



Single Pion Production in Neutrino-Nucleon Interaction

Minoo Kabirnezhad

National Center for Nuclear Research

Warsaw, Poland

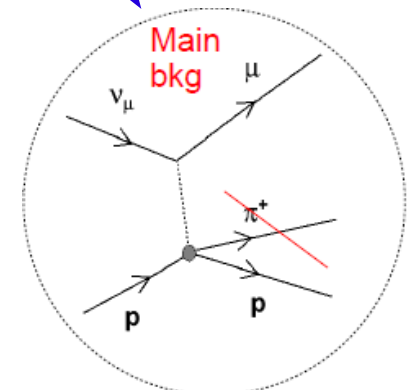
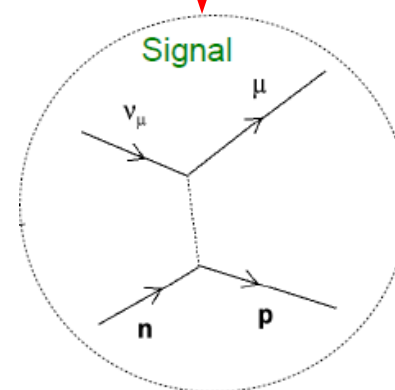
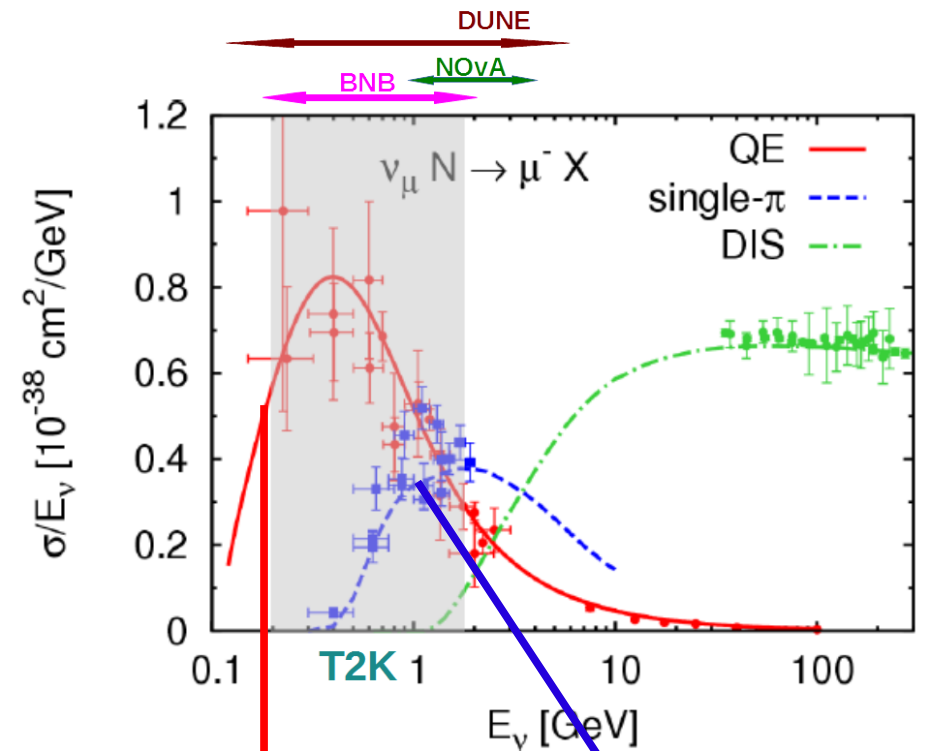
NuInt 2017

June 27th, 2017

Why Single Pion Production Model?

- Single pion production has significant contribution at intermediate energy.
- Good model is very crucial for precision neutrino oscillation measurements.
- Main goal is to develop a single pion production model **suitable** for Neutrino generators (NEUT)

- Fast enough
- Precise enough (with no assumption)
- Full kinematics
- Valid up to $W=2$ GeV (for NEUT)



Charged Current Single-Pion Production

$$\nu p \rightarrow \mu^- p \pi^+$$

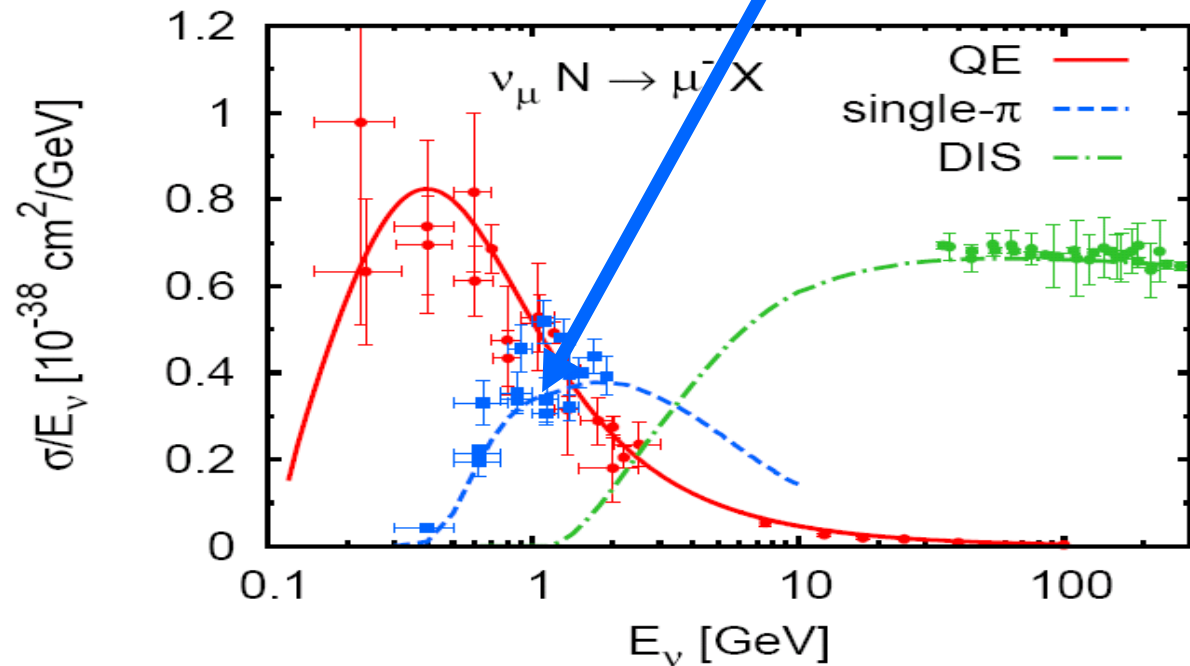
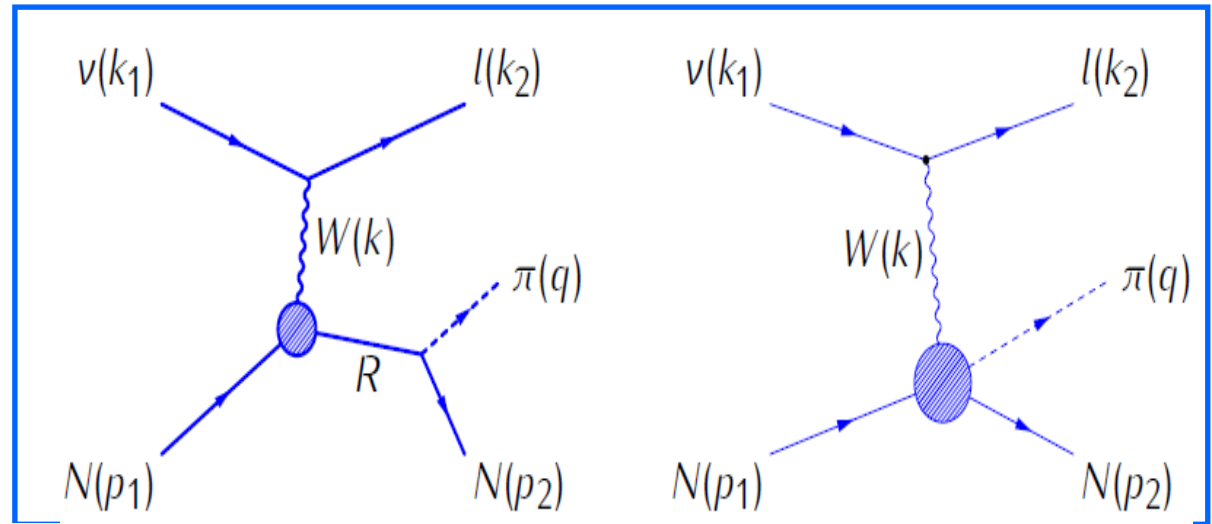
$$\nu n \rightarrow \mu^- p \pi^0$$

$$\nu n \rightarrow \mu^- n \pi^+$$

$$\bar{\nu} n \rightarrow \mu^+ n \pi^-$$

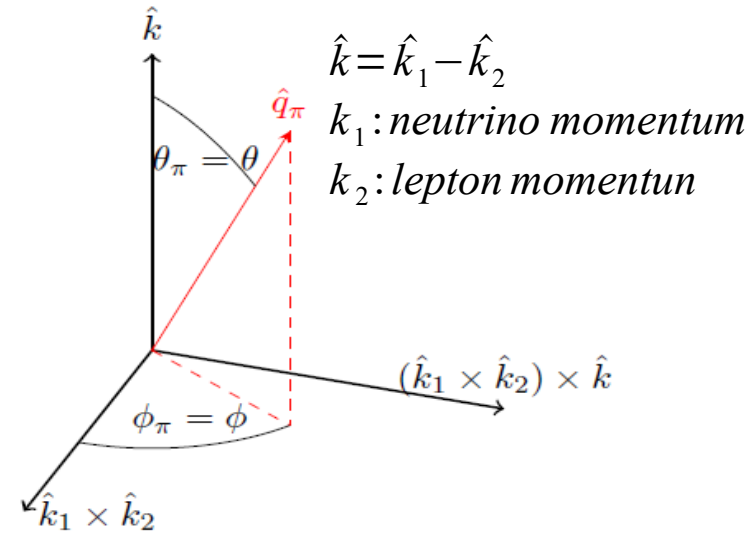
$$\bar{\nu} p \rightarrow \mu^+ n \pi^0$$

$$\bar{\nu} p \rightarrow \mu^+ p \pi^-$$



Rein Model

- Define a suitable framework; Adler frame
- Calculate both resonant and non-resonant interactions
- Add them coherently to include the interference effects

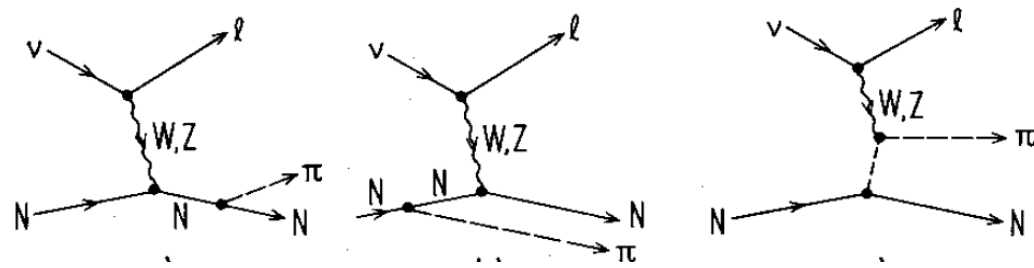


→ Resonant Interaction

- **Rein-Sehgal** model is based on **helicity amplitudes** derived in a relativistic quark model
- The helicity amplitudes depend on the spin projection of the initial and final particle.

→ Nonresonant Interaction




Born graphs based on linear sigma model.




→ The lepton is assumed to be massless

Rein-Sehgal model

Rein-Sehgal¹ (RS) is default model in the **NEUT**

-  Easy to be implemented in generators
-  Valid up to $W < 2$ GeV region
-  The RS model in NEUT includes the lepton mass and Graczyk-Sobczyk² form-factor

 Not a full kinematic model $d\sigma/dW dQ^2$
The helicity amplitudes are **not** a function of pion angles

 Pion angles is described by density matrix. NEUT **only** implemented the Δ resonance.

