

# Features of Quantum Information at Colliders

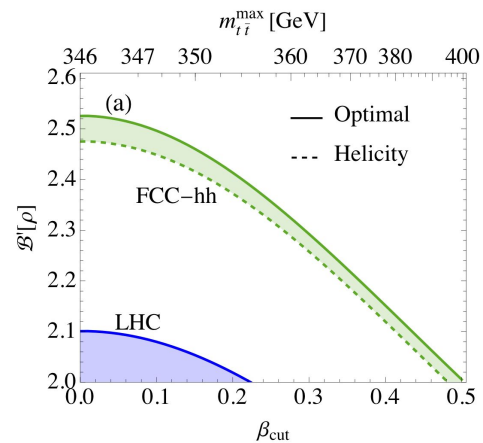
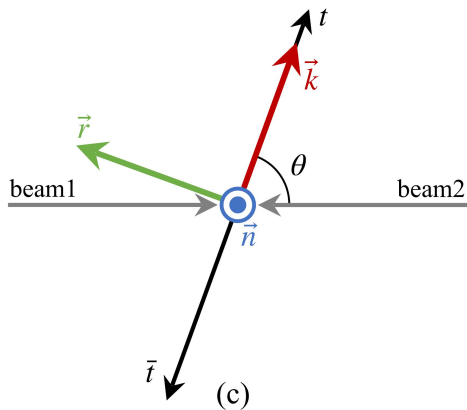
Matthew Low (University of Pittsburgh)

Physics Potential of Future Colliders, TRIUMF, Canada

Main References: 2310.17696, 2311.09166, 2407.01672

*with* Kun Cheng, Tao Han, Arthur Wu

# Outline



Low vs. High  
Energy



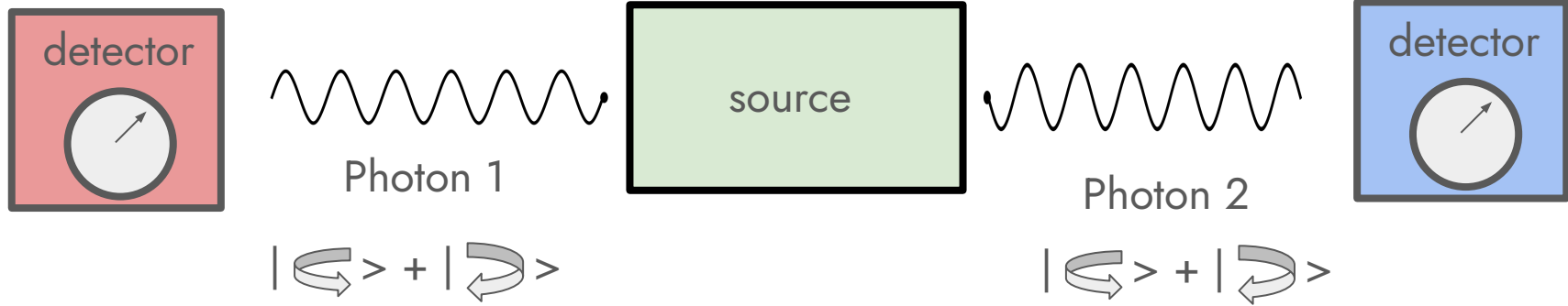
Bell Tests



Future Colliders

# Quantum Information at High Energies

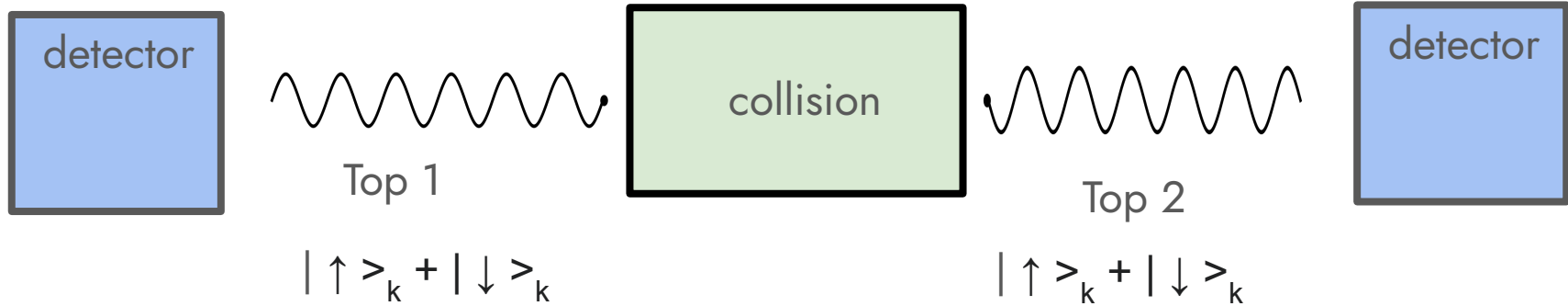
- Recall the traditional quantum experiments with two photons



