

# The Status of the Nuclear Shell Model

Still kicking after 70 years

Ragnar Stroberg

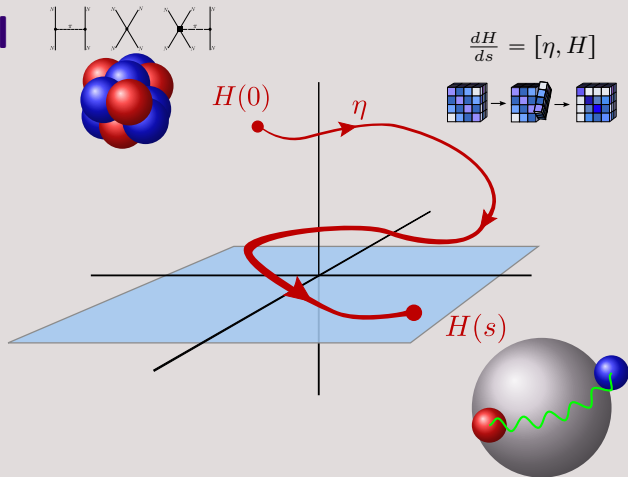
University of Washington

Theory Canada 14

University of British Columbia

Vancouver, BC

May 31, 2019

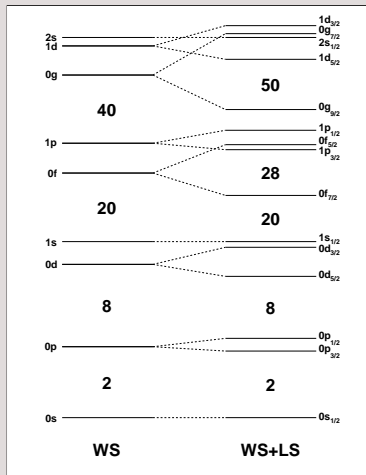
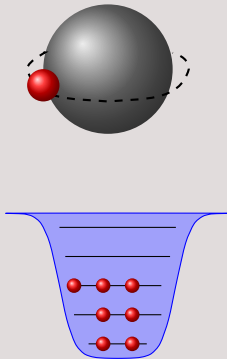
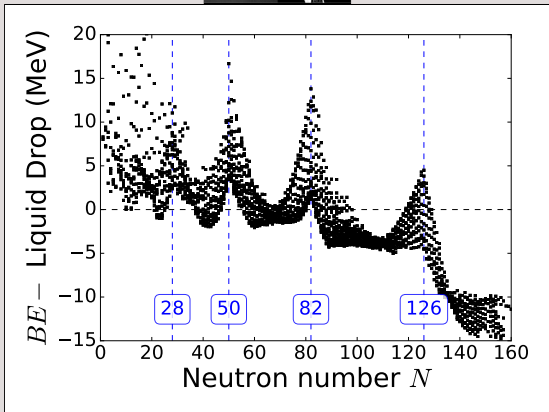


## 🌐 Background

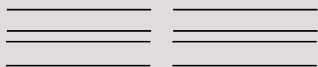
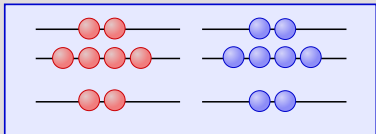
- The Mayer-Jensen shell model
- The interacting shell model
- Attempts to derive the residual interaction
- Phenomenology

## 🌐 Recent progress

- Forces from chiral EFT (with consistent 3N)
- Similarity renormalization group
- Valence-space in-medium SRG
- Some results and future work



Mayer 1948; Haxel, Jensen, and Suess 1949

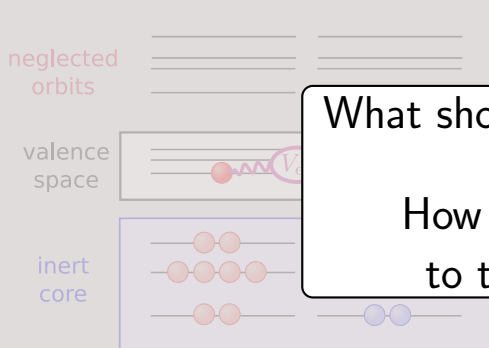
neglected  
orbitsvalence  
spaceinert  
core

$$H_{SM} = E_{core} + \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ijkl} a_i^\dagger a_j^\dagger a_l a_k$$

Diagonalize  $H_{SM}$  in space of all possible valence configurations.

$^{56}\text{Ni}$  in  $fp$  shell:

- $\sim 10^9$  possible configurations
- $\gtrsim 10^{11}$  nonzero matrix elements



$$H_{SM} = E_{core} + \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ijkl} a_i^\dagger a_j^\dagger a_l a_k$$

What should we use for  $V_{eff}$ ?

How is  $H_{SM}$  related to the "bare"  $H$ ?

space of all configurations.

- $\sim 10^9$  possible configurations
- $\gtrsim 10^{11}$  nonzero matrix elements

Operator  $P$  projects from large Hilbert space to smaller, tractable space.

Complementary projector  $Q = \mathbb{1} - P$

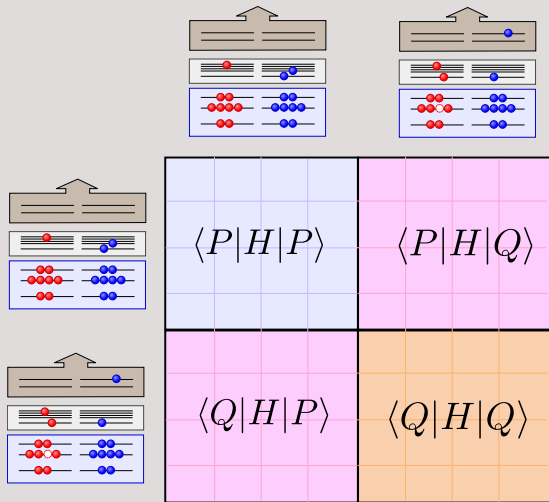
### Basic idea:

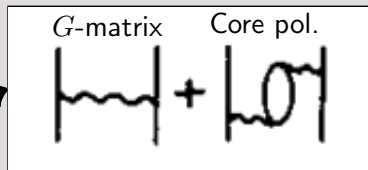
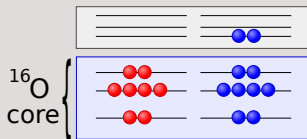
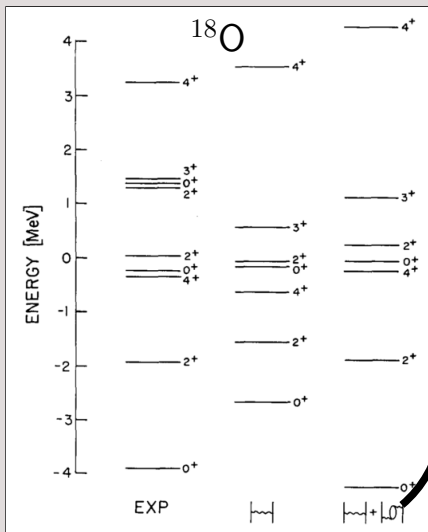
We want an *effective interaction*  $H_{\text{eff}}$  that acts only in the  $P$  space and reproduces the spectrum of  $H$  in the big space.

