

Scintillator Detectors

Darren R Grant — June 2026



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SIMON FRASER
UNIVERSITY

Material References/Sources*

Scintillation Detectors - Particle Detection via Luminescence (W. Kolanoski — https://www.desy.de/~schleper/lehre/Det_Dat/SS_2018/02_Scintillators.pdf)

Scintillator Detectors (S. Leoni — <https://www0.mi.infn.it/~sleoni/TEACHING/Nuc-Phys-Det/PDF/Lezione-partI-6-scint-det.pdf>)

Introduction to Radiation Detectors and Electronics - III. Scintillation Detectors (H. Spieler, 1998 — https://www-physics.lbl.gov/~spieler/physics_198_notes/PDF/III-Scint.pdf)

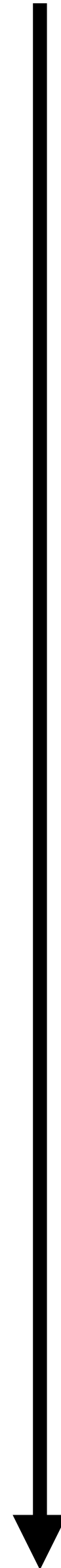
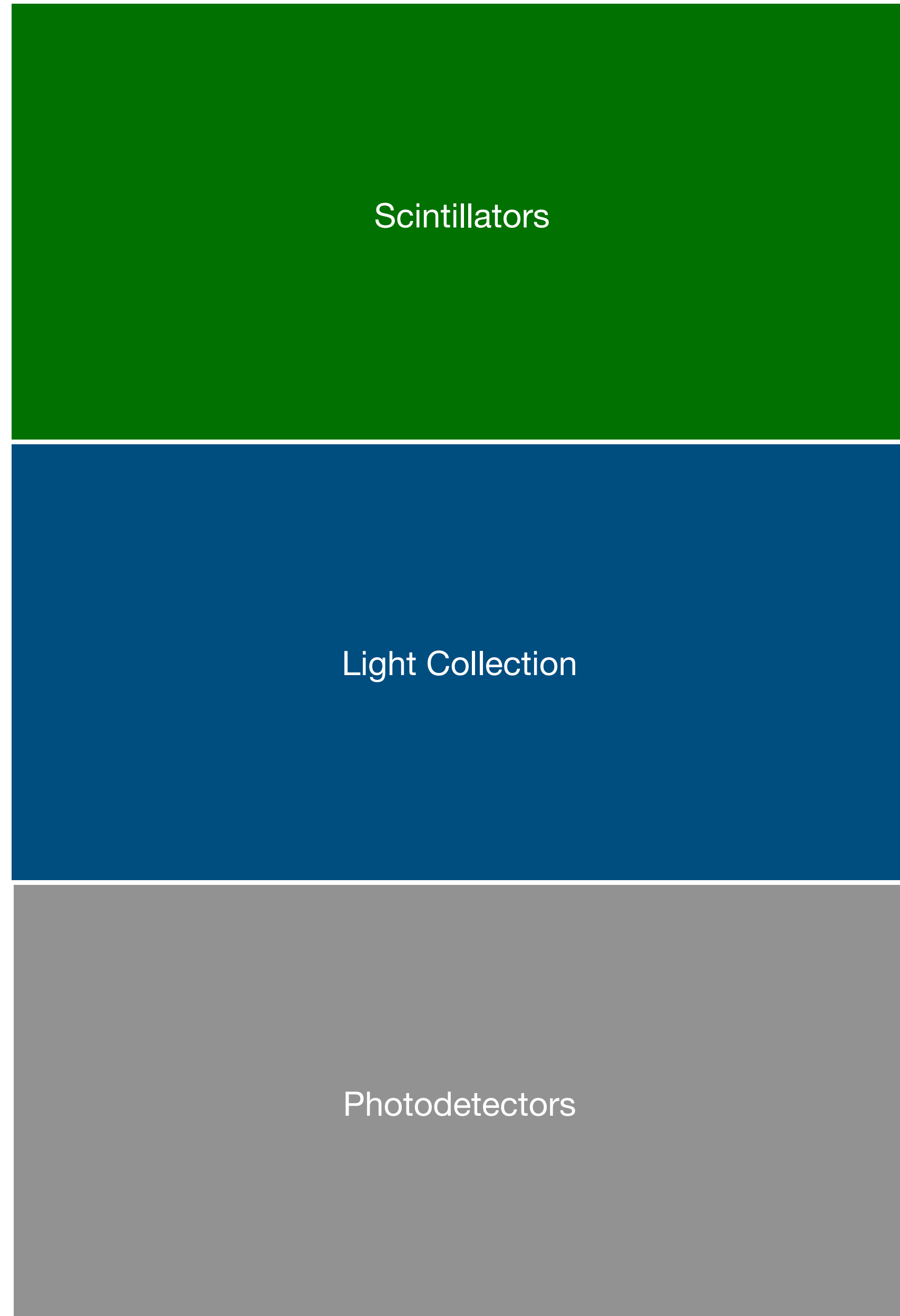
Scintillation Detectors and Their Applications (C. Woody, 2019 — <https://grids.triumf.ca/2019/lectures/Lecture7.pdf>)

Scintillation Detectors (S. Xella, 2009 — https://www.nbi.dk/~xella/lecture_16Feb2009.pdf)

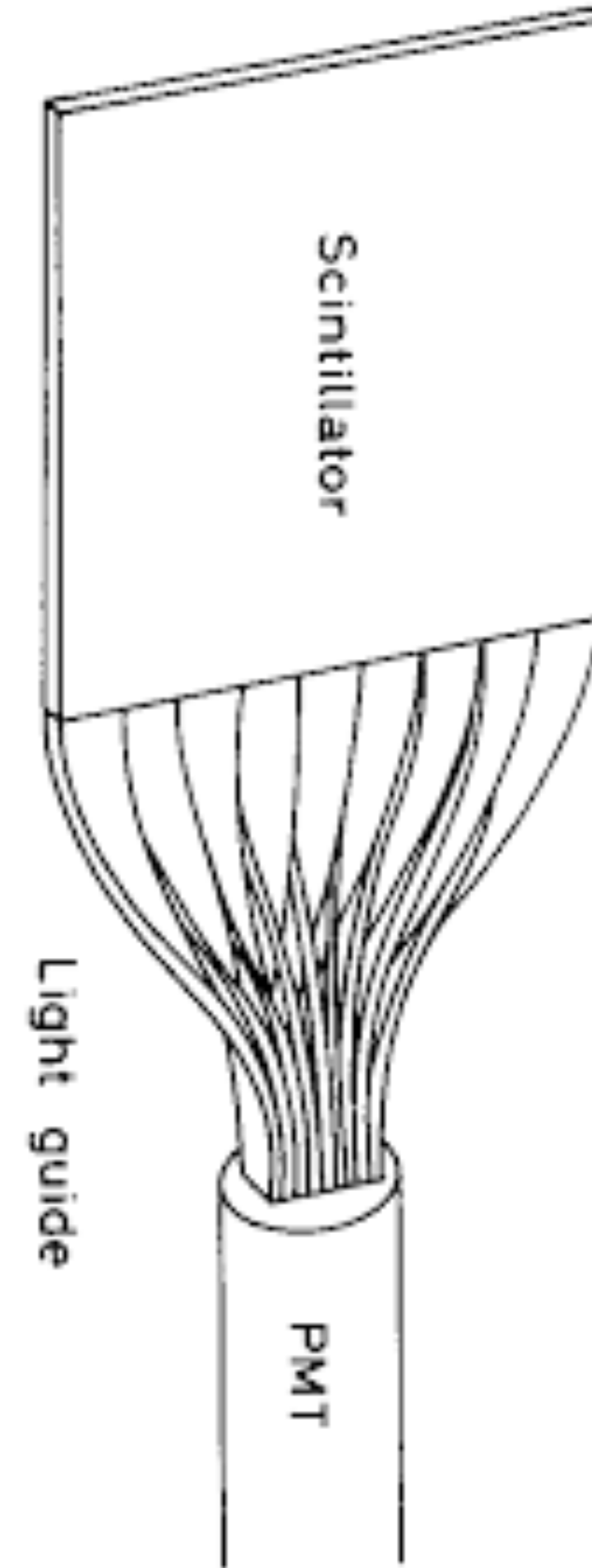
Scintillation Detectors (J. Yanez, 2025 — https://indico.triumf.ca/event/574/contributions/6081/attachments/4802/6978/scintillation_detectors_yanez_grids_2025.pdf)

Detector and detector systems for particle and nuclear physics I (J. Zmeskal, 2020 — https://www.oeaw.ac.at/fileadmin/Institute/SMI/PDF/Detectors_WS2020-21_part2.pdf)

Outline



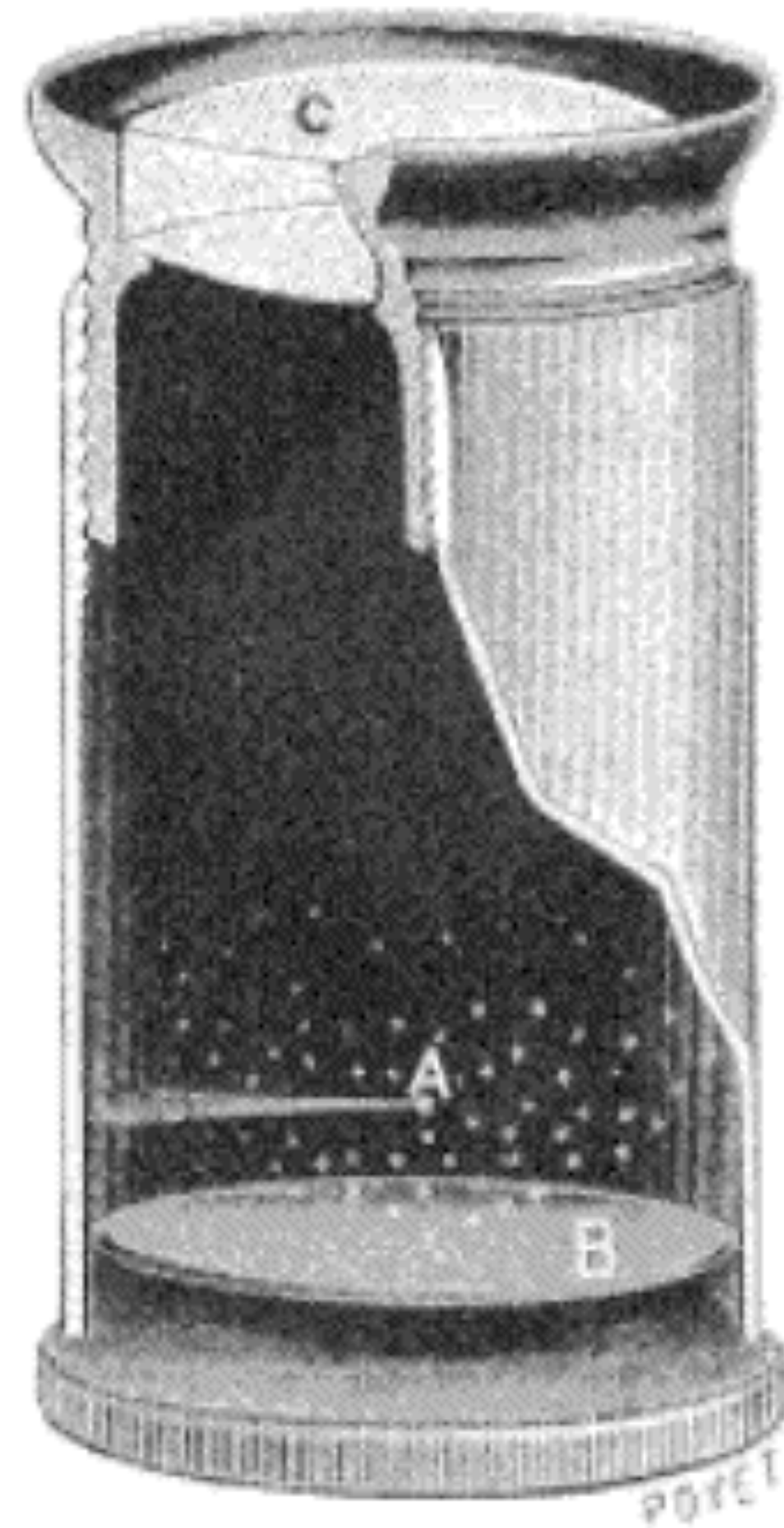
Outline



Scintillation (aside) - brief history

W. Crookes (1903) first demonstrates the “spintharoscope”, providing visual observations of individual scintillations (alpha particles impinging on a ZnS screen); precursor for scintillation counters.

Curran and Baker (1944) introduced the use of the photomultiplier tube in place of the naked eye, reviving the long-term use of scintillator detectors.

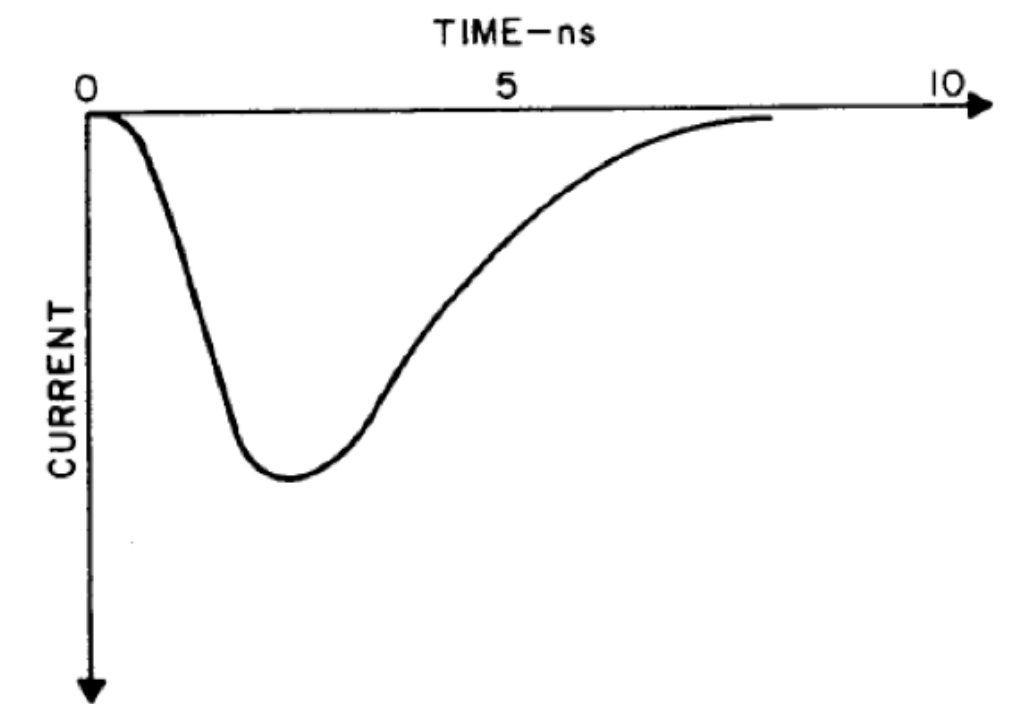
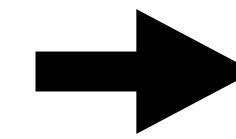
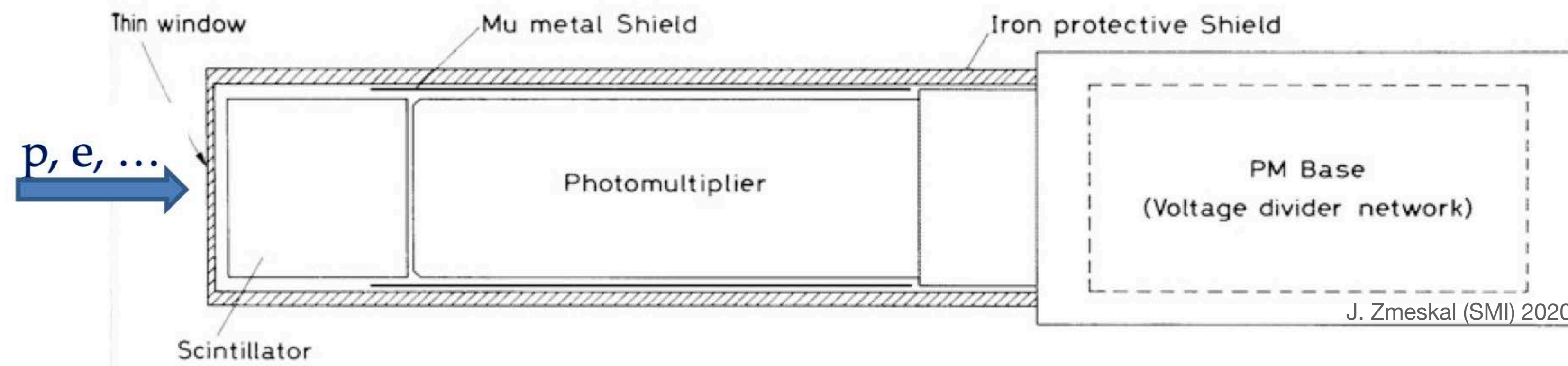


Artist's rendering of a spintharoscope from 1904.
Scale ~5cm. A: Ra-source; B: ZnS screen; C: Magnifying glass

Scintillation detectors - Introduction

General Principles

- Scintillator -> Luminescence via absorption of ionizing radiation (EM or particle)
 - Energy loss (dE/dx) is converted into light ($h\nu$).
- Light guide -> Collect and direct the scintillation light
- Photodetector -> Convert light to electrons and convert to digital signal



Use cases

- Tracking detectors -> Trajectory
- Calorimetry -> Energy deposition
- Time of flight -> Velocity
- Trigger and veto systems -> Signature

Organic

- Plastics, liquids, organic crystals
- Low density ($\sim 1 \text{ g/cm}^3$)
- Low Z
 - Requires interspersing high Z absorber material to achieve high stopping power for high energy γ 's
 - n detection by (n,p) interactions
- Up to 10,000 γ/MeV
- ns decay times
- Relatively inexpensive
- Moderately rad hard ($\sim 10 \text{ kGy/yr}$)

Inorganic

- Crystals
- Can high density ($> 8 \text{ g/cm}^3$)
- Typical have high Z
 - Leads to homogeneous detectors with very good energy resolution
 - Requires good light collection
 - Poor n detection efficiency
- Up to 50,000 γ/MeV
- ns to msec decay times
- Expensive
- Fairly radiation hard ($\sim 100 \text{ kGy/yr}$)

Noble gases

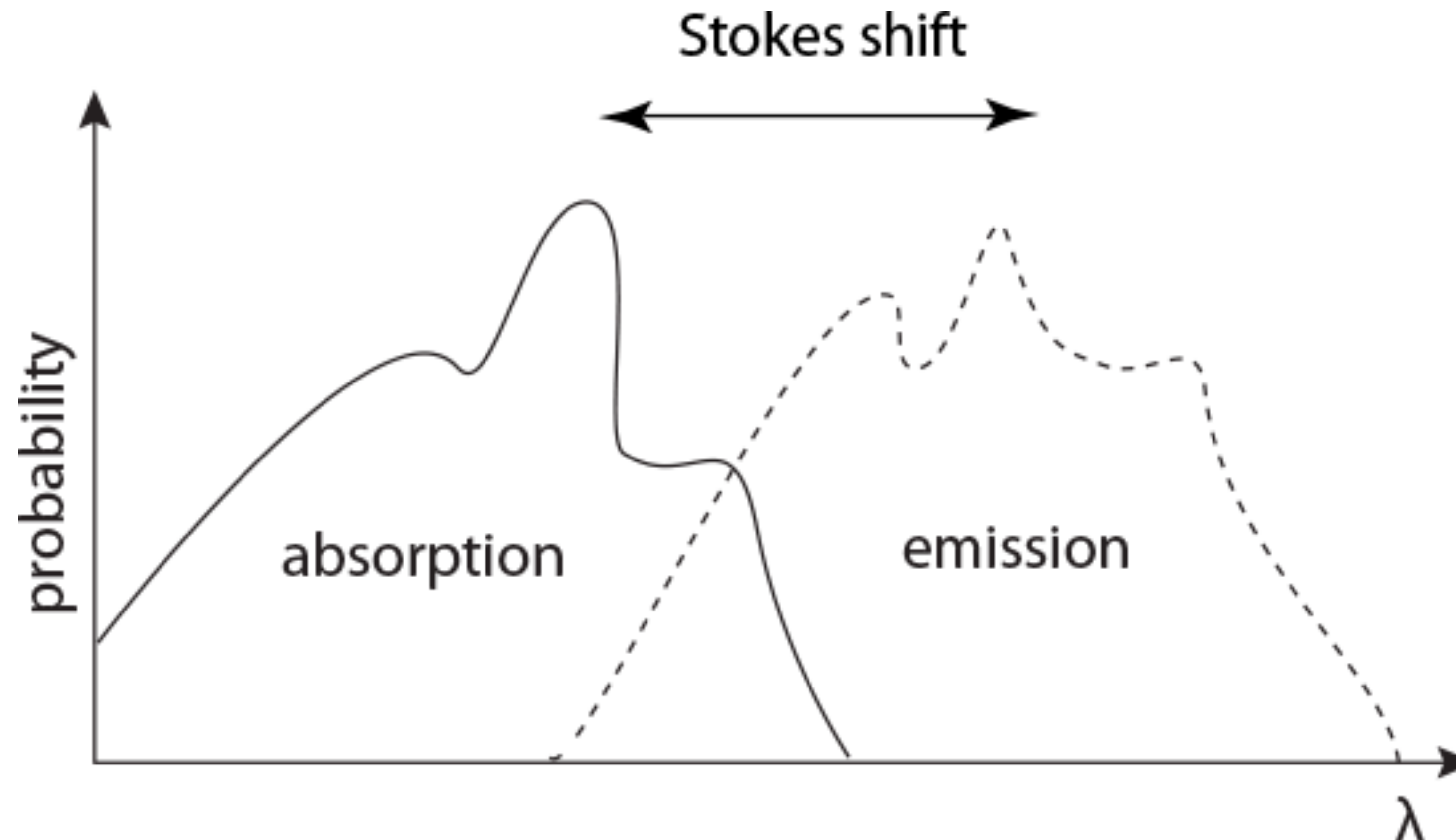
- LAr, LKr, LXe
- Requires working at cryogenic temperatures
- Moderate density and Z
 - ($\sim 1.4\text{-}3.0 \text{ g/cm}^3$)
- Scintillation in UV or VUV
- Light yield $\sim 50,000 \gamma/\text{MeV}$
- Produce ionization plus scintillation

Key properties

- High conversion efficiency and (linear) light yield
- (Short) luminescence decay time
- **Medium transparency to emitted light**

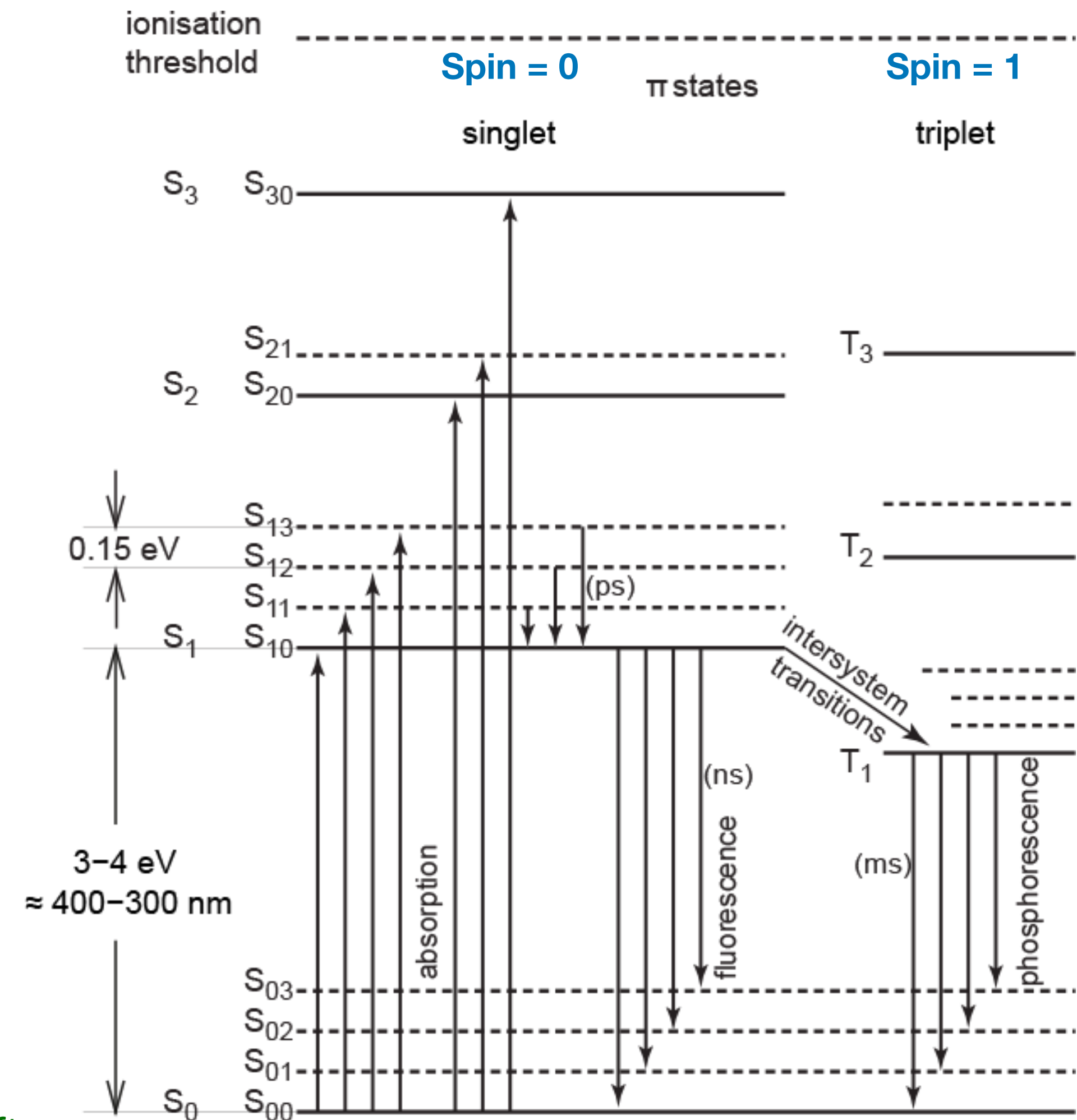
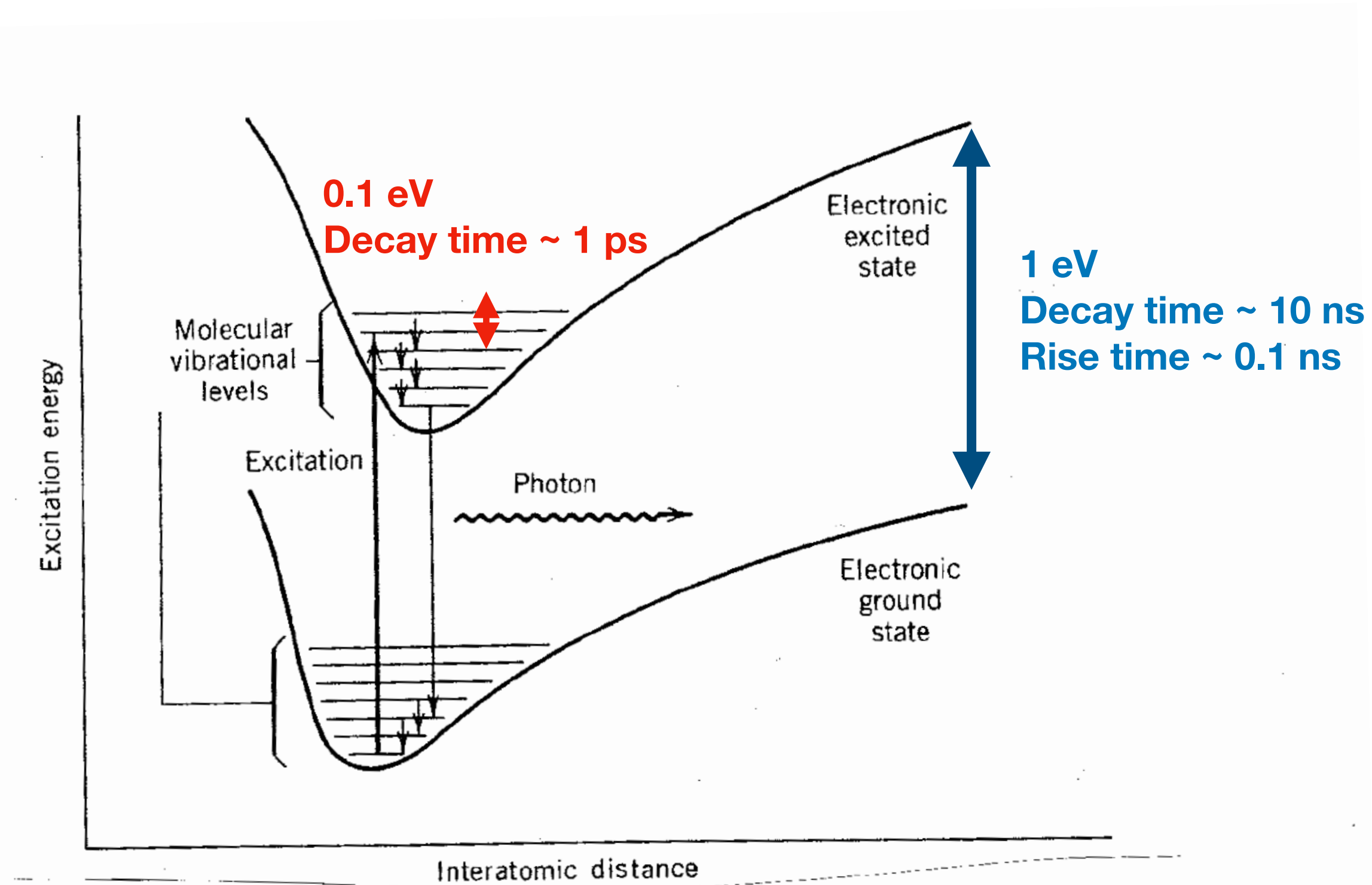
Scintillation medium should be transparent to the emitted light

- In principle the light may be reabsorbed by the scintillator.
- For light to escape, the emission energy must be lower than the absorption energy, a phenomenon known as Stokes shift



Scintillators - Organic (Mechanism)

Solid or liquid hydrocarbons where the excited electrons are not strongly involved in bonding of the material



Low Z

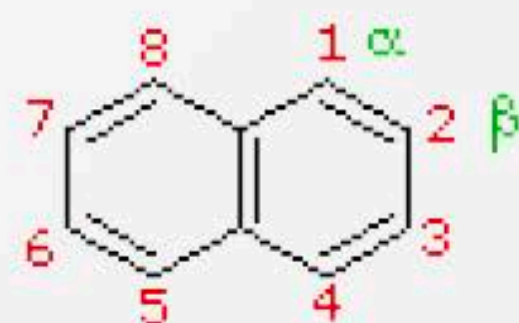
Low conversion efficiency ($\sim 8 - 10$ photons/keV)

Radiationless transitions (released via vibrations) produces **Stokes shift**

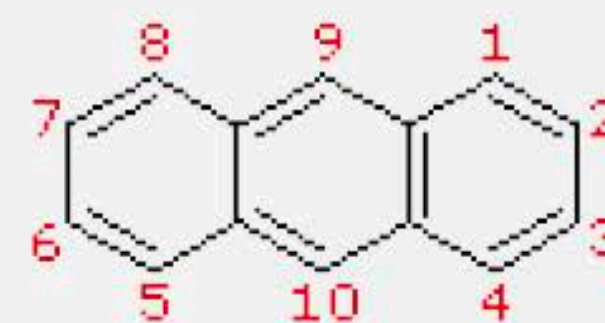
Fluorescence $\sim 10^{-8} \text{ s}$ (Fast)

Phosphorescence $\sim 10^{-6} \text{ s}$ (Slow - emission after intra-band transition)

crystals

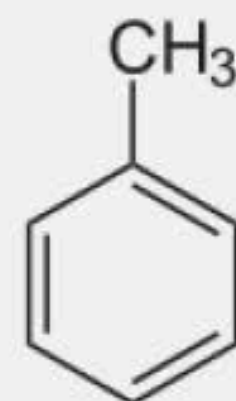


naphthalene
C₁₀H₈
m.p. 81° C

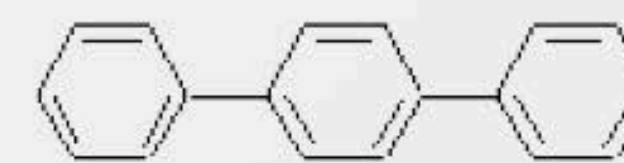


anthracene
C₁₄H₁₀
m.p. 217° C

liquids
(solutions)

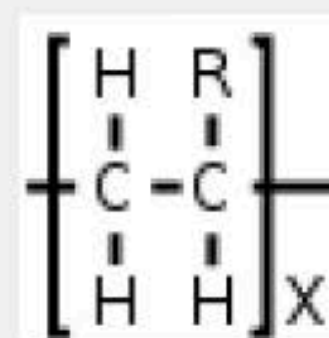


e.g. toluene

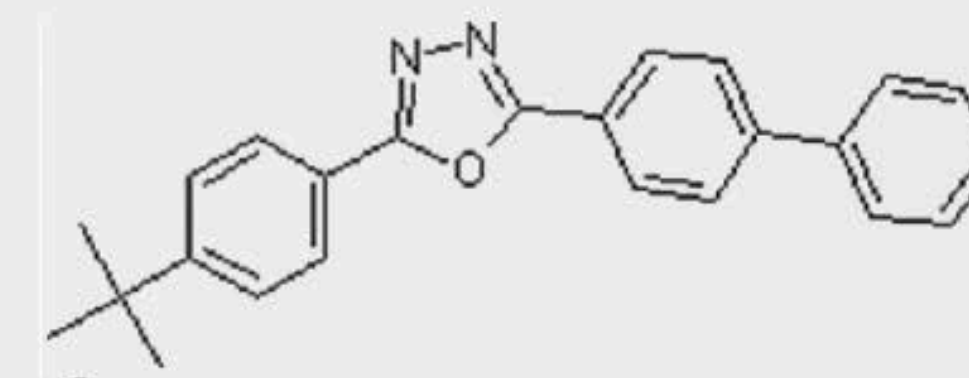


e.g. p-terphenyl

plastics
(polymerized
solutions)



e.g. polyvinyltoluene



e.g. Butyl-PBD

solvent

+

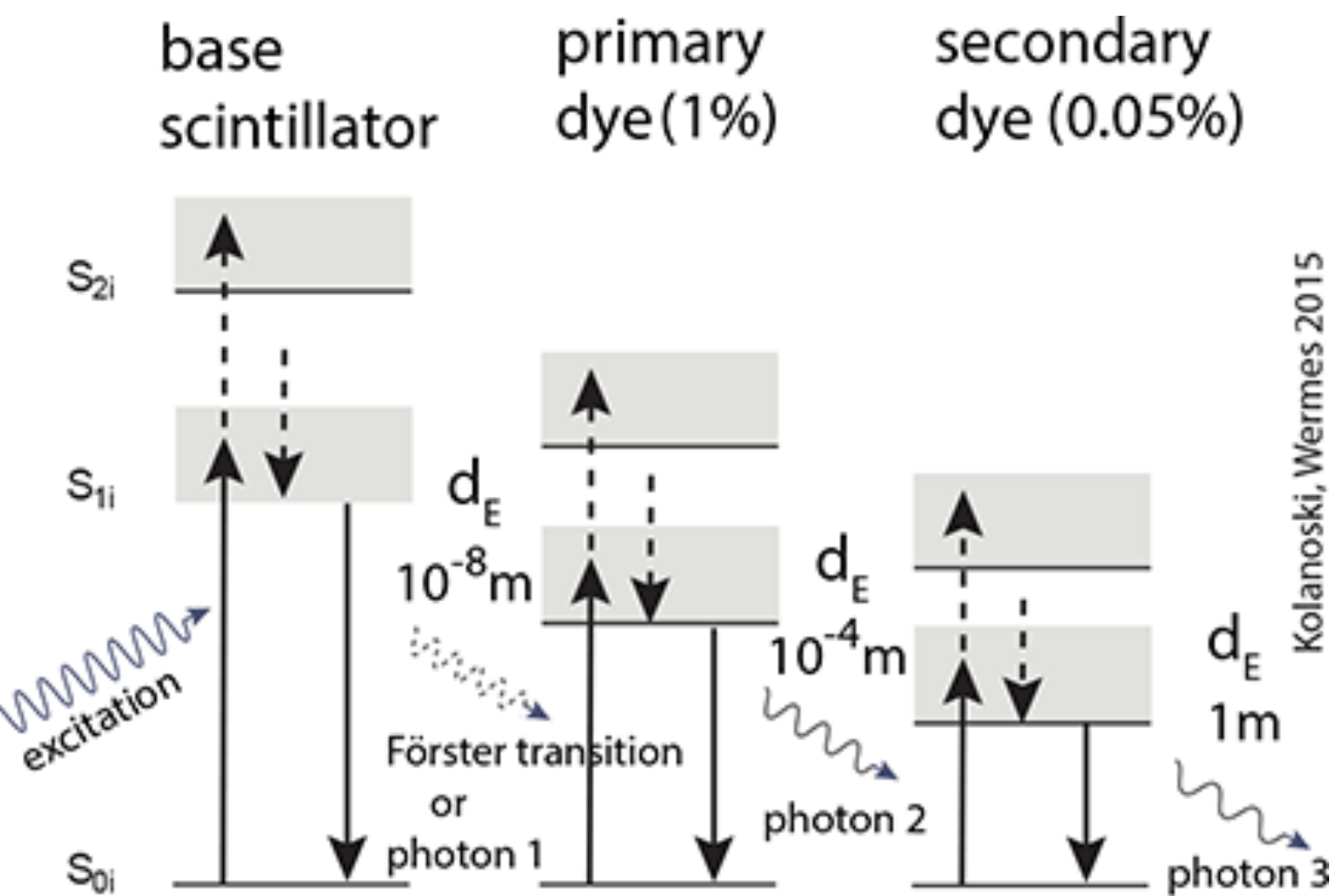
scintillator

High light yield but fragile and challenging to scale; rarely used

Solvents scintillate with low light yield; mix with a primary fluor (emitter) and secondary fluor (wavelength shifter). Impurities quench light output.

