

Probing Lepton Number Violation at Same-Sign Lepton Colliders

Douglas Tuckler

TRIUMF & Simon Fraser University

dtuckler@triumf.ca

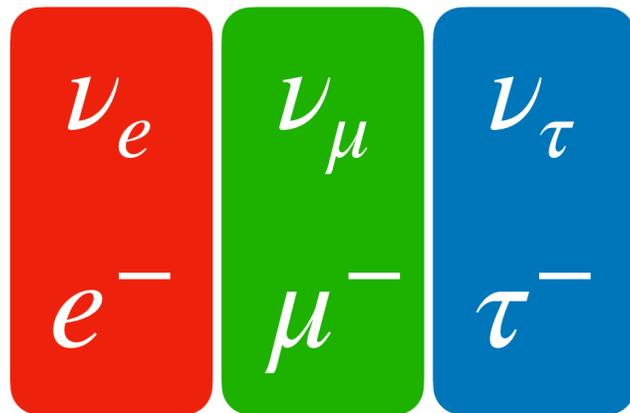
WNPPC 2025

Feb. 16, 2025

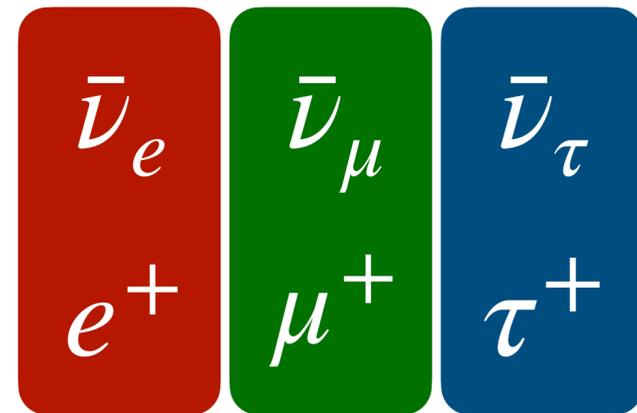
Lepton Number in the Standard Model

Lepton number (L) tracks the *number of leptons in a system*

Leptons: $L = +1$



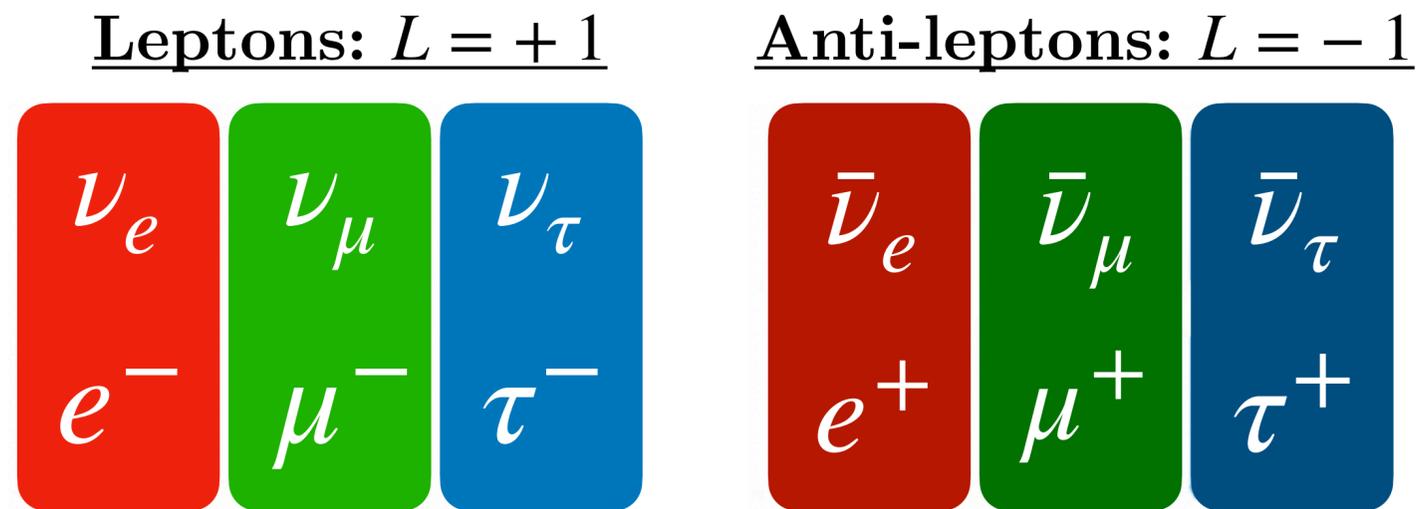
Anti-leptons: $L = -1$



Can also define *lepton family numbers* for each generation L_e , L_μ , L_τ so that total lepton number is $L = L_e + L_\mu + L_\tau$

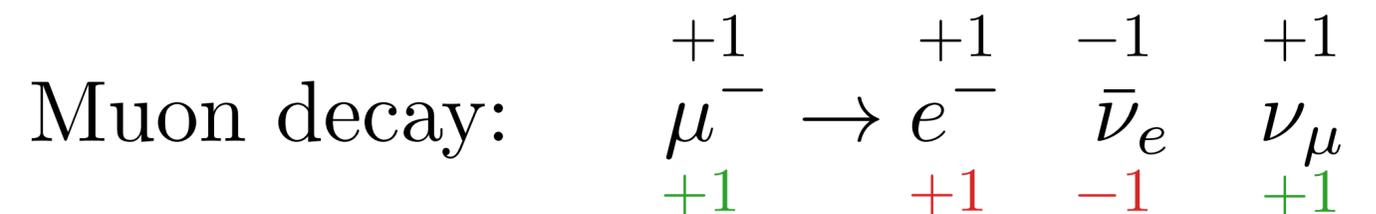
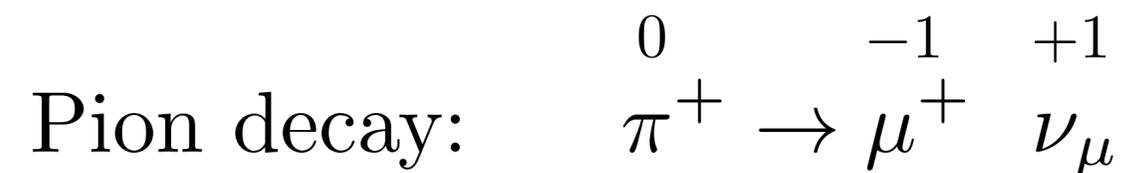
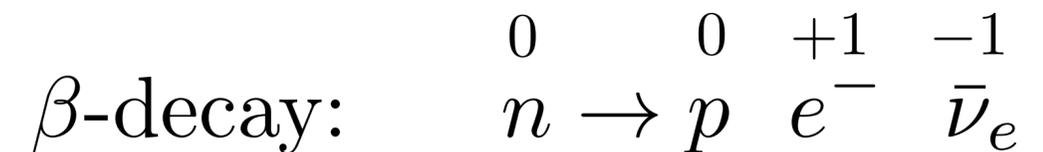
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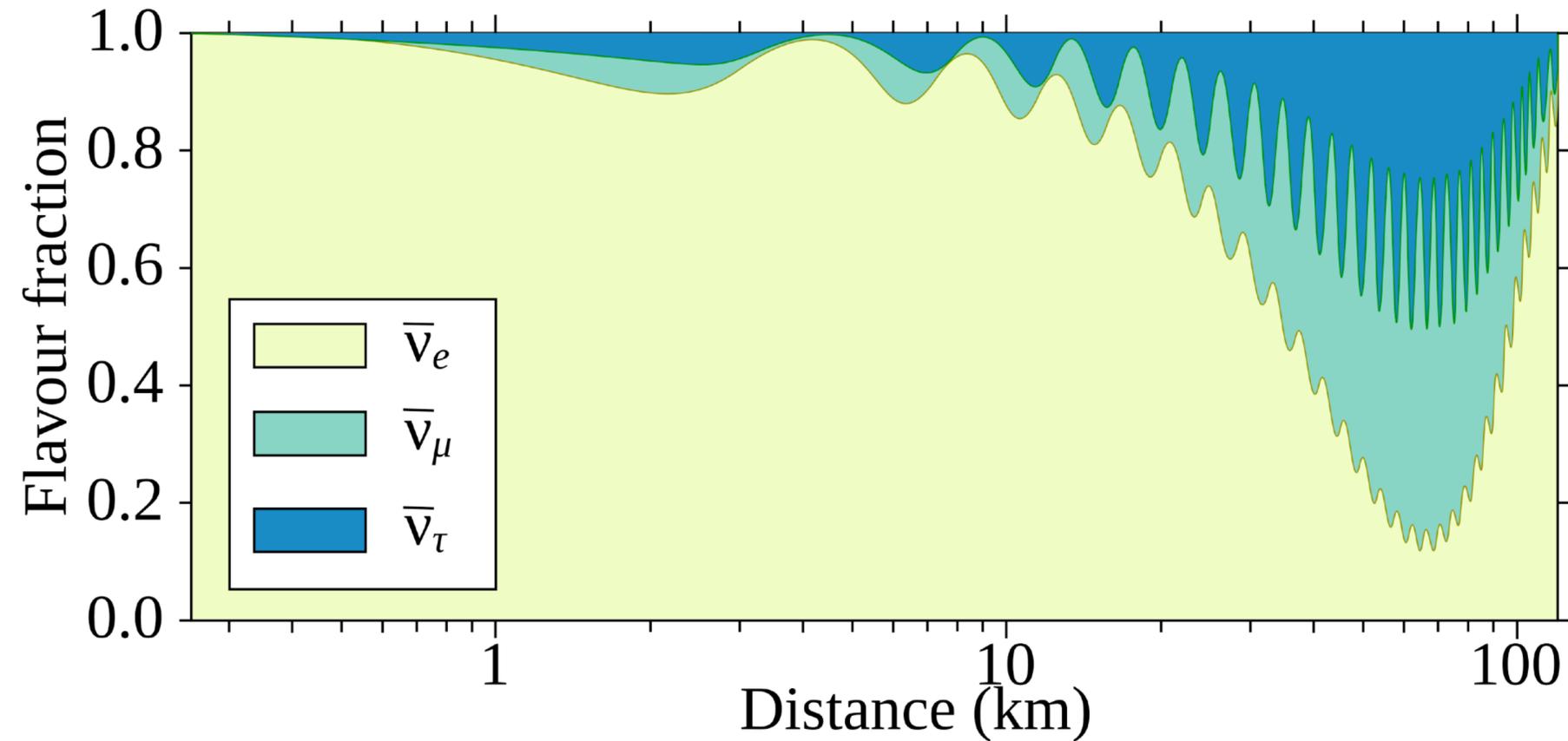
Lepton number is observed to be conserved
— *accidental symmetry of the SM*



Strong bounds on LNV processes. See [PDG — Tests of Conservation Laws](#)

However, there are fundamental reasons for lepton number violation in our universe.

Neutrino Oscillations



$$P(\nu_\alpha \rightarrow \nu_\beta) \simeq |U_\alpha|^2 |U_\beta|^2 \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$

Neutrino oscillations only possible if *neutrinos have mass!*

Some open questions:

- 1) Neutrino oscillations violate lepton flavor. Does this imply that lepton number is not conserved?
- 2) How do neutrinos get mass?
Neutrinos are massless in the SM.

