



INPP

INSTITUTE OF NUCLEAR & PARTICLE PHYSICS

@OHIO UNIVERSITY



PAINT26

Ab initio optical potentials derived from SA-NCSM densities for elastic scattering from Magnesium isotopes

Charlotte Elster

Thanks to collaborators:

G. Sargsyan, J. Fuentealba, K. Beyer

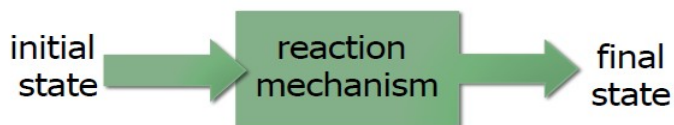
Supported by

Exploring Nuclei: Specifically Exotic nuclei are usually short lived

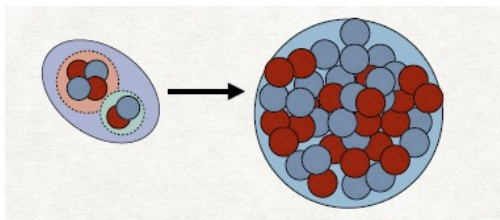
Thus one can only study them through reactions:

Have to be studied with reactions in inverse kinematics

direct reaction:



Many-body
problem



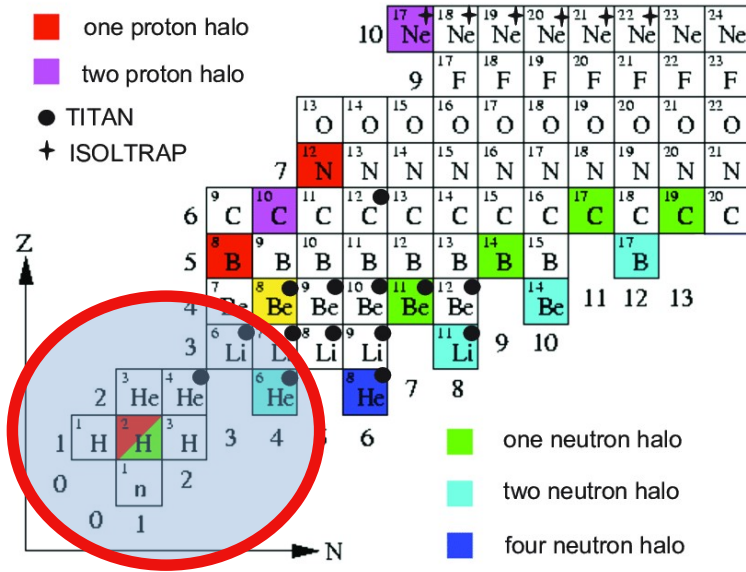
Quantum mechanical scattering problem

“idealist” approach: Just do it!



Not so fast!

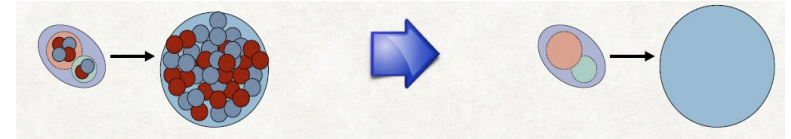
Chart of light nuclei



“Idealists” current domain

For heavier nuclei:

A more “pragmatic” approach is needed



Reduce the Many-Body to a Few-Body Problem

Isolate relevant degrees of freedom

Solve the few-body problem with effective interactions

Effective interactions from *ab initio* methods

Start from many-body Hamiltonian with 2 (and 3) body forces

Theoretical foundations laid by Feshbach and Watson in the 1950s

Feshbach:

effective nA interaction via Green's function from solution of many body problem using basis function expansion, e.g. SCGF, CCGF (current truncation to singles and doubles)

energy ~ 10 MeV

Watson:

- ▶ Multiple scattering expansion, e.g. spectator expansion
(current truncation to two active particles)

Spectator Expansion:

Siciliano, Thaler (1977)

Picklesimer, Thaler (1981)

Chinn, Elster, Thaler, Weppner
(1995)

Expansion in:

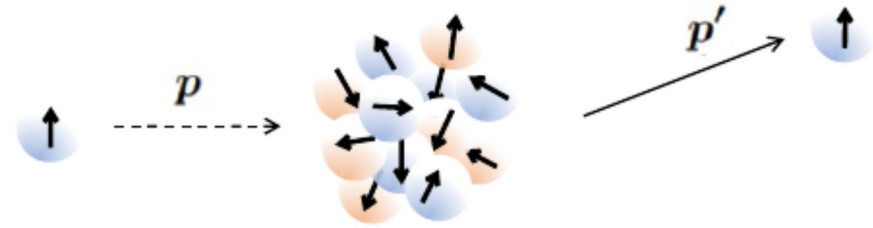
- ◆ *particles active in the reaction*
- ◆ *antisymmetrized in active particles*

Intended for "fast reaction", i.e. $\gtrsim 80$ MeV

Building up an *ab initio* framework for elastic scattering from nuclei

Ingredients:

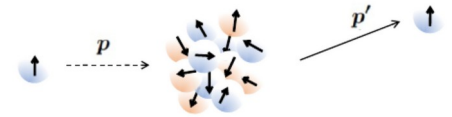
- 1) realistic nuclear interaction
must describe target and projectile
- 2) controllable structure framework
- 3) controllable reaction framework
- 4) a way to connect everything



	Two-nucleon force
LO (Q^0)	Weinberg '90
NLO (Q^2)	Ordonez, van Kolck '92
N ² LO (Q^3)	Ordonez, van Kolck '92
N ³ LO (Q^4)	Kaiser '00-'02
N ⁴ LO (Q^5)	Entem, Kaiser, Machleidt, Nosyk '15 Epelbaum, HK, Meißner '15

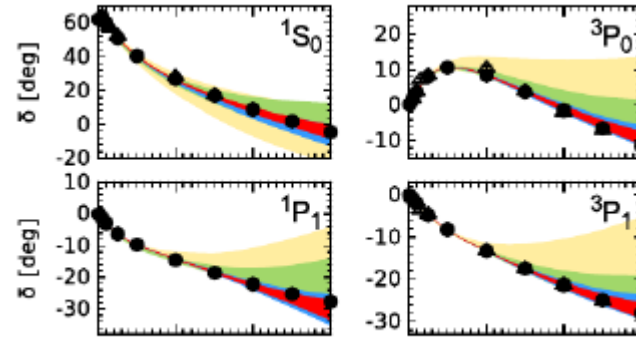
Chiral NN
Interaction

Building up an *ab initio* framework



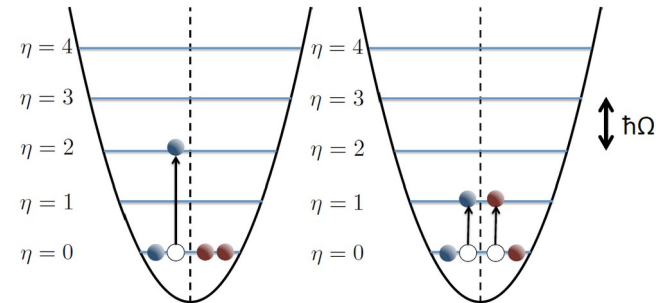
Ingredients:

- 1) realistic nuclear interaction
must describe target and projectile
- 2) controllable structure framework
here, no-core shell model (NCSM and SA-NCSM)
- 3) controllable reaction framework
here, spectator expansion
- 4) a way to connect everything



	$E(^3\text{H})$	$E(^3\text{He})$	$E(^4\text{He})$	$r_p(^4\text{He})$
NNLO	-8.249	-7.501	-27.759	1.43(8)
NNLO+NNN	-8.469	-7.722	-28.417	1.43(8)
Experiment	-8.482	-7.717	-28.296	1.467(13)

Ekström et al., PRL 110, 192502 (2013)



Building up an *ab initio* framework

Ingredients:

1) realistic nuclear interaction

must describe target and projectile

2) controllable structure framework

here, no-core shell model (NCSM and SA-NCSM)

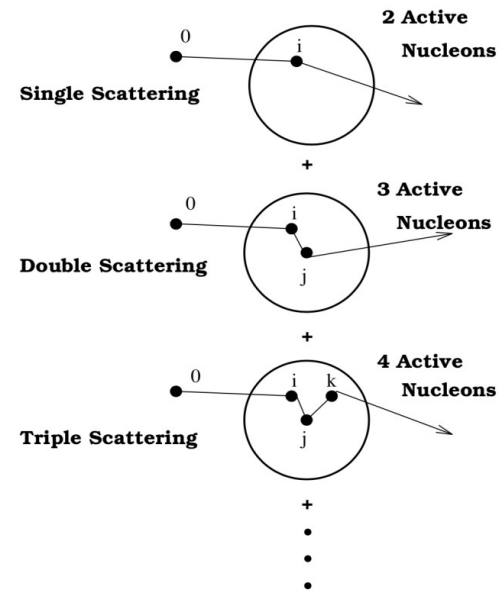
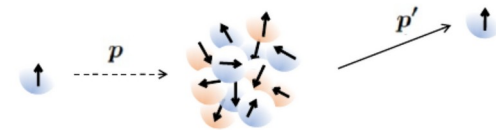
3) controllable reaction framework

here, spectator expansion of
Watson multiple scattering series

4) a way to connect everything

	$E(^3\text{H})$	$E(^3\text{He})$	$E(^4\text{He})$	$r_p(^4\text{He})$
NNLO	-8.249	-7.501	-27.759	1.43(8)
NNLO+NNN	-8.469	-7.722	-28.417	1.43(8)
Experiment	-8.482	-7.717	-28.296	1.467(13)

