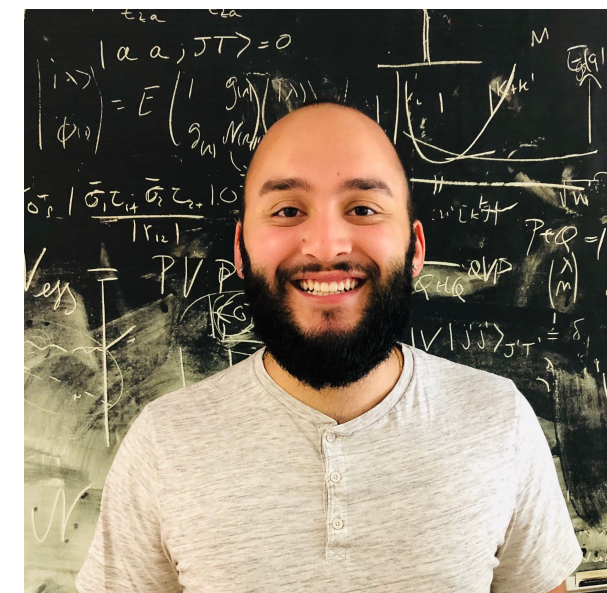
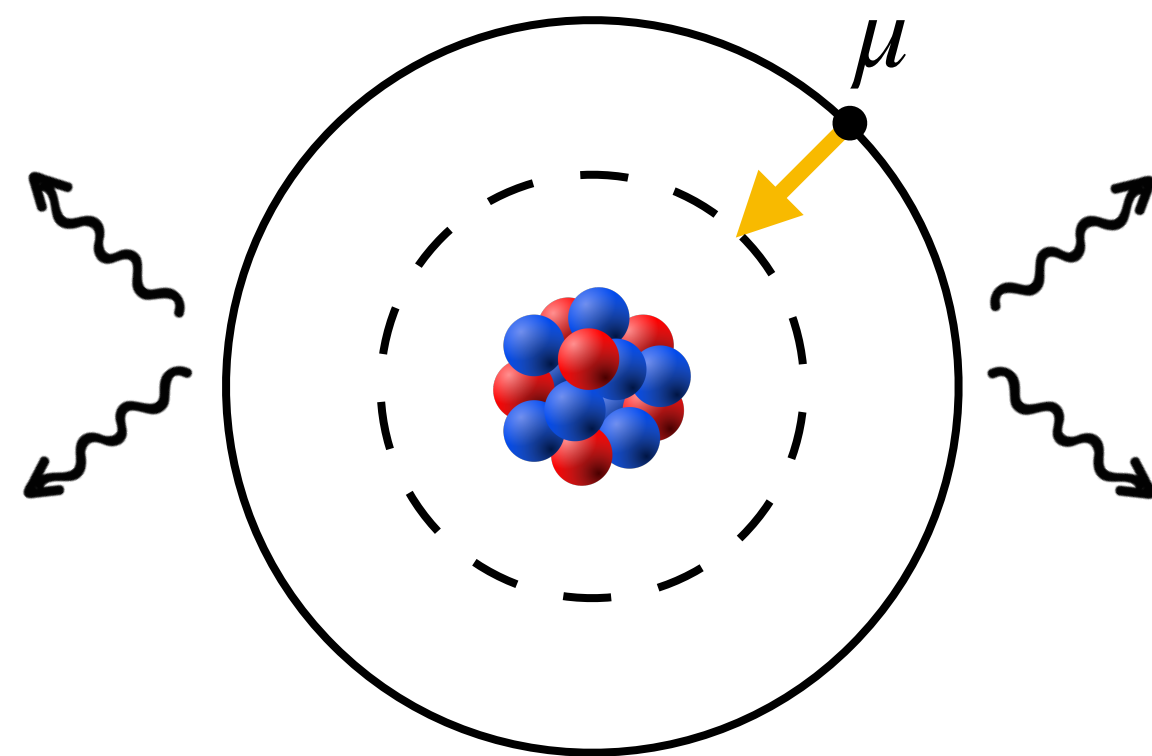


# Ab initio nuclear corrections to light muonic atoms

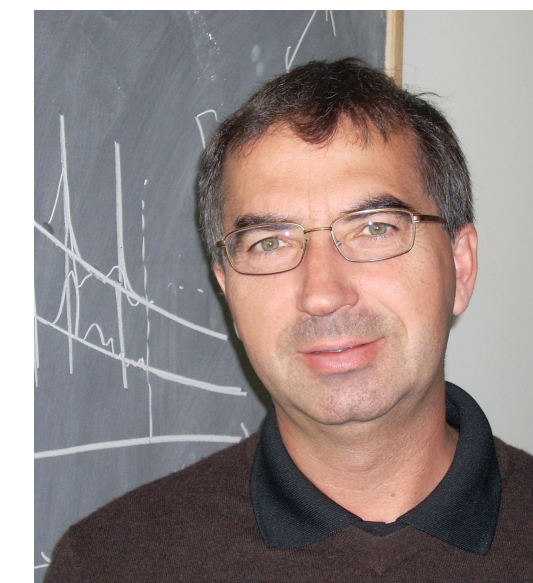
No-Core Shell Model calculations of polarizability for  ${}^6\text{Li}$  and  ${}^7\text{Li}$  atoms



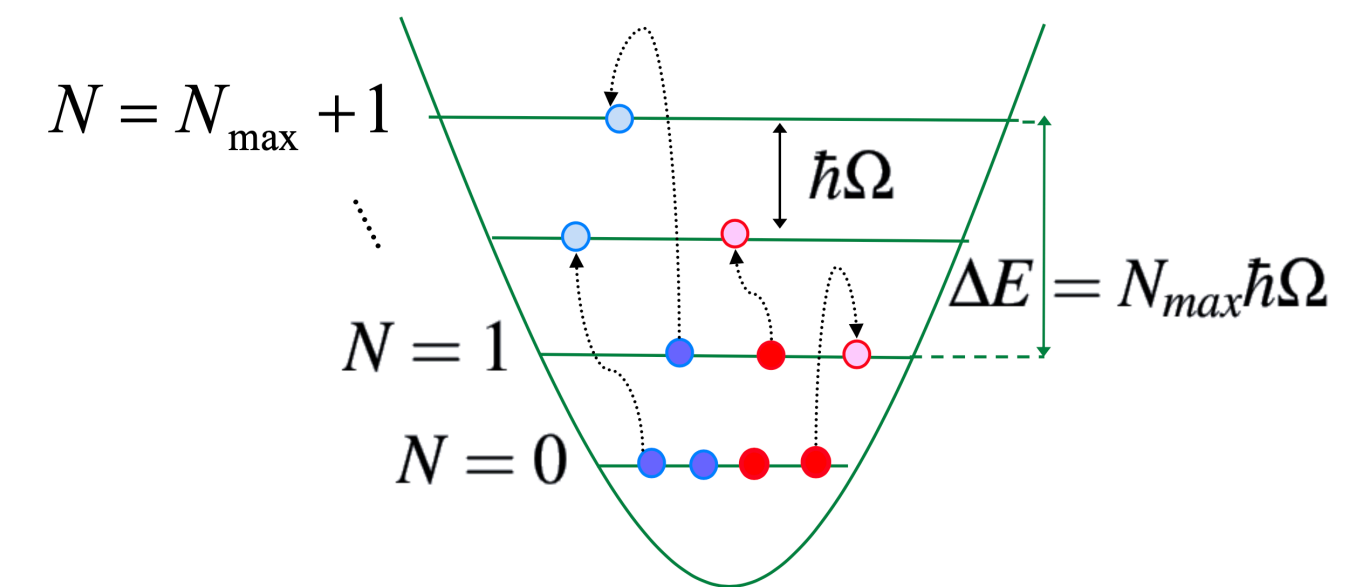
Michael Gennari



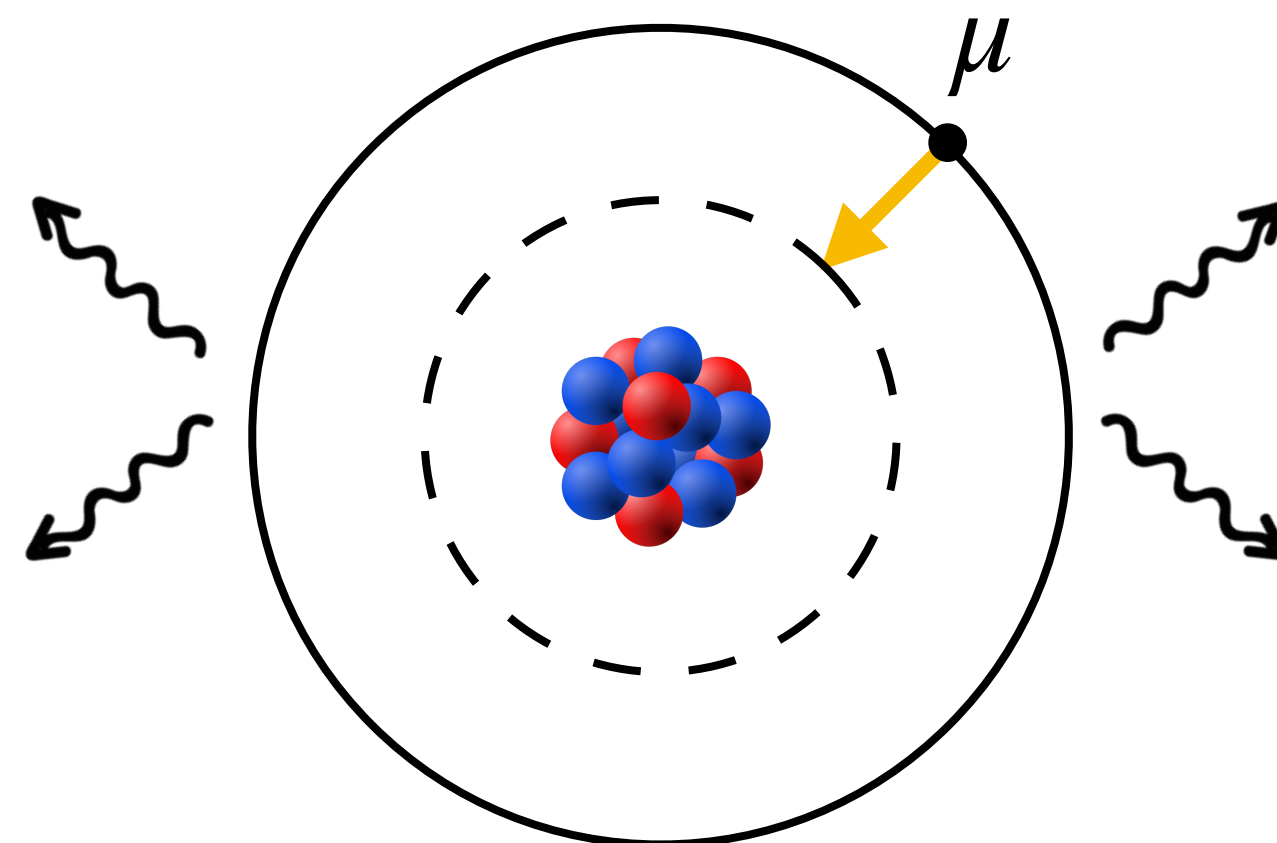
Diego Araujo Nájera



Petr Navratil

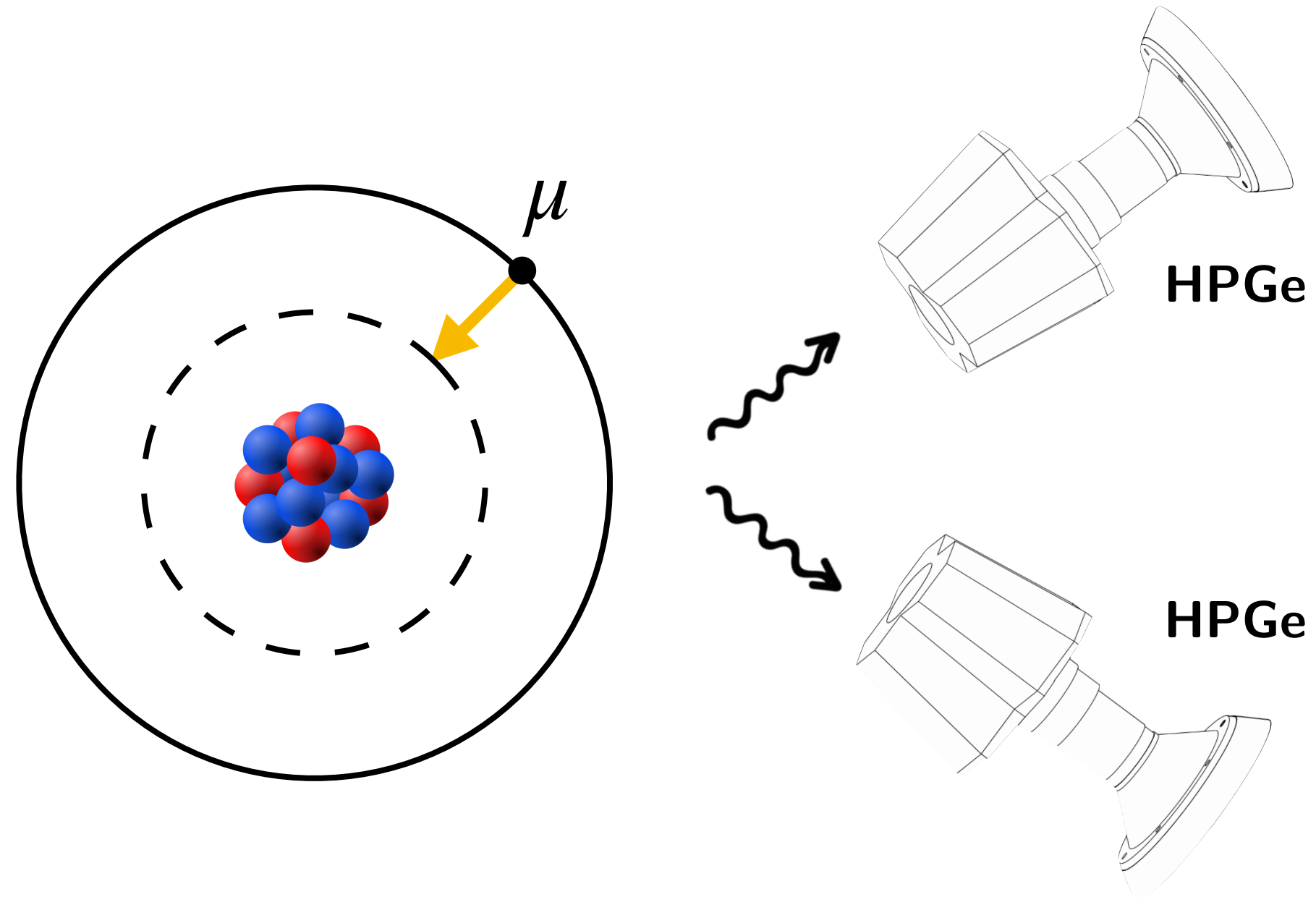


# Muonic atoms as a precision probe



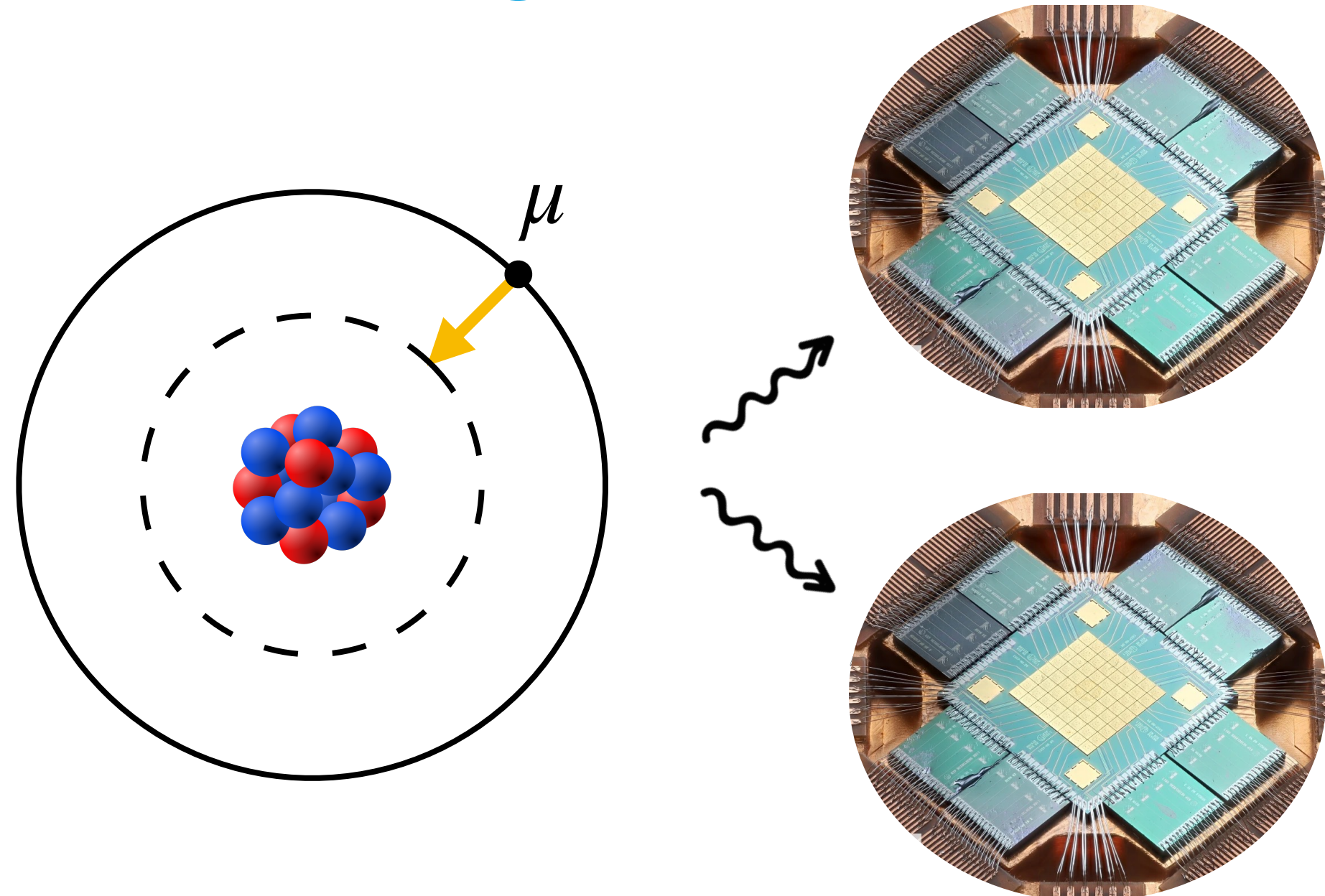
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Observing muonic atoms



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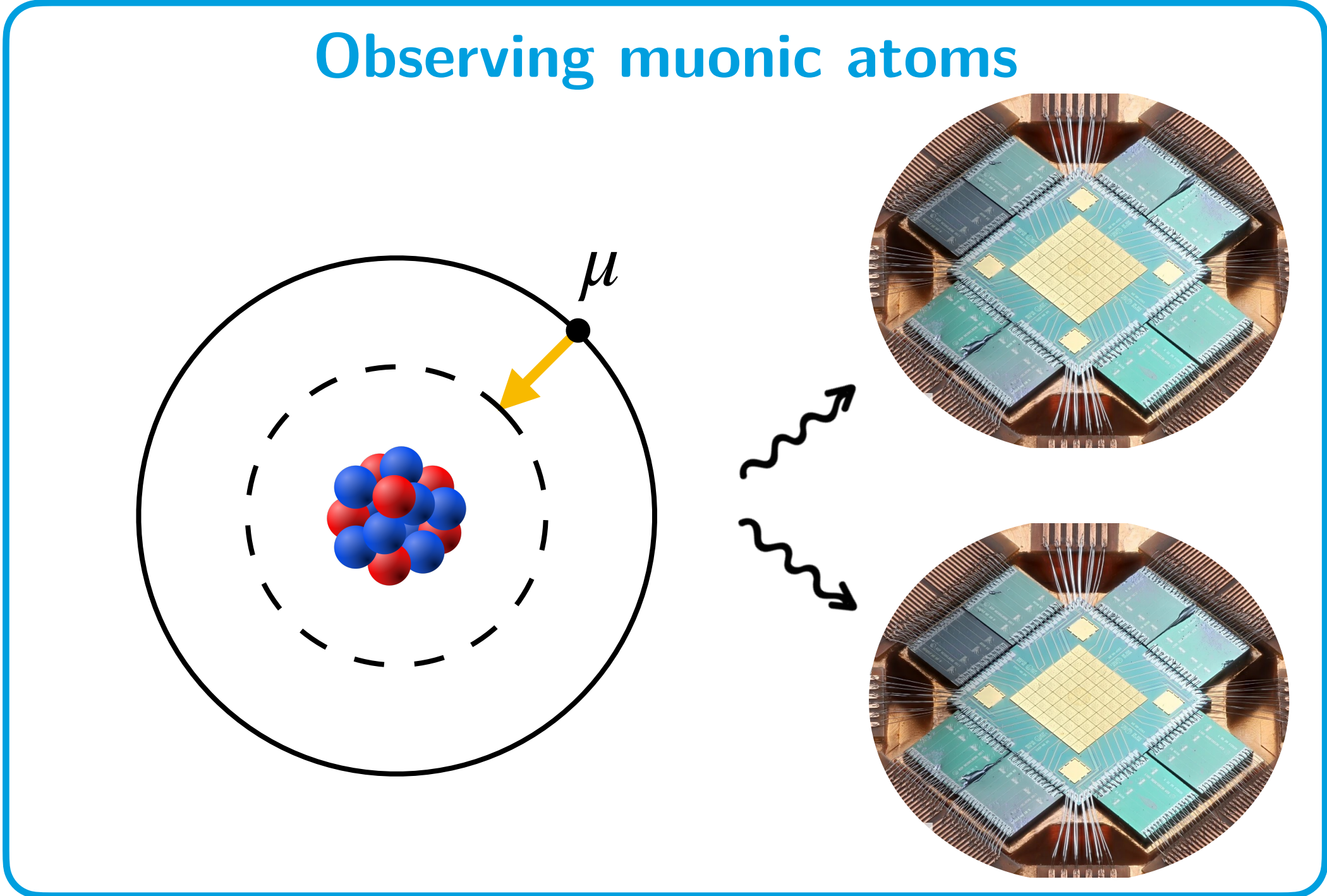


[Unger et al. J. Low Temp. Phys. (2024)]

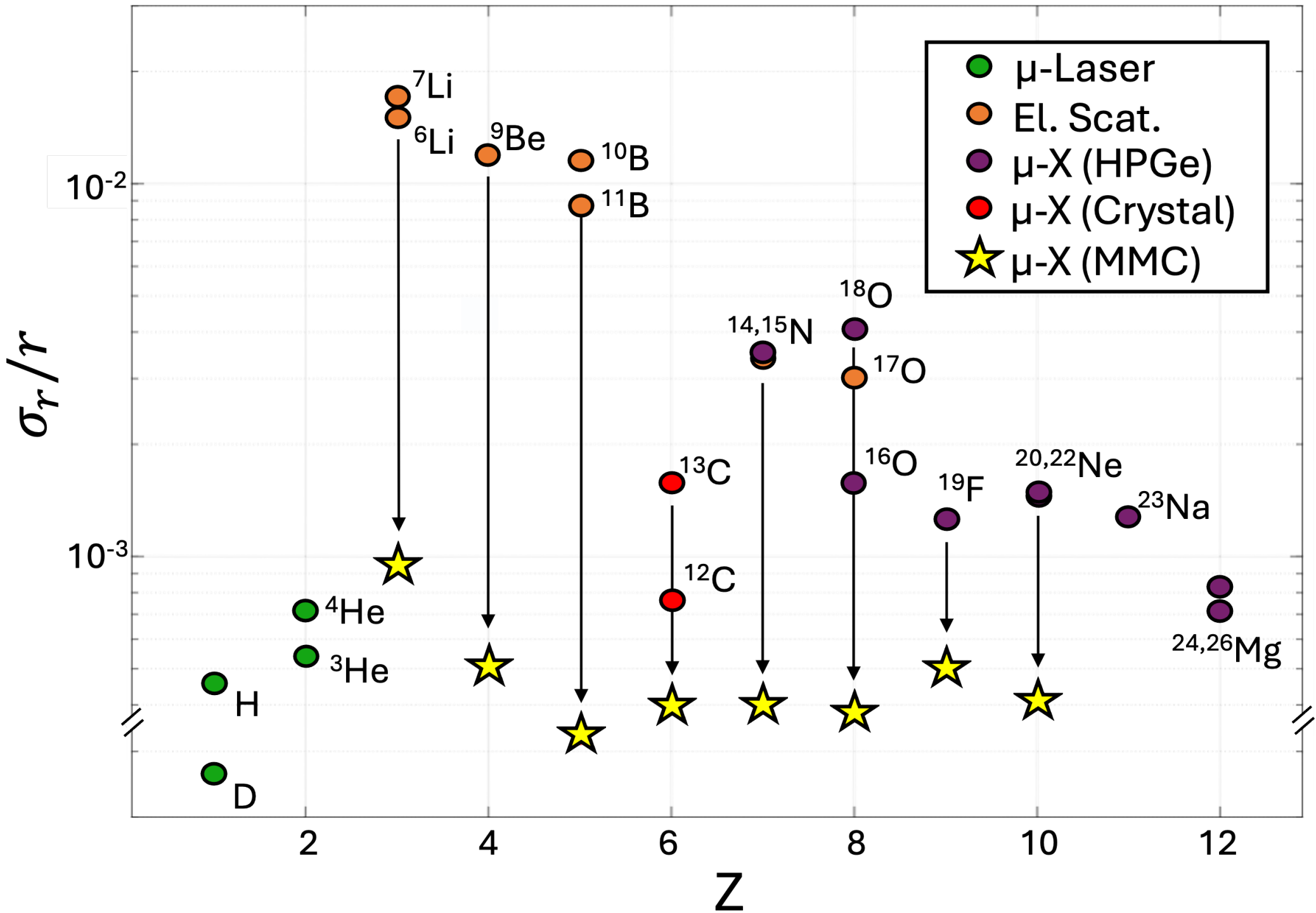
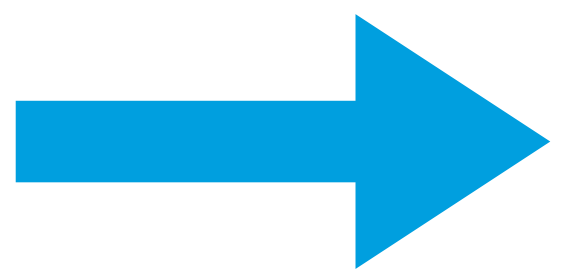
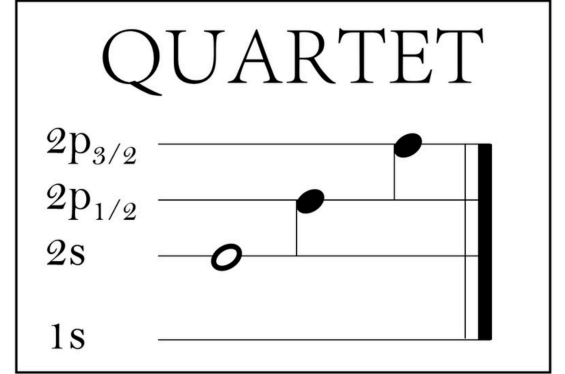
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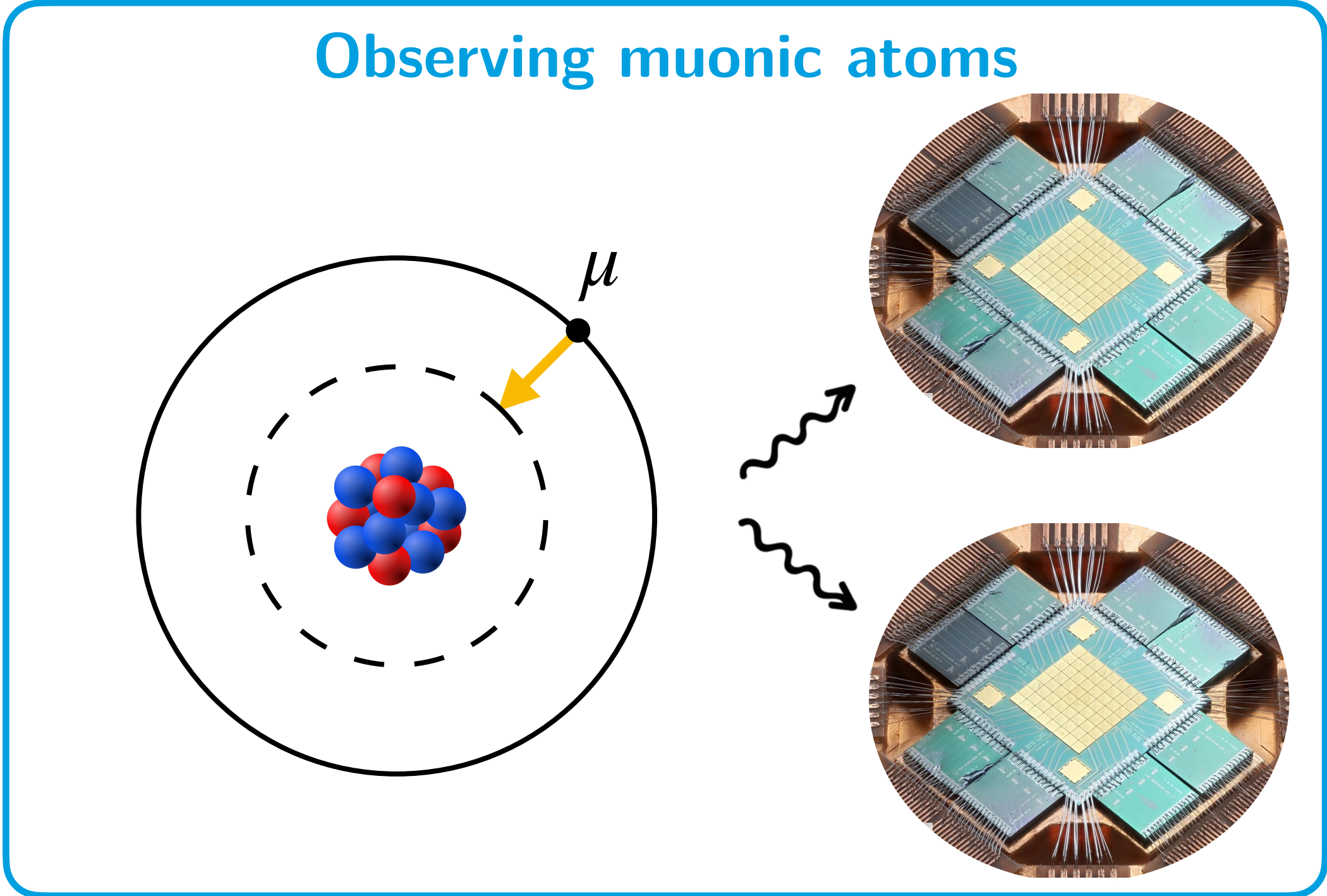