- A source creates *similarly-prepared* quantum states
- The quantum states includes *quantum correlations* between polarizations
- Detectors choose an *axis*, then detect *left* or *right* polarized

# Quantum Information at High Energies

- The set-up at colliders is a bit different



- Collisions create “similarly-prepared” quantum states
- The quantum states involves *quantum correlations* between *spins*
- Use decay products of top to *infer* spin

# Quantum Information at High Energies

- Consider the density matrix for two qubits

$$\rho = \frac{1}{4} \left( \mathbb{I}_4 + \sum_i B_i^+ \sigma_i \otimes \mathbb{I}_2 + \sum_j B_j^- \mathbb{I}_2 \otimes \sigma_j + \sum_{ij} C_{ij} \sigma_i \otimes \sigma_j \right)$$

- Polarization of first top (3 DOF):  $B_i^+$
- Polarization of second top (3 DOF):  $B_j^-$
- Spin correlations (9 DOF):  $C_{ij}$ 
  - Calculated before interest in entanglement

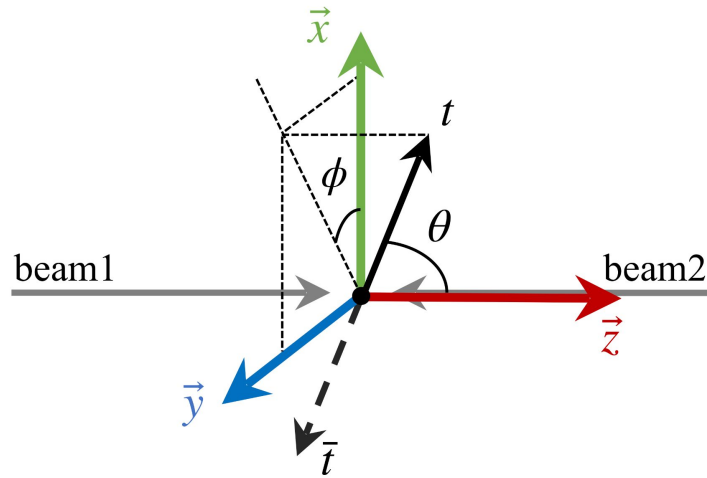
$$C_{33} = \begin{cases} -0.456 & (-0.389) & \text{Helicity at Tevatron} \\ +0.910 & (+0.806) & \text{Beamline at Tevatron} \\ +0.918 & (+0.913) & \text{Off - Diagonal at Tevatron} \\ +0.305 & (+0.311) & \text{Helicity at LHC(14 TeV),} \end{cases}$$

Parke [1202.2345](#)

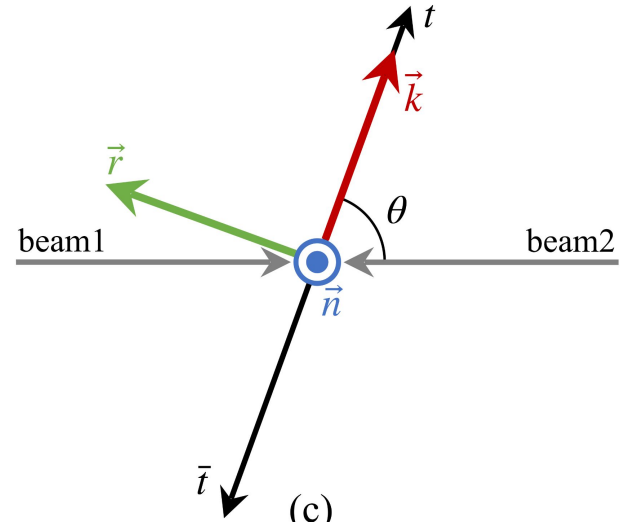
# Quantum Information at High Energies

- Beamline Basis (x,y,z)

Helicity basis (k,r,n)



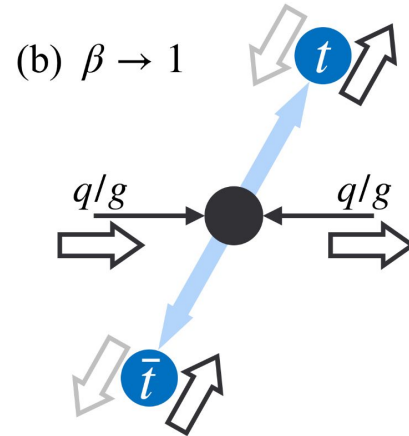
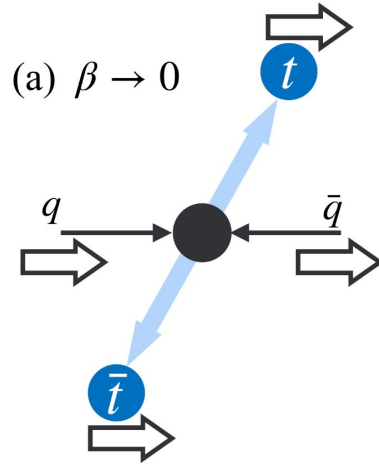
(a)



