

# Applications of *ab initio* nuclear theory to electroweak processes in atomic nuclei



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# Outline

1. Motivation
2. Challenges
3. Solution
4. Nuclear AIMs and Atomic AIMs
5. Summary



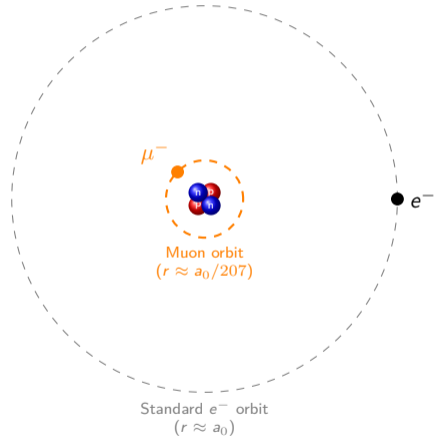
# Muonic Atom

## The Concept:

- An exotic atom where one (or more) electrons are replaced by a muon ( $\mu^-$ ).

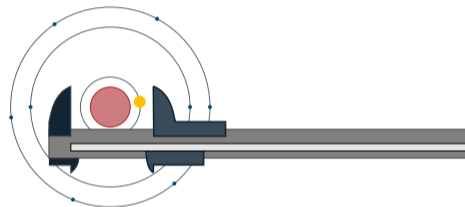
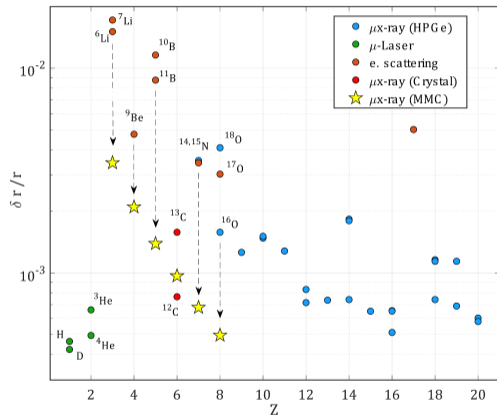
## Physical Implication:

- The muon orbits roughly 200 times closer to the nucleus.
- This makes muonic atoms incredible laboratories for probing nuclear structure and testing Quantum Electrodynamics (QED).





# QUARTET collaboration's experiment



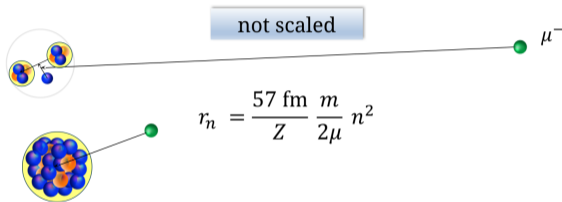
- aims to improve the measurement of absolute charge radii of light nuclei by a factor of 10



# Challenges

## A 2-scale problem

- Nuclear size  $\approx 10^{-5}$  a.u.
- Muonic 1s Bohr radius  $\approx 10^{-3}$  a.u.

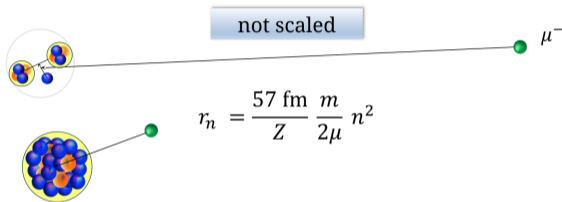




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Gauss' law must only be used with caution due to finite nuclear size

$\implies$  An accurate charge distribution extending up to several tens of fm is needed



## Proposed solution: No-Core Shell Model...



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- Drawback of NCSM: HO basis has an inherent Gaussian cutoff
- No-Core Shell Model with Continuum (NCSMC) provides a natural solution

$$\left| \Psi_{NCSMC}^{(A)} \right\rangle = \sum_{\lambda} c_{\lambda} |A\lambda J^{\pi} T\rangle + \sum_{\nu} \int dr \frac{\gamma_{\nu}(r)}{r} \hat{A}_{\nu} |\Phi_{\nu}\rangle$$