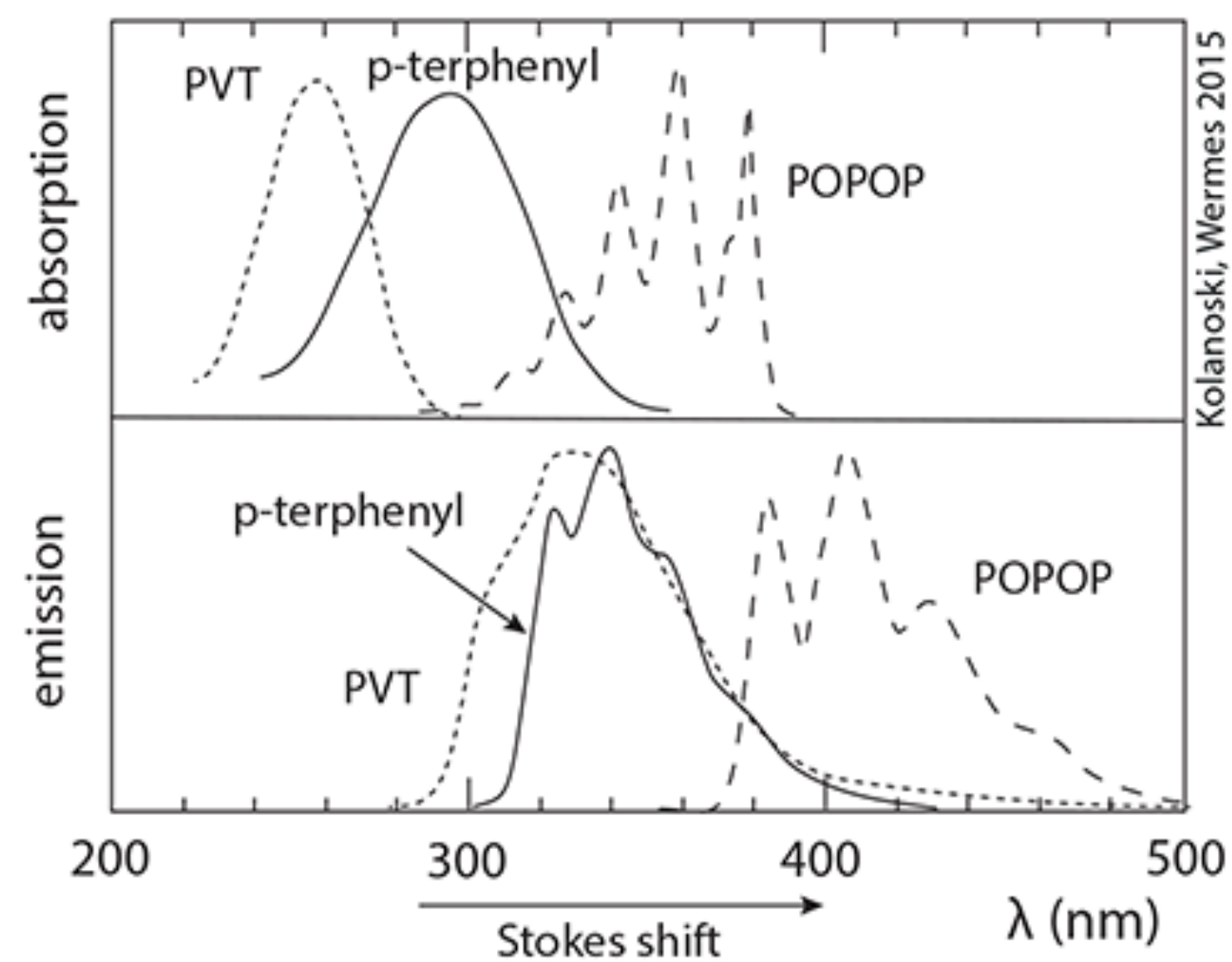
Cast in any shape for easy handling; Aging (solvents, grease, high temperature, ionizing radiation) impacts light yield

Scintillators - Organic (Liquid)

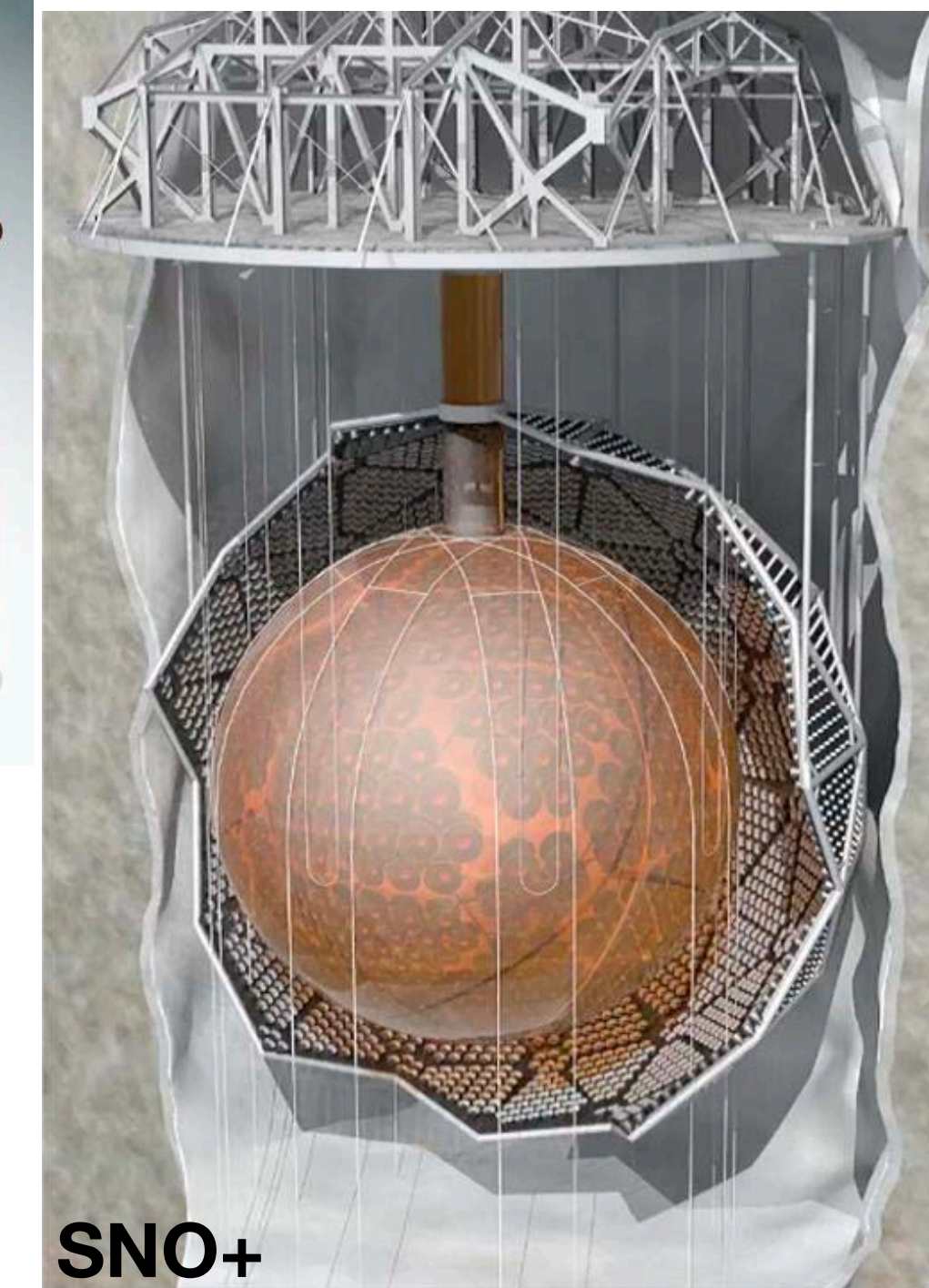
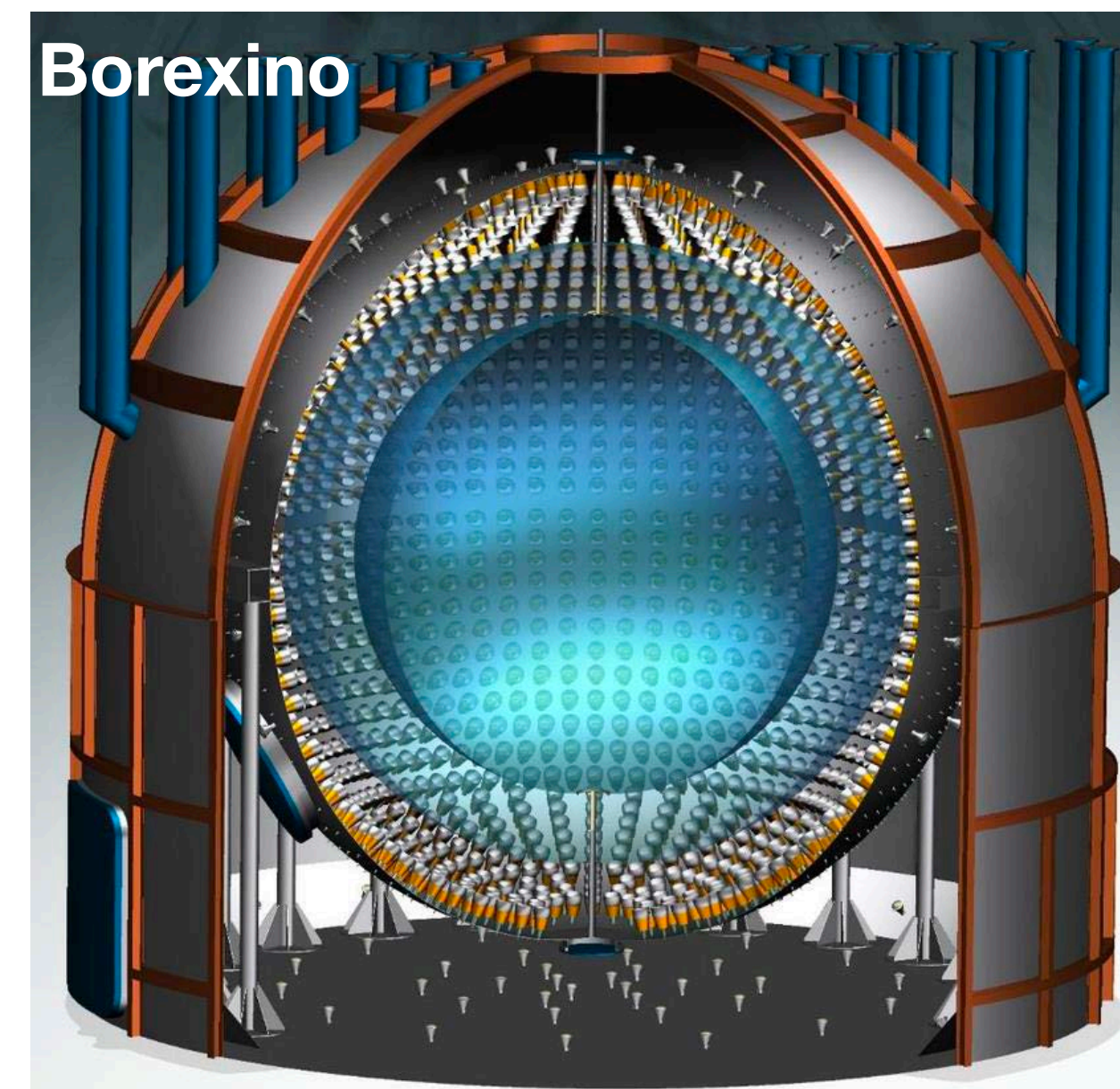


(a) Wavelength shifting: principle.

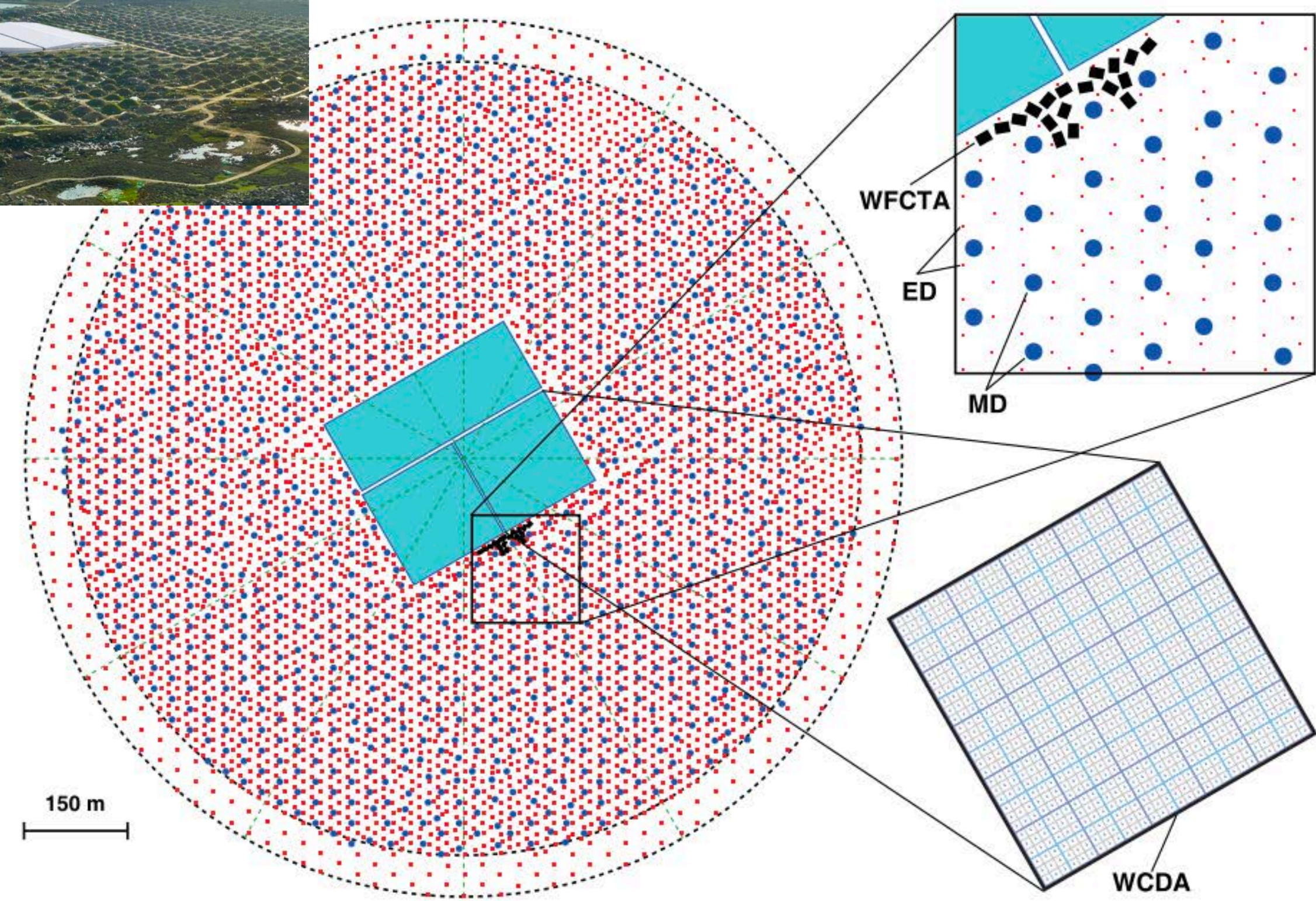
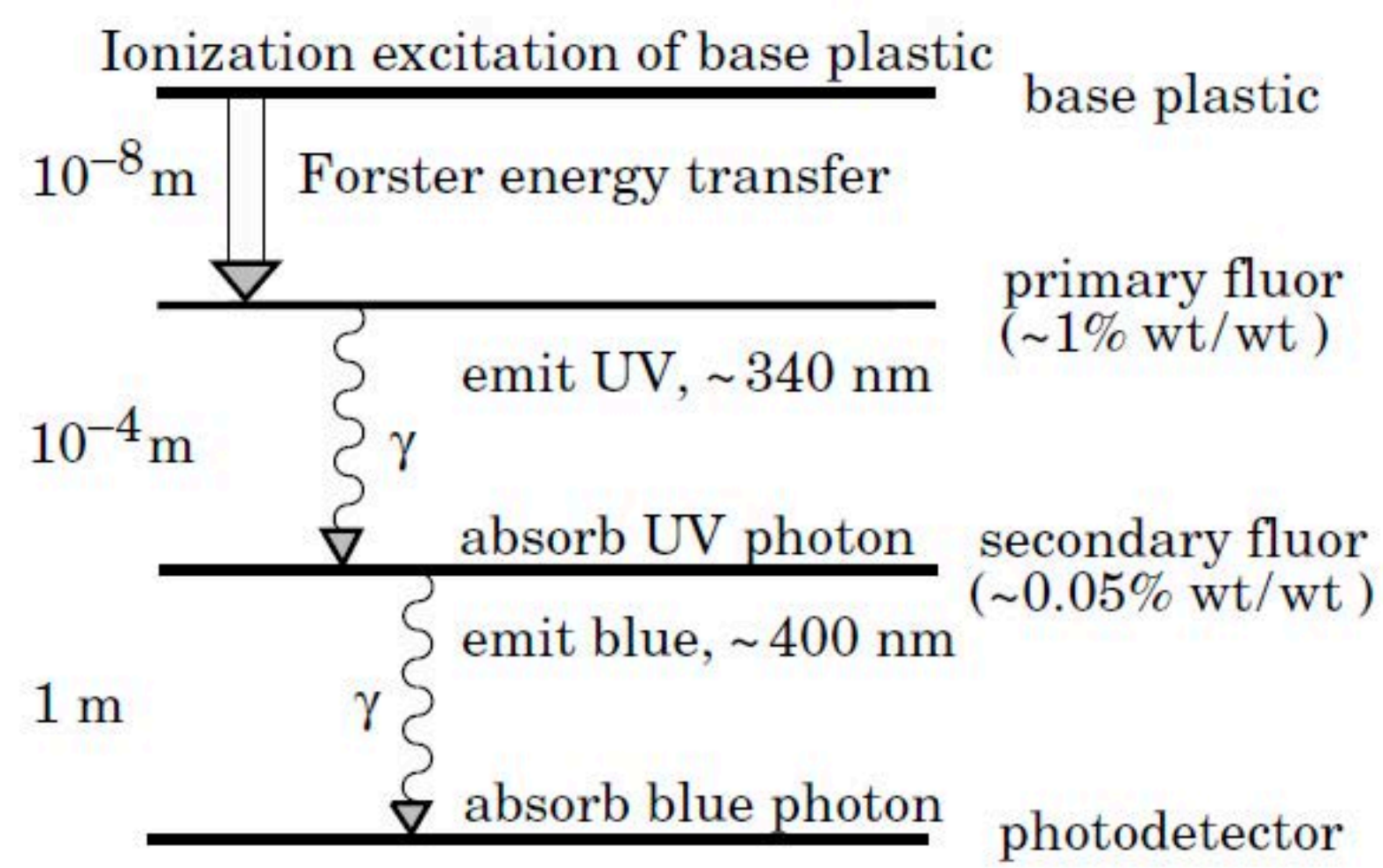
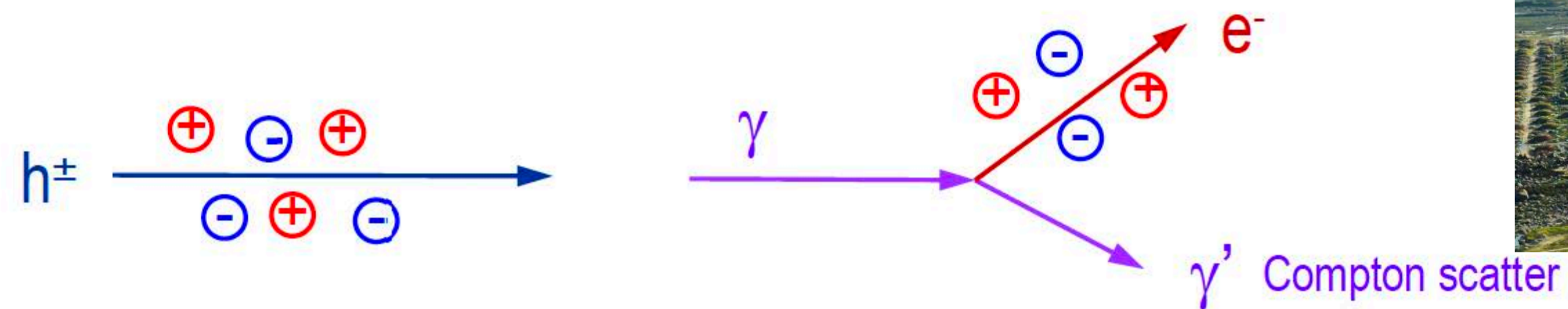
Solvents scintillate with low yield; mix base (PVT) with a primary fluor (emitter) and secondary fluor (wavelength shifter). Impurities quench light output.



(b) Absorption and emission spectra.

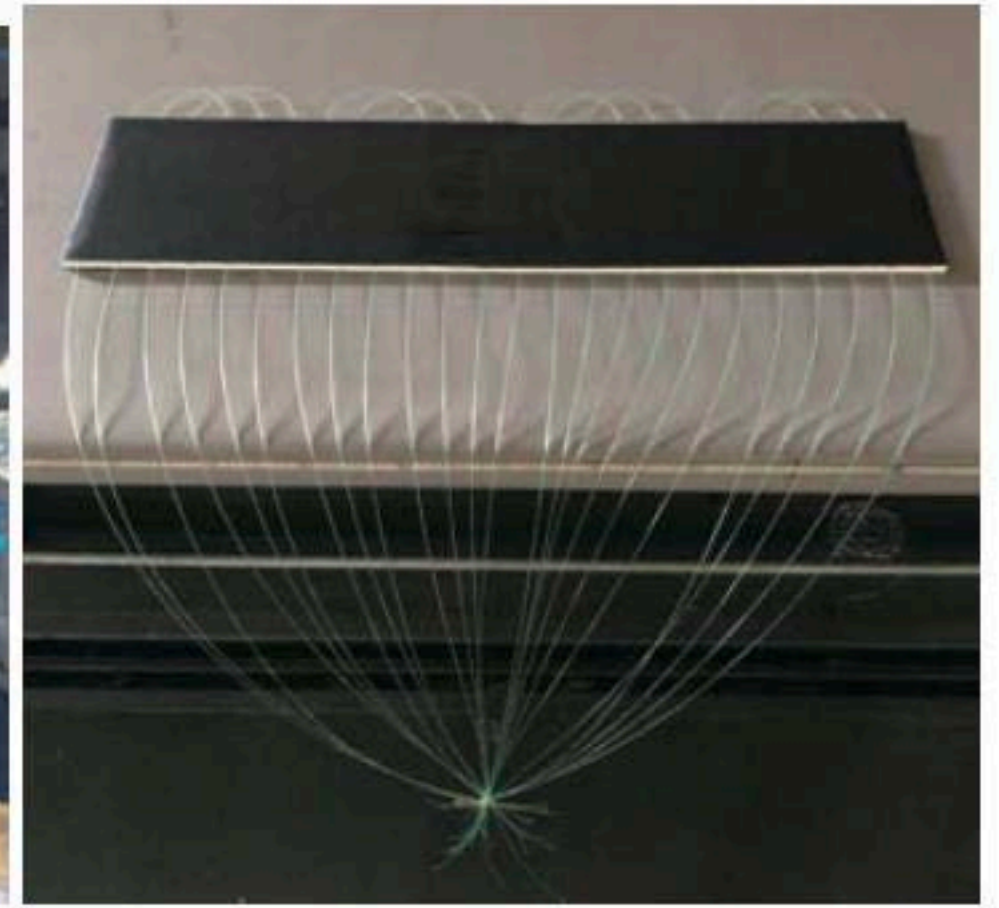
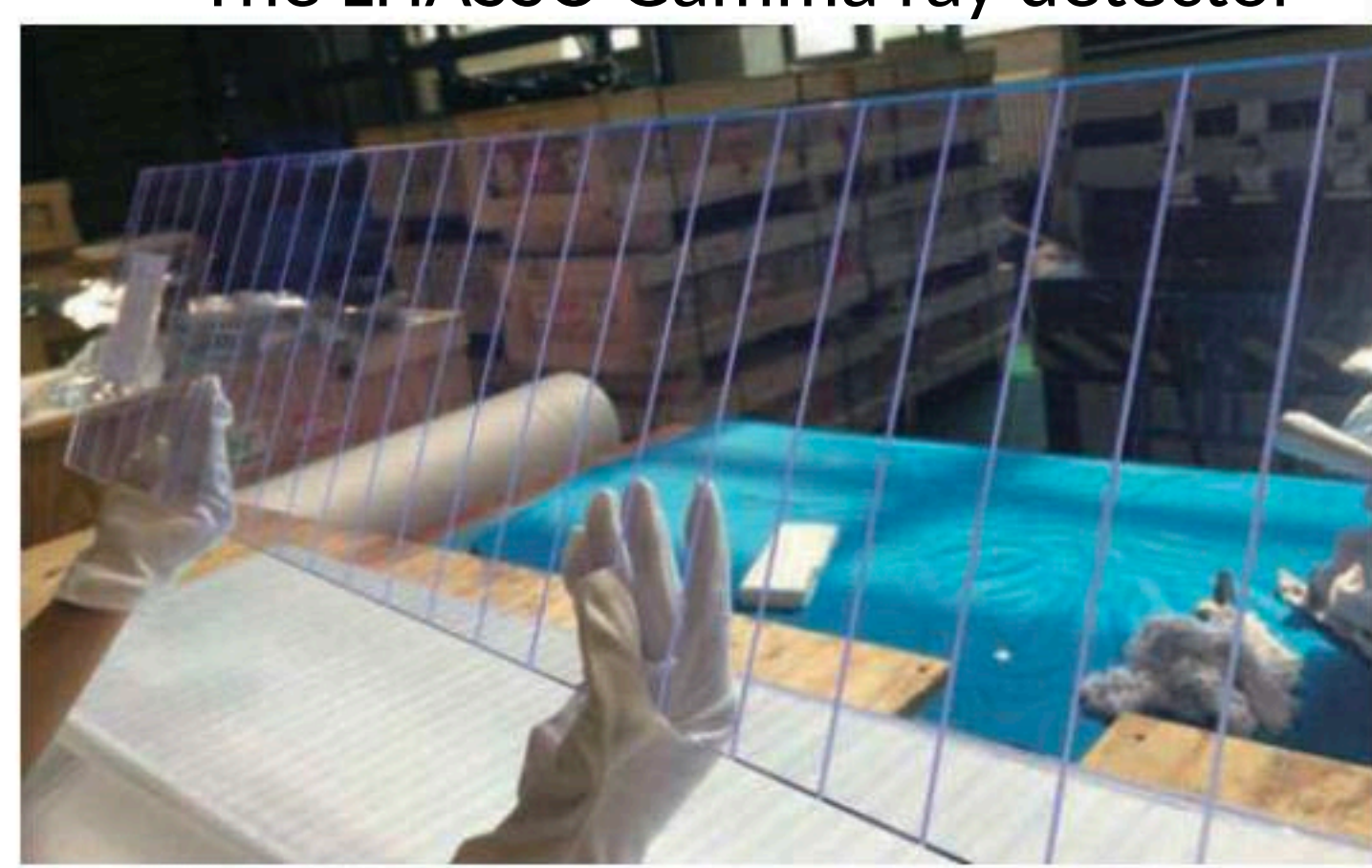


Scintillators - Organic (Plastic)



The LHAASO Gamma ray detector

- Charged particles produce ionization that causes excitations in the medium
- Gamma rays interact and produce electrons that can result in ionization
- Each step in the process involves efficiency losses; overall light yield of approximately 1 photon per 100 eV of energy deposition.



Scintillators - Organic (Light Yield)

Follows Birk's (semi-empirical) Law

$$\frac{dL}{dx} = \frac{S \frac{dE}{dx}}{1 + k_B \frac{dE}{dx}}$$

S is the scintillation efficiency

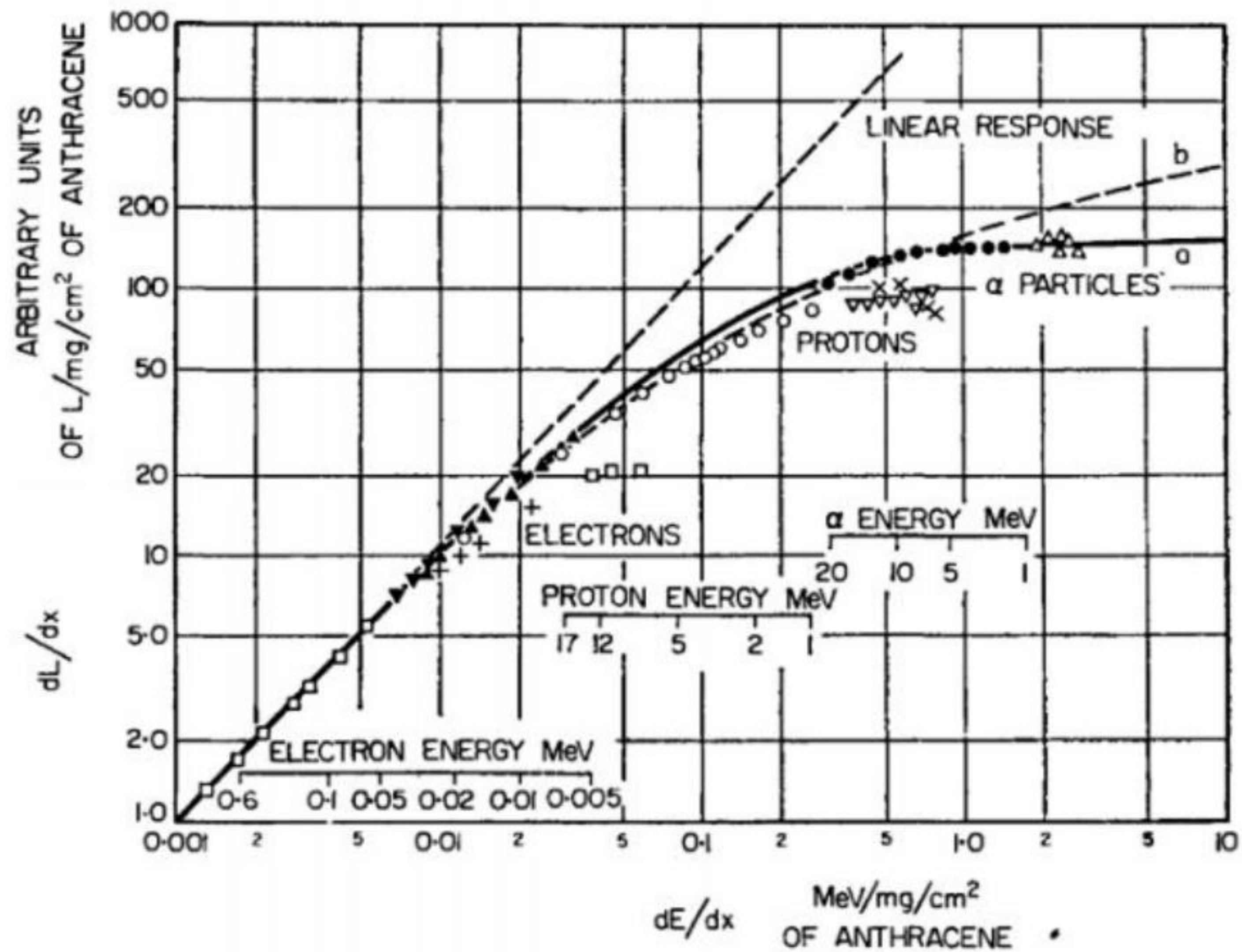
k_B is the quenching strength – ionization saturates

For minimum ionizing particles (GeV muons)

$$\frac{dL}{dx} = S \frac{dE}{dx}$$

For strongly ionizing radiation (MeV alpha particles)

$$\frac{dL}{dx} = \frac{S}{k_B}$$



The quenching strength must be determined experimentally.

