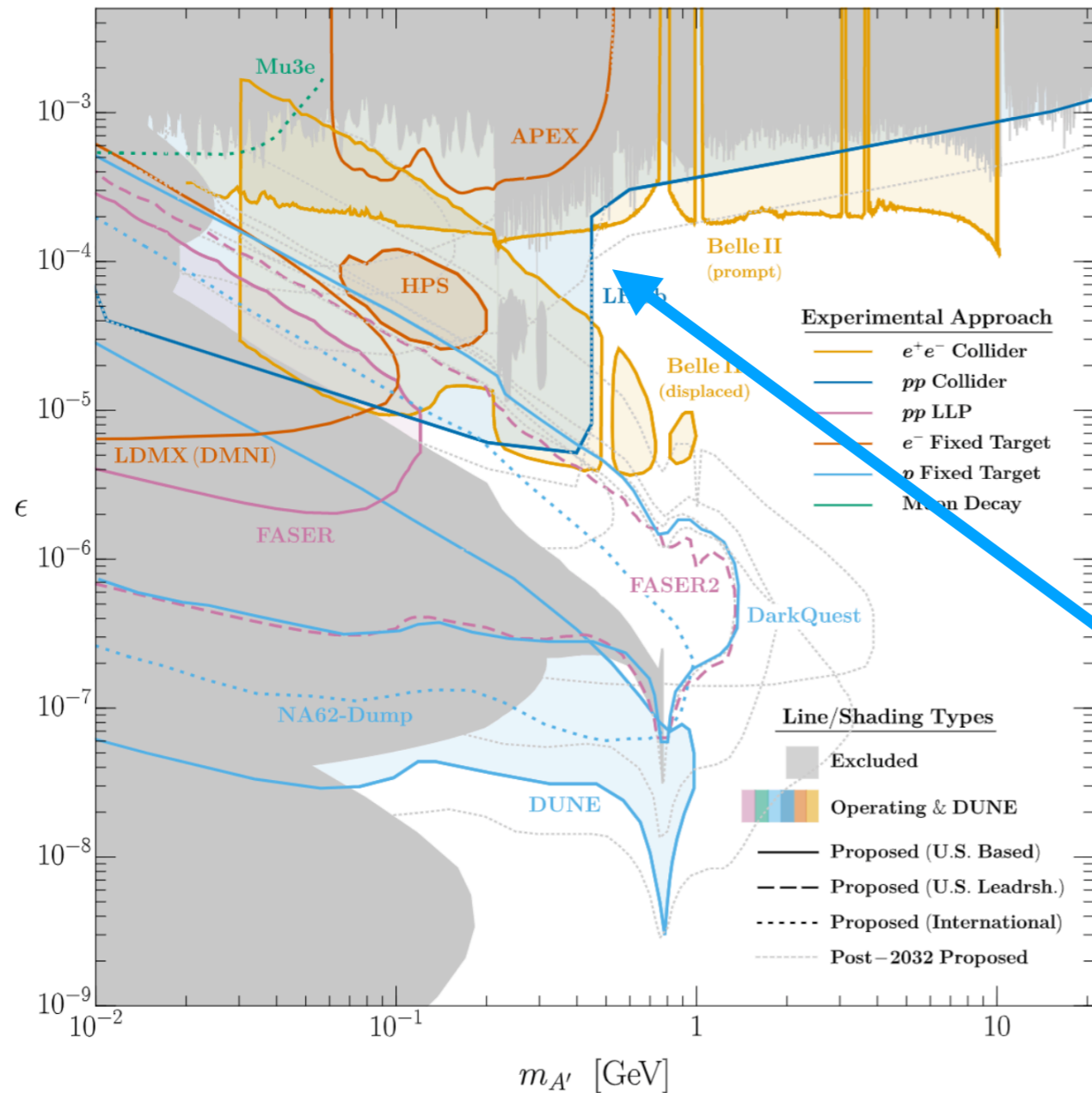


We will focus on invisible signatures for LHC to compare with

There are some cases that Light DM focuses not directly relevant

Future Connections

Weakly Coupled Dark Photons



Effort to highlight **weak coupled Dark Photon**

Coupling weak enough to be long-lived

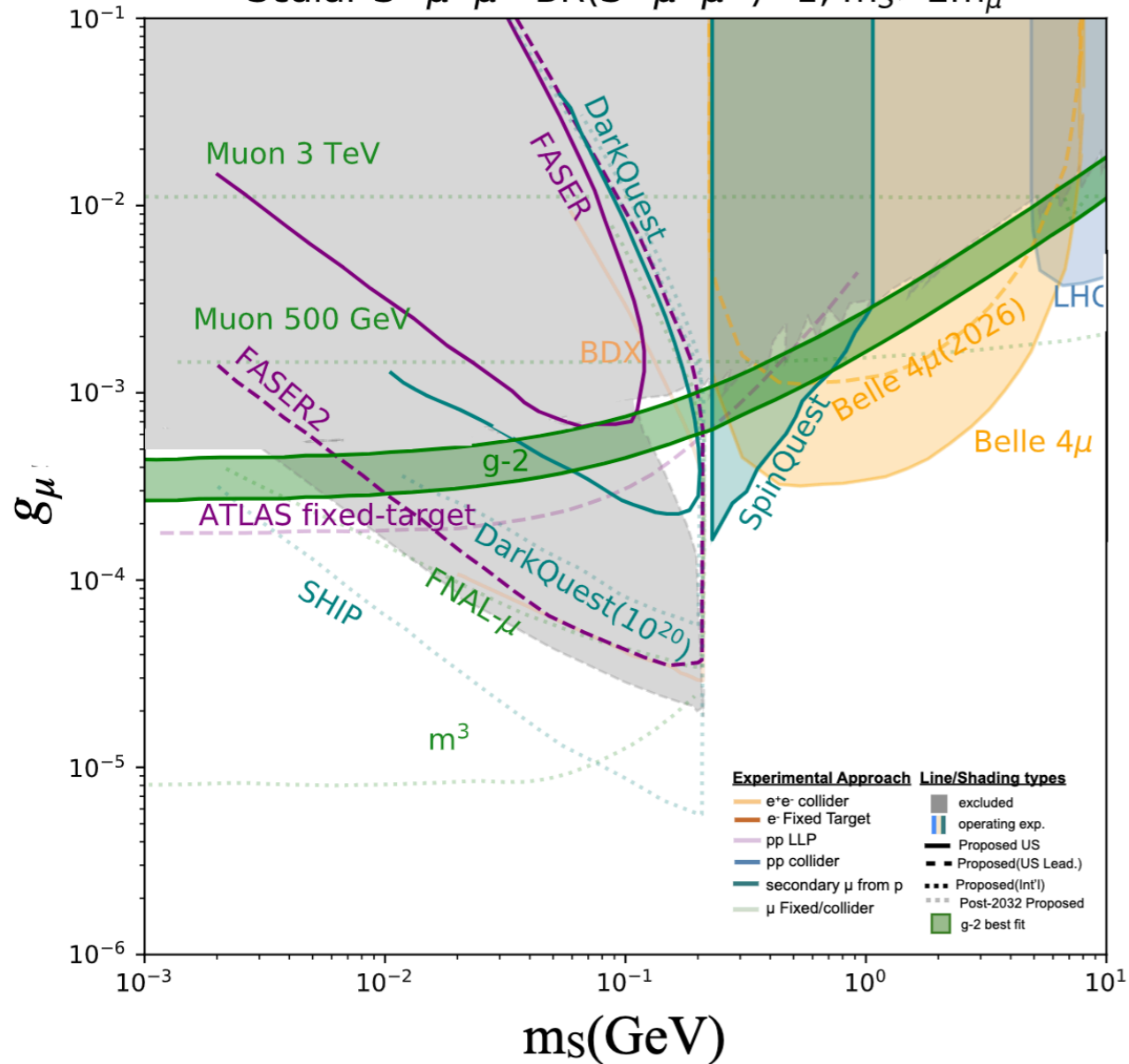
Potential to connect w/LL group

Light DM considered $g-2$ models highlights specific final states

Other Highlights

Dark Scalar Coupling to Muons

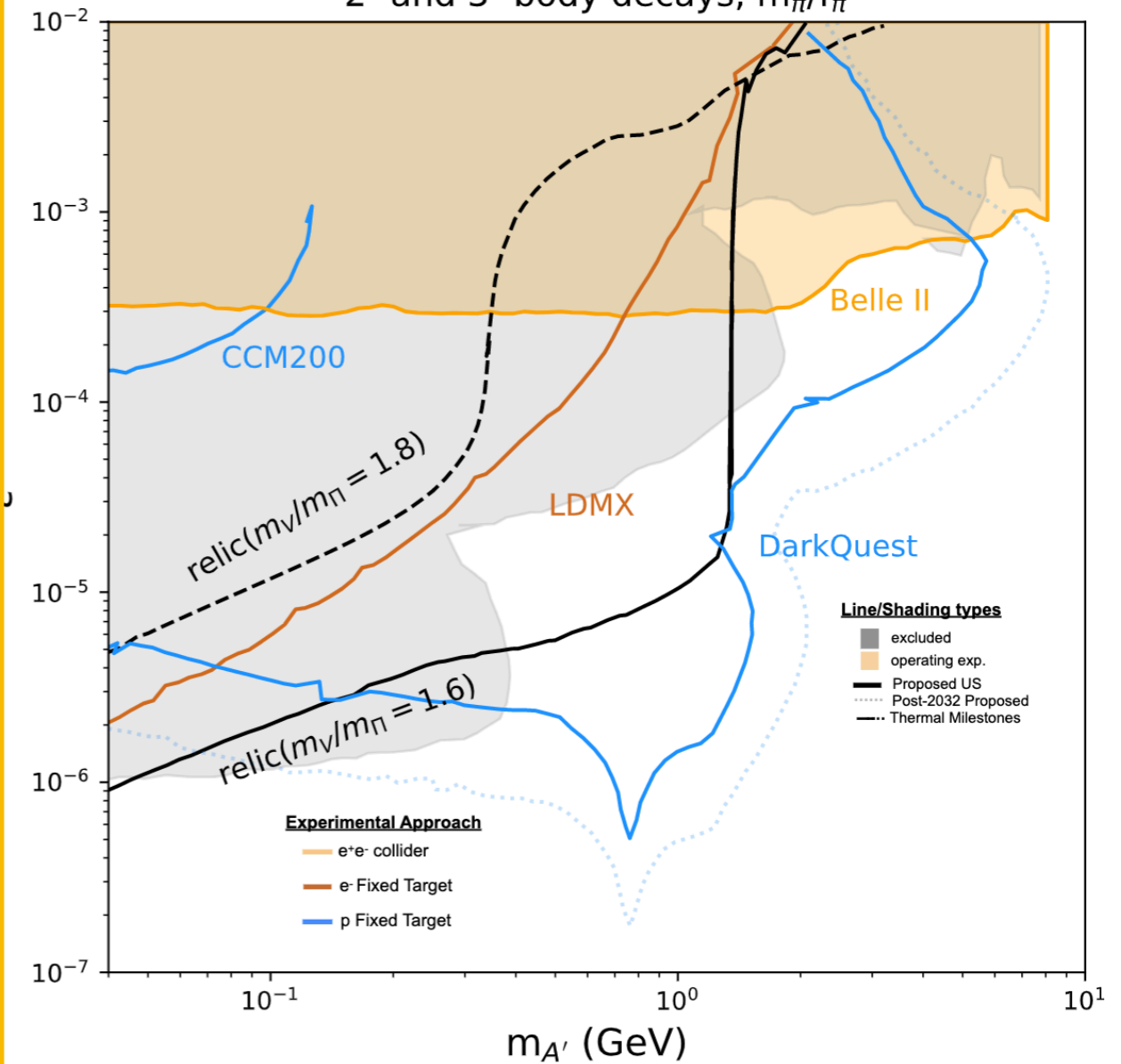
Scalar $S \rightarrow \mu^+ \mu^-$ $BR(S \rightarrow \mu^+ \mu^-) = 1$, $m_S > 2m_\mu$



Light DM considered g-2 models
highlights specific final states

SIMP Dark Photon Model

2- and 3- body decays, m_π/f_π



Light DM considered g-2 models
highlights specific final states

Comparisons w/PBC

$$\mathcal{L}_{\text{vector}} = -g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q - g_\ell \sum_{\ell=e,\mu,\tau} Z'_\mu \bar{\ell} \gamma^\mu \ell,$$

Adding Mixing with photon

$$g_q = g_\ell = \frac{\epsilon}{2e \cos \theta_W}$$

Portal	Coupling
Dark Photon, A_μ	$-\frac{\epsilon}{2 \cos \theta_W} F'_{\mu\nu} B^{\mu\nu}$
Dark Higgs, S	$(\mu S + \lambda S^2) H^\dagger H$
Axion, a	$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$
Sterile Neutrino, N	$y_N L H N$

LHC Spin 1 results are very similar to Dark Photon in PBC
 For the most part simple rescaling can allow for result comparisons

Dark Photon's have previously been discussed here <https://indico.cern.ch/event/729789/>

<https://arxiv.org/pdf/1901.09966.pdf>

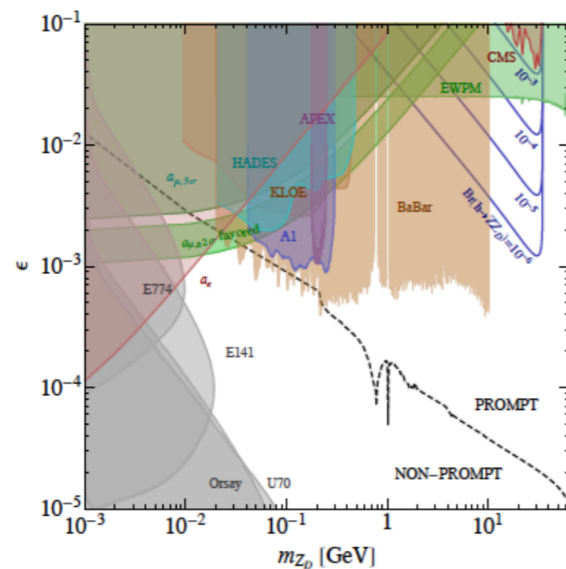
Actually Reconciling

- To reconcile the models we **wanted a Madgraph Model**
 - Started from here Dark Vector + Dark Higgs model here

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{\epsilon}{2 \cos \theta_W} F'_{\mu\nu} B^{\mu\nu} + g_{DM} \cos(\theta_a) Z_D \chi\chi + g_{DM} \sin(\theta_a) Z \chi\chi$$

We started with a Madgraph model with
Dark Photon to SM couplings
Also, includes Dark Higgs

Adding DM terms to the
model so we can probe
invisible decays



D. Curtin et al. ([Phys. Rev. D 90, 075004 \(2014\)](#))

In the following slides we will recast
the CMS monojet analysis and
projections to Dark Photon

**Just look at the invisible final state
(LDMX/Belle bounds at low mass)**

Analytic Form

- Additionally with model we can compare w/LHCDMWG
 - From the Lagrangian we can write

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{\epsilon}{2 \cos \theta_w} F'_{\mu\nu} B^{\mu\nu} + g_{DM} \cos(\theta_a) Z_D \chi\chi + g_{DM} \sin(\theta_a) Z \chi\chi$$

$$\begin{pmatrix} Z \\ Z_D \end{pmatrix} = \begin{pmatrix} \cos \theta_a & \sin \theta_a \\ -\sin \theta_a & \cos \theta_a \end{pmatrix} \begin{pmatrix} Z_0 \\ X \end{pmatrix} \quad \text{Taking usual mixing scenario}$$

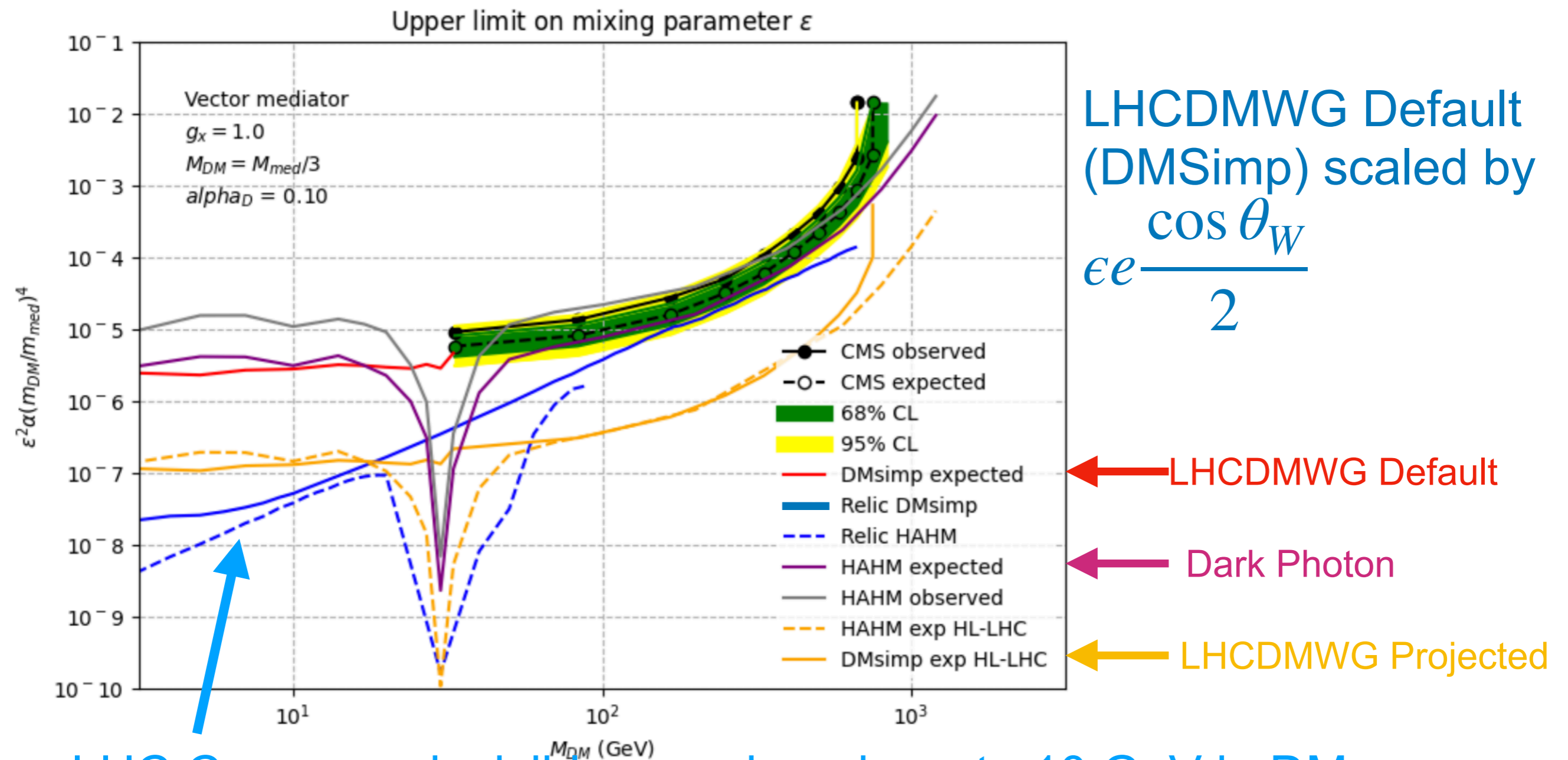
$$g_q = \frac{e \sin \theta_a}{2 \tan \theta_w} \approx e\epsilon \frac{1}{\Delta_z - 1} \frac{\cos \theta_w}{2}$$

$$\Delta_z = \left(\frac{M_{z'}}{M_z} \right)^2$$

Master Formula Allows us to translate between the two

The Result

- LHC Monojet Analysis is in MadAnalysis
 - Relic density computed with MadDM (maps well)

