

definitive proof of the solution. SNO.

sol: N_{theory} is correct.

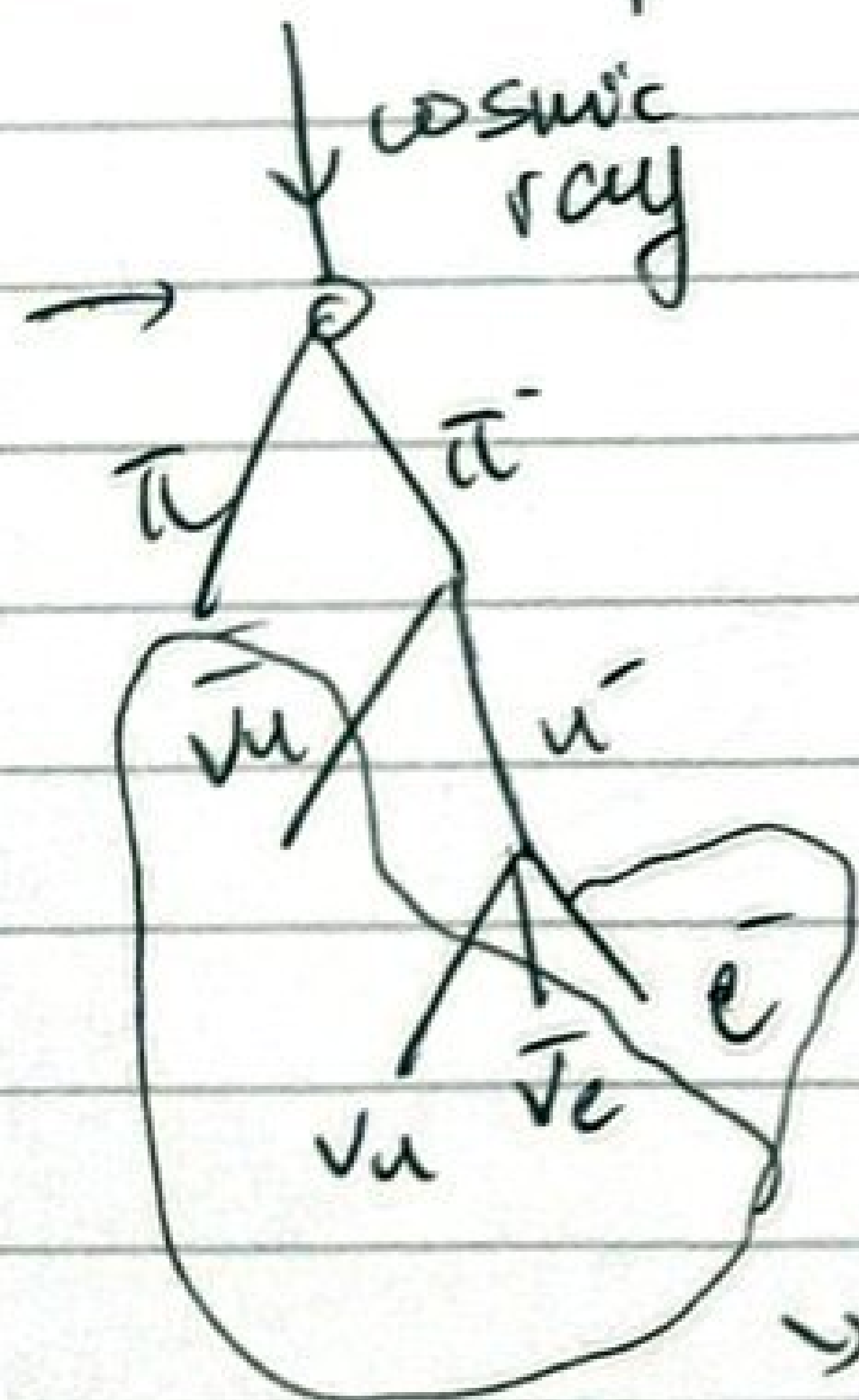
$N_{\text{exp}} \approx \frac{1}{3} N_{\text{th}}$ because $\sim \frac{1}{3}$ remained as ν_e , and $\frac{2}{3}$ of ν_e oscillated to $\nu_\mu + \nu_\tau$.

Why was SNO convincing?