Ekström et al., PRL 110, 192502 (2013)



Building up an *ab initio* framework

Ingredients:

1) realistic nuclear interaction

must describe target and projectile

2) controllable structure framework

here, no-core shell model (NCSM and SA-NCSM)

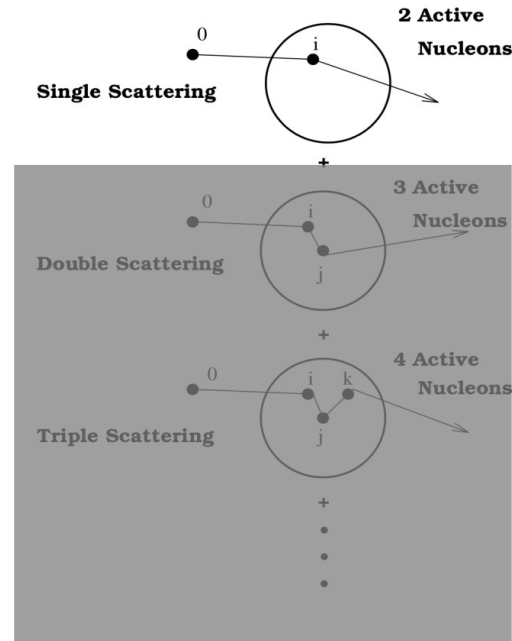
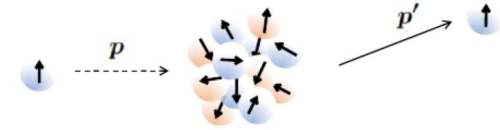
3) controllable reaction framework

here, spectator expansion of
Watson multiple scattering series

4) a way to connect everything

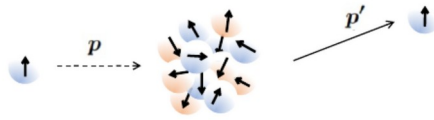
	$E(^3\text{H})$	$E(^3\text{He})$	$E(^4\text{He})$	$r_p(^4\text{He})$
NNLO	-8.249	-7.501	-27.759	1.43(8)
NNLO+NNN	-8.469	-7.722	-28.417	1.43(8)
Experiment	-8.482	-7.717	-28.296	1.467(13)

Ekström et al., PRL 110, 192502 (2013)



Leading order in
spectator
expansion can be
computed *ab initio*

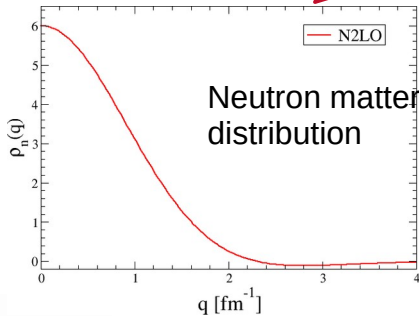
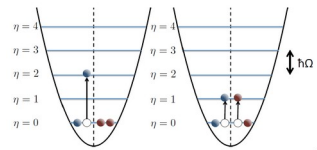
Framework for *ab initio* Elastic Scattering



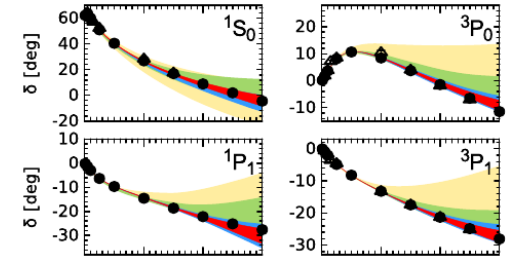
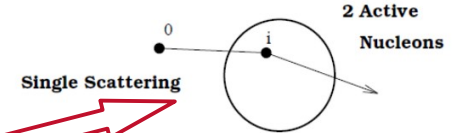
Same NN force in all parts

Reaction theory:
spectator expansion

Structure theory:
No-Core Shell Model



	Two-nucleon force
LO (Q^0)	Weinberg '90
NLO (Q^2)	Ordonez, van Kolck '92
N ² LO (Q^3)	Ordonez, van Kolck '92
N ³ LO (Q^4)	Kaiser '00-'02
N ⁴ LO (Q^5)	Entem, Kaiser, Machleidt, Nosyk '15 Epelbaum, HK, Meißner '15

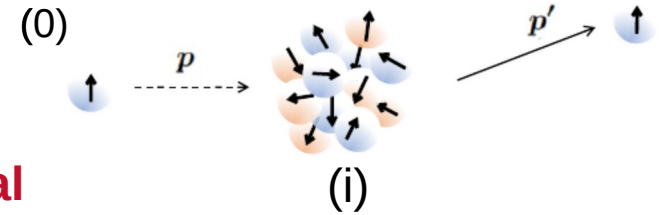


Calculations shown with NNLO_{opt} chiral interaction
A. Ekstrom et al. PRL 110, 192502 (2013)

Setting up the *ab initio* framework

$$\left(\begin{array}{c} \text{effective} \\ \text{interaction} \end{array} \right) = \left(\begin{array}{c} \text{thing that puts} \\ \text{them together} \end{array} \right) \times \left(\begin{array}{c} \text{reaction} \\ \text{information} \end{array} \right) \times \left(\begin{array}{c} \text{structure} \\ \text{information} \end{array} \right)$$

NN interaction represented by **Wolfenstein amplitudes**



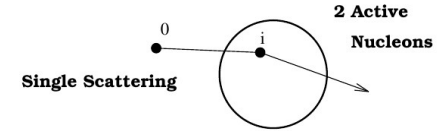
A: central
C: spin-orbit
M, G, H: tensor

$$\begin{aligned} \overline{M}(q, \mathcal{K}_{NN}, \epsilon) = & \underline{A(q, \mathcal{K}_{NN}, \epsilon)} \mathbf{1} \otimes \mathbf{1} \\ & + \underline{iC(q, \mathcal{K}_{NN}, \epsilon)} (\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{n}}) \otimes \mathbf{1} \\ & + \underline{iC(q, \mathcal{K}_{NN}, \epsilon)} \mathbf{1} \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{n}}) \\ & + \underline{M(q, \mathcal{K}_{NN}, \epsilon)} (\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{n}}) \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{n}}) \\ & + [G(q, \mathcal{K}_{NN}, \epsilon) - H(q, \mathcal{K}_{NN}, \epsilon)] (\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{q}}) \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{q}}) \\ & + [G(q, \mathcal{K}_{NN}, \epsilon) + H(q, \mathcal{K}_{NN}, \epsilon)] (\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{K}}) \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{K}}) \\ & + D(q, \mathcal{K}_{NN}, \epsilon) [(\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{q}}) \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{K}}) + (\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{K}}) \otimes (\boldsymbol{\sigma}^{(i)} \cdot \hat{\mathbf{q}})] \end{aligned}$$

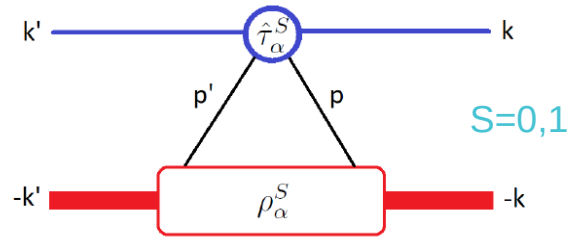
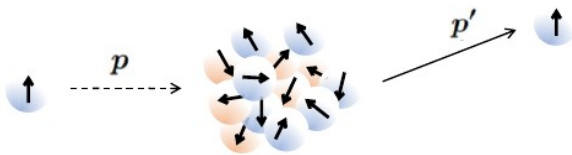
$$\begin{aligned} q &= p' - p \\ \mathcal{K} &= \frac{1}{2}(p' + p) \\ \hat{\mathbf{n}} &= \frac{\mathcal{K} \times q}{|\mathcal{K} \times q|} \end{aligned}$$

Most general structure of NN amplitudes consistent with invariance principles

Computing the leading order effective potential



$$\left(\text{effective interaction} \right) = \left(\text{thing that puts them together} \right) \times \left(\text{reaction information} \right) \times \left(\text{structure information} \right)$$



scatter off nucleus
with
 0^+ ground state

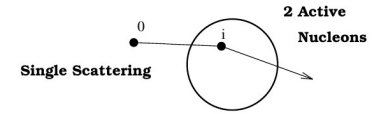
$$\hat{U}_p(q, \mathcal{K}_{NA}, \epsilon) = \sum_{\alpha=n,p} \sum_{S=0}^1 \int d^3\mathcal{K} \eta(q, \mathcal{K}, \mathcal{K}_{NA}) \hat{\tau}_{p,\alpha}^S \left(q, \frac{1}{2} \left(\frac{A+1}{A} \mathcal{K}_{NA} - \mathcal{K} \right); \epsilon \right) \rho_{\alpha}^S \left(\mathcal{K} - \frac{A-1}{A} \frac{q}{2}, \mathcal{K} + \frac{A-1}{A} \frac{q}{2} \right),$$

Watson optical potential:

$$U_p = \hat{U}_p - \hat{U}_p G_0(E) P U_p,$$

($N \neq Z$ nuclei treated exactly)