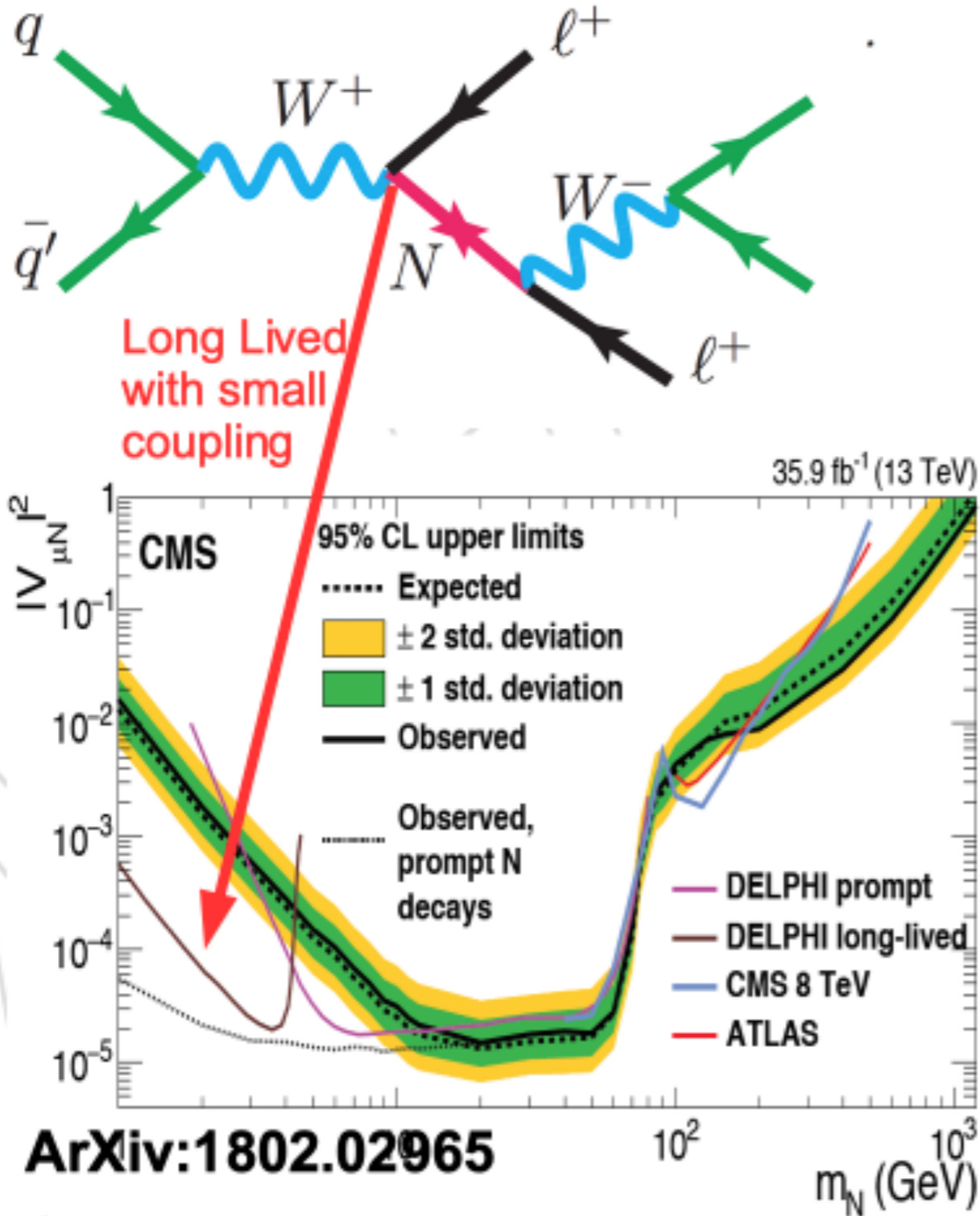


LHC Can cover invisible searches down to 10 GeV in DM mass

HNLS

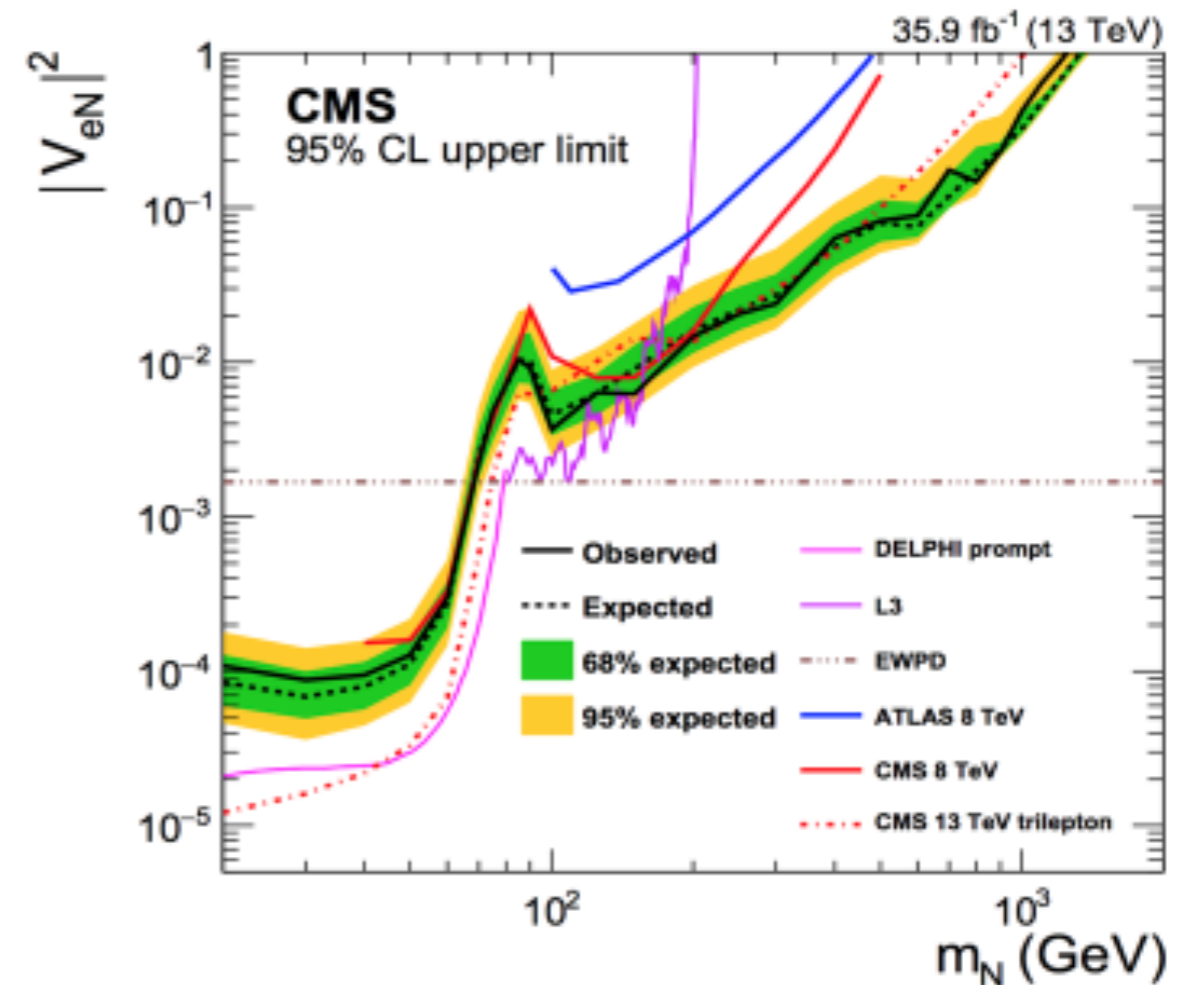
Expect Bounds at low mass to improve

Once long lived analysis is performed

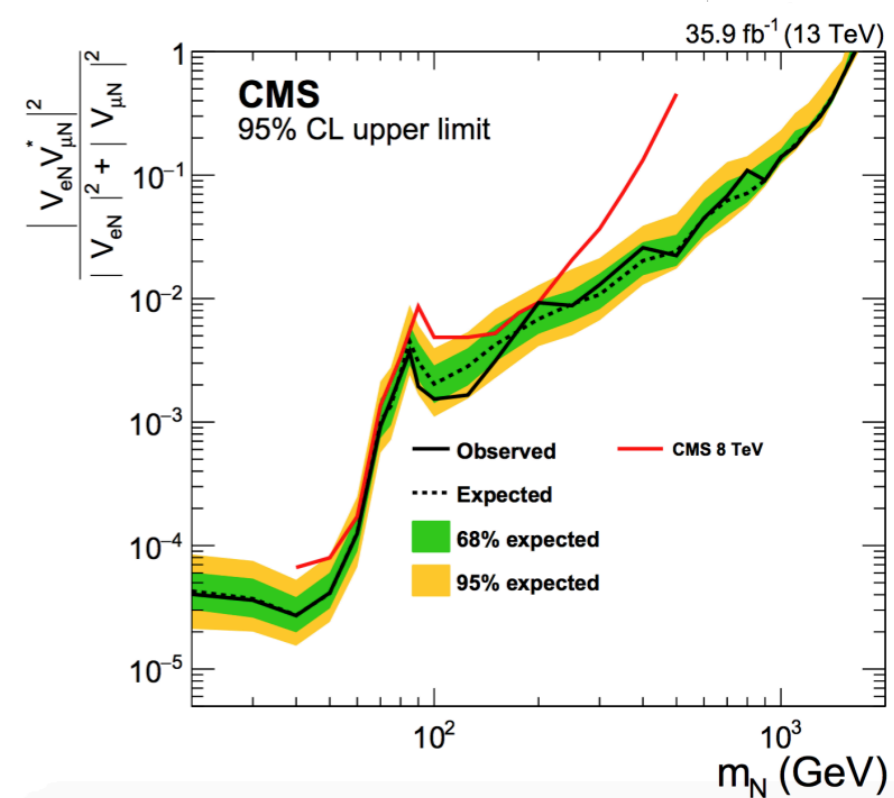
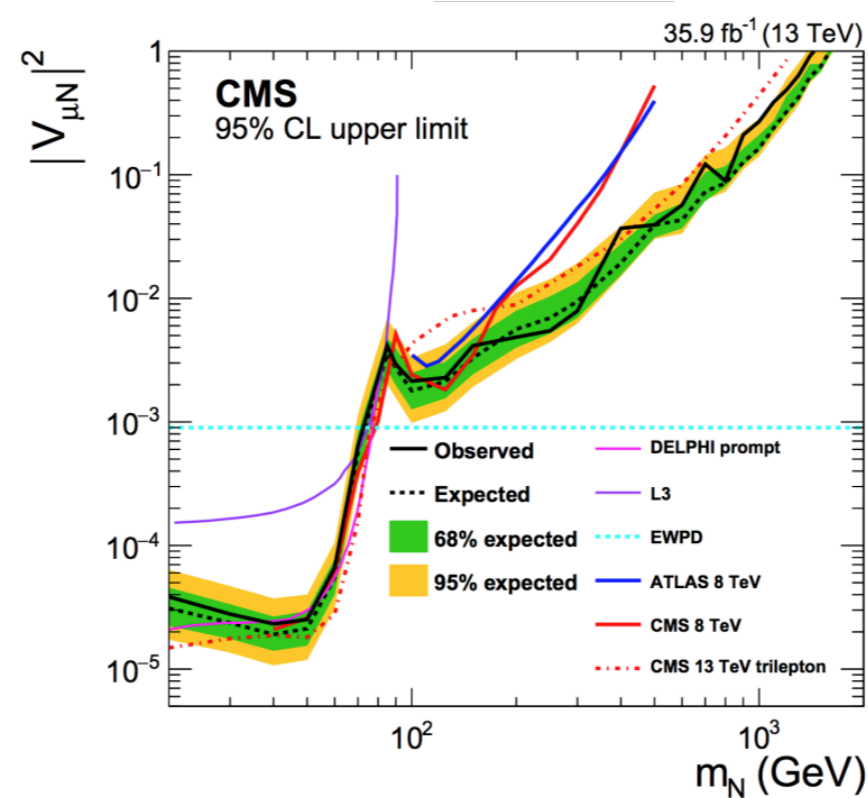
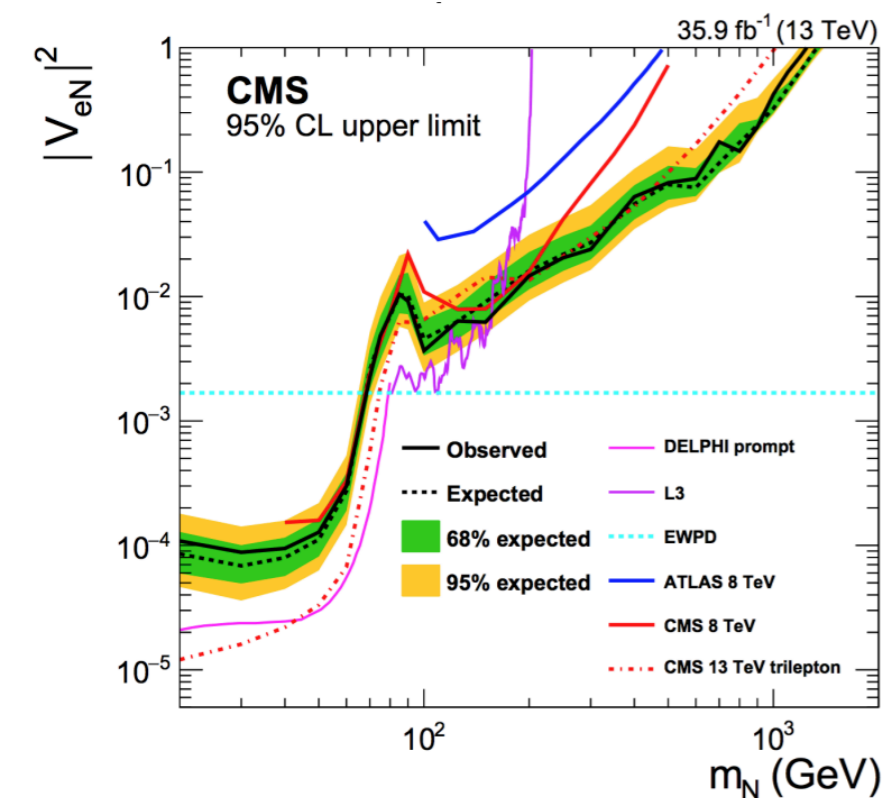
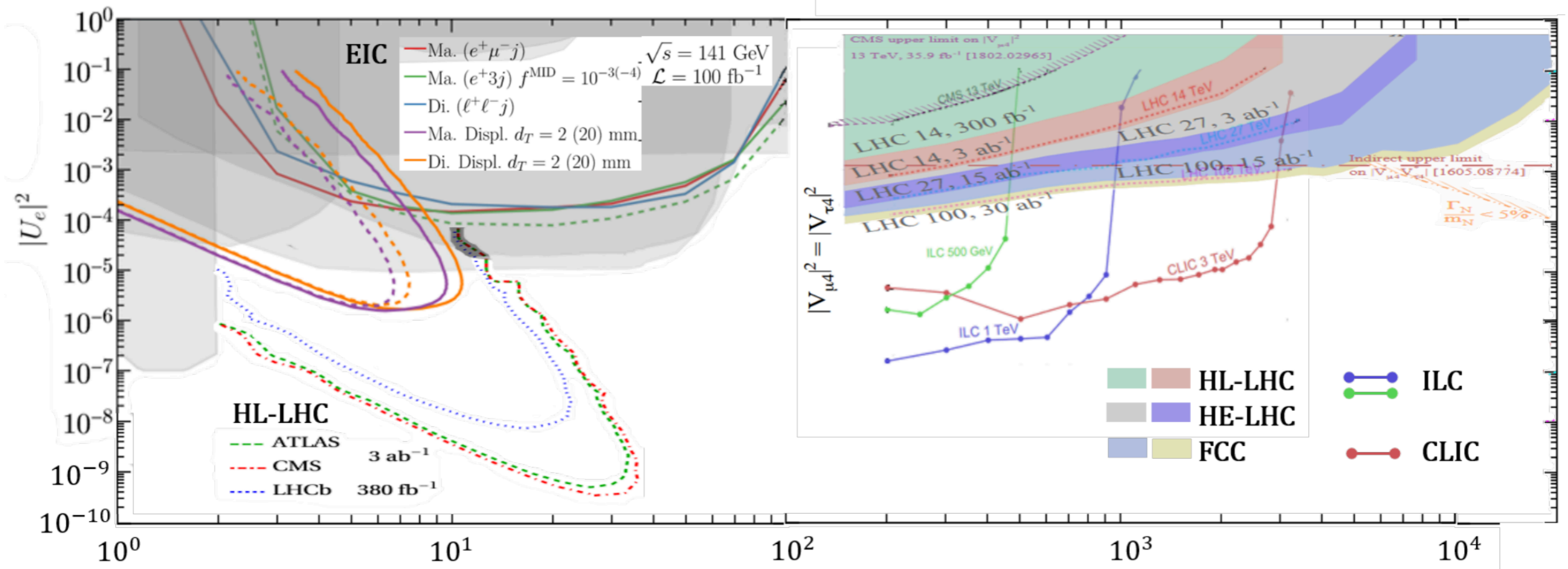