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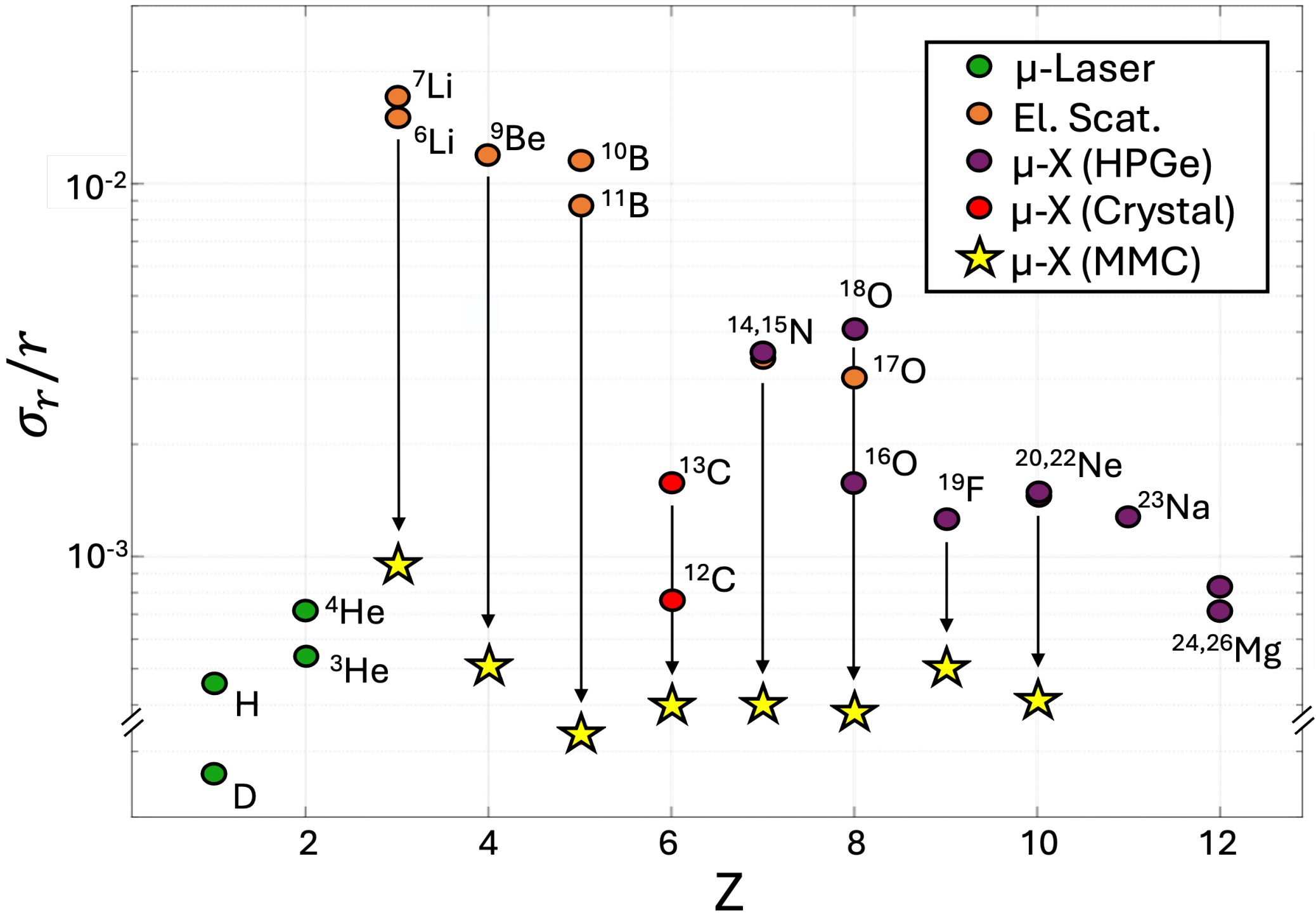
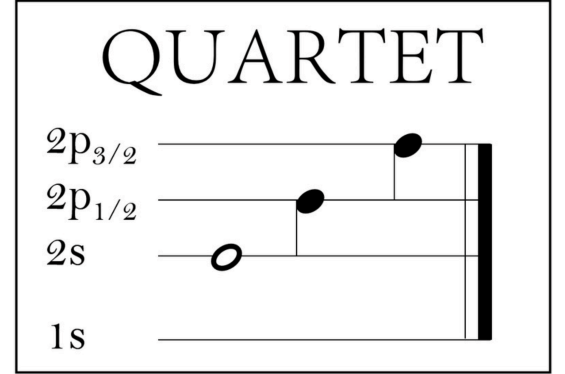
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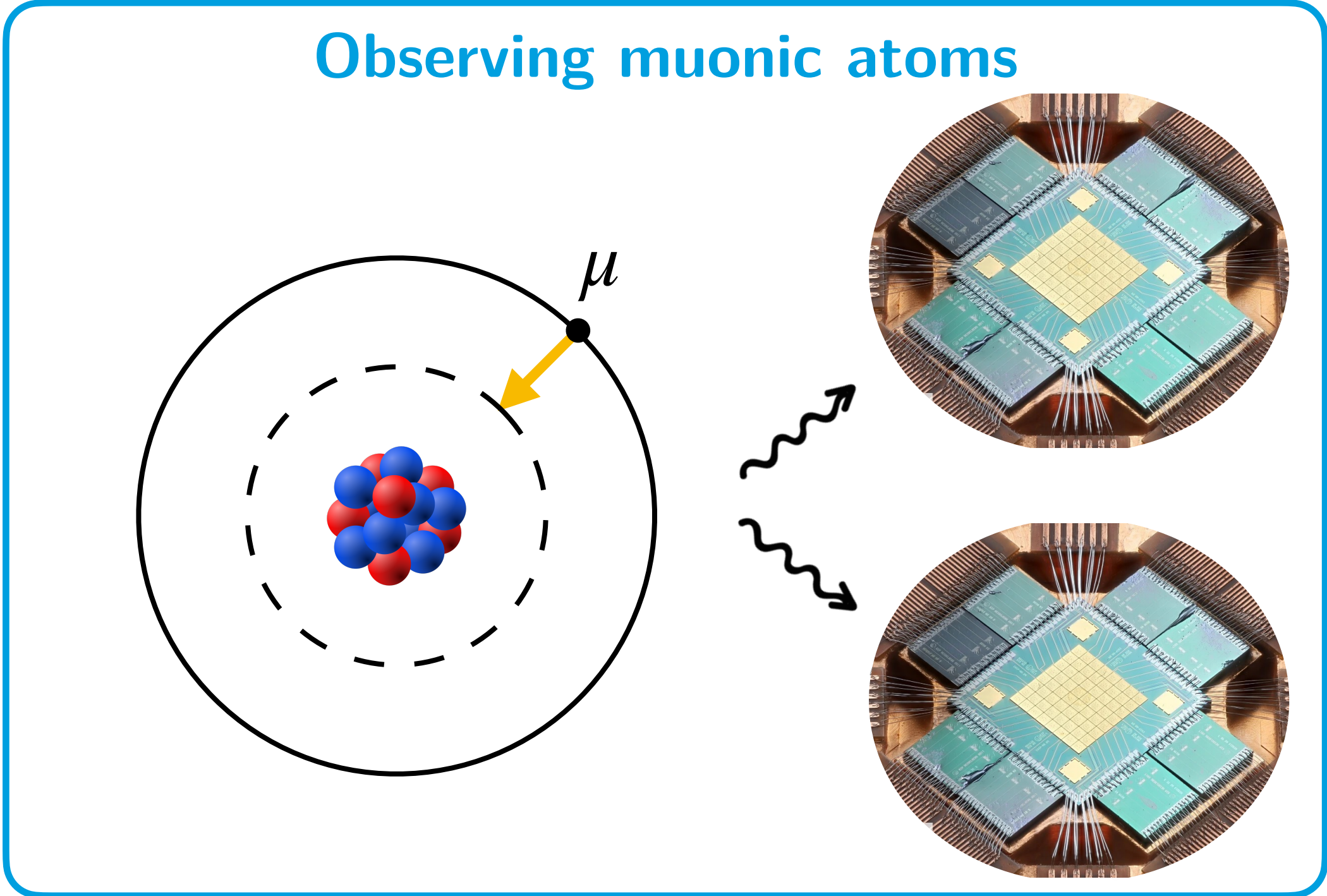


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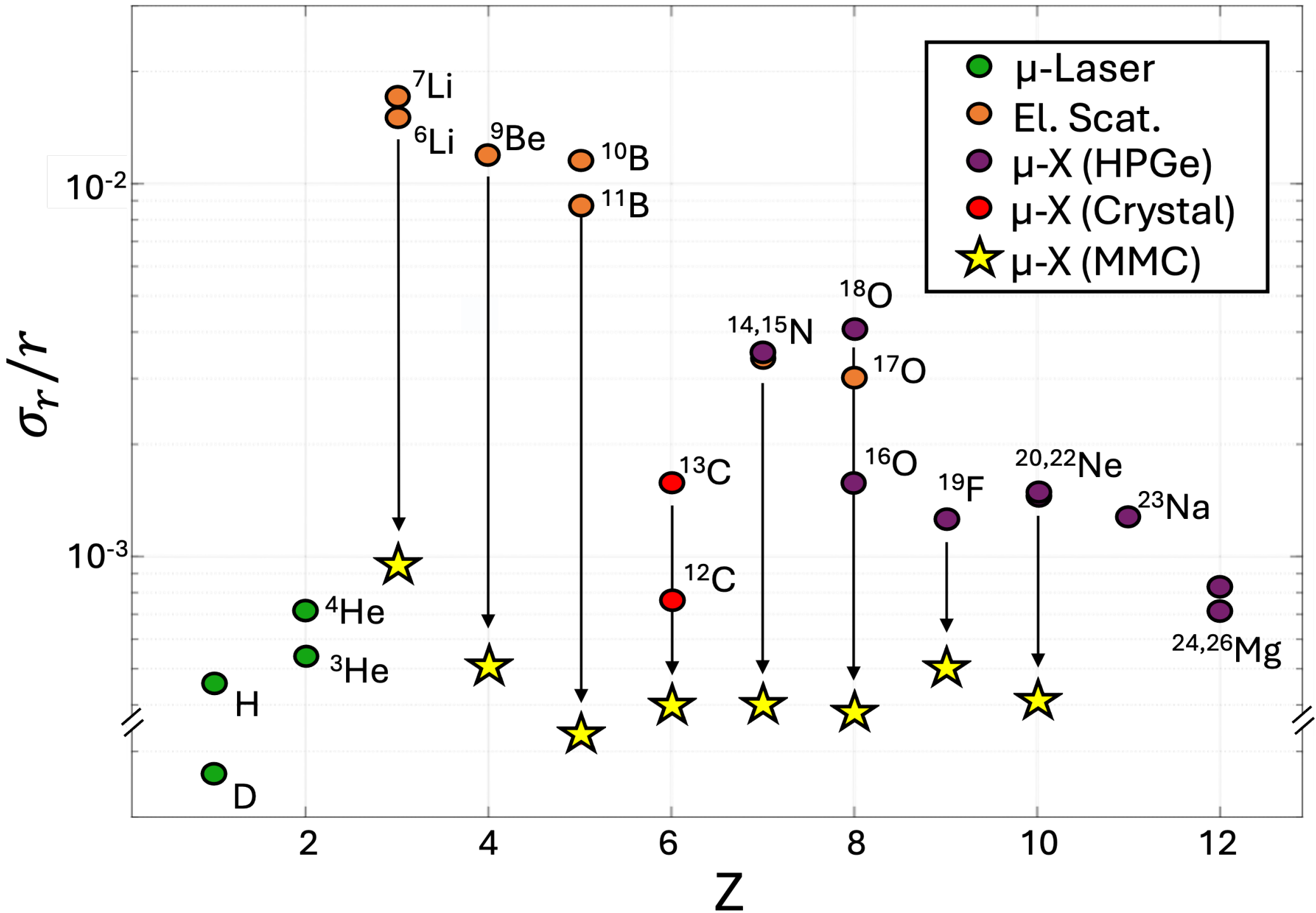
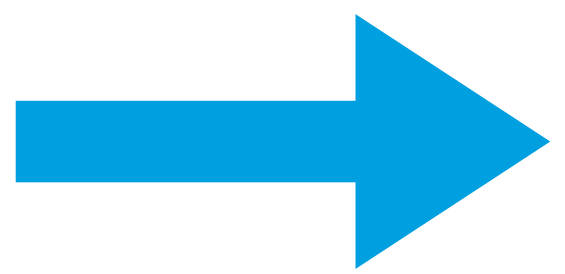
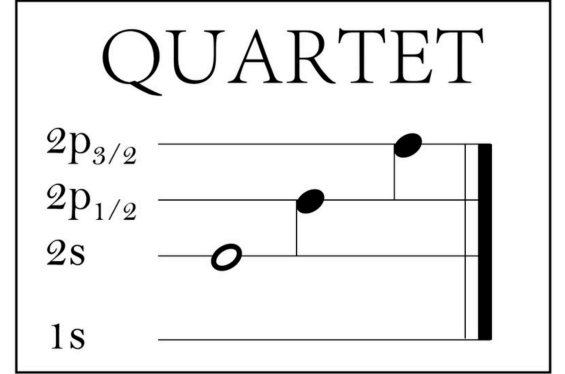
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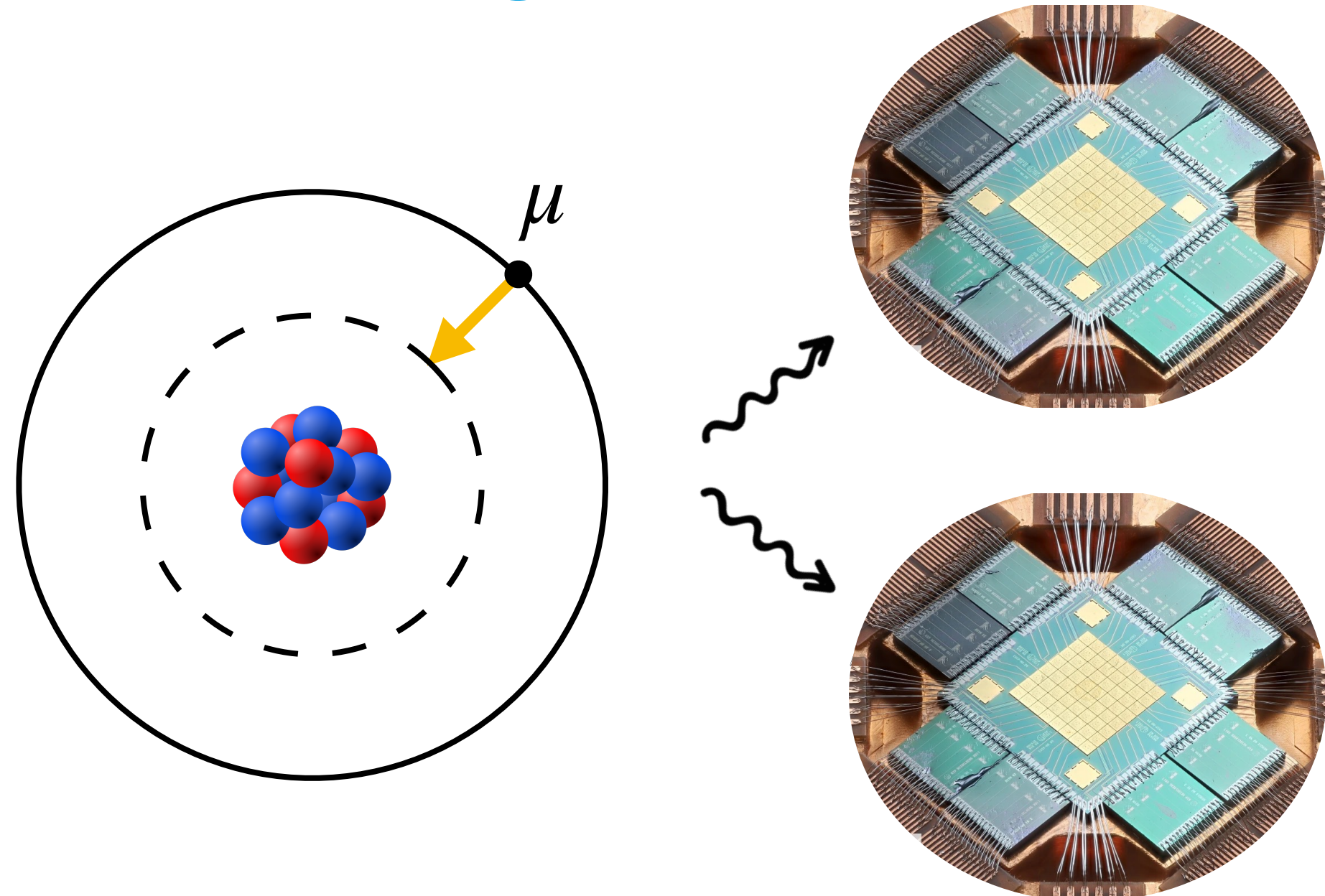
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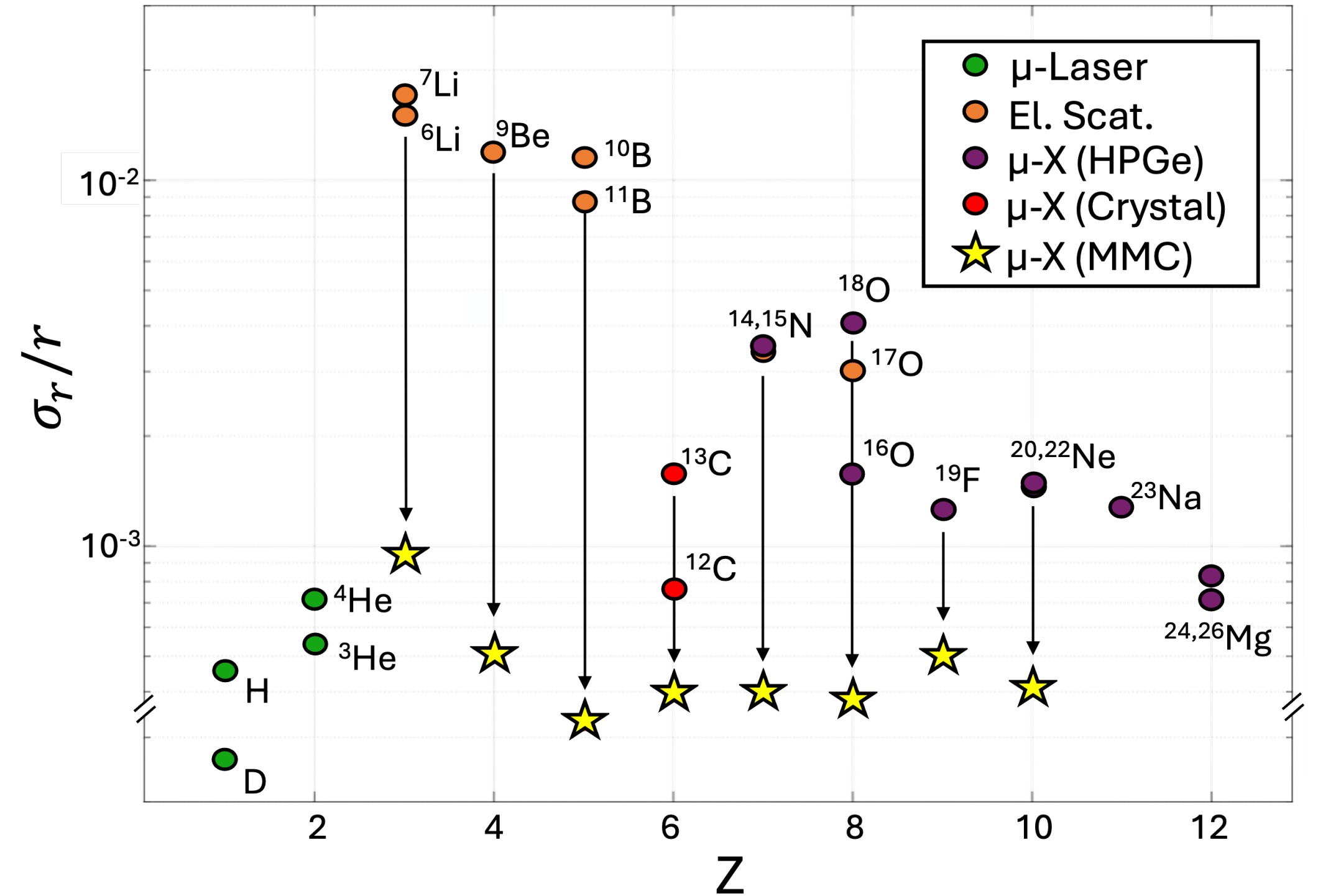
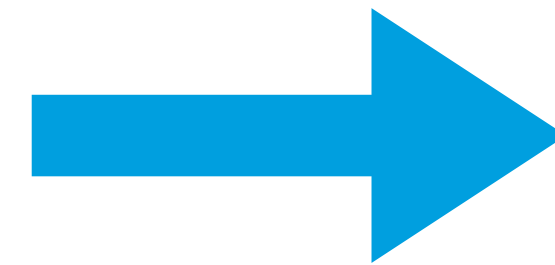
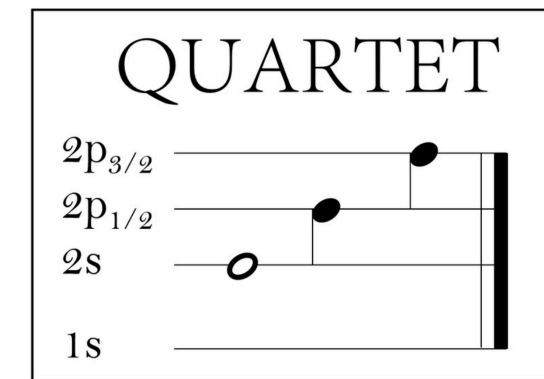
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# Muonic atoms as a precision probe

## Observing muonic atoms



[Unger et al. J. Low Temp. Phys. (2024)]



[QUARTET, Collaboration Meeting on Muonic X-ray (2025)]

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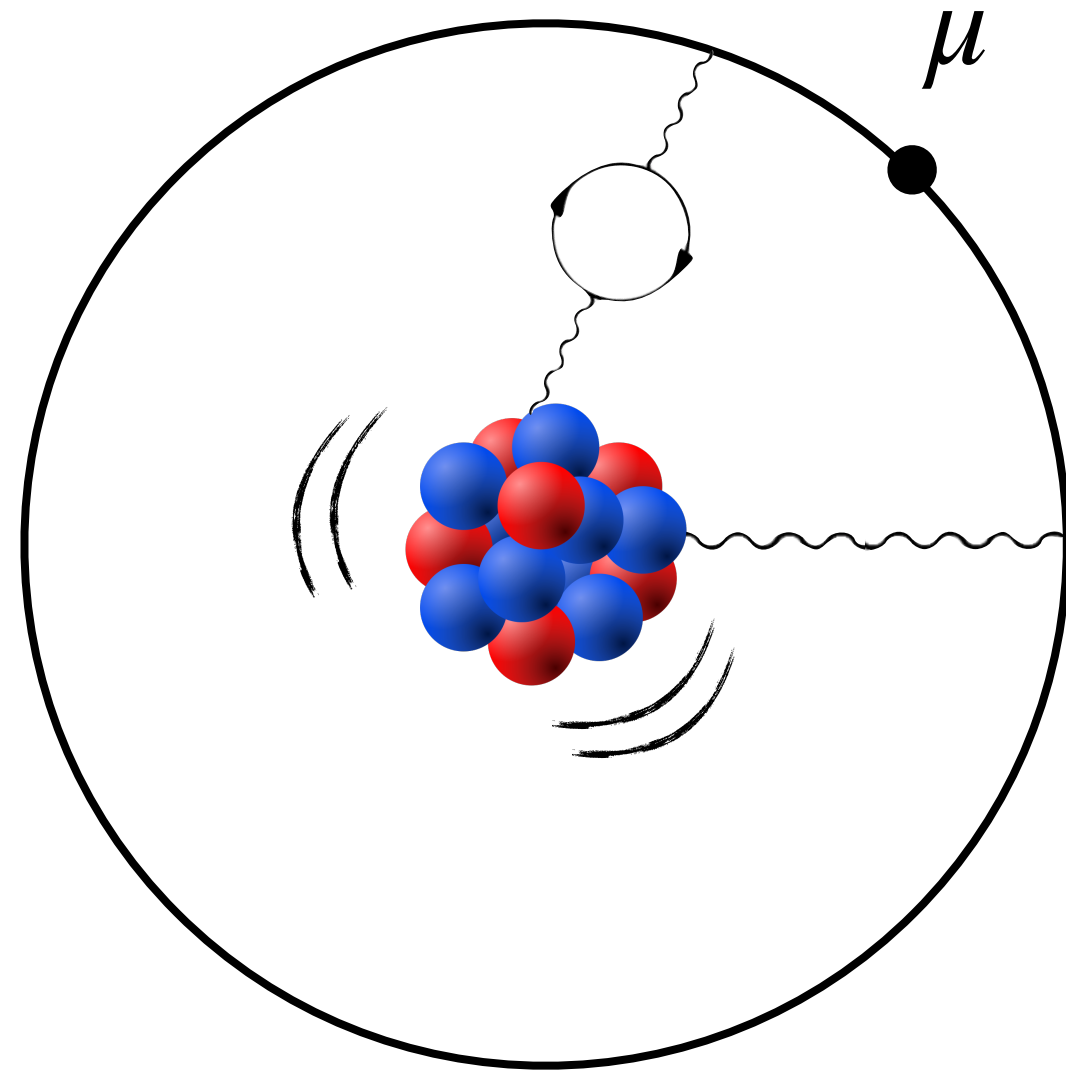
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**Beyond Nuclear:**  $\delta_C$  estimation for  $|V_{ud}|$  extraction

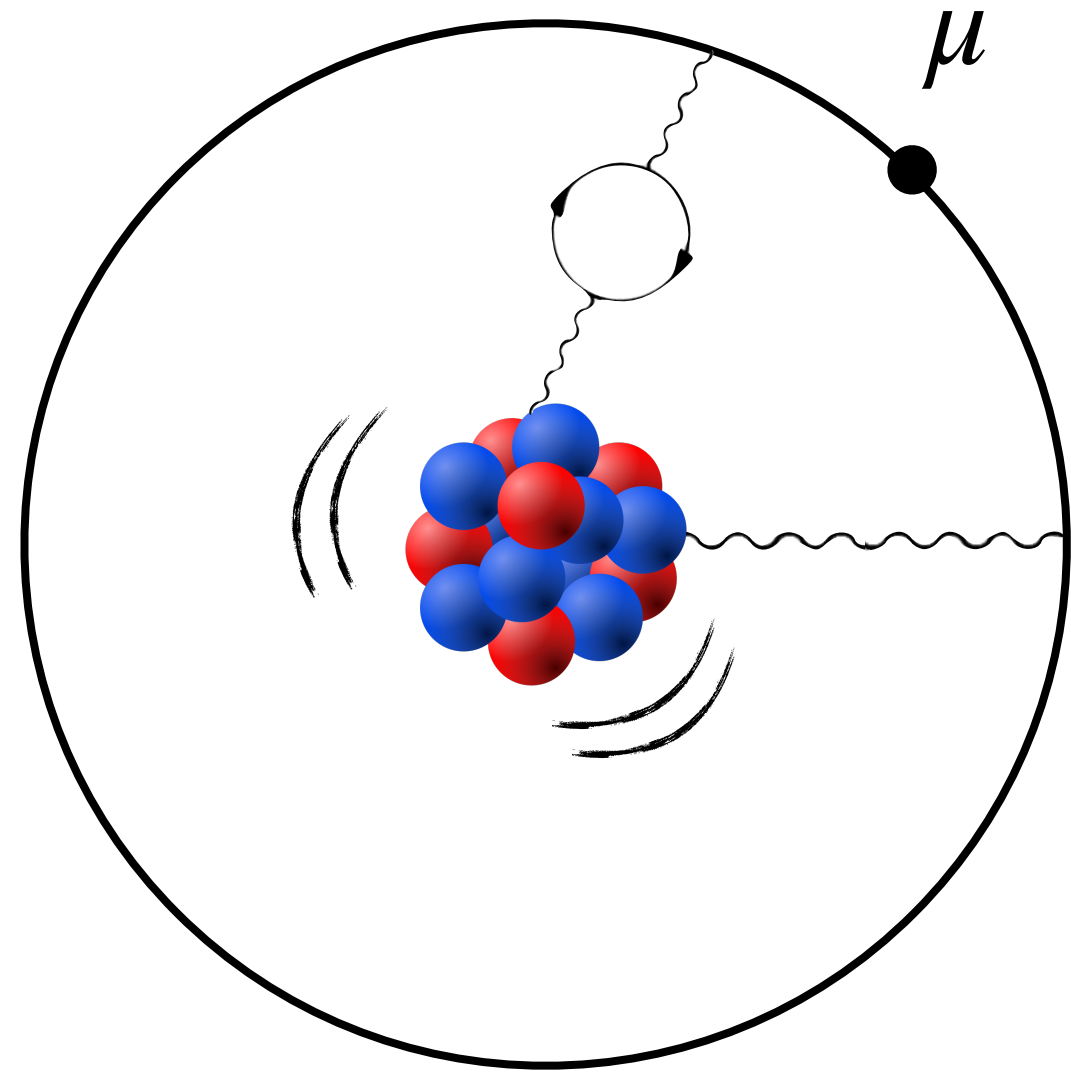
[V. Katyal et al. PRA (2025)] [C.-Y. Seng, M. Gorchtein, PLB (2023)]

[B. Ohayon, Atomic Data and Nuclear Data Tables (2025)]