Computing the leading order effective potential



$$\left(\text{effective interaction} \right) = \left(\text{thing that puts them together} \right) \times \left(\text{reaction information} \right) \times \left(\text{structure information} \right)$$

$$\begin{aligned} \hat{U}_p(\mathbf{q}, \mathcal{K}_{NA}, \epsilon) &= \sum_{\alpha=p,n} \int d^3\mathcal{K} \eta(\mathbf{q}, \mathcal{K}_{NA}, \epsilon) A_{p,\alpha} \left(\mathbf{q}, \frac{1}{2} \left(\frac{A+1}{A} \mathcal{K}_{NA} + \mathcal{K} \right), \epsilon \right) \rho_{\alpha}^{S=0}(\mathbf{P}', \mathbf{P}) \\ &+ i(\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{n}}) \sum_{\alpha=N,Z} \int d^3\mathcal{K} \eta(\mathbf{q}, \mathcal{K}_{NA}, \epsilon) C_{p,\alpha} \left(\mathbf{q}, \frac{1}{2} \left(\frac{A+1}{A} \mathcal{K}_{NA} + \mathcal{K} \right), \epsilon \right) \rho_{\alpha}^{S=0}(\mathbf{P}', \mathbf{P}) \\ &+ i \sum_{\alpha=N,Z} \int d^3\mathcal{K} \eta(\mathbf{q}, \mathcal{K}_{NA}, \epsilon) C_{p,\alpha} \left(\mathbf{q}, \frac{1}{2} \left(\frac{A+1}{A} \mathcal{K}_{NA} + \mathcal{K} \right), \epsilon \right) S_{n,\alpha}(\mathbf{P}', \mathbf{P}) \cos \beta \\ &+ i(\boldsymbol{\sigma}^{(0)} \cdot \hat{\mathbf{n}}) \sum_{\alpha=N,Z} \int d^3\mathcal{K} \eta(\mathbf{q}, \mathcal{K}_{NA}, \epsilon) M_{p,\alpha} \left(\mathbf{q}, \frac{1}{2} \left(\frac{A+1}{A} \mathcal{K}_{NA} + \mathcal{K} \right), \epsilon \right) S_{n,\alpha}(\mathbf{P}', \mathbf{P}) \cos \beta \end{aligned}$$

matter distribution = density

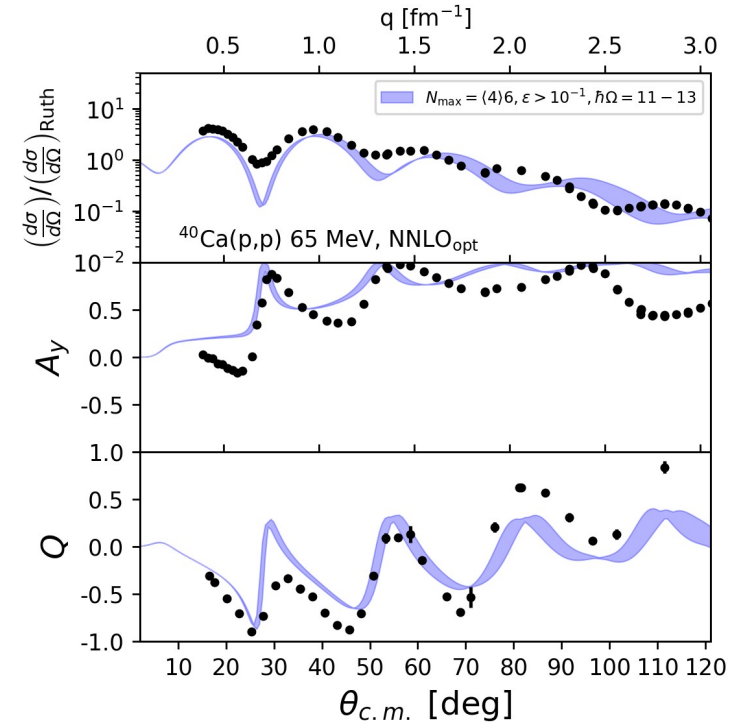
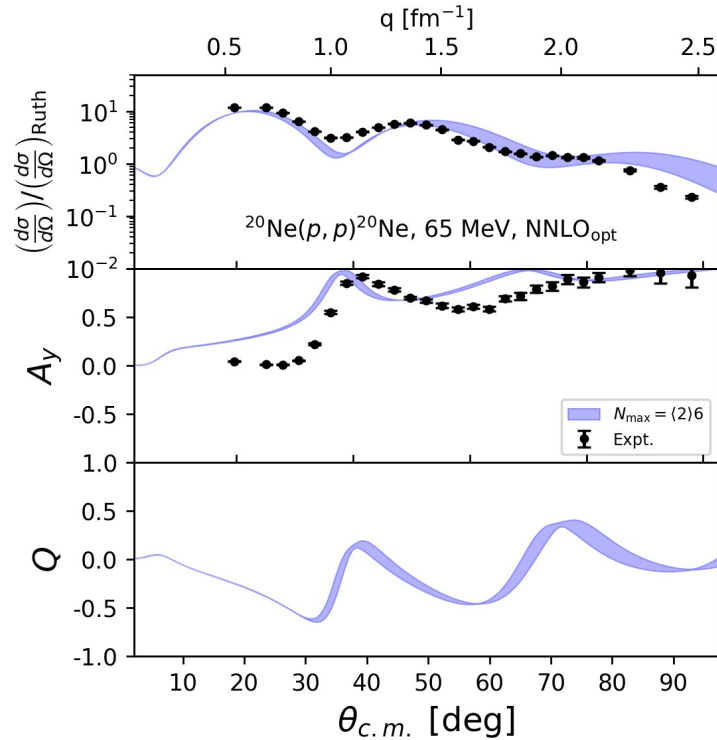
spin-projected momentum distribution

with $\mathcal{P}' = \left(\mathcal{K} - \frac{A-1}{A} \frac{\mathbf{q}}{2} \right)$ and $\mathcal{P} = \left(\mathcal{K} + \frac{A-1}{A} \frac{\mathbf{q}}{2} \right)$

Similar observation with SA-NCSM One-Body Densities

NNLO_{opt} chiral potential

$\hbar\omega=13-17$ MeV



$\hbar\omega=11-13$ MeV

Baker, Elster, Dytrych, Launey, PRC 110, 034605 (2024)

rms calc. 3.04 – 3.25 fm
rms exp. 3.48 fm

Ab initio leading order spectator seems to describe scattering data better for deformed nuclei

Why Magnesium Isotopes?




- Concentrate on Mg isotopes with 0^+ ground states
- Those isotopes are all deformed
- There is considerable experimental information on scattering from ^{24}Mg in our energy regime of interest
- ^{32}Mg of special interest (island of inversion)
- SA-NCSM can provide off-shell one-body densities for ^{24}Mg , ^{26}Mg , ^{28}Mg , ^{32}Mg

Why Magnesium Isotopes?

- Concentrate on Mg isotopes with 0+ ground states
- Those isotopes are deformed
- **There is considerable experimental information on scattering from ^{24}Mg** in our energy regime of interest
- ^{32}Mg of special interest (island of inversion)
- SA-NCSM can provide off-shell one-body densities for ^{24}Mg , ^{26}Mg , ^{28}Mg , ^{32}Mg

New development for phenomenology:

Uncertainty-quantified phenomenological optical potentials for single-nucleon scattering

C. D. Pruitt ^{*}, J. E. Escher , and R. Rahman [†]
Lawrence Livermore National Laboratory, Livermore, California 94550, USA

PHYSICAL REVIEW C **107**, 014602 (2023)

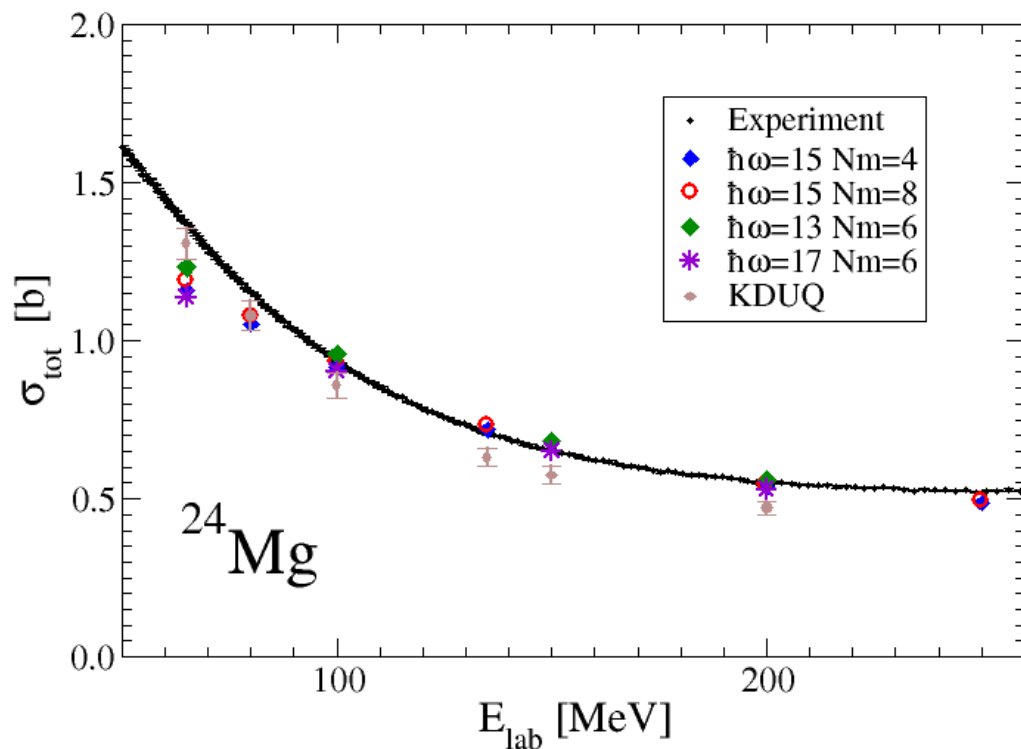
- Phenomenological Woods-Saxon based optical potentials fitted to stable nuclei
- **Extrapolation of phenomenological optical potentials with uncertainty quantification**
- **Ab initio is a theoretical well defined procedure for moving towards the drip-line with theory errors**