$$H|\Psi_i\rangle = E_i|\Psi_i\rangle$$

$$H_{\text{eff}}|\Phi_i\rangle = E_i|\Phi_i\rangle, \quad P|\Phi_i\rangle = |\Phi_i\rangle$$

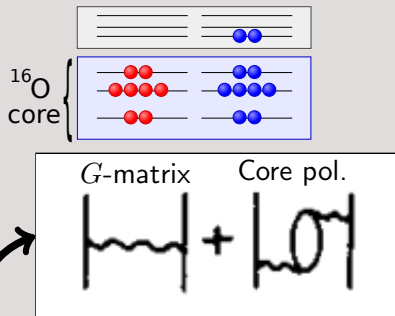
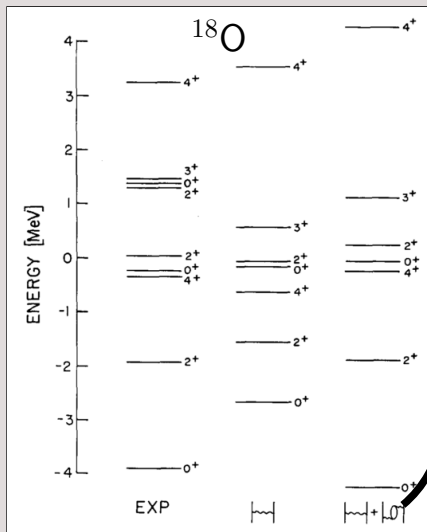
So we require:  $QH_{\text{eff}}P = 0$ .





$$V_{\text{eff}} = V + VQ \sum_{n=1}^{\infty} \left( \frac{1}{E_0 - H_0} QV \right)^n$$

Bertsch 1965; Kuo and Brown 1966; Barrett and Kirson 1970; Schucan and Weidenmüller 1973; Vary, Sauer, and Wong 1973



### Trouble:

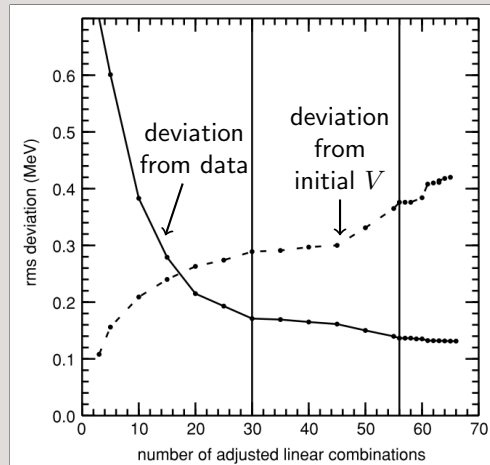
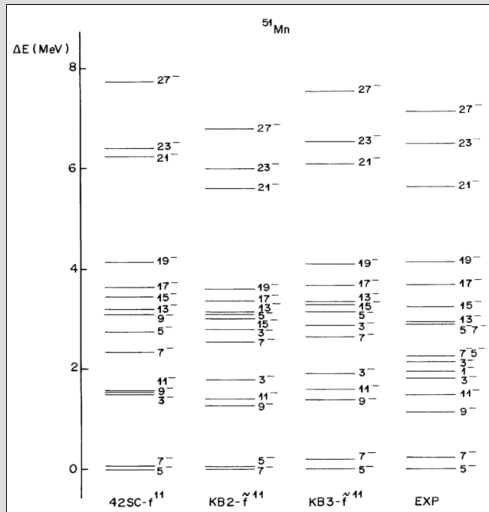
- 3rd order not small, makes it worse
- Perturbation series diverges
- Neglected virtual states important



## SHELL MODEL INTERACTION

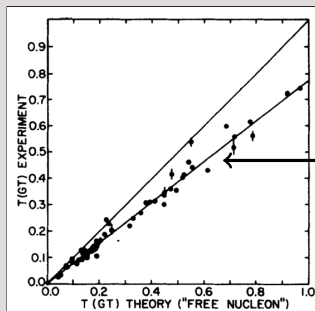
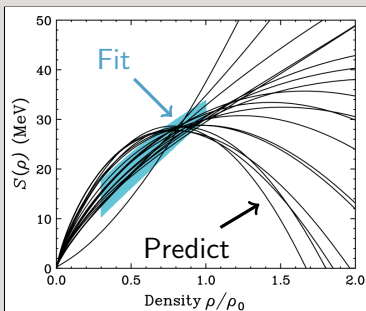
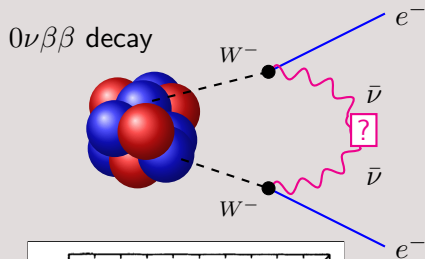
1. Take core energy from data
2. Fit single particle energies  $\epsilon_i$  to core+1 system
3. Adjust “monopoles” of  $V_{ijkl}$  to data (or treat as free parameters)
4. Scale  $V_{ijkl}$  by  $\approx A^{1/3}$
5. Use effective operators for electroweak observables, e.g.  $e_p = 1.5, e_n = 0.5$