# Generating effective potentials

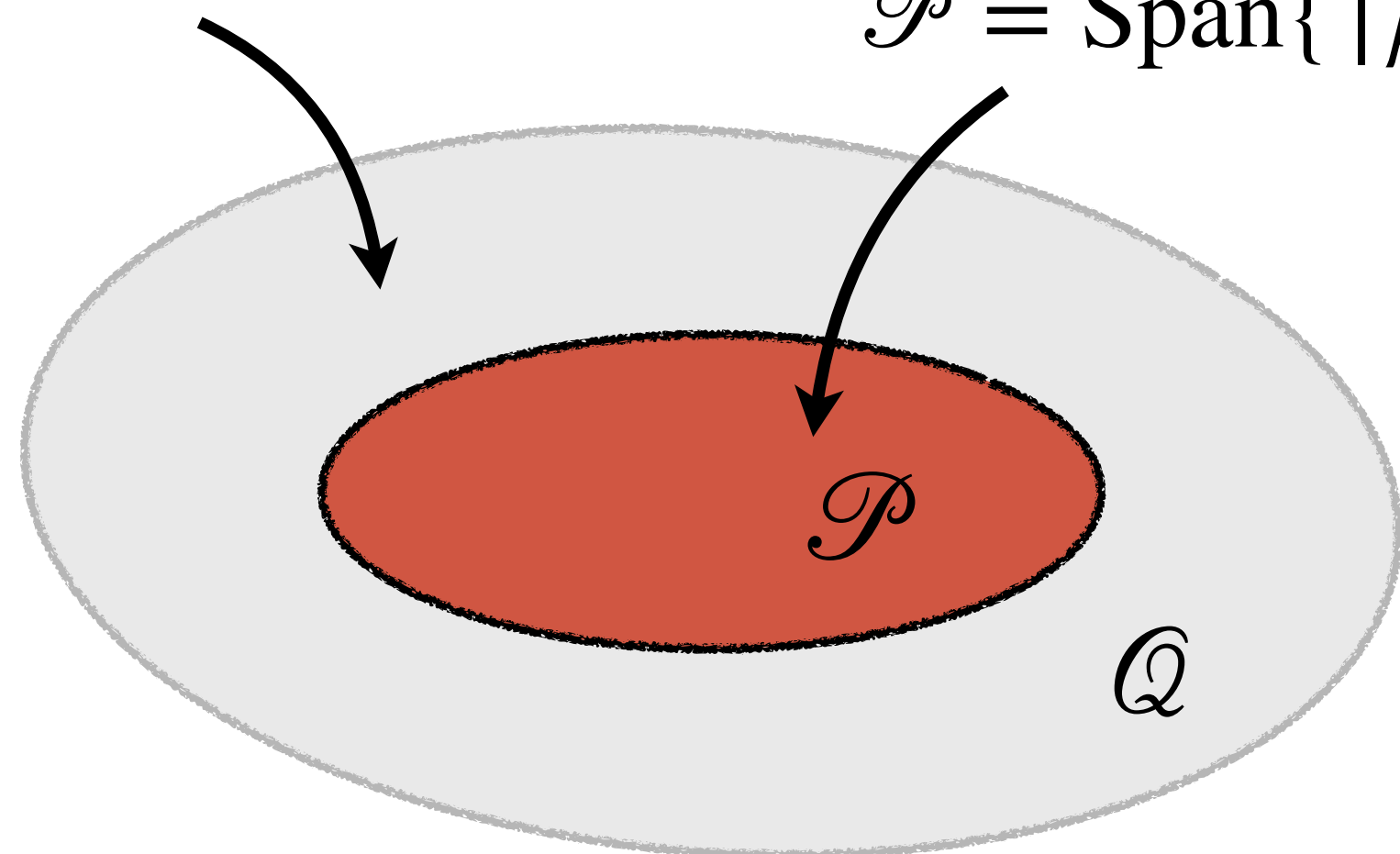


# Generating effective potentials

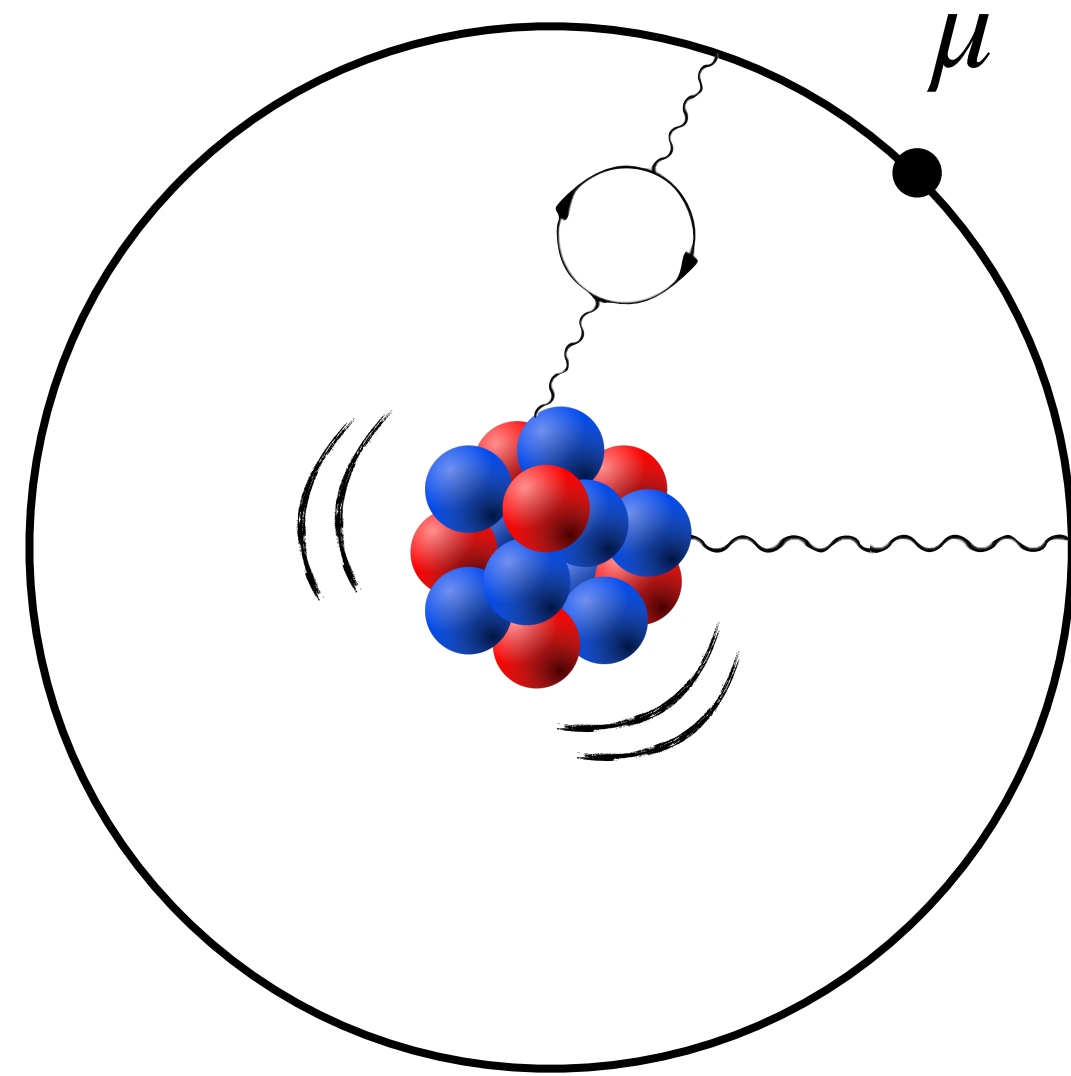


$$\mathcal{P} + \mathcal{Q} = \text{Span}\{ |\mu - {}^7\text{Li}\rangle \}$$

$$\mathcal{P} = \text{Span}\{ |\mu\rangle \} \otimes |\Psi_0\rangle$$



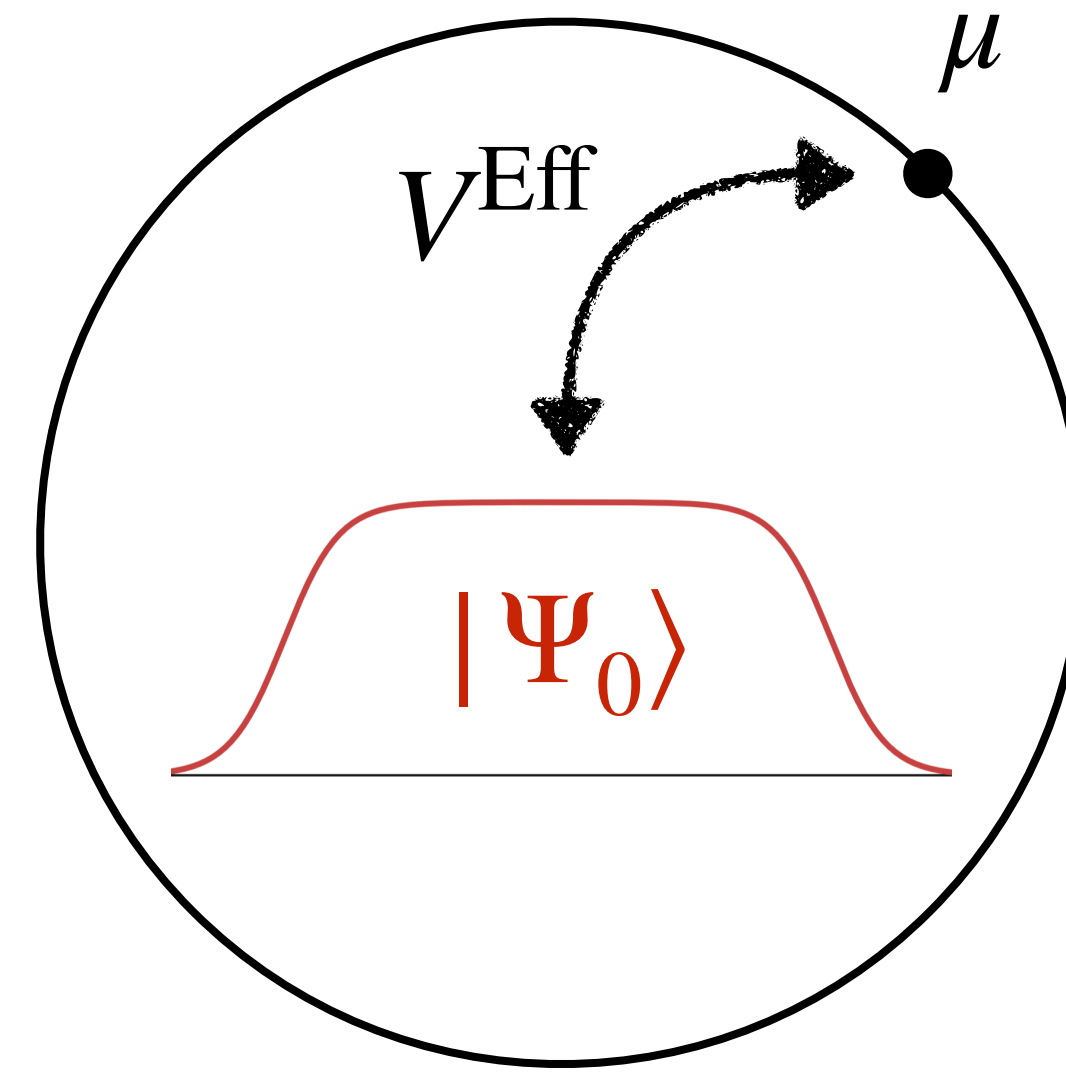
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$$V_{\mu}^{\text{Eff}} = V_{\mu N}^{\mathcal{P}} + \Delta V_{\mu N}$$

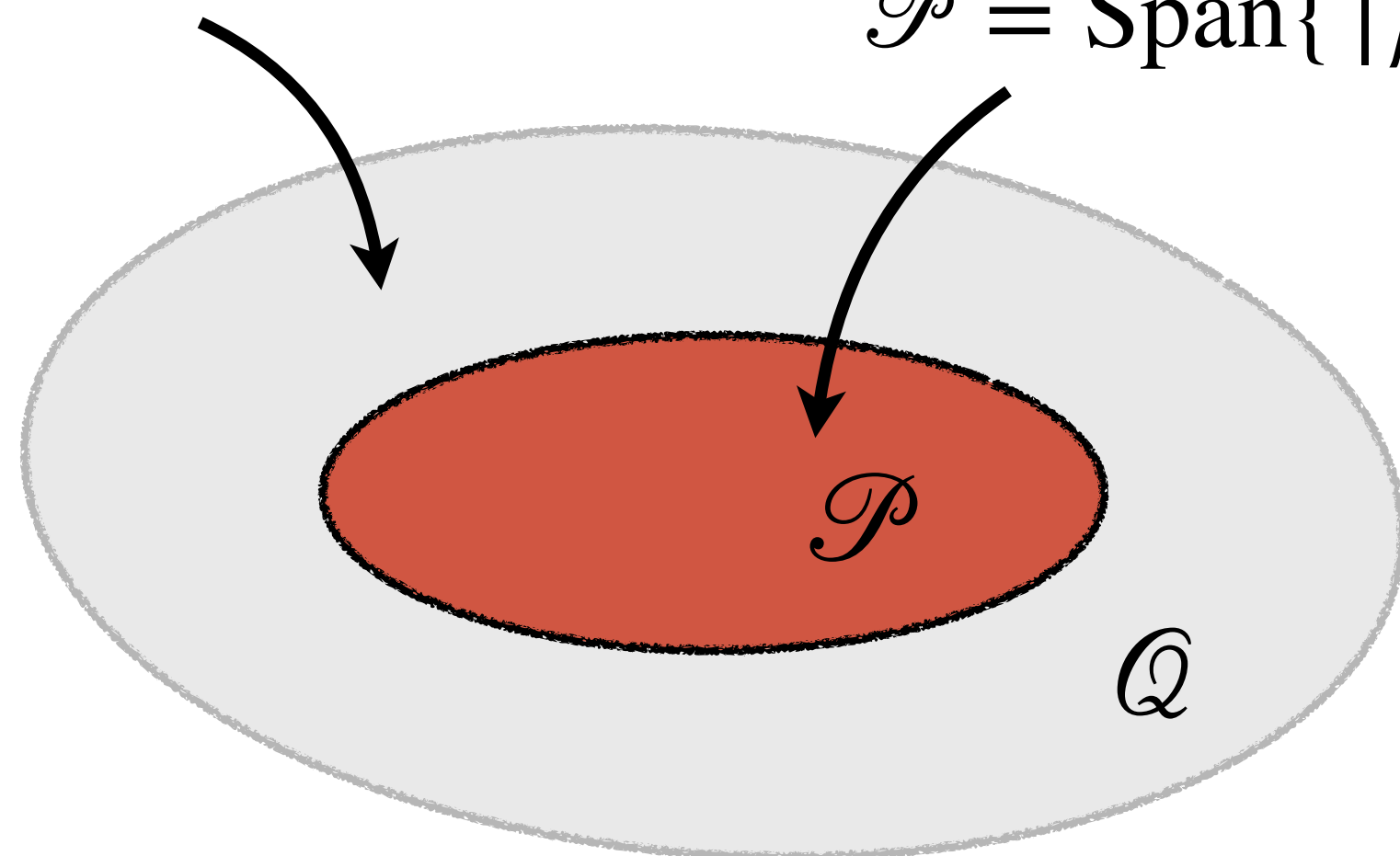


[Borie and Rinker, RMP (1982)]

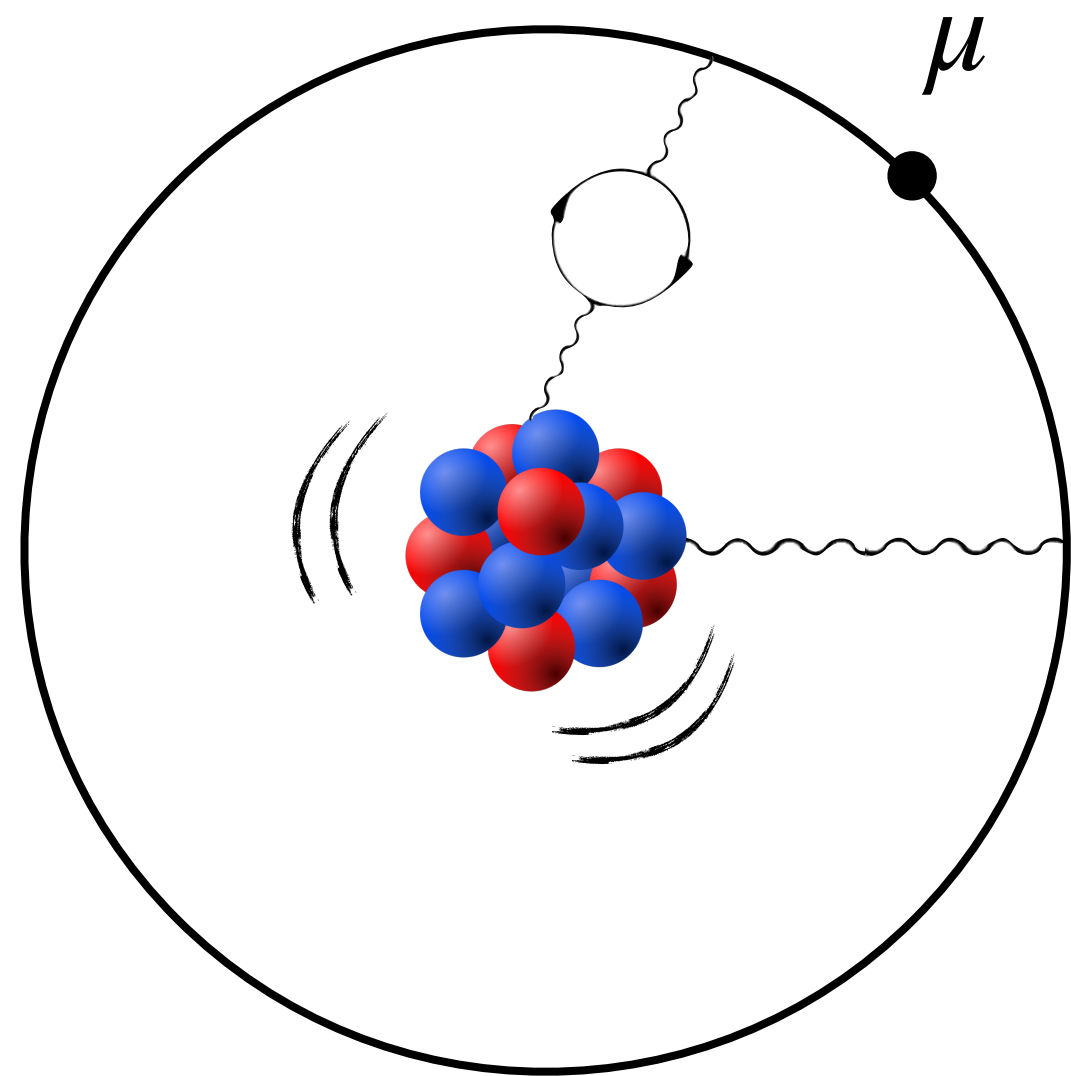


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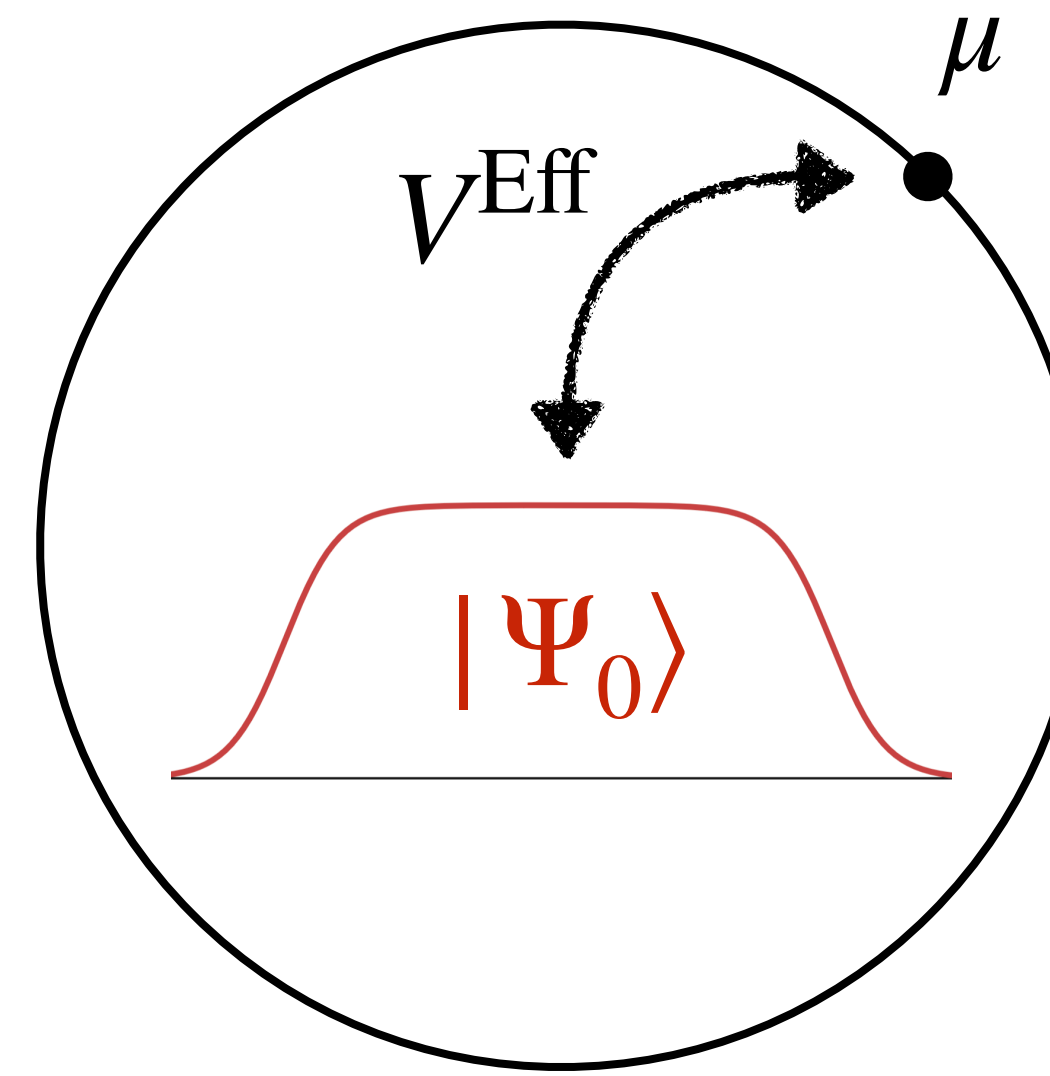
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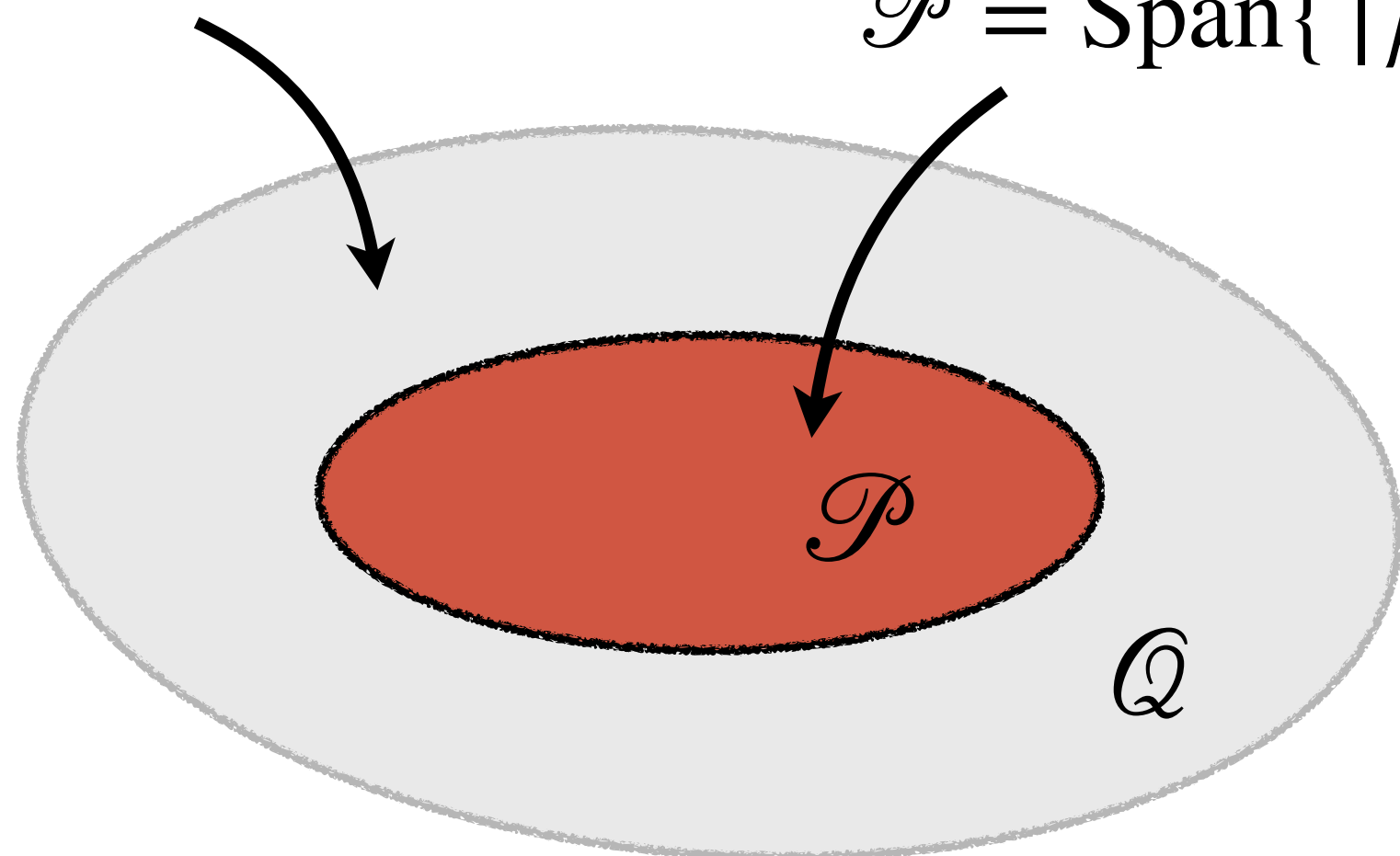


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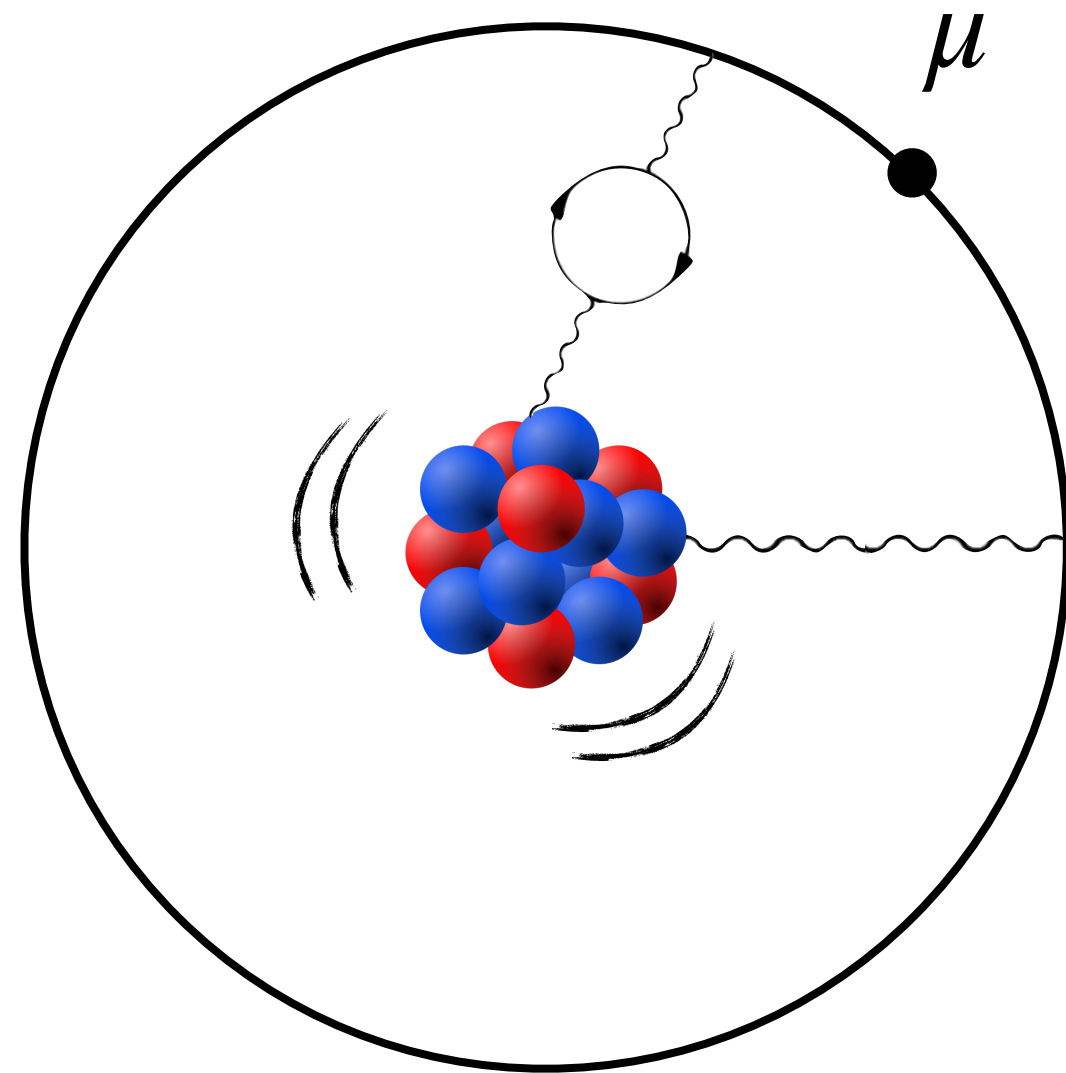


## Effective muonic system

- Perturbation theory at 2nd order:

$$\Delta V_{\mu N} = \sum_q \frac{V_{\mu N}^{\mathcal{P}} |q\rangle \langle q| V_{\mu N}^{\mathcal{P}}}{E_0 - E_q}$$

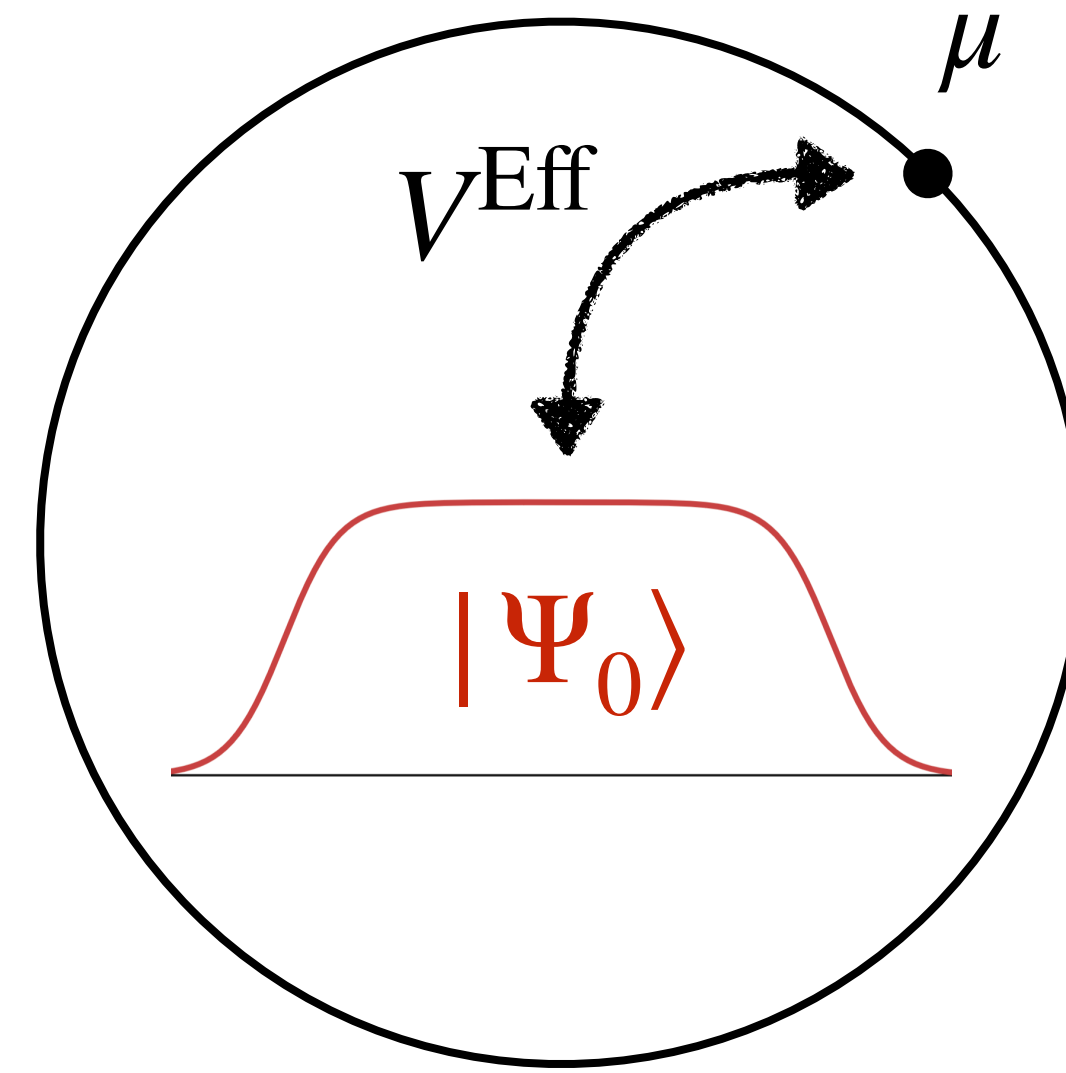
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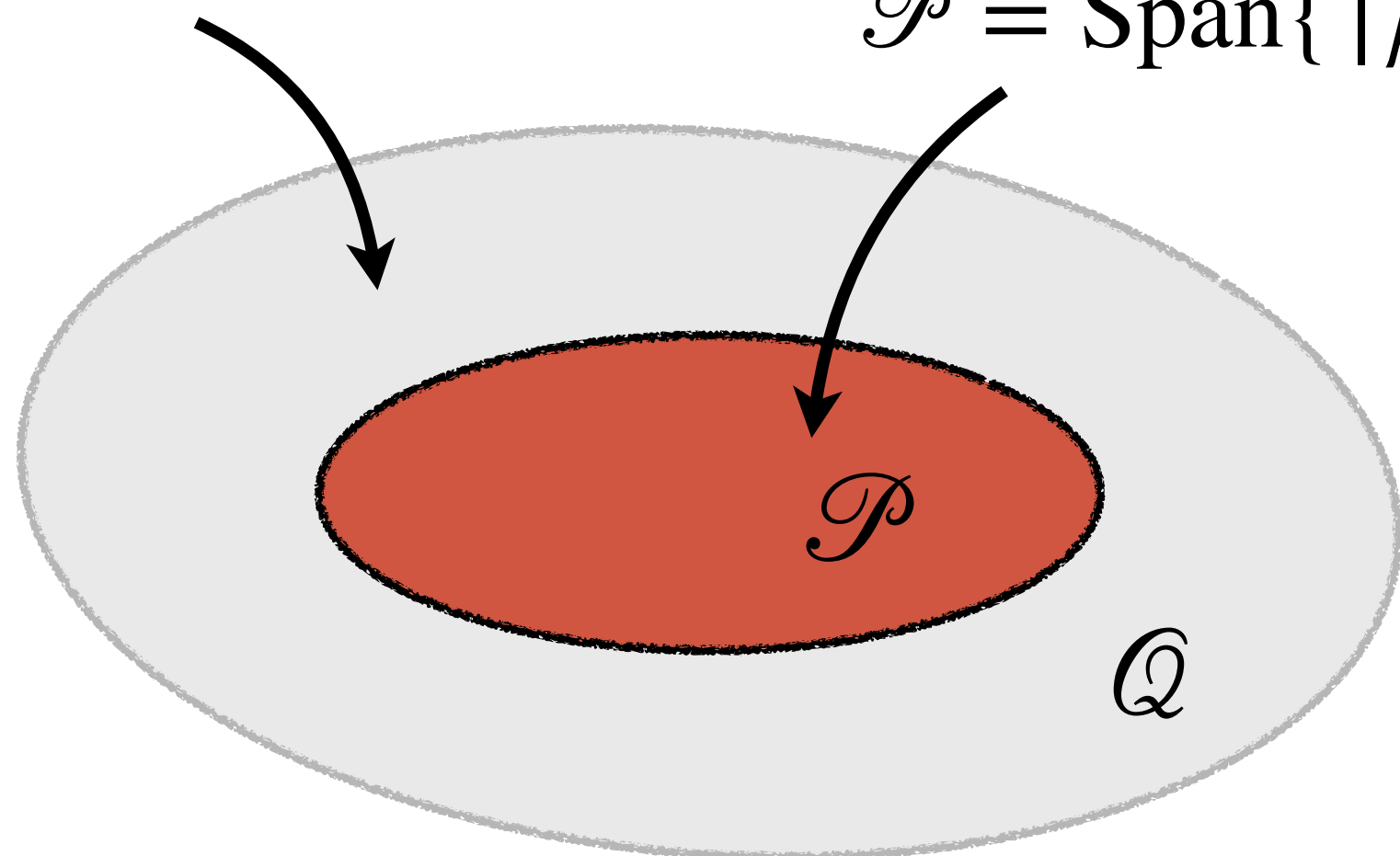