ArXiv:1802.02965

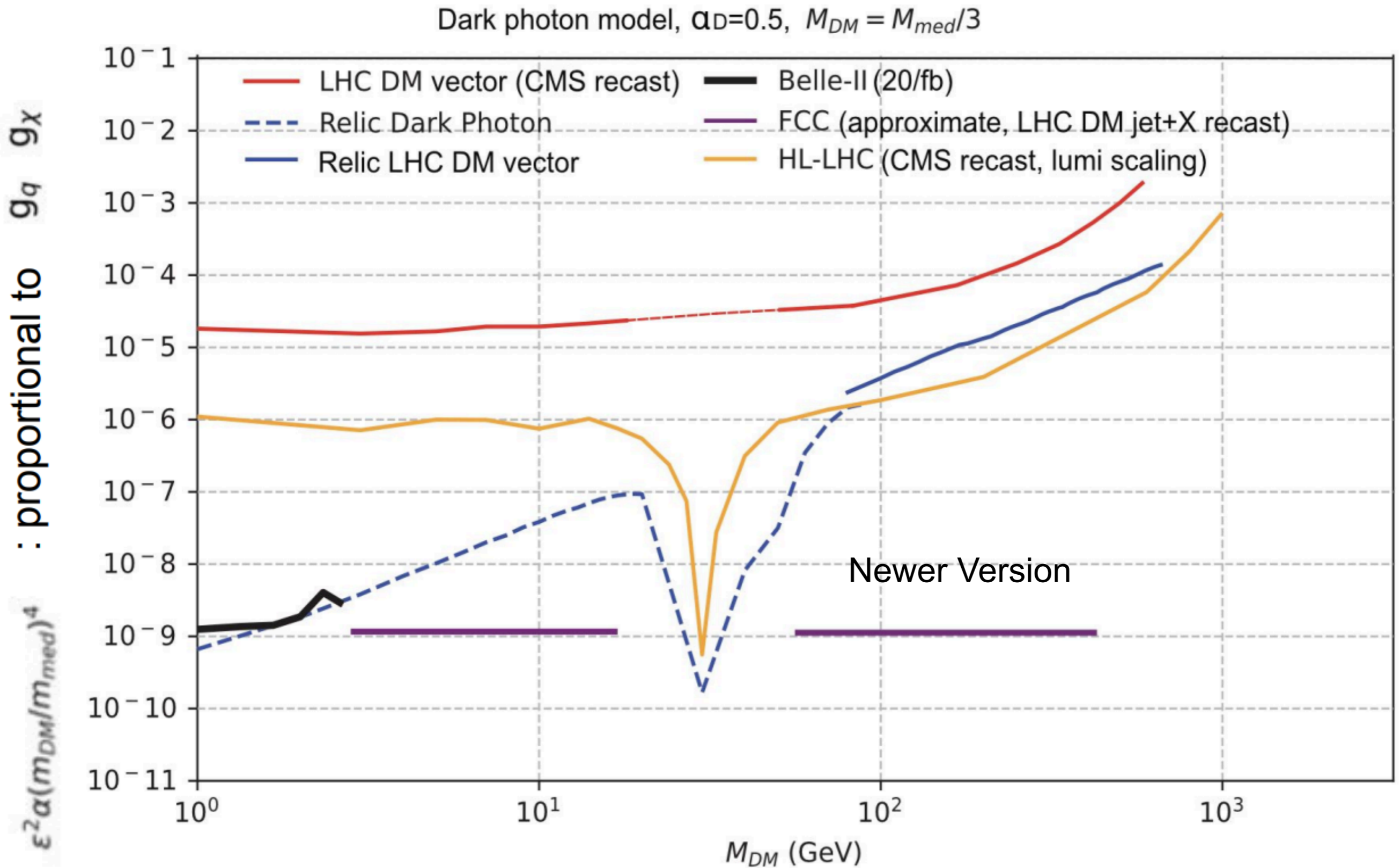
CMS-EXO-17-028



HNLS



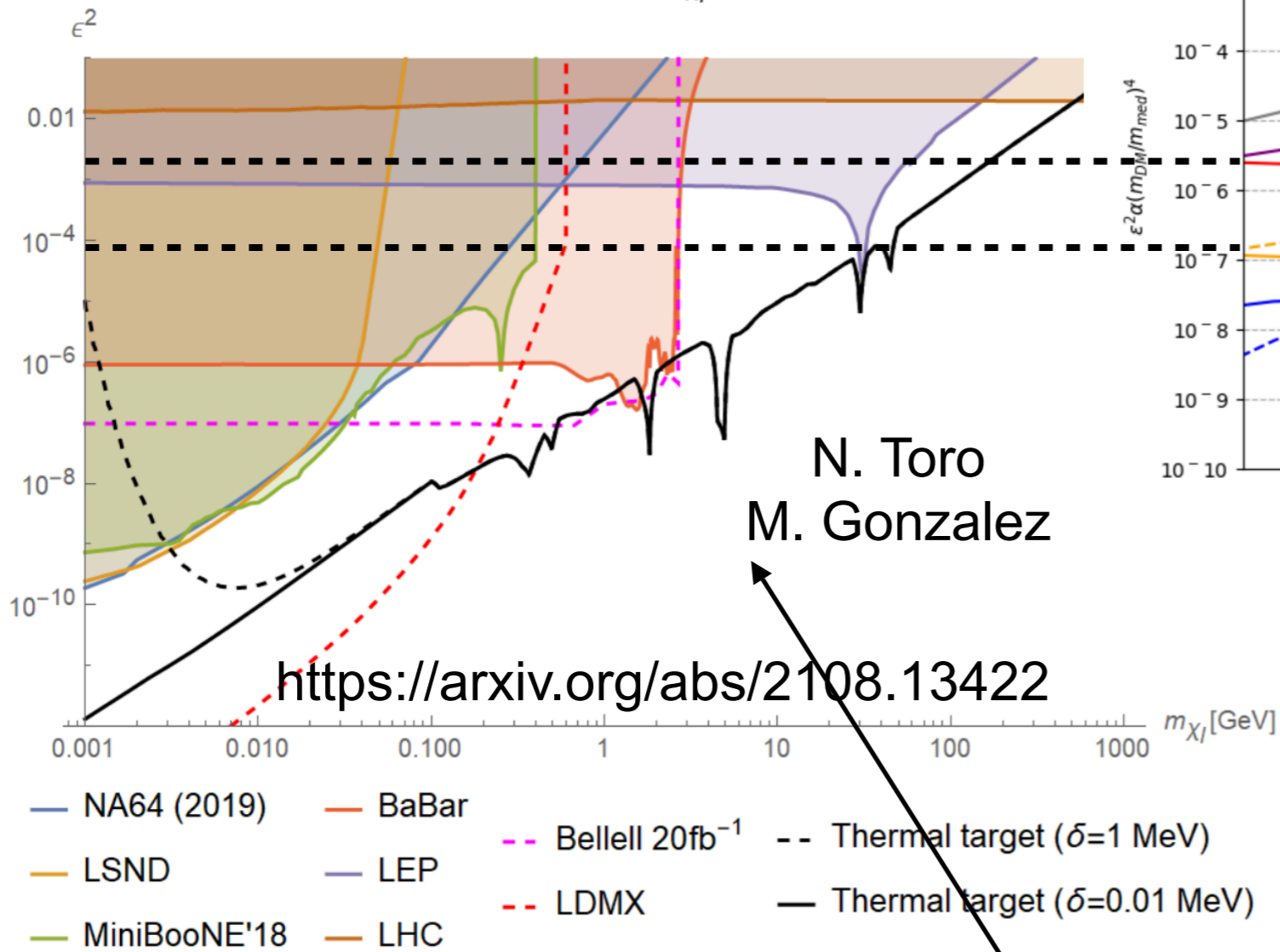
Dark Photon w/ $\alpha_D=0.5$



Check of Some Params

Relic calculations match pretty closely with other calculations

$$\alpha_D = 0.5, \quad m_{A'} = 3(m_{\chi_I} + \delta/2)$$

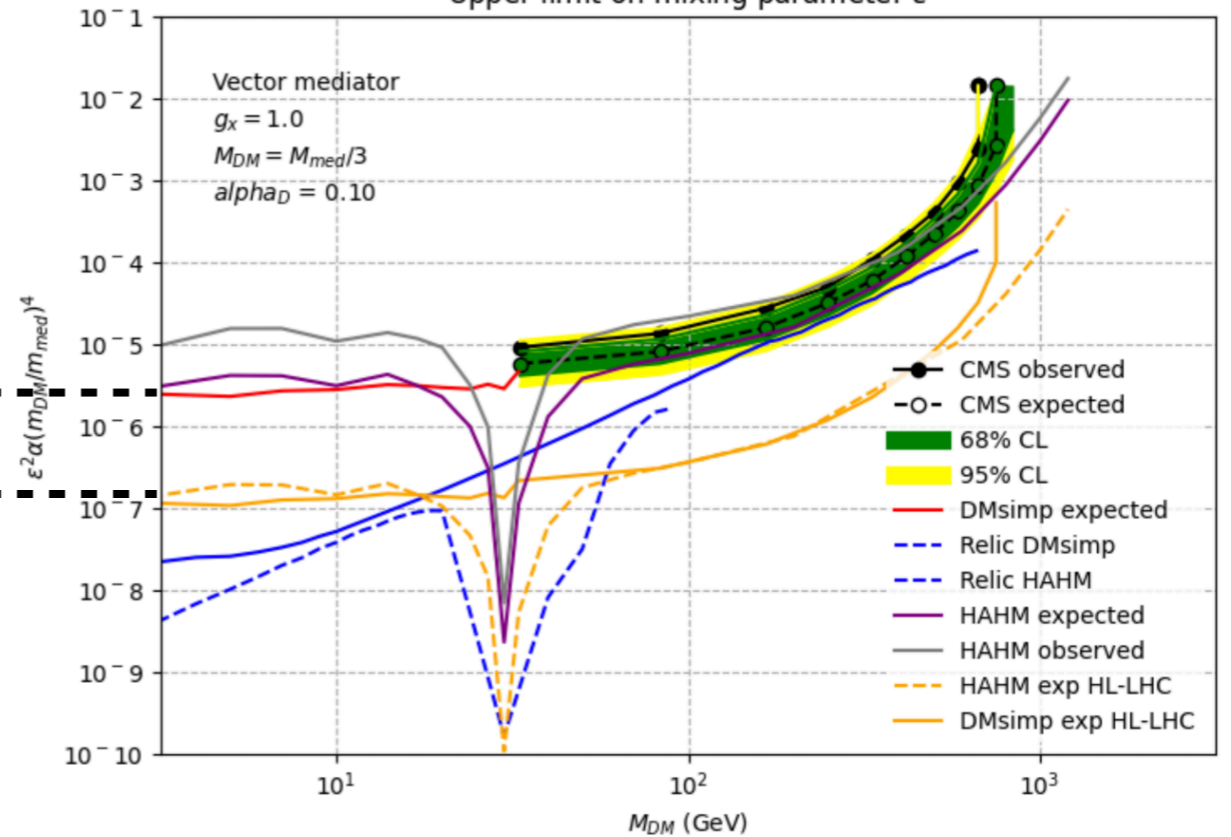


<https://arxiv.org/abs/2108.13422>

N. Toro
M. Gonzalez

Cross Check

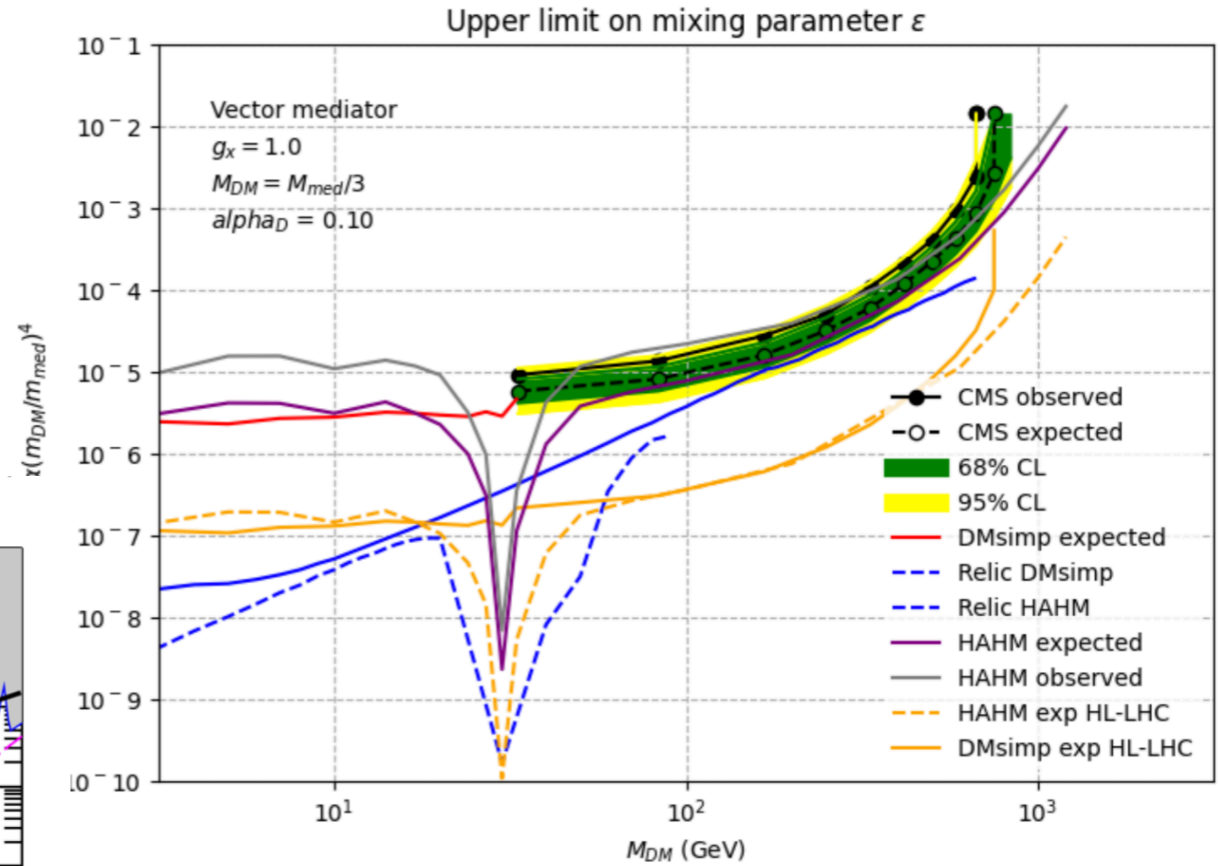
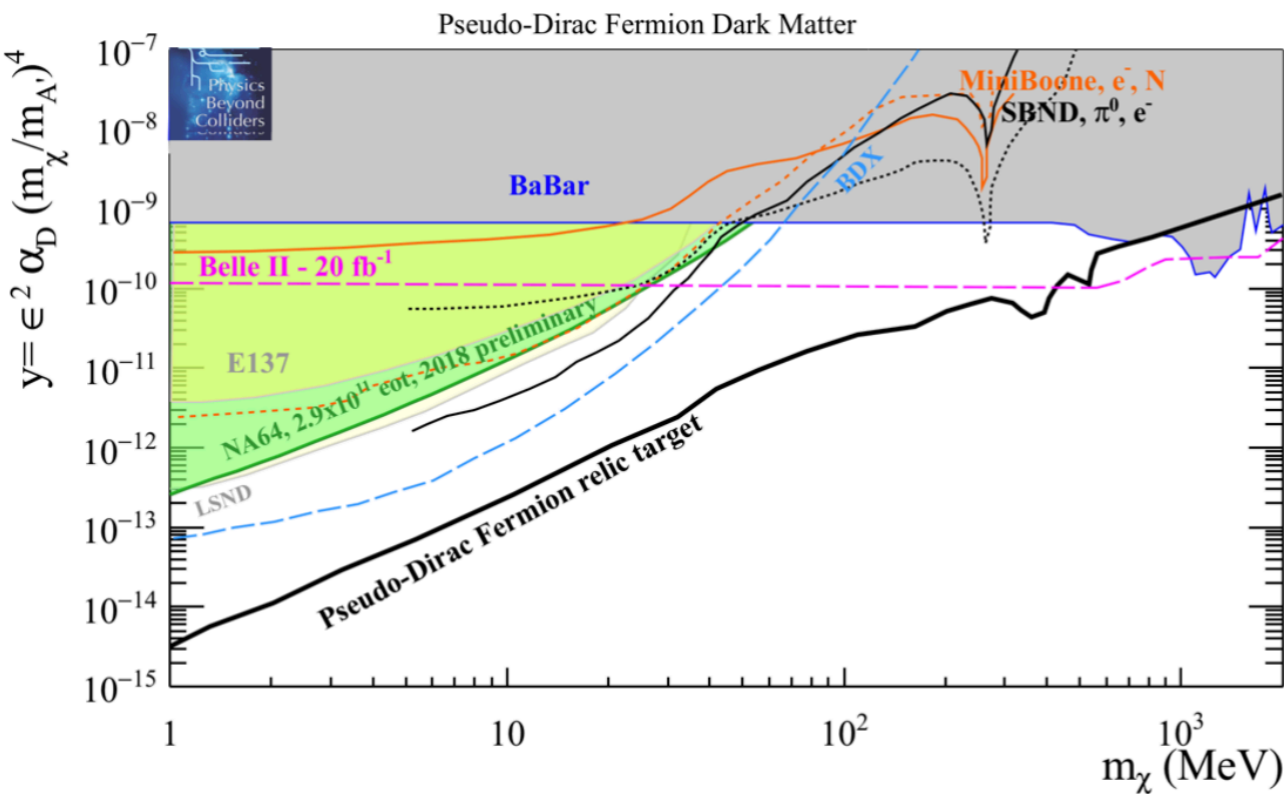
Upper limit on mixing parameter ϵ