The RS model is improved by including the pion angles in helicity amplitudes

Output of the modified RS model

$$d\sigma/dW dQ^2 d\Omega_\pi$$

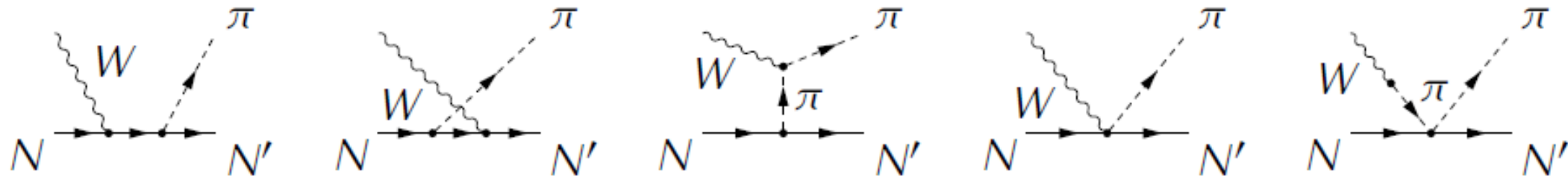
Resonance	M_R	Γ_0	χ_E
$P_{33}(1232)$	1232	117	1
$P_{11}(1440)$	1430	350	0.65
$D_{13}(1520)$	1515	115	0.60
$S_{11}(1535)$	1535	150	0.45
$P_{33}(1600)$	1600	320	0.18
$S_{31}(1620)$	1630	140	0.25
$S_{11}(1650)$	1655	140	0.70
$D_{15}(1675)$	1675	150	0.40
$F_{15}(1680)$	1685	130	0.67
$D_{13}(1700)$	1700	150	0.12
$D_{33}(1700)$	1700	300	0.15
$P_{11}(1710)$	1710	100	0.12
$P_{13}(1720)$	1720	250	0.11
$F_{35}(1905)$	1880	330	0.12
$P_{31}(1910)$	1890	280	0.22
$P_{33}(1920)$	1920	260	0.12
$F_{37}(1950)$	1930	285	0.40

1. D. Rein and L. M. Sehgal, Annals Phys. 133 (1981) 79. 5

2. K. M. Graczyk and J. T. Sobczyk, Phys. Rev. D 77 (2008) 053001.

non-resonant background is defined by a set of diagrams determined by HNV model based on non-linear sigma model.

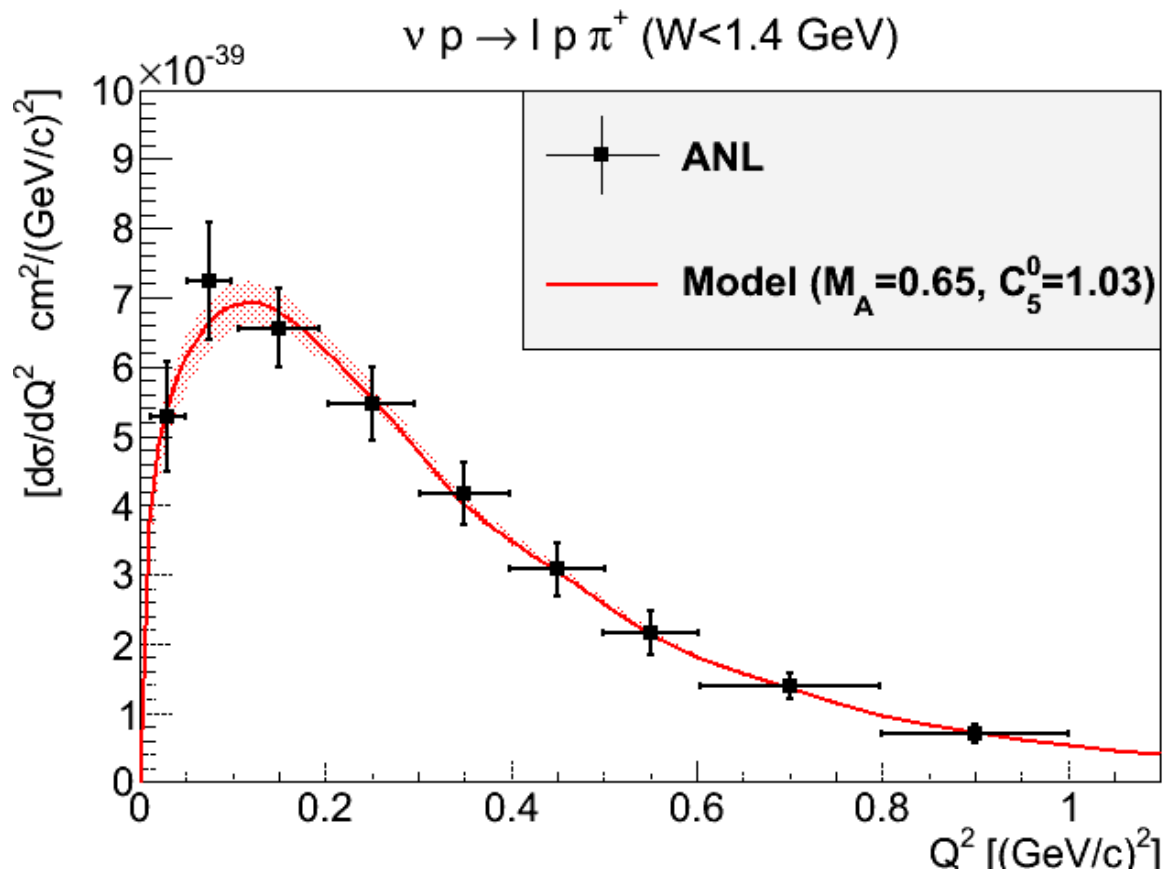
E. Hernandez, J. Nieves and M. Valverde,
Phys. Rev. D 76 (2007) 033005



- Helicity amplitudes of above diagrams are calculated in the Adler frame.
- the model is based on chiral symmetry and it is not reliable at high energy and high W .
- A practical solution (suggested by J. Nieves) is to multiply a form-factor by the virtual pion propagator.
- To preserve the conservation of vector current we have to multiply the same form-factor to other amplitudes
- The nonresonant background has no contribution at $W > 1.6$ GeV

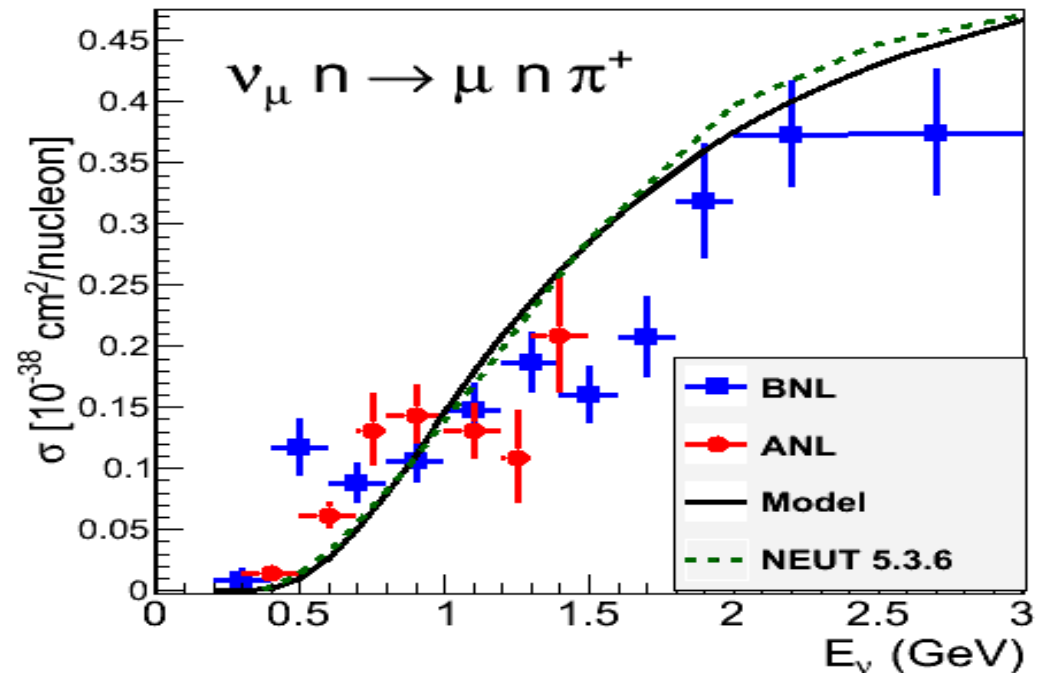
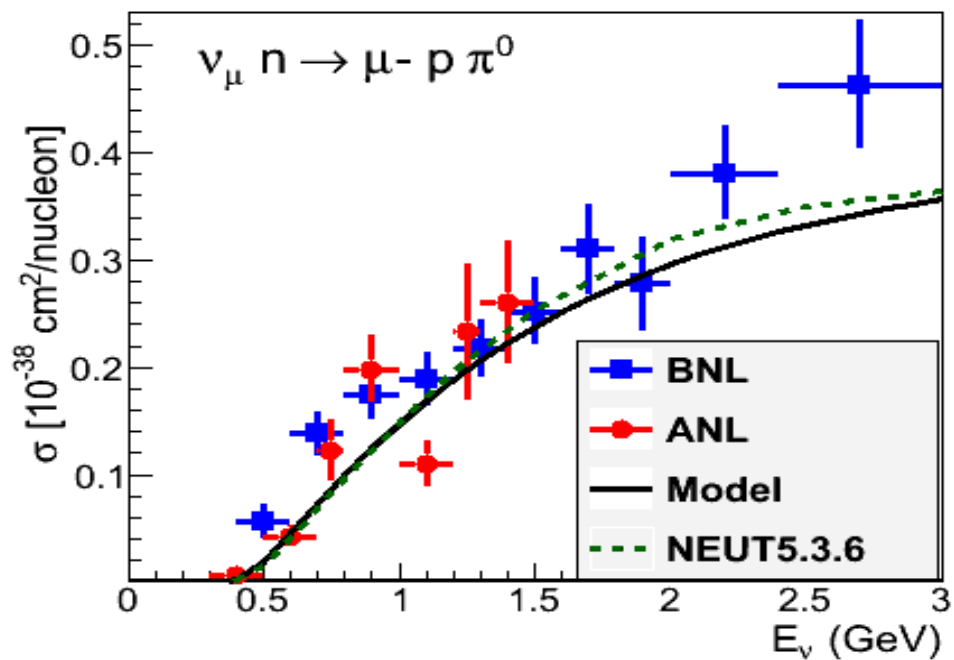
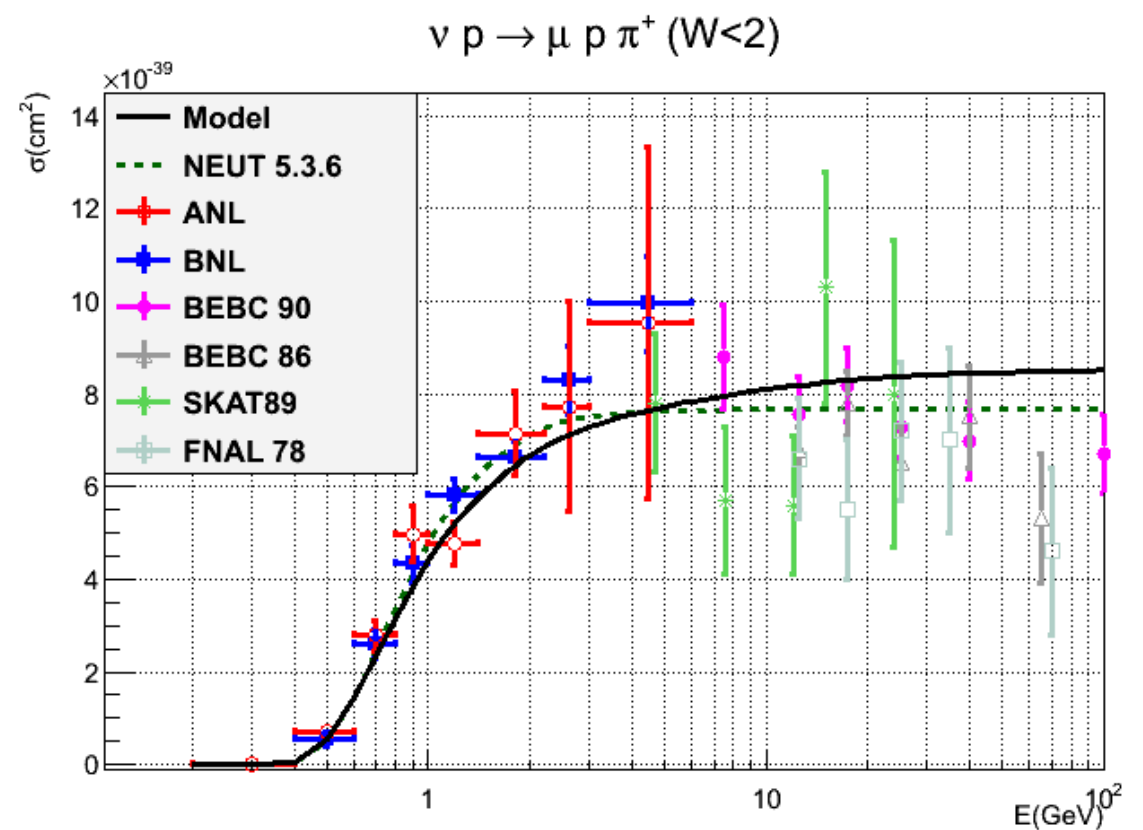
Verifying the model is difficult with limited data sets!