Neutrino Masses: The Seesaw Mechanism

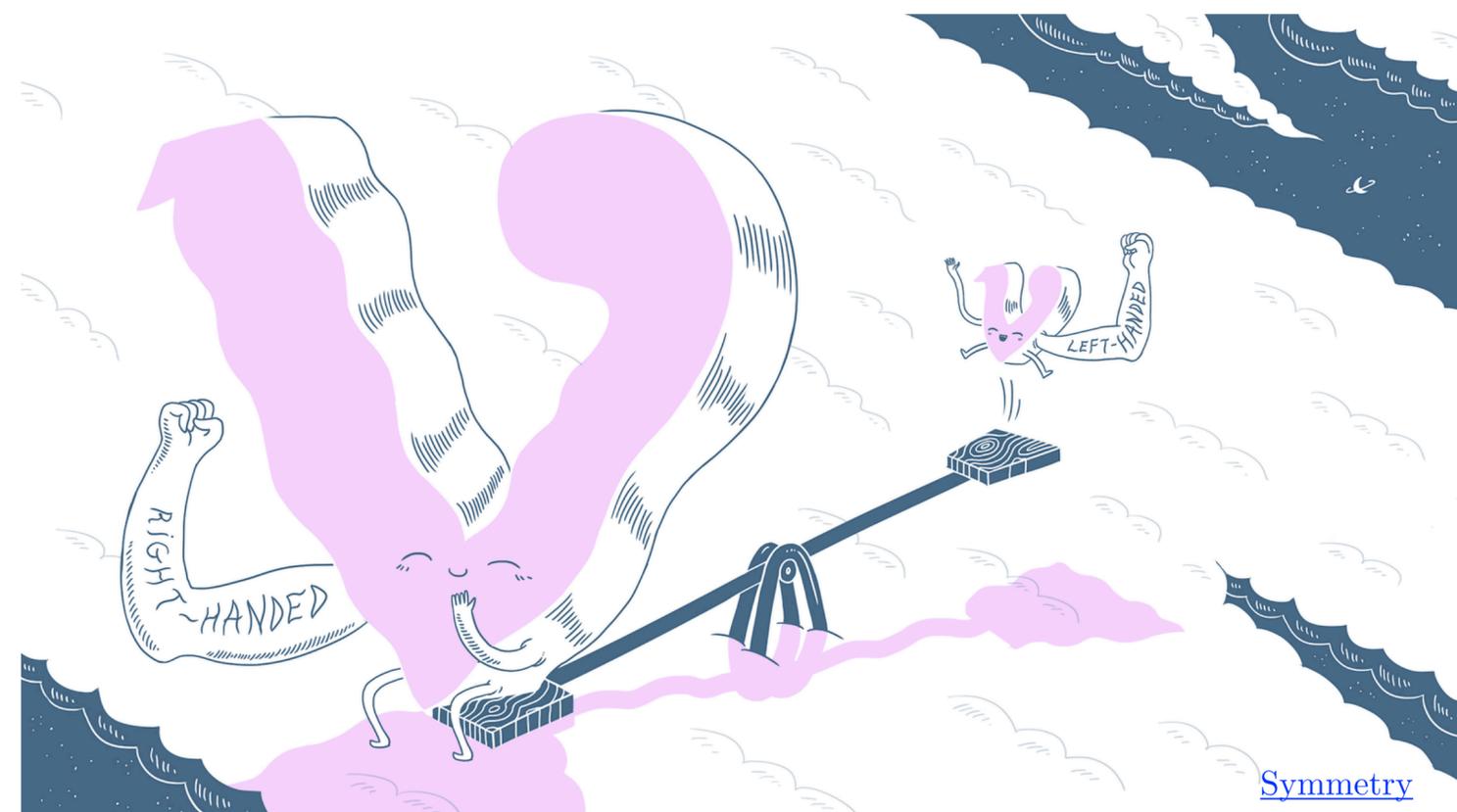
- Introduce right-handed neutrinos/sterile neutrinos/heavy neutral leptons N_R that are:
 - 1) SM singlets — no electric charge, but have lepton number $L = 1$
 - 2) Majorana fermions — their own anti-particles

$$\mathcal{L} \supset Y \bar{L} \tilde{H} N_R + \frac{1}{2} M_N \bar{N}_R N_R \xrightarrow{\text{EWSB}} \overset{\text{Dirac Mass}}{m_D} \bar{\nu}_L N_R + \overset{\text{Majorana Mass}}{M_N} \bar{N}_R N_R$$

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D \\ m_D & M_N \end{pmatrix} \rightarrow m_\nu \simeq \frac{Y^2 v^2}{M_N}$$

Neutrino masses *inversely proportional to Majorana mass*

$$m_\nu \sim 0.1 \text{ eV}, Y \sim 1 \implies M_N \sim 10^{14} \text{ GeV}$$



[Symmetry](#)

HNLs and Lepton Number Violation

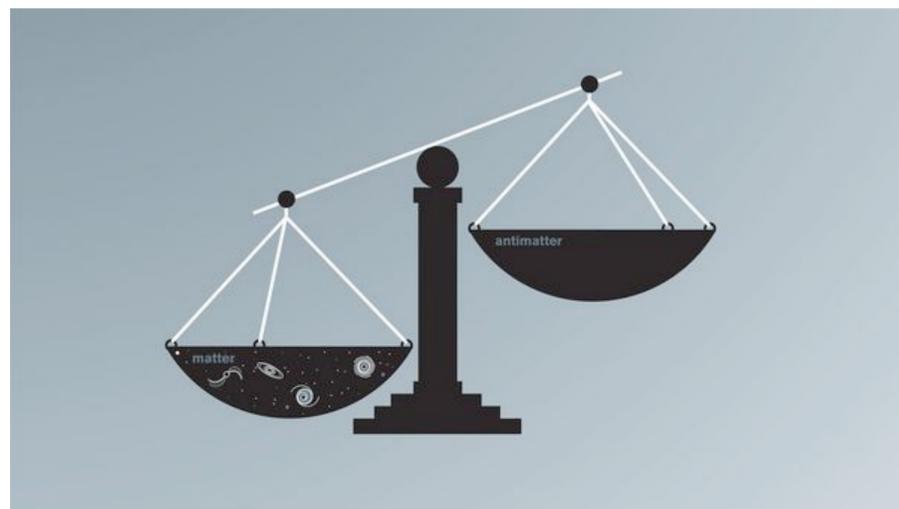
$$\mathcal{L} \supset Y \bar{L} \tilde{H} N_R + \frac{1}{2} M_N \bar{N}_R N_R \xrightarrow{\text{EWSB}} m_D \bar{\nu}_L N_R + \boxed{M_N \bar{N}_R N_R}$$

Marjorana Mass
Violates Lepton Number!

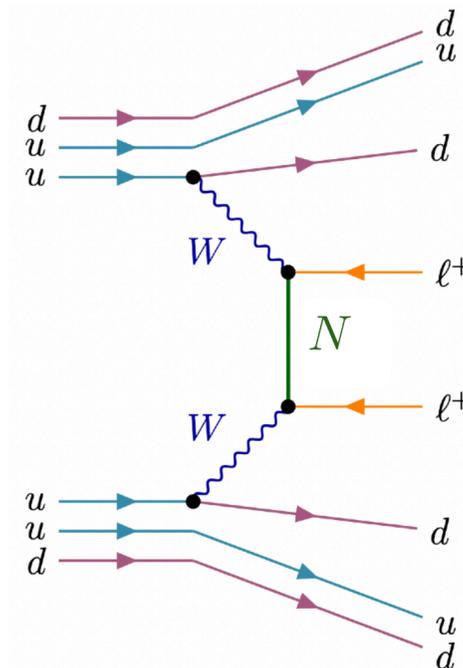
Implications of LNV by heavy neutral leptons:

1) Matter-Antimatter Asymmetry

$$\Gamma(N \rightarrow \ell) \neq \Gamma(N \rightarrow \bar{\ell})$$



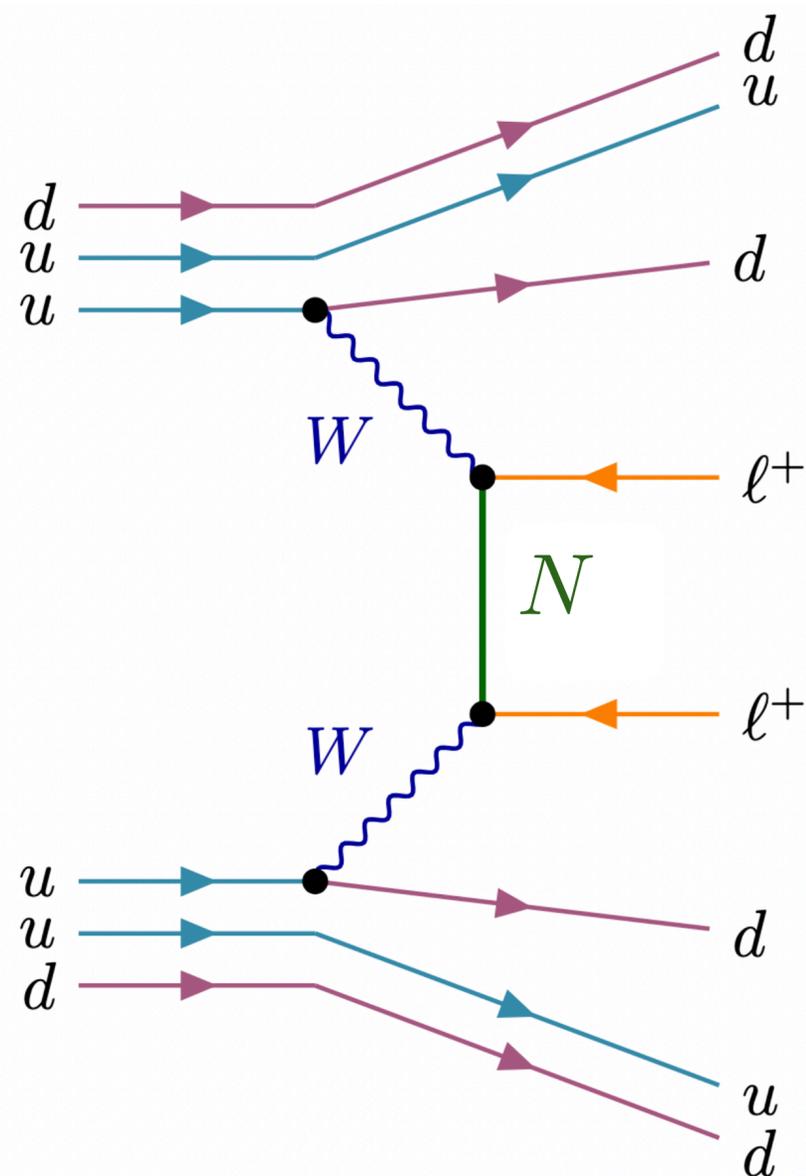
2) Neutrinoless Double Beta Decay



Connected to the question of whether neutrinos are Dirac or Majorana fermions

Neutrinoless Double Beta Decay

- **Observation of $0\nu\beta\beta$ would be clear evidence of LNV and physics beyond-the-SM!**
- Receives contributions from Majorana HNLs