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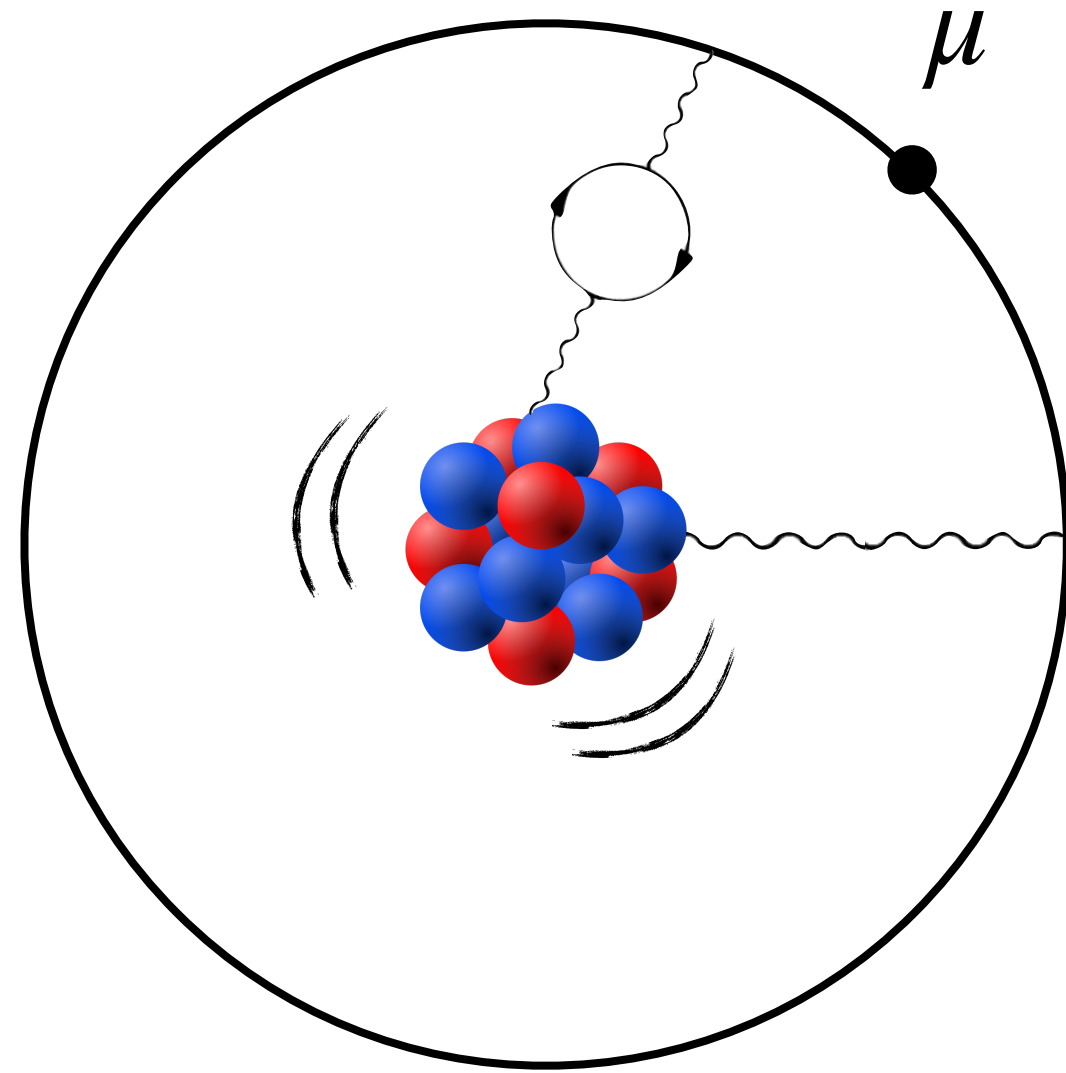
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- (i) Neglect intermediate excited muon
- (ii) "Plane-wave" approx for internal/external muons

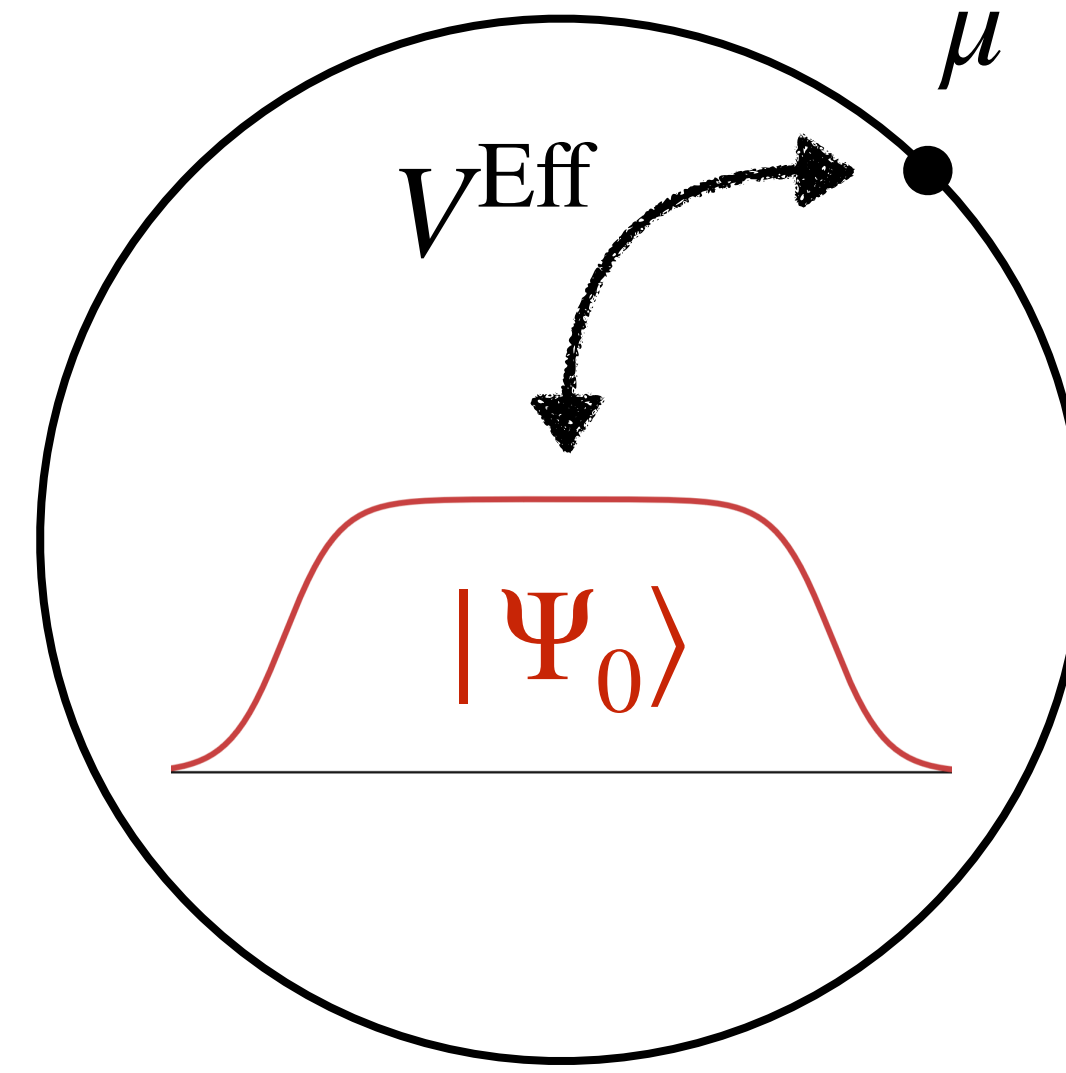
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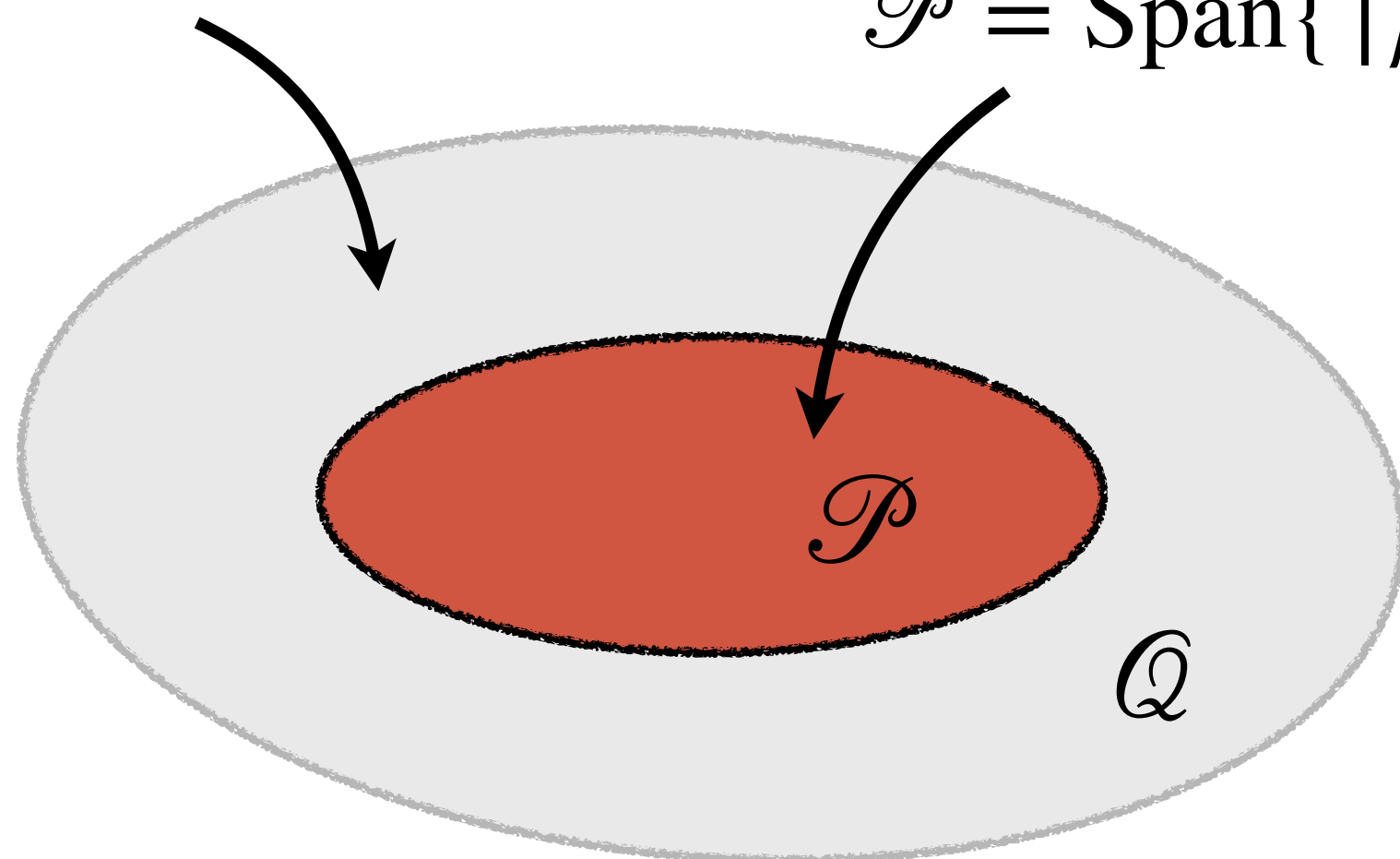


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- Future: control approx with **potential NRQED EFT**

[Peset et al., EPJA (2015)]

# The two-photon exchange nuclear correction

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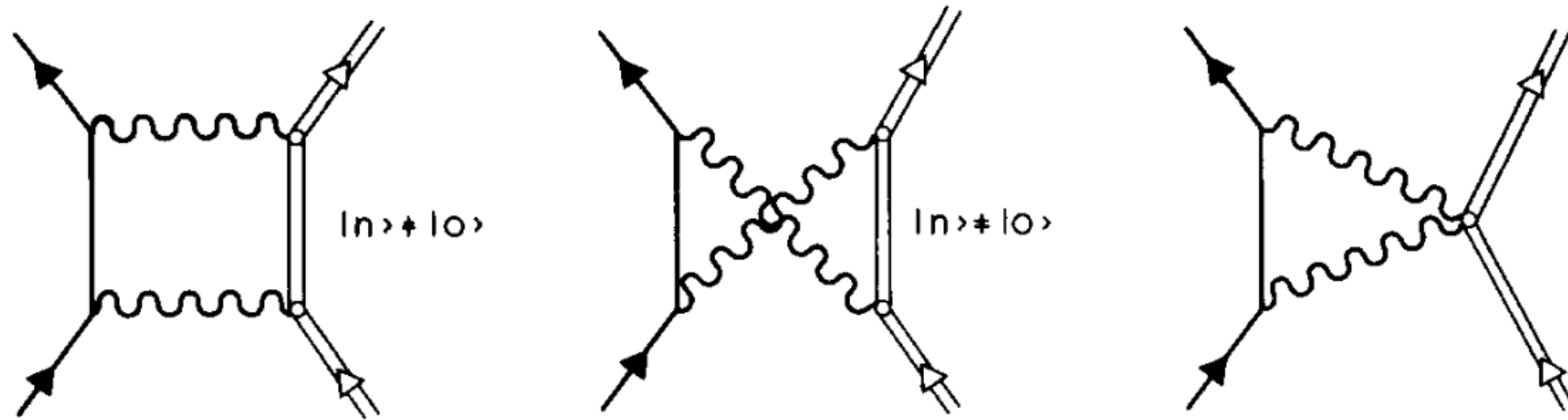
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Two-photon polarizability contributions



$$\Delta E_{nS} = -\frac{(4\pi Z\alpha)}{m_r} |\phi_{nS}(0)|^2 \text{Im} \int \frac{d^4q}{(2\pi)^4} D^{\mu\rho}(q) D^{\nu\tau}(-q) t_{\mu\nu}(q, k) T_{\rho\tau}(q, -q)$$

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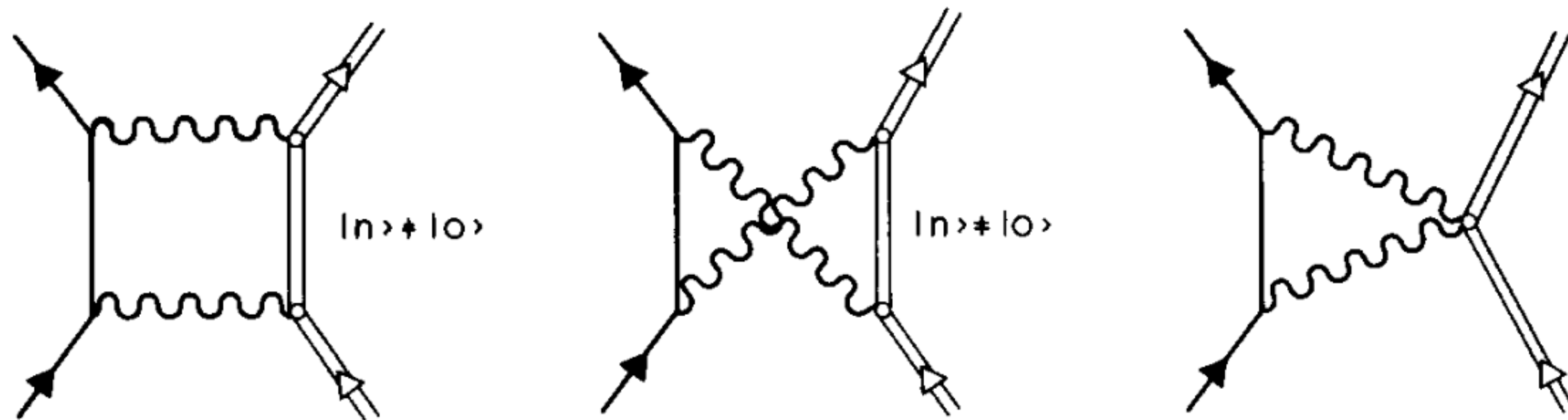
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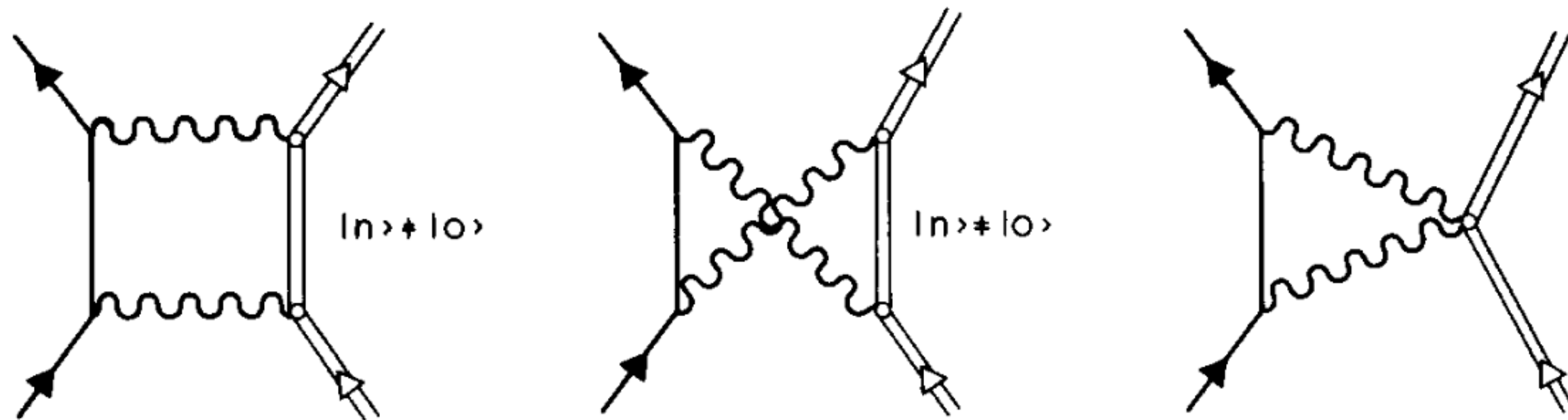
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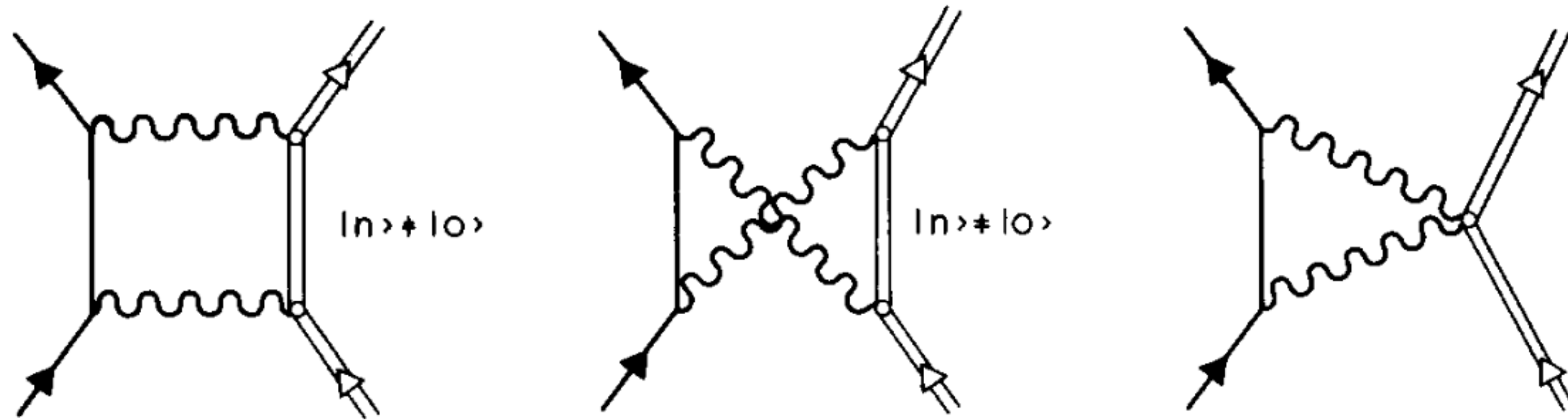
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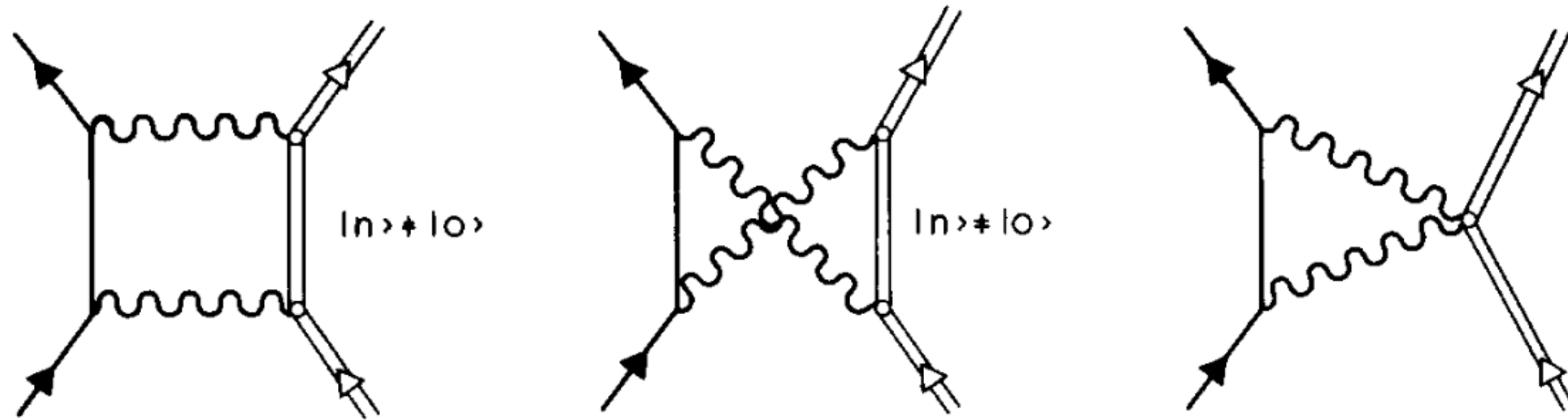
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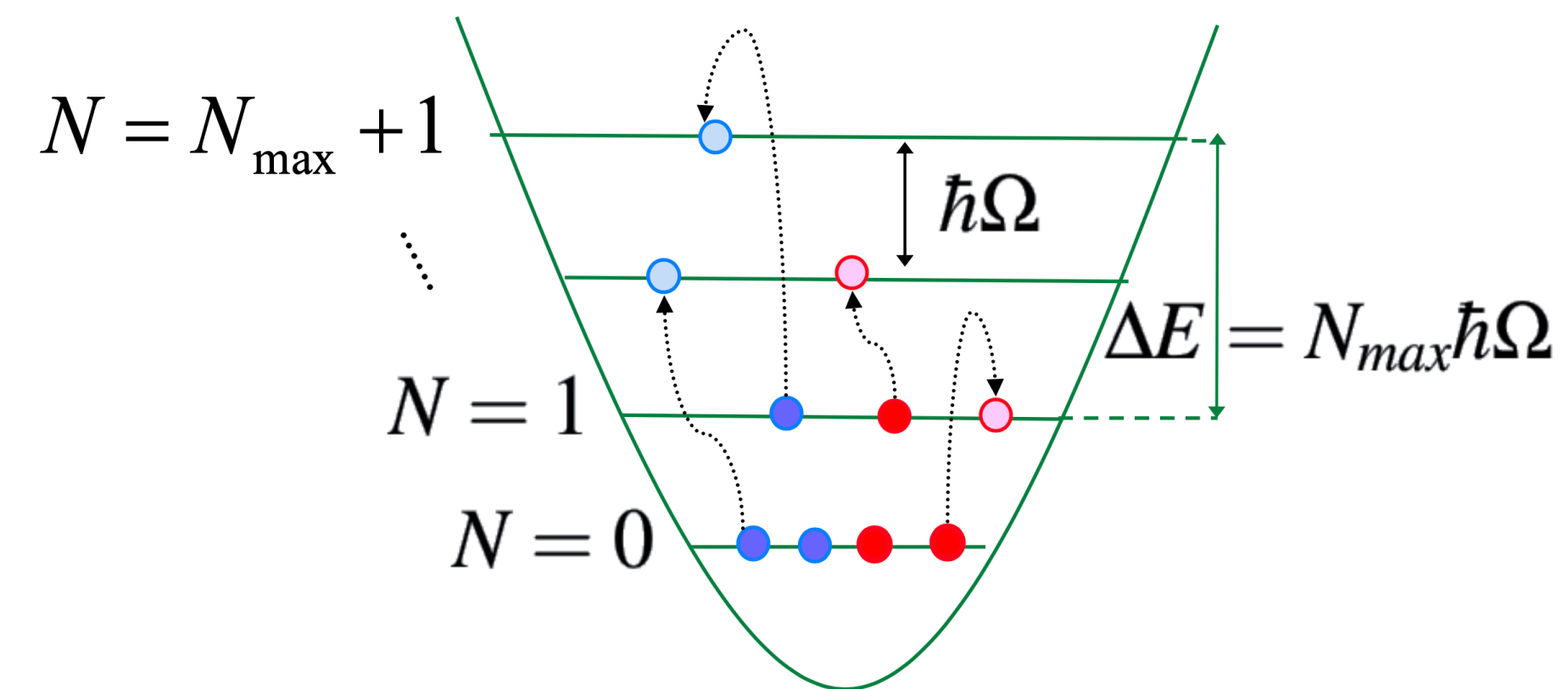
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- Response functions:
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# Ab initio nuclear corrections



# Nuclear physics modelling

## Model used for nuclear currents

### ⊙ Electromagnetic current modelling

- General one-body current for point-like particles
- Nucleon charge form factor  $F(q^2) = (1 + q^2/\Lambda^2)^{-2}$
- Non-relativistic reduction:
  - ➔  $\vec{j} \equiv$  charge convection + magnetic moment rotational

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### Multipole decomposition of nuclear currents

[Donnelly, Haxton, Atomic and Nuclear Data Tables (1979)]

- $M_{JM_j;TM_T}(q) \equiv \int d^3x \mathbf{M}_J^{M_j}(qx) J_0(x)_{TM_T}$
- $T_{JM_j;TM_T}^E(q) \equiv \int d^3x \left[ \frac{1}{q} \nabla \times \bar{\mathbf{M}}_{JJ}^{M_j}(qx) \right] \cdot \vec{J}(x)_{TM_T}$
- $T_{JM_j;TM_T}^M(q) \equiv \int d^3x \bar{\mathbf{M}}_{JJ}^{M_j}(qx) \cdot \vec{J}(x)_{TM_T}$

➔ Truncation at  $J = 3$

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- Non-relativistic reduction:
  - ➔  $\vec{j} \equiv$  charge convection + magnetic moment rotational

### Multipole decomposition of nuclear currents

[Donnelly, Haxton, Atomic and Nuclear Data Tables (1979)]

- $M_{JM_j;TM_T}(q) \equiv \int d^3x \mathbf{M}_J^{M_j}(qx) J_0(x)_{TM_T}$
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➔ Truncation at  $J = 3$

## Model used for nuclear many-body state

### Ab initio nuclear interaction [Entem et al. (2017)] [Somà et al. (2020)]

- Two  $\chi$ EFT interactions considered
- N4LO-E7 and N3LO

➔ Estimate interaction uncertainty

# Nuclear physics modelling

7

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### Many-body approximation

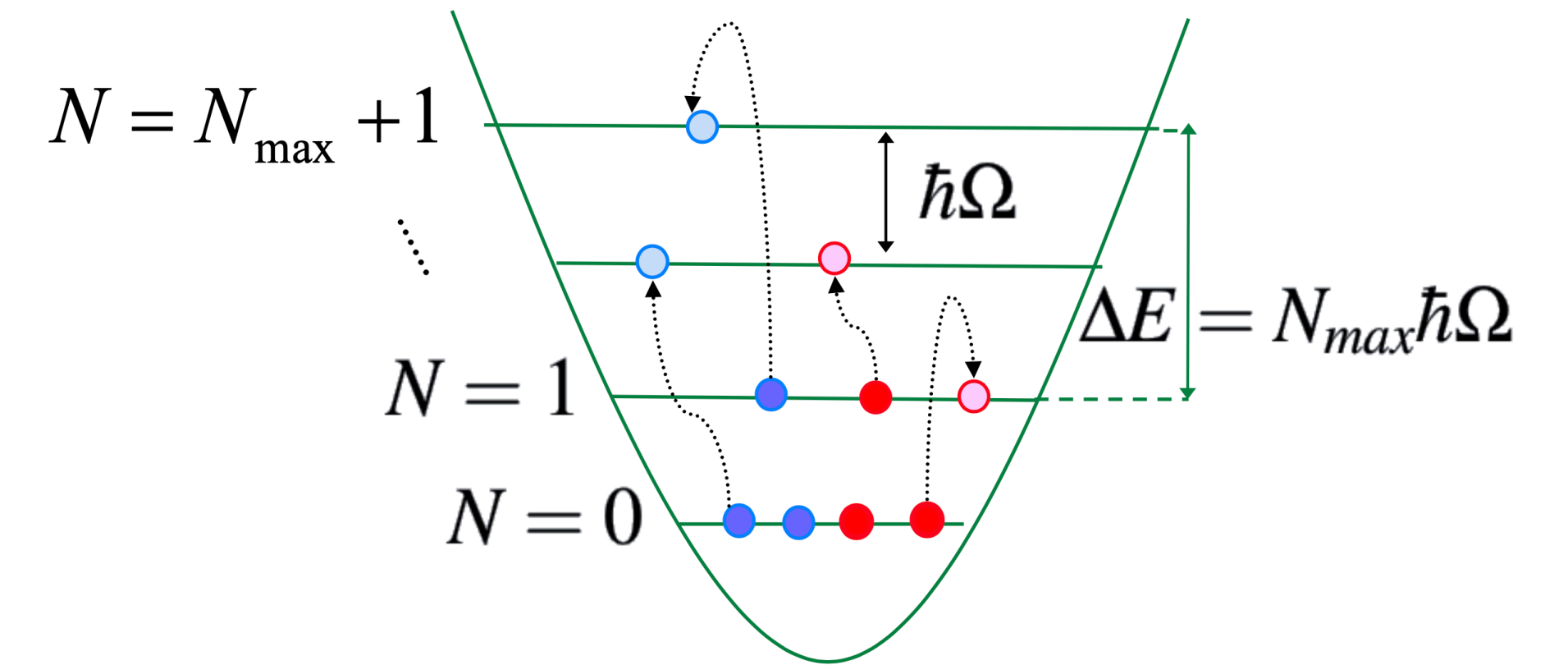
- No-Core Shell Model
- More details in next section

➔ Negligible many-body approximation uncertainty

# Ab initio No-Core Shell Model

Anti-symmetrized products of  
many-body HO states

8



# Ab initio No-Core Shell Model

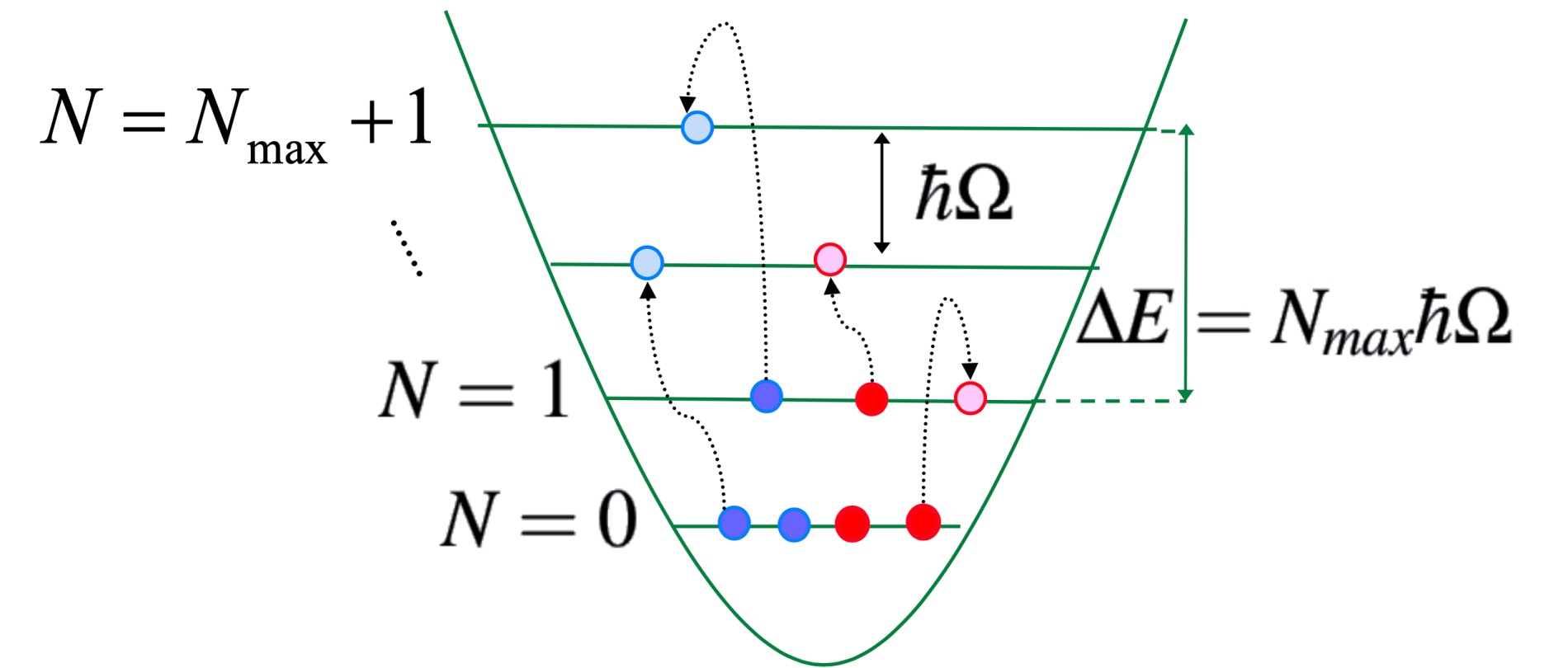
## Lanczos tridiagonalization algorithm [Lanczos (1950)]