- Existing bubble chamber data on “free” nucleon are old and with large error and it is very unlikely to be improved :(
- We extract the dipole axial form-factor from the bubble chamber data.

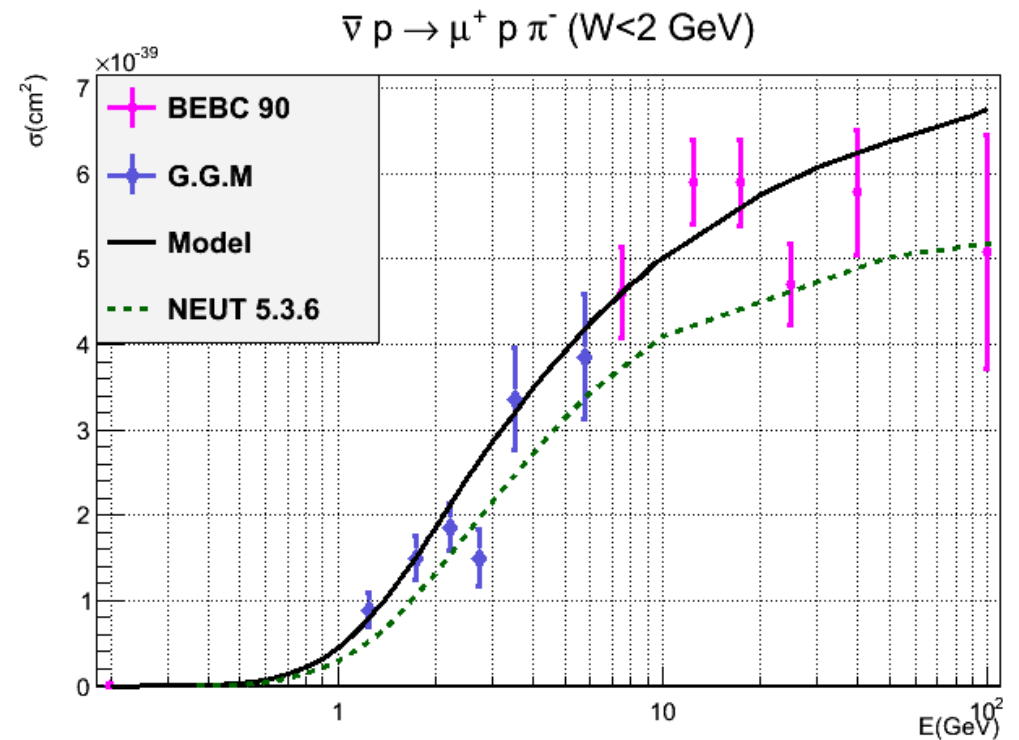
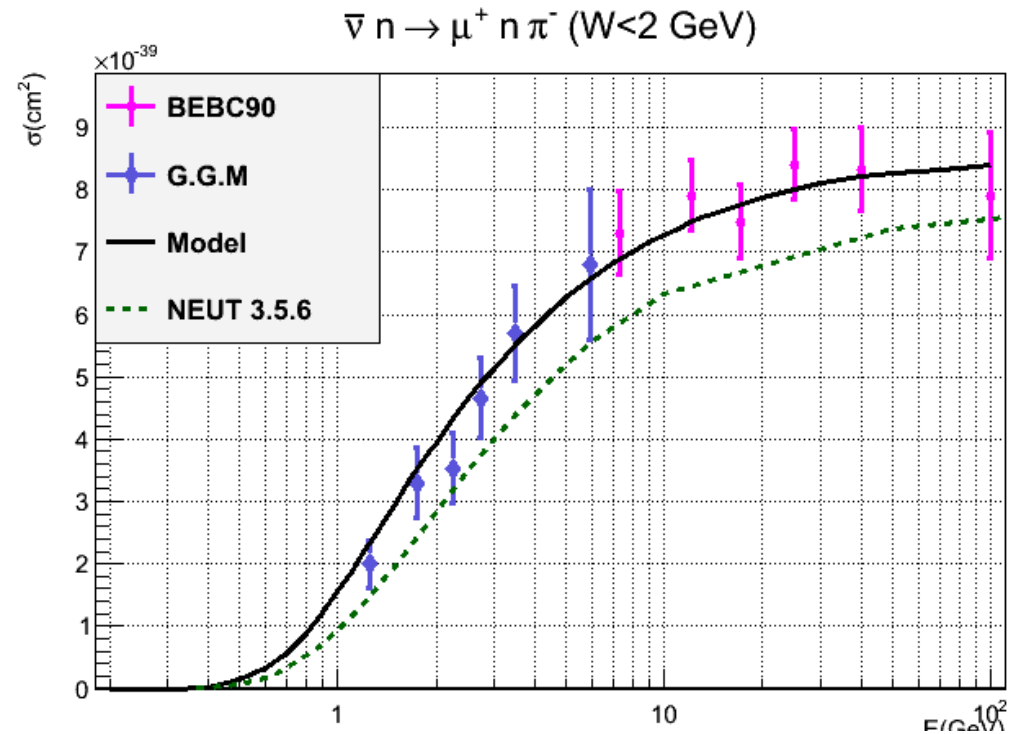


Total Cross-section for charged current neutrino interaction

ANL & BNL data is reanalyzed by
C. Wilkinson et al., Phys. Rev. D 90 (2014) 11,
112017



Total Cross-section for charged current anti-neutrino interactions

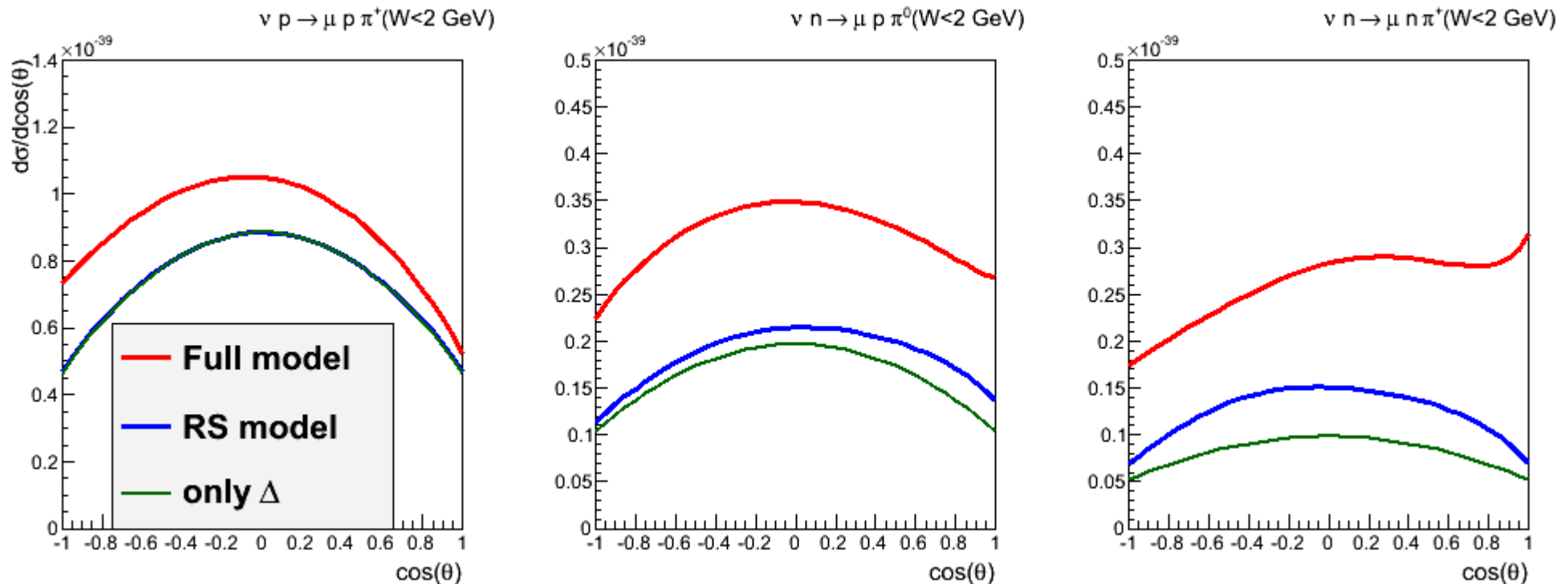


Data from

D. Allasia et al., Nucl. Phys. B 343 (1990) 285.
 T. Bolognese, J. P. Engel, J. L. Guyonnet and J. L.
 Riester, Phys. Lett. 81B (1979) 393.

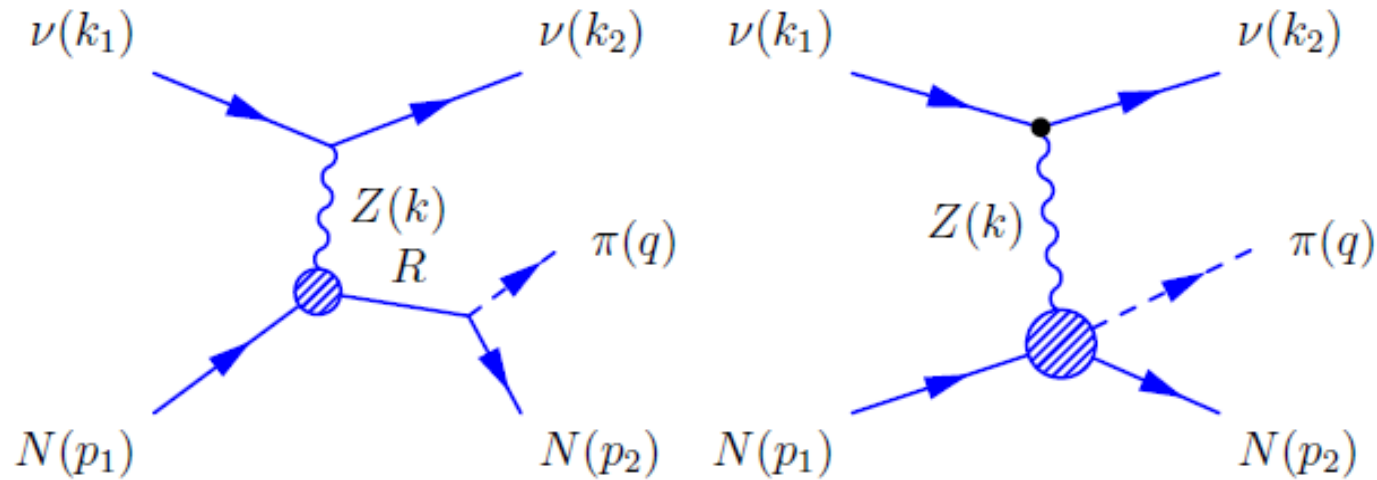
Angular Distribution In the Adler frame

For T2K energy



- “Only Δ ” has a symmetric distribution.
- Any deviation from the symmetric distribution comes from the interference effects.