(c)

# Quantum Information at High Energies

- Example:  $q\bar{q} \rightarrow t\bar{t}$



- Spin correlations intuitively *depend* on basis choice
- Quantum states *cannot depend* on basis choice

# Quantum Information at High Energies

- Consider reconstructing the density matrix from collider events
  - If  $\sigma_i \otimes \sigma_j$  is the same for each event

$$C_{ij} = \text{tr}(\rho(\sigma_i \otimes \sigma_j))$$

- If  $\sigma_i \otimes \sigma_j$  is the different for each event

$$\langle C_{ij} \rangle_a = \text{tr}(\rho(\sigma_i \otimes \sigma_j)_a)$$

- Rather than quantum states, at colliders we reconstruct “**fictitious states**”
  - Entangled fictitious state => entanglement (but numerical value not meaningful)
  - Bell non-local fictitious state => Bell non-locality (but numerical value not meaningful)

Cheng, Han, ML [2311.09166](#)

Cheng, Han, ML [2407.01672](#)



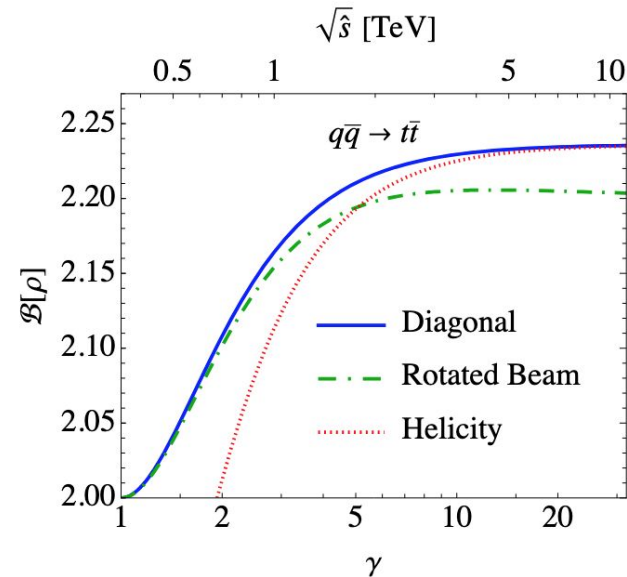
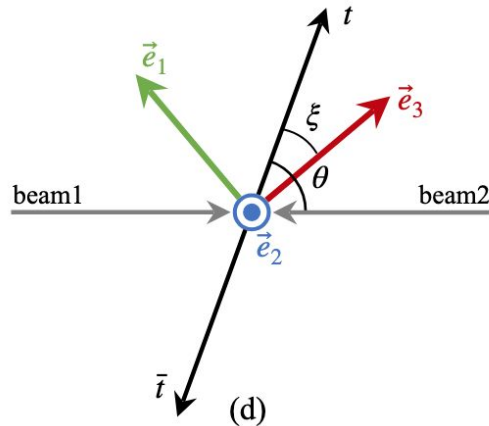
# Quantum Information at High Energies

- Fictitious states *depend* on the spin quantization axis
- There is an *optimal* direction to maximize:
  - Entanglement
  - Bell inequality violation

Cheng, Han, ML [2311.09166](#)

Cheng, Han, ML [2407.01672](#)