3 channels: $\left\{ \begin{array}{ll} \nu + d \rightarrow \nu + p + n & (\text{NC, measure } \phi_{\nu_{\text{total}}}) \\ \nu + d \rightarrow e + p + p & (\text{CC, measure } \phi_e) \\ \nu + e \rightarrow \nu + e & (\text{mix, " some combination of } \phi_e, \phi_\mu + \phi_\tau) \end{array} \right.$

$\Rightarrow$ the ~~was~~ well recognised SNO result plot.

atmospheric ν oscillation.

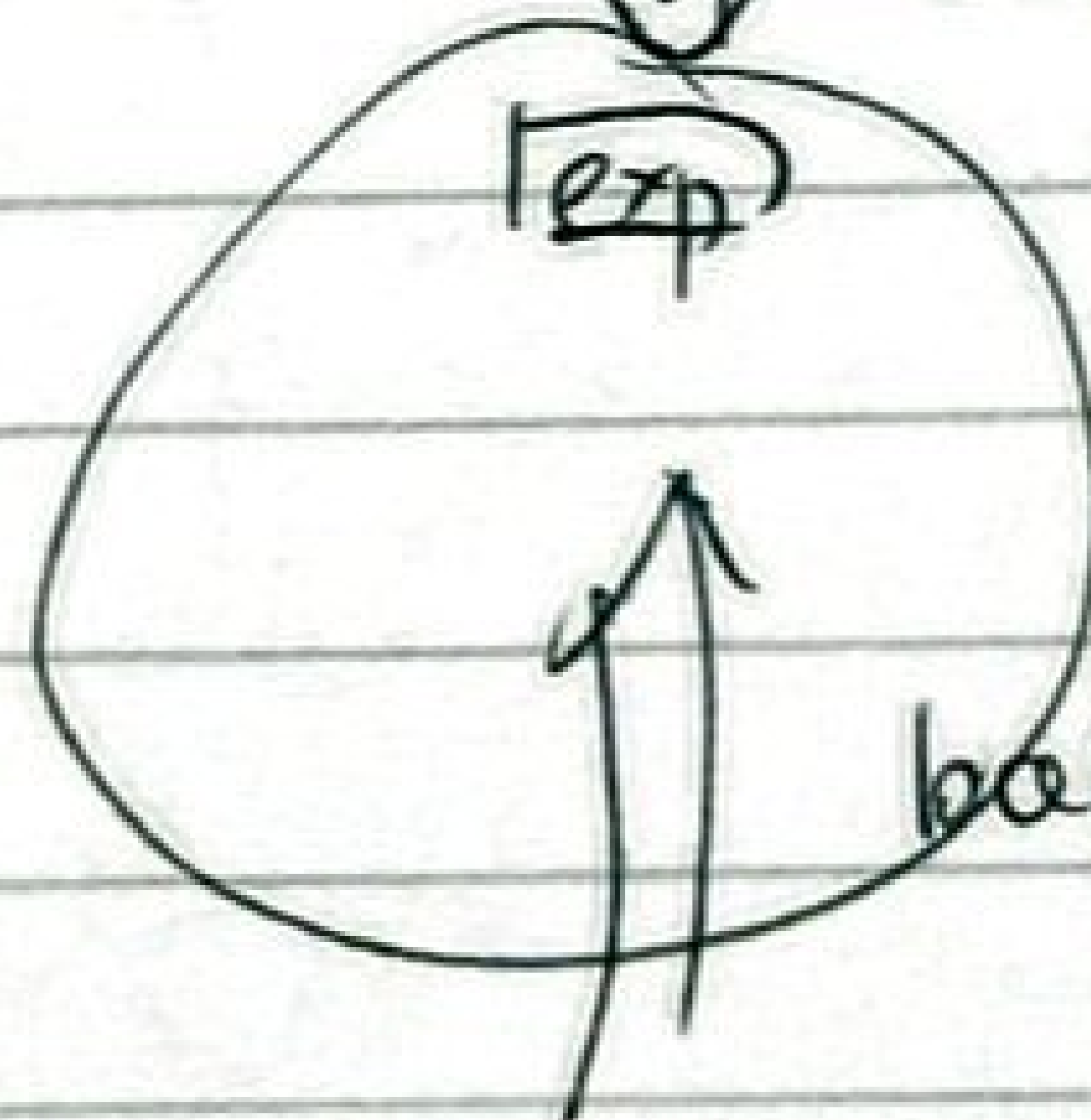


~~without~~ osci

At production. $\phi_\mu : \phi_e \approx 2:1$.