Neutral Current Single-Pion Production



$$\nu p \rightarrow \nu p \pi^0$$

$$\bar{\nu} p \rightarrow \bar{\nu} p \pi^0$$

$$\nu p \rightarrow \nu n \pi^+$$

$$\bar{\nu} p \rightarrow \bar{\nu} n \pi^+$$

$$\nu n \rightarrow \nu n \pi^0$$

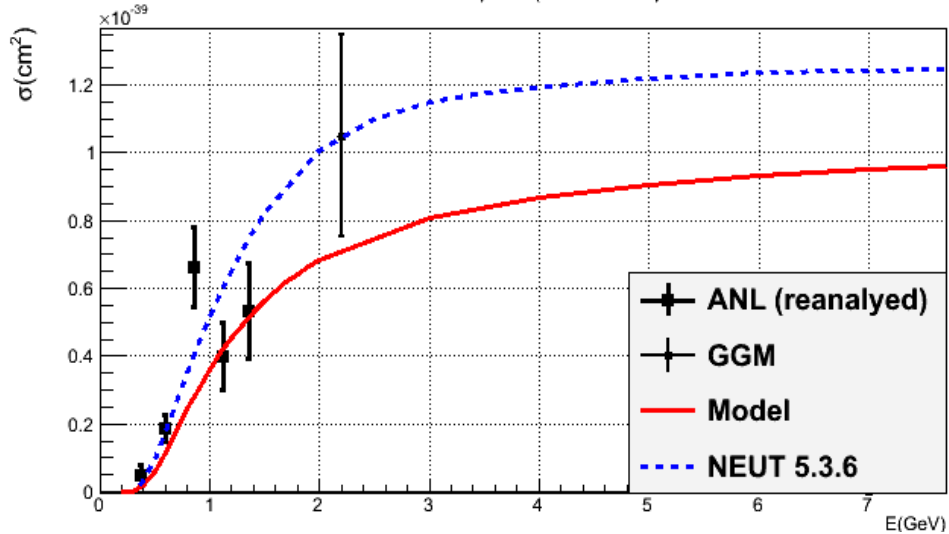
$$\bar{\nu} n \rightarrow \bar{\nu} n \pi^0$$

$$\nu n \rightarrow \nu p \pi^-$$

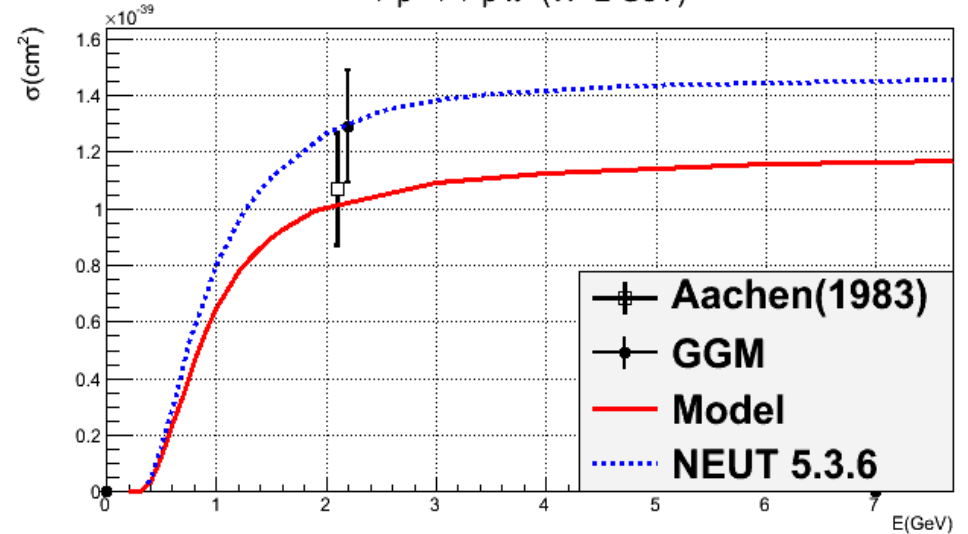
$$\bar{\nu} n \rightarrow \bar{\nu} p \pi^-$$

Total Cross-section for neutral current neutrino interactions

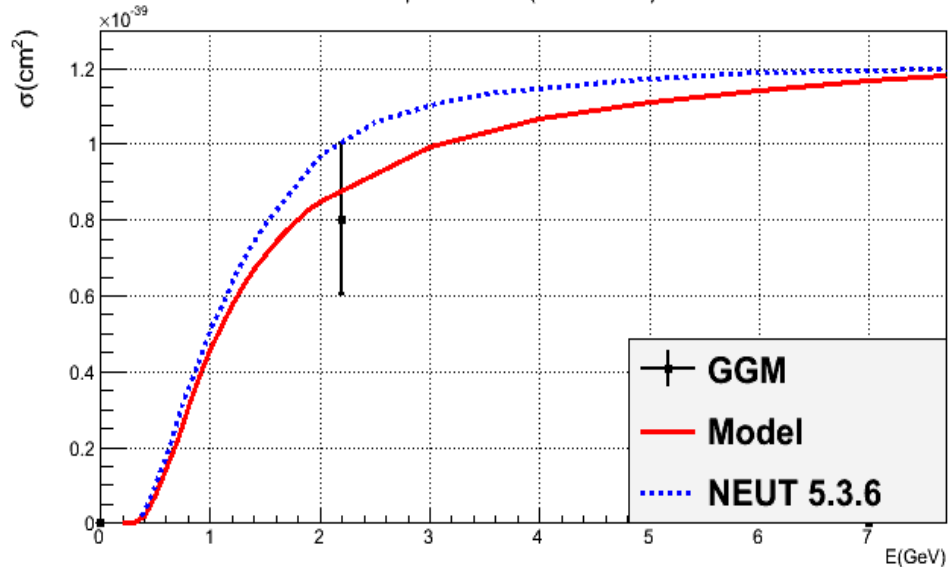
$\nu n \rightarrow \nu p \pi^- (W < 2 \text{ GeV})$



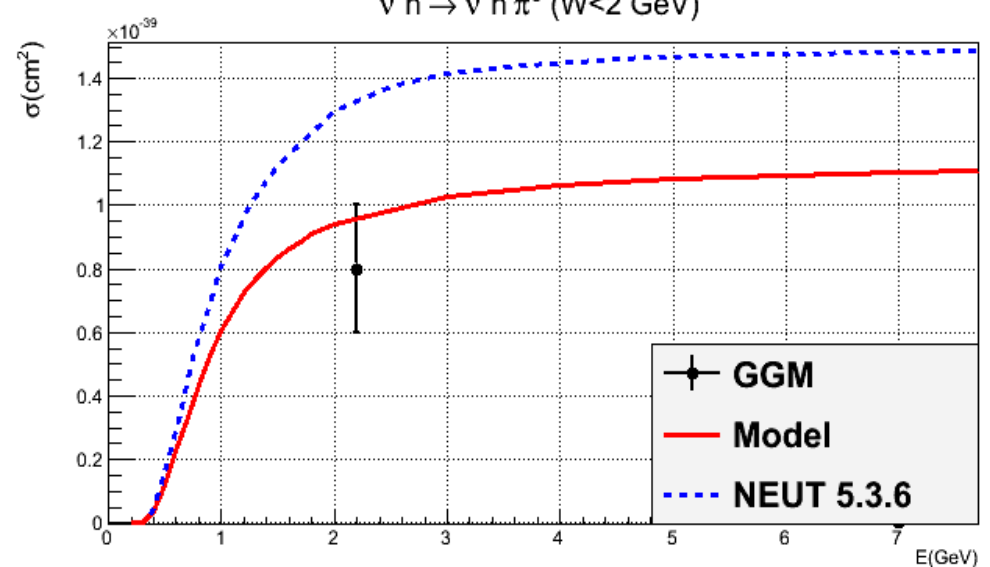
$\nu p \rightarrow \nu p \pi^0 (W < 2 \text{ GeV})$



$\nu p \rightarrow \nu n \pi^+ (W < 2 \text{ GeV})$



$\nu n \rightarrow \nu n \pi^0 (W < 2 \text{ GeV})$



The NEUT Implementation MK-Model

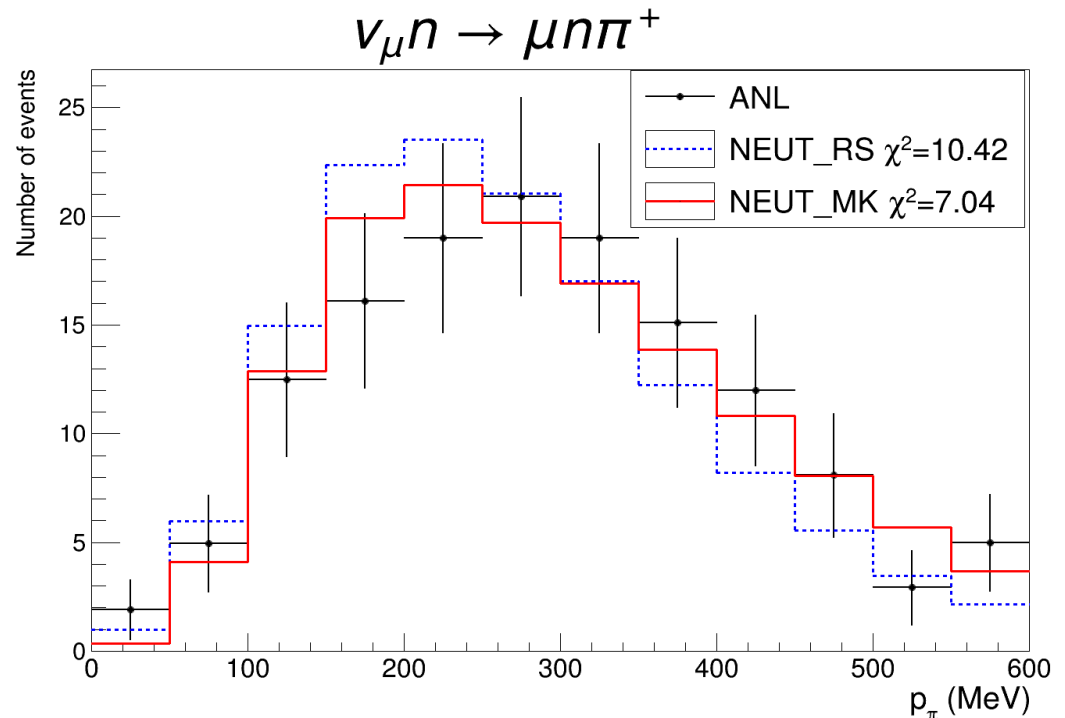
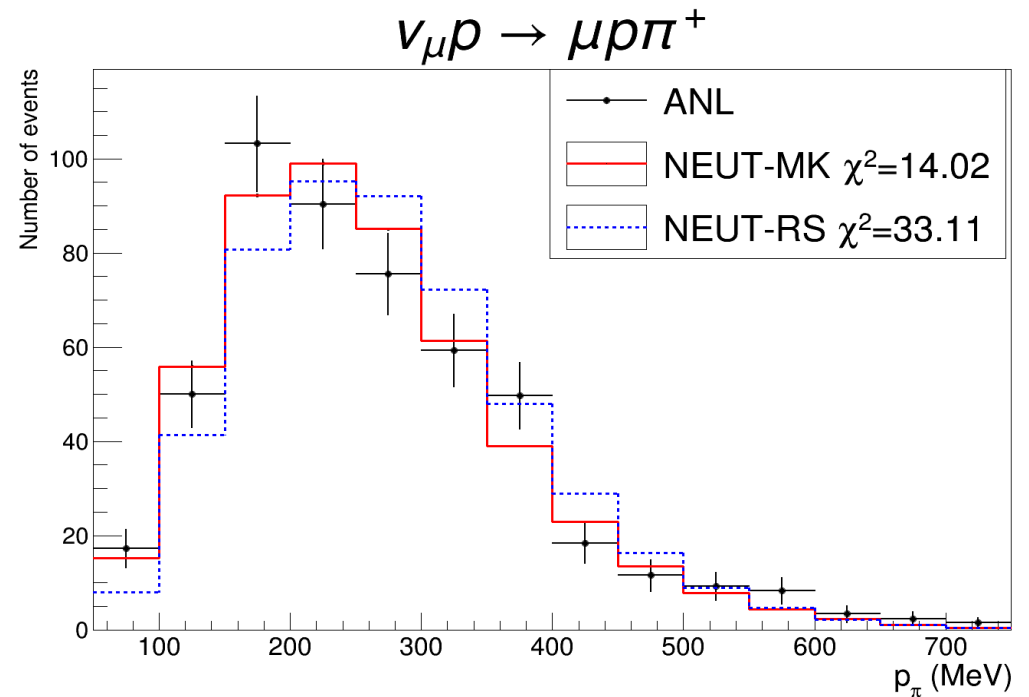
Clarence Wret and I
implemented the model in
NEUT