Scintillators - Organic (Signal Timing Shape)

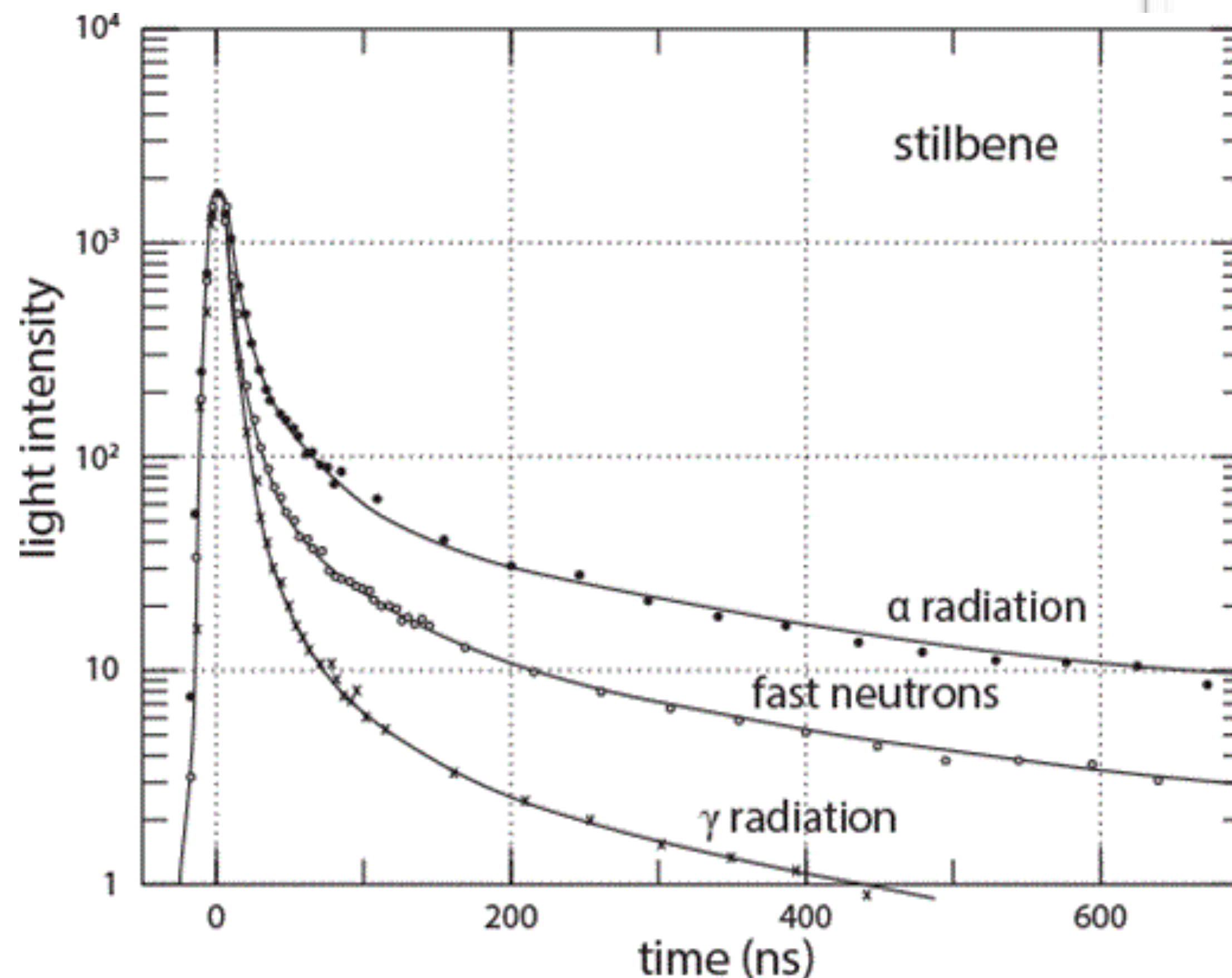
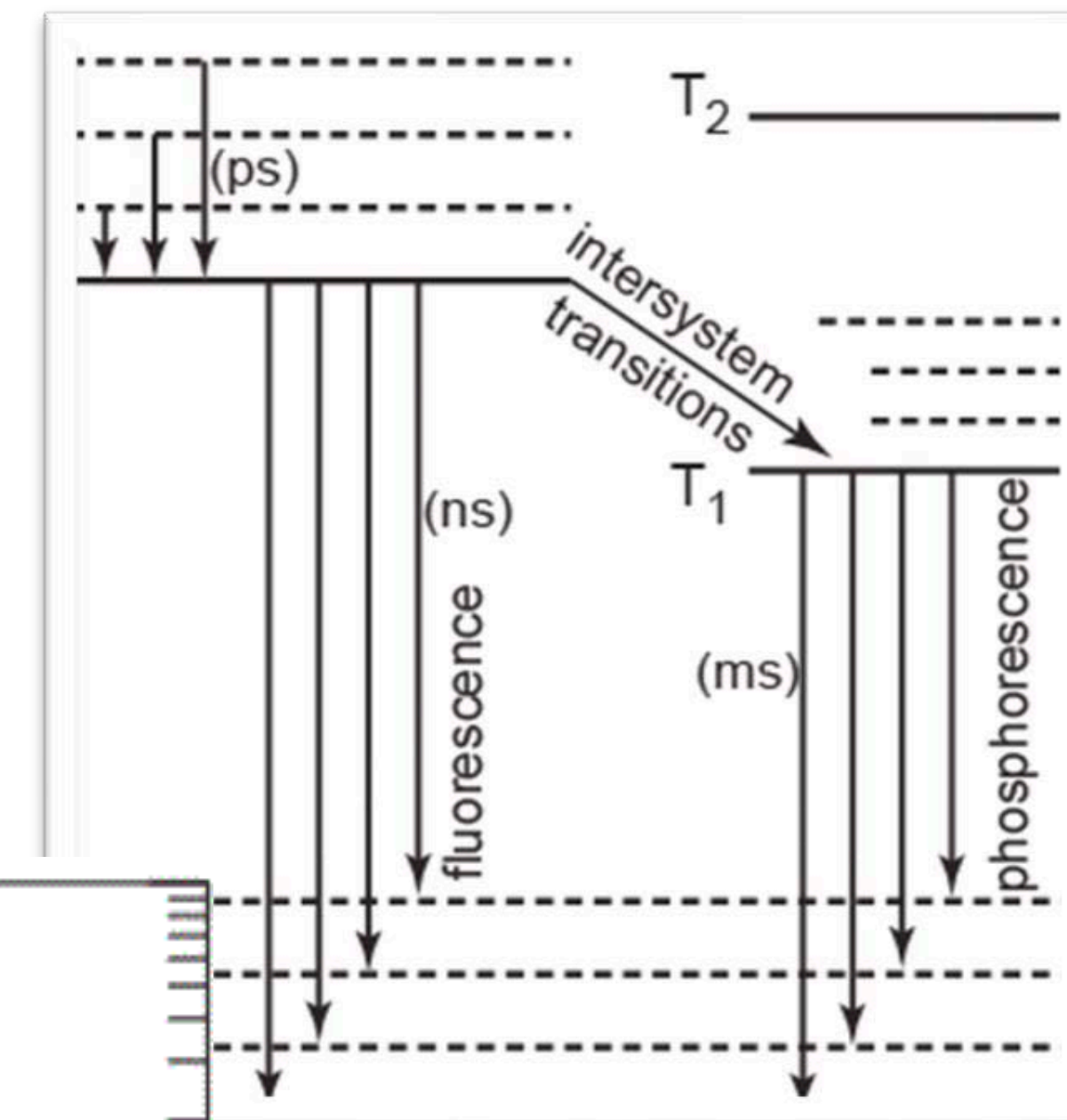
Prompt fluorescence emission follows a simple exponential decay; ie. For rise time:

$$I(t) = I_0 g(t) e^{-t/\tau_{dec}}$$

Many scintillators have a second component (with longer decay time)

$$I(t) = I_1 g(t) e^{-t/\tau_f} + I_2 e^{-t/\tau_s}$$

Relative contribution of I_1 and I_2 depend on the ionization density of the incident particles -> leverage for pulse shape discrimination to distinguish particle types.



Scintillators - Inorganic (Mechanism)

Relies on the crystal lattice structure

- Energy levels created by impurities in the band gap (luminescence centres)

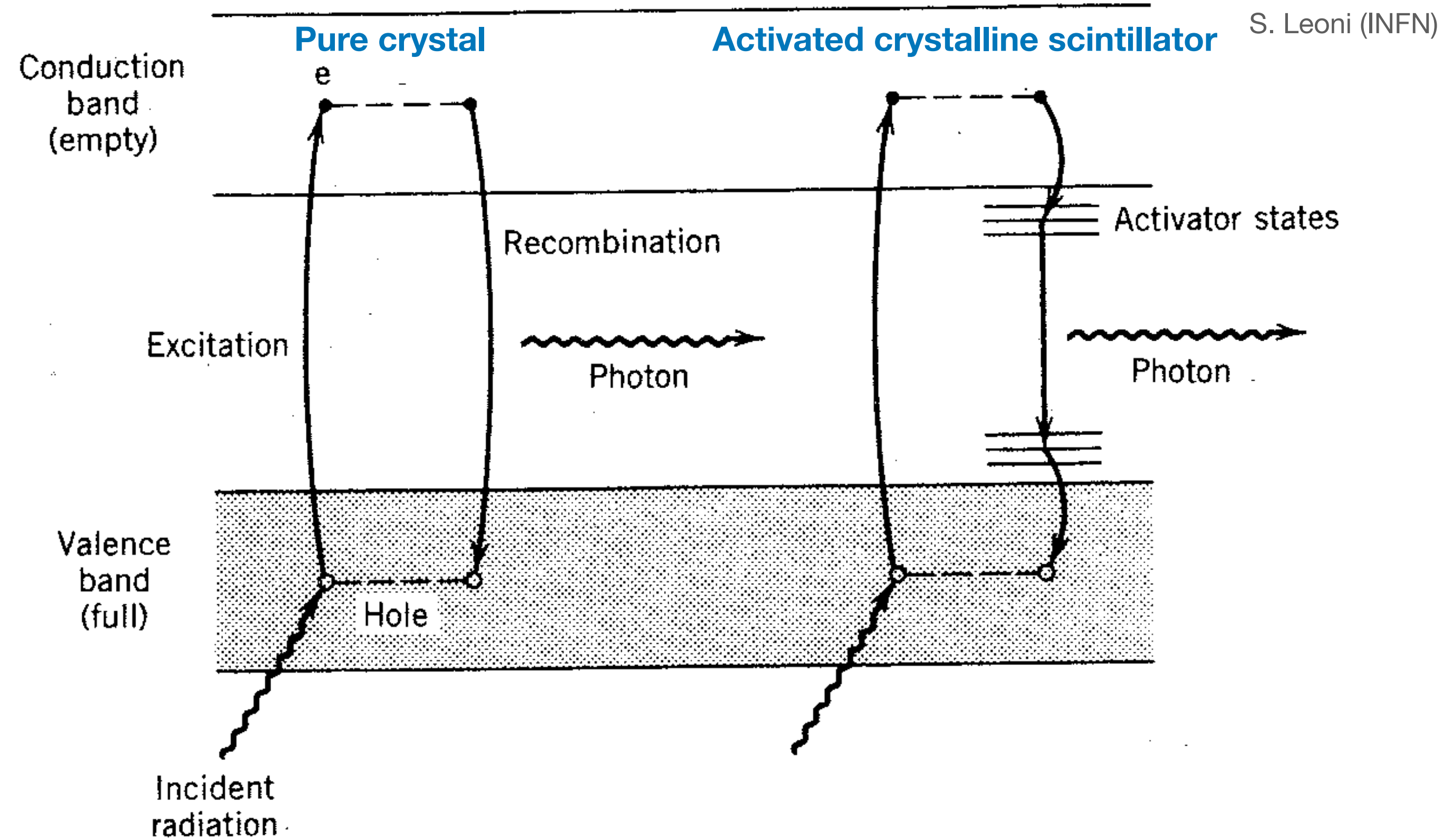
Excited electrons move between atomic states, from the valence to the conducting band

Doping material changes the energy structure of the lattice; used to minimize re-absorption from the crystal as the emitted light has a lower energy than the energy-gap

eg. NaI(Tl) \rightarrow 4 eV with decay time of \sim 230 ns and rise time of \sim 10 ns

High Z (high stopping power \rightarrow total absorption calorimeters)

High conversion efficiency \rightarrow \sim 40 photons/keV (\sim 40x better than plastic scintillators)



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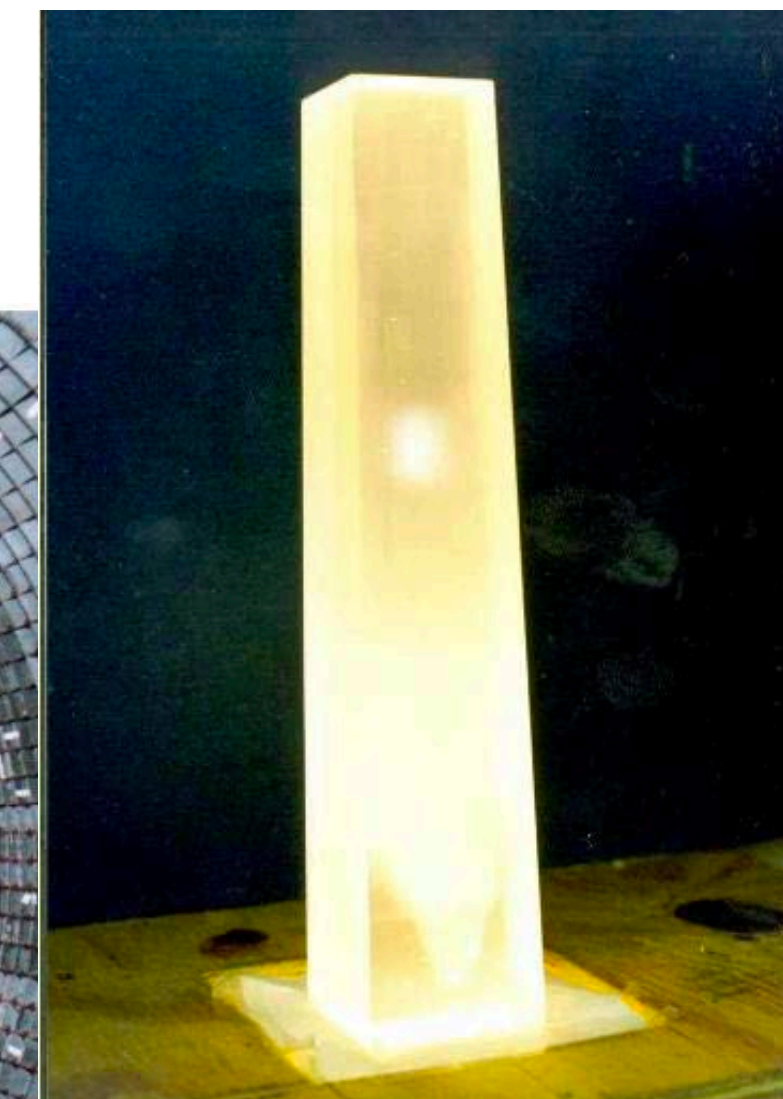
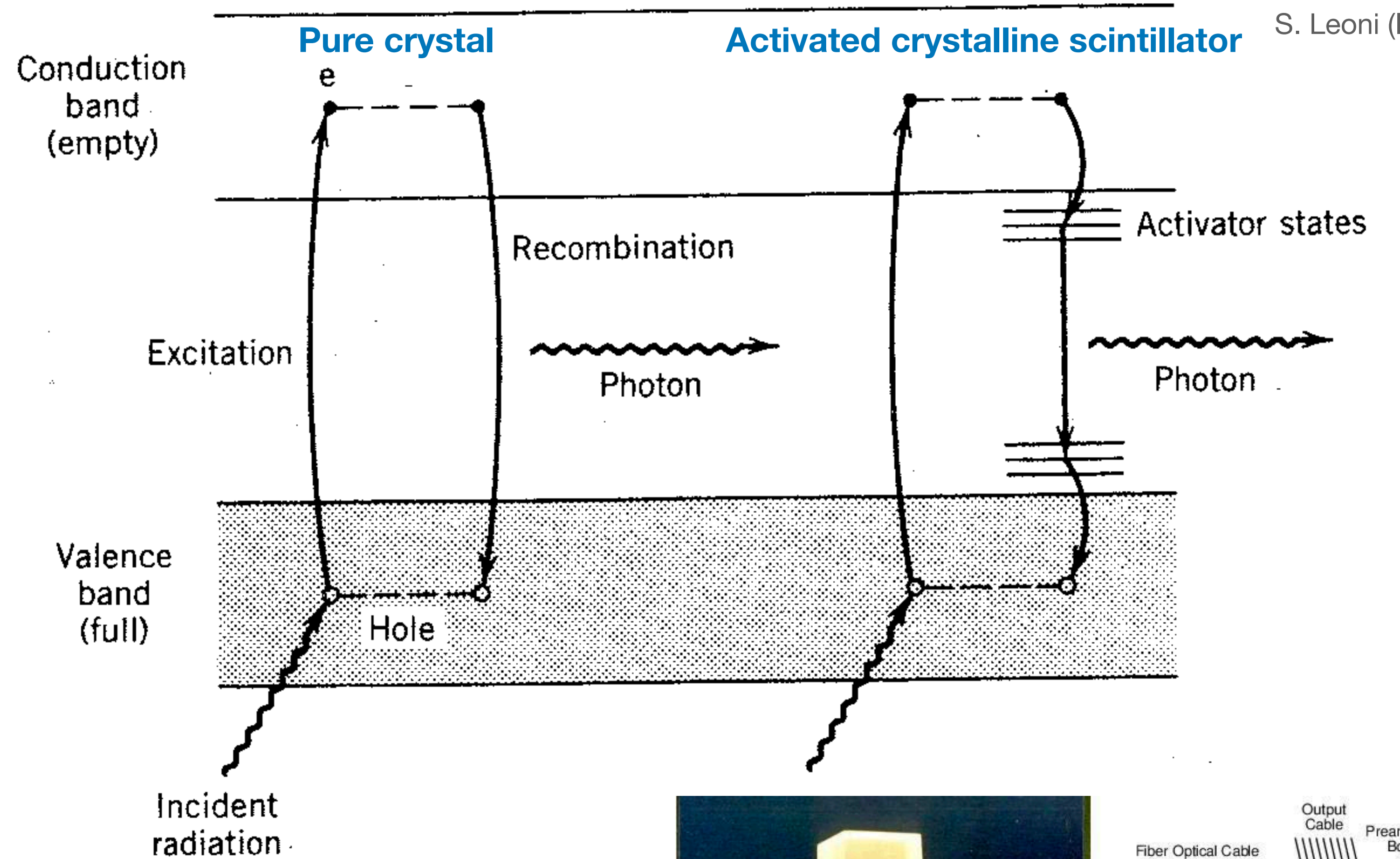
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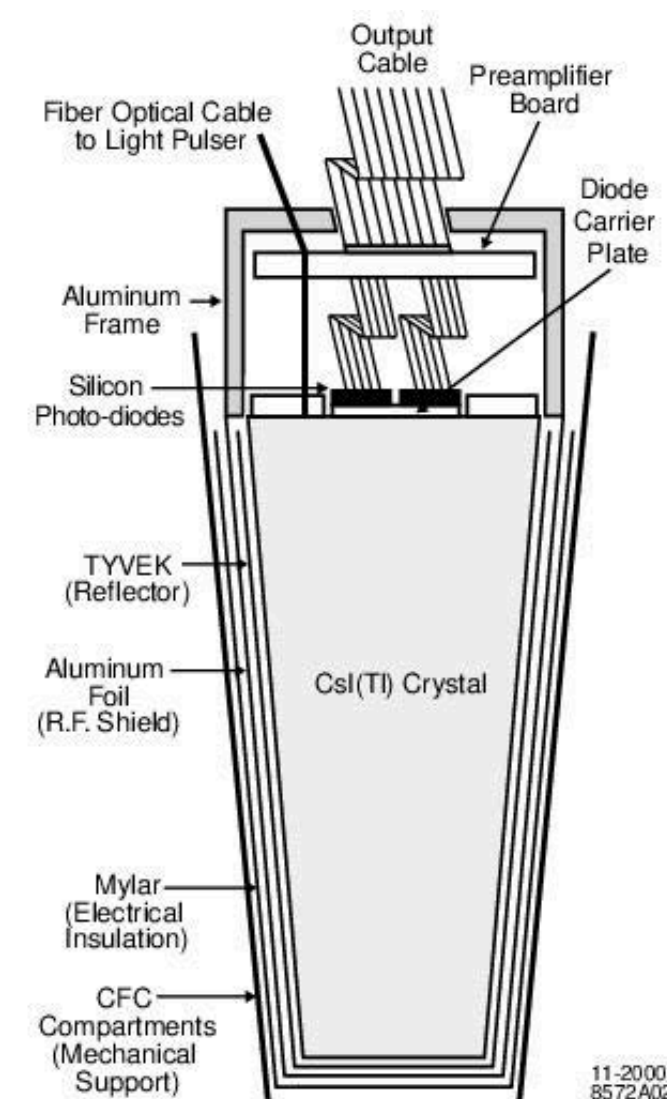
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BaBar calorimeter



Scintillators - Inorganic (Examples; Light Yield)

Light Yield depends on the absorbed energy and the electron-hole pair production average energy

$$L = \frac{E_{abs}}{W_{e/h}} S Q$$

S = transfer energy efficiency (thermalized pairs -> luminescence centres)

Q = quantum yield of the luminescence centre

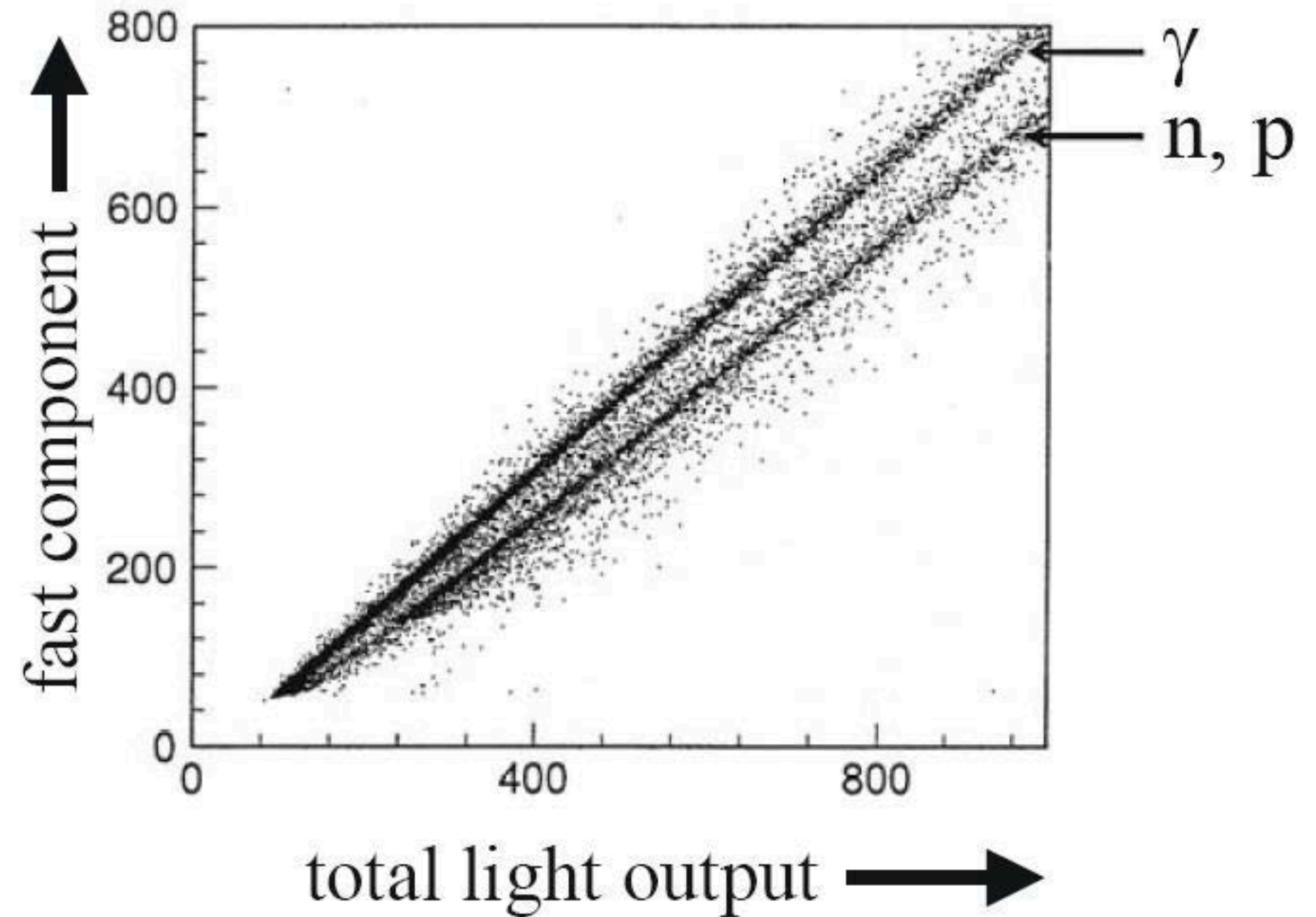
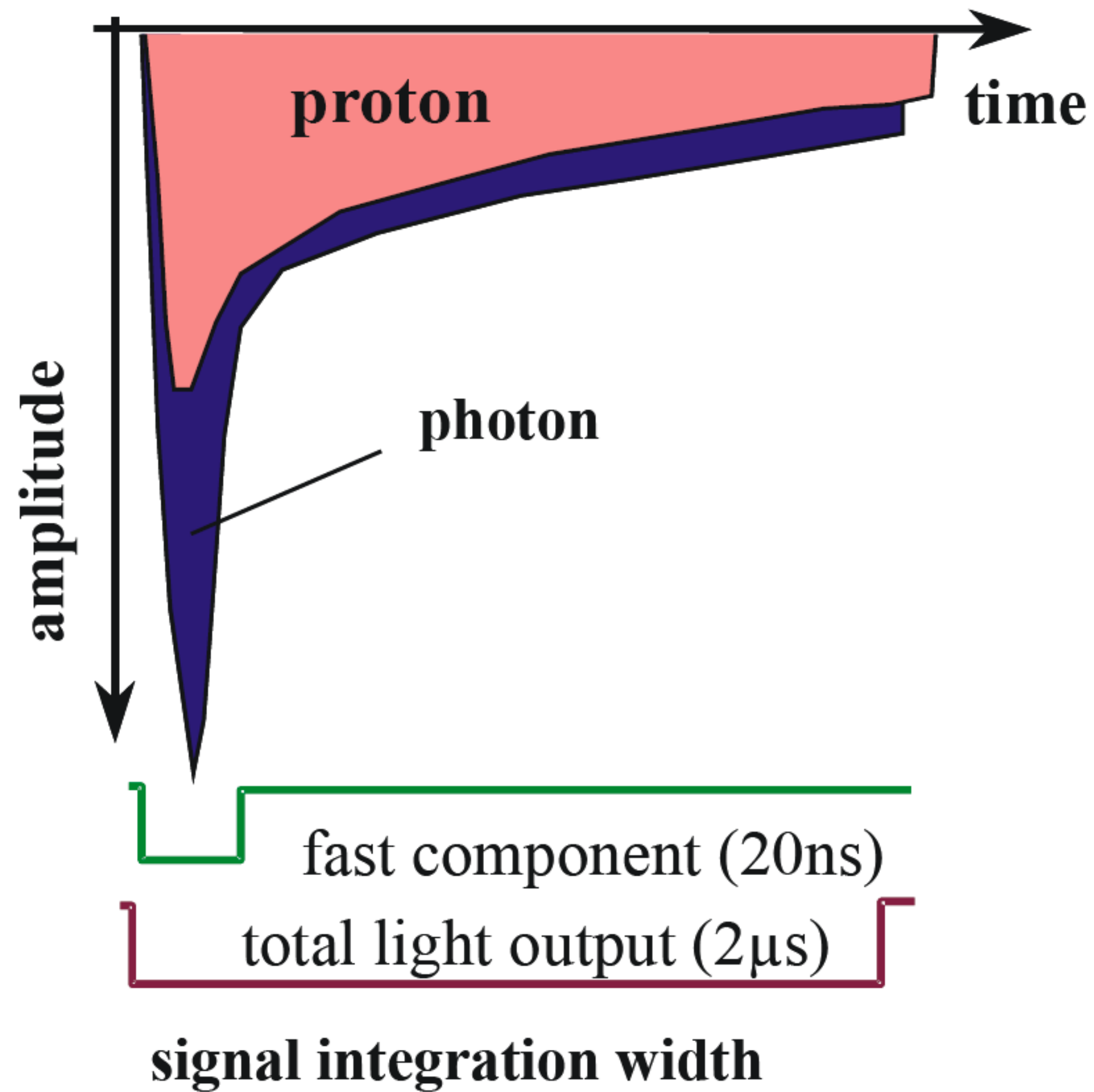
Maximum theoretical yield approximately 140 photons/keV*

Scintillator composition	Density (g/cm ³)	Index of refraction	Wavelength of max.Em. (nm)	Decay time Constant (μs)	Scinti Pulse height ¹⁾	Notes
Nal(Tl)	3.67	1.9	410	0.25	100	2)
Csl	4.51	1.8	310	0.01	6	3)
Csl(Tl)	4.51	1.8	565	1.0	45	3)
CaF ₂ (Eu)	3.19	1.4	435	0.9	50	
BaF ₂	4.88	1.5	190/220 310	0,0006 0.63	5 15	
BGO	7.13	2.2	480	0.30	10	
CdWO ₄	7.90	2.3	540	5.0	40	
PbWO ₄	8.28	2.1	440	0.020	0.1	
CeF ₃	6.16	1.7	300 340	0.005 0.020	5	
GSO	6.71	1.9	430	0.060	40	
LSO	7	1.8	420	0.040	75	
YAP	5.50	1.9	370	0.030	70	

1) Relative to Nal(Tl) in %; 2) Hygroscopic; 3) Water soluble

Scintillators - Inorganic (Signal Pulse Shape)

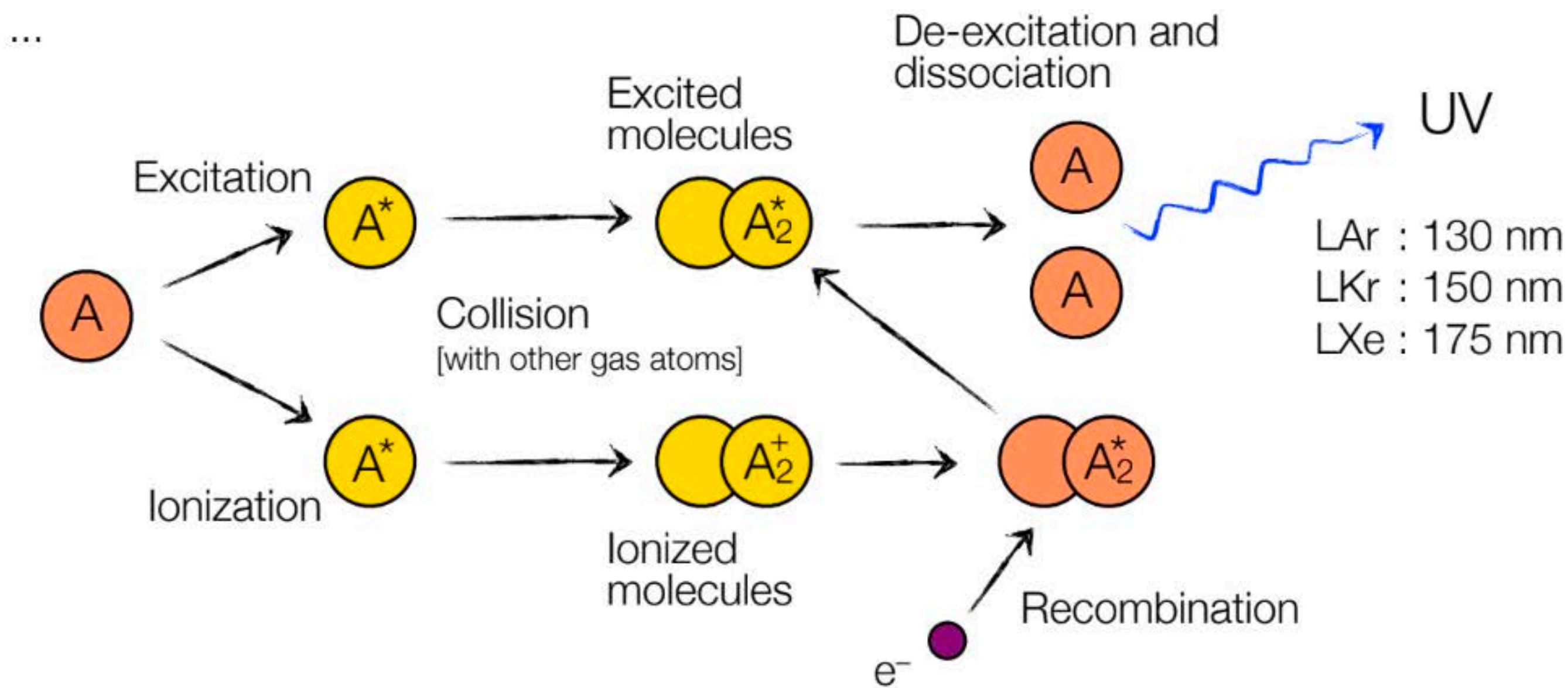
e.g. BaF₂ one of the fastest scintillators with an emission component with sub-ns decay time and a possible timing resolution of ~200 ps



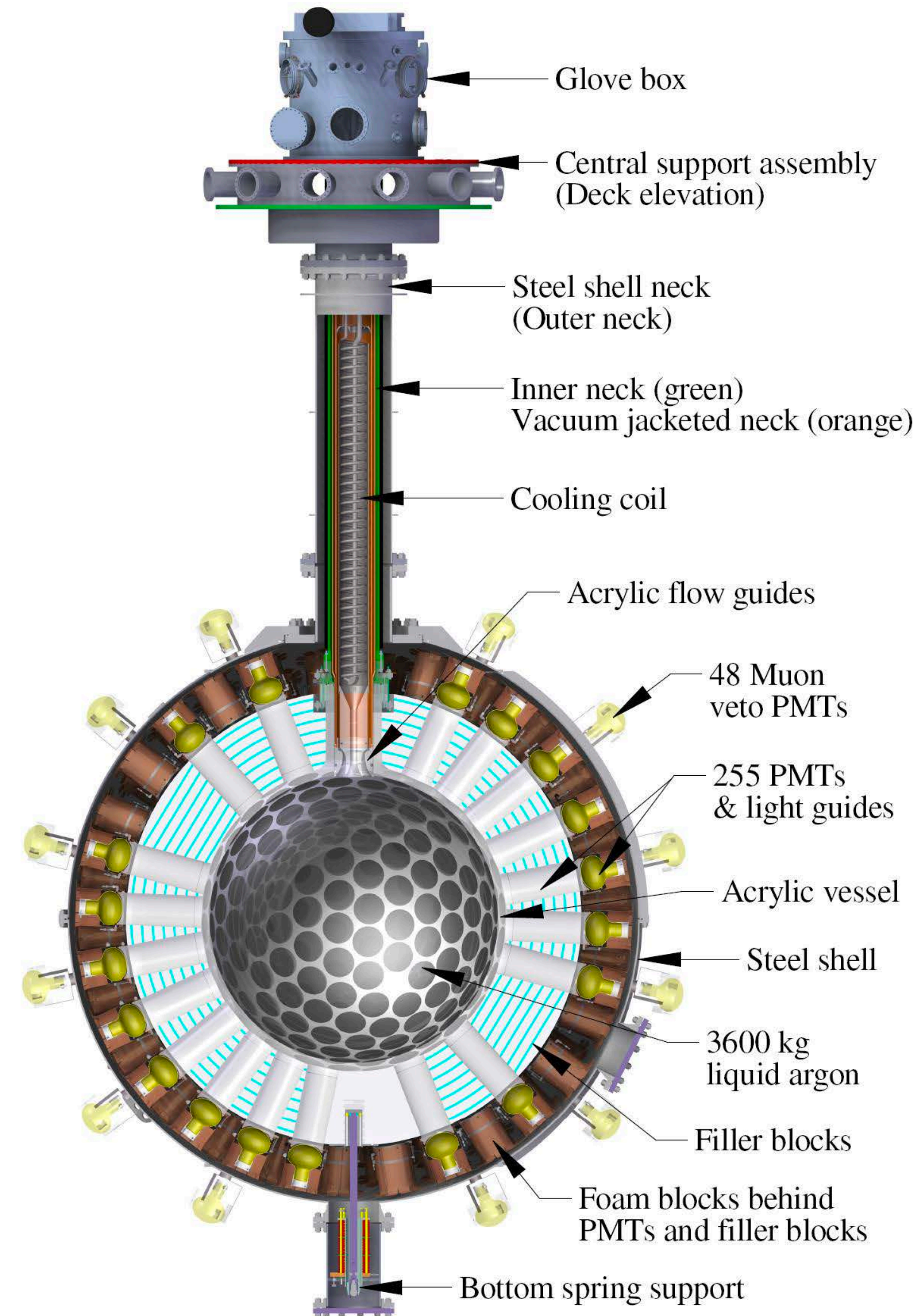
Scintillators - Noble Gas

High-purity inert gases may generation scintillation light via emission through decay of excited dimers to the ground state.