energy GeV +

baseline $\sim 10 \text{ km}$

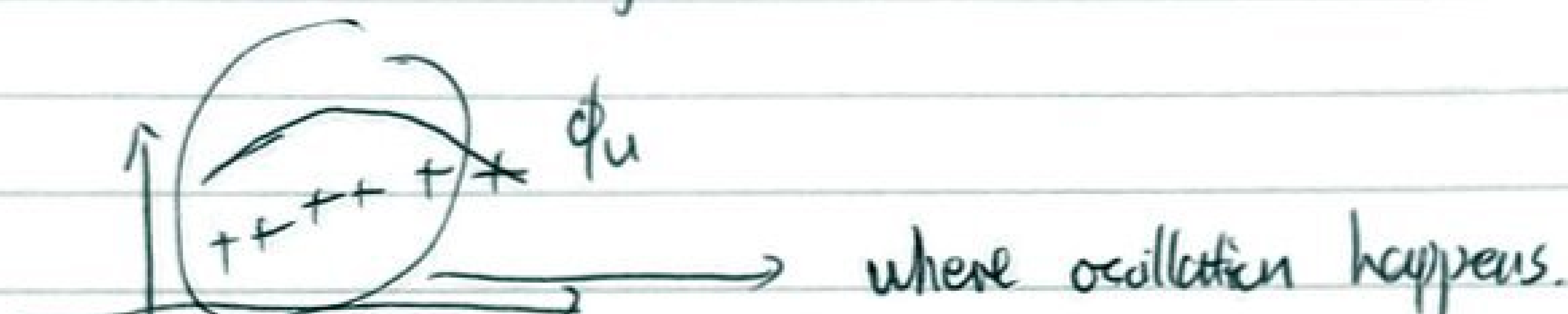
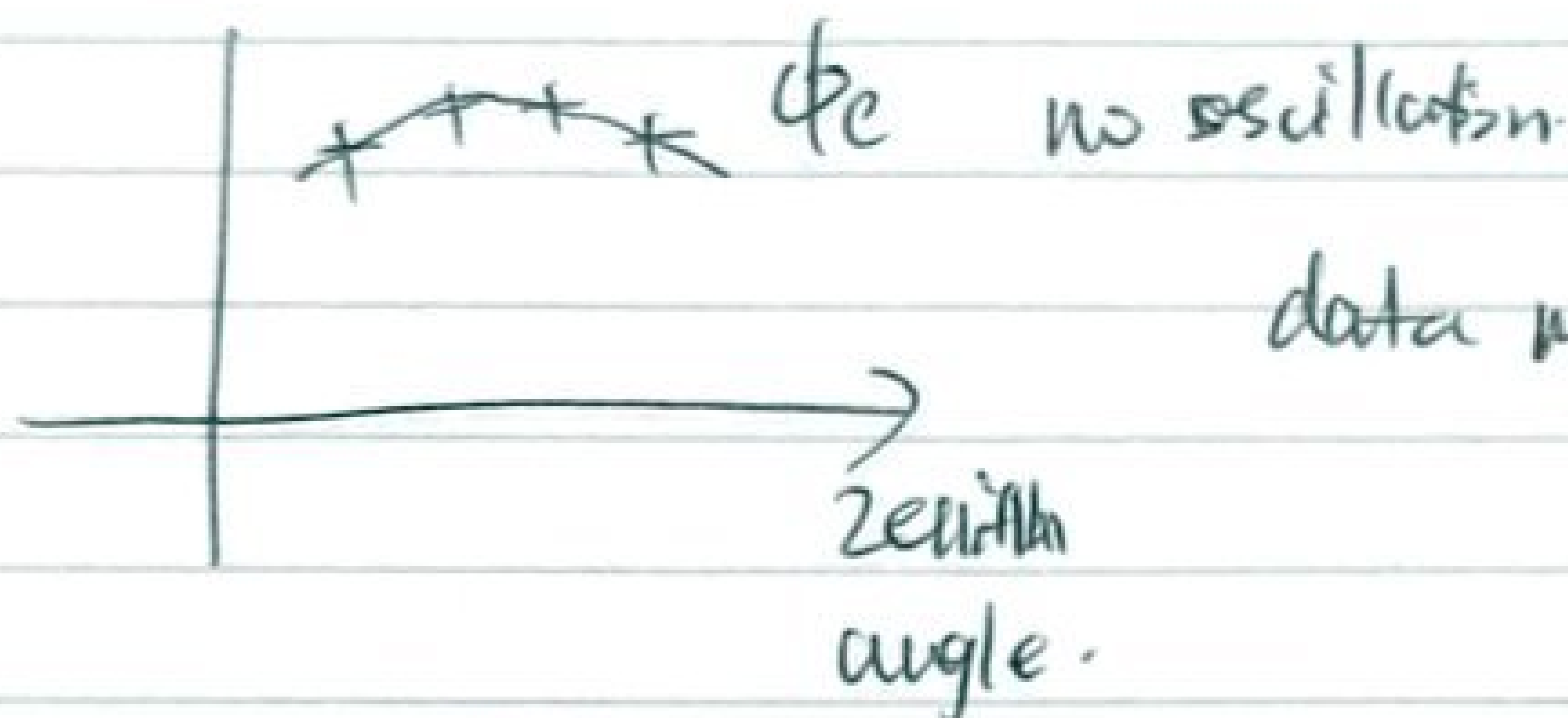