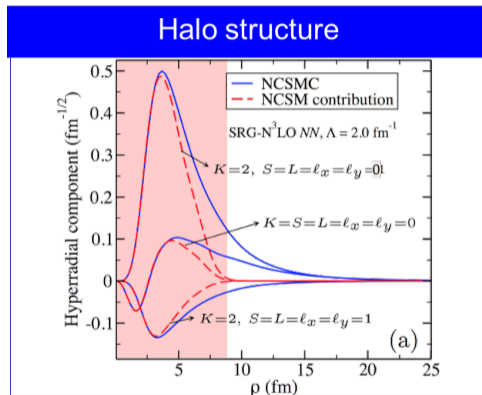
- $|A\lambda J^{\pi} T\rangle = \left| \text{single cluster} \right\rangle$ : NCSM basis
- $|\Phi_{\nu}\rangle = \left| \text{two clusters} \right\rangle$ : binary-cluster basis



## Example: ${}^6\text{He}$

A  $N_{\text{max}} = 12$  model space for NCSM corresponds to  $r \approx 8$  fm HO “spatial box”.

NCSMC extend much further providing correct asymptotics and internal part of the wave function.



*C. Romero-Redondo et al. (2016). Phys. Rev. Lett. 117*



# Example: light antiproton-nucleus system

## Light antiproton-nucleus systems at low energies with the *ab initio* NCSM/RGM method

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(Dated: February 11, 2026)

The availability of low-energy antiproton beams at the CERN Antiproton Decelerator has renewed interest in using antimatter as a probe of nuclear structure and in forming exotic antiprotonic few-body systems. In this work, we extend the *ab initio* No-Core Shell Model combined with the Resonating Group Method (NCSM/RGM), which was successfully applied to light-nucleus structure and reactions, to antiproton-nucleus dynamics at low energies. The NCSM/RGM formalism is adapted to antiproton projectiles by removing the requirement of antisymmetrization under exchange of target and projectile constituents, while retaining a fully microscopic description of the nuclear target and the relative motion. We focus on the lightest systems,  $\bar{p} + d$ ,  $\bar{p} + {}^3\text{H}$ , and  $\bar{p} + {}^3\text{He}$ , for which benchmarking against exact solutions of the Schrödinger equation enables stringent validation and helps disentangle methodological uncertainties—e.g., those associated with the choice of configurations included in the NCSM/RGM expansion—so that the dominant residual uncertainty can be attributed to the  $N\bar{N}$  interaction. We compute phase shifts, scattering lengths, cross sections, antiprotonic-atom level shifts and widths, nuclear quasi-bound energies, and annihilation densities. We find that the hard short-range components of the meson-exchange-based  $N\bar{N}$  interaction lead to slow convergence of the NCSM/RGM kernels expanded in a harmonic-oscillator basis, requiring exceptionally large model spaces and posing significant numerical challenges. We discuss practical strategies to mitigate these limitations and assess the impact of missing closed-channel configurations, which is a significant source of uncertainties in very light systems.

We can even compute atomic states for an antiprotonic system with 400 fm of channel radius.



# Nuclear AIMs and Atomic AIMs

- General matrix element of type

$$\langle \Psi_f^{(A)} | \hat{O} | \Psi_i^{(A)} \rangle$$

- Nuclear charge density operator

$$\hat{\rho}(\mathbf{r}) = e \sum_{i=1}^A \left( \frac{1}{2} - t_{i3} \right) \delta(\mathbf{r} - \mathbf{r}_i)$$



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- **Objective:** Collaboration with experts in Multi-Configuration Dirac-Fock methods [Cre: Paul Indelicato (LKB)] to use these nuclear calculations of charge and magnetic density as input

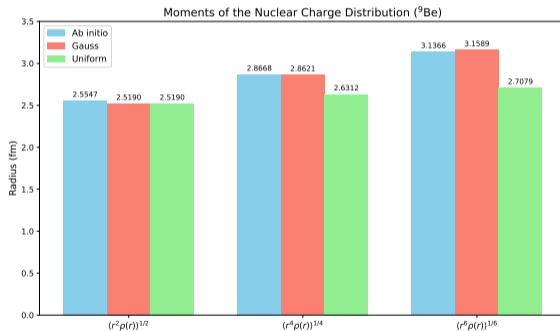


## Example: $^9\text{Be}$ muonic

Preliminary test calculations from P. Indelicato. The nuclear size effect is included in the Dirac equation potential.

1s binding energy (eV)

Ab initio	Gauss	Uniform
-43277.235	-44377.730	-44377.430





# Conclusion

## Summary:

- Muonic atom performs very well in probing nuclear structure
- NCSMC provides solutions that extend to large distances
- ...and can be used as MCDF input to achieve hyperfine splitting predictions

## Outlook:

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Thank you for your attention!