

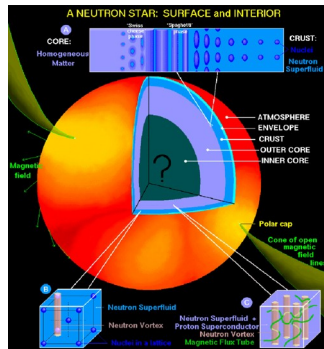
From light nuclei to neutron stars

Alex Gezerlis



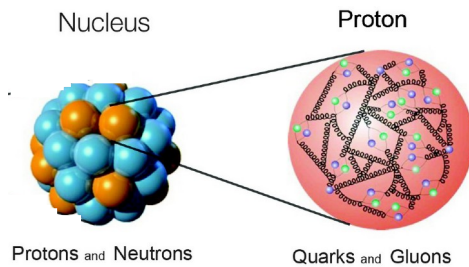
Workshop on progress in ab initio nuclear theory
TRIUMF, Vancouver BC
February 25, 2026

Outline

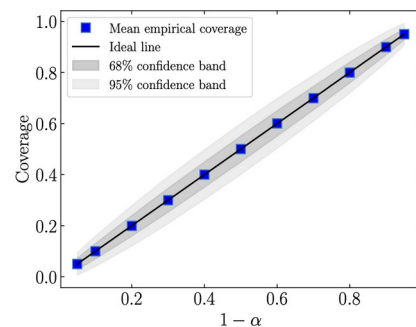


Credit: Dany Page

Motivation



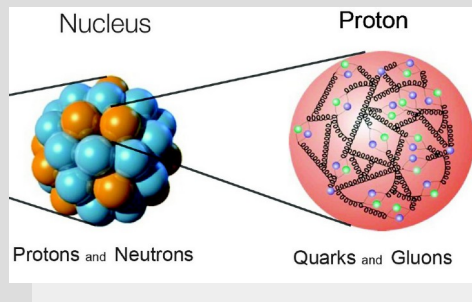
Nuclear methods



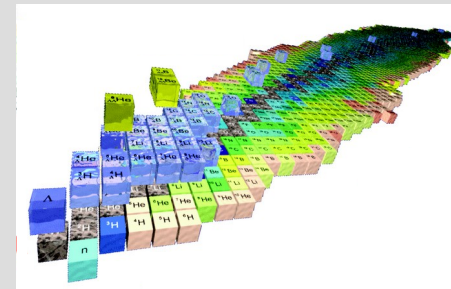
Recent results

Physical systems studied

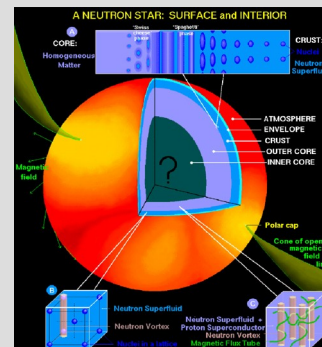
Nuclear forces



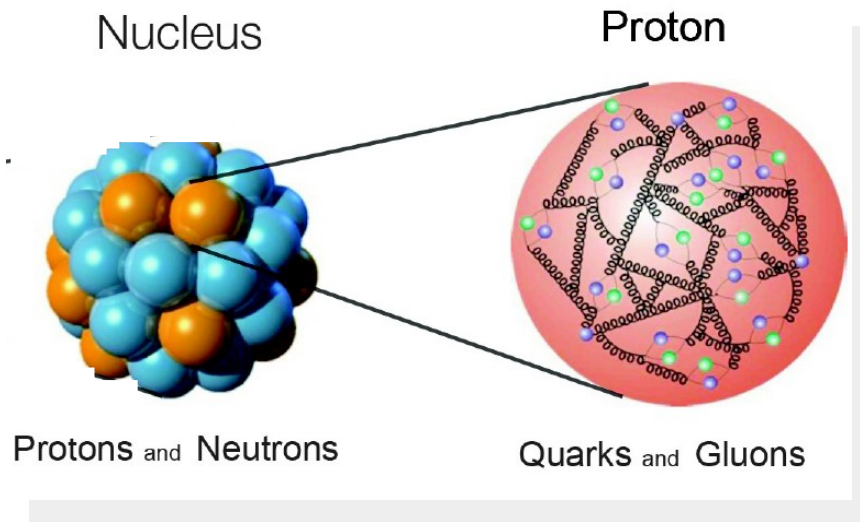
Nuclear structure



Nuclear astrophysics

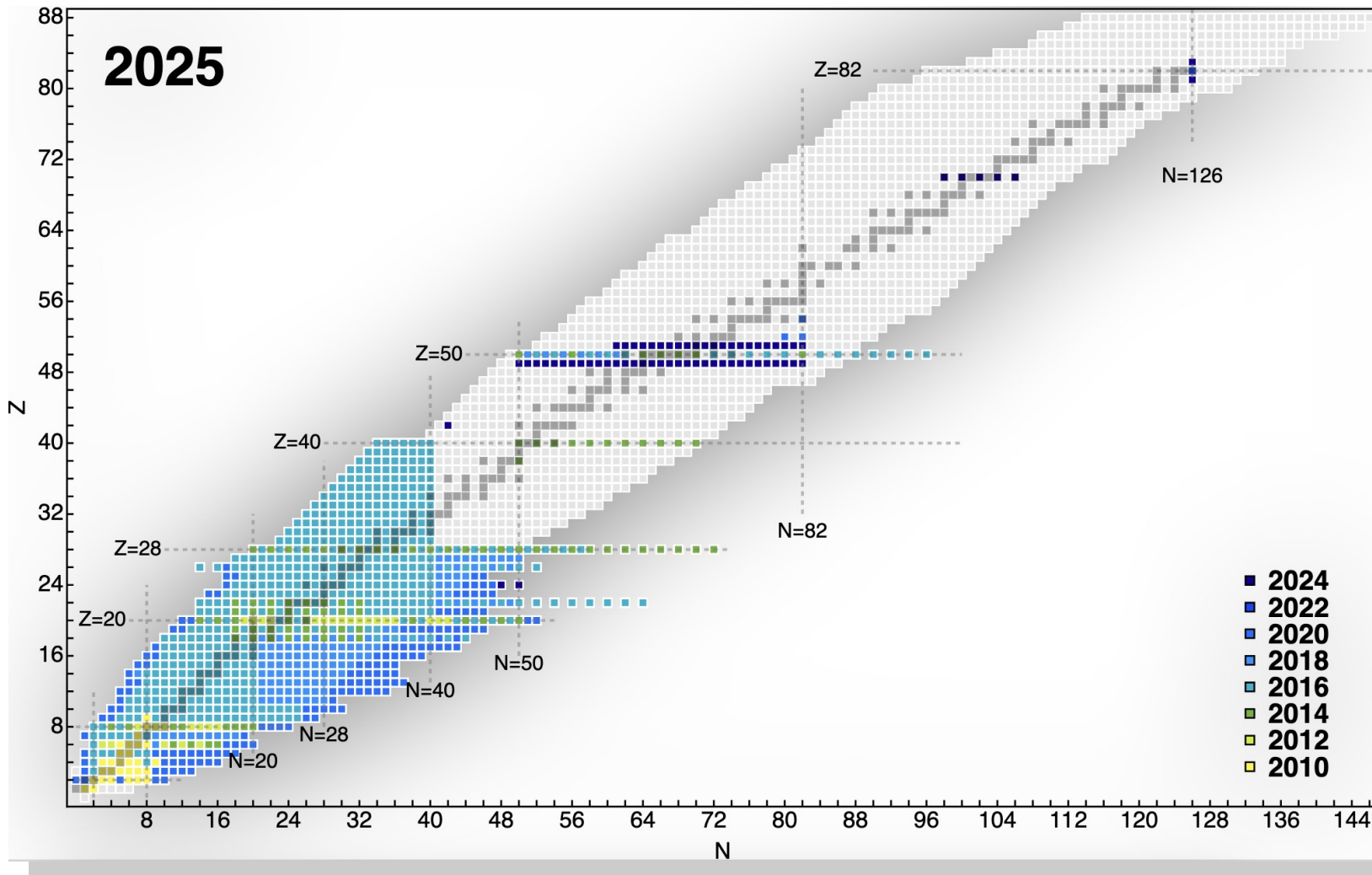


Key system: few nucleons



- No unique nuclear potential
- Preferable to use combination of phenomenological (high-quality) and more modern (conceptually clean) approach
- Desirable to make contact with underlying level
- New era, where practitioners design interactions themselves