Earth

baseline $\sim 10^4 \text{ km}$

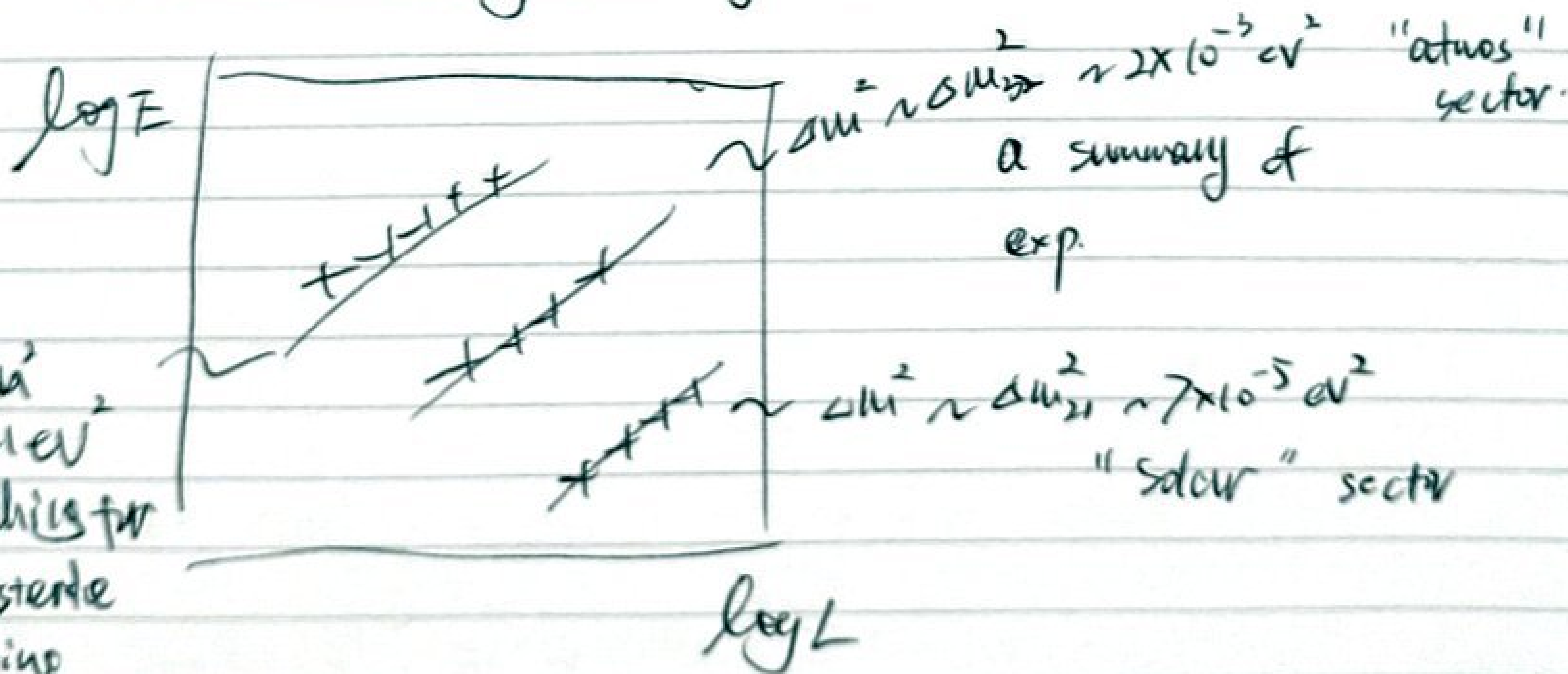
baseline vary from $10\text{ km} - 10^4\text{ km}$.