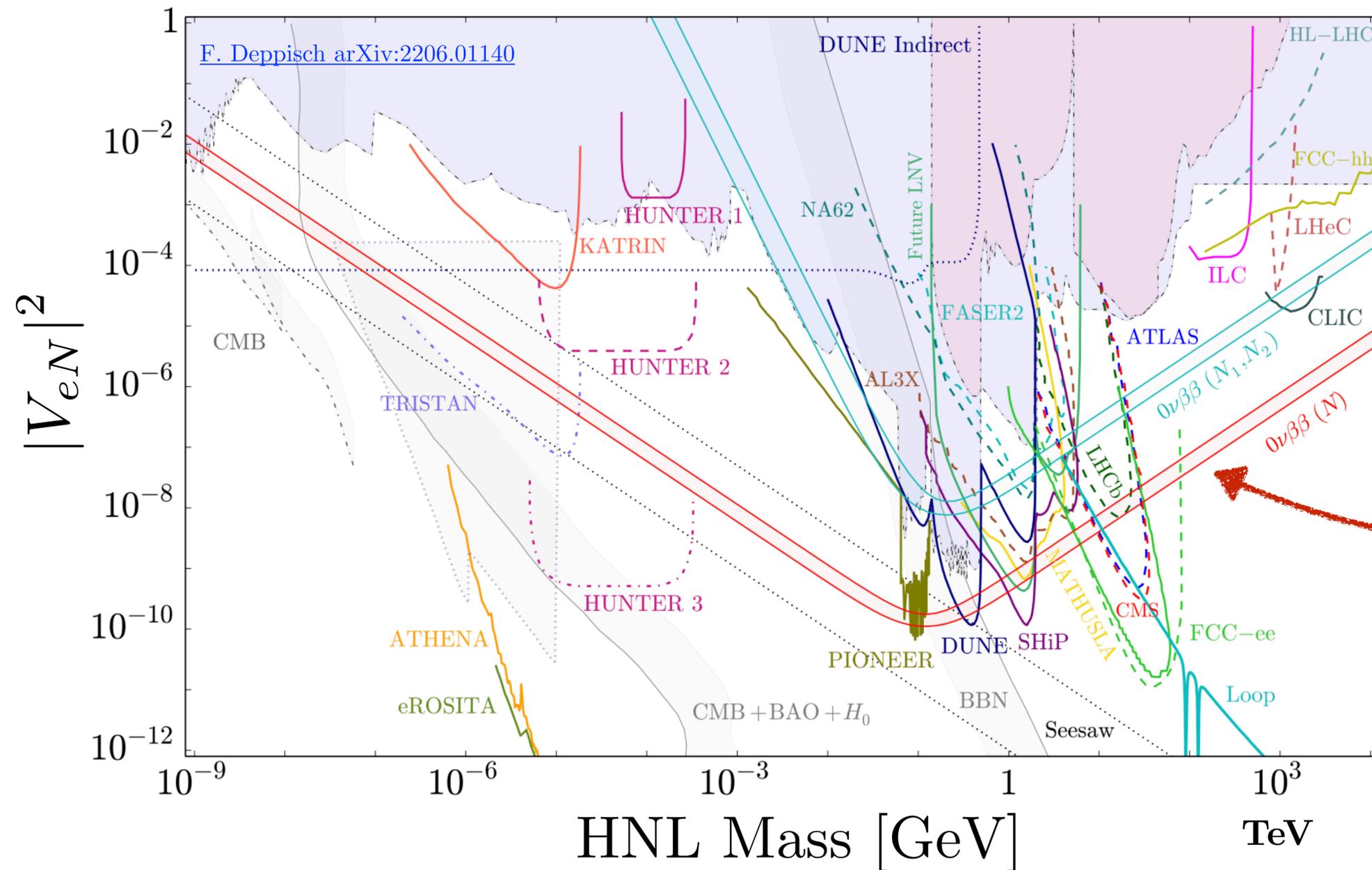
$$(T_{1/2}^{0\nu})^{-1} \sim \underbrace{G_{0\nu} g_A^4 m_p^2 |\mathcal{M}_N^{0\nu}|^2}_{\text{Nuclear physics stuff ?}} \left| \sum_{i=1}^3 \underbrace{\frac{U_{ei}^2 m_i}{\langle \mathbf{p}^2 \rangle}}_{\text{SM neutrino}} + \underbrace{\frac{V_{eN}^2}{M_N}}_{\text{HNL}} \right|^2$$

Existing Bounds on $0\nu\beta\beta$

Experiment	Isotope	$T_{1/2}^{0\nu}/10^{25}\text{yr}$ (90%C.L.)
GERDA	^{76}Ge	18.0
MAJORANA	^{76}Ge	8.30
CUORE	^{130}Te	2.61
KamLAND-Zen	^{136}Xe	22.86
EXO-200	^{136}Xe	3.50 <small>arXiv:2404.12316</small>

Neutrinoless Double Beta Decay

Electron-Mixed HNL

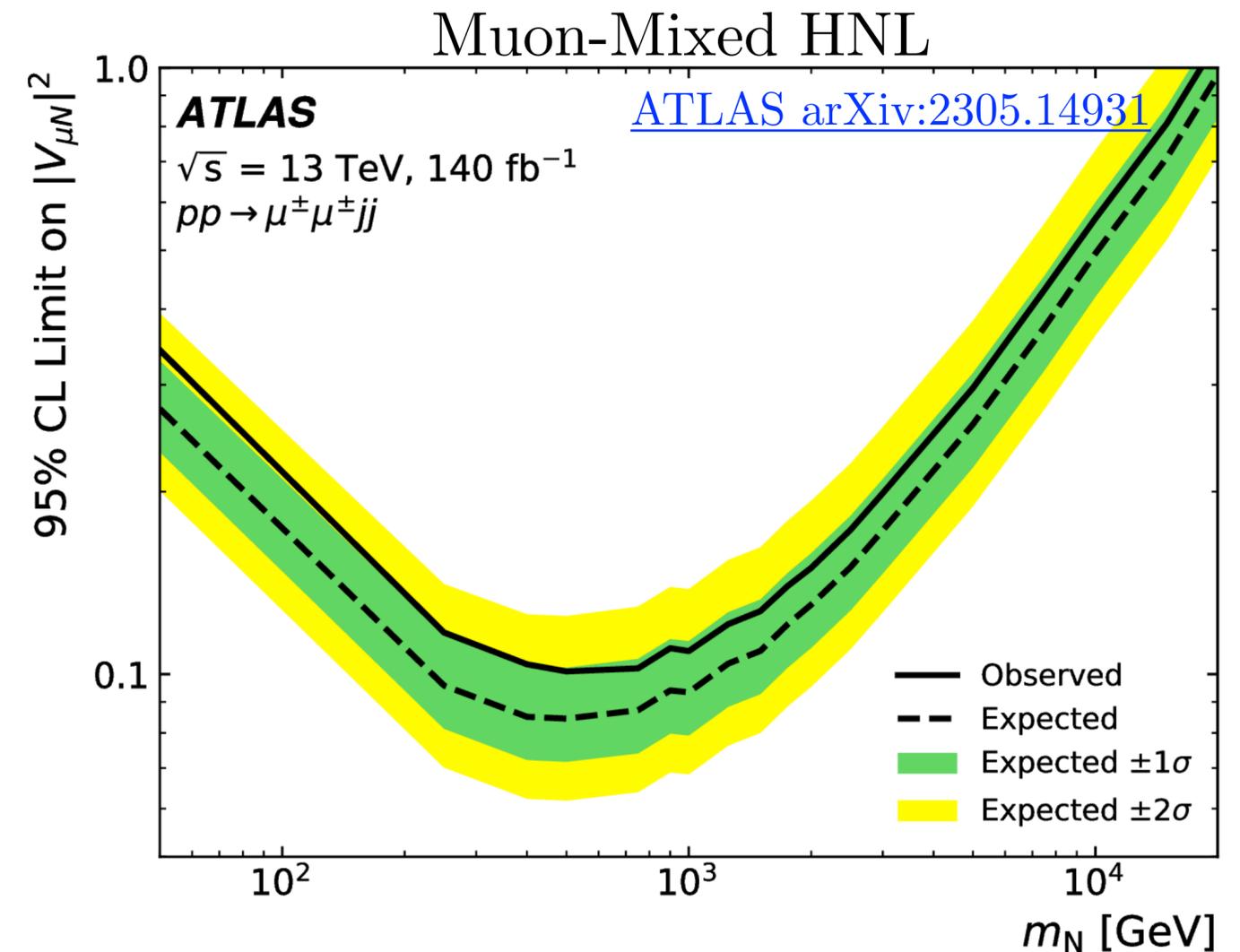
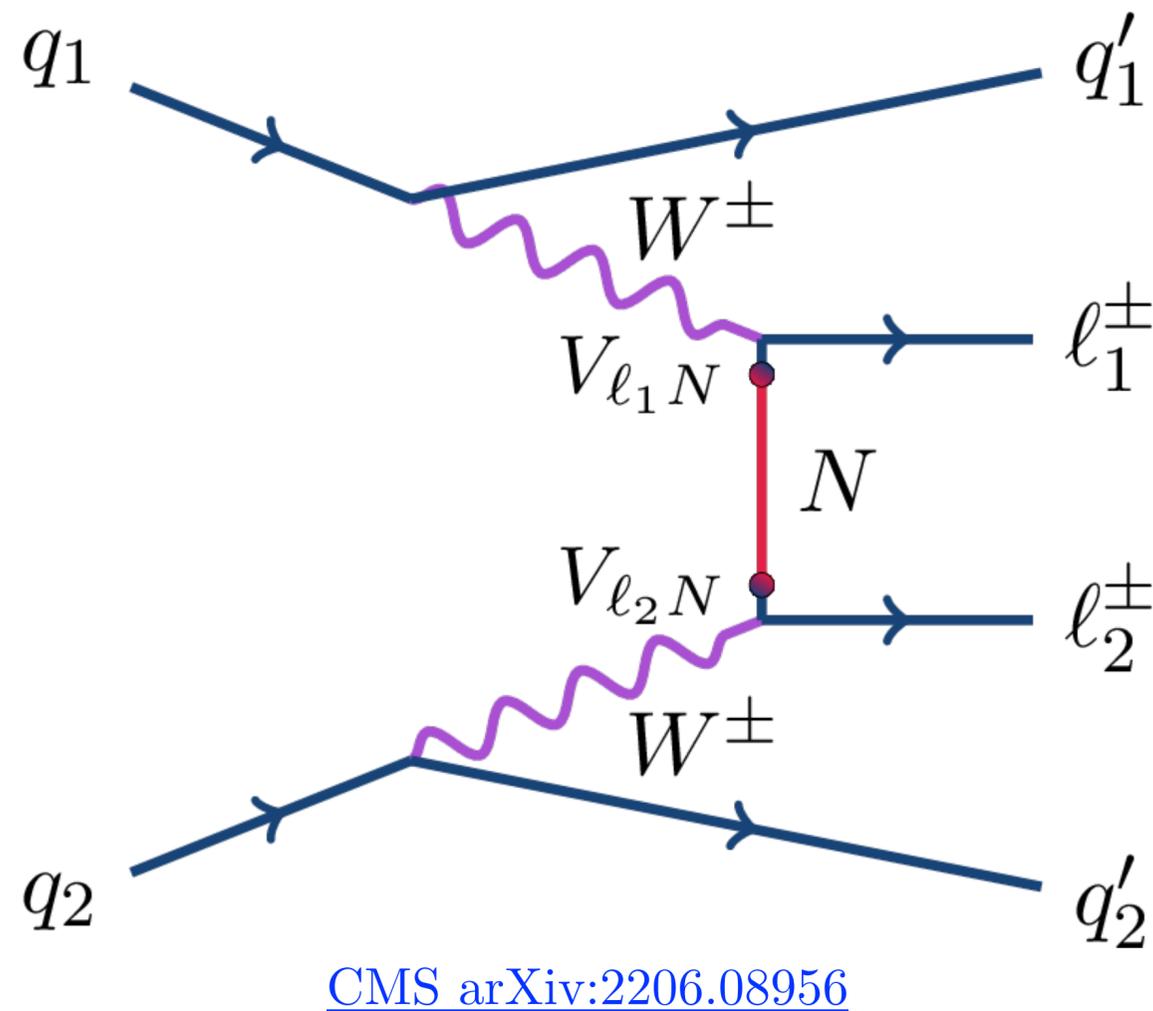


Probe of electron-mixed HNLs beyond the TeV-scale!