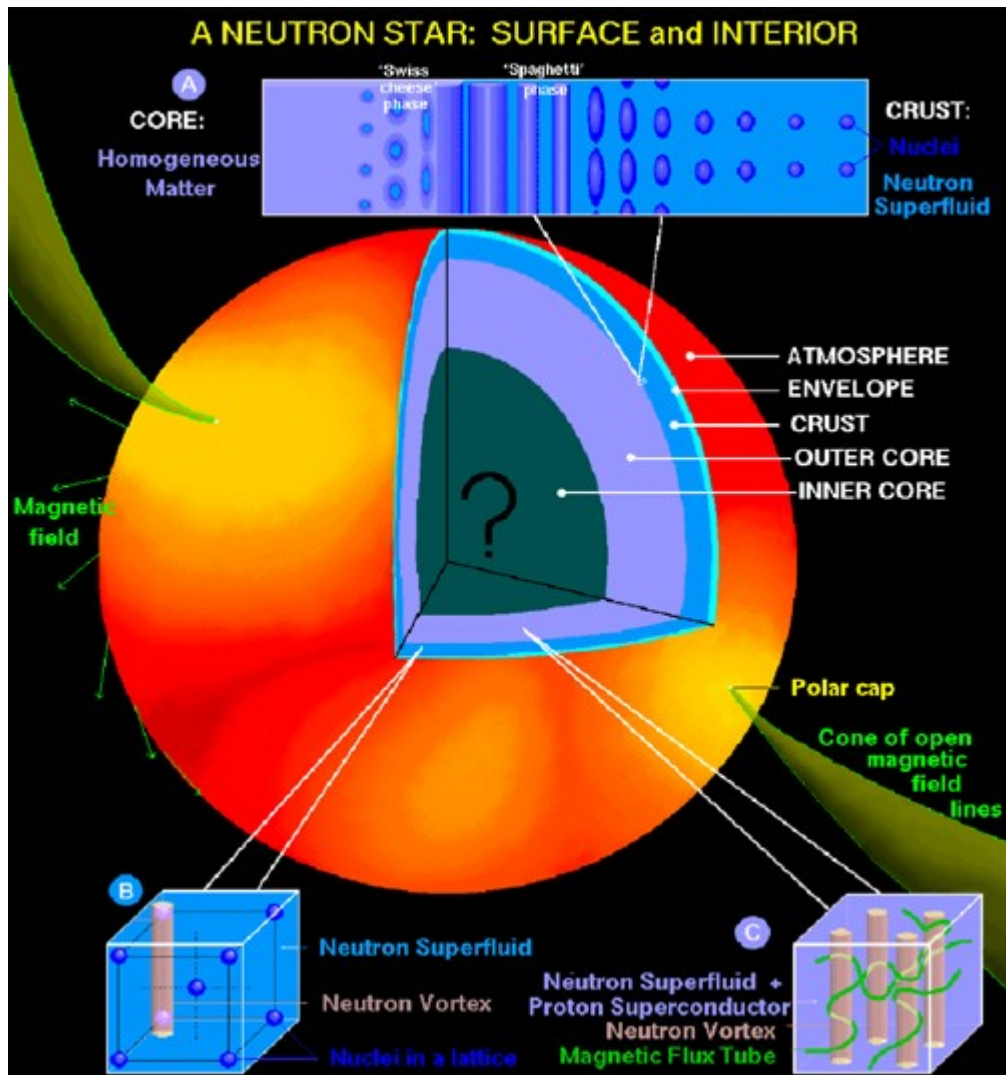
Key system: nuclei



Heiko Hergert's
propaganda plot

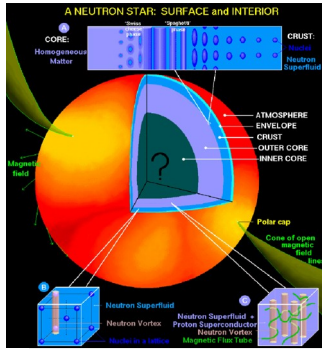
- Lots of recent progress
- Open-shell nuclei are the current frontier
- Goal is to study nuclei *from first principles* (when possible)

Key system: neutron stars



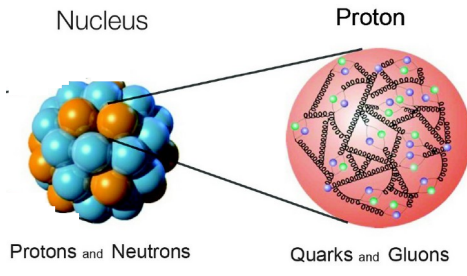
- Ultra-dense: 1.4 solar masses (or more) within a radius of 10 kilometres
- Terrestrial-like (outer layers) down to exotic (core) behaviour
- Observationally probed, i.e., not experimentally accessible
- Goal is to study neutron stars *from first principles* (when possible)

Outline

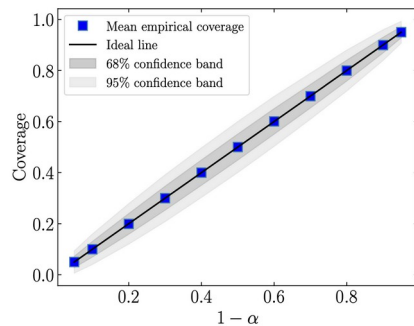


Credit: Dany Page

Motivation

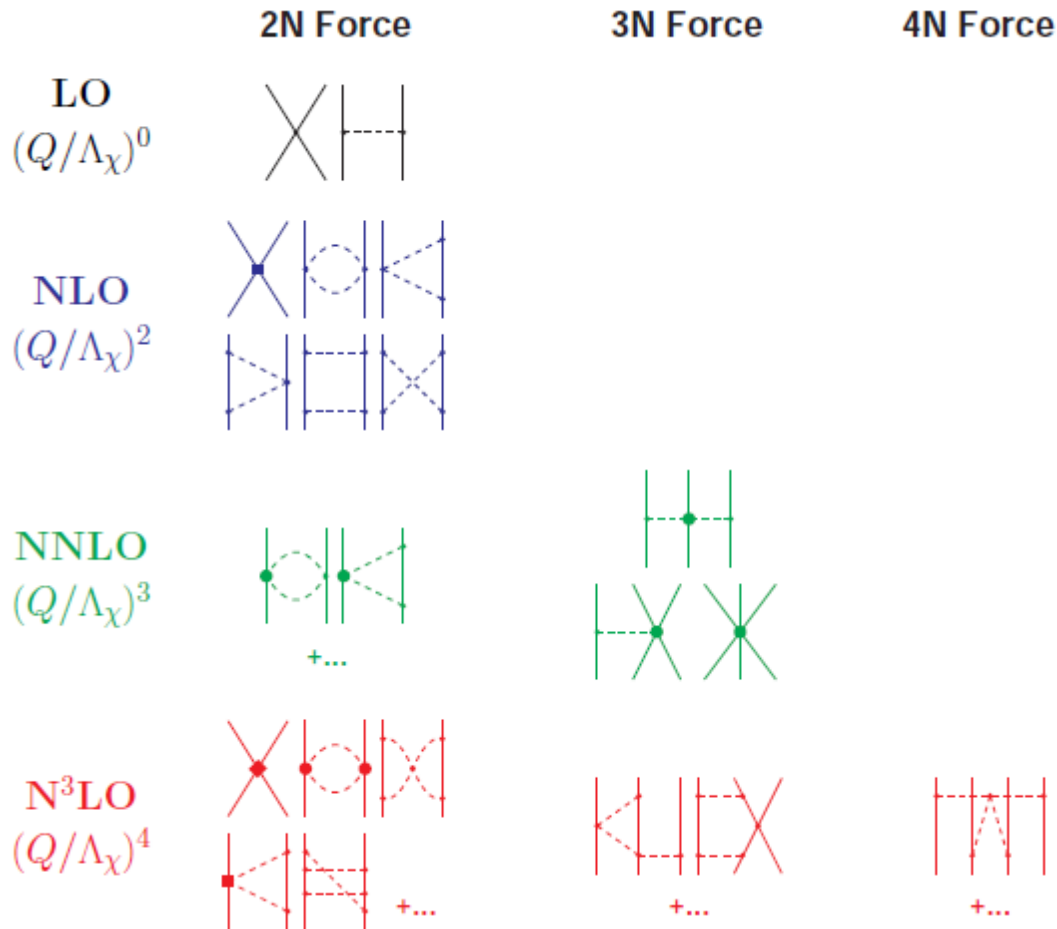


Nuclear methods



Recent results

Nuclear interactions



- Attempts to connect with underlying theory (QCD)
- Low-momentum expansion
- Naturally emerging many-body forces
- Low-energy constants from experiment or lattice QCD
- Now available in non-local, local, or semi-local varieties
- Power counting's relation to renormalization actively investigated

**But even with the interaction in place,
how do you solve the many-body problem?**

Nuclear many-body problem

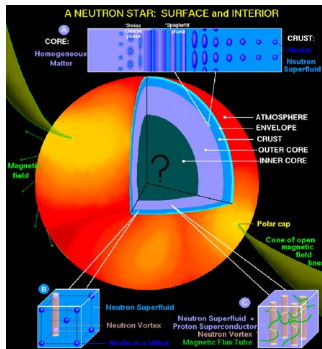
$$H\Psi = E\Psi$$

where

$$H = \sum_i K_i + \sum_{i<j} V_{ij} + \sum_{i<j<k} V_{ijk}$$

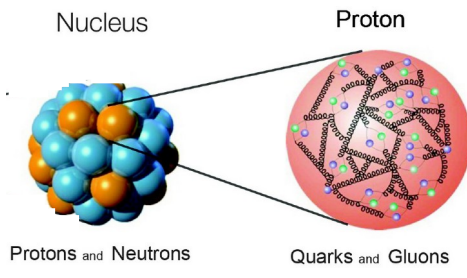
Wave function depends on coordinates, spin projections, and isospin projections, so we are faced with a large number of complex coupled second-order differential equations