VALIDATION

Start and anchor on ^{24}Mg (experiment available)

Neutron total cross section



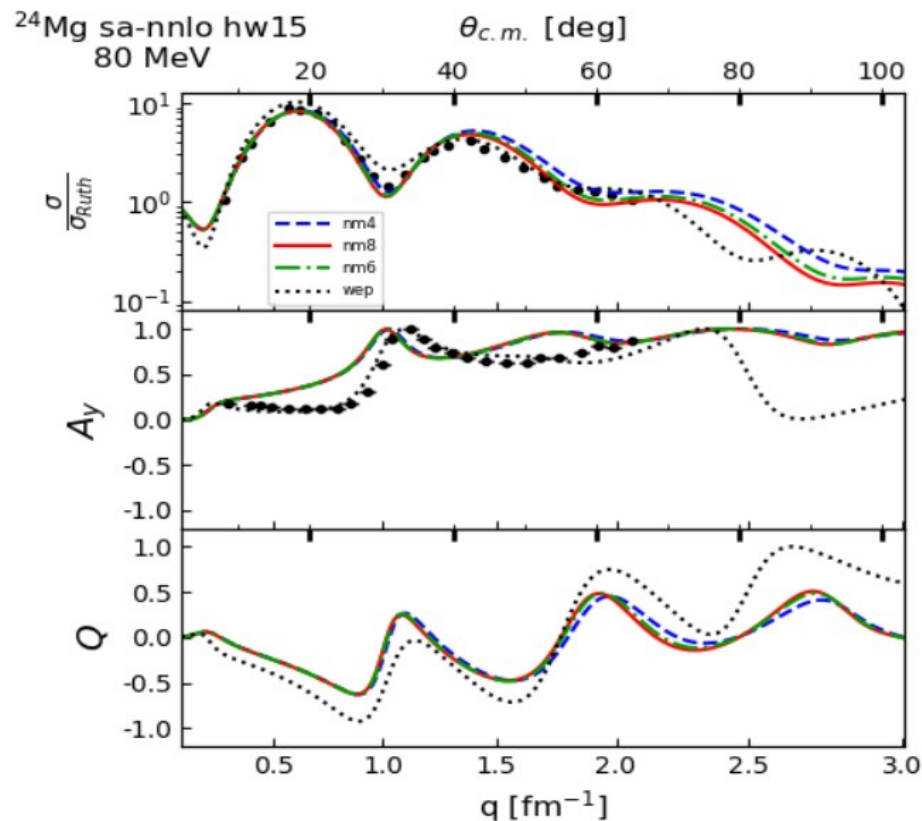
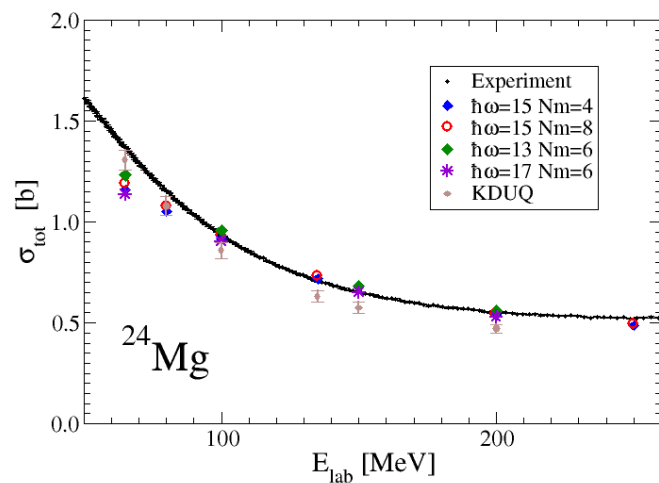
Ab initio calculation employs
NNLO_{opt} chiral potential

KDUQ: Koning-Delaroche
phenomenological optical potential
with Uncertainty Quantification



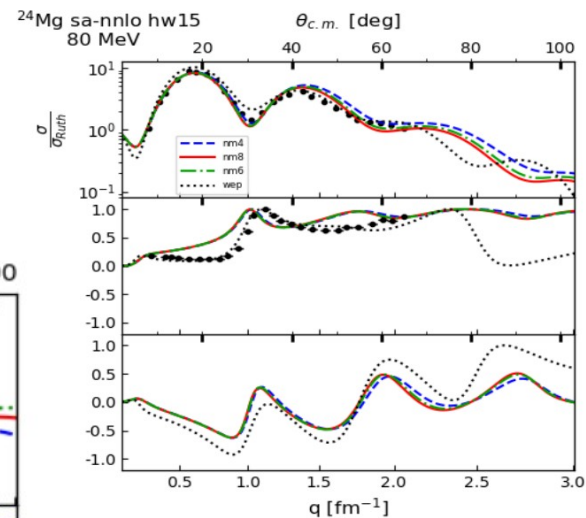
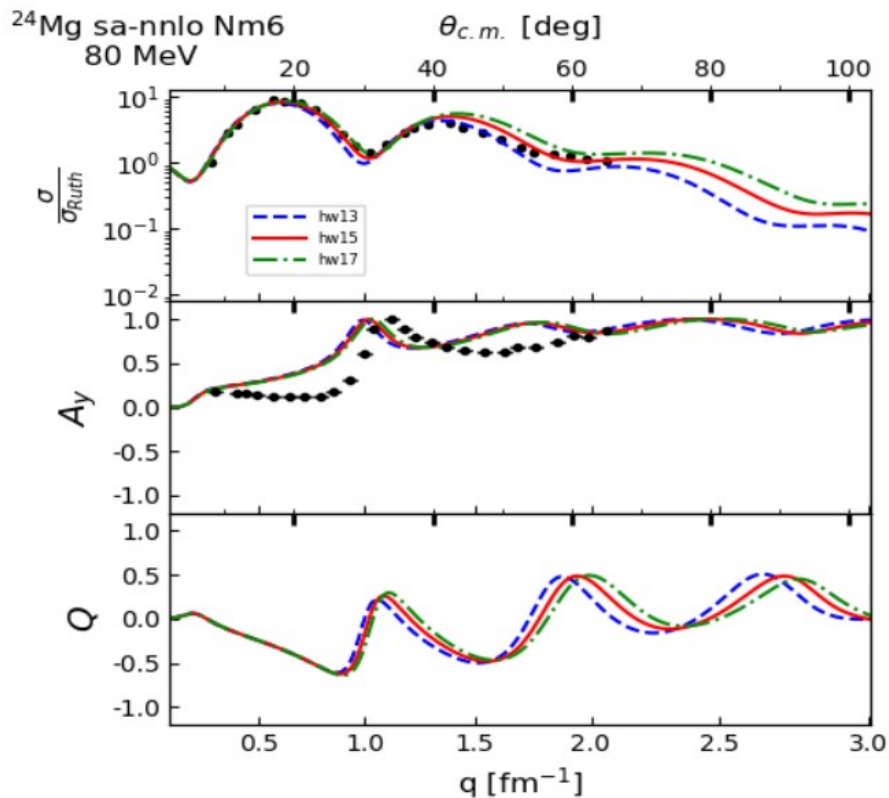
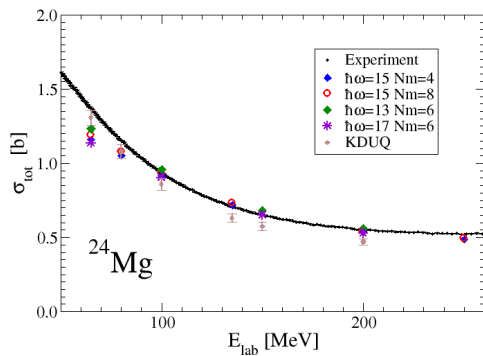
Start and anchor on ^{24}Mg