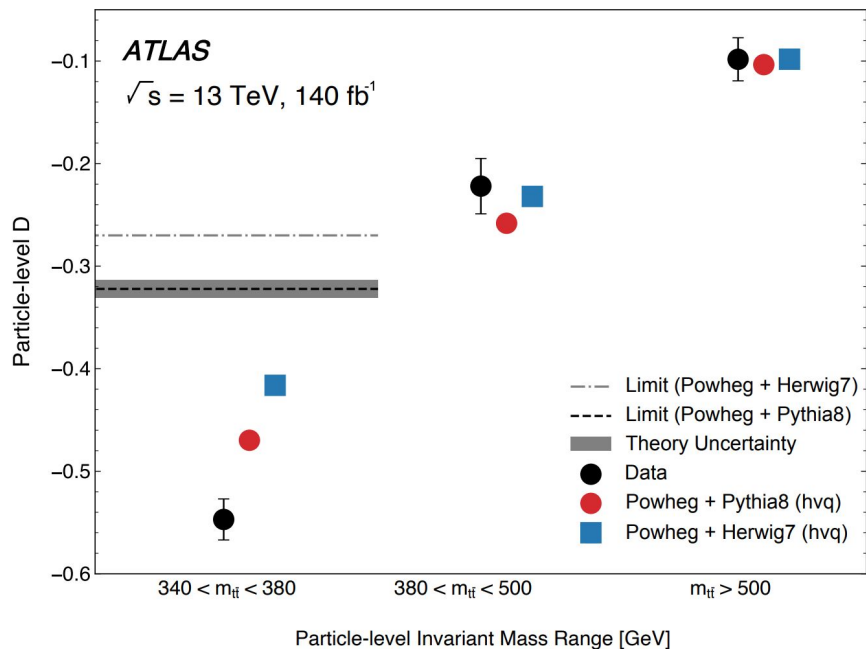
- Rotate by  $\xi$  from helicity basis



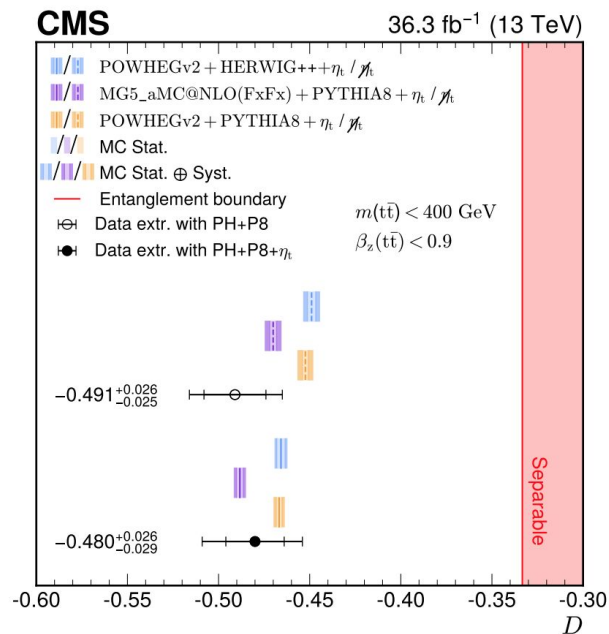
# Quantum Information at High Energies

- Entanglement already detected in leptonic top pair events

ATLAS [2311.07288](#)



CMS [2406.03976](#)



# Bell Tests

- Inequality that is satisfied by all local hidden variable theories
  - Bell (1964)
  - Clauser, Horne, Shimony, Holt (1969)
- “for experiments with entangled photons, establishing the violation of Bell inequalities and pioneering quantum information science.”
- For two qubits, the CHSH inequality is the Bell inequality

$$|\langle A_1 B_1 \rangle - \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle + \langle A_2 B_2 \rangle| \leq 2$$

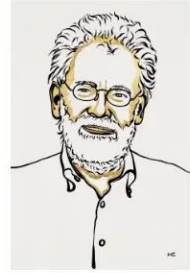
## The Nobel Prize in Physics 2022



Ill. Niklas Elmehed © Nobel Prize Outreach  
Alain Aspect  
Prize share: 1/3



Ill. Niklas Elmehed © Nobel Prize Outreach  
John F. Clauser  
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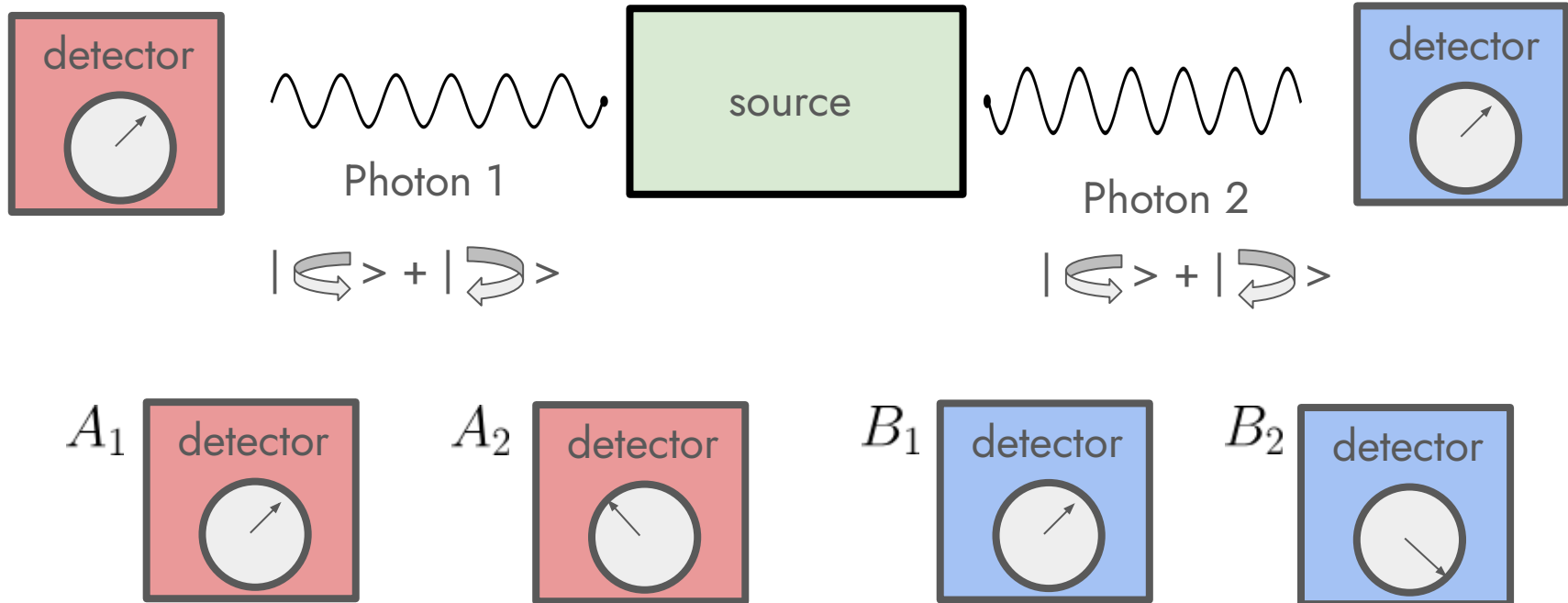