Outline

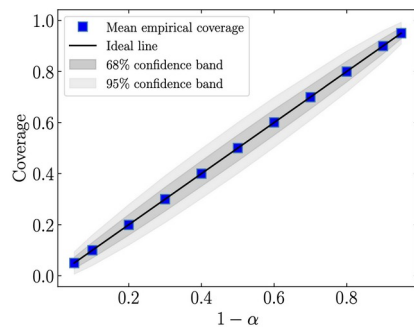


Credit: Dany Page

Motivation



Nuclear methods



Recent results

Recent results

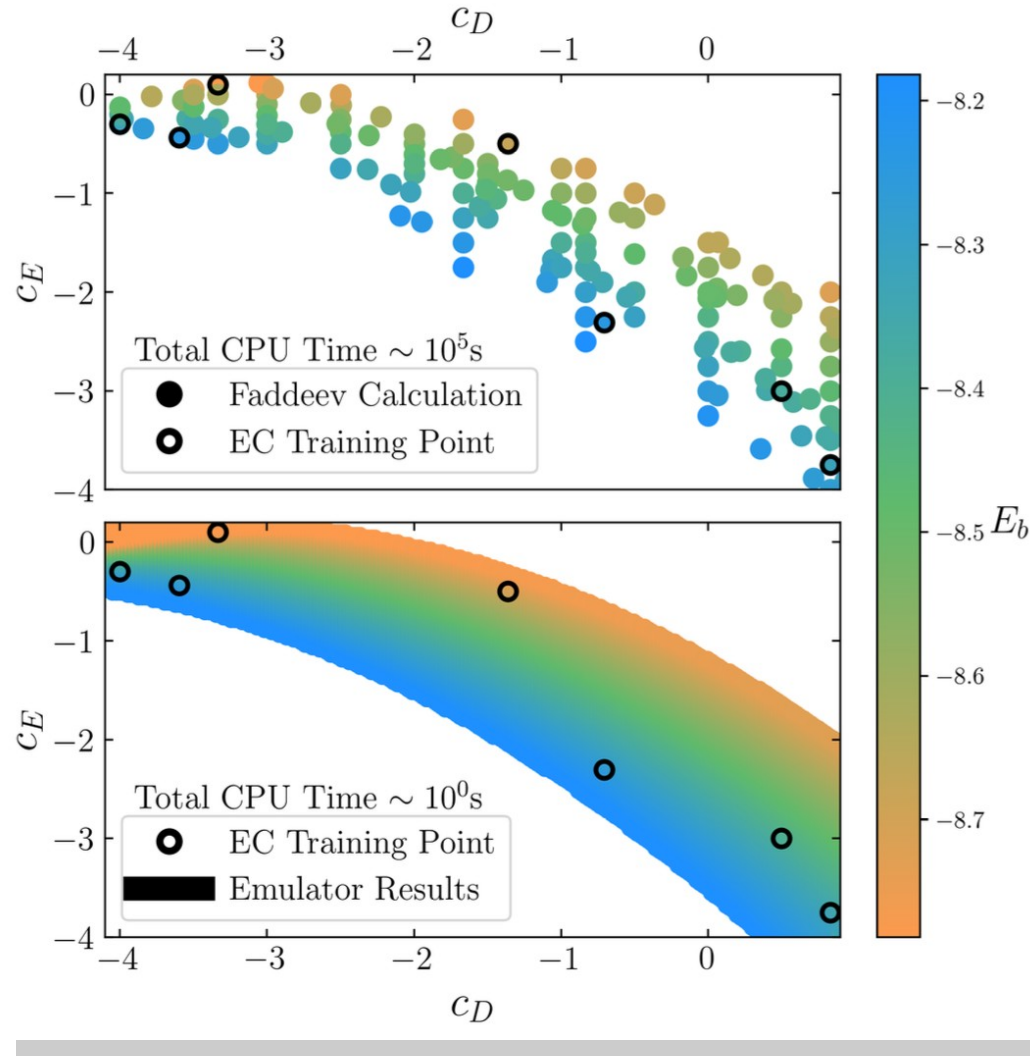
- **(QMC and emulators for light nuclei)**
- **Neural-network wave functions for light nuclei**
- **Conformal prediction for nucleon-nucleon scattering**

(QMC and emulators for light nuclei)

R. Curry, K. Hebeler, S. Gandolfi, A. Gezerlis,
A. Schwenk, R. Somasundaram, and I. Tews,
[arXiv:2510.15860](https://arxiv.org/abs/2510.15860)

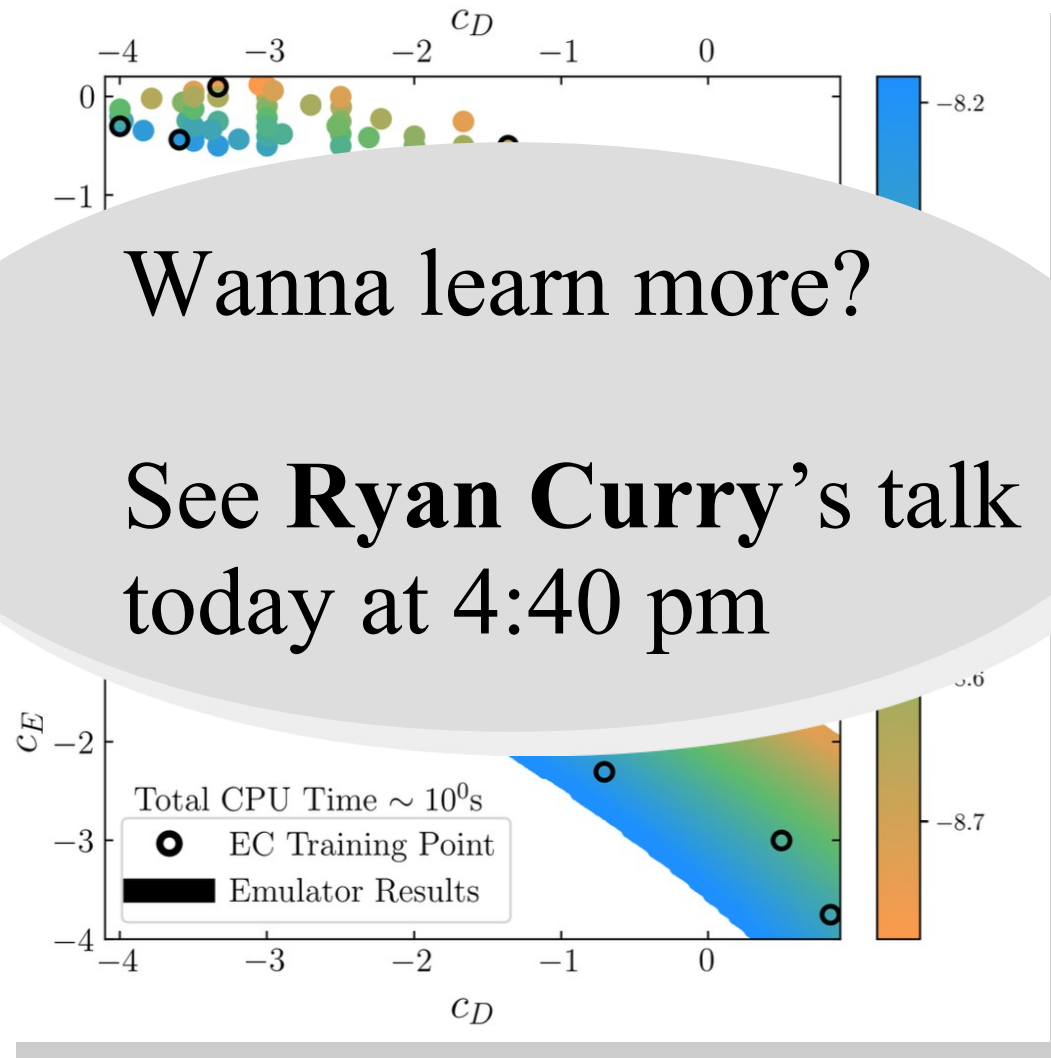
Varying the three-nucleon interaction

Emulators make the calculation so cheap that we can fill out our plots:



Varying the three-nucleon interaction

Emulators make the calculation so cheap that we can fill out our plots:



Neural-network wave functions for light nuclei

P. Weng, A. Gezerlis, and J. Holt, arXiv:2505.11442

Quantum Monte Carlo in one slide

Variational Monte Carlo

Encode parameters in trial wave function Ψ_V and use Rayleigh-Ritz principle to minimize expectation value of Hamiltonian

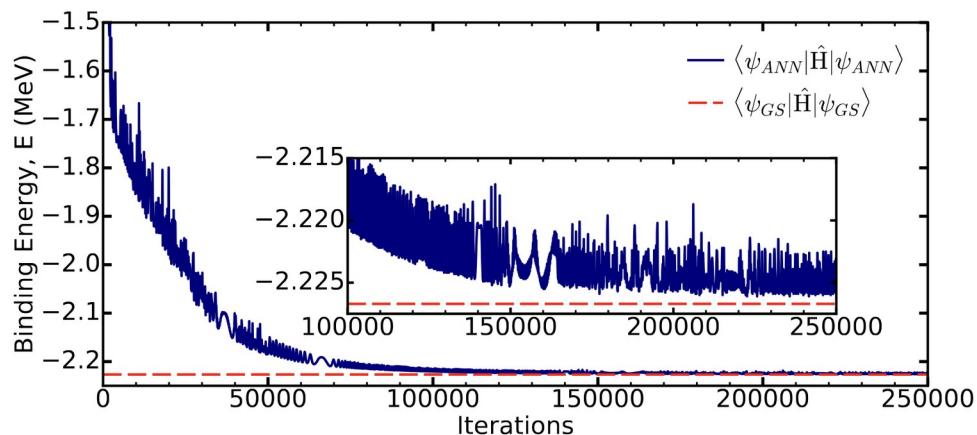
Diffusion Monte Carlo

Project out excited-state contributions, to reach the ground state

$$\Psi(\tau \rightarrow \infty) = \lim_{\tau \rightarrow \infty} e^{-(\mathcal{H} - E_T)\tau} \Psi_V$$