Neutron total cross section

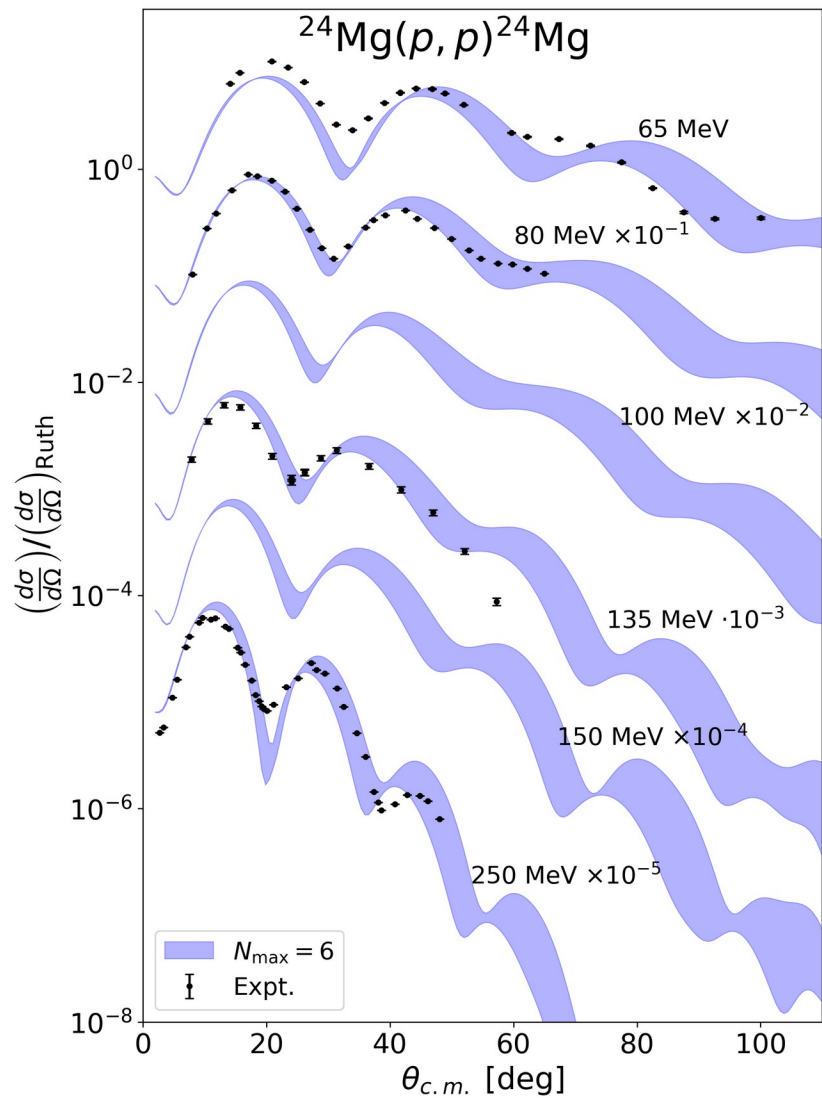


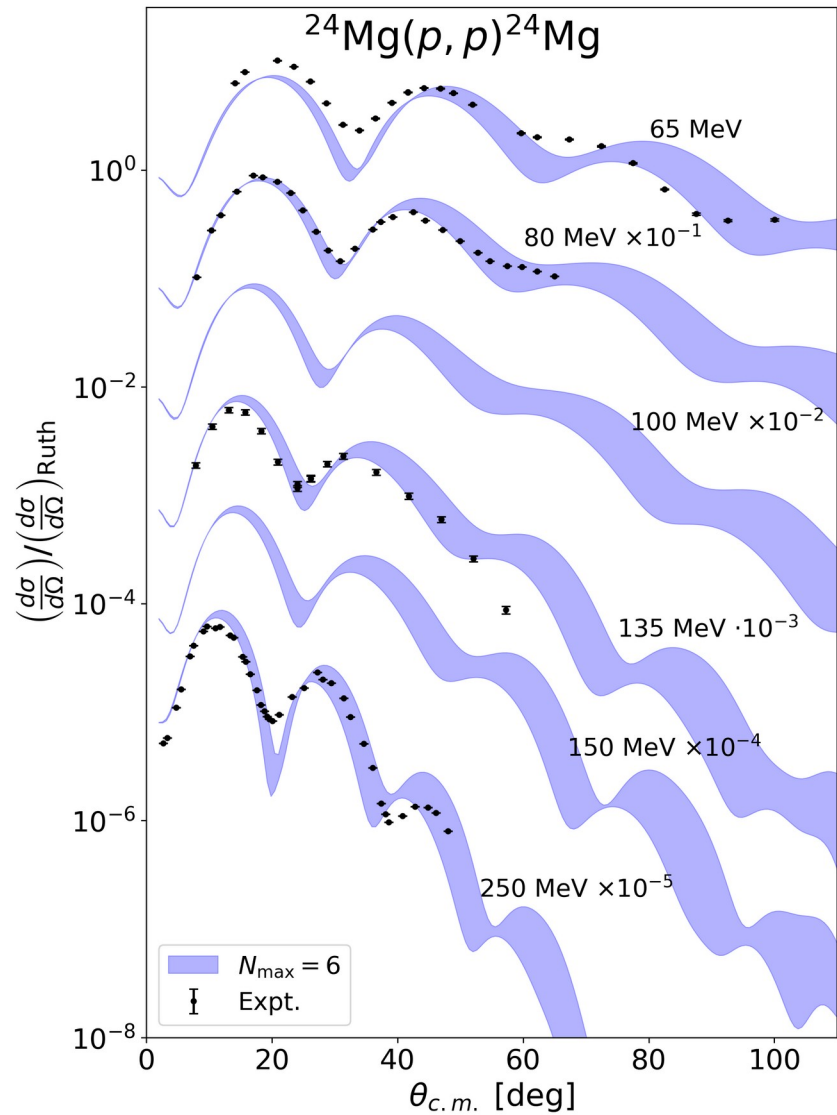
Start and anchor on ^{24}Mg

Neutron total cross section

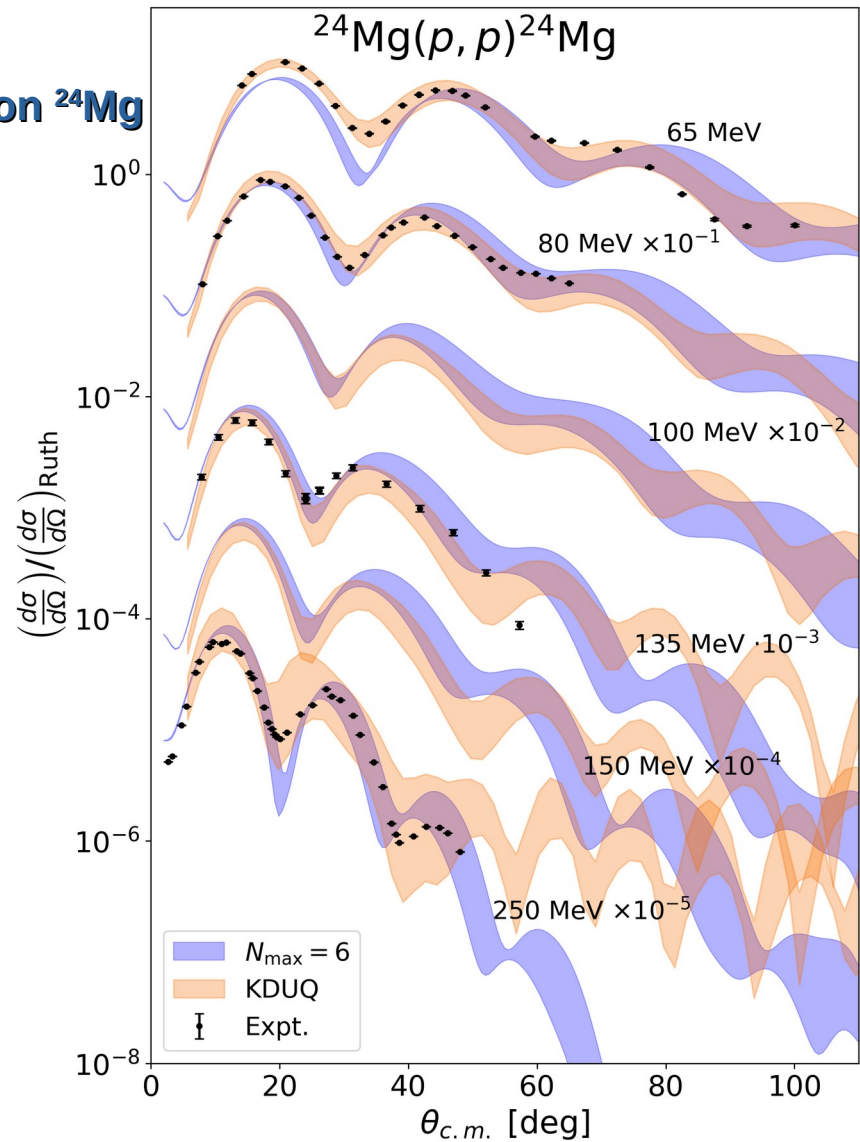


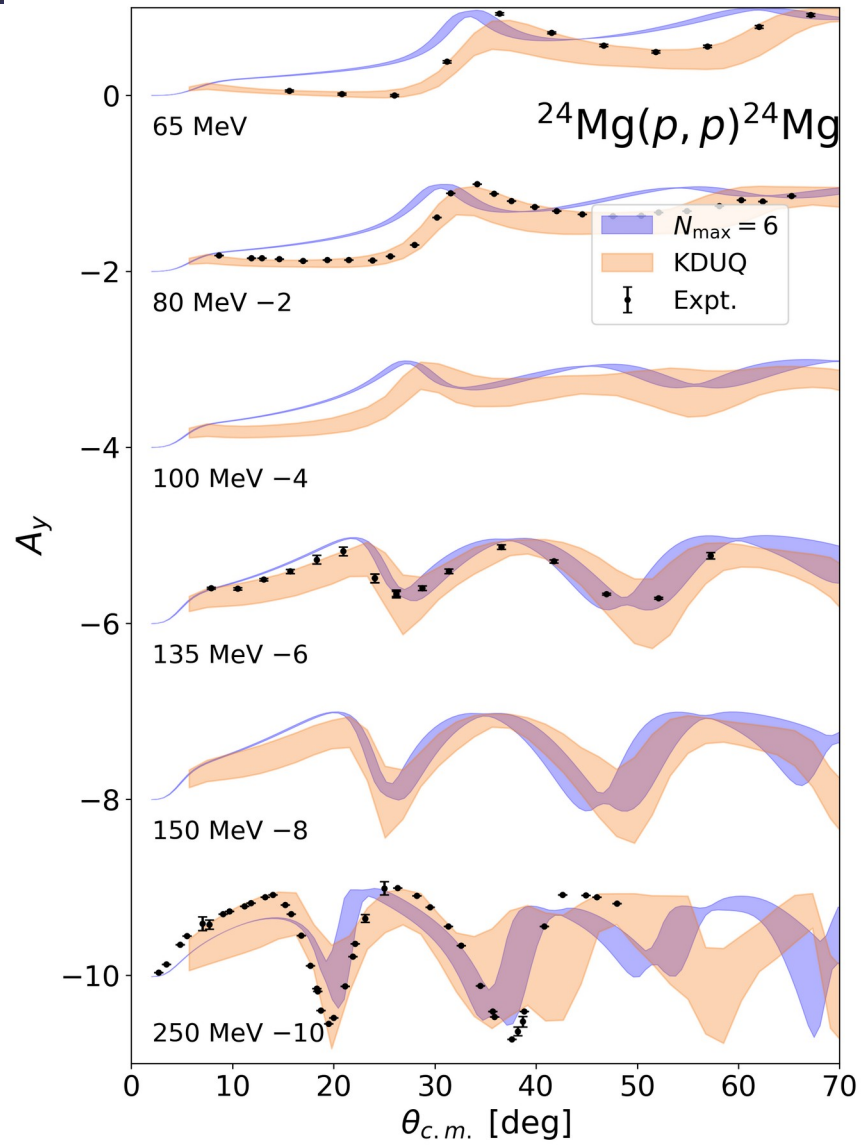
