## FROM SPIN TO STRUCTURE

Beam Spin Asymmetry in Exclusive Pion Electroproduction

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- Many unknowns in theory of strong force
- Meson electroproduction in Jefferson Lab Hall C probes hadron structure
- Use observables to study non-perturbative QCD in the transition regime



#### **Theoretical Approaches**





Compare observables to Regge and GPD models = test **relevant degrees of freedom** at given kinematics

### Beam Spin Asymmetry



3

$$A_{LU} = \frac{1}{P} \left( \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \right)$$

Difference in cross-sections based on helicity (±1) of the incident electron



A. Bacchetta et al. Phys. Rev. D 70, 117504 (2004). M. Diehl and S. Sapeta. Eur. Phys. J. C 41, 515-533 (2005).

## Beam Spin Asymmetry



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- Difference in cross-sections based on helicity (±1) of the incident electron
- Caused by interference between transversely and longitudinally polarized virtual photons  $\gamma^*$



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#### Beam Spin Asymmetry



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- Difference in cross-sections based on helicity (±1) of the incident electron
- Caused by interference between transversely and longitudinally polarized virtual photons
- Acceptances cancel in a ratio

My research: asymmetry analysis of the reaction:

$$e^- + p \longrightarrow e^{-\prime} + \pi^+ + n$$

Extraction of  $\sigma_{LT'}/\sigma_0$  over a range of kinematics

M. Diehl and S. Sapeta. Eur. Phys. J. C 41, 515-533 (2005).



 $\mathbf{Q^2}$ : 4-momentum of  $\gamma^*$ 

**x**<sub>*B*</sub>: momentum fraction of struck parton

-t: 4-momentum transfer from  $\gamma^*$  to meson

 $\epsilon: \ \mbox{longitudinal to transverse} \ \gamma^* \ \mbox{flux ratio} \label{eq:expansion}$ 





 $\mathbf{Q}^2$ : 4-momentum of  $\gamma^*$  $Q^2 = -(\mathbf{p}_e - \mathbf{p}'_e)^2$ 

 $\mathbf{x}_{B}$ : momentum fraction of struck parton

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 $\mathbf{Q}^2$ : 4-momentum of  $\gamma^*$  $Q^2 = -(oldsymbol{p}_e - oldsymbol{p}'_e)^2$ 

**x**<sub>B</sub>: momentum fraction of struck parton

$$x_B = \frac{Q^2}{2\boldsymbol{p}_p \cdot \boldsymbol{p}_{\gamma^*}}$$

-t: 4-momentum transfer from  $\gamma^*$  to meson

 $\epsilon :$  longitudinal to transverse  $\gamma^*$  flux ratio





**Q**<sup>2</sup>: 4-momentum of  $\gamma^*$  $Q^2 = -(\pmb{p}_e - \pmb{p}'_e)^2$ 

**x**<sub>*B*</sub>: momentum fraction of struck parton

$$x_B = \frac{Q^2}{2\boldsymbol{p}_p \cdot \boldsymbol{p}_{\gamma^*}}$$

-t: 4-momentum transfer from  $\gamma^*$  to meson

 $-t=-(oldsymbol{p}_{\gamma^*}-oldsymbol{p}_{\pi})^2$ 

 $\epsilon :$  longitudinal to transverse  $\gamma^*$  flux ratio





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## Welcome to Hall C!



#### **Hall C:** electron beam $\rightarrow$ fixed target $\rightarrow$ spectrometers



- $\blacksquare Spectrometers are magnetic and moveable \rightarrow choose charge, momentum, and angles to detect$
- Coincidence experiment: need simultaneous detection in High Momentum Spectrometer and Super HMS

## Welcome to Hall C!



#### $\blacksquare Hall C: electron beam \rightarrow fixed target \rightarrow spectrometers$





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## Particle Identification (PID)



- Spectrometer detector stack contains drift chambers for tracking, hodoscopes for triggering, threshold Cherenkovs and calorimeter for PID
- Fixed charge, momentum: PID via mass separation
- $\rightarrow\,$  Choose index of refraction to distinguish between particles



#### **Event Selection**



$$e^- + p \rightarrow e^{-\prime} + \pi^+ + n$$

- Select coincidences via  $t_{COIN} = t_{HMS} t_{SHMS} \approx 0$
- Neutron not detected ightarrow use missing mass  $m_X pprox m_N$

$$m_X^2 = (\boldsymbol{p}_e + m_p - \boldsymbol{p}_{e\prime} - \boldsymbol{p}_{\pi})^2$$