Emission typically in the far UV -> wavelength shifters



Pulse shape discrimination possible via measurement of ratio of single (fast) to triplet (slow) components that have separation for electromagnetic and nuclear recoils events.



Scintillators - Noble Gas

Gas/liquid	Z	Boiling point at 1 bar (K)	ρ (liquid) (g/cm ³)	$\lambda_{\text{peak}}^{\text{em}}$ (nm)	Photons per MeV	τ_{dec} (typ.) (ns)
<i>gases (1 atm, 20 °C)</i>						
N ₂	7	77	1.17×10^{-3}	390	100	2.5
CF ₄	8.6	145.3	3.93×10^{-3}	300/630	1200	6
Helium	2	4.2	1.66×10^{-4}	78	1100	<20
Neon	10	27.1	8.39×10^{-4}	~80		1.2
Argon	18	87.3	1.66×10^{-3}	127	18 200	6
<i>liquids</i>						
Helium	2	4.2	0.13	80	15 000	10
Neon	10	27.1	1.21	78	30 000	15
Argon	18	87.3	1.40	127	40 000	6
Krypton	36	119.8	2.41	147	25 000	3
Xenon	54	165.0	3.06	175	46 000	3
<i>solid (for reference)</i>						
NaI(Tl)	46.5		3.67	410	43 000	245

Very low Z
 Fast decay times
 Light yields that can reach similar
 levels to crystals

Organic Scintillators

Advantages

- Very fast
- Easily shaped
- Small temperature dependence
- Relatively inexpensive

Disadvantages

- Lower light yield (efficiency ~3%)
- Prone to radiation damage

Inorganic Scintillators

Advantages

- High light yield (efficiency ~13%)
- High density
- Good energy resolution

Disadvantages

- Crystal growth -> expensive
- Large temperature dependence

- In addition to Stokes Shift, attenuation lengths are impacted by: fluorine concentration and stabilizer additives (self-absorption); optical clarity and uniformity; surface quality

- Plastic scintillators age (exposure to solvent vapours, high temperature, flexing, irradiation) reducing the light yield. The surface is particularly fragile and can experience "crazing" (micro-fractures) in particular where there is surface contact with oils (including finger prints), solvents,...

- Long-lived luminescence (100s ns) may occur in plastic scintillators that does not follow an exponential decay.

- Radiation damage creates color centres that absorb more strongly blue/UV than longer wavelengths, reducing both light yield and attenuation length.

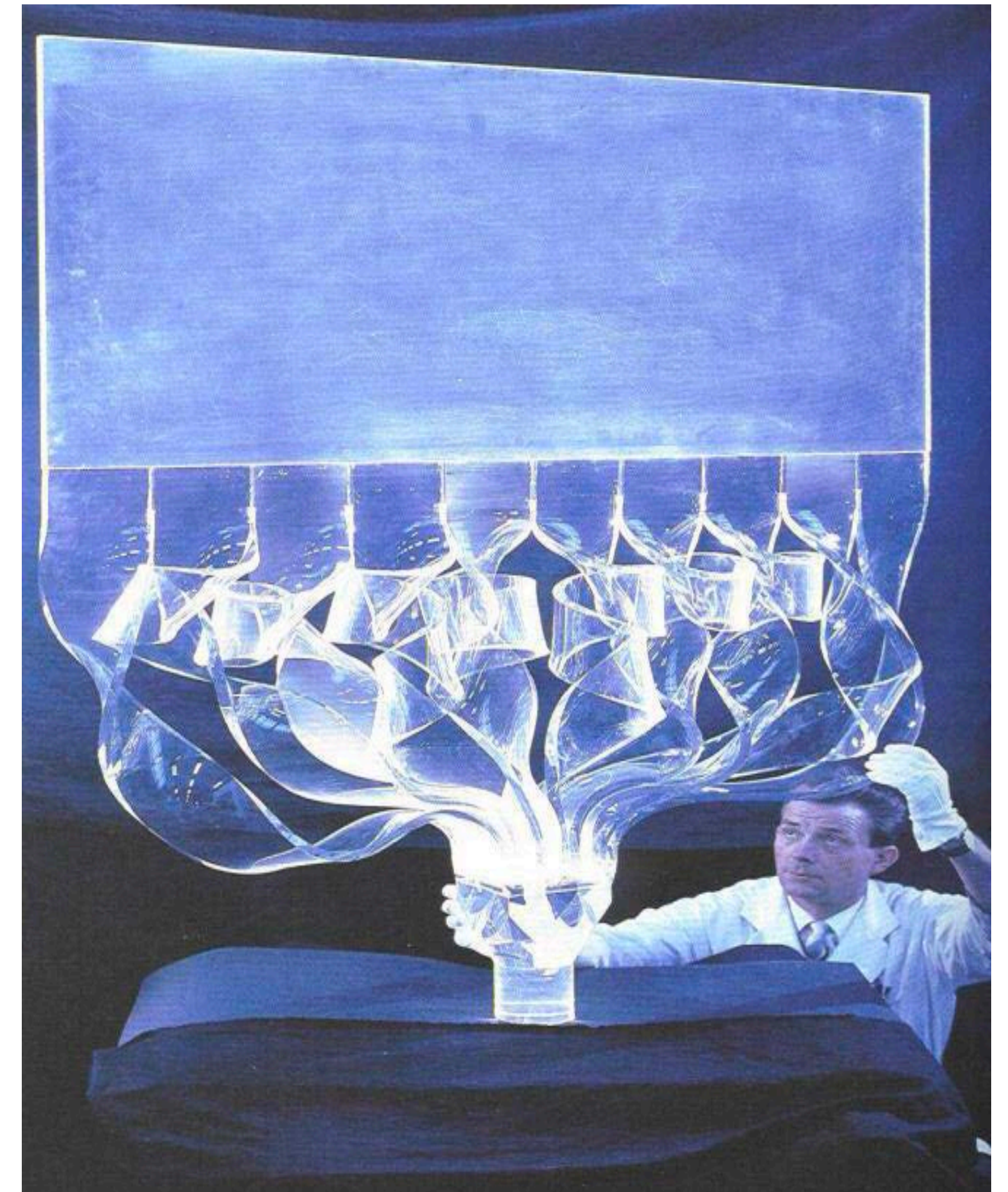
Scintillation light is produced isotropically

Establish the connection between the scintillator and the photoconductor

- high efficiency
- uniformity
- detector geometric constraints and matching to the readout device

General Approaches

- Reflectors
- Light focusing guides
- Wavelength shifters
- Optical fibres



Light Collection - Reflectors

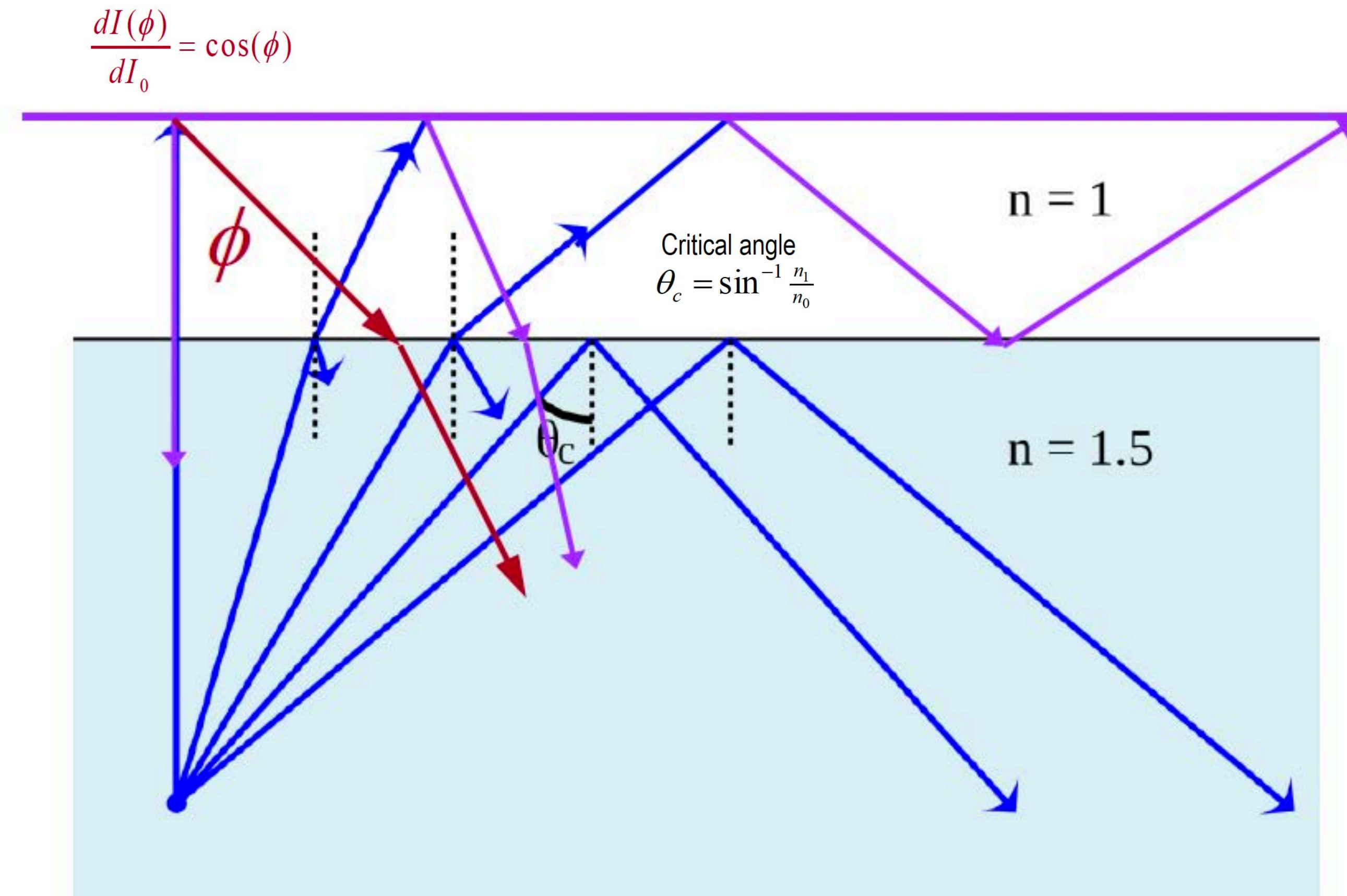
Total internal reflection (simplest approach can often be the best)

Escape fraction

- Can be reflected back into the scintillator by an external reflector
- An air gap between the scintillator surface and reflector preserves the internal reflection
- A specular reflector preserves the angles while a diffuse reflector (Lambert's Law) will change the angle of the exiting rays and improve efficiency

Trapped fraction (internal reflection)

- requires highly polished surfaces
- self-absorption can produce non-uniformities

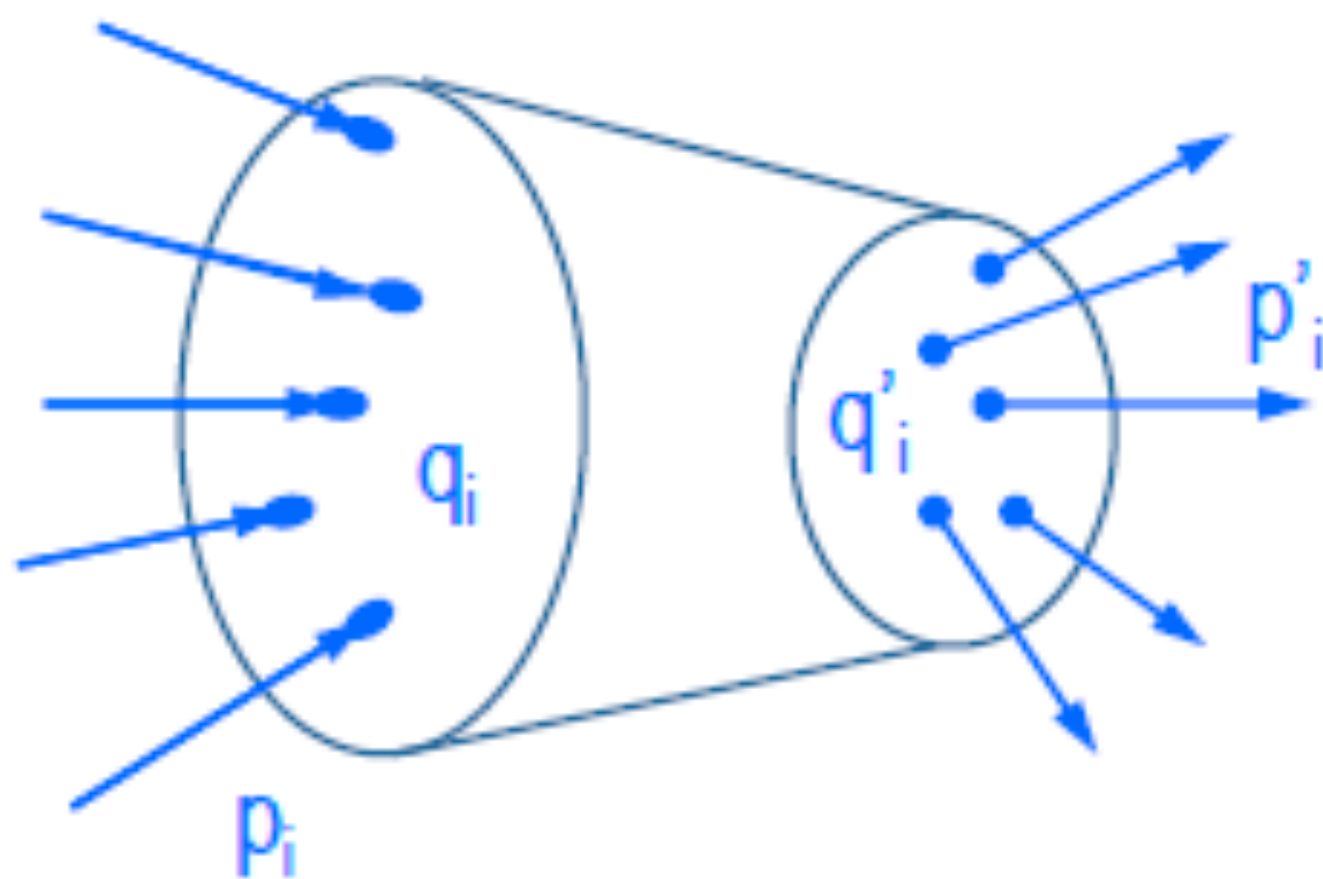


Light Collection - Focusing Guides

Light is concentrated while maintaining a constant phase space density with time

Louisville's Theorem

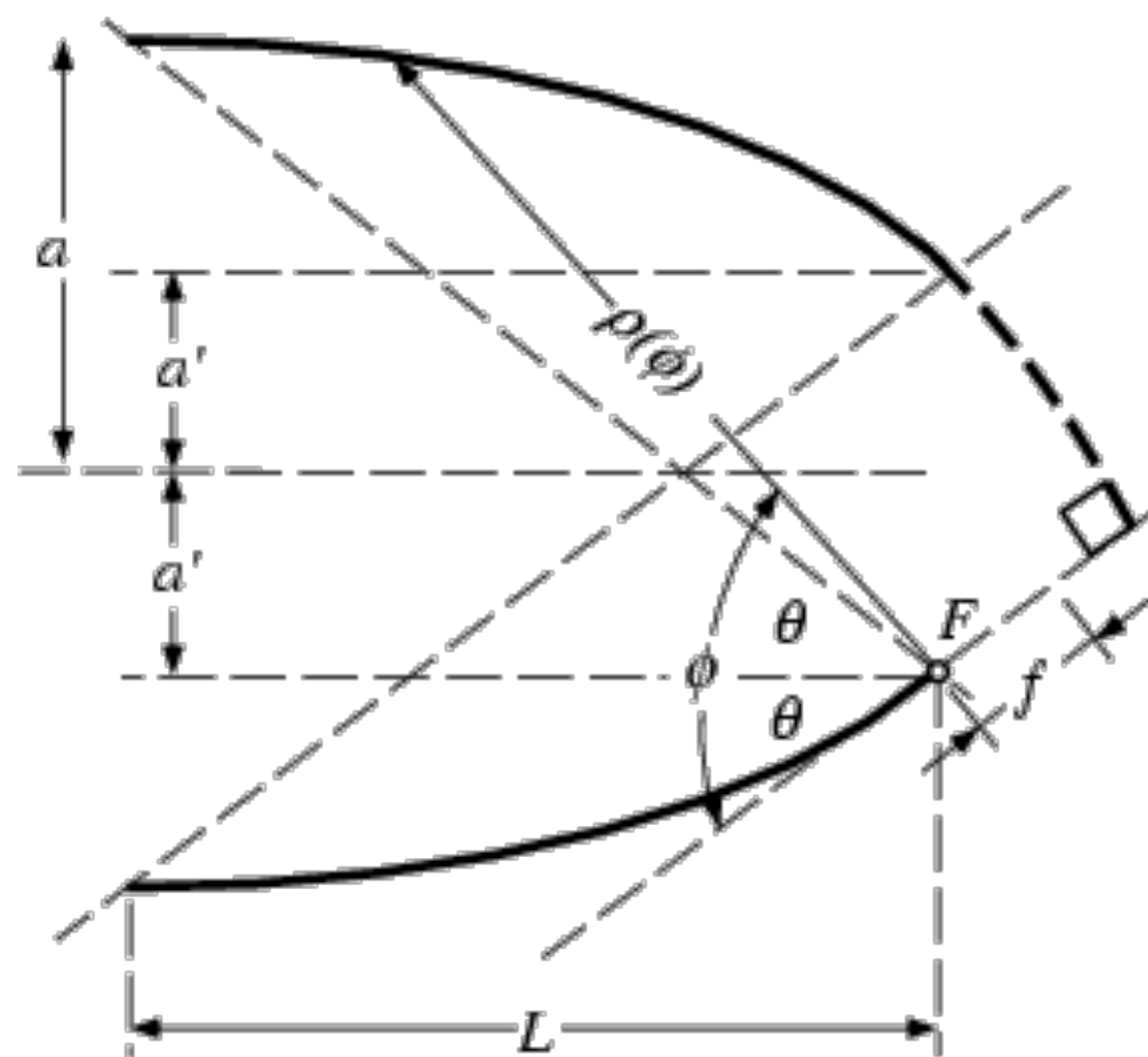
$$\frac{d\rho}{dt} = \sum_{i=1}^n \left(\frac{\partial \rho}{\partial q_i} \dot{q}_i + \frac{\partial \rho}{\partial p_i} \dot{p}_i \right) = 0$$



Area reduction increases divergence

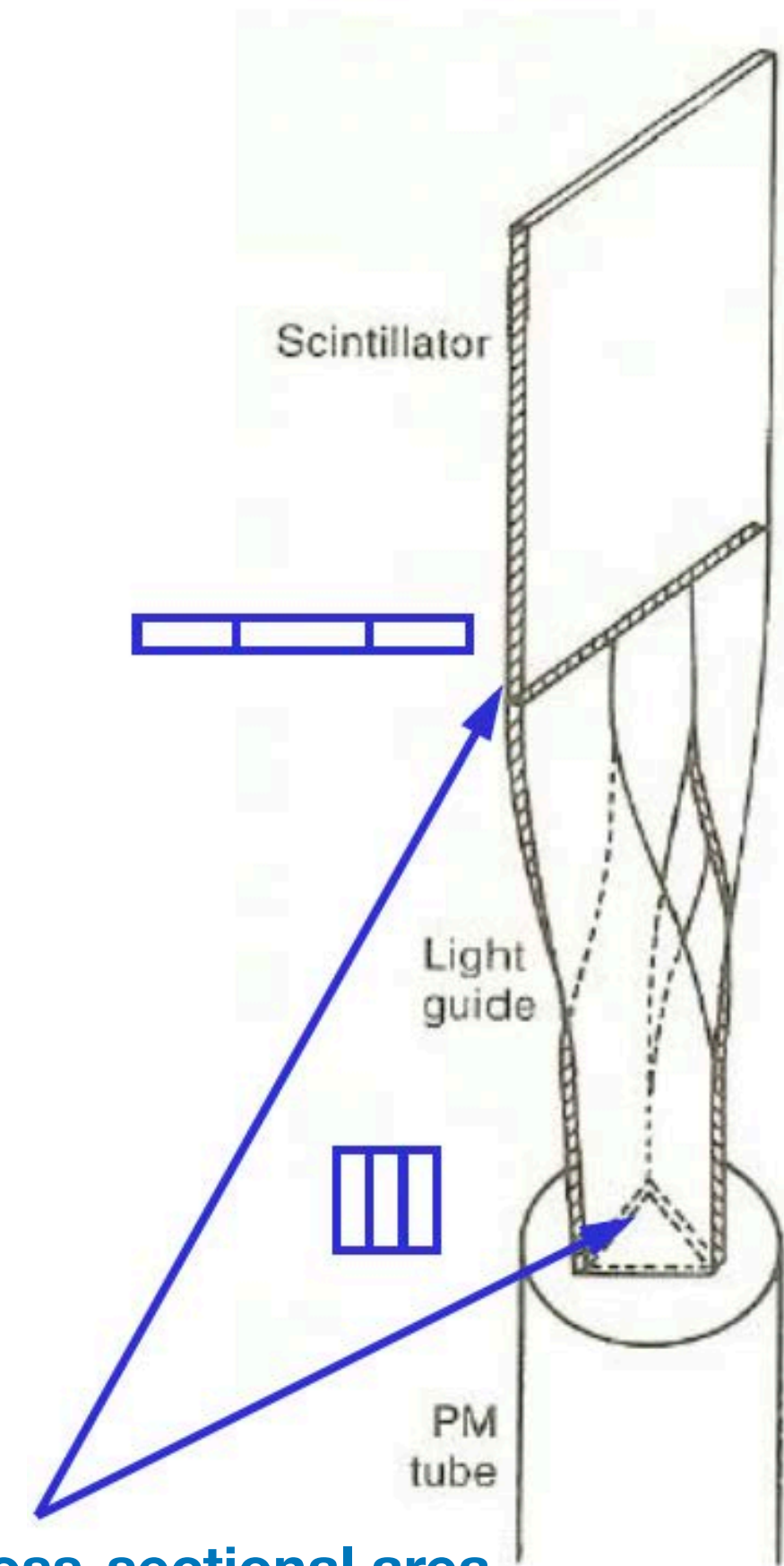
Winston Cone*

Off-axis parabola of revolution
Focuses rays on exit aperture



Preserve initial cross-sectional area

Adiabatic Light Guide



G.Knoll, Radiation Detection and Measurement

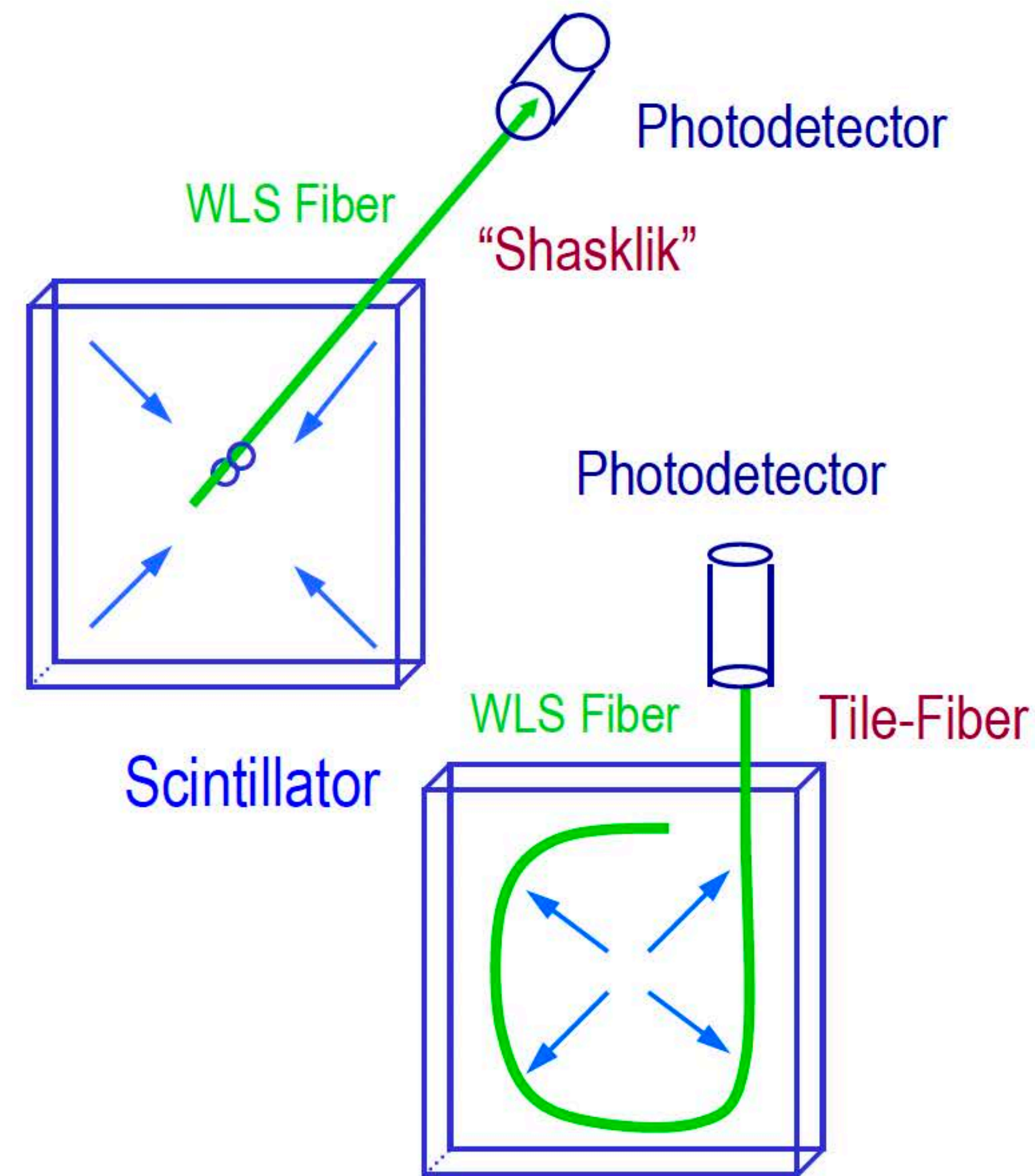
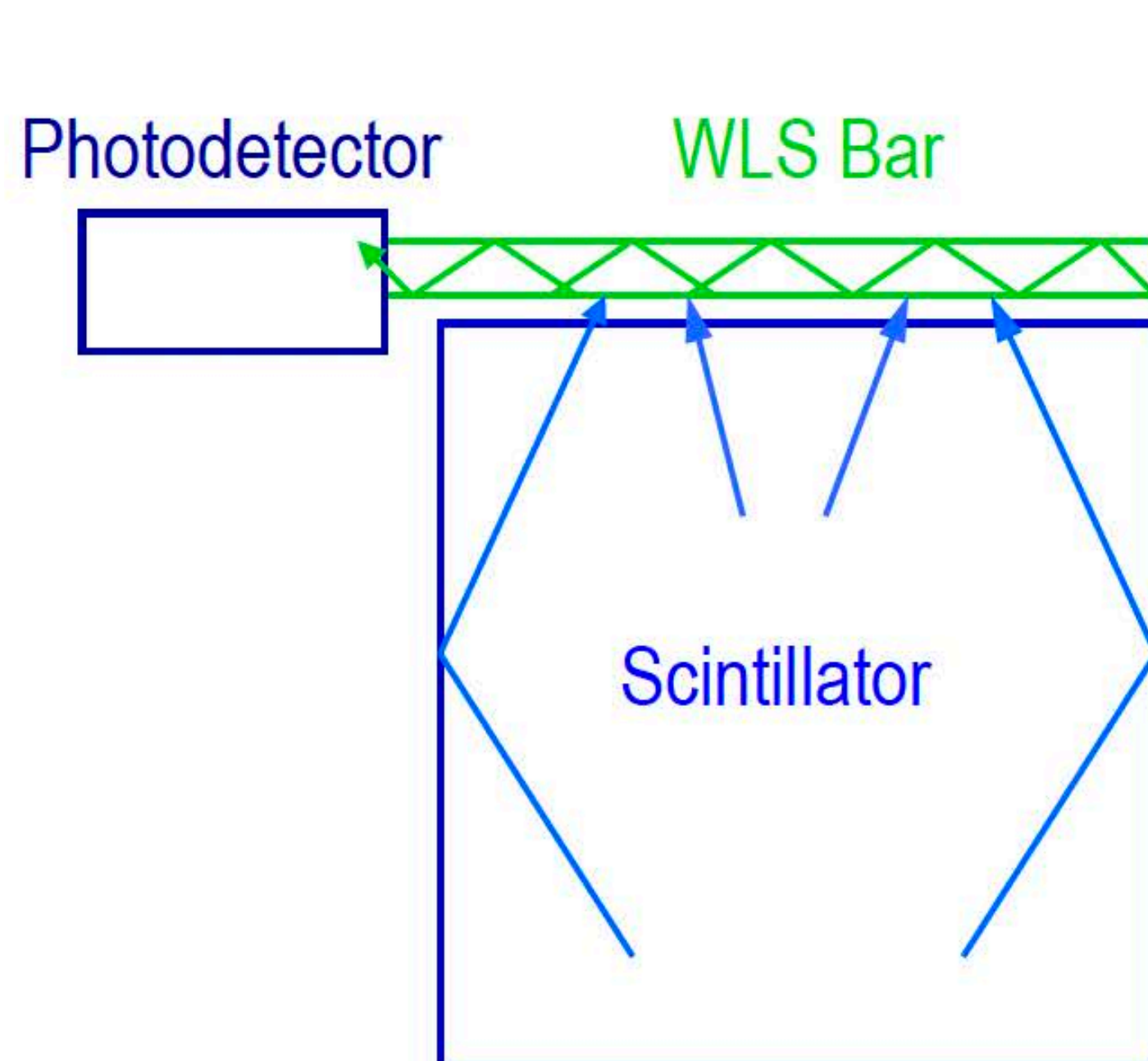
* R. Winston, J. Opt. Sci. Amer. 60 (1970) 245-247

Light Collection - Wavelength Shifters

Wavelength shifter absorbs scintillation light and reemits light as a longer wavelength and in a different direction

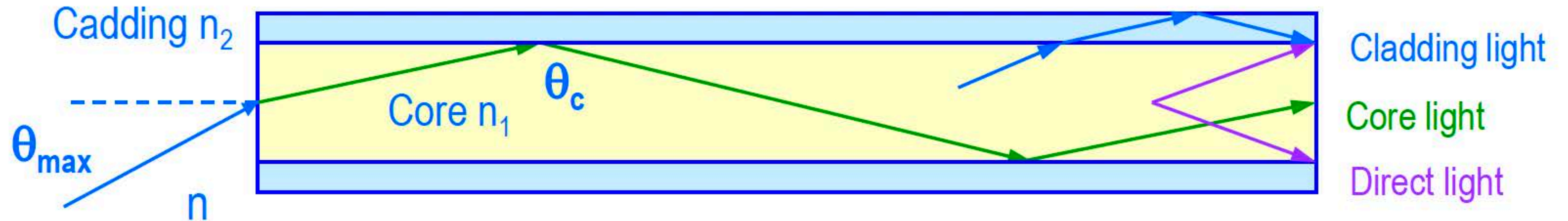
Wavelength shifter element acts as its own light guide, bringing the shifted light to the photodetector

The transfer process may suffer from inefficiencies (scintillator -> wavelength shifter -> absorption -> re-emission -> collection -> photodetector)



Light Collection - Optical Fibres

Fibres may be used both as light guides and as an active (scintillator or Cherenkov) detector material



Small overall trapping fraction

- The light is trapped by internal reflection between the core and cladding ($n_1 > n_2$)
- The Numerical Aperture defines the maximum angle that can be trapped

$$N.A. = n \cdot \sin \theta_{\max} = \sqrt{n_1^2 - n_2^2}$$

Attenuation length is wavelength dependent

Some direct light can escape; some escaping light can be trapped in the cladding

$$F_{fib} = \frac{1}{2} \left(1 - \frac{n_2}{n_1} \right)$$

G.Knoll, Radiation Detection and Measurement

$n_1 = 1.58$ (polystyrene)

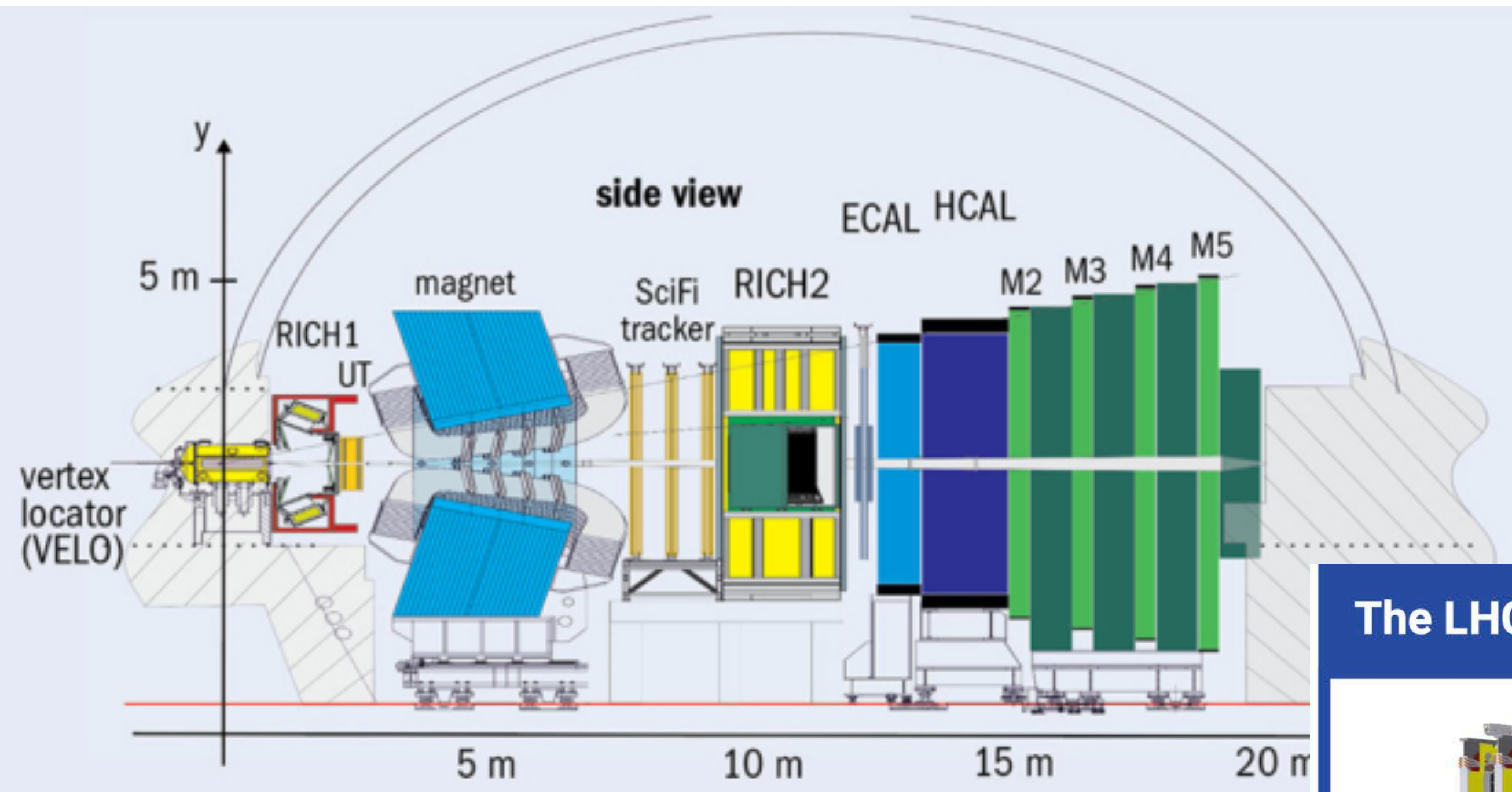
$n_2 = 1.49$ (PMMA)

$F_{fib} \sim 3\%$

$\times 2 \sim 6\%$ if use reflector or read out both ends

Light Collection - Optical Fibres

IceCube Fibre Optical Module



The LHCb Scintillating Fibre Tracker



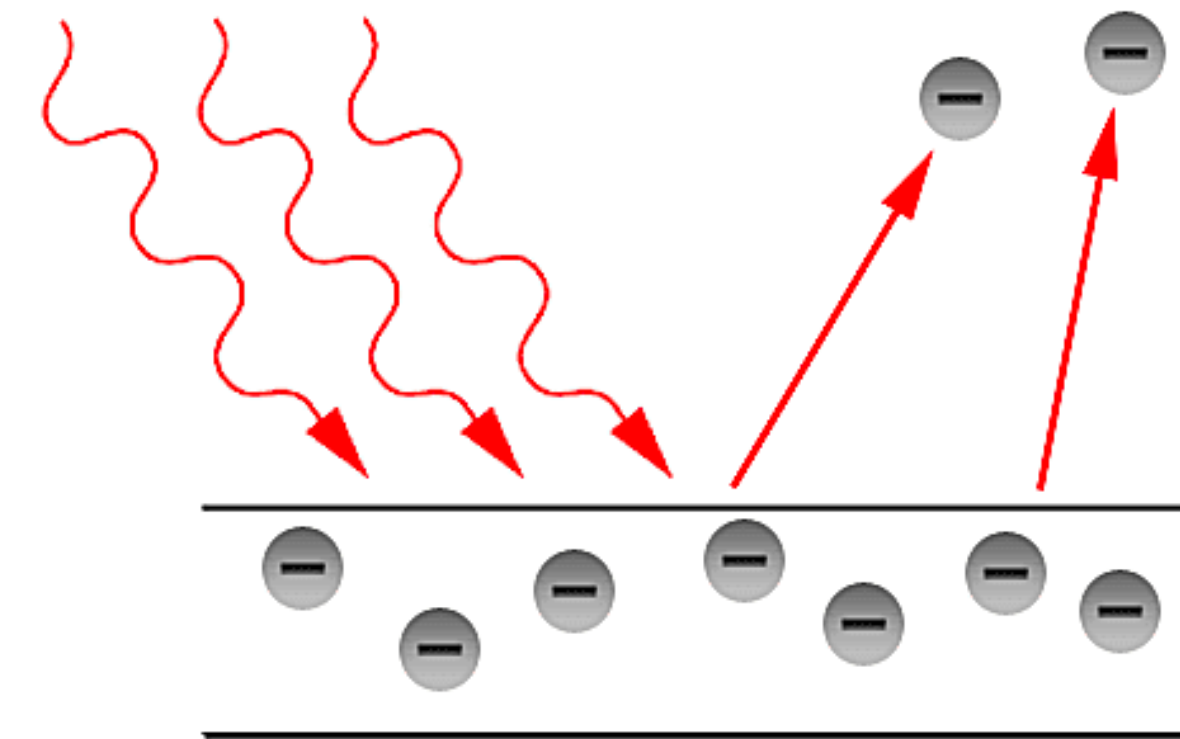
- Three stations with four layers each, covering a total active area of 340 m²
- Eight 2.5 m long six-layer fibre mats per module
 - 250 μm diameter scintillating fibres
 - 11 000 km of fibre used throughout the detector
- Readout by silicon photomultiplier (SiPM) arrays
 - 524 288 readout channels in total
 - Cooled to -40 °C to mitigate radiation damage
- Signal processing with 40 MHz readout electronics
 - Custom ASIC (PACIFIC) for analogue processing & digitisation with three comparators per channel
 - Online zero-suppression & clustering on FPGAs

Have reached the stage where we produced and collected light from an ionizing interaction we want to study
Final goal -> convert this to a digital signal.