signature:



$\nu_\mu \rightarrow \nu_e$. so one measures the reduced flux of ν_μ

what are actually oscillating?



most experiments are pretty well approximated by two-level system oscillation, e.g. i.e. $P = \sin^2 \theta \sin^2 \left(\frac{\Delta m^2 c^4}{4E} \right)$

reason: $\Delta m_{32}^2 \approx \Delta m_{31}^2 \gg \Delta m_{21}^2$

if $\frac{L}{E} \sim \frac{1}{\Delta m_{32}^2} \sim \frac{1}{\Delta m_{31}^2} \rightarrow$ 2 level oscillation.

if $\frac{L}{E} \sim \frac{1}{\Delta m_{21}^2} \rightarrow$ not participating.

$\therefore$ solar: $\theta_{12}, \Delta m_{21}^2$, two level.
 atmos: $\theta_{32}, \Delta m_{32}^2$, two level.
 } at leading order, to build intuition

Last missing piece: matter effect.

In vacuum, we used this:

$$i \frac{d}{dx} | \nu_i \rangle = \frac{m_i^2}{2E} | \nu_i \rangle$$

we can solve this in flavor basis instead.

$$i \frac{d}{dx} | \nu_f \rangle = U_{fi} \frac{m_i^2}{2E} U_{fi}^* | \nu_f \rangle$$

Or, in 2-flavor case:

$$i \frac{d}{dt} \begin{pmatrix} | \nu_e \rangle \\ | \nu_\mu \rangle \end{pmatrix} = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos 2\theta & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} \begin{pmatrix} | \nu_e \rangle \\ | \nu_\mu \rangle \end{pmatrix}$$

↑ solve this. you will get the usual $P = \sin^2 \theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$

Now, in the presence of (high-density) matter:

ν 's dispersion relation will get modified.
similarly to γ traveling in medium.

consider the presence of electron, & only CC interaction.

$$\mathcal{L} = -2\sqrt{2} G_F (\bar{\nu}_e \gamma_\mu \nu_e) (\bar{e}_L \gamma_\mu e_L) \quad \text{integrating out } W$$

$$\langle \bar{e}_L \gamma_\mu e_L \rangle = \frac{N_e}{2} \delta_{\mu 0}$$

$$\Rightarrow \text{EOM for neutrinos: } (i\partial - \sqrt{2} G_F N_e \gamma_0) | \nu_e \rangle = 0$$

$$\frac{\Delta m^2}{4E} \begin{pmatrix} -C_2 & S_2 \\ S_2 & C_2 \end{pmatrix} \text{ from above.} \quad E = p + \underbrace{\sqrt{2} G_F N_e}_{\text{potential } V}$$

$$\Rightarrow H = H_{\text{vacuum}} + \begin{pmatrix} V & 0 \\ 0 & 0 \end{pmatrix}$$

✓ with this H , you can solve the Schrödinger Equation,

assuming constant matter density.

$$\text{define } \tilde{3} = \frac{2\sqrt{2} G_F N_e}{\Delta m^2/E}$$

$$P_{\alpha\beta} = \frac{\sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right) (1 + \tilde{3}^2 - 2\tilde{3} \cos 2\theta)}{1 + \tilde{3}^2 - 2\tilde{3} \cos 2\theta}$$

↑ sensitive to the sign of Δm^2 .

we have only clearly measured matter effect from the Sun, where matter density is NOT constant.

how solar ν actually oscillate is the so-called MSW effect.

But, matter effect $\rightarrow$ the sign of Δm^2 holds.
so, we know $\Delta m_{21}^2 > 0$.

with the two level approximation, & matter effect for solar ν -
you can get good intuition for all oscillation exp
measure angles θ_{12} , θ_{13} & θ_{23} . and Δm_{21}^2 , Δm_{32}^2 .