Subtract random time sample, empty target background







## Asymmetry $Q^2=2.1$ GeV<sup>2</sup>, $x_B = 0.21$





## Asymmetry $Q^2 = 3.0 \text{ GeV}^2$ , $x_B = 0.40$





Systematic Errors



## 1. Fitting Error

- Difference in σ<sub>LT'</sub>/σ<sub>0</sub> extracted using full (solid line) or approximated (dashed) fit
- Has a direction → total systematic error asymmetric
- 2. Cut Dependence
  - $\sigma_{LT'}/\sigma_0$  varies with exact values of missing mass and coincidence time cuts
- 3. Beam Polarization
  - $\blacksquare \ P = 89^{+1}_{-3}\% \rightarrow \text{Propagate to} \\ \sigma_{LT'}/\sigma_0$



Results





S.V. Goloskokov, P. Kroll, Eur. Phys. J. C 65 137 (2010).
 B. Berthou et al, Eur. Phys. J. C 78 478 (2018).
 T. Vrancx, J. Ryckebusch & J. Nys, Phys. Rev C, 89 065202 (2014).
 S. Diehl et al., Phys. Lett. B 839, 137761 (2023).
 T. K. Choi, K.-J. Kong & B.-G. Yu, J. Korean Phys. Soc. 67, 1089-1094 (2015).

























## Comparison with Theory (2)





# Best overall agreement is **YCK1**: Regge model with GPD parametrization of nucleon electromagnetic form factors (EMFFs)

S.V. Goloskokov, P. Kroll, Eur. Phys. J. C 65 137 (2010).
 B. Berthou et al, Eur. Phys. J. C 78 478 (2018).
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Measurements of \$\sigma\_{LT'}\$/\$\sigma\_0\$ from KaonLT, CLAS, and CLAS12
 KaonLT data extends kinematic range with finer t-binning
 Combine to determine \$\mathbb{Q}^2\$ dependence at fixed \$(x\_B, -t)\$)



S. Diehl et al., Phys. Rev. Lett. 125, 182001 (2020).

## **Q**<sup>2</sup> **Dependence (New!)**





σ<sub>LT'</sub>/σ<sub>0</sub> from KaonLT, CLAS, and CLAS12 as a function of Q<sup>2</sup>
 Flat or weak Q<sup>2</sup> dependence

S. Diehl et al., Phys. Lett. B 839, 137761 (2023). S. Diehl et al., Phys. Rev. Lett. 125, 182001 (2020). T. K. Choi, K.-J. Kong & B.-G. Yu, J. Korean Phys. Soc. 67, 1089-1094 (2015).

## **Physics Interpretation**



CLAS12 conclusions: best fit is GK2 (GPD  $H_T$ \*2)

 $\rightarrow$  GPD picture applicable

KaonLT conclusions: best fit is YCK1 (Regge + GPD EMFF)



- Our measured  $\sigma_{LT'}/\sigma_0$  is not explained by only quark and gluon degrees of freedom
- Hadronic degrees of freedom may be more relevant
- **TYCK1** uses GPDs in EMFF parametrization  $\rightarrow$  hybrid approach?
- ★ Need model-independent tests of GPD picture (see next talk)

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 B. Berthou et al, Eur. Phys. J. C 78 478 (2018).
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- Measured  $A_{LU}$  in  $e + p \rightarrow e' + \pi^+ + n$  and extracted  $\sigma_{LT'} / \sigma_0$ from KaonLT data over range of kinematics
- No exact agreement with predictions, closest is YCK1 (Regge + GPD EMFF)
- Flat or weak  $Q^2$  dependence of  $\sigma_{LT'}/\sigma_0$
- Manuscript being prepared for Physics Letters B

Precision data of hadronic reaction observables critical for proton structure and the strong force!

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## **EXTRA SLIDES**

## What If...





■ No Q<sup>2</sup> dependence → overlay curves at same x<sub>B</sub>
 ■ Seems to show same -t dependence within uncertainties

## **CLAS12** Results





GK (default) GK ( $H_T$ \*1.5) GK ( $H_T$ \*2) JML (Regge) GK (no pion pole)

S. Diehl et al., Phys. Lett. B 839, 137761 (2023). arXiv:2210.14557

## **Quick Derivation**



**D**efine the beam spin asymmetry  $A_{LU}$  as:

$$A_{LU} = \frac{1}{P} \left( \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \right) = \frac{1}{P} \left( \frac{Y^+ - Y^-}{Y^+ + Y^-} \right)$$