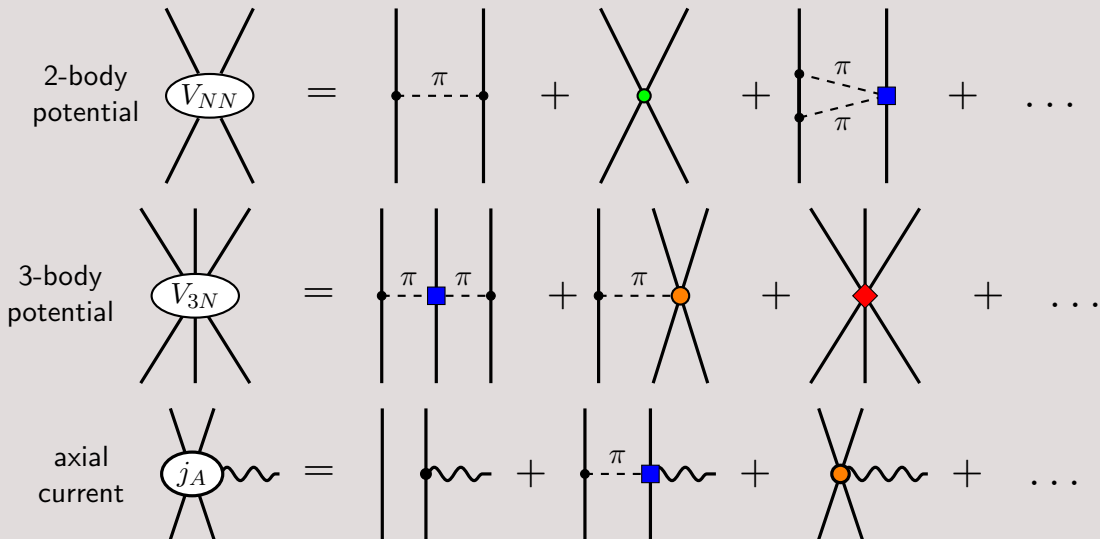
Poves and Zuker 1981; Brown et al. 2006

- Overfitting?
- Extrapolation?
- Lost connection to underlying interaction
- New valence space  
⇒ refit everything

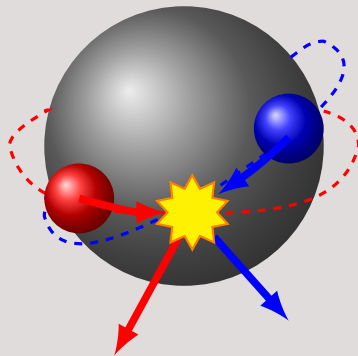
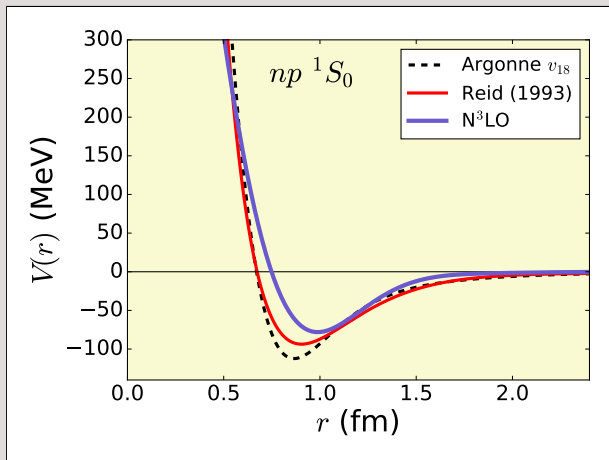


Bare Gamow-Teller operator needs to be "quenched" by factor  $\sim 0.75$

- Consistent  $NN$  and  $3N$  potential, and electroweak operators from chiral EFT
- Similarity Renormalization Group
- Non-perturbative many-body methods



Weinberg 1990; Kolck 1993; Ordóñez, Ray, and Kolck 1994; Epelbaum et al. 2002; Entem and Machleidt 2002; Navrátil 2007. ...



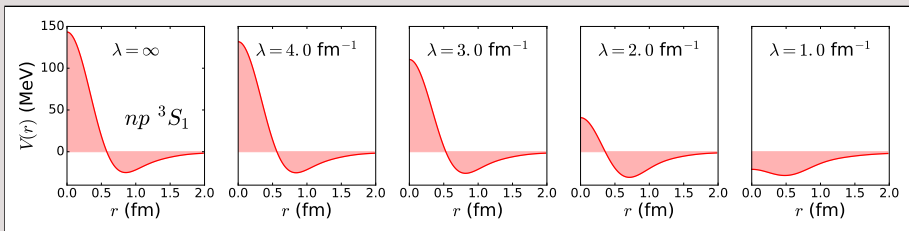
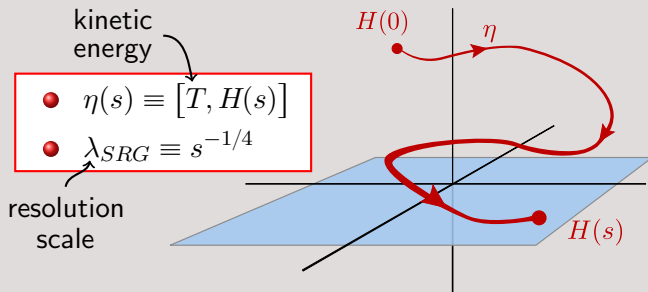
Short-range  
repulsion



Independent  
particles is a  
bad approx.

**SRG**

- $H(s) \equiv U(s)HU^\dagger(s)$
- $\frac{dU(s)}{ds} = \eta(s)U(s)$
- $\frac{dH(s)}{ds} = [\eta(s), H(s)]$



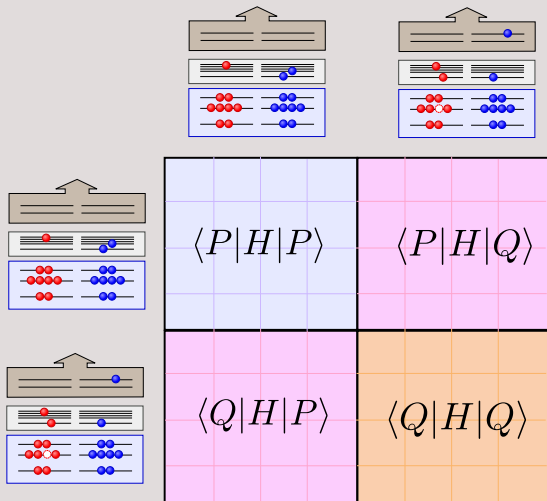
Glazek and Wilson 1994; Wegner 1994; White 2002; Bogner, Furnstahl, and Perry 2007

We want  $QHP = 0$ .

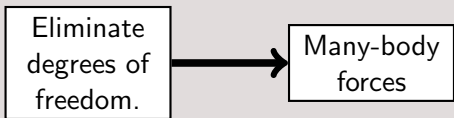
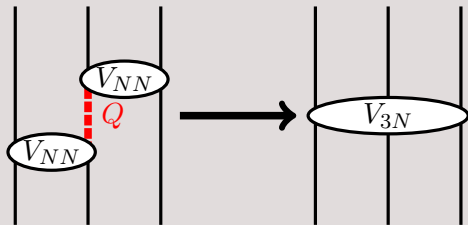
1. Partition  $H = H^d + H^{\text{od}}$   
with  
 $QH^dP = 0$

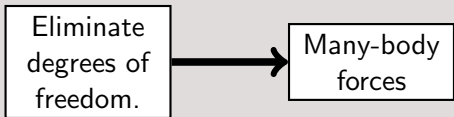
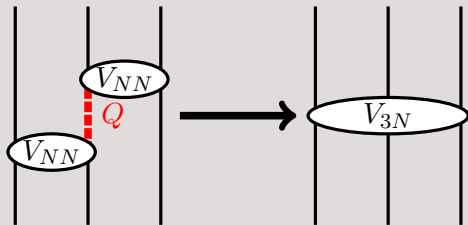
2. Define  
 $\eta(s) = \frac{H^{\text{od}}(s)}{\Delta}$

3.  $H^{\text{od}}(s) \rightarrow 0$  ,  $QHP \rightarrow 0$   
is a fixed point of the RG flow.