Data comparison can be done

- ◆ In any frame
- ◆ With any target
- ◆ With any selection

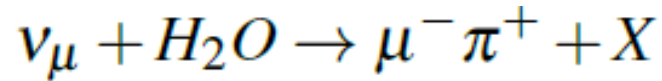
Data comparison with

NUISANCE ❤️



T2K data on Water

Pion kinematics



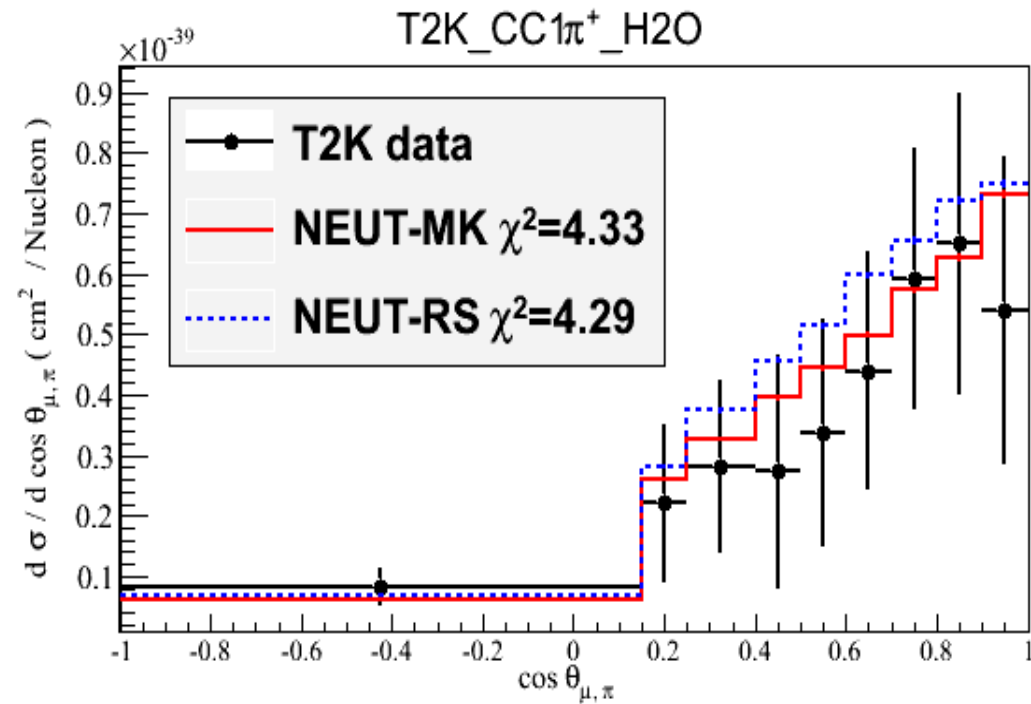
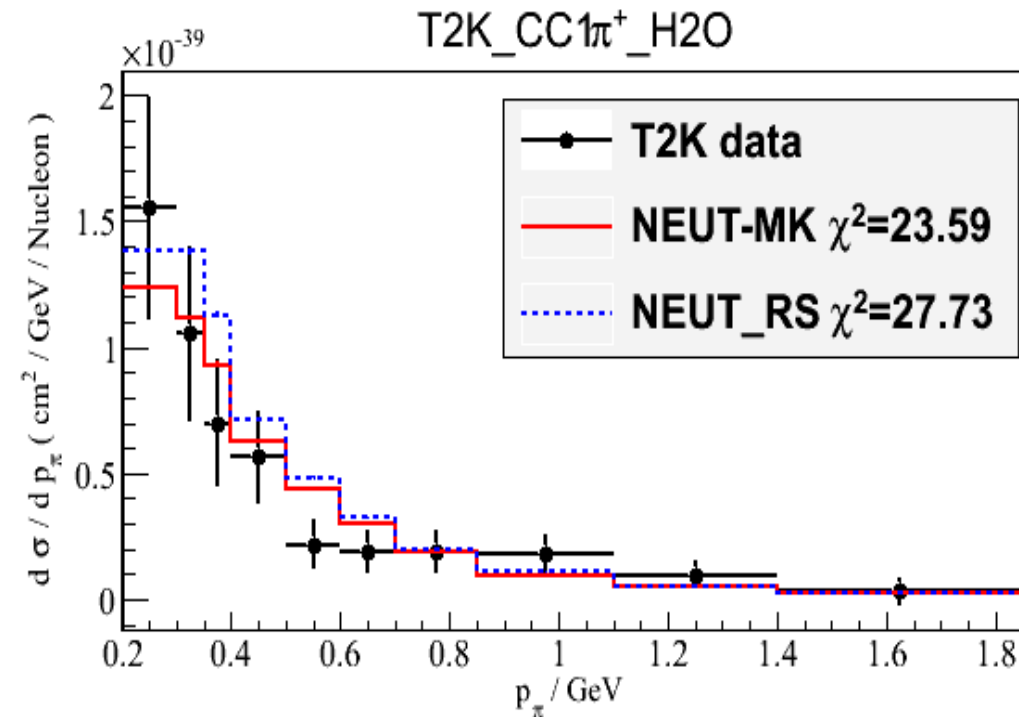
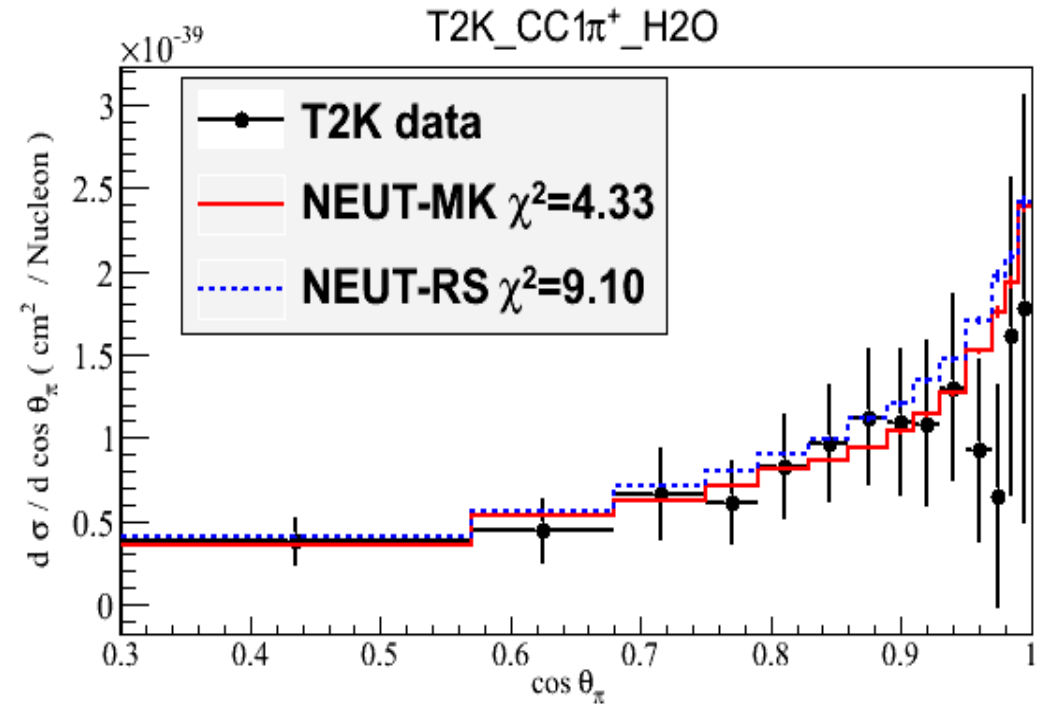
$$p_{\mu} > 200 \text{ MeV}/c$$

$$p_{\pi} > 200 \text{ MeV}/c$$

$$\cos \theta_{\mu} > 0.3$$

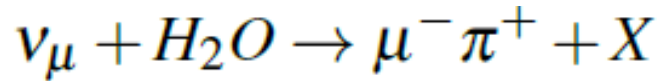
$$\cos \theta_{\pi} > 0.3$$

K. Abe et al. [T2K Collaboration], Phys. Rev. D 95 (2017) no.1, 012010.



T2K data on Water

Muon kinematics



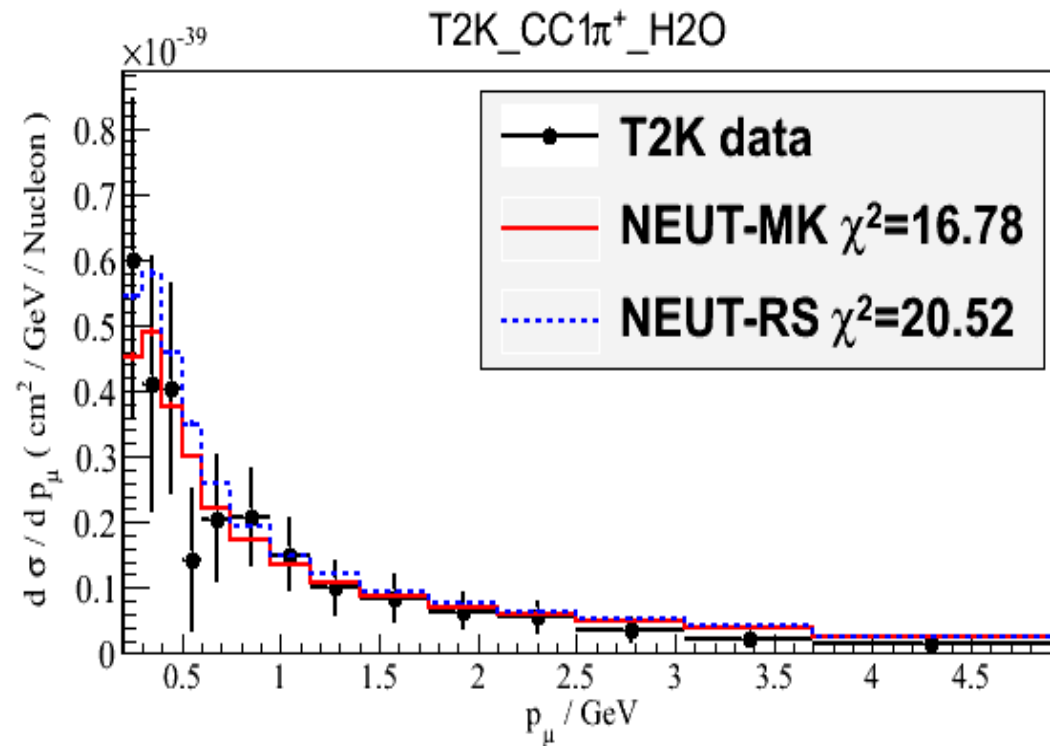
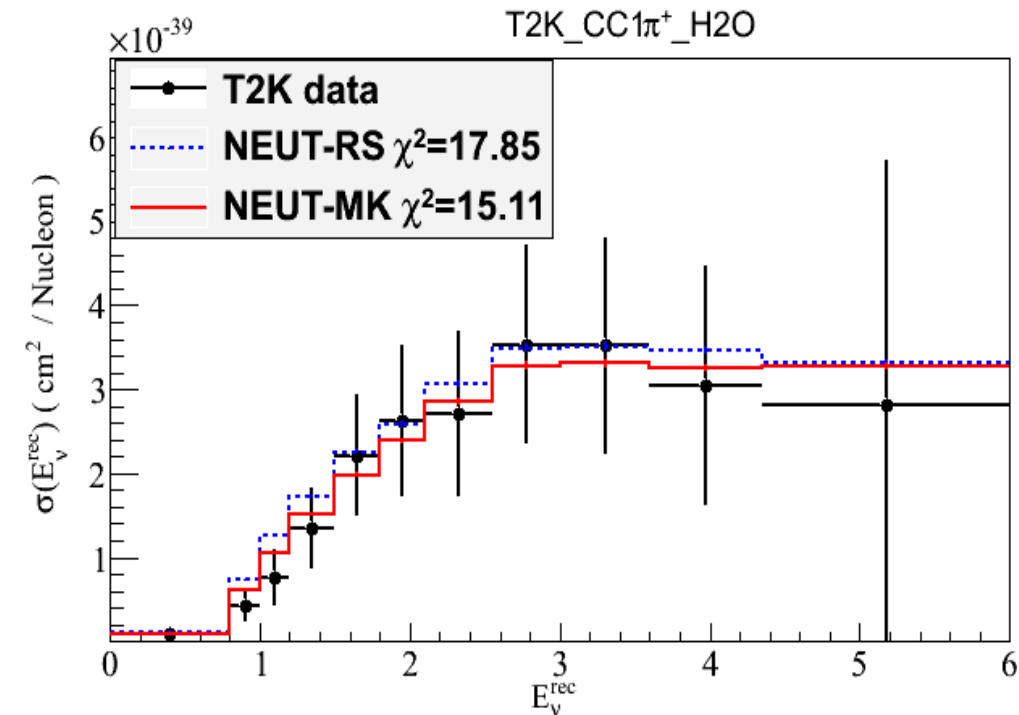
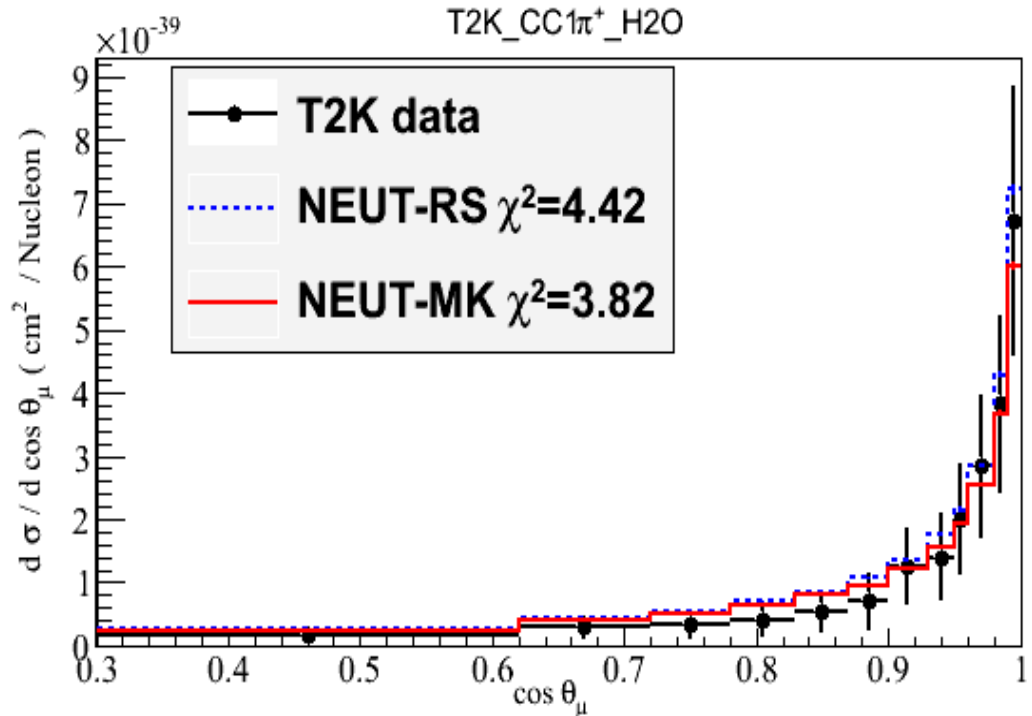
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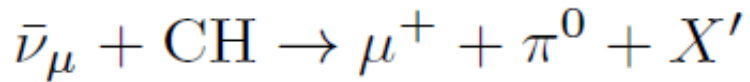
K. Abe et al. [T2K Collaboration], Phys. Rev. D 95 (2017) no.1, 012010.