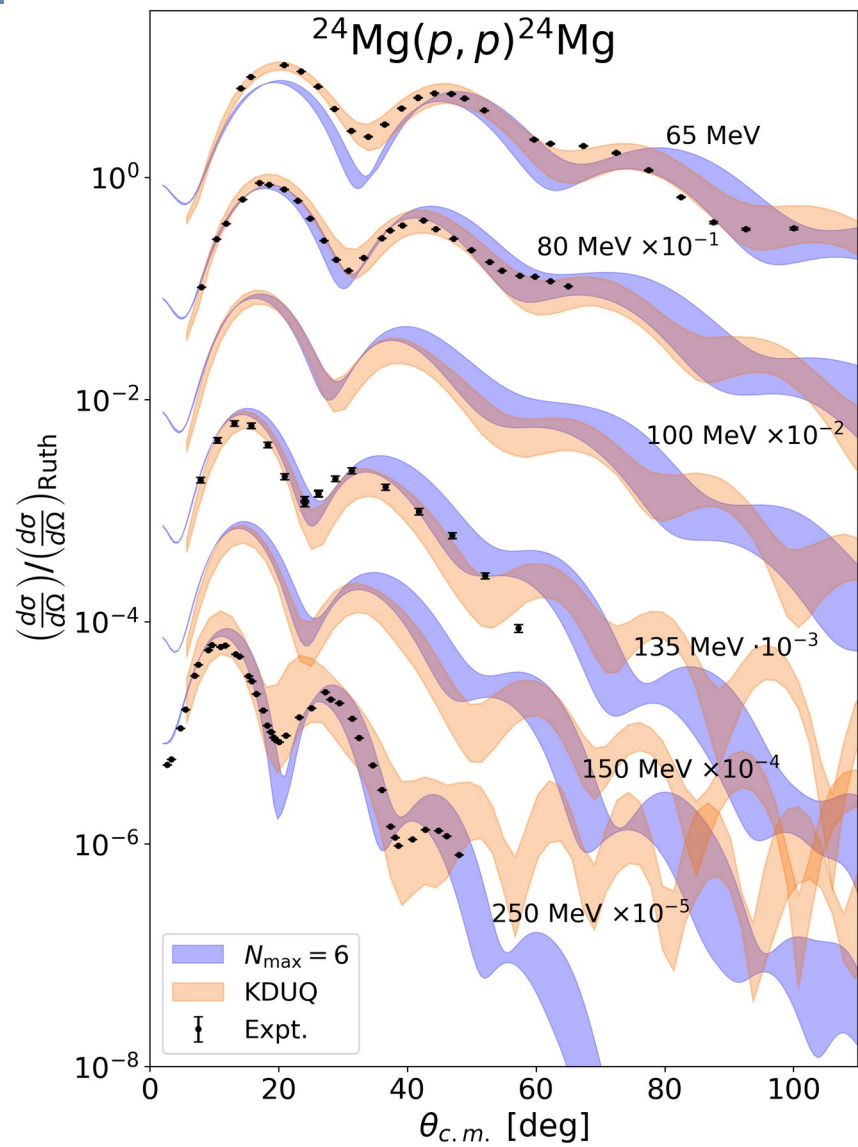
Start and anchor on ^{24}Mg



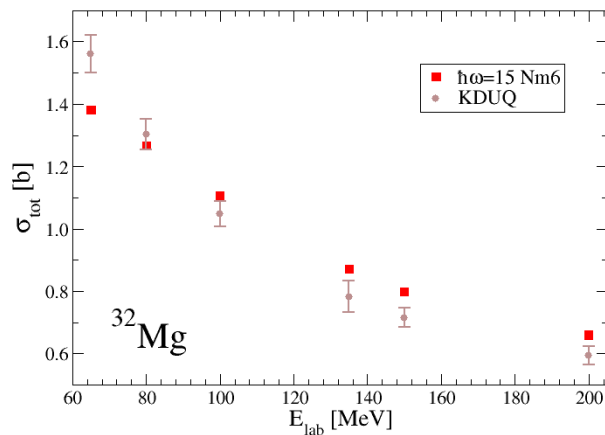
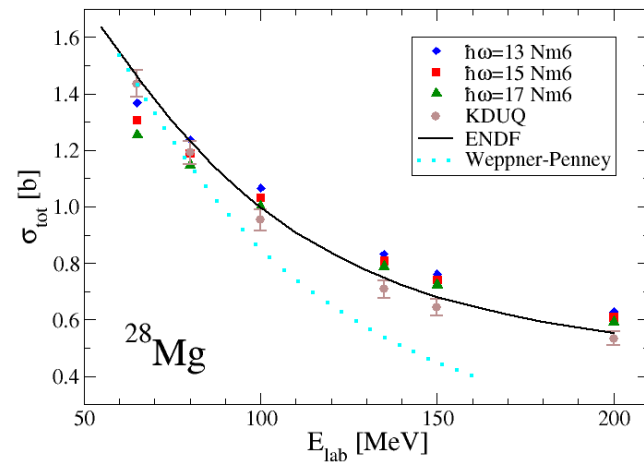
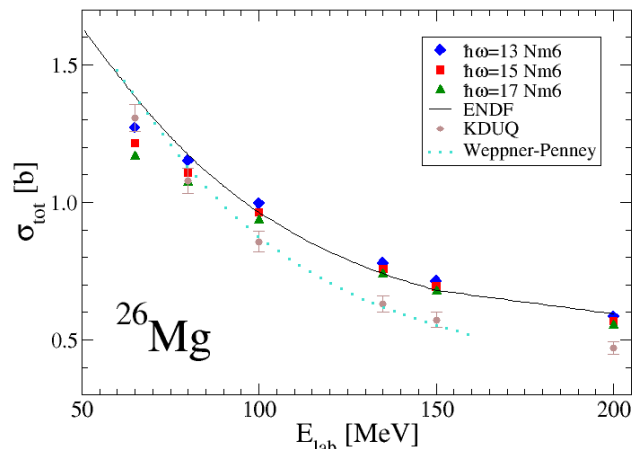
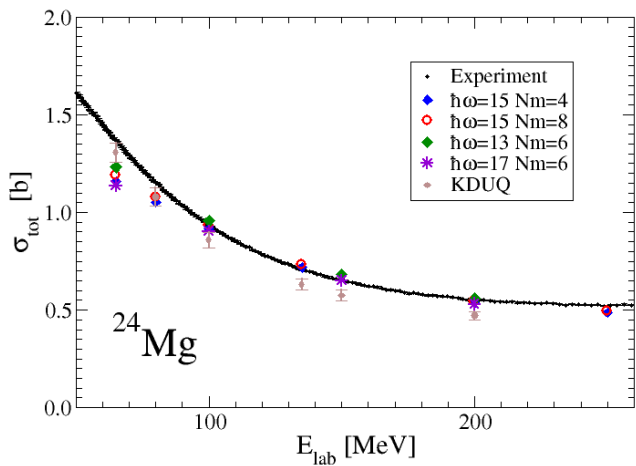