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Anton Zeilinger  
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# Bell Tests

- CHSH inequality

$$|\langle A_1 B_1 \rangle - \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle + \langle A_2 B_2 \rangle| \leq 2$$



# Bell Tests: example

$$|\langle A_1 B_1 \rangle - \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle + \langle A_2 B_2 \rangle| \leq 2$$

- Example:  $|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$

- Detector settings:  $A_1 = \sigma_x$      $B_1 = \frac{-1}{\sqrt{2}}(\sigma_x + \sigma_z)$   
 $A_2 = \sigma_z$      $B_2 = \frac{1}{\sqrt{2}}(\sigma_x - \sigma_z)$

- Result:    LHS =  $2\sqrt{2}$

- A different choice of detector settings may not violate the inequality

$$A_1 = \sigma_x \quad B_1 = \frac{1}{\sqrt{2}}(\sigma_x + \sigma_z)$$
$$A_2 = \sigma_z \quad B_2 = \frac{1}{\sqrt{2}}(\sigma_x - \sigma_z)$$

$$\text{LHS} = 0$$

# Bell Tests

- Bell's inequality (CHSH inequality)

$$|\langle A_1 B_1 \rangle - \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle + \langle A_2 B_2 \rangle| \leq 2$$

- All local theories, even with hidden variables, obey inequality
  - A quantum state may or may not violate CHSH
  - QM allows violation because spins *anti-commute*
- 
- At colliders, we don't measure spins *directly*
    - Only measure *momenta* of decayed particles
    - *Infer* spin value, assuming spin properties
    - Cannot "*test*" quantum mechanics at colliders

Abel, Dittmar, Dreiner 1992

# Bell Tests

- Bell's inequality (CHSH inequality)

$$|\langle A_1 B_1 \rangle - \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle + \langle A_2 B_2 \rangle| \leq 2$$

- At colliders, we don't compare hidden variable vs. QM
- Instead, we compare Bell-local QM with Bell-non-local QM
- At a future collider, we can imagine detectors with different capabilities
  - Partially polarized calorimeters?
  - Stern-Gerlach extensions?

Han, ML, Wu [2310.17696](#)

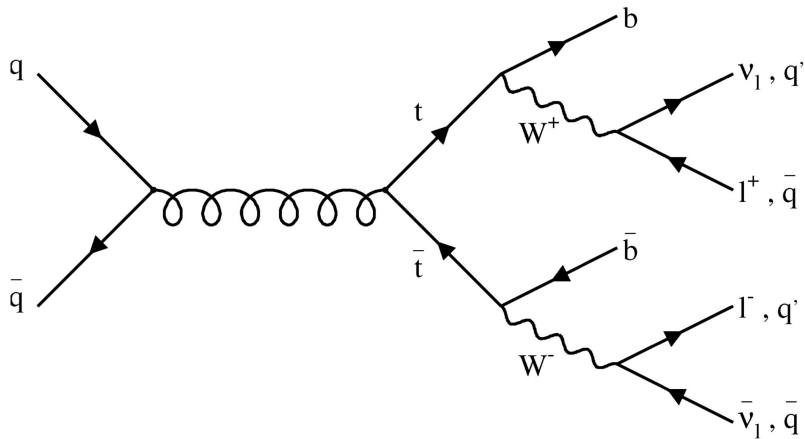
Abel, Dittmar, Dreiner 1992

We now suggest a modification of the ppCC experiment which might be a test of Bell's inequality. Through a conceptually simple change the ppCC experiment can be improved. Suppose one replaces the carbon-12 analyzers by polarized hydrogen ( $^1\text{H}$ ) targets and instead of measuring the scattering distribution, measure the total scattering cross-section or transmission rate. These targets then act as spin-filters, preferentially passing protons, whose spin are aligned with the target polarization direction <sup>#5</sup>. Then