Basic elements:

Photoelectric effect

- converts photons to photoelectrons
- performance driven by the quantum efficiency

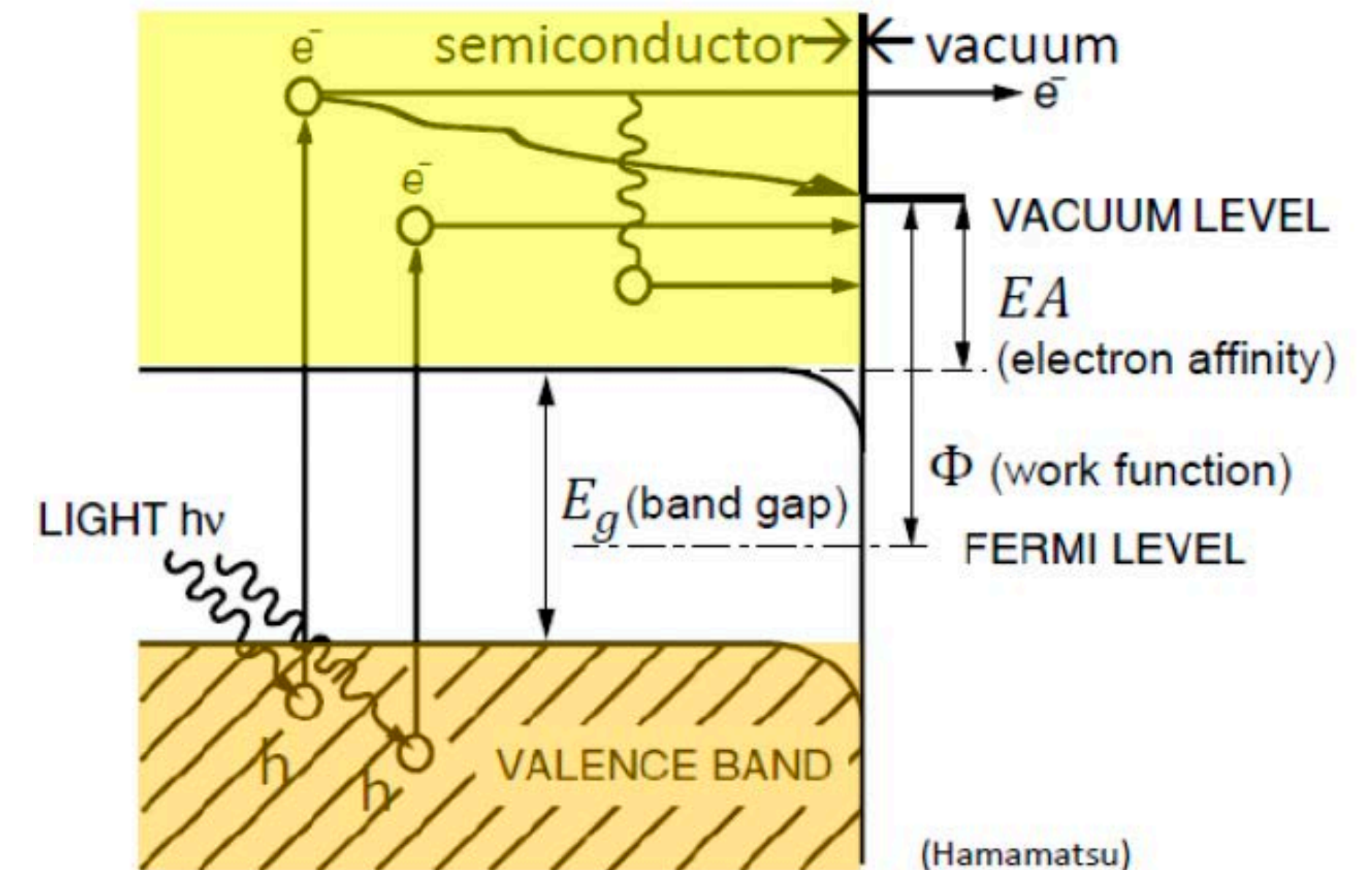
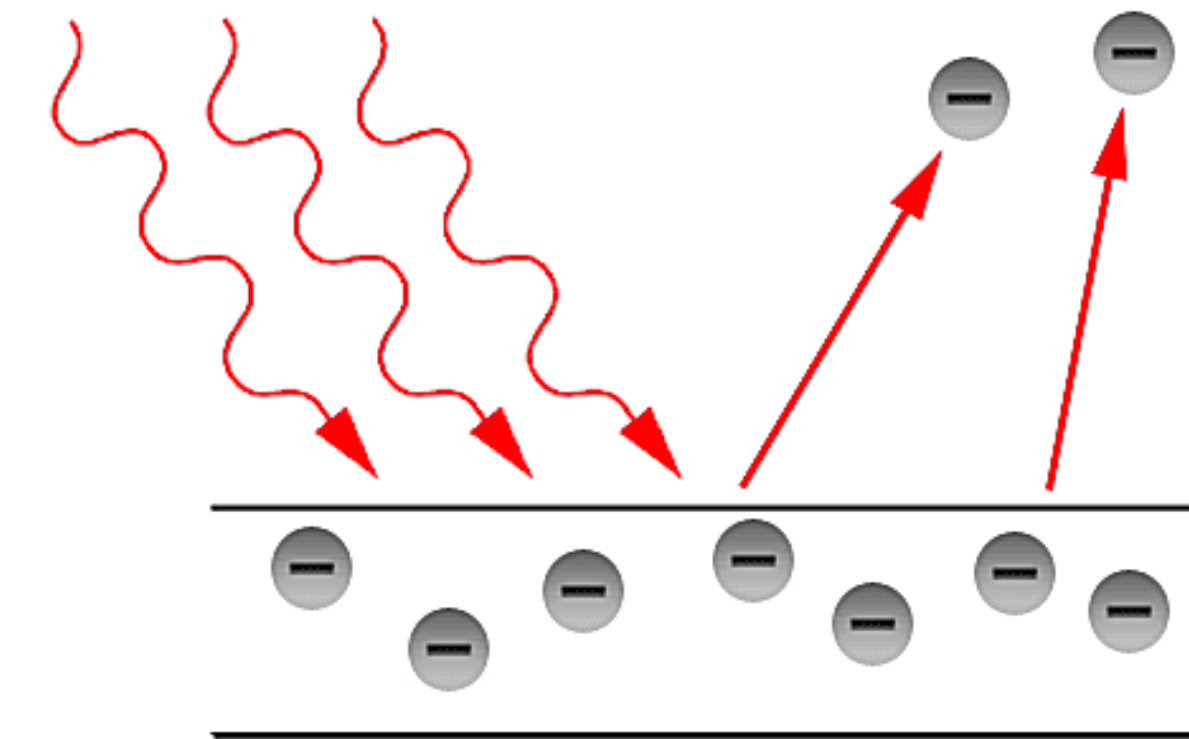


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Photon energy:

$$E_\gamma = h\nu = \frac{hc}{\lambda} \approx \frac{1239 \text{ eV} \cdot \text{nm}}{\lambda}$$

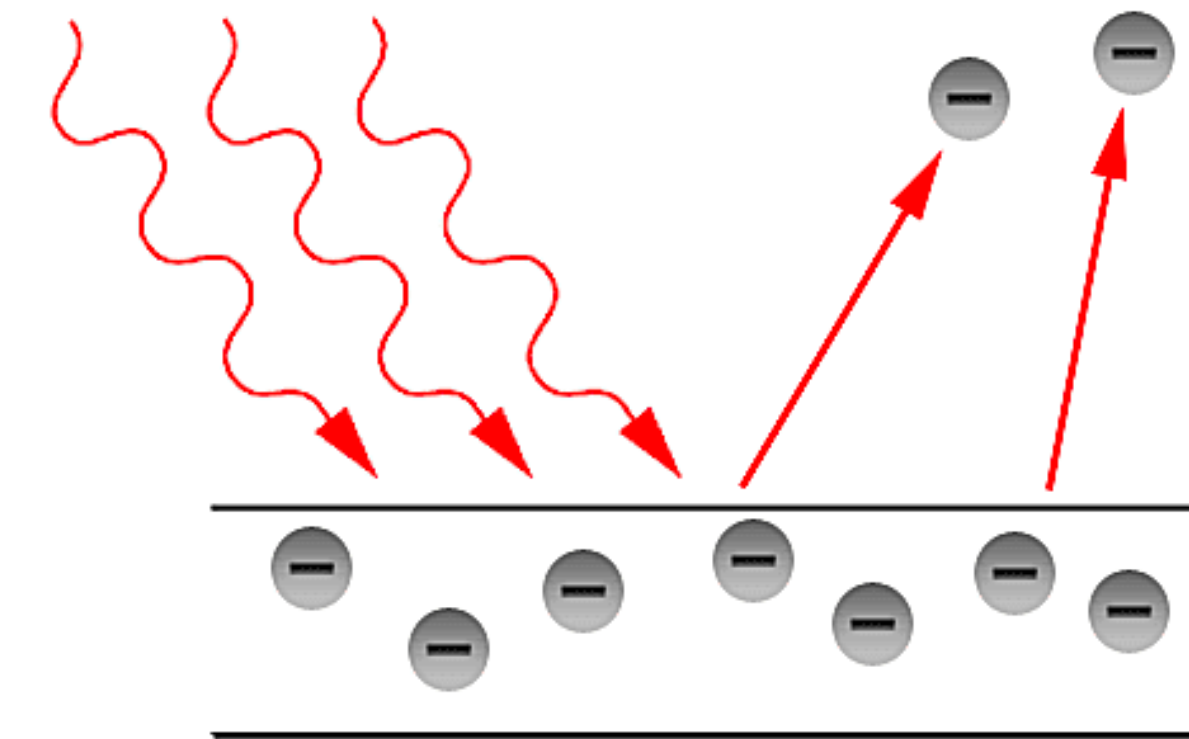
Visible range 400 nm – 780 nm
→ 3.1 - 1.6 eV

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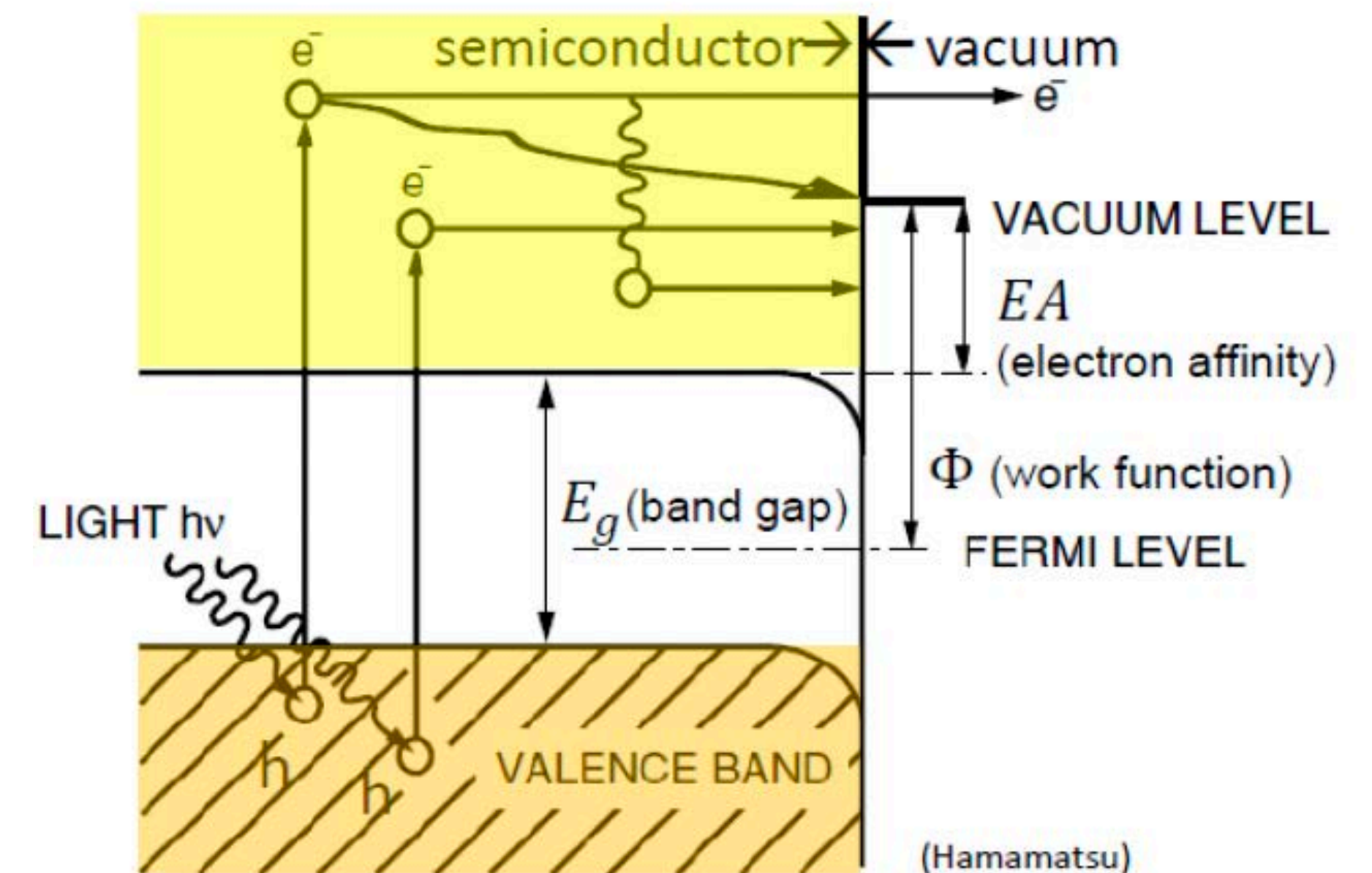
Charge amplification

- via gain or direct conversion (device dependent)
- driven by the signal-to-noise ratio

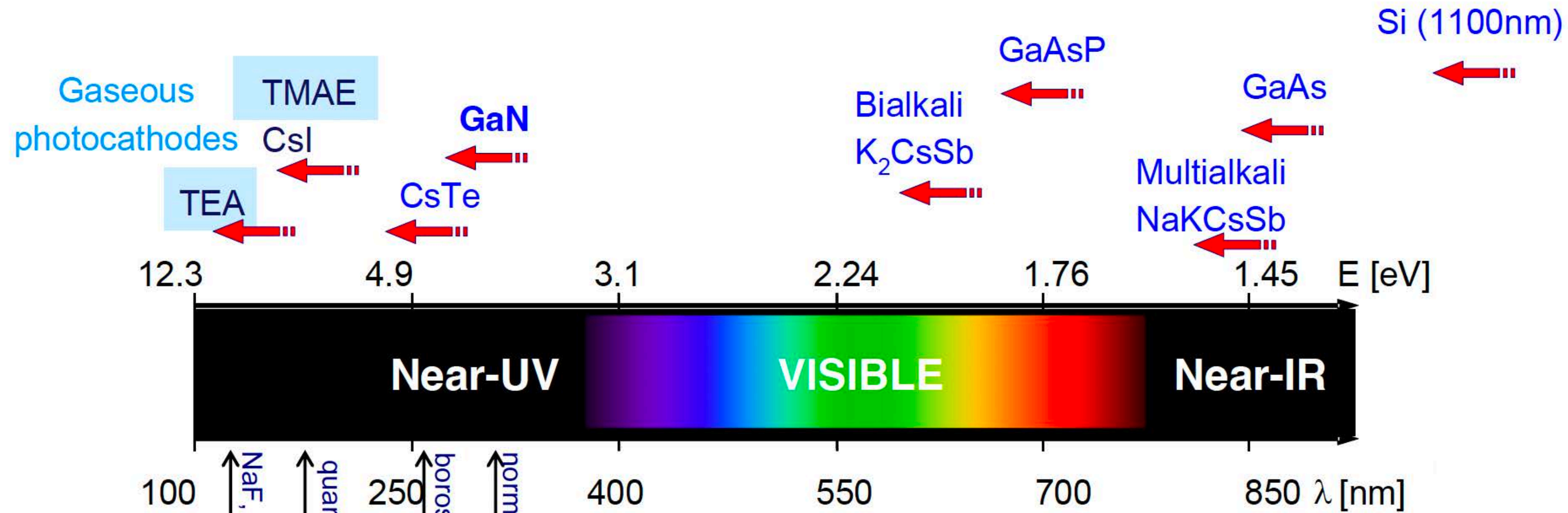
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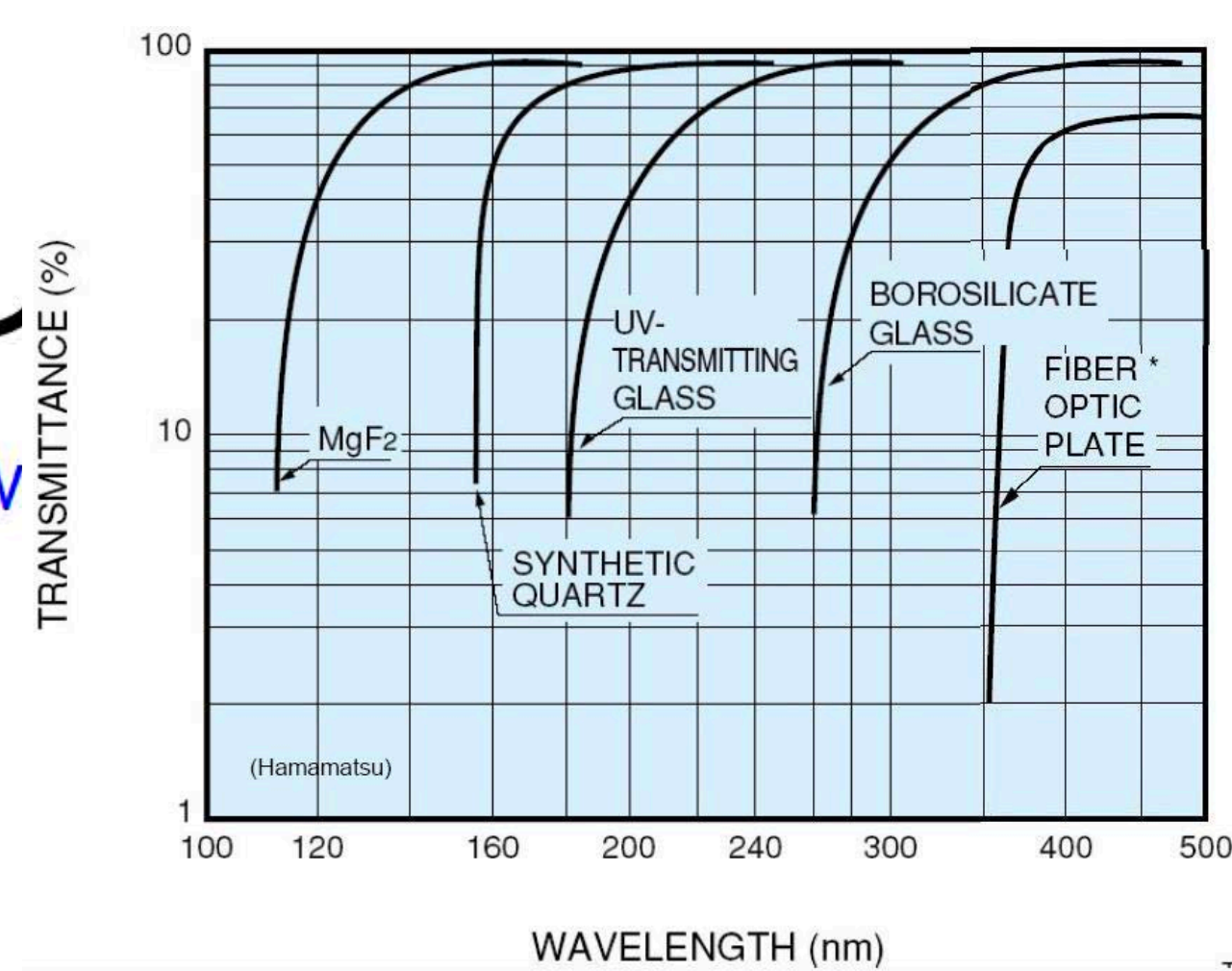


Photodetectors - Photocathodes



100 \uparrow NaF, MgF_2 , LiF, CaF_2
 \uparrow quartz
 250 \uparrow borosilicate glass
 \uparrow normal window glass

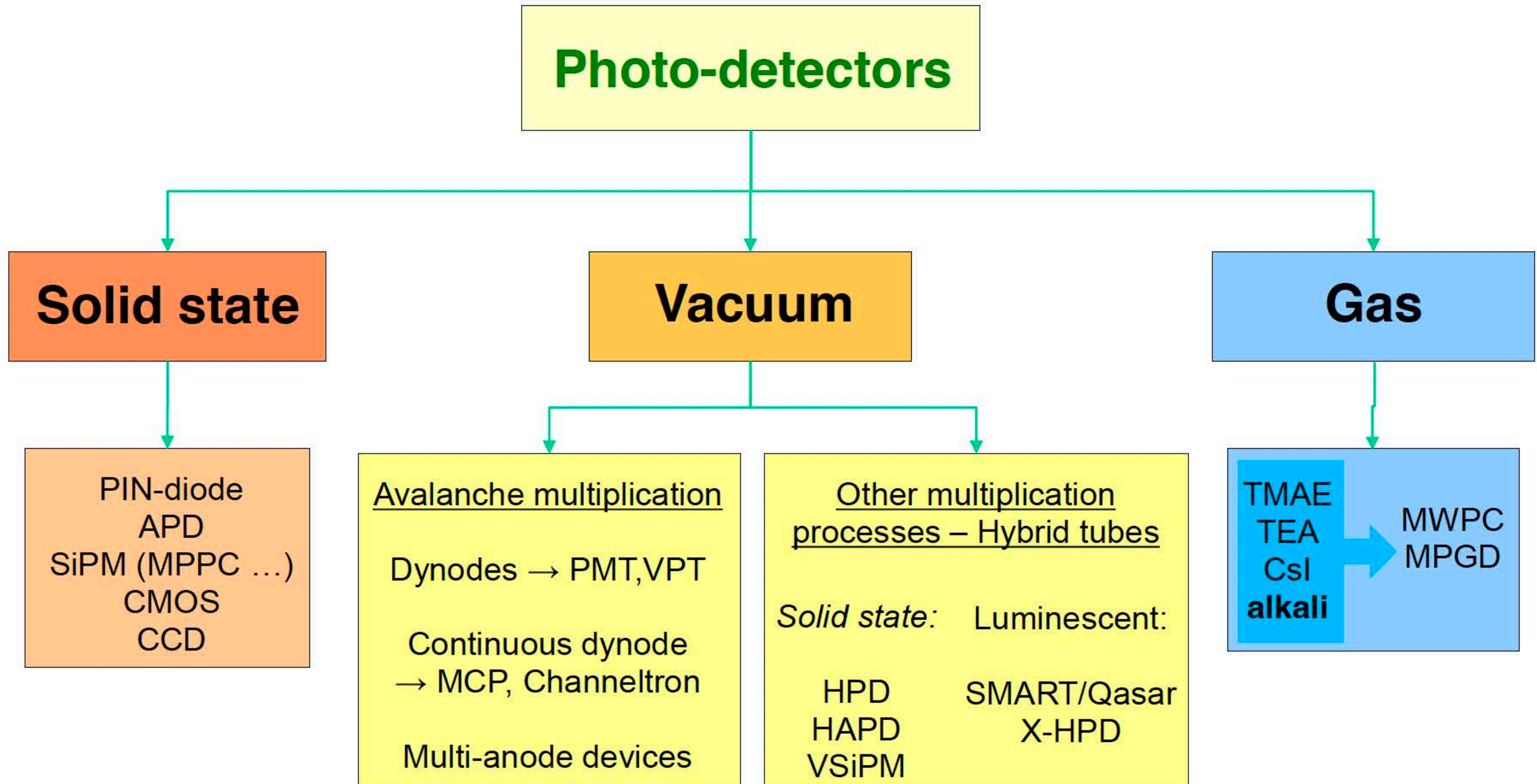
Cut-off limits of window



Nearly all photosensitive materials are very reactive (alkali metals)

Operation typically only in vacuum or high purity gas

Exceptions include CsI and Is



Photodetectors - Overview (properties)

Type	λ (nm)	$\epsilon_Q \epsilon_C$	Gain	Risetime (ns)	Area (mm ²)	1-p.e noise (Hz)	HV (V)	Price (USD)
PMT*	115–1700	0.15–0.25	10^3 – 10^7	0.7–10	10^2 – 10^5	10 – 10^4	500–3000	100–5000
MCP*	100–650	0.01–0.10	10^3 – 10^7	0.15–0.3	10^2 – 10^4	0.1–200	500–3500	10–6000
HPD*	115–850	0.1–0.3	10^3 – 10^4	7	10^2 – 10^5	10 – 10^3	$\sim 2 \times 10^4$	~ 600
GPM*	115–500	0.15–0.3	10^3 – 10^6	$O(0.1)$	$O(10)$	10 – 10^3	300–2000	$O(10)$
APD	300–1700	~ 0.7	10 – 10^8	$O(1)$	10 – 10^3	1 – 10^3	400–1400	$O(100)$
PPD	320–900	0.15–0.3	10^5 – 10^6	~ 1	1–10	$O(10^6)$	30–60	$O(100)$

Review of Particle Physics
Journal of Physics G, Vol 37, No 7A (2010)

PMT = Photomultiplier tube (sensitive to magnetic fields)

APD = Avalanche Photodiode

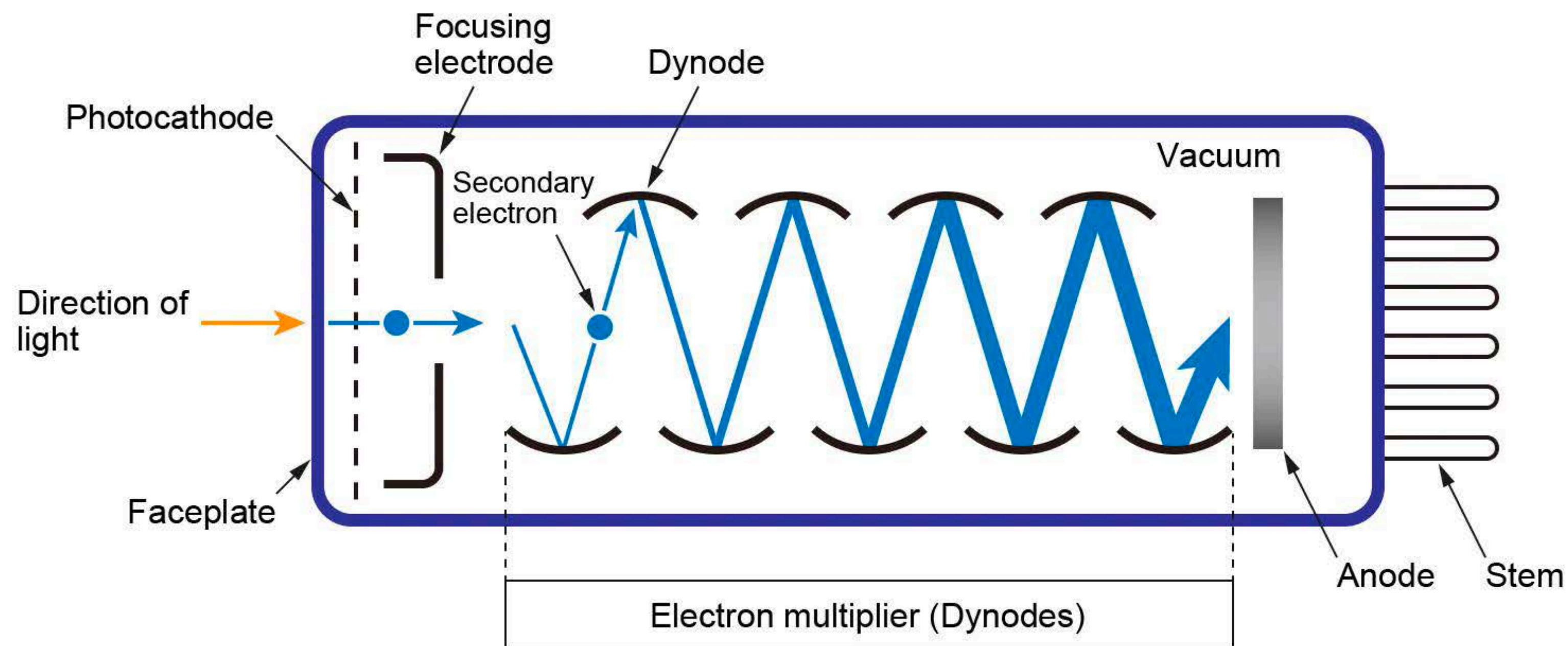
PPD = Pixellated Photon Detector (**SiPM**, GPMT, MPPC)

MCP = Multi-channel Plate Detector (dense, fast small diameter charge amplifiers)

HPD = Hybrid Photodiode Detector (vacuum PMT w/silicon sensor)

GPM = Gaseous Photon Detector (large area solid and gaseous photocathode)

Photodetectors - Photomultiplier tube (PMT)



Photon releases an electron from the photocathode (quantum efficiency)

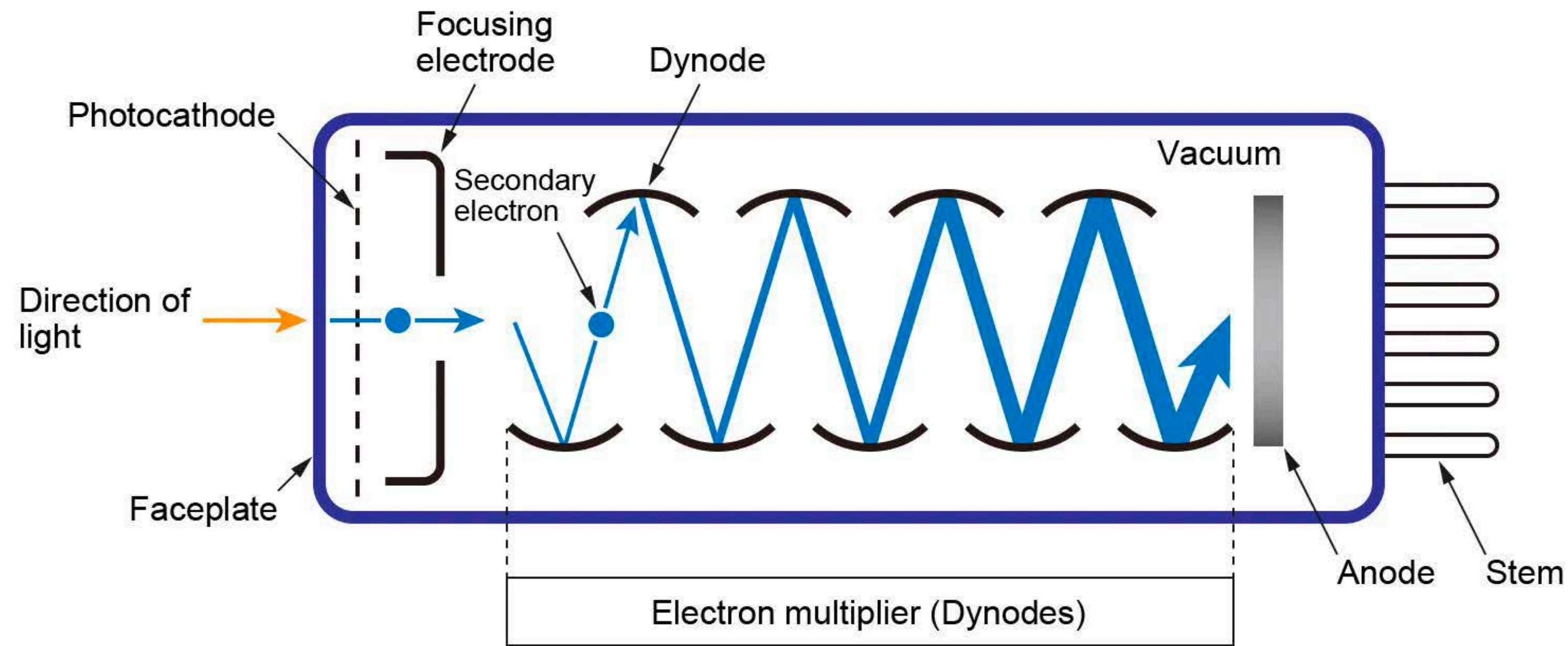
Collection of photoelectrons by the first dynode