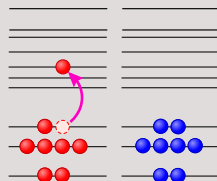






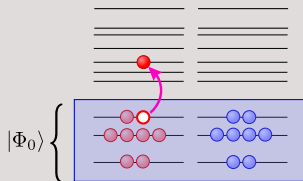


## Redefine the vacuum



16 particles

Many-particle forces

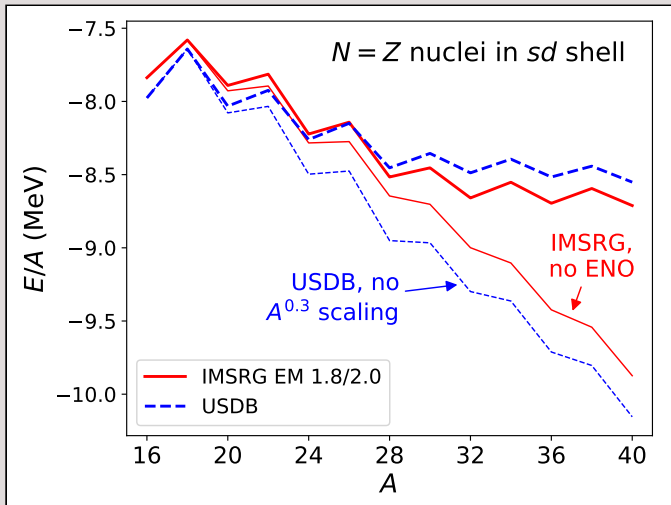


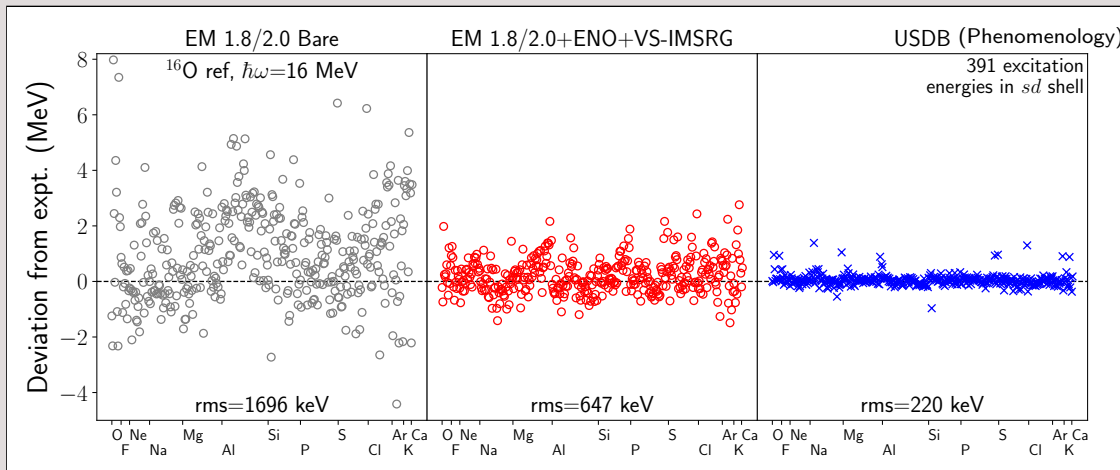
1 particle, 1 hole

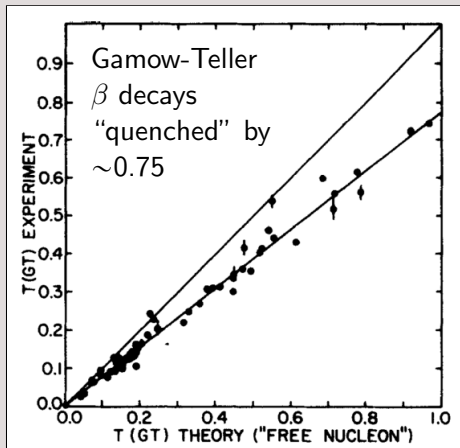
Many-quasiparticle forces

- Optimal reference depends on the nucleus being studied.
- Generate a different shell model interaction for *each nucleus*
- No loss of predictive power (no parameters to fit)

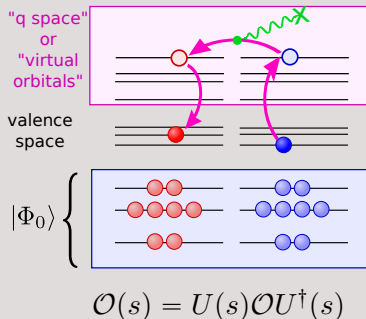
(USDB = phenomenology)



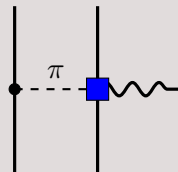


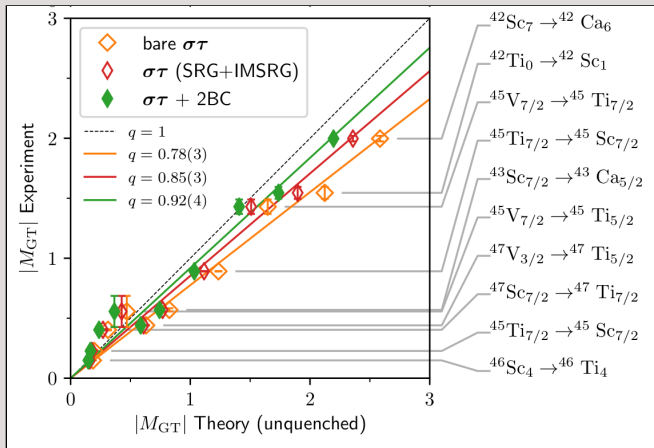
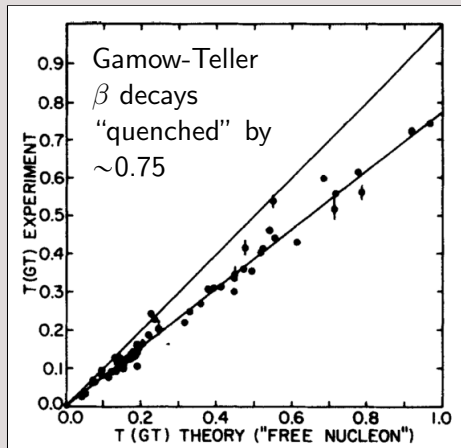