Only for electron-mixing. *What about muon- or tau-mixed HNLs?*

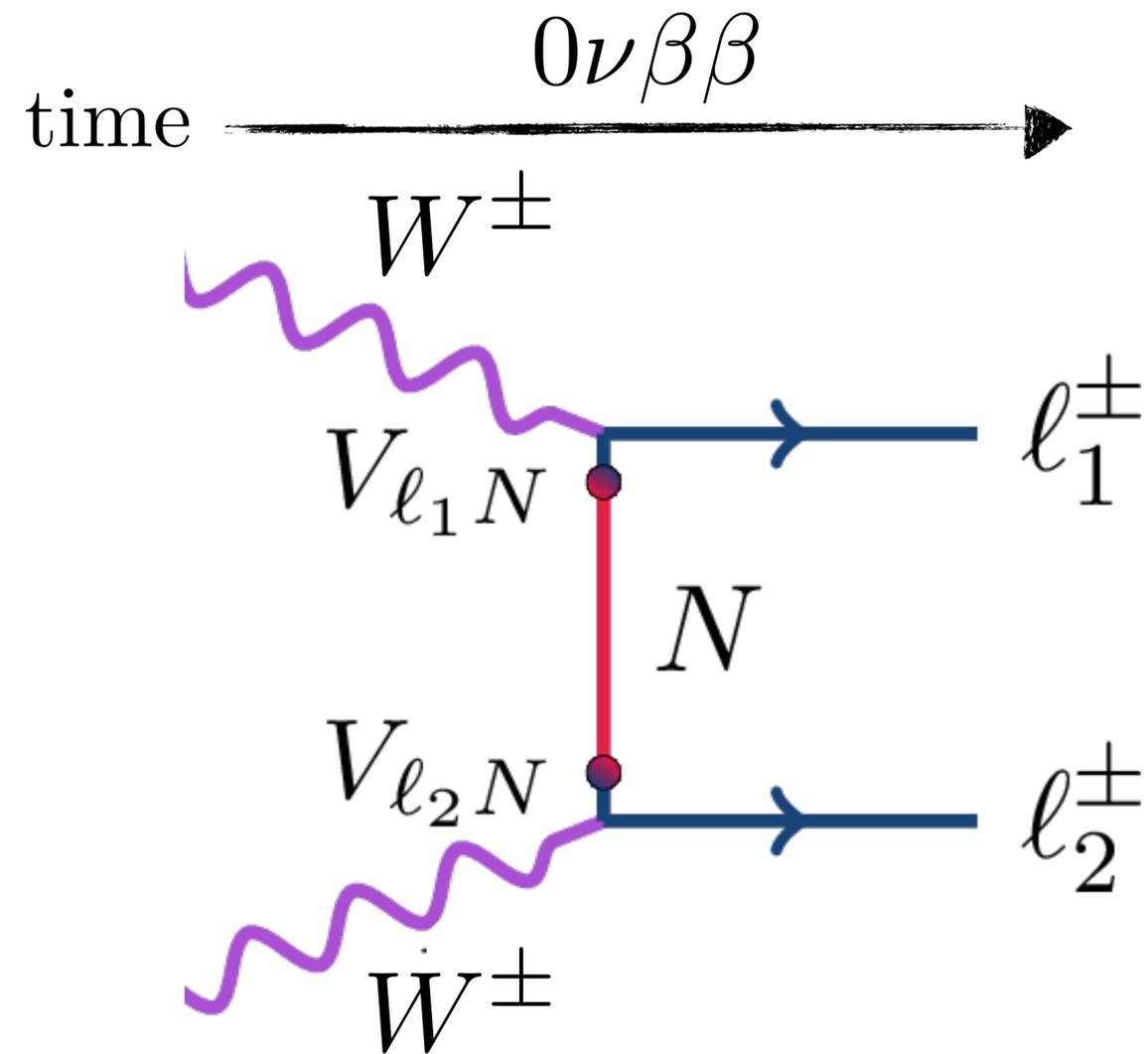
Same-Sign Leptons at the LHC

Same-sign final state leptons probe Majorana HNLs at the LHC

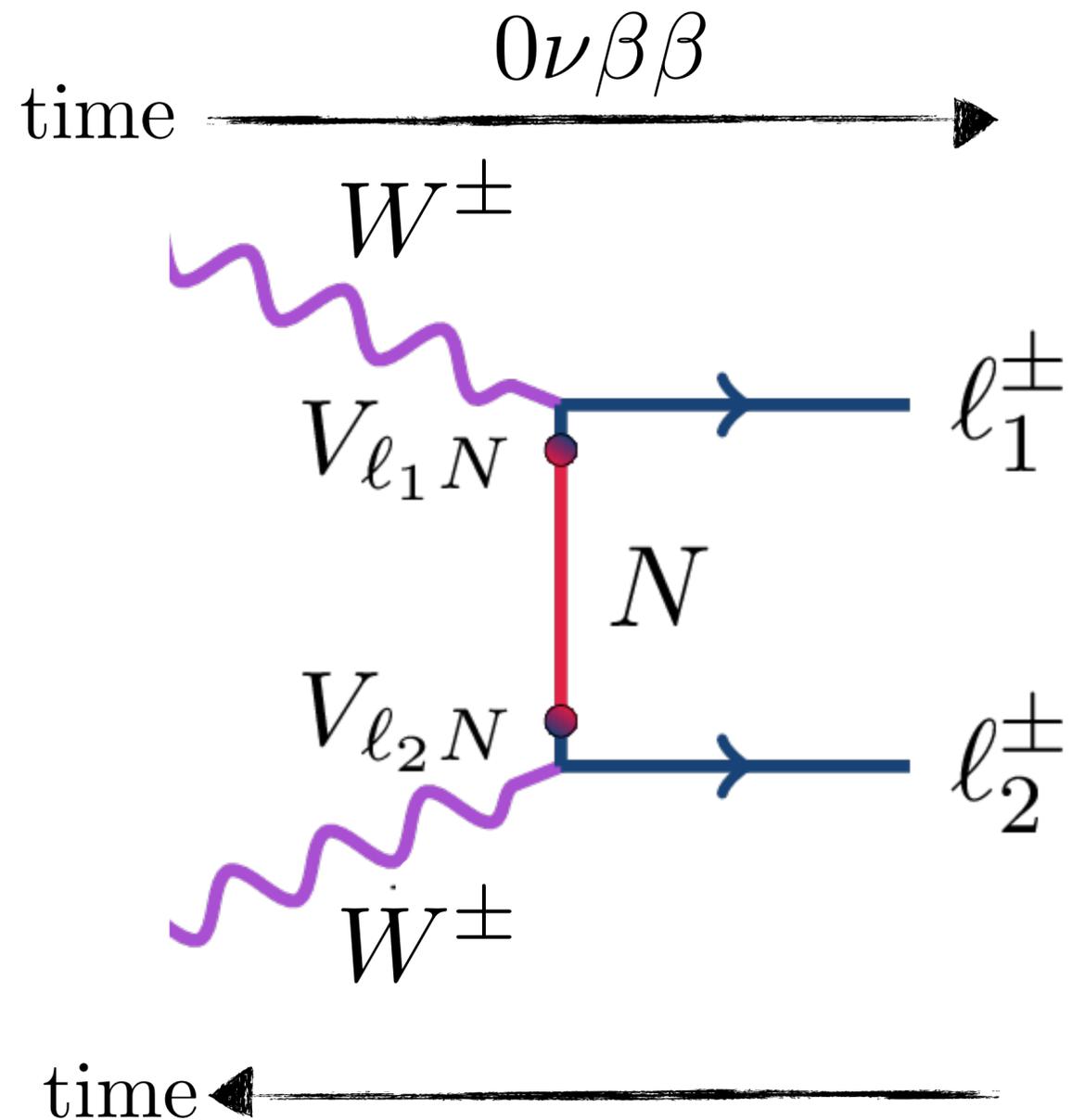


Not as sensitive as $0\nu\beta\beta$

HNLs at Same-Sign Lepton Colliders

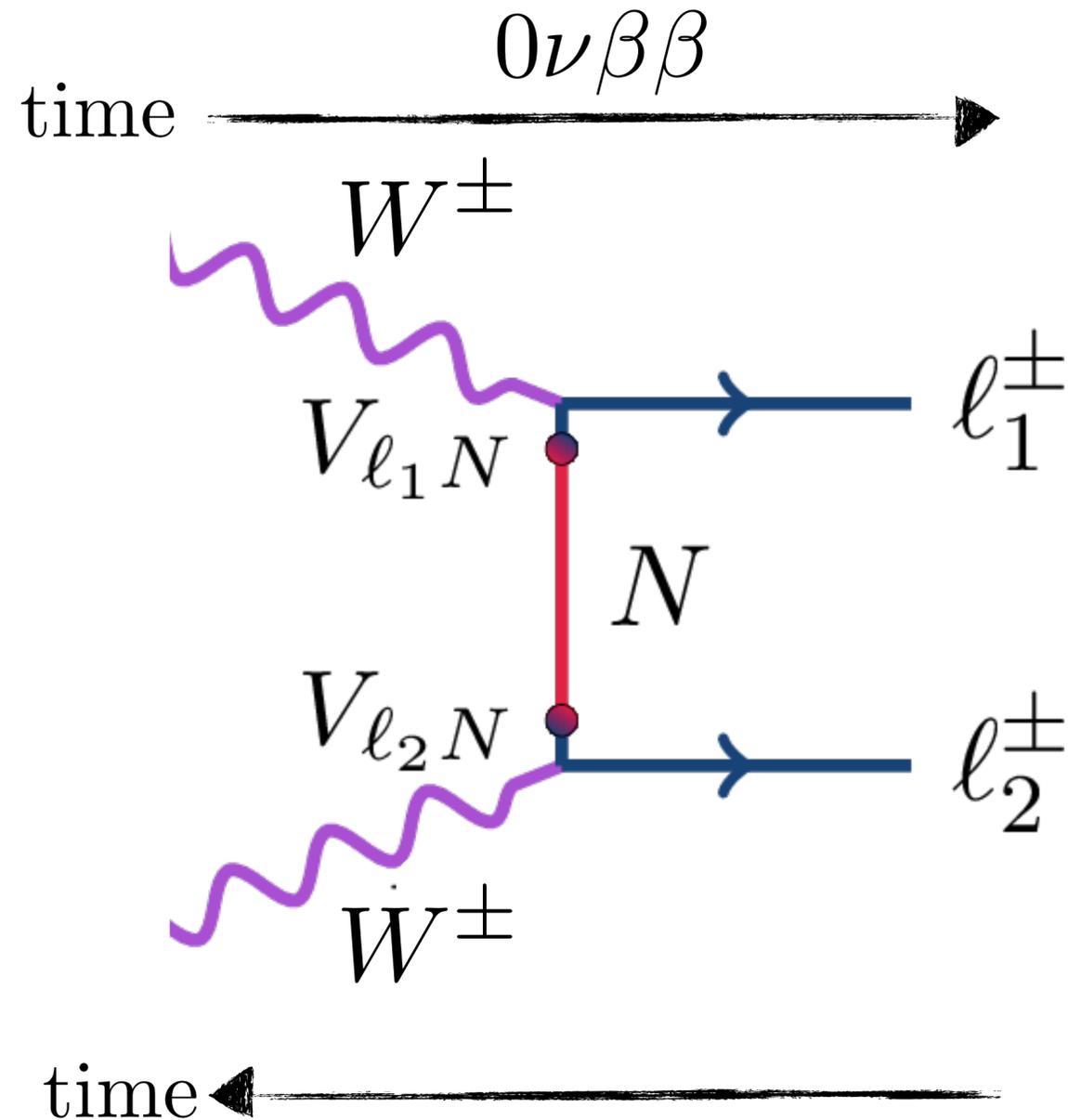


HNLs at Same-Sign Lepton Colliders



Same-Sign Lepton Colliders

HNLs at Same-Sign Lepton Colliders



Same-Sign Lepton Colliders

μ TRISTAN

arXiv > hep-ph > arXiv:2201.06664

arXiv:2201.06664

High Energy Physics – Phenomenology

[Submitted on 17 Jan 2022 (v1), last revised 21 Apr 2022 (this version, v2)]