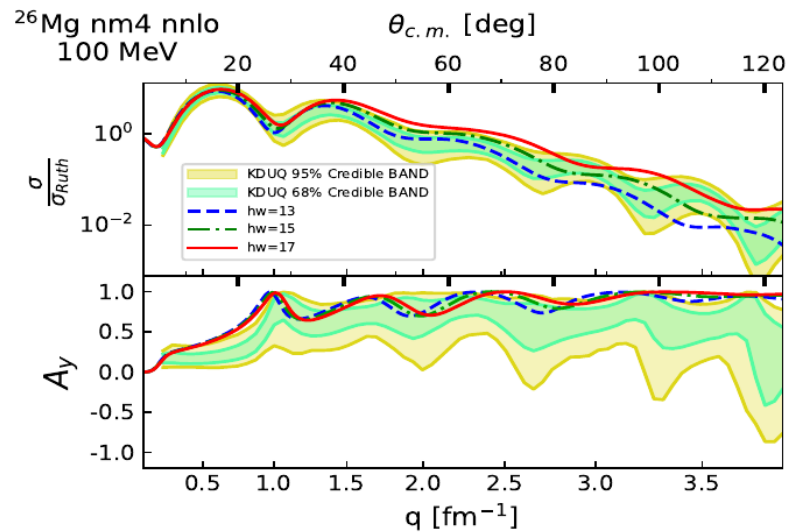
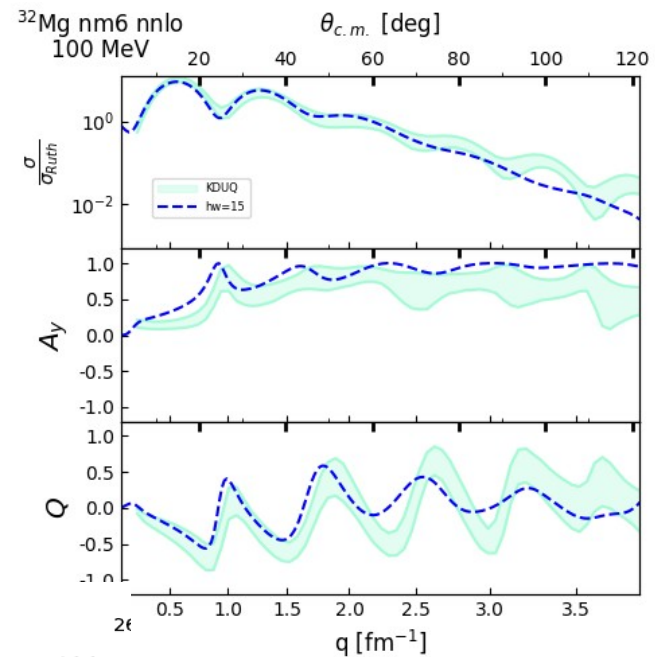
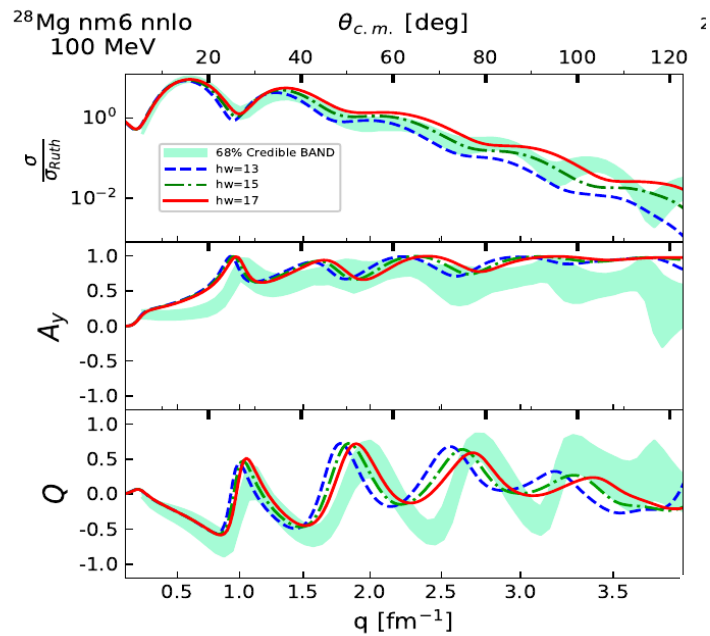
Start and anchor on ^{24}Mg





Neutron total cross sections for Mg isotopes





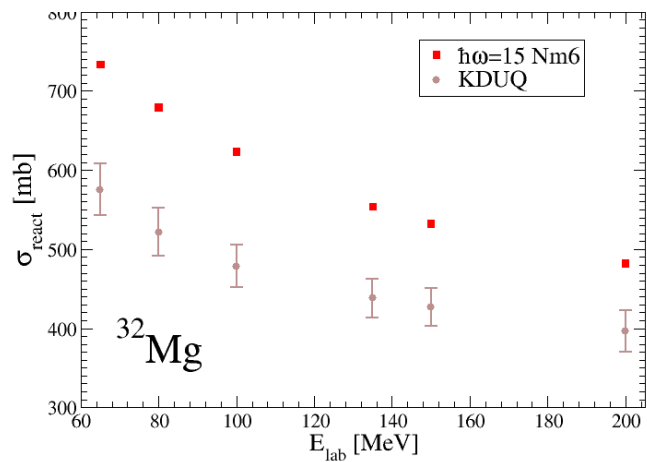
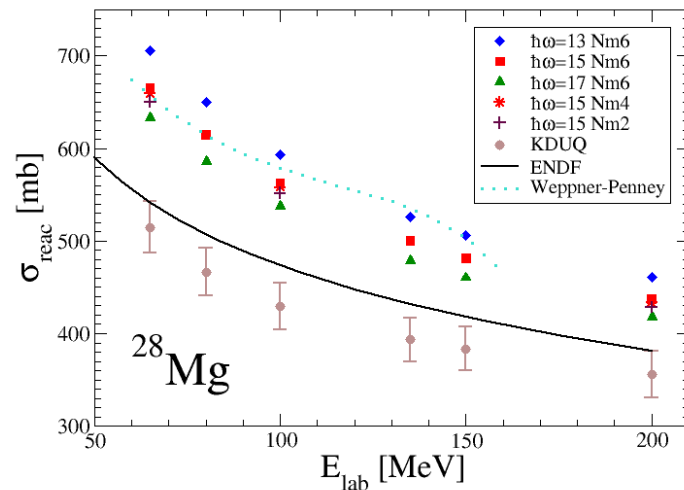
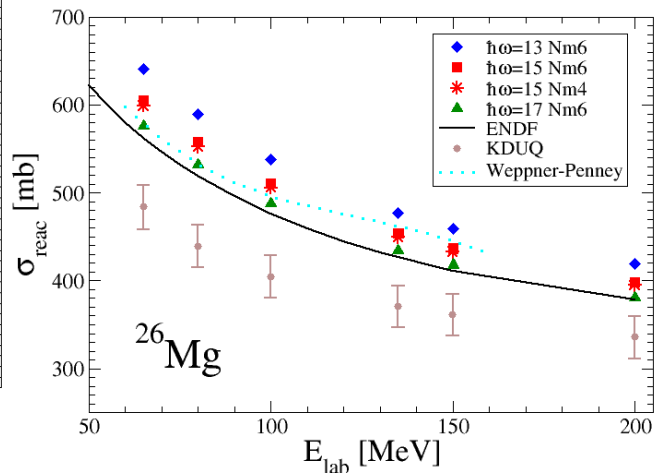
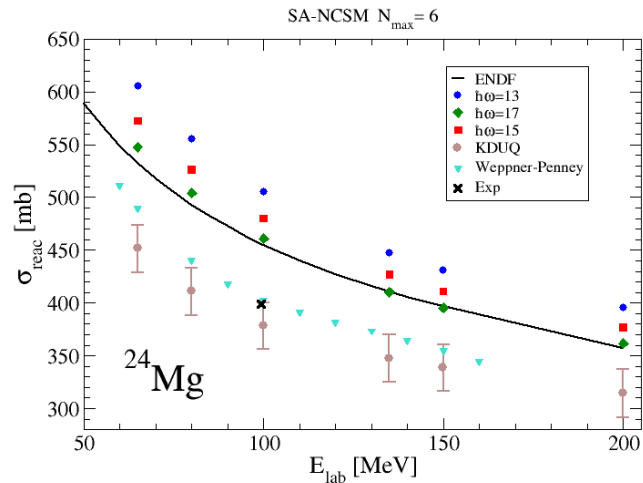
What we learned so far:

- **Consistent approach to p+A effective interaction in leading order multiple scattering expansion is possible.**
spin of projectile and struck target nucleon treated consistently
N \neq Z effects for Watson potential treated exactly
- Some indication that the leading order the spectator expansion describes elastic scattering data better for open-shell (deformed) and exotic nuclei than closed shell nuclei
(good for providing predictions for optical potential fits in exotic regime)
- **Carried out study of Mg isotopes:**
Calculations for ^{24}Mg compare favorably to experiment
Confidence in predicting ^{26}Mg , ^{28}Mg , and ^{32}Mg
- Test of extrapolations of phenomenological optical potentials towards the dripline
KDUQ extrapolates quite well (except reaction cross sections)

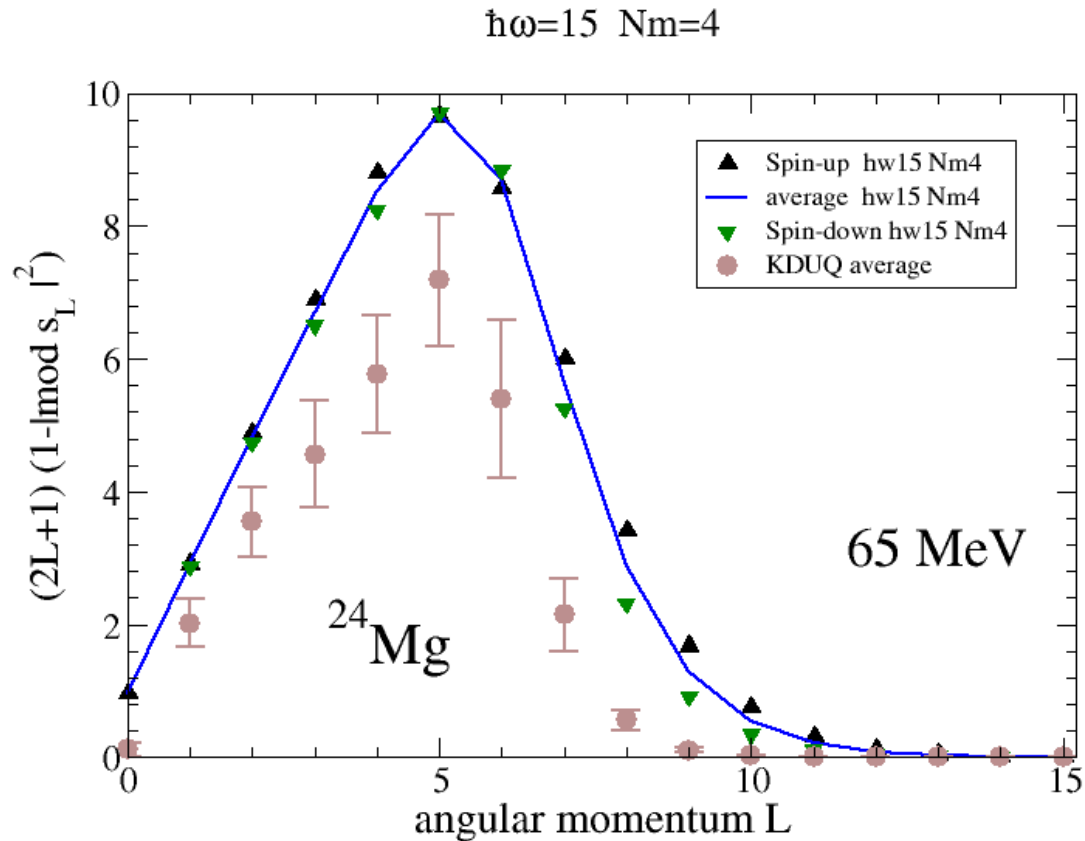


Backup slides

Proton reaction cross sections



Proton partial reaction cross sections



Diagonal densities (CoM)