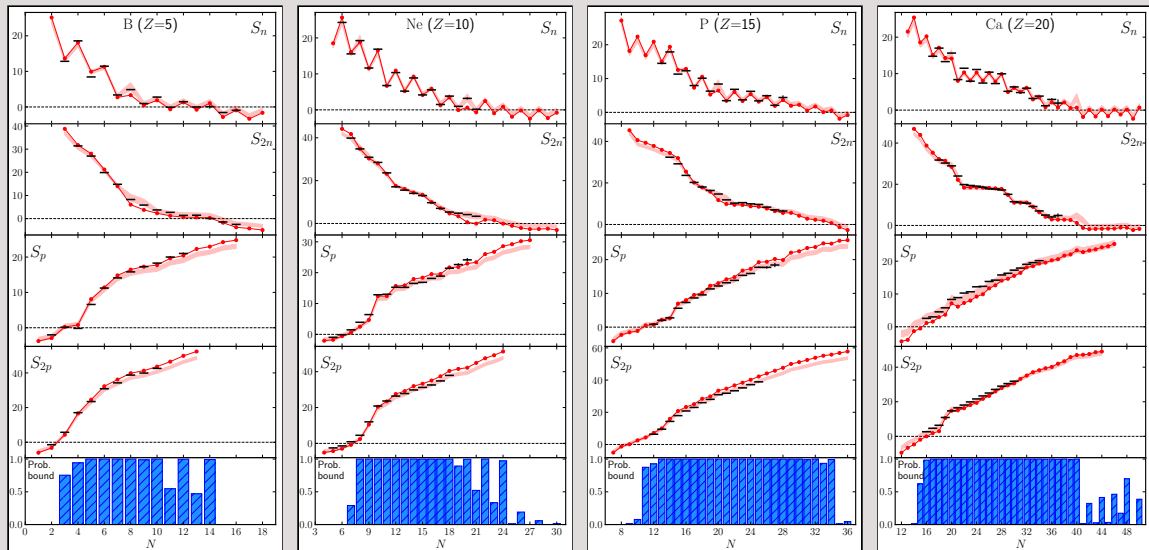
## Correlations



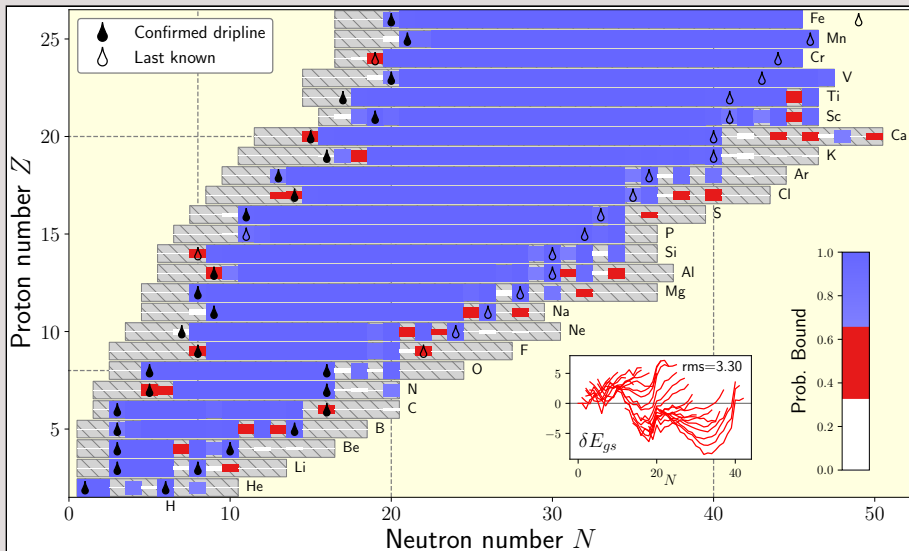
## Currents





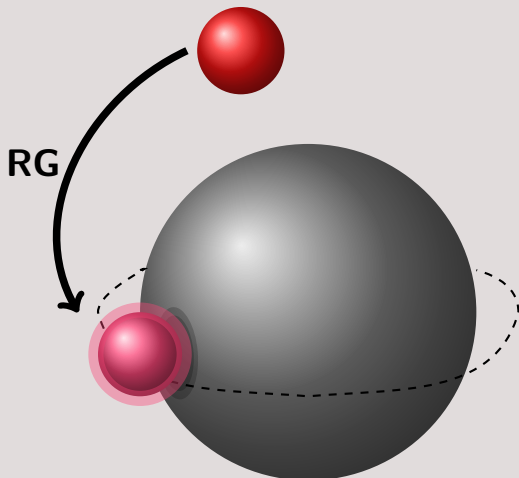


Holt et al. 2019

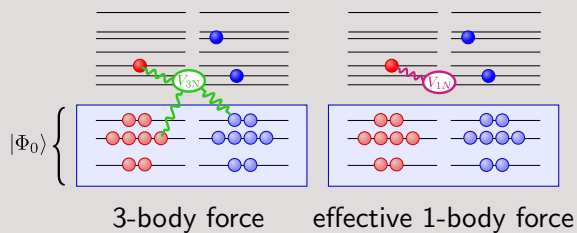




- RG connects free nucleons to shell model quasiparticles
- Old phenomenological adjustments can be traced to  $3N$  forces
- Some indications of universality in shell model interactions
- Strategy: Short distance  $\rightarrow$  RG flow, long distance  $\rightarrow$  diagonalization in valence space
- Diagonalization is exact, so bigger valence space should be better (see talk by T. Miyagi)

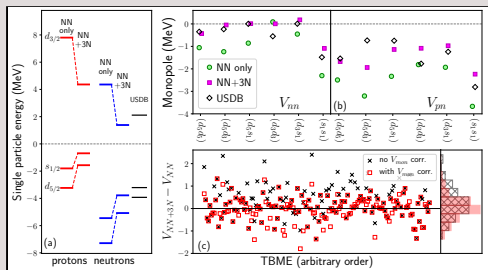
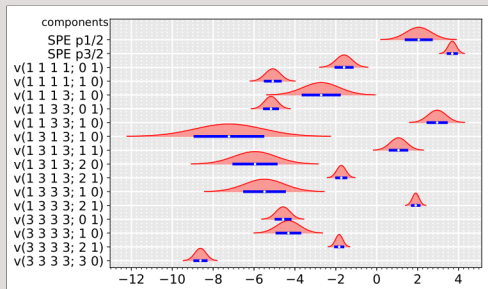