- Initialization: normalized pivot  $|\phi_1\rangle$
- Recursion:  $\alpha_i$ ,  $\beta_i$  and  $|\phi_i\rangle$ 
  - $\beta_{i+1}|\phi_{i+1}\rangle = H|\phi_i\rangle - \alpha_i|\phi_i\rangle - \beta_i|\phi_{i-1}\rangle$
  - $\alpha_i = \langle\phi_i|H|\phi_i\rangle$  and  $\beta_{i+1}$  st  $\langle\phi_{i+1}|\phi_{i+1}\rangle = 1$
- Output:
  - Lanczos basis and coefficients  $\{|\phi_i\rangle, \alpha_i, \beta_i\}$

$$\begin{pmatrix} \alpha_1 & \beta_2 & & & & \\ \beta_2 & \alpha_2 & \beta_3 & & & \\ & \beta_3 & \alpha_3 & \ddots & & \\ & & \ddots & \ddots & \beta_{k-1} & \\ & & & \beta_{k-1} & \alpha_{k-1} & \beta_k \\ & & & & \beta_k & \alpha_k \end{pmatrix}$$

→  $H$  in Lanczos basis

## Anti-symmetrized products of many-body HO states



# Ab initio No-Core Shell Model

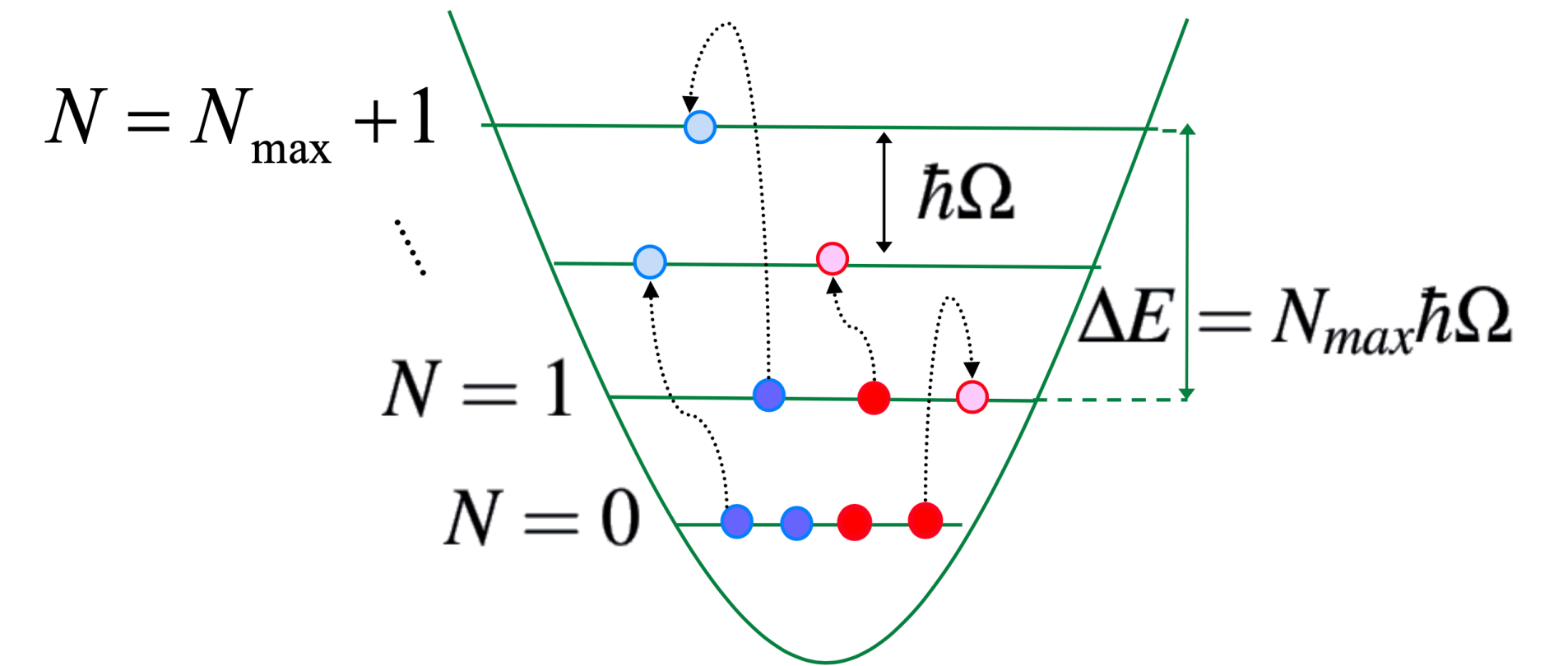
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## Anti-symmetrized products of many-body HO states



## Application to nuclear structure

- Efficient calculation of low-lying spectra
  - Selection rules  $\Rightarrow$  **Fast matrix-vector multiplication**
  - In practice:  $N_L \sim 100 - 200$  is sufficient
- Application to  ${}^{6-7}\text{Li}$ 
  - $N_L = 200$  for  $N_{\max} = 1$  to 9
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# Ab initio No-Core Shell Model

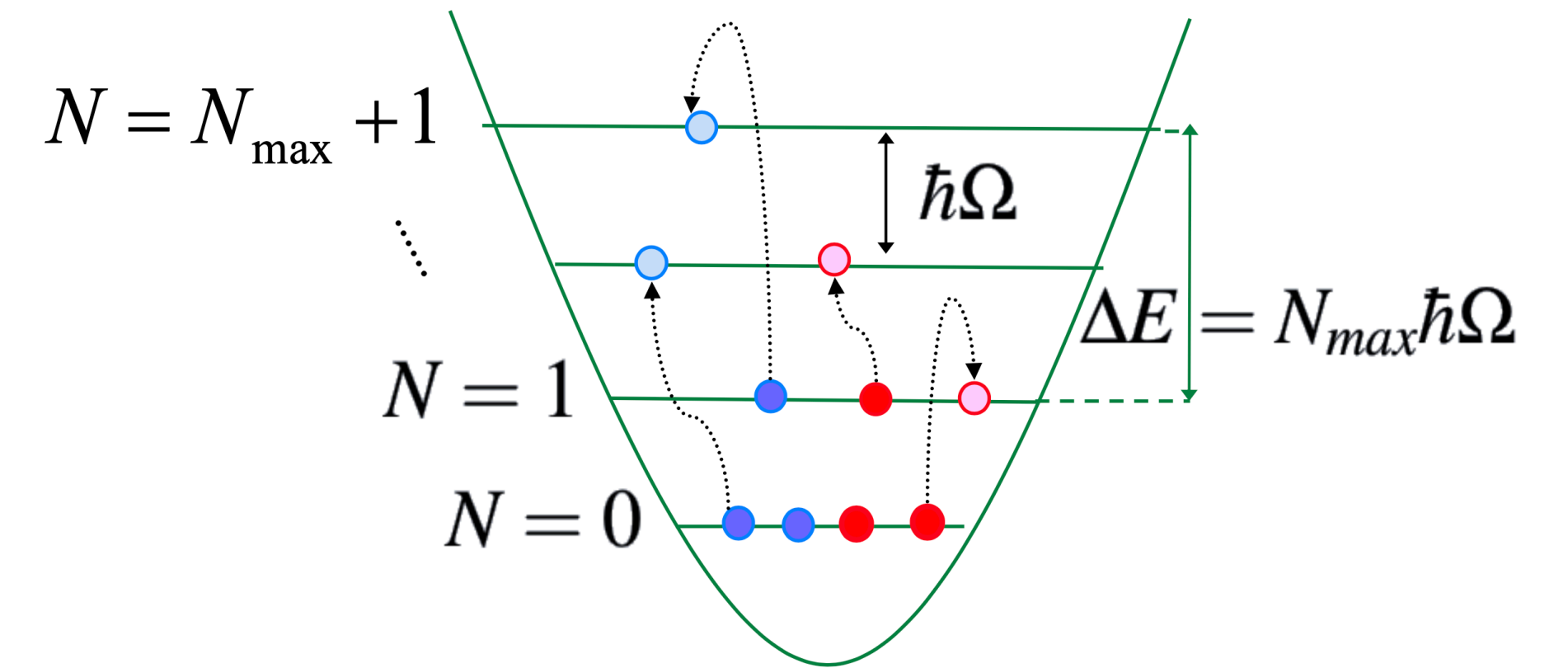
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## Lanczos-strength algorithm for $\delta_{\text{pol}}^A$

- Convergence problem
    - Often the **strength of  $S_O(\omega)$  is fragmented**
    - And only low-lying states converged
  - Lanczos strength algorithm
    - Pivot based on 1<sup>st</sup> Lanczos output:  $|\phi'_1\rangle = O|\Psi\rangle$
    - Recovers **exactly**  $\int d\omega \omega^n S_O(\omega)$  for  $n \leq 2N_L - 1$
- **One additional NCSM run per operator**

# Numerical results for ${}^6\text{-}{}^7\text{Li}$ isotopes

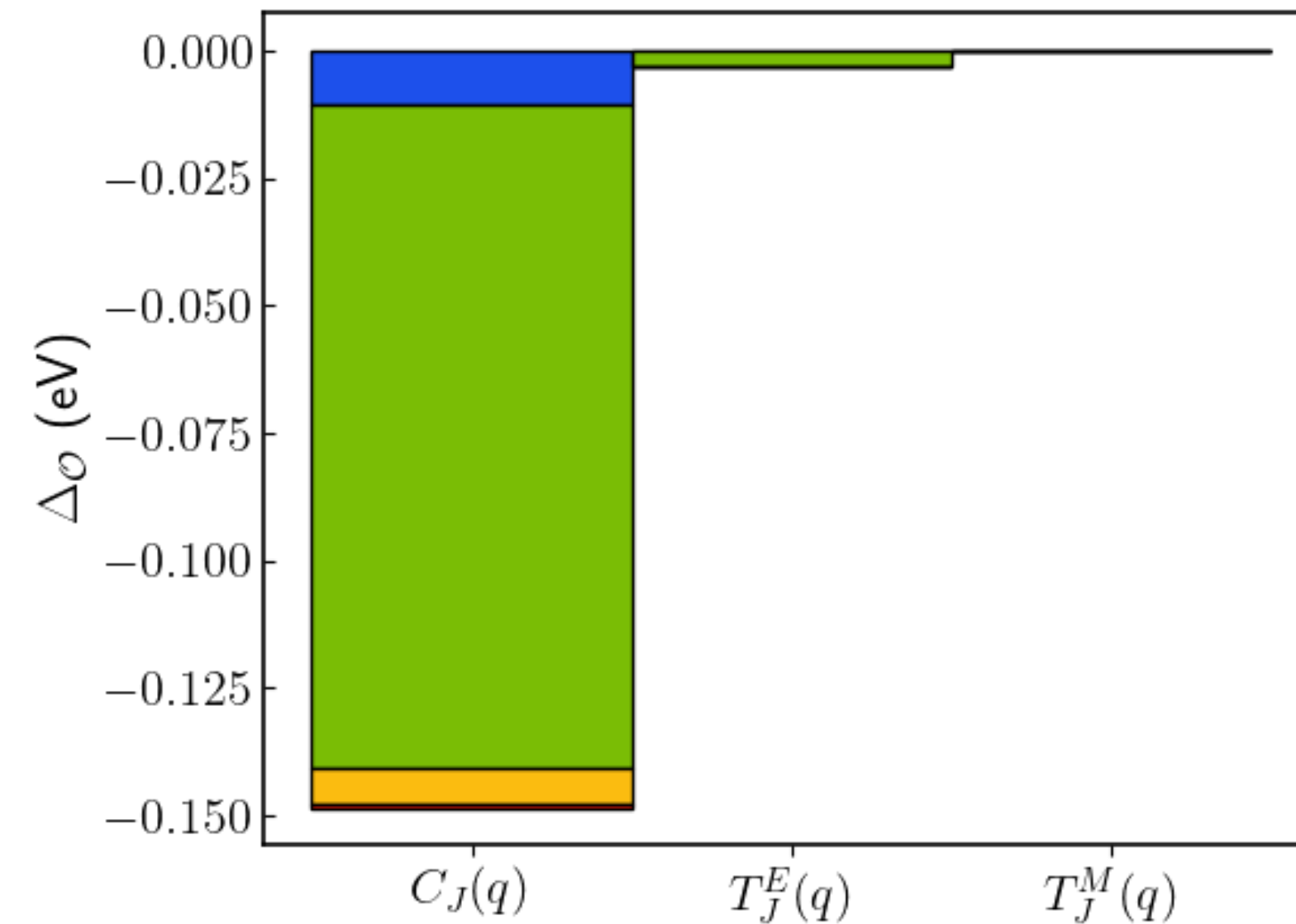
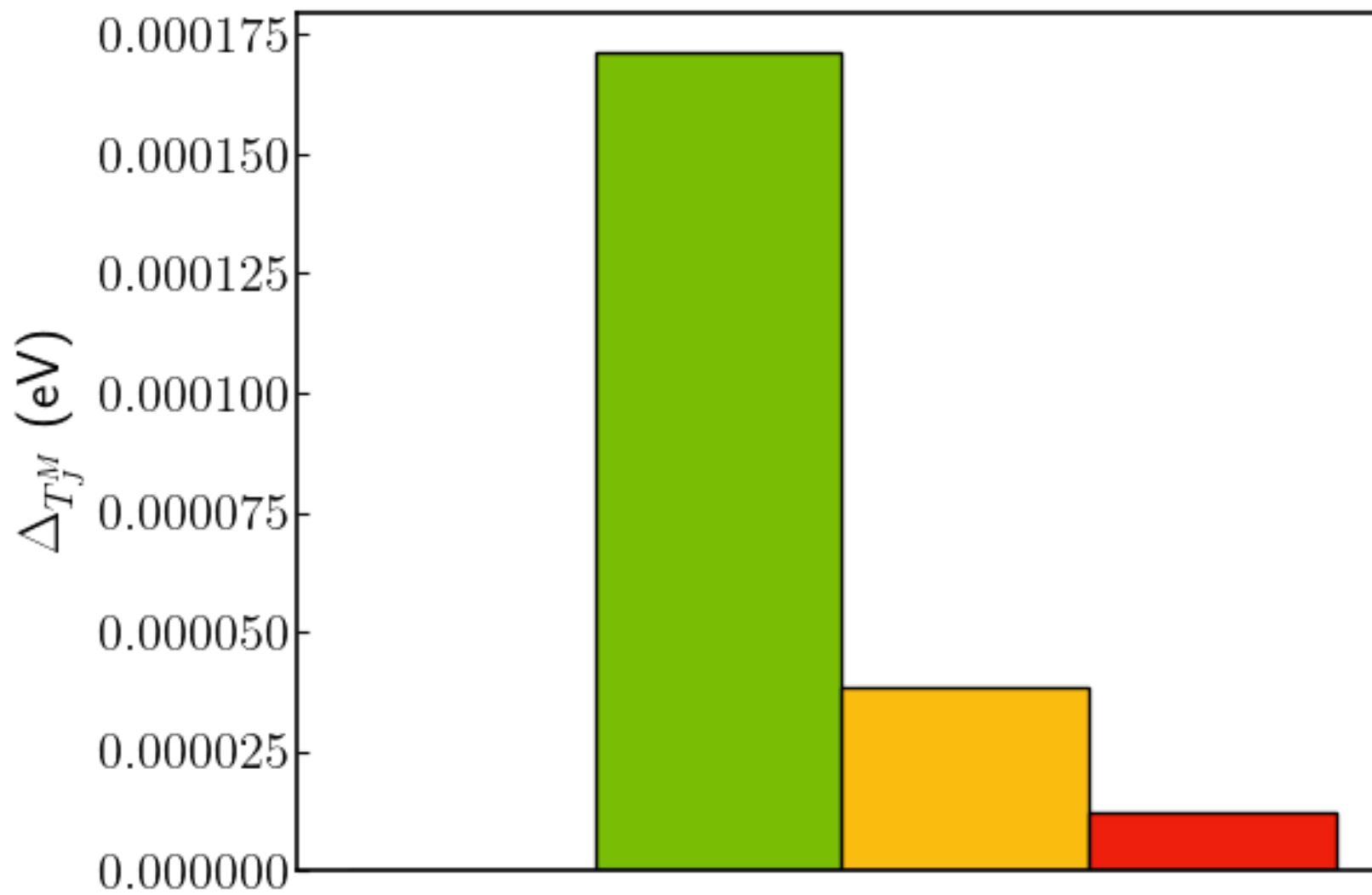
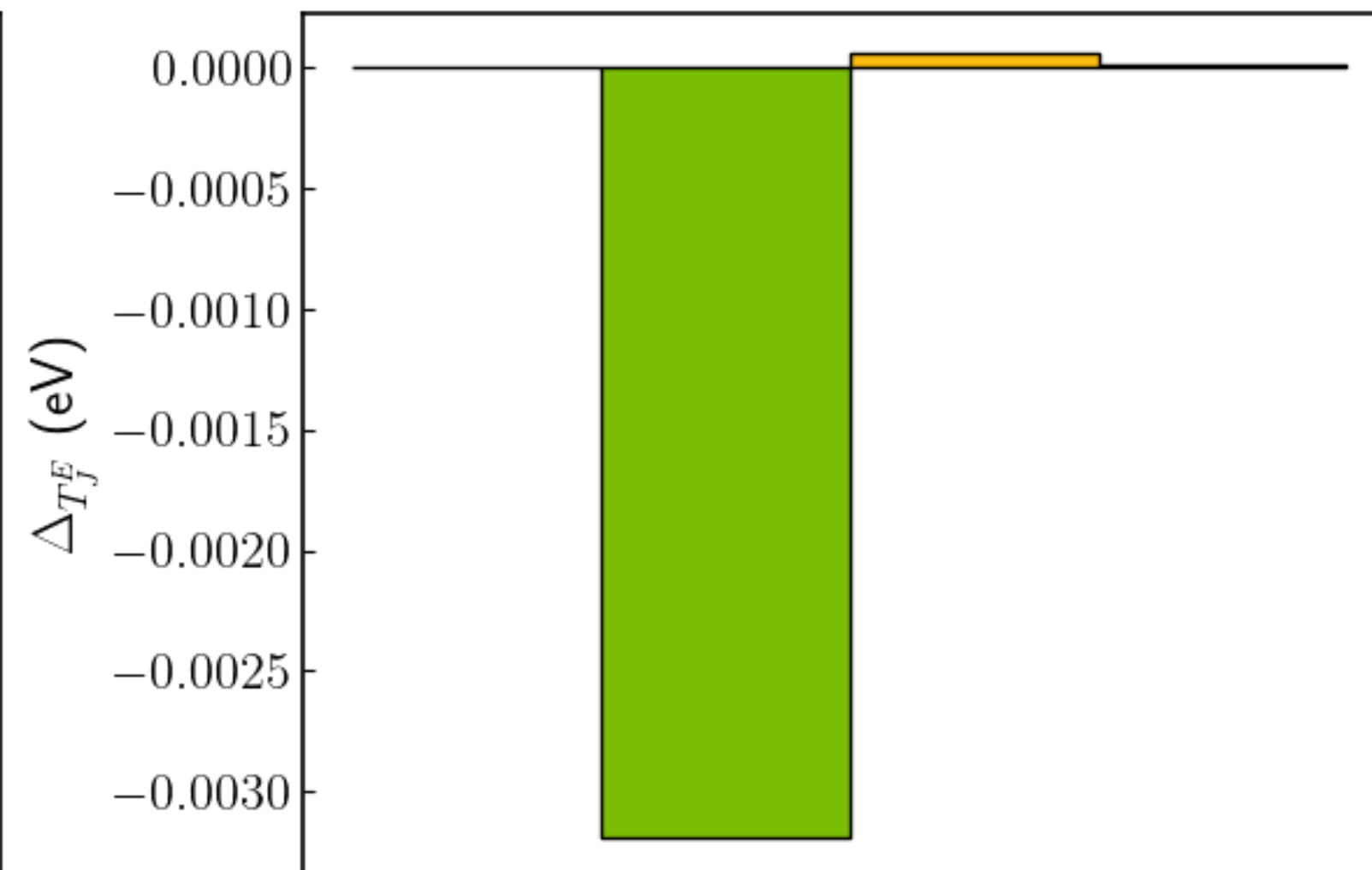
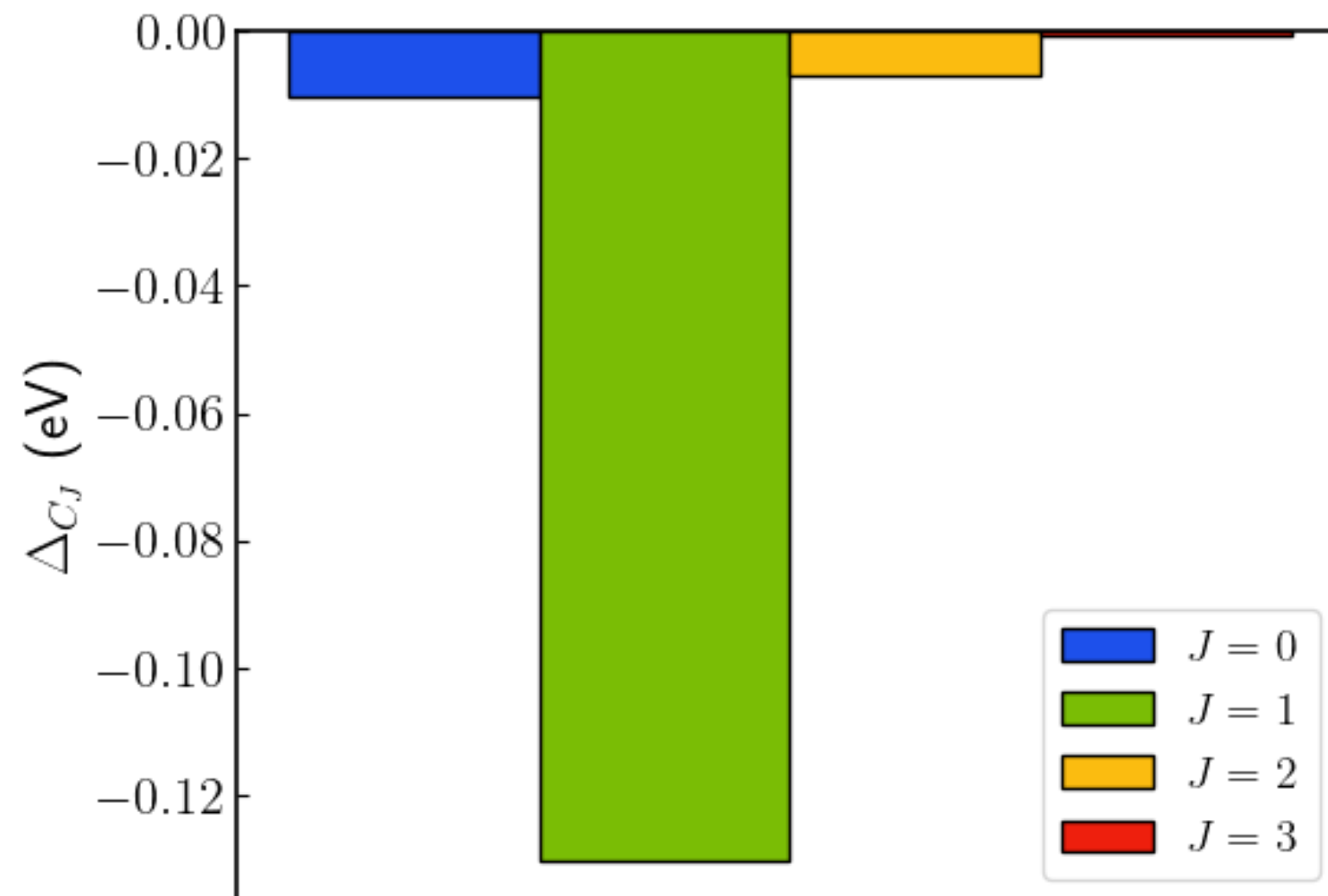


# A first test case for N4LO-E7 and $N_{\max} = 7$

## Numerical calculations

- ⦿  $q_{\max} = 700$  MeV and  $\Delta q = 10$  MeV
- ⦿ 10 different operators for  $J_{\max} = 3$
- ➔ **700 NCSM calculations at  $N_{\max} = 7$**

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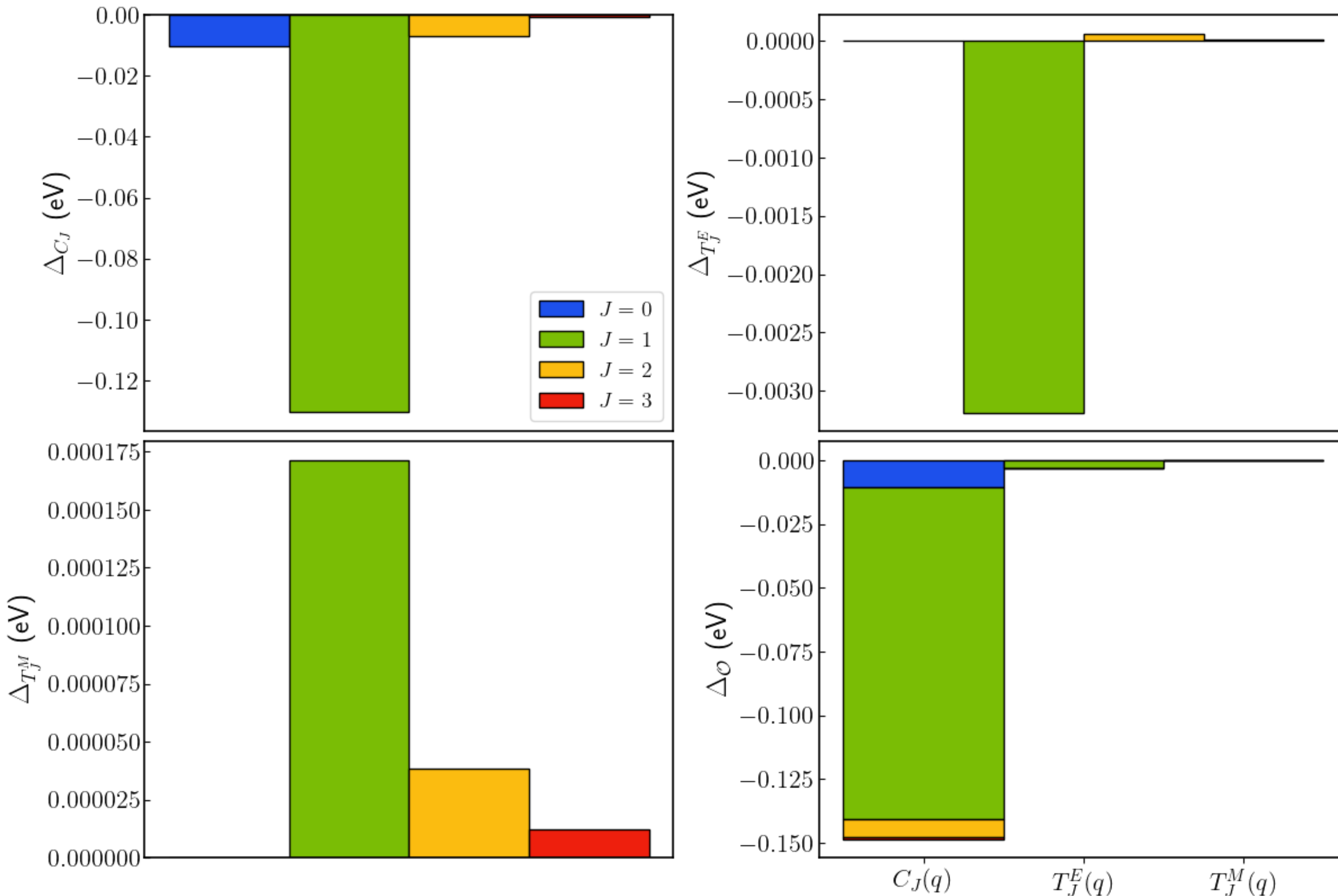


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10



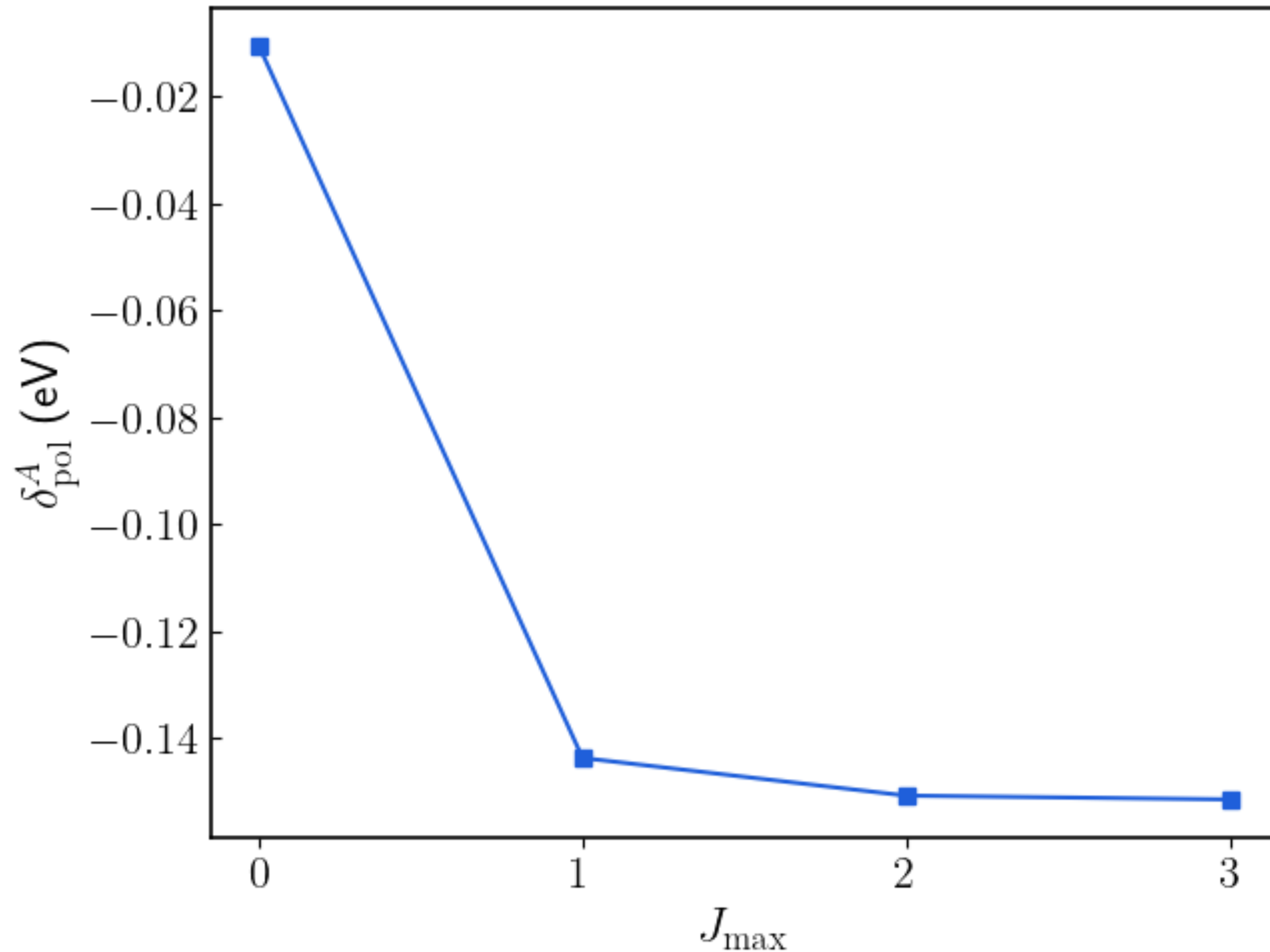
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### Observations

- ⊙ Contribution repartitions
  - Well-known **dipole** dominance
  - **Charge** contributions are dominant
- ⊙ Negligible contributions
  - TM is negligible for any  $J$
  - TE is relevant only for  $J = 1$
- ➔ **Only half the operators are relevant**