Bounds from LHC appear stronger than on left plot

Now Connecting them

Appears that we can now connect plots into one



LHC is complementary
 Similar goals on similar
 timescales

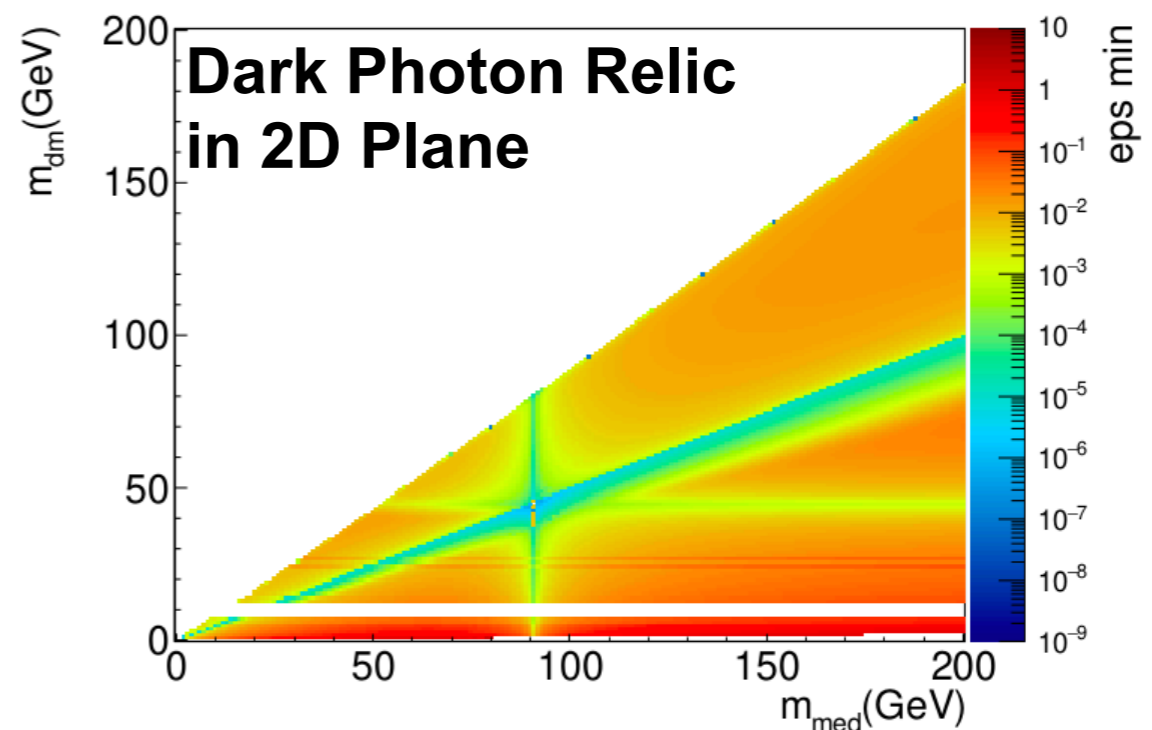
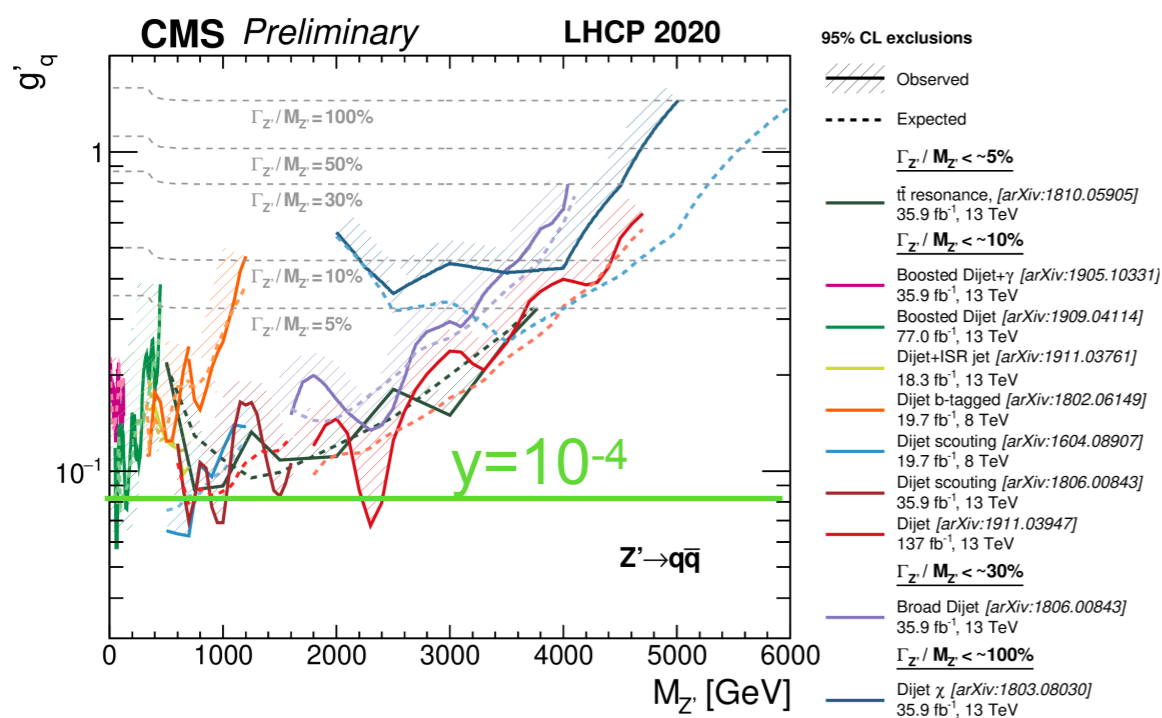
LHC goes left when mediator mass gets larger



Note Also that as $\frac{m_{DM}}{m_{med}}$ gets larger LHC DM searches are the only game in town

Additional Plots

- With Madgraph model we have some flexibility
 - MG mode has the full Higgs to dark photon couplings
 - Can envision adding the Higgs/Dark Higgs bounds
 - Visible searches provide bounds for heavy DM
- Since $g_q=0.01-0.1$ maps $y=10^{-7}-10^{-4}$ include jets/lepton bounds
 - $y > 10^{-4}$ we have largely excluded this up to 2 TeV



Comparisons w/PBC

$$\mathcal{L}_{\text{scalar}} = -g_{\text{DM}} \phi \bar{\chi} \chi - g_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} q$$

Enforcing a mixing with the Higgs
Higgs to Invisible dominates
bounds (adds VBF channel)

$$g_q = -\sin \theta$$

Portal	Coupling
Dark Photon, A_μ	$-\frac{\epsilon}{2 \cos \theta_W} F'_{\mu\nu} B^{\mu\nu}$
Dark Higgs, S	$(\mu S + \lambda S^2) H^\dagger H$
Axion, a	$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$
Sterile Neutrino, N	$y_N L H N$

DMWG presents results as a scalar w/o Higgs mixing

This eliminates the ϕ to SM vector boson coupling

However, Higgs to invisible is presented with Singlet Mixing model

Singlet Mixing Model

$$\mathcal{L} \supset -y_{\text{DM}} s \bar{\chi} \chi - \mu s |H|^2$$

What if we make a complete singlet scalar model?

Observed mass eigenstates

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

With vector boson interactions it will mix w/Higgs

$$\mathcal{L} \supset -y_{\text{DM}} (\sin \theta h_1 + \cos \theta h_2) \bar{\chi} \chi \quad \text{Higgs to Invisible}$$

$$+ (\cos \theta h_1 - \sin \theta h_2) \left(\frac{2M_W^2}{v} W_\mu^+ W^{-\mu} + \frac{M_Z^2}{v} Z_\mu Z^\mu - \sum \frac{m_f}{v} \bar{f} f \right)$$

Standard LHC Model w/MC....

To Map to PBC models

We need to fix DM coupling
and take it very large

Singlet Mixing Model

$$\mathcal{L} \supset -g_{\text{DM}} s \bar{\chi} \chi - \mu s |H|^2$$

What if we make a complete singlet scalar model?

Observed mass
eigenstates

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

With vector boson interactions it will mix w/Higgs

$$\mathcal{L} \supset -g_{\text{DM}} (\sin \theta h_1 + \cos \theta h_2) \bar{\chi} \chi \quad \text{Higgs to Invisible}$$

$$+ (\cos \theta h_1 - \sin \theta h_2) \left(\frac{2M_W^2}{v} W_\mu^+ W^{-\mu} + \frac{M_Z^2}{v} Z_\mu Z^\mu - \sum_f \frac{m_f}{v} \bar{f} f \right)$$

Singlet Mixing Model

$$\mathcal{L} \supset -g_{\text{DM}} s \bar{\chi} \chi - \mu s |H|^2$$

What if we make a complete singlet scalar model?

Observed mass eigenstates

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

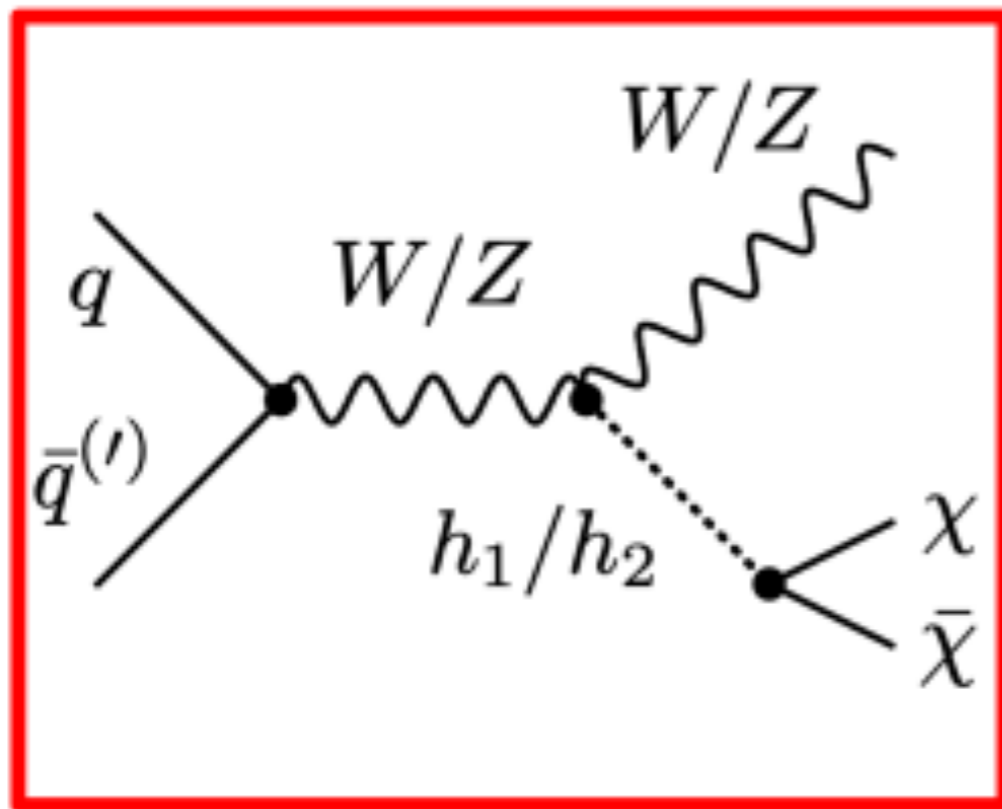
Modified Higgs Vector Boson Couplings

$$\mathcal{L} \supset -g_{\text{DM}} (\sin \theta h_1 + \cos \theta h_2) \bar{\chi} \chi$$

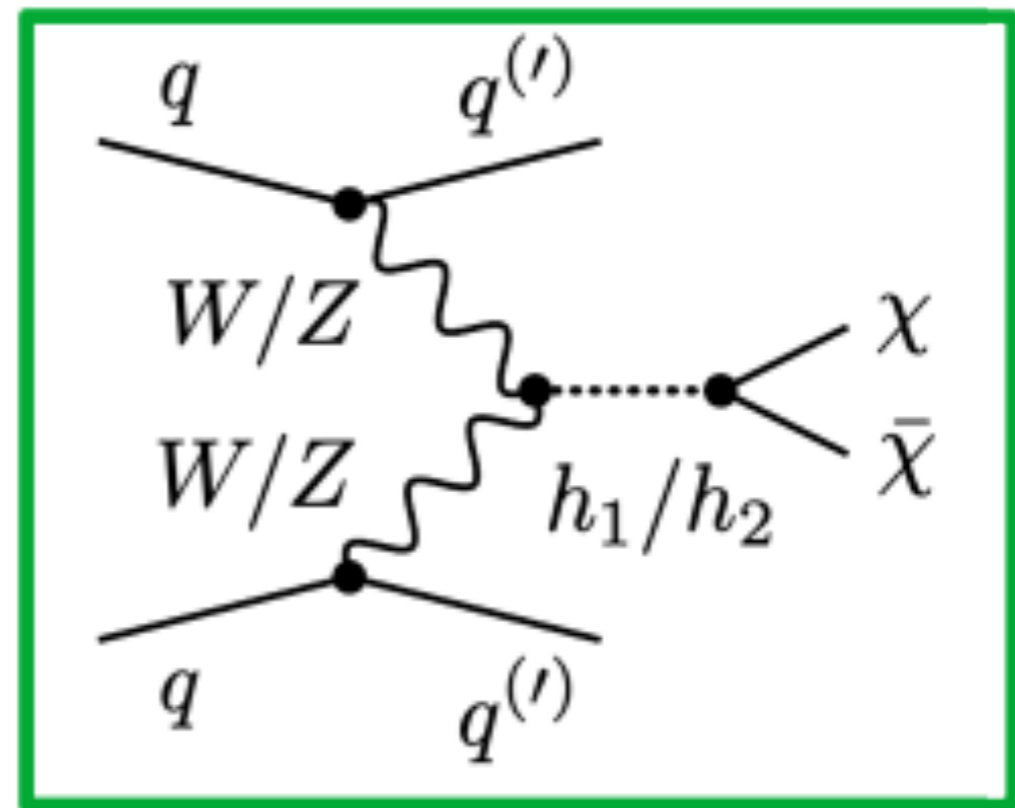
$$+ (\cos \theta h_1 - \sin \theta h_2) \left(\frac{2M_W^2}{v} W_\mu^+ W^{-\mu} + \frac{M_Z^2}{v} Z_\mu Z^\mu - \sum_f \frac{m_f}{v} \bar{f} f \right)$$

What are the scale of Modifications?

$$\Gamma(h_1 \rightarrow \chi\bar{\chi}) = \frac{g_{\text{DM}}^2 \sin^2 \theta m_{h_1}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{h_1}^2}\right)^{3/2}$$



Higgsstrahlung



VBF Higgs to invisible

What Drives Constraints

$$\Gamma(h_1 \rightarrow \chi\bar{\chi}) = \frac{g_{\text{DM}}^2 \sin^2 \theta m_{h_1}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{h_1}^2}\right)^{3/2}$$

Higgs to invisible bounds puts constraints a 10% bound equates to $\sin \theta < 0.002$ (note $g_{\text{DM}} = 1.0$)

Higgs boson coupling of 10% bound equates to $1 - \cos \theta < 0.1 \rightarrow \sin \theta < 0.3$

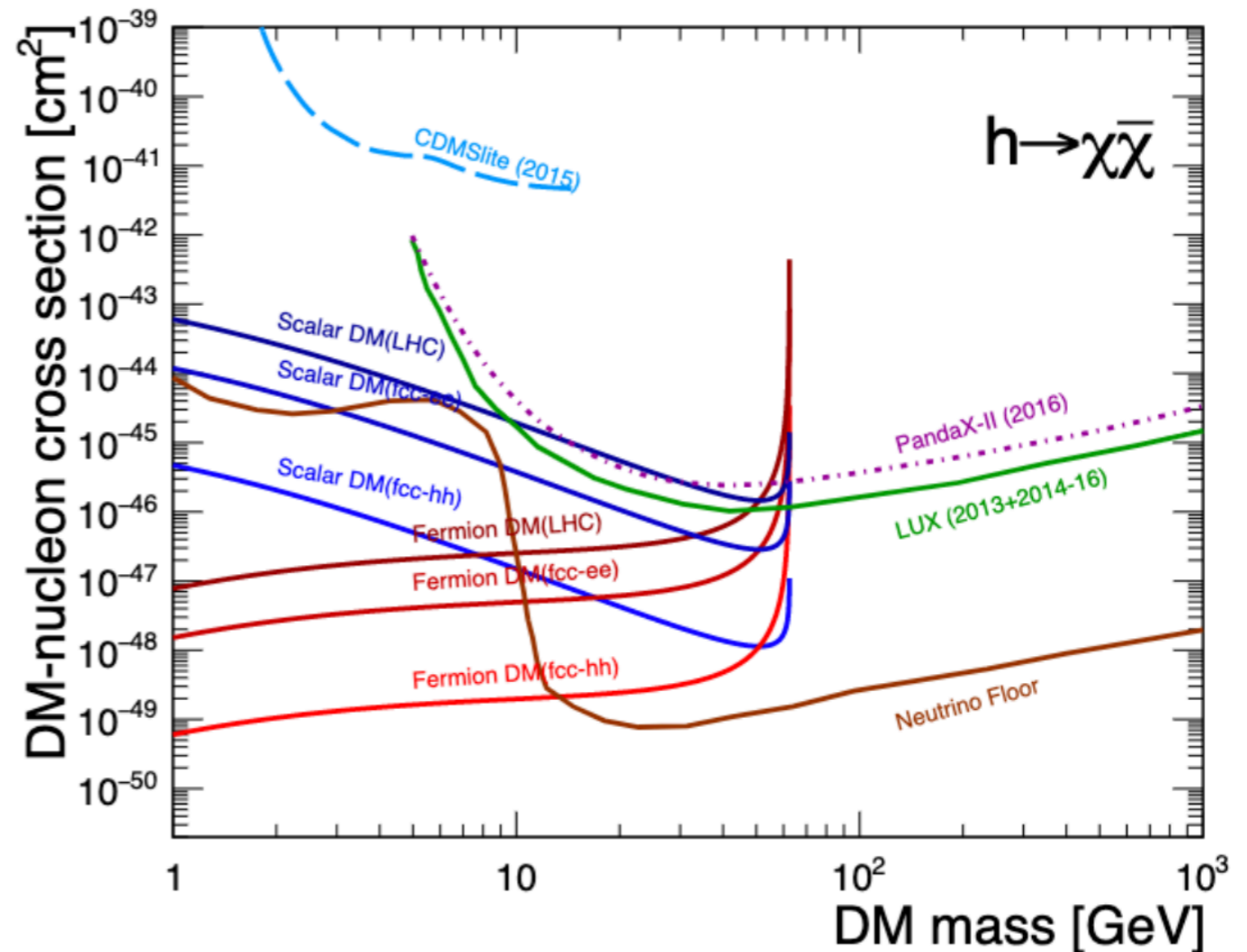
Both invisible decay and Couplings play a critical role

This model is effectively the same as the PBC model

Typically take $g_{\text{DM}} = y_{\text{DM}}$ makes Higgs to invisible less sensitive

Propagating Bounds

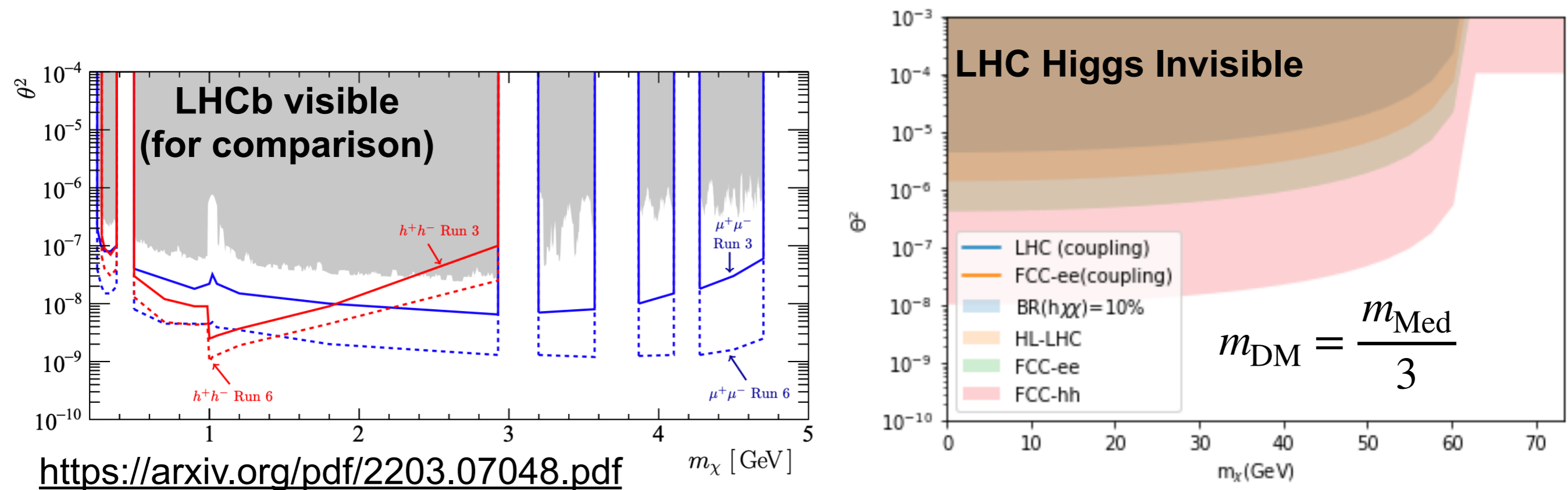
- Higgs to invisible Bounds
 - Current LHC $H(\text{inv}) > 0.1$
 - Future LHC $H(\text{inv}) > 0.02$
 - FCC-ee $H(\text{inv}) > 0.005$
 - FCC-hh $H(\text{inv}) > 0.0001$