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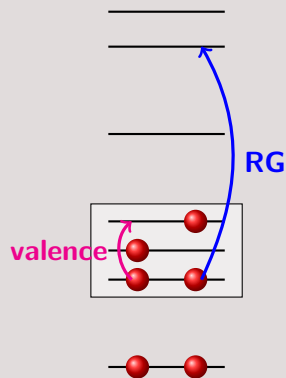
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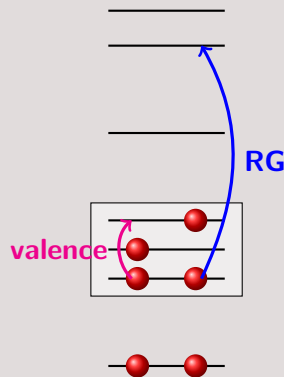


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## Summary

- The shell model is useful and successful because it picks the “right” d.o.f.
- Modern approaches can harness the virtues and insight of the old shell model, while avoiding some of the downsides
- Outstanding issues include: rigorous EFT power counting, quantification of many-body uncertainties, extended valence spaces, going beyond mass 100, the continuum . . .

### Collaborators:



A. Calci, J. Holt,  
P. Navrátil, P. Gysbers



NSCL/MSU S. Bogner, H. Hergert



LLNL S. Quaglioni, K. Wendt



ORNL/UT G. Hagen, G. Jansen, T. Morris,  
T. Papenbrock



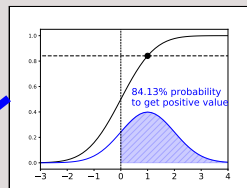
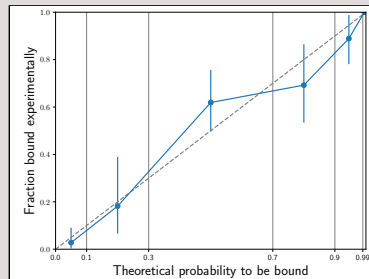
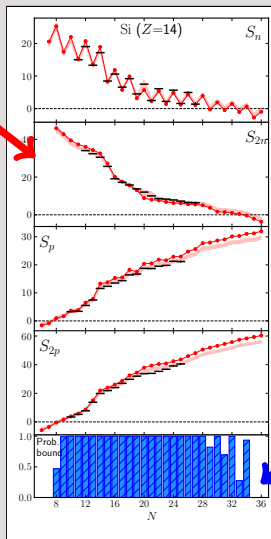
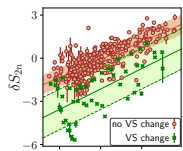
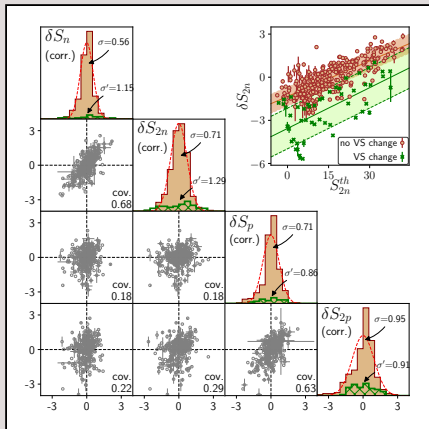
TU Darmstadt K. Hebeler, R. Roth, A. Schwenk



JGU Mainz J. Simonis

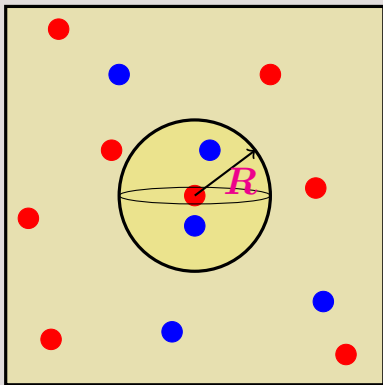
# Additional figures

# The "magic" EM1.8/2.0 interaction and the dripline





Cluster hierarchy  $H_{2b} > H_{3b} > H_{4b} \dots$  justified if  $\rho R^3 \ll 1$



Three vertical arrows point upwards, each with a black dot at its base. A wavy pink line connects the first two dots, with a bracket above it labeled  $R$ . To the right, a box contains the following text:

$$\rho \sim 0.16 \text{ fm}^{-3}$$

$$\Rightarrow R \ll 1.1 \text{ fm}$$

$$\langle H_{3b} \rangle \sim \int d^3 r_1 d^3 r_2 d^3 r_3 |\psi_1|^2 |\psi_2|^2 \underbrace{|\psi_3|^2}_{\rho} \underbrace{V(r_{12}, r_{13})}_{\sim V_0 \theta(R-r)}$$

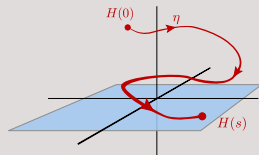
$$\sim \left( \frac{4\pi}{3} R^3 \rho \right)^2 V_0 \int d^3 r_1 |\psi_1|^2$$

# Similarity Renormalization Group for the effective interaction

$$H = E_0 + \sum_{ij} f_{ij} \{a_i^\dagger a_j\} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\} + \frac{1}{36} \sum_{\substack{ijk \\ lmn}} W_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}$$

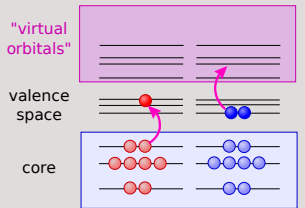
$$H = \text{red circle} + \text{red circle with 1 line} + \text{red circle with 2 lines} + \text{red circle with 3 lines}$$

$$\eta = \text{blue square with 1 line} + \text{blue square with 2 lines} + \text{blue square with 3 lines}$$



Flow equation:

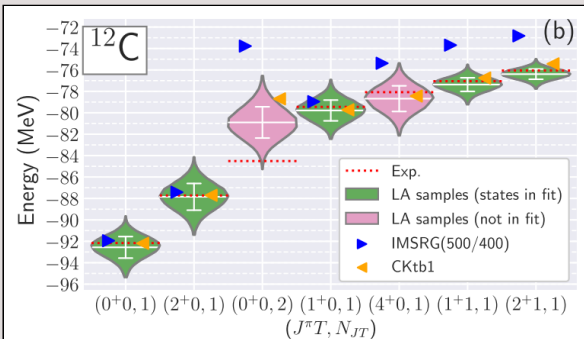
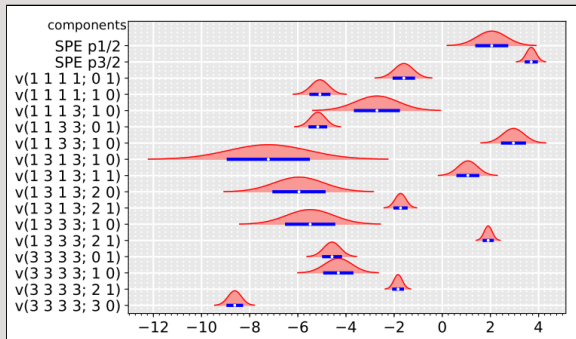
$$\frac{dH}{ds} = [\eta(s), H(s)] \Rightarrow \begin{cases} \frac{dE_0}{ds} = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \dots \\ \frac{df}{ds} = \text{diagram 4} + \text{diagram 5} + \text{diagram 6} + \dots \\ \frac{d\Gamma}{ds} = \text{diagram 7} + \text{diagram 8} + \text{diagram 9} + \dots \end{cases}$$

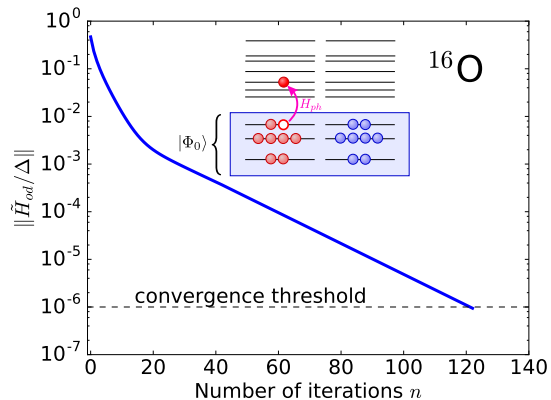
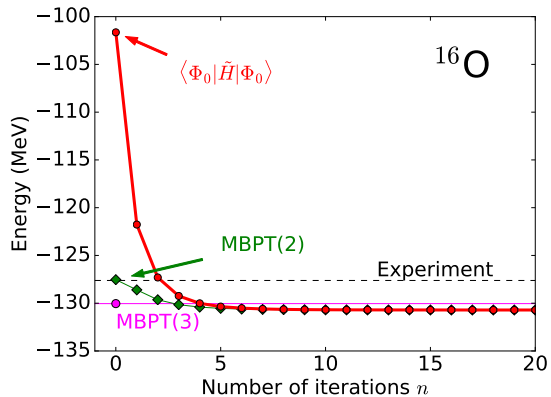


$$H^{\text{od}} \sim QHP \rightarrow 0$$

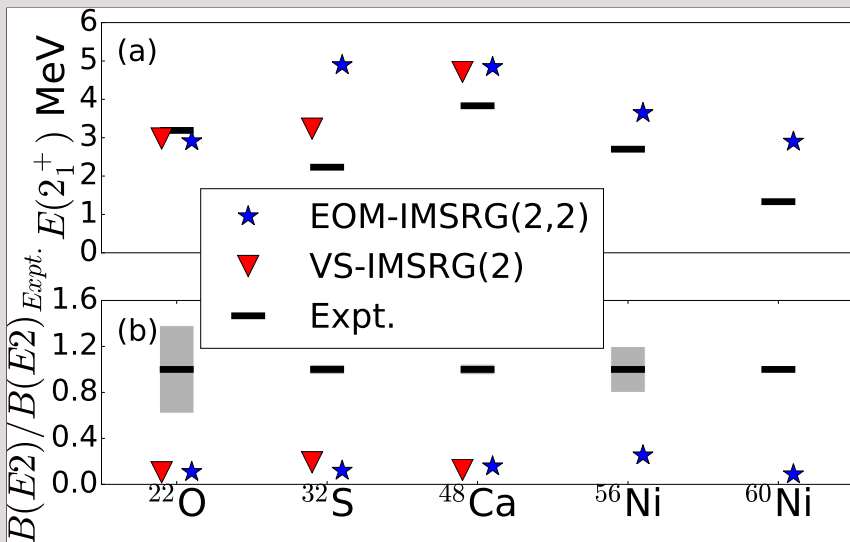
Glazek1993; Wegner 1994; Tsukiyama, Bogner, and Schwenk 2011; Bogner et al. 2014

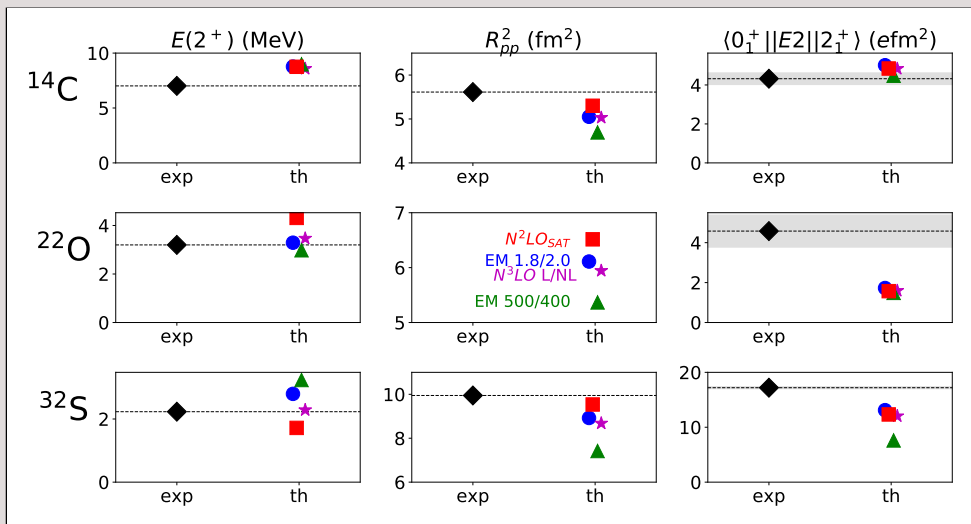
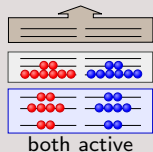
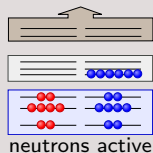
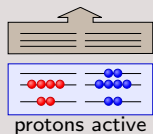
## A new twist on phenomenology