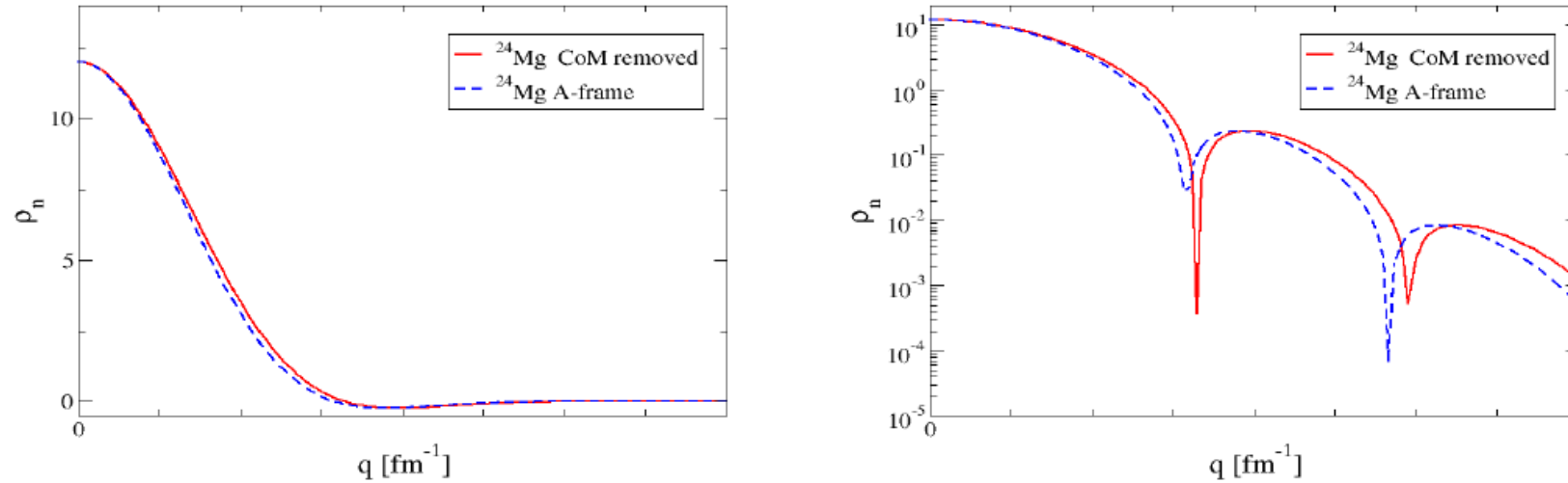


Figure 3: The diagonal one-body neutron densities for ^{24}Mg for $\hbar\omega=15$ and $N_{\text{max}}=6$. Shown is the diagonal density for a calculation where the CoM is removed vs a calculation in the A-frame.

Diagonal densities (neutron)

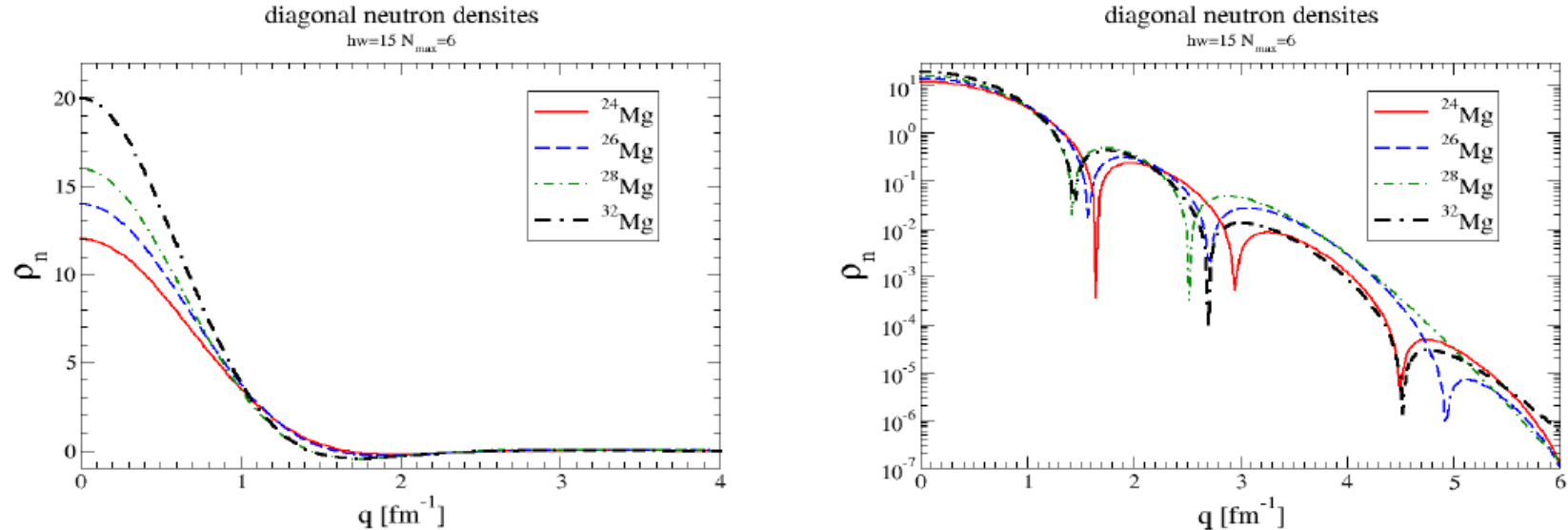


Figure 35: The diagonal neutron densities for $\hbar\omega=15$ and $N_{\max}=6$ for the Mg-isotopes we considered. All calculations are based on the NNLO_{opt} potential and SA-NCSM densities. They are plotted in linear as well as logarithmic scale.

Diagonal densities (proton)

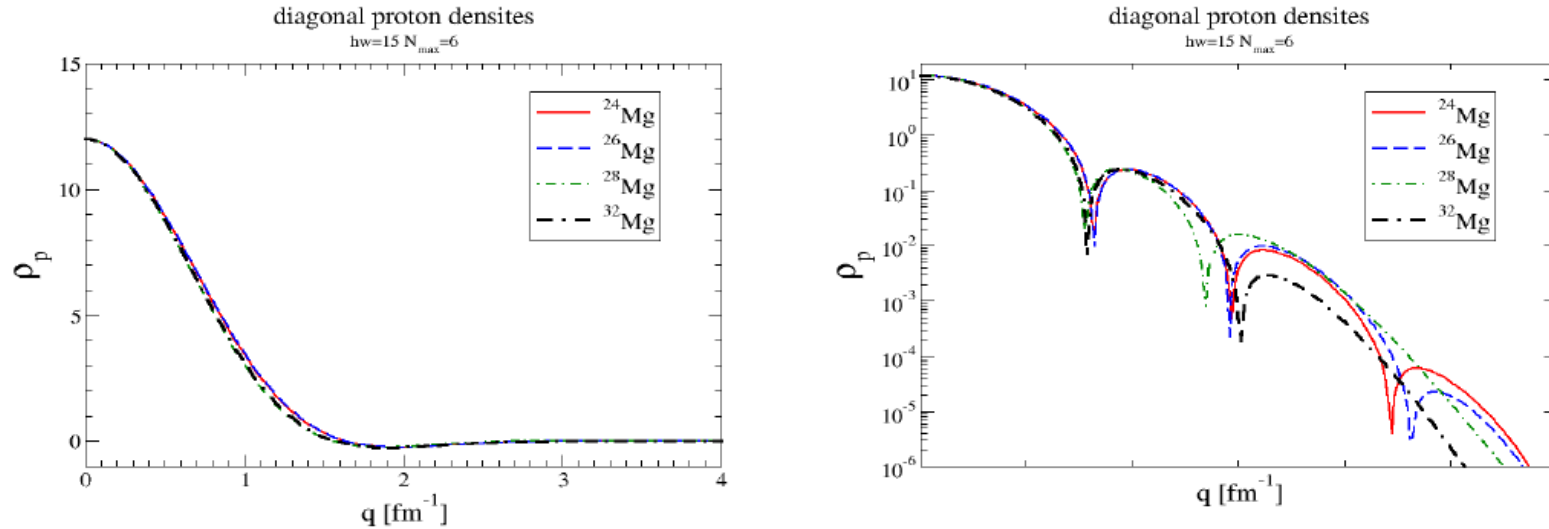


Figure 36: The diagonal neutron densities for $\hbar\omega=15$ and $N_{\max}=6$ for the Mg-isotopes we considered. All calculations are based on the NNLO_{opt} potential and SA-NCSM densities. They are plotted in linear as well as logarithmic scale.