# Future Colliders

- Top pair production
  - The top is a qubit and the anti-top

(See talk by Dorival Gonçalves)



$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_v} = \frac{1}{2} \left( 1 + |\vec{B}| \kappa_v \cos \theta_v \right)$$

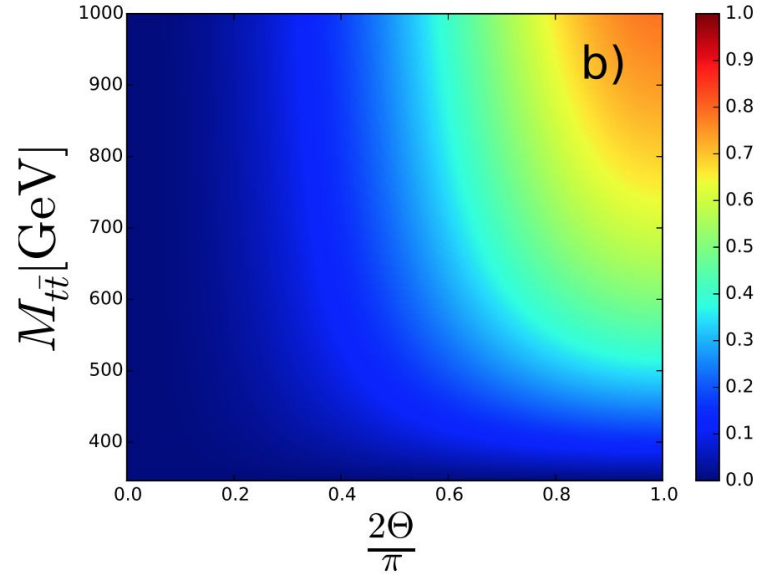
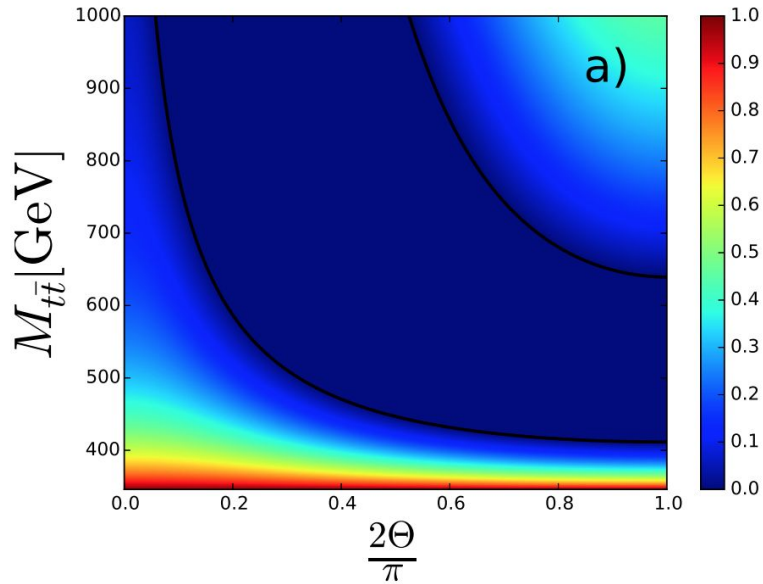
↑  
Spin analyzing power

- For the leptonic decay ( $bb\ell\nu\ell\nu$ ), the lepton spin analyzing power is 1.0
- For the semi-leptonic decay ( $bb\ell\nu qq$ ), the hadronic spin analyzing power is 0.6