- Current projections of Higgs to invisible similar to Direct Detection
 - Sensitivities comparable in the low DM mass region
 - LHC exceed neutrino floor for light DM

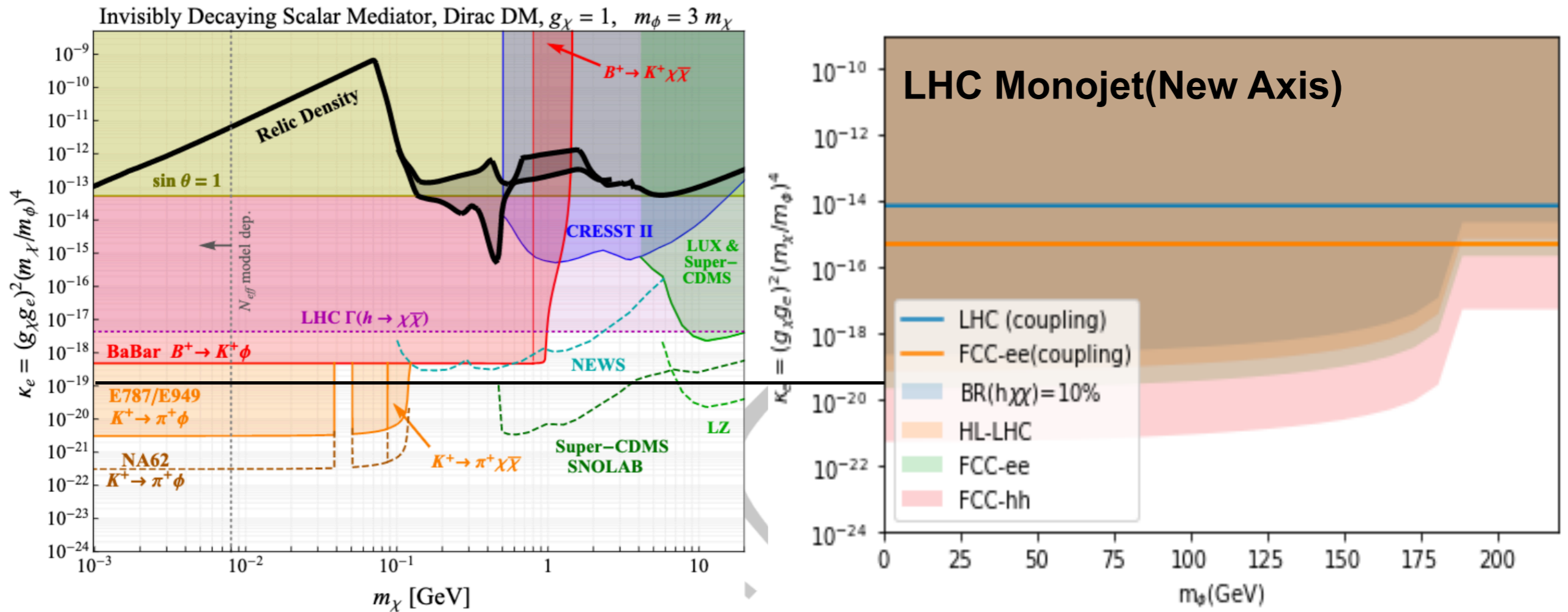
Comparing Standard Plot

- Often the scalar portal is presented in terms of θ^2
 - LHC bounds have clear and large sensitivity



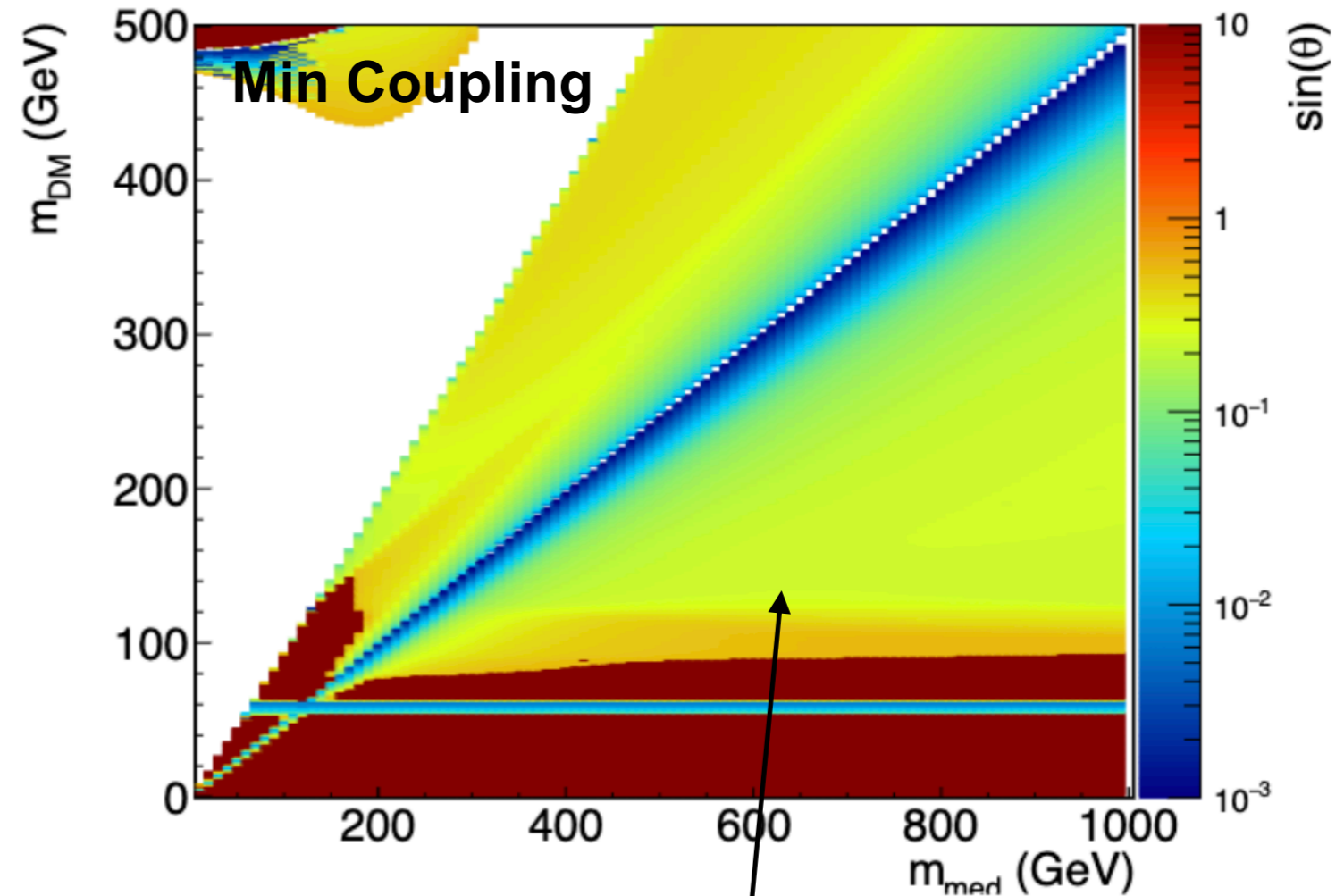
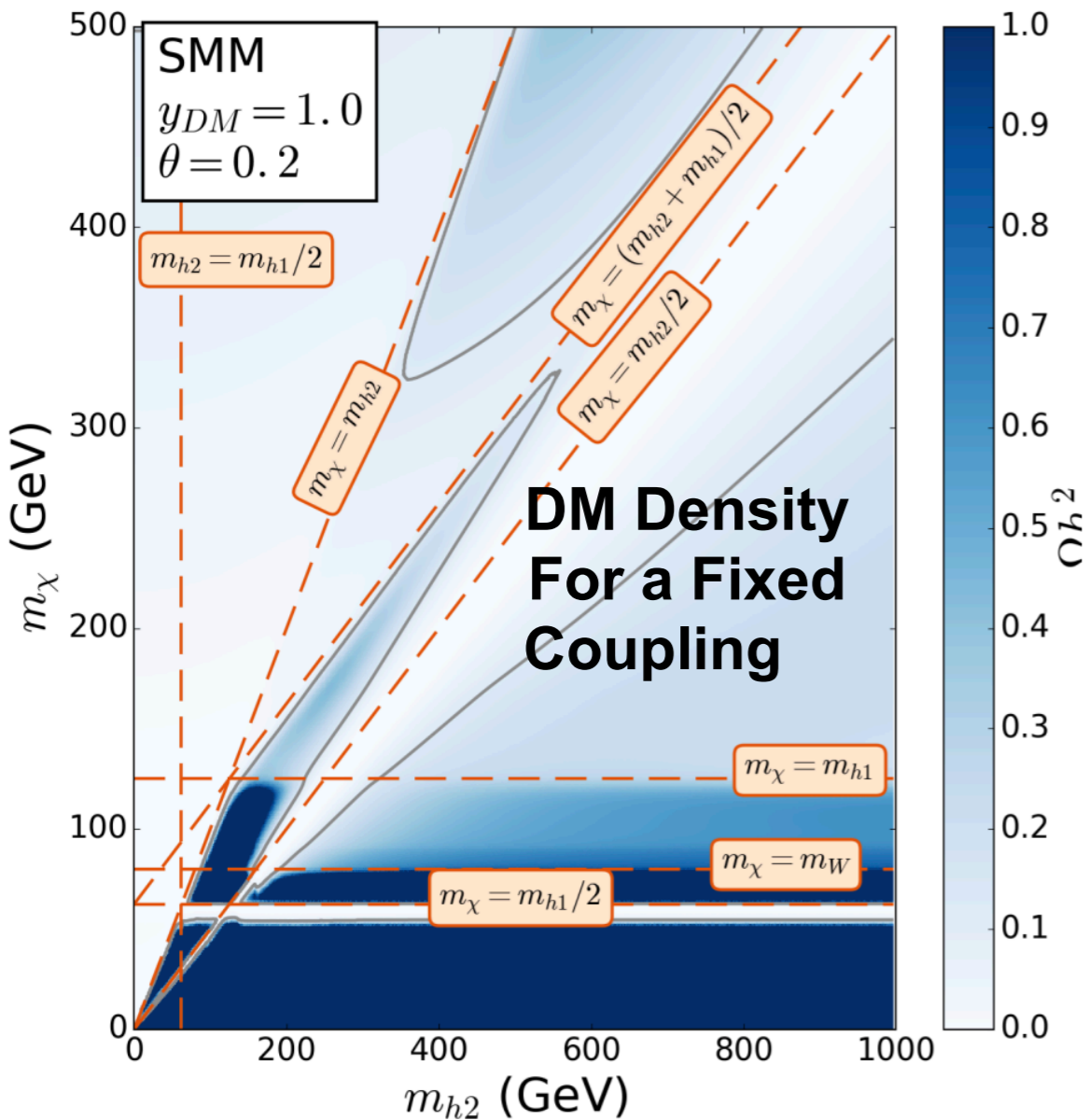
- Bounds for Monojet(invisible) comparable to visible bounds
 - Covers a variety of important final states

Scalar DM Bounds



- LHC Higgs to invisible dominates the scalar DM bounds
 - Additionally Higgs couplings bounds also impact bounds
 - Overall extends sensitivity beyond range of light DM models

Relic Density



- Overall minimum coupling bound is very large
 - Mostly constrained by a 5% Higgs coupling measurement
 - A 5% Higgs coupling bound is an equivalent bound on $\sin \theta < 0.1$

Comparisons w/PBC

$$\mathcal{L}_{\text{pseudo-scalar}} = -ig_{\text{DM}}\phi\bar{\chi}\gamma_5\chi - ig_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q}\gamma_5 q,$$

Pseudoscalar mediator again similar
Interpretation of couplings also similar

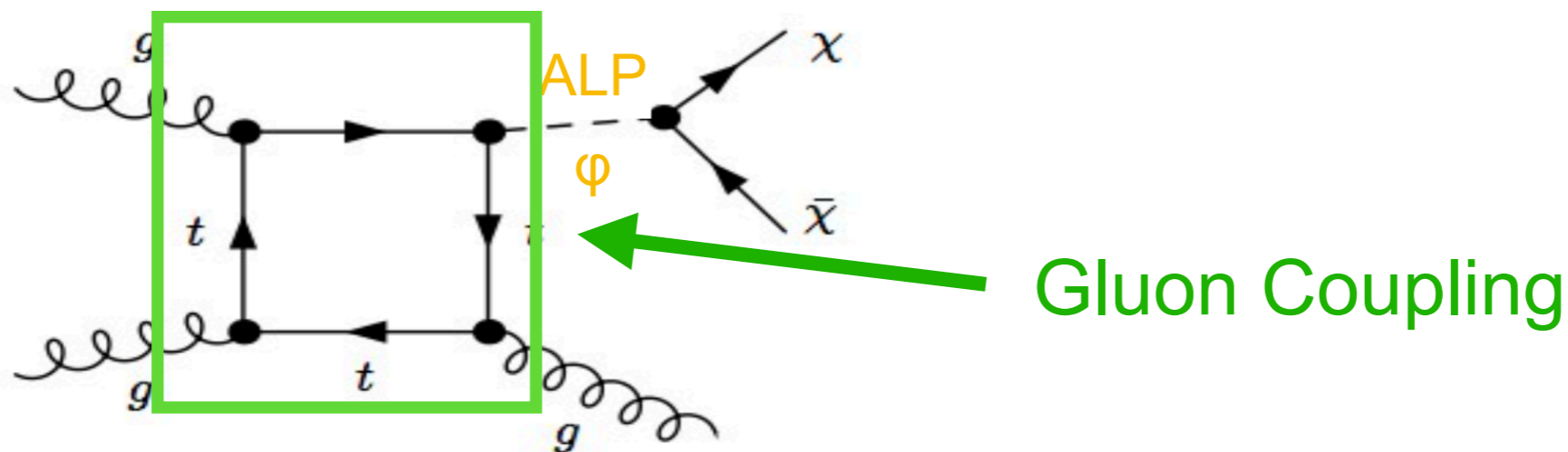
$$g_q = \frac{v}{f_a}$$

Portal	Coupling
Dark Photon, A_μ	$-\frac{\epsilon}{2\cos\theta_W} F'_{\mu\nu} B^{\mu\nu}$
Dark Higgs, S	$(\mu S + \lambda S^2) H^\dagger H$
Axion, a	$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$
Sterile Neutrino, N	$y_N L H N$

DMWG tends to present pseudoscalar results in two ways:
A single mediator (as a simplified model)
A mediator within a 2HDM

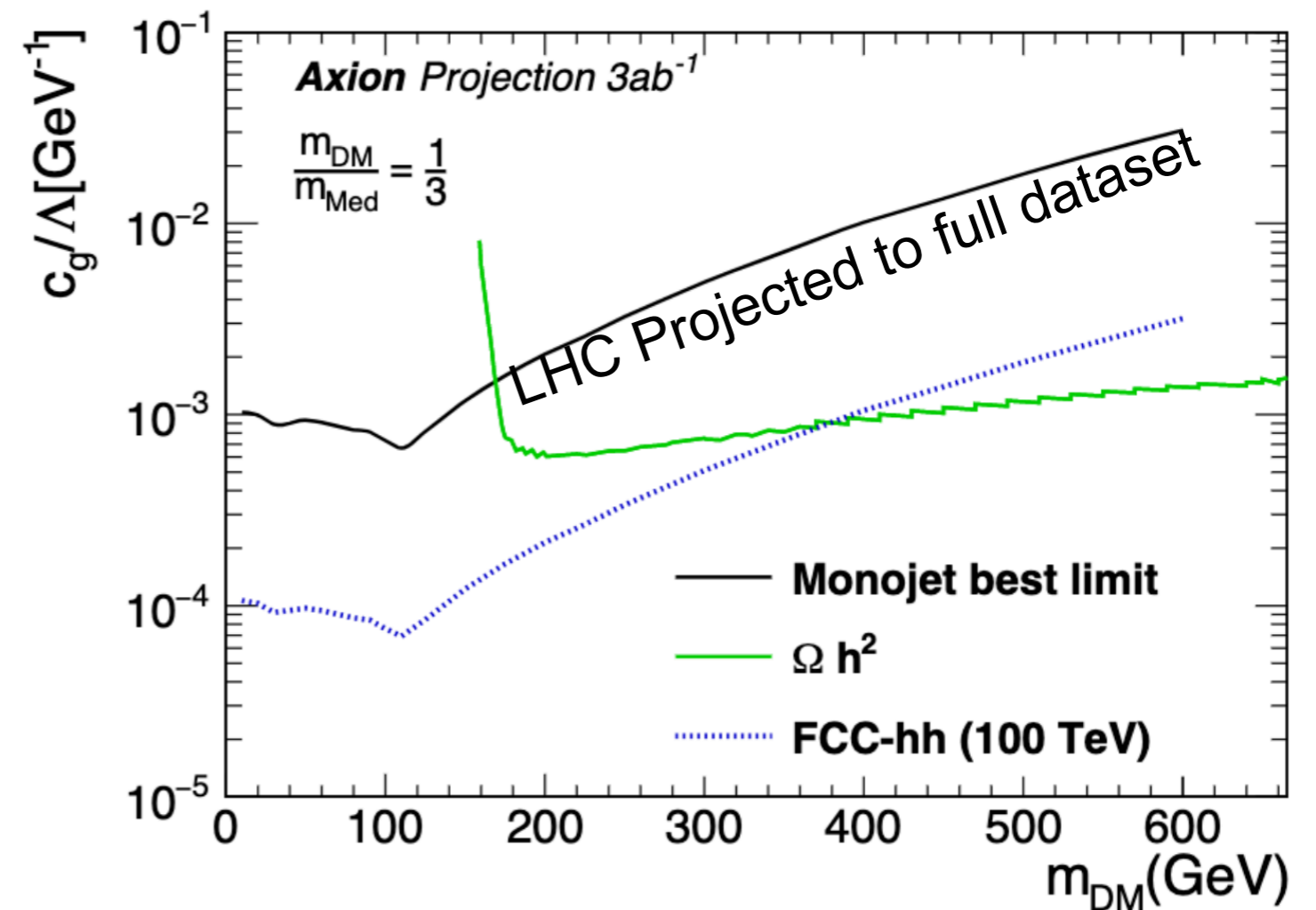
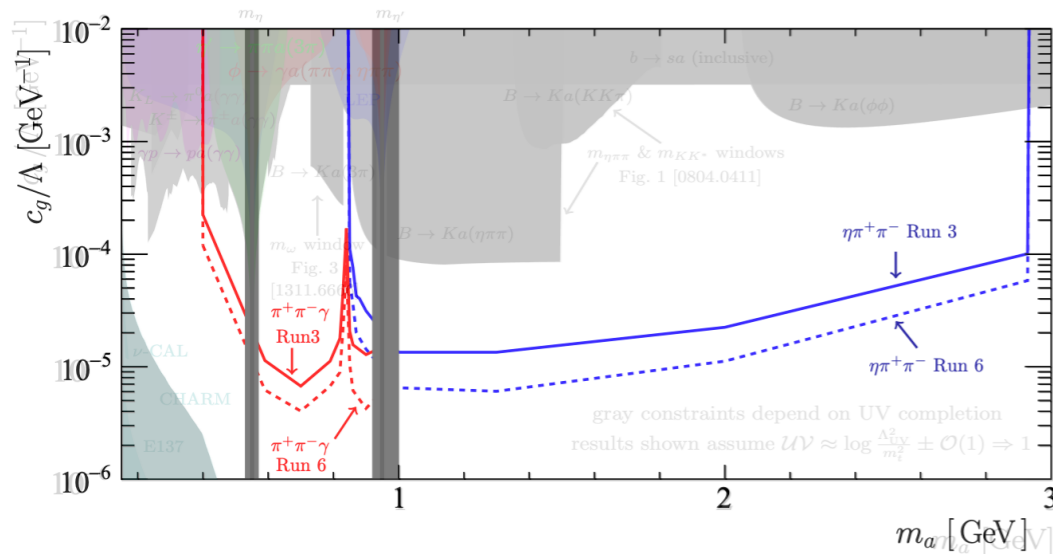
Axion Portal is a recast