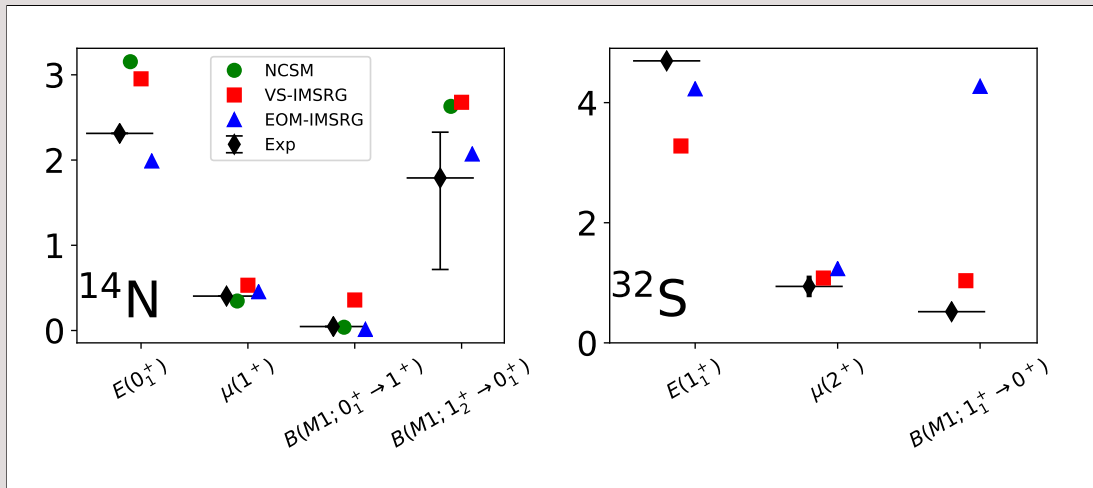


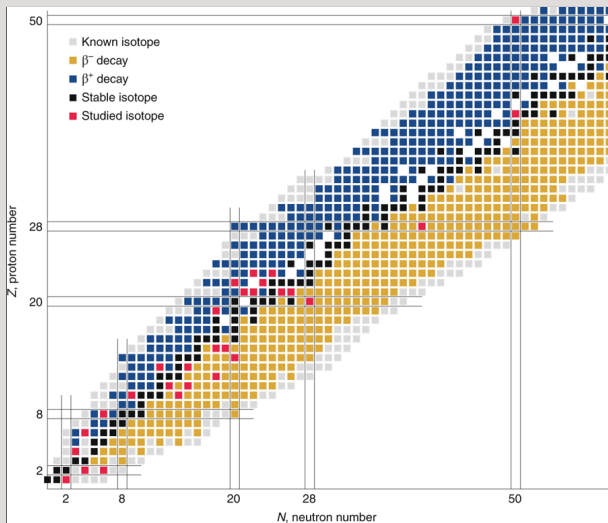
Tsukiyama 2010; Tsukiyama, Bogner, and Schwenk 2011; Hergert et al. 2013





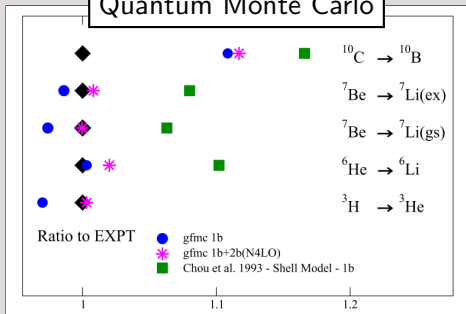
Entem and Machleidt 2003; Navrátil 2007; Gazit, Quaglioni, and Navrátil 2009; Ekström et al. 2015; Simonis et al. 2017



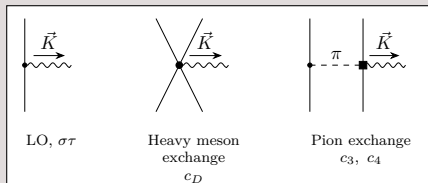
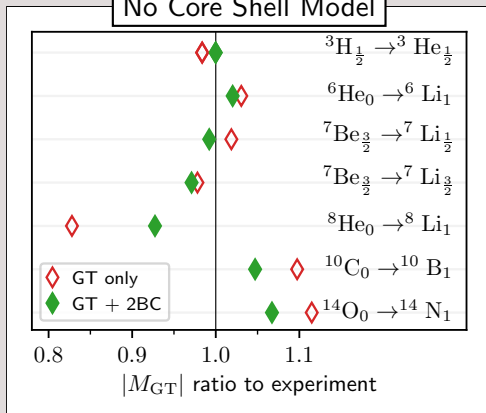




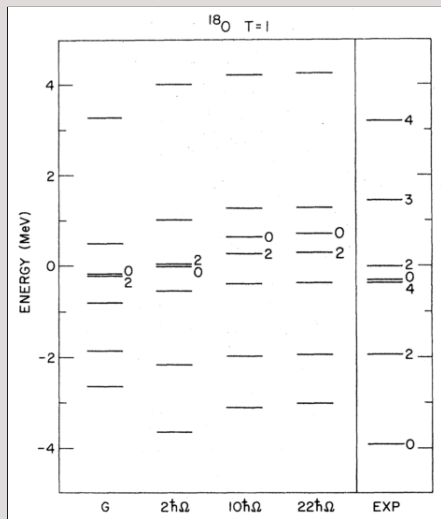
## Quantum Monte Carlo

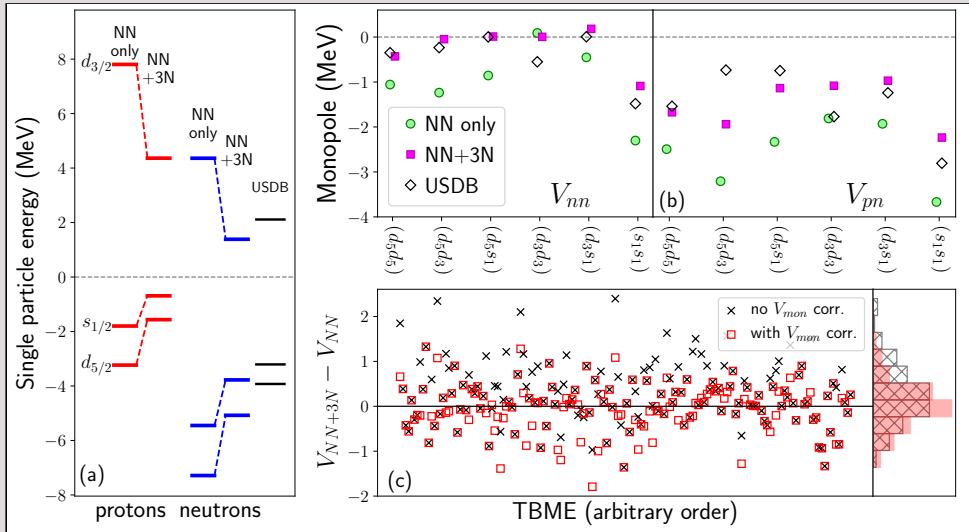


## No Core Shell Model



Pastore et al. 2017; Gysbers et al. 2019





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