MINERvA data 2016

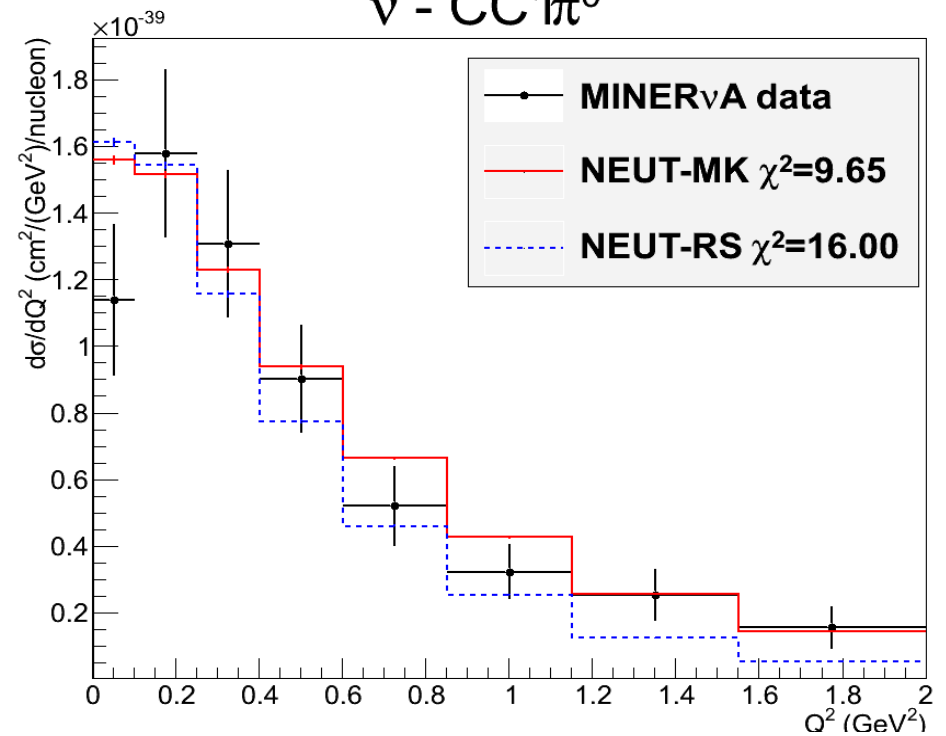
CH target

$W < 1.8$ GeV

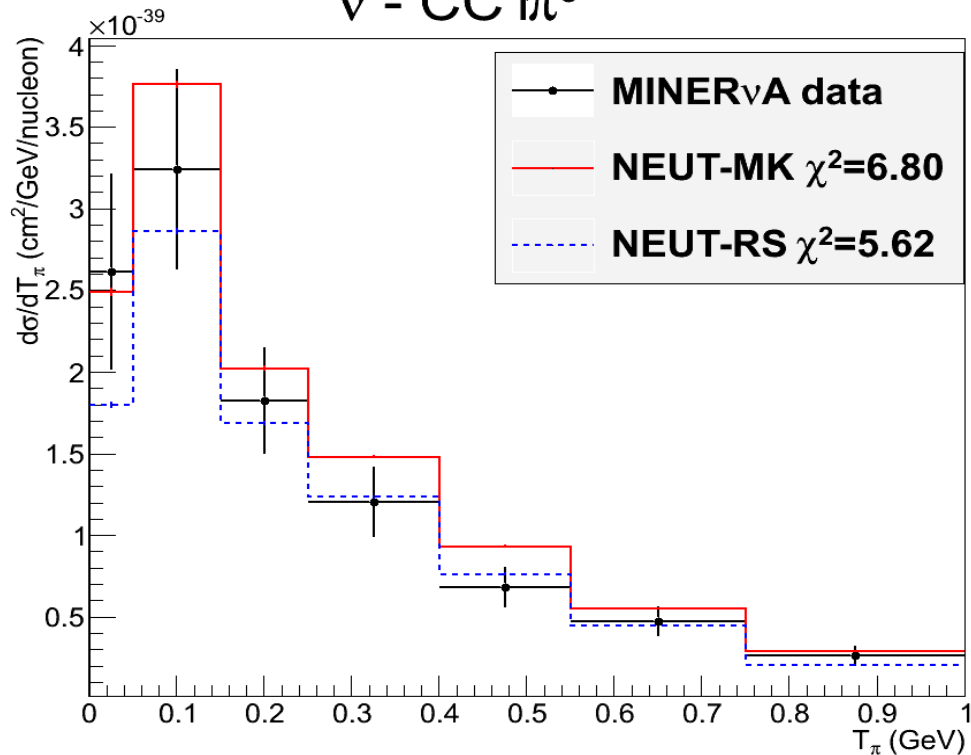


C. L. McGivern et al. [MINERvA Collaboration], Phys.Rev. D 94 (2016) no.5, 052005

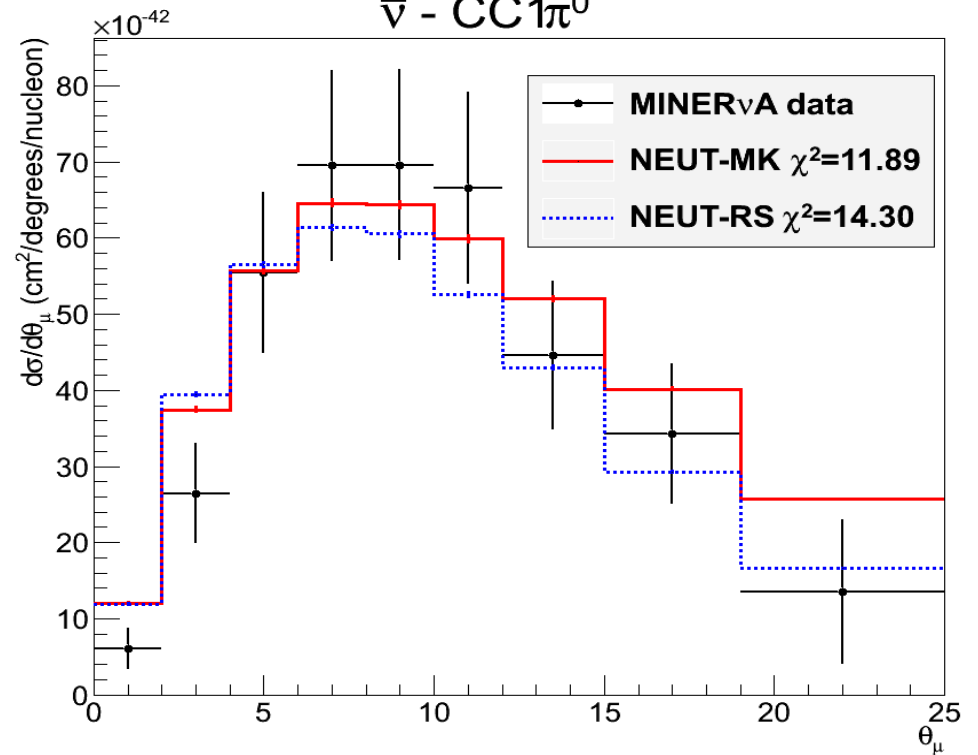
$\bar{\nu} - \text{CC}1\pi^0$



$\bar{\nu} - \text{CC}1\pi^0$



$\bar{\nu} - \text{CC}1\pi^0$



Summary

- The model consists of resonant and nonresonant interactions, **including the interference effects.**
- The **lepton mass** is included
- The model shows good agreements with bubble chamber data.
- The model is implemented in NEUT and it is satisfactory:
 - ✓ It is as fast as the Rein-Sehgal model.
 - ✓ It is valid up to $W=2$ GeV.
 - ✓ It is a full kinematics model $d\sigma/dW dQ^2 d\Omega_\pi$
- It will soon be tested in T2K analyses, hopefully improved data/MC agreement and reducing systematic uncertainties propagated to the oscillation measurements.



Thank you!

Thanks **J .Zmuda** and **R. GONZALEZ JIMENEZ**
for validating my results.

Backup

General framework: how to calculate the helicity amplitudes?

$$\begin{aligned} \mathcal{M}_{CC}(vN \rightarrow l_\lambda N' \pi) &= \frac{G_F}{\sqrt{2}} \cos \theta_C \langle N' \pi | \epsilon_\lambda^\rho J_\rho | N \rangle \\ &= \frac{G_F}{\sqrt{2}} \cos \theta_C \langle N' \pi | C_{L_\lambda} e_L^\rho J_\rho + C_{R_\lambda} e_R^\rho J_\rho + C_\lambda e_\lambda^\rho J_\rho | N \rangle \end{aligned}$$

Hadron current $J_\alpha = J_\alpha^V - J_\alpha^A$

Lepton current

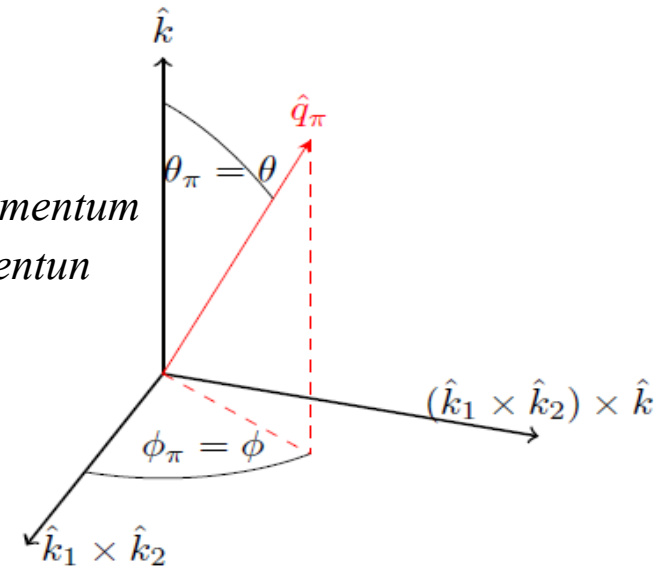
$$\epsilon^\alpha = \bar{u}_l(k_2) \gamma^\alpha (1 - \gamma_5) u_\nu(k_1)$$

can be interpreted as the intermediate gauge boson's polarization vector.

$$\epsilon_\lambda^\rho = C_{L_\lambda} e_L^\rho + C_{R_\lambda} e_R^\rho + C_\lambda e_\lambda^\rho$$

Four different polarizations $e_{\lambda_k}^\rho$

$\hat{k} = \hat{k}_1 - \hat{k}_2$
 k_1 : neutrino momentum
 k_2 : lepton momentum



$$e_L^\rho = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & -i & 0 \end{pmatrix}$$

$$e_R^\rho = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & -1 & -i & 0 \end{pmatrix}$$

$$e_\lambda^\rho = \frac{1}{\sqrt{|(\epsilon_\lambda^0)^2 - (\epsilon_\lambda^3)^2|}} \begin{pmatrix} \epsilon_\lambda^0 & 0 & 0 & \epsilon_\lambda^3 \end{pmatrix}$$

Hadronic Current

- Dirac equation allows us to have 16 independent Lorentz covariance
- Conservation of vector current reduce the number of $O(V_i)$ to six.

$$J_V^\rho e^{\lambda_k} = \sum_{i=1}^6 V_i \bar{u}_N(p_2) O^{\lambda_k}(V_i) u_N(p_1)$$

$$J_A^\rho e^{\lambda_k} = \sum_{i=1}^8 A_i \bar{u}_N(p_2) O^{\lambda_k}(A_i) u_N(p_1)$$

Dirac Spinors

$O(V)$ and $O(A)$ are 4*4 matrices in terms of Dirac matrices and particle's 4-momenta.

we can rewrite the right hand side of above eq. in terms of 2*2 matrices

$$J_V^\rho e^{\lambda_k} = \sum_{i=1}^6 \mathcal{F}_i \chi_2^* \Sigma_i^{\lambda_k} \chi_1$$

$$J_A^\rho e^{\lambda_k} = \sum_{i=1}^8 \mathcal{G}_i \chi_2^* \Lambda_i^{\lambda_k} \chi_1$$

Pauli spinors

Helicity Amplitudes

$$\tilde{F}_{\lambda_2, \lambda_1}^{\lambda_k} = \langle N\pi | e_{\lambda_k}^\rho J_\rho^V | N \rangle$$

$$\tilde{G}_{\lambda_2, \lambda_1}^{\lambda_k} = \langle N\pi | e_{\lambda_k}^\rho J_\rho^A | N \rangle$$

Can be defined by knowing the helicity of incident and outgoing nucleons and gauge boson's polarization. 16 helicity amplitudes for each vector and axial.