μ TRISTAN

Yu Hamada, Ryuichiro Kitano, Ryutaro Matsudo, Hiromasa Takaura, Mitsuhiro Yoshida

Compact Linear Collider

SOME THOUGHTS ON e^-e^- COLLISIONS IN CLIC

D. SCHULTE

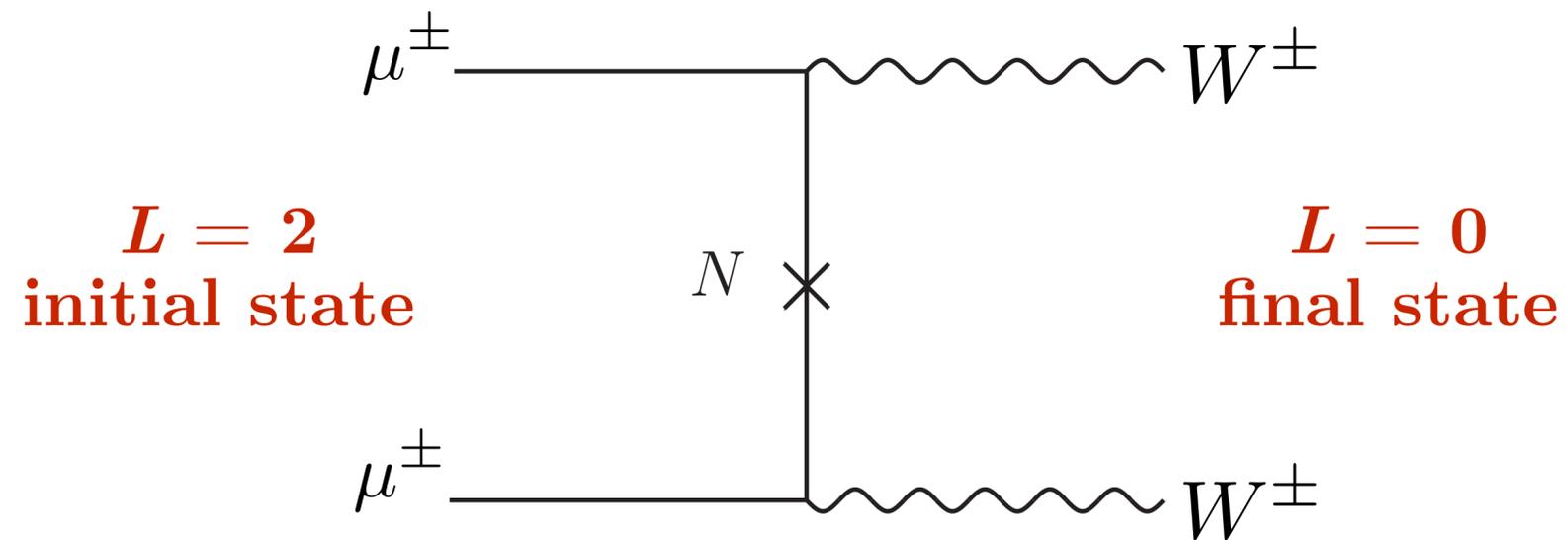
CERN, Geneva, Switzerland

Linear e^+e^- colliders have the potential to also be used to realise e^-e^- collisions. Some preliminary thoughts about the realisation of this at CLIC are presented. Luminosity and some background estimates are also given. [CLIC Note 512 \(2002\)](#)

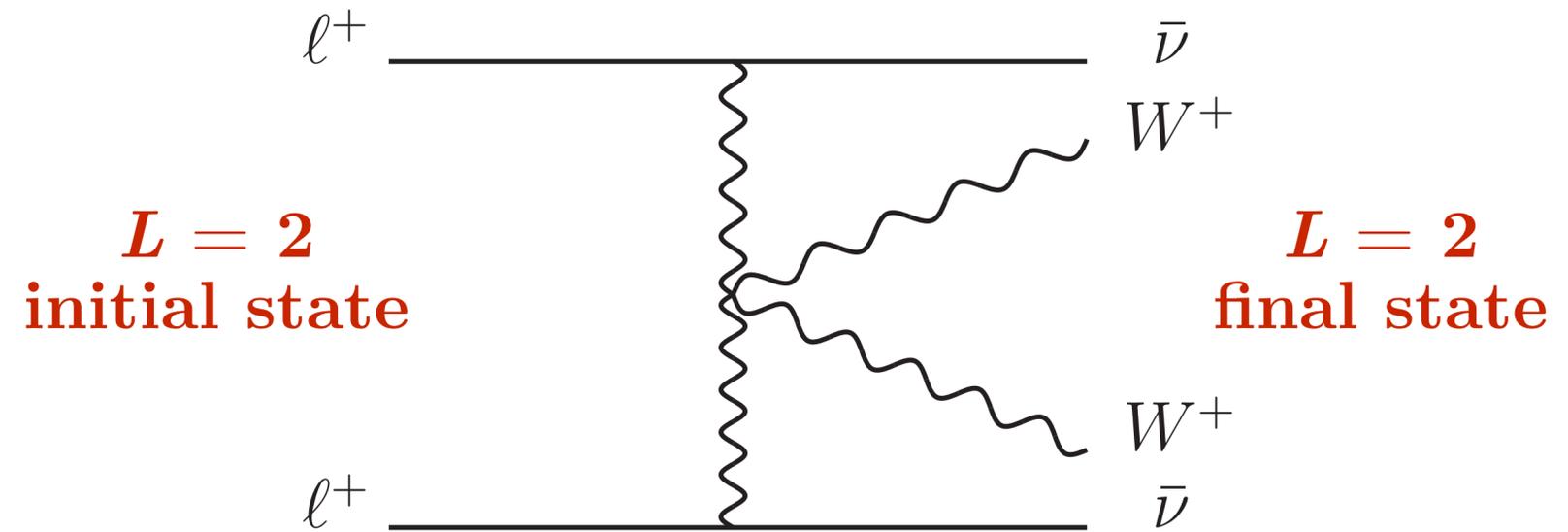
[CLIC Note 512 \(2002\)](#)

Advantages of Same-Sign Lepton Collider

HNL Signal



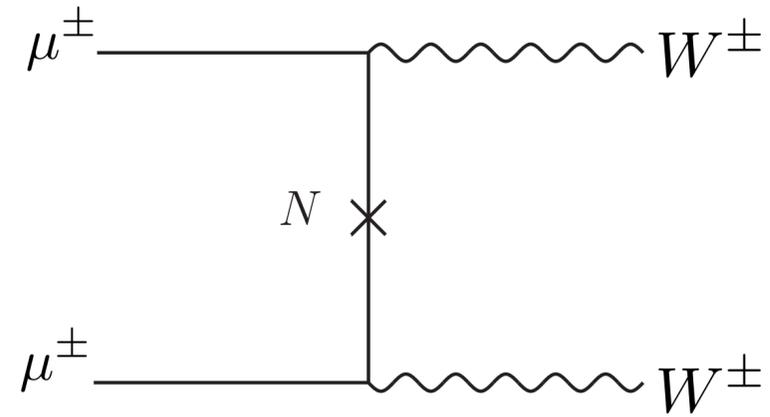
SM Background



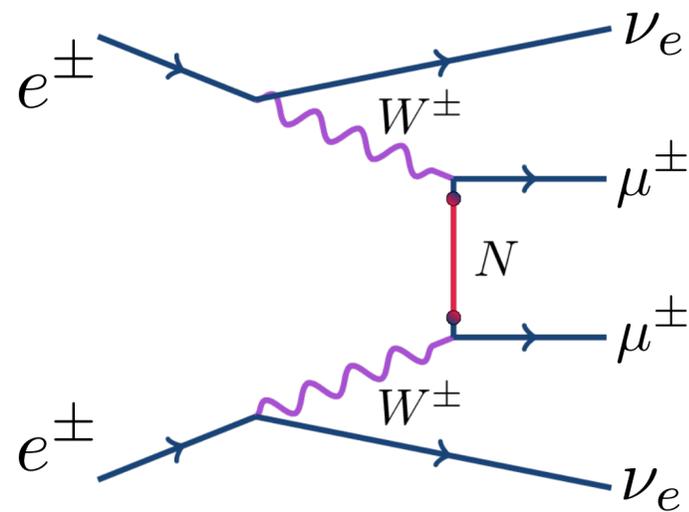
$\Delta L = 2$ transition is not allowed in the SM. Background processes are required to have additional neutrinos in the final state and are highly suppressed. This essentially a zero background search!

HNLs at Same-Sign Lepton Colliders

*W^\pm Pair
Production*

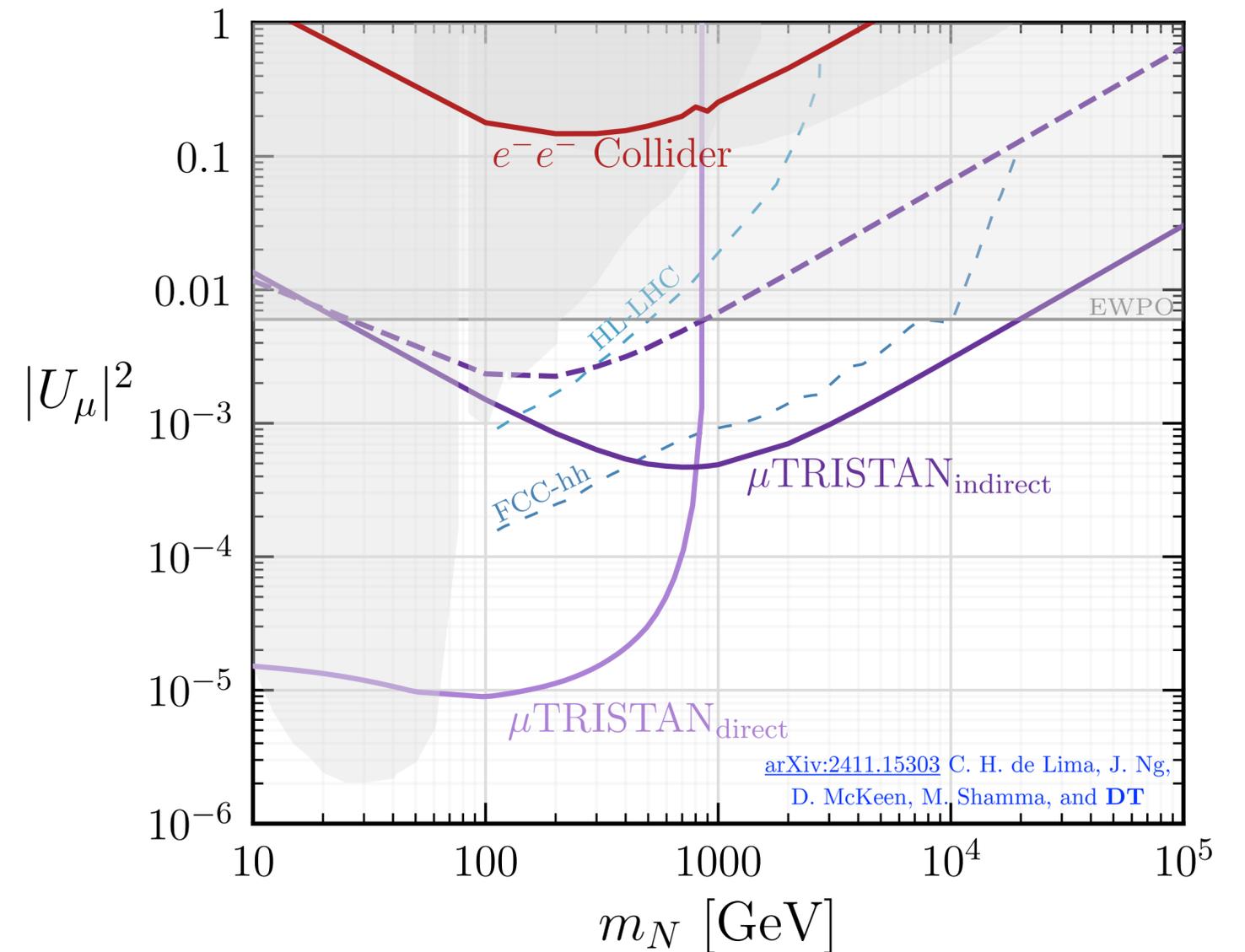


*W^\pm Boson
Scattering*



Probe HNLs with a different flavor initial state lepton!