- We can translate directly into the axion like portal
 - Governed by one formula $\frac{c_g}{\Lambda} = \frac{g_q}{v}$
 - Assumes Gluon coupling comes from a yukawa loop
 - Also LHC model assumes yukawa coupling(not need)
 - ▶ Photon coupling not considered in this setup
- With the model used by LHC DM WG gluon coupling is a loop



Axion Portal result

- Bounds written in ALP notation are quite strong
 - Relic density bound exists whend mediator mass is higher



LHCDMWG & FIP

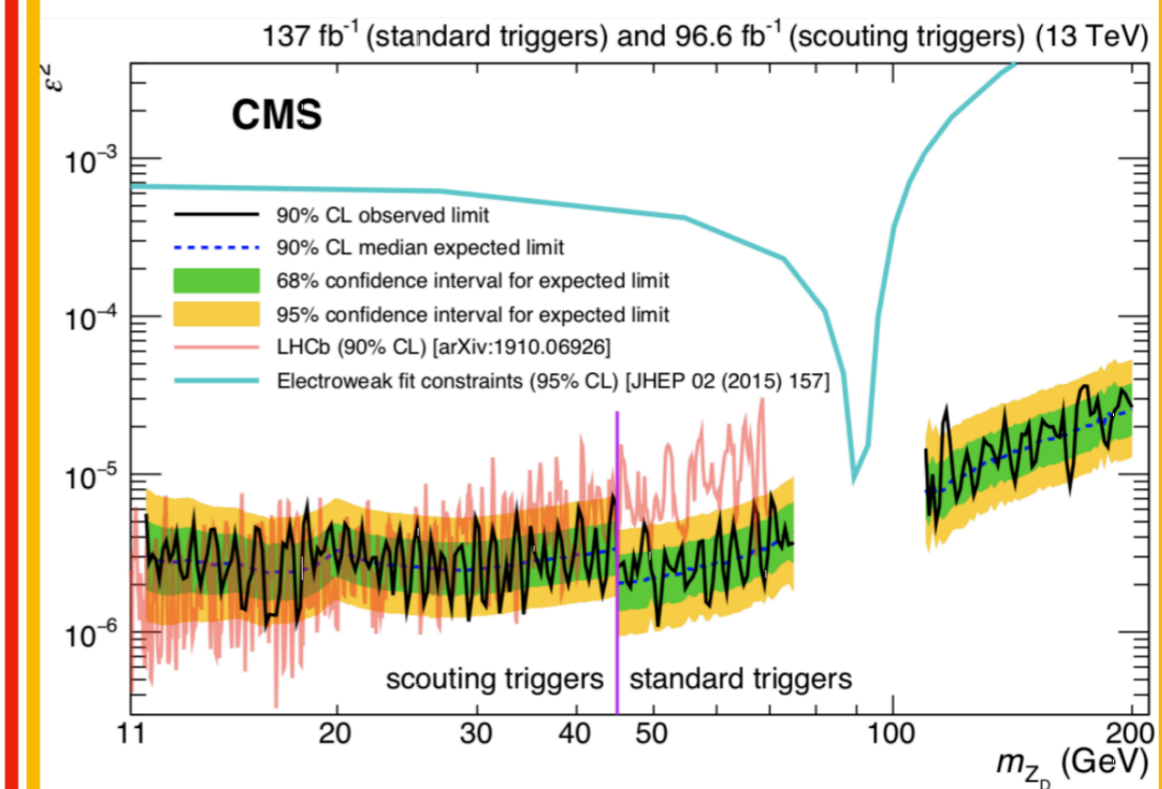
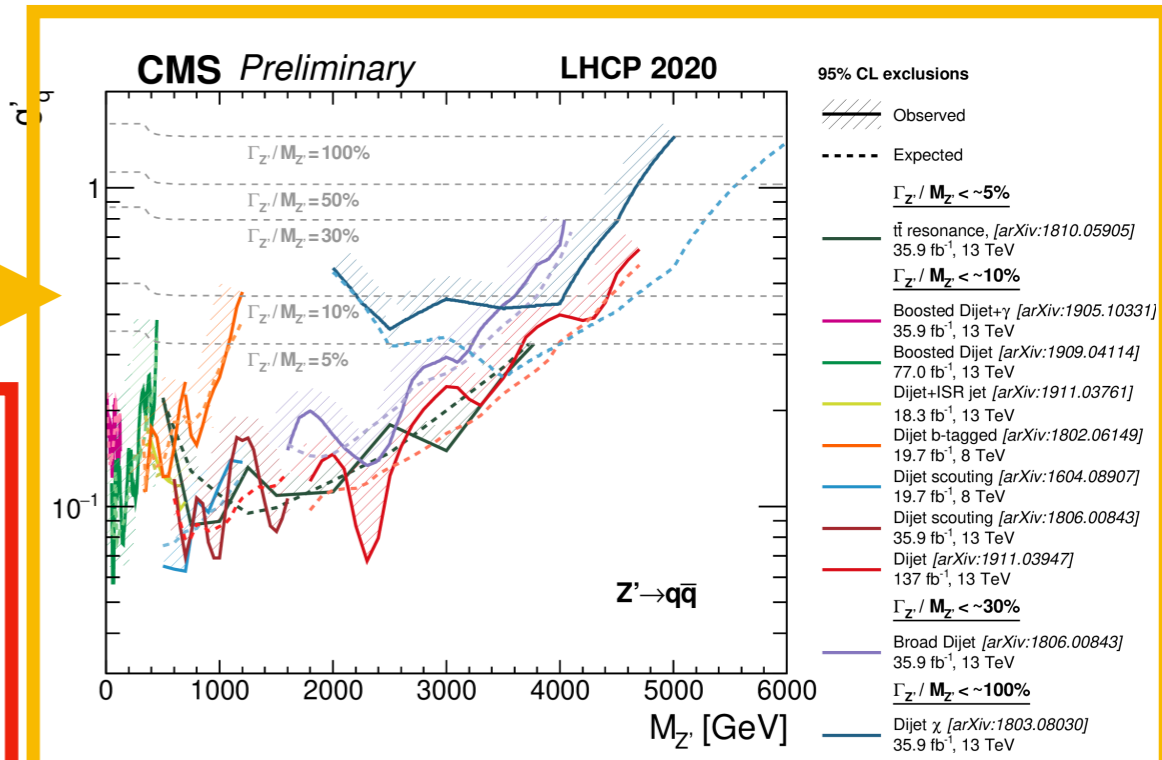
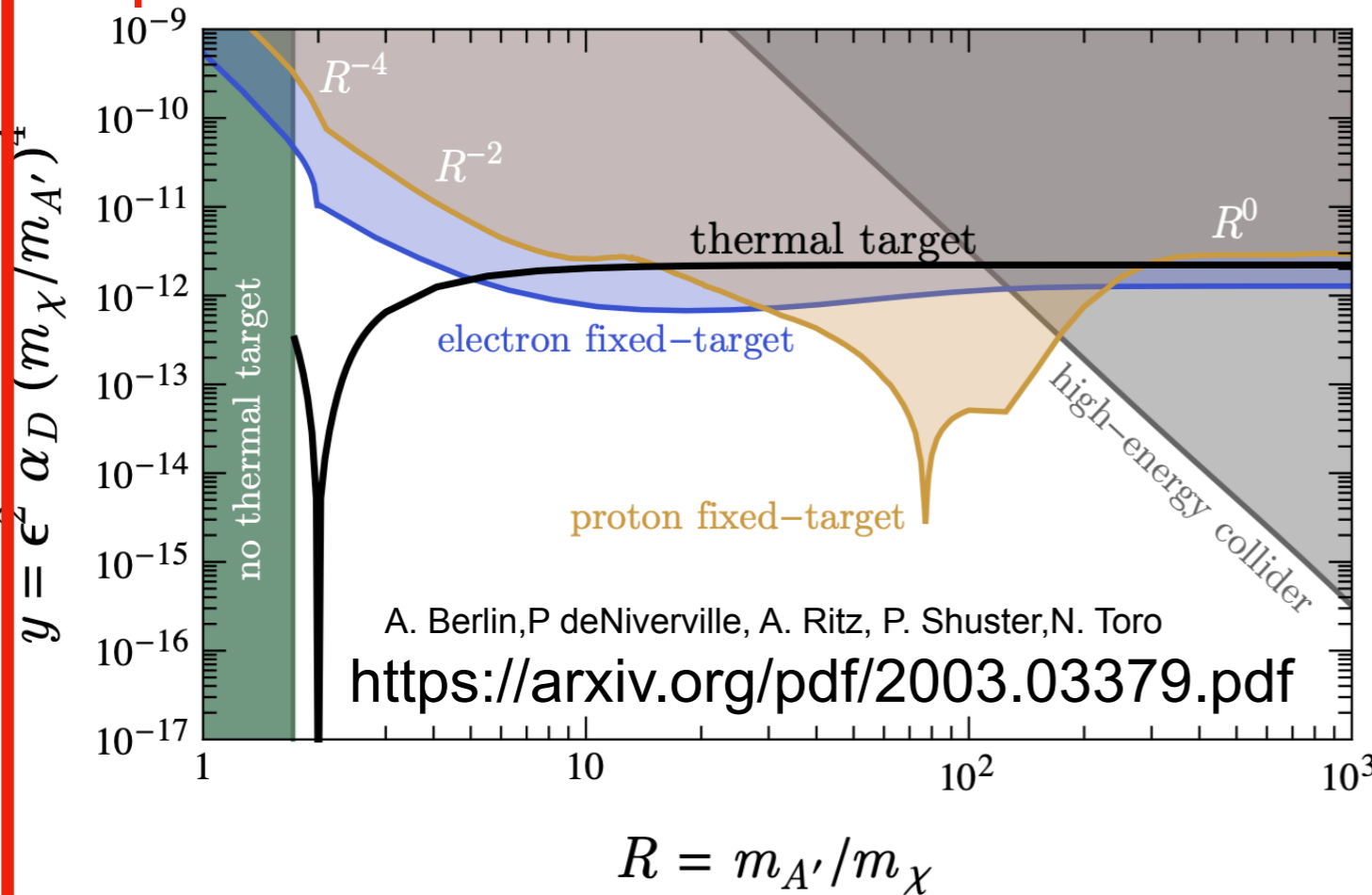
- The LHC is the only collider in town above 10 GeV
 - There is a lot it can say about Dark Matter
 - ▶ Particular in context of Higgs and heavy mediators
 - LHCDMWG is the forum for DM interpretations of the LHC
- Light Dark Sector group focuses on specific models
 - There is a large overlap of these models with LHC DM WG
 - We now have a model to enable Dark Photon Interpretations
 - Reconciled ALP and Dark Higgs Portals
 - ▶ Madgraph models exist for both
 - Part of a greater dark sectors effort underway
- New interpretations/models will motivate new directions at LHC

Other Points to keep in mind

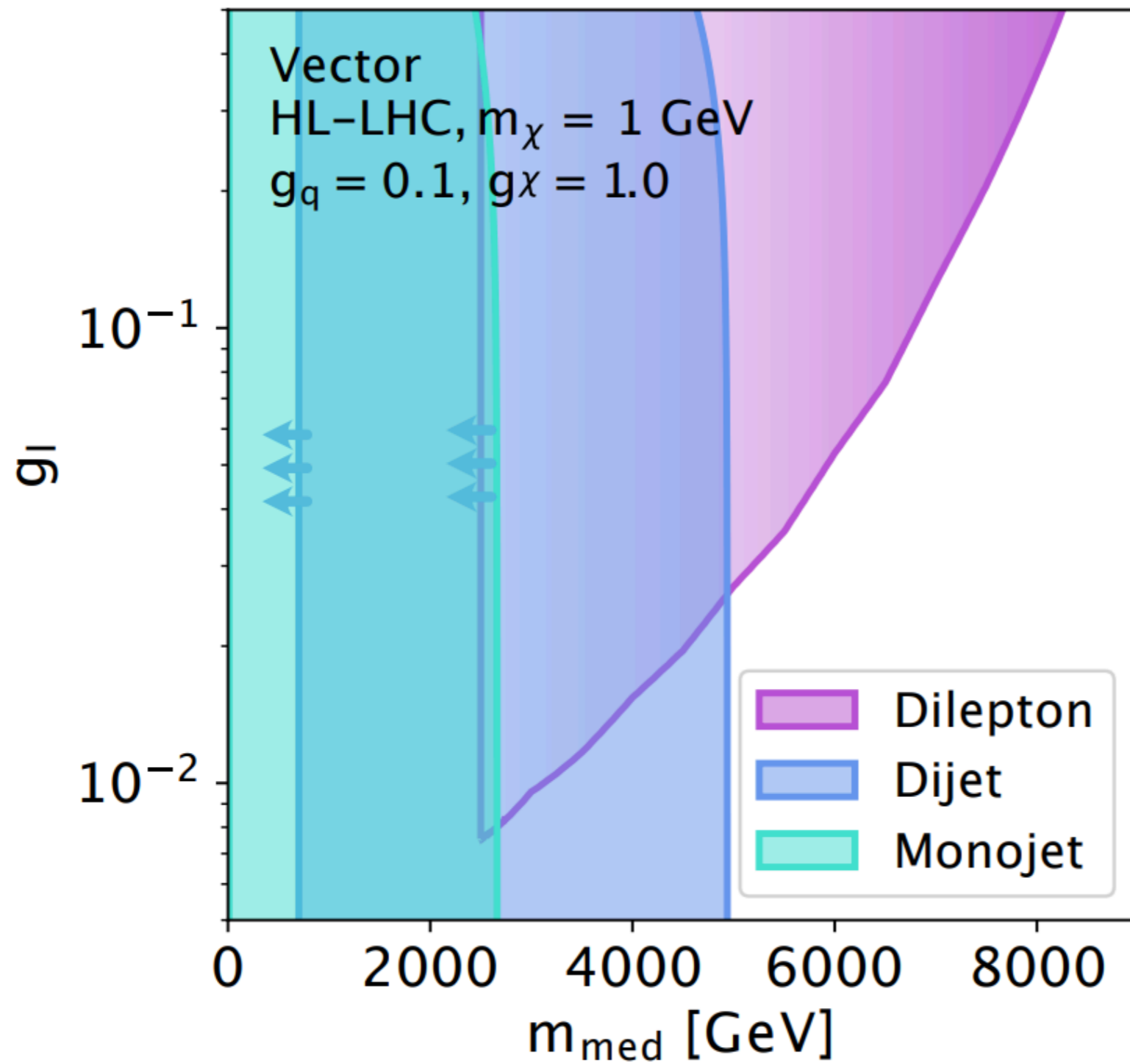
Visible Results for Quark and Lepton final states can be added into the mix



There are other ways to present LHC results on the same plot w/light DM experiments