Cross-Section

One needs to calculate the helicity amplitudes for resonant and nonresonant interactions.

$$\frac{d\sigma(\nu N \rightarrow l N \pi)}{dk^2 dW d\Omega_\pi} = \frac{G_F^2}{2} \frac{1}{(2\pi)^4} \frac{|\mathbf{q}|}{4} \frac{-k^2}{(k^L)^2} \sum_{\lambda_2, \lambda_1} \left\{ \begin{aligned} & \left| C_{L-} (\tilde{F}_{\lambda_2 \lambda_1}^{eL} - \tilde{G}_{\lambda_2 \lambda_1}^{eL}) + C_{R-} (\tilde{F}_{\lambda_2 \lambda_1}^{eR} - \tilde{G}_{\lambda_2 \lambda_1}^{eR}) + C_- (\tilde{F}_{\lambda_2 \lambda_1}^{e-} - \tilde{G}_{\lambda_2 \lambda_1}^{e-}) \right|^2 \\ & + \left| C_{L+} (\tilde{F}_{\lambda_2 \lambda_1}^{eL} - \tilde{G}_{\lambda_2 \lambda_1}^{eL}) + C_{R+} (\tilde{F}_{\lambda_2 \lambda_1}^{eR} - \tilde{G}_{\lambda_2 \lambda_1}^{eR}) + C_+ (\tilde{F}_{\lambda_2 \lambda_1}^{e+} - \tilde{G}_{\lambda_2 \lambda_1}^{e+}) \right|^2 \end{aligned} \right\}$$

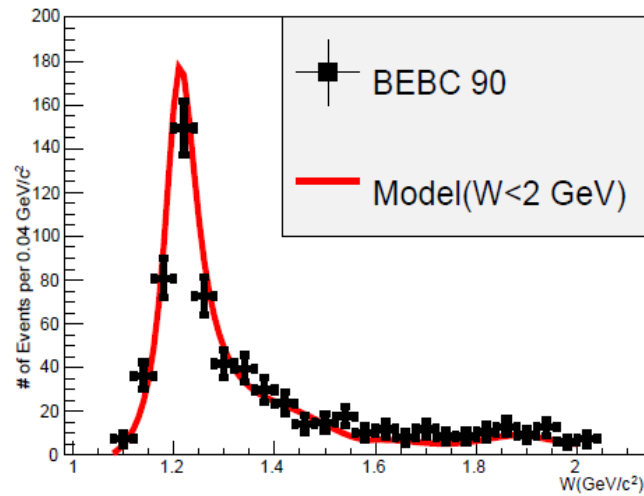
- Helicity amplitudes are a functions of W , $Q^2 = -k^2$ and pion angles in the Adler frame (θ, φ)
- The cross-section is used in several papers, but the lepton mass was ignored.

W -distribution

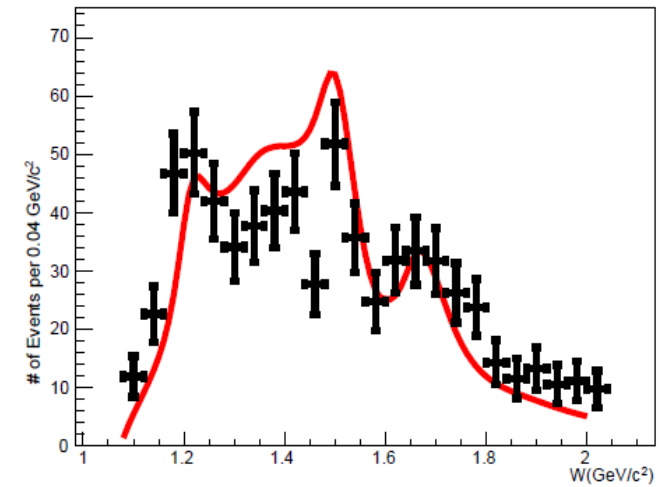
BEBC data

Hydrogen target

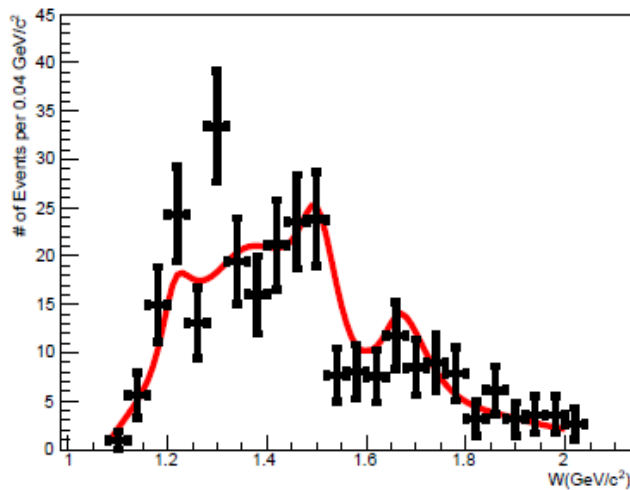
$\nu p \rightarrow \mu^- p \pi^+$



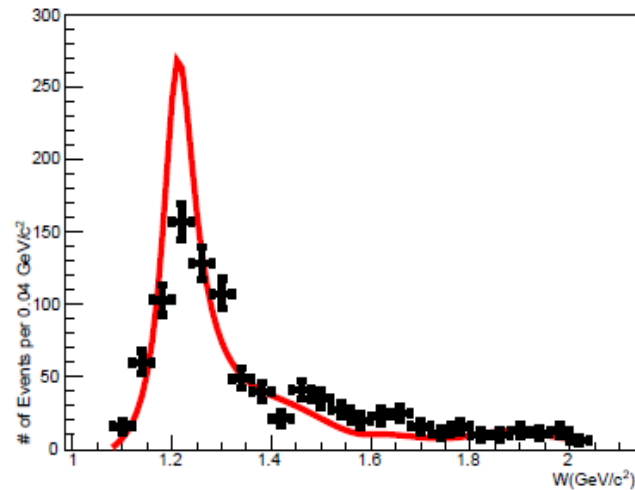
$\bar{\nu} p \rightarrow \mu^+ p \pi^-$



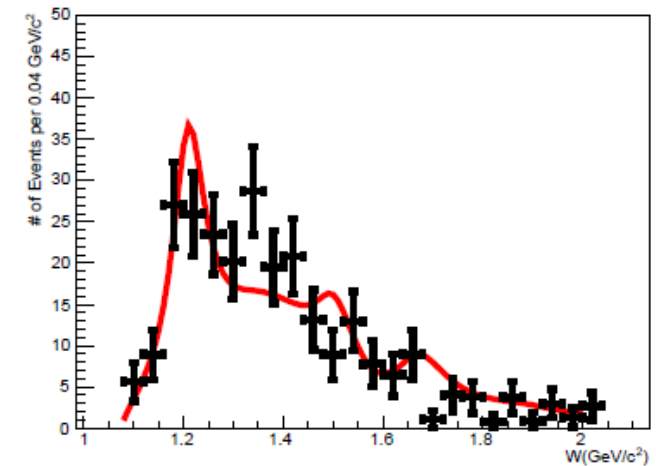
$\nu n \rightarrow \mu^- n \pi^+$



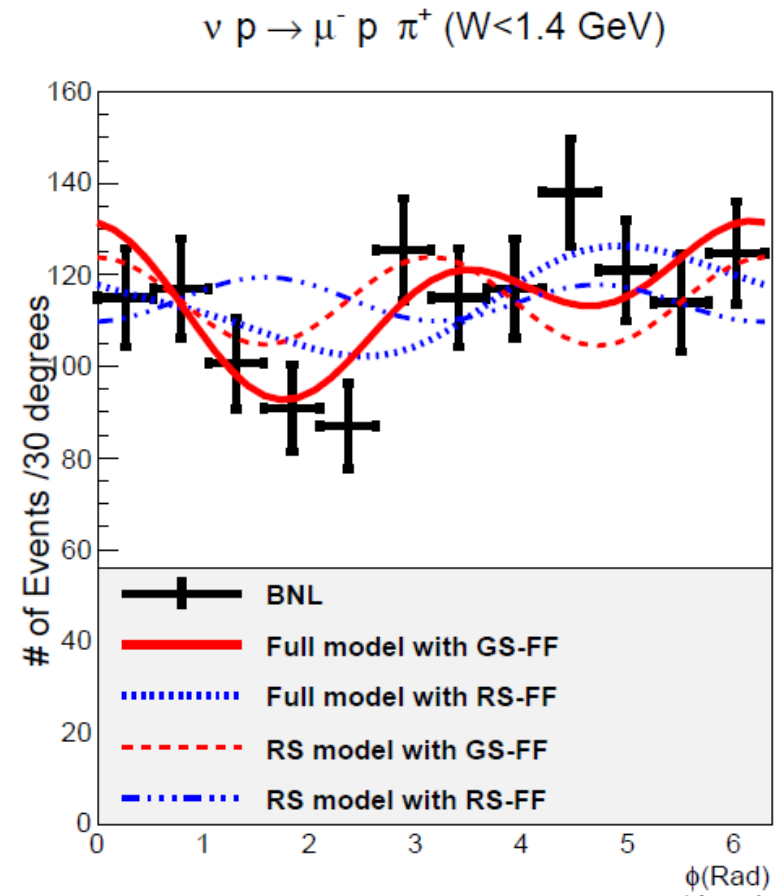
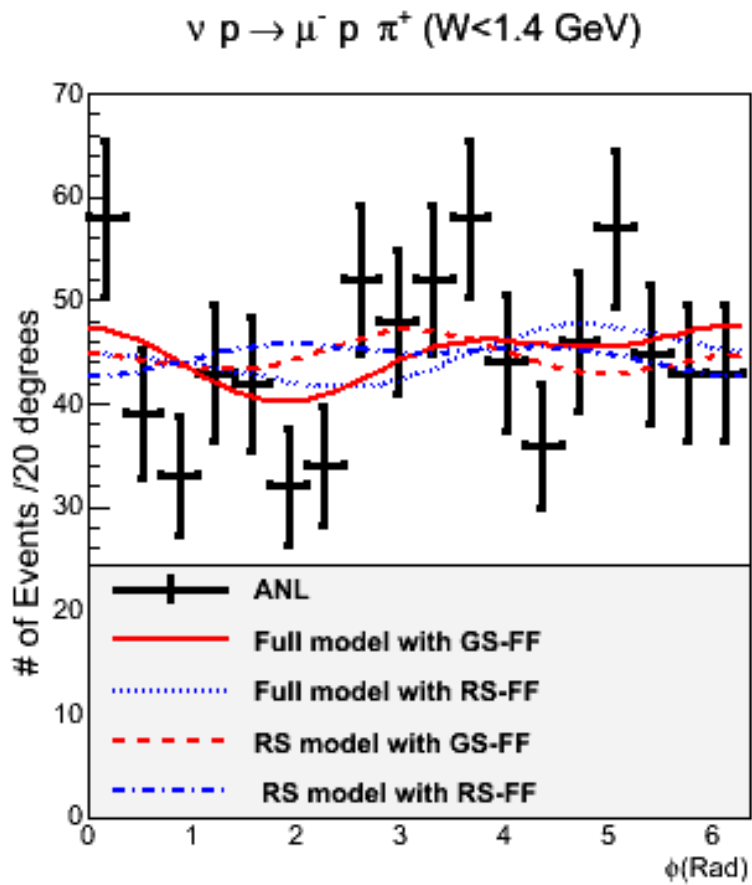
$\bar{\nu} n \rightarrow \mu^+ n \pi^-$