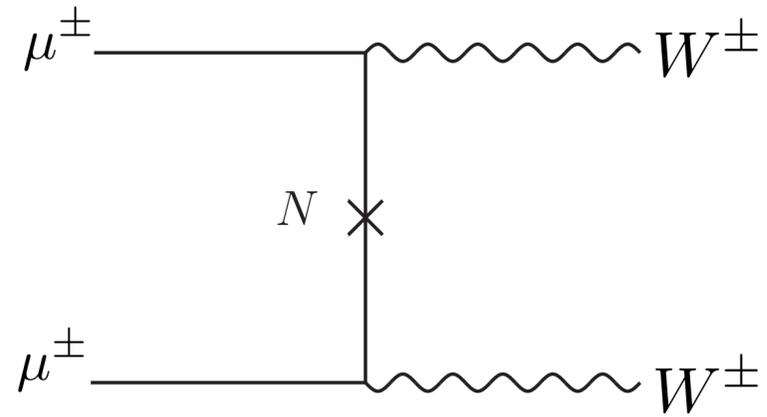
Muon-Mixed HNL



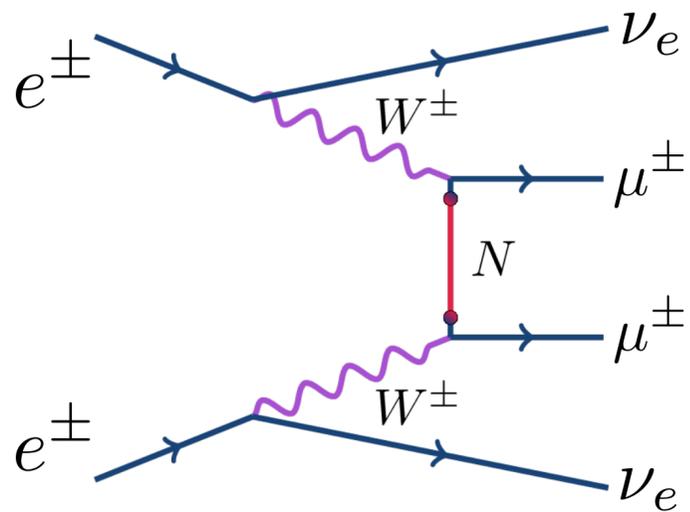
Dashed: $\sqrt{s} = 250$ GeV, Solid: $\sqrt{s} = 1$ TeV

HNLs at Same-Sign Lepton Colliders

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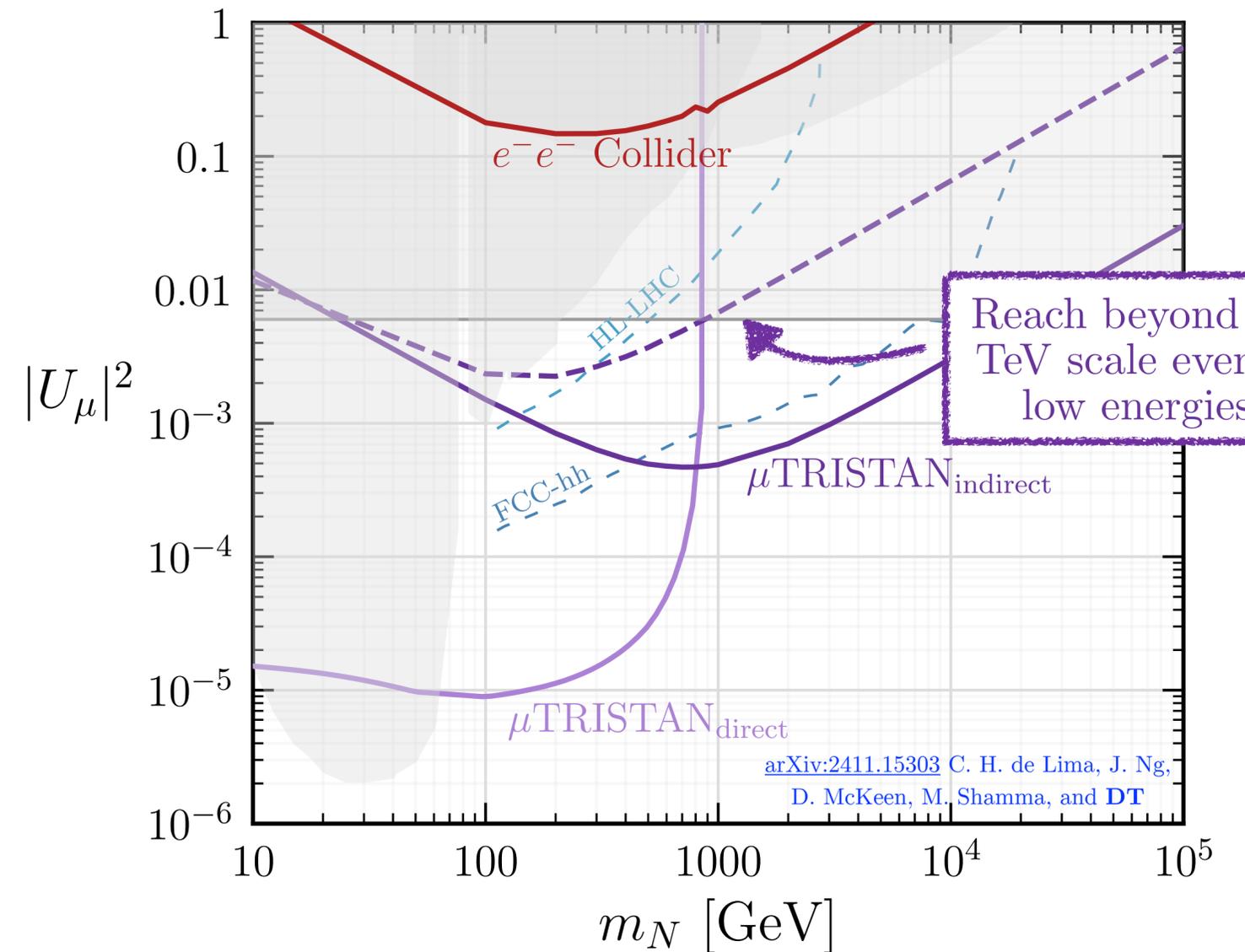


*W^\pm Boson
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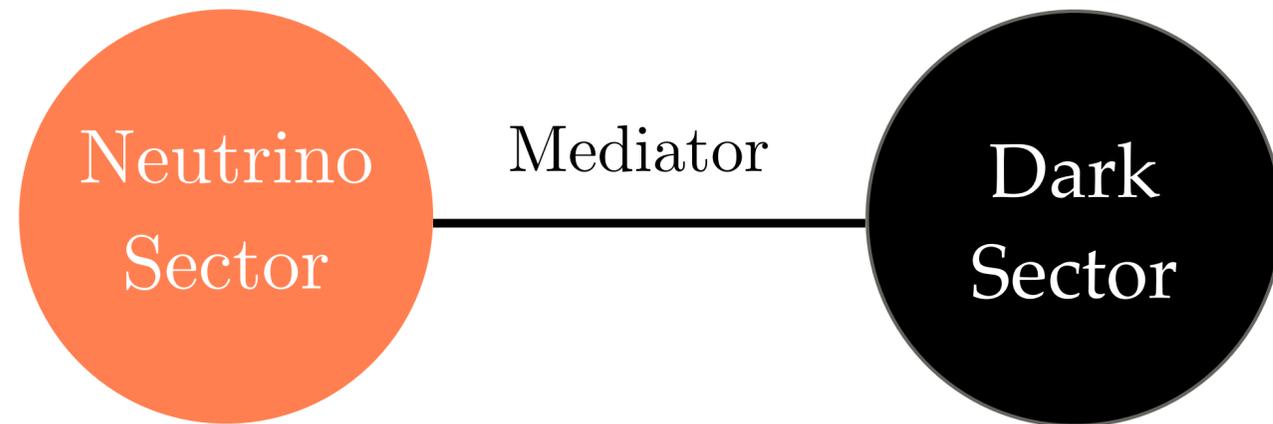
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Muon-Mixed HNL



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Lepton Number Violation and Dark Matter

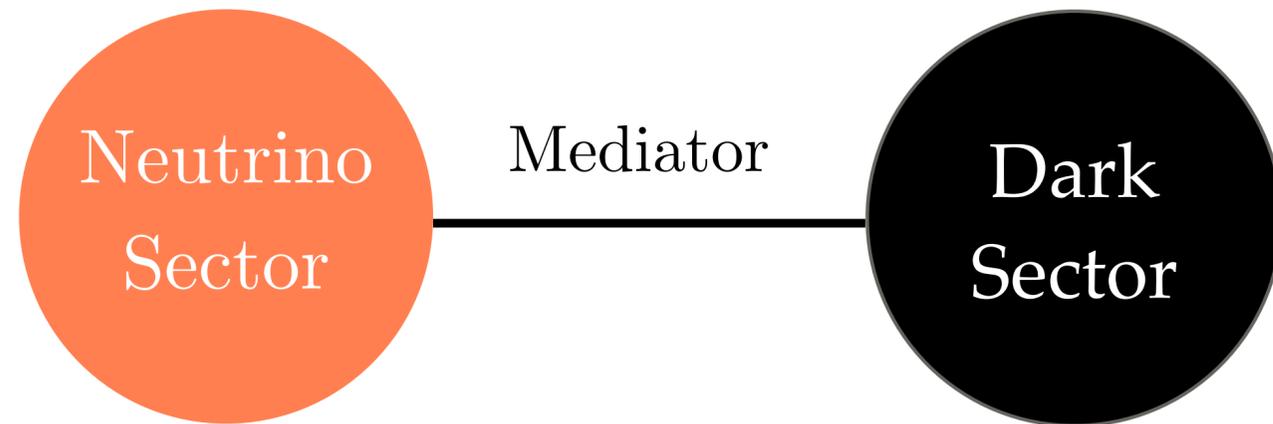


Introduce a new scalar that
interacts only with the SM neutrinos

$$\mathcal{L} \supset \frac{1}{2} \lambda_{\alpha\beta} \nu_{\alpha} \nu_{\beta} \phi$$