# Checking convergence in $J_{\max}$

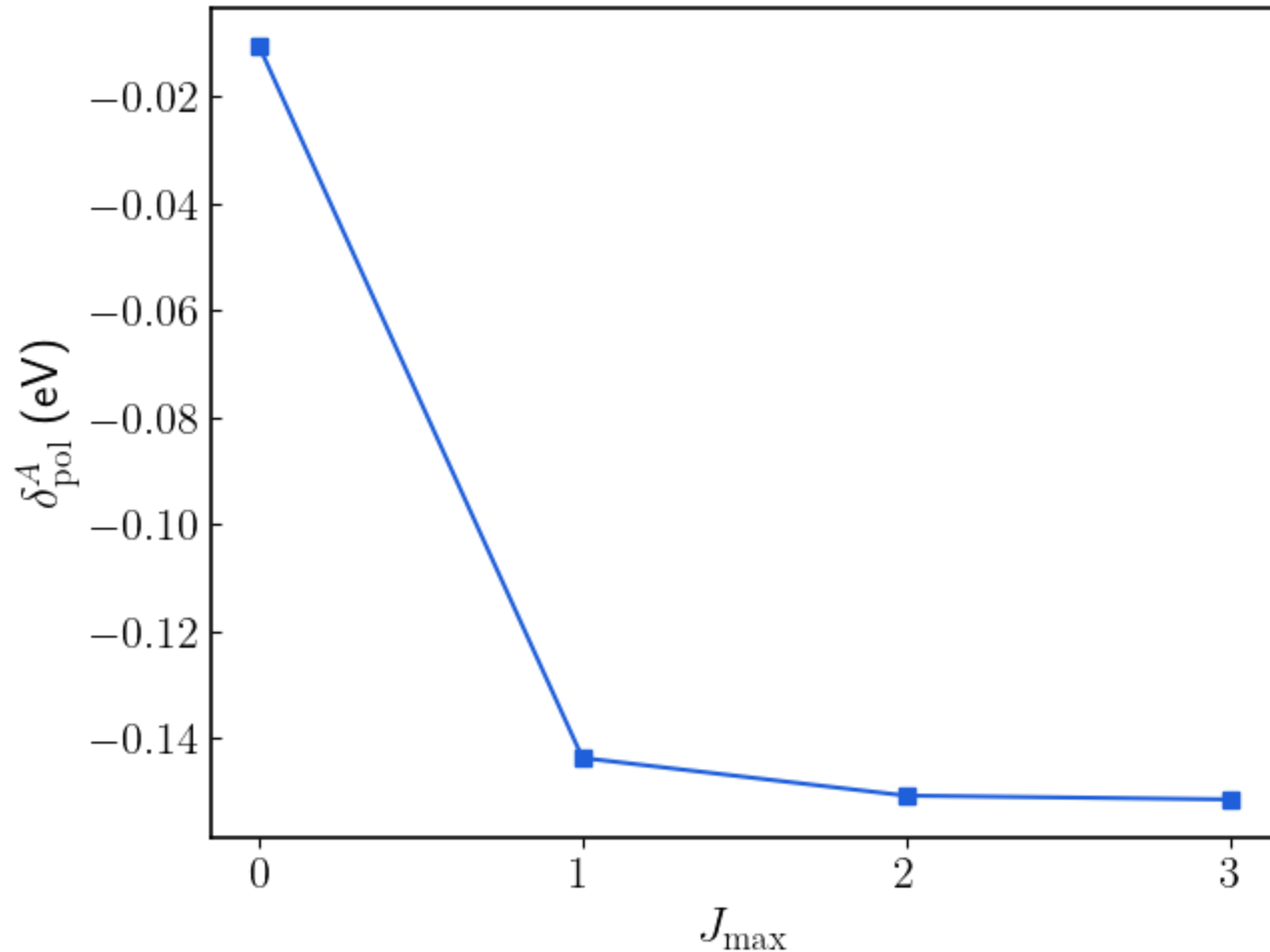


## Results

- ⦿ Here shown for  $N_{\max} = 7$  and N4LO-E7
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11



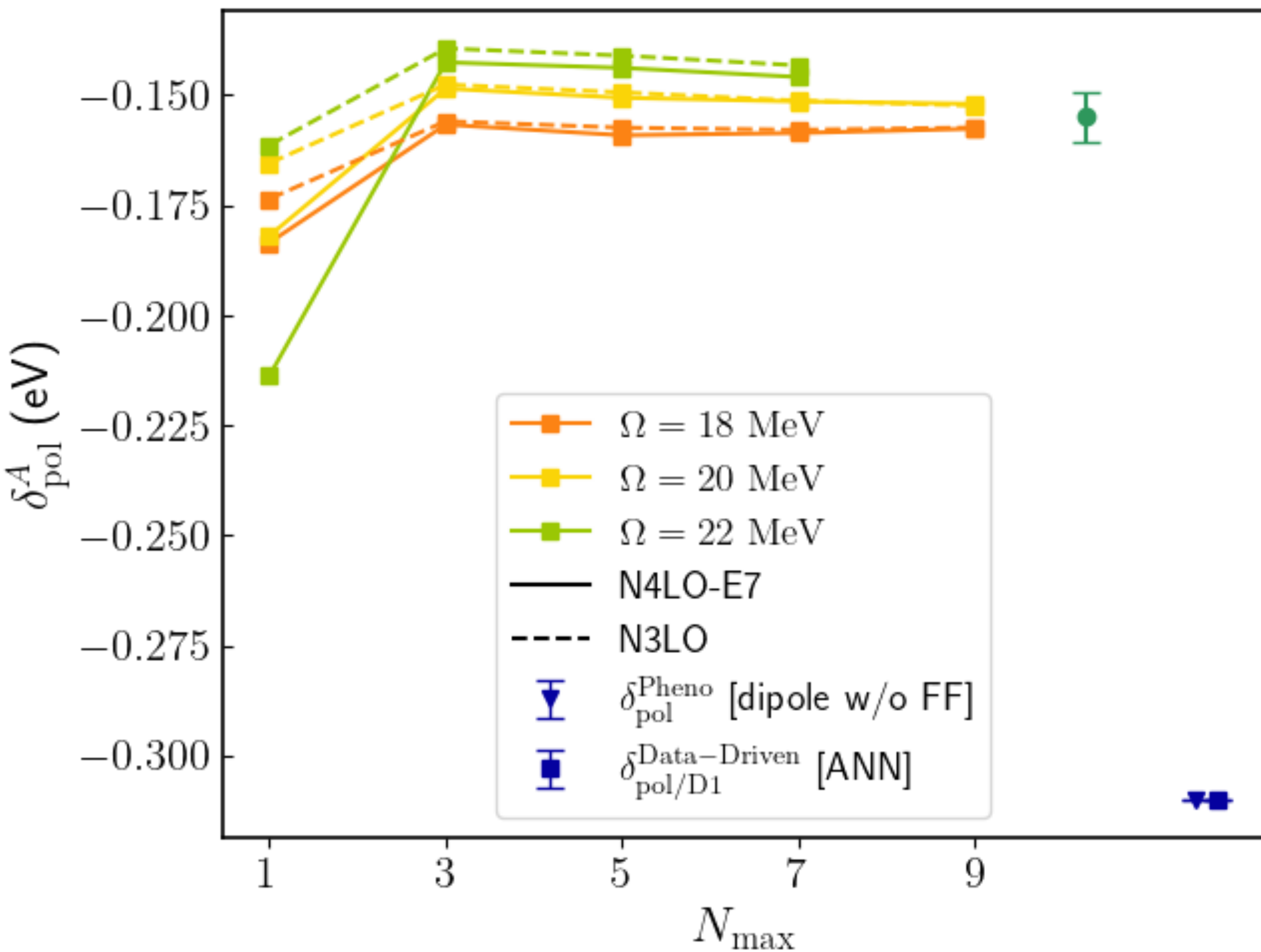
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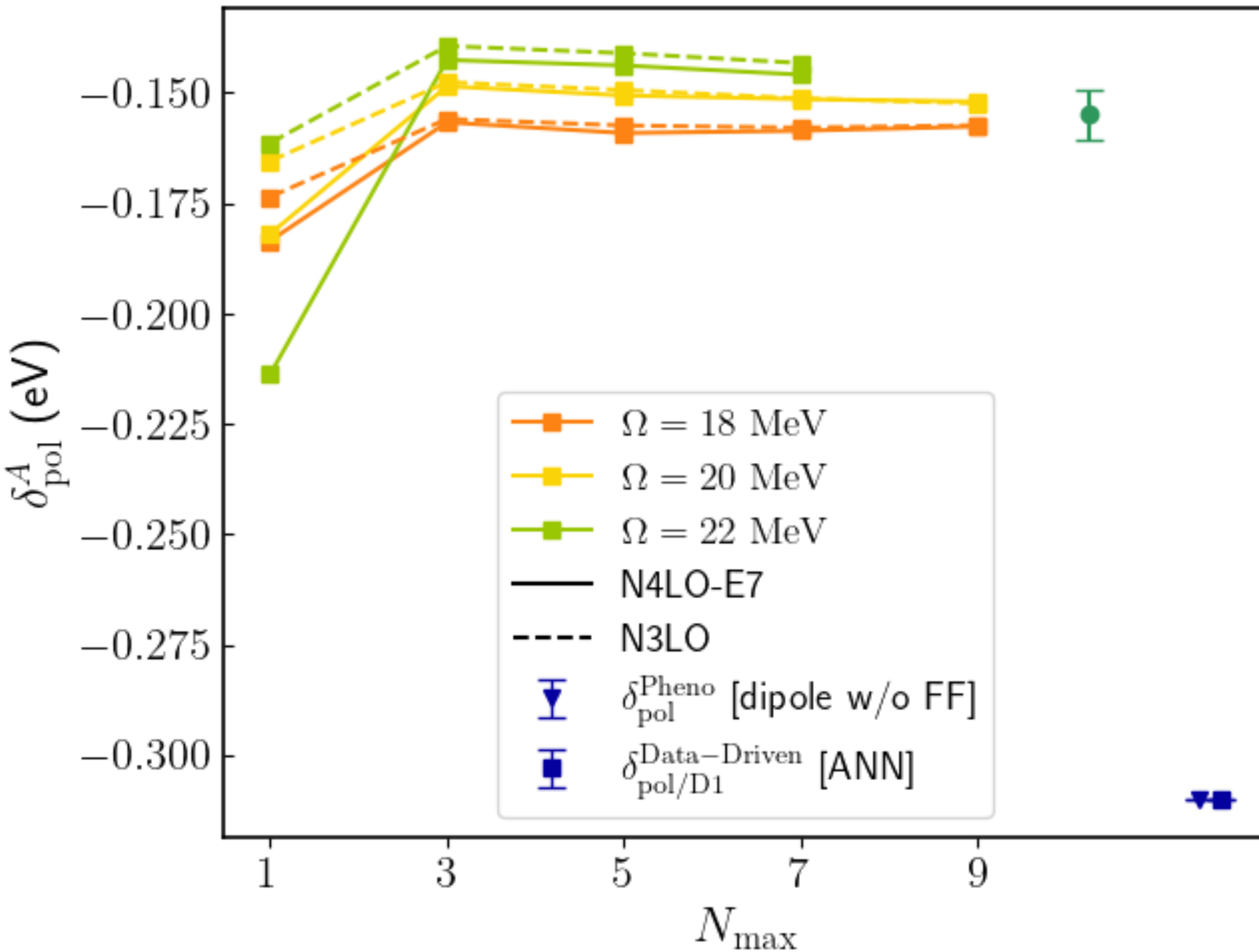
$$\epsilon_{J_{\max}} \lesssim 1 - 10 \text{ meV}$$

**Multipole truncation  $\Rightarrow$  Negligible uncertainty**

# Dependence on $(\Omega, N_{\max})$ and the interaction



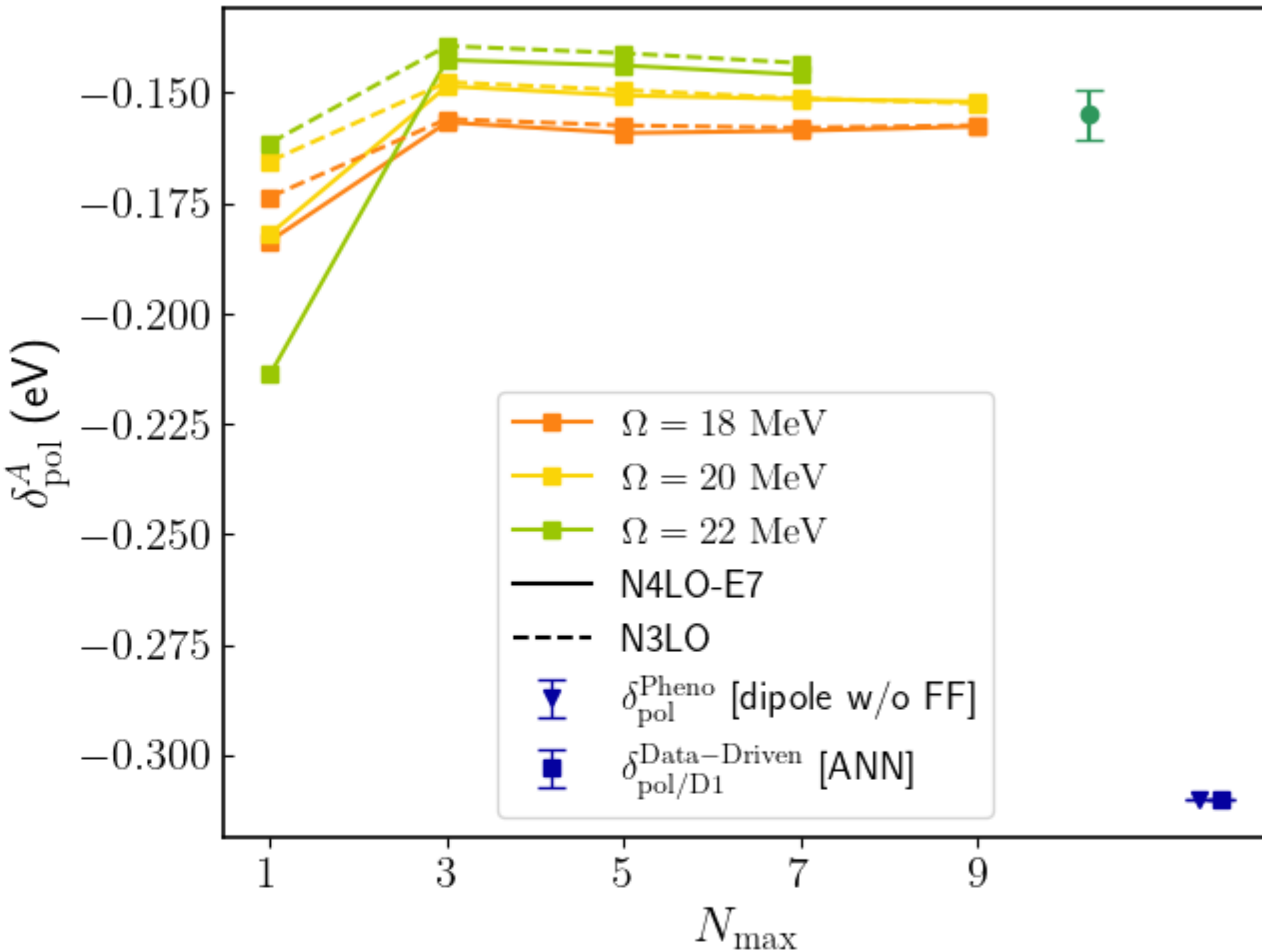
# Dependence on $(\Omega, N_{\max})$ and the interaction



## Preliminary numerical results

- Model-space dependence
  - Optimal frequency around 20 MeV
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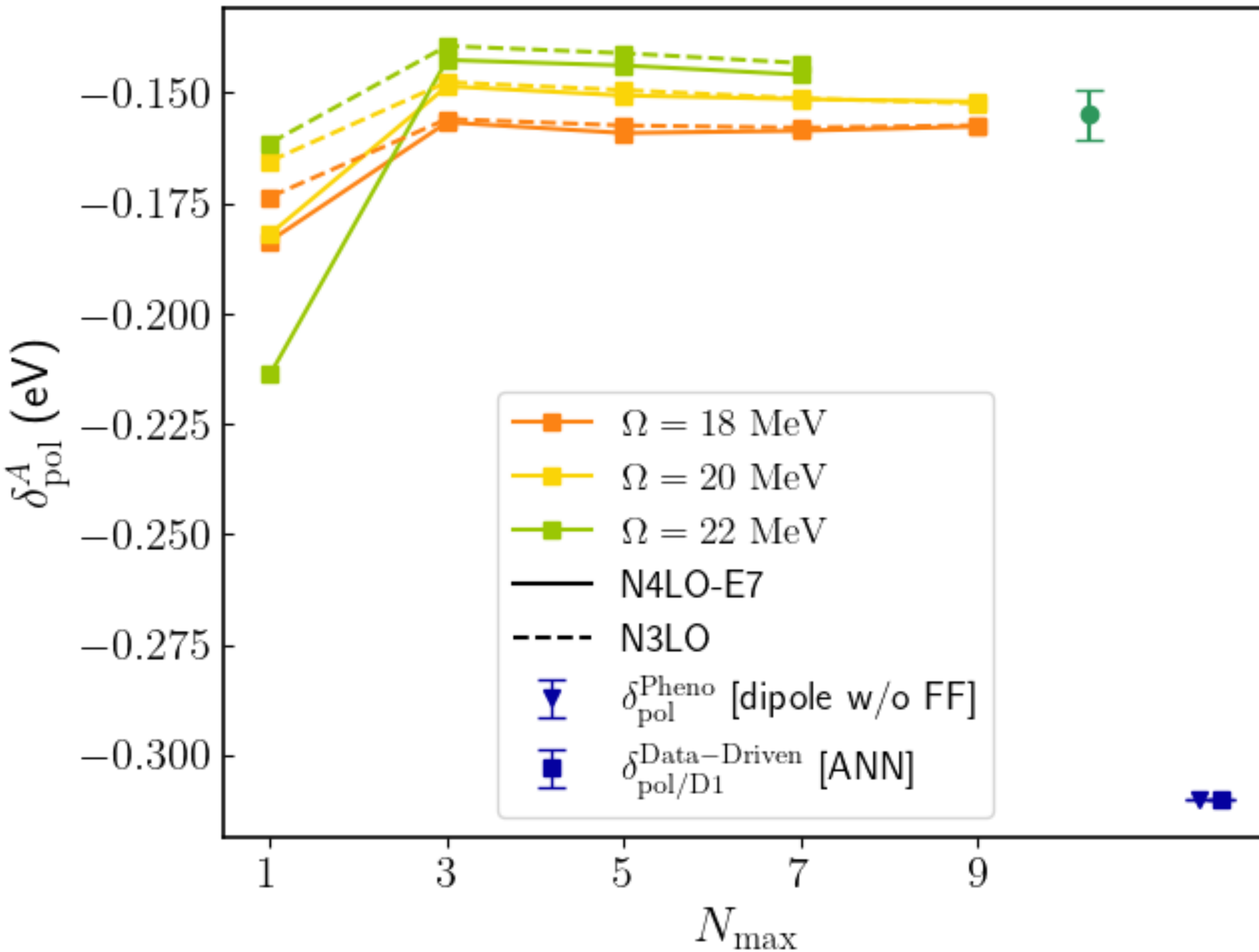
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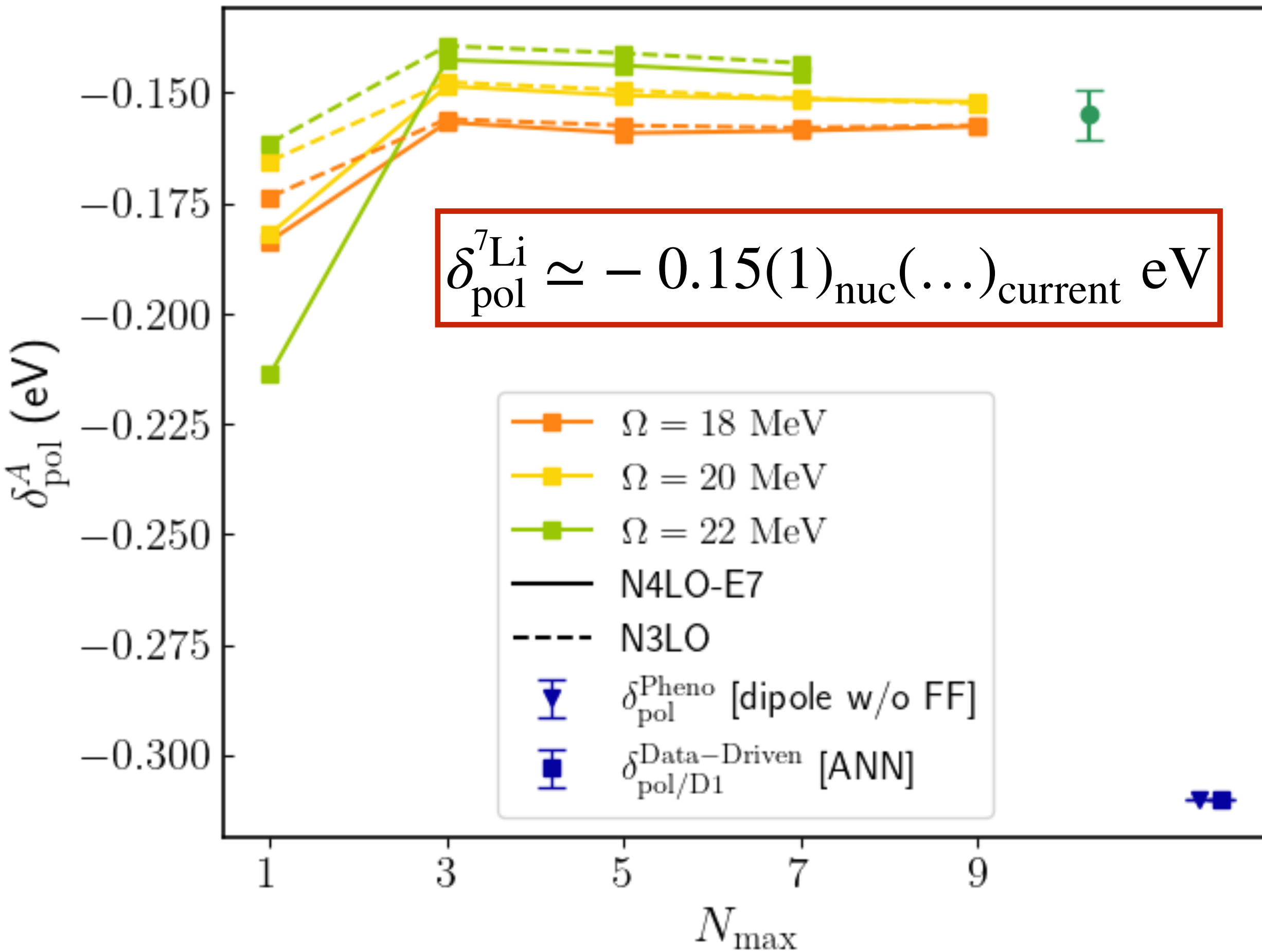
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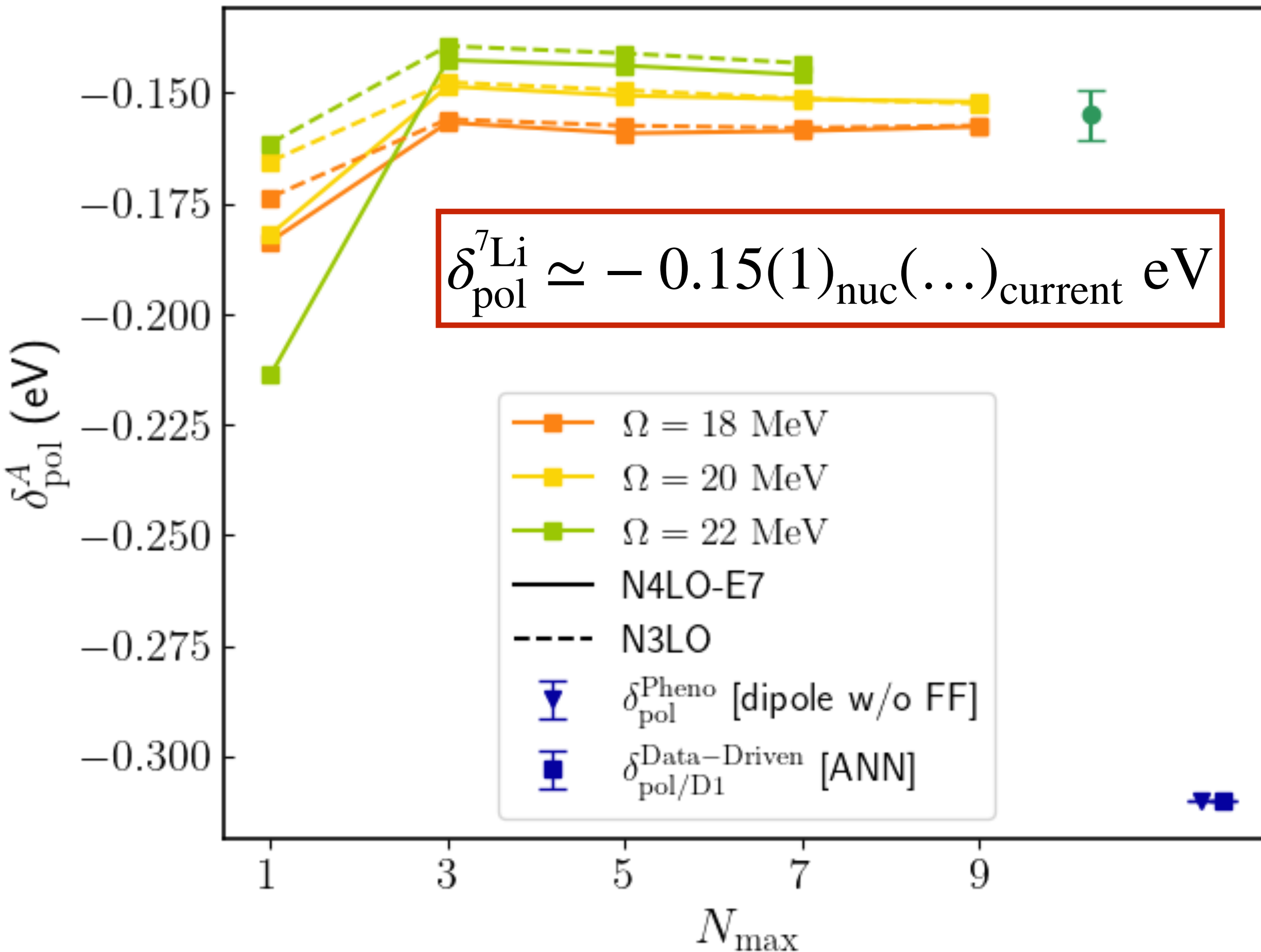
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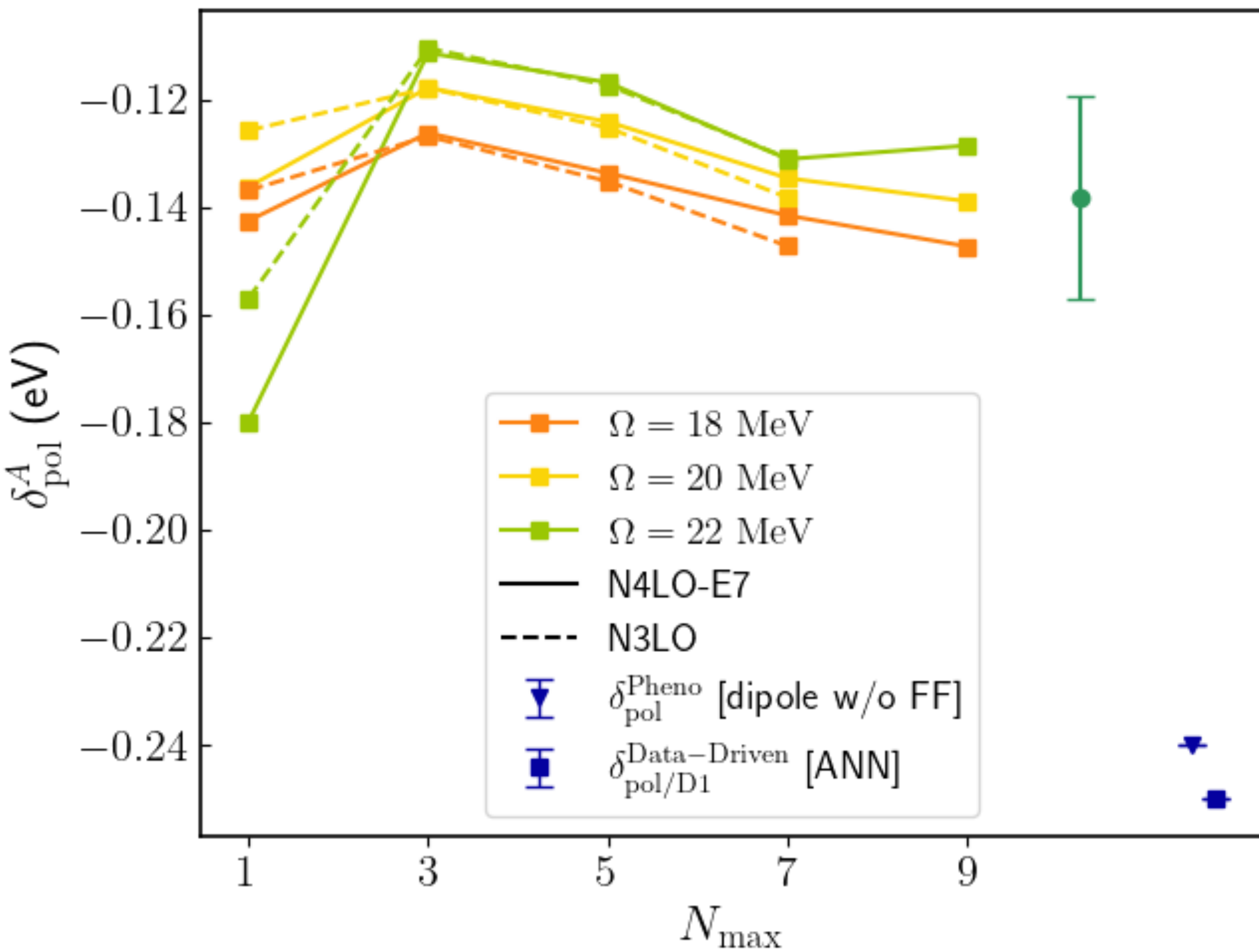


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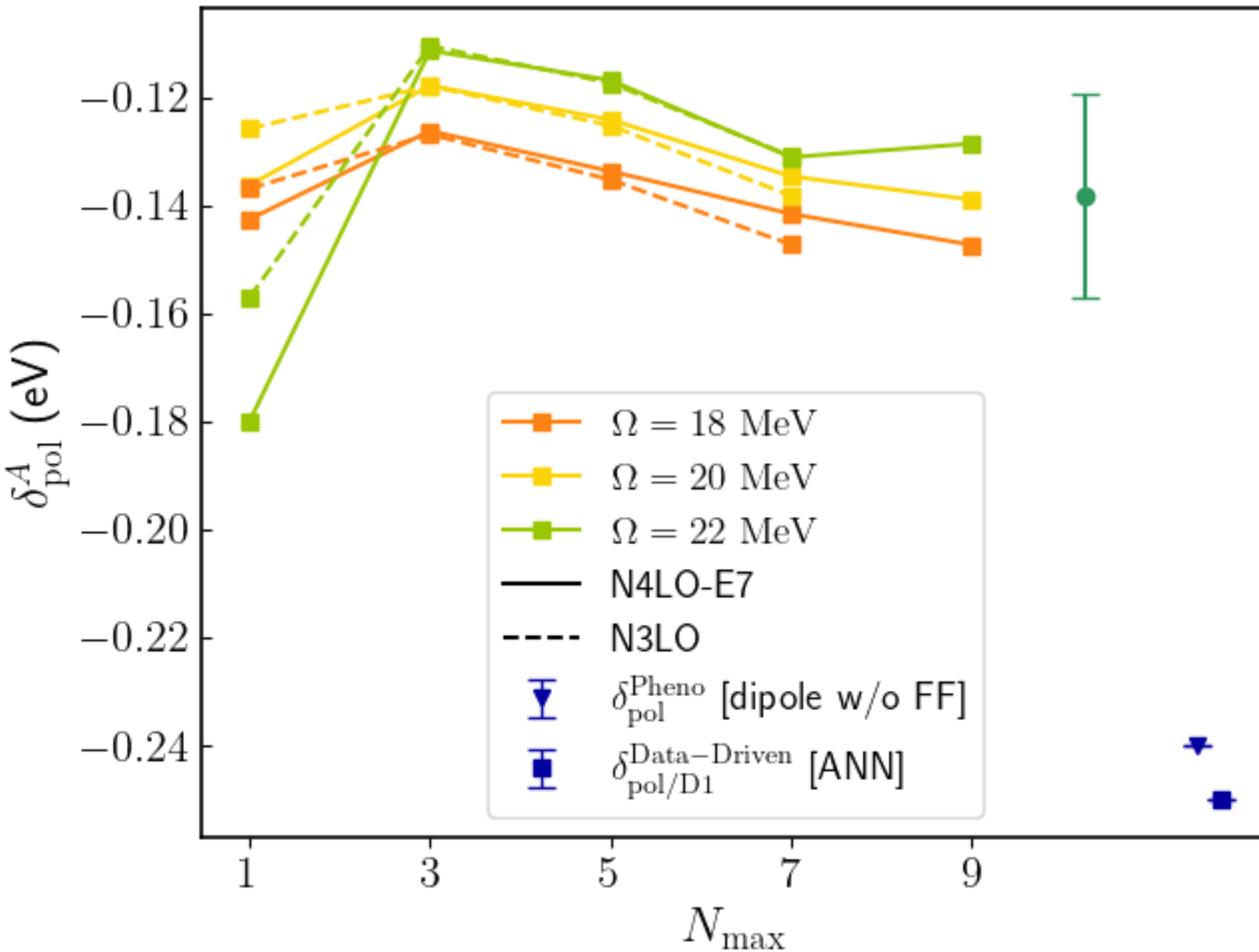
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**0.05 eV precision seems reachable for nuclear structure corrections!**