Earlier work

Deuteron



J. W. T. Keeble and A. Rios,
Phys. Lett. B **809**, 135743 (2020)

Light nuclei

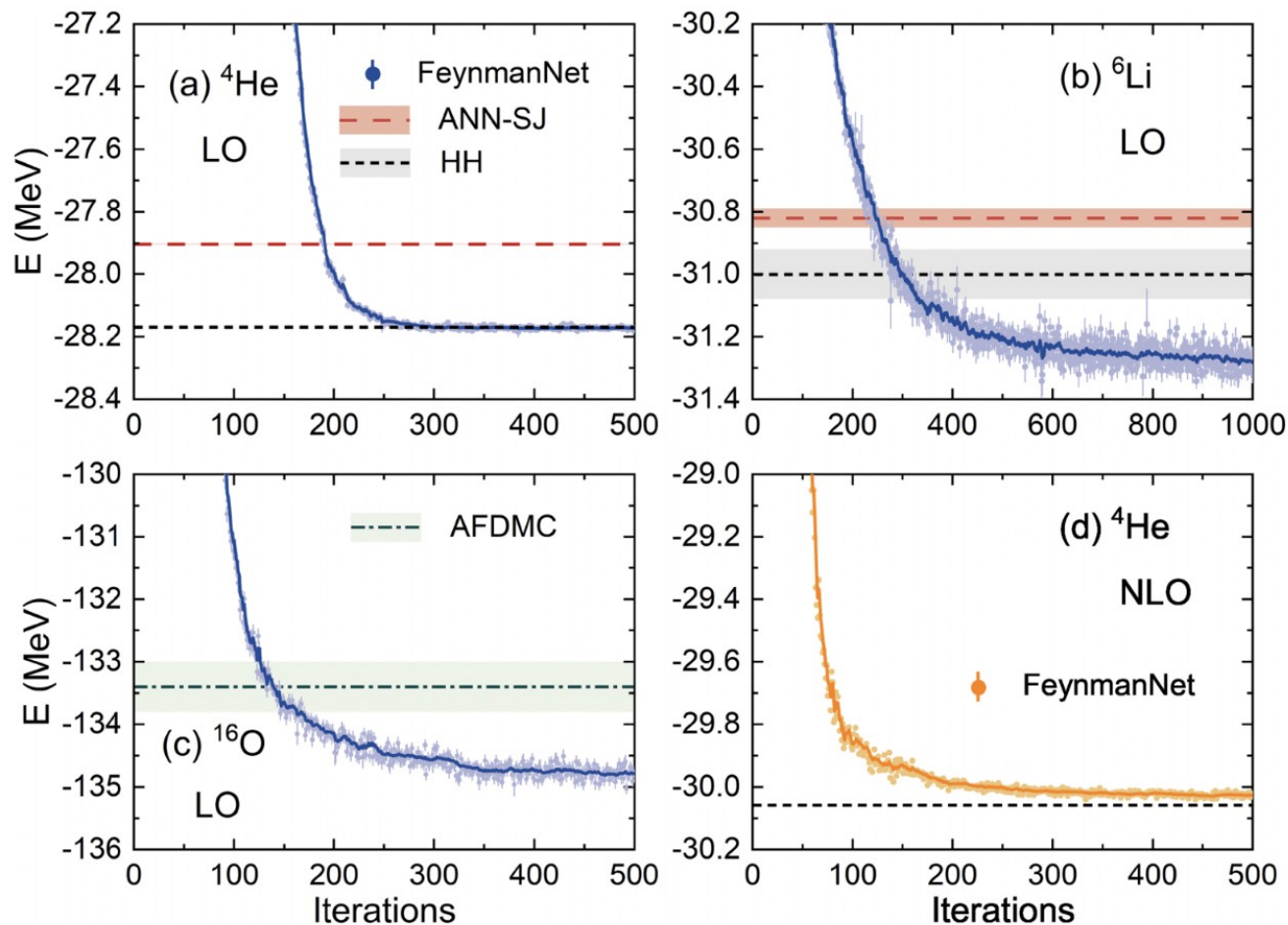
	Λ	VMC-ANN	VMC-JS	GFMC	GFMC _c
${}^2\text{H}$	4 fm^{-1}	-2.224(1)	-2.223(1)	-2.224(1)	-
	6 fm^{-1}	-2.224(4)	-2.220(1)	-2.225(1)	-
${}^3\text{H}$	4 fm^{-1}	-8.26(1)	-7.80(1)	-8.38(2)	-7.82(1)
	6 fm^{-1}	-8.27(1)	-7.74(1)	-8.38(2)	-7.81(1)
${}^4\text{He}$	4 fm^{-1}	-23.30(2)	-22.54(1)	-23.62(3)	-22.77(2)
	6 fm^{-1}	-24.47(3)	-23.44(2)	-25.06(3)	-24.10(2)

C. Adams, G. Carleo, A. Lovato, N. Rocco,
Phys. Rev. Lett. **127**, 022502 (2021)

N.B. Limited to pionless Hamiltonian

Earlier work

More on light nuclei



Y. L. Yang and P. W. Zhao, Phys. Rev. C **107**, 034320 (2023)
N.B. Also limited to pionless Hamiltonian

Neural networks for light nuclei

Spin-isospin correlations

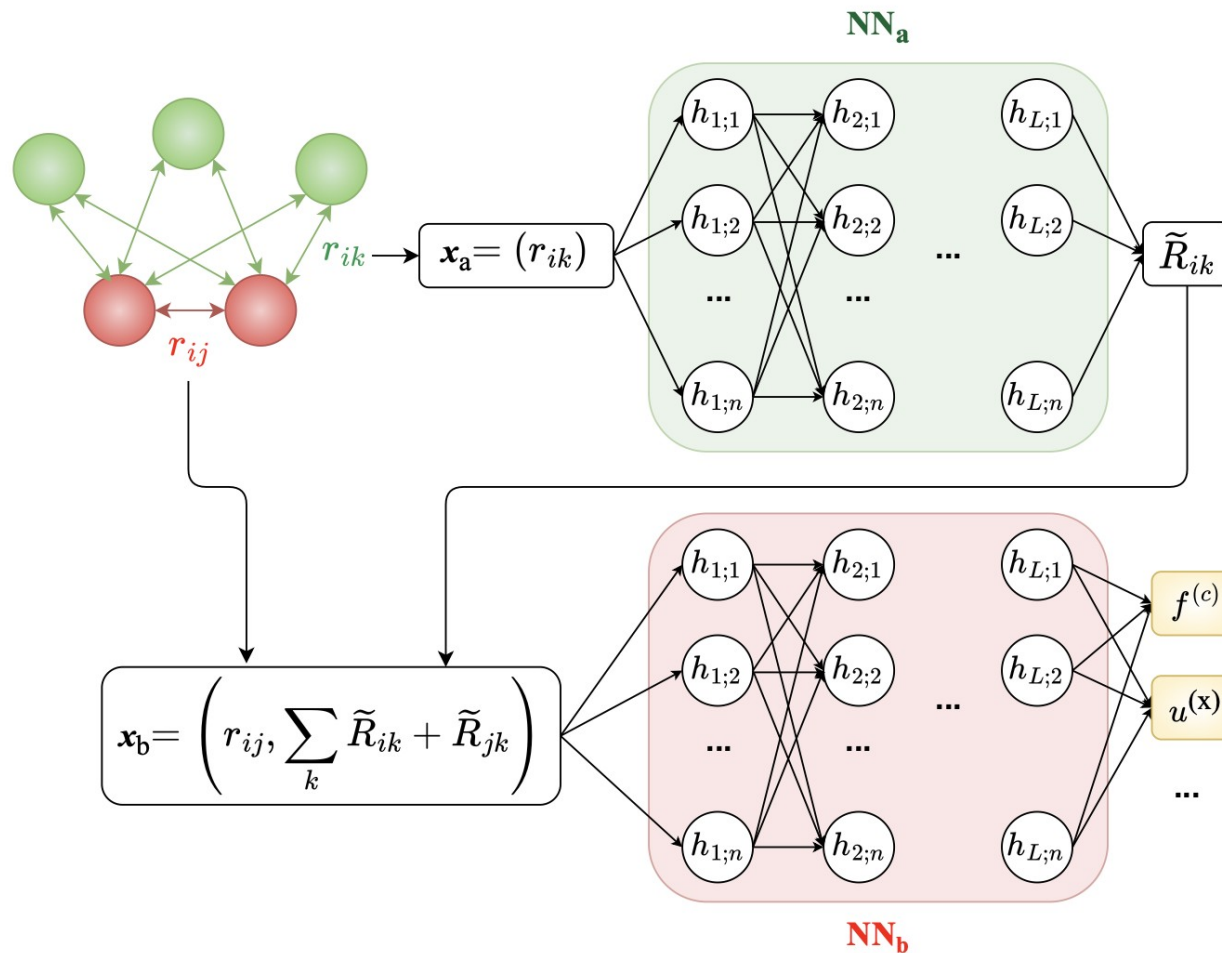
$$|\psi\rangle = \mathcal{S} \prod_{i < j} \left(1 + \sum_{\mathbf{x}} u_{ij}^{(\mathbf{x})} \hat{O}_{ij}^{(\mathbf{x})} \right) f_{ij}^{(c)} |\Phi\rangle$$

$$|\psi\rangle \rightarrow \left(1 + \sum_{i < j < k} \sum_{\text{cyc}} \sum_{\mathbf{x}} \epsilon^{(\mathbf{x})} \hat{V}_{ijk}^{(\mathbf{x})} \right) |\psi\rangle$$