Secondary emission from (N) dynodes in the chain (gain is a function of incoming electron energy)

Amplification factors of $O(e6)$ achievable



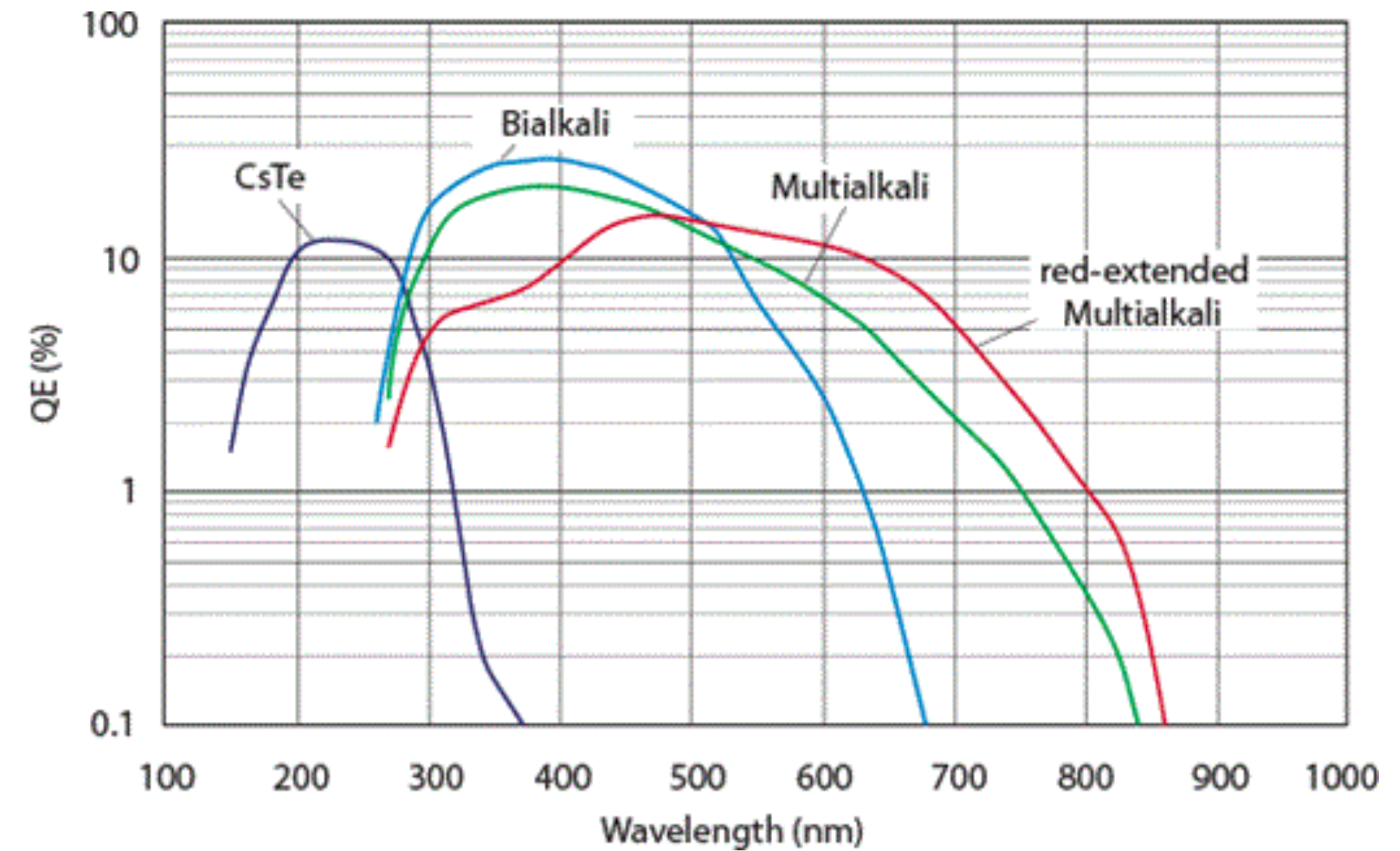
Photodetectors - Photomultiplier tube (PMT)



Monoalkali - CsI, CsTe, Kir

Bialkali - SbRbCs₃

Multialkali - SbNaKCs₃



Photodetectors - Silicon (Avalanche Photodiodes - APDs)

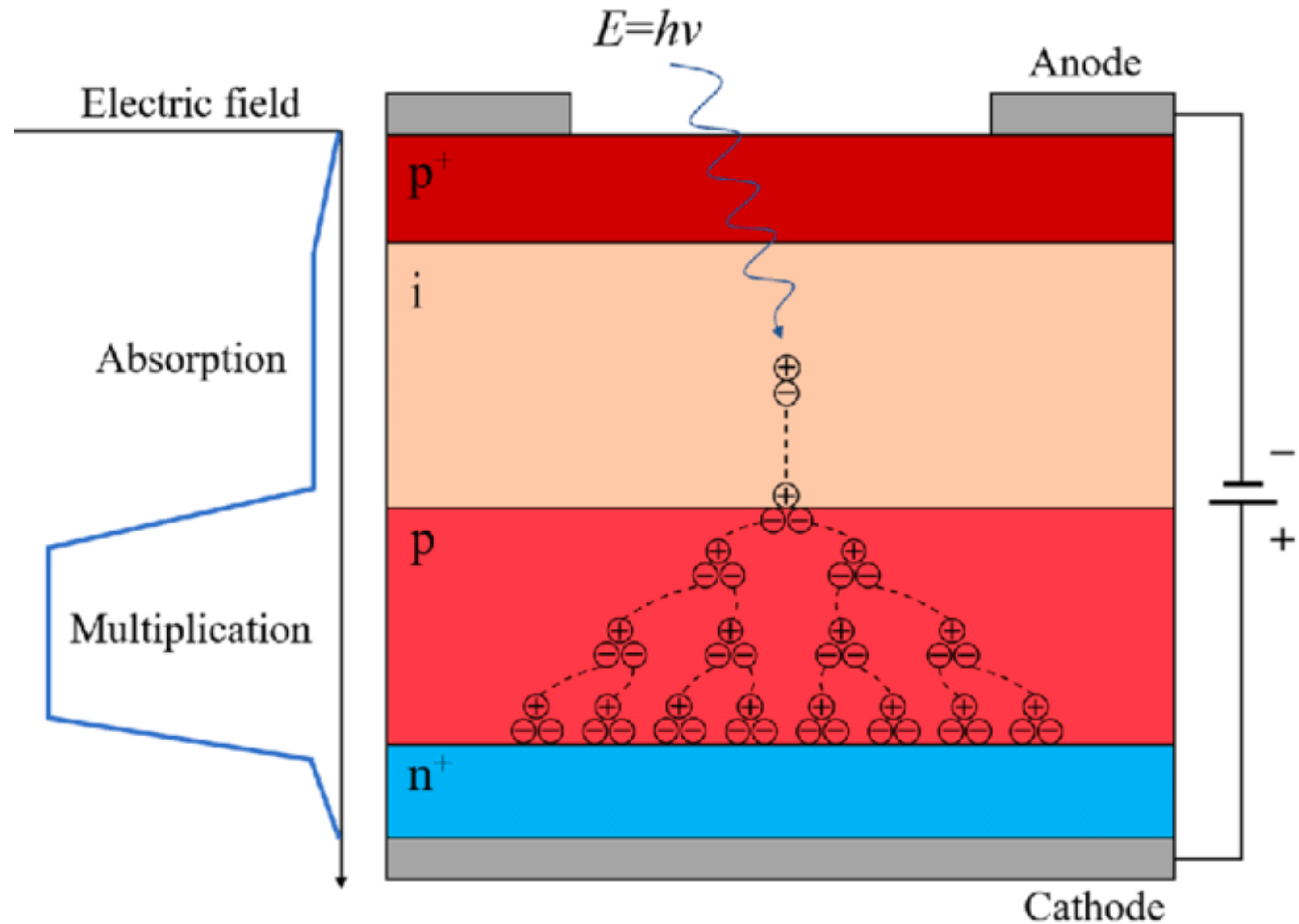
Formed utilizing multiple p-n junctions ->

A modification of the traditional P (hole-rich)-I-N (electron-rich) photodiode where the (I)ntrinsic region typically acts as a large charge storage region. Operation in reverse bias this region instead behaves as a small capacitor = "open switch")

In reverse bias, high doping creates a strong electric field

Electrons moved to the conduction band drift to the high-field region where an avalanche amplification takes place (50 - 500x)

Signal is obtained by an electric (can be further amplified).



Photodetectors - Silicon (Avalanche Photodiodes - APDs)

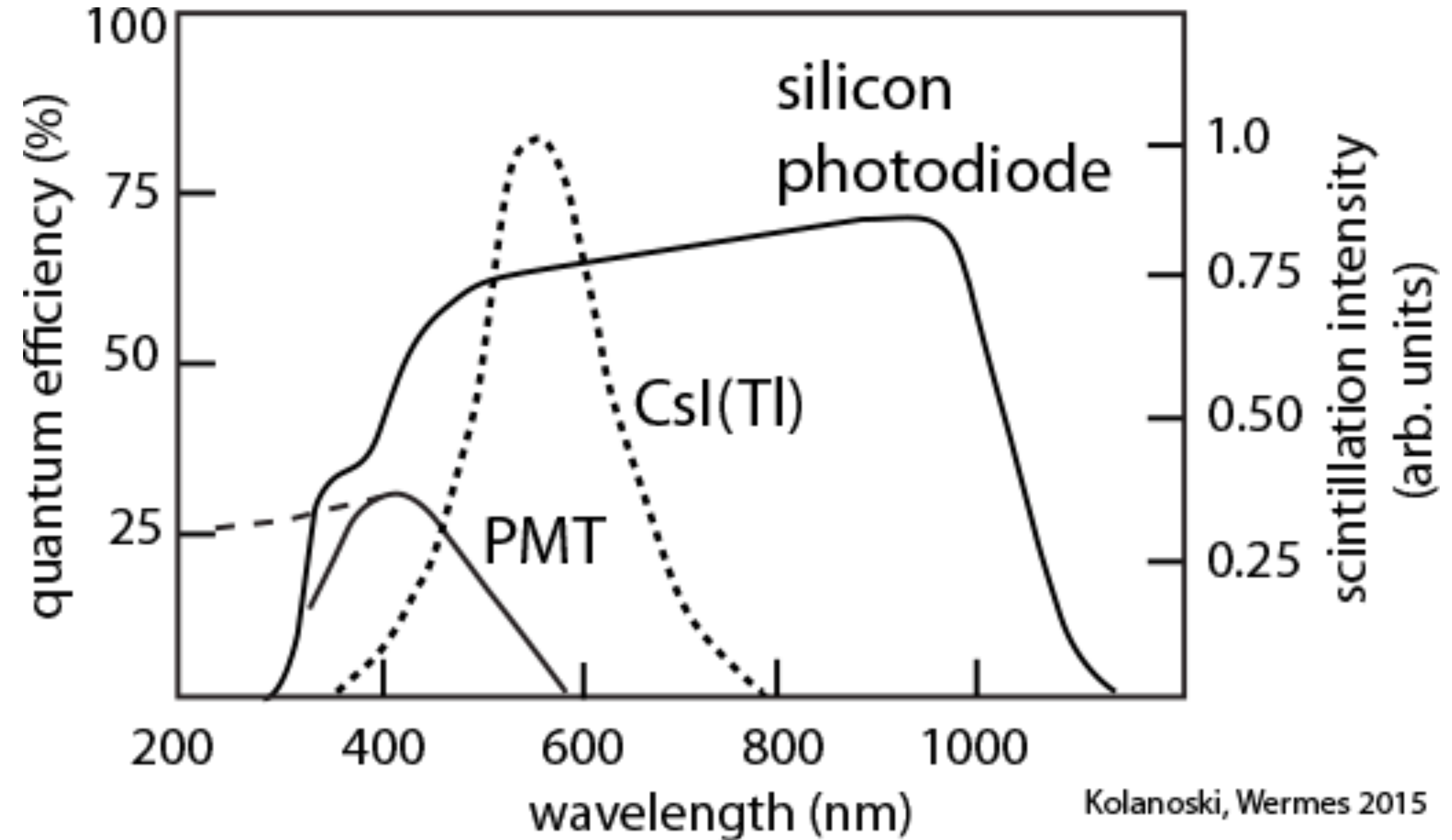
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Kolanoski, Wermes 2015

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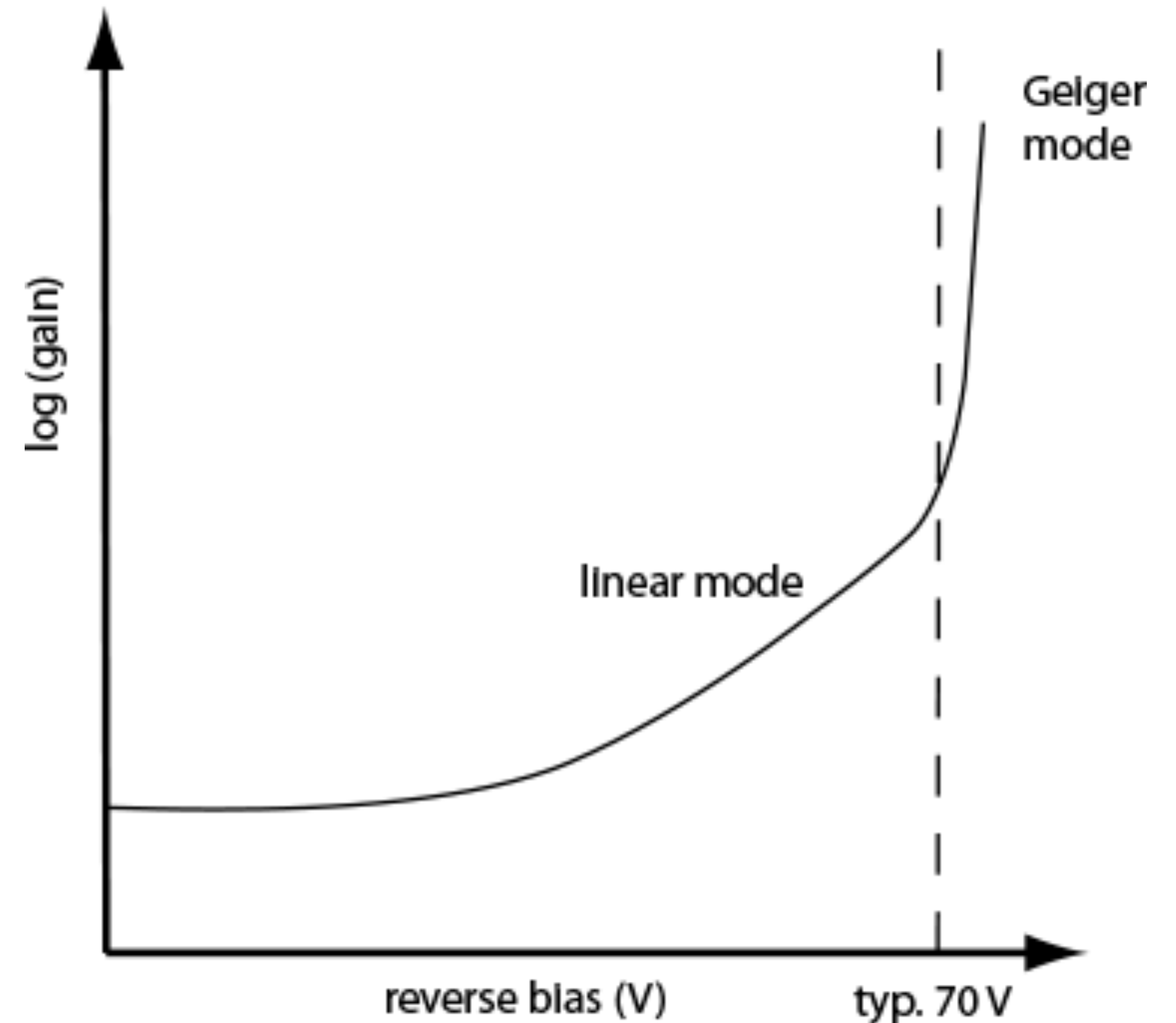
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Photodetectors - Silicon (SiPMs)

Arrays of Single Photon Avalanche Diodes (SPADs) biased slightly above the breakdown voltage (Geiger mode) -> single particle (including photons) can trigger an avalanche.

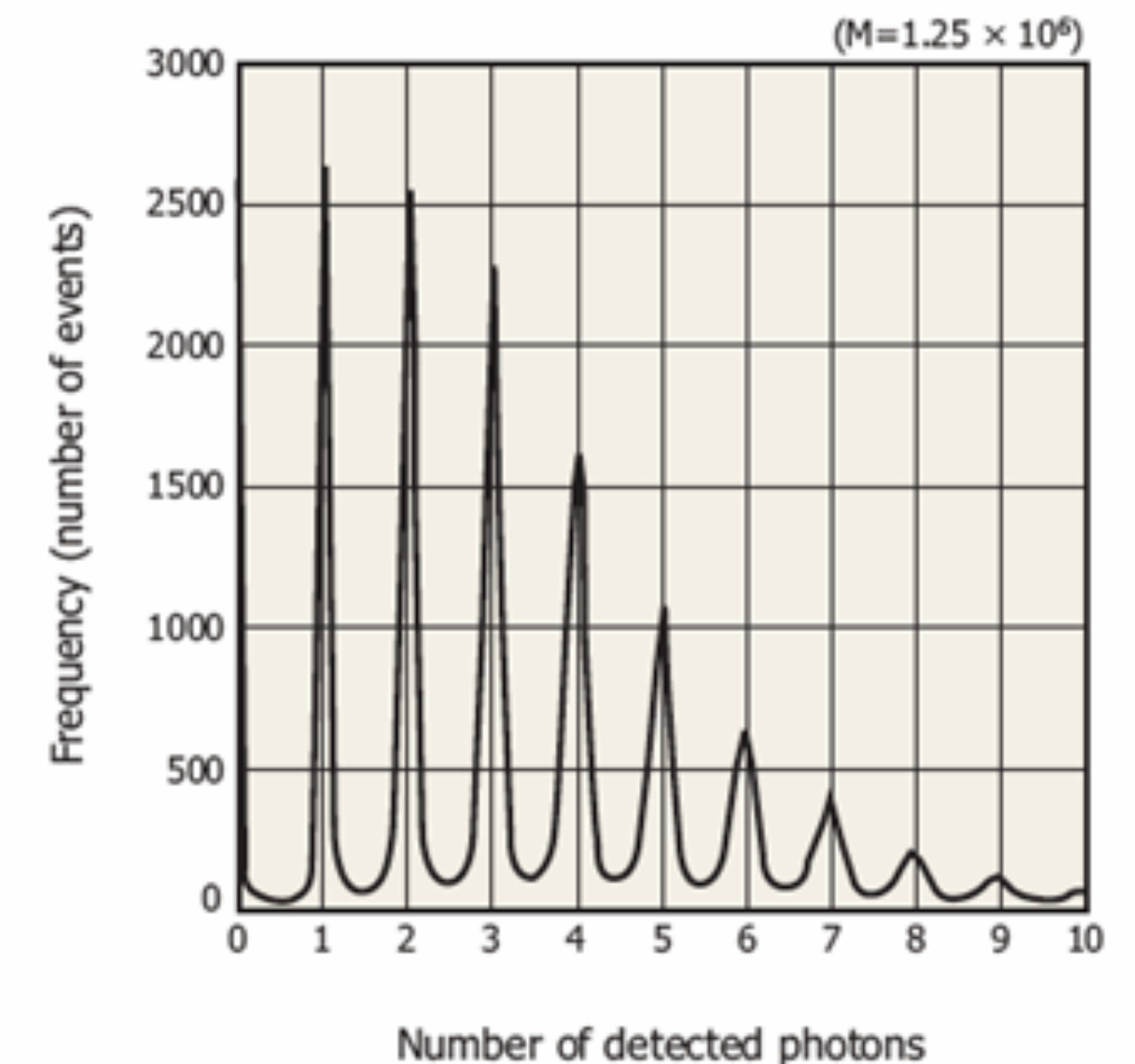
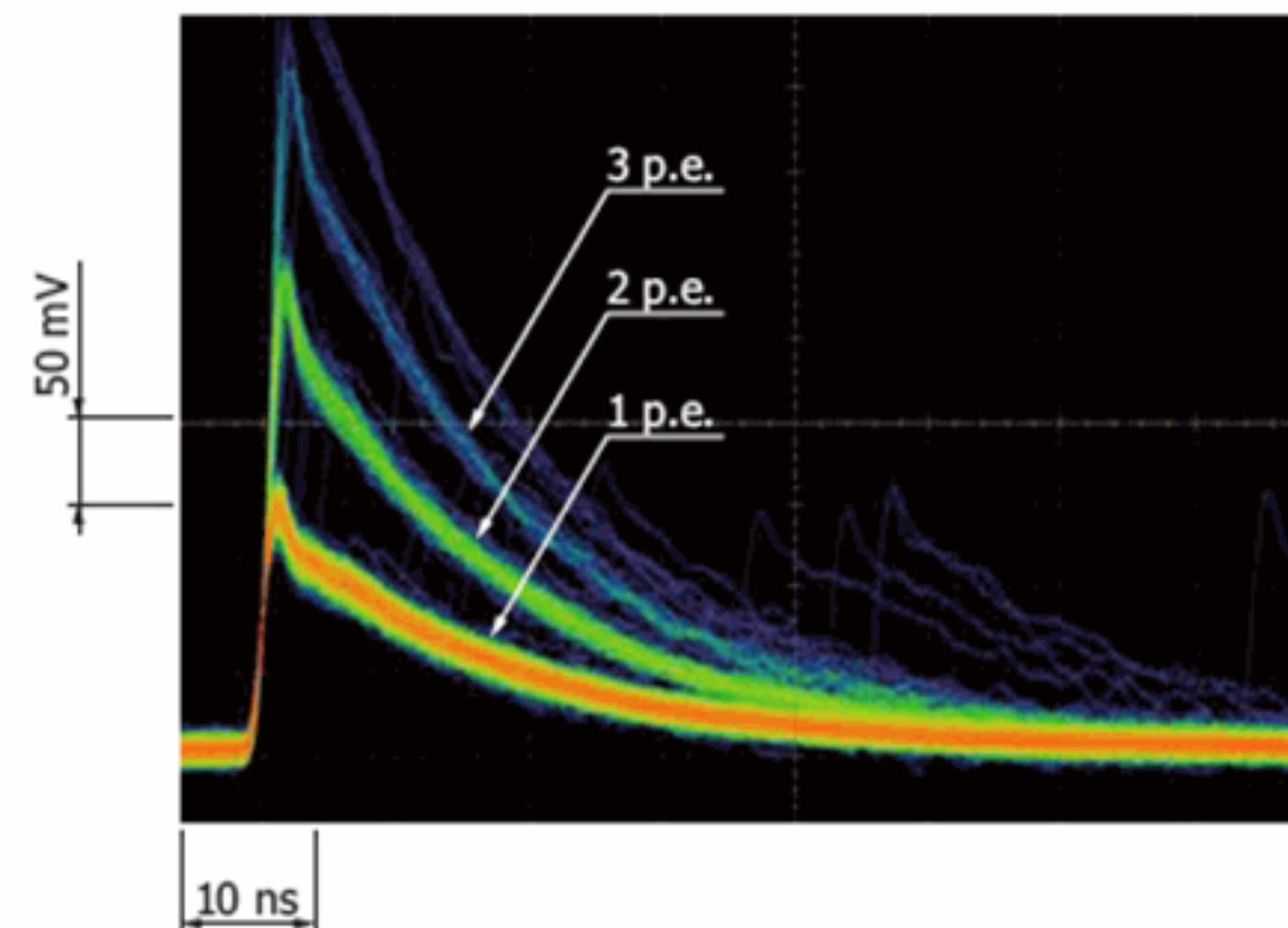
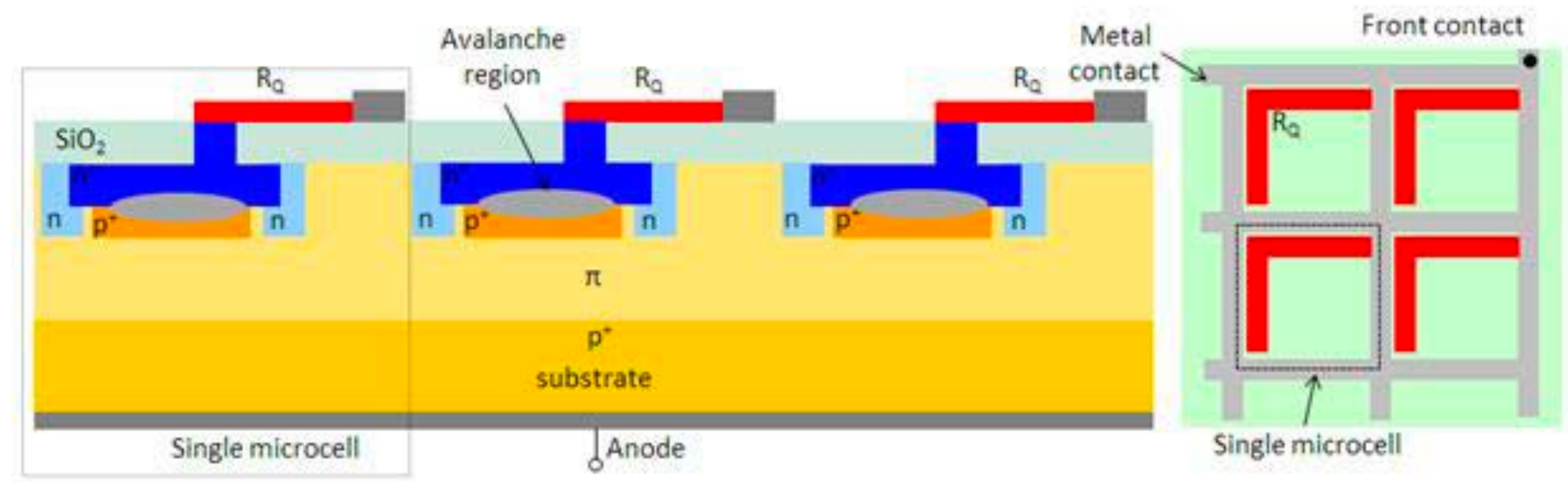
Cells are binary (yes/no); signal proportional to light intensity

Advantages

- High gains (similar to PMTs) can be realized
- Excellent single photon resolution
- Insensitive to magnetic fields

Disadvantages

- Non-linear output at high incident flux (pixel saturation)
- High noise rates (thermal carriers) ~ 100 kHz/sq-mm
- Large temperature dependence
- Susceptible to neutron damage

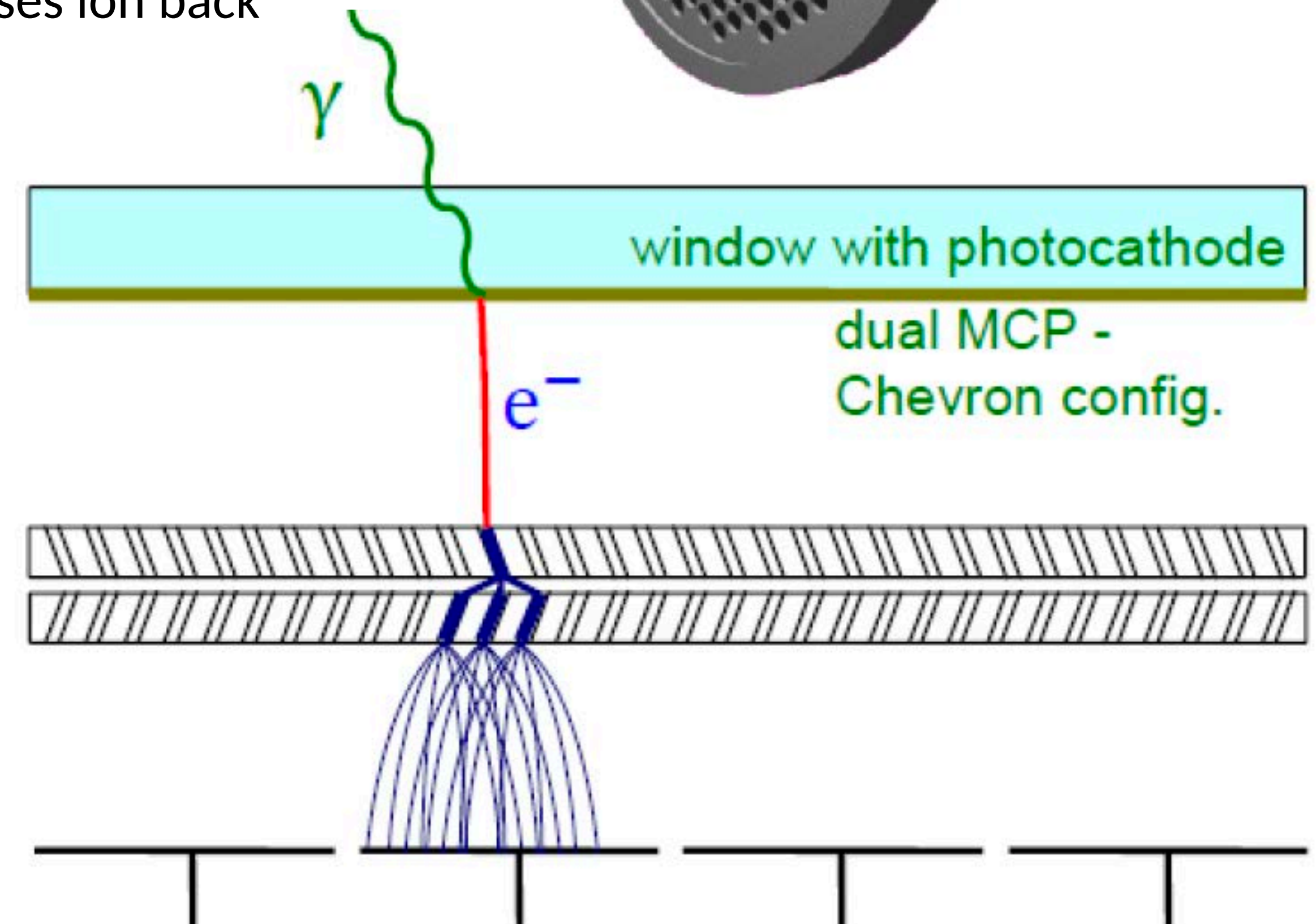
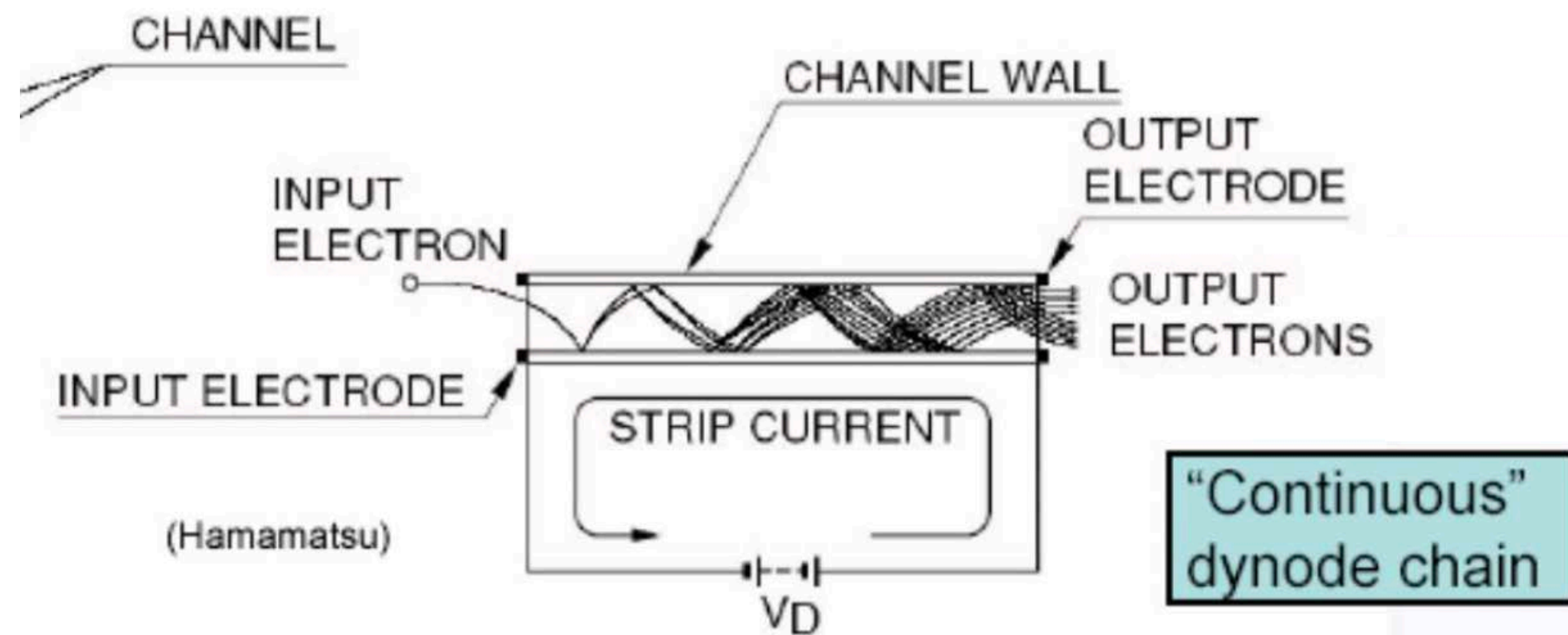
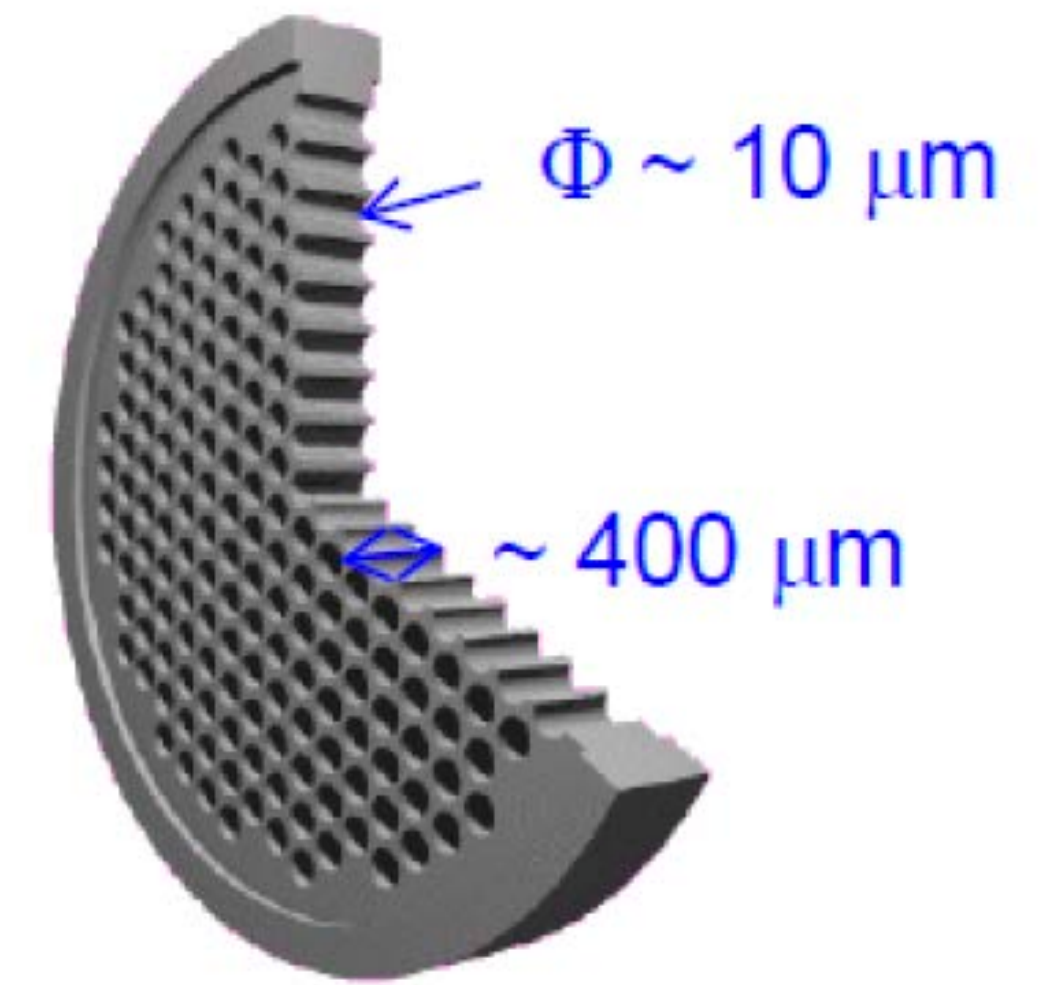


Photodetectors - Silicon (Multi-channel Plate Detectors - MCPs)

A thin (Pb-)glass plate with array of holes forming a continuous dynode structure.