# Model-space dependence for ${}^6\text{Li}$



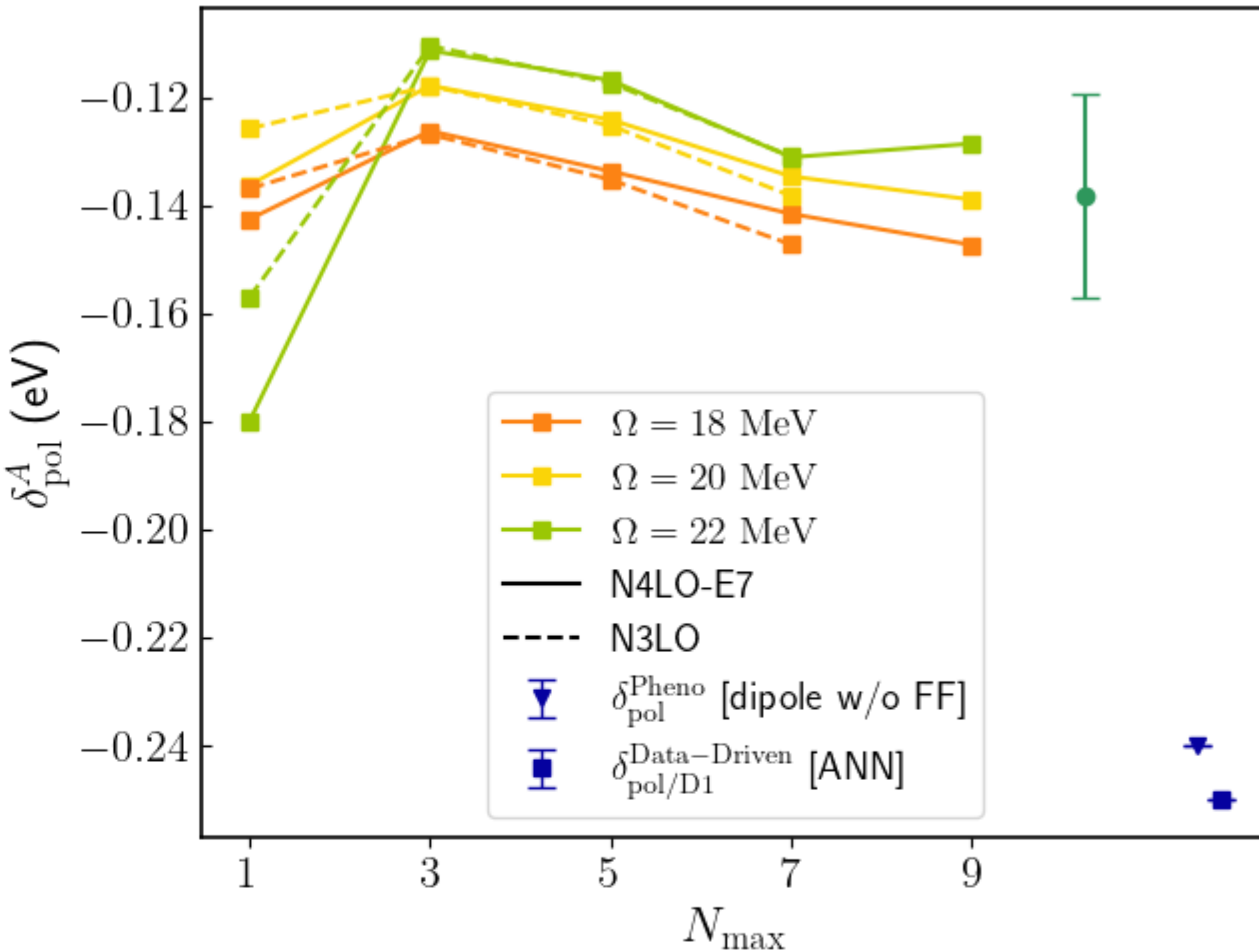
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- Critical role of latest supercomputer generation

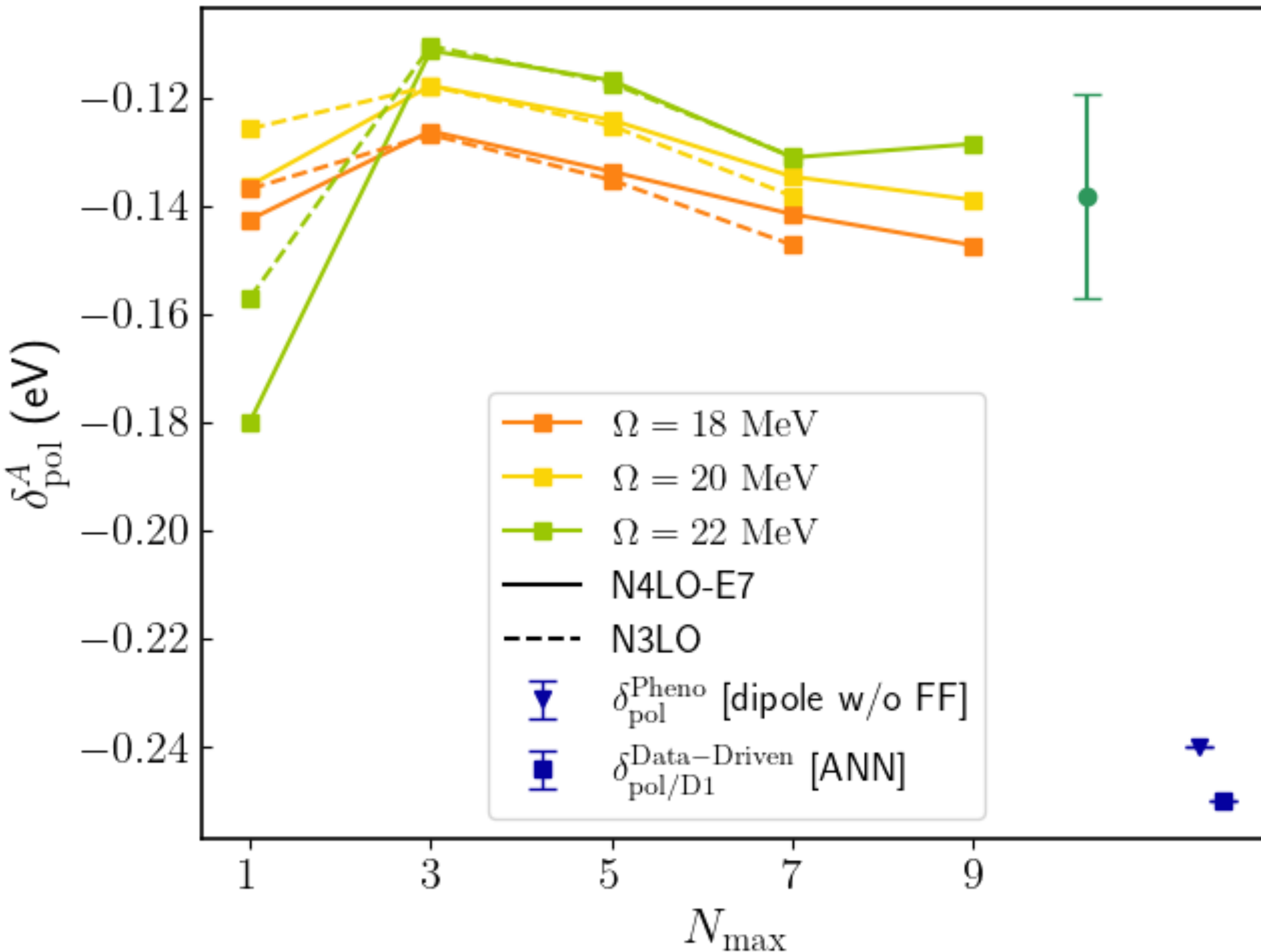
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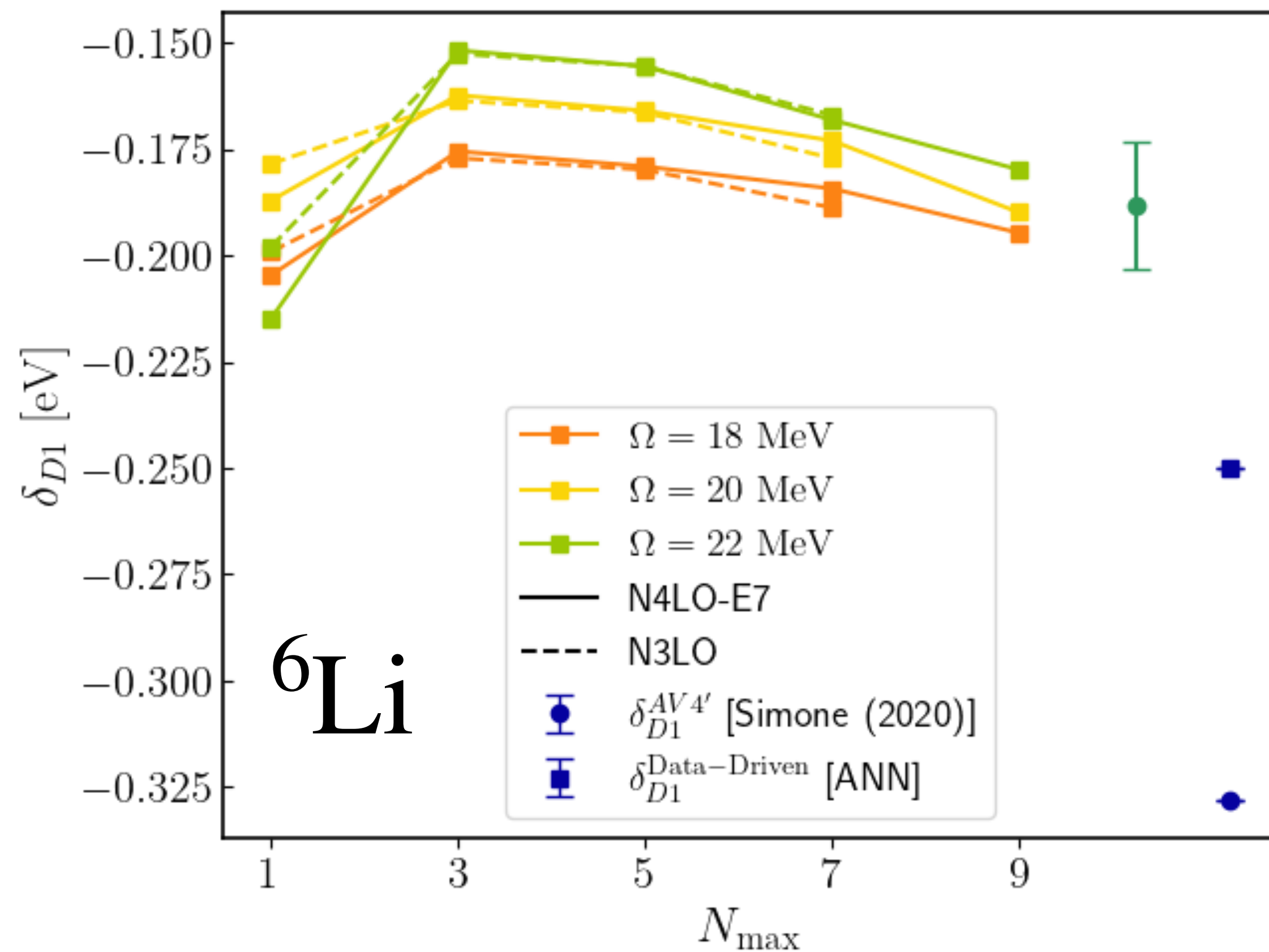
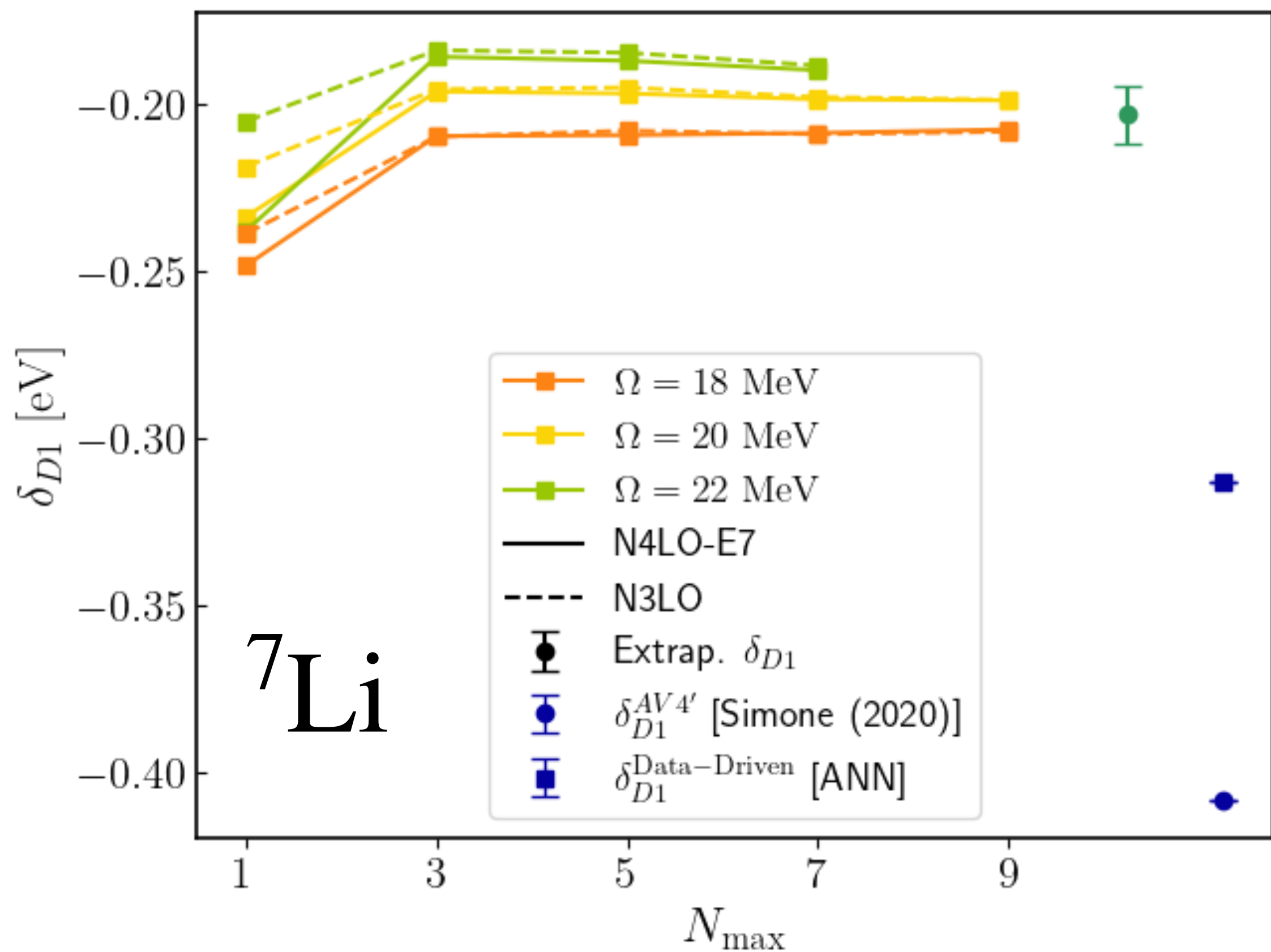


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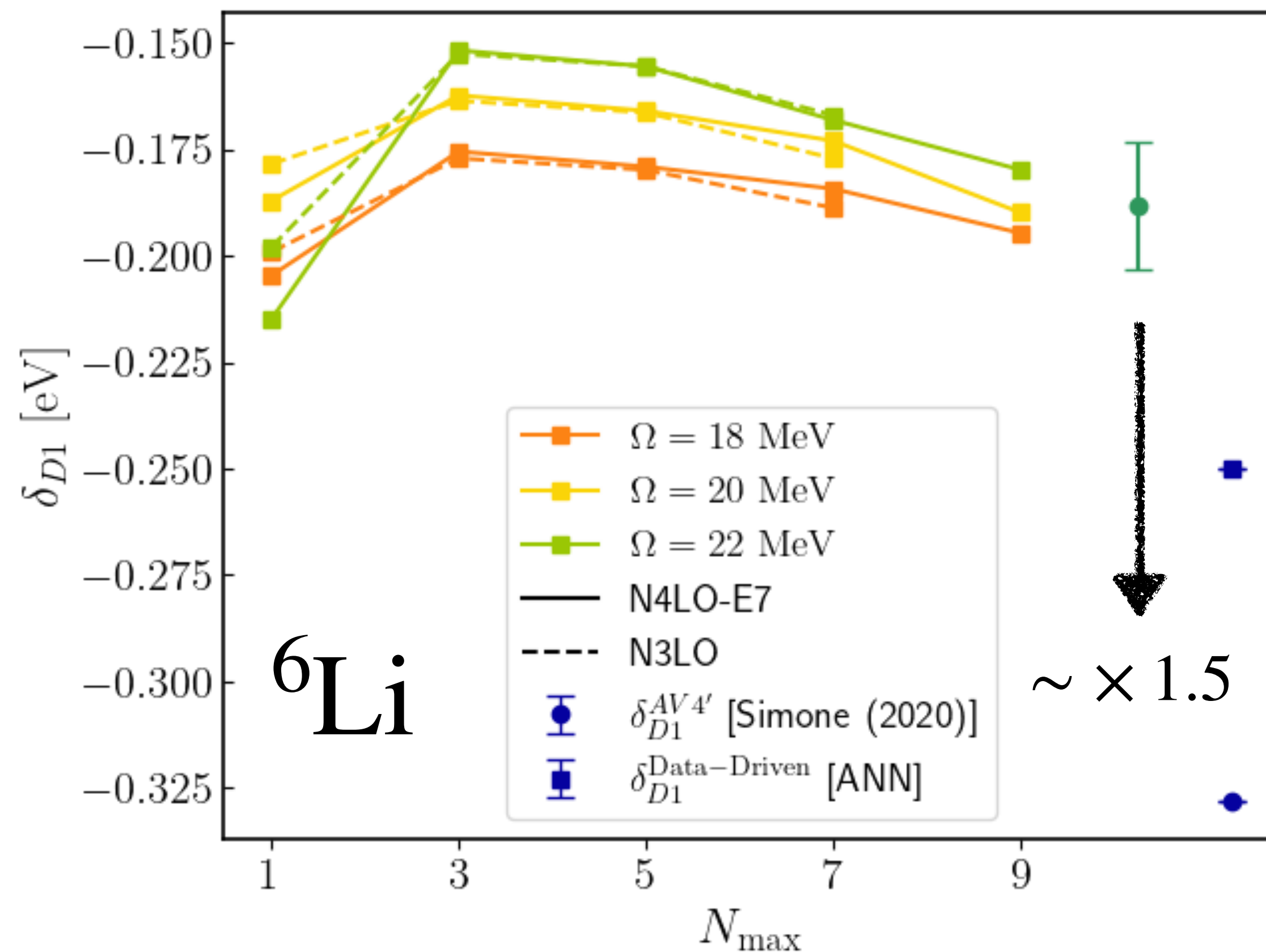
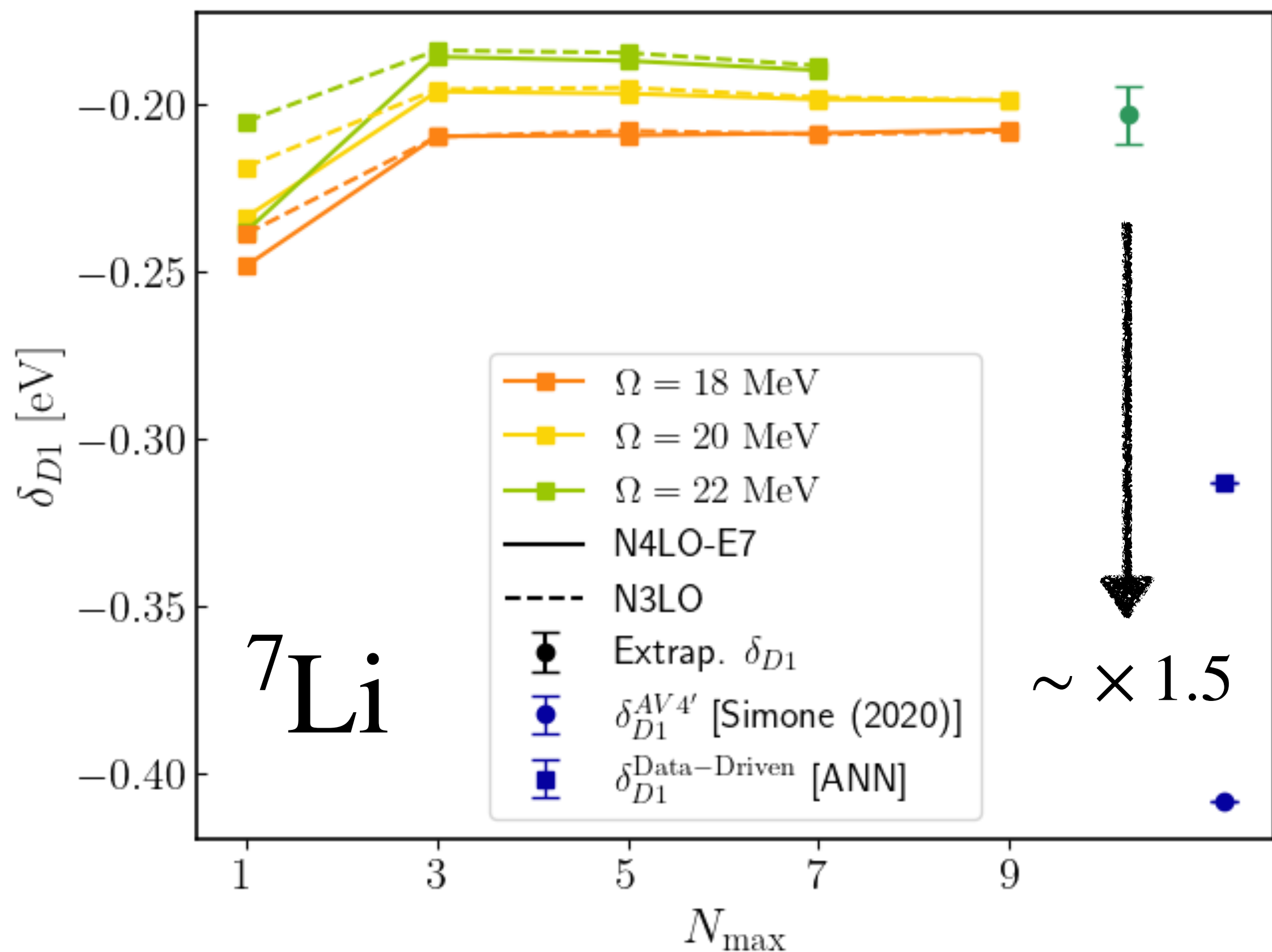
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**On-going results to be finalized** :  $\delta_{\text{pol}}^{{}^6\text{Li}} \sim -0.14 (\dots)_{\text{nuc}} (\dots)_{\text{current}} \text{ eV}$

# Benchmarking with pheno estimations



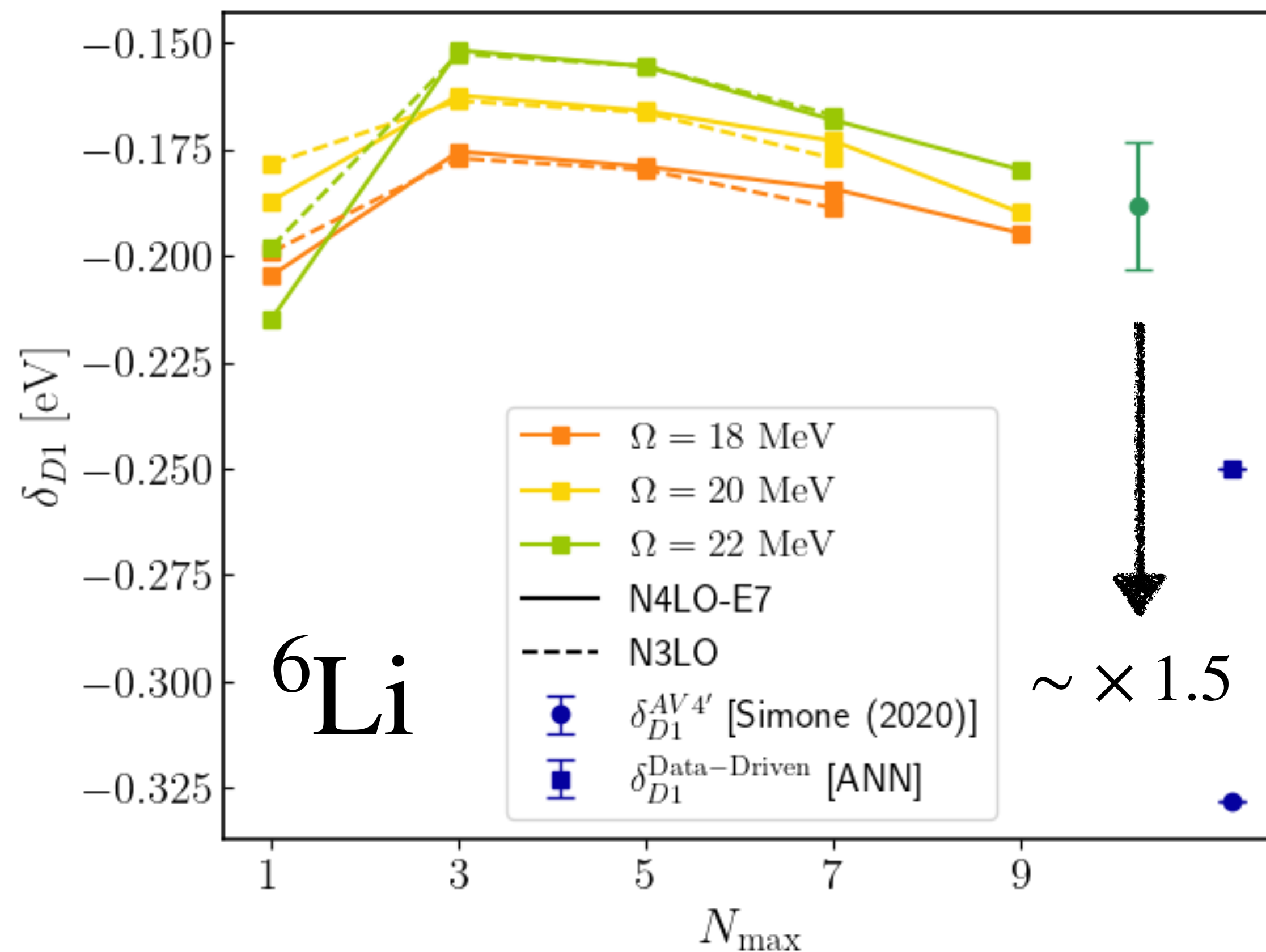
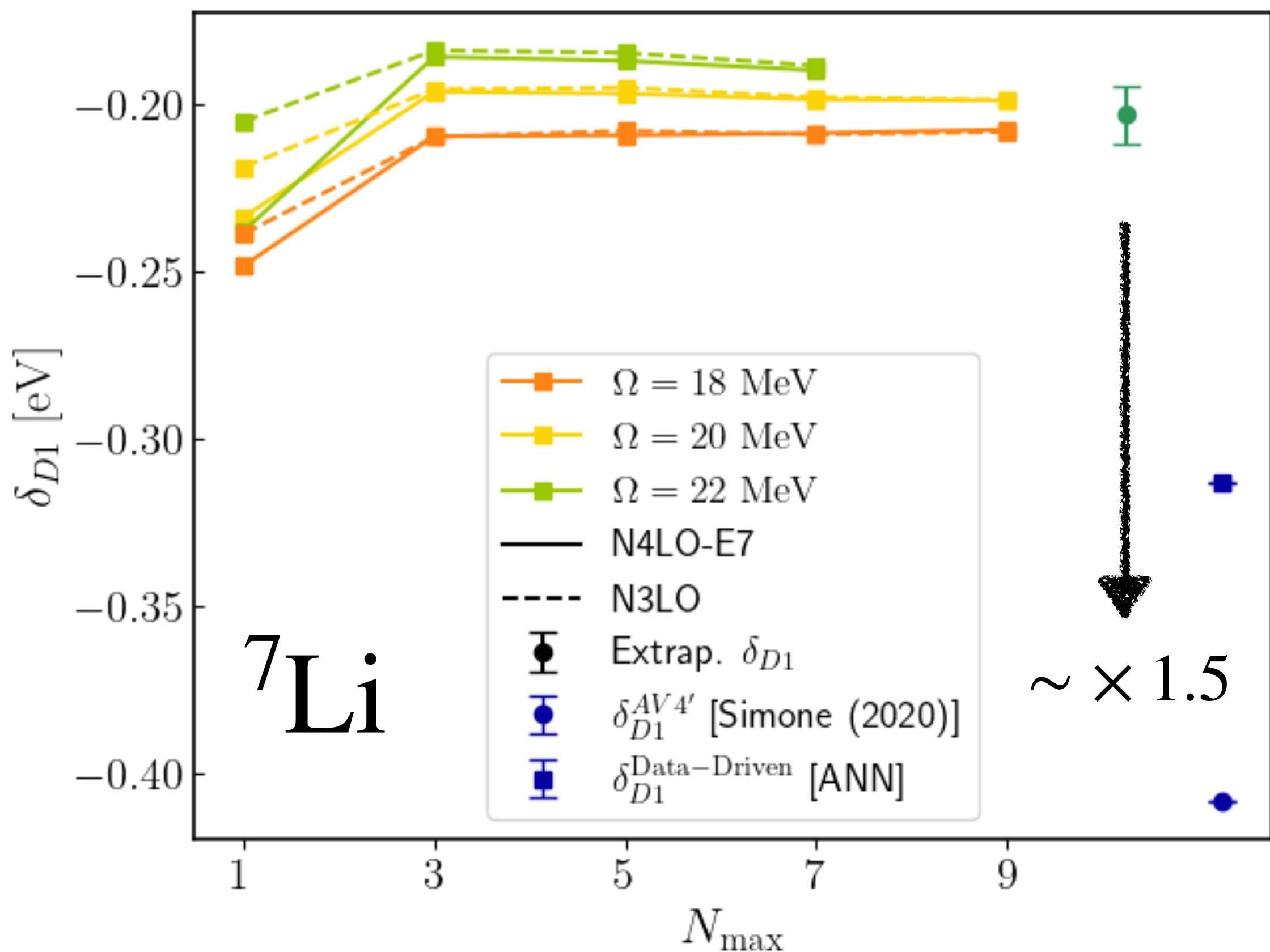
# Benchmarking with pheno estimations



**On-going investigation of discrepancy:** *might* come from non-trivial  $(M_i, M)$  dependence

$$S_X(\omega, q) \equiv \sum_{J \geq 0} \sum_{N \neq 0} |\langle N | O_{X,J}(q) | \Psi \rangle|^2 \delta(E_N - E_0 - \omega)$$

# Benchmarking with pheno estimations



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**NCSM calculations**  
 $\times (2J_i + 1)(2J + 1)$