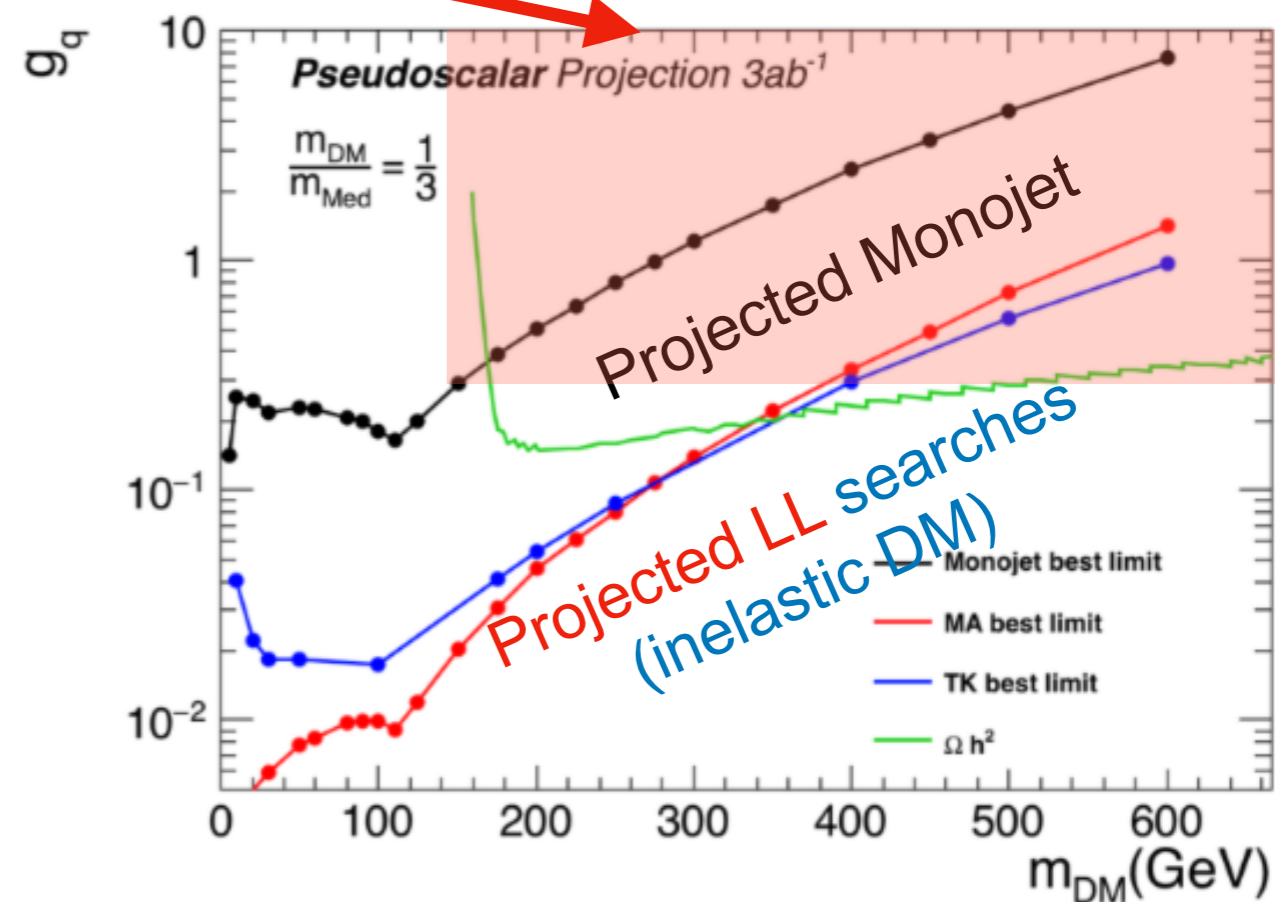
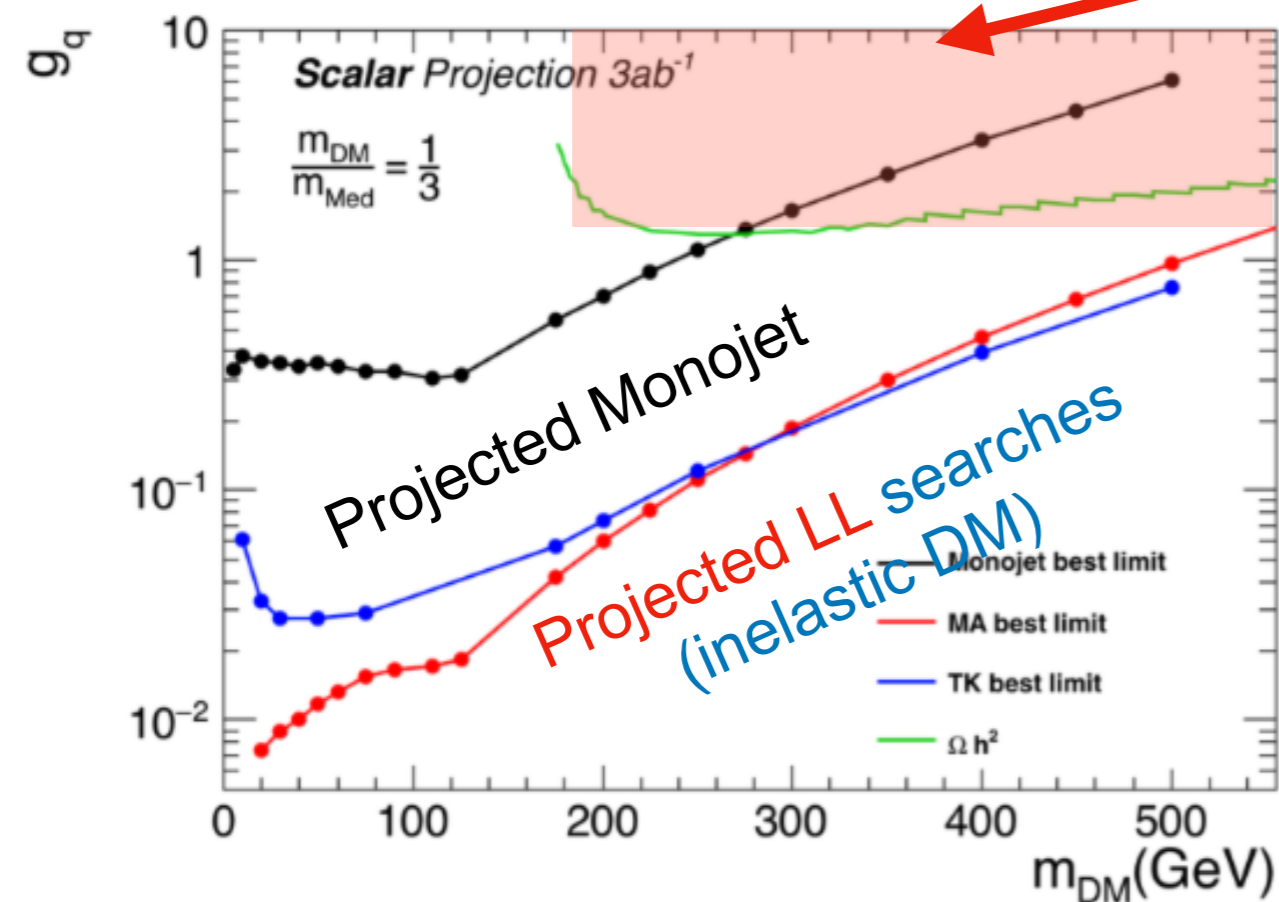
LHC Lepton Projections



Scalar/Pseudoscalar

- Heavy (pseudo)scalar models contend w/ relic bounds
 - Addition of Higgs to invisible also complicates this
 - Its **very hard to have a scalar/ALP without heavier objects**
 - Typically need a 2HDM or Higgs Mixing

Region that would not overclose DM



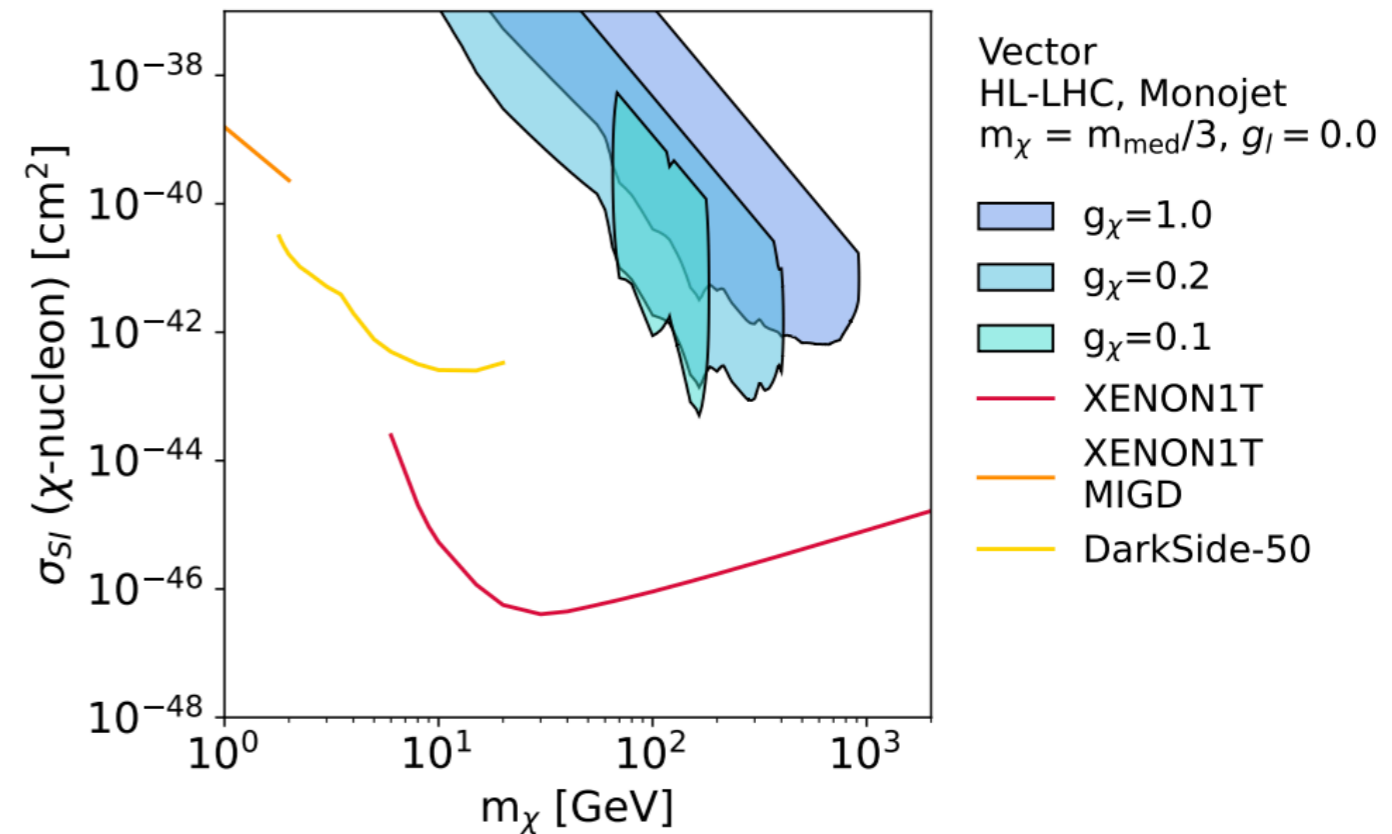
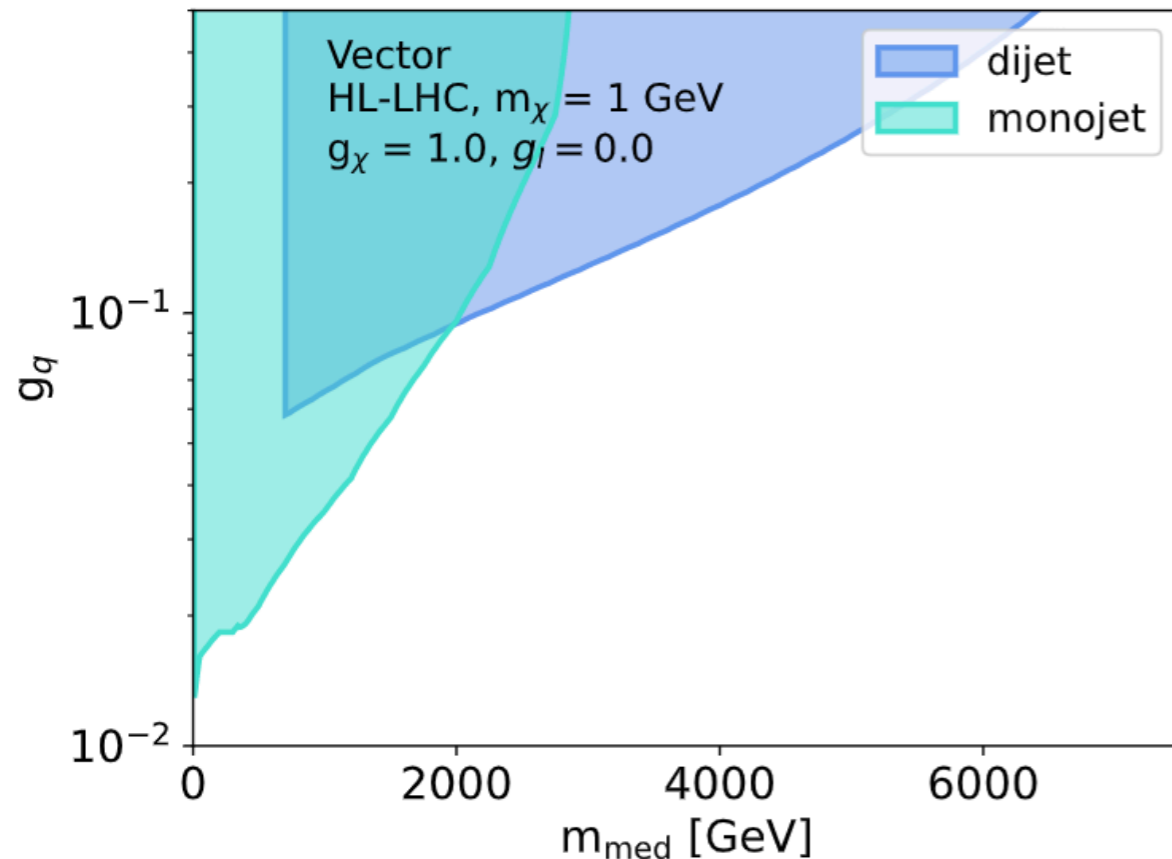
Floating the couplings ¹⁰²

- Floating the couplings gives us a new set of bounds
 - In practice varying couplings doesn't change bounds
 - However to make direct detection bounds coupling fixed
 - Monojet and dijet can probe couplings below $g_q = 0.1$

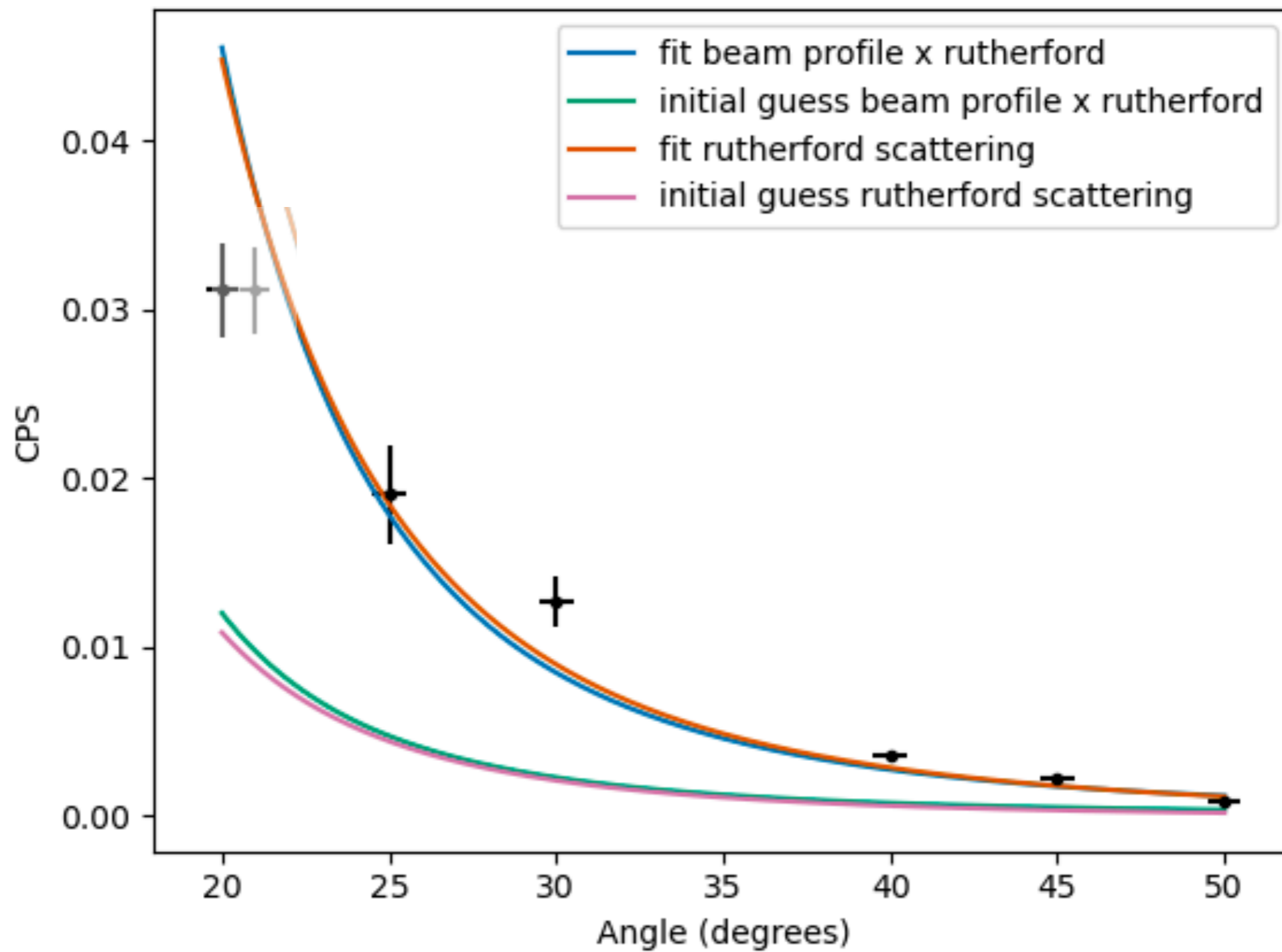
Float

$$\sigma_{\text{SI}} = \frac{f^2(g_q) g_{\text{DM}}^2 \mu_{n\chi}^2}{\pi M_{\text{med}}^4} \text{ Fix}$$

Fixing the Dark Matter Mass



gold scattering :P



LHC Default Models

- LHC has had 4 default models
 - Motivated by standard LHC signatures and comparison with ID/D
 - Additionally had benchmark coupling choices $g_q=0.25$ and $g_{DM}=1.0$

Spin 1

$$\mathcal{L}_{\text{vector}} = -g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q - g_\ell \sum_{\ell=e,\mu,\tau} Z'_\mu \bar{\ell} \gamma^\mu \ell,$$

$$\mathcal{L}_{\text{axial-vector}} = -g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \gamma_5 \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma_5 q - g_\ell \sum_{\ell=e,\mu,\tau} Z'_\mu \bar{\ell} \gamma^\mu \gamma_5 \ell$$

No Interference with the Z boson

Spin 0

$$\mathcal{L}_{\text{scalar}} = -g_{\text{DM}} \phi \bar{\chi} \chi - g_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} q,$$

$$\mathcal{L}_{\text{pseudo-scalar}} = -ig_{\text{DM}} \phi \bar{\chi} \gamma_5 \chi - ig_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} \gamma_5 q,$$