for N²LO chiral Hamiltonian

P. Weng, A. Gezerlis, and J. Holt, arXiv:2505.11442

Neural networks for light nuclei



Neural networks for light nuclei

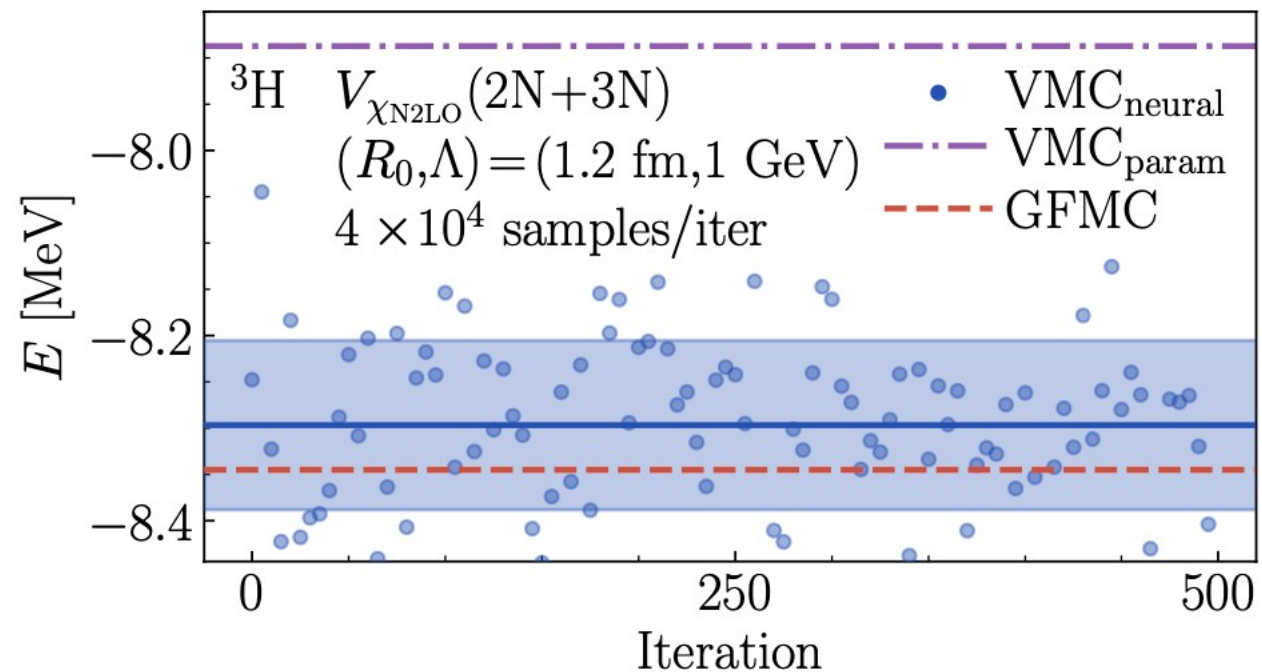
Nearly reproduces GFMC results
already at the VMC level

$E = E_k + V_{\chi\text{N}^2\text{LO}}(2\text{N})$				
	R_0 [fm]	E_{neural} [MeV]	E_{GFMC} [MeV]	$ \Delta E / E_{\text{GFMC}} $
${}^3\text{H}$	1.0	-7.338 ± 0.008	-7.554 ± 0.007	2.9%
	1.1	-7.500 ± 0.006	-7.625 ± 0.005	1.6%
	1.2	-7.678 ± 0.005	-7.740 ± 0.005	0.8%
${}^2\text{H}$	1.0	-2.217 ± 0.005	-2.21 ± 0.02	0.3%
	1.2	-2.212 ± 0.004	-2.20 ± 0.03	0.5%

P. Weng, A. Gezerlis, and J. Holt, arXiv:2505.11442

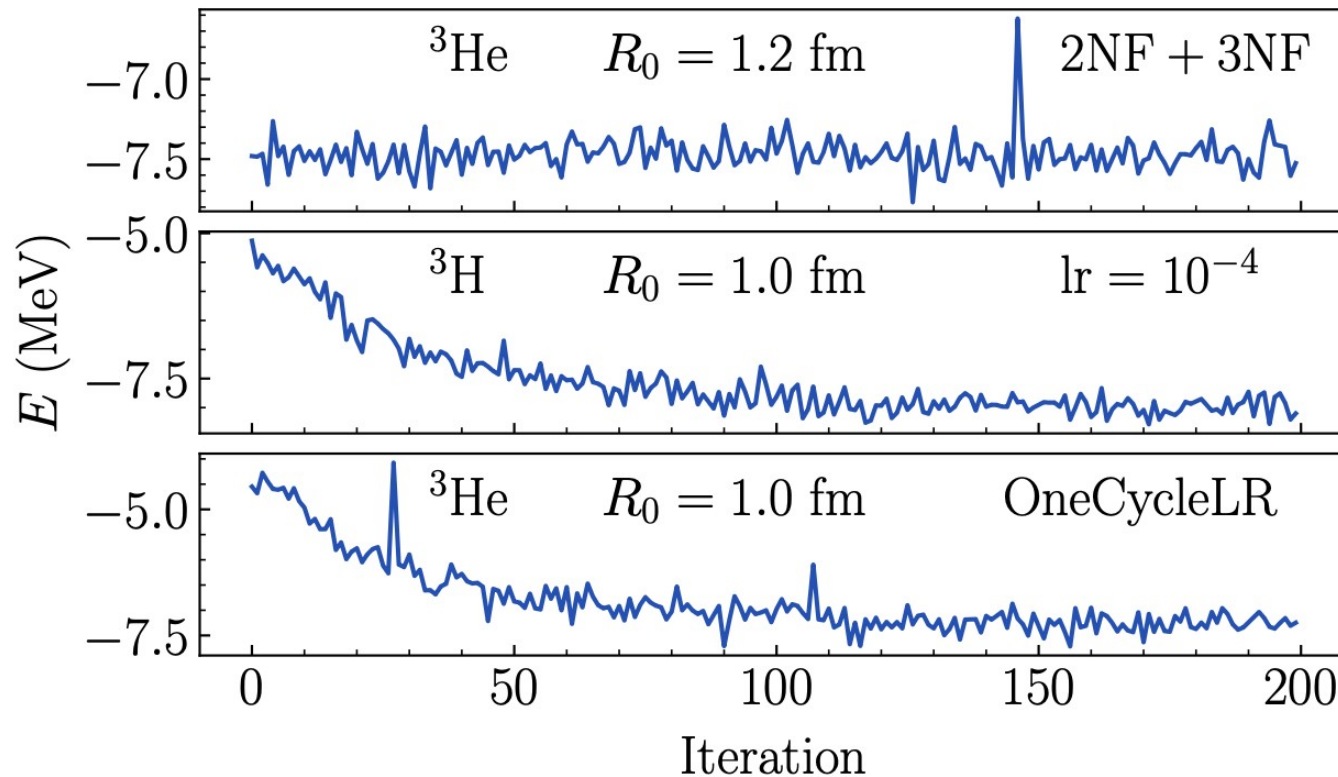
Neural networks for light nuclei

Dramatic improvement over
standard/parametric VMC
employed before, e.g., AFDMC



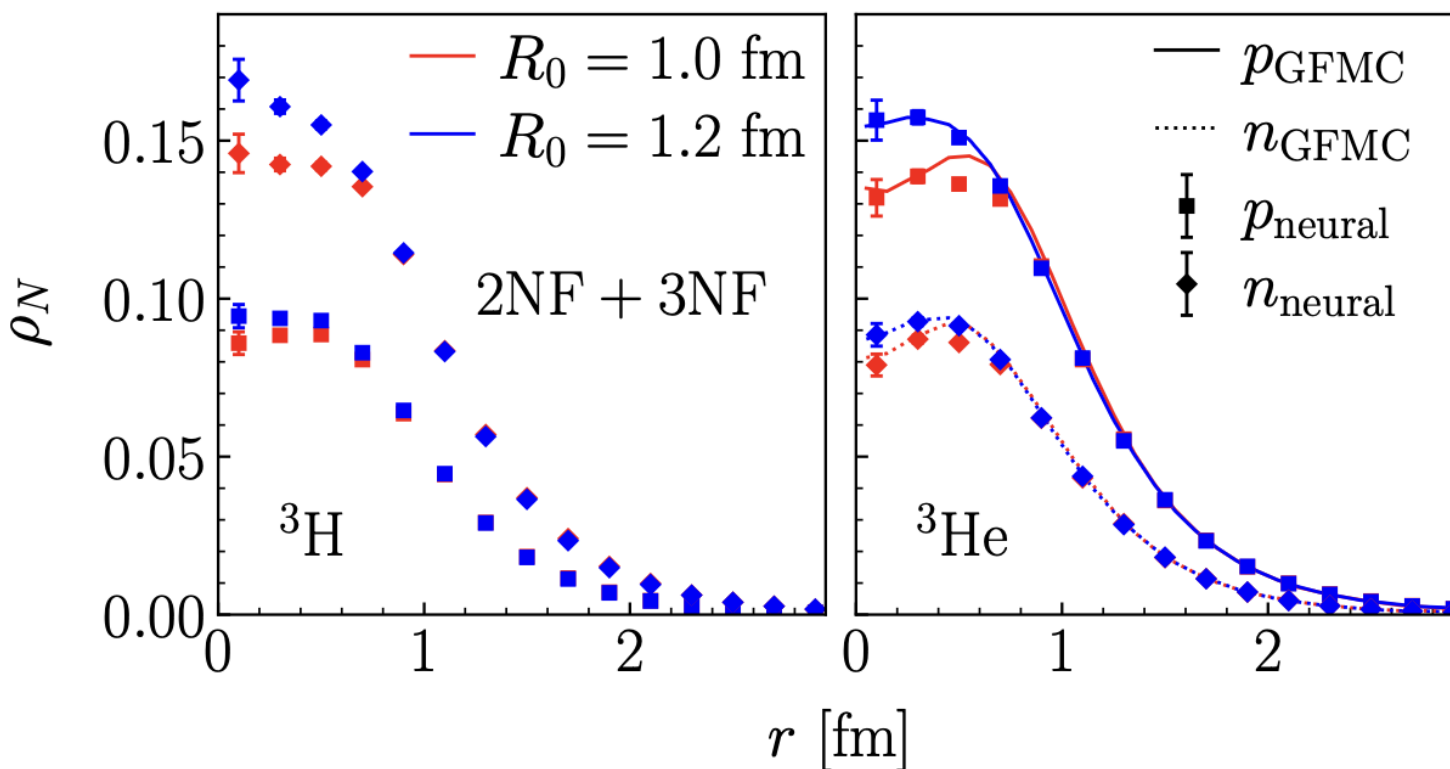
Neural networks for light nuclei

Transfer training: use one potential/nucleus as a starting point for other studies