$\nu n \rightarrow \mu^- p \pi^0$



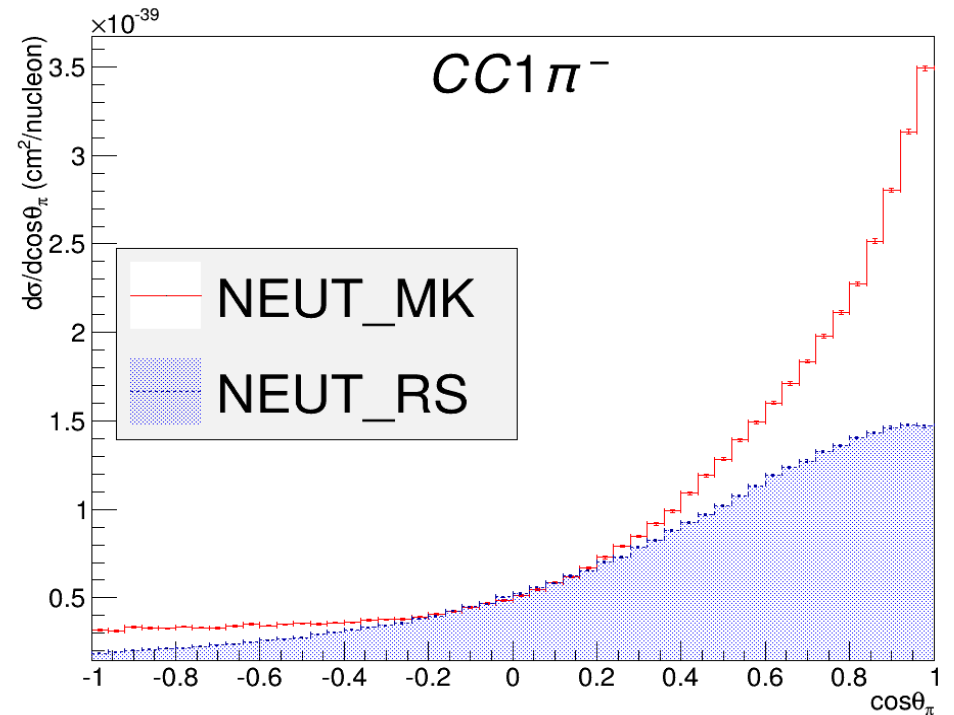
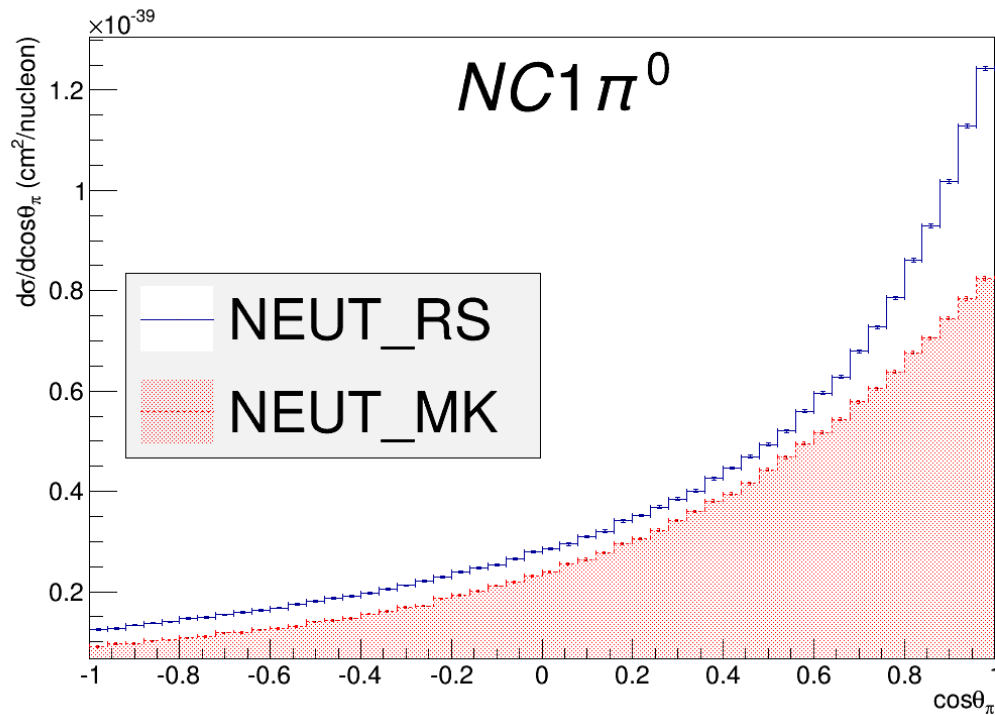
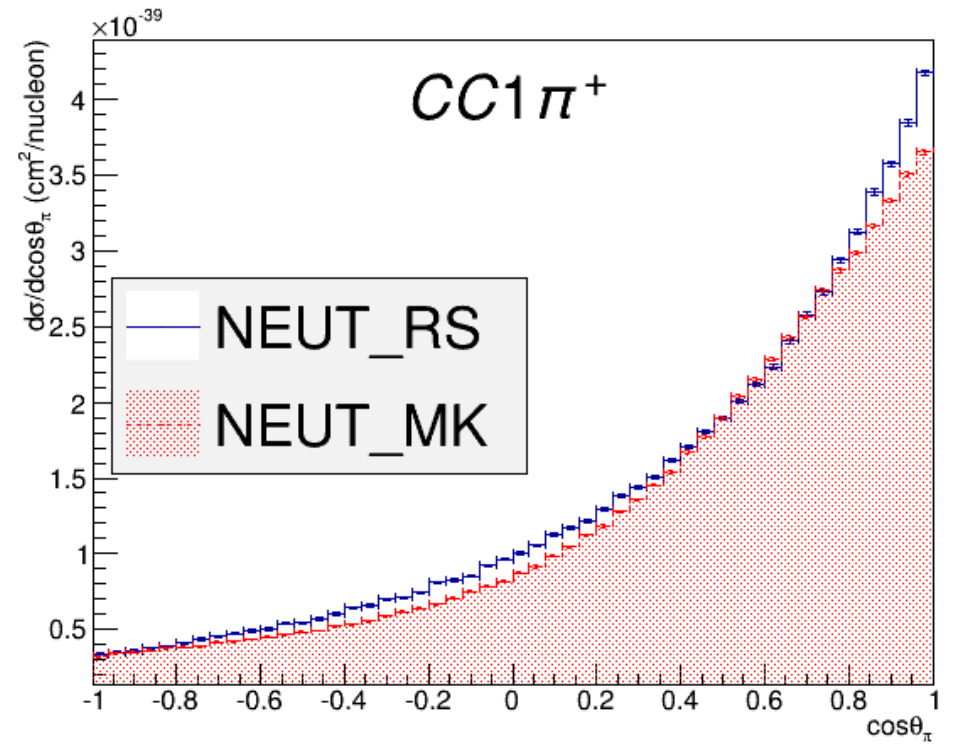
Pion azimuthal angle



Pion polar angle in the lab frame

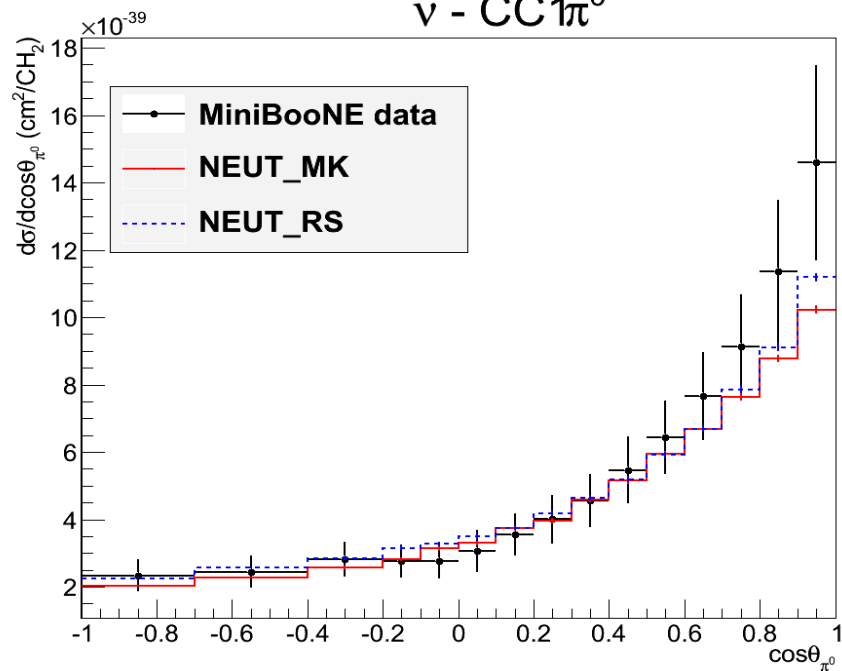
Nucleon target

Different samples

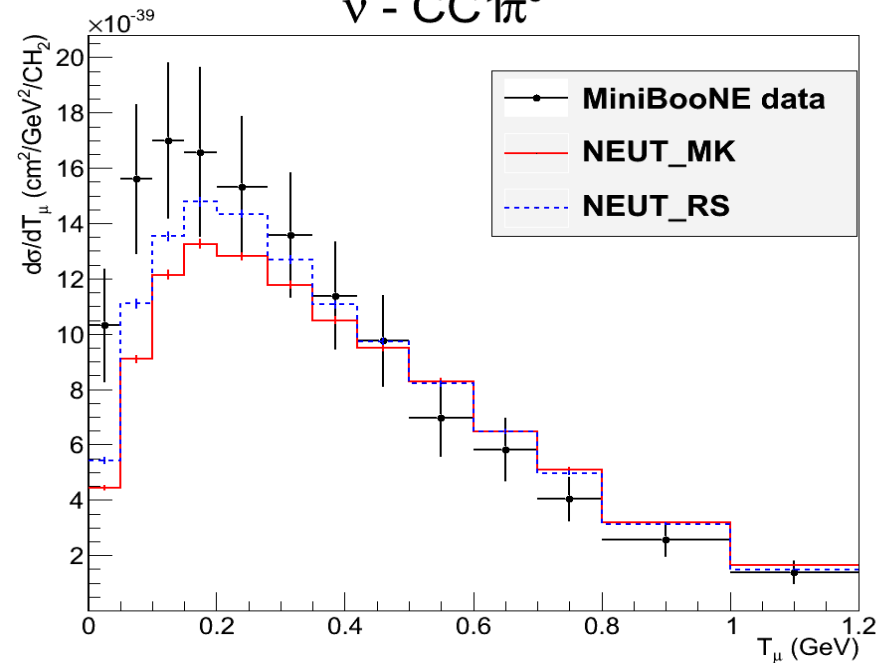


MiniBooNE data on CH2

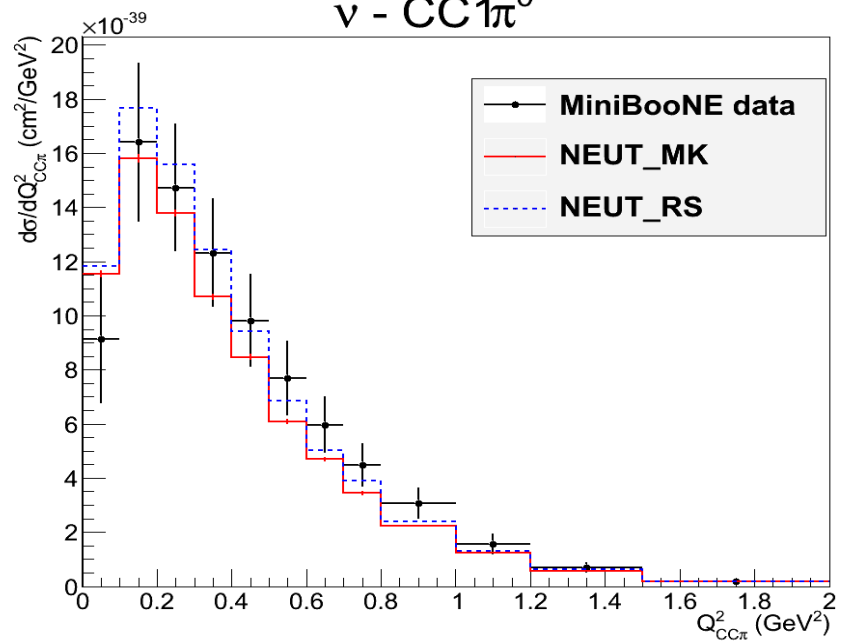
ν - CC1 π^0



ν - CC1 π^0



ν - CC1 π^0



ν - CC1 π^0