Next gen exp DUNE Hyper-K. aim to measure $\delta\phi$ only
both use $P_{\nu e}$ & $P_{\nu \mu}$ need 3-flavor & matter effect.

in constant matter density (here it's a good approximation for Earth surface).

$$\Delta_{31} \equiv \Delta m_{31}^2 L / 4E$$

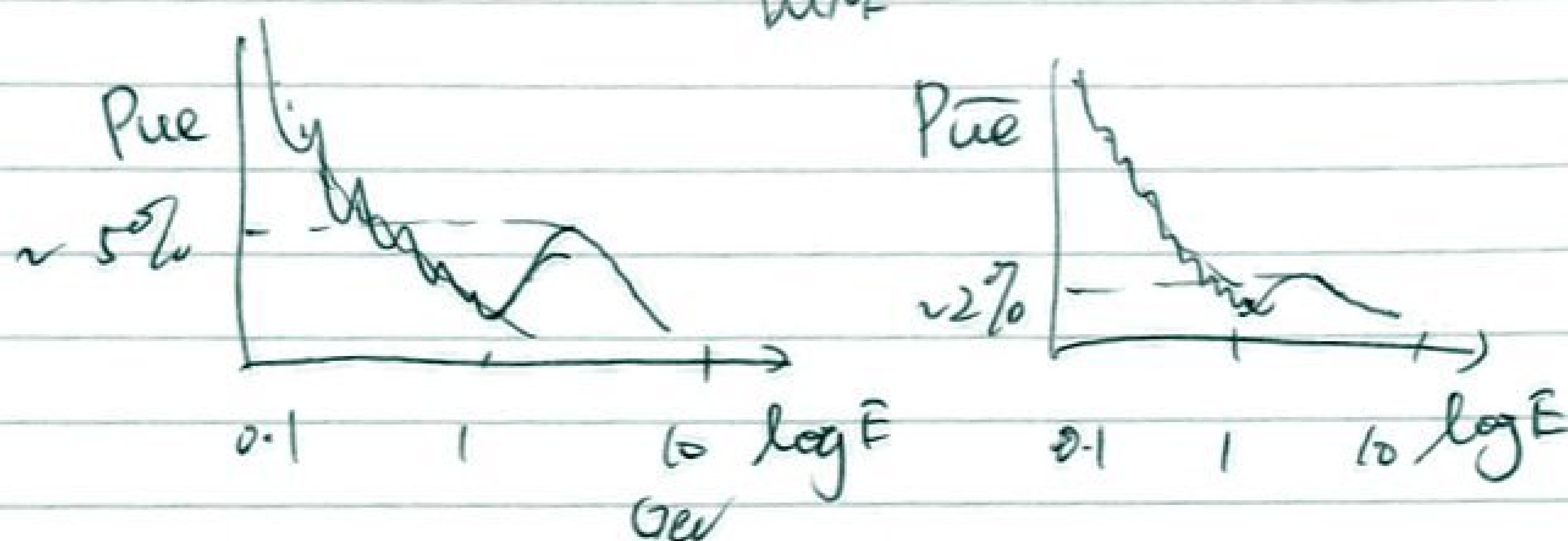
$$(DUNE) \quad P_{\nu e} \approx \sin^2 \theta_{23} \sin^2 \theta_{13} \Delta_{31}^2 \frac{\sin^2(\Delta_{31} - VL/2)}{(\Delta_{31} - VL/2)^2}$$

$$+ \sin^2 \theta_{23} \sin^2 \theta_{13} \sin^2 \theta_{12} \frac{\sin(\Delta_{31} - VL/2)}{\Delta_{31} - VL/2} \cdot \Delta_{31} \Delta_{21} \frac{\sin(VL/2)}{VL/2} \times \cos(\Delta_{31} + \delta_{CP})$$

note: "long baseline" ν_e need to be not tiny.

ν_L term together with $\Delta_{31} \Rightarrow$ ~~looks~~ sign of Δ_{31}
 δ_{CP} shows up in 2nd term.

$\Rightarrow$ MH. δ_{CP} both measured in P_{ue} & $P_{\bar{u}\bar{e}}$.



$\nu_{\mu}E$, Hyper-K aim to measure $\delta_{CP} \sim 10^\circ$
which causes $\sim 5\%$ change in peak height.