Neural networks for light nuclei

Point-particle distributions:
$$\rho_N(r) = \frac{1}{4\pi r^2} \left\langle \psi \left| \sum_{i=1}^A \frac{1 + \zeta \tau_{z,i}}{2} \delta(r - |\mathbf{r}_i - \mathbf{R}_{\text{cm}}|) \right| \psi \right\rangle$$



Conformal prediction for nuclear physics

H. Yousefi Dezdarani, R. Curry, and A. Gezerlis, Phys. Rev. C, **113**, 014004 (2026)

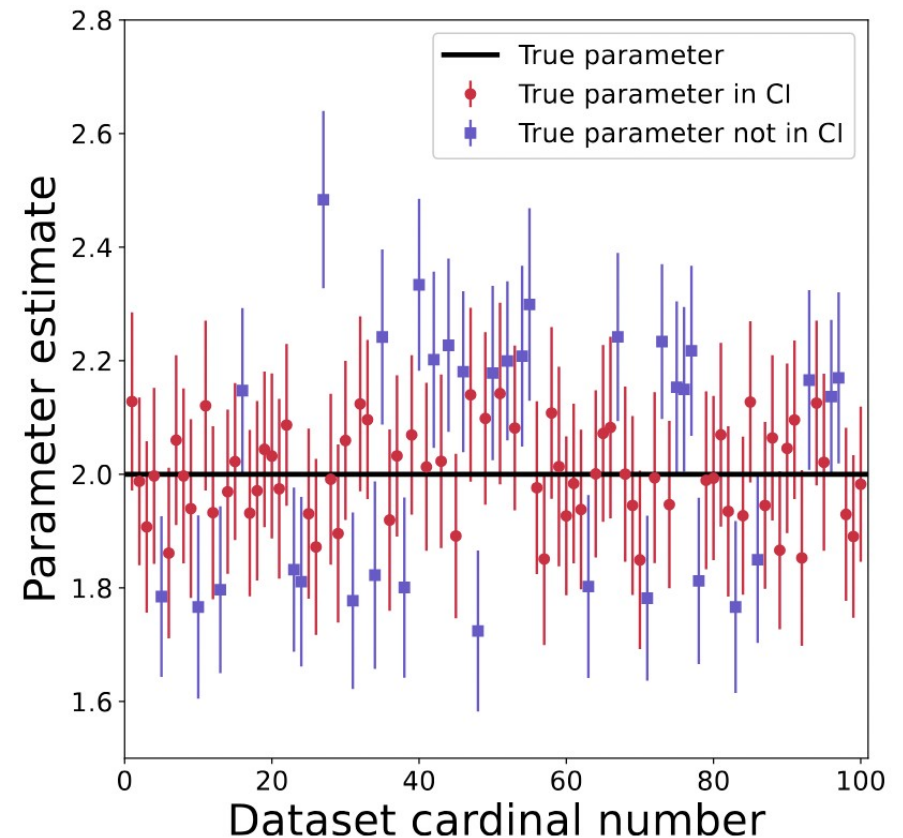
H. Yousefi Dezdarani, R. Curry, C. L. Armstrong, and A. Gezerlis, *in preparation*

A hint of the philosophy of statistics

(Frequentist) confidence interval

E.g., maximize likelihood and take an $n\sigma$ error bar around that point.

Confidence interval *does not* imply degree of belief about our single dataset, but *does* provide guaranteed coverage across datasets.



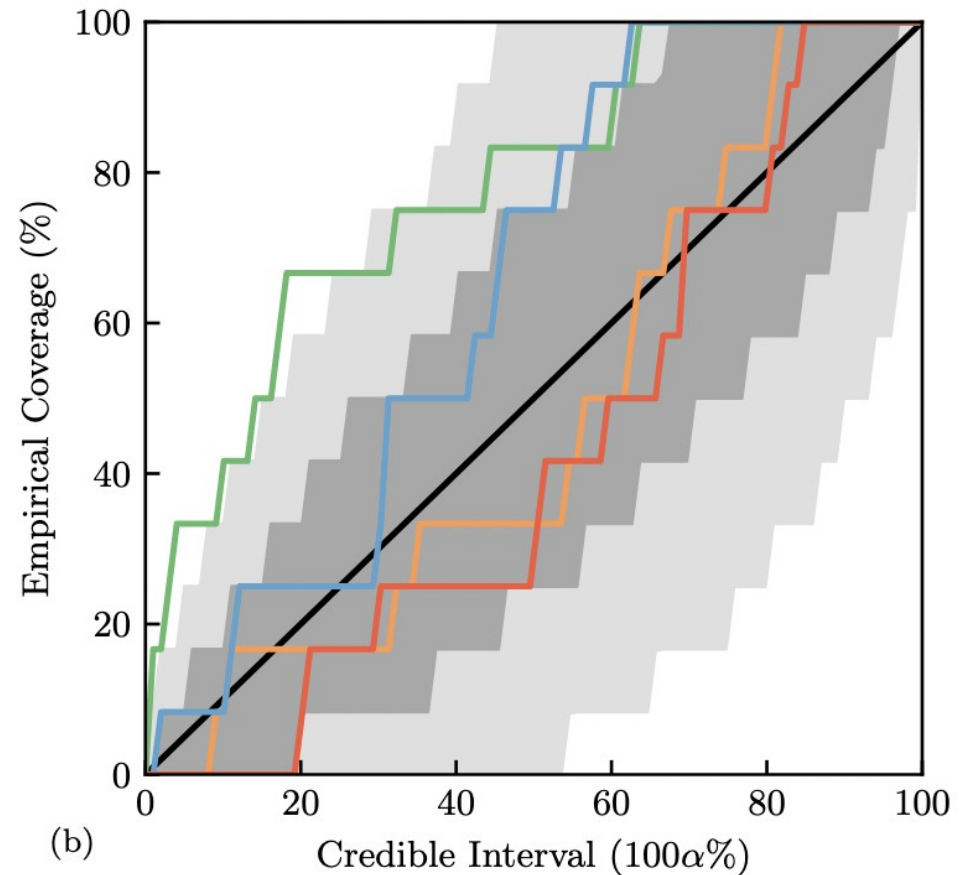
A hint of the philosophy of statistics

(Bayesian) credible interval

E.g., maximize posterior and take an $n\sigma$ error bar around that point.

Credible interval *does* imply degree of belief about our single dataset, but *does not* provide guaranteed coverage across datasets.

N.B. BUQEYE collaboration interprets coverage order-by-order



A hint of the philosophy of statistics

Summarizing attractive features

- (Frequentist) confidence interval has guaranteed coverage across datasets
- (Bayesian) credible interval combines prior and likelihood to encapsulate degree of belief about our single dataset

Can you get the best of both worlds?

- Yes (*contra* Betteridge's law of headlines)
- Conformal prediction is a tool that post-processes any pre-trained model to produce guaranteed coverage

Conformal prediction

- Distribution-free and model-agnostic uncertainty-quantification method
- Provides finite-sample prediction intervals with guaranteed coverage
- It accomplishes this by employing the quantile function (inverse of the cumulative distribution function) in an ingenious way:

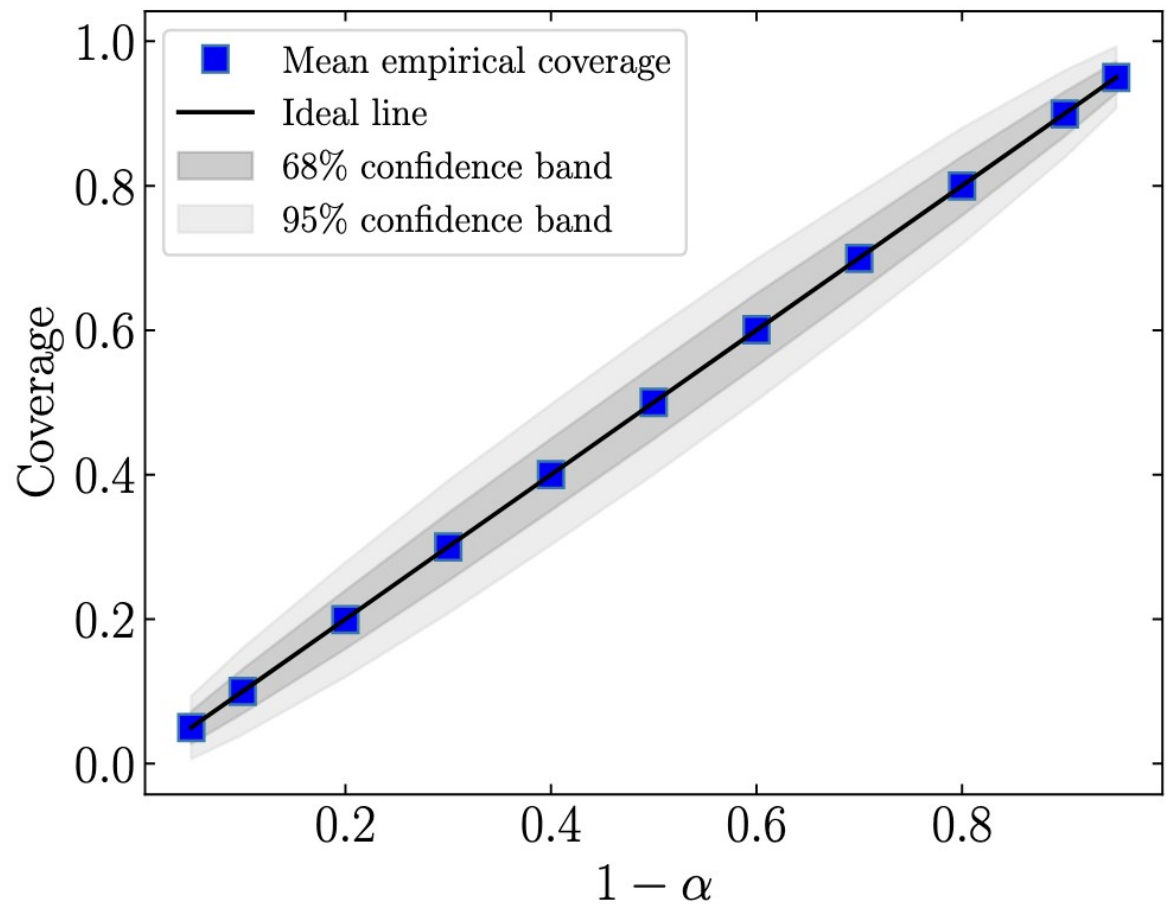
$$C(X_{n+1}) = [Q_Y \left(\frac{\alpha}{2} \mid X_{n+1} \right) - q, Q_Y \left(1 - \frac{\alpha}{2} \mid X_{n+1} \right) + q]$$