Neutrino state has $L = 2$
 $\Rightarrow \phi$ must have $L = -2$

Lepton Number Violation and Dark Matter



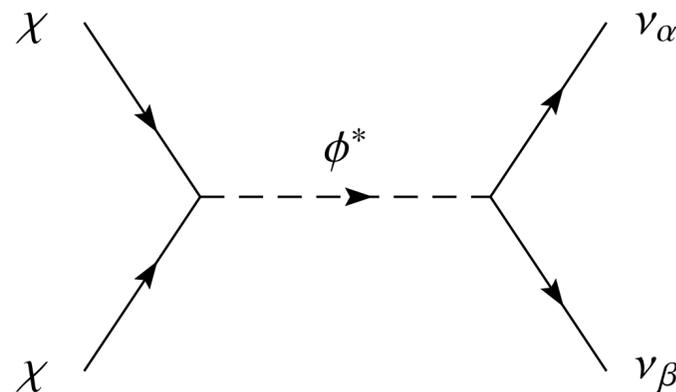
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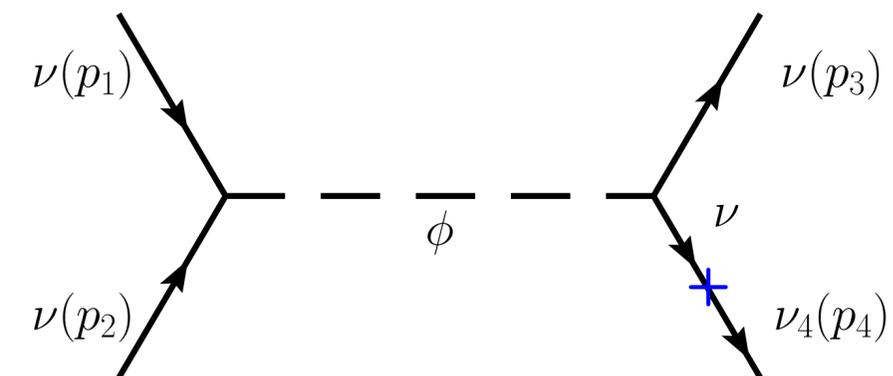
Thermal Dirac Fermion DM

$$\mathcal{L} = \frac{1}{2} y_{\chi} \bar{\chi} \chi \phi + \text{h.c.}$$

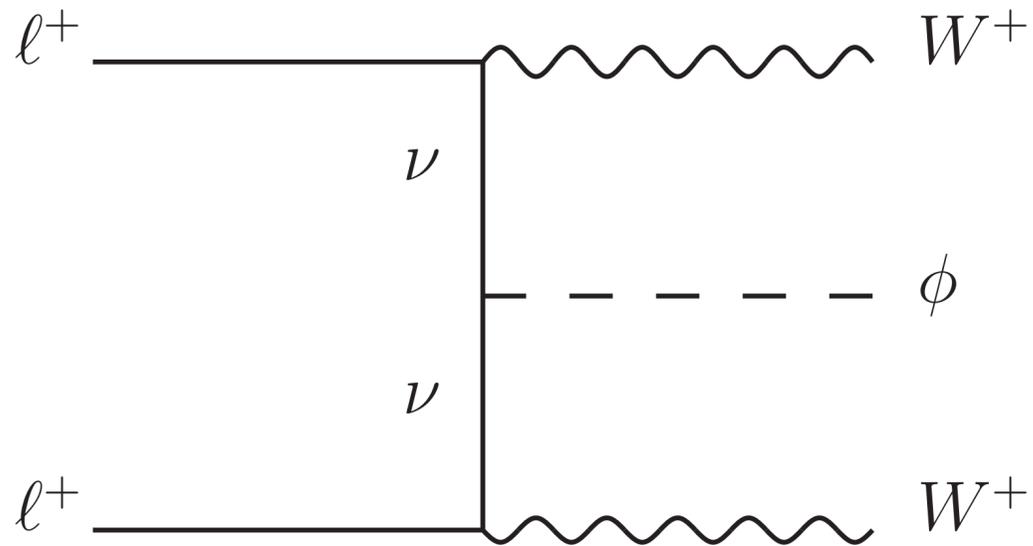


Sterile Neutrino DM

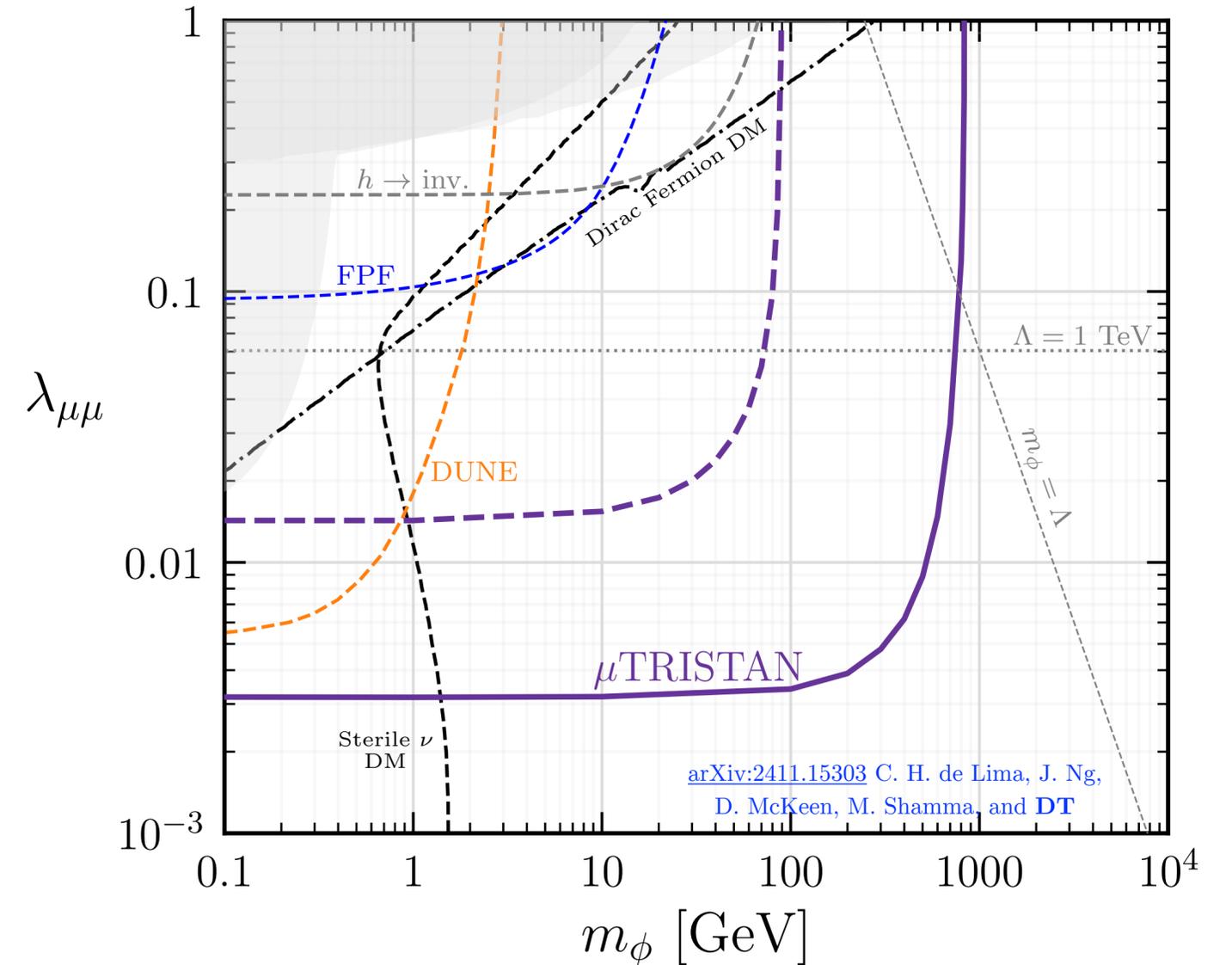
$$\nu_4 = \nu_s \cos \theta + \nu_s \sin \theta$$