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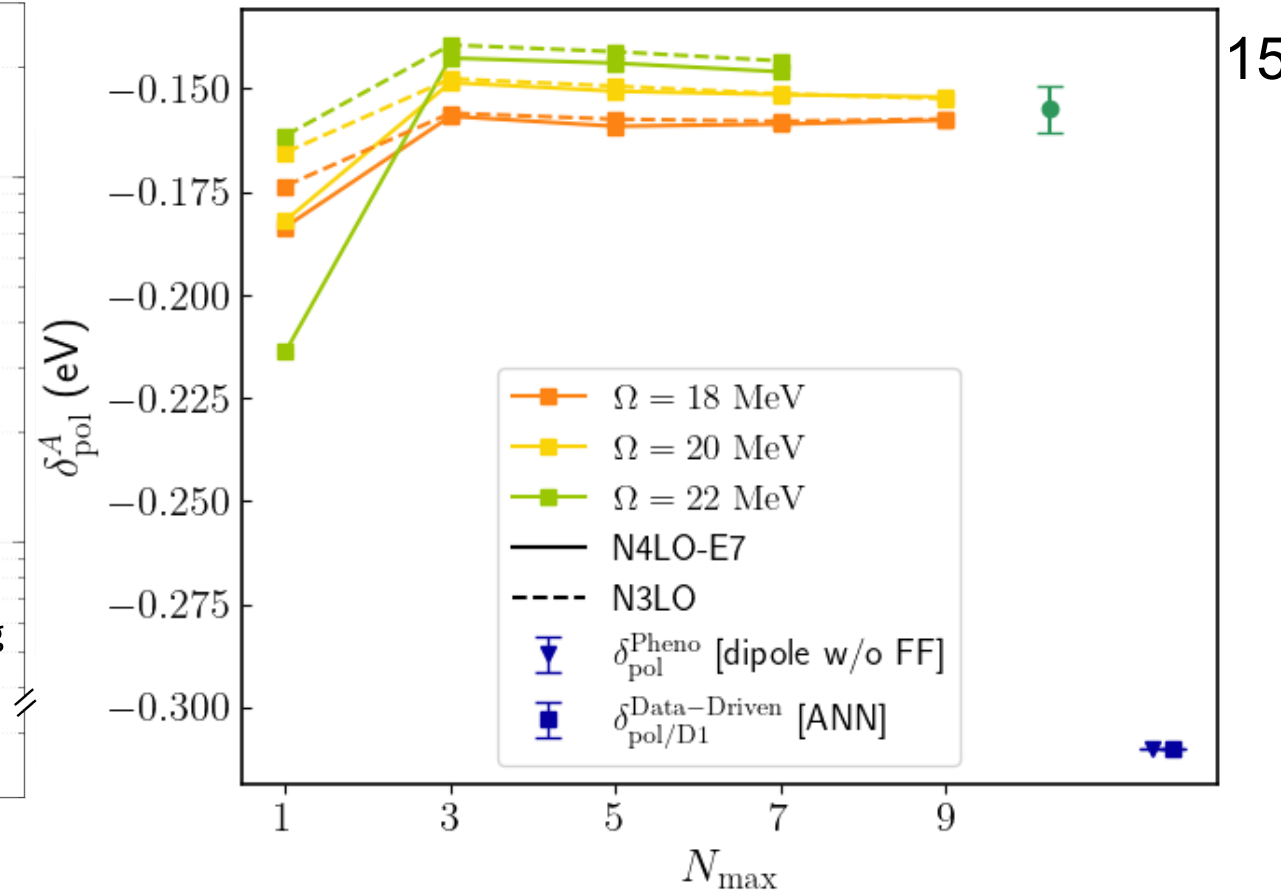
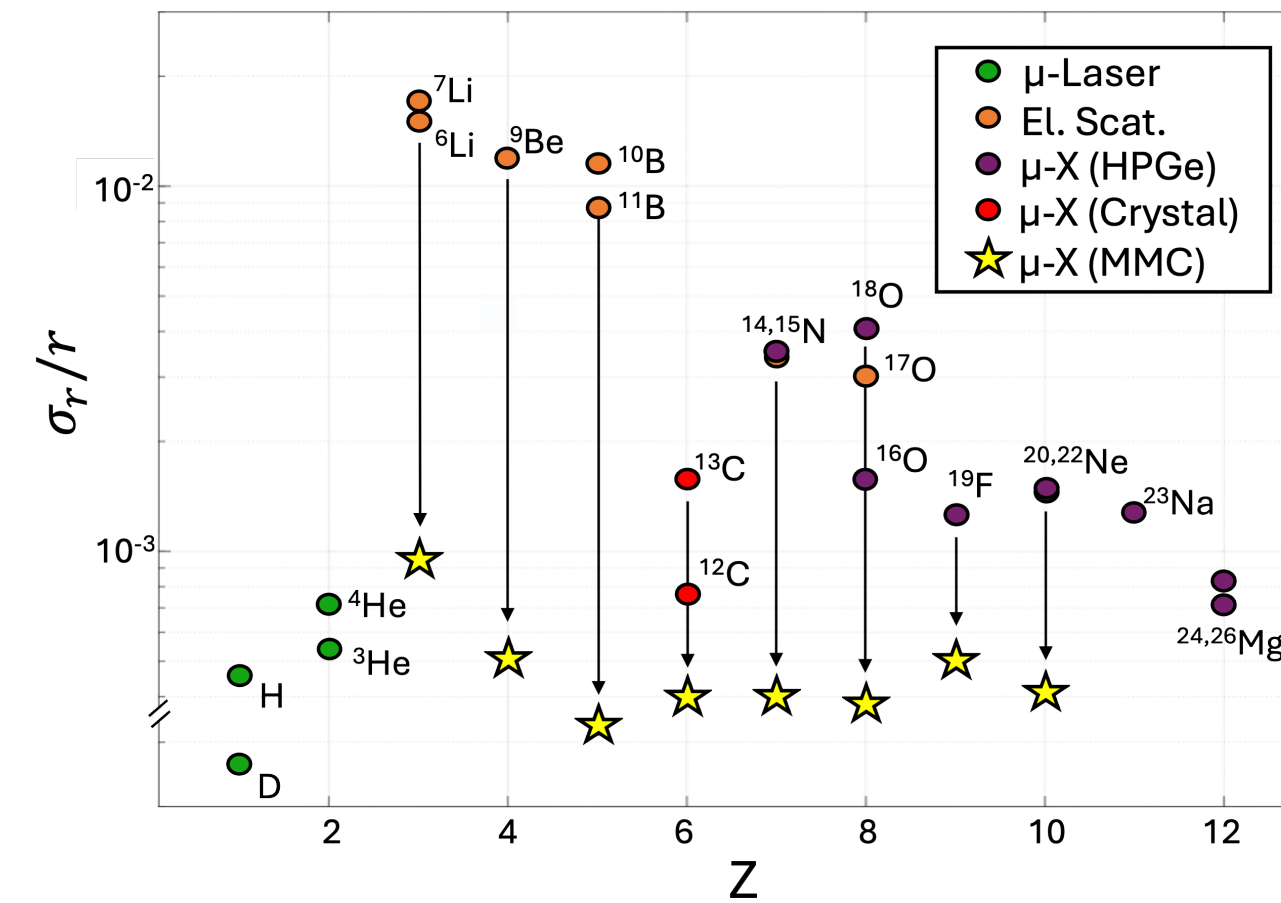
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# Conclusion

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## Summary

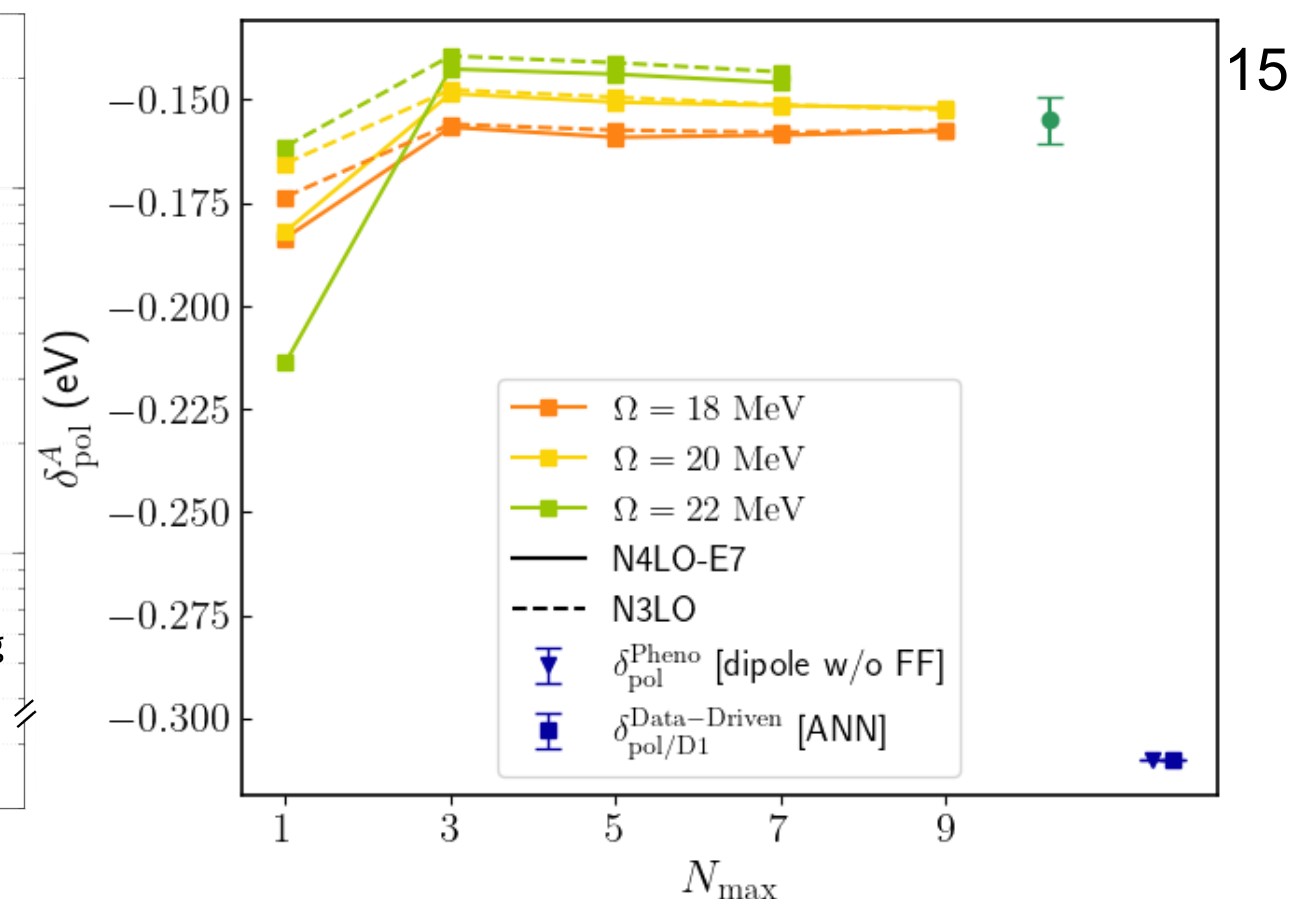
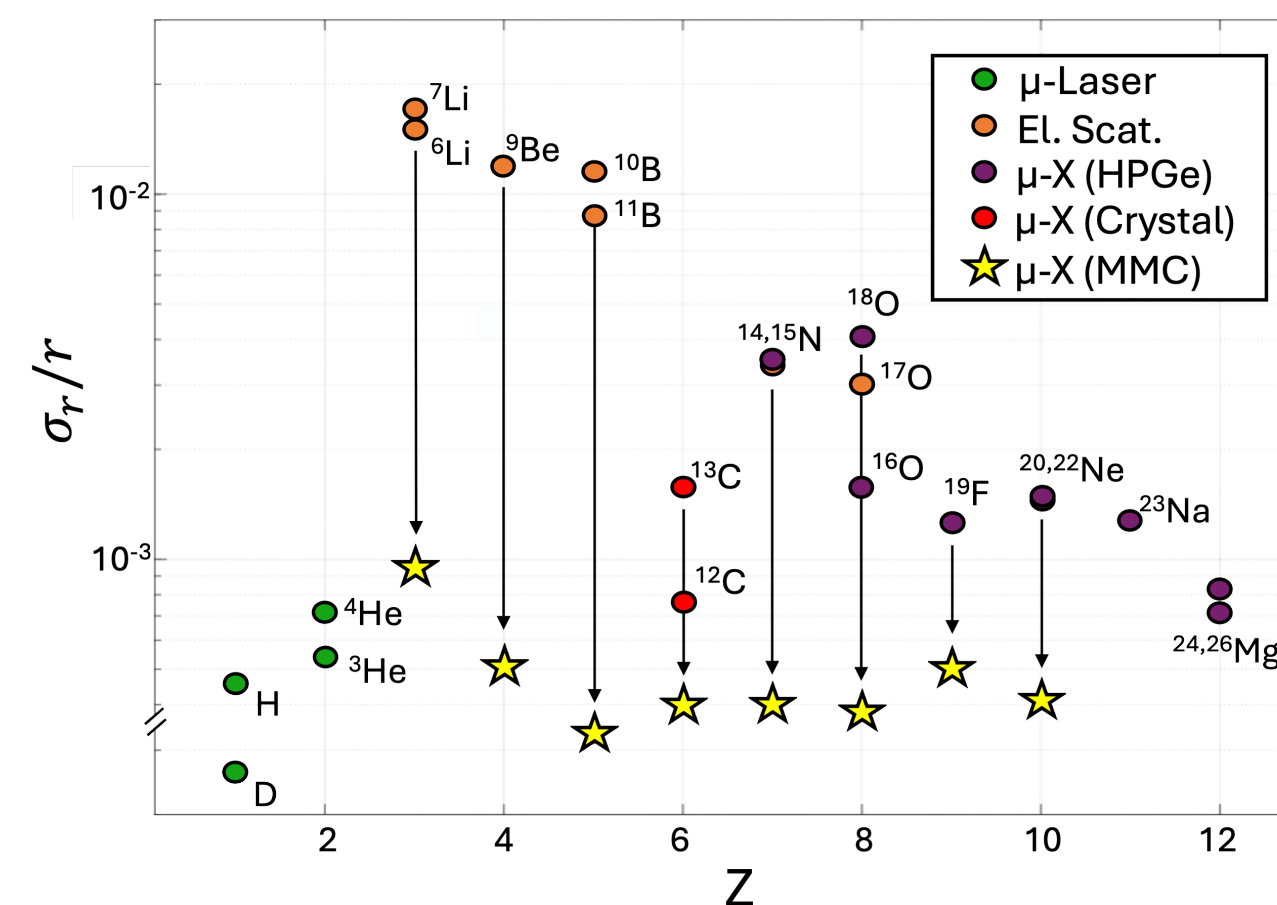
- Muonic atoms: a precision probe for nuclear physics
  - Radii extraction: reference point + isotope-shift
  - Precise reference point: muonic atoms
- ➔ **QUARTET**: 10x exp. improvement for  $Z \lesssim 10$



# Conclusion

## Summary

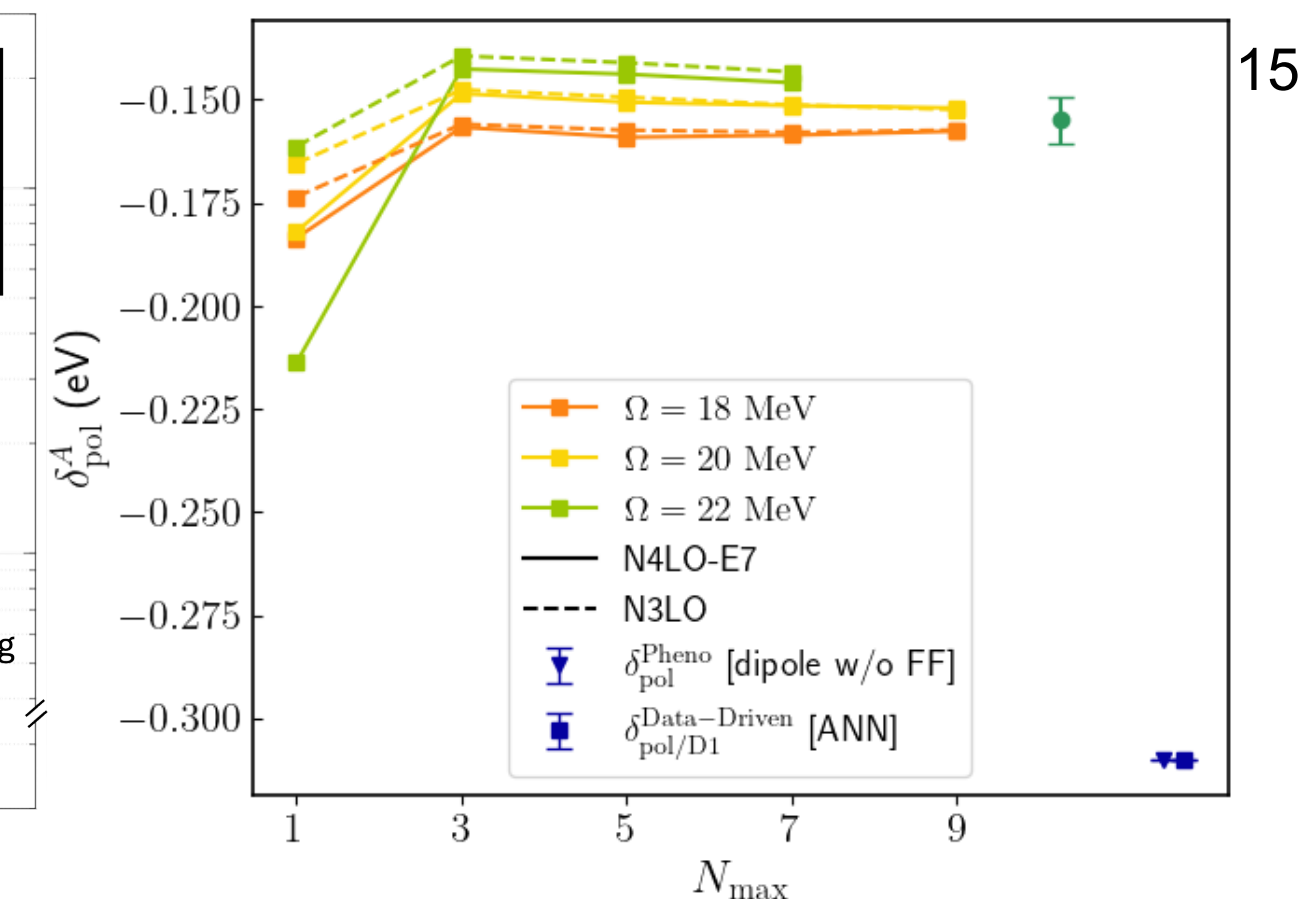
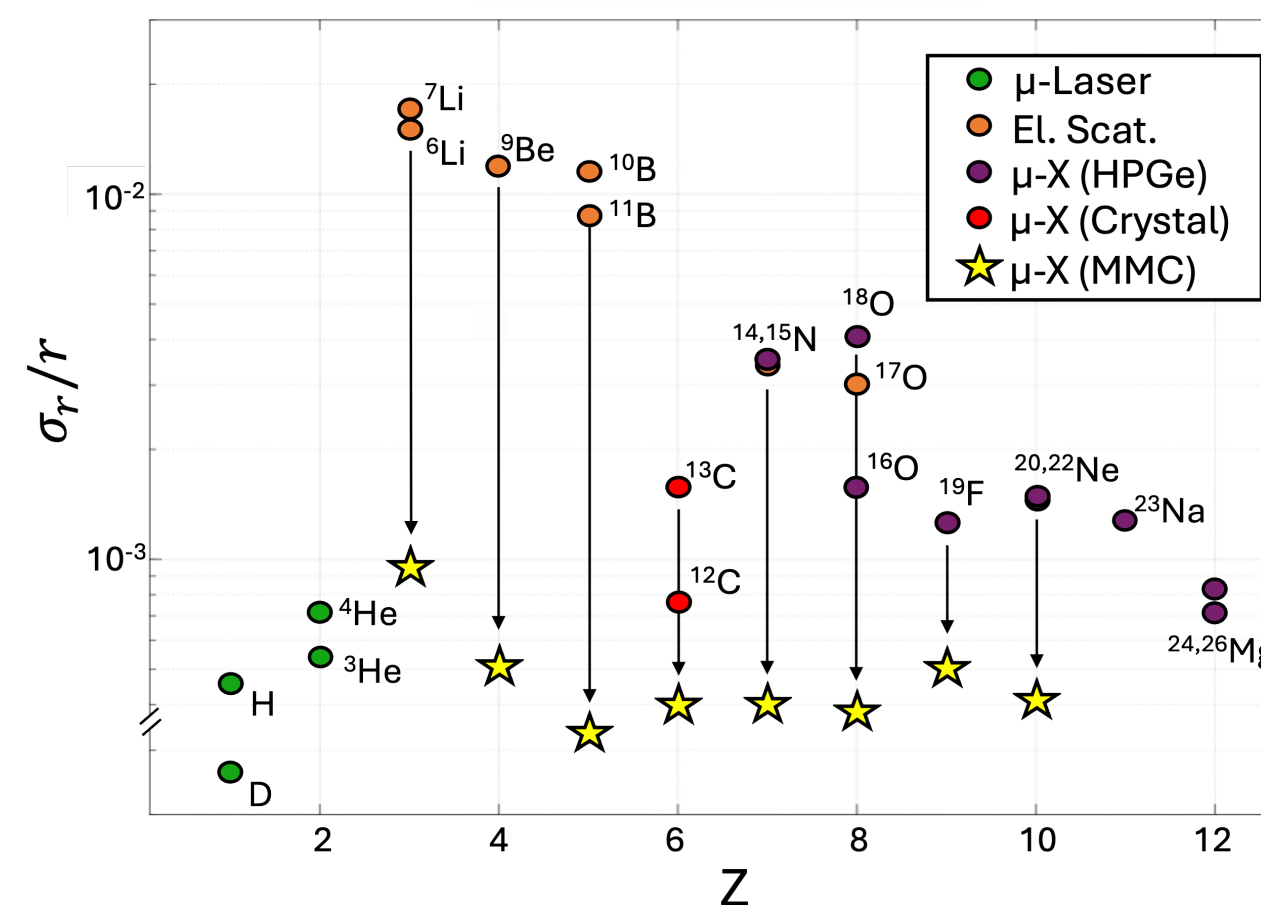
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# Conclusion

## Summary

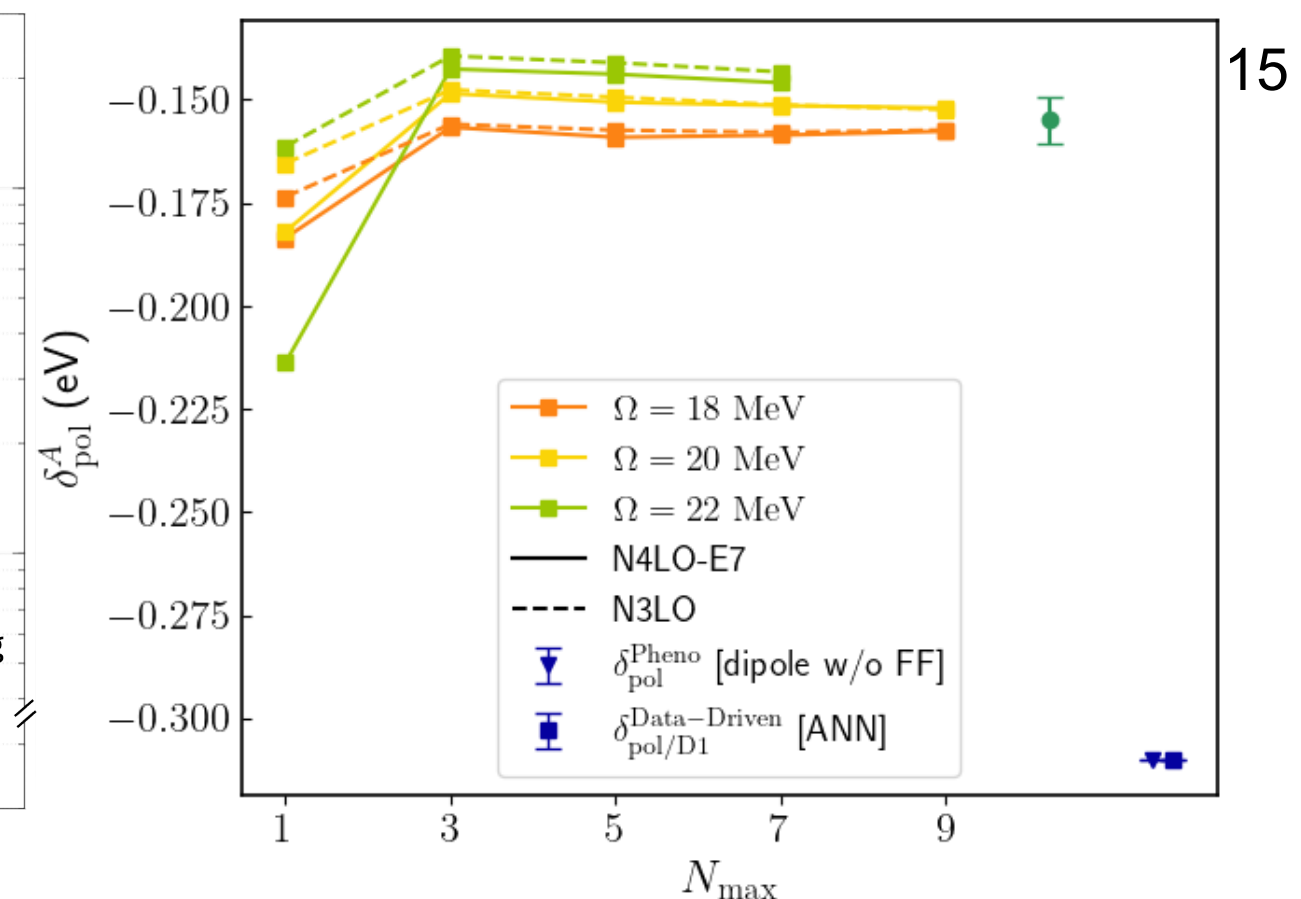
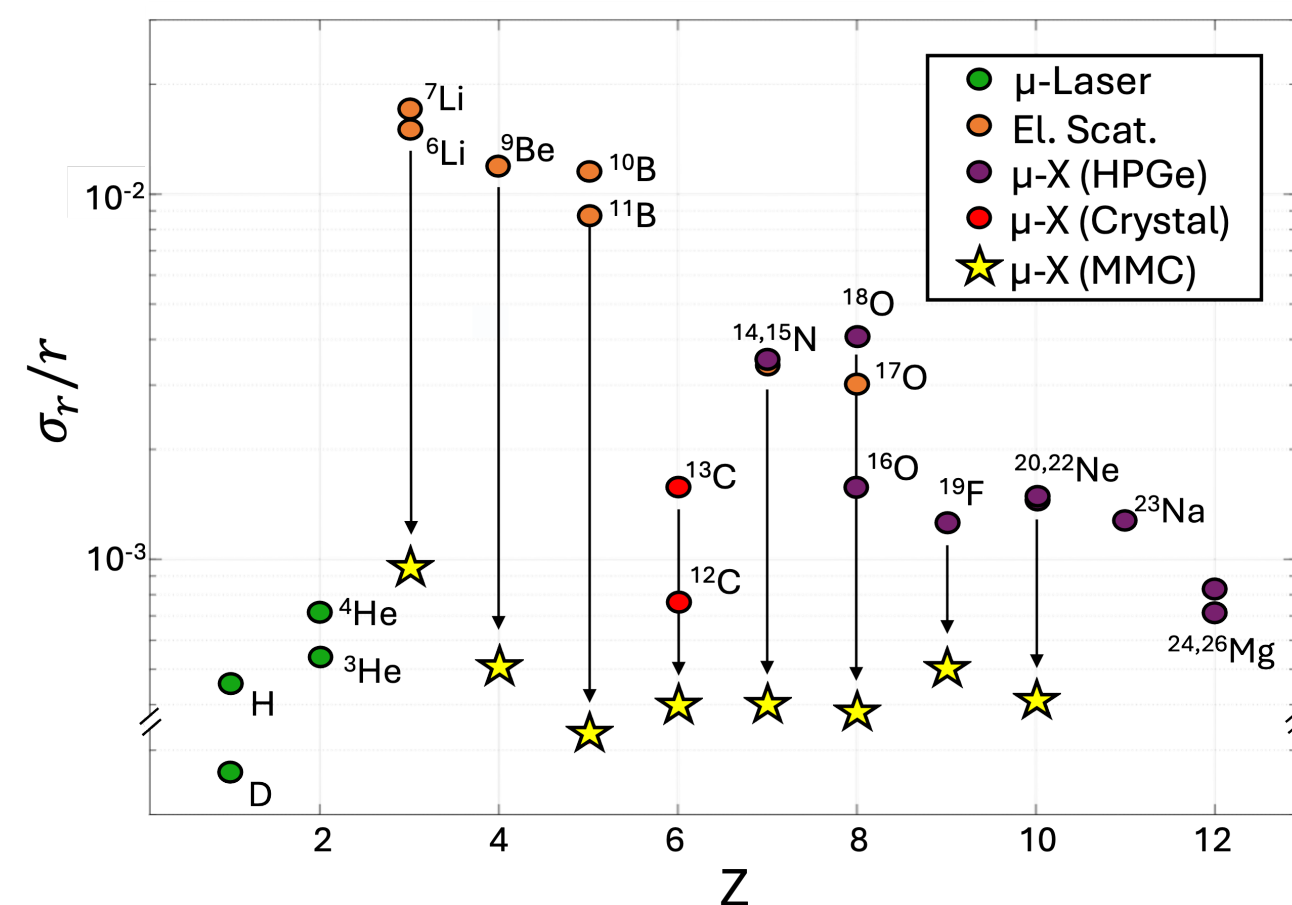
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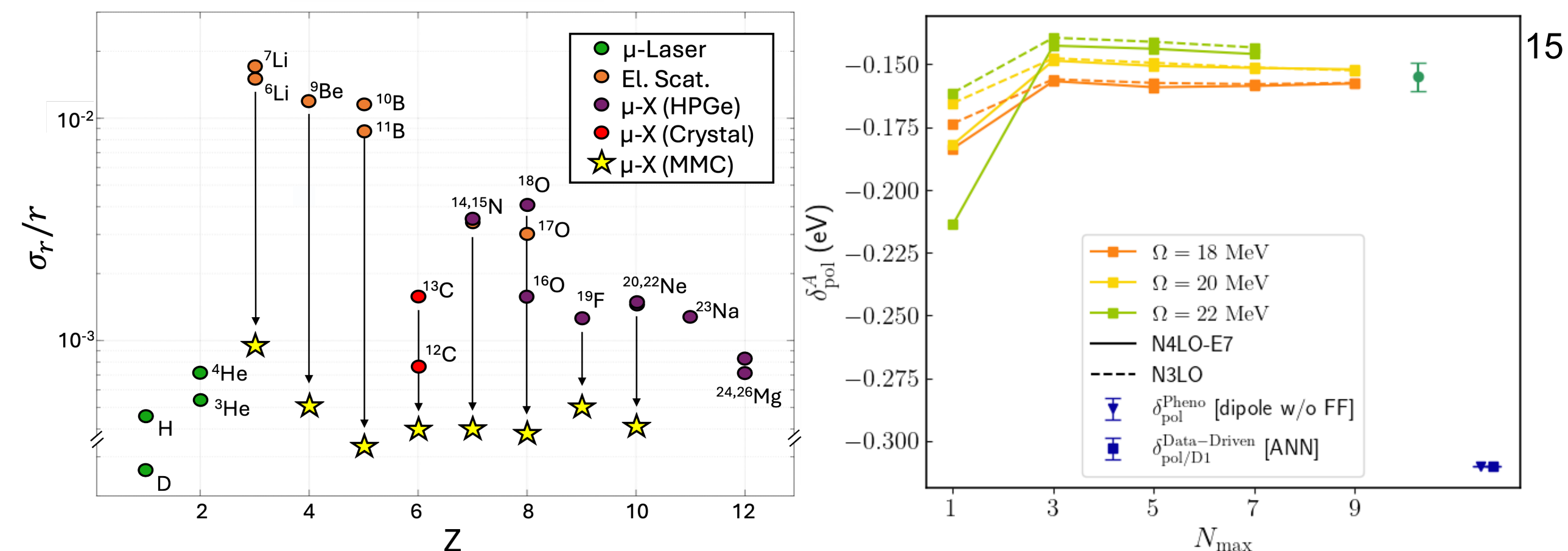
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- Completing on-going ab initio calculation
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  - Extension to  $^9\text{Be}$ ,  $^{10-11}\text{B}$   $\Rightarrow$  **See Diego's poster!**
- Towards better controlling theoretical uncertainty
  - Shifting from pheno towards EFT approach
  - EFT based on **potential-NRQED** for  $Z > 1$

[Peset et al., EPJA (2015)]

# THANK YOU!

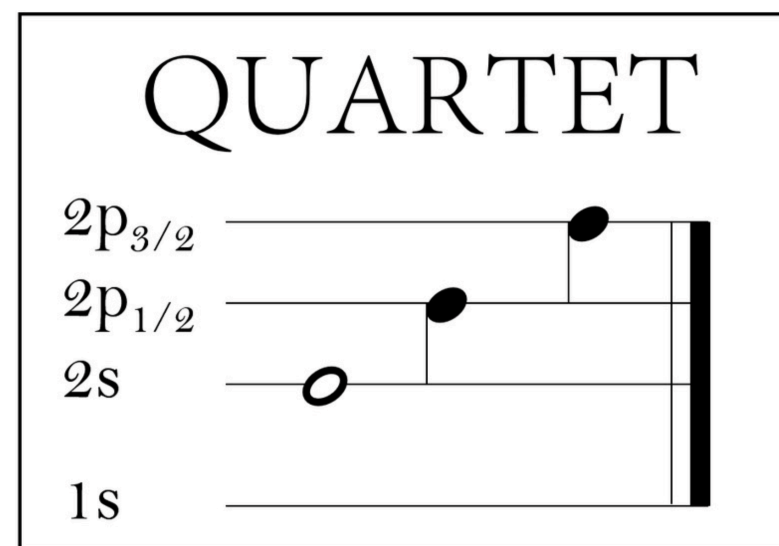


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UNIVERSITÄT  
DARMSTADT



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DARMSTADT

**STRONGINT Group**



**QUARTET Collaboration**



**Petr Navratil**  
**Diego Araujo Nájera**



JOHANNES GUTENBERG  
UNIVERSITÄT MAINZ

**Michael Gennari**      **Sonia Bacca**  
**Mikhail Gorchtein**      **Tim Egert**