$$\text{where } q = Q_S(1 - \alpha)$$

Conformal prediction: scattering

Application: BUQEYE pointwise model

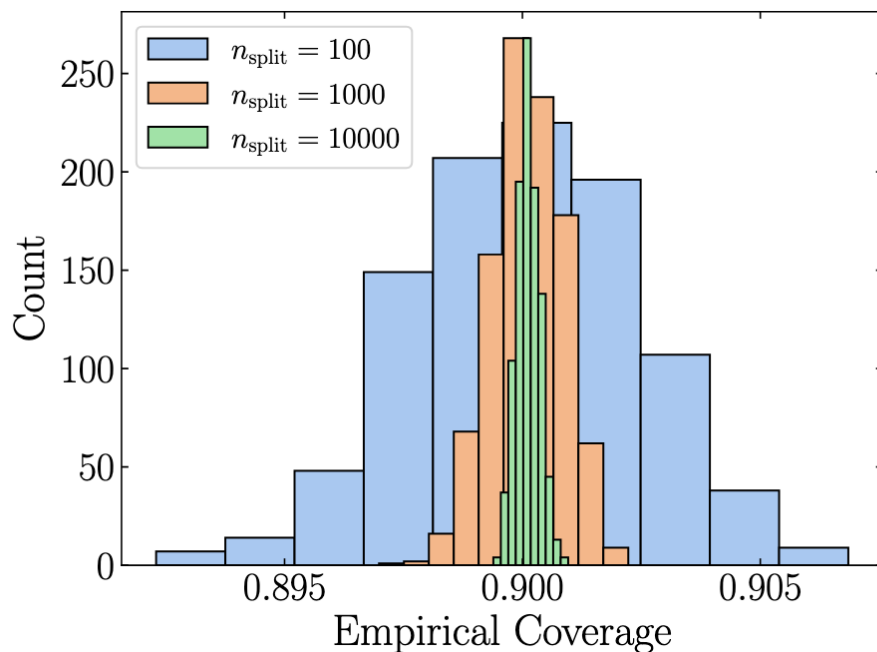
- Two-nucleon total cross section at $E = 50$ MeV
- Empirical coverage over 4000 independent trials
- There is near-perfect alignment between empirical coverage and ideal line



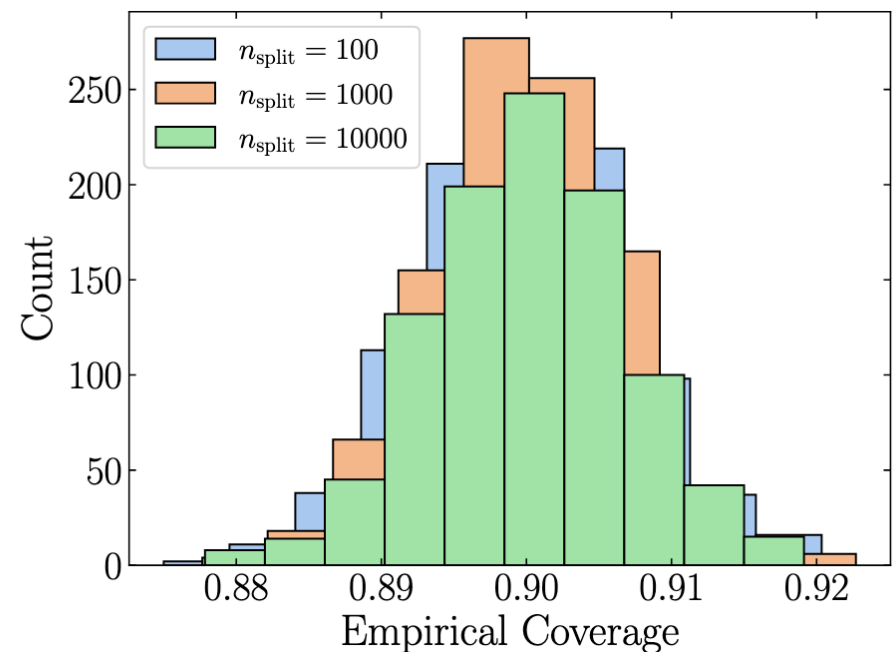
Conformal prediction: scattering

Application: BUQEYE Gaussian-process model

We drew posterior samples using the BUQEYE open-source code



conformal

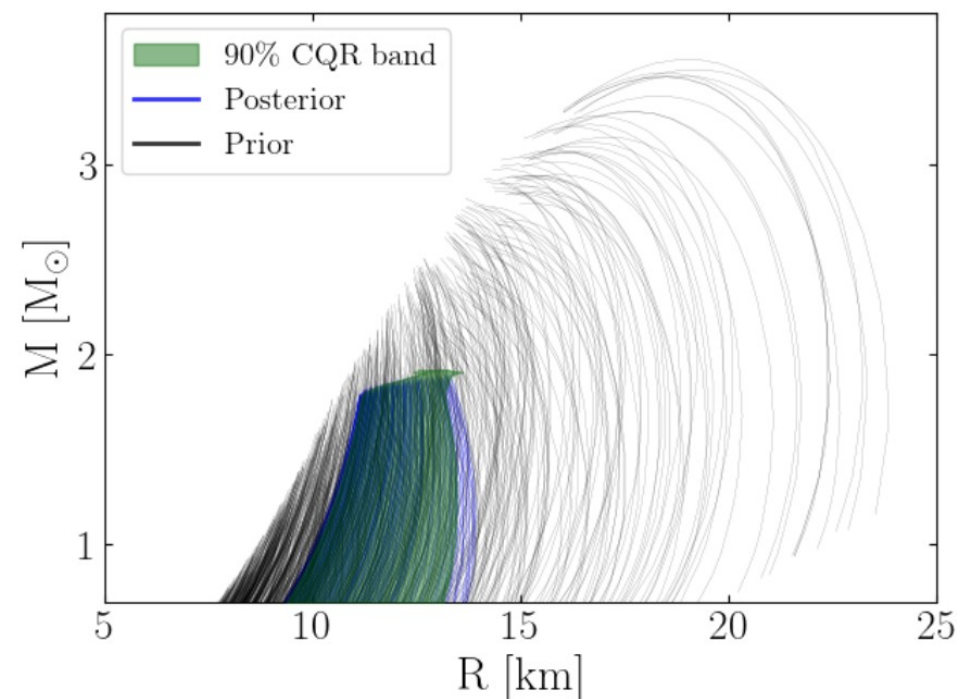
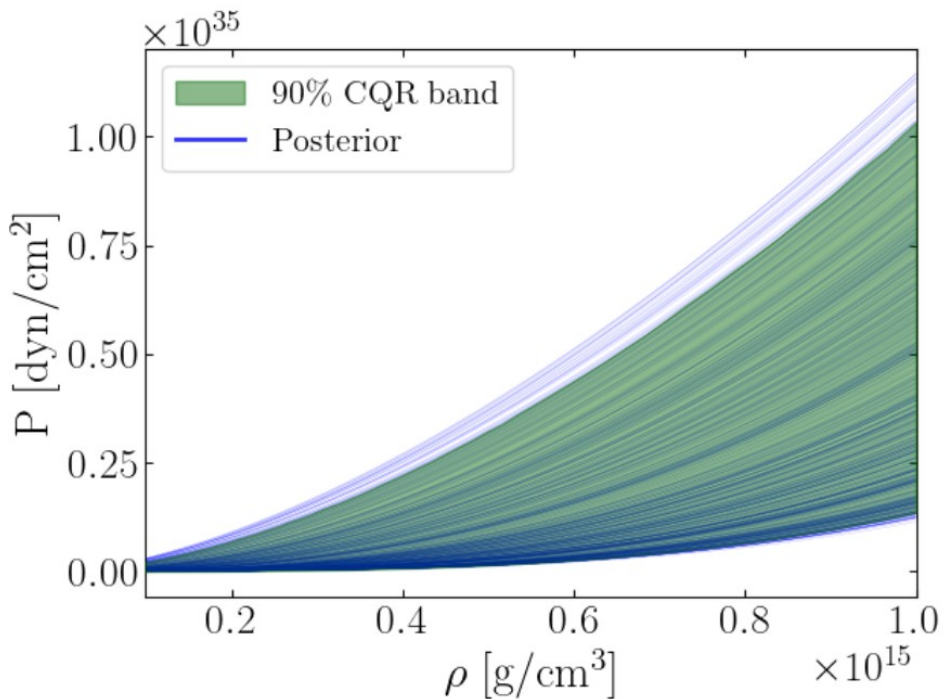