The avalanche process is highly contained providing a very small transit time spread -> very fast timing and excellent timing resolutions (possible to identify arrivals of individual photoelectrons)

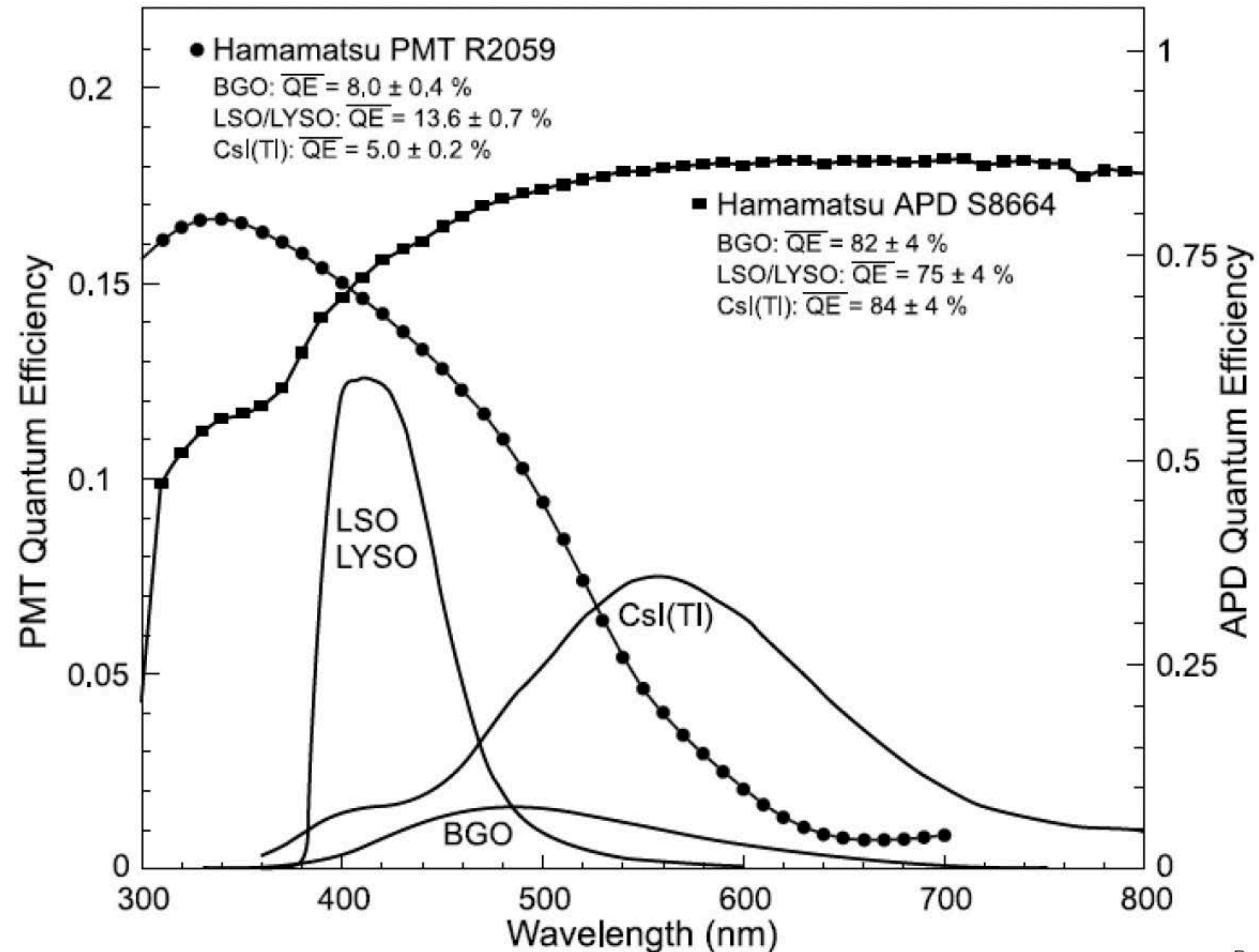
“Chevron configuration” improves radiation hardness and suppresses ion back streaming -> improved gain



Scintillator Detector Systems - Quantum Efficiency

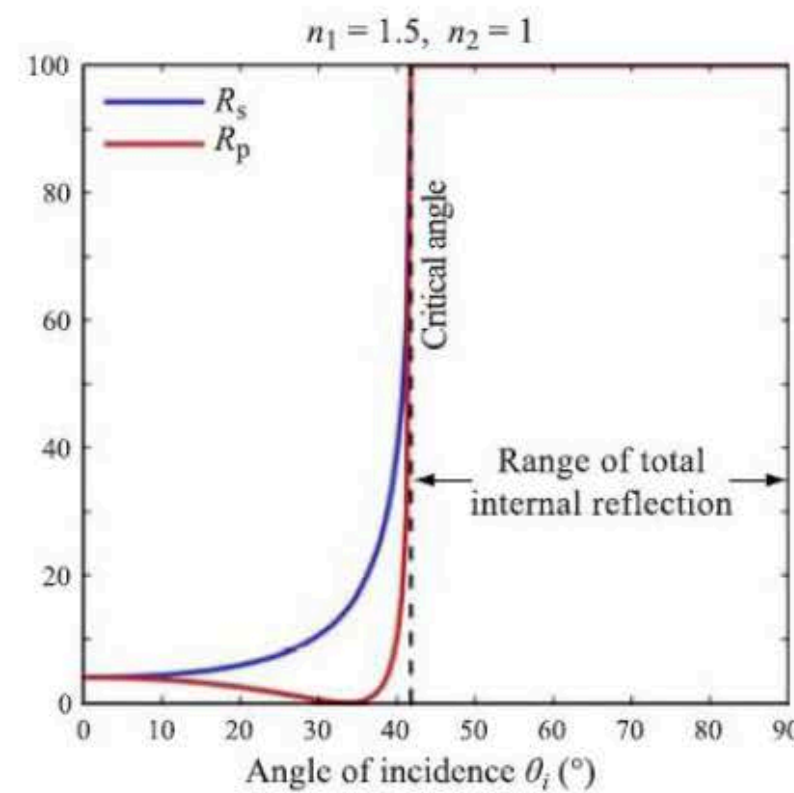
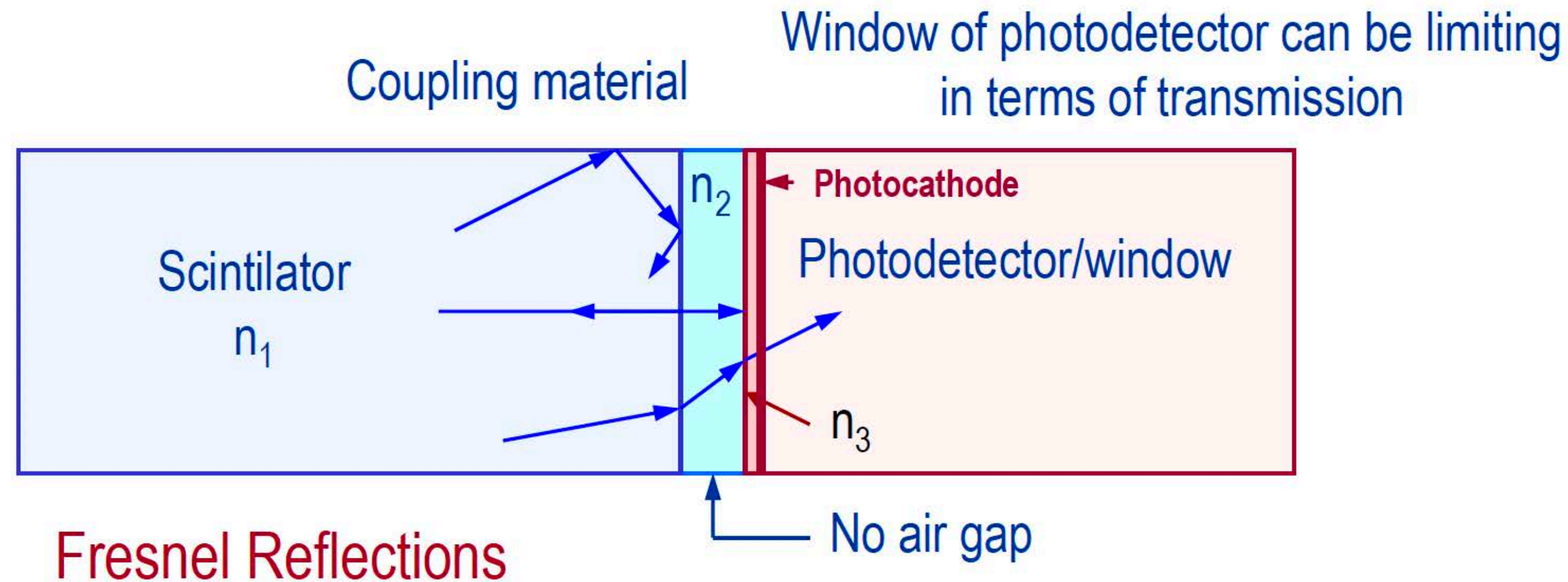
Optimization of the detector elements to harmonize the system

Emission Weighted Quantum Efficiency: $\overline{QE} = \int QE_{PD}(\lambda) \cdot I_{em}(\lambda) d\lambda$



Scintillator Detector Systems - Optical Coupling

Optimization of the detector elements to harmonize the system



Wikipedia

Normal incidence

$$R = \left[\frac{n_2 - n_1}{n_2 + n_1} \right]^2$$

$$n_1 = 1.5, n_2 = 1.0$$

$$R \sim 4\%$$

Coupling materials ($n \sim 1.4$):

- Optical grease/glue
- Silicon "cookies"

$$\Rightarrow n_2 \sim n_3$$

However, for heavy crystals $n_1 > n_2$

Scintillator Detector Systems - Applications

Nuclear and particle physics:

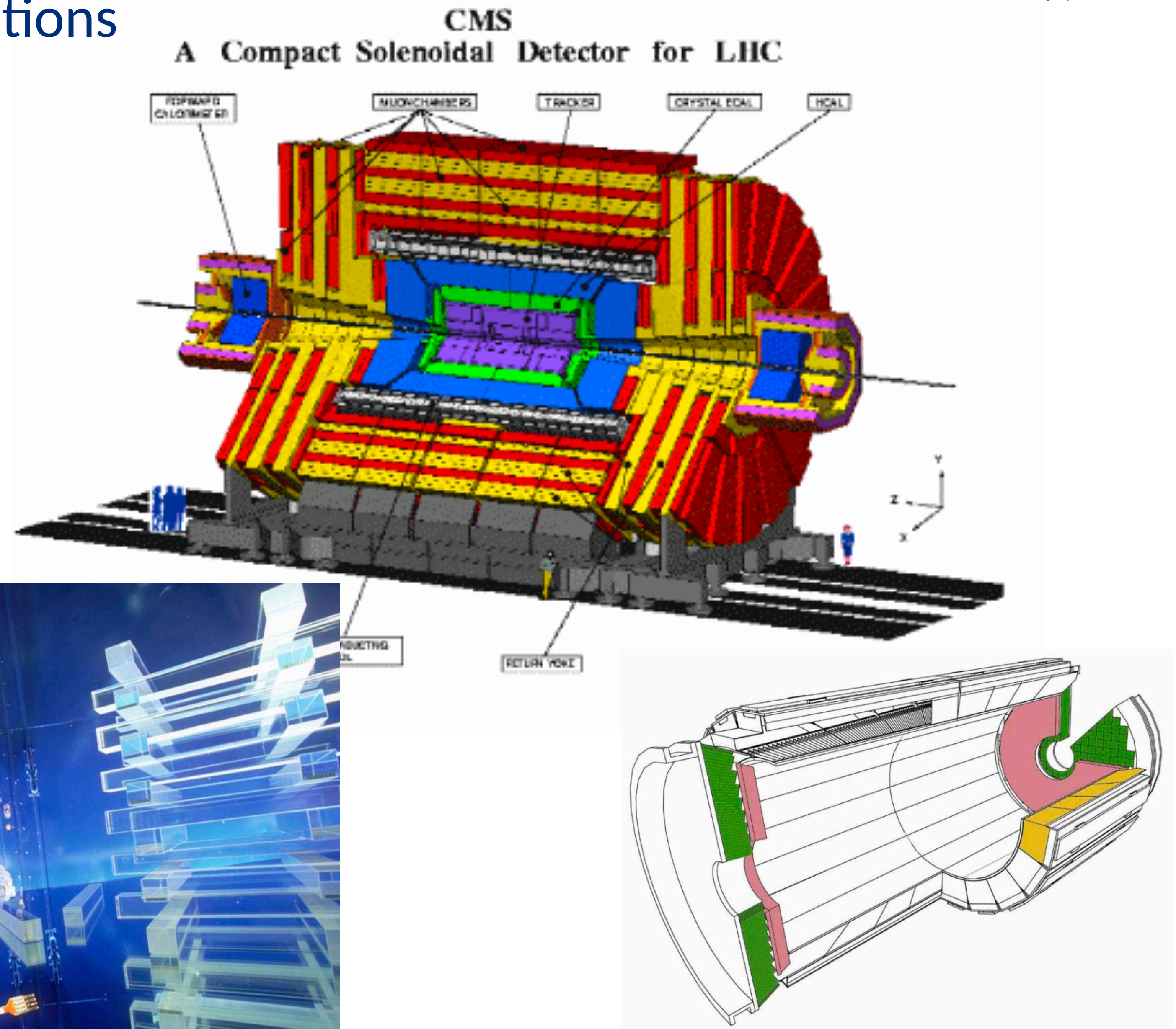
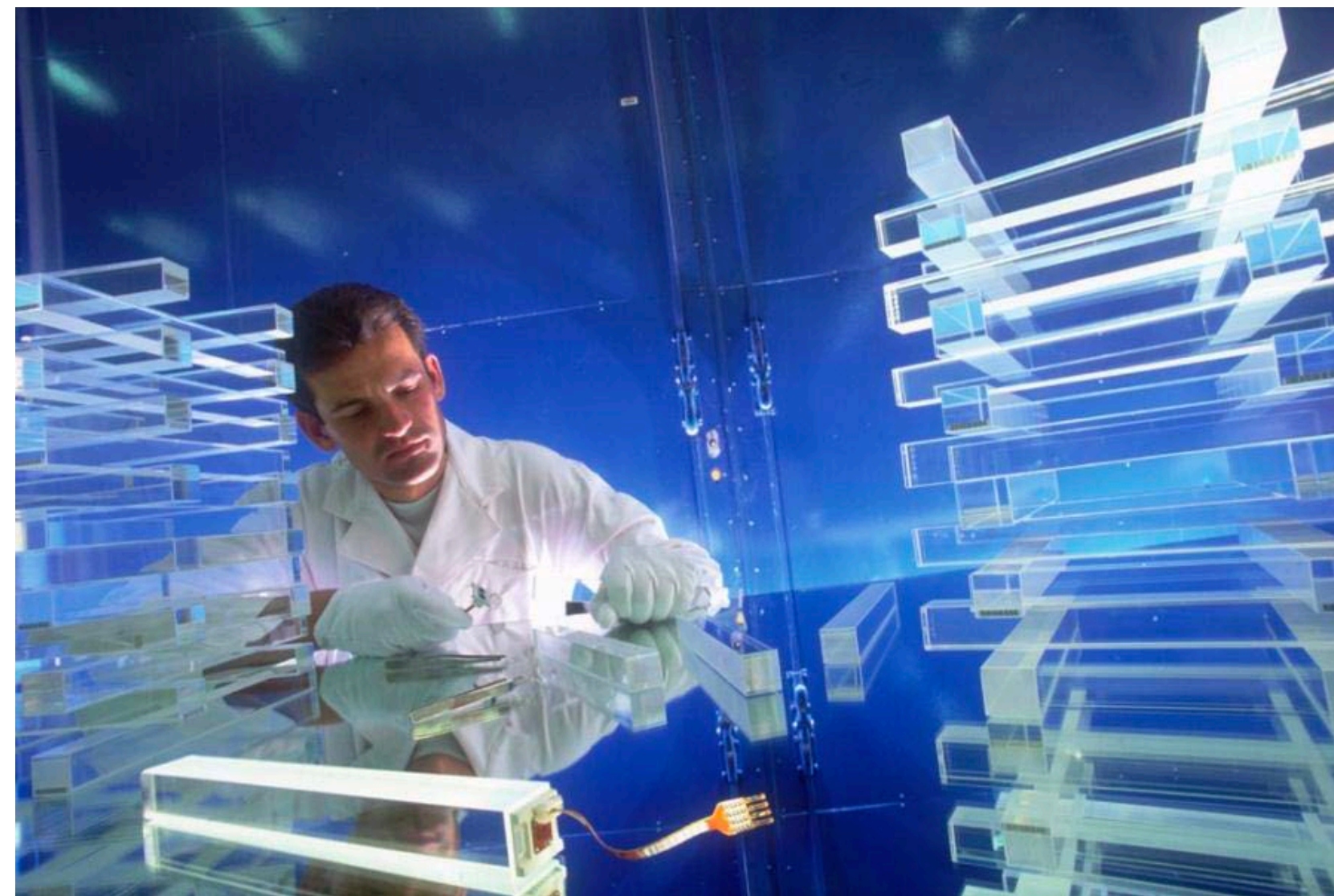
- Calorimetry
- Tracking
- Particle ID (Time of flight)

Astrophysics:

- Dark Matter

Medical Imaging

- PET



Order of 75k PbWO₄ crystals crystals in the endcap (vacuum phototriodes) and barrel (APD) sampling calorimeters

Scintillator Detector Systems - Applications

Nuclear and particle physics:

- Calorimetry
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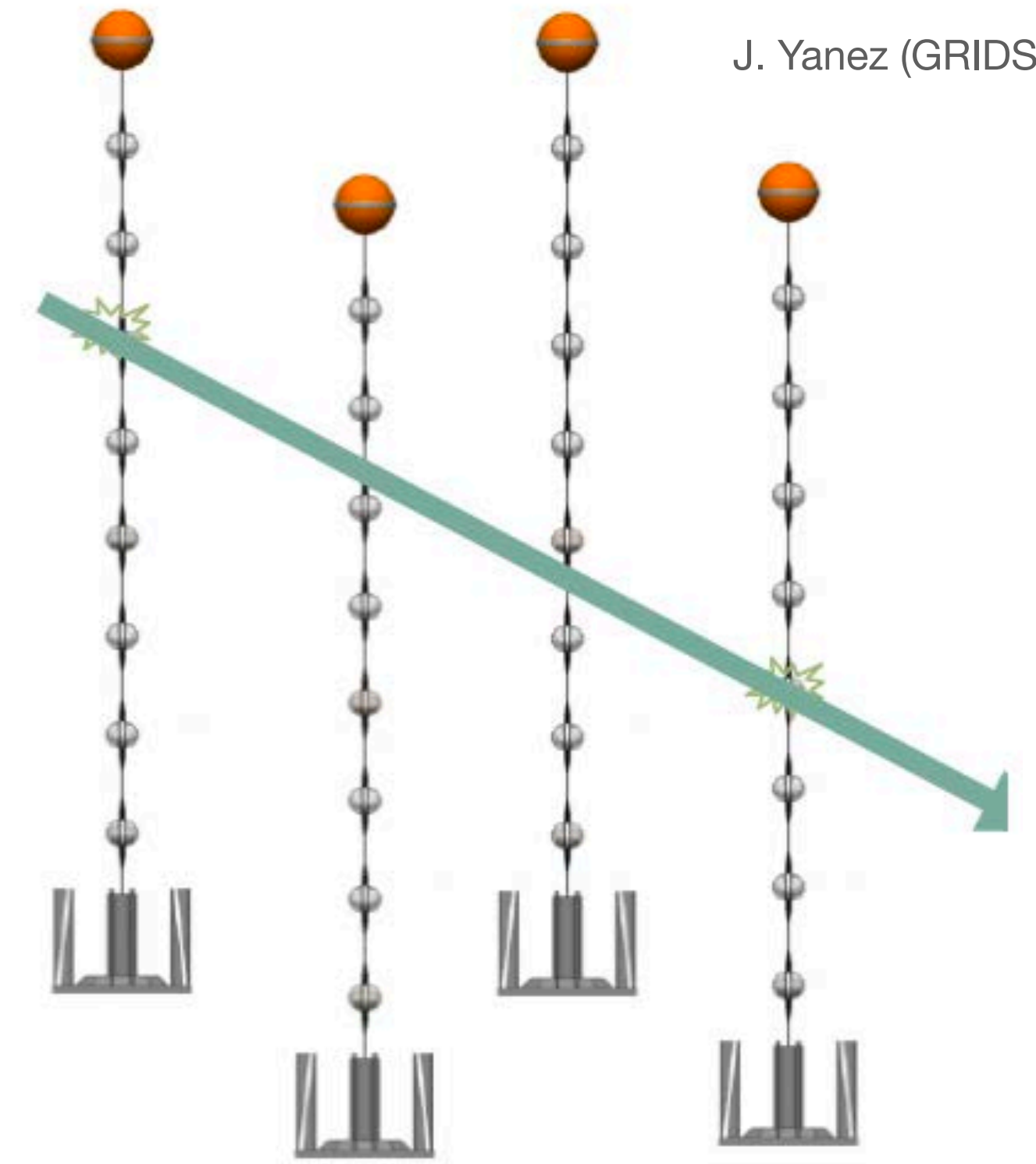
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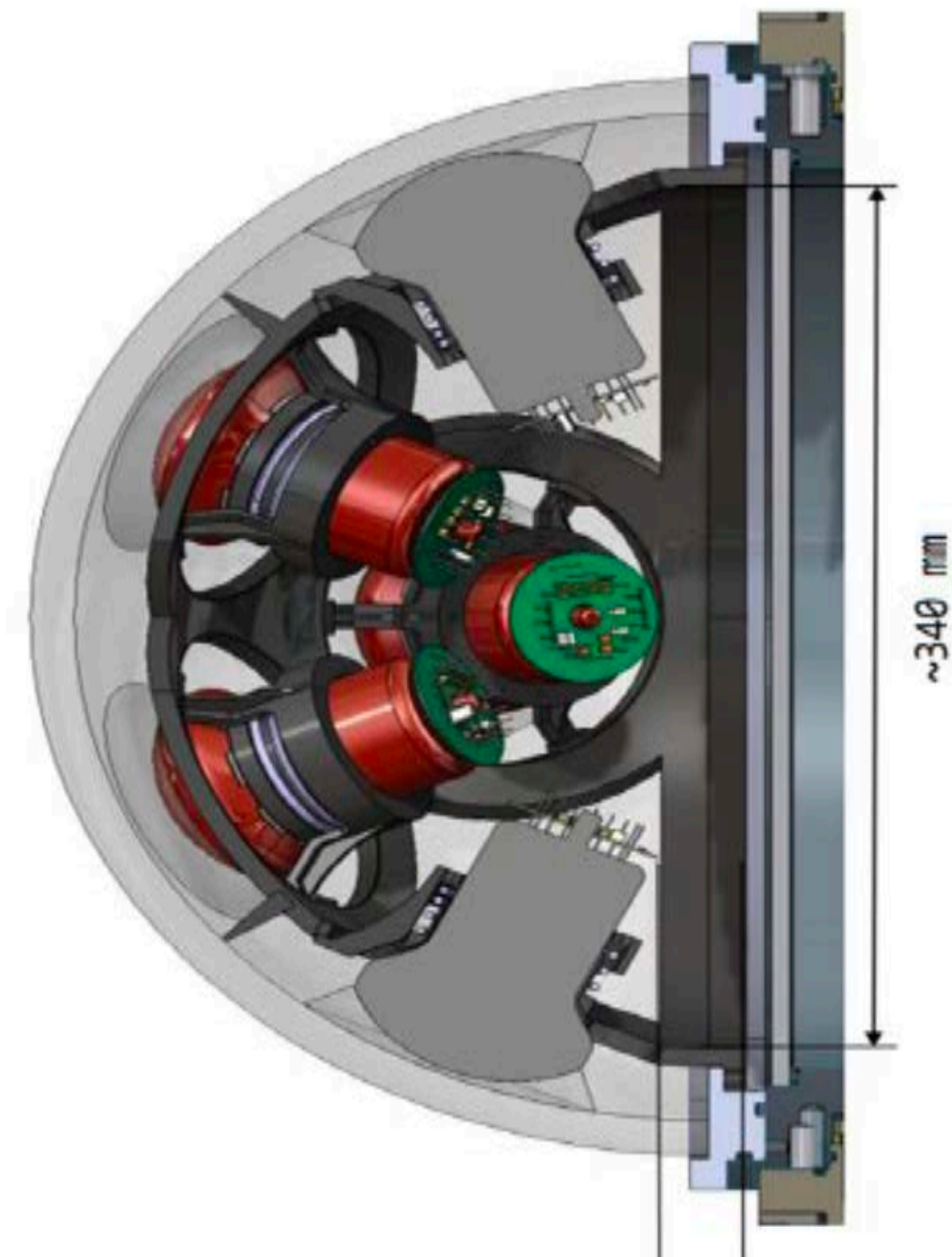
Medical Imaging

- PET

P-ONE



Atmospheric muon calibration w. Plastic scintillator and SiPM readout.



Scintillator Detector Systems - Applications

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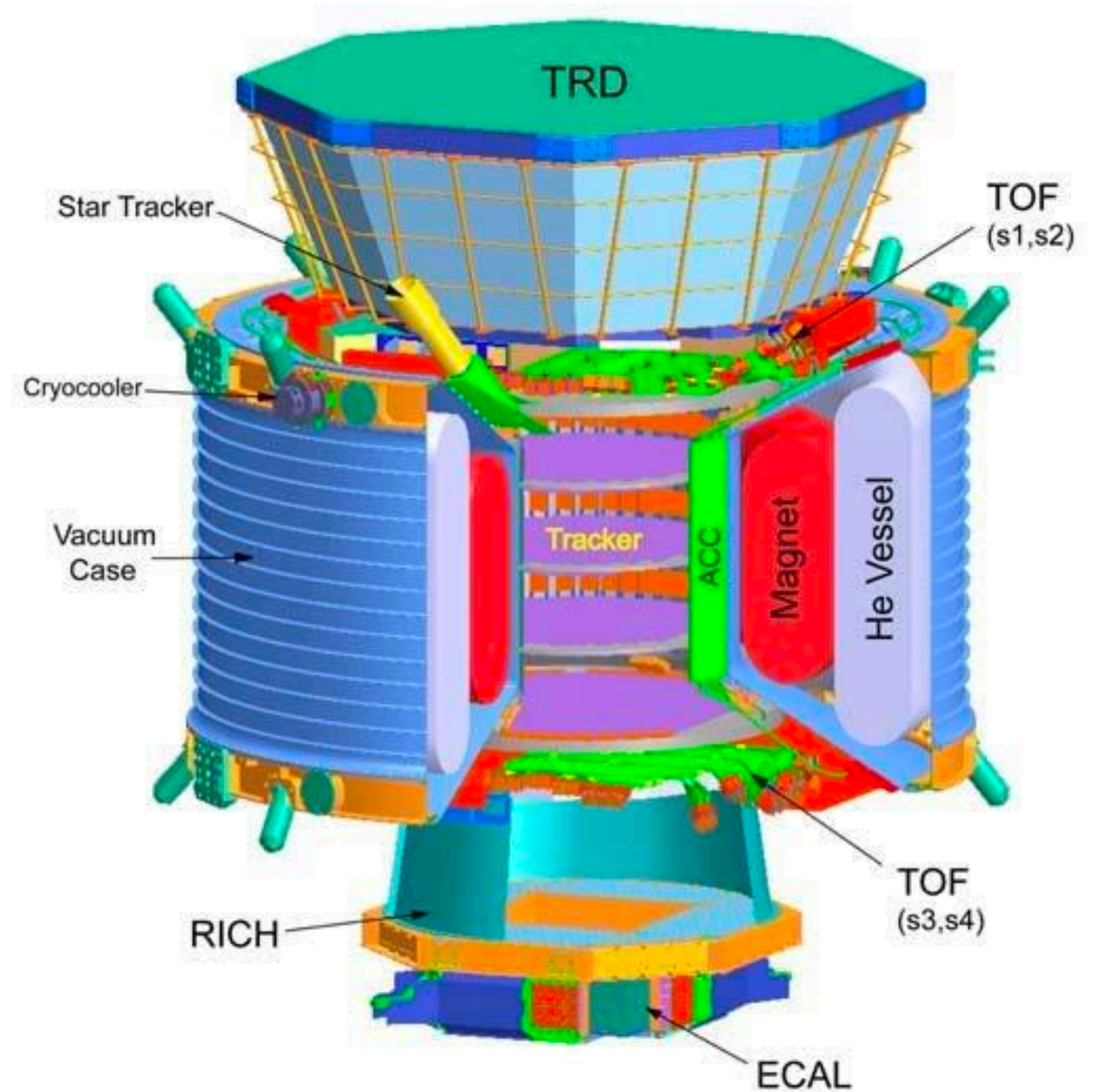
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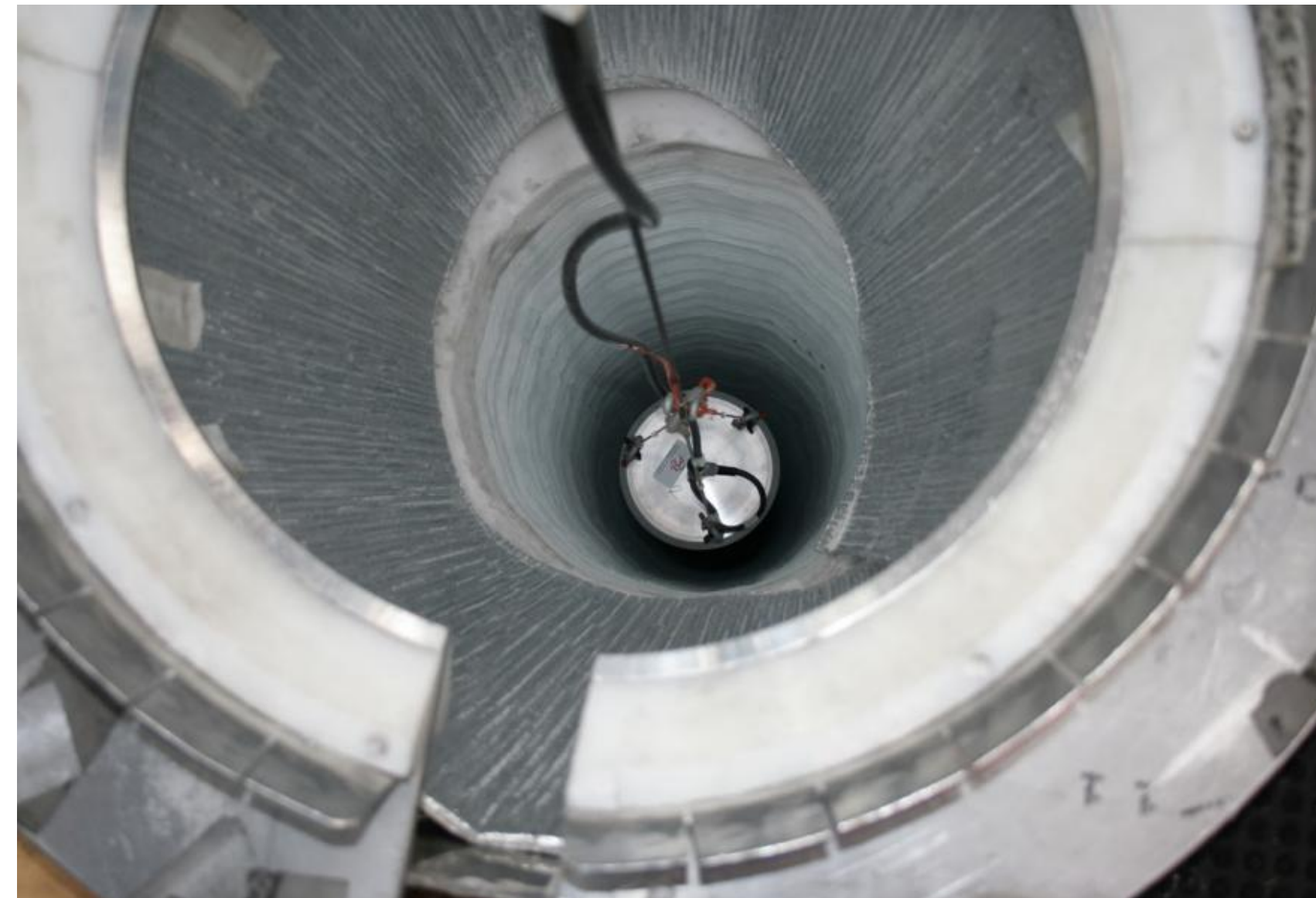
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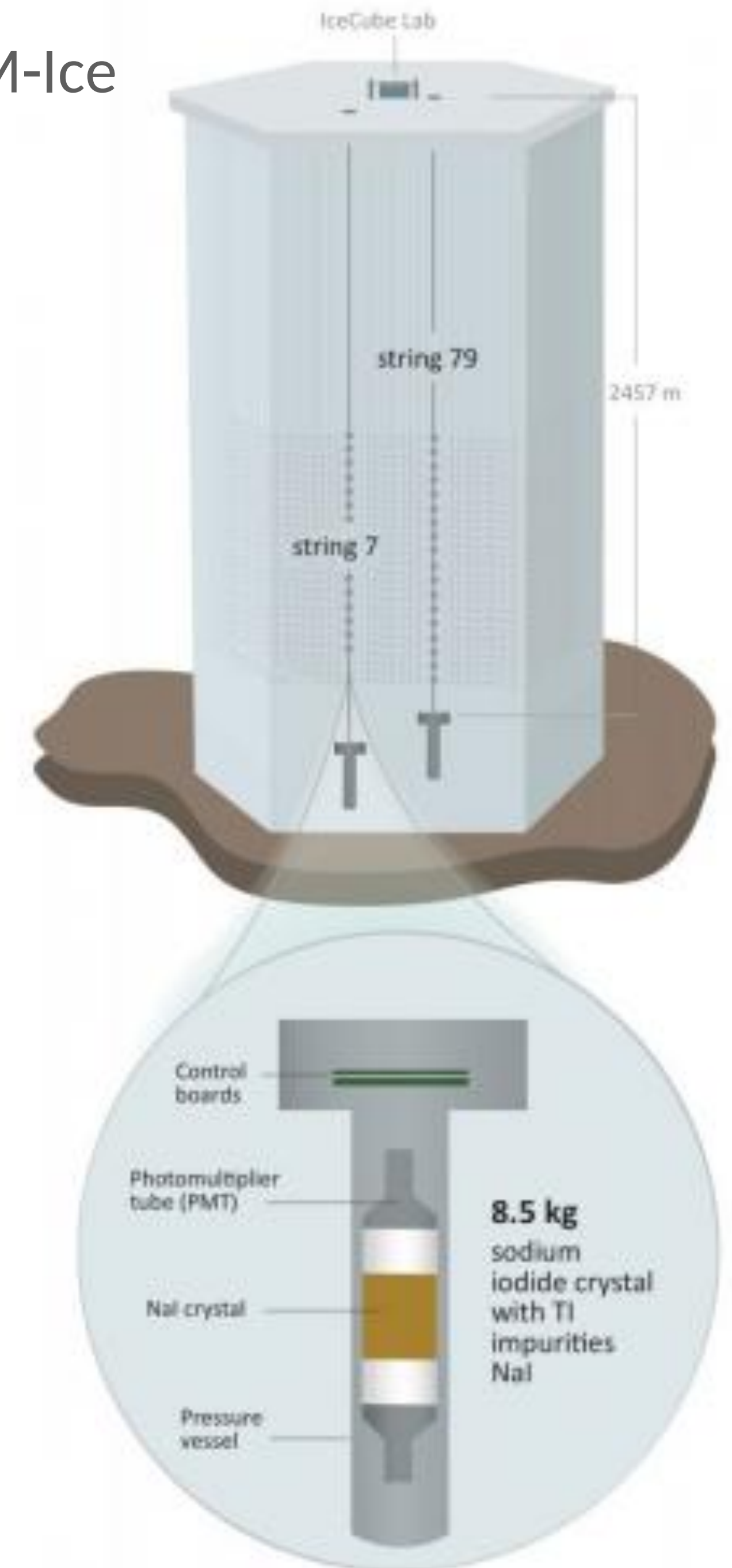
- PET