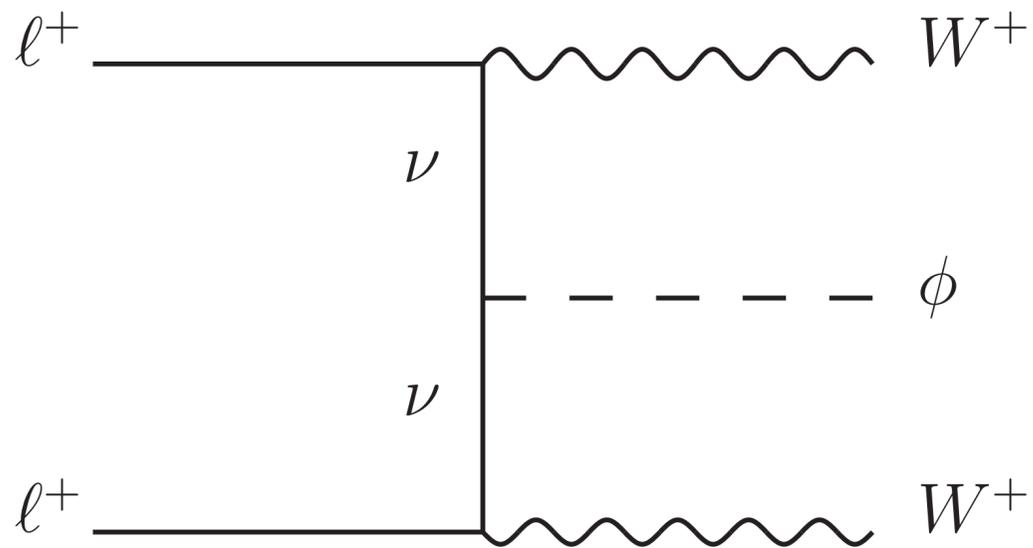
Neutrinophilic Scalars



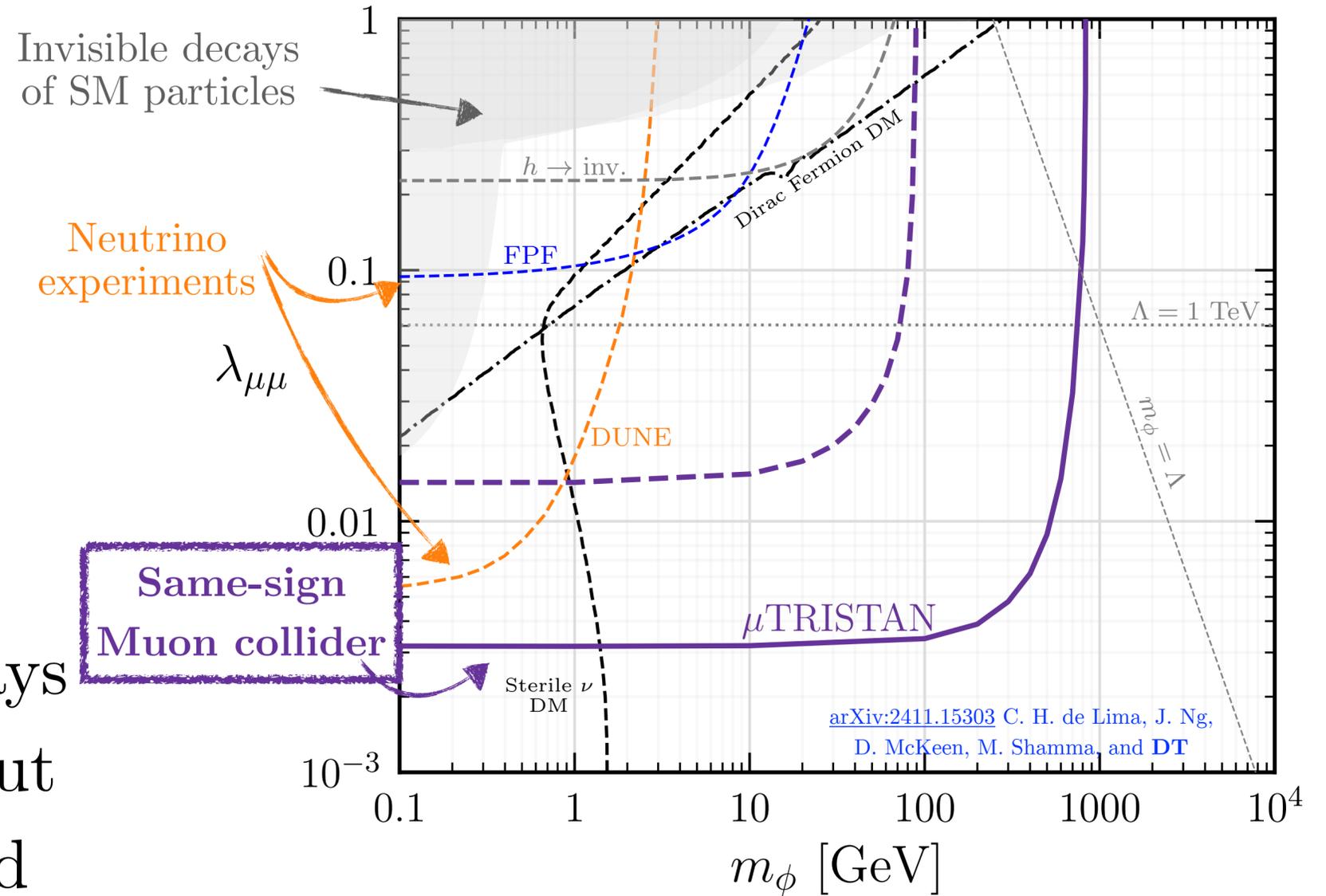
Since the scalar has $L = -2$ but decays invisibly, this is an *apparent LNV* but lepton number is actually conserved



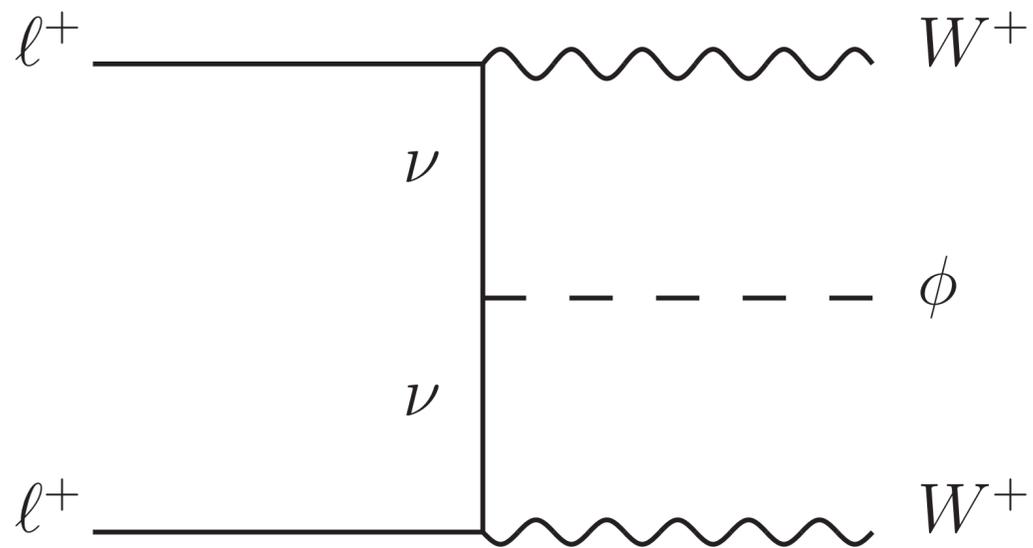
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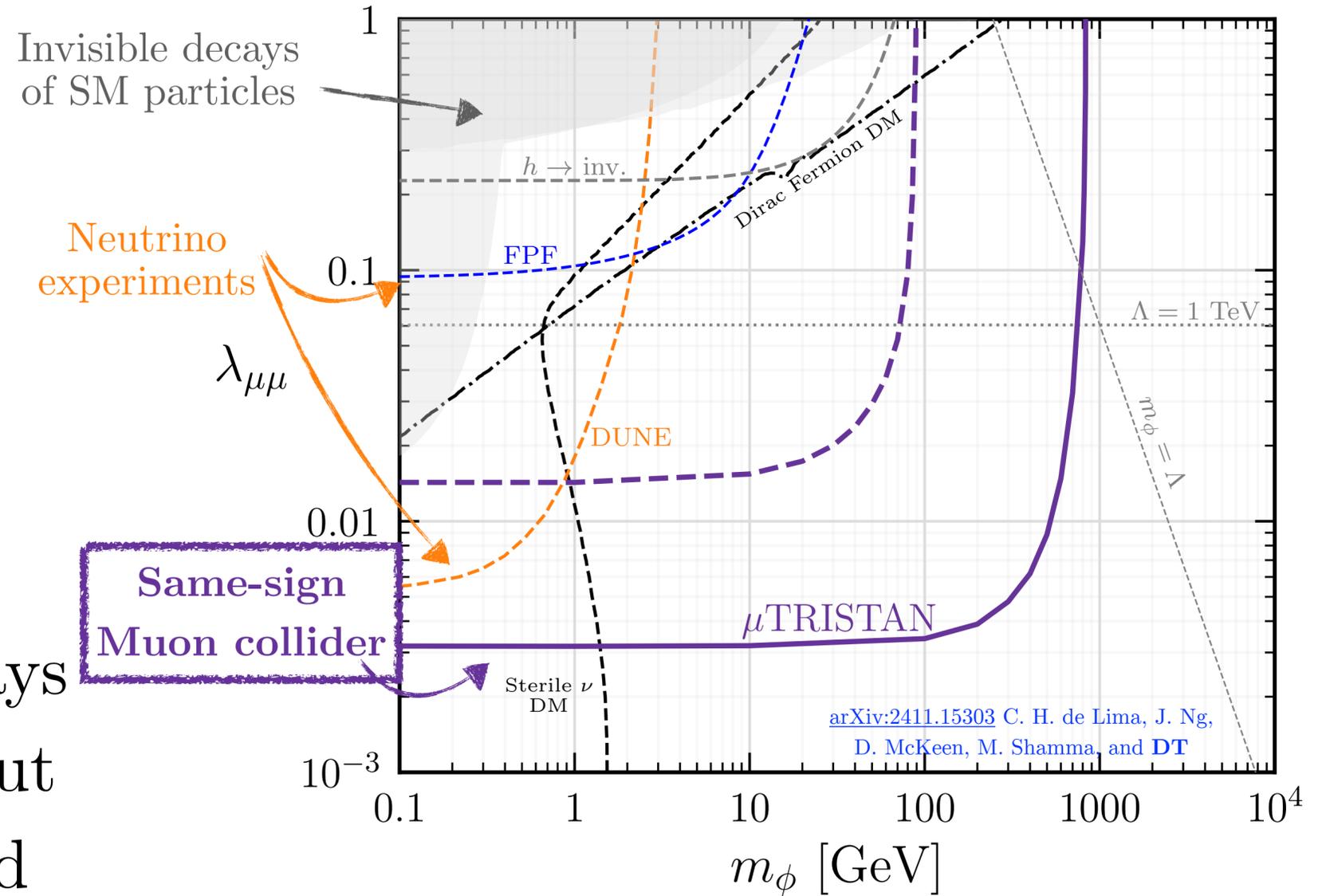
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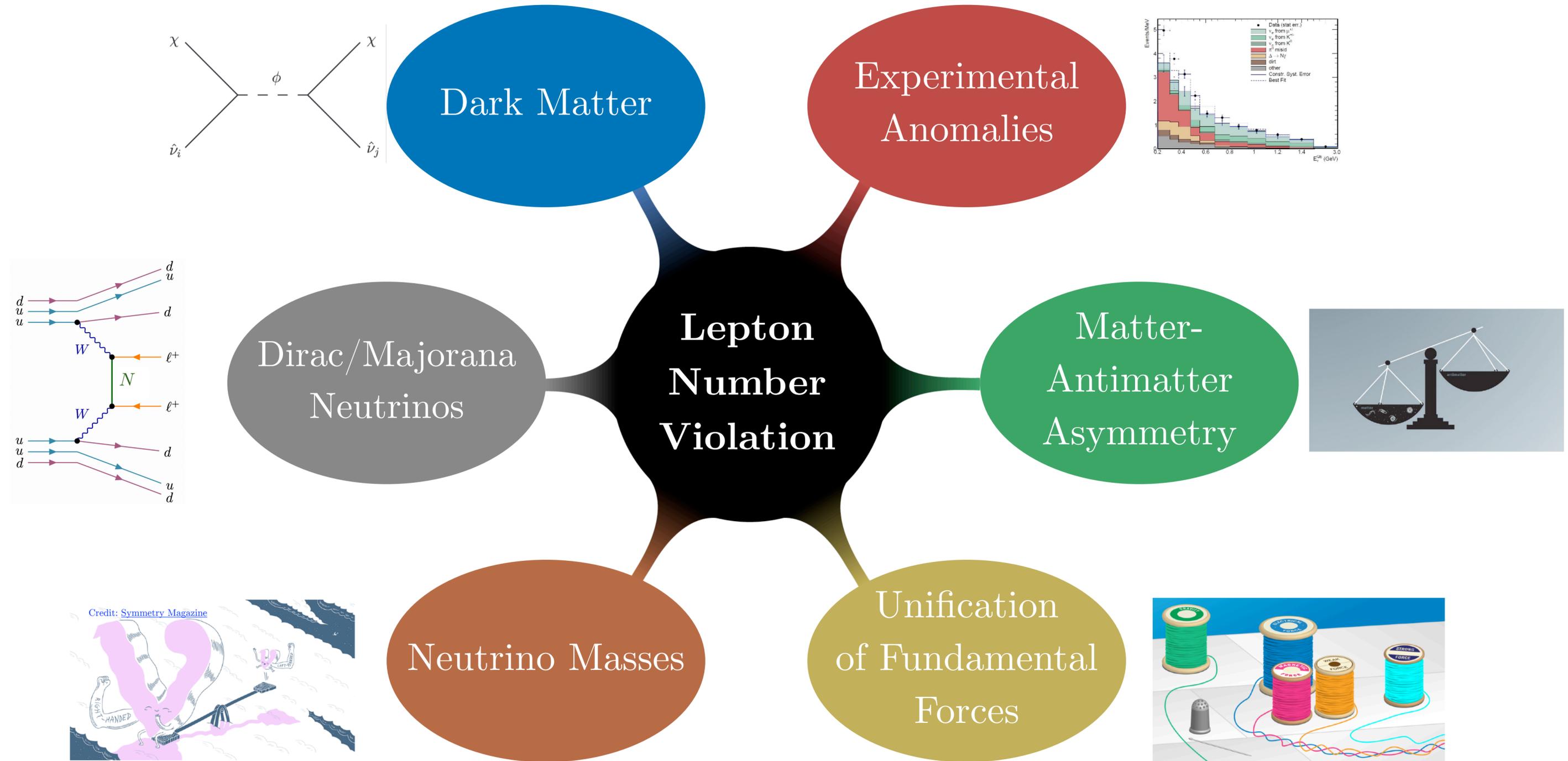


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Same-sign muon collider can cover probe the full dark matter parameter space!

Summary



Thanks!
Questions?