Polarized cross-section in Rosenbluth equation appears when separating events by helicity:

$$2\pi \frac{d^2\sigma}{dtd\phi} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{LT}}{dt} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi + h\sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{LT'}}{dt} \sin\phi$$

Beam spin asymmetry provides much cleaner access to  $\sigma_{LT'}$ :

$$A_{LU} = \frac{\sqrt{2\epsilon(1-\epsilon)}\frac{\sigma_{LT'}}{\sigma_0}\sin\phi}{1+\sqrt{2\epsilon(1+\epsilon)}\frac{\sigma_{LT}}{\sigma_0}\cos\phi + \epsilon\frac{\sigma_{TT}}{\sigma_0}\cos2\phi}$$

M. Diehl and S. Sapeta. Eur. Phys. J. C 41, 525-533 (2005).

## **SHMS Focal Plane Detectors**





Photo by N. Heinrich.

Detector	Purpose
Ar/Ne Cherenkov	Not installed
Drift chambers	Tracking
Hodoscopes	Triggering, tracking
$C_4F_8O$ Cherenkov	Particle identification
Aerogel Cherenkov	Particle identification
Shower counters	Calorimetry

S. Ali et al, to be published (2022).



## -t Binning



Sum all events at one  $(Q^2, x_B)$  and separate into -t bins with similar numbers of events



## **Combining SHMS Settings**





## GPDs and Hard/Soft Factorization





GPD models rely explicitly on hard/soft factorization: process can be written as convolution of hard (perturbative) scattering and soft (non-perturbative) object

- GPDs encode 3D nucleon structure information: extraction is of high interest
- **QCD** predicts factorization at **"sufficiently high"**  $Q^2$
- Experimental data needed to identify onset of factorization
- R. Devenish and A. Cooper-Sarkar, Deep Inelastic Scattering (Oxford University Press, 2004). Figure by N. Heinrich.

## **Regge Models**



## **Regge trajectories**

Linear relationship observed between mass  $M^2$  and spin  $\alpha$  for baryons of the same quark content

 Feynman propagator replaced with Regge propagator

- Exchange of a series of particles along a Regge trajectory
- My reaction: mesons exchanged, typically π and ρ propagators
- Cutoff is a free parameter in many models

T. Regge. Introduction to complex orbital momenta. Nuovo Cim, 14:951-976, 1959



W. Li, Exclusive Backward-Angle Omega Meson Electroproduction, Ph.D. thesis, University of Regina (2017).

## **Model Details**



- Vrancx-Ryckebusch (**VR**): exchange of  $\pi(140)$ ,  $\rho(770)$ , and  $a_1(1260)$  **Regge** trajectories
- Goloskokov-Kroll (GK): uses twist-2 longitudinal (*Ẽ*, *H̃*) and twist-3 transverse (*E<sub>T</sub>*, *H<sub>T</sub>*) GPDs, with pion pole contributions. GK1: default GK model
  - **GK2:** modification  $H_T \rightarrow H_T * 2$ , as seen in CLAS12 BSA paper
- Yu-Choi-Kong (YCK): Regge trajectories, including tensor meson a<sub>2</sub>(1320) and axial mesons a<sub>1</sub> and b<sub>1</sub>(1235), with pion pole contributions.

YCK are co-authors on this paper.

**YCK1:** nucleon EMFFs mediated by GPDs **YCK2:** nucleon EMFFs use dipole parametrization

T. Vrancx, J. Ryckebusch & J. Nys, Phys. Rev C, 89 065202 (2014). arXiv:1310.7715
S.V. Goloskokov, P. Kroll, Eur. Phys. J. C 65 137 (2010). arXiv:1106.4897
T. K. Choi, K.-J. Kong & B.-G. Yu, J. Korean Phys. Soc. 67, 1089-1094 (2015). arXiv:1508.00969
S. Diehl et al., Phys. Lett. B 839, 137761 (2023). arXiv:2210.14557



- CEBAF produces polarized beam up to 12 GeV
- Polarization flipped at 30 Hz in a pseudo-random sequence
  - I Mott polarimeter at injector gives source polarization:  $90 \pm 1\%$
- Spin precession calculation shows Hall C receives 99% of the source polarization
- Final value P = 89<sup>+1</sup><sub>-3</sub>%: Uncertainty from the beam energy uncertainty (3.6 MeV) and the range of possible linac energy imbalance

Thanks to Steve Wood for polarization calculations and Dave Gaskell for uncertainty estimate