No Interference with the higgs boson

Only quark couplings guaranteed in interpretation

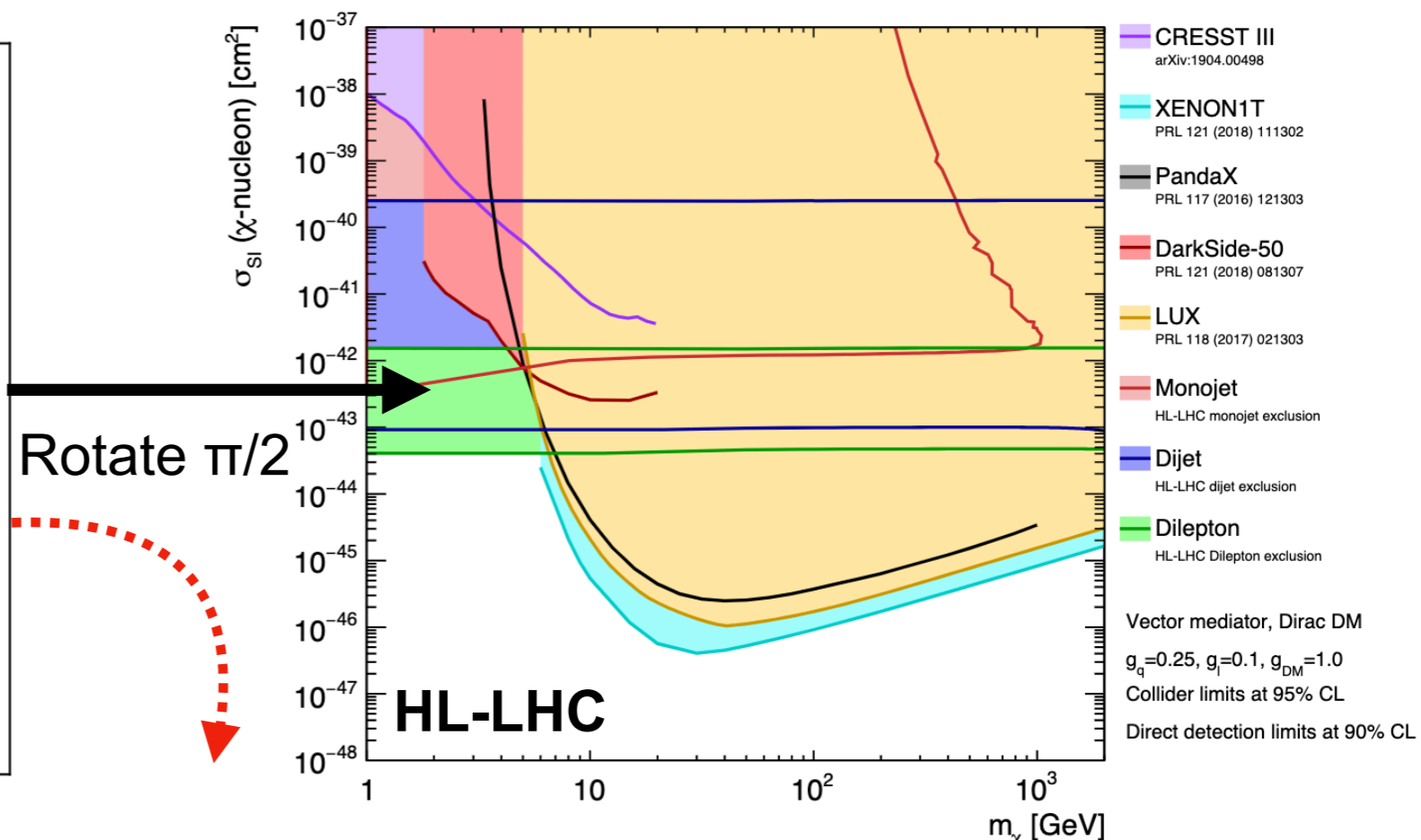
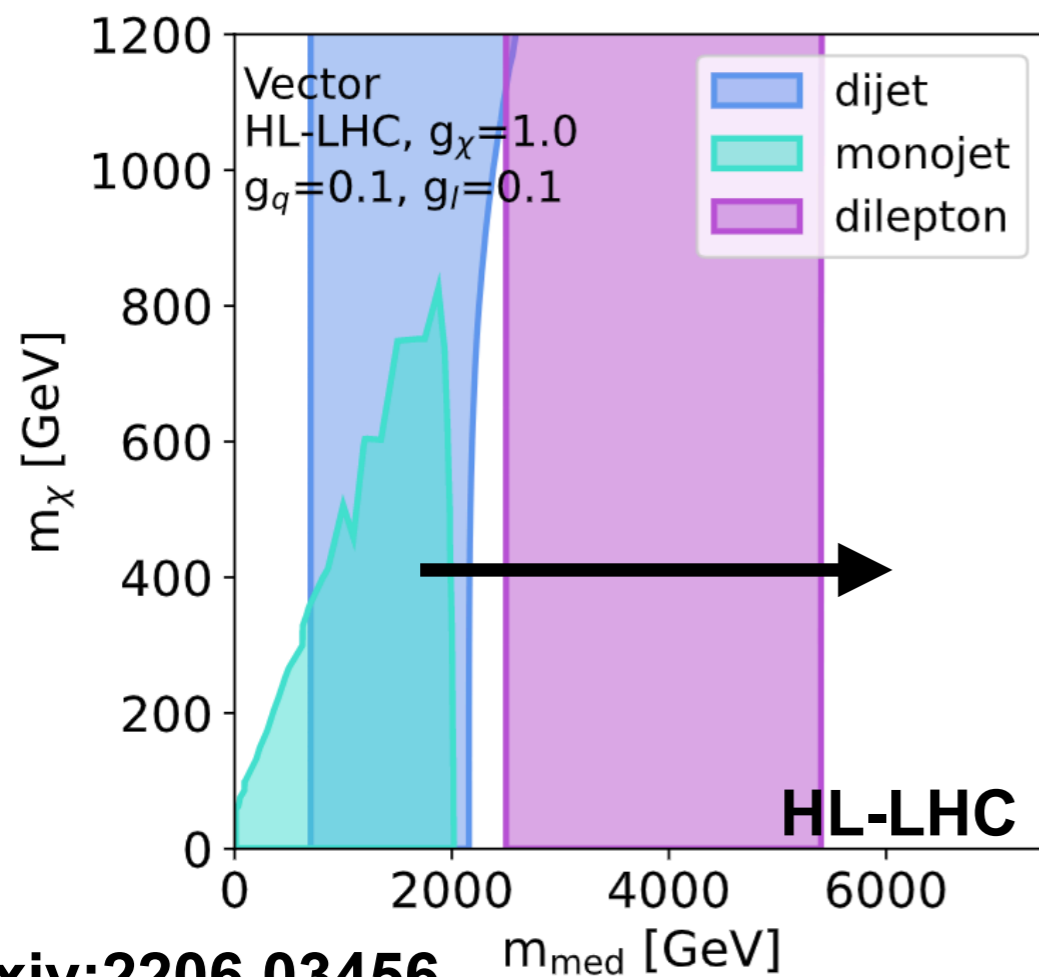
These remain the main ways to interpret DM at LHC

LHC Model Presentation

- Traditionally presented models in mass vs mass plane
 - With fixed couplings
 - Idea was to see how high a mass we could achieve

Wanted to see how high we could go

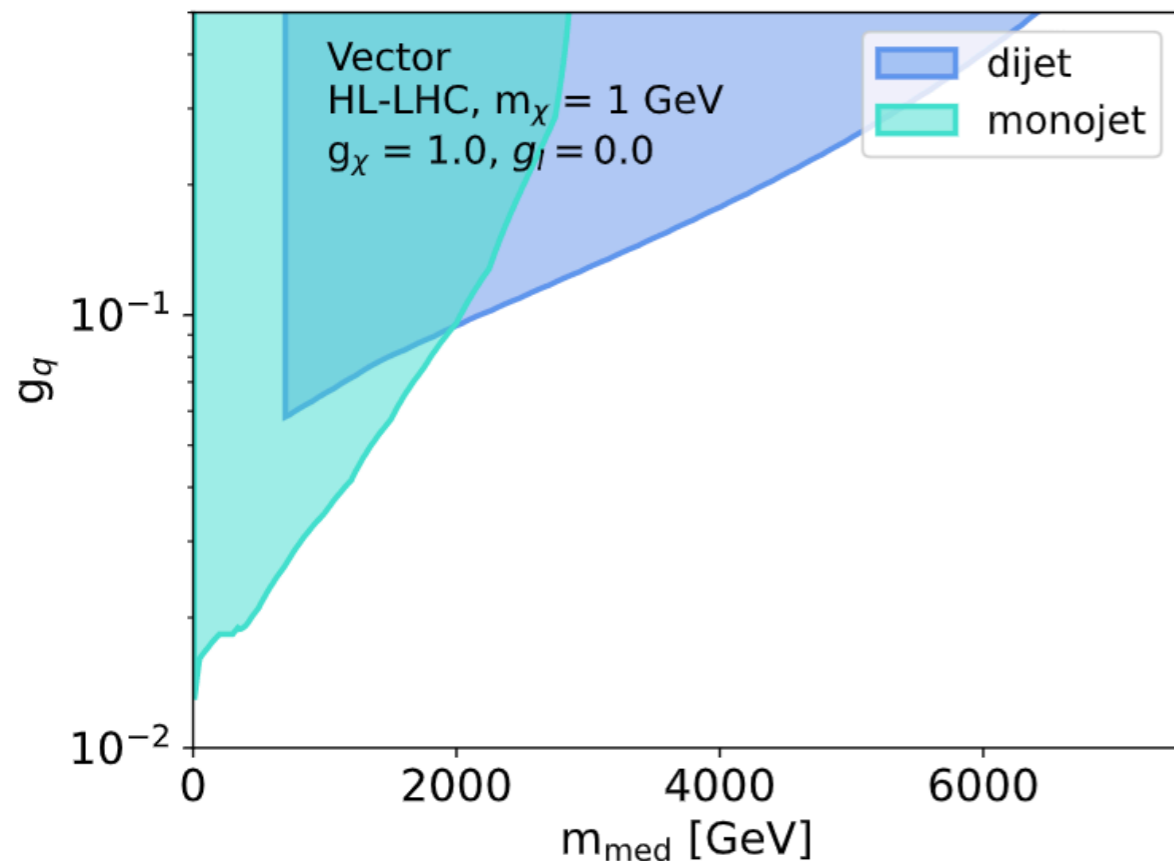
$$\sigma_{\text{SI}} = \frac{\text{Fixed } f^2(g_q)g_{\text{DM}}^2 \mu_{n\chi}^2}{\pi \text{Floated } M_{\text{med}}^4}$$



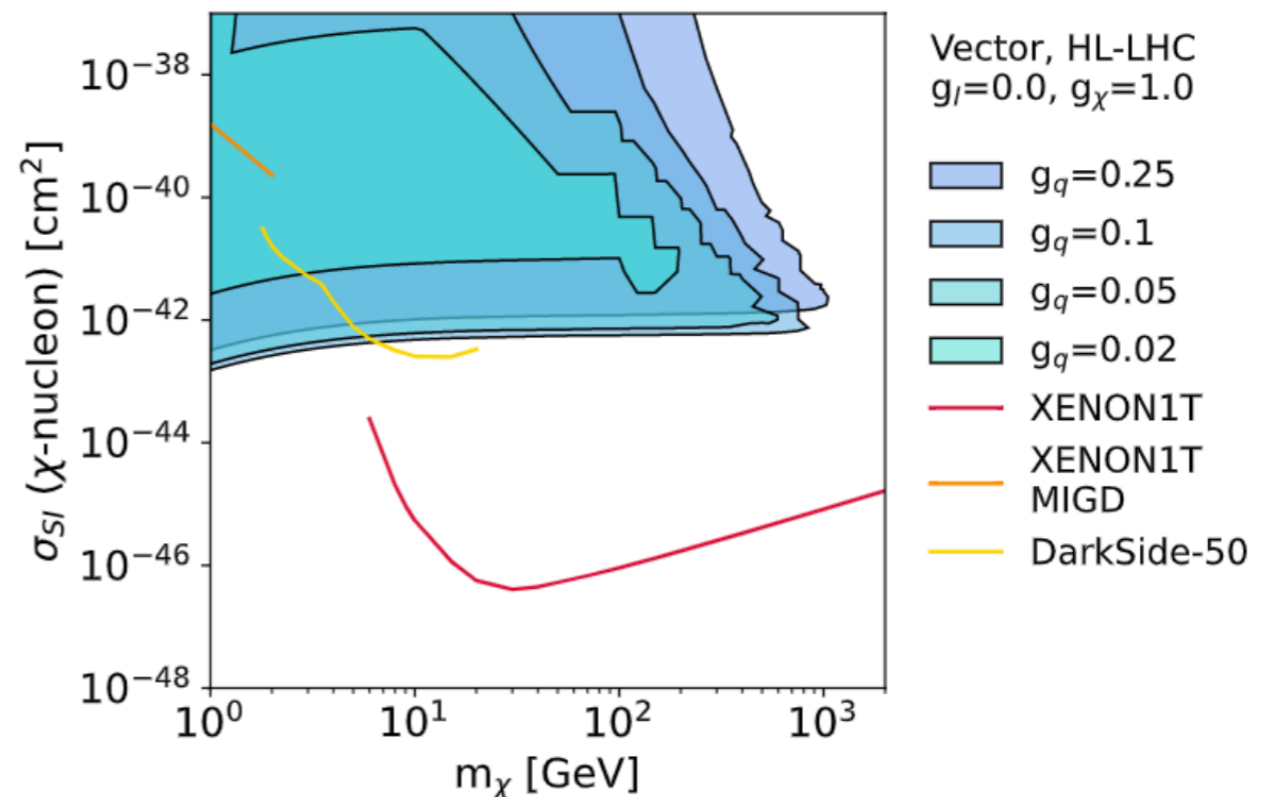
Floating the couplings ¹⁰⁶

- Floating the couplings gives us a new set of bounds
 - In practice **varying couplings doesn't change bounds much**
 - However to make direct detection bounds coupling fixed
 - **Monojet and dijet can probe couplings below $g_q = 0.1$**

Fixing the Dark Matter Mass

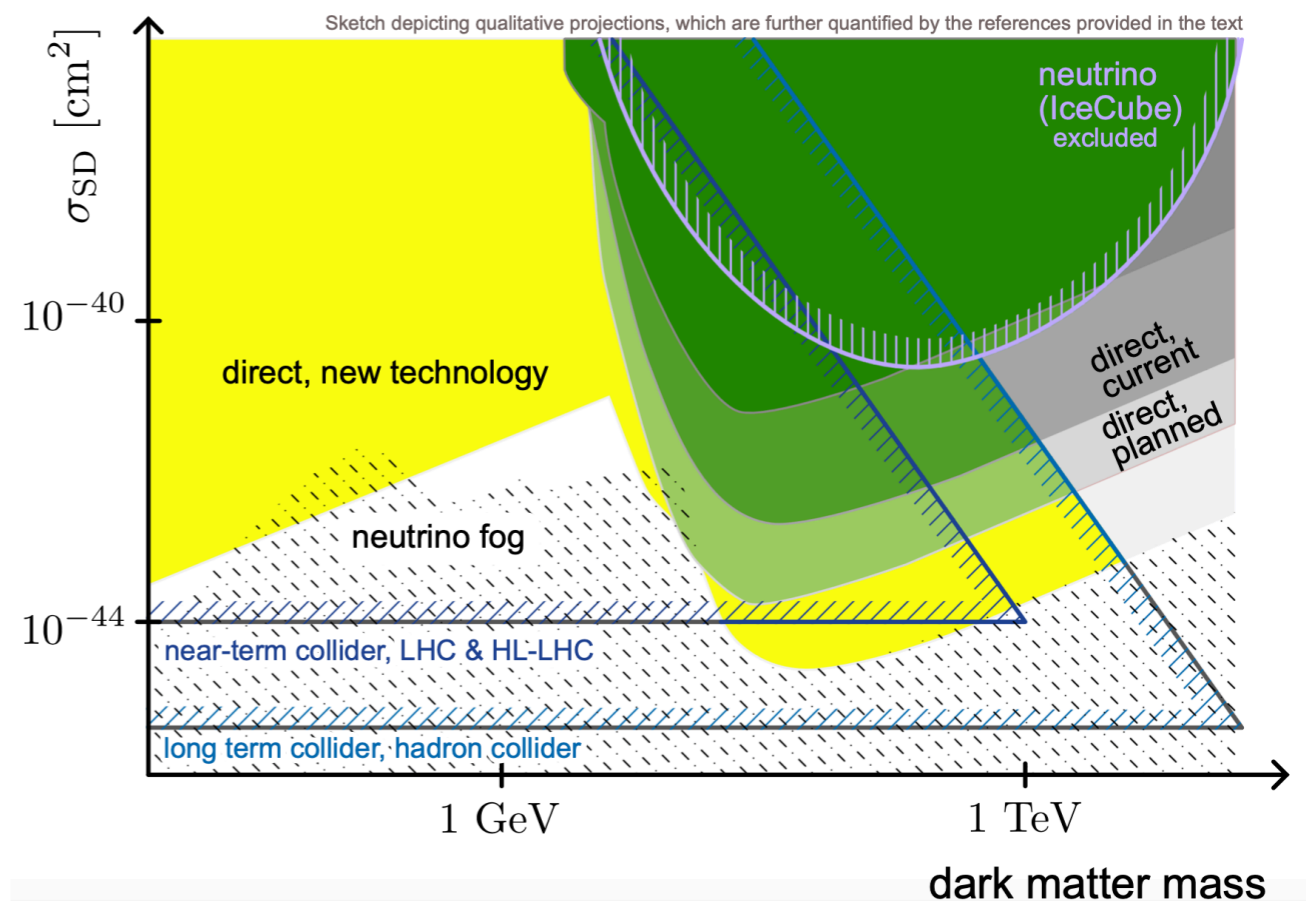
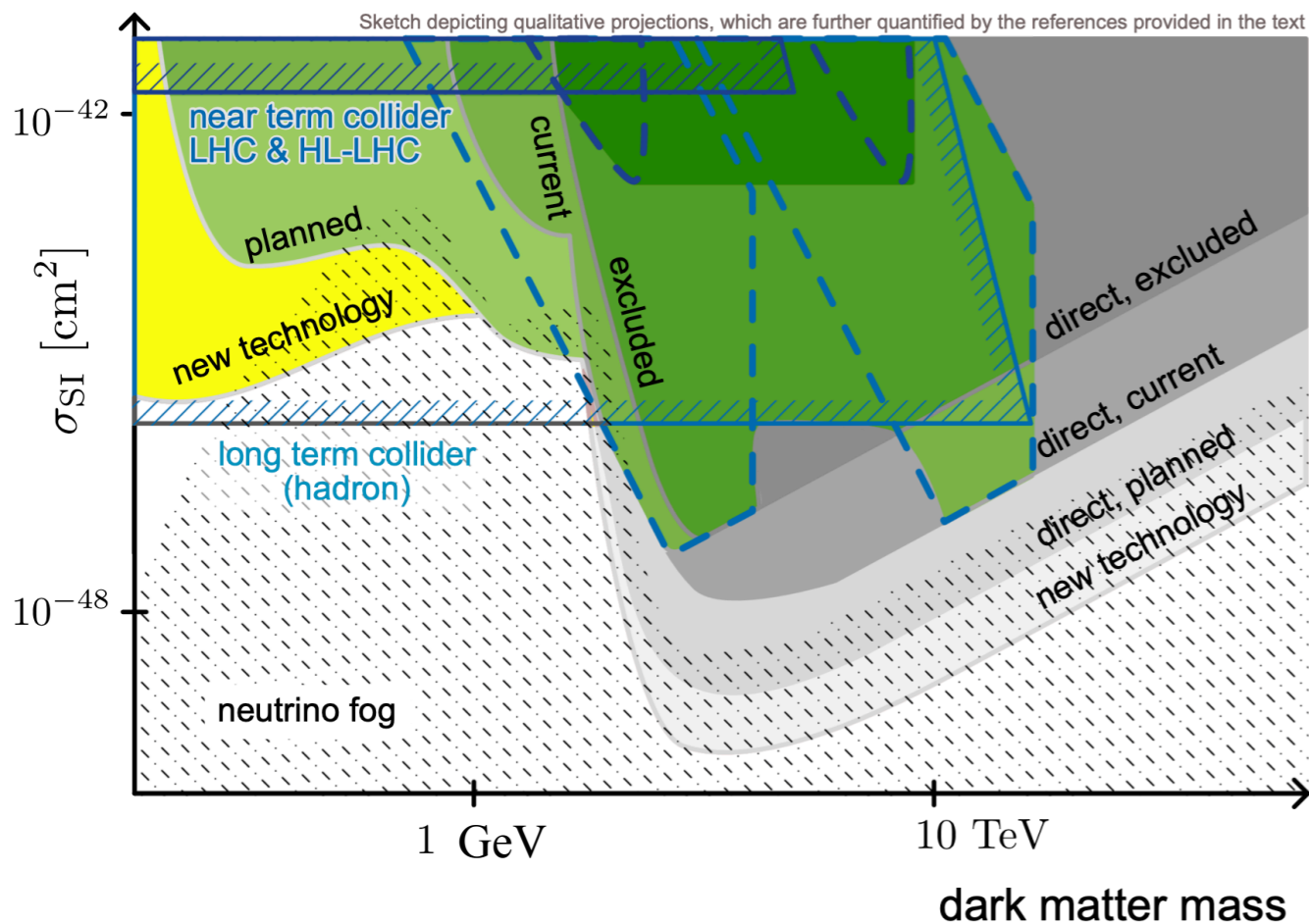


$$\sigma_{\text{SI}} = \frac{f^2(g_q) g_{\text{DM}}^2 \mu_{n\chi}^2}{\pi M_{\text{med}}^4} \quad \text{Float}$$



Ultimate Bounds

- Given these variations we can standardize these



- From high mass invisible studies draw general conclusions
 - Varying coupling bounds doesn't dramatically change LHC
 - The LHC can **provide complementarity to Direct Detection**

Translating to y