Bayesian

Conformal prediction: neutron stars

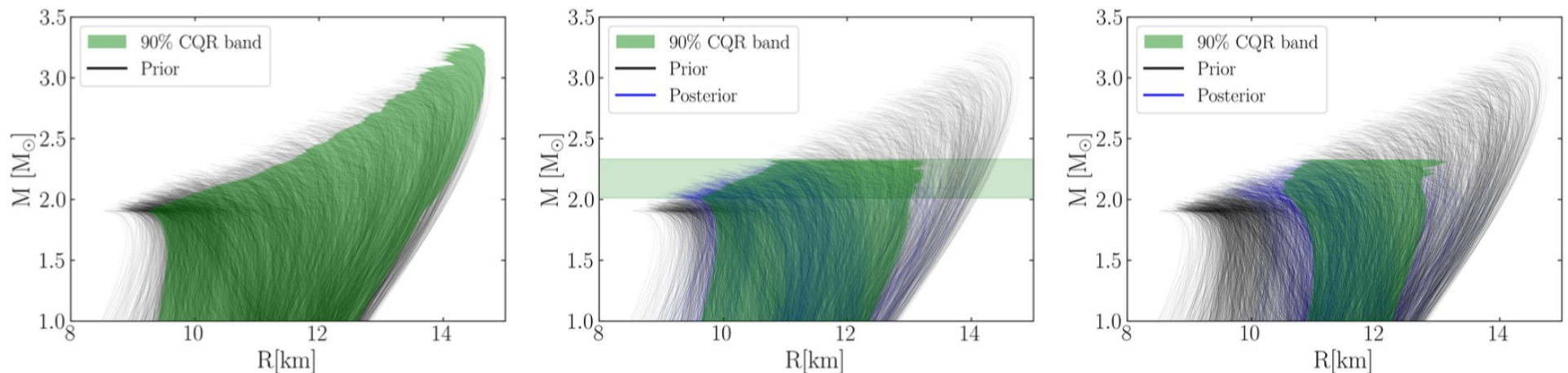
Application: TOV equations for polytrope

$$\frac{dP(r)}{dr} = -\frac{G m(r) \rho(r)}{r^2} \left[1 + \frac{P(r)}{\rho(r)c^2} \right] \left[1 + \frac{4\pi r^3 P(r)}{m(r)c^2} \right] \left[1 - \frac{2Gm(r)}{c^2 r} \right]^{-1}$$
$$\frac{dm(r)}{dr} = 4\pi r^2 \rho(r)$$



Conformal prediction: neutron stars

Application: multimessenger analysis (NMMA collaboration)



Post-processing carried out using individual EOS samples

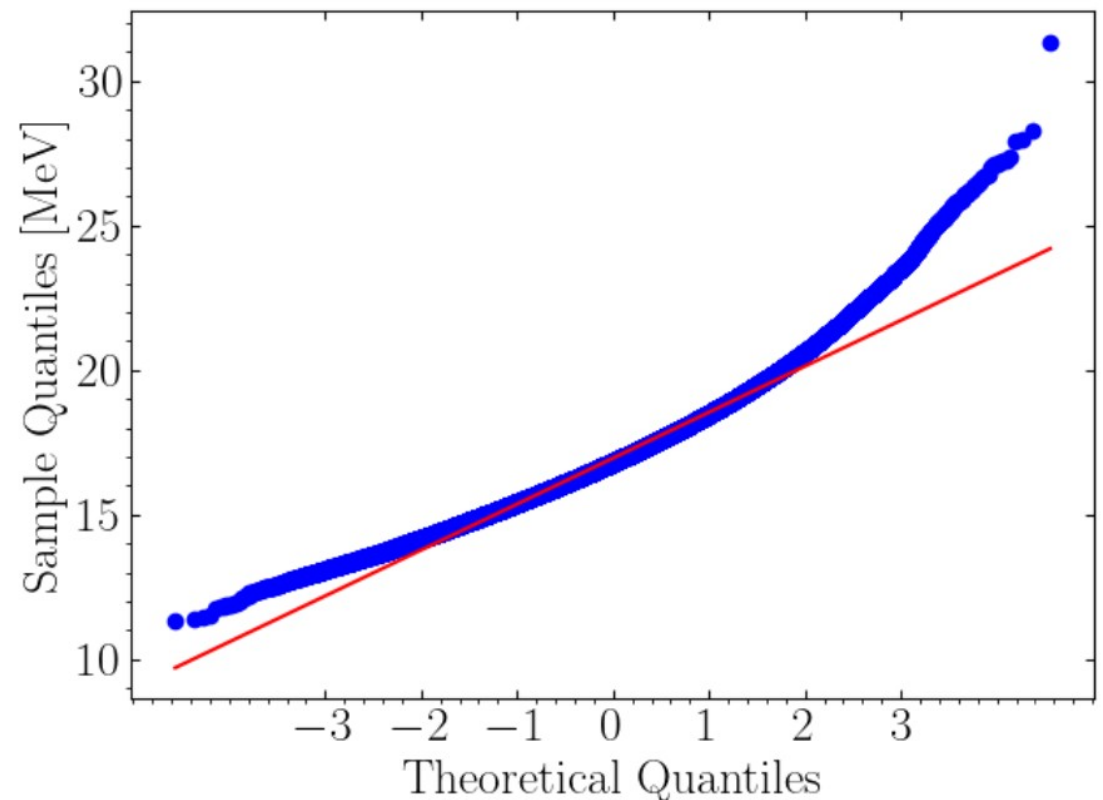
T. Dietrich *et al*, *Science* **370**, 1450 (2020)

H. Yousefi Dezdarani, R. Curry, C. L. Armstrong, and A. Gezerlis, *in preparation*

Conformal prediction: neutron stars

Application: QMC neutron matter equation of state (1)

- Deviations from the diagonal reflect a non-normal distribution
- Non-gaussianity motivates use of distribution-free method

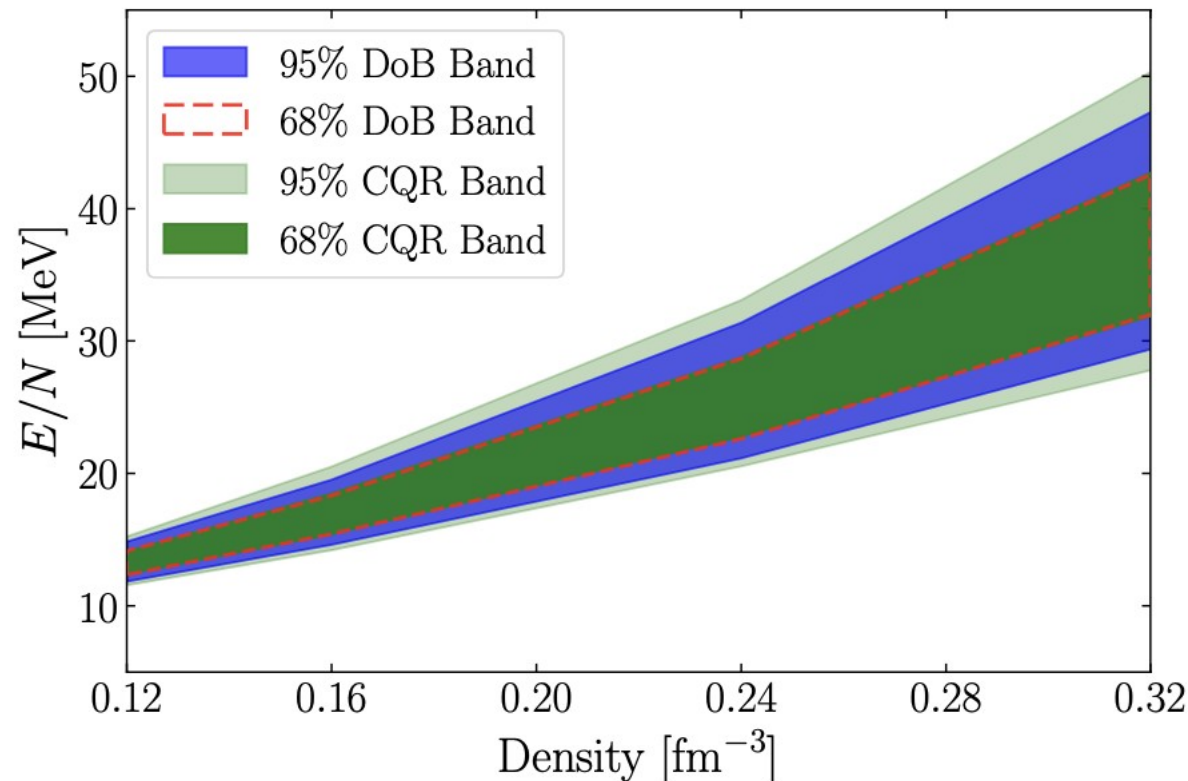


C. L. Armstrong *et al*, Phys. Rev. Lett. **135**, 142501 (2025)

H. Yousefi Dezdarani, R. Curry, C. L. Armstrong, and A. Gezerlis, *in preparation*

Conformal prediction: neutron stars

Application: QMC neutron matter equation of state (2)



DoB = Bayesian
CQR = conformal

Conformal prediction
interval wider than
Bayesian one

C. L. Armstrong *et al*, Phys. Rev. Lett. **135**, 142501 (2025)

H. Yousefi Dezdarani, R. Curry, C. L. Armstrong, and A. Gezerlis, *in preparation*

Conclusions

- We used PMM and EC emulators together with QMC studies of light nuclei
- We used chiral Effective Field Theory interactions in neural-network studies of light nuclei
- We applied conformal prediction to nuclear scattering and neutron star EOSs

Acknowledgments

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MINISTÈRE DE LA RECHERCHE ET DE L'INNOVATION

Collaborators

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Texas A&M

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TU Darmstadt

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