High-purity NaI(Tl) crystals read-out by PMTs and deployed deep in the Antarctic ice sheet

Tl doping -> impurities on the crystal lattice that behave as activator sites for scintillation

DM-Ice



Scintillator Detector Systems - Applications

Nuclear and particle physics:

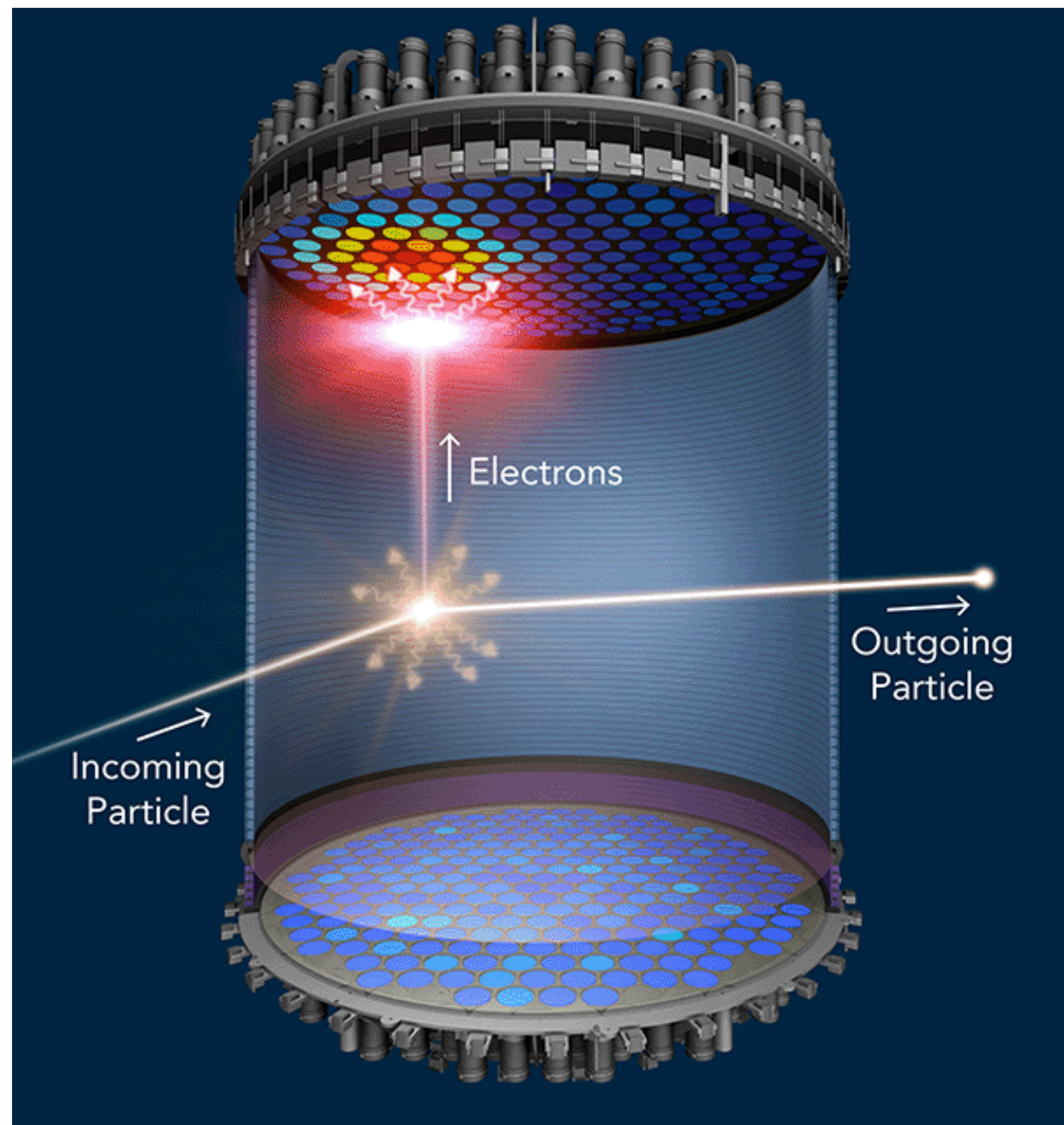
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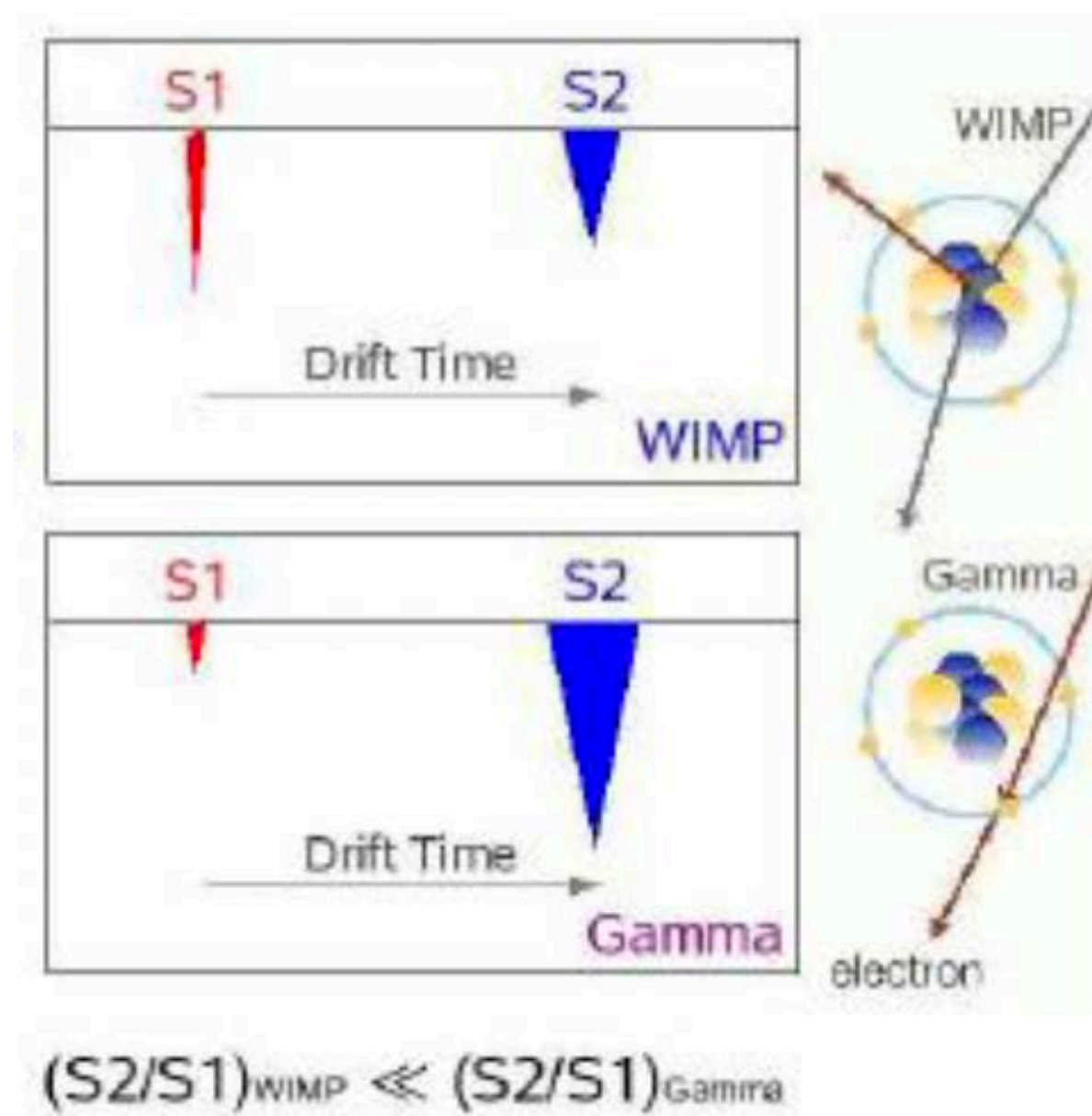
- Dark Matter

Medical Imaging

- PET



Dual-phase time projection chamber



WIMP interaction with nucleus -> nuclear recoil -> large scintillation signal in LXe, small ionization charge

Background (EM) interacts with electrons -> small scintillation signal, large charge

S1 = initial scintillation signal

S2 = ionization charge drifted to the gas phase via large electric field, amplified to produce 2nd delayed scintillation signal

Scintillator Detector Systems - Applications

Nuclear and particle physics:

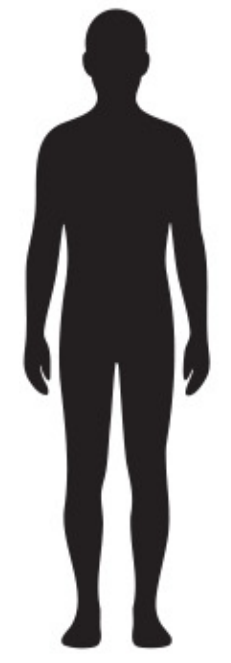
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Astrophysics:

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Medical Imaging

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XENON100, 62 kg



LUX, 375 kg



LZ, 10 ton



XLZD, 60–80 ton

* ask Thomas

Scintillator Detector Systems - Applications

Nuclear and particle physics:

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Ring of Photon Detectors

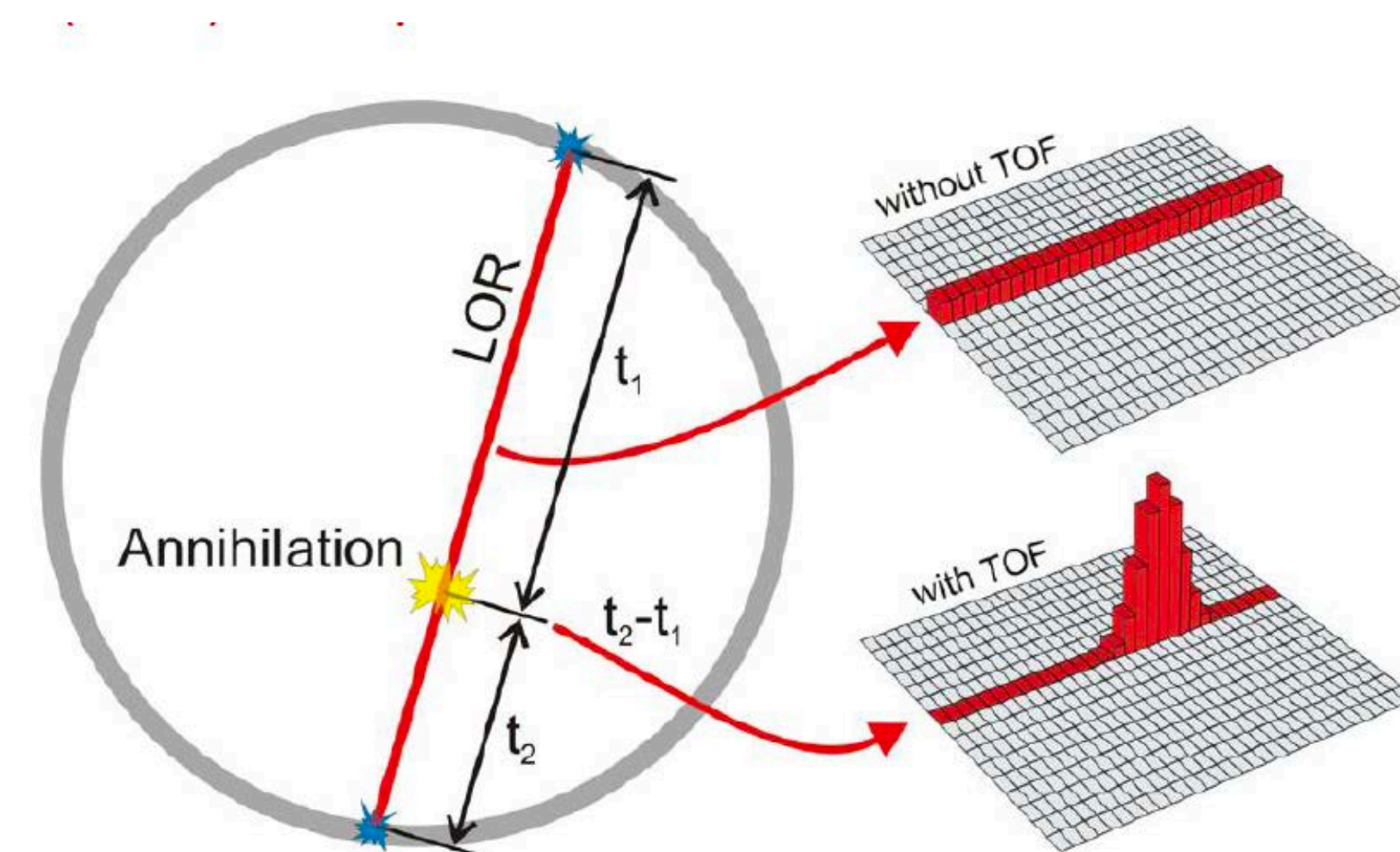
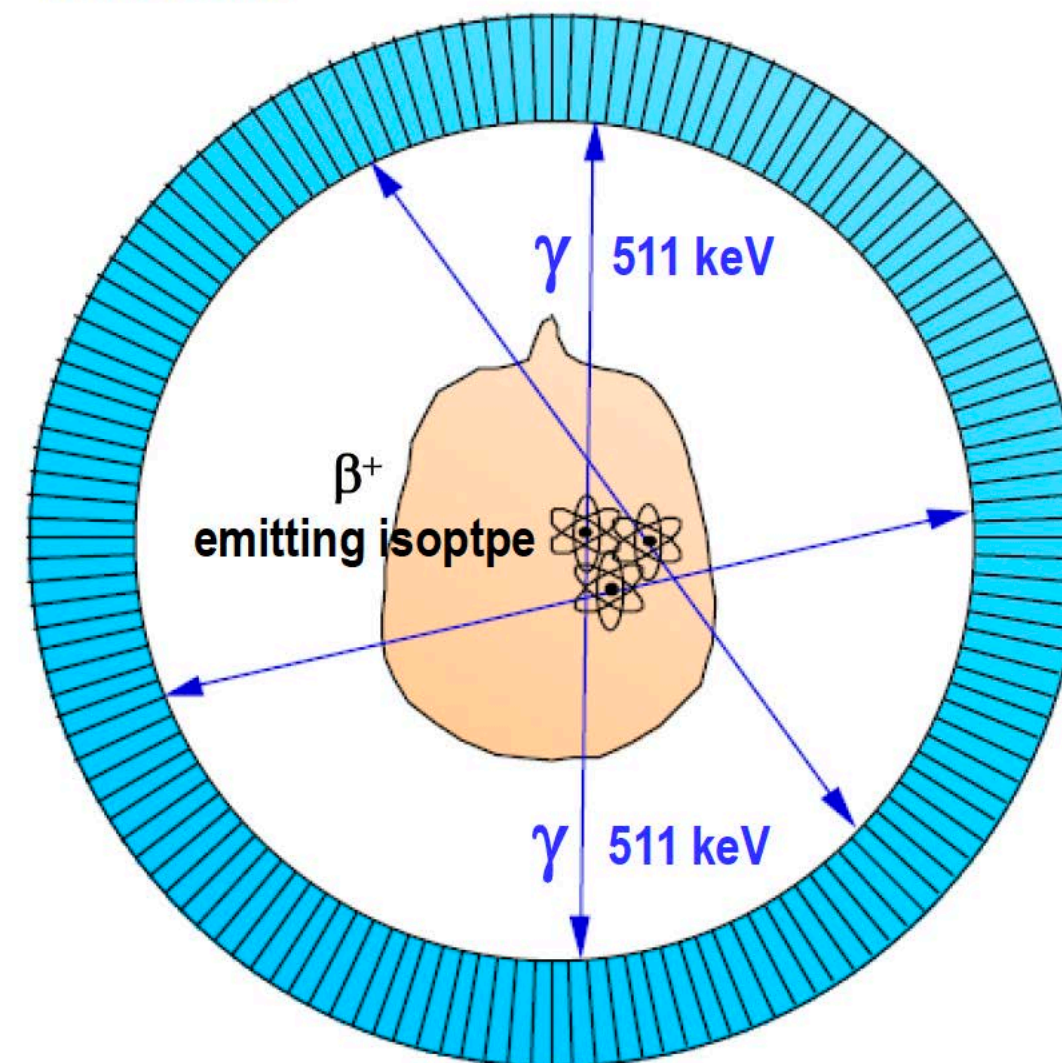
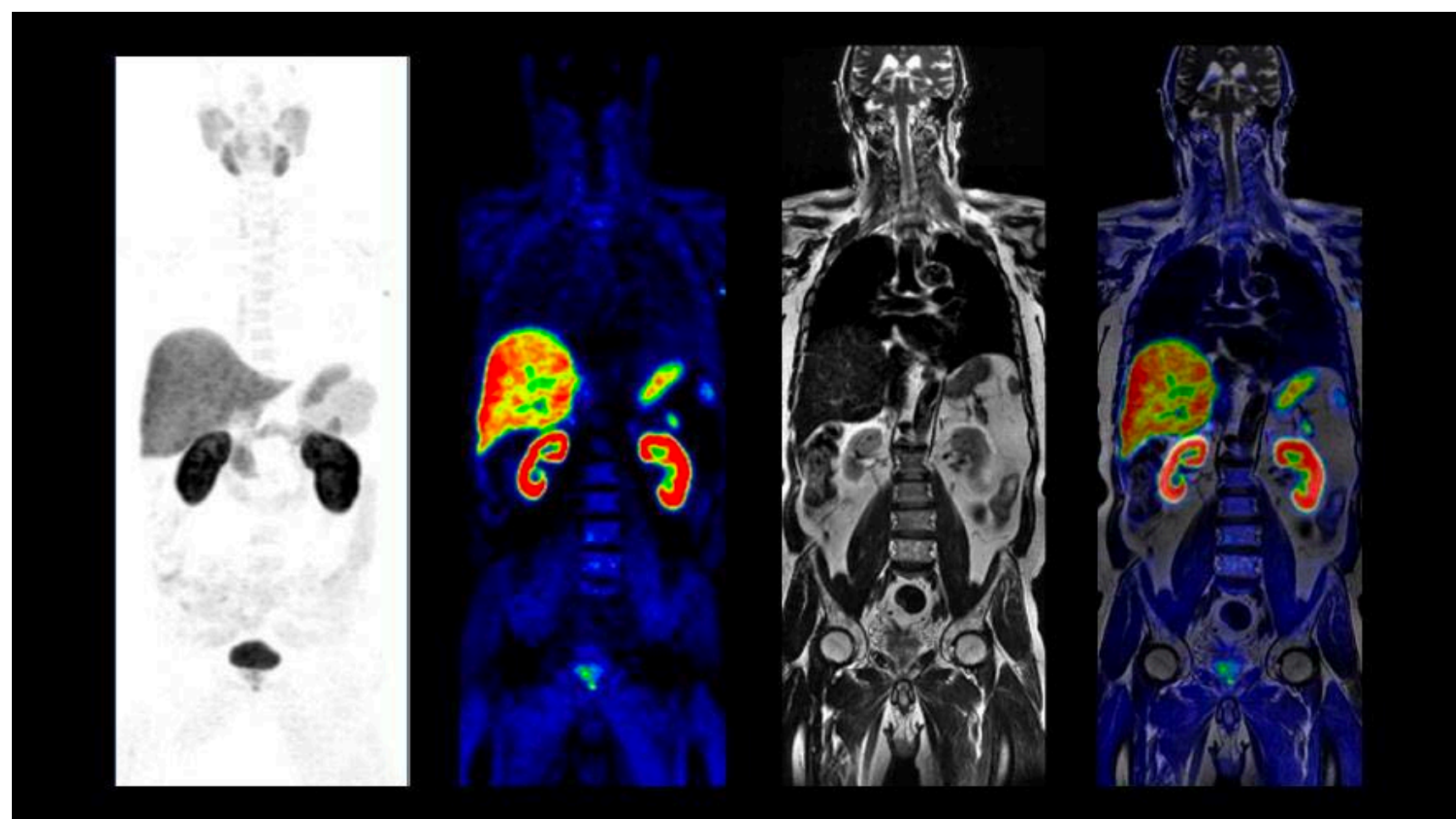


Figure courtesy of W.Moses



Timing is used to mitigate large random backgrounds
Requires fast, bright crystals

Reaching timing resolutions below 10 ps would improve image quality, reduce patient exposure and scan time (in particular for whole-body imaging).



$\sigma_t < 100$ ps
 $\sigma_E < 4\%$ @ 511 keV

Summary

Scintillators are an incredibly versatile particle detector medium

Adaptability provides opportunities for broad use across many areas

Photodetectors are a continuing area of exciting development

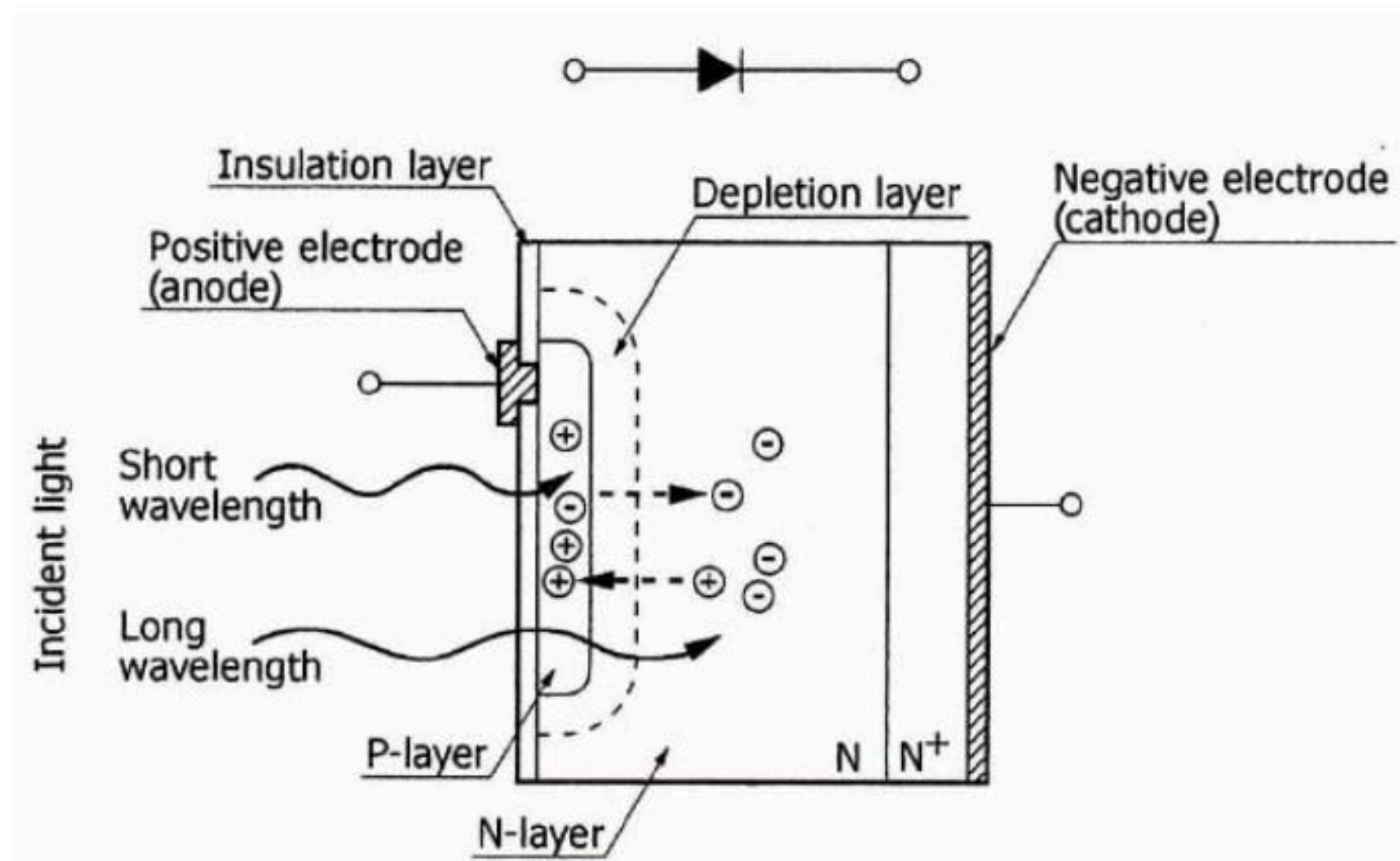
Detector design balances all of the factors that will impact performance:

- Interactions under observation
- Harmonizing the detector elements
- Optimizing signal to noise
- Geometry and scalability
- Environment
- ...

Backup slides

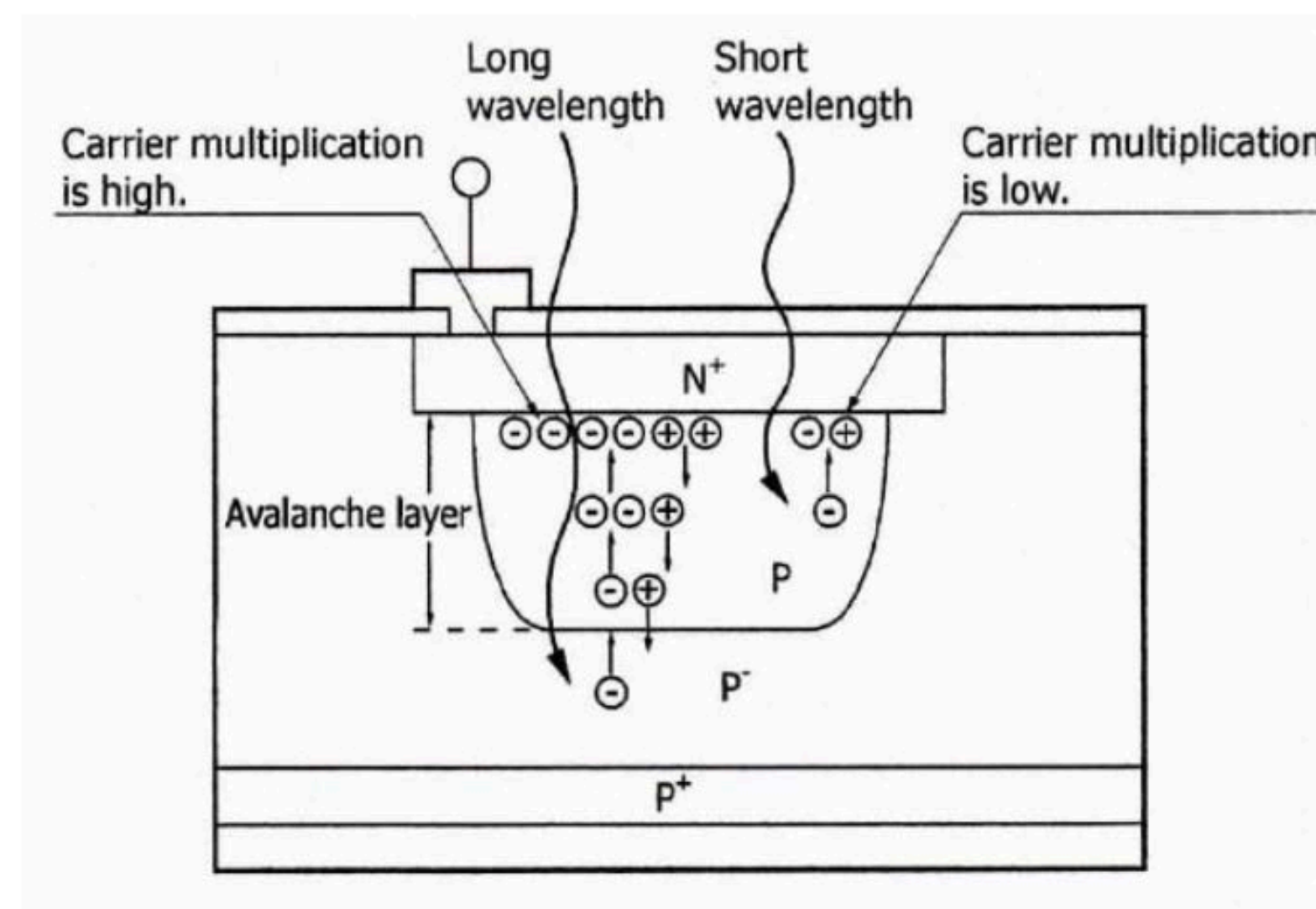
Photodetectors - Silicon (APDs)

PIN photodiodes



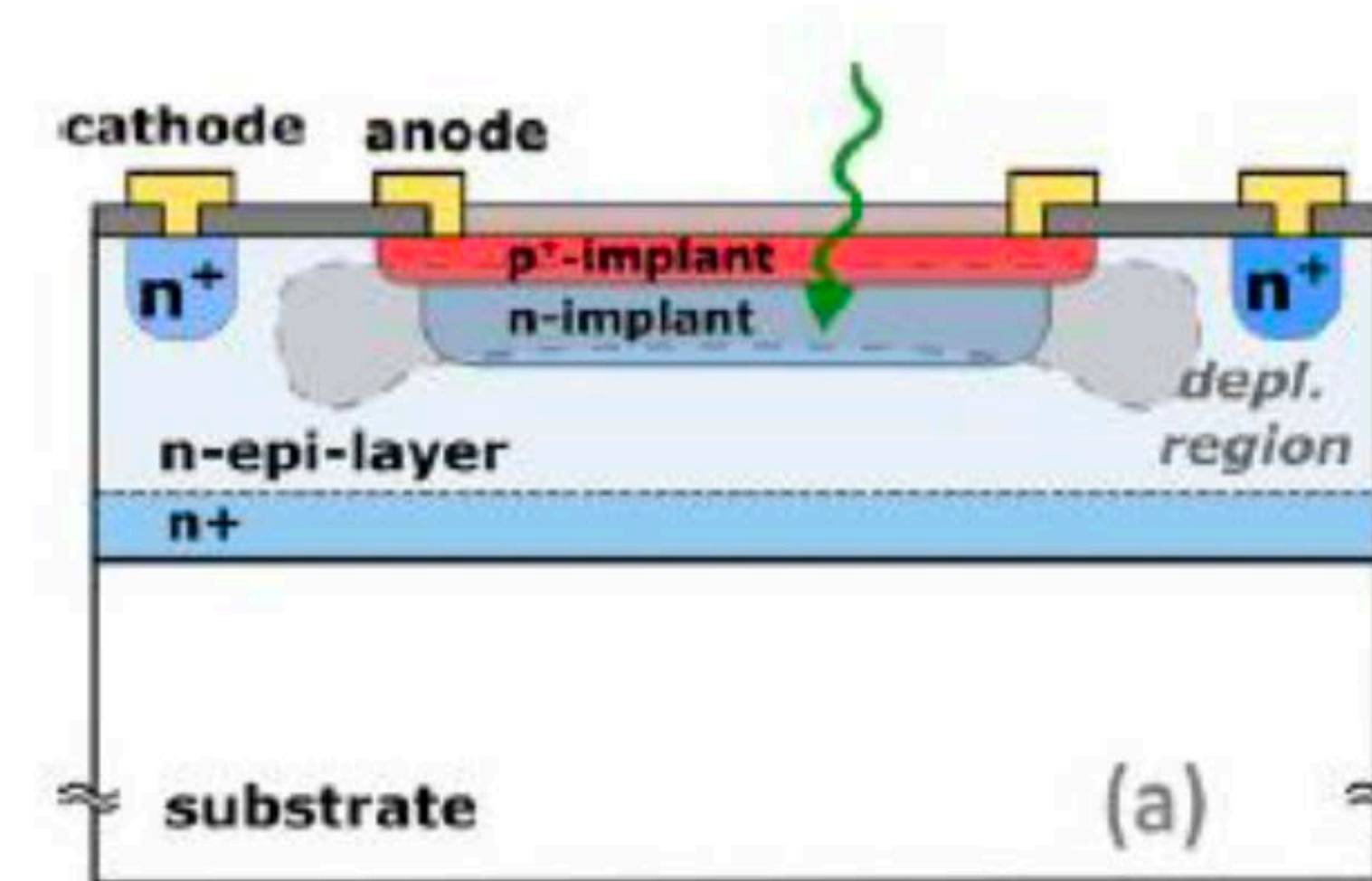
Gain = 1
Linear output

Avalanche Photodiodes (APDs)



Gain = 50 - 500
Bias voltage < Breakdown voltage
Linear output

Geiger APDs (GAPDS)



Gain = 50 - 500
Bias voltage > breakdown voltage
"Quantized" output