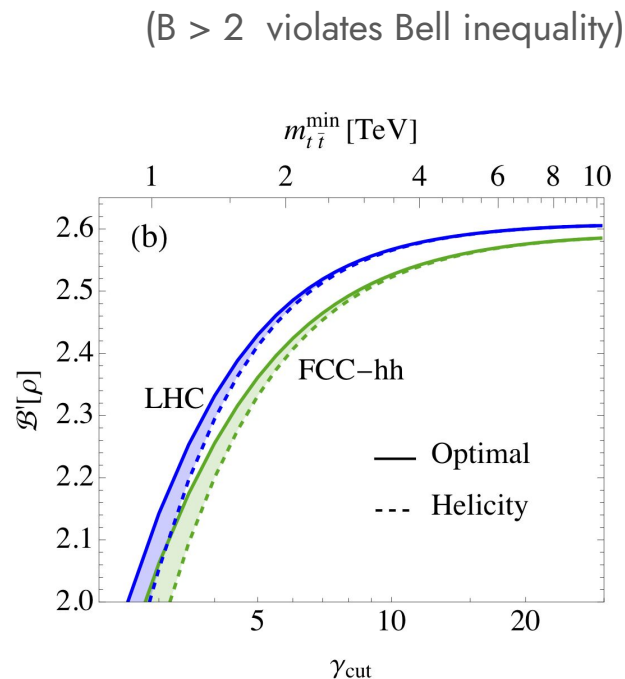
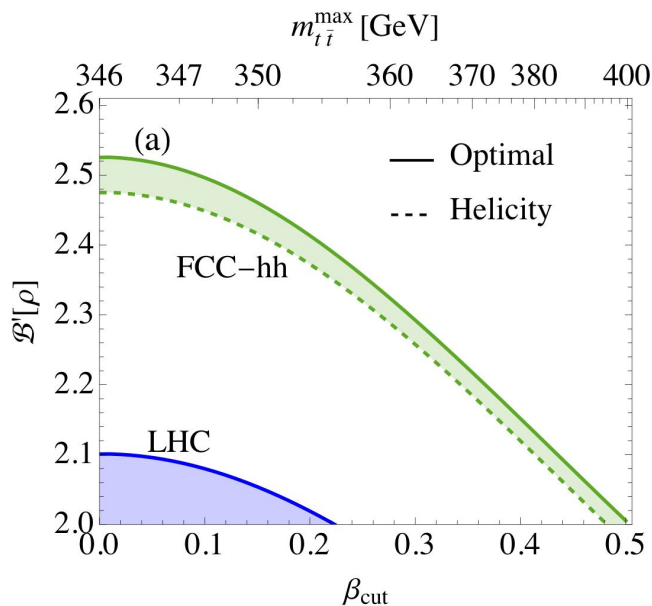
# Future Colliders

- Top pair production
  - The top is a qubit and the anti-top
  - Density matrix is a mixed state of qq-initiated and gg-initiated



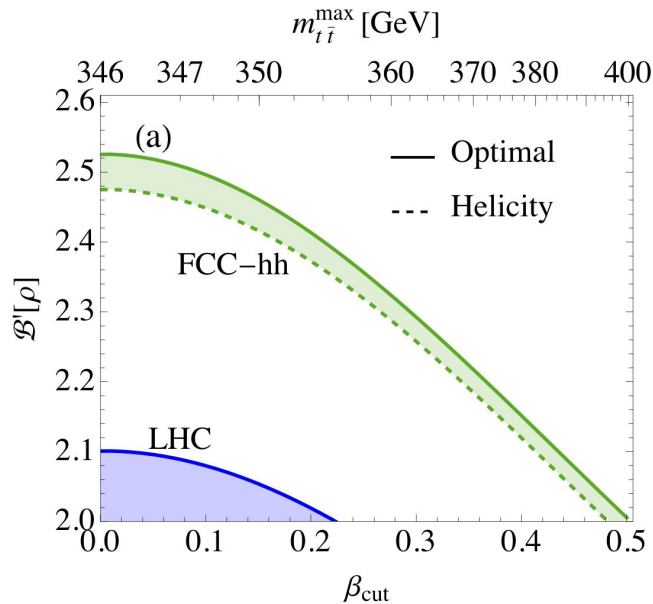
# Future Colliders

- Top pair production
  - FCC-hh (100 TeV) has a different gg/qq fraction
  - Increases signal at threshold
  - Decreases signal at high- $p_T$

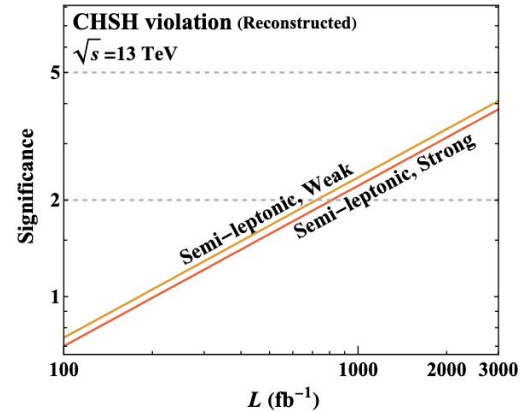


# Future Colliders

- Top pair production
  - FCC-hh (100 TeV) has a different gg/qq fraction
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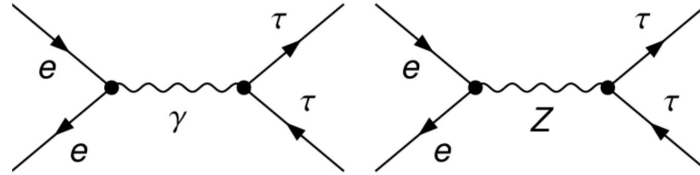


Prediction at HL-LHC



# Future Colliders

- Tau pair production
  - FCC-ee would be sensitive to tau pair production



- Electroweak production leads to non-zero polarization

$$C = \begin{pmatrix} 0.4878 & 0 & 0 \\ 0 & -0.4878 & 0.0011 \\ 0 & 0.0011 & 1 \end{pmatrix} \quad B^+ = B^- = \begin{pmatrix} 0 \\ 0.0001 \\ 0.2194 \end{pmatrix}$$

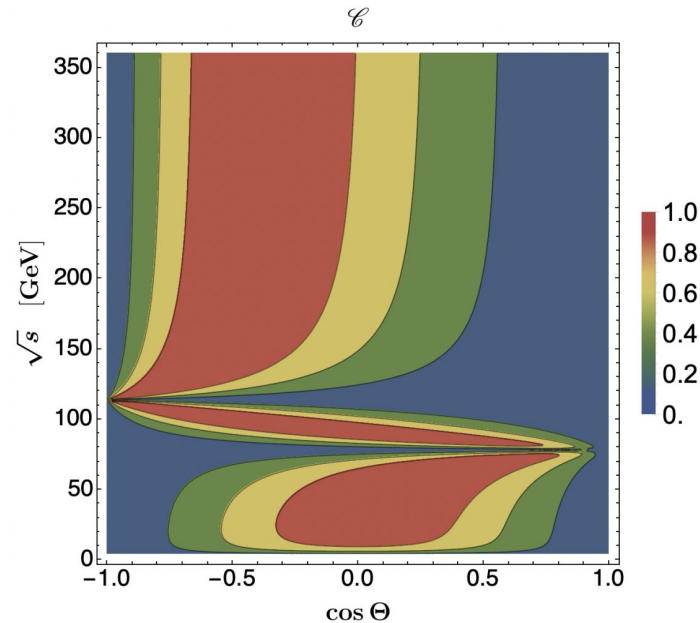
- Ideal decay is  $\tau \rightarrow \pi \nu$

Decay Product( $\tau^-$ )	Branch(%)	Spin Analyzing Power
$\nu_\tau \pi^-$	$10.82 \pm 0.05$	$1.000 \pm 0.005$
$\nu_\tau \pi^- \pi^+ \pi^-$	$9.31 \pm 0.05$	$-0.148 \pm 0.006(\pi^+)$ $-0.038 \pm 0.005(\pi^-)$
$\nu_\tau \mu^- \bar{\nu}_\mu$	$17.39 \pm 0.04$	$-0.341 \pm 0.005$
$\nu_\tau e^- \bar{\nu}_e$	$17.82 \pm 0.04$	$-0.336 \pm 0.005$

# Future Colliders

Fabbrichesi, Marzola [2405.09201](#)

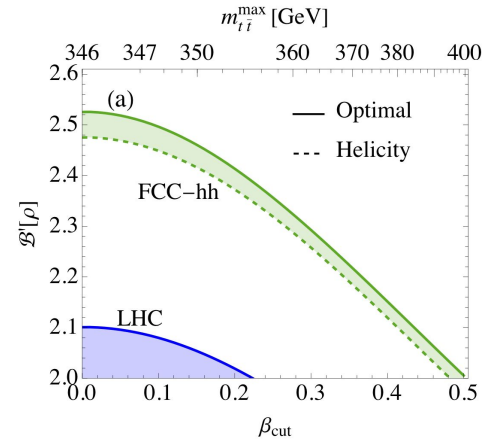
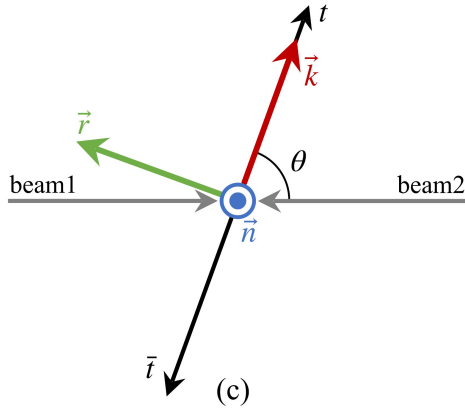
- Tau pair production
  - The entanglement as a function of scattering angle and collision energy



- Entangled ( $>0$ )
- Separable ( $=0$ )

- Statistical significance  $\gg 5\sigma$  for entanglement and for Bell inequality violation

# Summary



Low vs. High  
Energy



Bell Tests



Future Colliders